

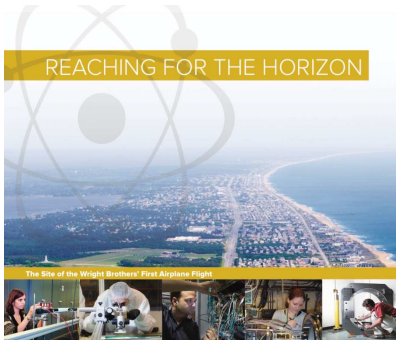
Extracting Resonance Parameters
from $\gamma^* p \rightarrow n\pi^+$ at CLAS

Kijun Park

Nov. 13-18, 2016

- 1 Introduction
- 2 Physics Result Highlight
- 3 New Interesting Results !
- 4 Summary

Long Range Plan 2015

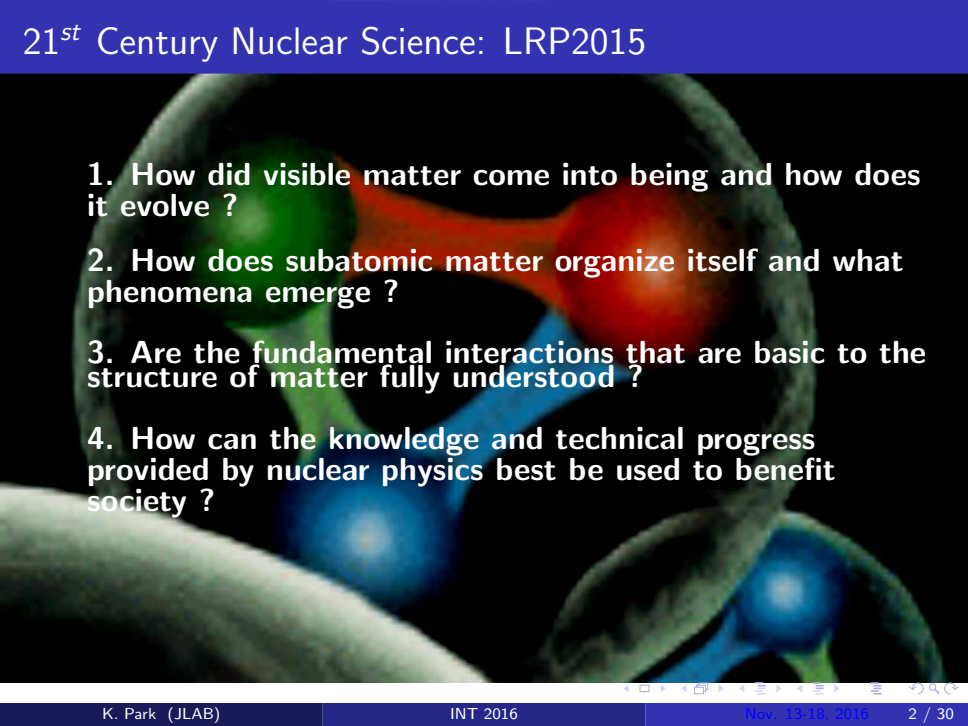


The 2015 LONG RANGE PLAN for NUCLEAR SCIENCE

<http://science.energy.gov/np/reports>



1. Fully utilize programs at existing & under construction facilities (JLab12, RHIC, NSCL, FRIB,...)
2. Sustain a targeted program in fundamental symmetries & neutrino research
3. Invest in a ton-scale neutrino-less double beta decay experiment
4. Construct a high-energy high-luminosity polarized **EIC** with highest priority following the completion of FRIB
5. Invest in mid- and small-scale projects at universities and laboratories

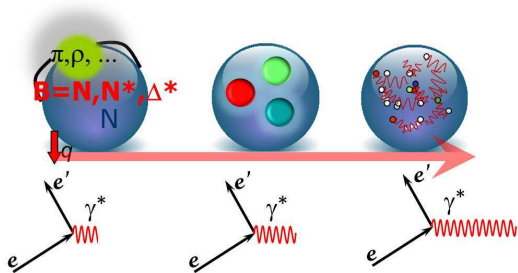
- 
1. How did visible matter come into being and how does it evolve ?
 2. How does subatomic matter organize itself and what phenomena emerge ?
 3. Are the fundamental interactions that are basic to the structure of matter fully understood ?
 4. How can the knowledge and technical progress provided by nuclear physics best be used to benefit society ?



How can we approach these questions ?

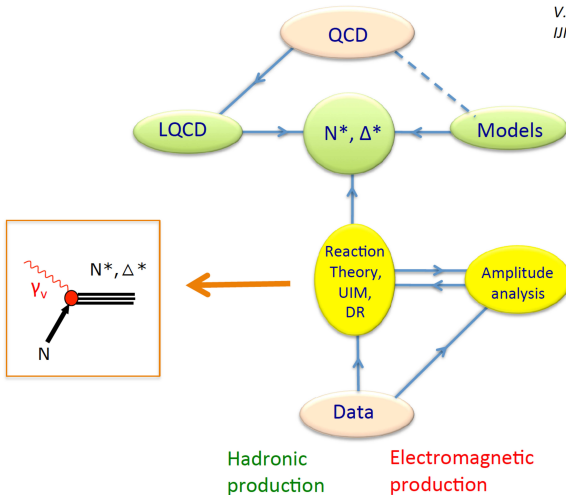
The most challenging problems in Hadron Physics

