

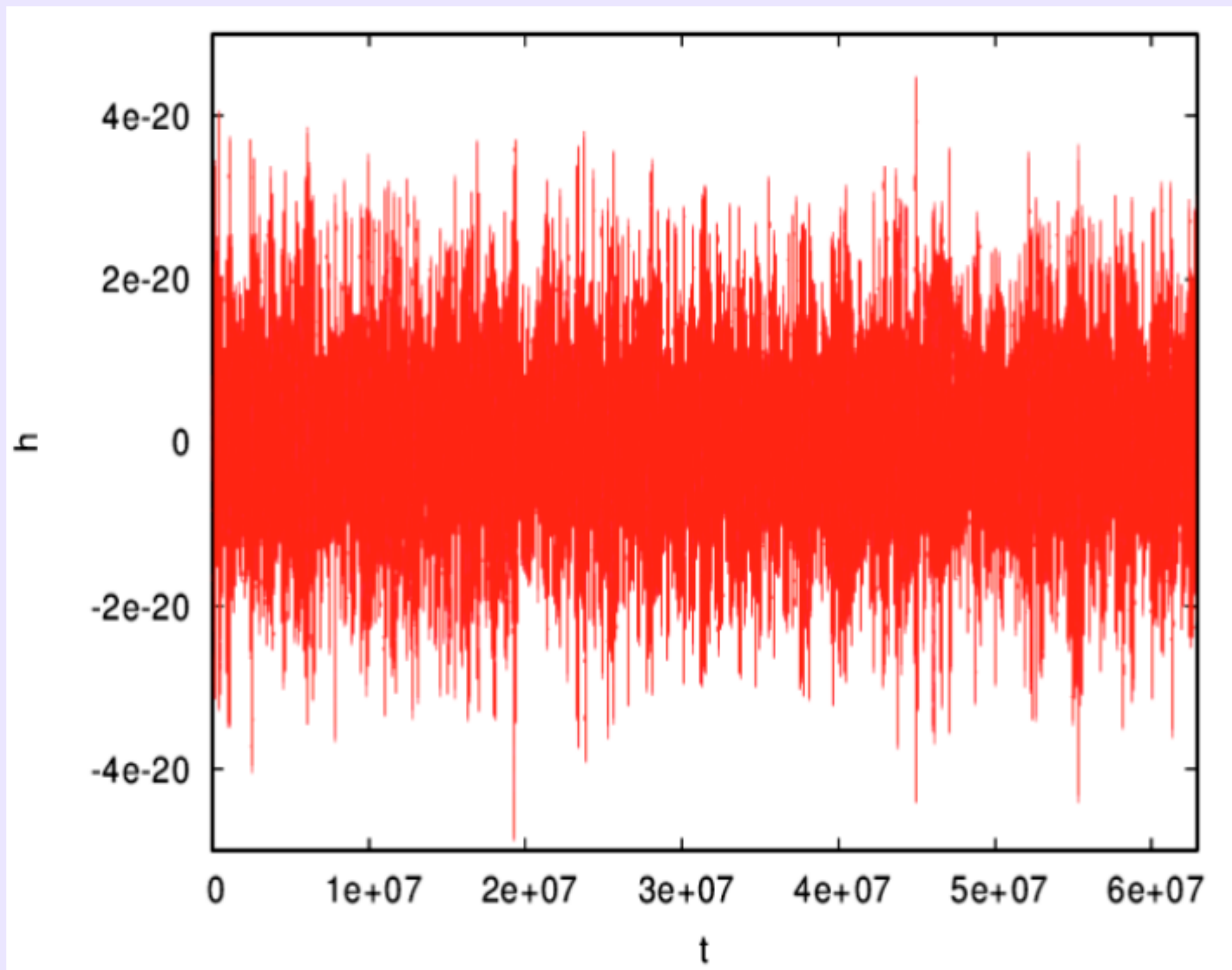
Probing Neutron Star Interiors with Gravitational Waves

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Montana State University

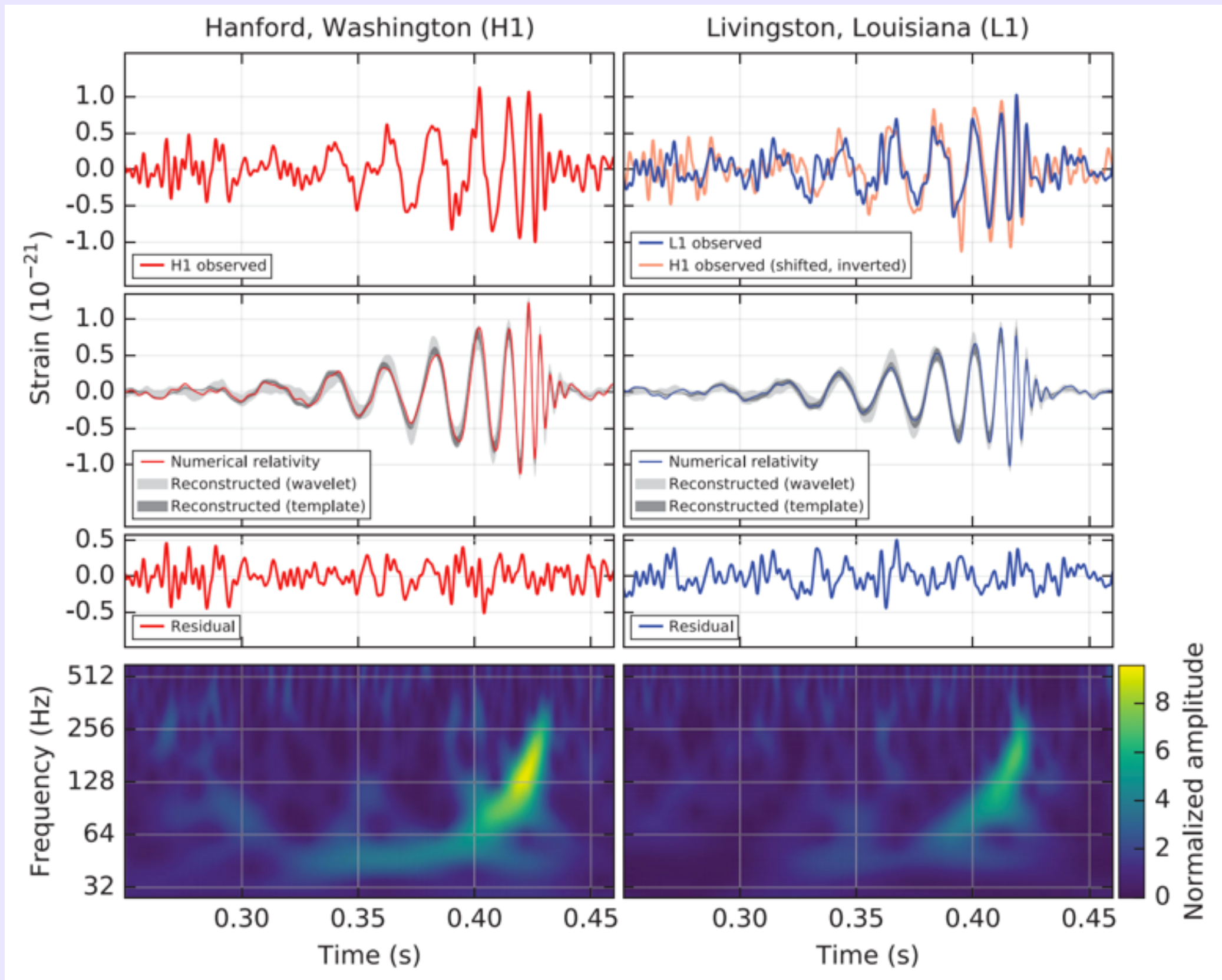
with Kent Yagi, Antoine Klein, Neil Cornish, and Nico Yunes

INT, 2016

Can you spot the Gravitational Wave?



Exciting Times for Gravitational Wave Astrophysics!



From now on...

- Gravitational Waves and Data Analysis
- Neutron Star Compact Binaries
- Model Selection
- Results

Gravitational Waves

Definition: Wave-like perturbation of the gravitational field

$$\square h_{\mu\nu} = T_{\mu\nu}$$

Generation: Accelerating masses (changing quadrupole and higher multipole moments)

$$h_{ij} \sim \frac{1}{R} \frac{d^2 Q_{ij}}{dt^2}$$

Amplitude: Small

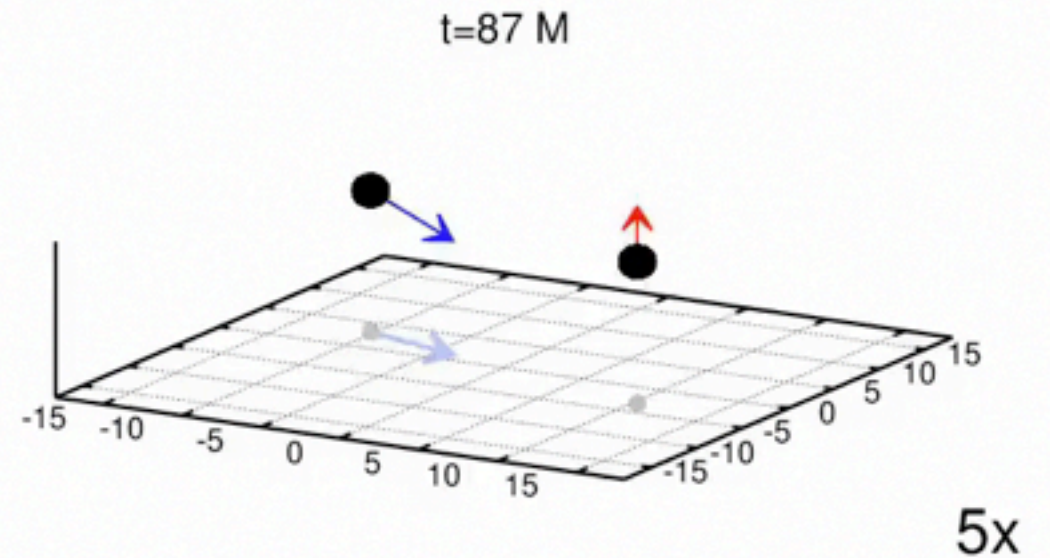
$$h \sim \frac{G}{c^4} \frac{mu^2}{R} \sim 10^{-22}$$

Propagation: Light speed, weakly interacting

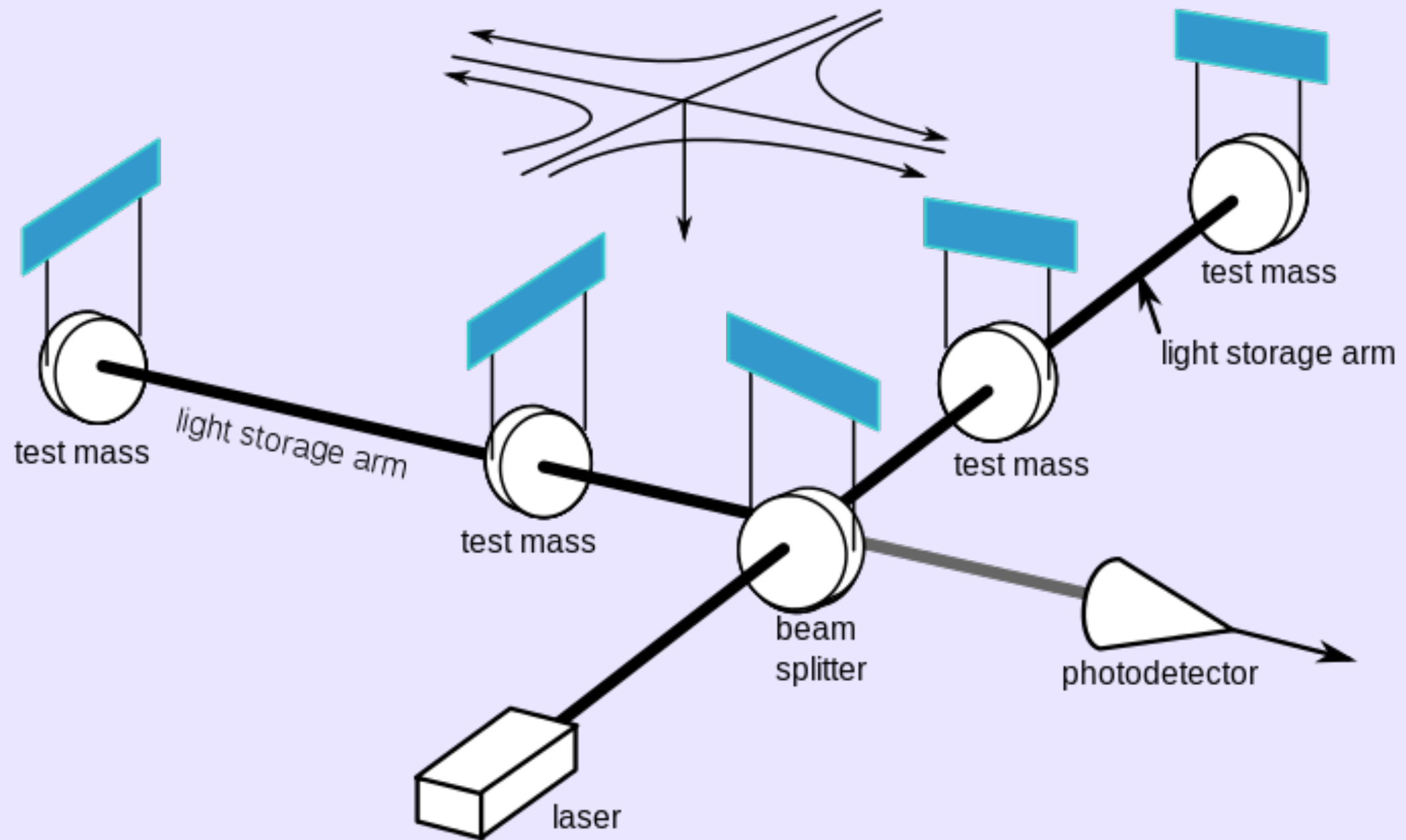
Spectrum: Kepler 3rd Law: $f \sim \sqrt{\frac{m}{r_{12}^3}} \sim \frac{1}{m}$, $E_{rad} \sim \%m$

Example: for GW150914, $E_{GW} \sim 3M_{\odot} \sim 10^3 E_{SN} \sim 0.6 E_{GRB}$

[Lousto, RIT Group]



