



Exotic Glue in the Nucleus?

Double Helicity Flip Gluon Operators from Lattice QCD

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Collaborator: Will Detmold

April 27, 2016

Outline

- 1 Motivation
- 2 Double Helicity Flip Gluon Structure Function: $\Delta(x, Q^2)$
- 3 Lattice Study
- 4 Preliminary Results: ϕ meson
- 5 Question for Discussion
- 6 Summary

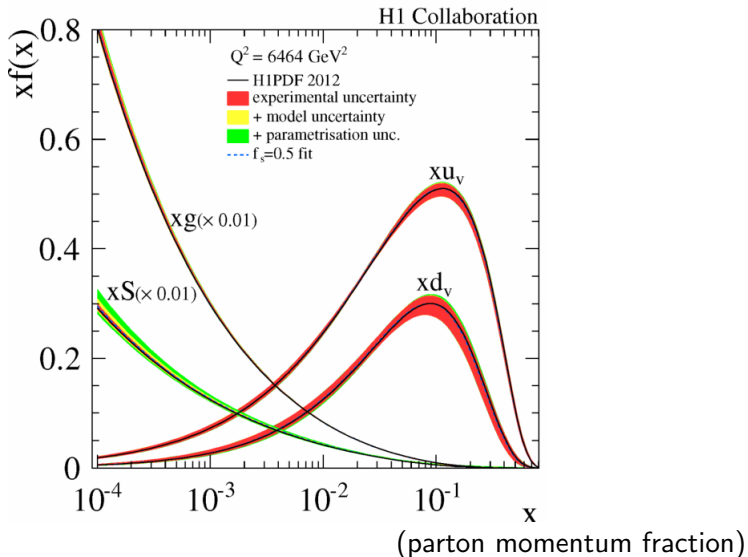
Motivation

Understanding gluons in hadron and nuclear structure is

- Important
 - ▶ e.g., Dominance of gluon PDF at low x

Gluons are Important in Hadron Structure

Parton distribution function in the proton



Motivation

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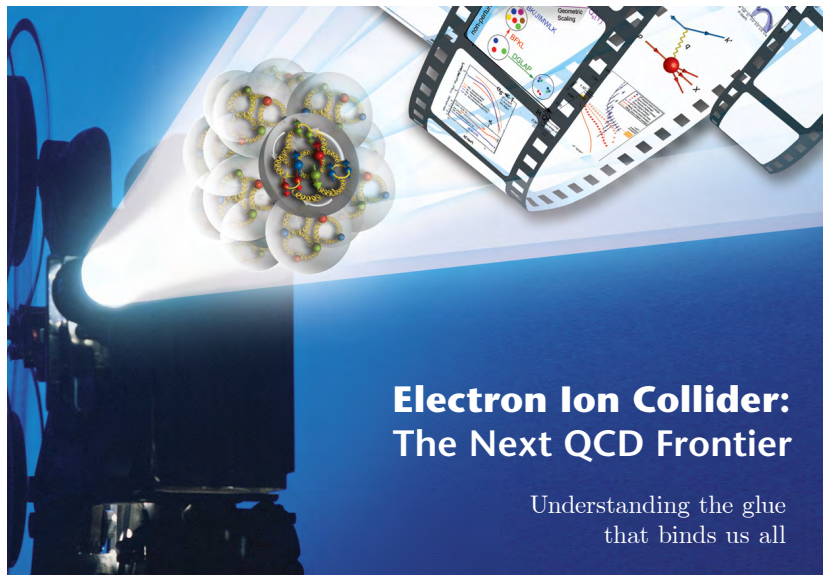
- Important
 - ▶ e.g., Dominance of gluon PDF at low x
- Hard
 - ▶ No direct expt. measurement of glue in a nucleus (yet)
 - ▶ Gluon probed only indirectly in electron scattering from hadrons/nuclei (does not couple to photon)
 - ▶ Drell-Yan more direct but messy

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- Topical
 - ▶ Electron-Ion Collider
 - ▶ JLab 12 GeV (lesser extent)

Significant Experimental Progress Expected



**Electron Ion Collider:
The Next QCD Frontier**

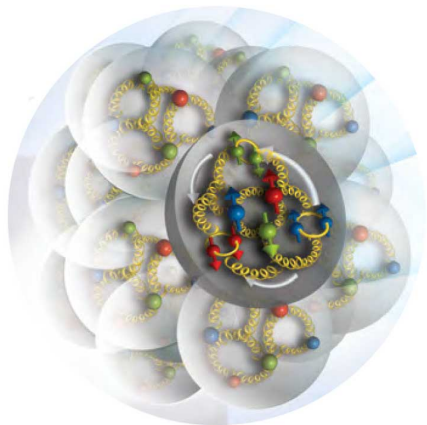
Understanding the glue
that binds us all

Motivation

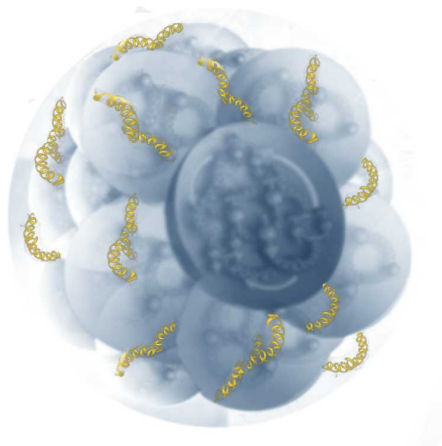
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'Exotic' Glue in the Nucleus



'Exotic' Glue in the Nucleus



'Exotic' Glue

Contributions to gluon observables that are not from nucleon degrees of freedom.