- Non-perturbative **DCSB** generates more than 98% of **dress quark masses** as well as **dynamical structure**
 - although, Higgs mechanism $< 2\%$ in N , N^* masses
- Quark-gluon **confinement** in baryons emerges from QCD
 - dressed quarks, meson-baryon cloud, dressed gluon,...
- Study of the excited states of the nucleon is important step in the development of a fundamental understanding of strong interaction —
[N. Isgur, V. Burkert (2000)]
- The most fundamental question: “ **WHAT ARE THE RELEVANT DEGREE-OF-FREEDOM AT VARYING DISTANCE SCALE ?** ”

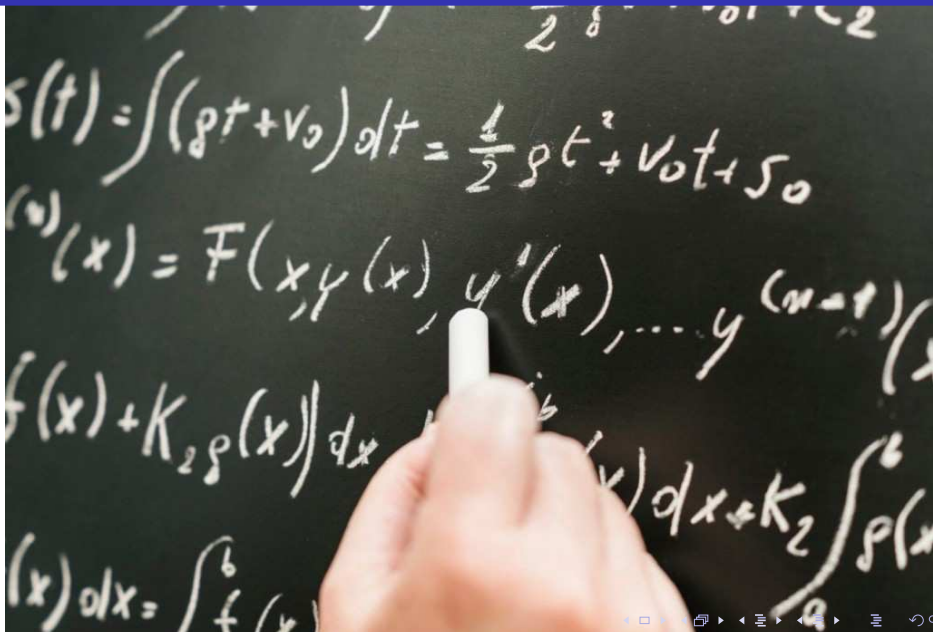


Analysis Chain

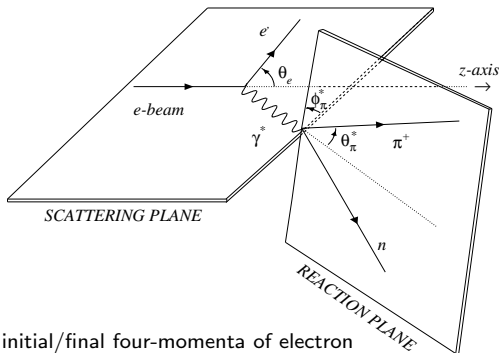
V.B., T.-S.H. Lee
IJMP E13 (2004) 1035



Let me talk a little bit about kinematics



Reaction, $\vec{e}p \rightarrow e'\pi^+n$



- $k_{i,f}$: the initial/final four-momenta of electron
- $E_{i,f}$: the initial/final energy of electron
- θ_e : the electron scattering angle
- $p_{\gamma,i}$: the virtual photon/target four-momenta
- $W^2 = (p_\gamma + p_i)^2 = M_p^2 + 2M_p\nu - Q^2$
- ν : transferred energy $= E_i - E_f = \frac{p_i \cdot p_\gamma}{M_p}$
- Q^2 : virtuality of the exchanged photon $= -(k_i - k_f)^2 = 4E_i E_f \sin^2(\theta_e/2)$
- θ_π^* : the angle between the virtual photon and the hadron (π^+)
- ϕ_π^* : the angle between the electron scattering plane and the hadronic production plane

Formalism, $\vec{e}p \rightarrow e'\pi^+n$

- assume: one photon exchange approximation

$$\frac{d^5\sigma}{dE_f d\Omega_e d\Omega_\pi^*} = \Gamma_\nu \cdot \frac{d^2\sigma}{d\Omega_\pi^*}$$

where,

$$\Gamma_\nu: \text{virtual photon flux: } \frac{\alpha}{2\pi^2 Q^2} \frac{(W^2 - M_p^2) E_f}{2M_p E_e} \frac{1}{1-\epsilon},$$

$$\epsilon: \text{virtual photon polarization: } \left(1 + 2 \left(1 + \frac{\nu^2}{Q^2}\right) \tan^2 \frac{\theta_e}{2}\right)^{-1}$$

$$\frac{d^2\sigma}{d\Omega_\pi^*} = \frac{p_\pi^*}{k_\pi^*} \left(\sigma_0 + h \sqrt{2\epsilon(1-\epsilon)} \sigma'_{LT} \sin \theta_\pi^* \sin \phi_\pi^* \right)$$