Gravitational Wave Detectors

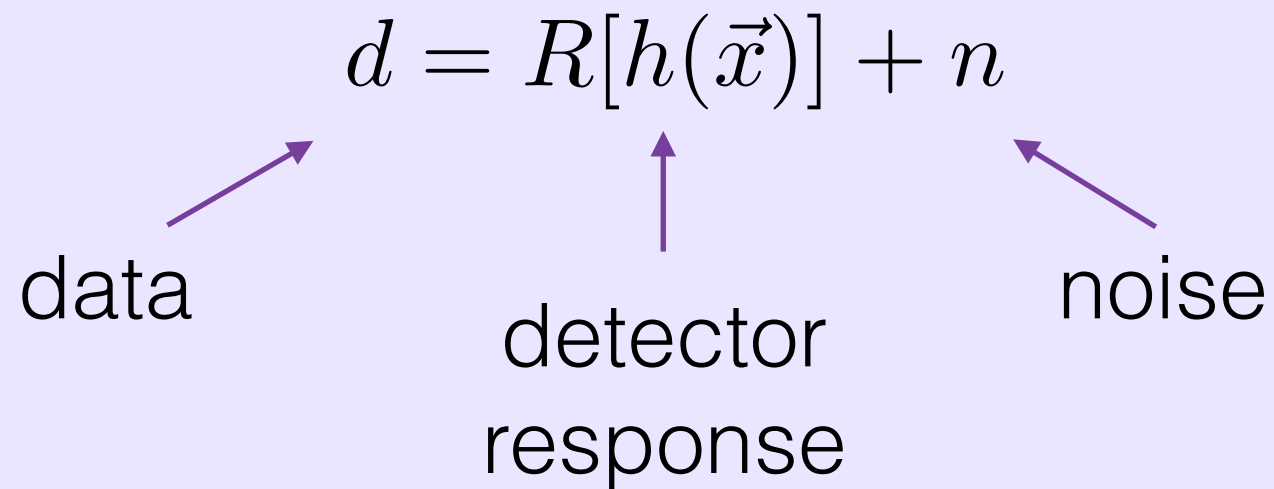


$$h = \frac{\Delta L}{L} \Rightarrow \Delta L = 10^{-19} m$$

Gravitational Wave Data Analysis

$$d = R[h(\vec{x})] + n$$

data detector response noise



Fit the data with a theoretical model for the GW signal

- 1) Get data d
- 2) Select a model $h'(\vec{x})$
- 3) Calculate the residual $r = d - R[h'(\vec{x})]$
- 4) Is this just noise? $p(r) = p(n) = p(d|h'(\vec{x}))$


Likelihood

Bayesian Probability Theory

Degree of belief interpretation of probability

Initial Understanding + New Observations = Updated Understanding

$$p(\vec{x})$$

Prior

$$p(d|\vec{x})$$

Likelihood

$$p(\vec{x}|d)$$

Posterior

+

=

Bayes' Theorem

$$p(\vec{x}|d, M) = \frac{p(\vec{x}|M)p(d|\vec{x}, M)}{p(d|M)}$$

Evidence

$$p(d|M) = \int p(\vec{x}|M)p(d|\vec{x}, M)d\vec{x}$$

Bayesian Model Selection

Probability of model M $p(M|d) \sim p(M)p(d|M)$

Odds ratio

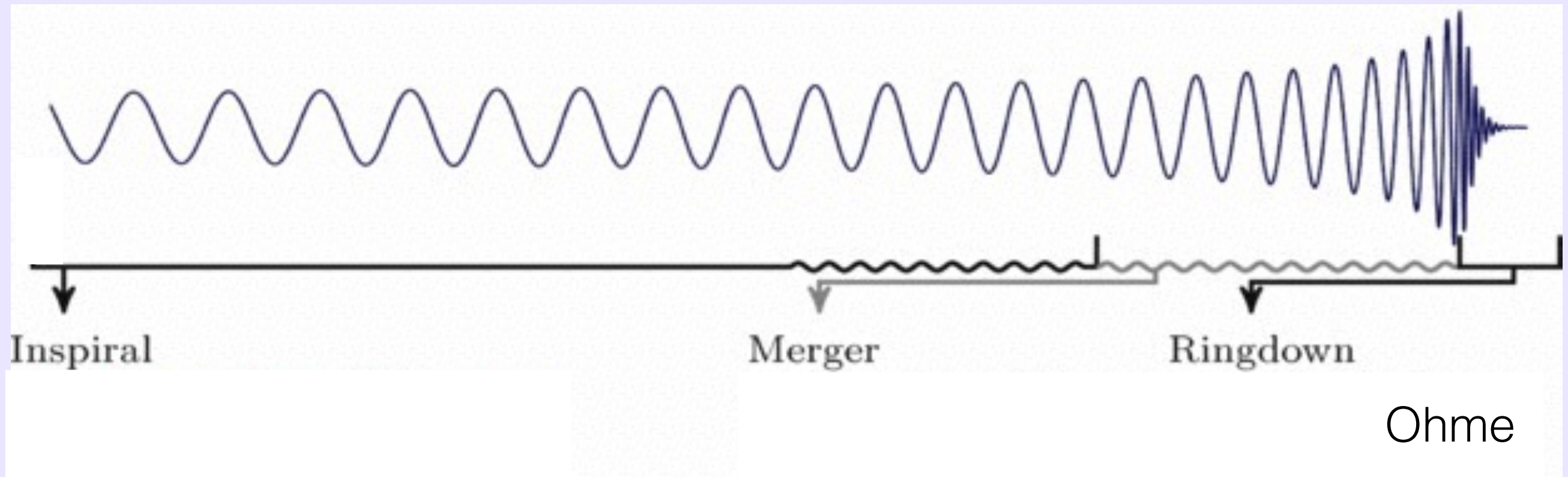
$$\mathcal{O}_{ij} = \frac{p(M_i|d)}{p(M_j|d)}$$

$$= \frac{p(M_i)p(d|M_i)}{p(M_j)p(d|M_j)}$$

**Bayes
Factor**

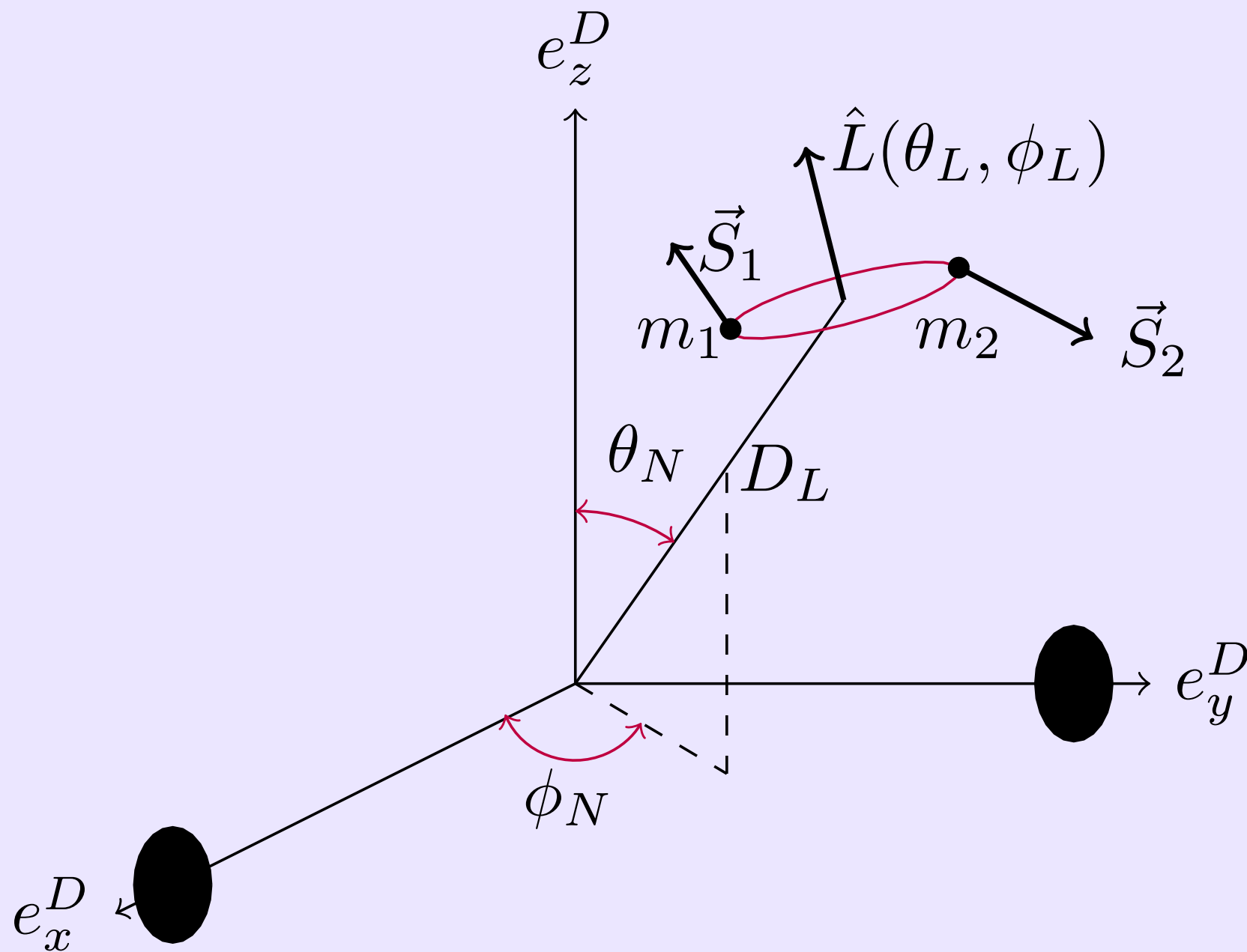
*(the evidence
from before)*

Quasicircular Compact Binary Inspirals



NS/NS:	[20 mins, 10,000 cycles, (10-a few k)Hz]	~400km
BH/NS:	[(3-7) mins, (1,000-5,000) cycles, (10-merger)Hz]	~(500-700)km
BH/BH:	[secs - mins, (100-700) cycles, (10-merger)Hz]	~(600-1000)km

Compact Binary System

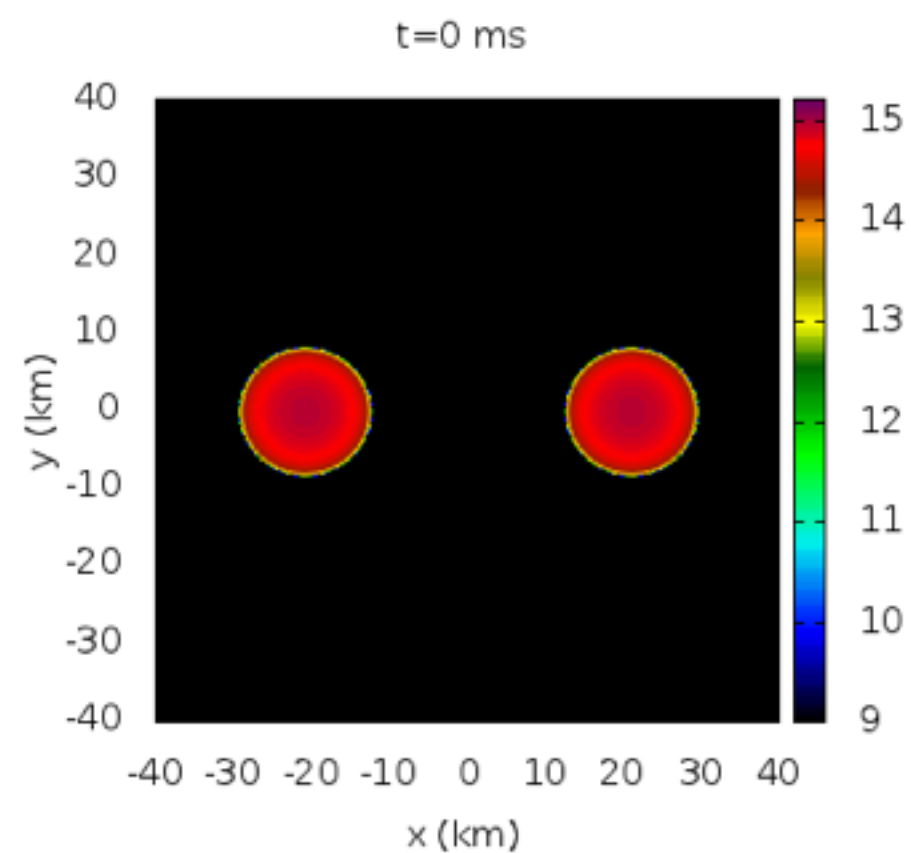
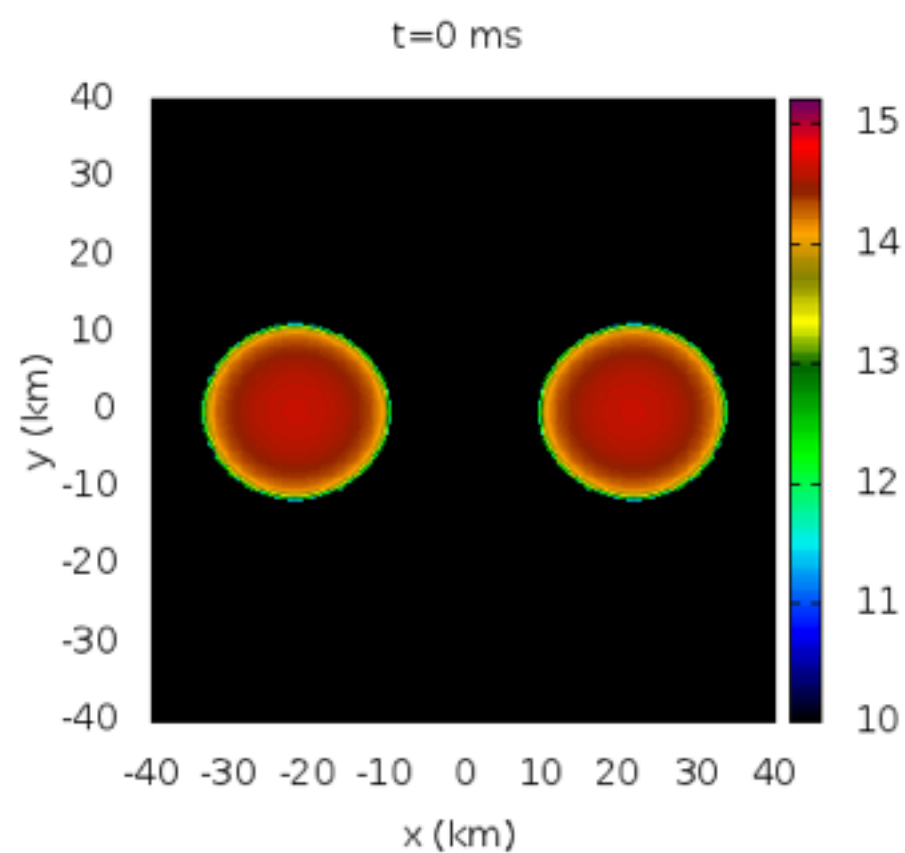


$$\vec{x} = (m_1, m_2, \vec{S}_1, \vec{S}_2, D_L, \theta_N, \phi_N, \theta_L, \phi_L, t_c, \phi_c)$$

