

Applying Twisted Boundary Conditions for Few-Body Nuclear Systems

C. Körber & T. Luu

12th, April 2016 | Christopher Körber | PhD Student | Forschungszentrum Jülich & University of Bonn
INT — Nuclear Physics from lattice QCD

Outline

- ▶ ***Motivation***
- ▶ ***Implementation***
- ▶ ***Explanation***
- ▶ ***Methodology***
- ▶ ***Numerical Results***
 - ***2-Body***
 - ***3-Body***

Finite Volume (FV)

As a tool for

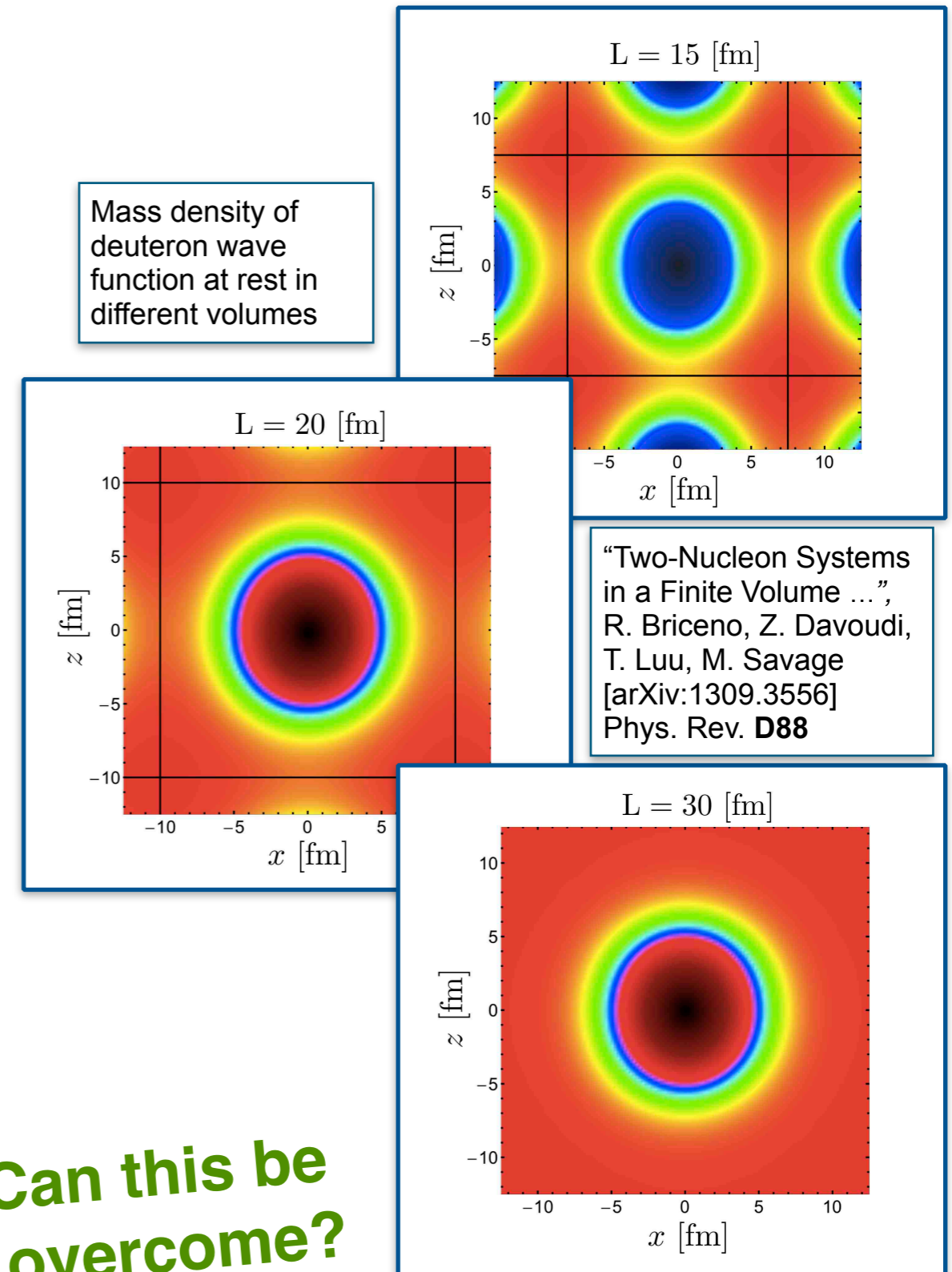
- ▶ **Non-perturbative physics**
Evaluation of path integral
- ▶ **Many-body physics**
Scalability of HMC algorithms
- ▶ **Physical finite systems**
E.g. carbon nanotubes

But...

- ▶ **Systematic “errors”...**
... since observables depend on the volume (L)
- ▶ **Broken symmetries...**
... since rotation symmetry is reduced to ‘just’ cubic rotation symmetry

- ▶ **Multiple calculations...**
... at multiple volumes must be performed to extrapolate to infinite volume

Mass density of
deuteron wave
function at rest in
different volumes



“Two-Nucleon Systems
in a Finite Volume ...”,
R. Briceno, Z. Davoudi,
T. Luu, M. Savage
[arXiv:1309.3556]
Phys. Rev. **D88**

**Can this be
overcome?**

Twisted Boundaries

A 'knob' for reducing FV effects

"...Volume Dependence of Light Hadron Masses", NPLQCD [arXiv:1104.4101], Phys.Rev. **D84**

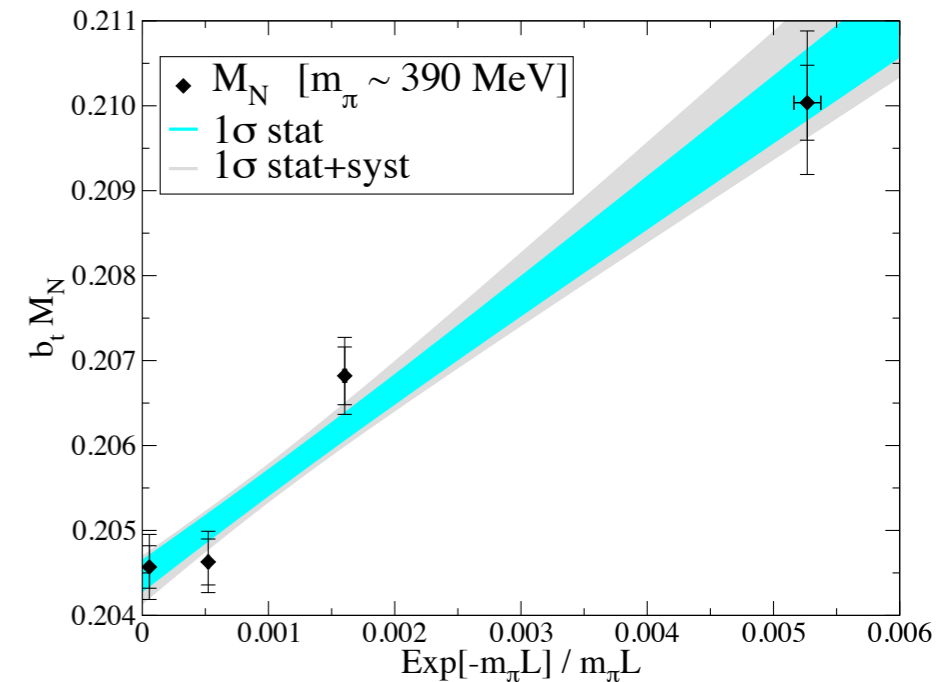
► **Reduce FV effects by increasing volume**

Bound states scale exponentially in volume...

$$\mathcal{O}(e^{-\lambda L})$$

... but the computational cost increases as well

$$\mathcal{O}(N_L^3 \times N_t)$$



Nucleon mass volume (L) dependence.

► **Change the boundary conditions**

Usually LQCD and NLEFT calculations utilize periodic boundary conditions (PBs)

$$\psi(\vec{r} + \vec{e}_i L) = \psi(\vec{r})$$

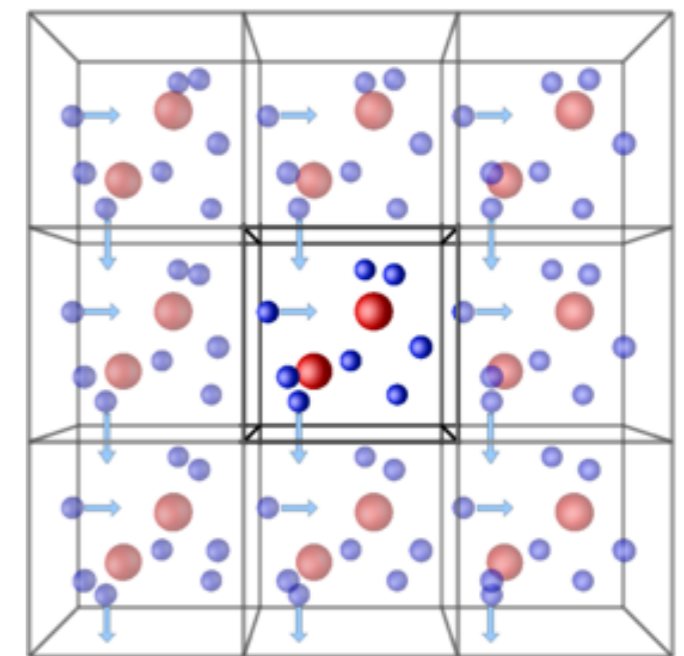
Anti-periodic boundary conditions

$$\psi(\vec{r} + \vec{e}_i L) = -\psi(\vec{r})$$

Twisted boundary conditions

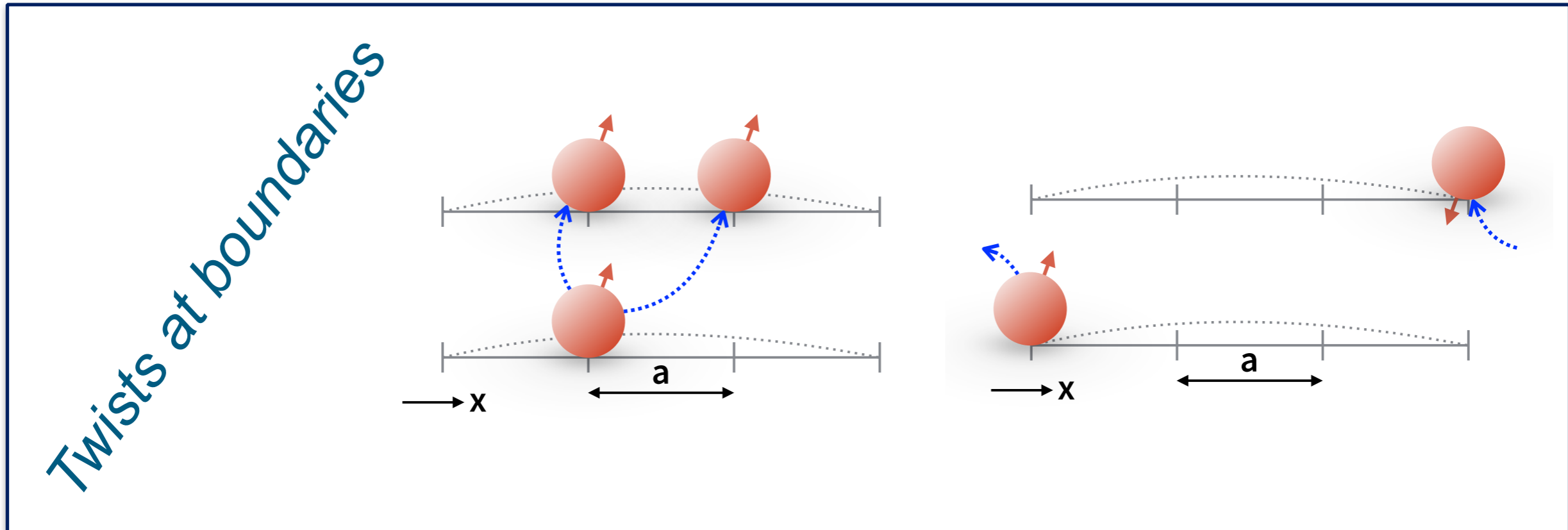
$$\psi(\vec{r} + \vec{e}_i L) = e^{i\vec{\phi} \cdot \vec{e}_i} \psi(\vec{r})$$

'i-Periodic':
 $e^{i\pi/2} = i$



[isaacs.sourceforge](http://isaacs.sourceforge.net)

TBCs — Implementation

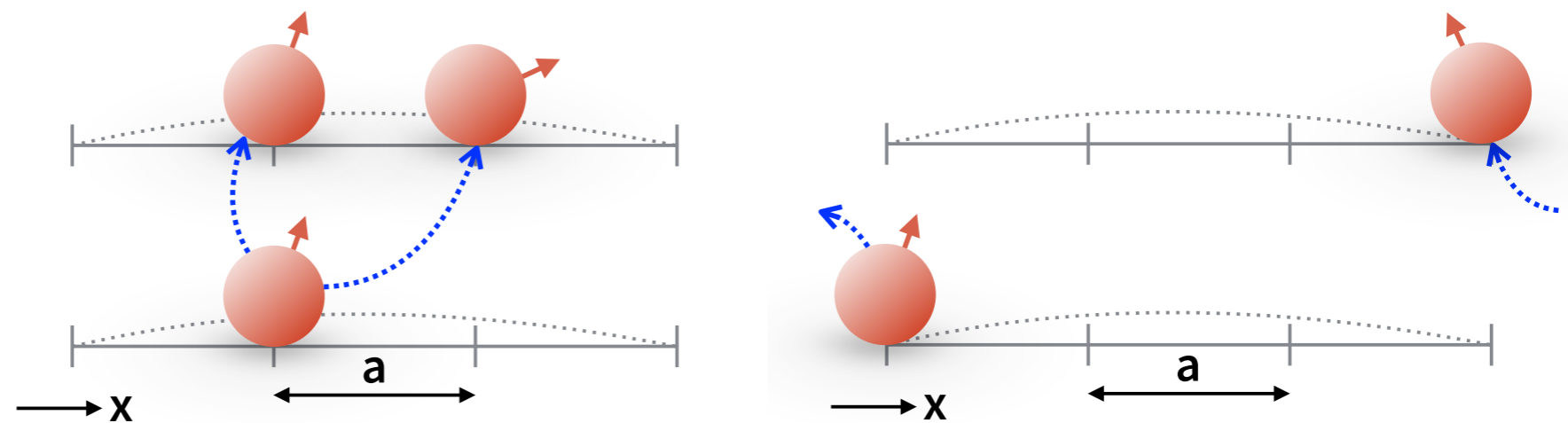


$$\sum_{\vec{n}' \in \mathbb{Z}^3} a^\dagger(\vec{n}' + \vec{l}) a(\vec{n}') |\vec{n}\rangle = \begin{cases} |\vec{n} + \vec{l}\rangle & , \quad \vec{n} + \vec{l} \in L^3 \\ |(\vec{n} + \vec{l})_{L^3}\rangle e^{-i\vec{\phi} \cdot \vec{l}} & , \quad \vec{n} + \vec{l} \notin L^3 \end{cases}$$

