

Uncertainty Quantification Challenges From Modern *Ab Initio* Nuclear Theory

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Technische Universität Darmstadt
Institut für Kernphysik

Bayesian Methods in Nuclear Physics
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European Research Council
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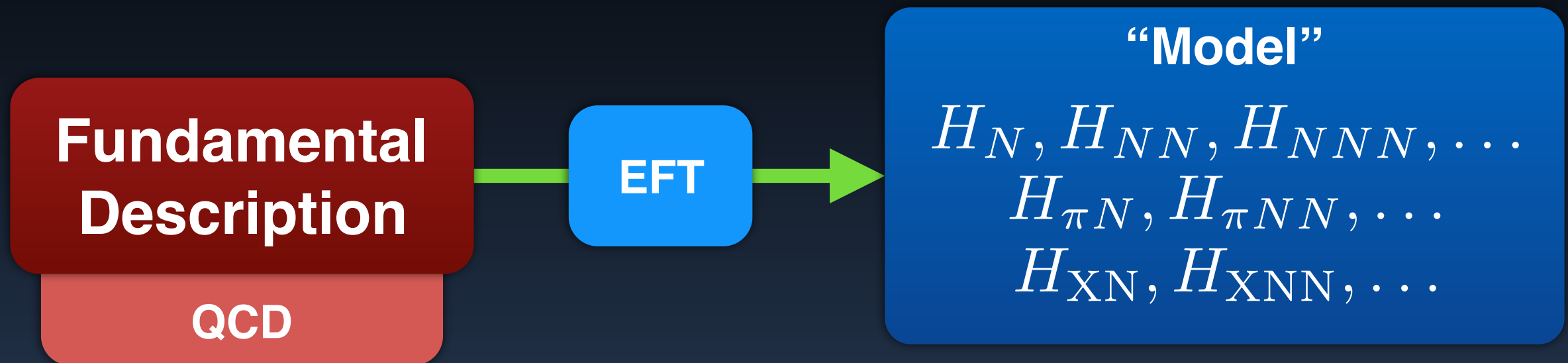
TECHNISCHE
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DARMSTADT

The Modern *Ab Initio* Approach to Low Energy Nuclear Physics

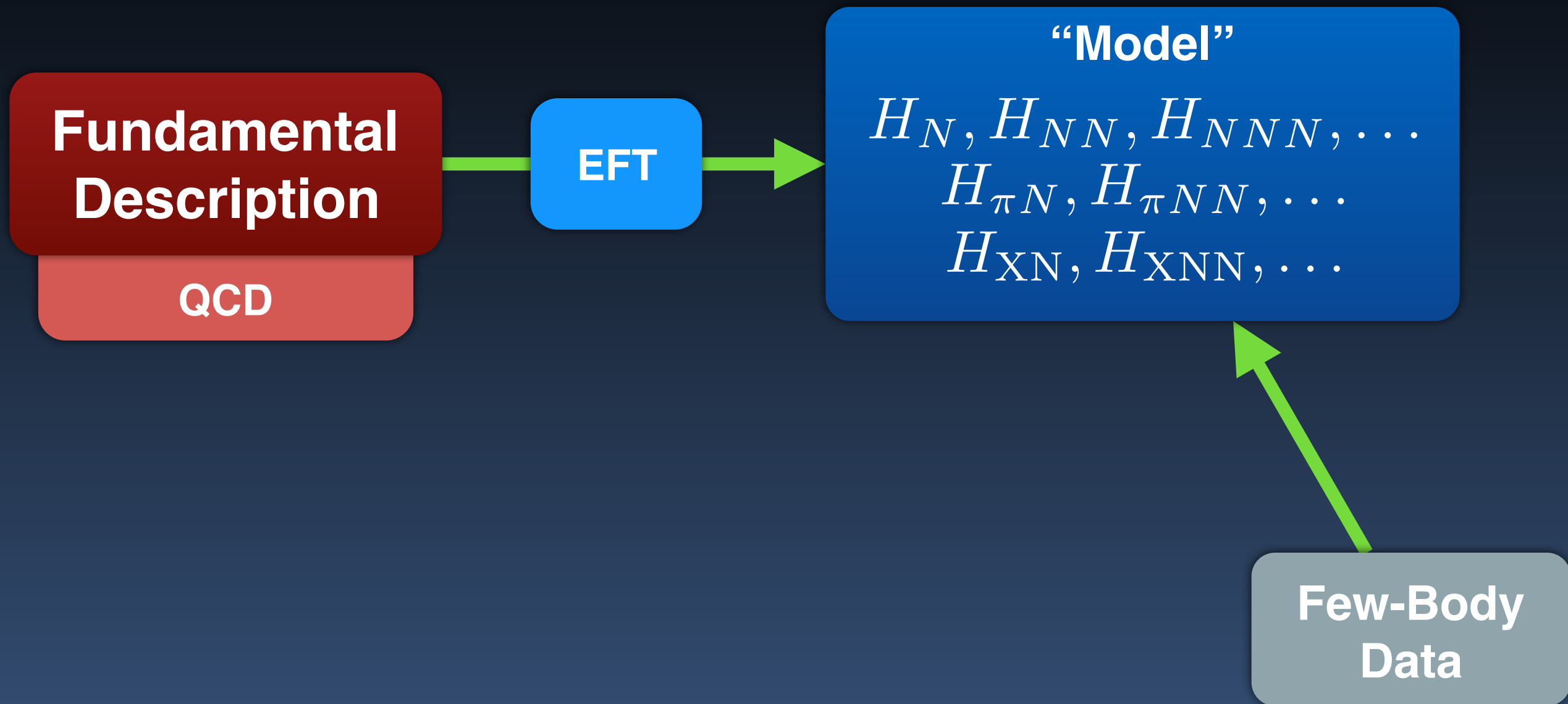
**Fundamental
Description**

QCD

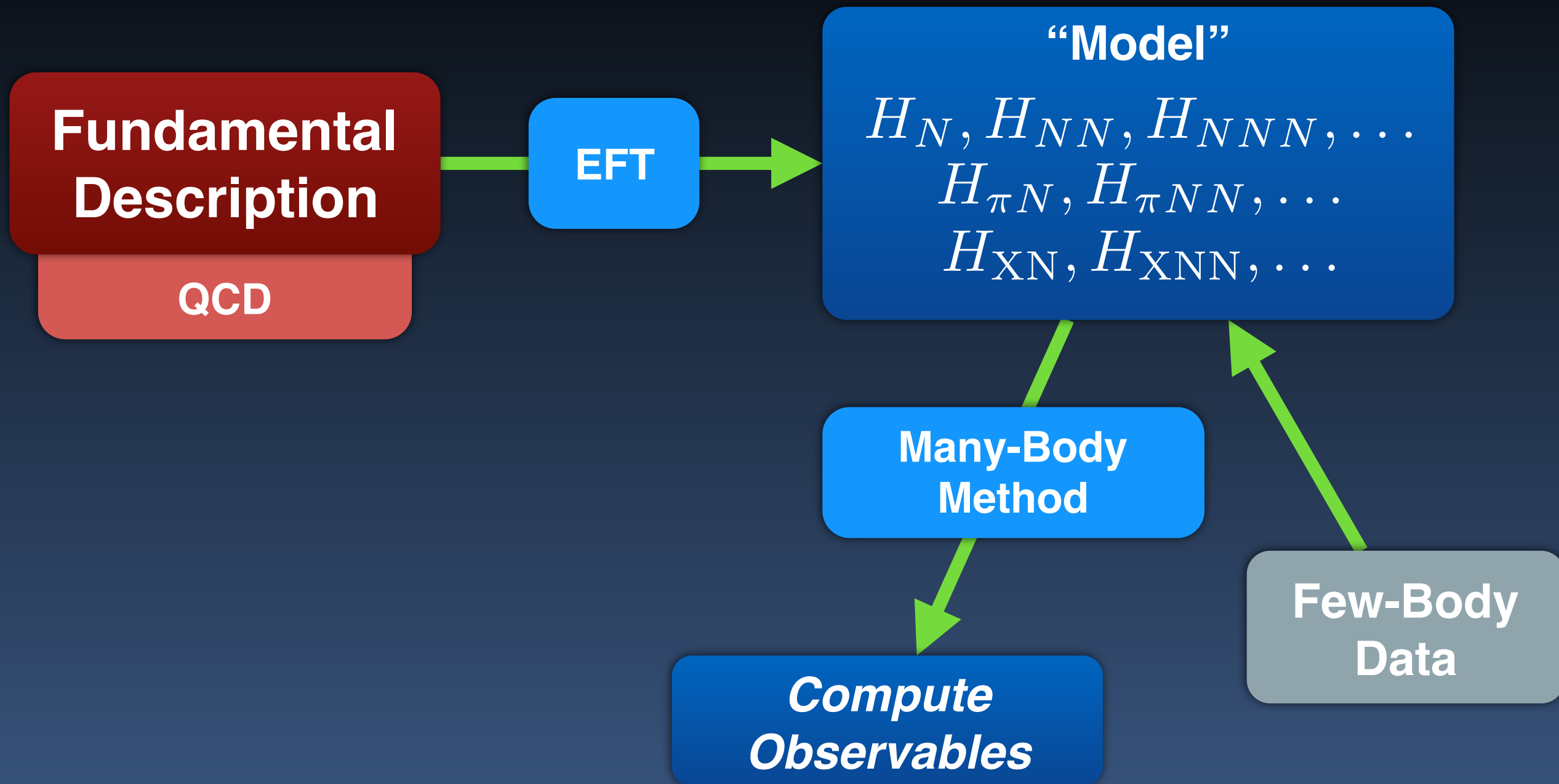
The Modern *Ab Initio* Approach to Low Energy Nuclear Physics



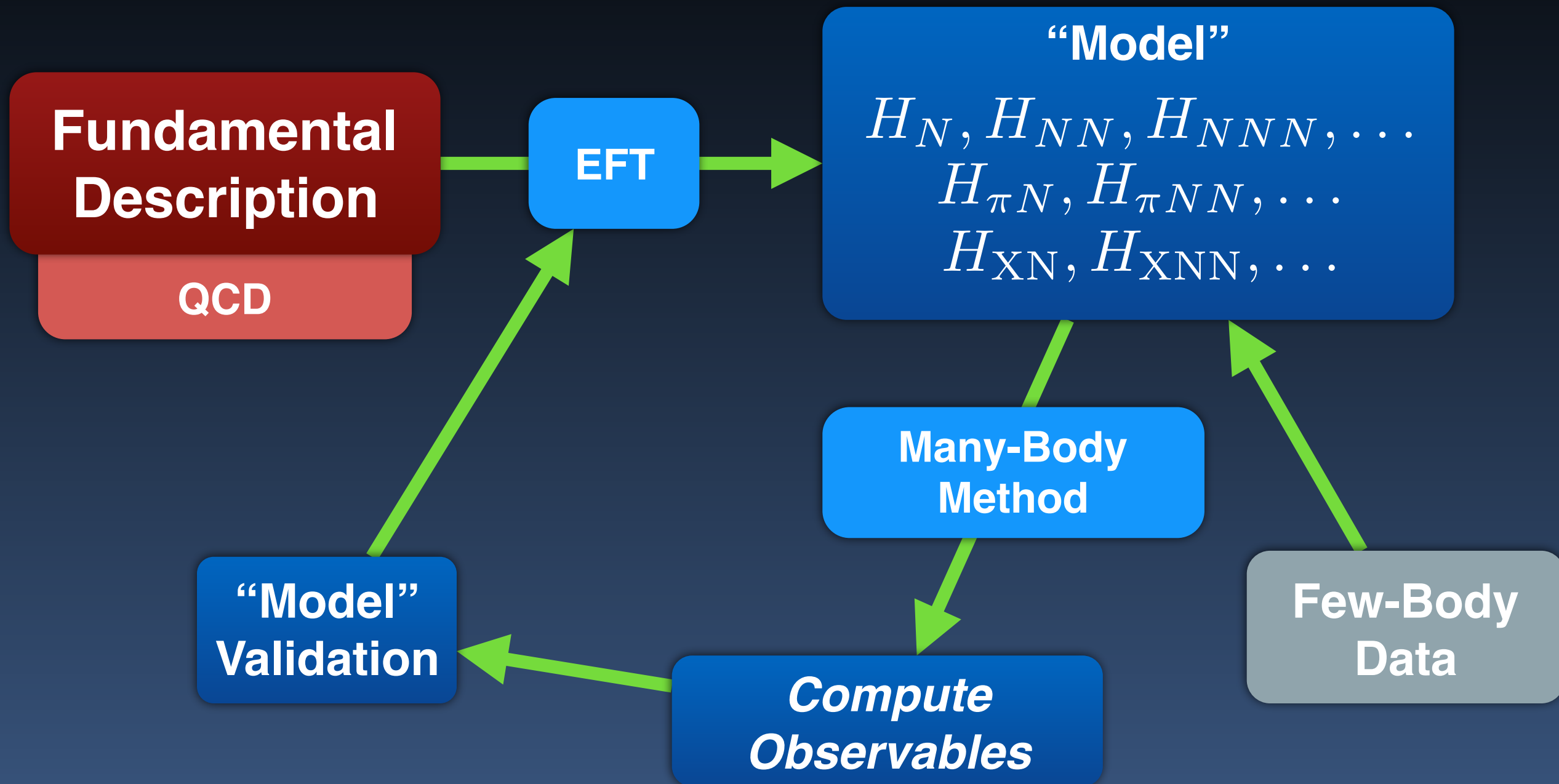
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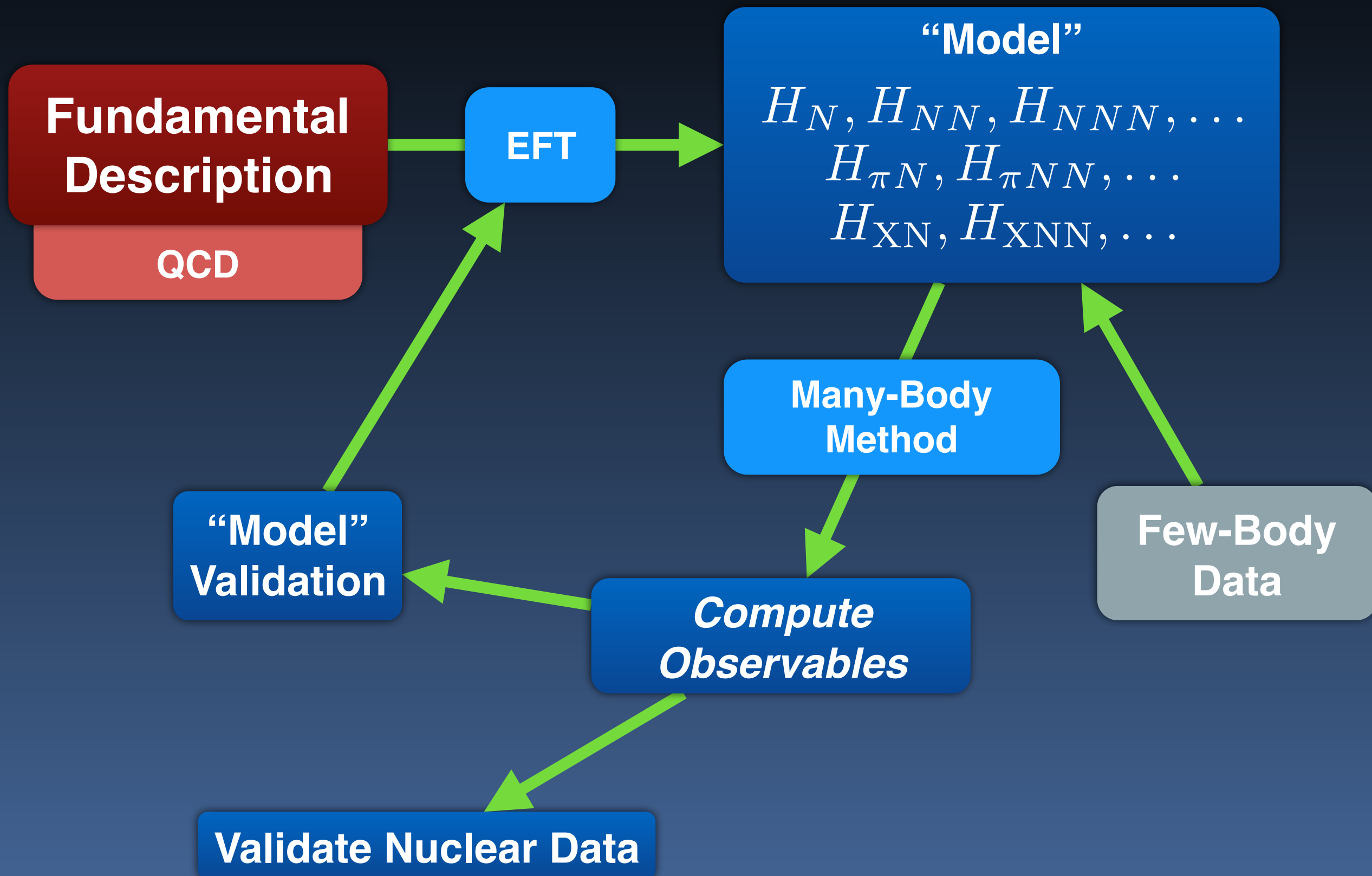
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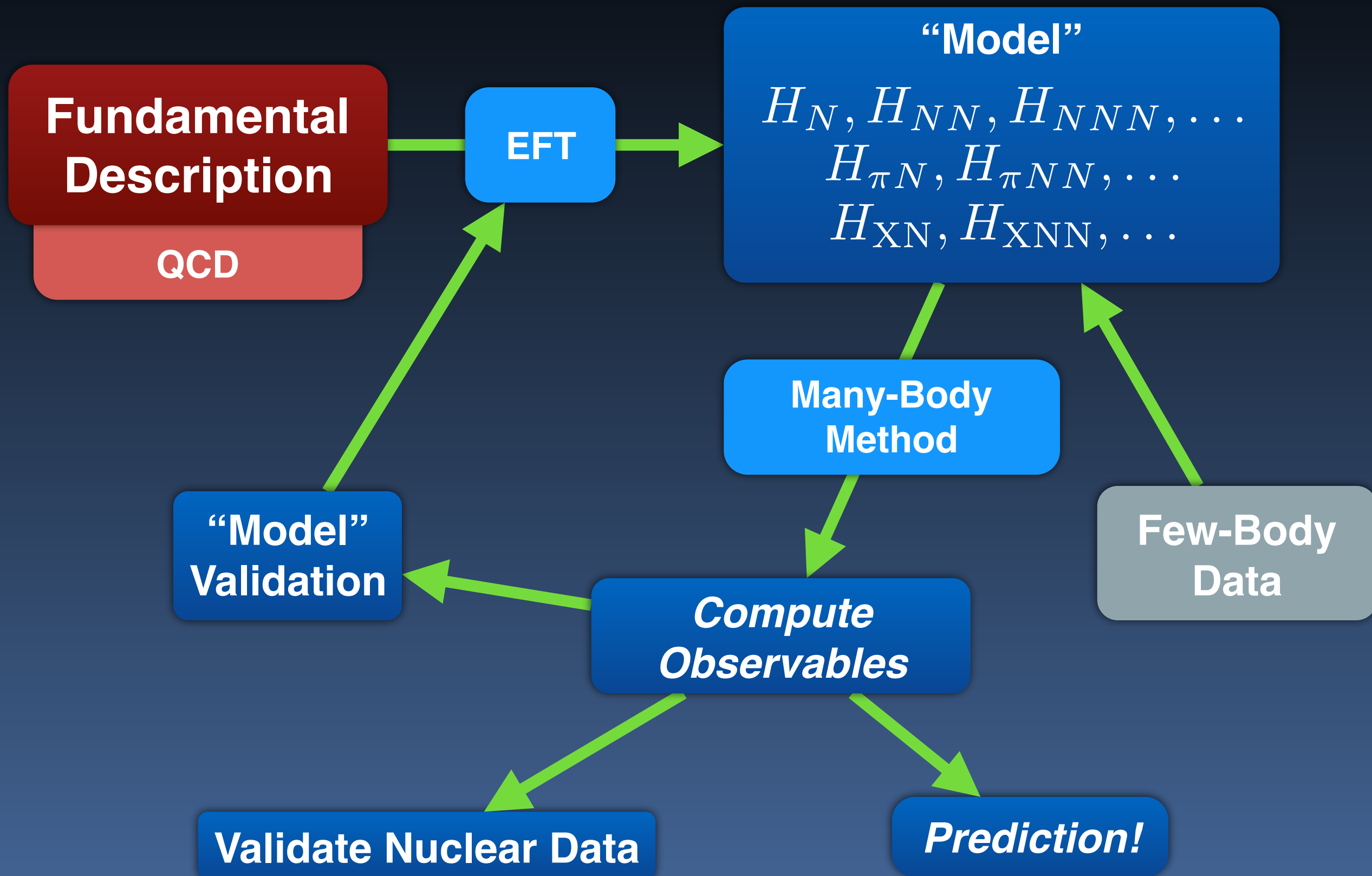
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Nuclear forces from chiral effective field theory

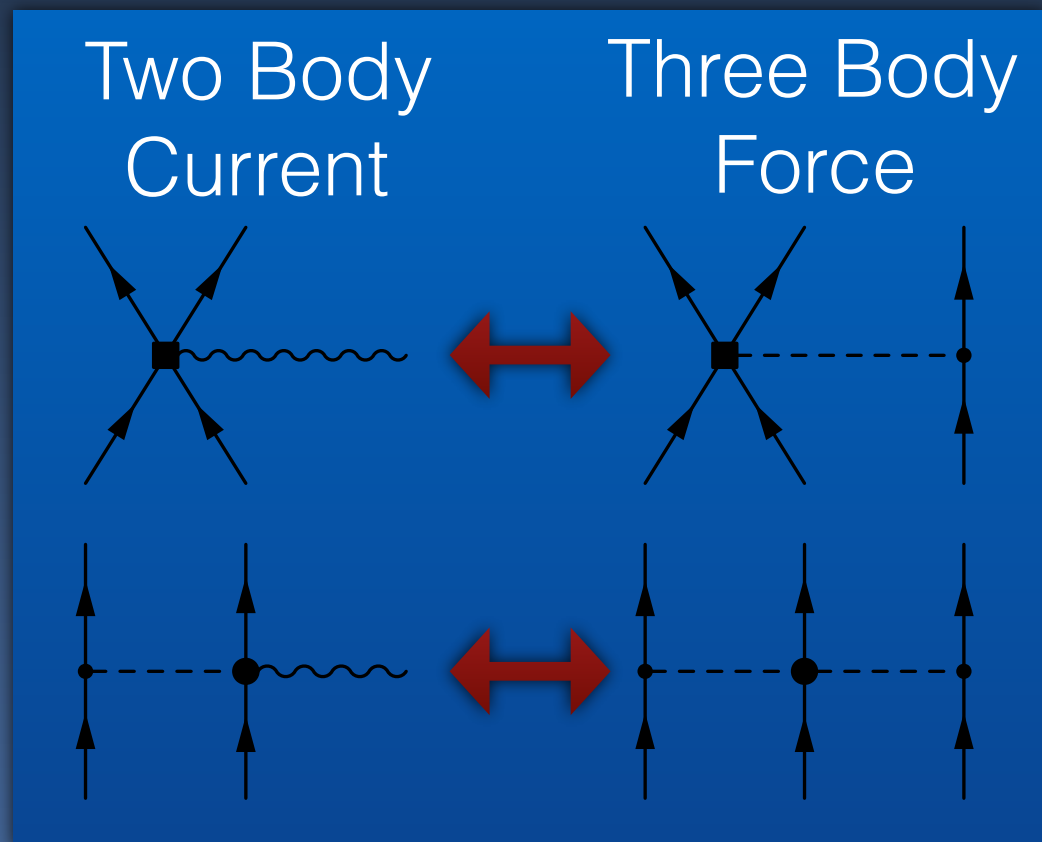
[Weinberg; van Kolck; Epelbaum et al.; Entem & Machleidt; ...]

