1. Compute:

- (a) Tr (a)
- (b) $\operatorname{Tr}(\gamma_5)$
- (c) $\operatorname{Tr}(\gamma_5 \not a)$
- (d) Tr (a b)
- (e) Tr $(\gamma_5 \not \approx \not b)$
- (f) Tr (\$\$ \$\$ \$\$ \$\$ \$\$ \$\$ \$\$
- (g) $\operatorname{Tr}(\gamma_5 \not a \not b \not c)$
- (i) Tr $(\gamma_5 \not \approx \not \gg \not \ll \not q)$

where $a \equiv a_{\mu}\gamma^{\mu}$. Do not use a particular basis for the gamma matrices, but just the anti-commutation relations $\{\gamma^{\mu}, \gamma^{\nu}\} = -2\eta^{\mu\nu}, \{\gamma_5, \gamma^{\mu}\} = 0, \gamma_5^2 = 1.$

Hints: For some of these objects it is useful to use the cyclic property of a trace; for others it can help to insert $(\gamma_5)^2 = 1$ and make use of $\{\gamma_5, \gamma^{\mu}\} = 0$. The trace of four γ 's and a γ_5 will involve the totally antisymmetric ϵ tensor: $\epsilon^{0123} = 1 = -\epsilon_{0123}$. Note that the 4-index ϵ tensor (in contrast to the 3-index version) is *anti*-cyclic, $\epsilon^{bcda} = -\epsilon^{abcd}$.

- 2. Define the sixteen matrices $\Gamma_a = \{1, \gamma_5, \gamma^{\mu}, \gamma^{\mu}\gamma_5, \sigma^{\mu\nu}\}$ with $a = 1 \dots, 16$, accounting for all the different values of the μ and ν indices. Show that $\operatorname{Tr} \Gamma_a^{\dagger} \Gamma_b = 4\delta_{ab}$. You can assume γ^0 is Hermitian, while γ^i is anti-Hermitian, which is consistent with $\{\gamma^{\mu}, \gamma^{\nu}\} = -2\eta^{\mu\nu}$. Explain therefore why a complete set of bilinears on can make from $\overline{\Psi}$ and Ψ are the sixteen quantities $\overline{\Psi}\Gamma_a\Psi$. How do these tensors transform under the Lorentz group? (eg as scalar, four vector...).
- 3. Consider the spinors in the chiral basis in which γ_5 is diagonal:

$$\gamma^{\mu} = \begin{pmatrix} 0 & \sigma^{\mu} \\ \bar{\sigma}^{\mu} & 0 \end{pmatrix} , \qquad \gamma_5 = \begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix} , \qquad (1)$$

where $\sigma^{\mu} = \{1, \vec{\sigma}\}$ and $\bar{\sigma}^{\mu} = \sigma_2(\sigma^{\mu})^T \sigma_2 = \{1, -\vec{\sigma}\}$, and

$$v_1(\mathbf{p}) = \frac{\not p + m}{\sqrt{2(m+\omega)}} \begin{pmatrix} 0\\1\\0\\-1 \end{pmatrix}, \quad v_2(\mathbf{p}) = \frac{\not p + m}{\sqrt{2(m+\omega)}} \begin{pmatrix} -1\\0\\1\\0 \end{pmatrix}$$
(3)

where $p^0 = \omega = \sqrt{m^2 + |\mathbf{p}|^2}$ with \mathbf{p} being the 3-momentum, and $\not p = p_\mu \gamma^\mu$. Also define spinors $\bar{u}_r = u_r^{\dagger} \gamma^0$ and $\bar{v}_r = v_r^{\dagger} \gamma^0$.

(a) Show that

$$(\not\!\!\!p + m)u_r(\mathbf{p}) = (-\not\!\!\!p + m)v_r(\mathbf{p}) = 0 \tag{4}$$

(b) Define the spin operator $S_z = \sigma^{12}$, where $\sigma^{\mu\nu} = \frac{i}{4} [\gamma^{\mu}, \gamma^{\nu}]$; show that in the particle rest frame $(\mathbf{p} = 0)$

$$S_z u_1(0) = \frac{1}{2} u_1(0) , \quad S_z u_2(0) = -\frac{1}{2} u_2(0) , \quad S_z v_1(0) = -\frac{1}{2} v_1(0) , \quad S_z v_2(0) = \frac{1}{2} v_2(0) .$$
 (5)

(c) Prove Srednicki eq. (38.21):

$$\bar{u}_{s'}(\mathbf{p})\gamma^{\mu}u_s(\mathbf{p}) = 2p^{\mu}\delta_{s's} , \qquad \bar{v}_{s'}(\mathbf{p})\gamma^{\mu}v_s(\mathbf{p}) = 2p^{\mu}\delta_{s's}$$
(6)

(d) Prove Srednicki eq (38.23):