Exotic glue operator:
operator in nucleon = 0
operator in nucleus $\neq 0$

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Double Helicity Flip Gluon Structure Function: $\Delta(x, Q^2)$

Jaffe and Manohar (1989)

Leading-twist, double-helicity-flipping structure function $\Delta(x, Q^2)$ sensitive to exotic glue in the nucleus

- Clear signature for exotic glue in nuclei with spin ≥ 1 :
NO analogous twist-2 quark PDF \rightarrow unambiguous
- Experimentally measurable (JLab LOI 2016)
- Moments are calculable on the lattice

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First Lattice Study

- First moment of $\Delta(x, Q^2)$ in spin-1 ϕ (or ρ) meson

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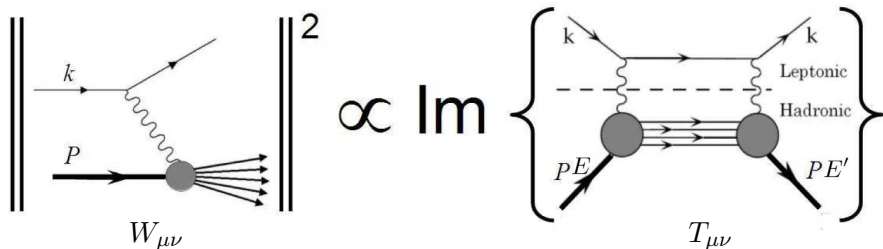
PRELIMINARY

Double Helicity Flip Gluon Structure Function: $\Delta(x, Q^2)$

Hadronic tensor for inelastic lepton scattering from a polarized spin-one target:

$$\frac{1}{4\pi} \int d^4x e^{iq \cdot x} \langle p, E' | [j_\mu(x), j_\nu(0)] | p, E \rangle$$

Optical theorem relates squared amplitude of DIS process to imaginary part of forward scattering amplitude:

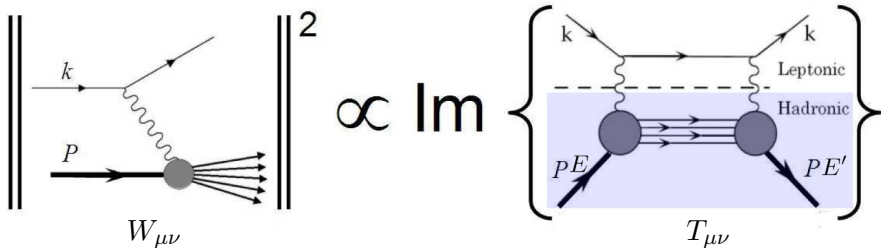


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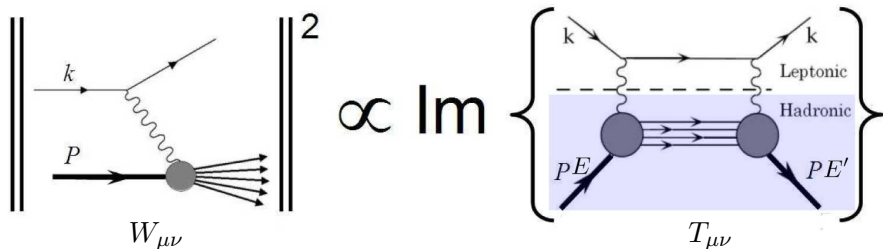
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Target Polarization Vector
↙ ↘

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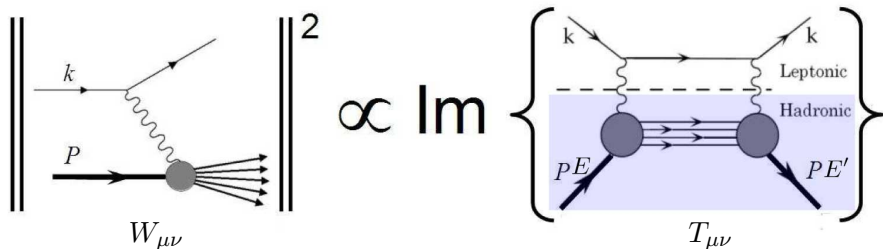
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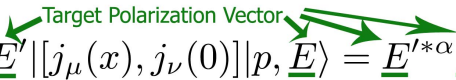
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Helicity projection:

$$W_{\mu\nu, \alpha\beta} = \sum_{hH, h'H'} P(hH, h'H')_{\mu\nu, \alpha\beta} A_{hH, h'H'}$$

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Helicity projection operators
Target helicity

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Double helicity flip amplitude:

$$\Delta(x, Q^2) = A_{+-,-+} = A_{-+,-+}$$

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Double helicity flip piece of structure function:

$$W_{\mu\nu, \alpha\beta}^{\Delta=2} = \frac{1}{2} \left\{ \left[\left(E_\mu^{*\prime} - \frac{q \cdot E^{*\prime}}{\kappa\nu} \left(p_\mu - \frac{M^2}{\nu} q_\mu \right) \right) \left(E_\nu - \frac{q \cdot E}{\kappa\nu} \left(p_\nu - \frac{M^2}{\nu} q_\nu \right) \right) + (\mu \leftrightarrow \nu) \right] \right. \\ \left. - \left[g_{\mu\nu} - \frac{q_\mu q_\nu}{q^2} + \frac{q^2}{\kappa\nu^2} \left(p_\mu - \frac{\nu}{q^2} q_\mu \right) \left(p_\nu - \frac{\nu}{q^2} q_\nu \right) \right] \left[E^{*\prime} \cdot E + \frac{M^2}{\kappa\nu^2} q \cdot E^{*\prime} q \cdot E \right] \right\} \Delta(x, Q^2)$$

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Double Helicity Flip Structure Function

Double Helicity Flip Gluon Structure Function: $\Delta(x, Q^2)$

Relate to matrix elements of operator (next slide) using

Operator Product Expansion

$$\mathcal{T}_{\mu\nu}(q) \equiv i \int d^4x e^{iq \cdot x} T(j_\mu(x) j_\nu(0))$$

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Showing only double helicity flip part: tower of gluonic operators

$$\frac{1}{2} \mathcal{T}_{\mu\nu}(q) = \dots + \sum_{n=2,4,\dots} \frac{2^n q^{\mu_1} \dots q^{\mu_n}}{(Q^2)^n} C_n(Q^2) \mathcal{O}_{\mu\nu\mu_1 \dots \mu_n}$$

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where

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Symmetrize and trace subtract in μ_1, \dots, μ_n

$$\mathcal{O}_{\mu\nu\mu_1\dots\mu_n} = \underline{S} \left[G_{\mu\mu_1} \overleftrightarrow{D}_{\mu_3} \dots \overleftrightarrow{D}_{\mu_n} G_{\nu\mu_2} \right]$$