We can use EFT to build nuclear forces!

- Systematically improvable

$$\langle O(Q) \rangle = \langle O_{\text{EFT}}^{(\nu)}(Q) \rangle + \mathcal{O}\left(\frac{Q}{\Lambda}\right)^{\nu+1}$$

- Connects different sets of strong-interaction phenomena:
 $\pi\pi\text{N}$, NN , NNN



	2N forces	3N forces	4N forces
LO ($\frac{Q^0}{\Lambda^0}$)			
NLO ($\frac{Q^2}{\Lambda^2}$)			
N ² LO ($\frac{Q^3}{\Lambda^3}$)			
N ³ LO ($\frac{Q^4}{\Lambda^4}$)			
	+ ...	+ ...	+ ...

Nuclear forces from chiral effective field theory

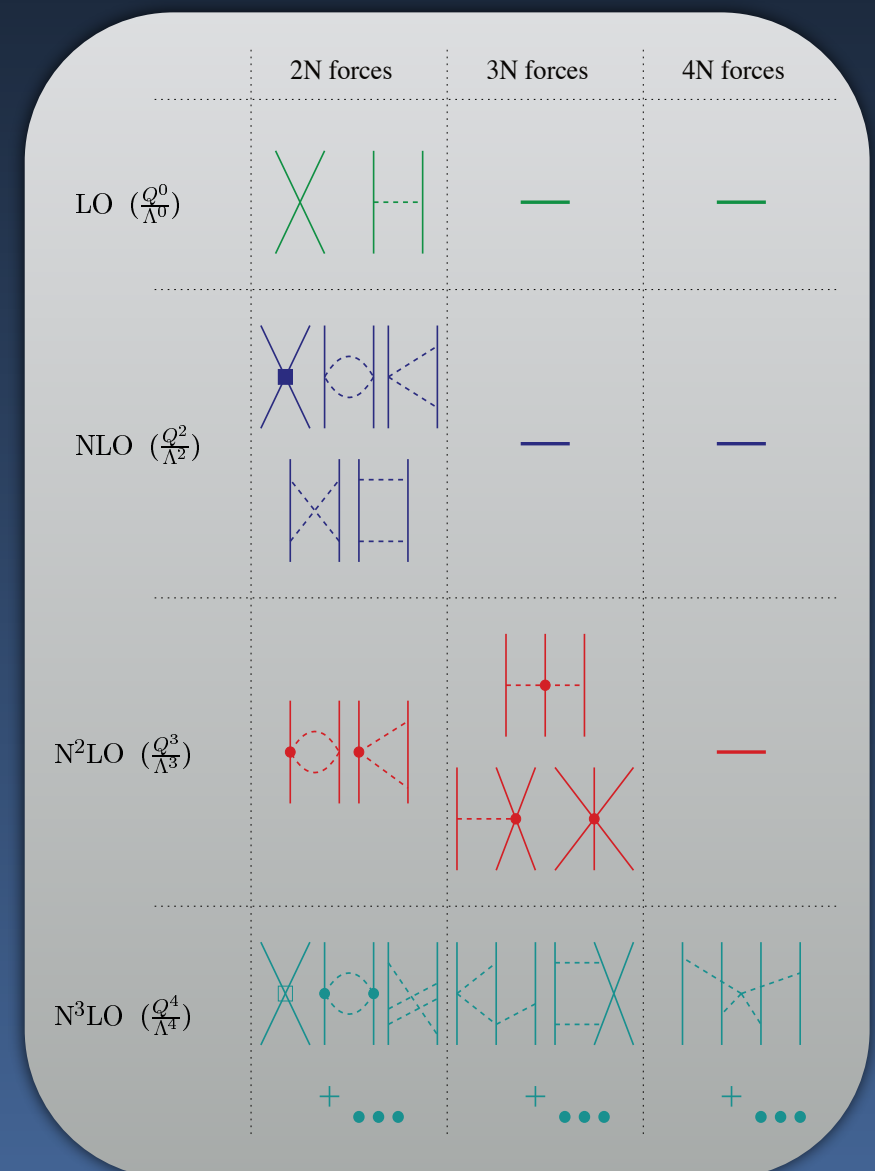
[Weinberg; van Kolck; Epelbaum et al.; Entem & Machleidt; ...]

$$P(\theta; \Lambda, \text{Data}) \propto \exp \left[-\frac{1}{2} \sum_i \frac{\left(y_i^{(\text{data})} - y_i^{(\text{EFT})}(\theta, \Lambda) \right)^2}{\sigma_i^2} \right]$$

$$y_{i,2\text{B}/3\text{B}}^{(\text{EFT})}(\theta, \Lambda) = m(\theta, \Lambda) + u_{\text{EFT}}$$

θ = Set of Parameters (LECs)

Λ = Regularization Specification



Nuclear forces from chiral effective field theory

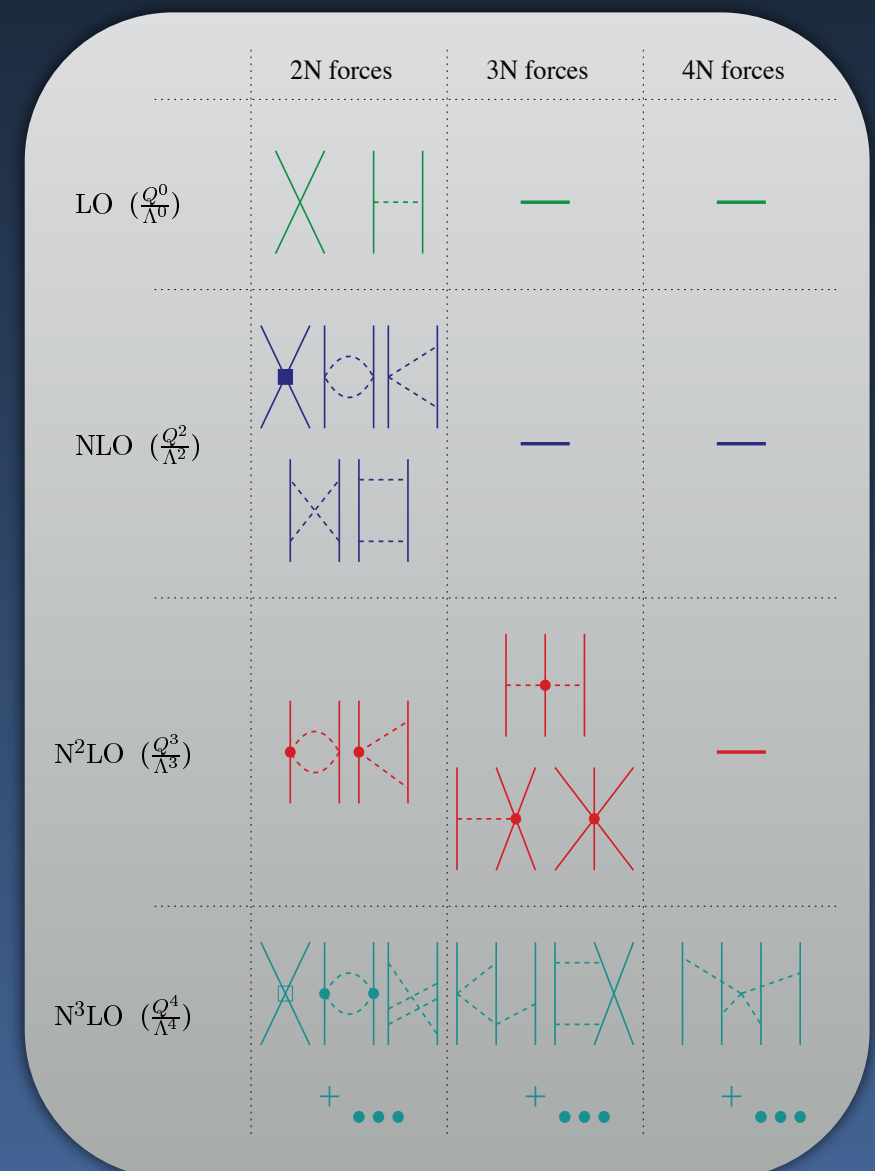
[Weinberg; van Kolck; Epelbaum et al.; Entem & Machleidt; ...]

$$P(\theta; \Lambda, \text{Data}) \propto \exp \left[-\frac{1}{2} \sum_i \frac{\left(y_i^{(\text{data})} - y_i^{(\text{EFT})}(\theta, \Lambda) \right)^2}{\sigma_i^2} \right]$$

$$y_{i,2B/3B}^{(\text{EFT})}(\theta, \Lambda) = m(\theta, \Lambda) + u_{\text{EFT}}$$

Open Problems

- How to choose regularization?
- How to correctly order power counting?
- What data should we fit?



Two Main Classes of Ab Initio Many-Body Methods

Quasi Exact

Quantum Monte Carlo

Direct Basis Expansion

Exact many-body correlations

Scale factorially in number
of particles

Reference State

Coupled Clusters Expansion

In-medium SRG

(Indirect Basis Expansion)

Approximate many-body
correlations

Scale polynomially in
number of particles

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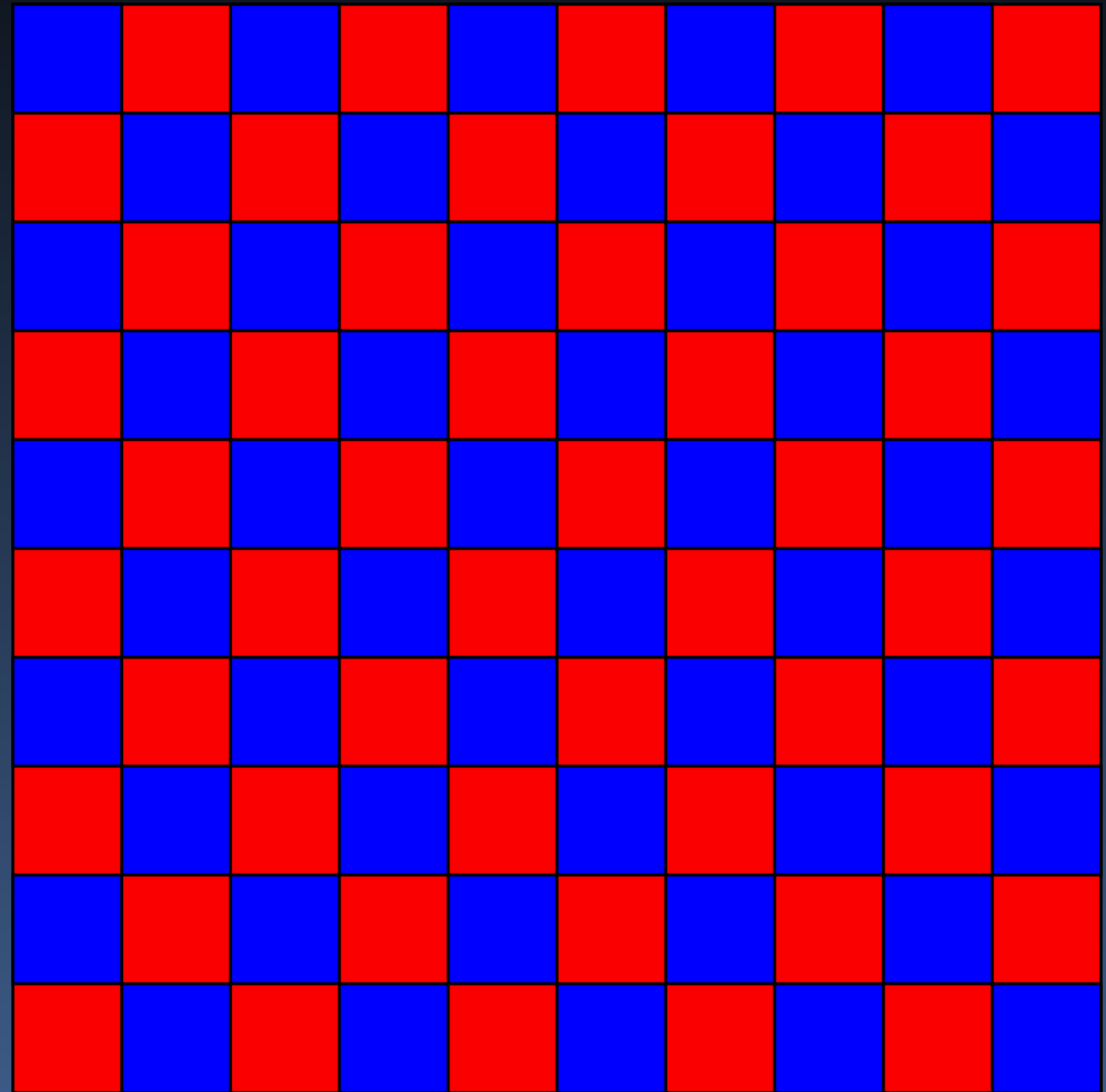
Approximate many-body
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Scale polynomially in
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Direct Basis Expansion

For only a few (< 12) particles, one can construct the many-body Hamiltonian as a large matrix

Many-body correlations are included exactly!

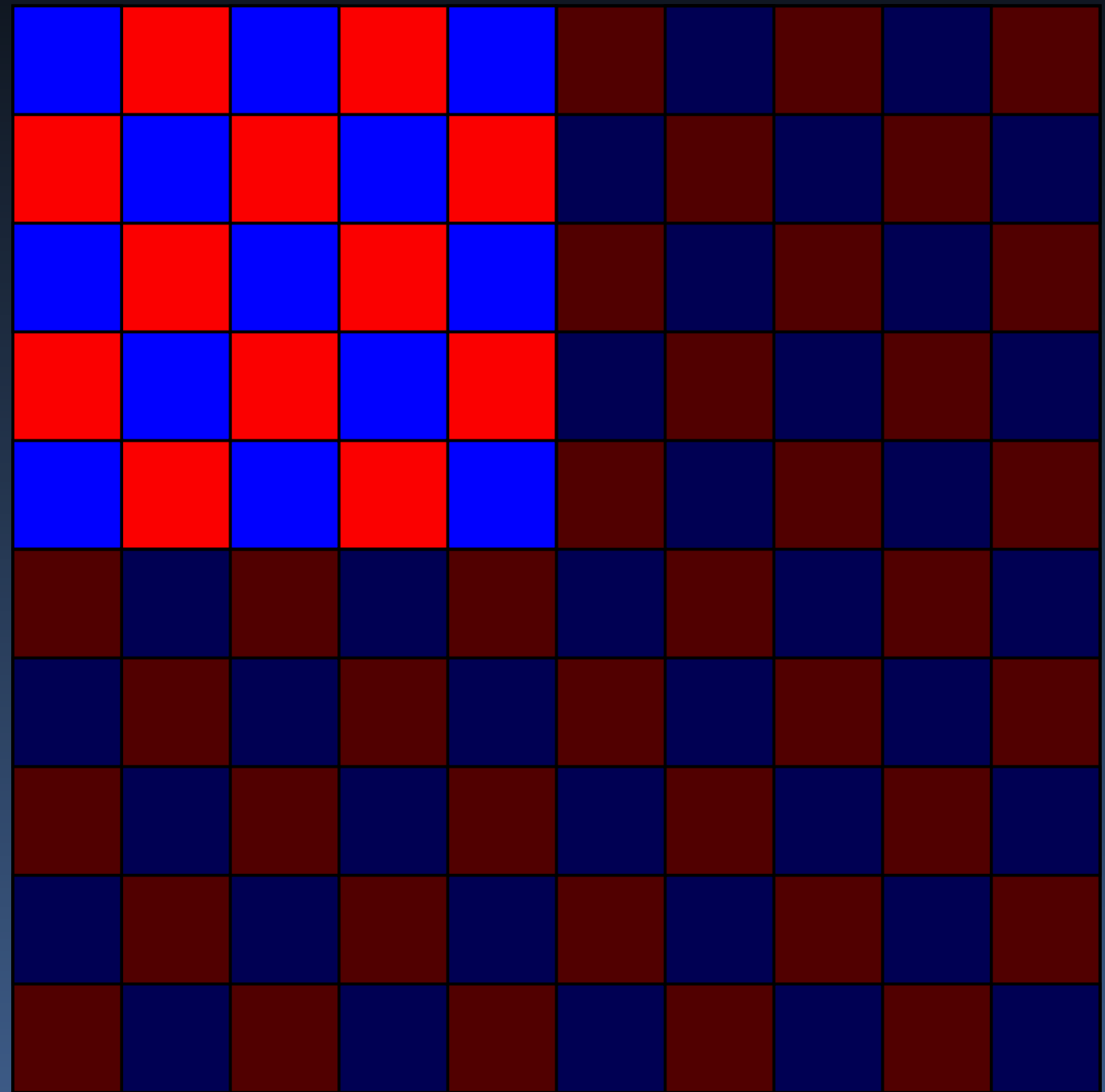


Quasi-Exact Many-Body Methods

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Uncertainty enters from the basis truncation

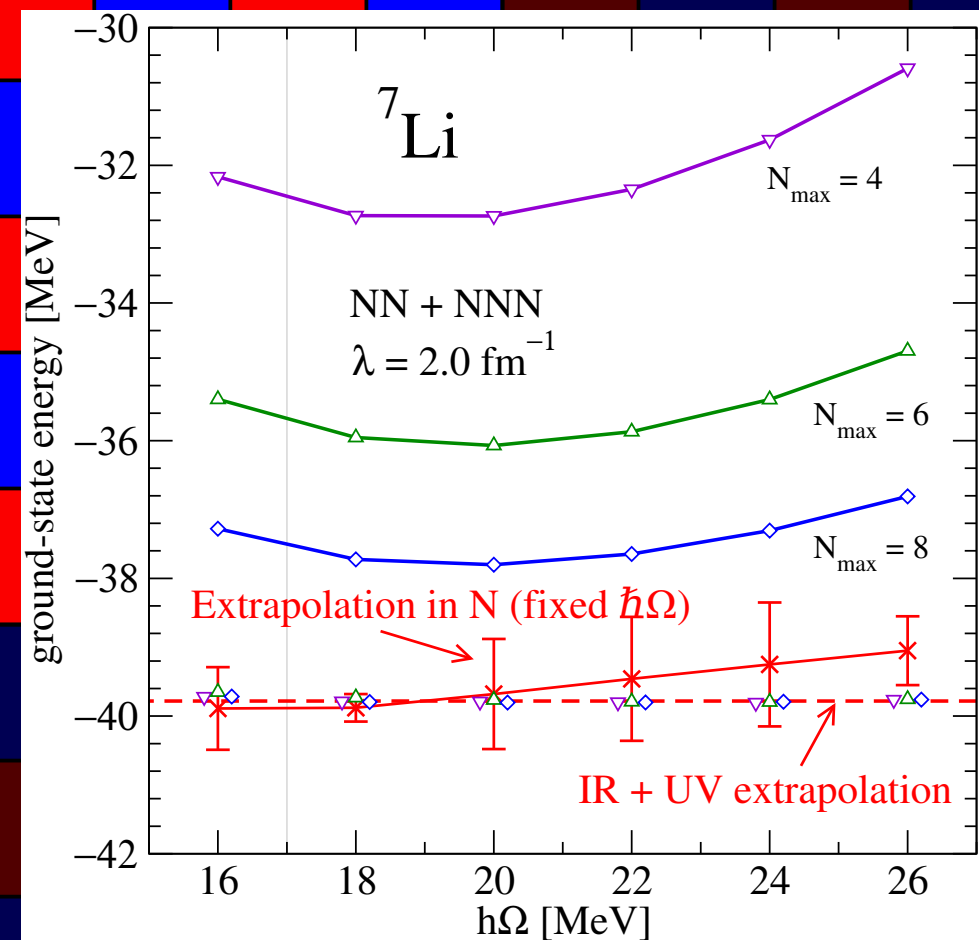


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E. D. Jurgenson et al., PRC 87 054312 (2013)

Quantifying Errors in Finite Oscillator Basis

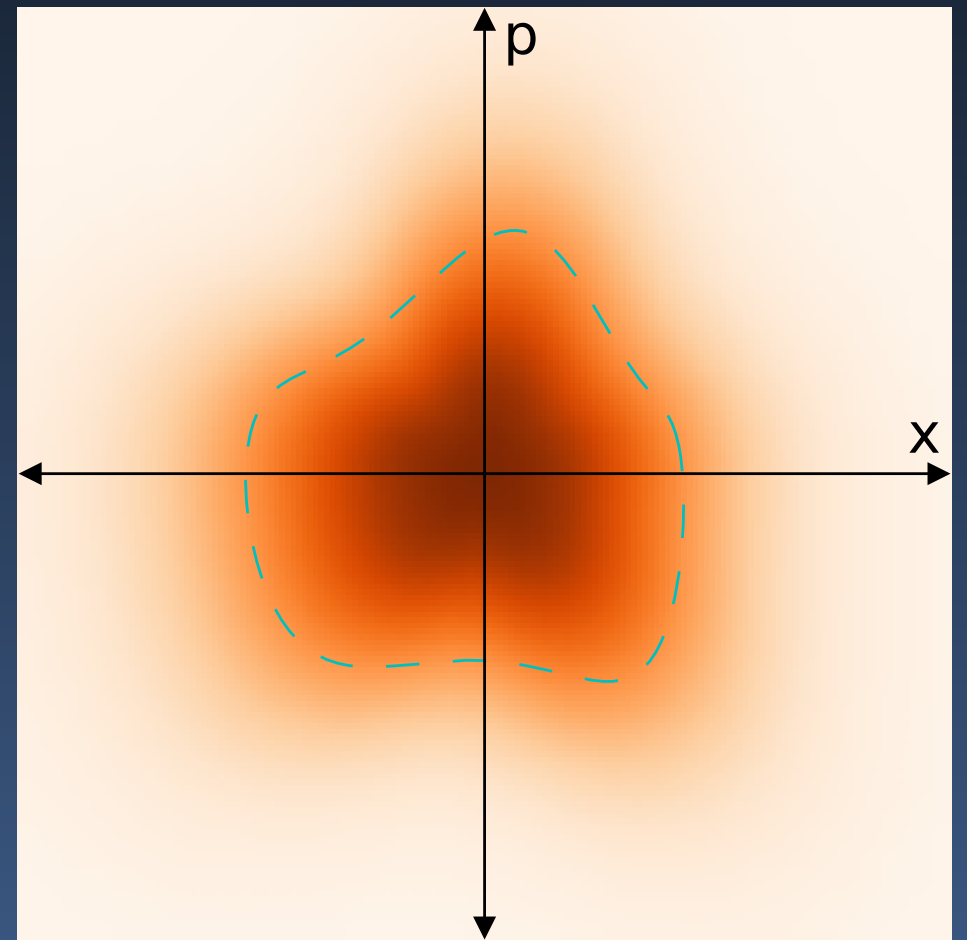
R. J. Furnstahl et al., PRC 86 031301(R)