$$\sigma_0 = \sigma_U + \epsilon \sigma_{TT} \sin^2 \theta_\pi^* \cos 2\phi_\pi^* + \sqrt{2\epsilon(1+\epsilon)} \sigma_{LT} \sin \theta_\pi^* \cos \phi_\pi^*$$

where,

h : beam helicity state

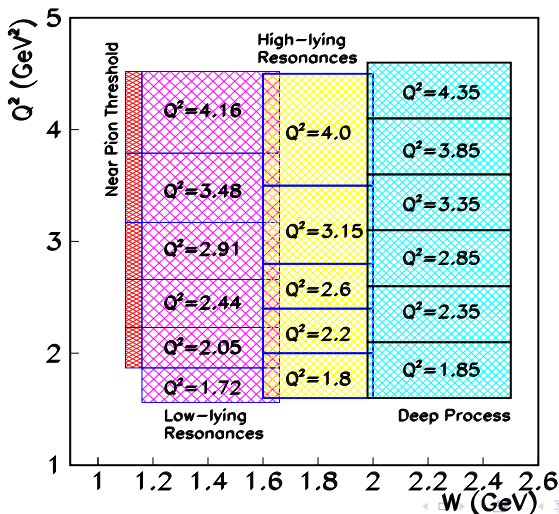
σ_0 : unpolarized cross-section

$$\sigma_U = \sigma_T + \epsilon \sigma_L$$

Kinematics is completely defined by five variables (Q^2 , W , θ_π^* , ϕ_π^* , and ϕ_e)

Kinematic coverage, $E_0=5.49, 5.75$ GeV, $P_e \sim 70\%$

- Kinematic range $W(\text{excitation})$, $Q^2(\text{resolution})$ of $\gamma^* p \rightarrow n\pi^+$
- From the **near pion threshold** to **Deep Process** regime



Let me talk about the highlighted results briefly...



- Two different approaches: **UIM, DR**

UIM

- BG UIM is built from nucleon exchange in s -, u - and π, ω, ρ exchange in t - channel
- Unitarization of multipole amplitudes in the K -matrix approximation
- Resonance contributions are parameterized in the unified BW form with energy dependence

DR

- Fixed- t dispersion relation for the invariant amplitude
- Re -Amplitude to Born-term (nucleon exchange in s -, u -, π exchange in t -channel)
- Integral Im -Amplitude with the isospin structure

- **Two model-uncertainties**

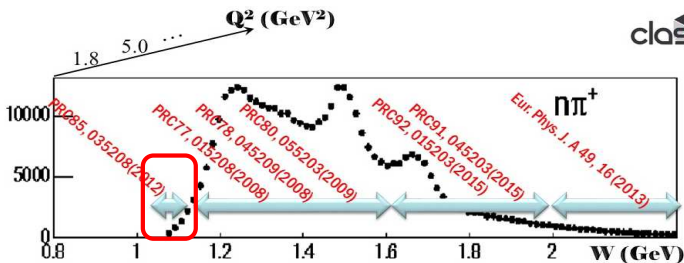
- 1/ BG determination in the UIM and Born term in DR
- 2/ A width and mass of resonances from PDG

- **Take into account...**

- 1/ All(13) **** and *** states in the 1st, 2nd, 3rd
- 2/ $\Delta(1905)F_{35}$, $\Delta(1950)F_{37}$ in 4th resonance region

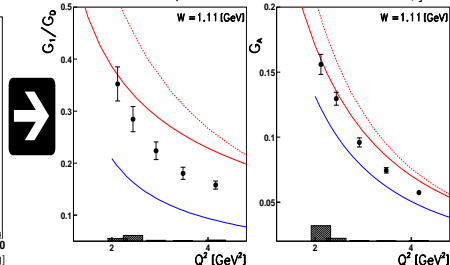
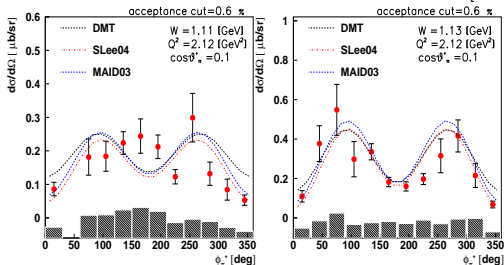
- Same BR from PDG2012

Near threshold ($W < 1.15$ GeV)



- Generalized form factor (G_1) and Axial Form Factor (G_A) near pion threshold
- Multipole fit vs. LCSR, Both are consistent result in lowest W

[red solid: LCSR+FF, dash: pure LCSR, blue solid: MAID07 ↓]

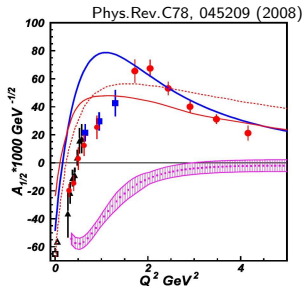
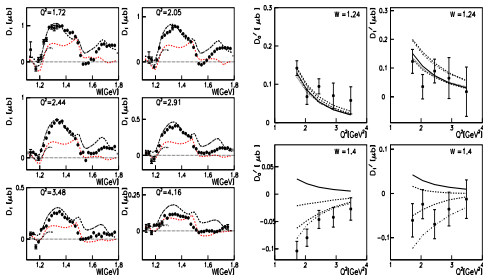