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Twist at each step



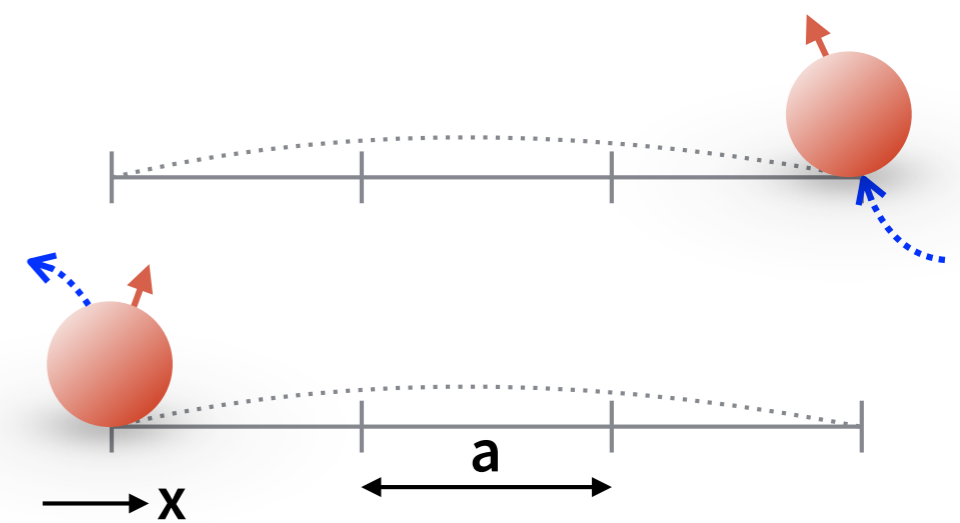
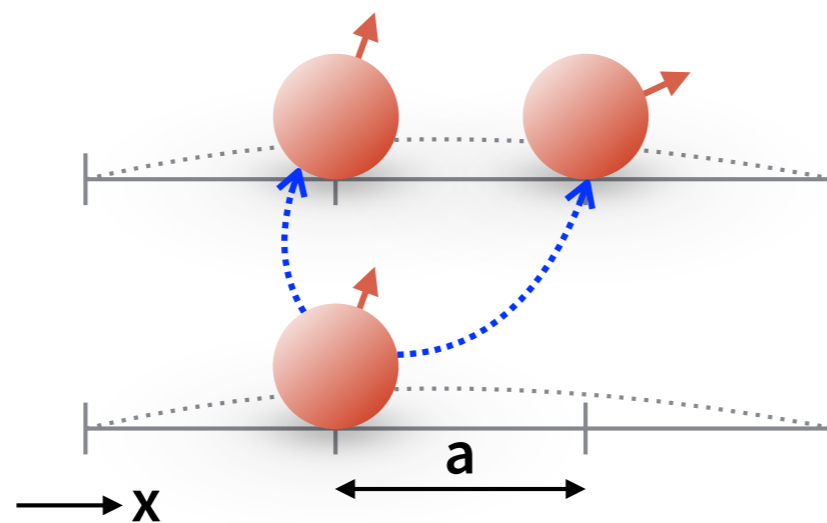
TBCs — Implementation

Hermitian operators stay hermitian

$$\langle \vec{m}_1^{\vec{\phi}_1}, \dots; \vec{m}_N^{\vec{\phi}_N} | O | \vec{n}_1^{\vec{\phi}_1}, \dots; \vec{n}_N^{\vec{\phi}_N} \rangle$$

$$= \langle \vec{m}_1^{\vec{0}}, \dots; \vec{m}_N^{\vec{0}} | O | \vec{n}_1^{\vec{0}}, \dots; \vec{n}_N^{\vec{0}} \rangle \exp \left(i \sum_{i=1}^N \vec{\phi}_i / N_L \cdot (\vec{n}_i - \vec{m}_i) \right)$$

Twist at each step



TBCs — Summary

- ▶ **Twists are implemented at each step**
Well defined momenta (translation invariance)

$$\vec{p} \mapsto \frac{2\pi}{L} \vec{n}_p + \frac{\vec{\phi}}{L} \quad \vec{n}_p \in \mathbb{Z}^3$$

- ▶ **Twists are connected to particle “hops”**
For each d.o.f. you could use a different twist

$$\vec{\phi}_1, \vec{\phi}_2, \dots, \vec{\phi}_N$$

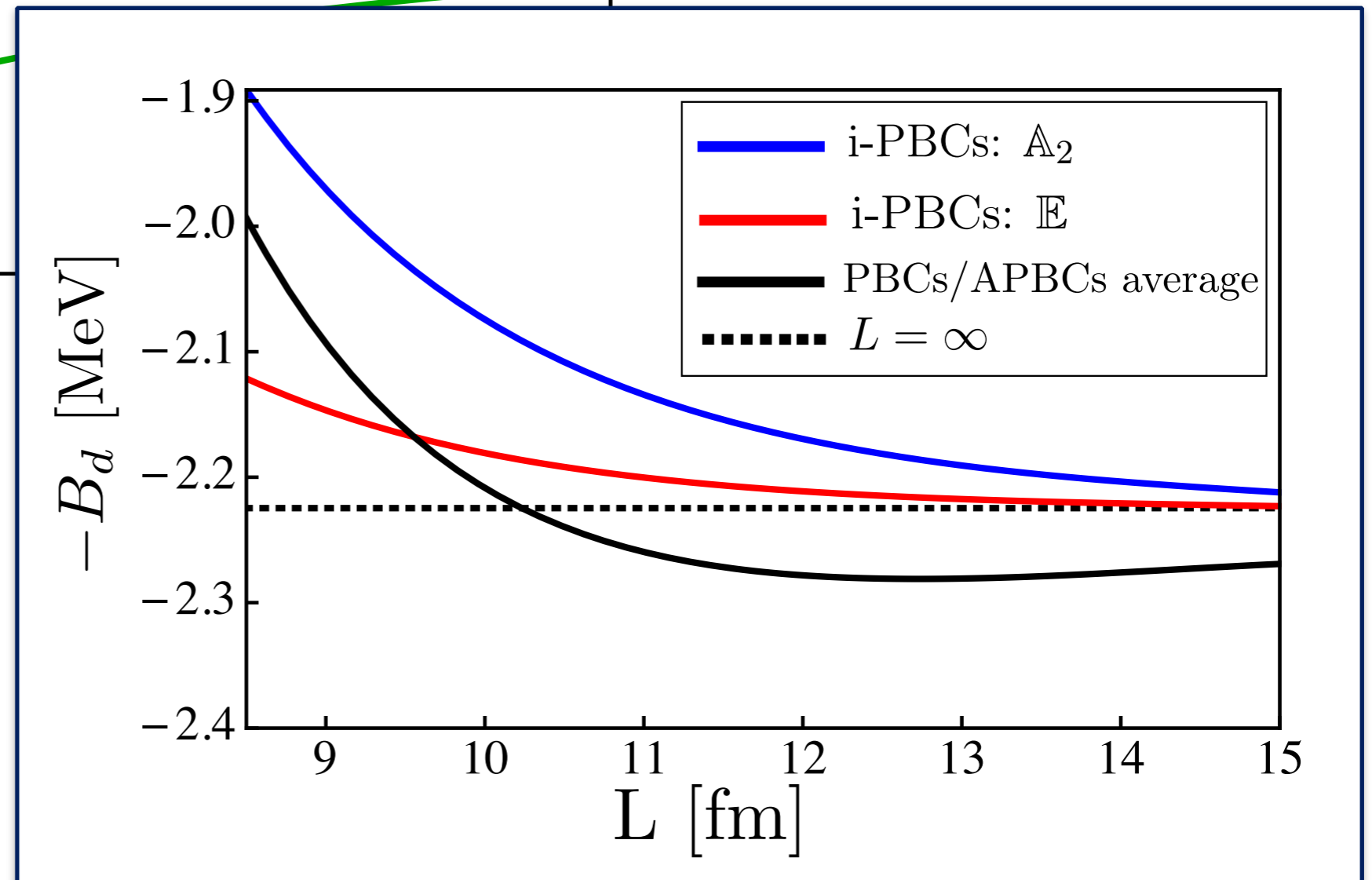
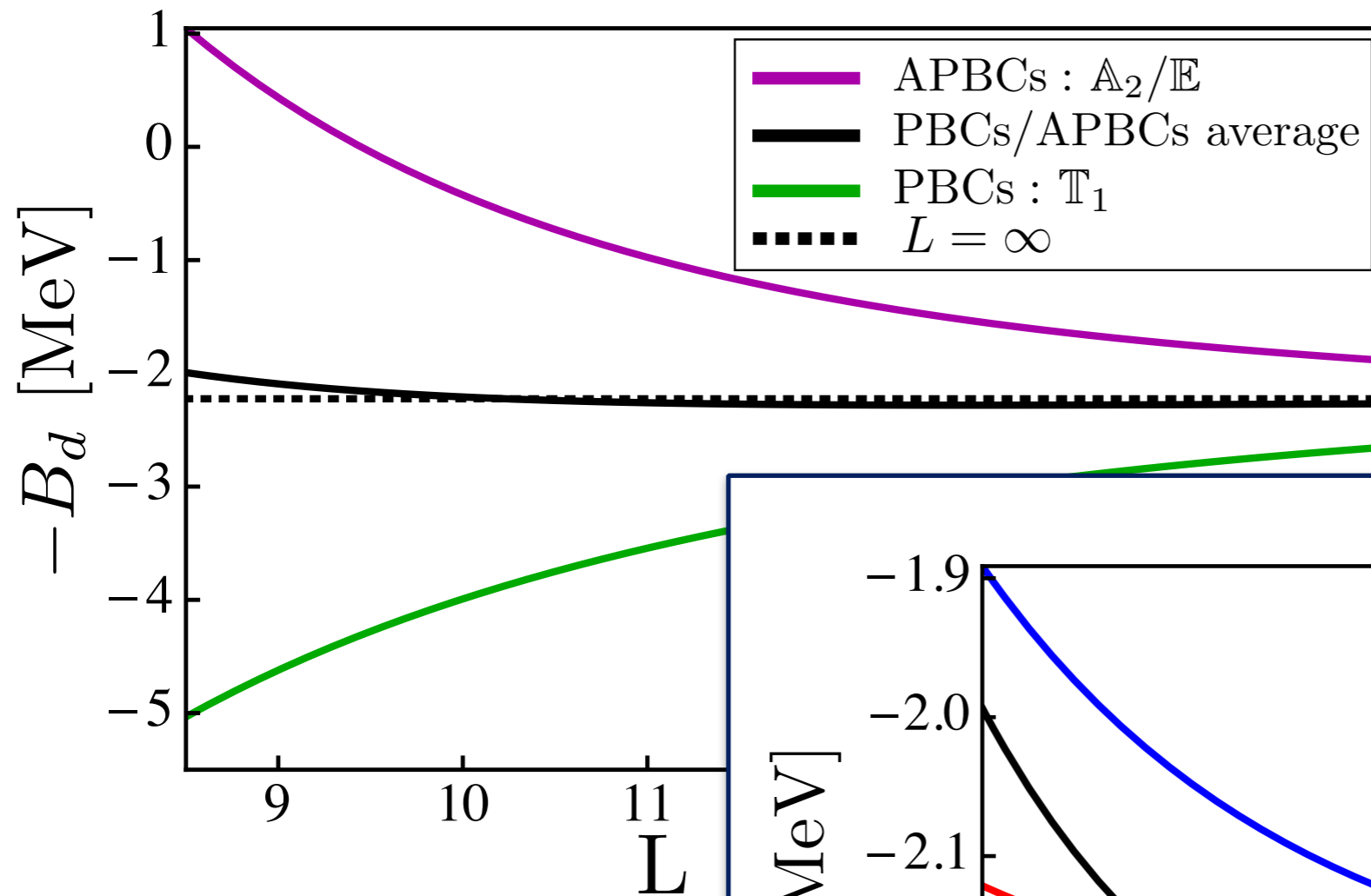
- ▶ **Twists can change CMS motion**
Constrained on twist conserve CMS motion

$$\sum_i \vec{\phi}_i = \vec{0}$$

Analytic Deuteron Results

Why should you use twisted boundaries?

Analytic Deuteron Results



“Lüscher’s Formula with a twist”,
 R. Briceno, Z. Davoudi, T. Luu, M. Savage
 [arXiv:1311.7686], Phys.Rev. **D89**

TBCs — How does this work?

For two-body systems

Schrödinger formalism

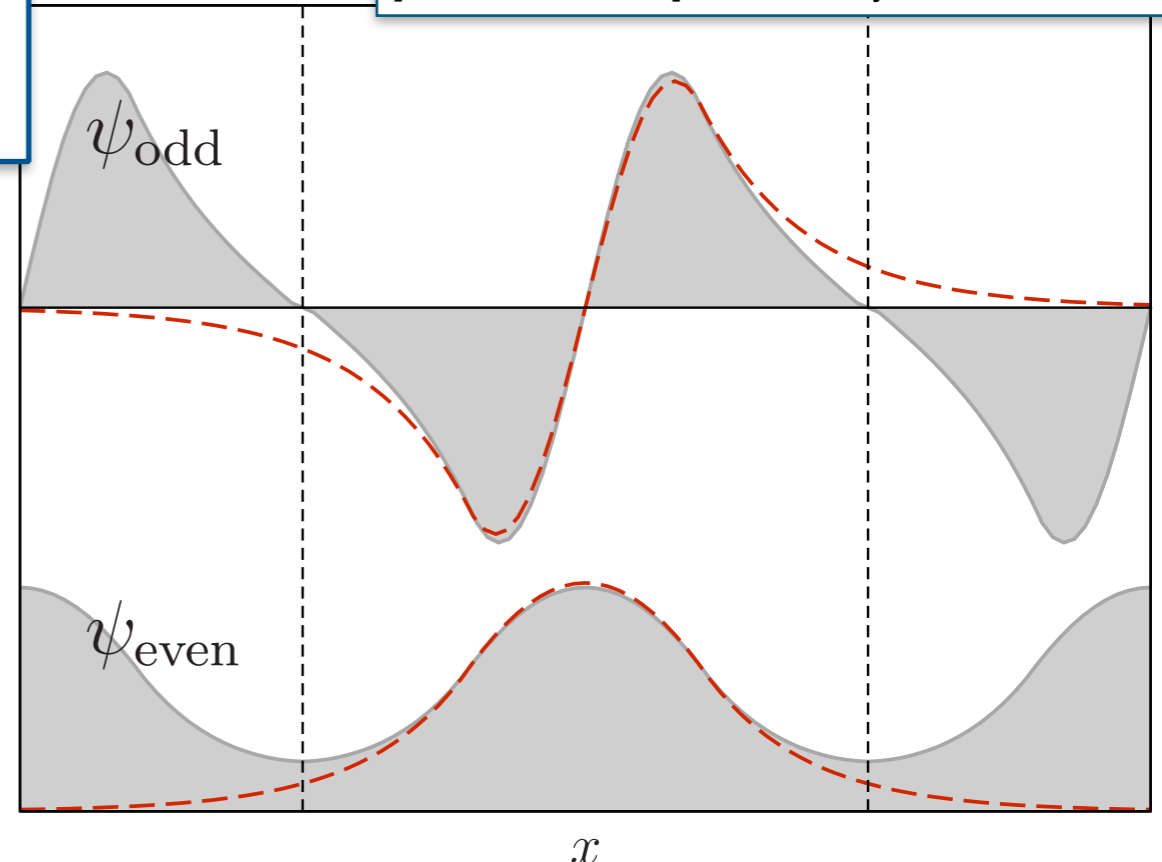
$$\Delta E_L(L, \phi) := \langle \psi_L | \hat{H}_L - E_\infty | \psi_L \rangle$$

“Volume Dependence of the Energy Spectrum in Massive Quantum Field Theories. 1. Stable Particle States”,
M. Lüscher
Commun.Math.Phys. **104** (1986) 177