S. A. Coon et al. PRC 86 054002

$$E - E_{\text{exact}} = \Delta E =$$

$$\Delta E_{\text{IR}}(L) + \Delta E_{\text{mix}}(L, \Lambda) + \Delta E_{\text{UV}}(\Lambda)$$

$$u_i(r) = \sum_{n=0}^{N_{\text{max}}} c_{i,n} \phi_n(r)$$



Quantifying Errors in Finite Oscillator Basis

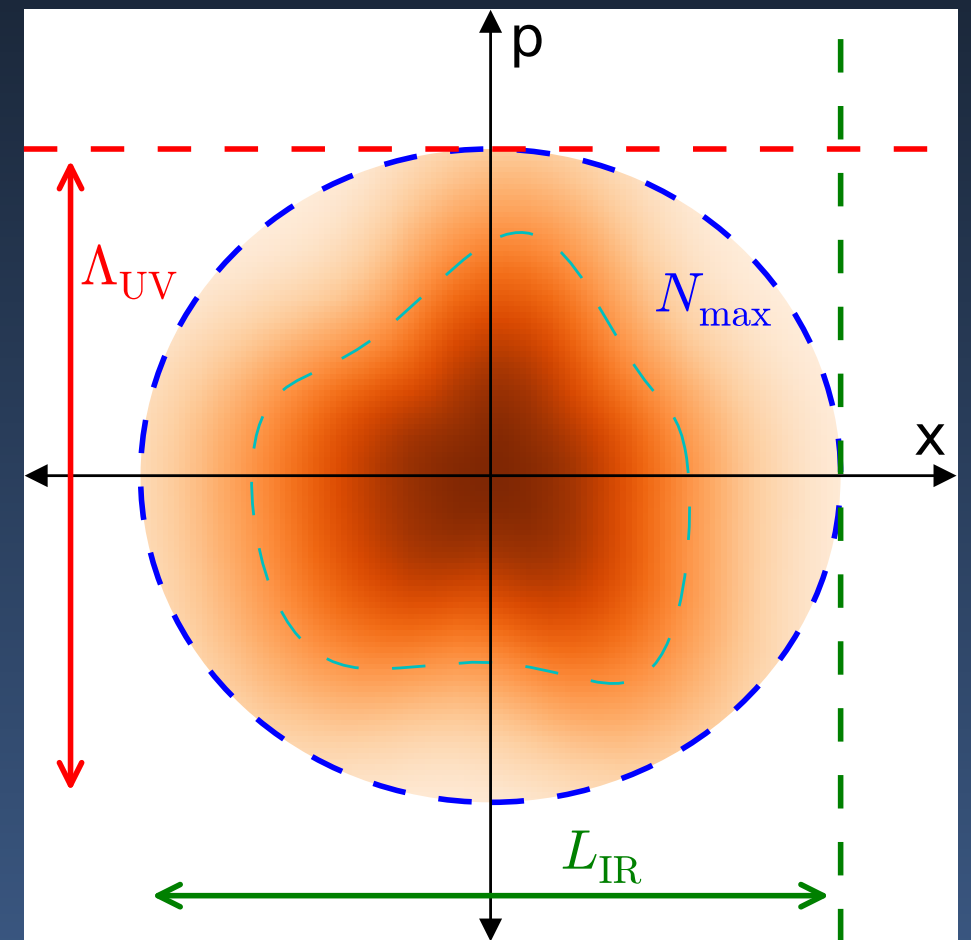
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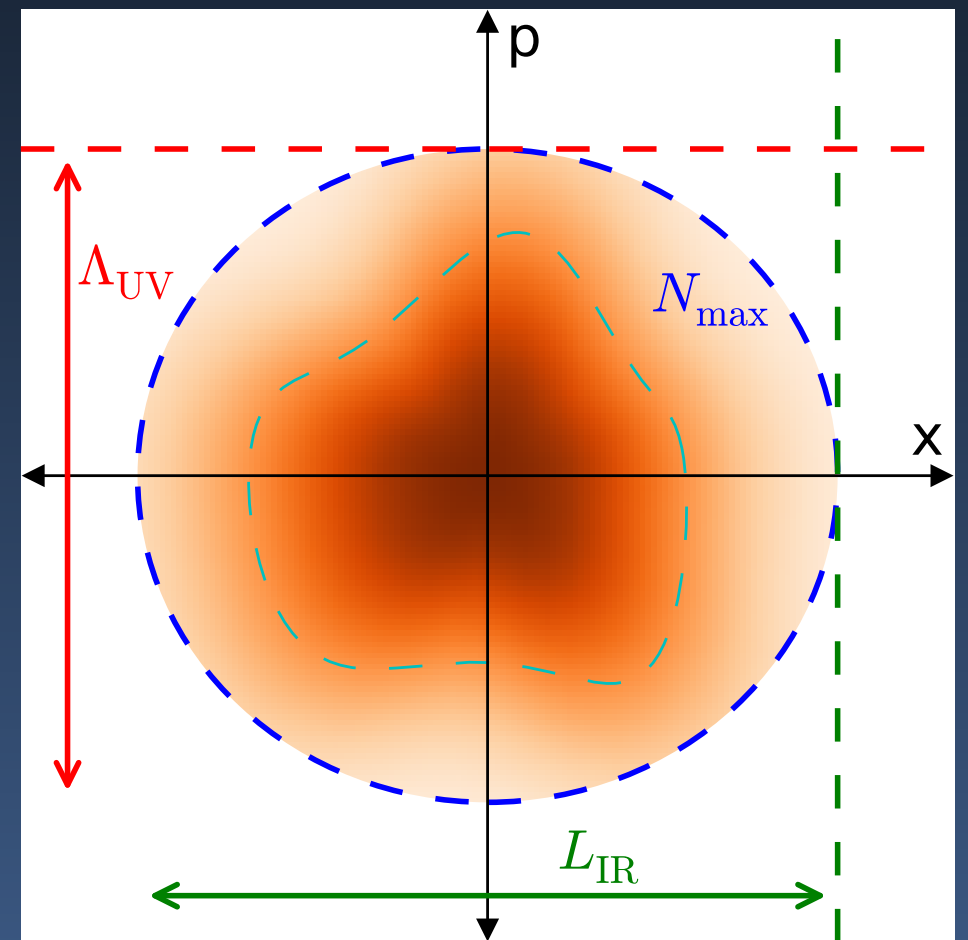
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The Infrared Error Term:

$$u_i(r) = \sum_{n=0}^{N_{\text{max}}} c_{i,n} \phi_n(r)$$



Quantifying Errors in Finite Oscillator Basis

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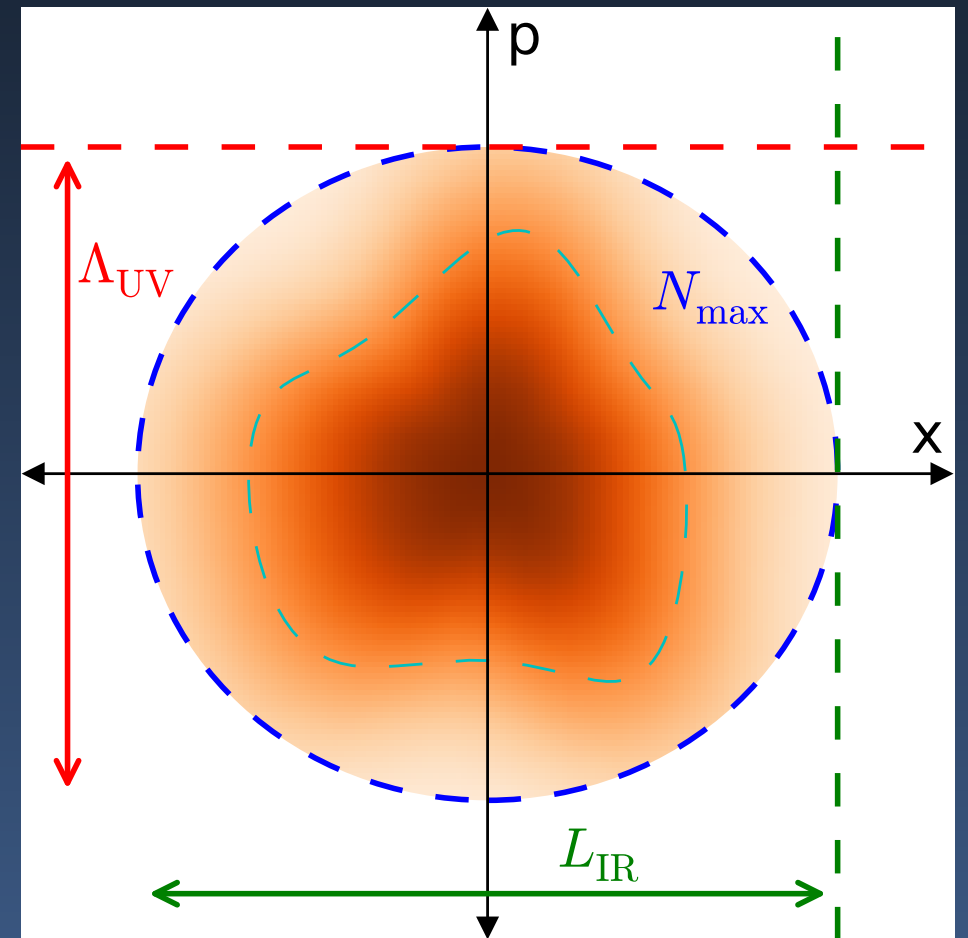
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The Infrared Error Term:

- Universal for short range Hamiltonians

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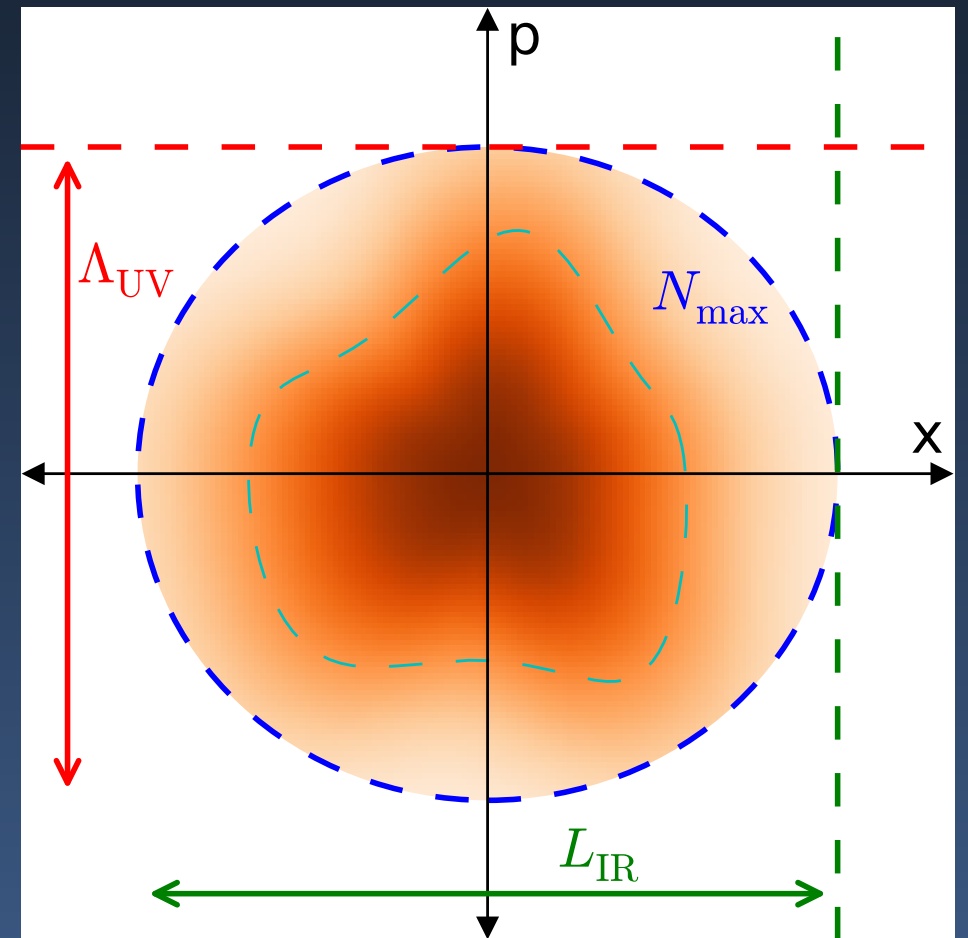
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The Infrared Error Term:

- Universal for short range Hamiltonians
- Can be systematically corrected

$$u_i(r) = \sum_{n=0}^{N_{\text{max}}} c_{i,n} \phi_n(r)$$



Quantifying Errors in Finite Oscillator Basis

R. J. Furnstahl et al., PRC 86 031301(R)

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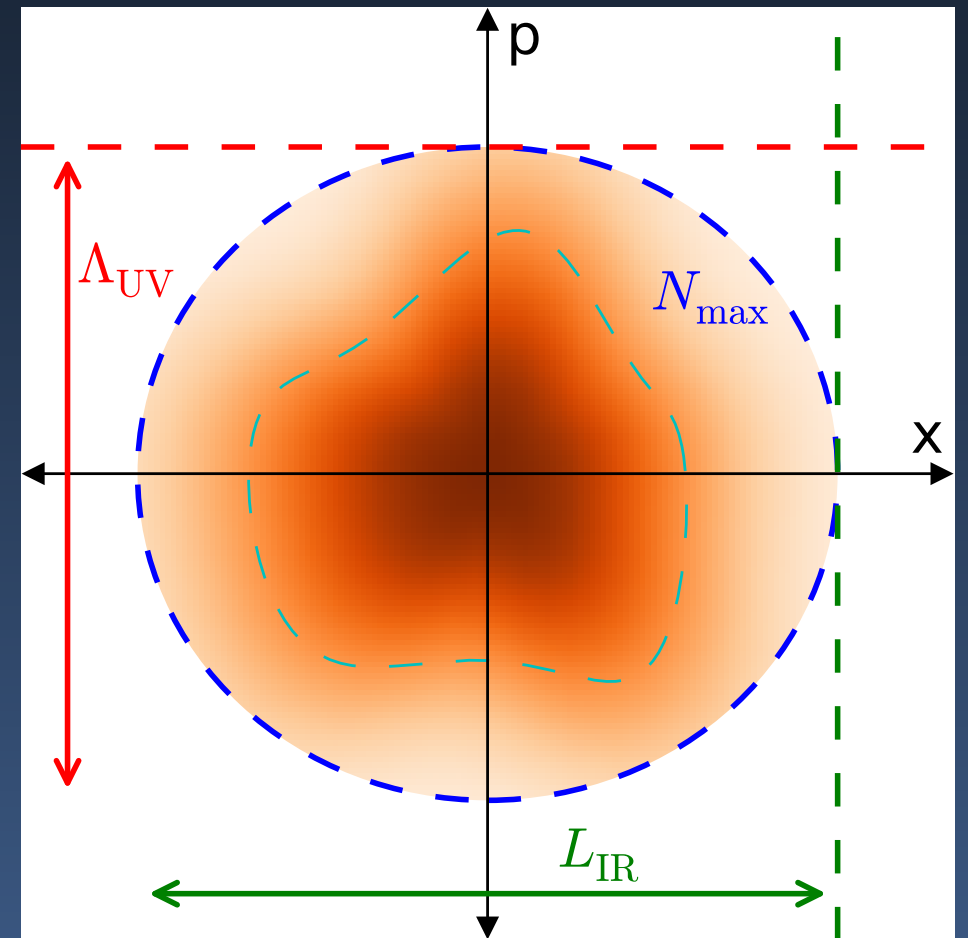
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- Universal for short range Hamiltonians
- Can be systematically corrected
- Resembles error from putting system into a infinite well

$$u_i(r) = \sum_{n=0}^{N_{\text{max}}} c_{i,n} \phi_n(r)$$



Quantifying Errors in Finite Oscillator Basis

R. J. Furnstahl et al., PRC 86 031301(R)

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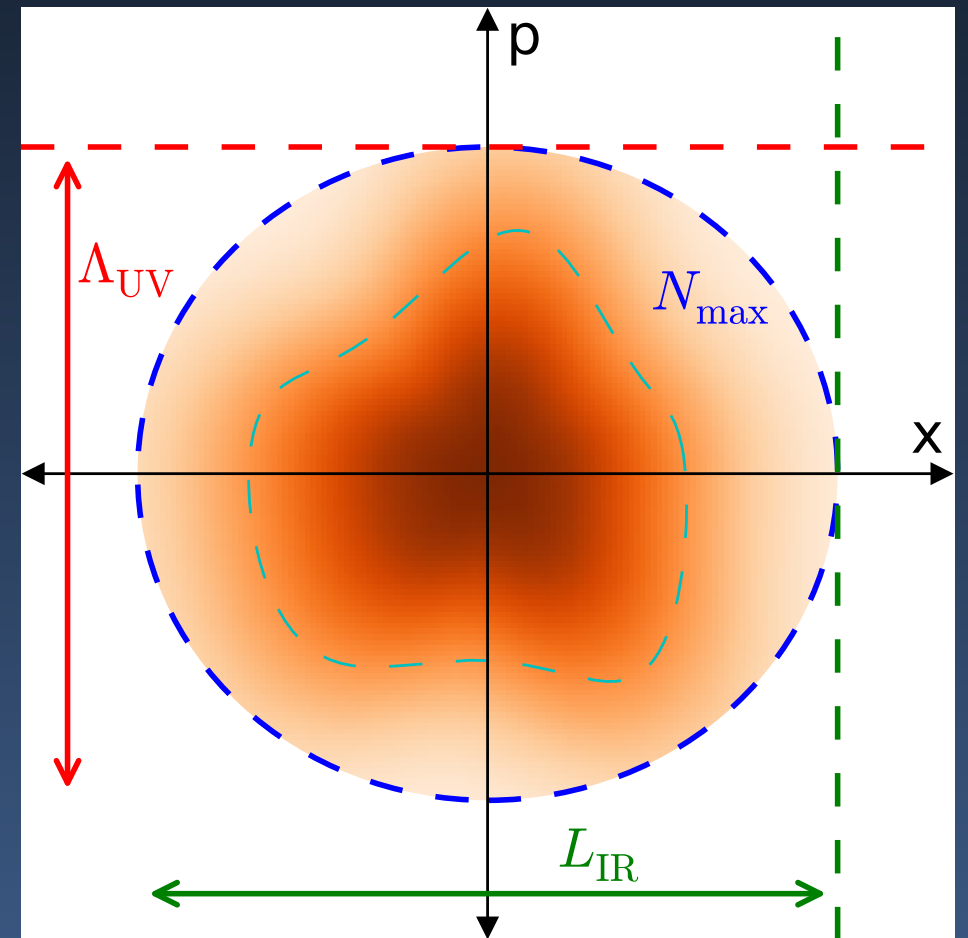
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The Infrared Error Term:

- Universal for short range Hamiltonians
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Requires Calculation to be UV converged

$$u_i(r) = \sum_{n=0}^{N_{\text{max}}} c_{i,n} \phi_n(r)$$



Quantifying Errors in Finite Oscillator Basis

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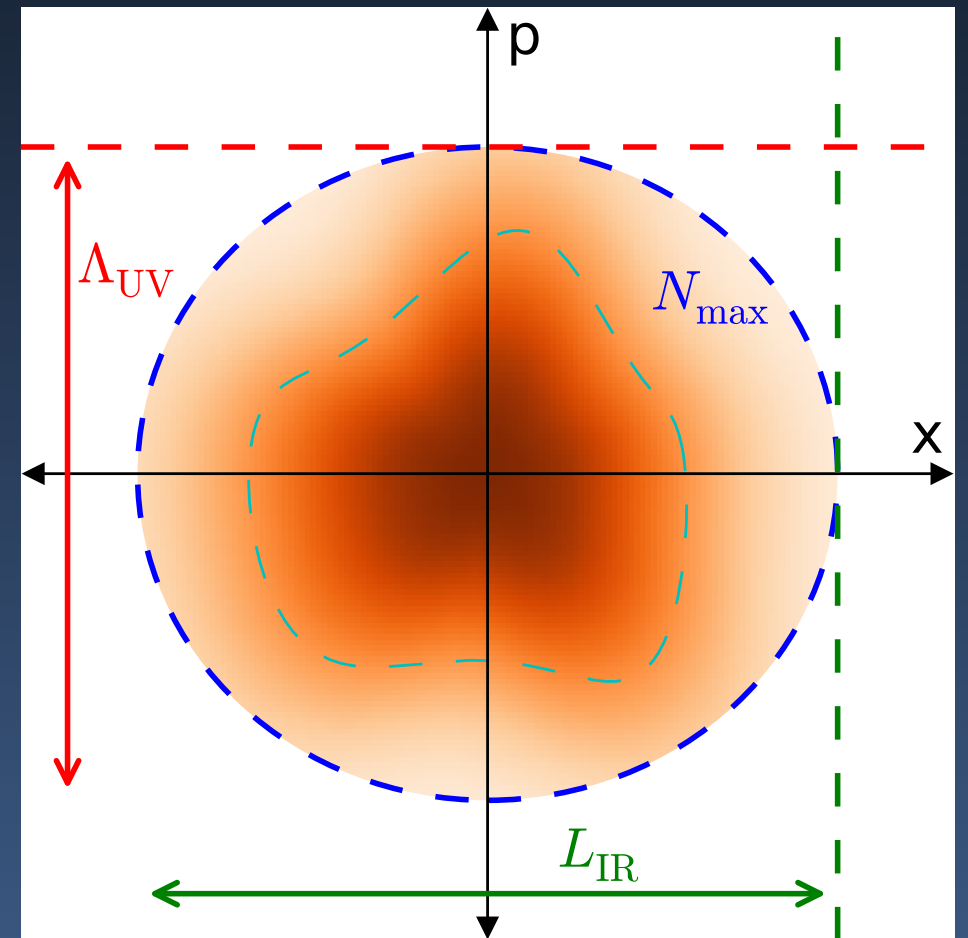
S. N. More et al., PRC 87 044326

R. J. Furnstahl et al., PRC 89 044301

$$\langle r \rangle / \langle r \rangle_{\infty} = 1 + \sum_{m=1}^{\infty} P_m(\beta) e^{-m\beta}$$

$$\langle E \rangle / \langle E \rangle_{\infty} = 1 + \sum_{m=1}^{\infty} a_m \frac{p_m(\beta)}{p_m(-\beta)} e^{-m\beta}$$

$$u_i(r) = \sum_{n=0}^{N_{\text{max}}} c_{i,n} \phi_n(r)$$



This is only one part of the Basis Truncation Error!!!

