

Phase conversion dissipation in multi-component stars

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Alford, Han & Schwenzer,
[arXiv:1404.5279](https://arxiv.org/abs/1404.5279)

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The Phases of Dense Matter

INT Program INT-16-2b Week III

Dense matter in neutron stars

Is there quark matter in NS?

properties

observables

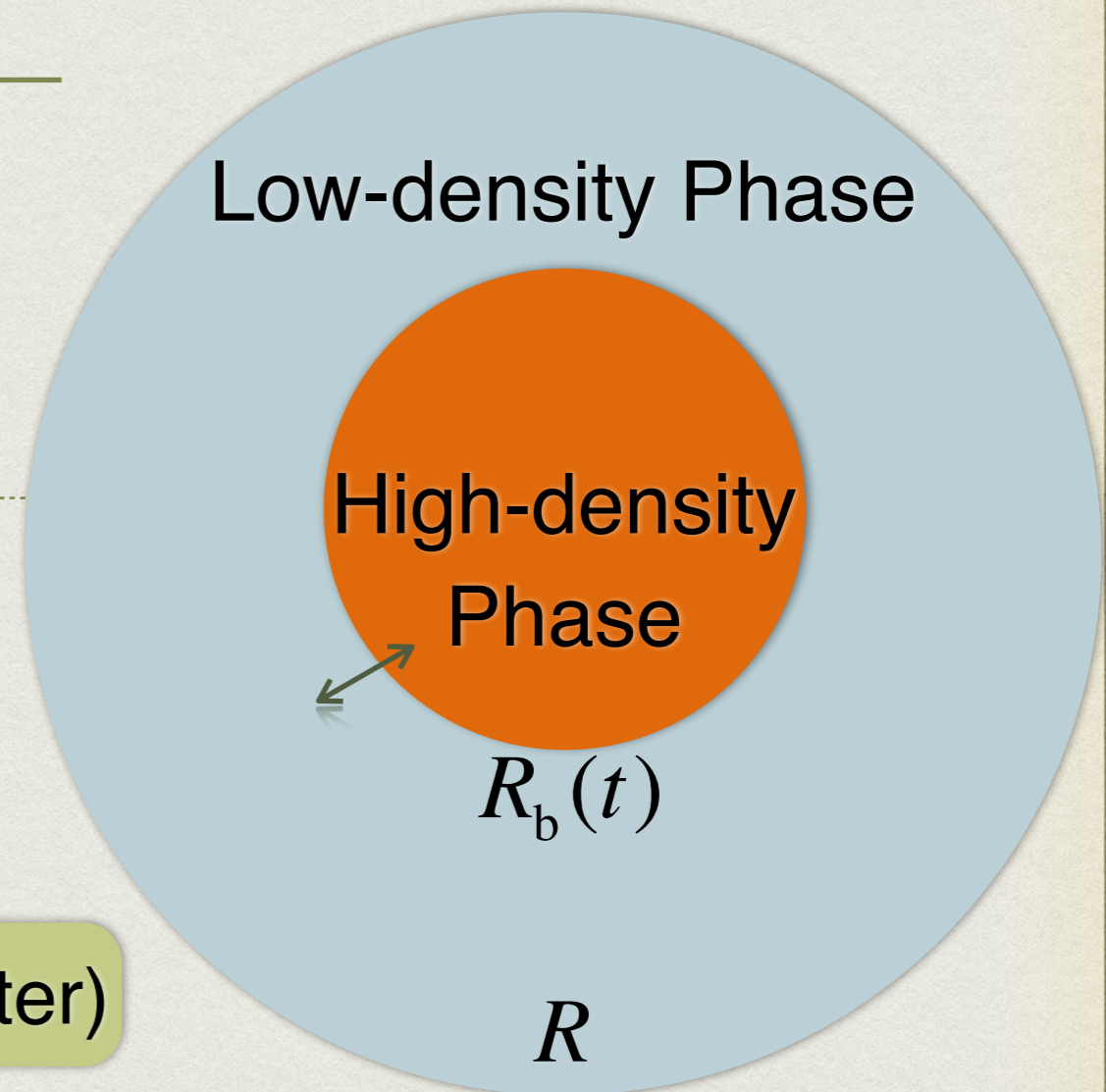
equations of state

mass, radius, moment of inertia...

thermal & transport properties, vortex pinning

cooling, spin down, glitches, neutrinos, GW, magnetic field...

Dissipation (in uniform matter)



- Shear viscosity from particle scattering (strong/EM interaction)
- Bulk viscosity from particle transformation (weak interaction)

Dissipation at an interface

-Ekman layer damping from shear rubbing of a fluid core along a solid crust

Phase conversion dissipation

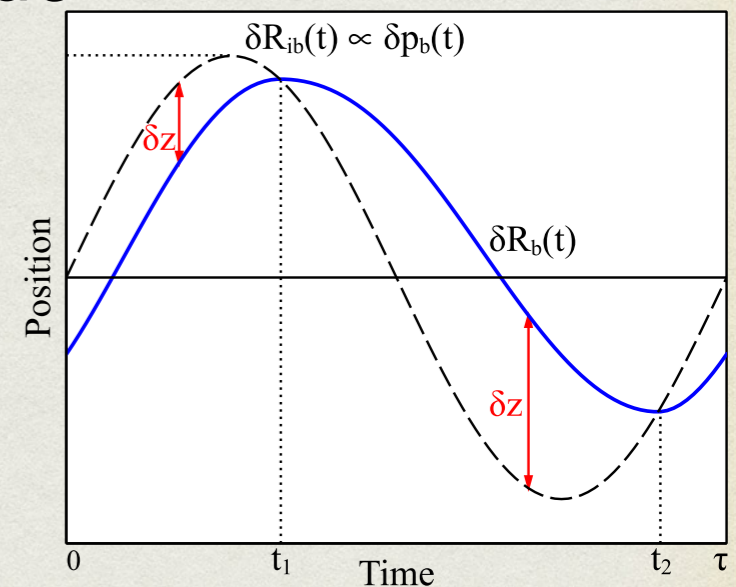
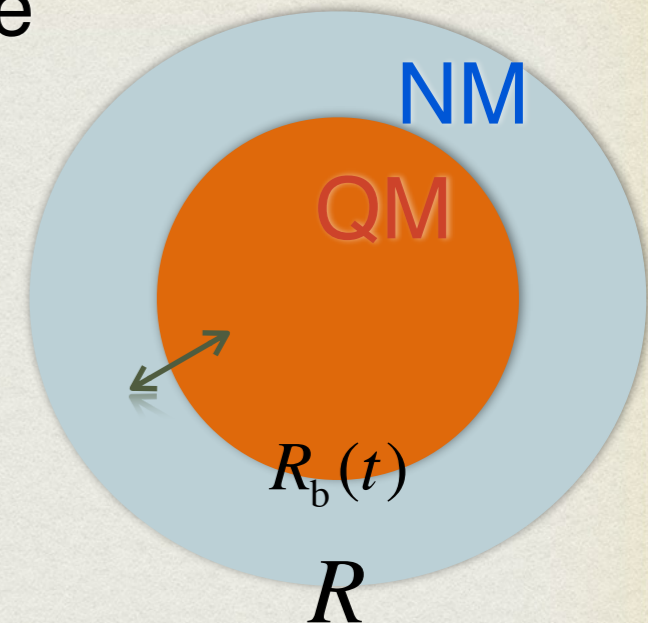
-Between fluids in different phases with first-order transition separated by a sharp interface