$\vec{e}p \rightarrow e'\pi^+n$ for low lying N^* ($W = 1.15 - 1.69$ GeV)

- Transition Form Factors for $N(1440)1/2^+$ (old conv: $P_{11}(1440)$)
- $A_{1/2}$ shows a sign change in $Q^2 \sim 0.8$ GeV²
- $S_{1/2}$ is large at low Q^2 and drop off smoothly with increasing Q^2
- A complex interplay btw inner core of quarks in the first radial excitation and external MB cloud
- Quark core in DSEQCD (thick blue curve), MB cloud contribution (purple band)

$N\pi$ loops MB, running quark mass (red solid curve)

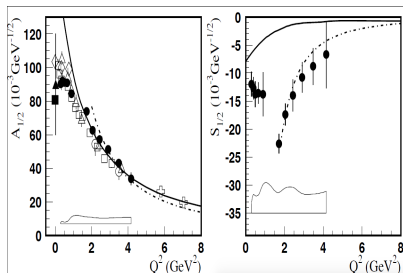
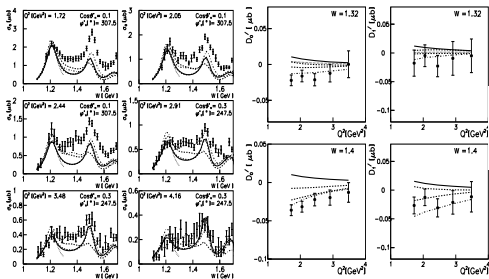
$N\sigma$ loops MB, fixed constituent quark mass (red dashed curve) [\Downarrow including single π and 2π data]



$\vec{e}p \rightarrow e'\pi^+n$ for low lying N^* ($W = 1.15 - 1.69$ GeV)

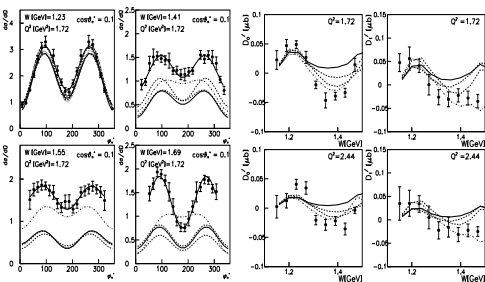
- **Transition Form Factors for $N(1535)1/2^-$ (old conv: $S_{11}(1535)$)**
- $\beta_{N\eta}^{PDG} = 0.45 - 0.60 \rightarrow \beta_{N\pi}^{PDG} = 0.485$ & $\beta_{N\eta}^{PDG} = 0.460$, excellent agreement
- Sensitive to long. as well (strong interference S_{11} - P_{11})
- Previously Opposite sign of $S_{1/2}$! \rightarrow Impossible to change in quark model (LFRQM failed for $S_{1/2}$!)
 \rightarrow Combined with the difficulties in the description of
 (1) large width of $S_{11}(1535) \rightarrow \eta N$
 (2) large $S_{11}(1535) \rightarrow \phi N, \Lambda K$ couplings
 \rightarrow It shows that $3q$ picture for $S_{11}(1535)$ should be complemented ! [I.Aznuryan]

[\downarrow solid: LFRQM, dash-dot: LCSR]

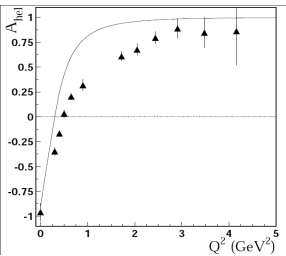
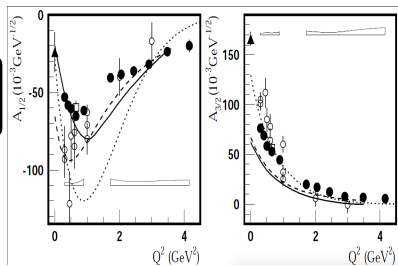


$\vec{e}p \rightarrow e'\pi^+n$ for low lying N^* ($W = 1.15 - 1.69$ GeV)

- Transition Form Factors for $N(1520)3/2^-$ (old conv: $D_{13}(1520)$)
- $A_{1/2}$ is large at high Q^2 , $A_{3/2}$ is small at high Q^2



[▲: RPP(2014), ○, □: DESY, NINA data]



$$\frac{A_{1/2}^{D13}}{A_{3/2}^{D13}} = \frac{-1}{\sqrt{3}} \left(\frac{Q^2}{\alpha} - 1 \right)$$

$$A_{hel} = \frac{A_{1/2}^2 - A_{3/2}^2}{A_{1/2}^2 + A_{3/2}^2}$$

- Asymptotic Q^2 behavior of A_{hel} vs. Q^2
- NRQ simple harmonic oscillator model (solid line) with spin, orbit flip amplitudes
- $A_{1/2} \ll A_{3/2}$ at low Q^2 , $A_{3/2} \ll A_{1/2}$ at high Q^2