Quantifying Errors in Finite Oscillator Basis

$$y_i^{(\text{thy})} = m(\theta, \Lambda) + u_{\text{EFT}} + v_{\text{Basis}}$$

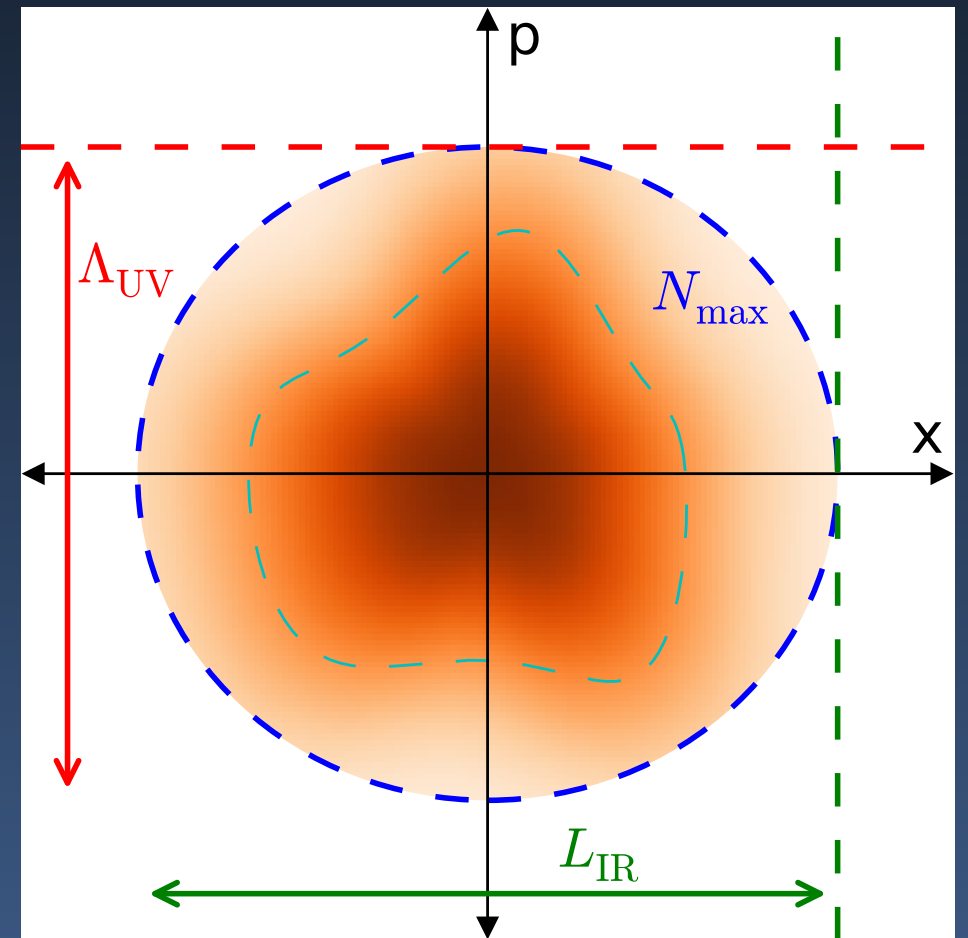
$$v_{\text{Basis}} = v_{\text{IR}} + v_{\text{UV}} + v_{\text{mixed}}$$

v_{IR}

$$\langle r \rangle / \langle r \rangle_{\infty} = 1 + \sum_{m=1}^{\infty} P_m(\beta) e^{-m\beta}$$

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Coupled Clusters Expansion

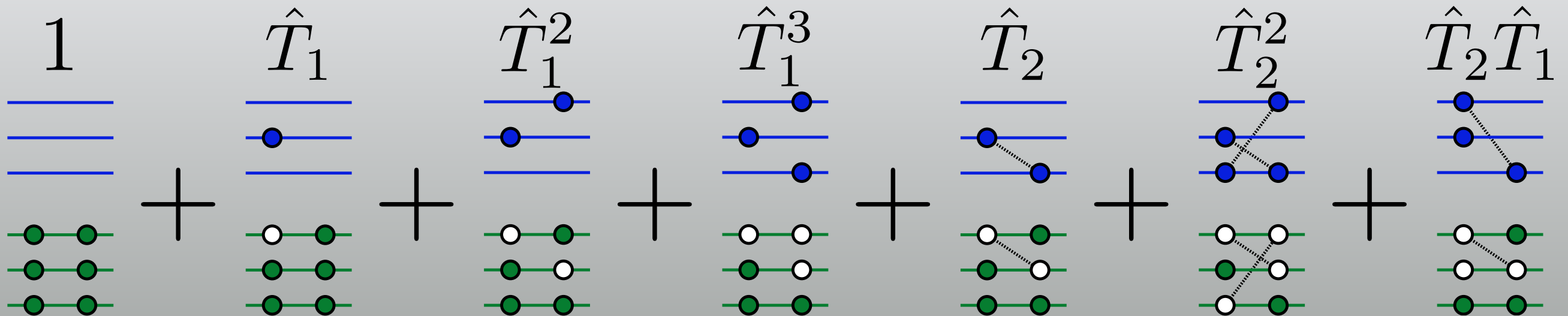
$$\hat{T} = \text{[diagram 1]} + \text{[diagram 2]} + \text{[diagram 3]} + \dots$$

The diagram shows the expansion of the cluster operator \hat{T} as a sum of terms. Each term is represented by a diagram of particles (red and blue spheres) and their interactions. The first term shows three particles (one red, two blue) with no interactions. The second term shows three particles with a single interaction line connecting two of them. The third term shows three particles with two interaction lines forming a triangle. The expansion continues with higher-order terms indicated by ellipses.

$$|\Psi\rangle = e^{\hat{T}} |\Phi\rangle$$

$$0 = \langle \Phi_i^a | e^{-\hat{T}} \hat{H} e^{\hat{T}} | \Phi \rangle$$

$$0 = \langle \Phi_{ij}^{ab} | e^{-\hat{T}} \hat{H} e^{\hat{T}} | \Phi \rangle$$



Coupled Clusters Expansion

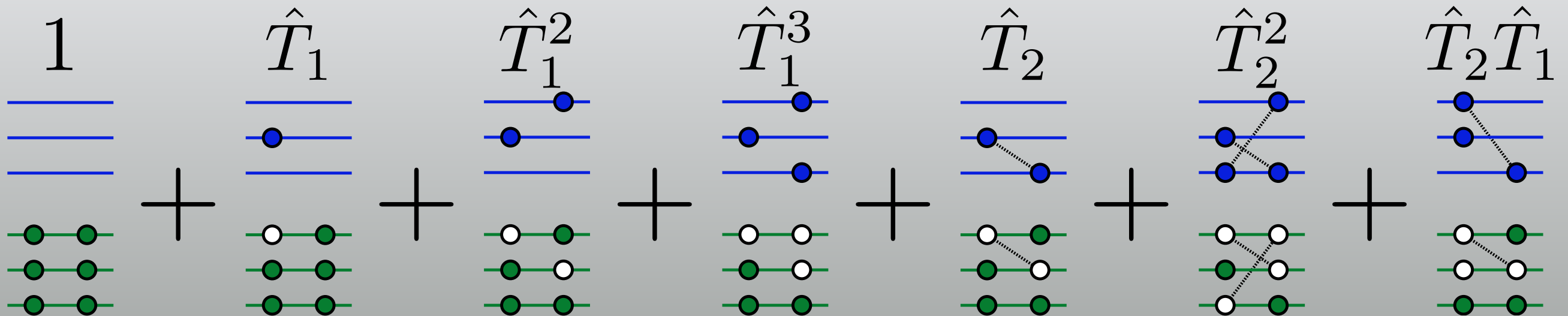
$$\hat{T} = \text{[diagram 1]} + \text{[diagram 2]} + \text{[diagram 3]} + \dots$$

The diagram shows the expansion of the cluster operator \hat{T} as a sum of terms. The first term is a single cluster with one blue and one red particle. The second term is a pair of clusters, one blue and one red. The third term is a pair of clusters, one blue and one red, with a red 'X' over it, indicating it is excluded. The expansion continues with an ellipsis.

$$|\Psi\rangle = e^{\hat{T}} |\Phi\rangle$$

$$0 = \langle \Phi_i^a | e^{-\hat{T}} \hat{H} e^{\hat{T}} | \Phi \rangle$$

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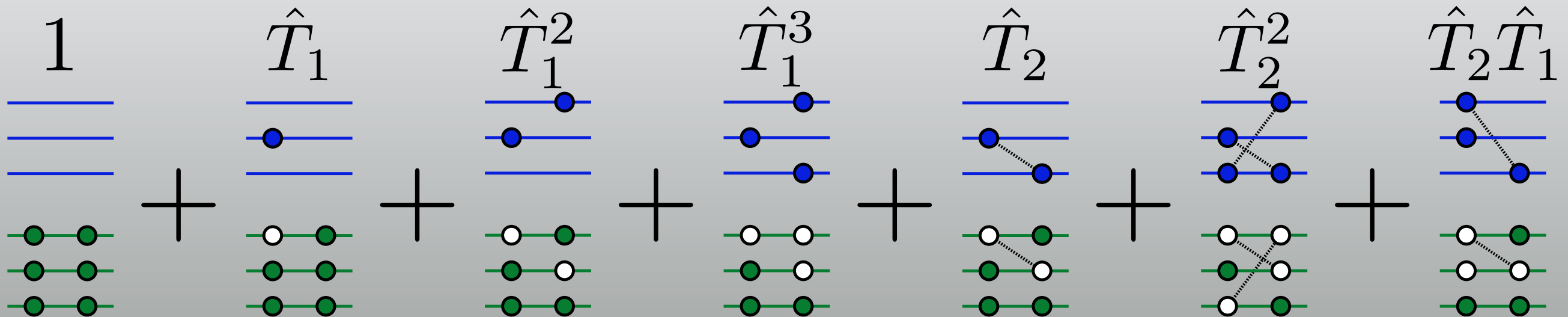
Reference State Methods (Coupled Clusters)

$$y_i^{(\text{thy})} = m(\theta, \Lambda) + u_{\text{EFT}} + v_{\text{Basis}} + w_{\text{M.B. corr}}$$

$$|\Psi\rangle = e^{\hat{T}} |\Phi\rangle$$

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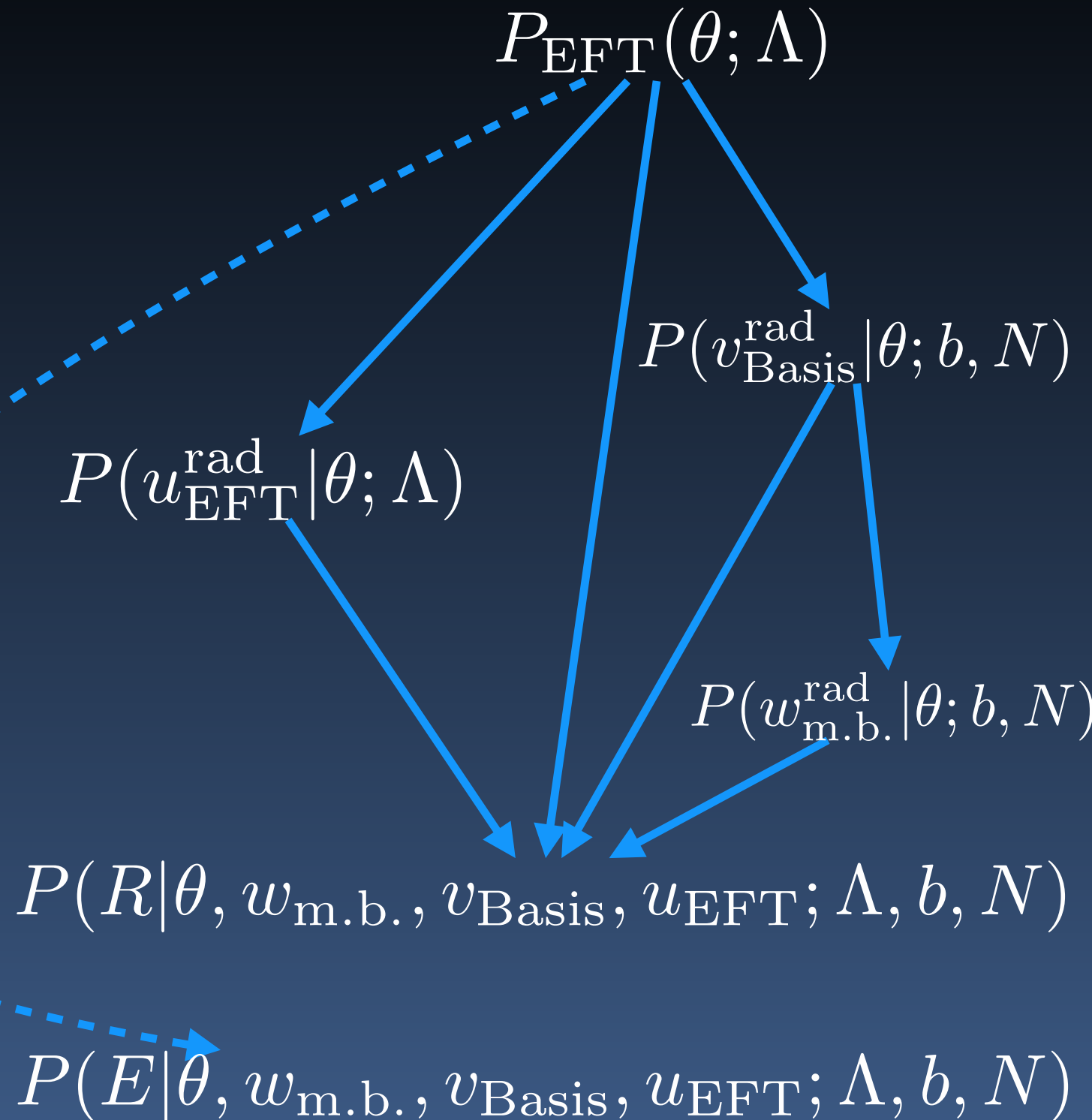


Sub Conclusions

Uncertainty enters in many places modern *ab initio* calculations

**Underlying interactions model
Basis systematics
Missing many-body correlations**

The relative sizes of each contribution can be observable dependent



Neutron and weak-charge distributions of the ^{48}Ca nucleus

Nature Physics 12, 186–190 (2016) [10.1038/nphys3529](https://doi.org/10.1038/nphys3529)

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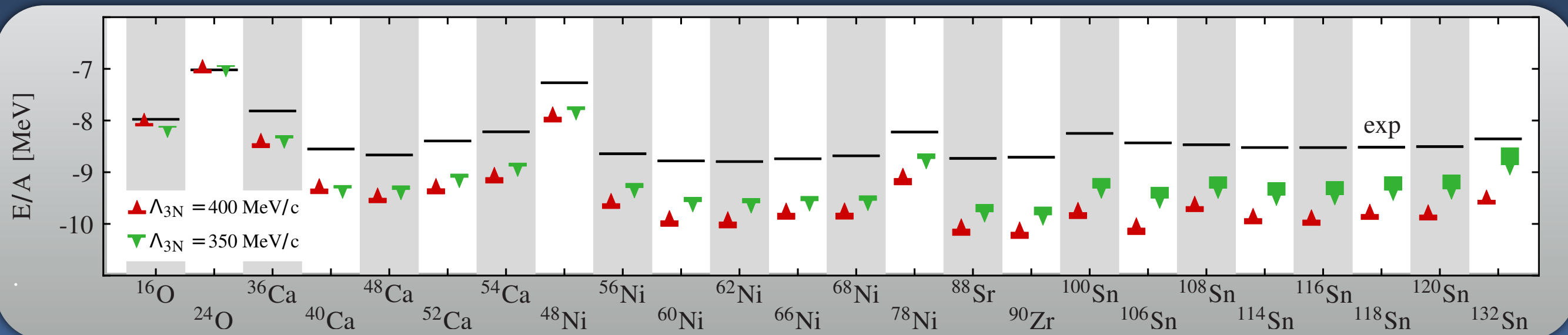
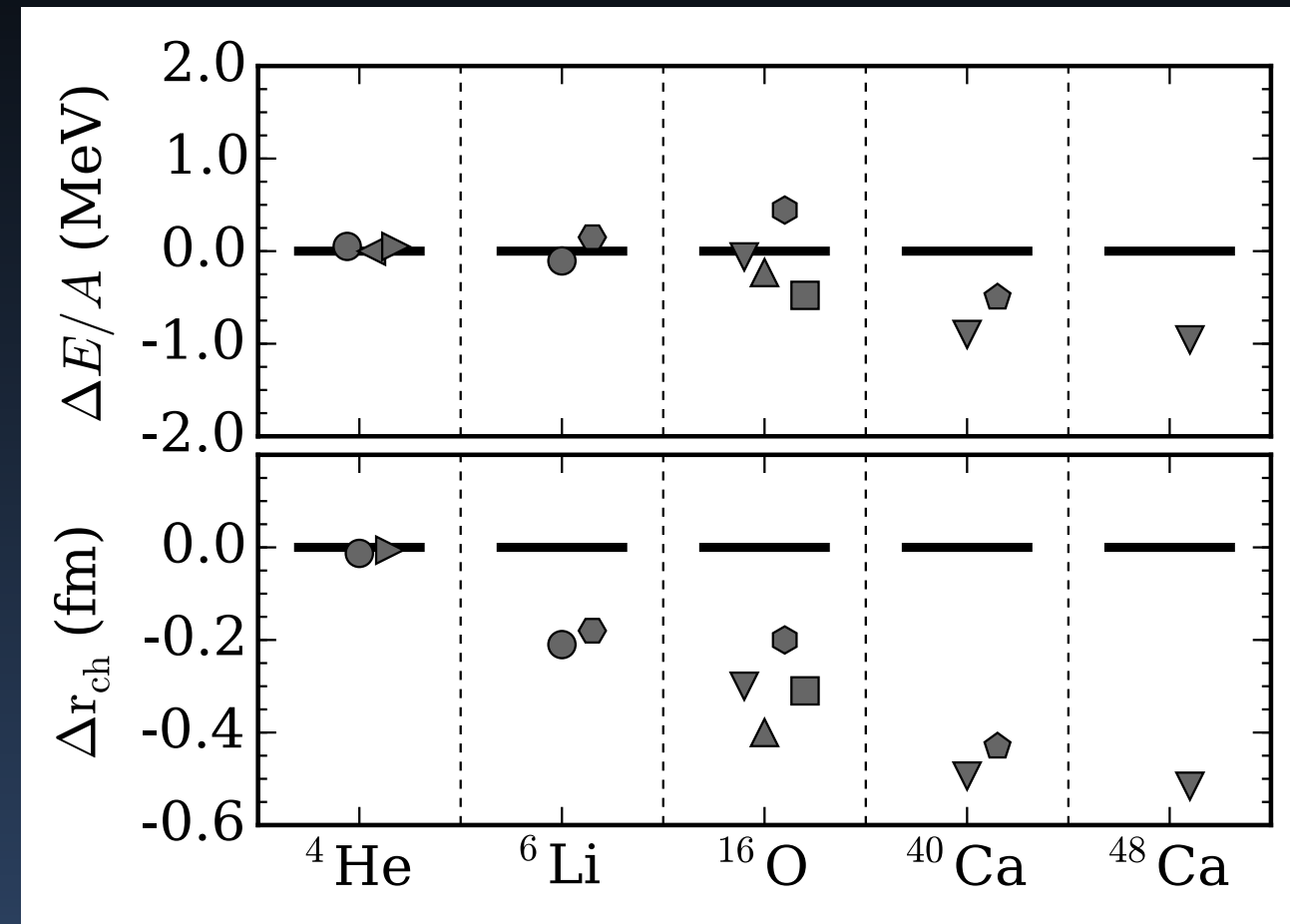
¹¹Technische Universität Darmstadt, Germany

¹²ExtreMe Matter Institute EMMI, Darmstadt, Germany

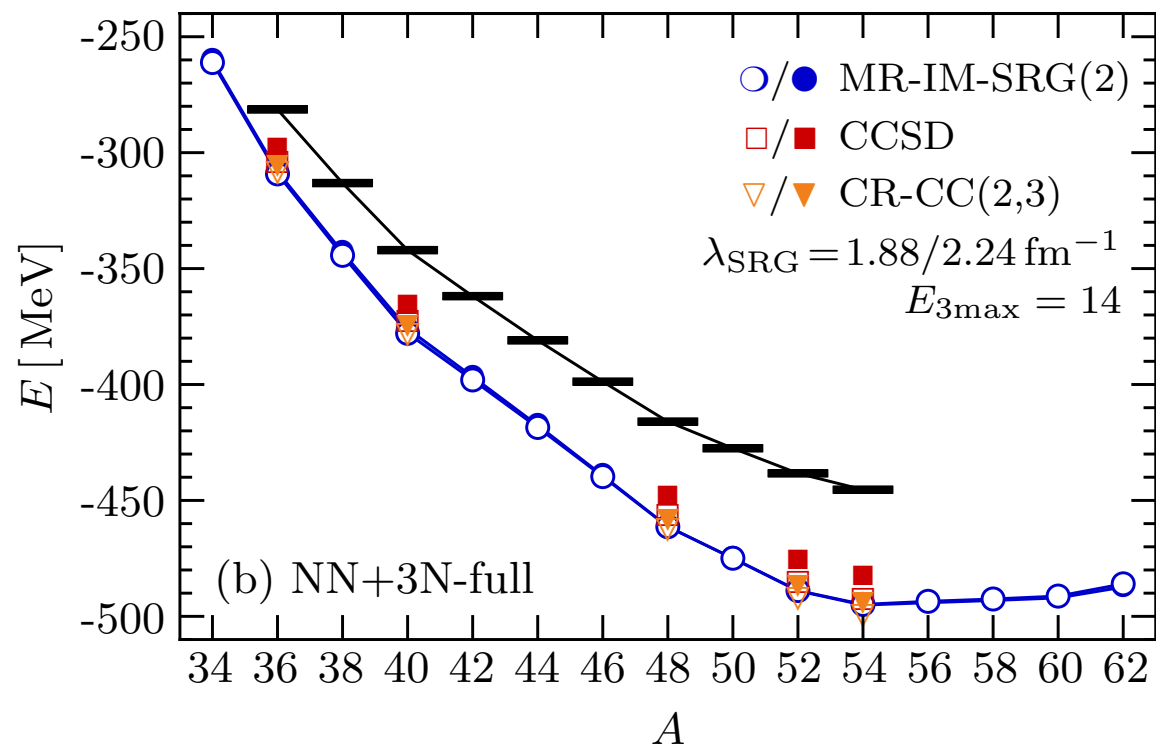
Problems in Many-body Systems

Effective Field Theory fit to
A=2, 3, 4 body systems.