-Quark/hadron conversion

1) flavor-changing process $d \leftrightarrow s$ out of equilibrium due to global oscillations

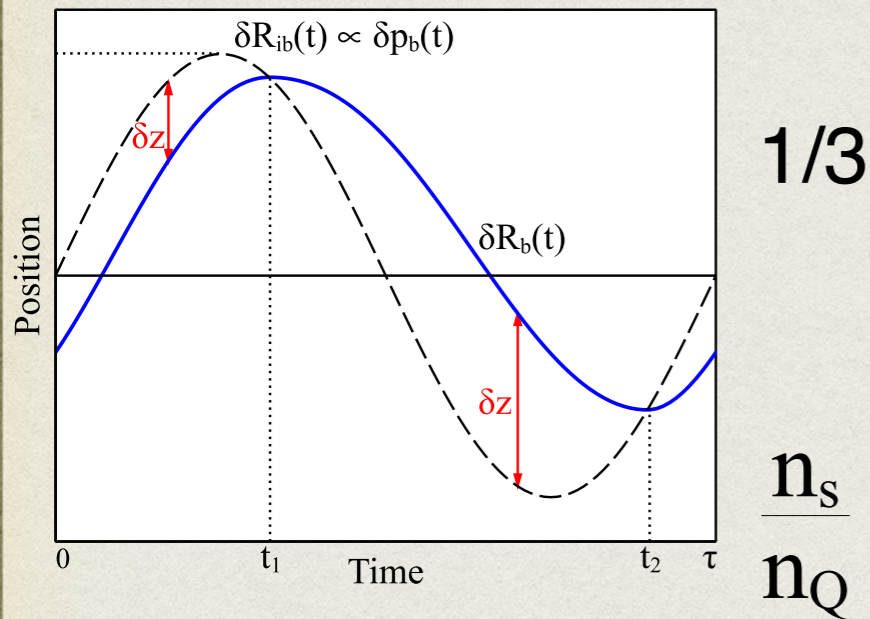
2) instantaneous restoration \Leftrightarrow
phase boundary moves arbitrarily fast (no diss.)

3) finite rate of weak interaction and flavor diffusion
→ a phase lag in system response
→ dissipates energy

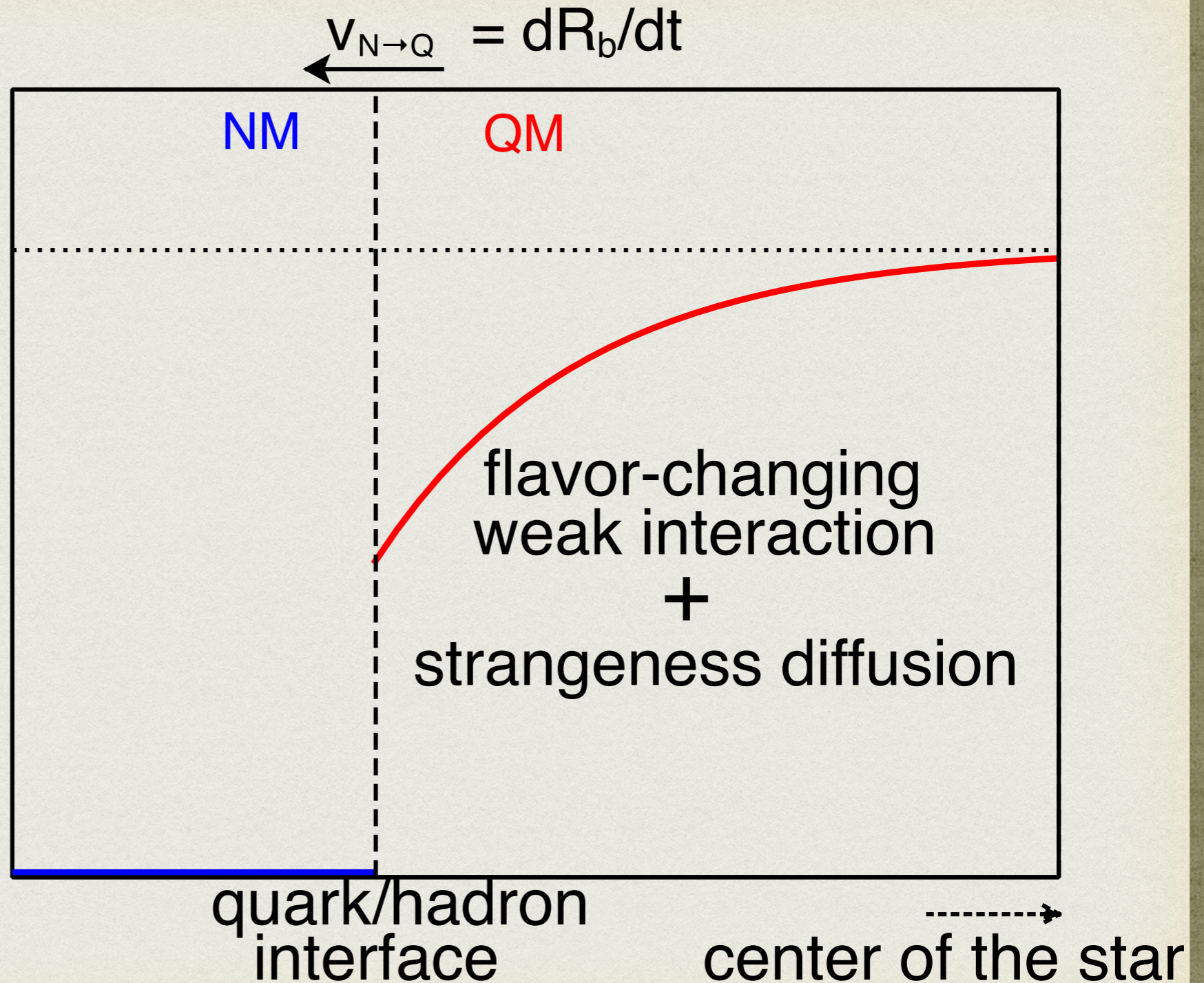


Toy model formalism

$$(udd) \leftrightarrow (uds)$$



- steady-state transport; boundary conditions
- no acceleration/ deceleration effects; no turbulence
- no leptons



Dissipated energy (pdV work)

out of phase

$$dW = -dS \left(\frac{n_N}{n_Q} - 1 \right) \int_0^\tau \delta p_b(t) \frac{d\delta R_b(t)}{dt} dt$$

-actual front speed $n \rightarrow q$

$$\left. \frac{d\delta R_b}{dt} \right|_{N \rightarrow Q} \propto \sqrt{D_Q / \tau_Q}$$