Finite volume wave as copies of infinite volume wave (PB)

$$\langle \vec{r} | \psi_L \rangle = \sum_{\vec{n} \in \mathbb{Z}^3} \langle \vec{r} + \vec{n}L | \psi_\infty \rangle + \mathcal{O}(e^{-\kappa L})$$

“Non-relativistic bound states in a finite volume”,
H. W. Hammer, S. König, D. Lee
[arXiv:1109.4577] Annals Phys. **327**



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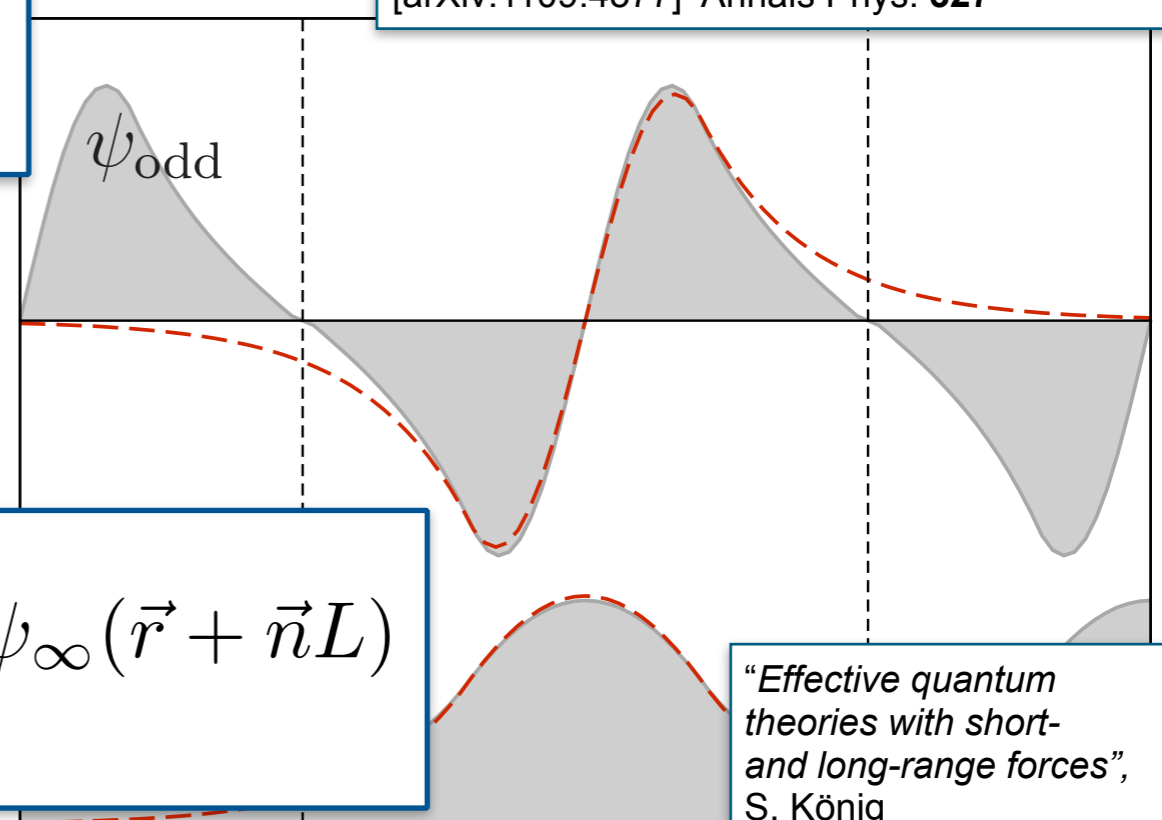
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Energy shift corresponds to different box overlap

$$\Delta E_L^{(LO)}(L) := \sum_{|\vec{n}|=1} \int d^3 \vec{r} \psi_\infty^*(\vec{r}) V(r) \psi_\infty(\vec{r} + \vec{n}L)$$



“Effective quantum theories with short- and long-range forces”,
S. König
Dissertation -
University of Bonn (2013)

Periodic boundary result: S-Wave

$$\Delta E_L^{(LO)}(L) := \sum_{|\vec{n}|=1} \int d^3\vec{r} \psi_\infty^*(\vec{r}) V(r) \psi_\infty(\vec{r} + \vec{n}L) = -3|\mathcal{A}|^2 \frac{e^{-\kappa L}}{\mu L}$$

$$V(\vec{r}) = 0 \Rightarrow \psi_\infty(\vec{r}) = \mathcal{A} \sqrt{\frac{1}{4\pi}} \frac{e^{-\kappa r}}{r} \quad \forall R < r < L, \quad \kappa^2 = -2\mu E_\infty$$

Periodic boundary result: S-Wave

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Twisted boundaries

$$\langle \vec{r}^{\vec{\phi}} | \psi_L \rangle = \sum_{\vec{n} \in \mathbb{Z}^3} \langle \vec{r} + \vec{n}L | \psi_\infty \rangle e^{-i\vec{n} \cdot \vec{\phi}} + \mathcal{O}(e^{-\kappa L})$$

Twisted boundaries result

“Topological phases for bound states moving in a finite volume”,
S. Bour, H. Hammer, S. König, D. Lee, U. Meißner
[arXiv:1107.1272] Phys.Rev. **D84**

$$\begin{aligned} \Delta E_L^{(LO)}(L, \vec{\phi}) &:= \sum_{|\vec{n}|=1} \int d^3\vec{r} \psi_\infty^*(\vec{r}) V(r) \psi_\infty(\vec{r} + \vec{n}L) e^{-i\vec{n} \cdot \vec{\phi}} \\ &= -|\mathcal{A}|^2 \frac{e^{-\kappa L}}{\mu L} \sum_{i=1}^3 \cos(\phi_i) \end{aligned}$$

Formalism and Methodology

For 2-Body Level

Sources of errors

*(Numerical deviation
from theoretical prediction)*

$$\hat{H}_L |\psi\rangle = E_L |\psi\rangle, \quad V_L(\vec{n}) = \frac{c}{a^3} \delta_{\vec{n}, \vec{0}}$$

▶ ~~Contact interaction estimate~~
(Fitting on Lattice)

▶ Solving procedure
(Lanczos like iteration)

▶ Resolution not sufficient
(Discretization errors)

*Uncertainty of solving
procedure*

$$=: \delta E_S \simeq 10^{-4} [MeV]$$

**Same spacing for all
computations (?)**

Momenta in FV:

$$\vec{r} \in \mathbb{R}^3 \mapsto \vec{r} \in L^3 \quad \Rightarrow \quad \vec{p} \in \mathbb{R}^3 \mapsto \frac{2\pi}{L} \vec{n}_p, \quad \vec{n}_p \in \mathbb{Z}^3$$

Momenta in discrete space:

Discretization (one step derivative):

$$\partial_{x,a}^2 f(\vec{r}) := \frac{1}{a^2} (f(\vec{r} + a\vec{e}_x) - 2f(\vec{r}) + f(\vec{r} - a\vec{e}_x))$$

Dispersion relation (one step derivative):

$$\hat{p}_x^2 |\vec{p}\rangle \mapsto \frac{2(1 - \cos(p_x a))}{a^2} |\vec{p}\rangle$$

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Momenta in discrete space:

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$$\partial_{x,a}^2 f(\vec{r}) := \frac{1}{a^2} \left(e^{i\phi_x a/L} f(\vec{r} + a\vec{e}_x) - 2f(\vec{r}) + e^{-i\phi_x a/L} f(\vec{r} - a\vec{e}_x) \right)$$

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How does discretization affect momenta?

▶ *Contact interactions*

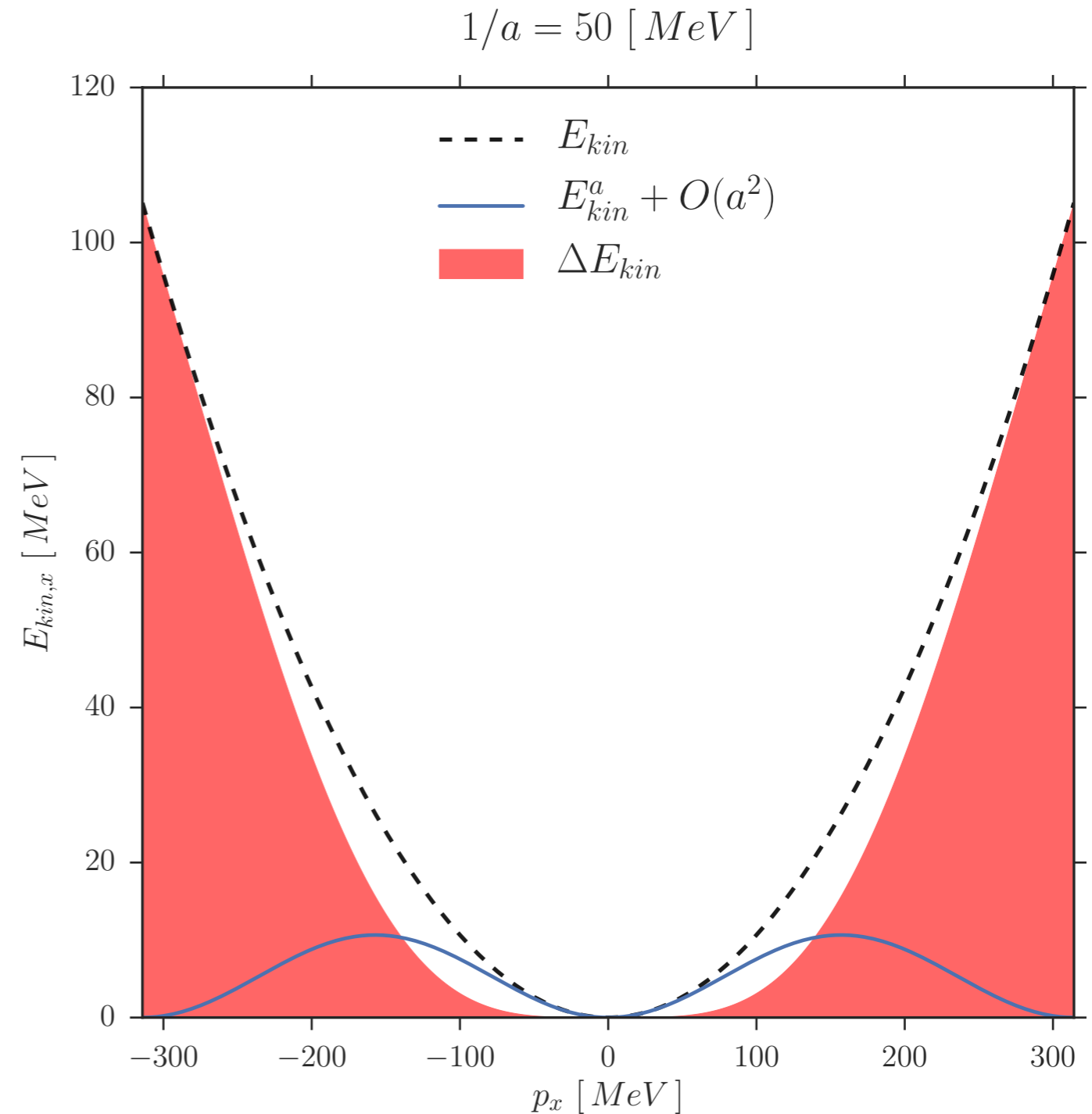
LECs are cutoff dependent (fitting)

How does discretization affect momenta?

► **Contact interactions**

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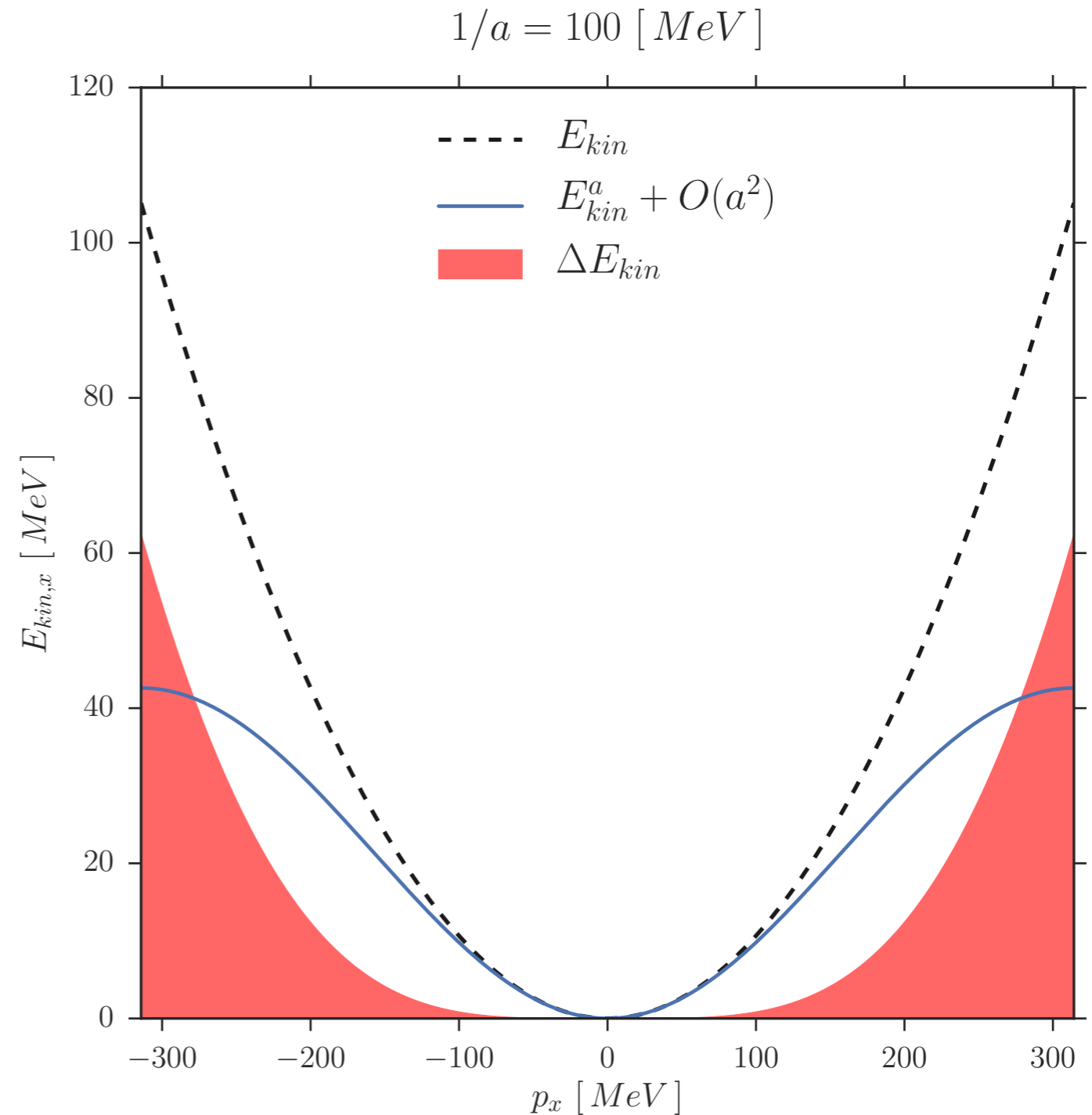