Coalescing Neutron Stars and Nuclear Physics

MS1

AP4

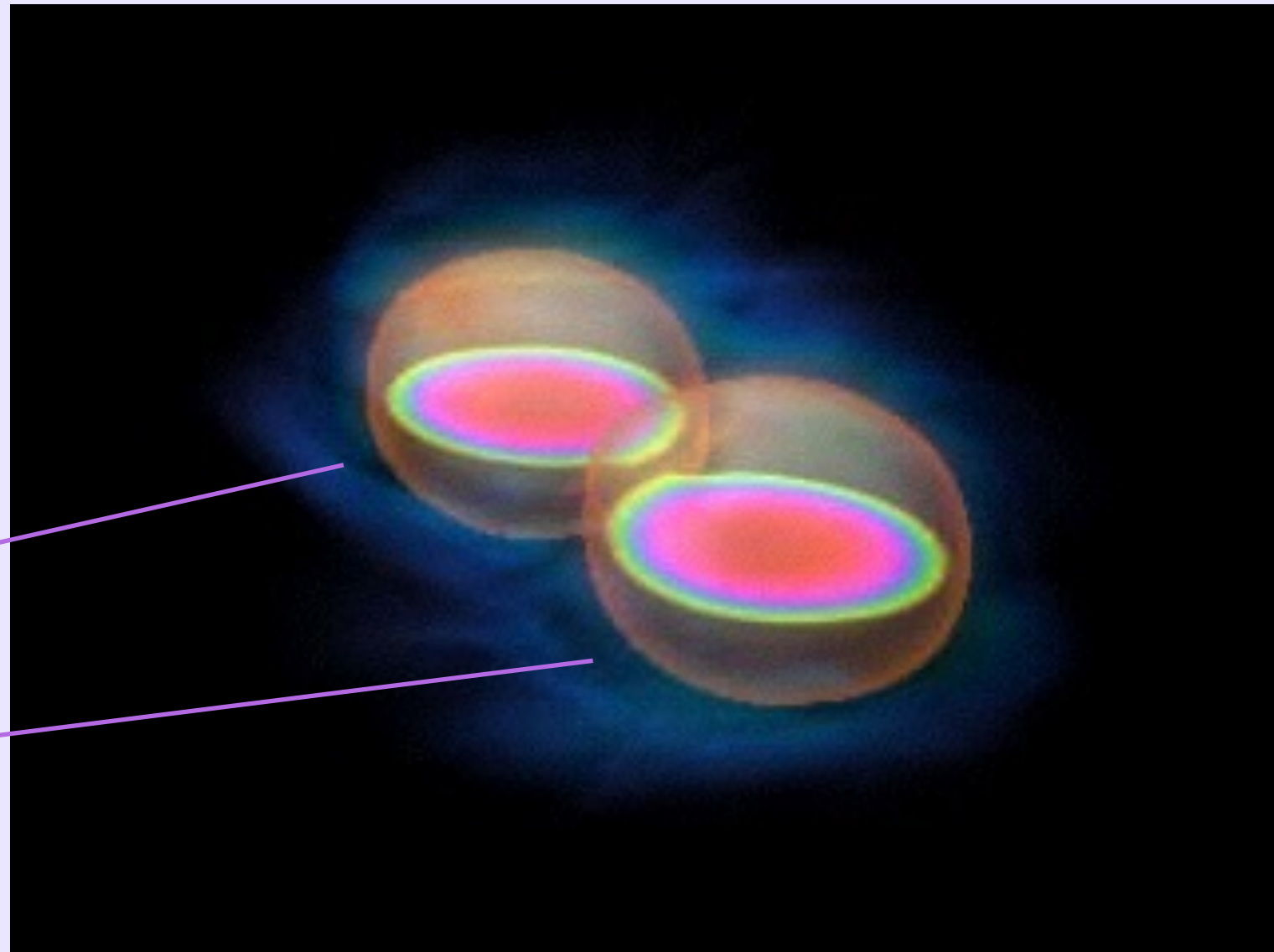


Hotokezaka et al.

Neutron Star Inspirals

Tidal deformability

$$Q_{ij} = -\lambda \mathcal{E}_{ij}$$



Calder

Tidal Deformability

$$\vec{x} = (m_1, m_2, \vec{S}_1, \vec{S}_2, D_L, \theta_N, \phi_N, \theta_L, \phi_L, t_c, \phi_c)$$

$$+ \{ \lambda_i(m_i, \text{EoS}), Q_i(m_i, \text{EoS}) \}$$

= 19 parameters

Tidal Deformability

We can measure the tidal deformability*
with a few bright sources

Read et al. (2009)
Del Pozzo et al. (2013)
Wade et al. (2014)
Agathos et al. (2015)
Lackey and Wade (2015)

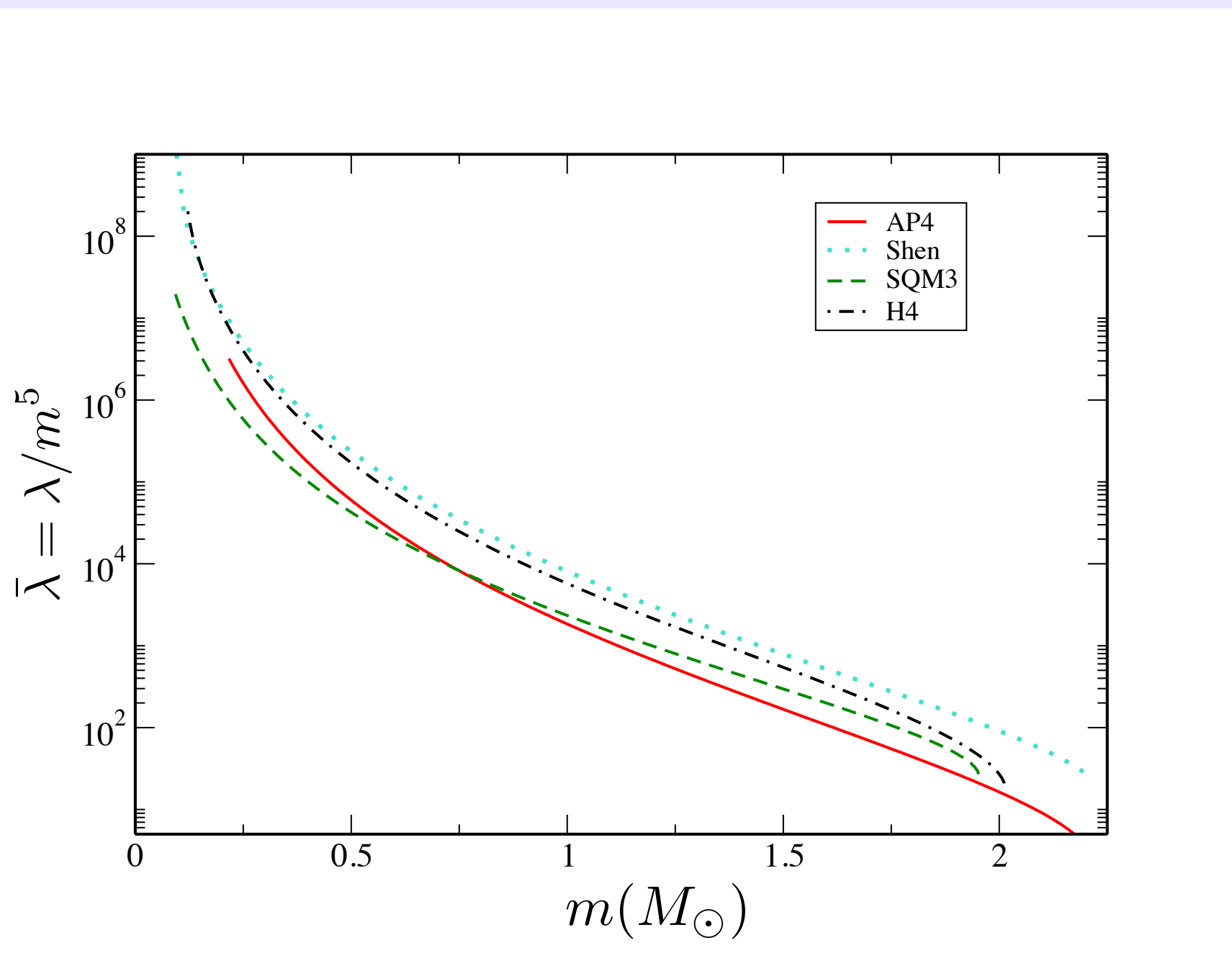
Equation of State

EoS	Method/Model	Composition
AP4	variational	n, p, e, μ
GCR	variational	n
SV	SHF	n, p, e, μ
SIG, SkI4	SHF	n, p, e, μ
DBHF ⁽²⁾ (A)	BHF	n, p, e, μ
MPa	BHF	n, p, e, μ
G4, GA-FSU2.1	RMF	n, p, e, μ
SIG-YBZ6-S $\Lambda\Lambda$ 3, SkI4-YBZ6-S $\Lambda\Lambda$ 3	SHF	n, p, e, μ , H
NIY5KK*	BHF	n, p, e, μ , H
MPaH	BHF	n, p, e, μ , H
H4	RMF	n, p, e, μ , H
SIG178	SHF	n, p, e, μ , K
SV222	SHF	n, p, e, μ , K
GA-FSU2.1-180	RMF	n, p, e, μ , K
ALF4, ALF5	variational	n, p, e, μ , π , Q
GCR-ALF	variational	n, Q
SQM3	MIT bag	Q (u, d, s)

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**Can GWs Distinguish NS Binaries
with Different Internal Composition ?**

Model Selection

We need to calculate the evidence and the odds ratio

$$\mathcal{O}_{ij} = \frac{p(M_i) p(d|M_i)}{p(M_j) p(d|M_j)}$$


How much we believe in each model before acquiring the data.

Based on our previous experience, observational evidence, and theoretical understanding of the Universe.

Which of two competing model fits the data at hand better.

Bayes Factor

$$\text{BF} = \frac{p(d|M_i)}{p(d|M_j)}$$

When is the BF 'large enough'?

BF	Interpretation
<1	Negative
<3	Barely worth mentioning
<10	Strong
<100	Very Strong
>100	Desicive

Jeffreys scale
of BF
interpretation

The Evidence (or the ratio)

- Laplace Approximation
- Schwarz-Bayes Information Criterion
- Reversible Jump MCMC
- Thermodynamic Integration
- Nested Sampling
- Savage-Dickey Density Ratio

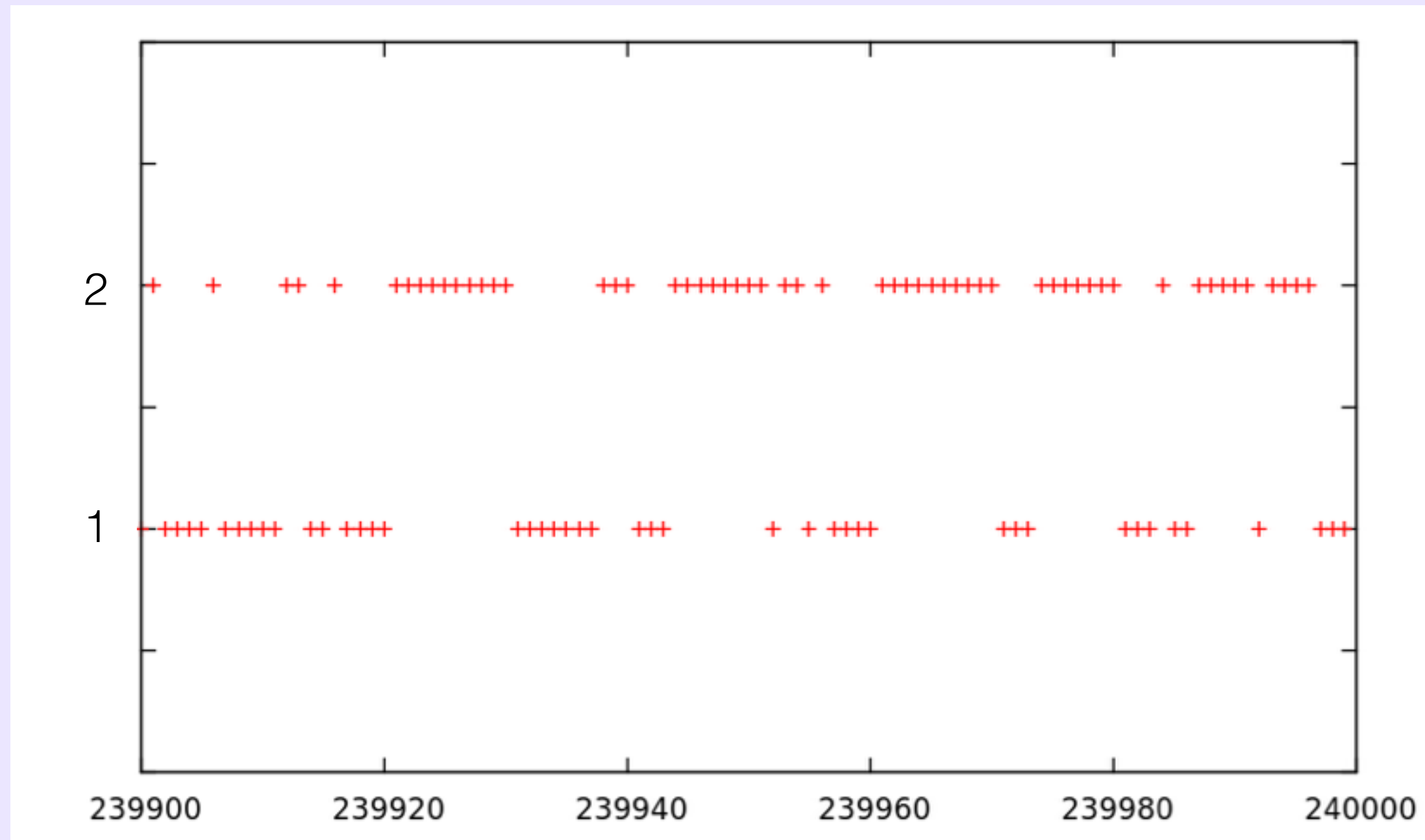
Model Parameters

$$\vec{x} = (m_1, m_2, \vec{S}_1, \vec{S}_2, D_L, \theta_N, \phi_N, \theta_L, \phi_L, t_c, \phi_c)$$

$$+ \{\text{EoS}\}$$

15 *continuous* parameters, and 1 *discrete*

Reversible Jump Markov Chain Monte Carlo



$$\text{Bayes Factor} = \frac{\# \text{ of iterations in model 1}}{\# \text{ of iterations in model 2}}$$

Errors (with RJMCMC)

$$\text{Bayes Factor} = \frac{\# \text{ of iterations in model 1}}{\# \text{ of iterations in model 2}}$$