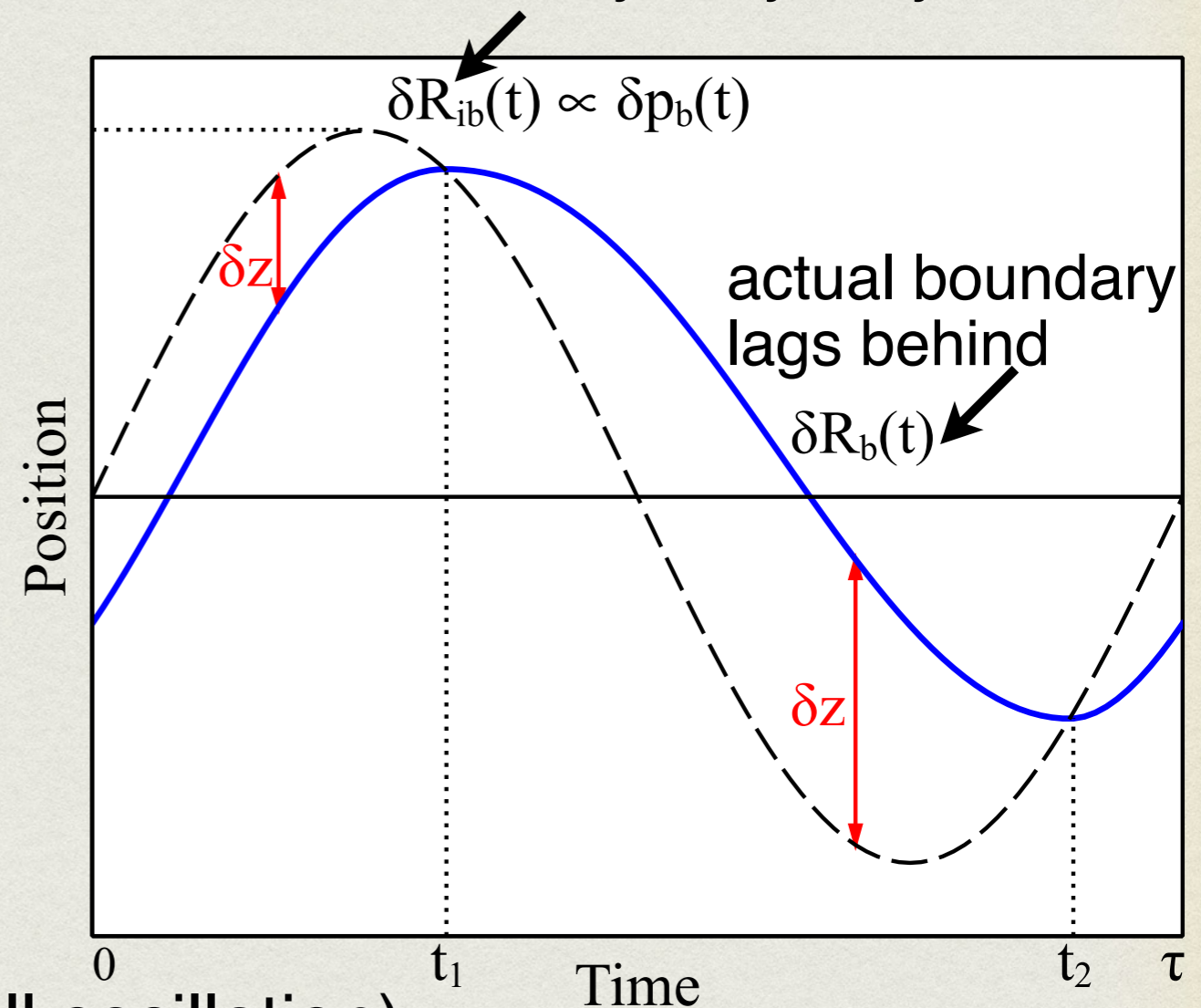
-similarly for the other half cycle

-two regimes

subthermal: $\Delta\mu \ll T \ll \mu$ (small oscillation)

suprathermal: $T \ll \Delta\mu \ll \mu$ (large oscillation)

“ideal boundary” fully in sync



R-mode saturation

-integrating over the whole star

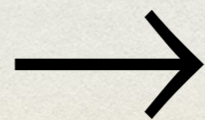
$$W \propto \begin{cases} \alpha^3, \text{subthermal} \\ \alpha^2, \text{suprathermal} \end{cases}$$

-non-linearly increasing with amplitude (resonant behavior)

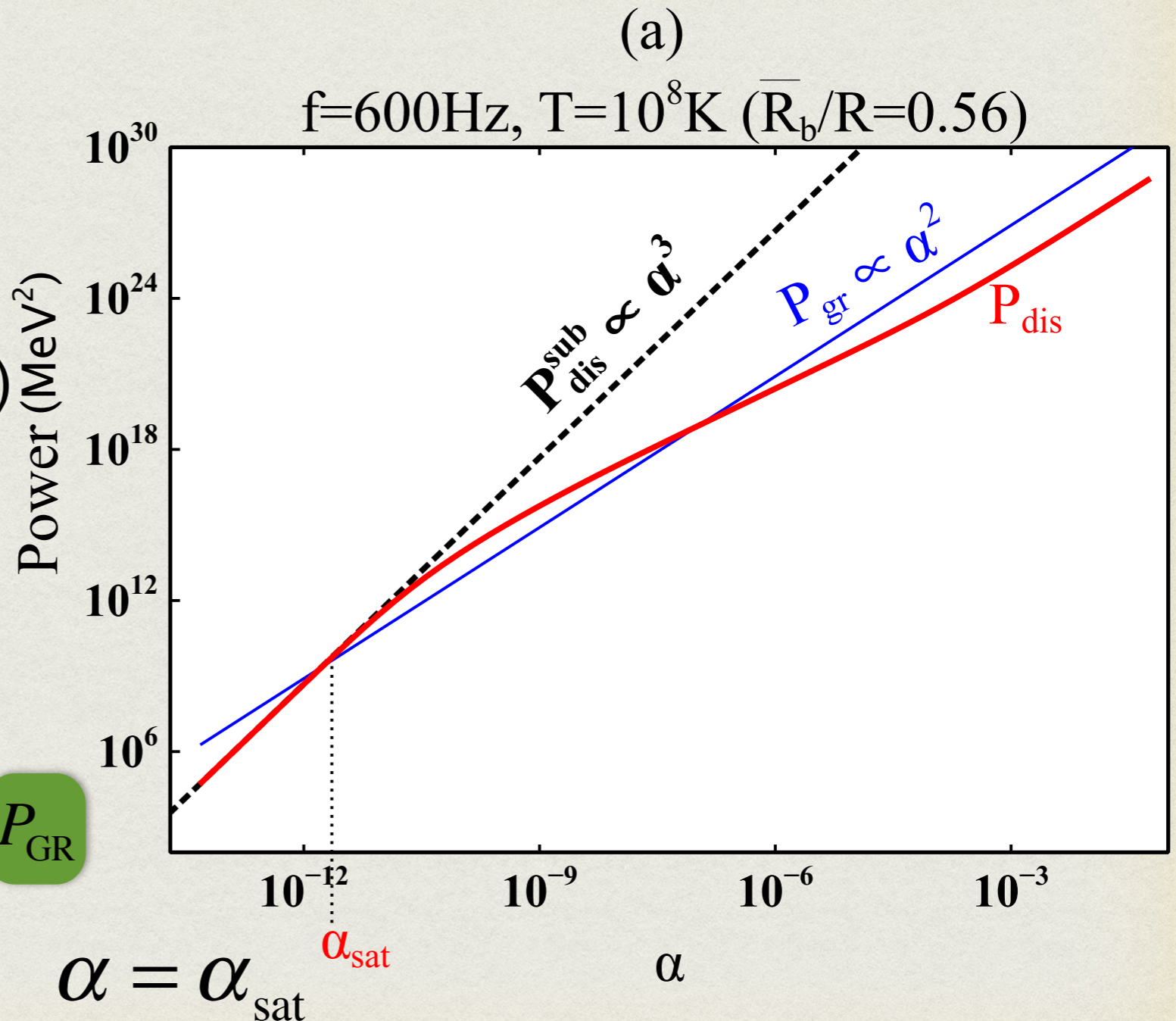
-dissipated power balances gravitational radiation power:

r-mode saturated

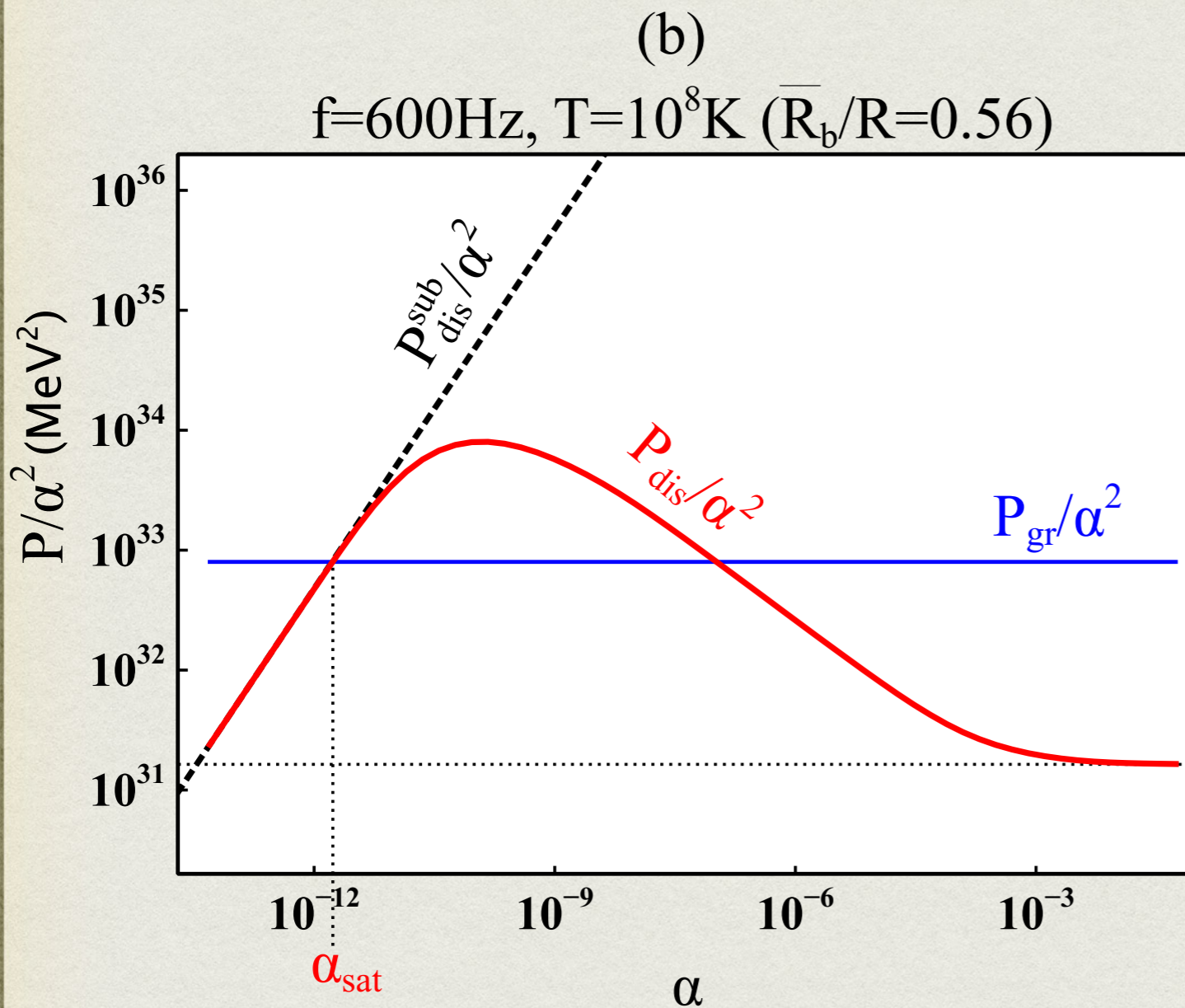
$$P_{\text{dis}} = W_{\text{dis}}(\alpha) \cdot f = \left. \frac{d\tilde{E}(\alpha)}{dt} \right|_{\text{GR}} = P_{\text{GR}}$$



$$\alpha = \alpha_{\text{sat}}$$



Phase conversion dissipation



-Varying parameters e.g. the core size, rotation frequency or mass of the star will shift the curves

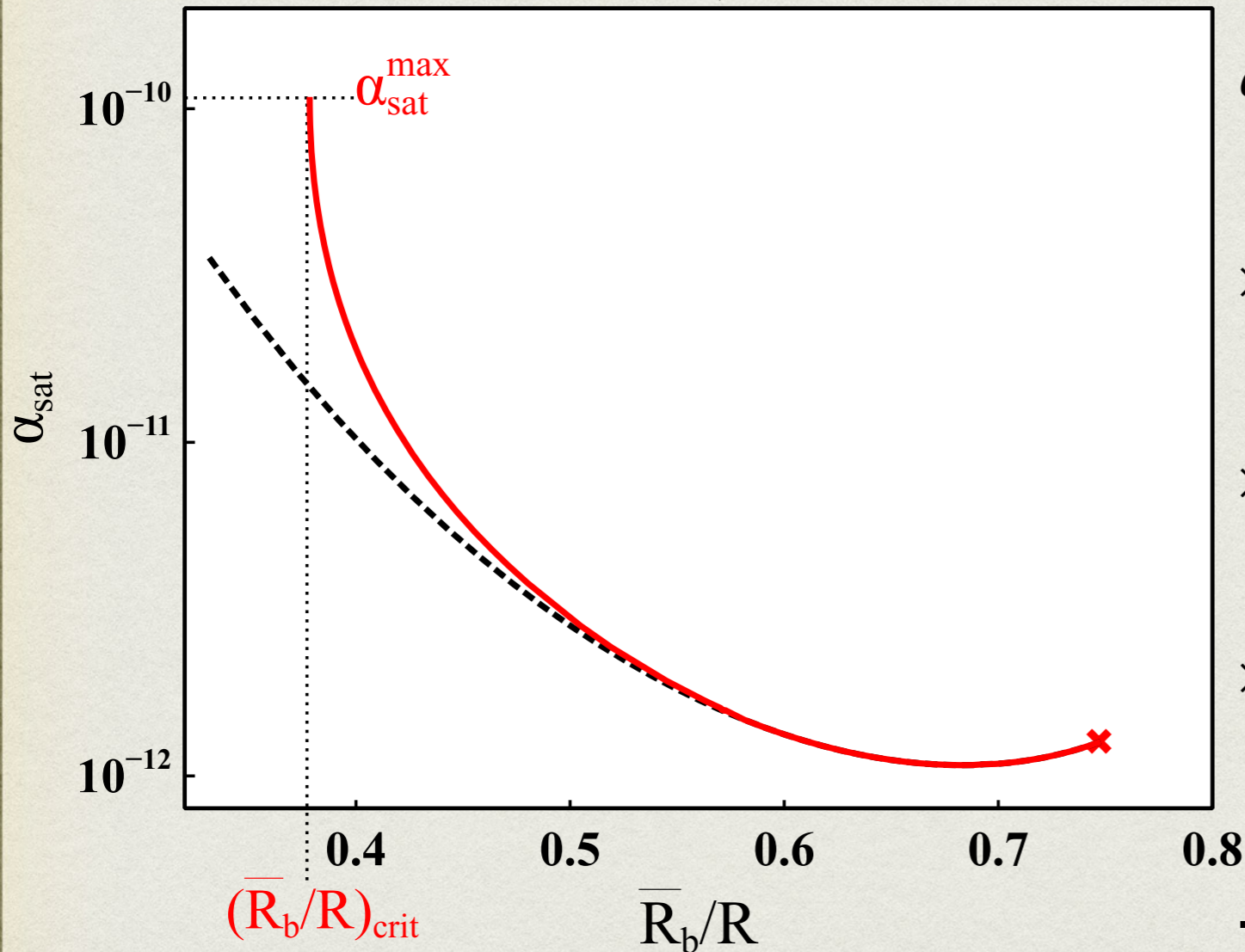
-There is no saturation if damping is insufficient to counteract driving (blue line above the maximum of red)