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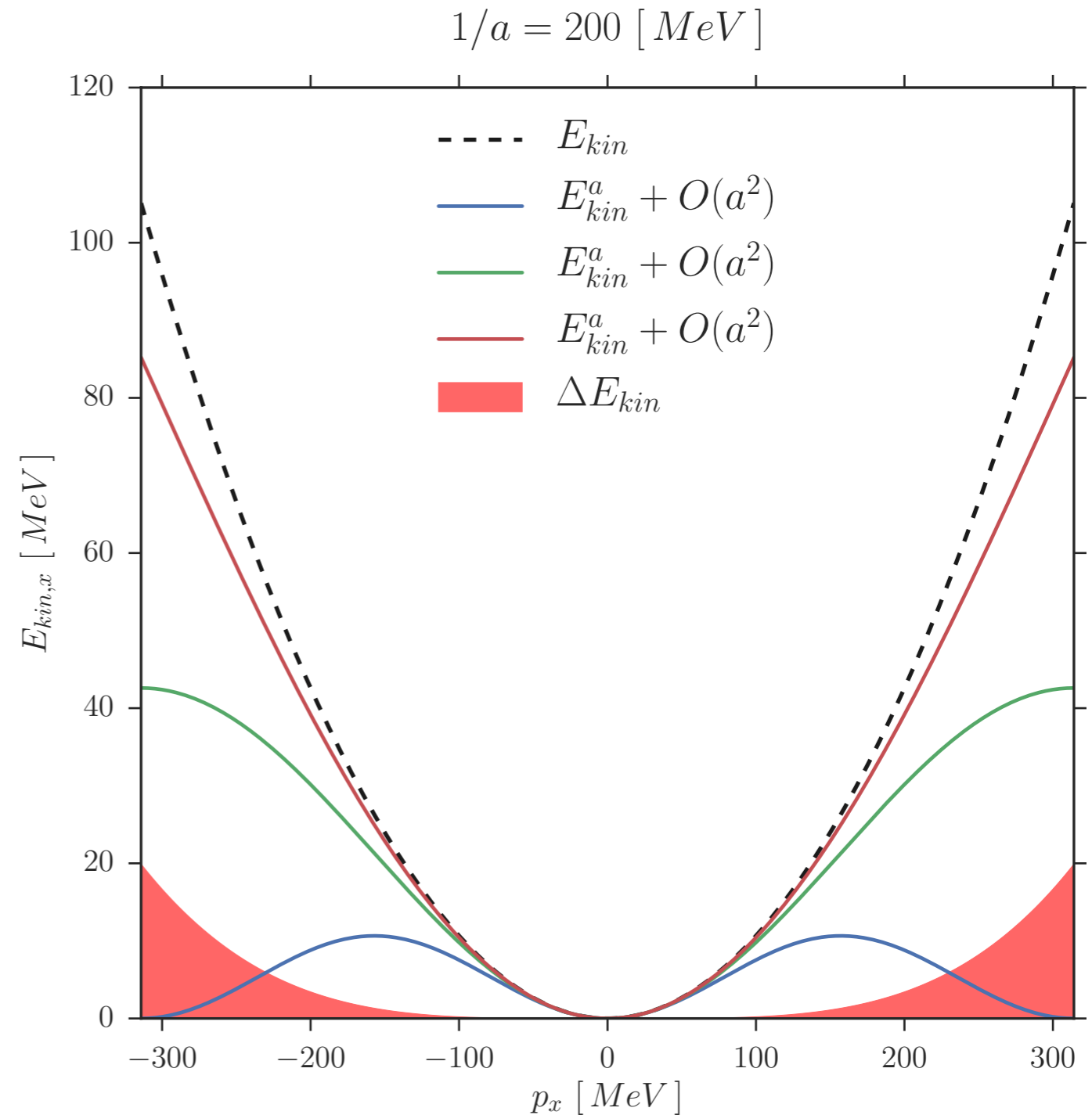


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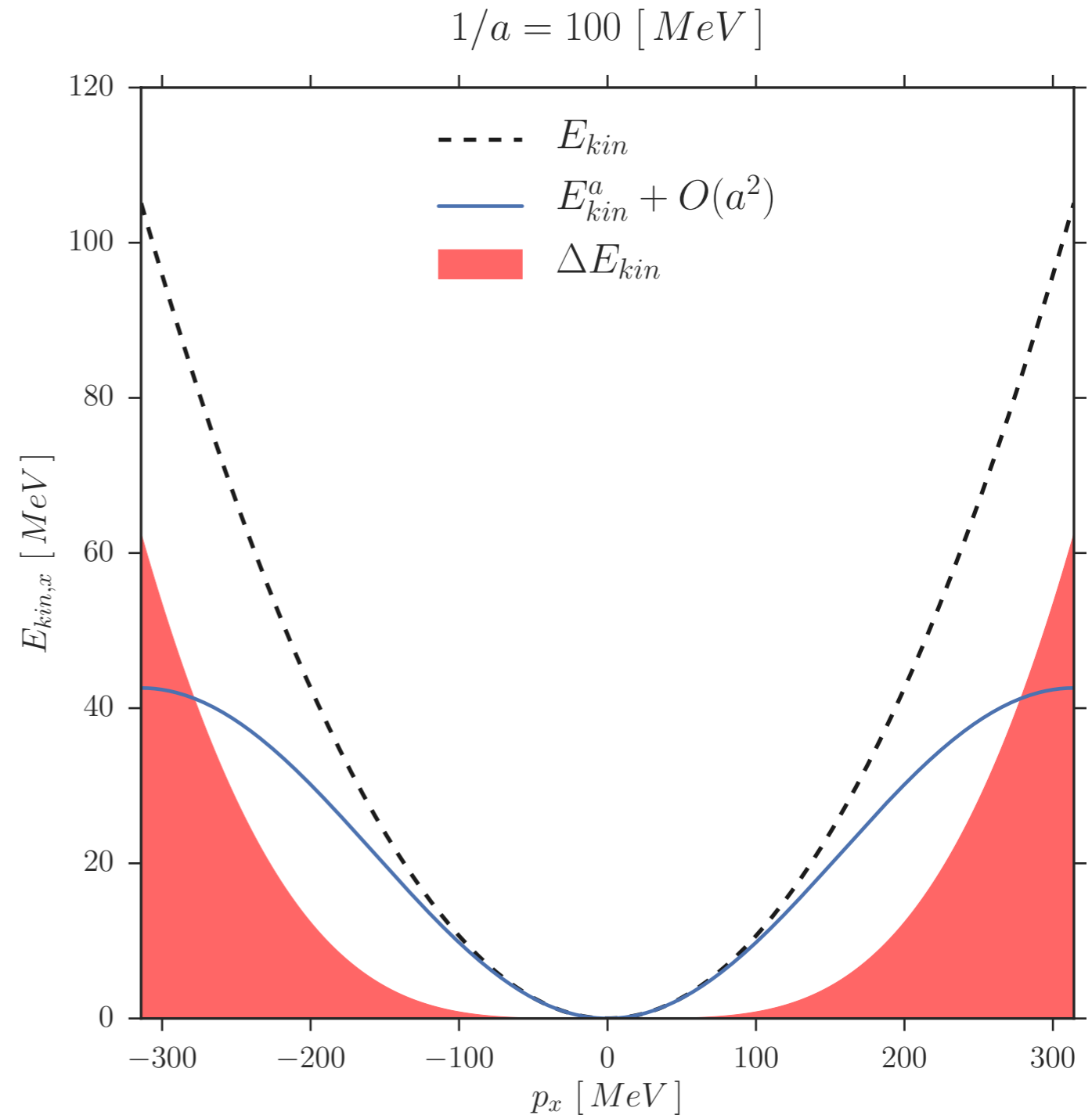


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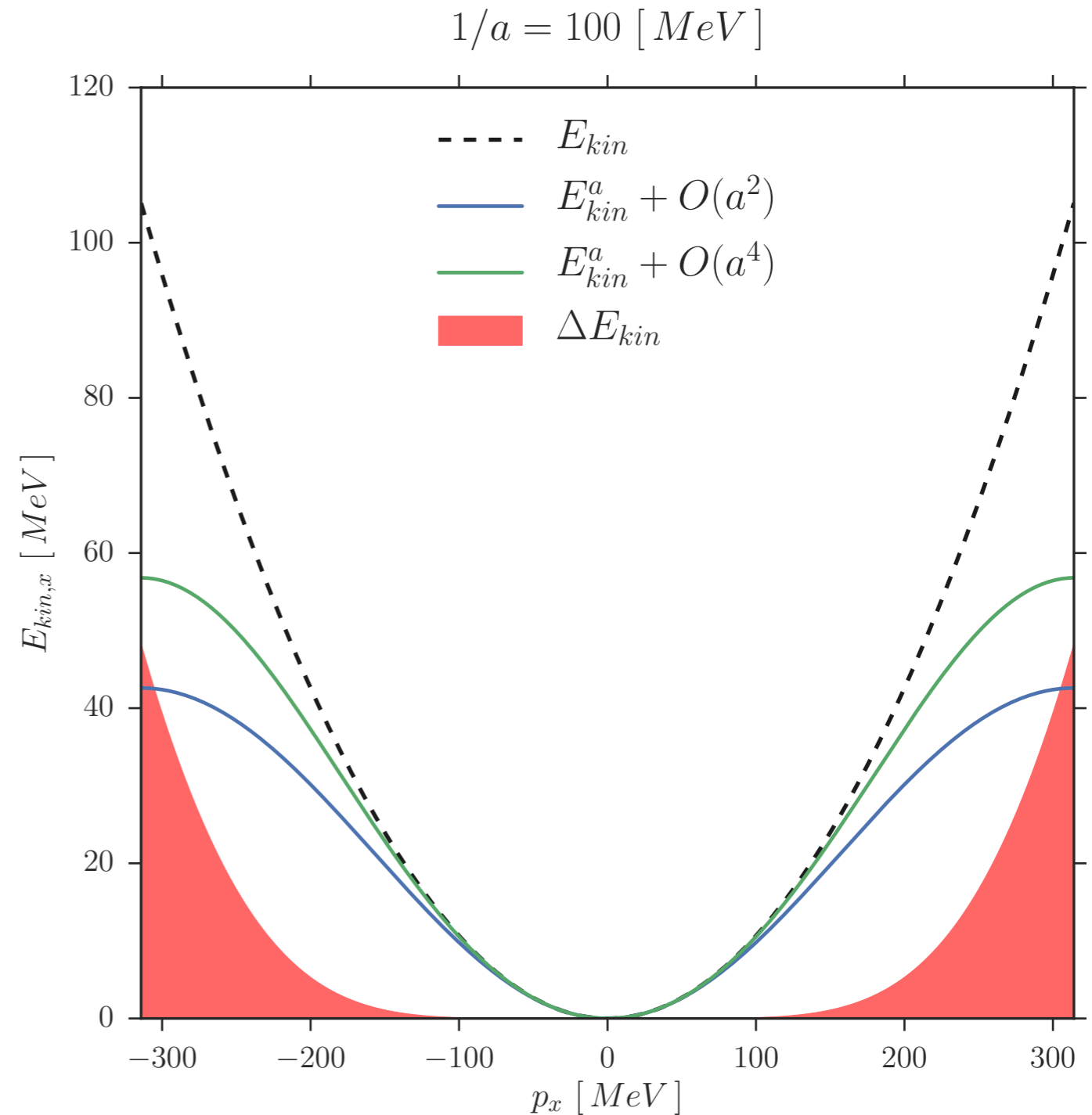
$$\partial_{x,a}^{(N)^2} f(\vec{r}) = \frac{1}{a^2} \sum_{n=0}^N (-1)^{n+1} \omega_n (f(\vec{r} + an\vec{e}_x) + f(\vec{r} - an\vec{e}_x))$$

How does discretization affect momenta?

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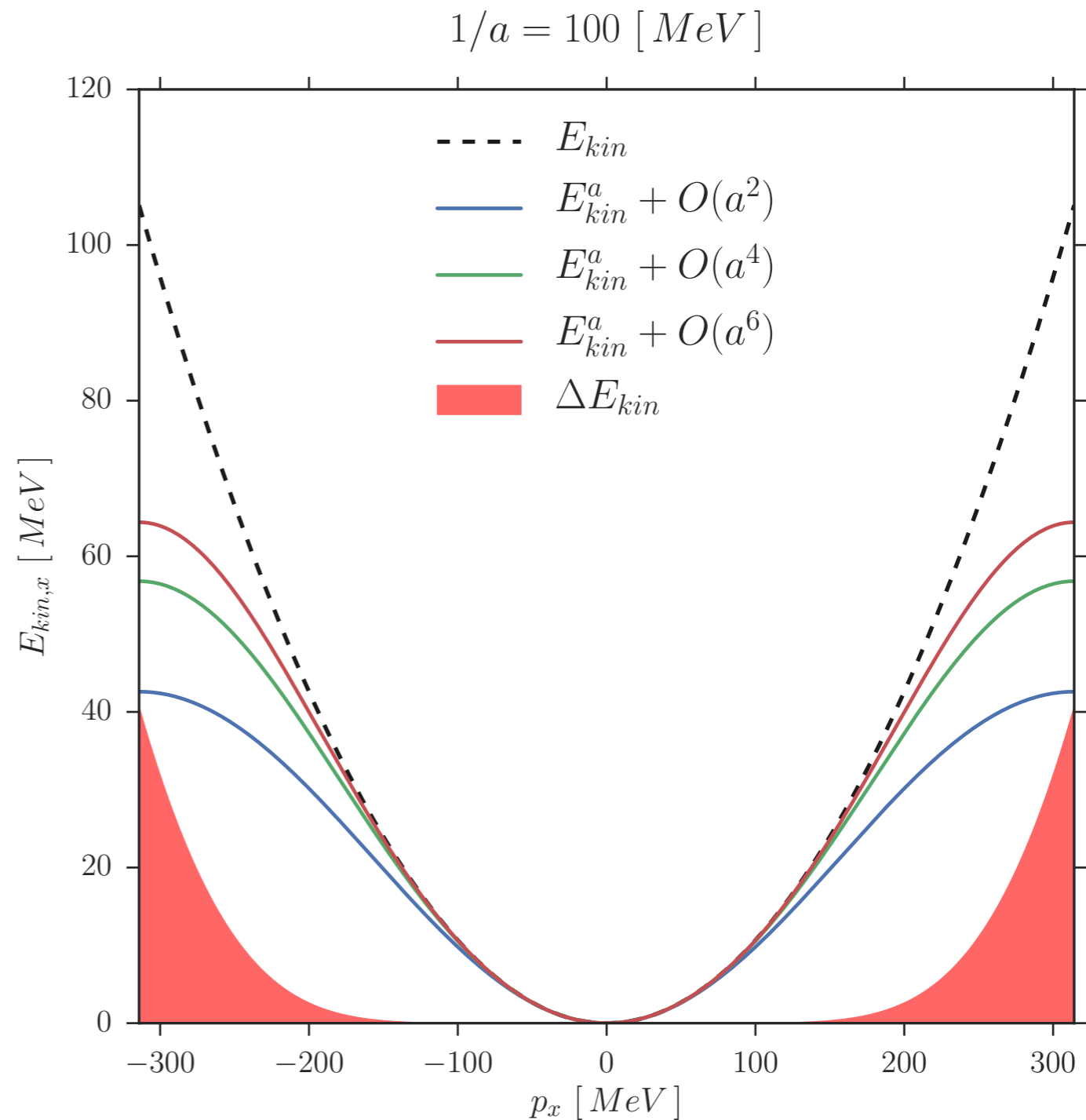
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Twisted discrete dispersion: $N_L = 3$

$$a = 1.97 \text{ fm}$$

$$\vec{n}_p \sim \vec{n}'_p \Leftrightarrow \vec{n}_p = \vec{n}'_p + N_L \vec{e}$$

$$\vec{p}^{\vec{\phi}} = \frac{2\pi}{L} \vec{n}_p + \frac{\vec{\phi}}{L}$$

$$\Delta E := \text{mean} \left(\frac{E_{kin} - E_{kin}^{L, \vec{\phi}, a}}{E_{kin}} \right)$$