For well-mixed chains

$$\text{Var}(\text{BF}) = \text{BF}^2 \left(\frac{N_1 - N_{12}}{N_1 N_{12}} + \frac{N_2 - N_{21}}{N_2 N_{21}} \right)$$

Prior

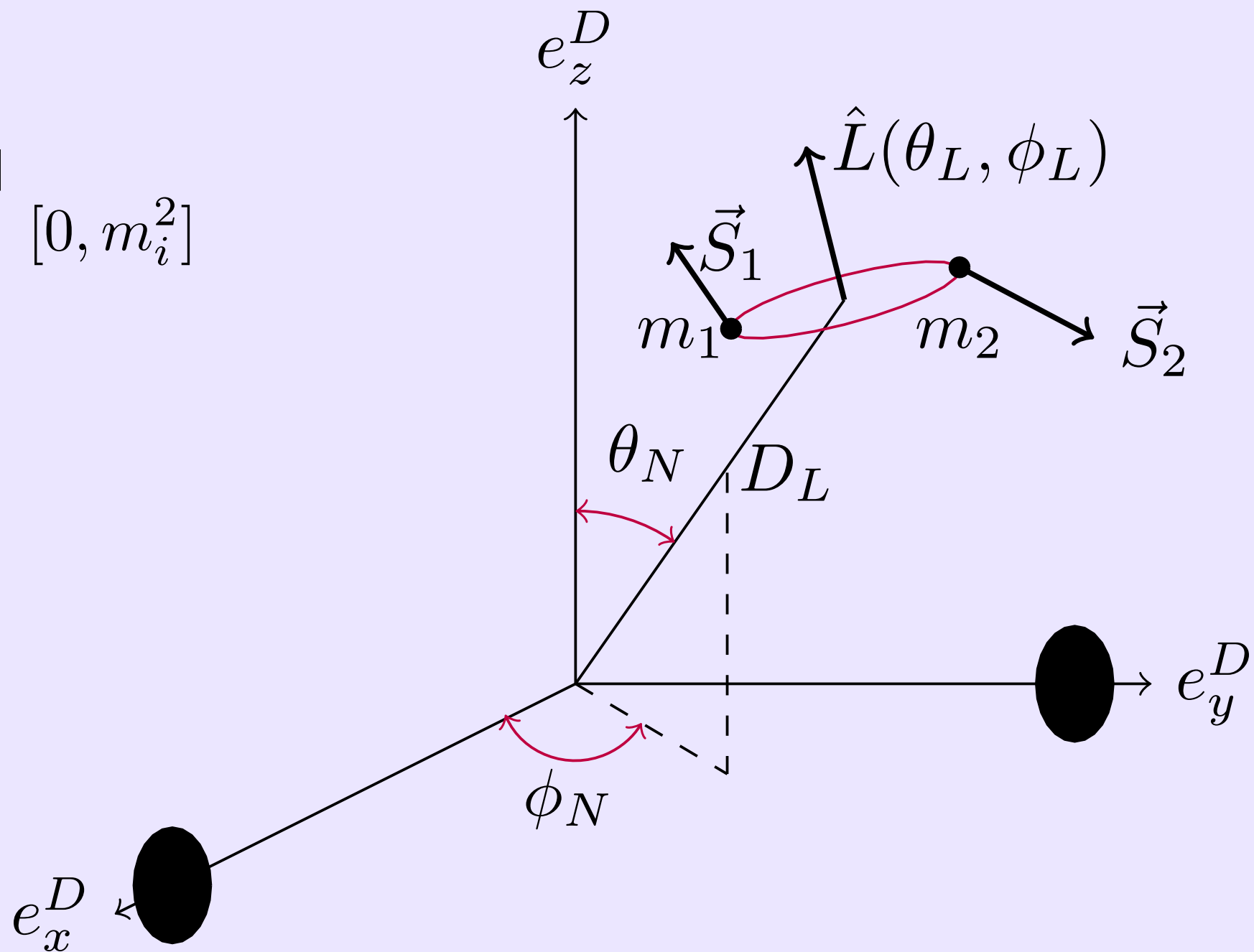
m_1
 m_2 Uniform in
[0.1, 3.2] M_\odot

\vec{S}_2
 \vec{S}_1 Uniform in
direction and
magnitude in $[0, m_i^2]$

D_L Uniform in
volume

θ_N
 ϕ_N Uniform in
the sky

θ_L
 ϕ_L Uniform in
direction



Likelihood: the Noise Model

$$p(d|h) = p(d - R[h]) = p(n)$$

Correlated *Gaussian* noise

$$p(n_1 \dots n_N) = \frac{1}{\sqrt{\det(2\pi C)}} e^{-\frac{1}{2} n_i C_{ij}^{-1} n_j}$$

Stationary noise

$$C_{f_i f_j} \sim \delta_{ij} S(f_i)$$

Likelihood: the Noise Model

Easier to evaluate

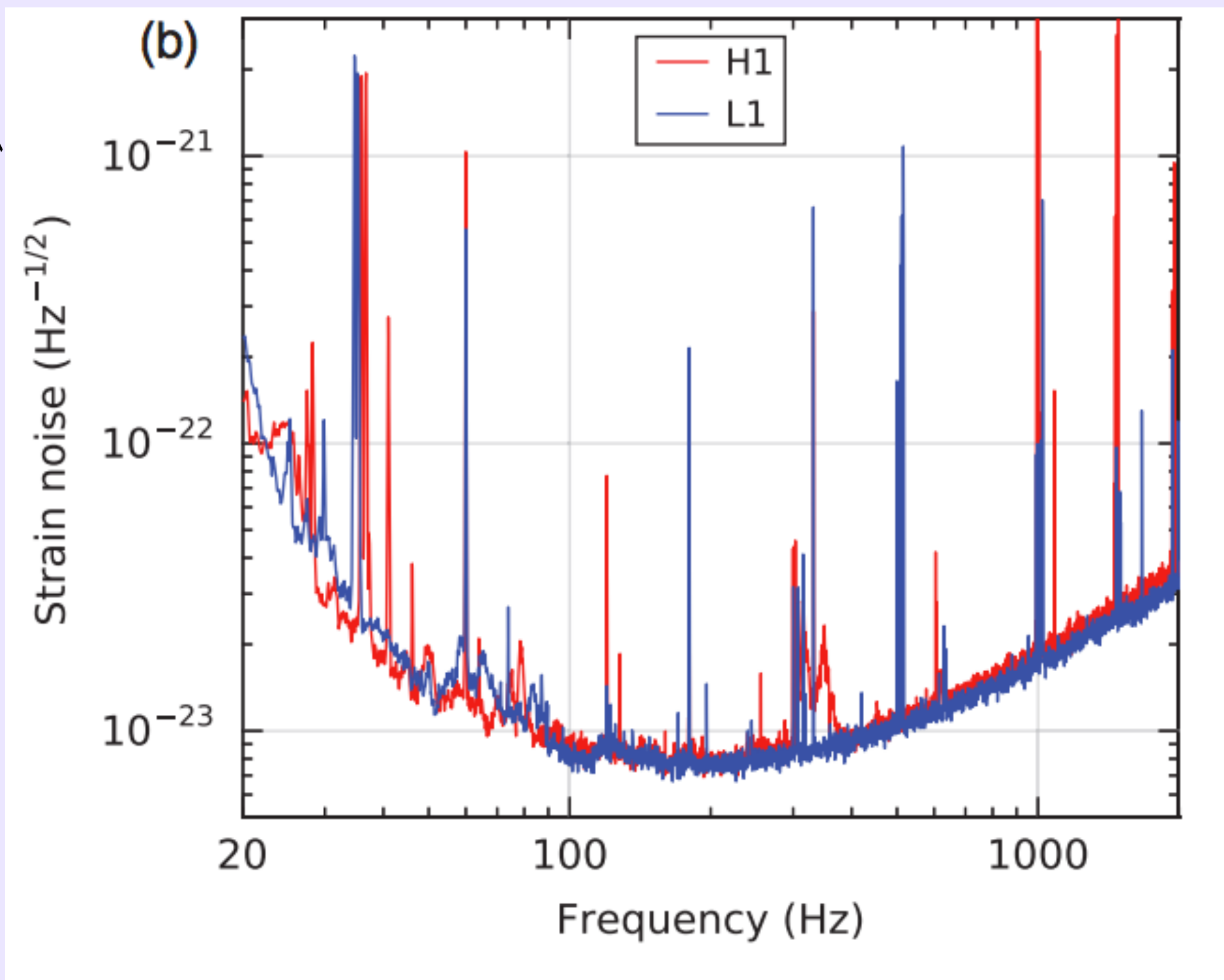
$$n_i C_{ij}^{-1} n_j = (n|n) \sim \int \frac{\tilde{n}(f) \tilde{n}^*(f)}{S(f)} df$$

Our noise model

$$p(d|\vec{x}) \sim e^{-\frac{(d-h(\vec{x})|d-h(\vec{x}))}{2}}$$

Noise

related to $S(f)$



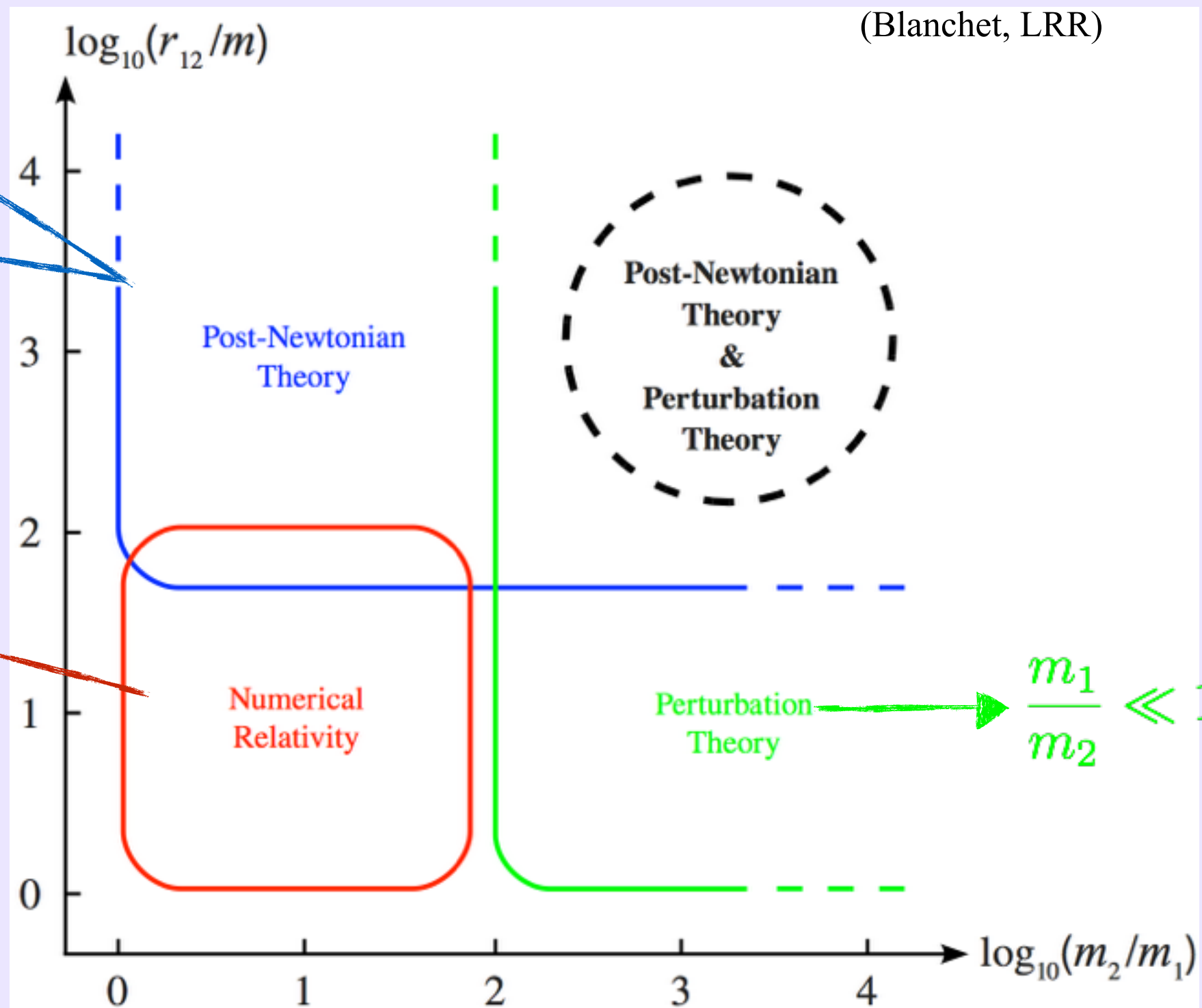
Building Models

$$G_{\mu\nu} = \frac{8\pi G}{c^4} T_{\mu\nu}$$

$$\frac{Gm_{\text{tot}}}{c^2 r_{12}} \ll 1$$

$$\frac{v}{c} \ll 1$$

(Blanchet, LRR)



Models: Inspiral GW

$$\vec{x} = (m_1, m_2, \vec{S}_1, \vec{S}_2, D_L, \theta_N, \phi_N, \theta_L, \phi_L, t_c, \phi_c) + \{\text{EoS}\}$$

- GW described by \vec{x} , $(m_1, m_2) \leq M_{\text{max}}(\text{EoS})$
- No GW, otherwise