Selection Rules in Symmetric Quark Model

- The first orbital excitation states
 $|70, {}^28, 1, 1, J\rangle - S_{11}(1535)(****), D_{13}(1520)(****)$
 $|70, {}^48, 1, 1, J\rangle - S_{11}(1650)(****), D_{13}(1700)(***), D_{15}(1675)(****)$
- **Moorhouse selection rule** (Moorhouse, PRL16, 772 (1966))
 $\gamma + p(|56, {}^28; 0, 0, 1/2\rangle) \leftrightarrow N^*(|70, {}^48\rangle)$: vanishing TME for charge operator
 $\gamma + n(|56, {}^28; 0, 0, 1/2\rangle) \leftrightarrow N^*(|70, {}^48\rangle)$
- Λ selection rule (Zhao, PRD74, 094014 (2006))
 $N^*|70, {}^48\rangle \leftrightarrow K(K^*) + \Lambda$
- Faiman-Hendry selection rule (Faiman, Hendry, PR173, 1720 (1968))
 $\Lambda^*|70, {}^48\rangle \leftrightarrow N(|56, {}^28; 0, 0, 1/2\rangle) + \bar{K}$

Moorhouse selection rule must be violated !

Spin-dependent potential from one-gluon-exchange and $SU(6) \otimes O(3)$ symmetry breaking, interaction H_{hyper} is introducing mass splitting and configuration mixing in $SU(6)$ multiplets

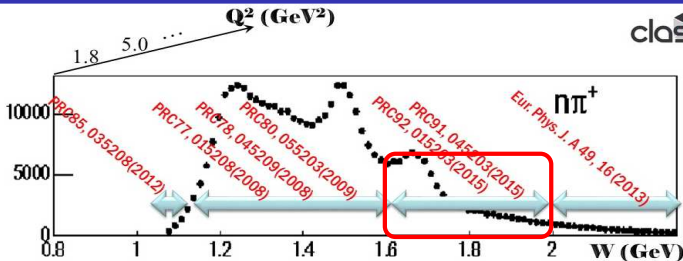
$$H_{hyper} = \frac{2\alpha_s}{3m_i m_j} \left[\frac{8\pi}{3} S_i \cdot S_j \delta^3(r_{ij}) + \frac{1}{r_{ij}^3} \left(\frac{3(S_i \cdot r_{ij})(S_j \cdot r_{ij})}{r_{ij}^2} - S_i \cdot S_j \right) \right]$$

TABLE I. Violations of some $SU(6)$ rules.

Quantity	SU(6) (Relative values)	This calculation (Relative values)	Experiment (Various units)
$A_{3/2}^n(D_{15} \rightarrow n\gamma)$	$-\alpha$	$-\alpha$	-60 ± 33^a
$A_{1/2}^n(D_{15} \rightarrow n\gamma)$	-0.71α	-0.71α	-33 ± 25^a
$A_{3/2}^p(D_{15} \rightarrow p\gamma)$	0	$+0.31\alpha$	$+20 \pm 13^a$
$A_{1/2}^p(D_{15} \rightarrow p\gamma)$	0	$+0.22\alpha$	$+19 \pm 14^a$
$A(D_{15} \rightarrow \bar{K}N)$	β	β	$+0.41 \pm 0.03^b$
$A(D_{05} \rightarrow \bar{K}N)$	0	-0.28β	-0.09 ± 0.04^c
$\langle \sum e_i r_i^2 \rangle_p$	γ	γ	$+0.82 \pm 0.02^d$
$\langle \sum e_i r_i^2 \rangle_n$	0	-0.16γ	-0.12 ± 0.01^e

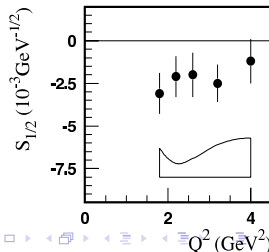
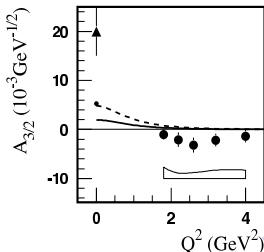
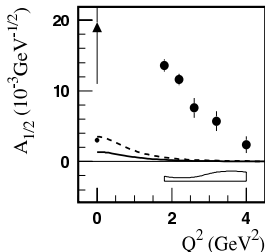
[Reference:Isgur, Karl, PRL 41, 1269 (1978).]

$\vec{e}p \rightarrow e'\pi^+n$ for high lying N^* ($W = 1.65 - 2.0$ GeV)

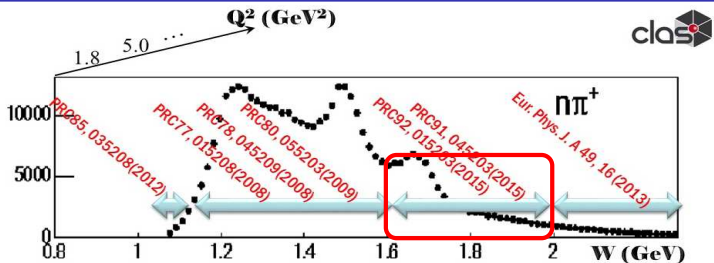


- Transition Form Factors for $N(1675)5/2^-$ (old conv: $D_{15}(1675)$)
- SQTM, Moorhouse selection rule: suppression Transverse Amplitudes