-If saturation occurs, most likely in the low-amplitude regime:

analytical calculation

Low-amplitude approximation

$f=600\text{Hz}, T=10^8\text{K}$



$$\alpha_{\text{sat}}^{\text{approx}} \approx 4.2 \times 10^{-11} \gamma \left(\frac{\tilde{D}_N}{1.5 \text{MeV}} \right) \left(\frac{\tau_N}{2 \times 10^{-8} \text{s}} \right)^{-1}$$

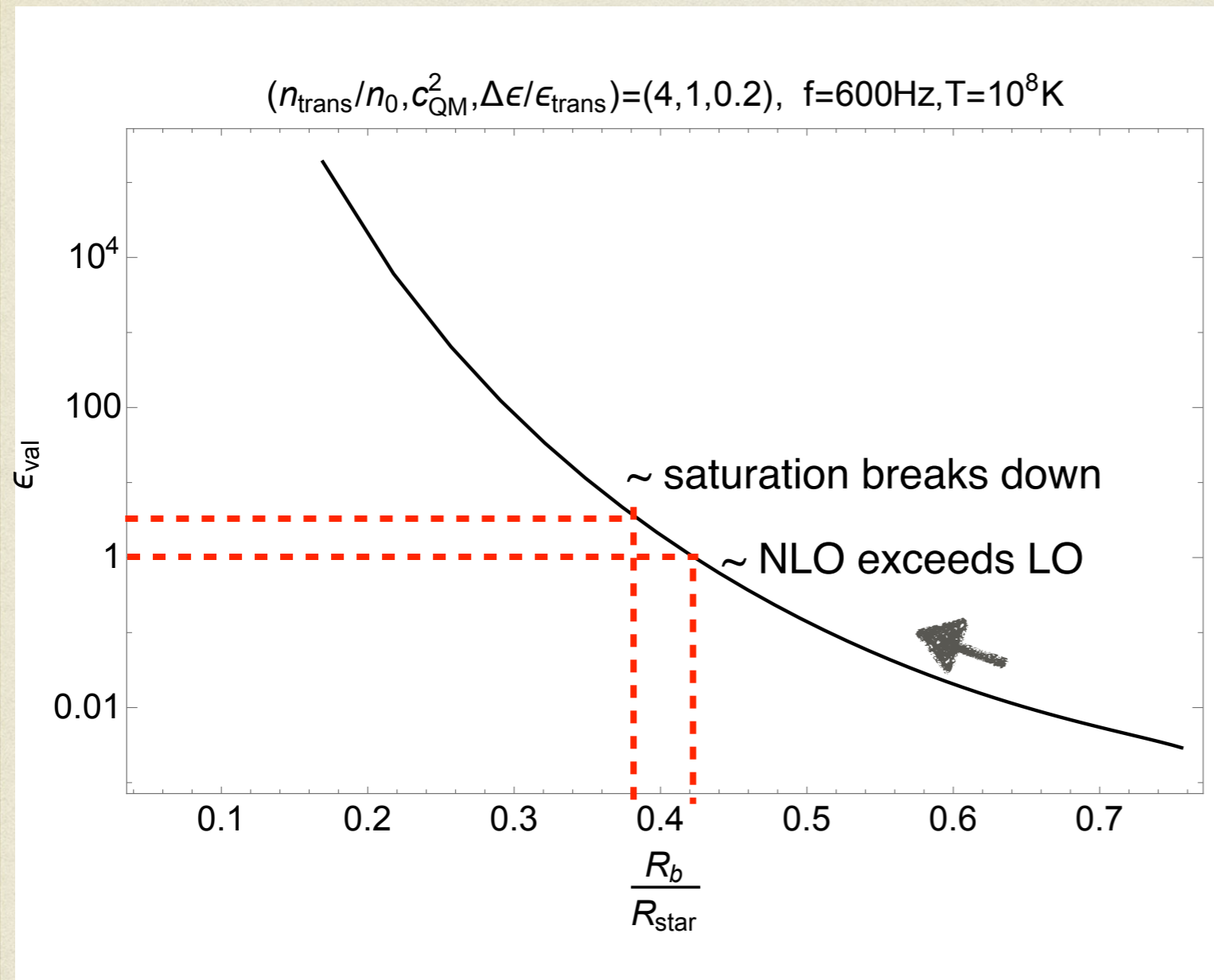
$$\times \left(\frac{b_Q}{1/3} \right)^{-2} \left(\frac{n_N}{2n_0} \right)^{-4} \left(\frac{\chi_K^N}{(100 \text{MeV})^2} \right)^3 \left(\frac{g_b}{g_u} \right)^3$$

$$\times \left(\frac{\varepsilon_{\text{crit}}^Q}{2\varepsilon_{\text{crit}}^N} \right)^3 \left(\frac{\varepsilon_{\text{crit}}^N}{600 \text{MeV-fm}^{-3}} \right)^3 \left(\frac{M}{1.4 M_\odot} \right)^2$$

$$\times \left(\frac{\tilde{J}}{0.02} \right)^2 \left(\frac{f}{1 \text{kHz}} \right)^{-1} \left(\frac{R}{10 \text{km}} \right) \left(\frac{\bar{R}_b/R}{0.4} \right)^{-8}$$

-Very accurate at low amplitude, but does not capture the sudden weakening at small core size

Range of validity



-Compute the next-to-leading order contribution

$$P_{\text{dis}}(\alpha) = P_0(\alpha) + P_1(\alpha) + \dots$$

$$P_0(\alpha) \propto \alpha^3, P_1(\alpha) \propto \alpha^5$$

$$P_1(\alpha_{\text{sat}}^{\text{LO}}) \leq \epsilon P_0(\alpha_{\text{sat}}^{\text{LO}})$$