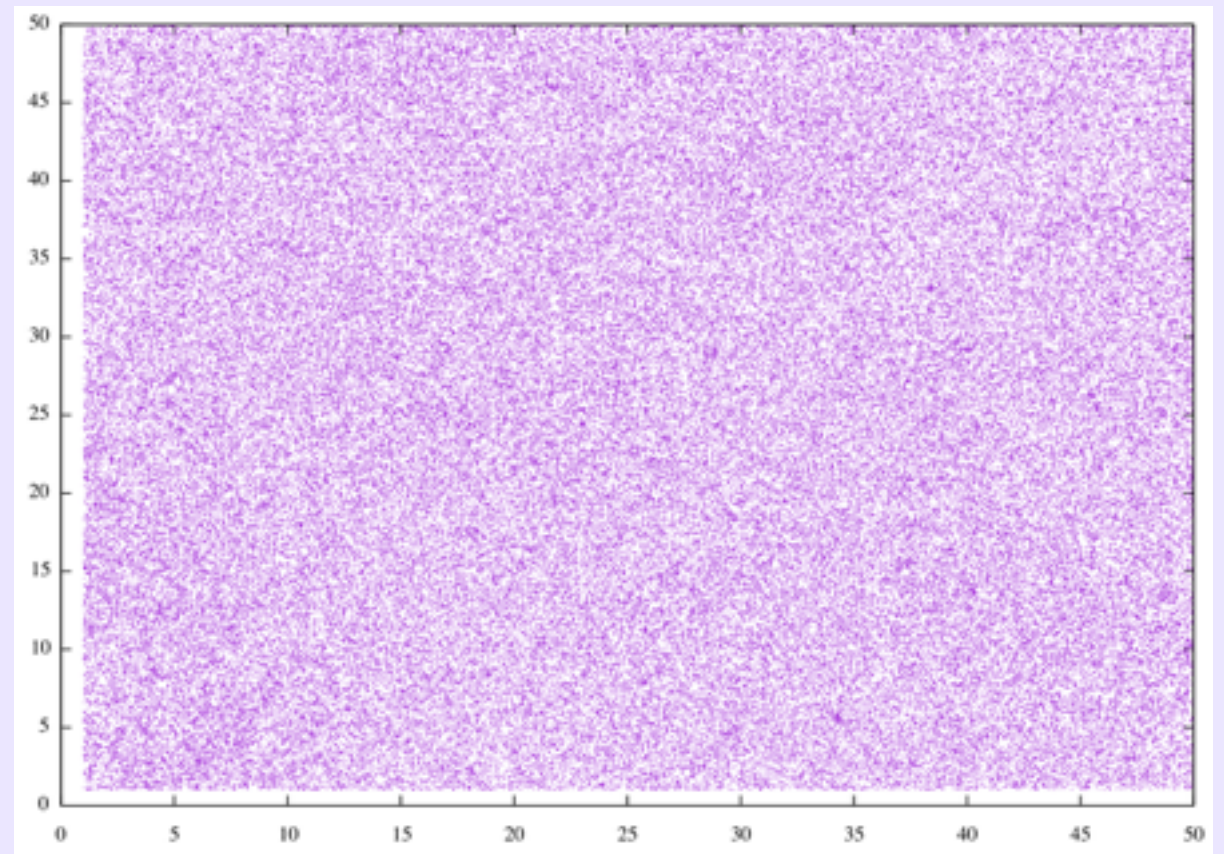
Reasonably fast to evaluate

Proposal Distribution

Prior

Provides access to the entire prior volume

(essential to pass the constant likelihood test)



Jiggle

Search around the current position

$$\vec{\theta} + \vec{\epsilon}$$

Proposal Distribution

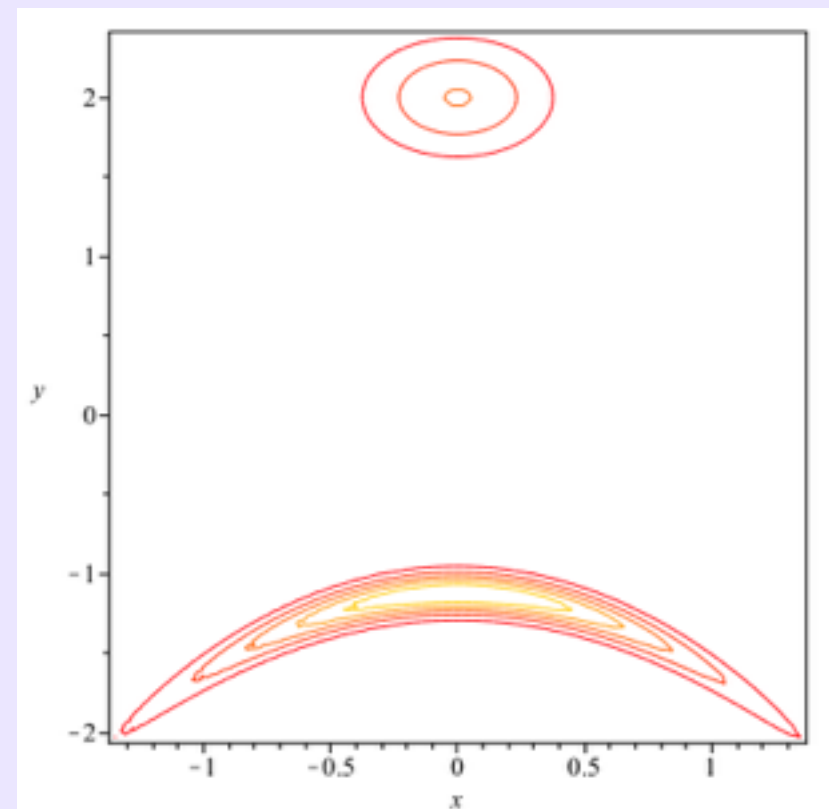
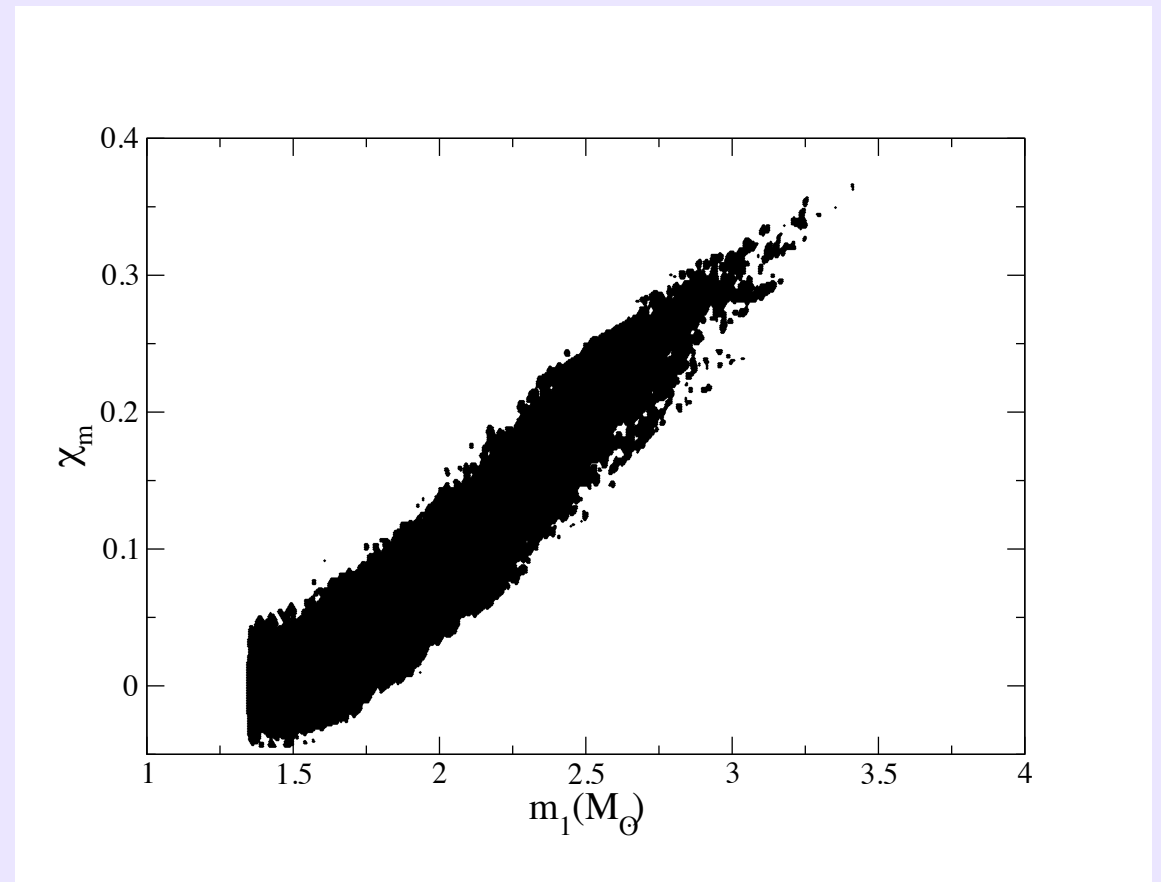
Fisher

Jump along the eigendirections (scaled by the eigenvalues) of the *Fisher* Information Matrix

$$F_{ij} = (h_{,i} | h_{,j})$$

Langevin

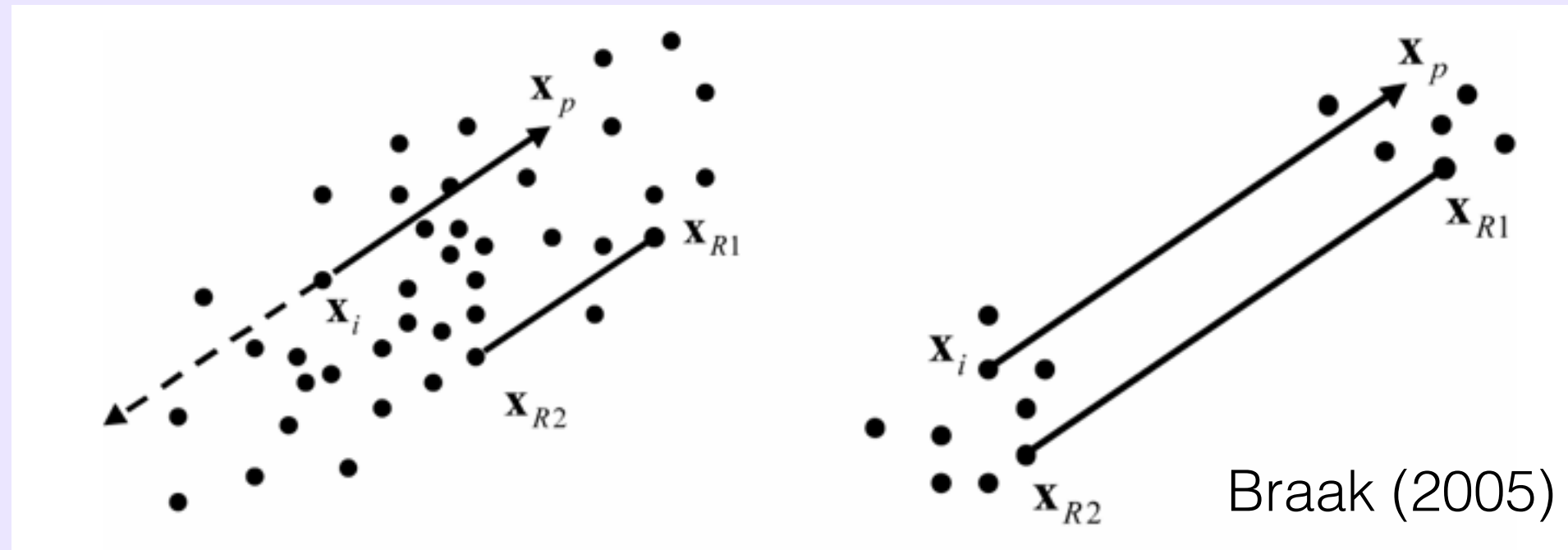
jump along the likelihood gradient



Proposal Distribution

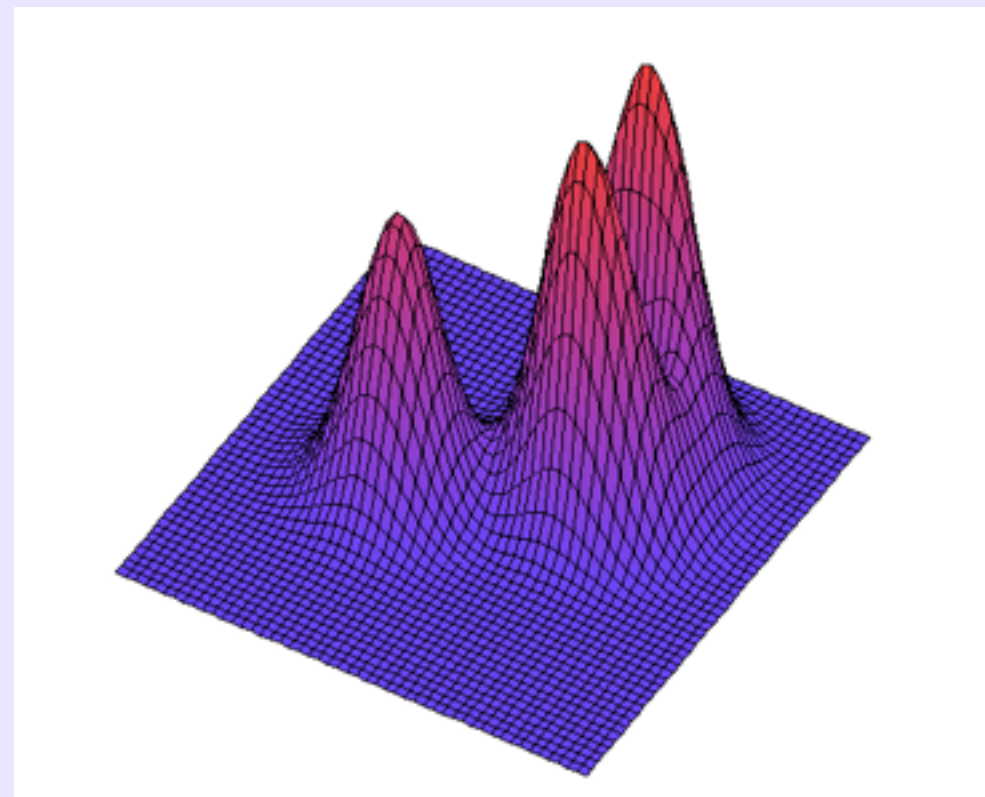
Differential Evolution

(technically it is not memoryless)