Double Helicity Flip Gluon Structure Function: $\Delta(x, Q^2)$

$$\begin{aligned} & \langle pE' | \mathcal{O}_{\mu\nu\{\mu_1\dots\mu_n\}} - \text{Tr}[pE] \rangle \\ &= (-2i)^{n-2} S [(p_\mu E'_{\mu_1}{}^* - p_{\mu_1} E'_\mu{}^*)(p_\nu E_{\mu_2} - p_{\mu_2} E_\nu) \\ & \quad + (\mu \leftrightarrow \nu)] p_{\mu_3} \dots p_{\mu_n} A_n(Q^2) \dots, \end{aligned}$$

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Dispersion relation for helicity flip part of $\mathcal{T}_{\mu\nu}$ (previous slide) and analytic continuation give **moments**:

$$\int_0^1 dx x^{n-1} \Delta(x, Q^2) = \frac{\alpha_s(Q^2)}{3\pi} \frac{A_n(Q^2)}{n+2}, \quad n = 2, 4, 6, \dots,$$

Unpolarized scattering: symmetric piece of hadronic tensor $W_{\mu\nu}$, \rightarrow even n

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Double Helicity Flip Gluon Structure Function: $\Delta(x, Q^2)$

Parton model interpretation

For a target in the infinite momentum frame polarized in the \hat{x} direction perpendicular to its momentum,

$$\Delta(x, Q^2) \propto \int_x^1 \frac{dy}{y^3} (g_{\hat{x}}(y, Q^2) - g_{\hat{y}}(y, Q^2))$$

$g_{\hat{x}, \hat{y}}(y, Q^2)$: probability of finding a gluon with momentum fraction y linearly polarized in the \hat{x} , \hat{y} direction

“How much more momentum of transversely polarized particle carried by gluons aligned rather than perpendicular to it in the transverse plane”

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Lattice Operators

First moment of $\Delta(x, Q^2)$



Matrix elt. of $\mathcal{O}_{\mu\nu\mu_1\mu_2} = S[G_{\mu\mu_1}G_{\nu\mu_2}]$

Lattice Operators

First moment of $\Delta(x, Q^2)$



Matrix elt. of $\mathcal{O}_{\mu\nu\mu_1\mu_2} = S[G_{\mu\mu_1}G_{\nu\mu_2}]$

- Relate $\mathcal{O}_{\mu\nu\mu_1\mu_2}$ to Euclidean operator
- Find linear combs. of Euclidean operator (with different indices) that
 - 1 Transform irriducibly under appropriate representations of $H(4)$
 - 2 Don't mix with same or lower-dimensional quark or gluon operators
 - ▶ 3 irreps. of dimension 2, 6, 2, i.e., 10 basis vectors
- Lattice simulation of matrix element in ϕ meson (spin-1)

Lattice Details

Luscher-Weisz gauge action with a clover-improved quark action with one level of stout link smearing

L/a	T/a	β	am_l	am_s
24	64	6.1	-0.2800	-0.2450
a (fm)	L (fm)	T (fm)	m_π (MeV)	m_K (MeV)
0.1167(16)	2.801(29)	7.469(77)	450(5)	596(6)
m_ϕ (MeV)	$m_\pi L$	$m_\pi T$	N_{cfg}	N_{src}
1040(3)	6.390	17.04	1042	10^5

- All ϕ polarization states ($\{1, 2, 3\}$ or $\{+, -, 0\}$)
 - ▶ on-diagonal
 - ▶ off-diagonal
- Momenta up to (1,1,1) in lattice units
- HYP smearing of gauge fields in operator (2-5 steps)

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- Momenta up to (1,1,1) in lattice units
- HYP smearing of gauge fields in operator (2-5 steps)

Lattice Details

Luscher-Weisz gauge action with a clover-improved quark action with one level of stout link smearing

L/a	T/a	β	am_l	am_s
24	64	6.1	-0.2800	-0.2450
a (fm)	L (fm)	T (fm)	m_π (MeV)	m_K (MeV)
0.1167(16)	2.801(29)	7.469(77)	450(5)	596(6)
m_ϕ (MeV)	$m_\pi L$	$m_\pi T$	N_{cfg}	N_{src}
1040(3)	6.390	17.04	1042	10^5

- All ϕ polarization states ($\{1, 2, 3\}$ or $\{+, -, 0\}$)
 - ▶ on-diagonal
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Luscher-Weisz gauge action with a clover-improved quark action with one level of stout link smearing

L/a	SYSTEMATICS IGNORED			am_s
24	• Quark mass effects			-0.2450
a (fm)	• Volume effects			m_K (MeV)
0.1167(16)	• Discretization effects			596(6)
m_ϕ (MeV)	• Renormalization (for now)			N_{src}
1040(3)	6.390	17.04	1042	10^5

- All ϕ polarization states ($\{1, 2, 3\}$ or $\{+, -, 0\}$)
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Extraction of A_2

$$\begin{aligned} & \langle pE' | \mathcal{O}_{\mu\nu\{\mu_1\dots\mu_n\}} - \text{Tr}|pE \rangle \quad \text{Symmetrize and trace subtract in } \mu_1, \dots, \mu_n \\ & = (-2i)^{n-2} \underline{S} \left[(p_\mu E'_{\mu_1} - p_{\mu_1} E'_\mu) (p_\nu E_{\mu_2} - p_{\mu_2} E_\nu) \right. \\ & \quad \left. + (\mu \leftrightarrow \nu) \right] p_{\mu_3} \dots p_{\mu_n} \boxed{A_n(Q^2)} \dots, \\ & \qquad \qquad \qquad \text{Reduced Matrix Element} \end{aligned}$$

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 & \quad \left. + (\mu \leftrightarrow \nu) \right] p_{\mu_3} \dots p_{\mu_n} \boxed{A_n(Q^2)} \dots, \\
 & \hspace{15em} \text{Reduced Matrix Element}
 \end{aligned}$$

$$\left[\frac{C_{3\text{pt}}^{EE'}}{C_{2\text{pt}}^{EE'}} \right] (t_{\text{sink}}, \tau) \propto \boxed{A_2}, \quad 0 \ll \tau \ll t_{\text{sink}}$$

factors of m and p

Extraction of A_2 : 3pt/2pt ratio

Some choice of irrep. and basis vector

ϕ momentum $(0, 0, 0)$

$$\begin{array}{c} \rho_0 \\ \rho_+ \\ \rho_- \end{array} \begin{pmatrix} \rho_0 & \rho_+ & \rho_- \\ \frac{2m^2 A_2}{\sqrt{3}} & 0 & 0 \\ 0 & -\frac{m^2 A_2}{\sqrt{3}} & 0 \\ 0 & 0 & -\frac{m^2 A_2}{\sqrt{3}} \end{pmatrix}$$