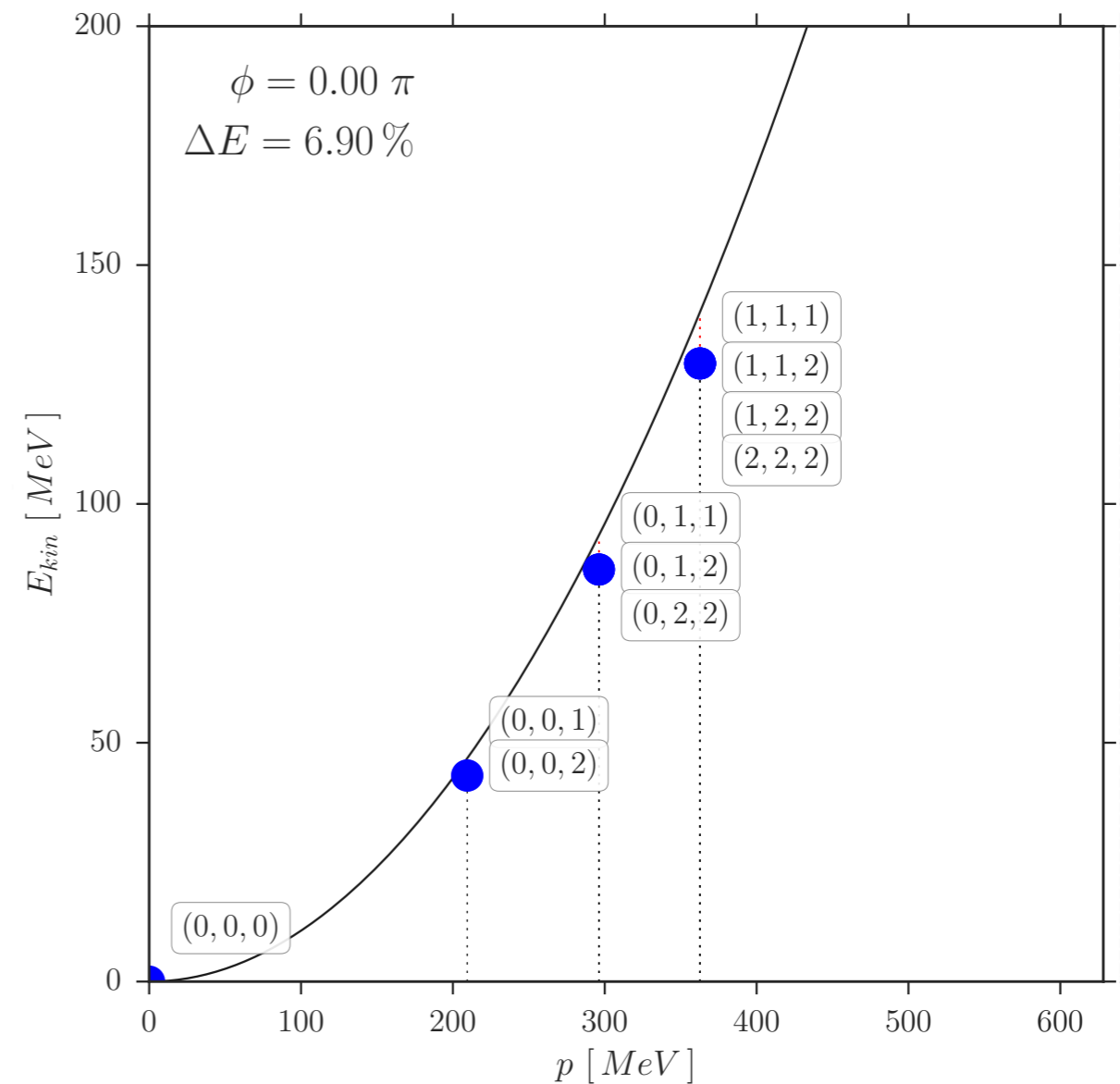
Legend



*Infinite volume
continuum dispersion*



*Finite volume
discrete dispersion*

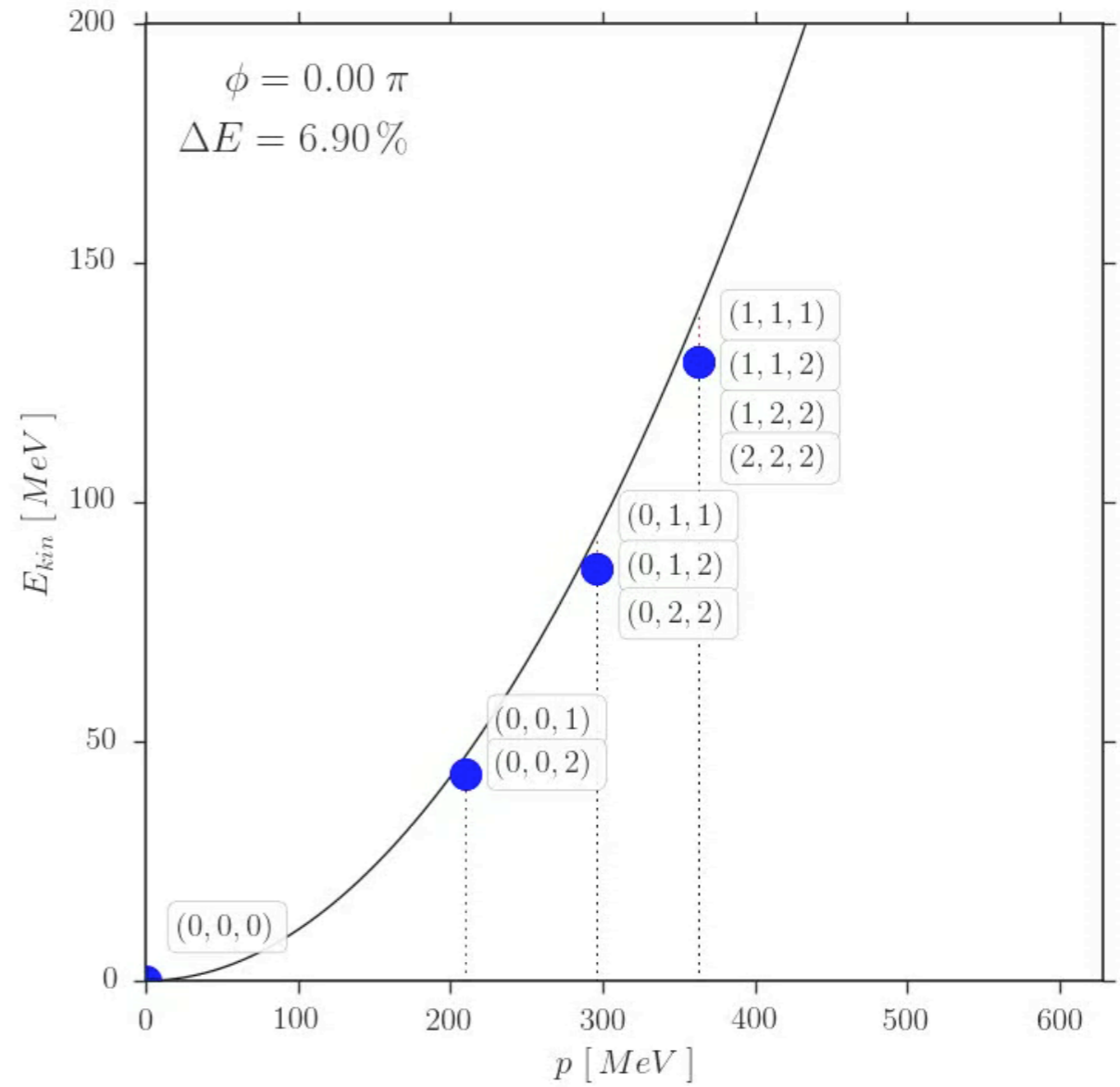
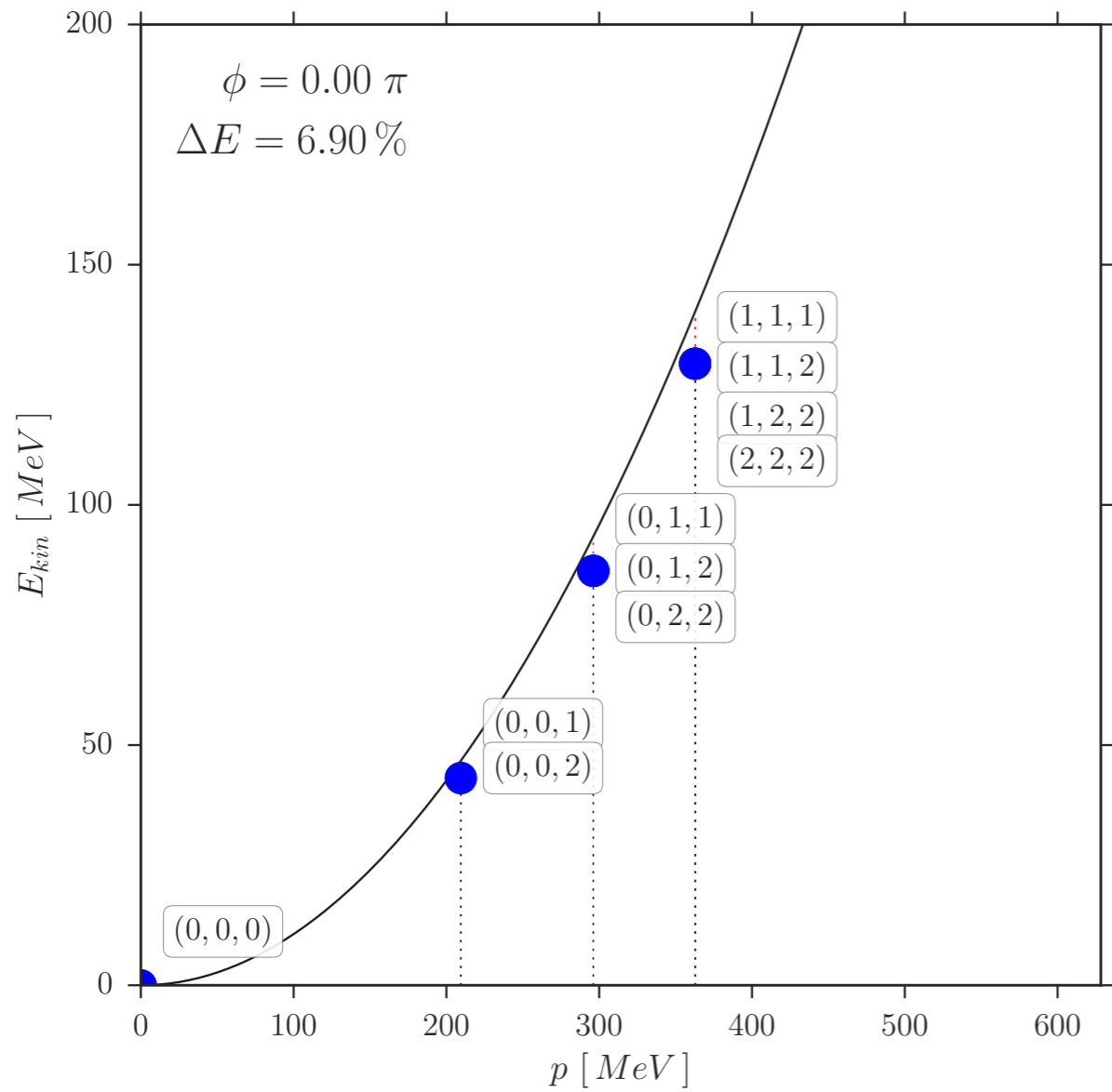