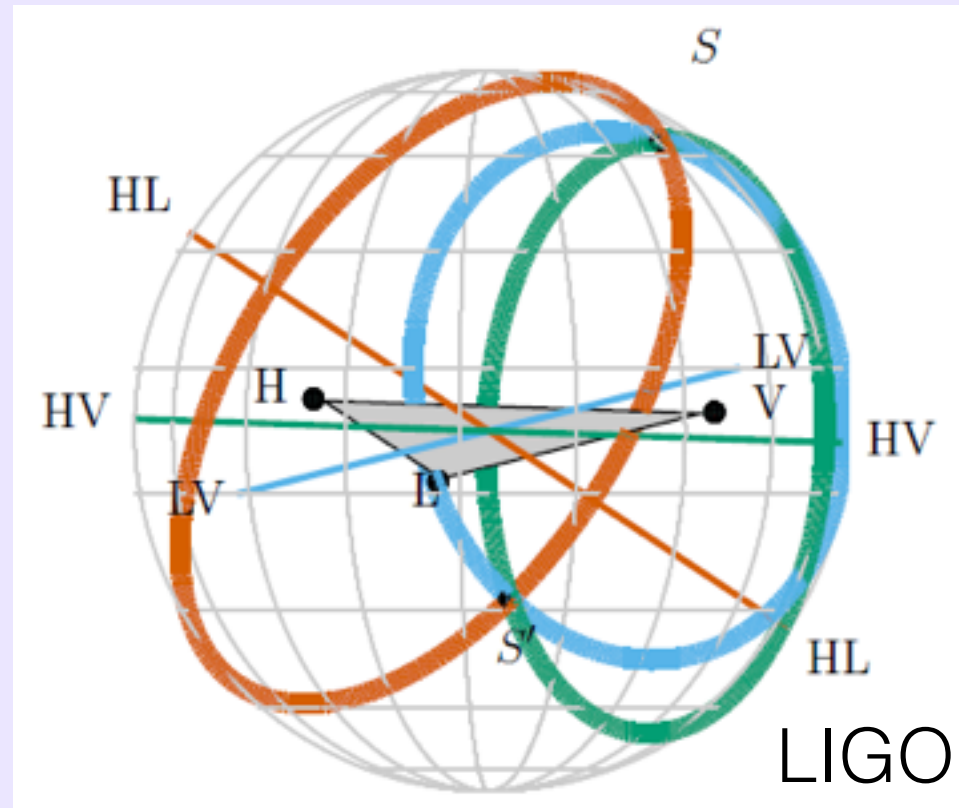
Model jumps

Pilot runs

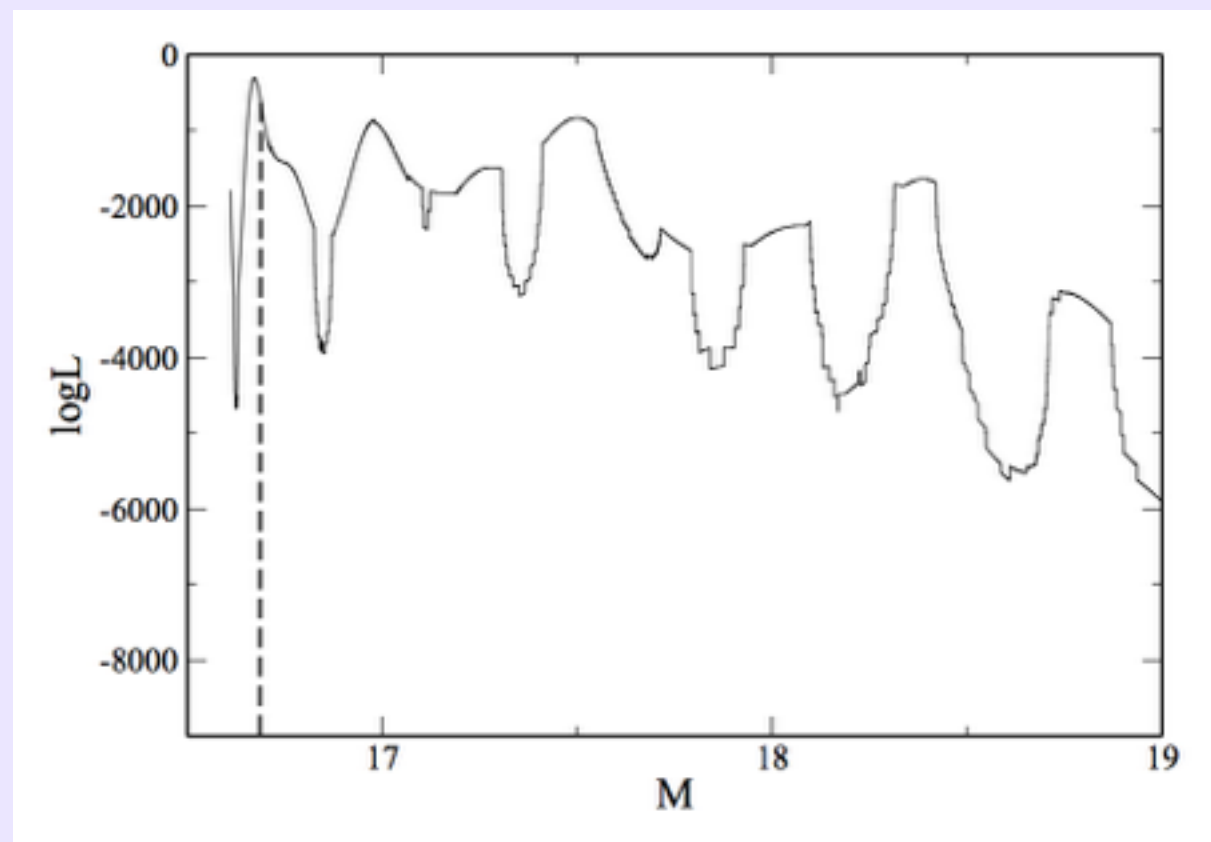


Proposal Distribution

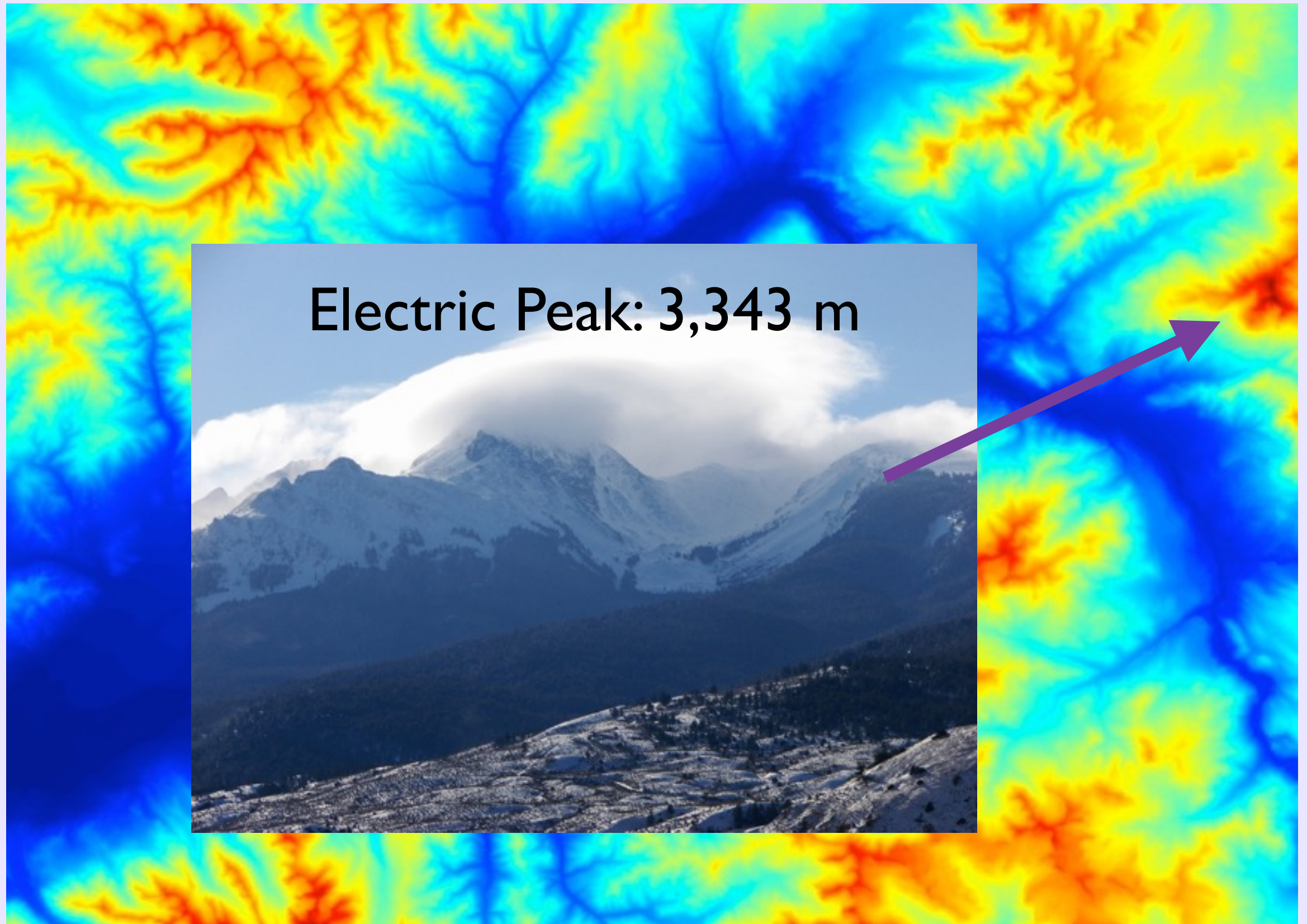
Sky jumps



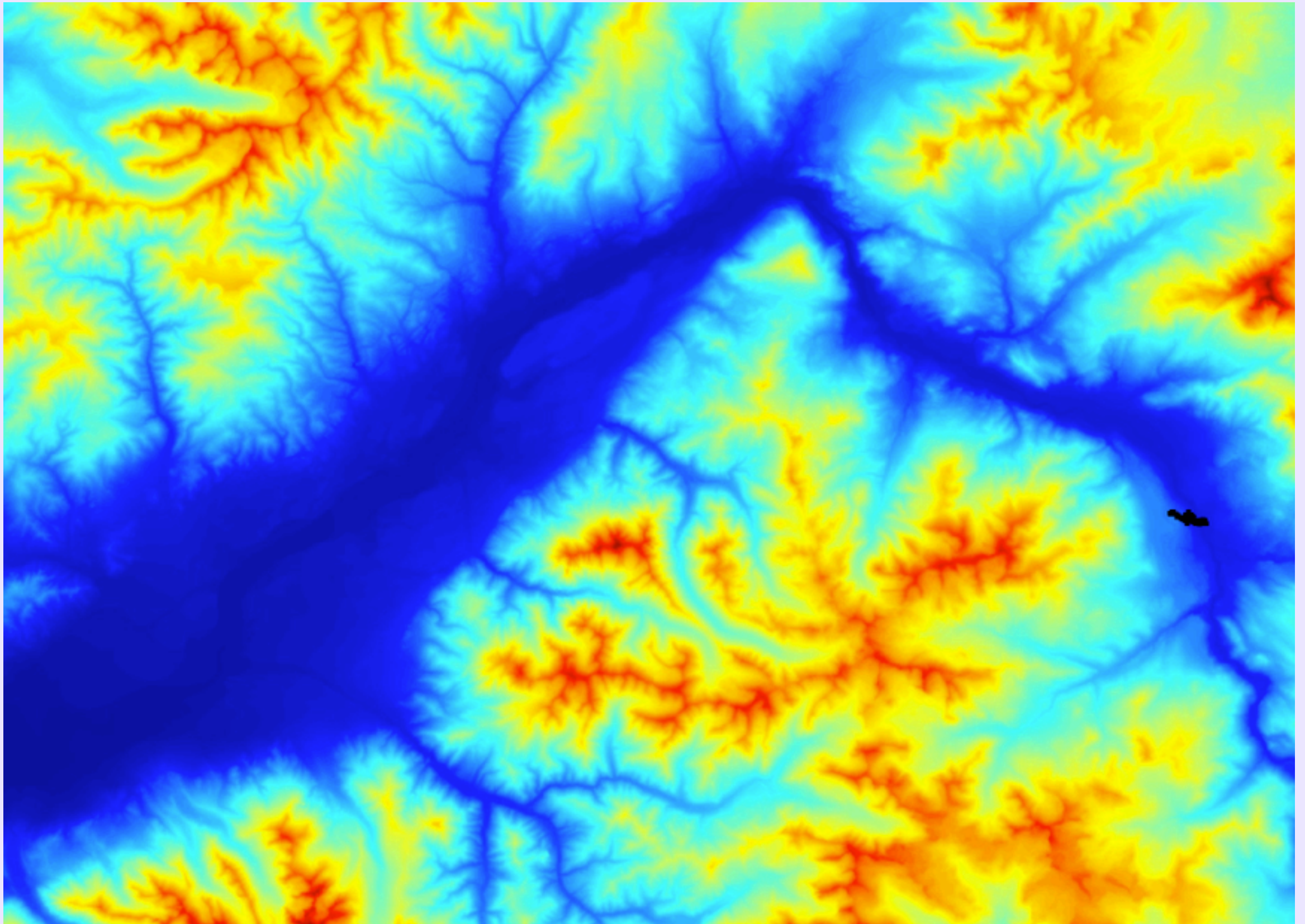
Customized



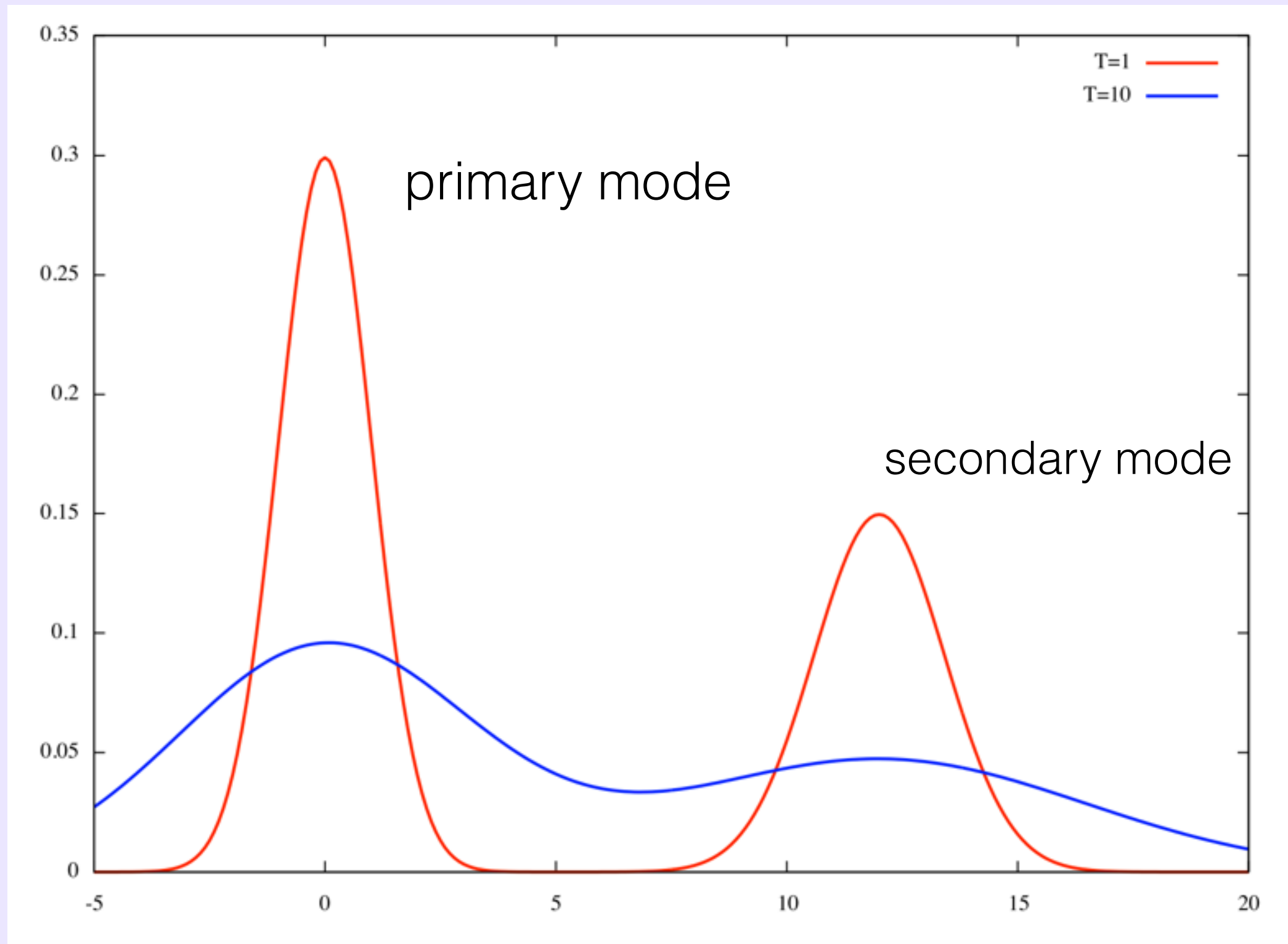
Finding the Highest Peak in Gallatin Range