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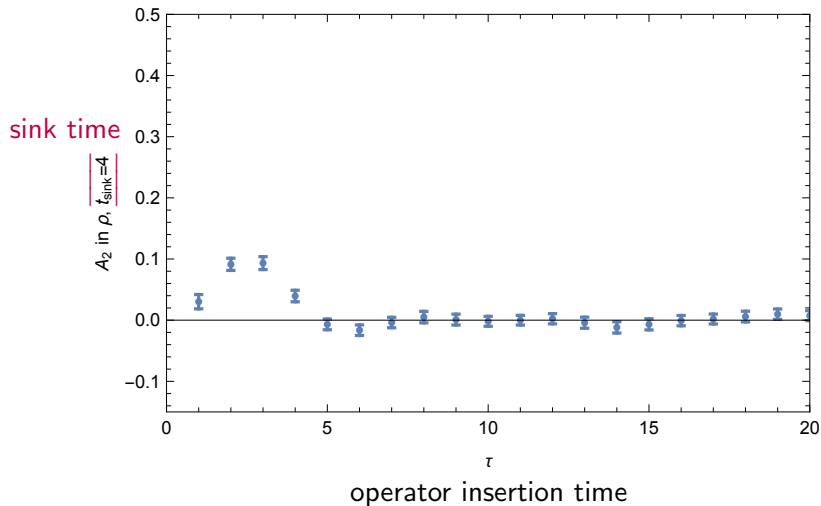
Extraction of A_2 : 3pt/2pt ratio

Some choice of irrep. and basis vector

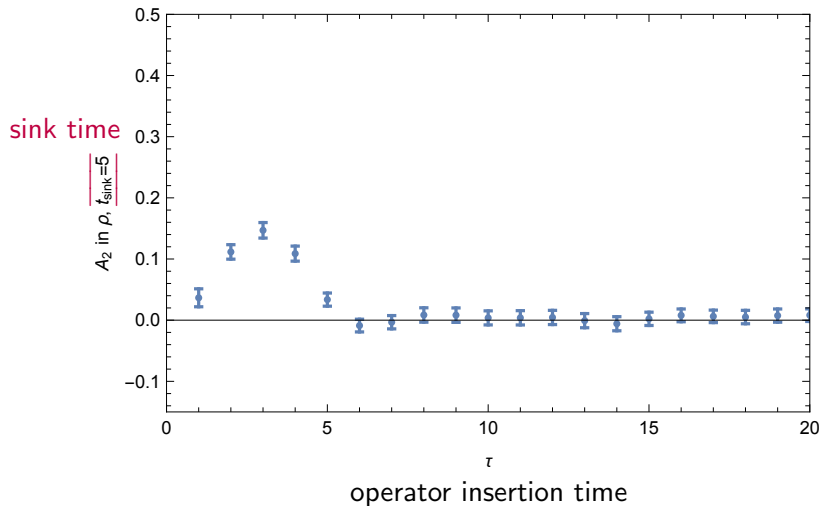
ϕ momentum (p, p, p)

$$\begin{array}{c}
 \rho_0 \\
 \rho_+ \\
 \rho_-
 \end{array}
 \left(
 \begin{array}{ccc}
 \rho_0 & \rho_+ & \rho_- \\
 \frac{2(m^3 + \sqrt{m^2 + 3p^2}m^2 + 4p^2m + 2p^2\sqrt{m^2 + 3p^2})A_2}{\sqrt{3}(m + \sqrt{m^2 + 3p^2})} & \frac{(1-i)p^2(m + 2\sqrt{m^2 + 3p^2})A_2}{\sqrt{6}(m + \sqrt{m^2 + 3p^2})} & \frac{(1+i)p^2(m + 2\sqrt{m^2 + 3p^2})A_2}{\sqrt{6}(m + \sqrt{m^2 + 3p^2})} \\
 \frac{(1+i)p^2(m + 2\sqrt{m^2 + 3p^2})A_2}{\sqrt{6}(m + \sqrt{m^2 + 3p^2})} & -\frac{(m^3 + \sqrt{m^2 + 3p^2}m^2 + 4p^2m + 2p^2\sqrt{m^2 + 3p^2})A_2}{\sqrt{3}(m + \sqrt{m^2 + 3p^2})} & -\frac{2ip^2(m + 2\sqrt{m^2 + 3p^2})A_2}{\sqrt{3}(m + \sqrt{m^2 + 3p^2})} \\
 \frac{(1-i)p^2(m + 2\sqrt{m^2 + 3p^2})A_2}{\sqrt{6}(m + \sqrt{m^2 + 3p^2})} & \frac{2ip^2(m + 2\sqrt{m^2 + 3p^2})A_2}{\sqrt{3}(m + \sqrt{m^2 + 3p^2})} & \frac{(m^3 + \sqrt{m^2 + 3p^2}m^2 + 4p^2m + 2p^2\sqrt{m^2 + 3p^2})A_2}{\sqrt{3}(m + \sqrt{m^2 + 3p^2})}
 \end{array}
 \right)$$