Large extrapolations do
not go hand in hand
with predictive power!

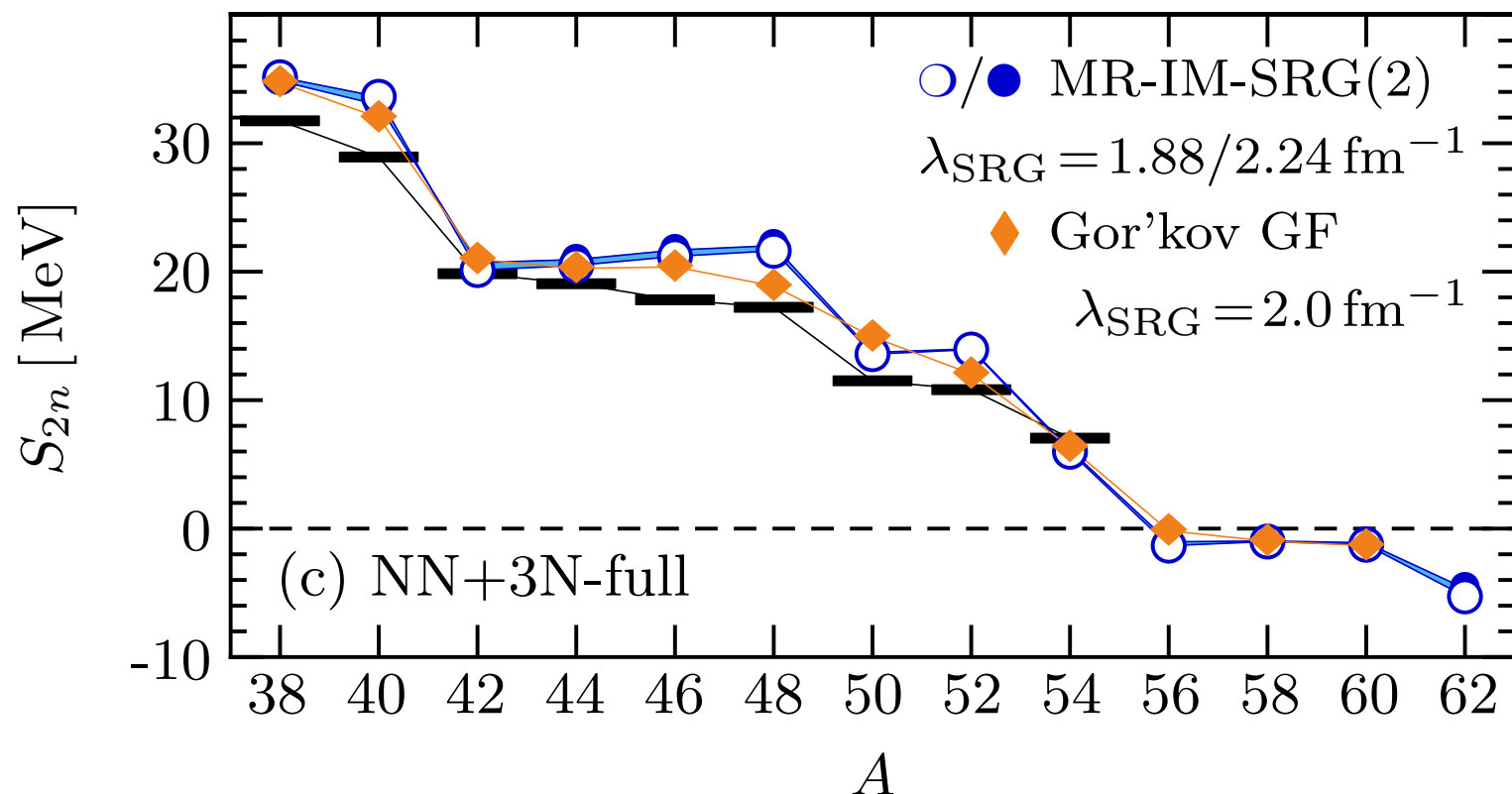


Catastrophic Errors?



EFT potentials have a catastrophic systematic error as the size of the nucleus increases

$$S_{2N}(N, Z) = E(N - 2, Z) - E(N, Z)$$

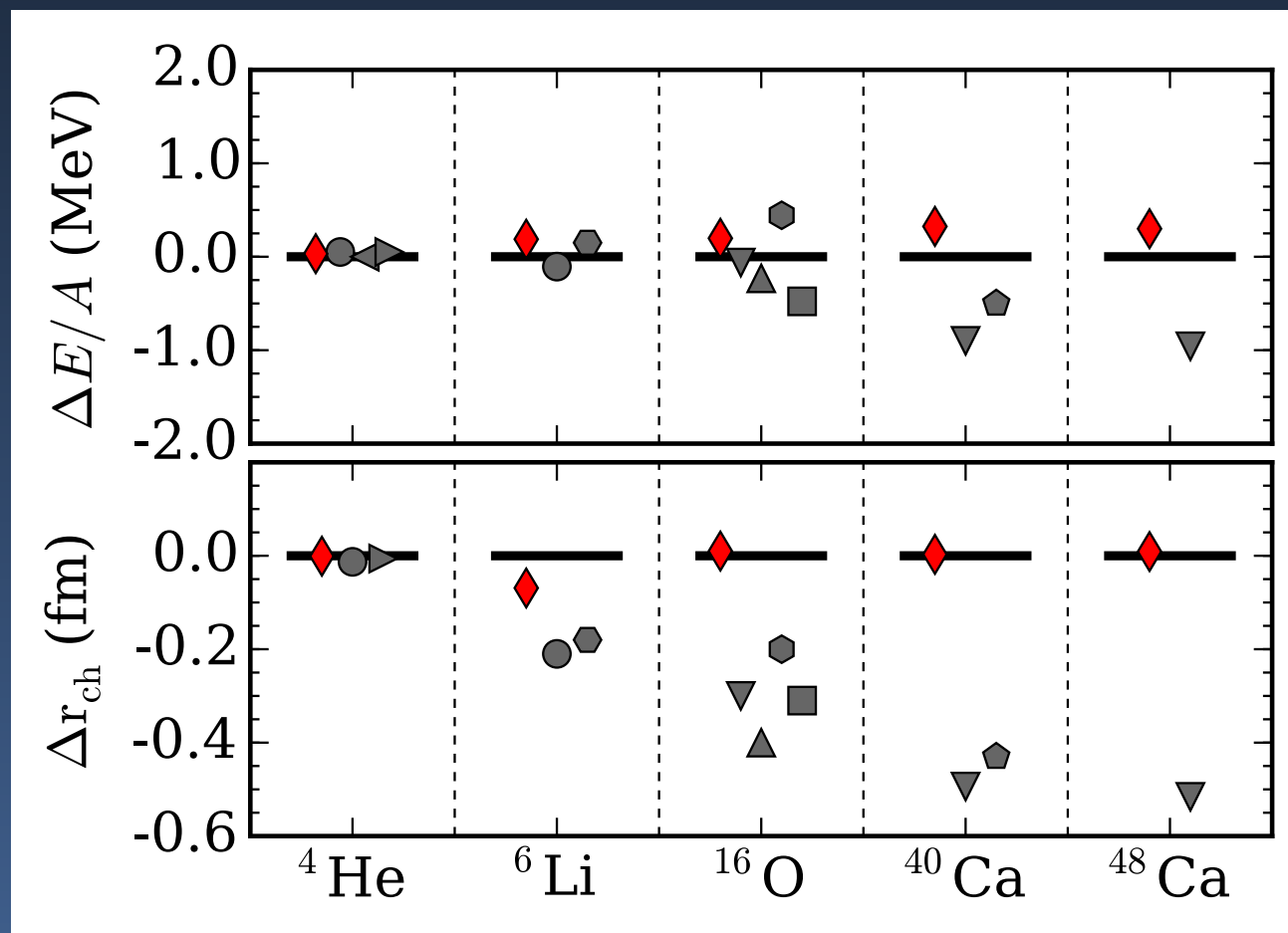


Differential Quantities seem unaffected by this systematic error!

How to Address This Issue?

Fit effective field theory to light nuclei and very low energy scattering data.

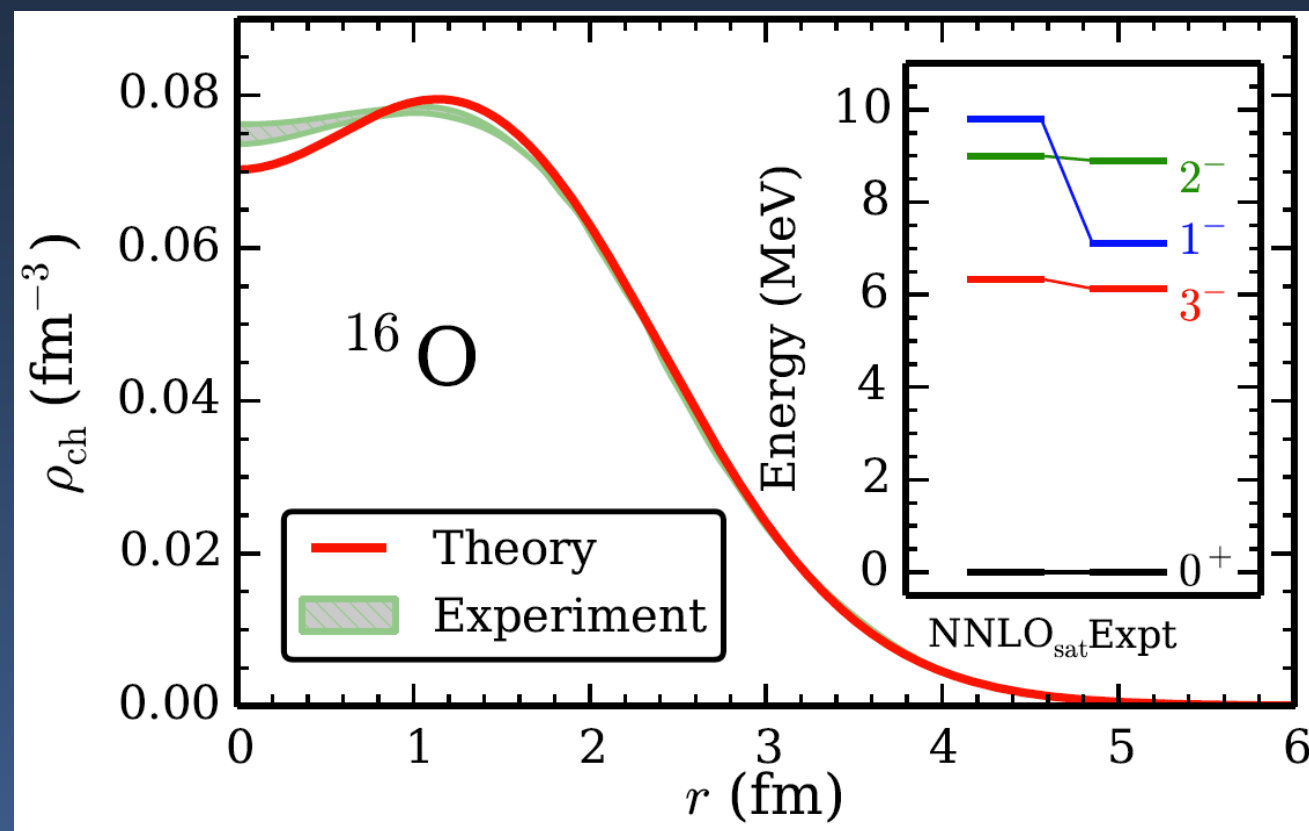
Ekström, Jansen, KAW, et. al., PRC(R) 91, 051301(2015)



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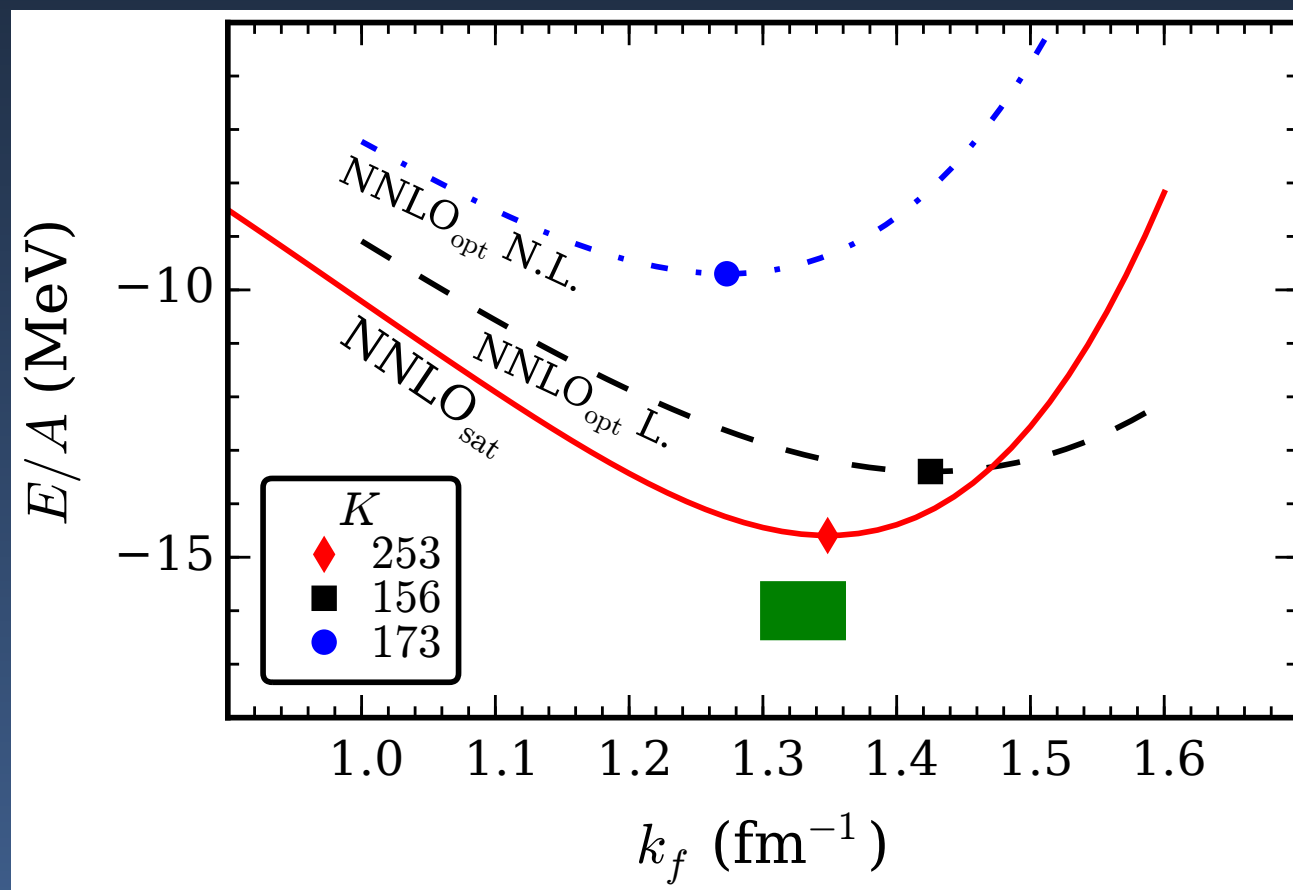
Ekström, Jansen, **KAW**, et. al., PRC(R) 91, 051301(2015)



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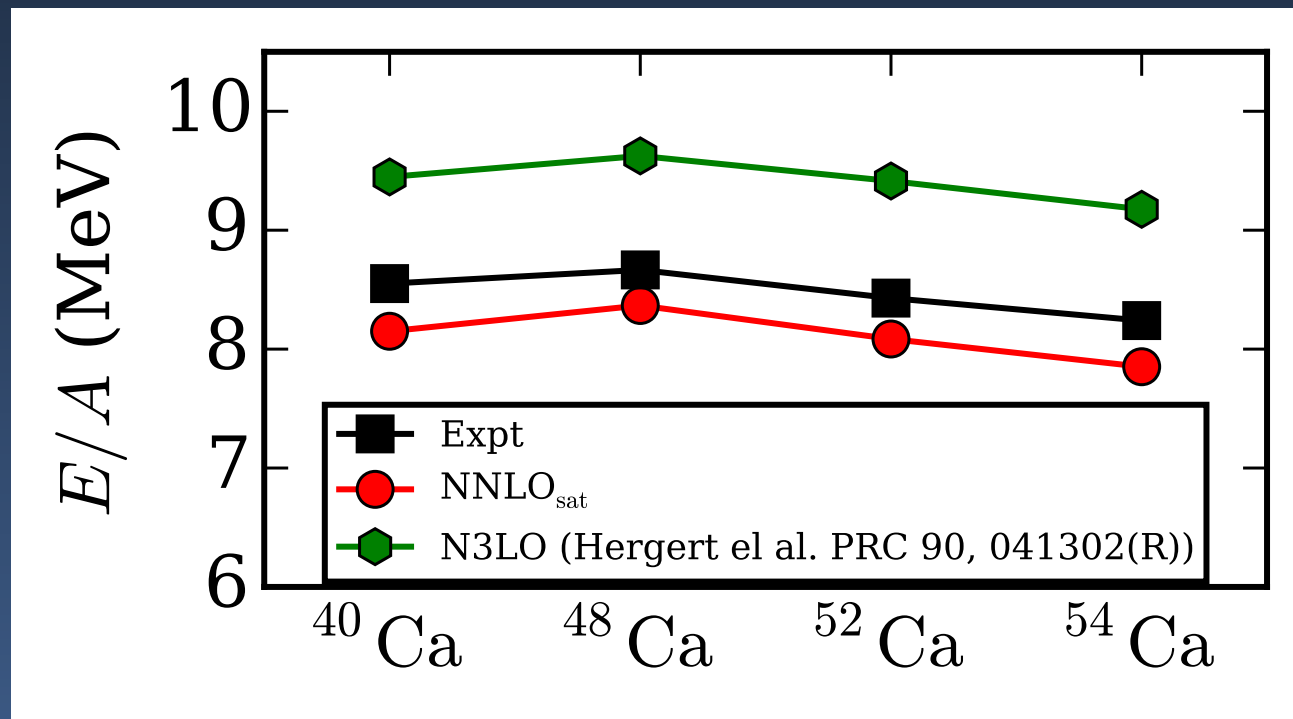
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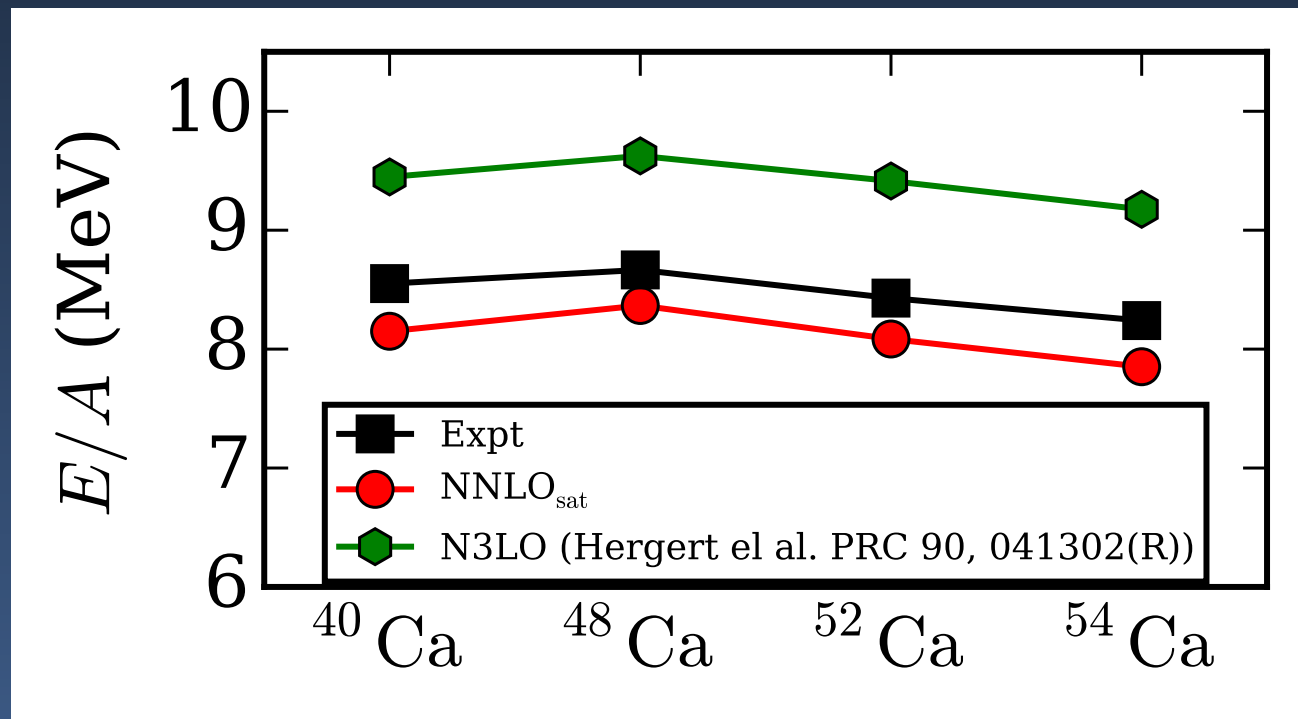
Ekström, Jansen, **KAW**, et. al., PRC(R) 91, 051301(2015)



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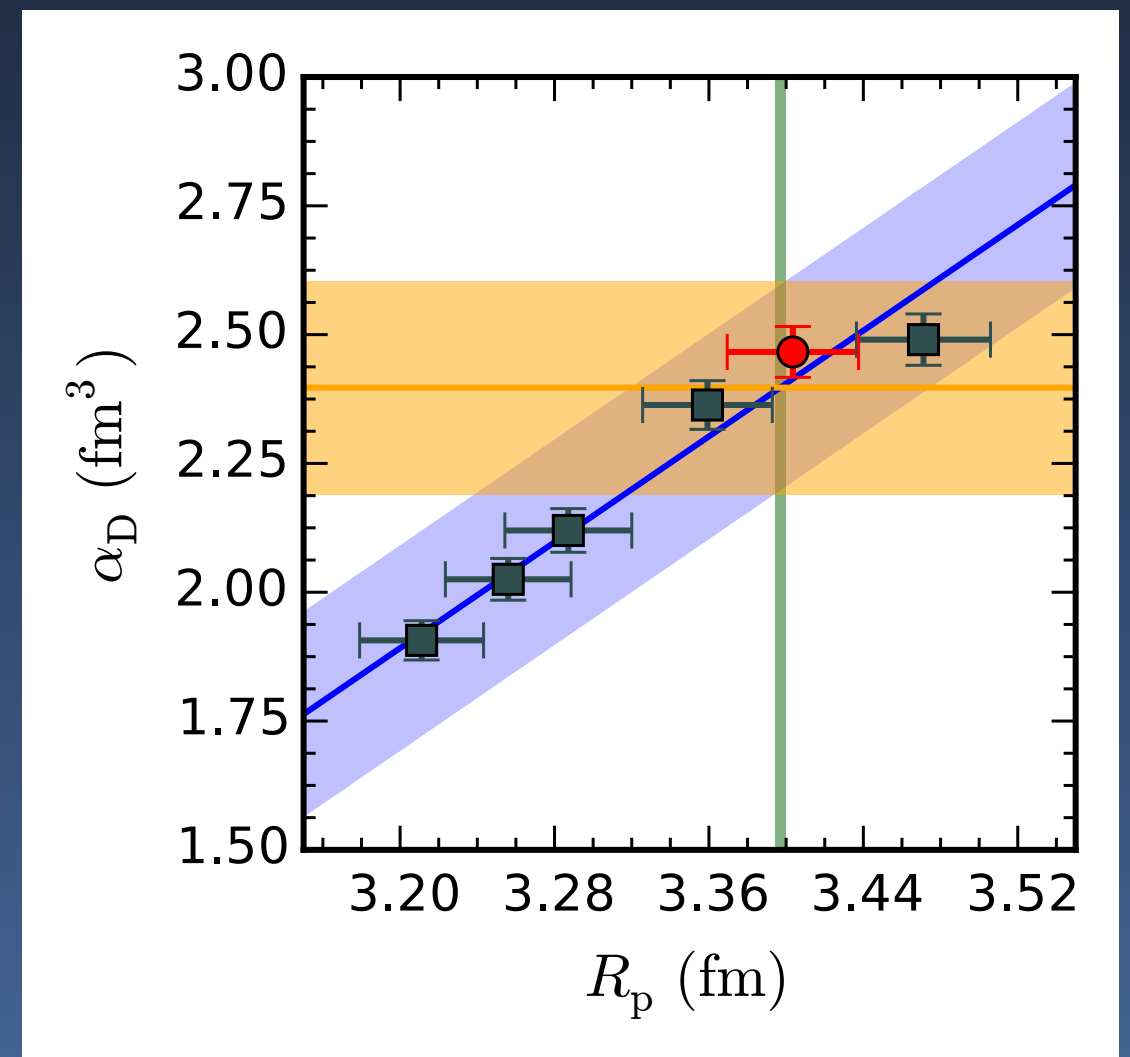
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Exploit correlations between predictions!