MCMC

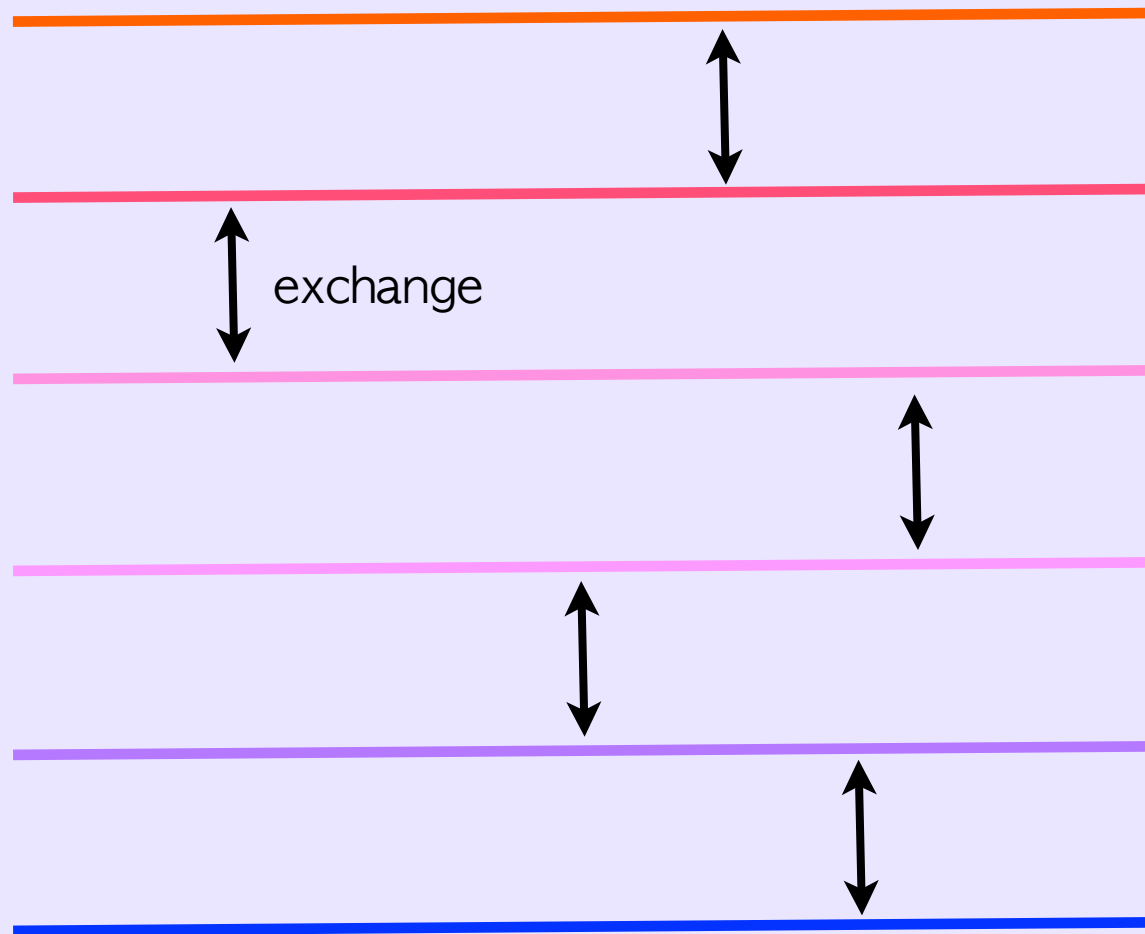


Multi-Modal Distributions



$$p(d|\vec{x}) \rightarrow p(d|\vec{x})^{1/T}$$

Exchanges



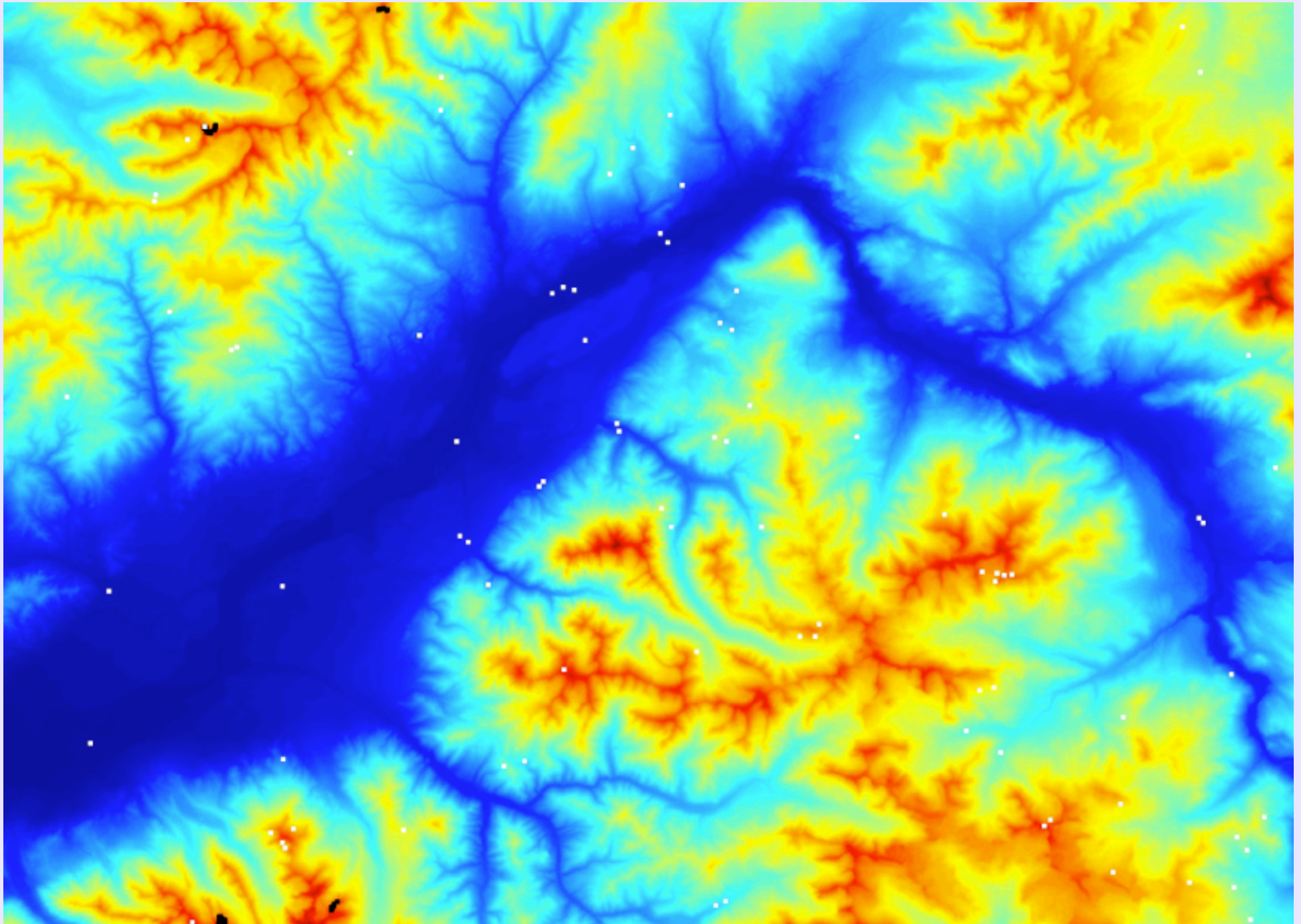
Wide exploration

Good solutions

Limited exploration

Kirkpatrick, Gelatt, Vecchi (1983)
Swendsen, Wang (1986)

Parallel Tempering

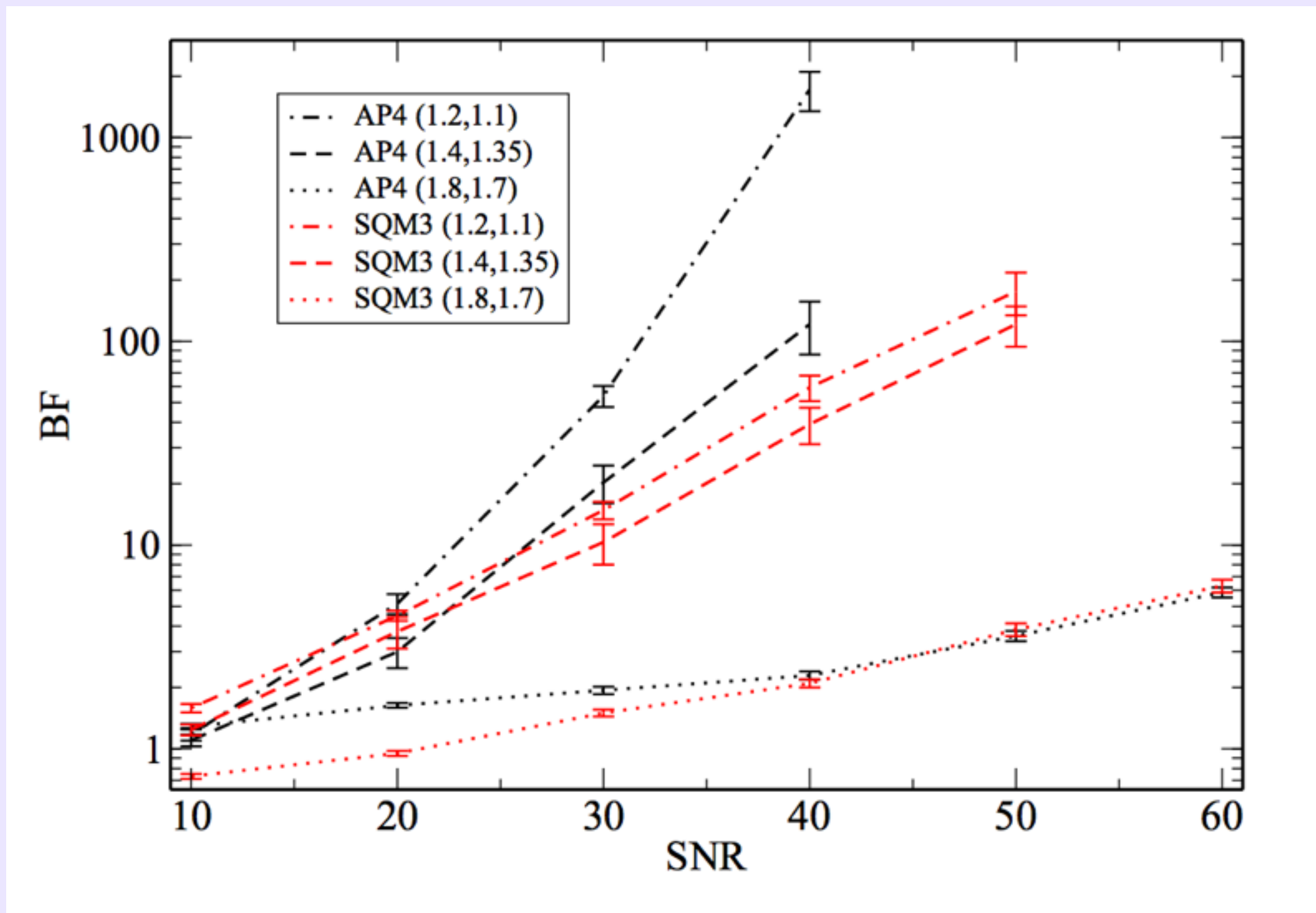


Our Analysis

EoS	Method/Model	Composition
AP4	variational	n, p, e, μ
GCR	variational	n
SV	SHF	n, p, e, μ
SGI, SkI4	SHF	n, p, e, μ
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ALF4, ALF5	variational	n, p, e, μ , π , Q
GCR-ALF	variational	n, Q
SQM3	MIT bag	Q (u, d, s)

If an EoS **with kaons** fits the data better than an otherwise identical EoS **without kaons**, then we have **detected** kaons in a NS interior