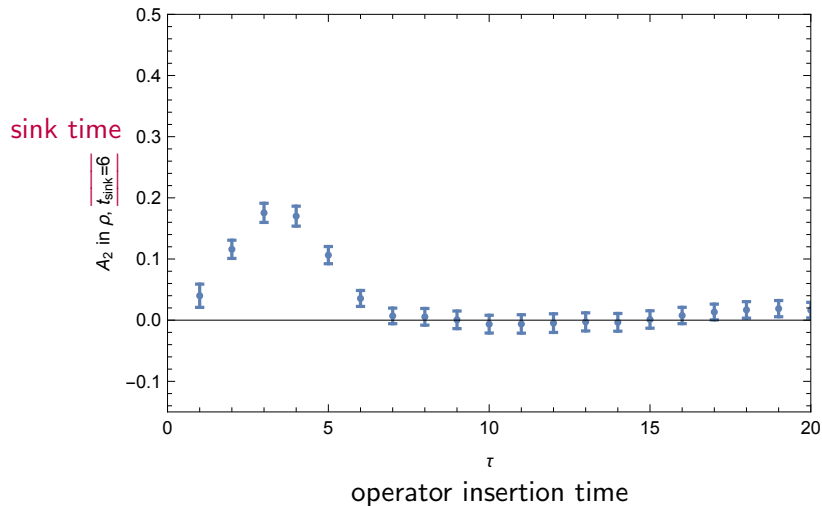
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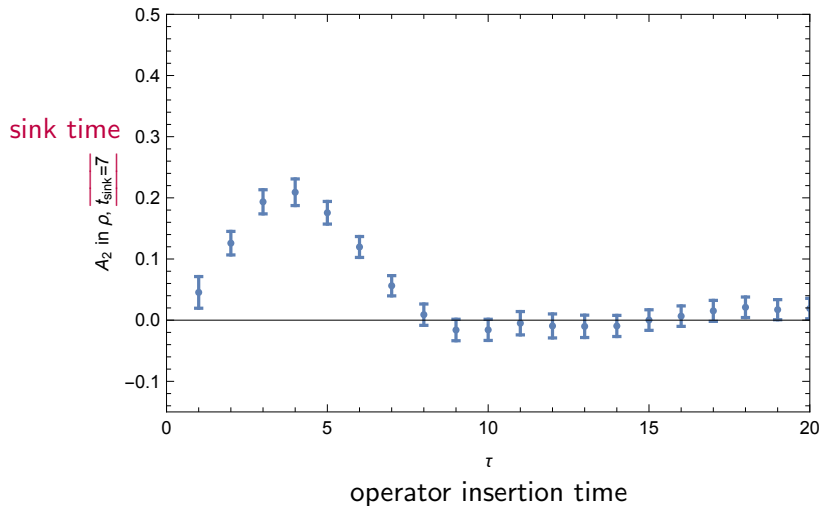
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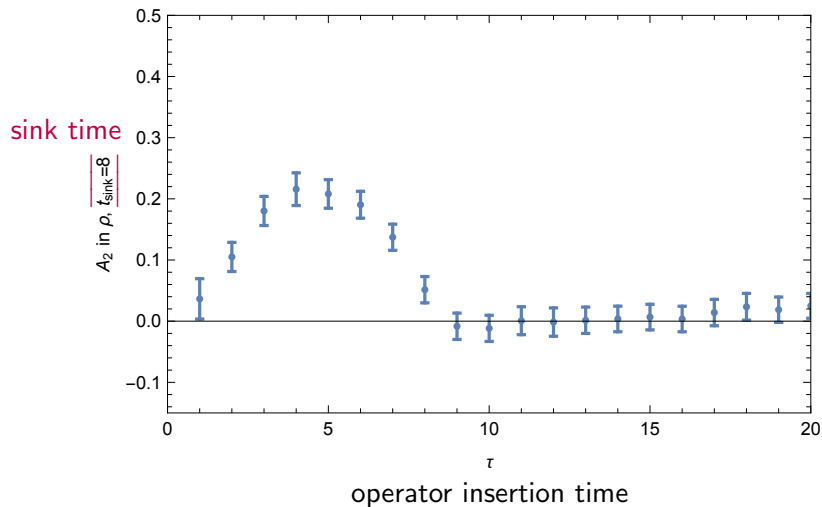
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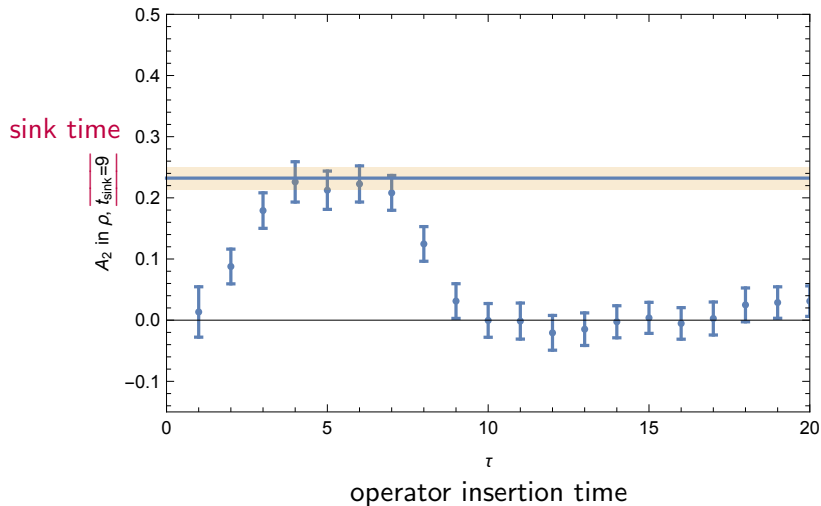
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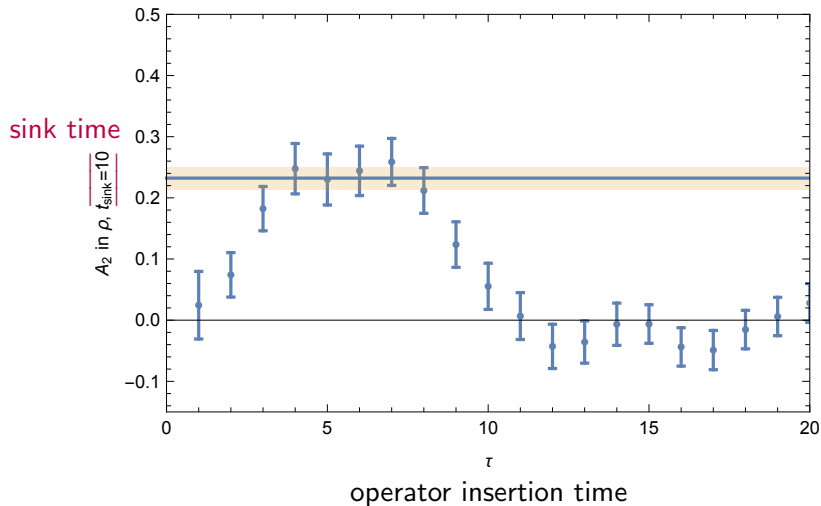
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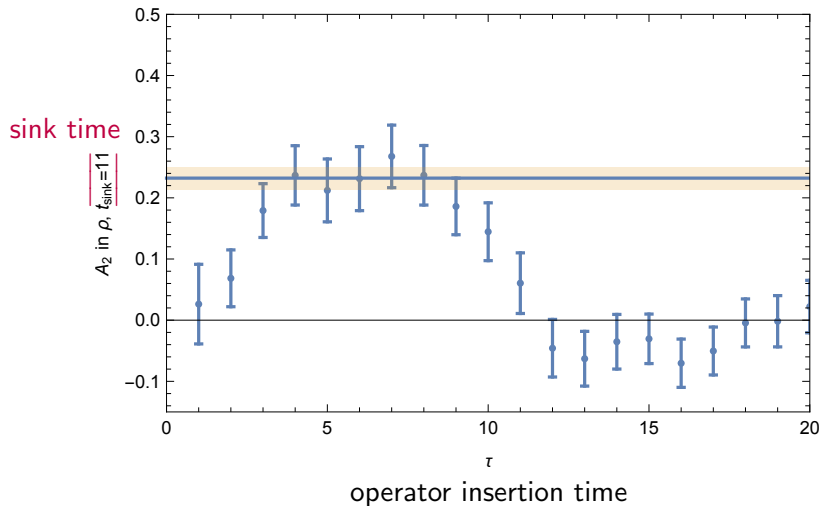
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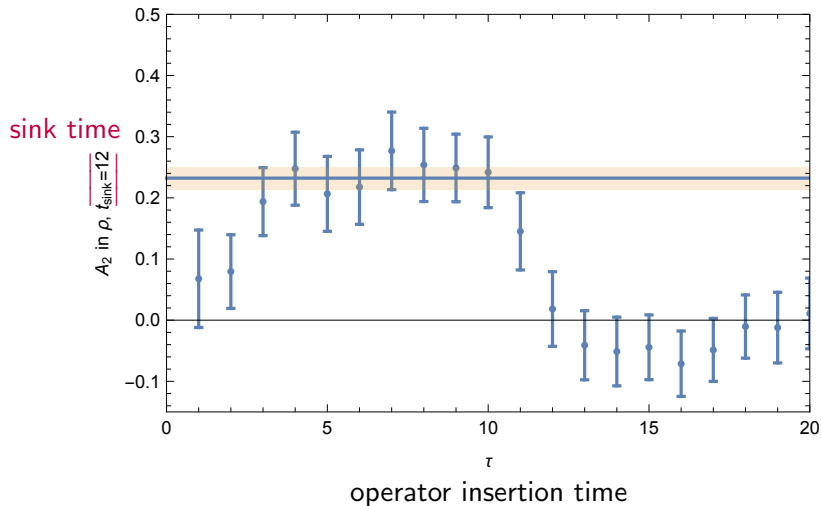
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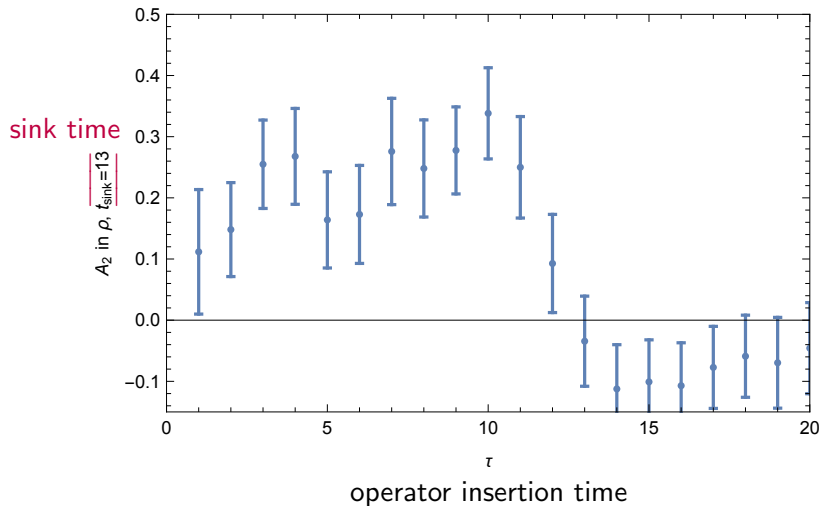
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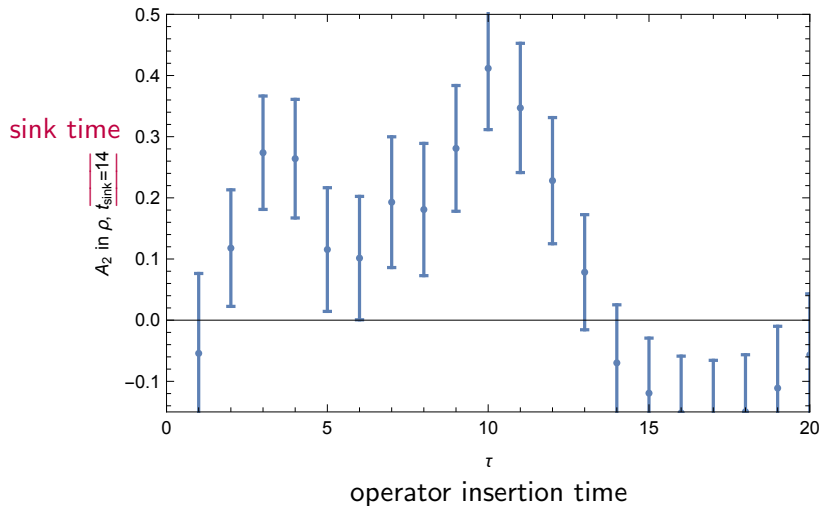
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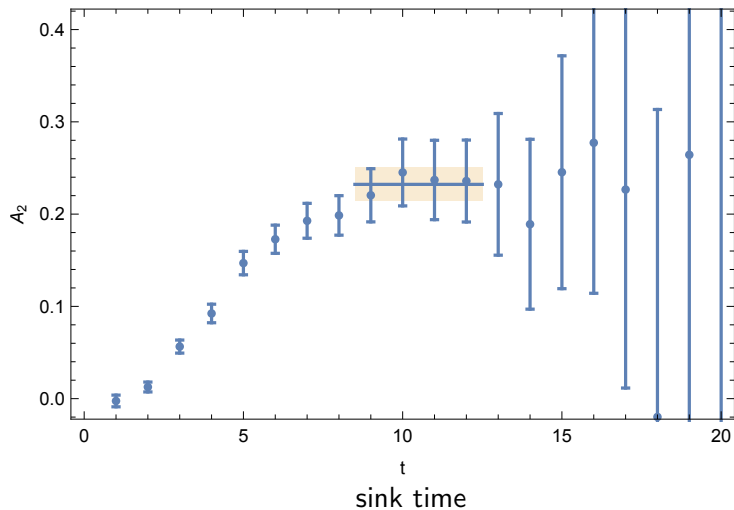
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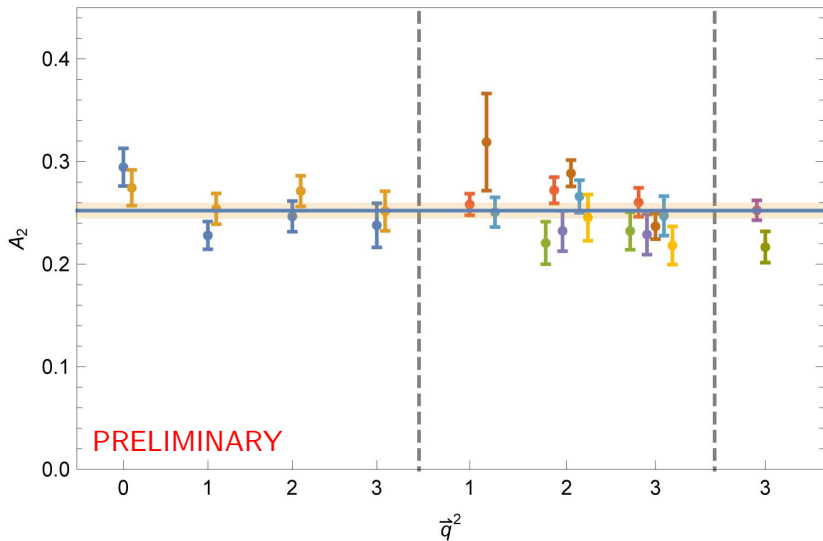
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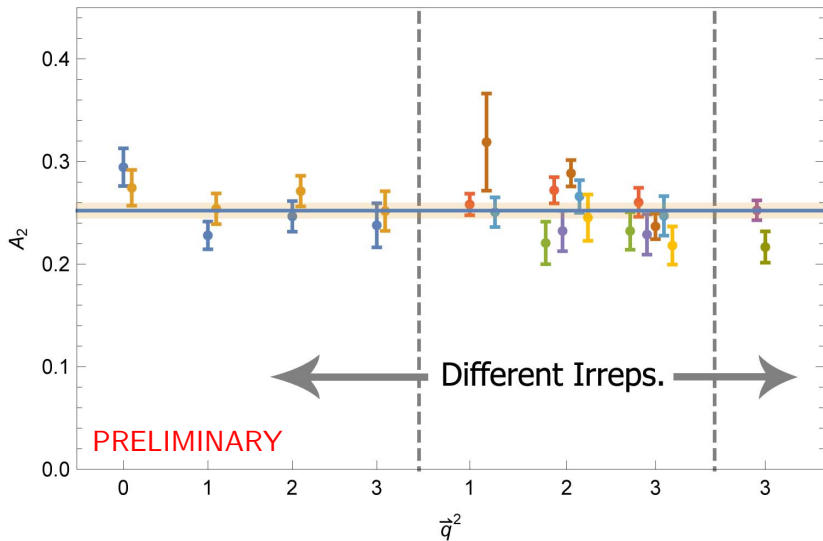
Outline

