

1. With  $\sigma^{\mu\nu} = (i/2)[\gamma^\mu, \gamma^\nu]$ , compute how  $\bar{\Psi}\sigma^{\mu\nu}\Psi$  transforms under  $C$ ,  $P$  and  $T$ .
2. On problem set #1 we saw that all fermion bilinears can be expressed as a linear combination of the sixteen bilinears  $\bar{\Psi}\Gamma_a\Psi$  where  $\Gamma_a$  is one of the matrices  $\{1, \gamma_5, \gamma^\mu, \gamma^\mu\gamma_5, \sigma^{\mu\nu}\}$  and  $\text{Tr } \Gamma_a^\dagger \Gamma_b = 4\delta_{ab}$ .
  - (a) Express  $\bar{\Psi}\sigma_{\mu\nu}\gamma_5\Psi$  in terms of the  $\bar{\Psi}\Gamma_a\Psi$ . (Hint: the indices really do tell you how the object transforms under the Lorentz group, so you cannot expect  $\bar{\Psi}\sigma^{\mu\nu}\gamma_5\Psi$  to contain a term like  $\bar{\Psi}\Psi$ .)
  - (b) *Revised.* Define the operator  $S_3 = \int d^3x \bar{\Psi}\sigma^{12}\Psi$ . What is its action on electron states  $|\mathbf{p}, s, \mp\rangle$ , where  $\mp$  refers to electric charge (eg, “−”=particle, “+”= antiparticle),  $s = \pm$  is the spin index, and  $\mathbf{p}$  is the 3-momentum assumed to be nonrelativistic:  $|\mathbf{p}| \ll m$ . You can work out  $S_3$  in the chiral basis, using the  $u$  and  $v$  spinors given in problem set #1, but using indices  $s = 1, -1$  instead of  $s = 1, 2$  so that you can express the  $s$  dependence of your answer simply.
3. Consider a theory of a massless Dirac fermion,  $\mathcal{L} = \bar{\Psi}i\not{\partial}\Psi$ . We know that this is equivalent to a theory of two uncoupled massless Weyl fermions, and therefore should have an independent phase rotation symmetry for each one. In the Dirac language, the two phase symmetries are  $\Psi \rightarrow e^{i\alpha}\Psi$  and  $\Psi \rightarrow e^{i\beta\gamma_5}\Psi$ . This latter symmetry is called a chiral symmetry.
  - (a) Show that  $\Psi \rightarrow e^{i\beta\gamma_5}\Psi$  is a symmetry of the massless Dirac theory, in a basis independent manner.
  - (b) What is the Noether current  $j_5^\mu$  corresponding to this symmetry?
  - (c) Express the conserved charge  $Q_5 = \int d^3x j_5^0$  in terms of the  $b$  and  $d$  creation and annihilation operators. Can you interpret what this chiral charge is physically?
4. Consider the following Lagrangian for two real scalar fields  $\phi_{1,2}$  interacting with a Dirac fermion,

$$\begin{aligned} \mathcal{L} = & \bar{\Psi}(i\not{\partial} - m)\Psi + \frac{1}{2}(-\partial_\mu\phi_1\partial^\mu\phi_1 - \partial_\mu\phi_2\partial^\mu\phi_2 - m_1^2\phi_1^2 - m_2^2\phi_2^2) \\ & + g\phi_1\bar{\Psi}\Psi + ih\phi_2\bar{\Psi}\gamma_5\Psi. \end{aligned} \tag{1}$$

- (a) Show that the interaction terms are Hermitian if coupling constants  $g$  and  $h$  are real.
- (b) Show that this Lagrangian is invariant under  $C$ ,  $P$  and  $T$ , and make a table for how  $\phi_1$  and  $\phi_2$  must transform under these three discrete symmetries.
- (c) Suppose if instead of just the scalar mass terms, we add to  $\mathcal{L}$  a more general scalar interaction  $-V(\phi_1, \phi_2)$ , where  $V$  is a local hermitian polynomial consisting only of terms quartic in scalar fields (no derivatives). Which of the following symmetries can you break with the proper choice of  $V$ :  $C$ ,  $P$ ,  $T$ ,  $CP$ ,  $CT$ ,  $PT$ ,  $CPT$ ?
- (d) Consider the special case of the above Lagrangian with  $m = 0$ ,  $m_1 = m_2 = \mu$ ,  $h = g$ . For these values of the couplings the theory has an exact chiral  $U(1)$  symmetry where  $\Psi \rightarrow e^{i\alpha\gamma_5}\Psi$  and the scalar fields also transform. How must the scalars transform for this to be a symmetry? If you add  $V(\phi_1, \phi_2)$  to this theory (consisting only of quartic terms, as above), what is the most general form of  $V$  that respects this chiral  $U(1)$  symmetry?

5. Prove eq. 42.10 in Srednicki:

$$\int \frac{d^4 p}{(2\pi)^4} \frac{e^{ip(x-y)} f(p)}{p^2 + m^2 - i\epsilon} = i\theta(x^0 - y^0) \int \widetilde{dp} e^{ip(x-y)} f(p) + i\theta(y^0 - x^0) \int \widetilde{dp} e^{-ip(x-y)} f(-p) \quad (2)$$

where in the  $\int \widetilde{dp}$  integrals,  $p^0 = \omega = \sqrt{|\mathbf{p}|^2 + m^2}$ , and

$$\theta(\tau) = \begin{cases} 1 & \tau > 0 \\ 0 & \tau < 0 \end{cases} . \quad (3)$$

Be clear about how you are evaluating contour integrals.