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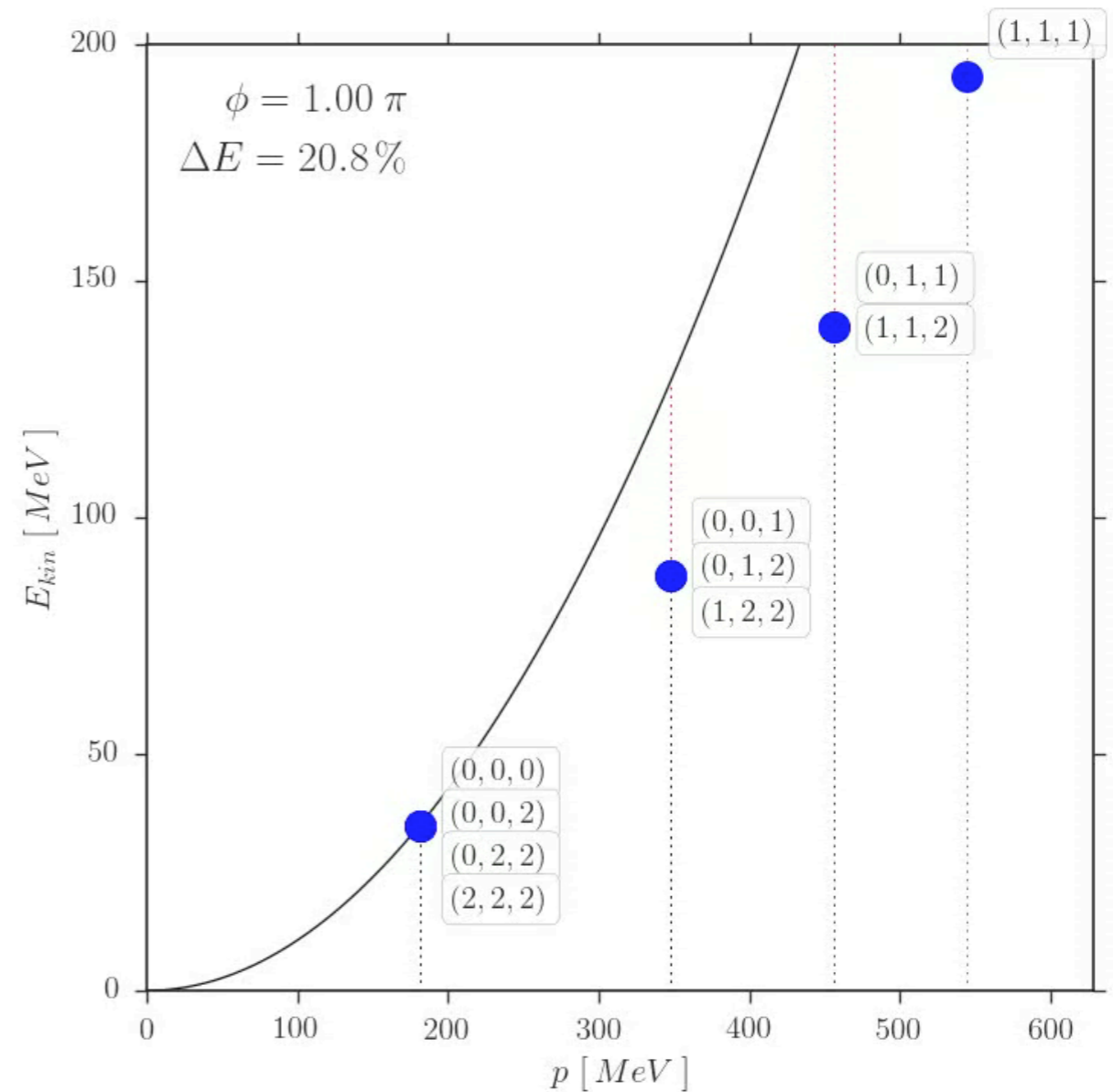
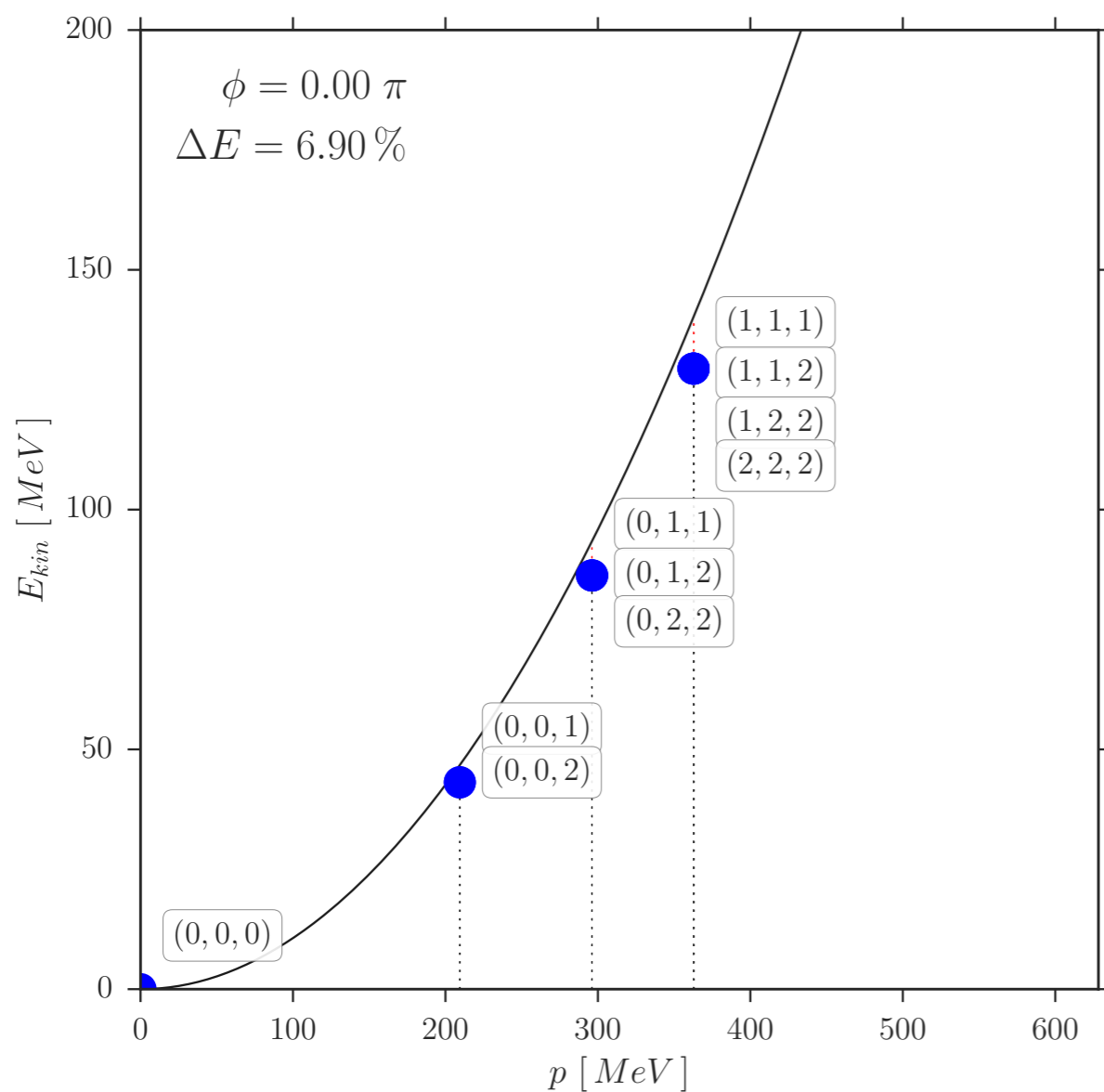


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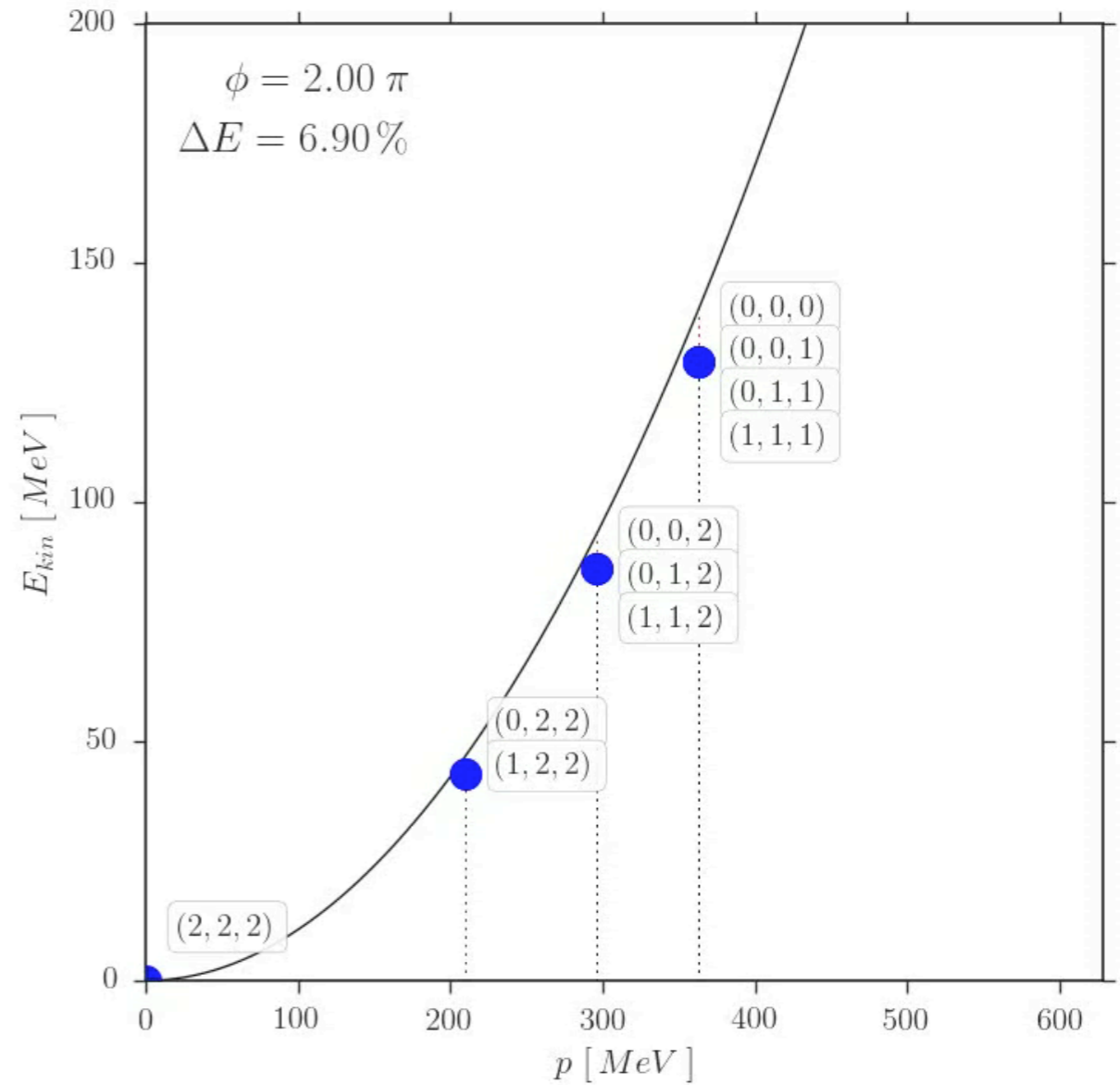
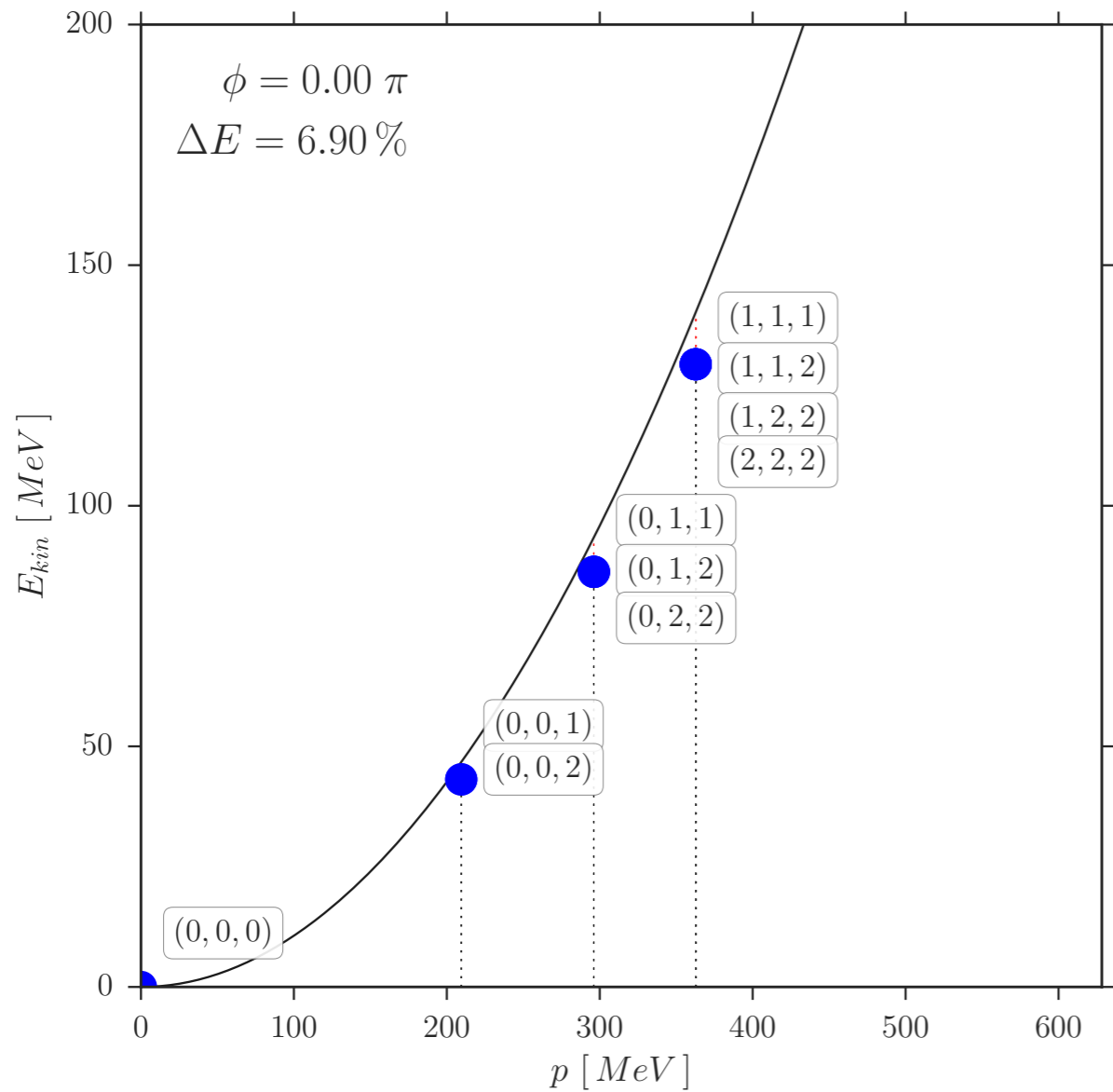


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Sources of errors

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- ▶ ~~Contact interaction estimate~~
(Fitting on Lattice)

Uncertainty of solving procedure

- ▶ Solving procedure
(Lanczos like iteration)

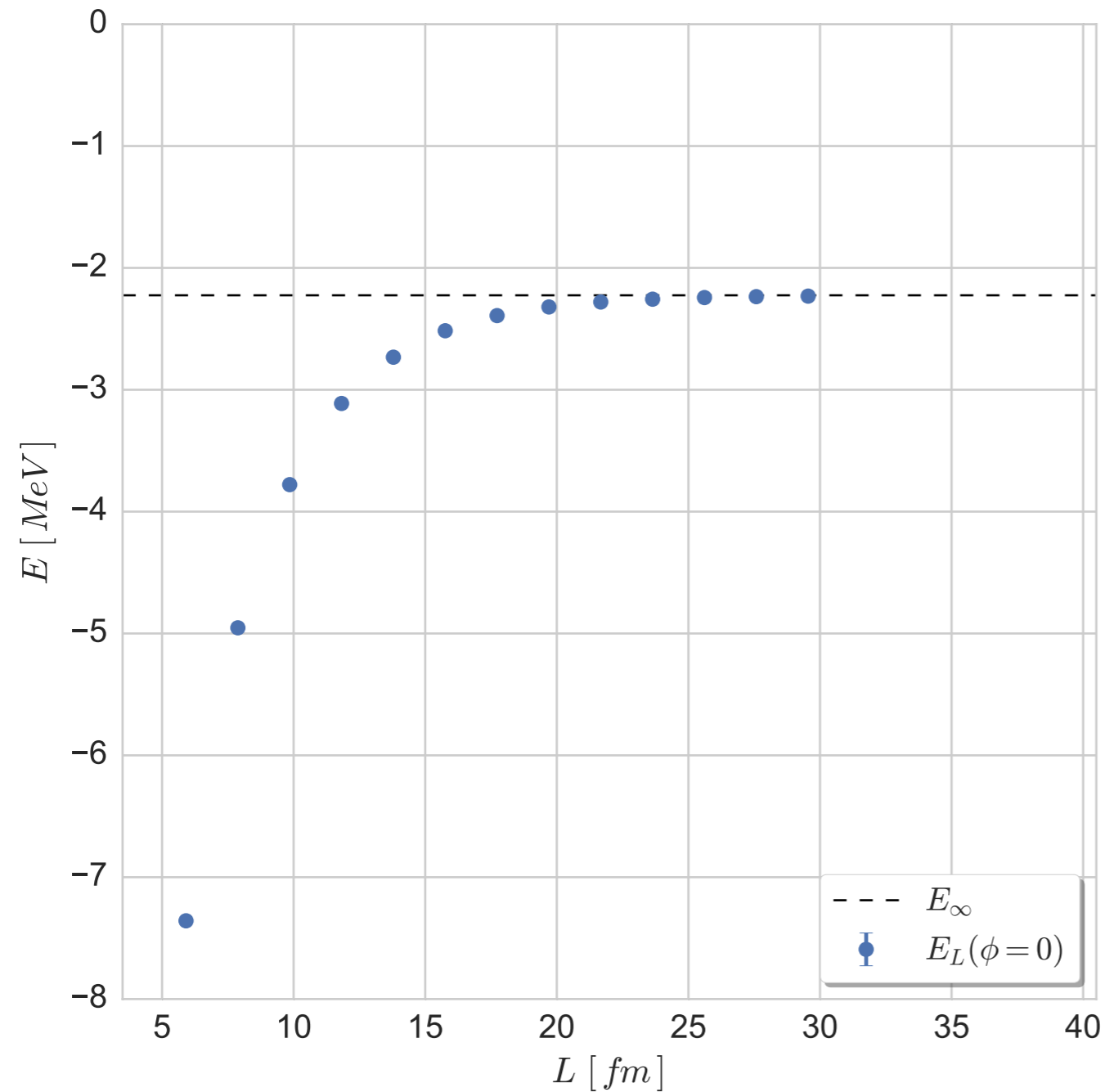
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- ▶ Resolution not sufficient
(Discretization errors)

Same spacing for all computations (?)