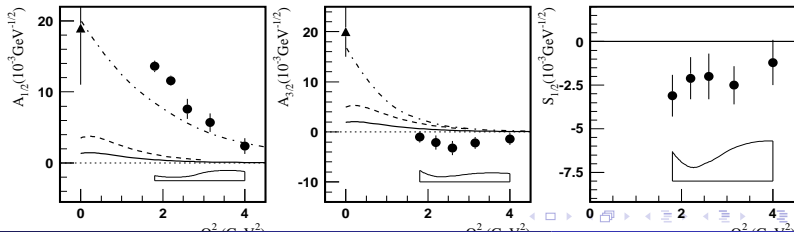
● Solid: M.M.Gianini/E.Santopinto (hQCM)
 dash: D.Merten& U.Loring(2003), Solid-dot($Q^2 = 0$): I.G.Aznauryan(LFRQ)



$\vec{e}p \rightarrow e'\pi^+n$ for high lying N^* ($W = 1.65 - 2.0$ GeV)

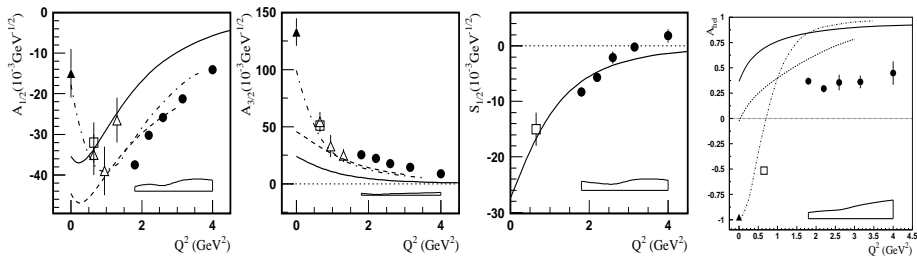


- Non-quark contributions dominance, A strong coupling $A_{1/2}$ for $Q^2 < 4$ GeV²
- Significant *MB contribution from the dynamical coupled-channel model* (dash-dot: B. Julia-Diaz, T-S.H. Lee, A. Matsuyama)
- A strong suppression of $A_{3/2}$ for $Q^2 > 1.8$ GeV²



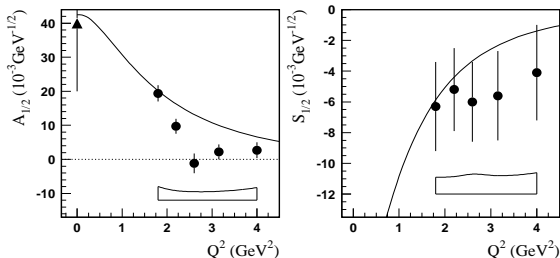
$\bar{e}p \rightarrow e'\pi^+n$ for high lying N^* ($W = 1.65 - 2.0$ GeV)

- **Transition Form Factors for $N(1680)5/2^+$ (old conv: $F_{15}(1680)$)**
- ▲ RPP(PDG:2014), \triangle V.Mokeev& I.G.Aznauryan(2013), \square I.G.Aznauryan(2005)
- Solid: M.M.Gianini/E.Santopinto (hQCM), dash-dot: Z.Lee& F.Close(1990), dash: D.Merten& U.Loring(2003)
- All models estimates amplitudes larger $A_{1/2}$ (lower $A_{3/2}$) than data
- *MB* contribution should be taken into account ?



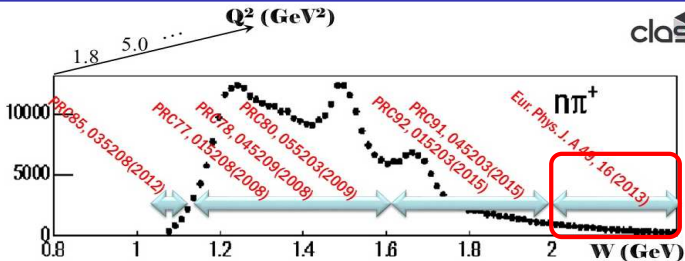
$\vec{e}p \rightarrow e'\pi^+n$ for high lying N^* ($W = 1.65 - 2.0$ GeV)

- Transition Form Factors for $N(1710)1/2^+$ (old conv: $P_{11}(1710)$)
- Finite size of $A_{1/2}$ for $Q^2 < 2.5$ GeV²
- Finite size and negative of $S_{1/2}$ for all given Q^2 GeV²



Solid: M.M.Gianini
E.Santopinto (hQCM)

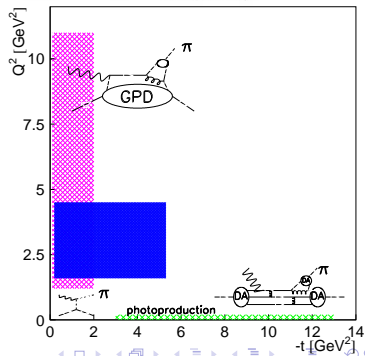
Deep Inelastic Process ($W > 2.0$ GeV)