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- 3 Lattice Study
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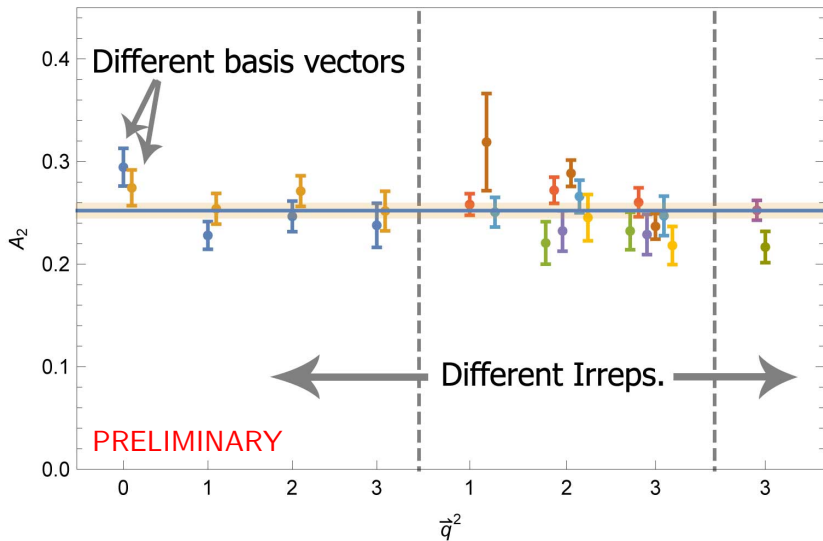
UNRENORMALISED reduced matrix element: ϕ meson