- ▶ Box too small
(Finite Volume)

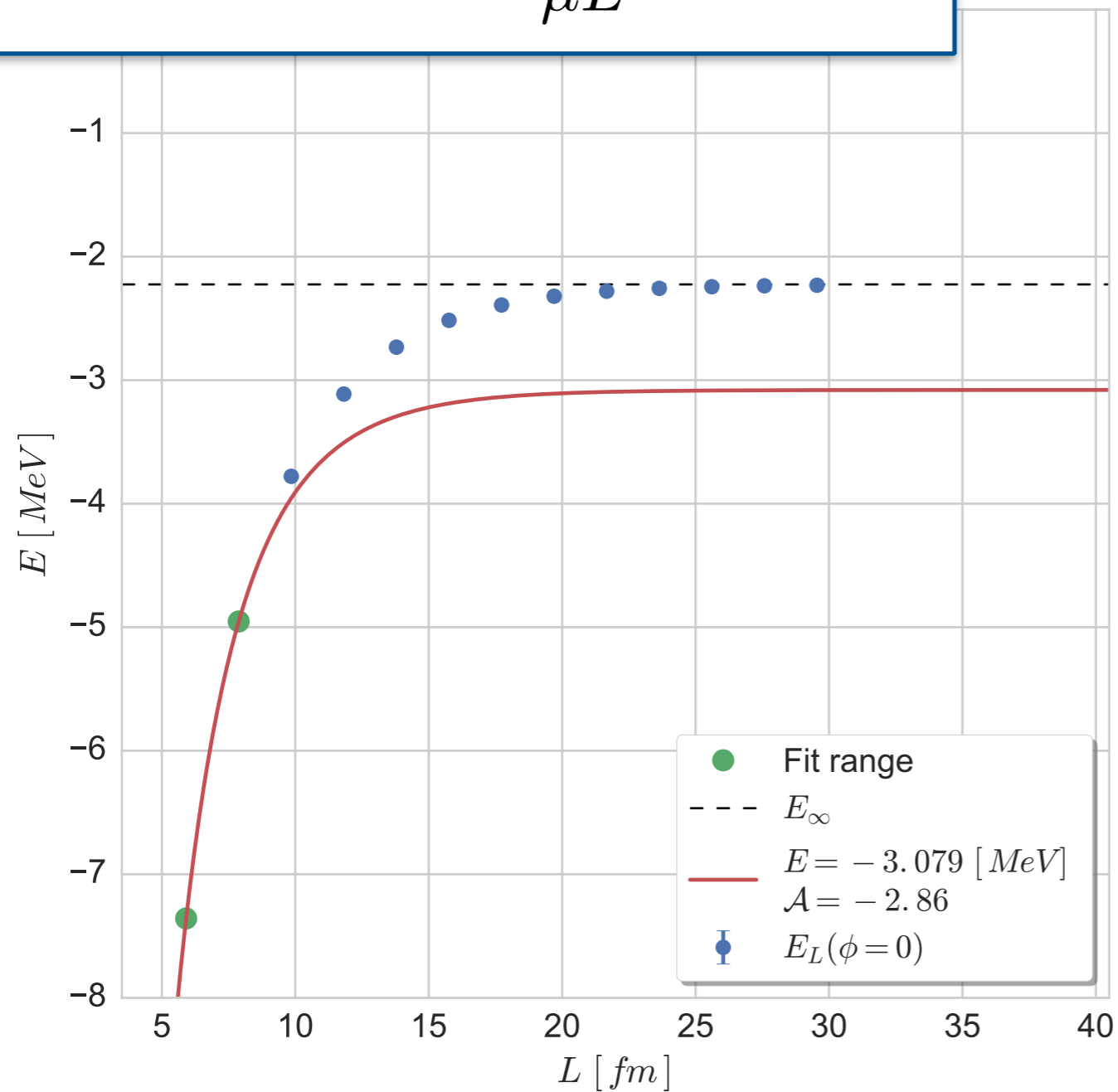
Size of NLO errors?



► **Goal:**
Reliable results for small boxes

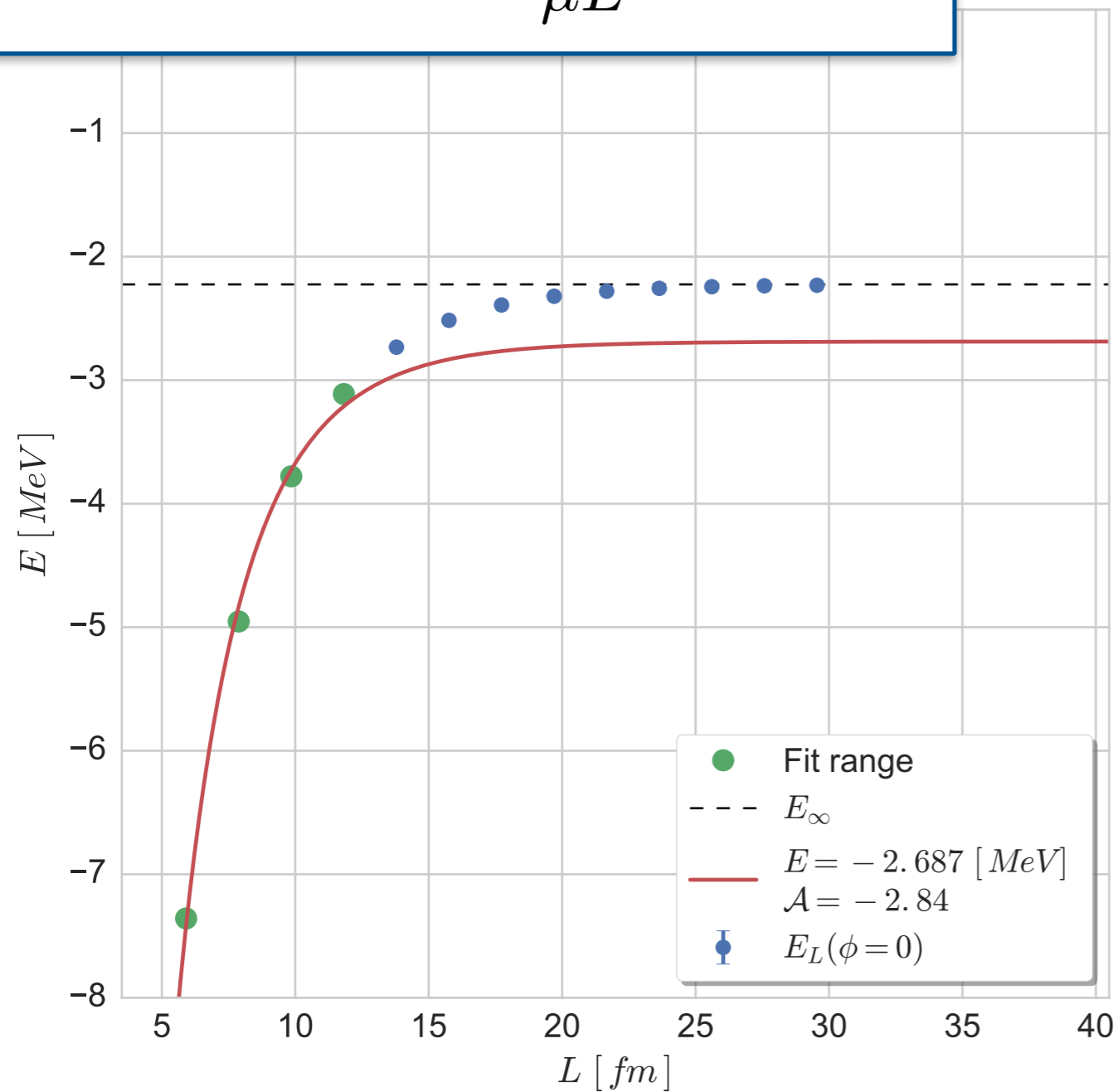
► **Data uncertainties**
Errors:
Numerical errors

$$E_L(E, \mathcal{A}) \stackrel{!}{=} A_0 \frac{e^{-\kappa(E)L}}{\mu L} \mathcal{A} + E$$