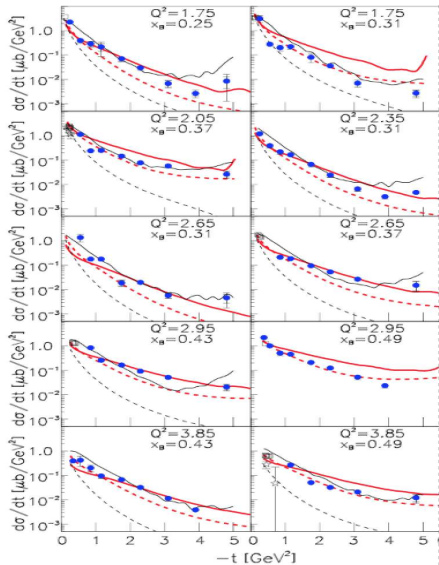
- Transition between hadronic and partonic picture of strong interaction
- GPD
→ Correlations of longitudinal momentum fraction with transverse spatial position
- DVMP: $N(e, e'NM)$, $M = \pi, \rho, \phi, \dots$
→ Connection to the transversity GPD
- Kinematic variables → x_{BJ} , Q^2 and $-t$

Blue box →

[K.Park, et al., Eur. Phys. J. A49 16, (2013)]

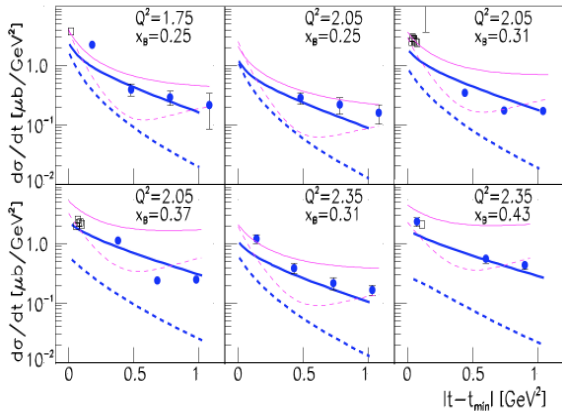


Hard exclusive forward, large-angled $\gamma^* p \rightarrow n\pi^+$

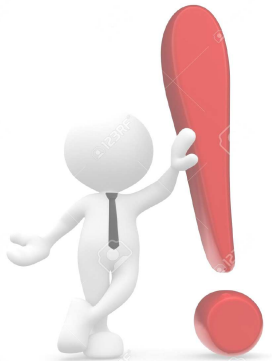


- Solid ($d\sigma/dt$), dashed curves ($d\sigma_L/dt$)
- Red curves: J. M. Laget, Regge-model
- Blue curves: M. Kaskulov, Hybrid (hadron-parton) model

Hard exclusive forward, large-angled $\gamma^* p \rightarrow n\pi^+$



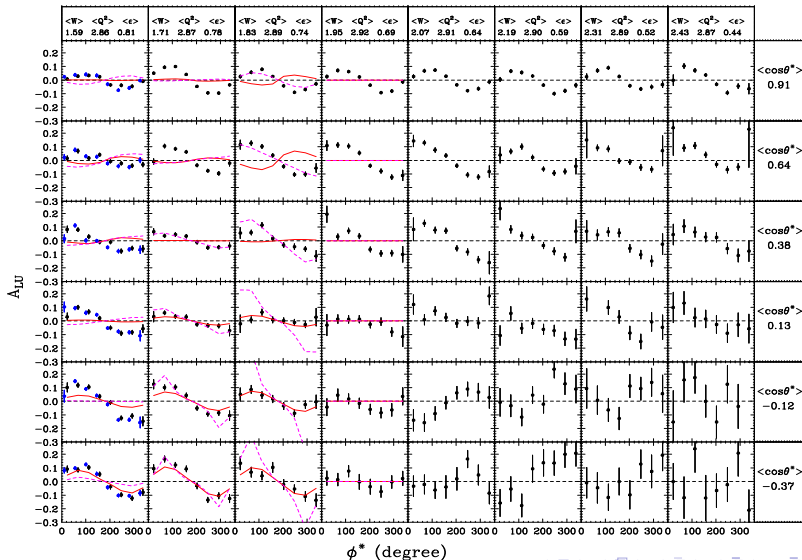
- Solid ($d\sigma/dt$), dashed curves ($d\sigma_L/dt$)
- **Magenta curves:**
M. Kaskulov, Duality model
→ Transverse: resonance excitation
→ Longitudinal: t -channel meson exchange
- **Blue curves:** G-K :
Transversity of GPDs
→ Partonic model (handbag diagram) (But w/o adjusting J_{lab} kinematics)

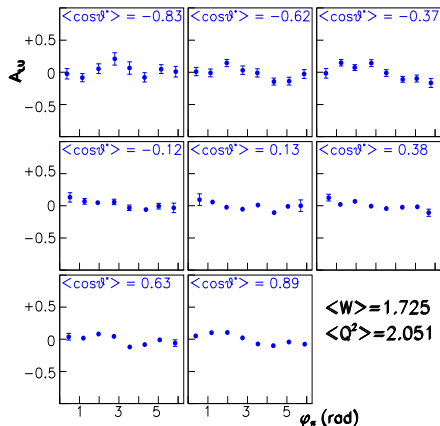
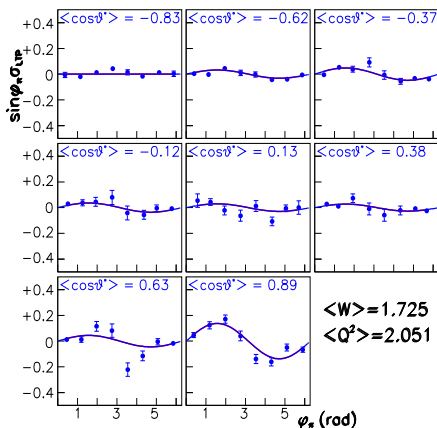


New upcoming results 2017 !!!
under CLAS⁶ Analysis Review !

