1. Compute the expectation value of the Wilson loop  $W_C = e^{ig \oint_C A_\mu d\ell^\mu}$  where C is a closed path lying in the x - y plane and  $A_\mu$  is a U(1) gauge field, assuming that there is a gaussian random field magnetic field pointing in the z direction with  $\langle B_z(\mathbf{r})B_z(\mathbf{r}')\rangle = \sigma^2 \delta^2(\mathbf{r} - \mathbf{r}')$ , the vectors  $\mathbf{r}, \mathbf{r}'$  lying in the x - y plane; how does  $\langle W_C \rangle$  vary as one changes C?

Note: if a variable x is described by an unnormalized probability distribution P(x), then  $\langle f(x) \rangle = \int dx P(x) f(x) / \int dx P(x)$ . In this case the unnormalized probability distribution for  $B_z(\mathbf{r})$  is given by  $\exp(-k \int d^2 r B_z(\mathbf{r})^2)$ , and the integral over x is replaced by a path integral over  $B_z(\mathbf{r})$ . You need to figure out how to perform the required integral, and how to relate the constant k to the parameter  $\sigma$ .

- 2. In class (and in the posted handout) we analyzed a large-N NJL model idimensionally reduced from  $4 \rightarrow 3$  dimensions to illustrate fermion condensation spontaneously breaking a  $U(1)_A$  symmetry. A key ingredient was the introduction of an auxiliary field  $\phi = (\sigma + i\pi)/\sqrt{2}$ . Here I would like you to construct a similar model, but one exhibiting spontaneous breaking of an  $SU(2)_L \times SU(2)_R$  symmetry down to  $SU(2)_V$  by introducing an auxiliary field  $\Phi = \sigma \mathbf{1} + i\pi_a \tau_a$ , where **1** is the unit matrix,  $\tau_a$  are the Pauli matrices with a = 1, 2, 3, and  $\sigma, \pi_a$  are all **real** scalar fields.
  - (a) Start with the Lagrangian

$$\mathcal{L} = \frac{N}{4g} \operatorname{Tr} \Phi^{\dagger} \Phi + \sum_{n=1}^{N} \sum_{i=1,2} \overline{\psi}_{ni} \partial \!\!\!/ \psi_{ni} + \sum_{n=1}^{N} \sum_{ij=1,2} \left( \overline{\psi}_{L,ni} \Phi_{ij} \ \psi_{R,nj} + \text{h.c.} \right) , \qquad (1)$$

where  $\psi$  is a four-component Dirac fermion,  $\psi_R = \frac{1+\gamma_5}{2}\psi$ ,  $\psi_L = \frac{1-\gamma_5}{2}\psi$ , and the  $\gamma$ -matrices are the usual  $4 \times 4$  matrices we use in four dimensions. What makes the theory 3-dimensional is that we assume  $\psi$  and  $\Phi$  are independent of one spatial coordinate, and so momentum integrals are over  $\int d^3p$  rather than  $\int d^4p$ . What is the full symmetry of this theory? Show explicitly how the  $\psi_L$ ,  $\psi_R$ ,  $\sigma$  and  $\pi_a$  fields transform under infinitesimal symmetry transformations.

- (b) What is the fermion Lagrangian if you integrate out the  $\Phi$  field?
- (c) Show that in the large N limit there can be a nontrivial vacuum with  $\langle \sigma \rangle = f \neq 0$  and  $\langle \pi_a \rangle = 0$ , using dim reg and  $\overline{MS}$  to deal with divergences and renormalization, and determining the condition on g for this ground state to be favored. Sketch f versus g.
- (d) What is the symmetry of this ground state? How many Goldstone bosons should there be? How do they transform under the unbroken symmetry?