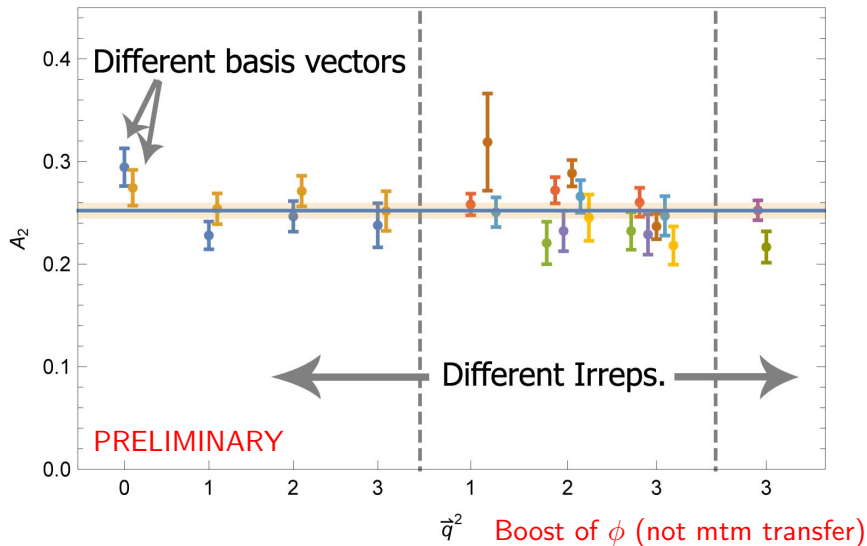
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CLEAR NON-ZERO SIGNAL