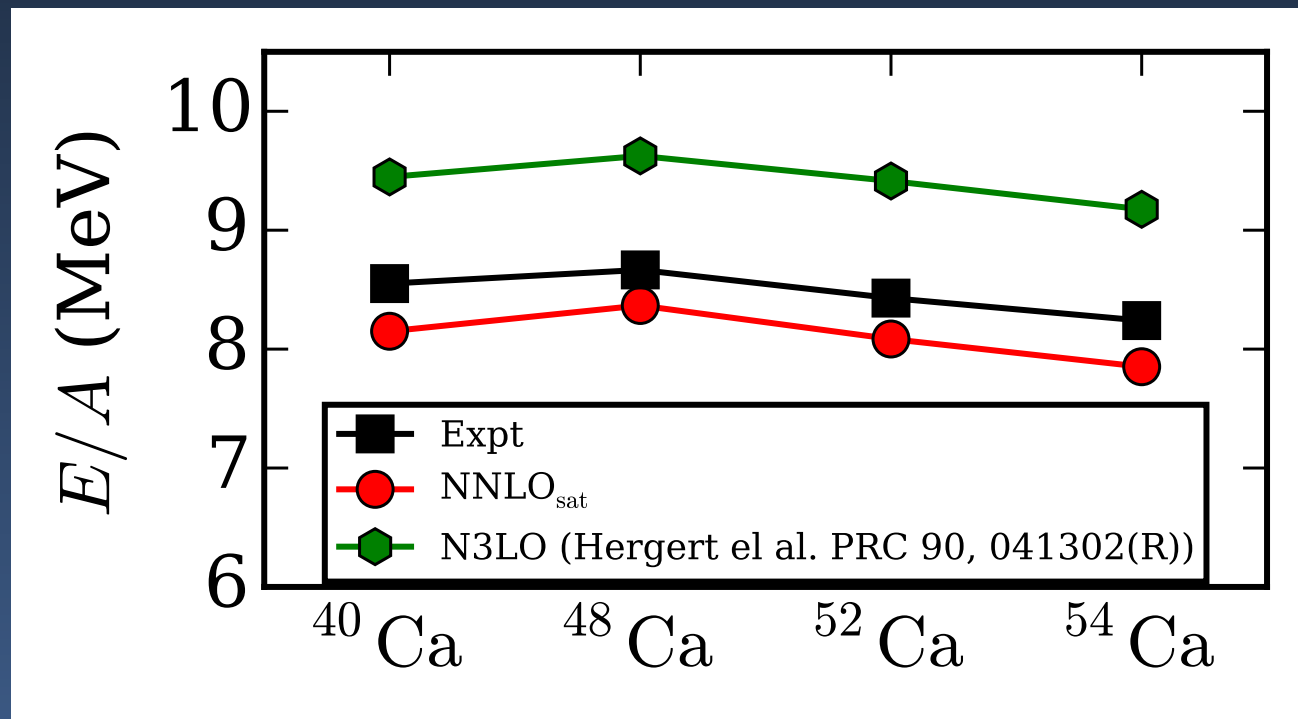
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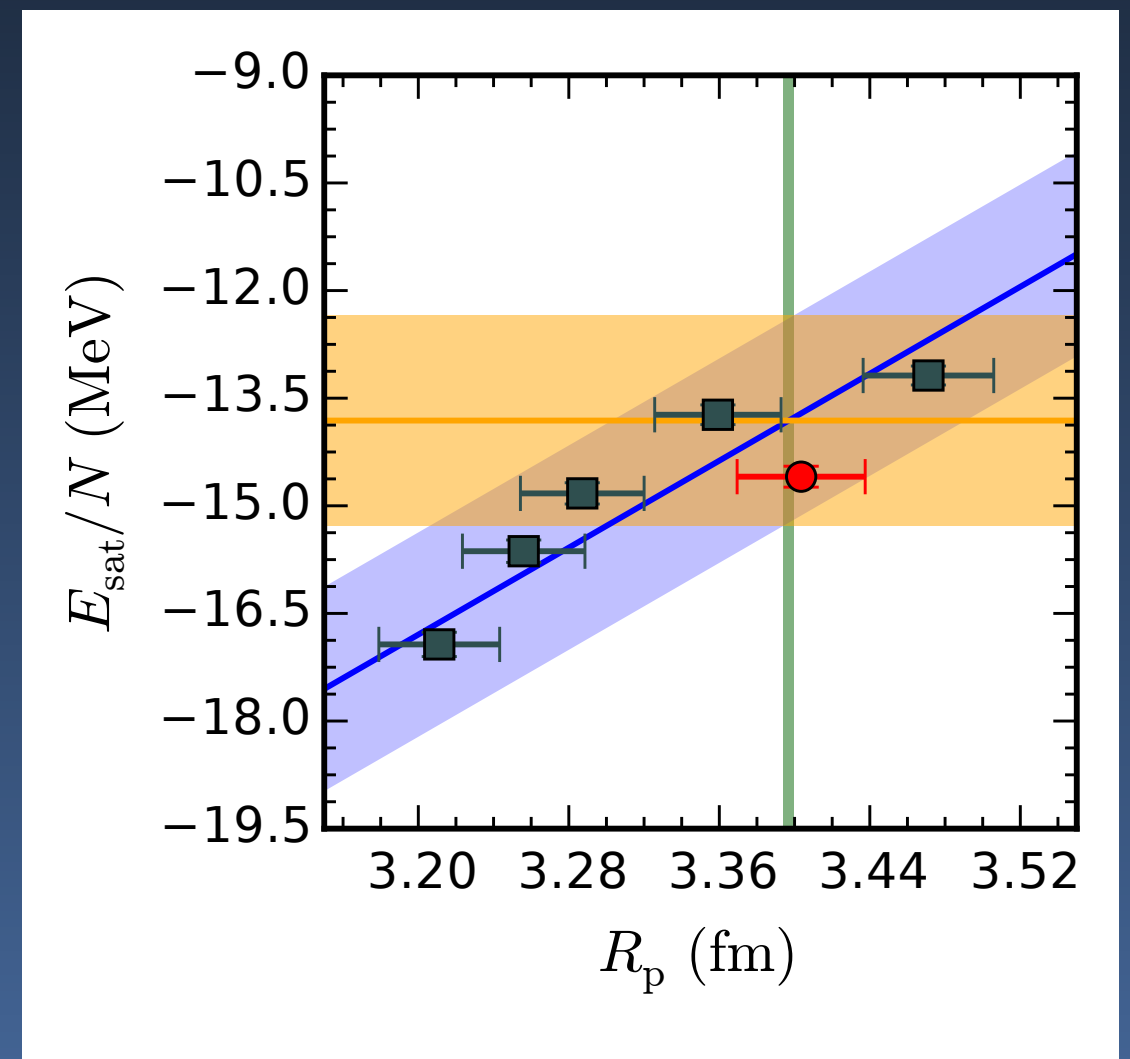
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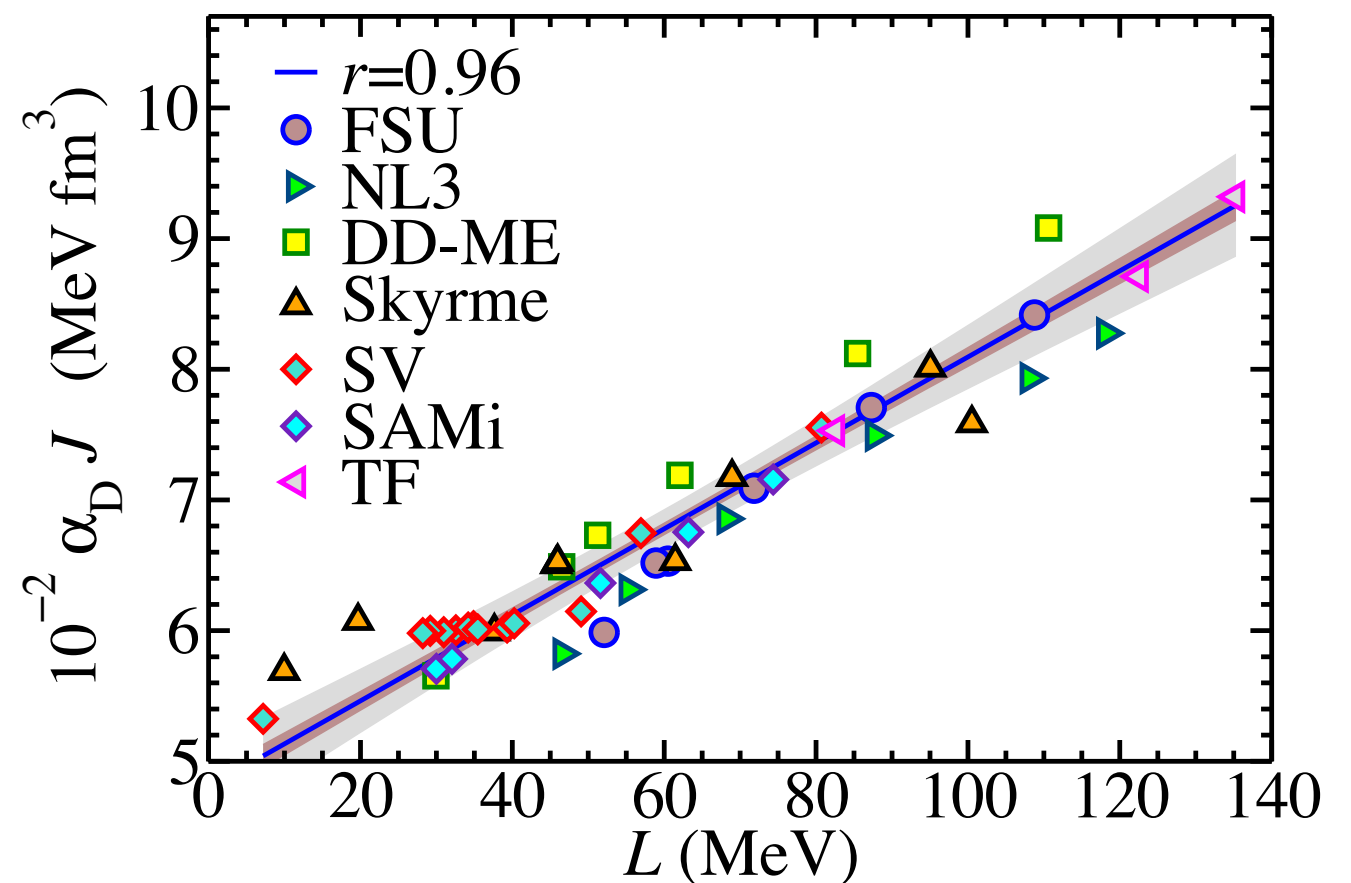
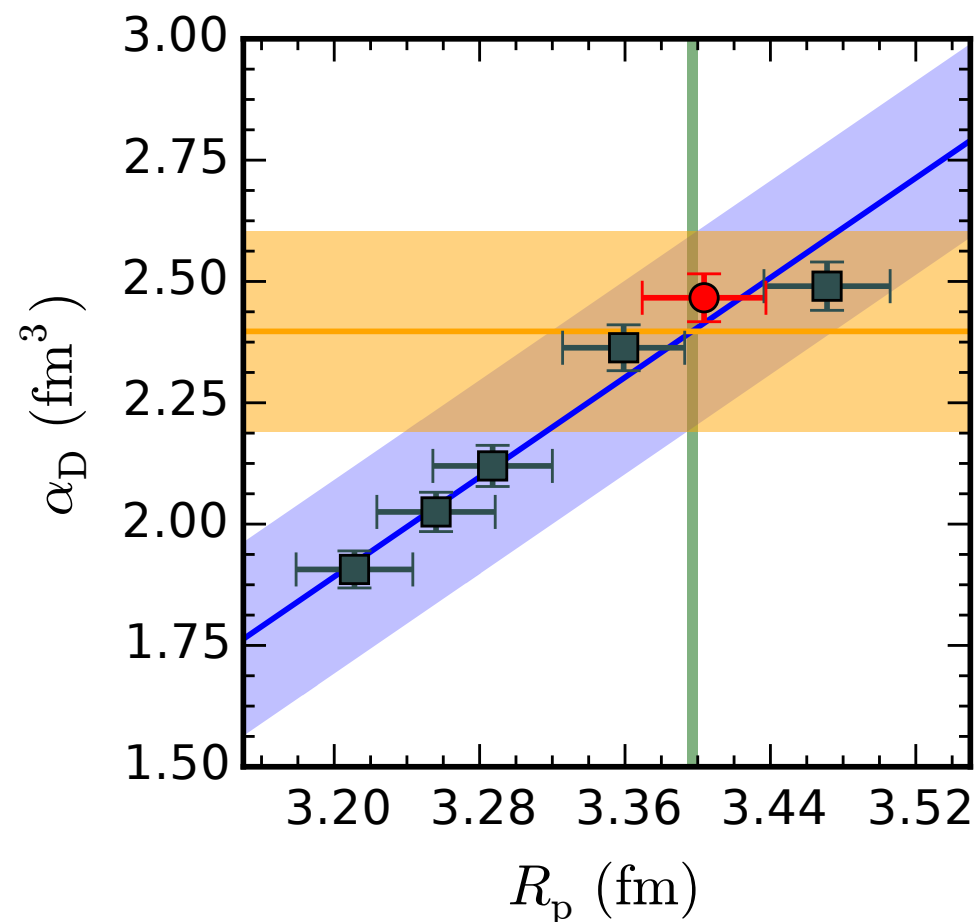
Exploiting Correlations Between Observables

Observables are correlated in heavy systems
(e.g. droplet model)

X. Roca-Maza, et. al., PRC 88, 024316 (2013)

What about lighter systems?

$$\alpha_D J = \frac{\pi e^2 A}{54} R^2 \left(1 + \frac{5}{3} \frac{L}{J} \epsilon_A \right)$$

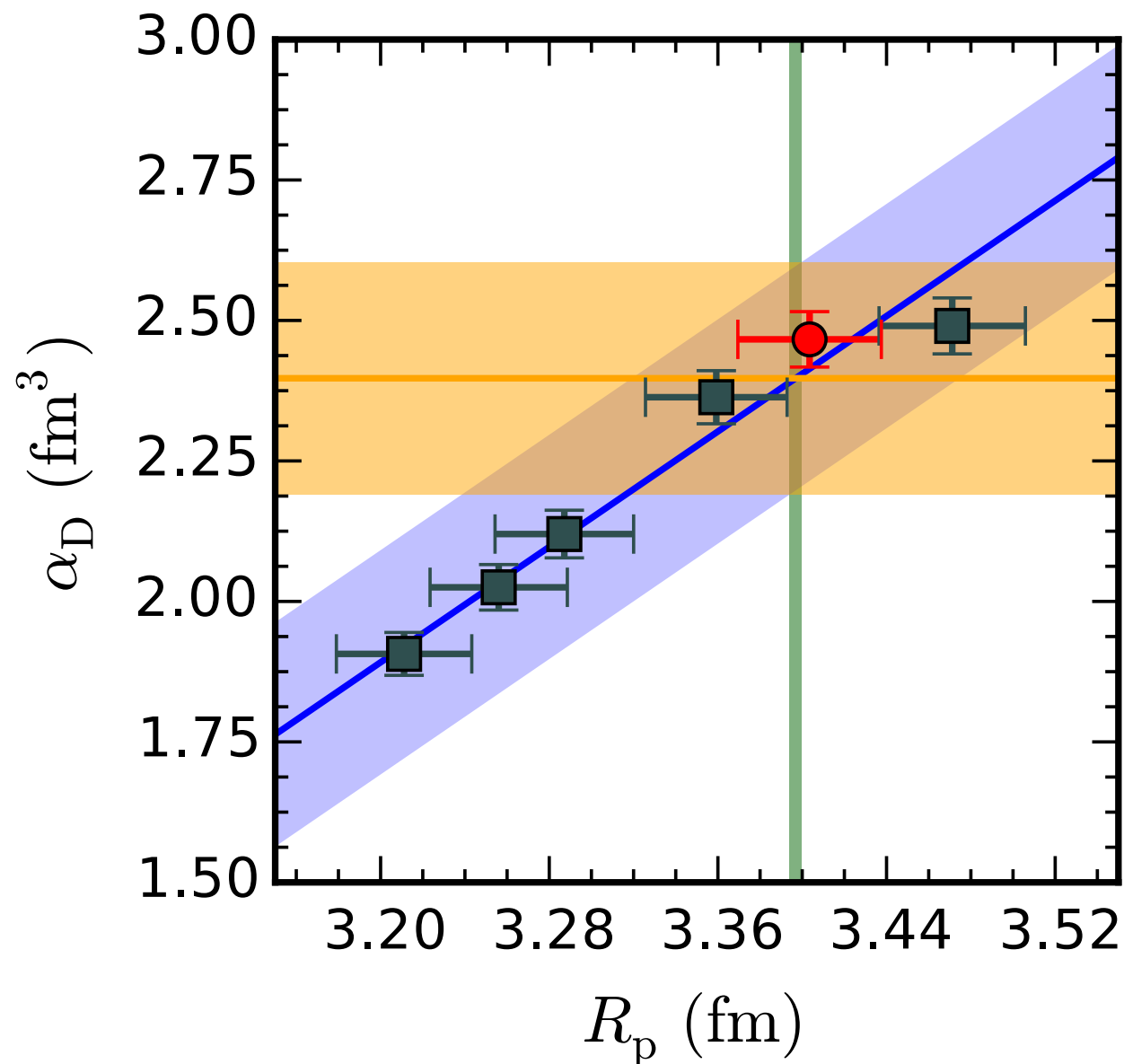


Exploiting Correlations Between Observables

Observables are correlated in heavy systems

(e.g. droplet model)

$$\alpha_D J \approx (aR^2 + b) \approx (a_{\alpha_D} R + b_{\alpha_D})(a_J R + b_J) + \dots$$

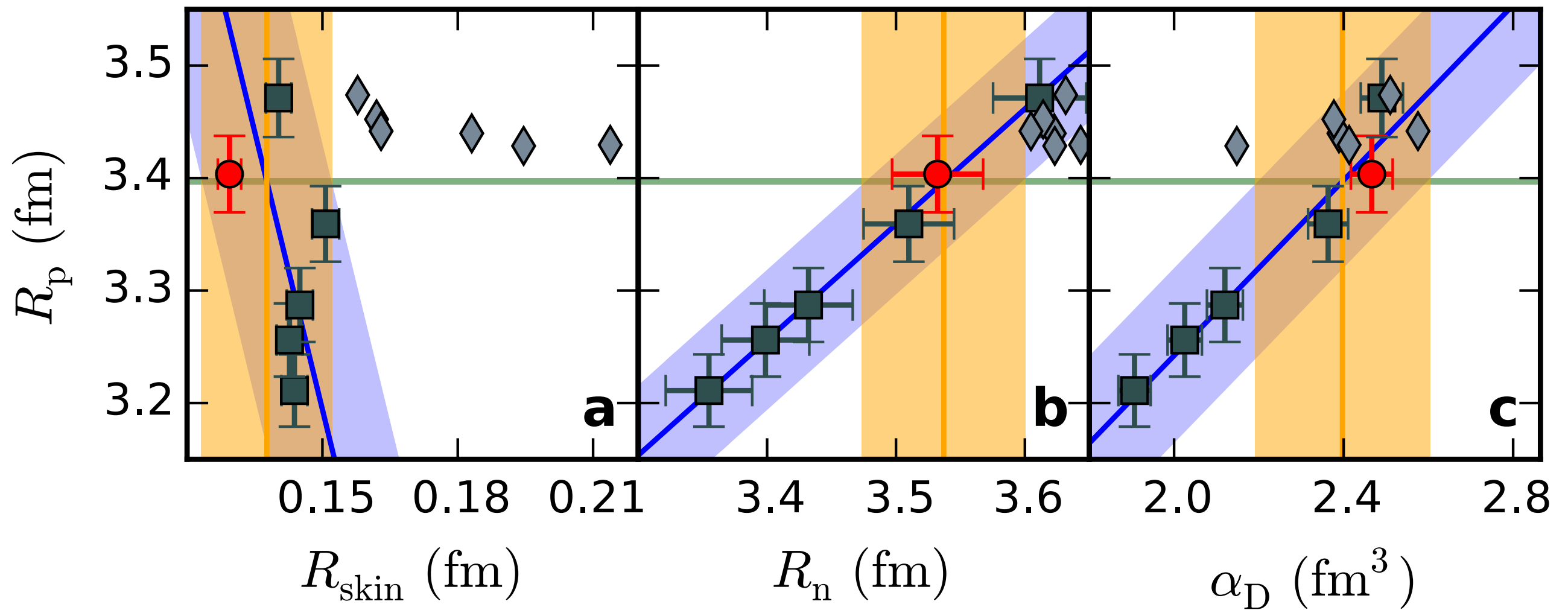


EFT in lighter systems has only a few parameters.