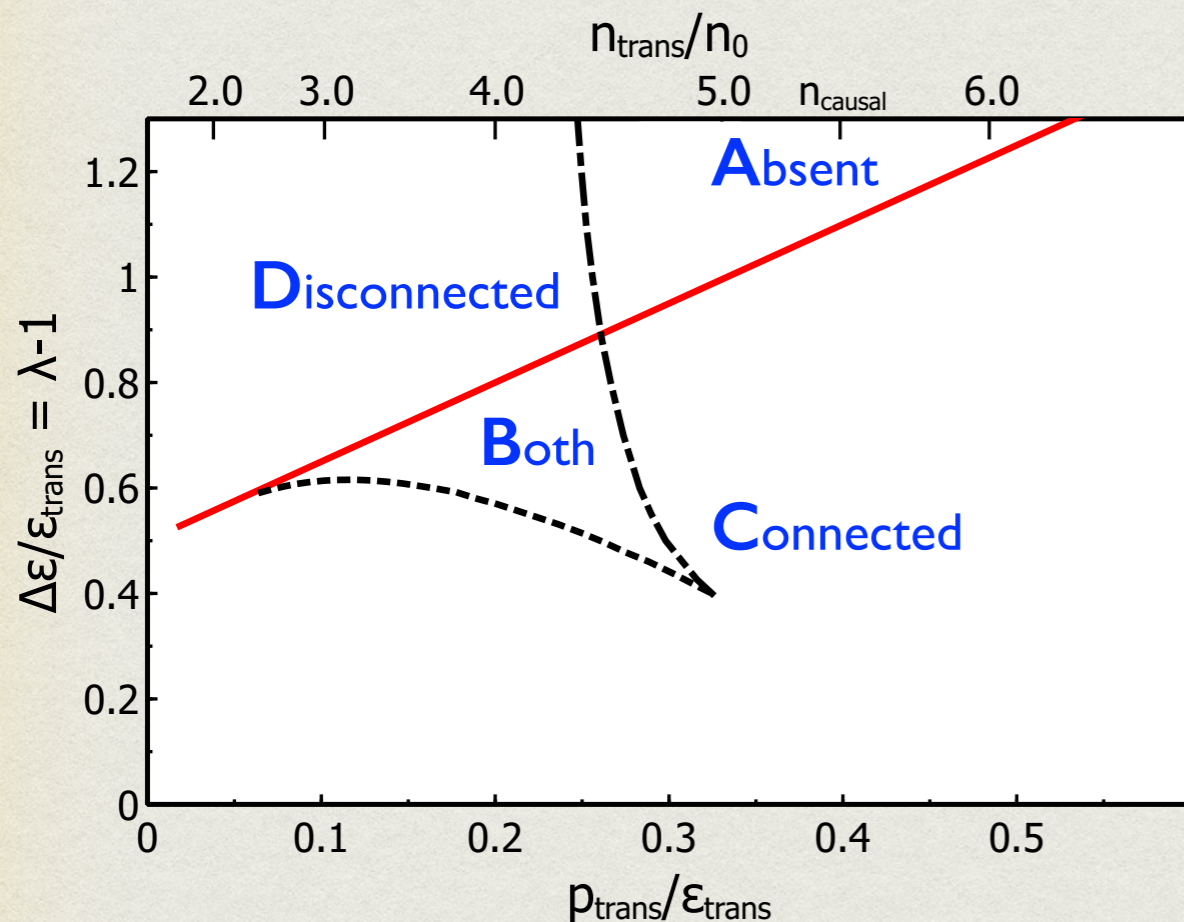
-Bound on parameter values

$$\epsilon \geq 2.96 \left(\frac{\gamma - 1}{0.5} \right)^{-2} \left(\frac{\epsilon_{\text{crit}}^{\text{Q}}}{2\epsilon_{\text{crit}}^{\text{N}}} \right)^2 \left(\frac{g_{\text{b}}}{g_{\text{u}}} \right)^2 \left(\frac{M}{1.4M_{\odot}} \right)^4$$

$$\times \left(\frac{\tilde{J}}{0.02} \right)^4 \left(\frac{f}{1\text{kHz}} \right)^6 \left(\frac{R}{10\text{km}} \right)^2 \left(\frac{\bar{R}_{\text{b}} / R}{0.4} \right)^{-14}$$