Proof of principle: similar signal in a nucleus \Leftrightarrow exotic glue

CLEAR NON-ZERO SIGNAL

Proof of principle: similar signal in a nucleus \Leftrightarrow exotic glue

SYSTEMATICS IGNORED

- Quark mass effects
- Volume effects
- Discretization effects
- Renormalization

Renormalisation:

- Gradient flow instead of HYP smearing
- Gradient-flowed lattice perturbation theory
- Non-perturbative renormalisation
(set gauge-fixed gluon matrix elt. = tree level value)

Soffer bound analogue

Explore gluon structure of ϕ meson more generally

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Soffer bound for transversity quark distributions:

$$|\delta q(x)| \leq \frac{1}{2} (q(x) + \Delta q(x))$$

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Soffer bound for transversity quark distributions:

$$|\delta q(x)| \leq \frac{1}{2} (q(x) + \Delta q(x))$$

The equation features two annotations: a pink arrow pointing to $\delta q(x)$ labeled "Transversity" and a green arrow pointing to $q(x)$ labeled "Spin-independent".

Soffer bound analogue

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Diagram annotations:

- A pink arrow points from the word "Transversity" to the $\delta q(x)$ term.
- A green arrow points from the word "Spin-independent" to the $q(x)$ term.
- A blue arrow points from the word "Spin-dependent" to the $\Delta q(x)$ term.

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Annotations in the image:
- A pink arrow points to $\delta q(x)$ with the label "Transversity".
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Direct analogue for leading moments of gluon distributions:

$$|A_2| \leq \frac{1}{2} B_2$$

Soffer bound analogue

Explore gluon structure of ϕ meson more generally

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Annotations:
 - Transversity (pink arrow pointing to $\delta q(x)$)
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Direct analogue for leading moments of gluon distributions:

$$G_{\mu\mu_1} G_{\nu\nu_2} \boxed{A_2} \leq \frac{1}{2} B_2$$

Annotation:
 - $G_{\mu\mu_1} G_{\nu\nu_2}$ (pink arrow pointing to $\boxed{A_2}$)

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Direct analogue for leading moments of gluon distributions:

$$G_{\mu\mu_1} G_{\nu\nu_2} \boxed{A_2} \leq \frac{1}{2} \boxed{B_2} G_{\mu_1\alpha} G_{\mu_2}^{\alpha}$$

Annotations:
 - $G_{\mu\mu_1} G_{\nu\nu_2}$ (pink) points to $\boxed{A_2}$ (pink box)
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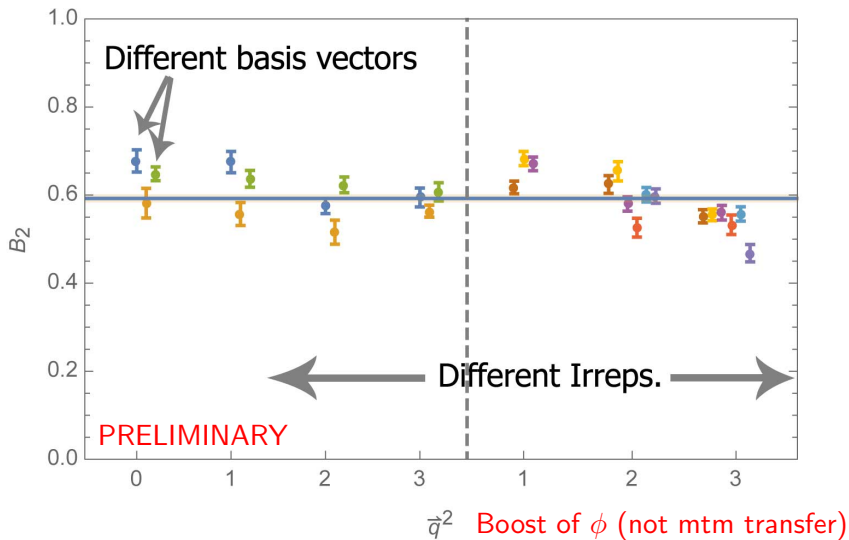
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Direct analogue for leading moments of gluon distributions:

$$G_{\mu\mu_1} G_{\nu\nu_2} \boxed{A_2} \leq \frac{1}{2} \boxed{B_2} \quad \tilde{G}_{\mu_1\alpha} G_{\mu_2}^\alpha \rightarrow 0$$

Annotations: $G_{\mu\mu_1} G_{\nu\nu_2}$ (pink), $\boxed{A_2}$ (pink box), $G_{\mu_1\alpha} G_{\mu_2}^\alpha$ (green), $\boxed{B_2}$ (green box), $\tilde{G}_{\mu_1\alpha} G_{\mu_2}^\alpha \rightarrow 0$ (blue)

UNRENORMALISED reduced matrix element: ϕ meson



Soffer bound analogue

If we imagine approx. the same renormalisation for A_2 and B_2 :

$$G_{\mu\mu_1} G_{\nu\nu_2} \downarrow \boxed{A_2} \leq \frac{1}{2} \boxed{B_2} \uparrow G_{\mu_1\alpha} G_{\mu_2}{}^\alpha$$

$$0.25 \leq \frac{1}{2} 0.6$$

First two moments of quark distributions: Soffer bound saturated to 80%
(lattice QCD, Diehl *et al.* 2005)

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Given lattice results for light nuclei
 \Rightarrow extrapolate to heavier nuclei?

Recall:

- Leading twist, double helicity flip
- Zero in spin < 1 states (nucleons, pions)
i.e., gluon contributions from non-nucleonic and pionic degrees of freedom

Experimental possibilities:

- JLab 2016 proposal: **Nitrogen**
- EIC: possibly deuteron and other lighter nuclei!

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in the ϕ (or ρ) meson

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e.g., Soffer bound analogue

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BUT: SYSTEMATICS IGNORED
 \Rightarrow no physically meaningful number (yet)

UNRENORMALISED reduced matrix element: ϕ meson, $n\text{HYP}=2$