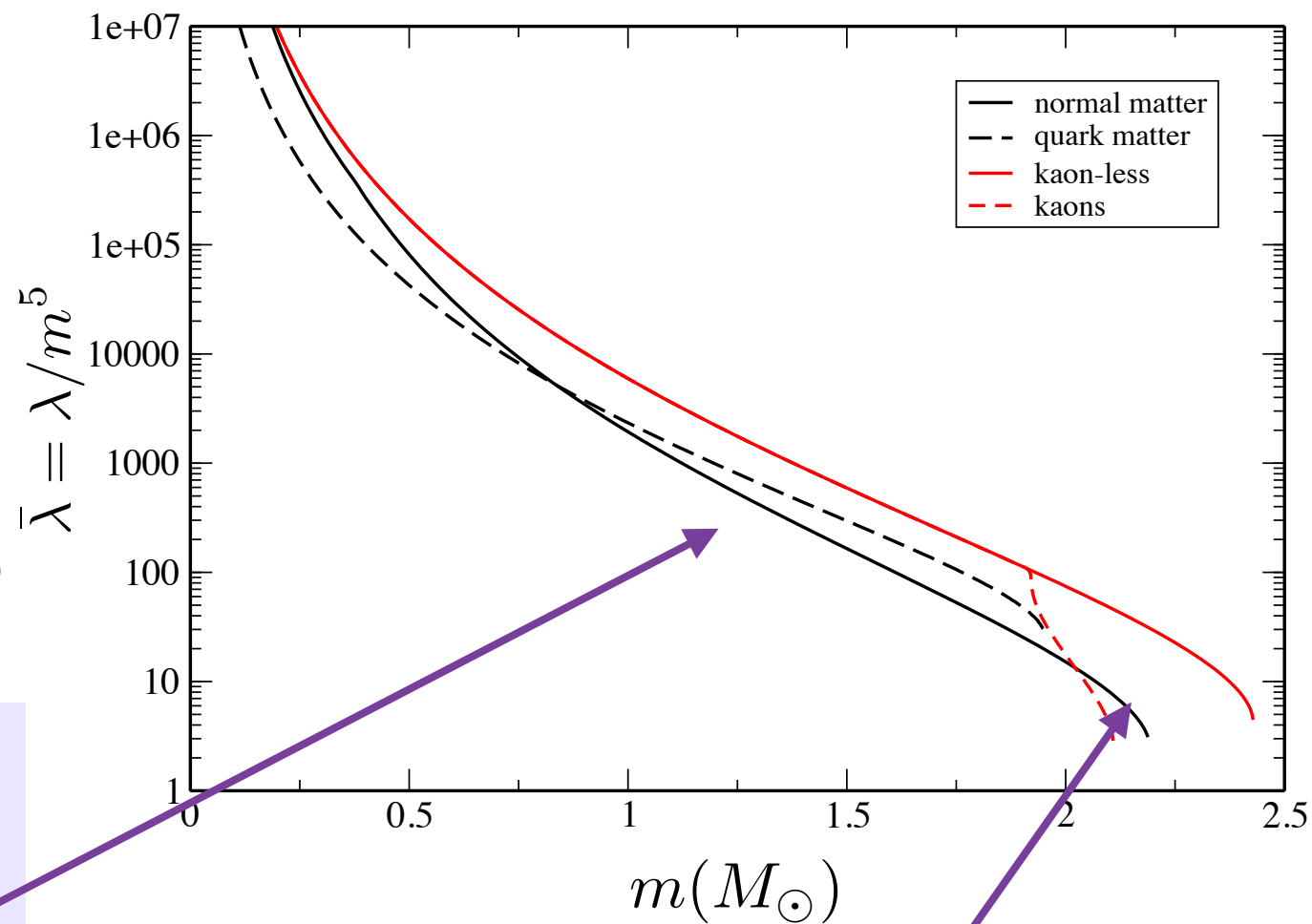
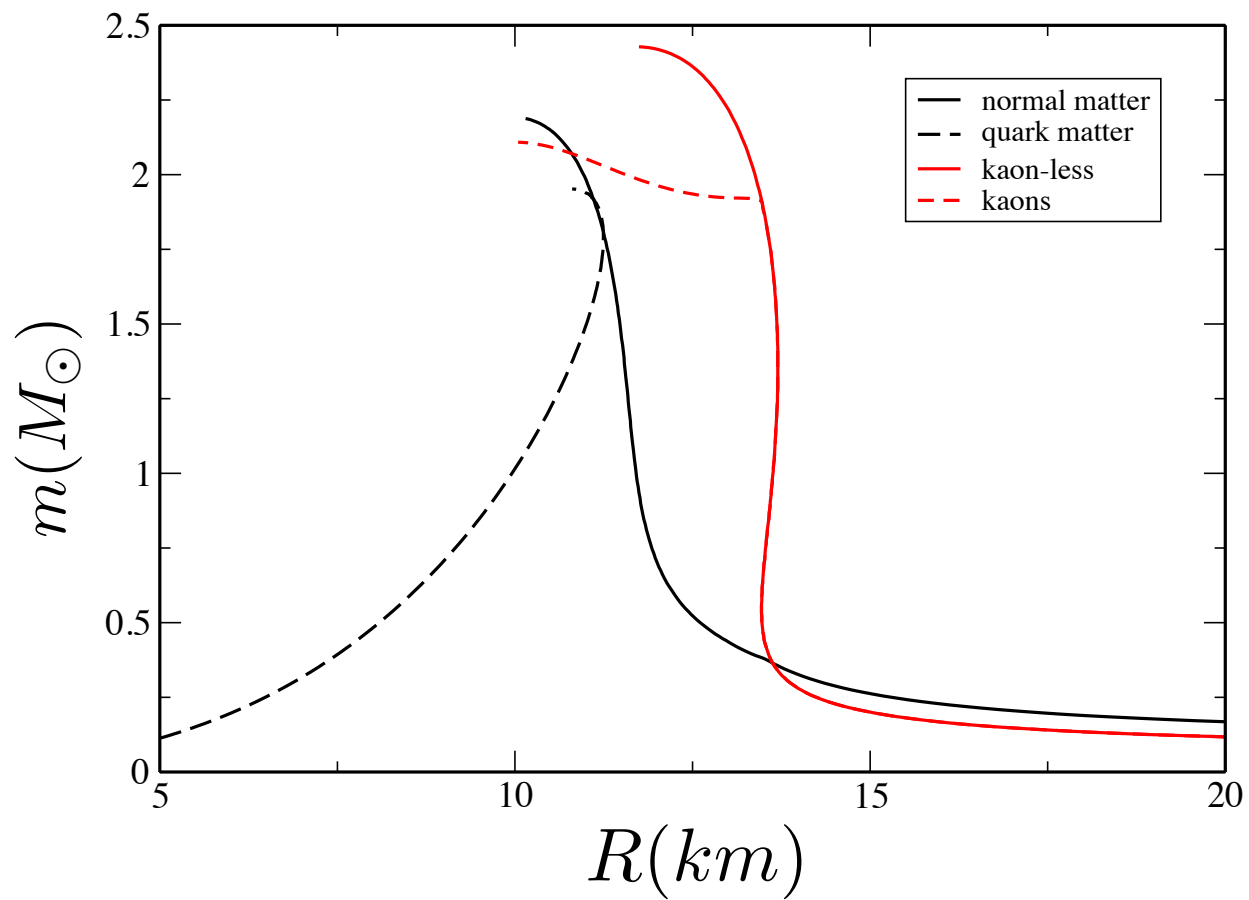
Bayes Factors



Results

	Strange Quark Stars	Hybrid Quark Stars	Kaons	Hyperons
aLIGO	Yes!	Maybe	Unlikely	Unlikely
SNR	20	30-40	50-60	50-60
mass	$(1.2, 1.5)M_{\odot}$	$1.4M_{\odot}$	$2M_{\odot}$	$2M_{\odot}$

Mass Matters



Directly measure low mass effects

Partial information on high mass effects

Errors

Systematic Errors: our models might be wrong

- General Relativity might be wrong

- Perturbative models not accurate enough

- Models not accurate astrophysically

- Unknown noise contribution

- Detector Calibration

Statistical Errors: finite signal strength

- Width of the Posterior

- Noise Realization

- Marginalization

Further work

P

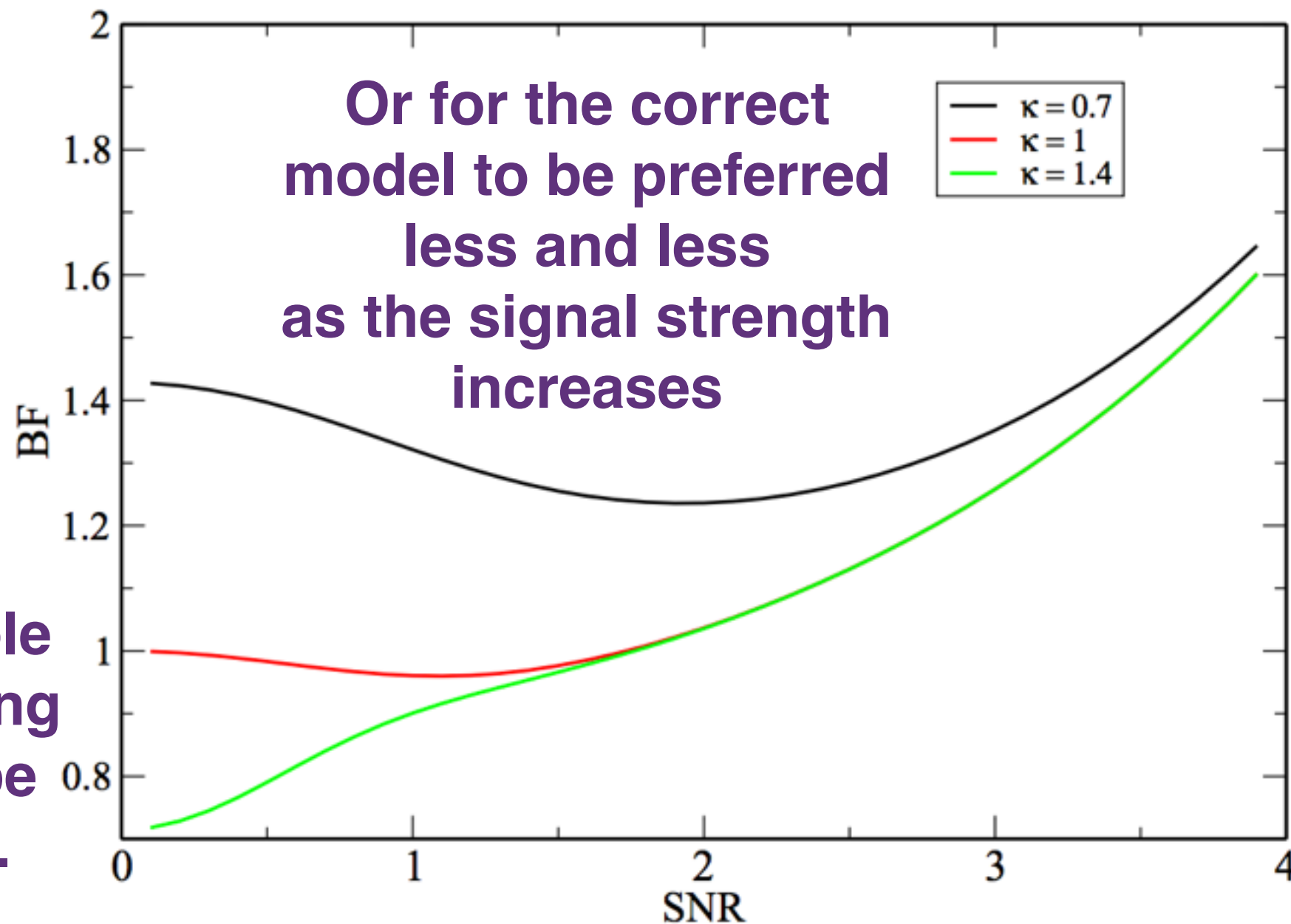
Further meaningful comparisons
Inspiral phase: improve modeling
Merger phase: modeling

S

Efficient trans-model jumps
Exploration of disfavored models
Thermodynamic Integration
Merger phase: unmodeled search

Thank you!

The Edge of the Prior



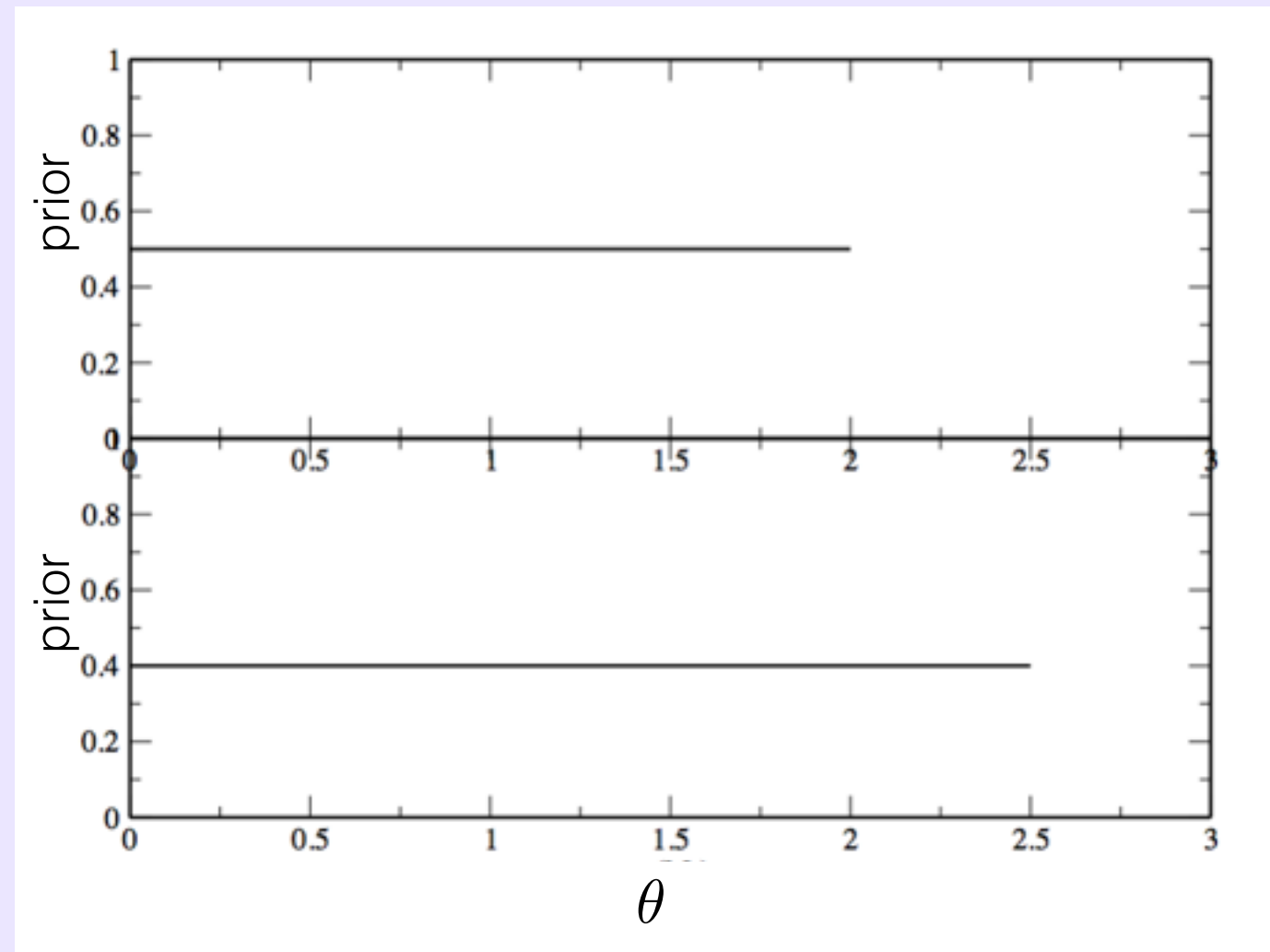
It is possible for the wrong model to be preferred.

Occam Penalty

A model that requires more parameters to fit the data is penalized

$$\frac{\delta\theta}{\Delta\theta}$$

But what if it's the denominator that changes between the various models?



Toy Model

We get N data from a signal $d(f) = f$

Two competing models

- $h_1 = af \quad a \in (0, 2)$
- $h_2 = af^{1.5} \quad a \in (0, 2\kappa)$

Likelihood

$$L_i = \frac{1}{\sqrt{2\pi\sigma}} \exp \left\{ - \sum^N \frac{[d(f) - h_i(f)]^2}{2\sigma^2} \right\}$$

Same Dimensionality, Different Prior Volume

$$h_1 = af \quad a \in (0, 2)$$

$$h_2 = af^{1.5} \quad a \in (0, 2\kappa)$$