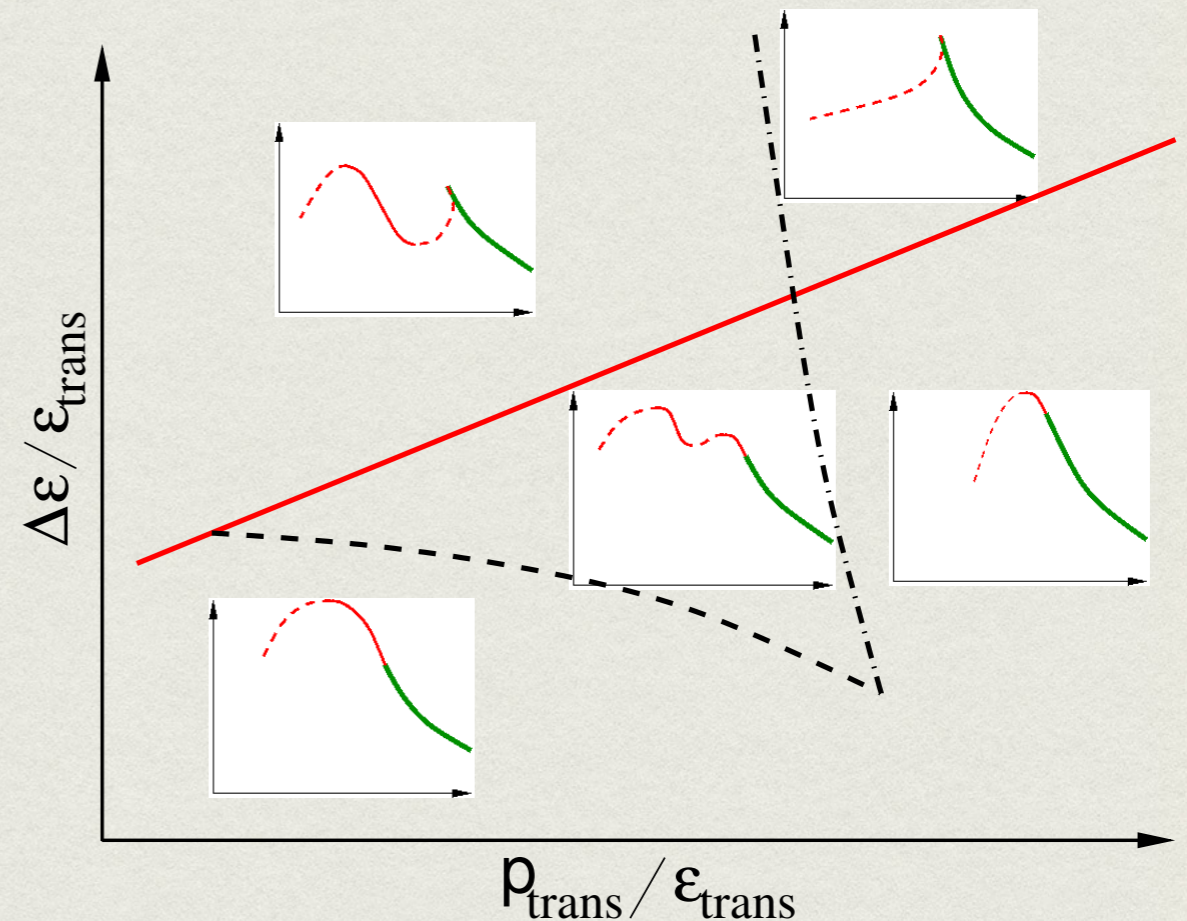
Hybrid star configurations

Soft NM (HLPS) + QM $c_{\text{QM}}^2 = 1$



-Below the red line in regions B and C, there is a connected hybrid star branch

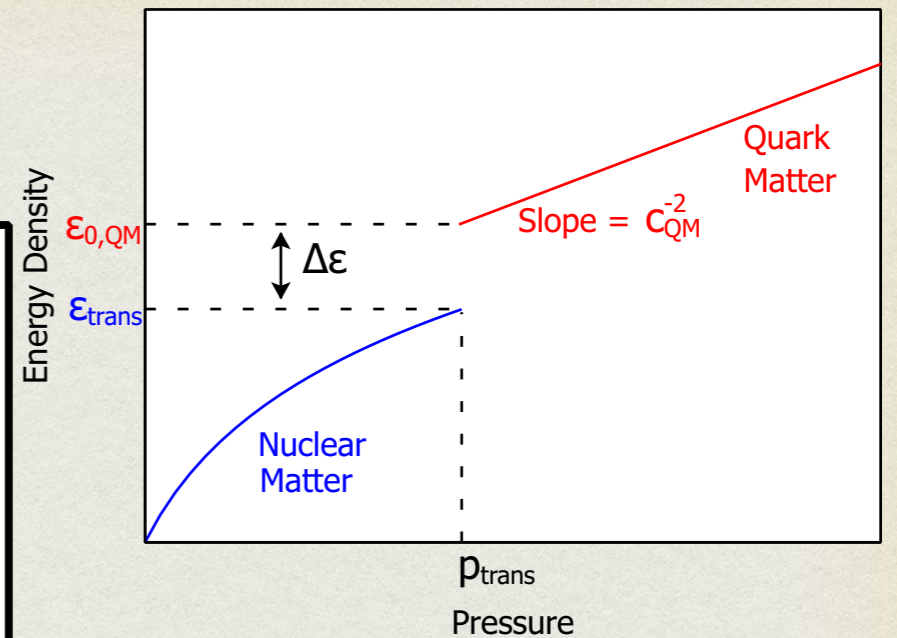
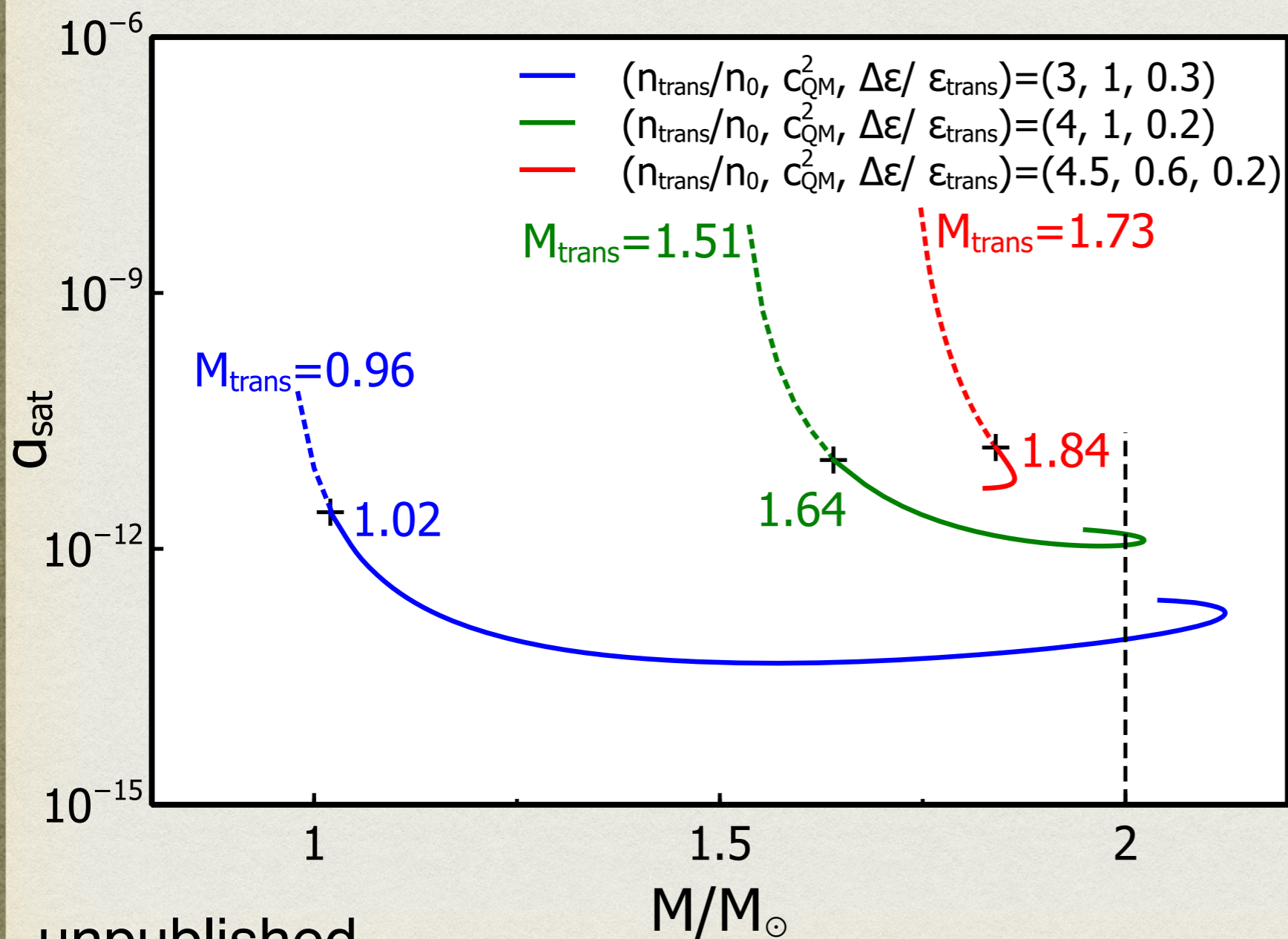
-with CSS parametrization