$\vec{e}p \rightarrow e'\pi^+n$, A_{LU} for $W = 1.6 - 2.4$ GeV

Curves: solid-MAID2007, dashed-JANR, Blue: Phys. Rev. C77 015208, (2008), Black points: current work P. Bosted/K. Park

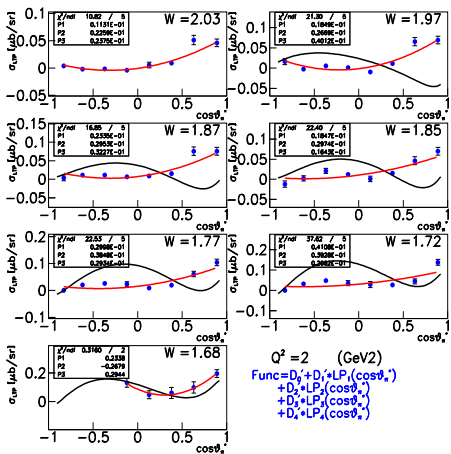


A_{LU}  $\sin \phi^* \sigma_{LT'}$ Blue curves: $\sin \phi^*$ fit

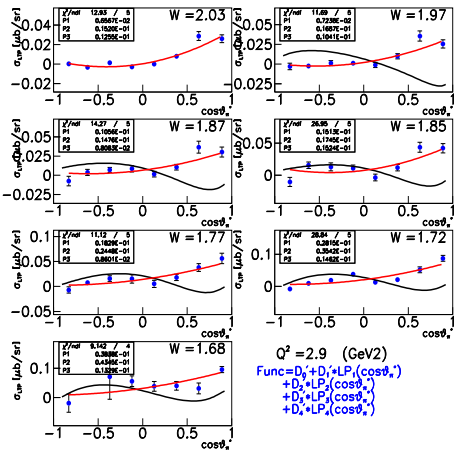
$\sigma_{LT'}$ vs. $\cos\theta_{\pi^*}$, MAID2007, PRELIMINARY

Red curves: Legendre fit, Black curves: MAID2007

$Q^2 = 2.0 \text{ GeV}^2$



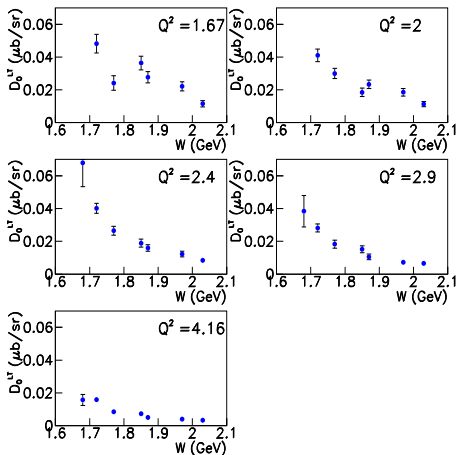
$Q^2 = 2.90 \text{ GeV}^2$



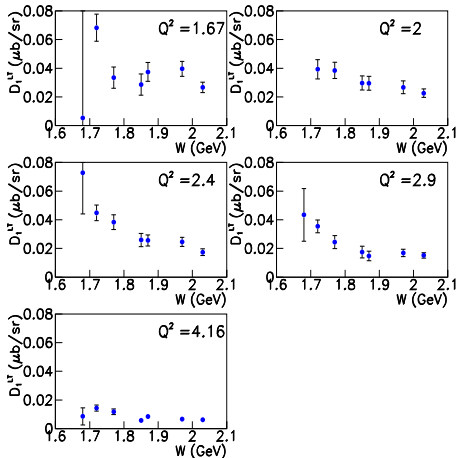
Moments $D^{LT'}$ vs. W , PRELIMINARY

Observation: interesting behavior above $W > 1.8$ GeV

Moments $D_0^{LT'}$



Moments $D_1^{LT'}$



- Extraction of the transition form factor have been carried out through the differential cross-sections/asymmetries measurements for $\gamma^* p \rightarrow n\pi^+$ data for nearly full range of kinematics, **near threshold** $< W < \mathbf{DIS}$ regime, $Q^2 = 1.6-4.5 \text{ GeV}^2$.
- Precision data for $\gamma^* p \rightarrow n\pi^+$ from CLAS allows to extract the helicity amplitudes for various reonance states, $N(1440)1/2^+$, $N(1520)3/2^-$, $N(1535)1/2^-$, $N(1675)5/2^-$, $N(1680)5/2^+$, and $N(1710)1/2^+$
- Coupled-channel analysis (including $p\pi^0, p\pi^+\pi^-, \dots K\Lambda, K\Sigma\dots$) is crucial in particular high W and this will improve considerably our knowledge on N^* -state electro-couplings.

→ **Coupled-Channel Analysis and A Full Mass Spectrum data help us to map out nucleon structure with an effective degree of freedom.**