Many-body observables must be correlated.

Neutron Radius and Skin of ^{48}Ca

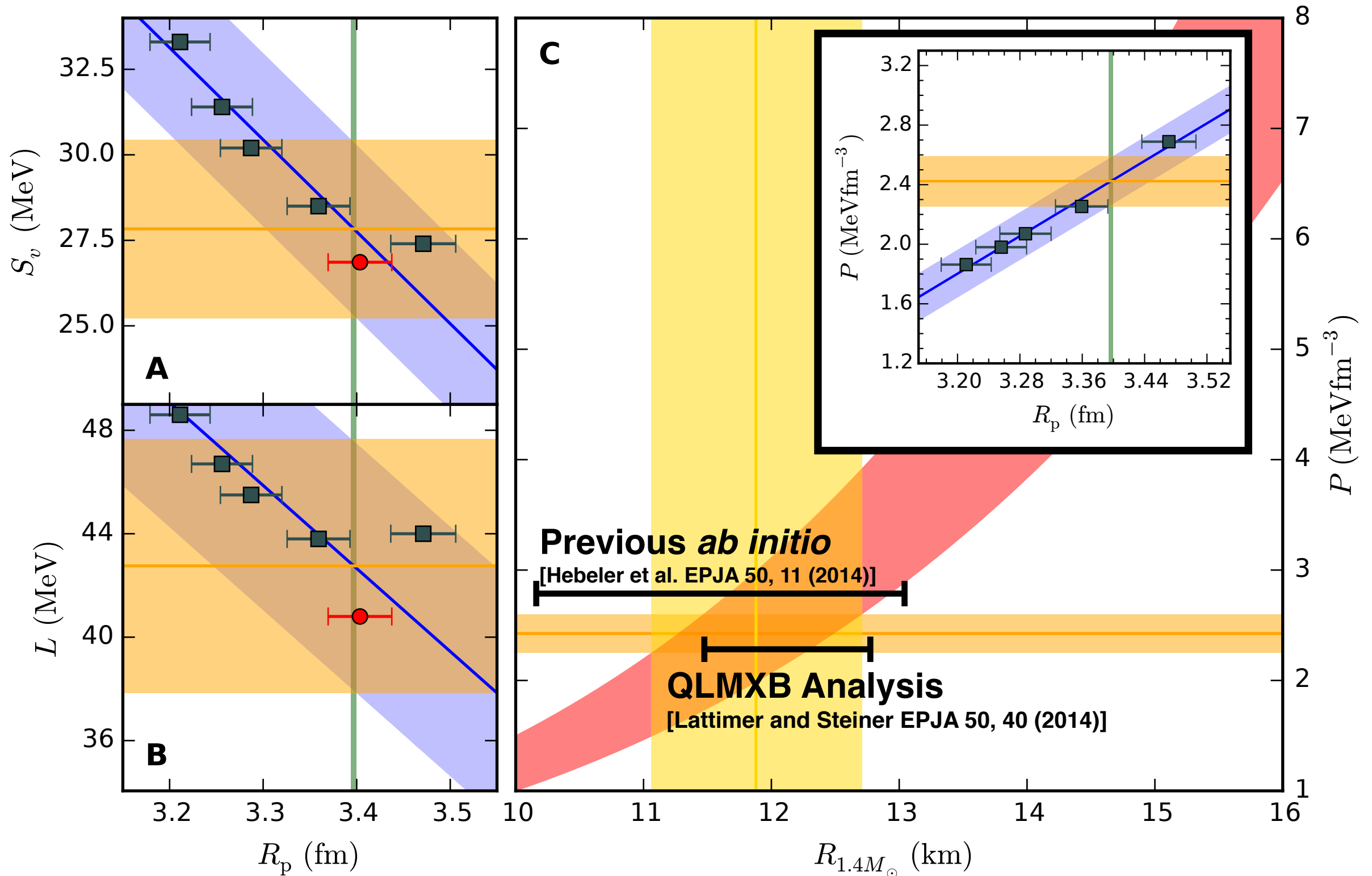
Hagen et al., Nature Physics 12, 186–190 (2016) [10.1038/nphys3529](https://doi.org/10.1038/nphys3529)



Ab initio gives a significantly thinner skin than DFT.
Ab initio and DFT give consistent prediction for α_D .

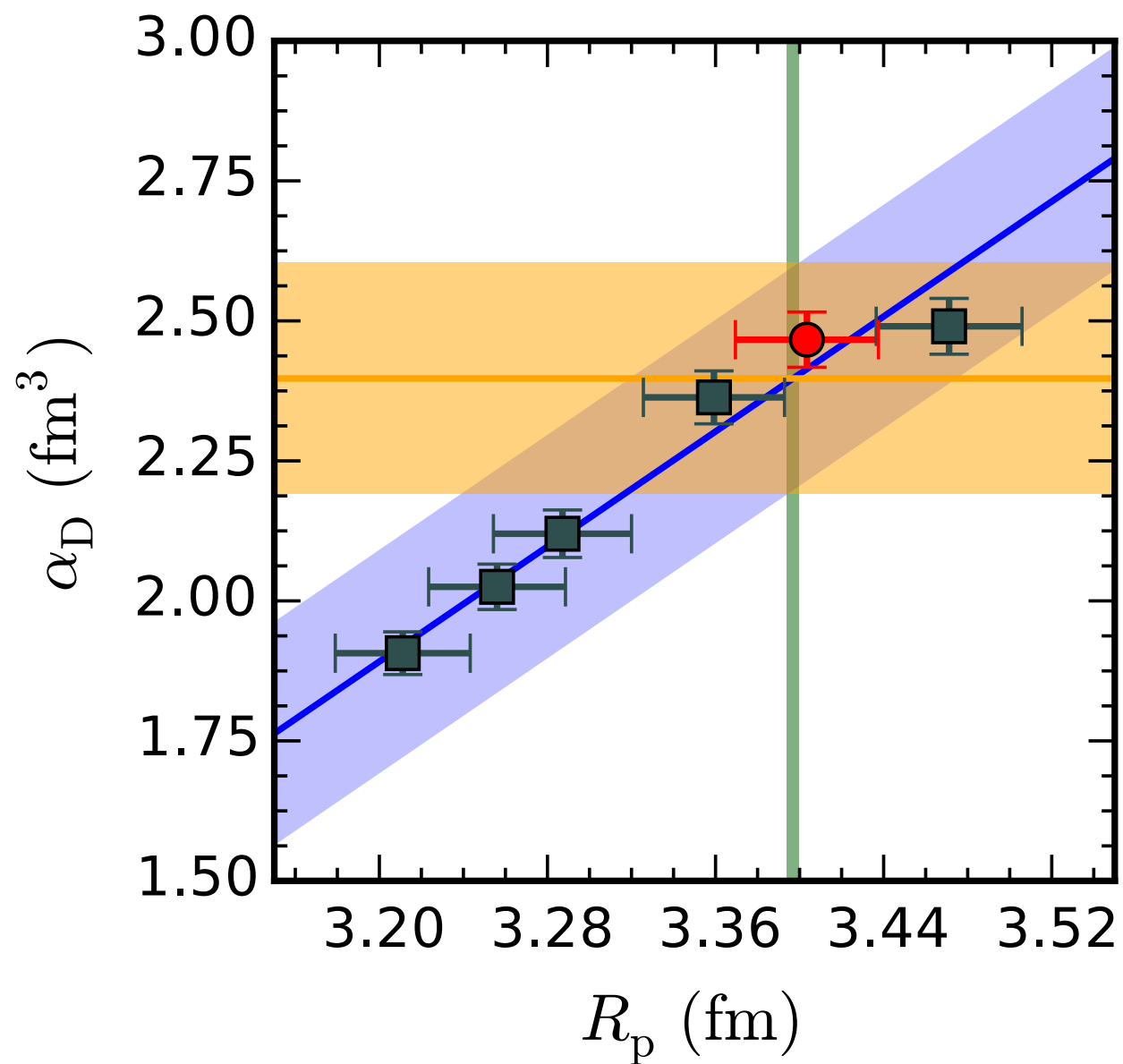
Implications for Neutron Stars

Hagen et al., Nature Physics 12, 186–190 (2016) [10.1038/nphys3529](https://doi.org/10.1038/nphys3529)



Exploiting Correlations Between Observables

Observables are correlated in nuclear systems



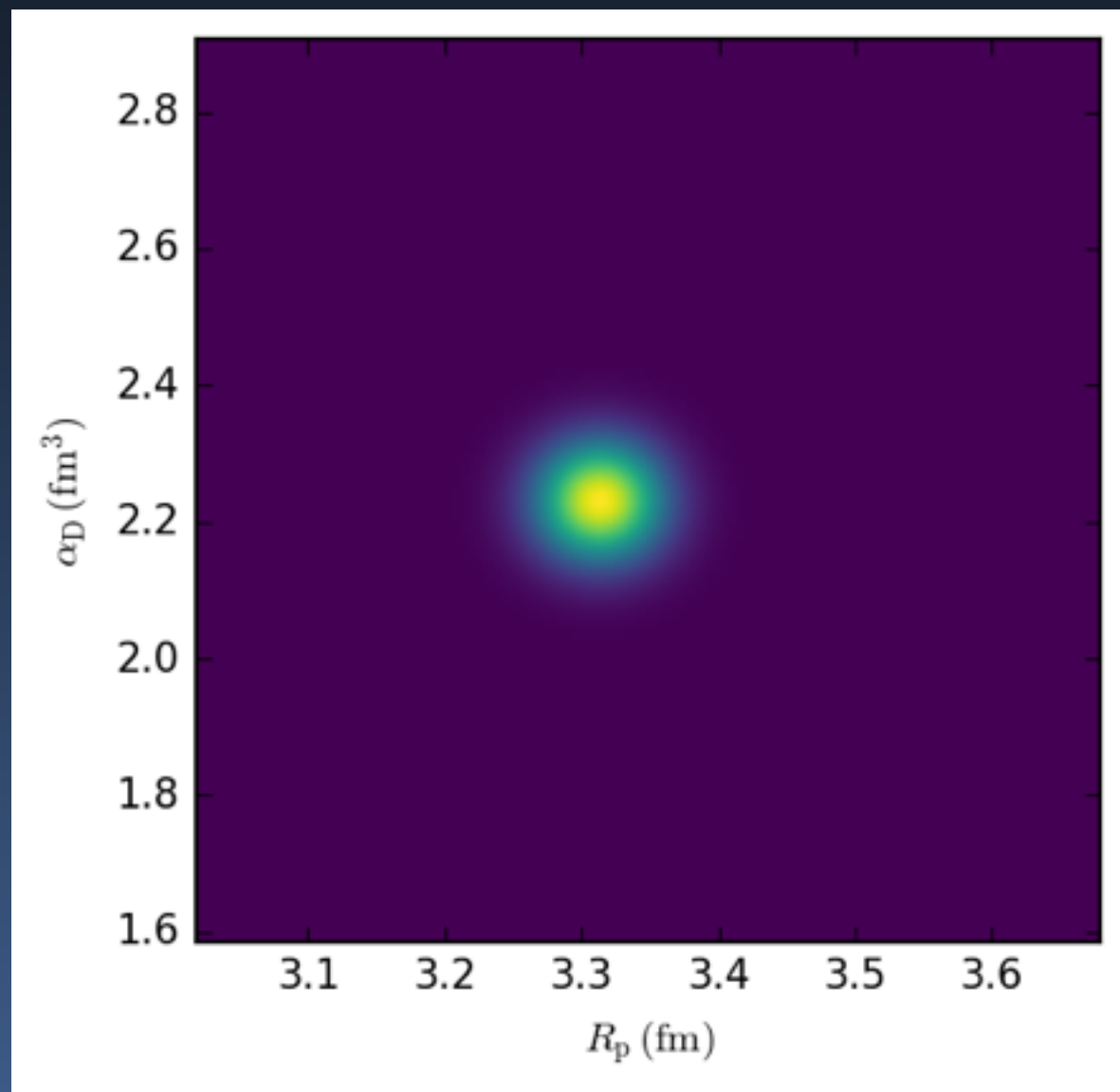
Explicit form of these correlations is unknown!

EFT in lighter systems has only a few parameters.

Many-body observables must be correlated.

But The Errors Are Correlated

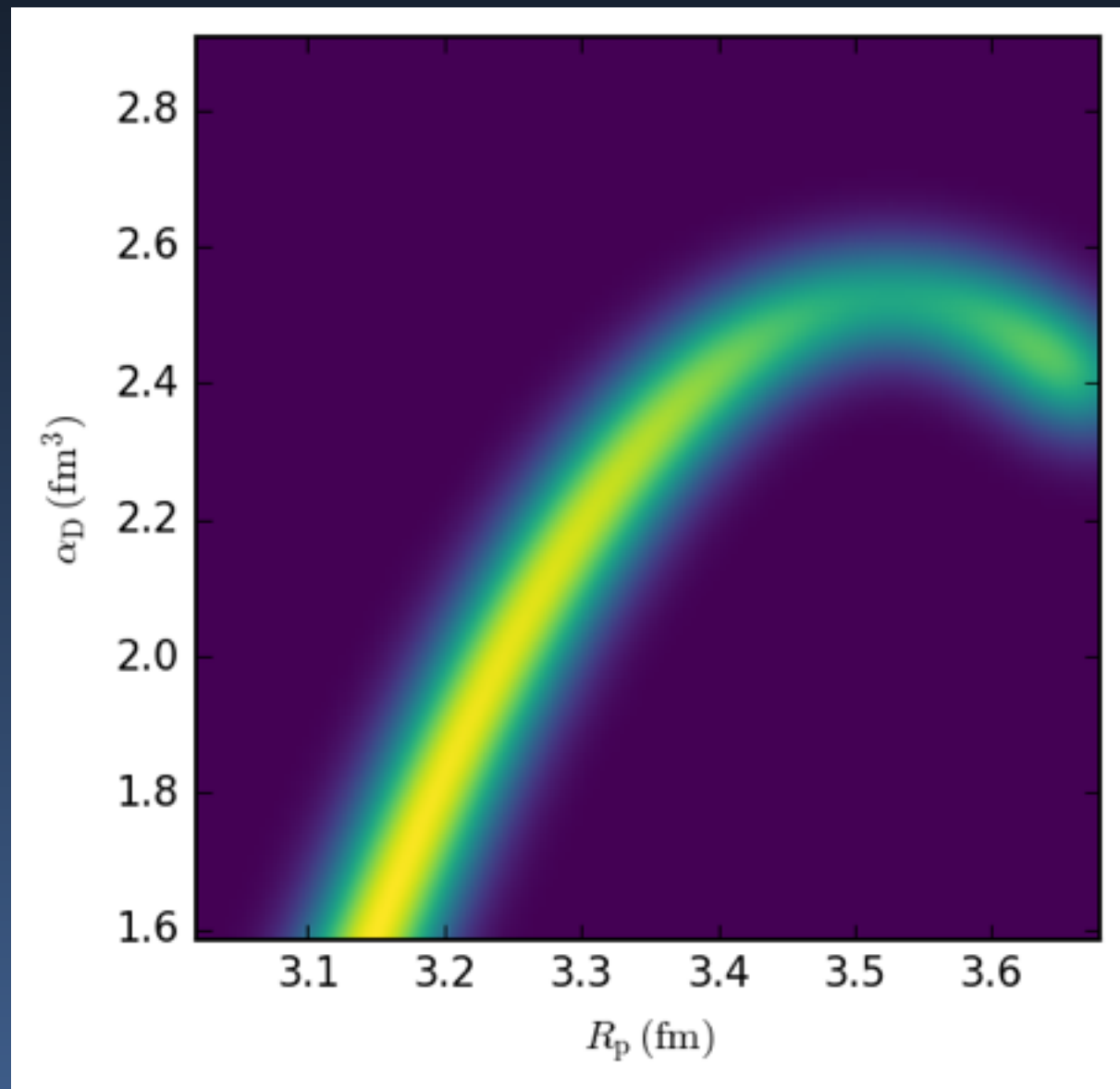
$$P_{\text{thy}}(O, X|m)$$



A single model (potential) gives us limited information about the joint probability distribution of two many-body observables

But The Errors Are Correlated

$$P_{\text{thy}}(O, X) = \int dm P(m) P_{\text{thy}}(O, X | m)$$



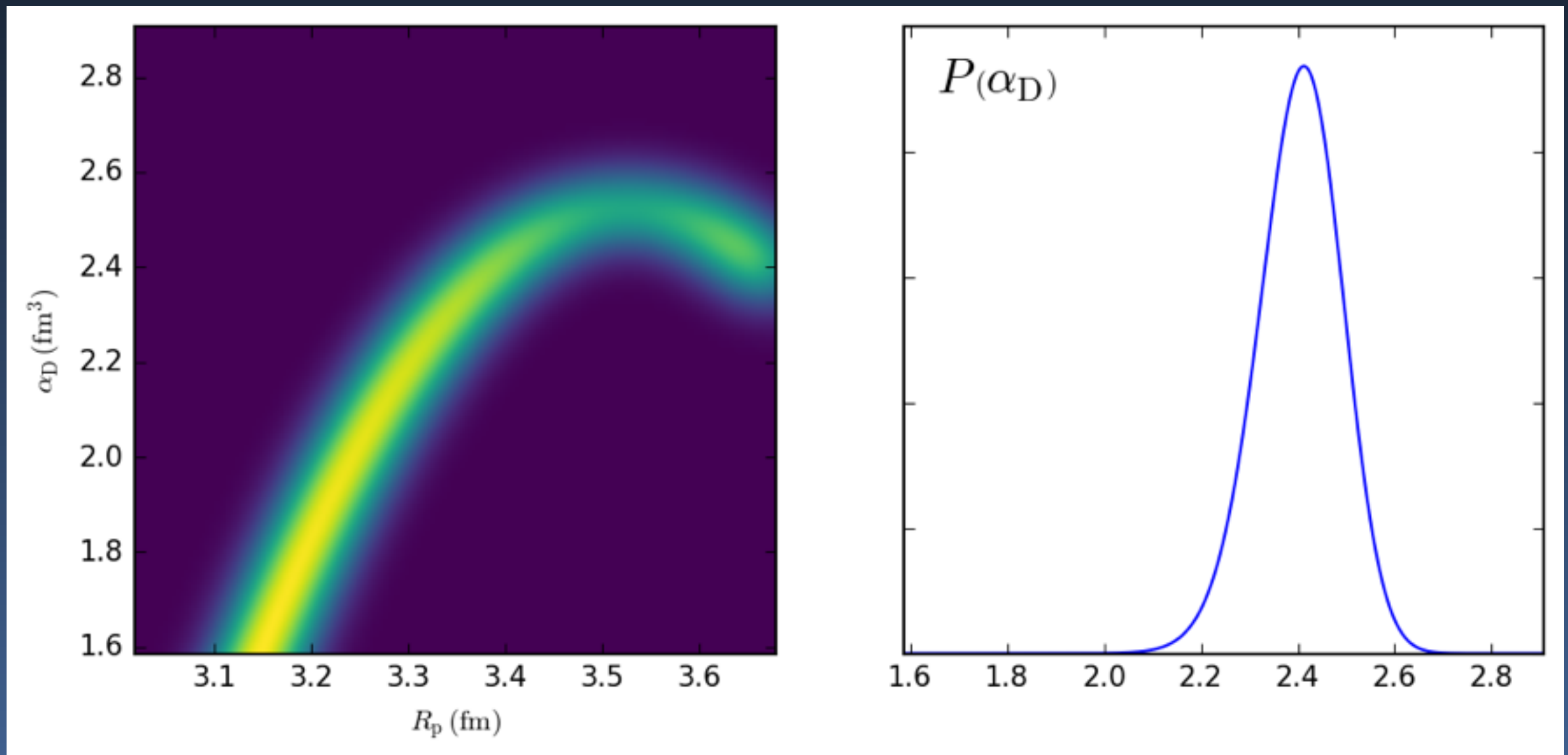
A single model (potential) gives us limited information about the joint probability distribution of two many-body observables

We can marginalize (integrate) over a family of models to better define our joint probability distribution.

But The Errors Are Correlated

We can marginalize (integrate) experimental data to further constrain our distribution.

$$P(O) = \int dX P_{\text{thy}}(O|X) P_{\text{expt}}(X)$$



Can we do this?

Can we compute:

$$P_{\text{thy}}(O, X|m)$$

$$P_{\text{thy}}(O, X) = \int dm P(m) P_{\text{thy}}(O, X|m)$$

$$P(O) = \int dX P_{\text{thy}}(O|X) P_{\text{expt}}(X)$$

Can we do this?

Can we compute:

$$P_{\text{thy}}(O, X|m) = \mathcal{N}(\{O_m, X_m\}, \Sigma_{O_m, X_m}) \quad \text{More or less}$$

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$$P_{\text{thy}}(O, X|m) = \mathcal{N}(\{O_m, X_m\}, \Sigma_{O_m, X_m}) \quad \text{More or less}$$

$$P_{\text{thy}}(O, X) = \frac{1}{n} \sum_{m=1}^n P_{\text{thy}}(O, X|m) \quad \text{Almost?}$$

$$P(O) = \int dX P_{\text{thy}}(O|X) P_{\text{expt}}(X)$$

Can we do this?

Can we compute:

$$P_{\text{thy}}(O, X|m) \approx \mathcal{N}(O_m, \sigma_{O_m})\mathcal{N}(X_m, \sigma_{X_m}) \text{ More or less}$$

$$P_{\text{thy}}(O, X) = \frac{1}{n} \sum_{m=1}^n P_{\text{thy}}(O, X|m) \quad \text{Almost?}$$

$$P(O) = \int dX P_{\text{thy}}(O|X) P_{\text{expt}}(X)$$

Can we do this?