-In regions B and D, there is a disconnected hybrid star branch

Dependence on QM EoS

$f=600\text{Hz}, T=10^8\text{K}$



QM EoS parameters

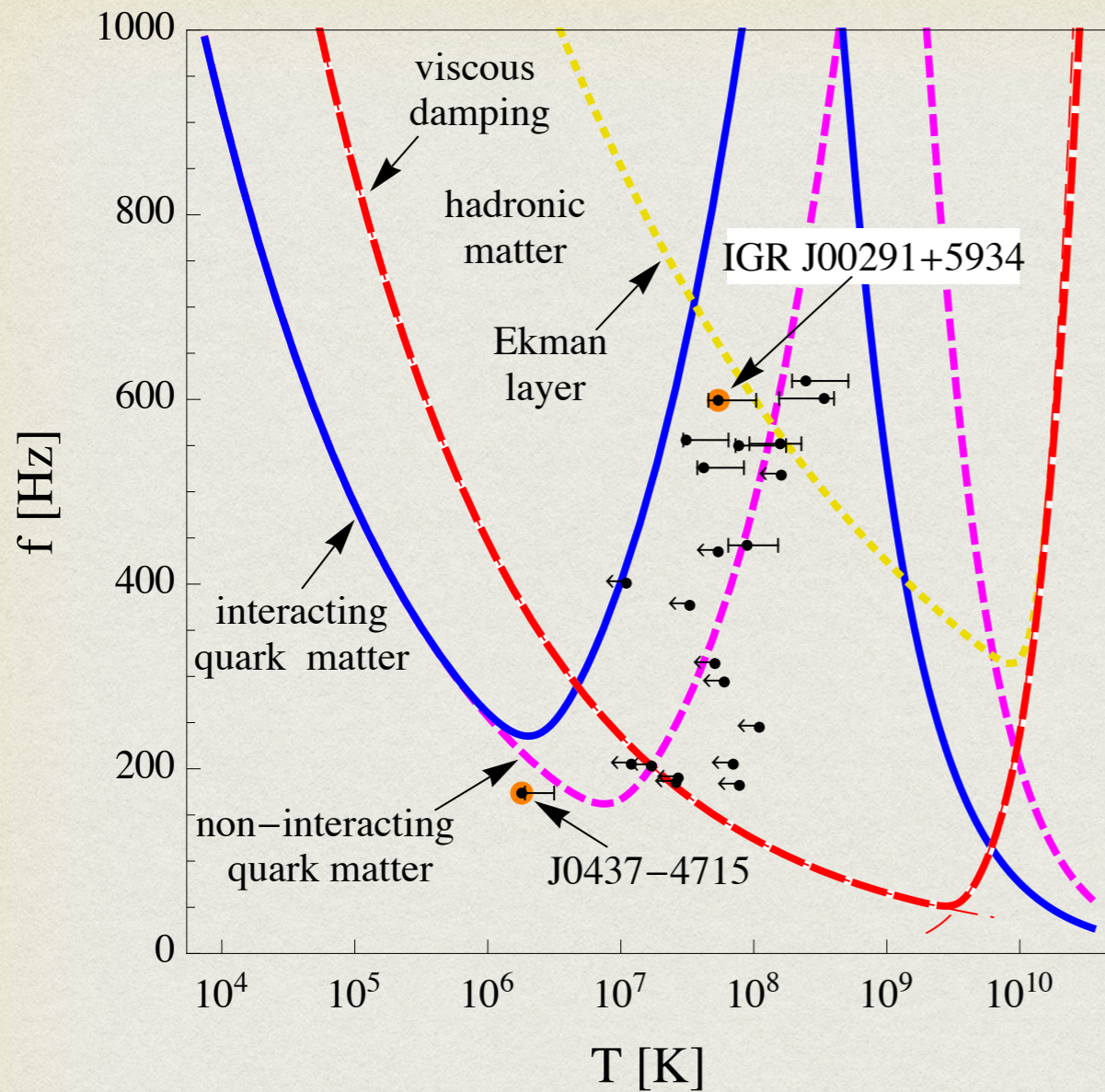
$$p_{\text{trans}} / \epsilon_{\text{trans}}$$

$$\Delta\epsilon / \epsilon_{\text{trans}}$$

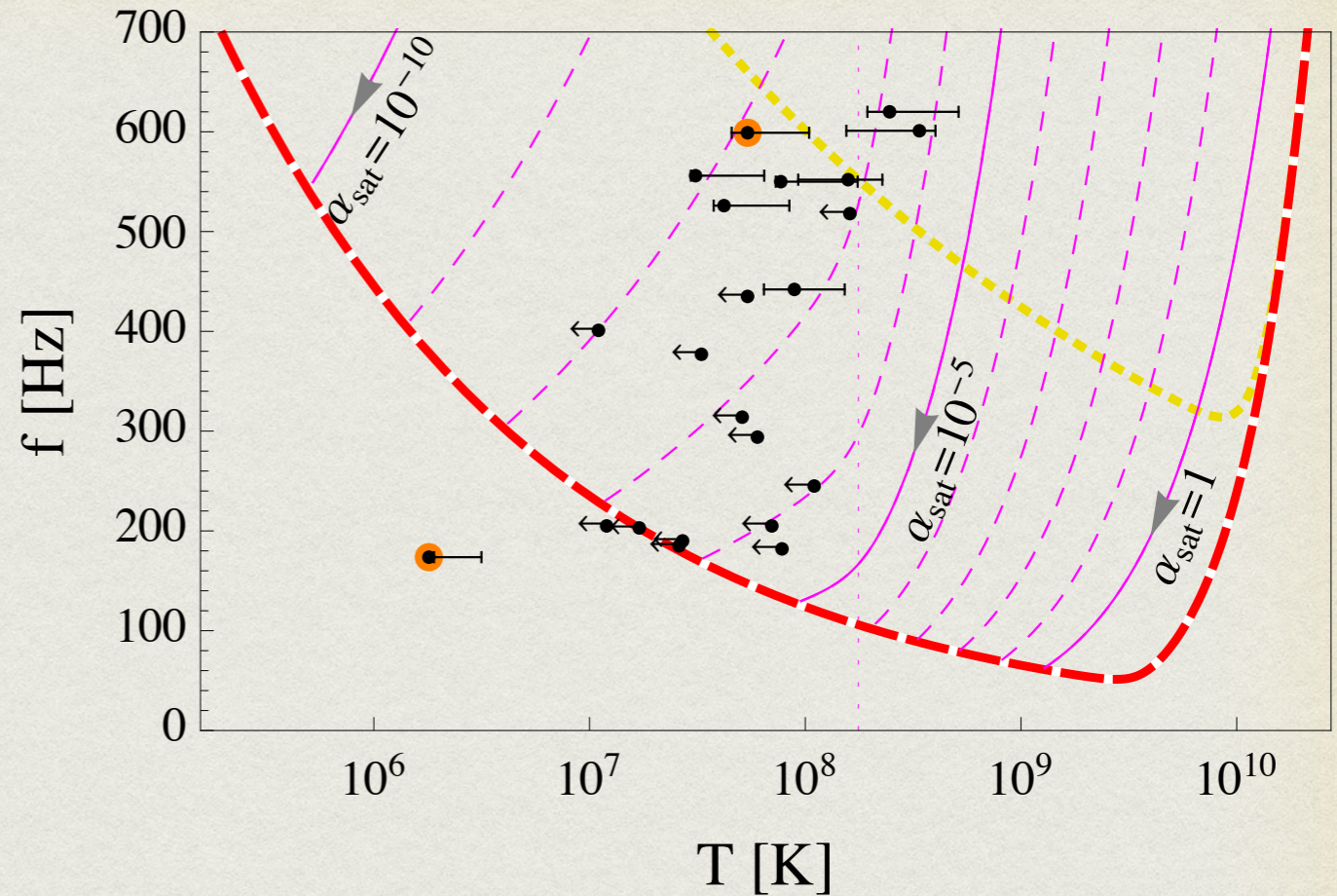
$$c_{\text{QM}}^2$$

unpublished

Astrophysical interpretation



Alford & Schwenzer, Phys. Rev. Lett. 113, 251102



Alford & Schwenzer, MNRAS 446.4 (2015)

-Damping vs. saturation
(no r-mode) (tiny r-mode)

-Thermal & spin evolution

Discussion

Toy model simplifications

- EoS parametrization
- (steady-state) laminar approach
- diffusion coefficient estimates in nuclear/quark matter
- density oscillations due to r-mode in non-uniform matter

Outlook

- vary nuclear EoSs; more quark matter EoS parameter space
- other 1st-order phase transitions with sharp interface
- mixed phases:
 - 1) domains of charged hadronic and quark matter coexist
 - 2) nuclear pasta in the inner crust
- hydrodynamic instability

Summary

A new saturation mechanism for oscillations in hybrid stars

-Phase conversion dissipation: a reasonable first estimate of the damping through a sharp boundary moving periodically in a steady state

Extremely low saturation amplitude can be achieved

$-\alpha_{\text{sat}} \approx O(10^{-12} \sim 10^{-10})$: consistent with millisecond pulsar frequency data

-At low amplitude analytic prediction is given

-Saturation amplitude relies on multiple parameters, among which the quark core size is the most sensitive one

Interpretation of observations

[arXiv:1404.5279](https://arxiv.org/abs/1404.5279)

-MSRPs could form a separate population of very cold sources and would not show *detectable* r-mode GW emission

THANK YOU!

Q & A