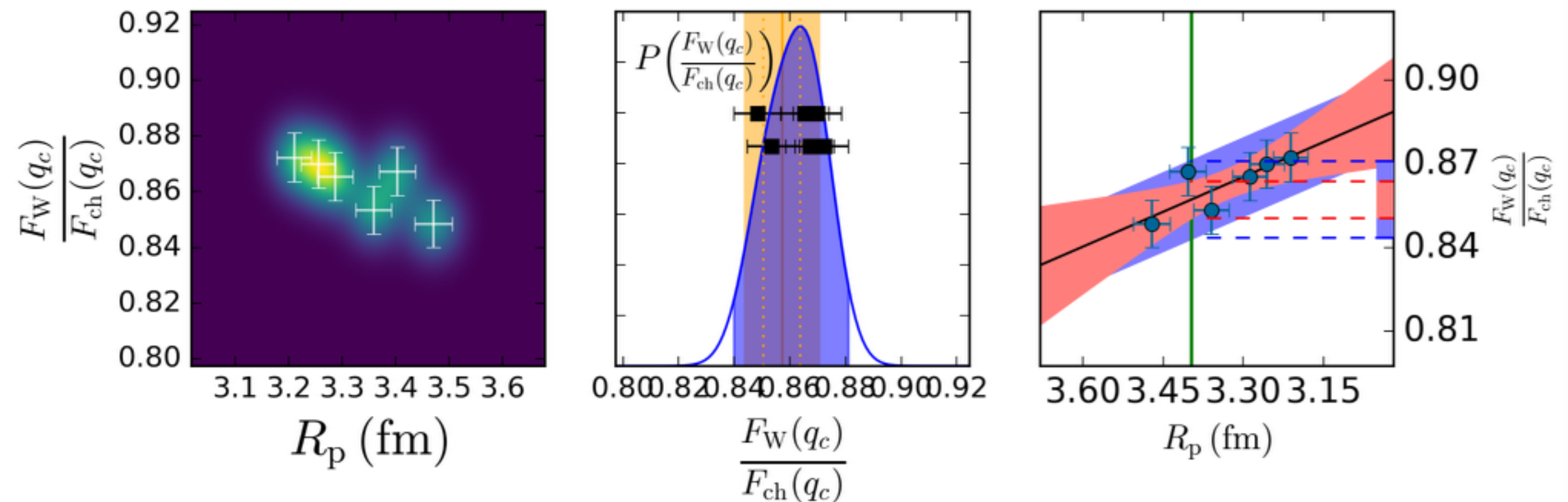
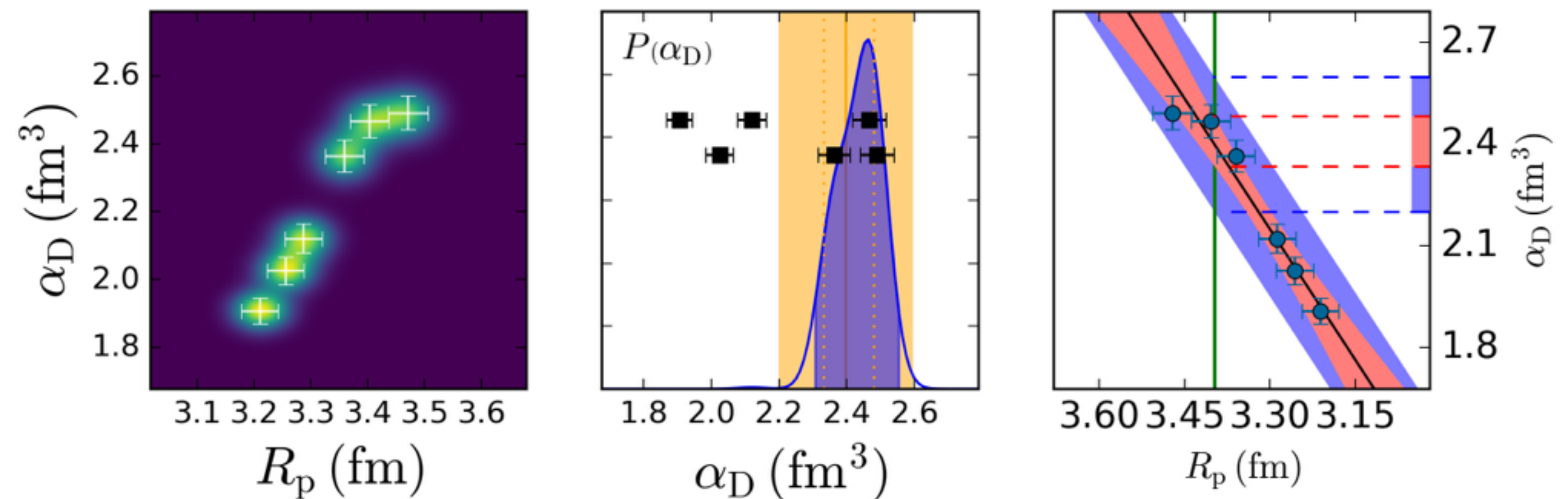
Can we compute:

$$P_{\text{thy}}(O, X|m) \approx \mathcal{N}(O_m, \sigma_{O_m})\mathcal{N}(X_m, \sigma_{X_m}) \text{ More or less}$$

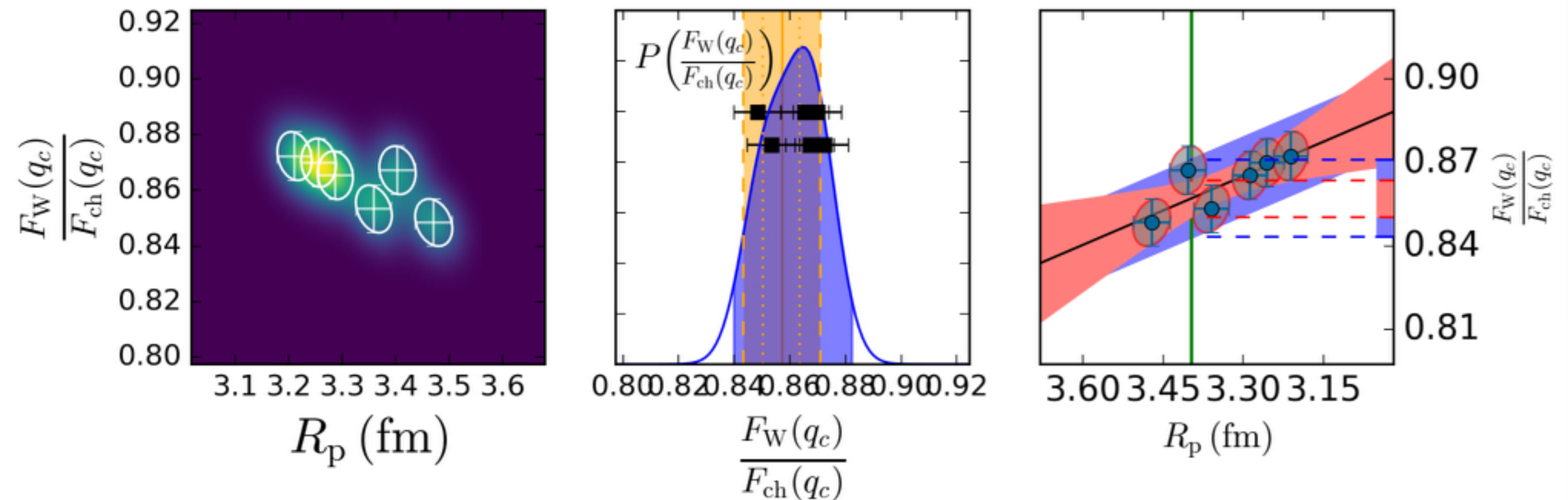
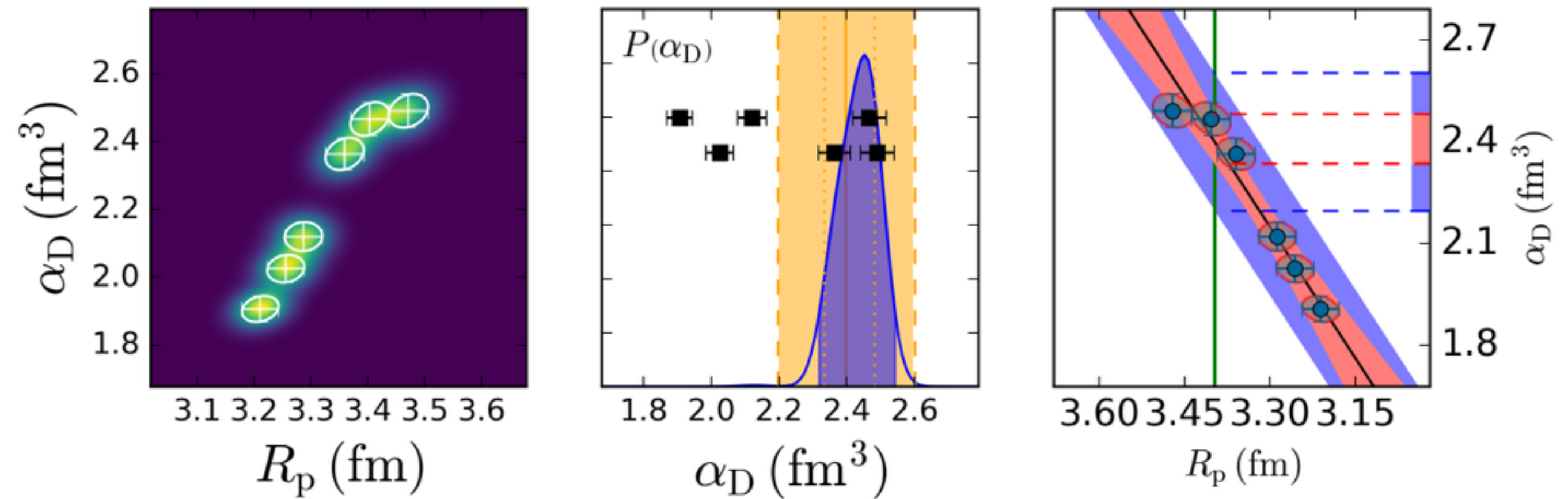
$$P_{\text{thy}}(O, X) = \frac{1}{n} \sum_{m=1}^n P_{\text{thy}}(O, X|m) \quad \text{Almost?}$$

$$P(O) = \int dX P_{\text{thy}}(O|X) P_{\text{expt}}(X) \quad \text{Yes!}$$

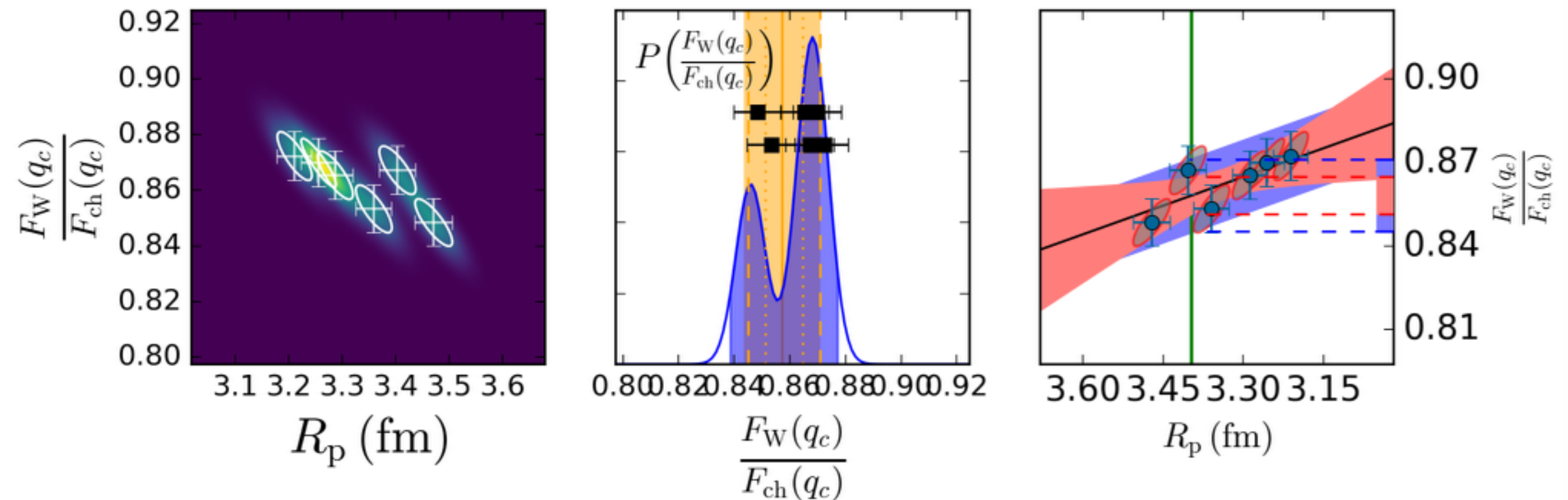
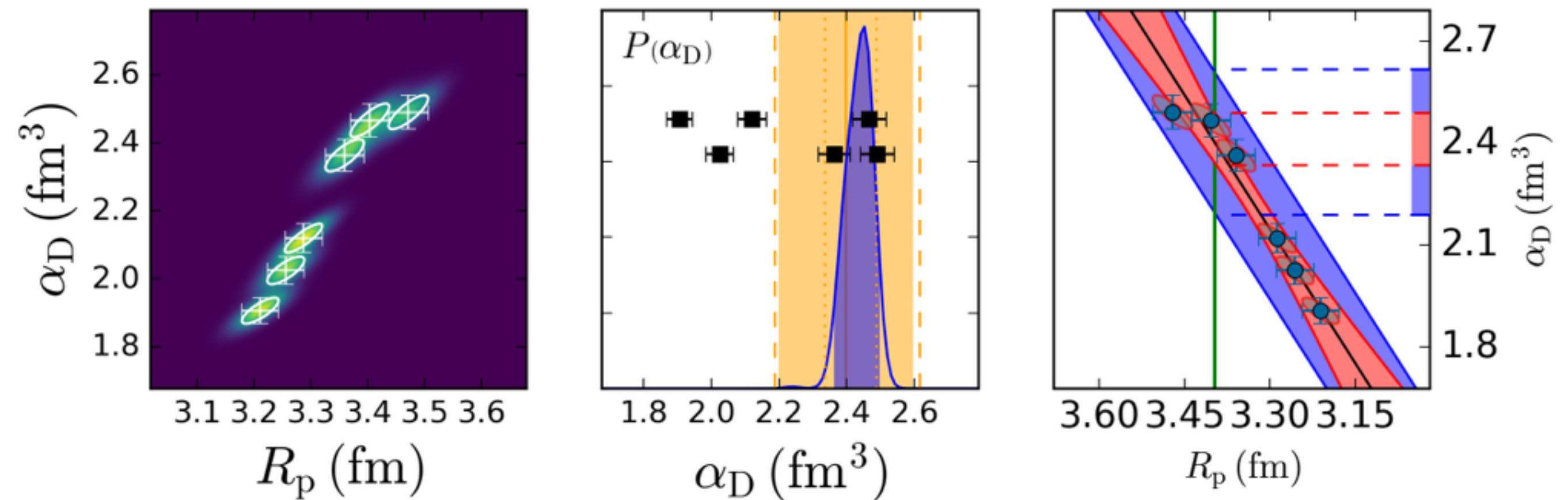
Dipole Polarizability and Weak Form Factor



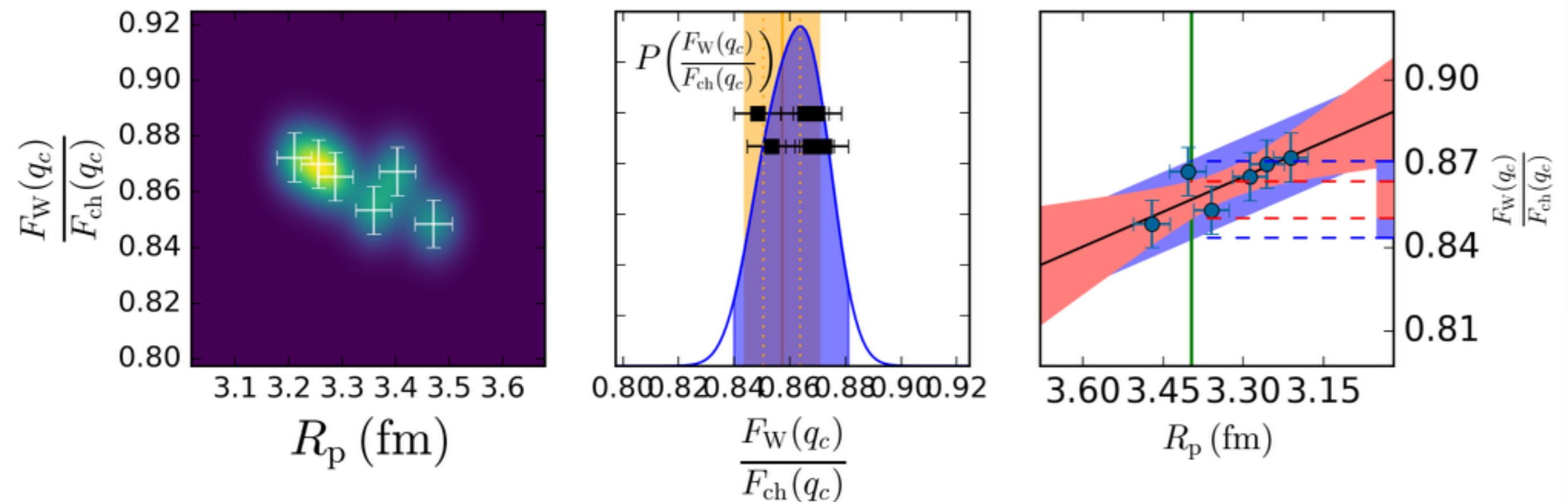
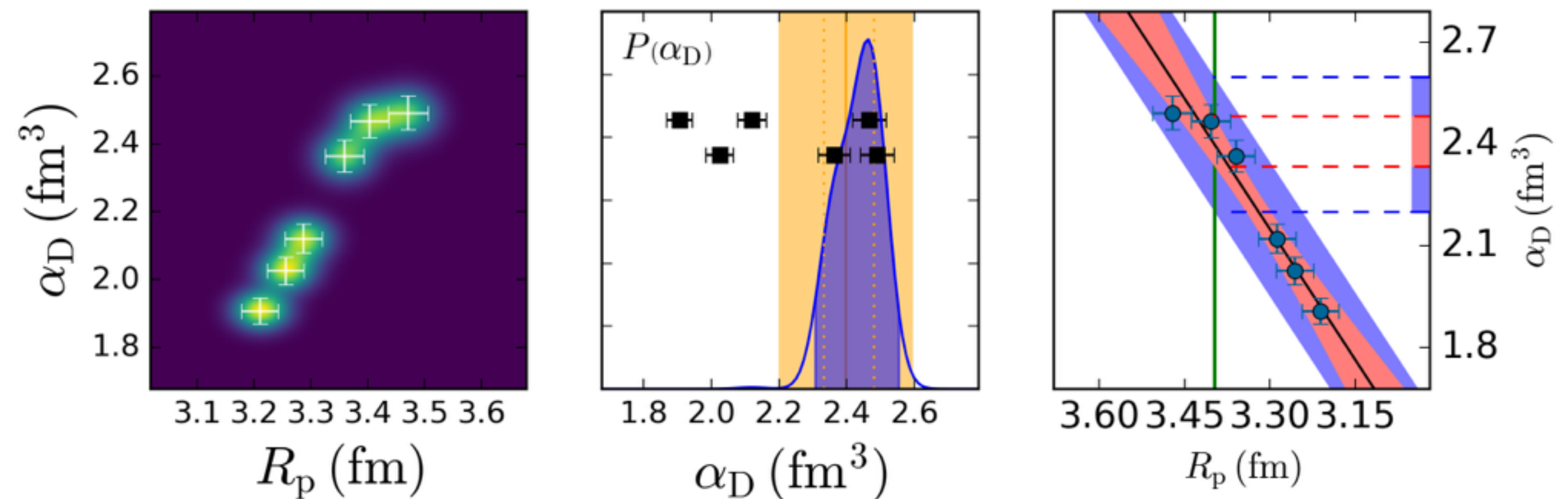
With Small Random Correlation Coefficient



With Large Random Correlation Coefficient



Dipole Polarizability and Weak Form Factor



Summary

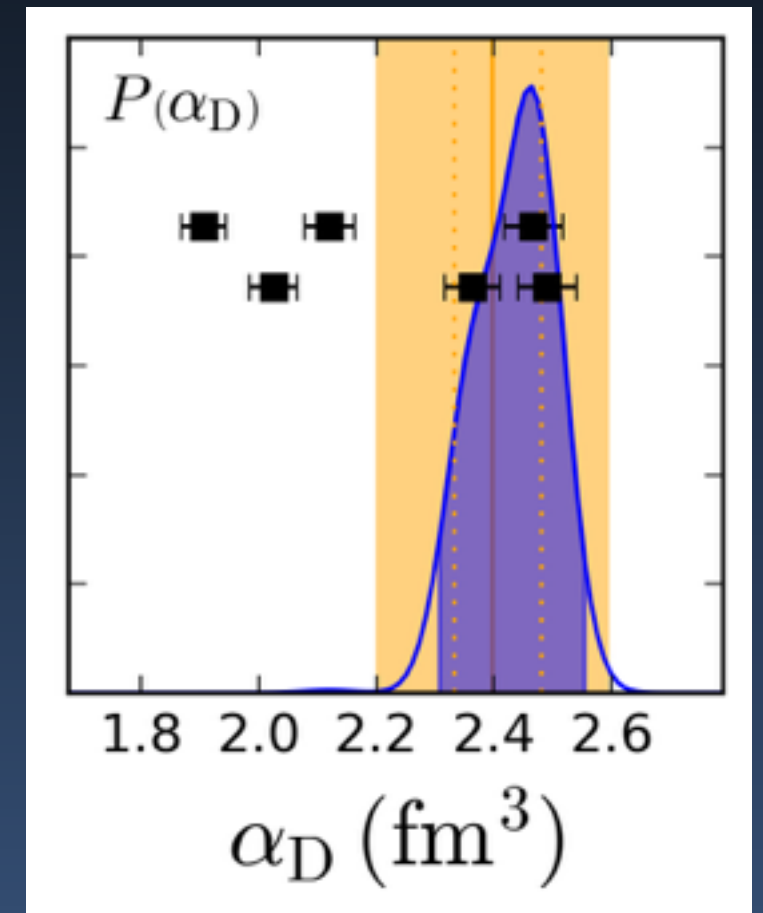
Chiral Effective Field Theory provides a tool for generating microscopic potentials.

Large systematic errors in many-body calculation tend to spoil this tool (EFT/model-independence)

Correlations between observables can be exploited to fix these issues

Need to exploit these correlations in as model independent manner as possible

Need to be able to fully propagate uncertainties from few-body input to many-body calculations (derivatives w.r.t model parameters)



Thank You