- ▶ **Goal:**
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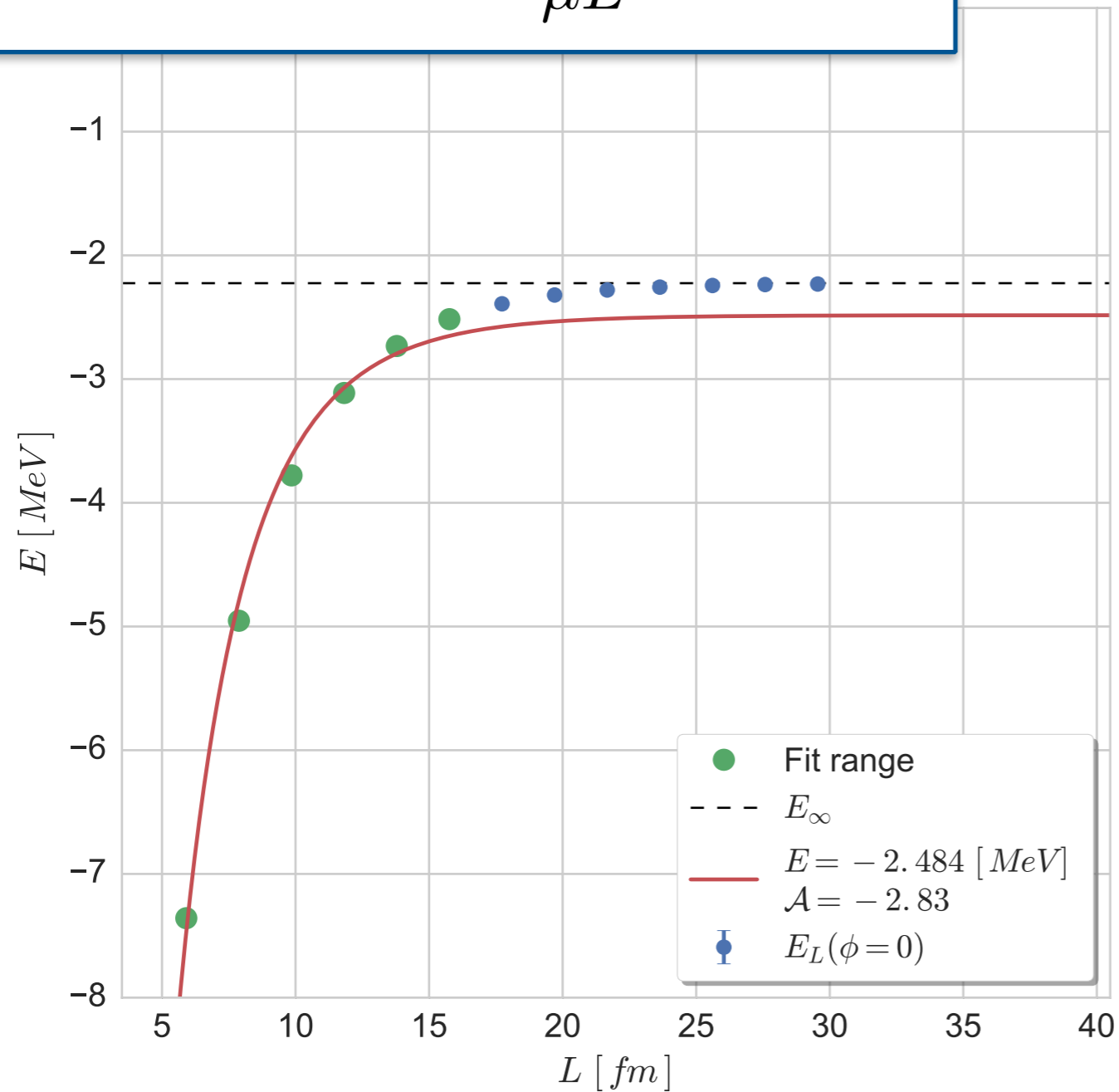
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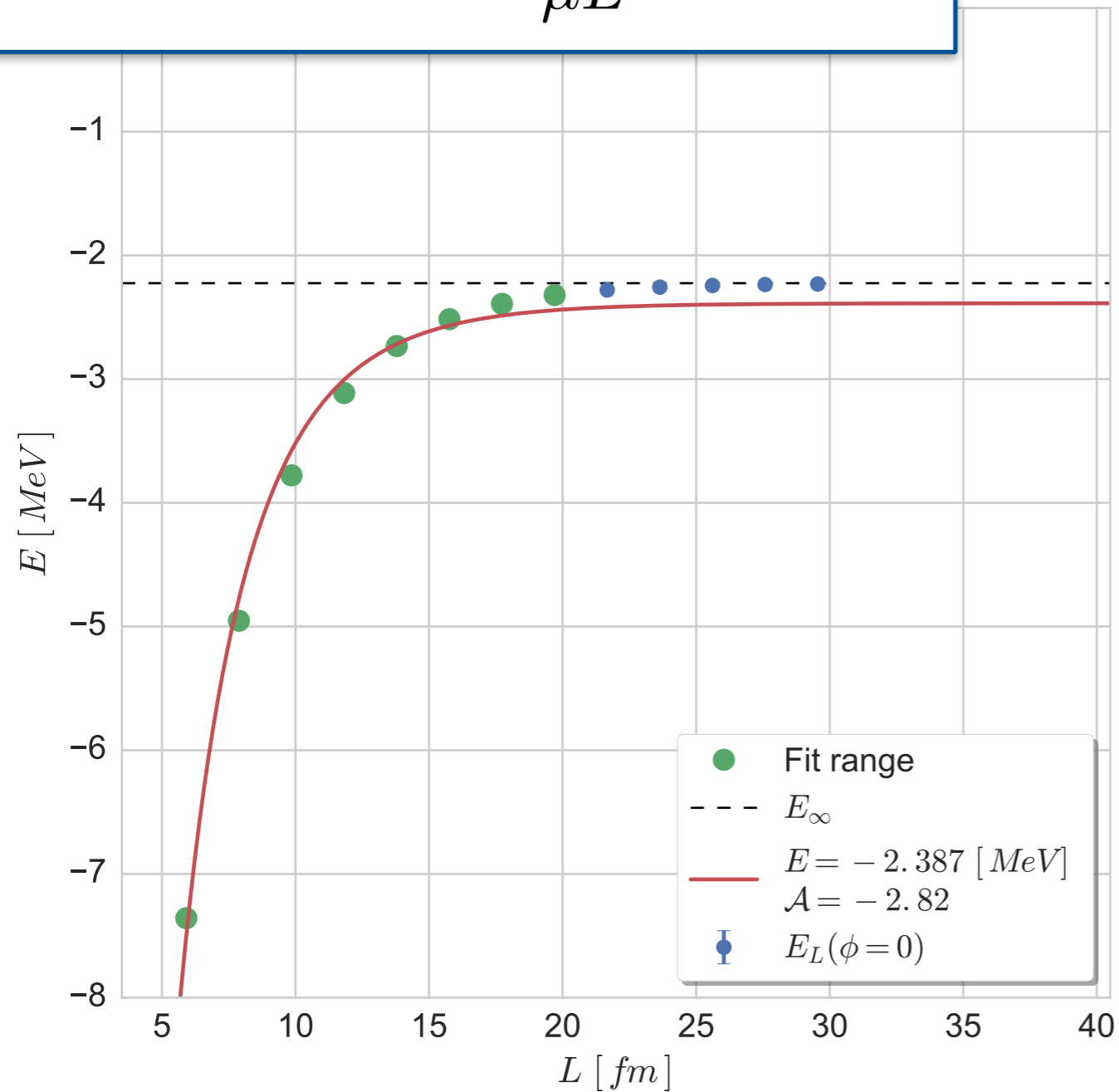
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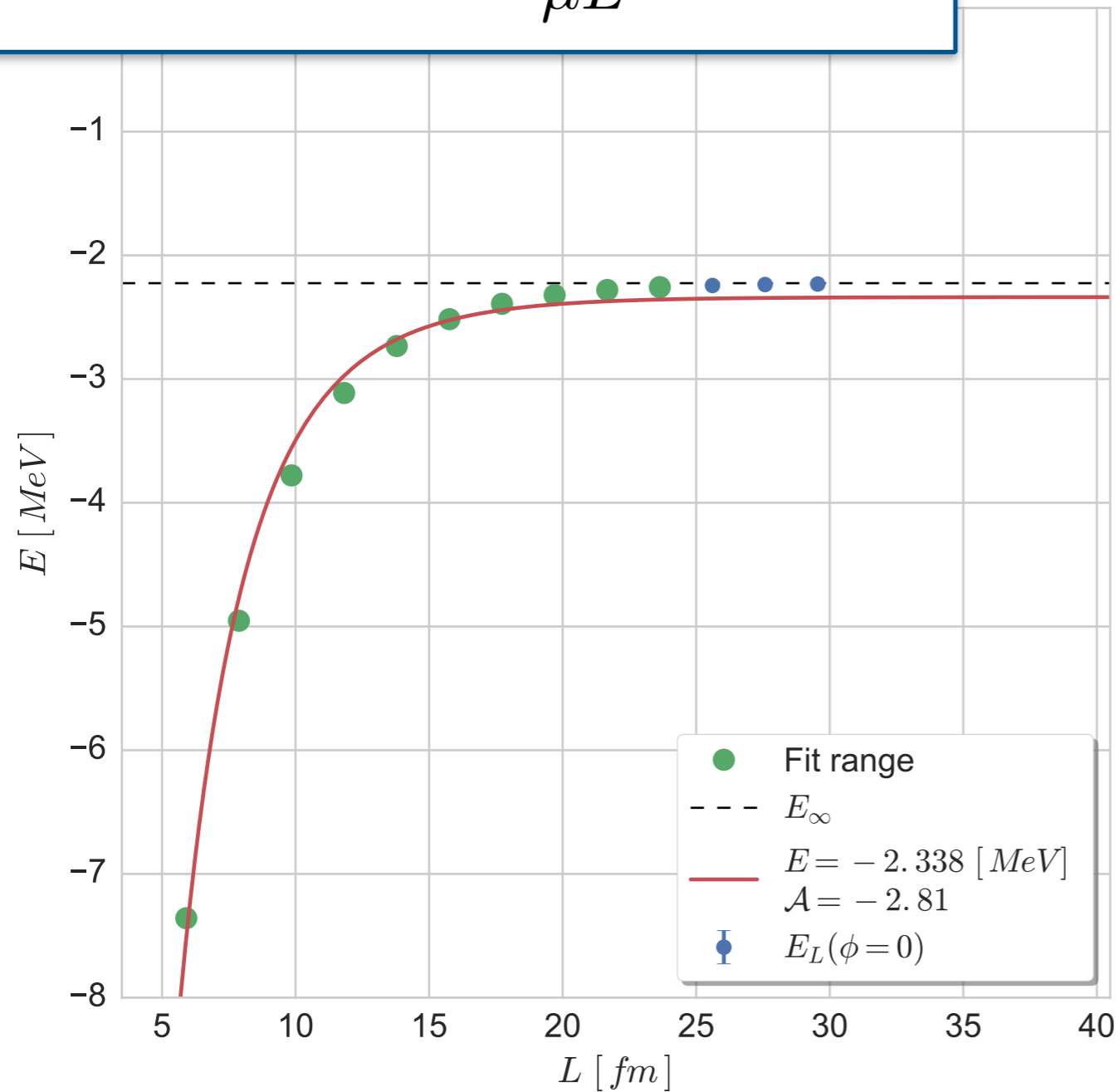
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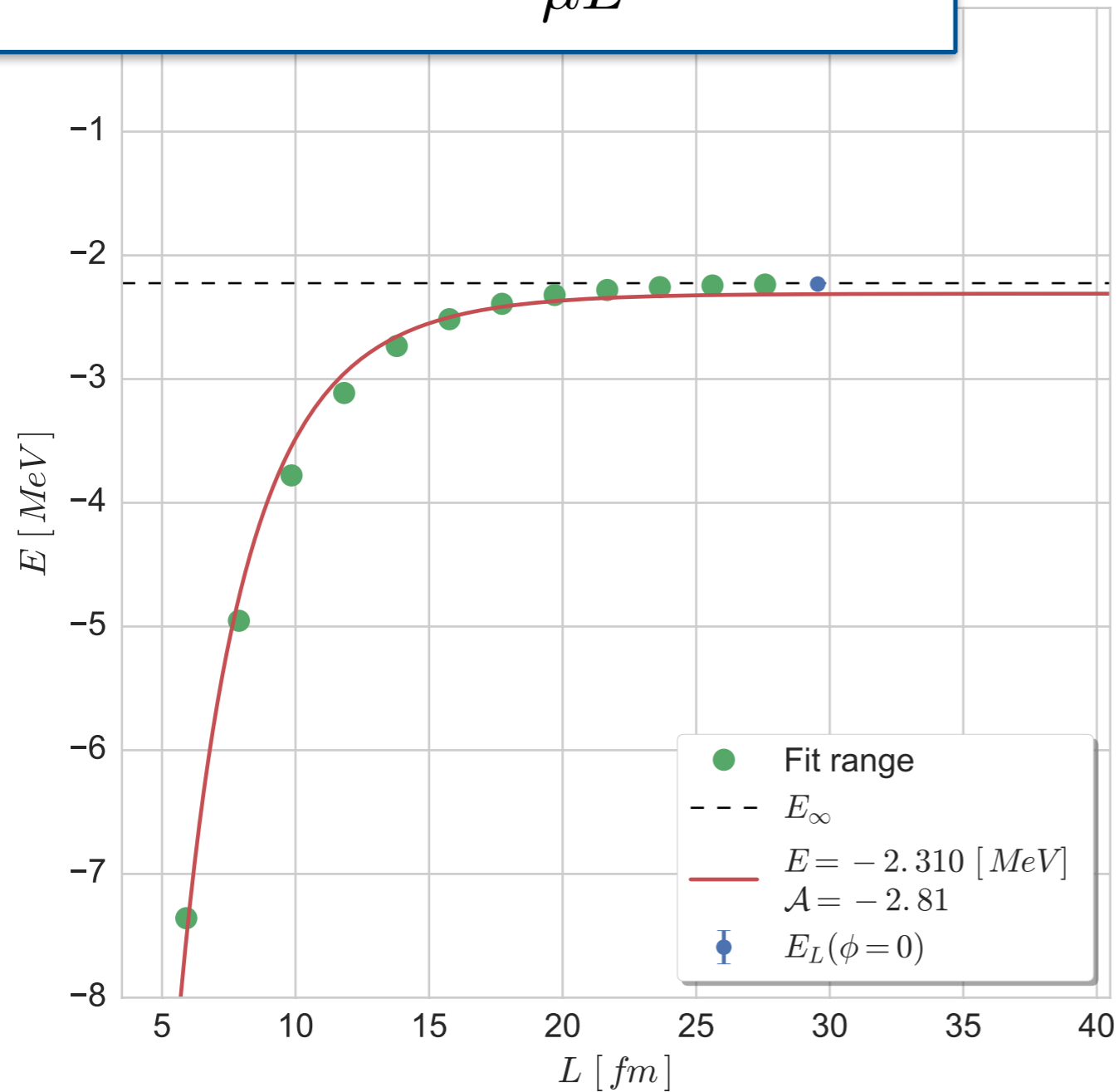
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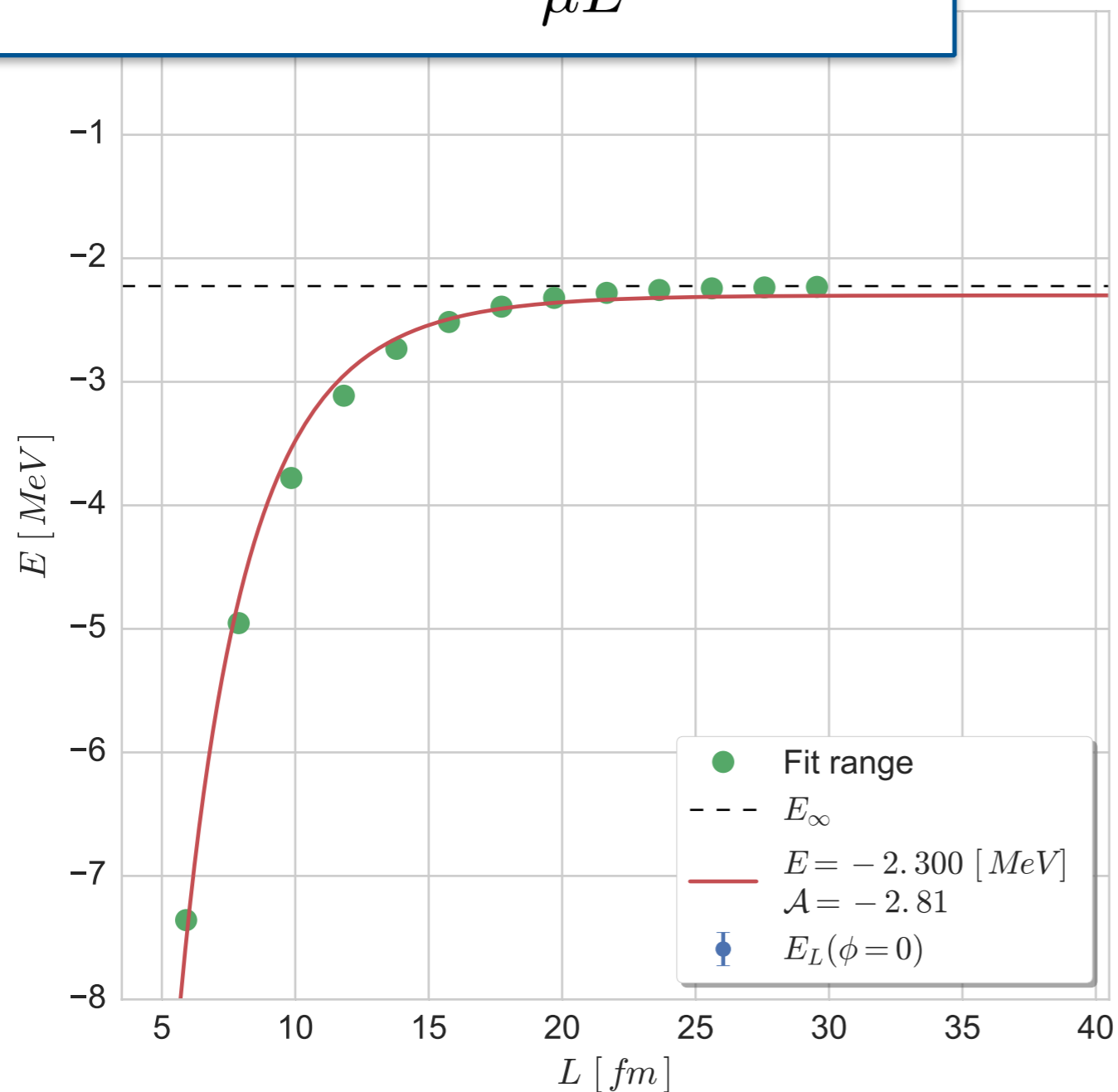
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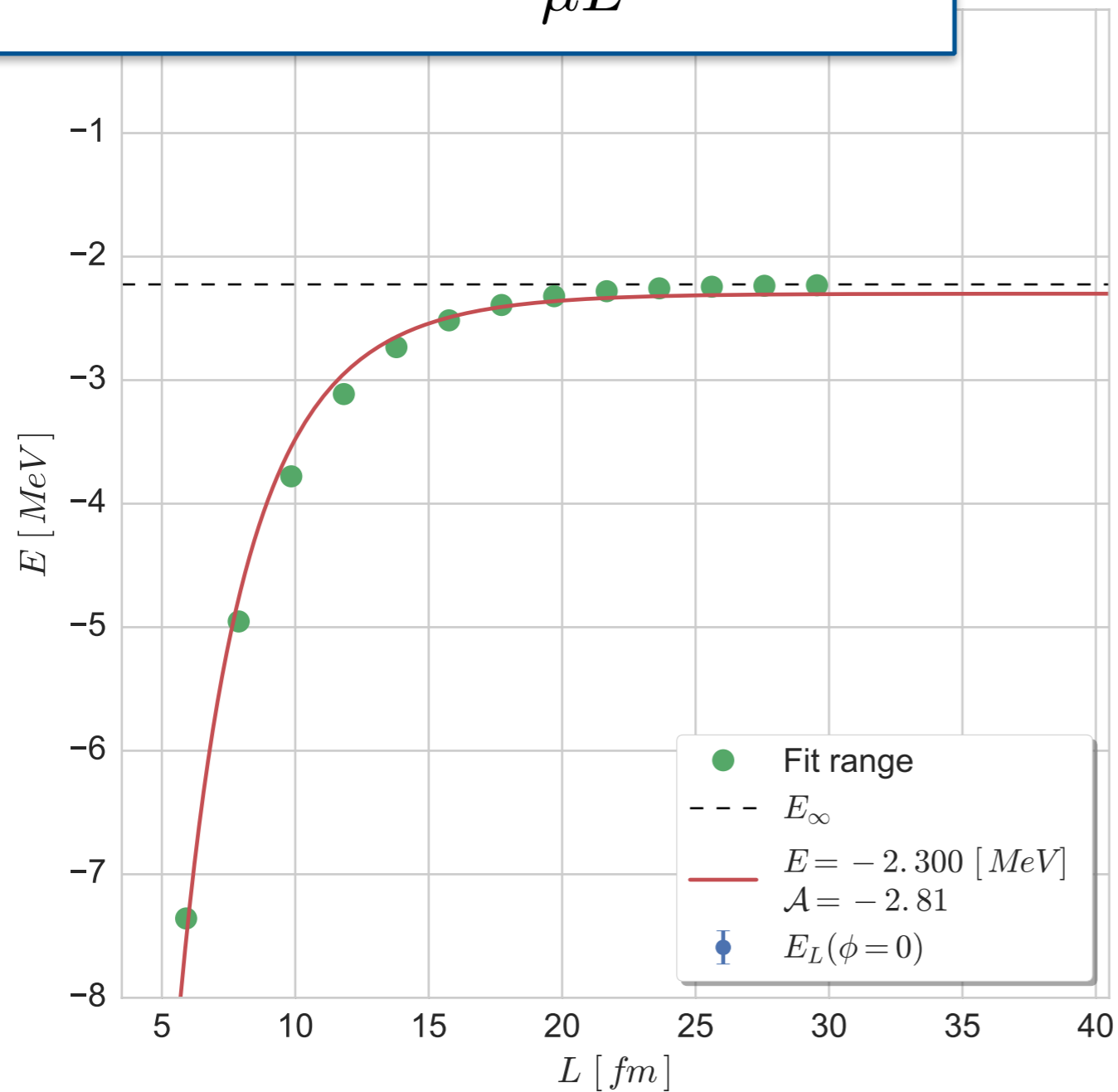
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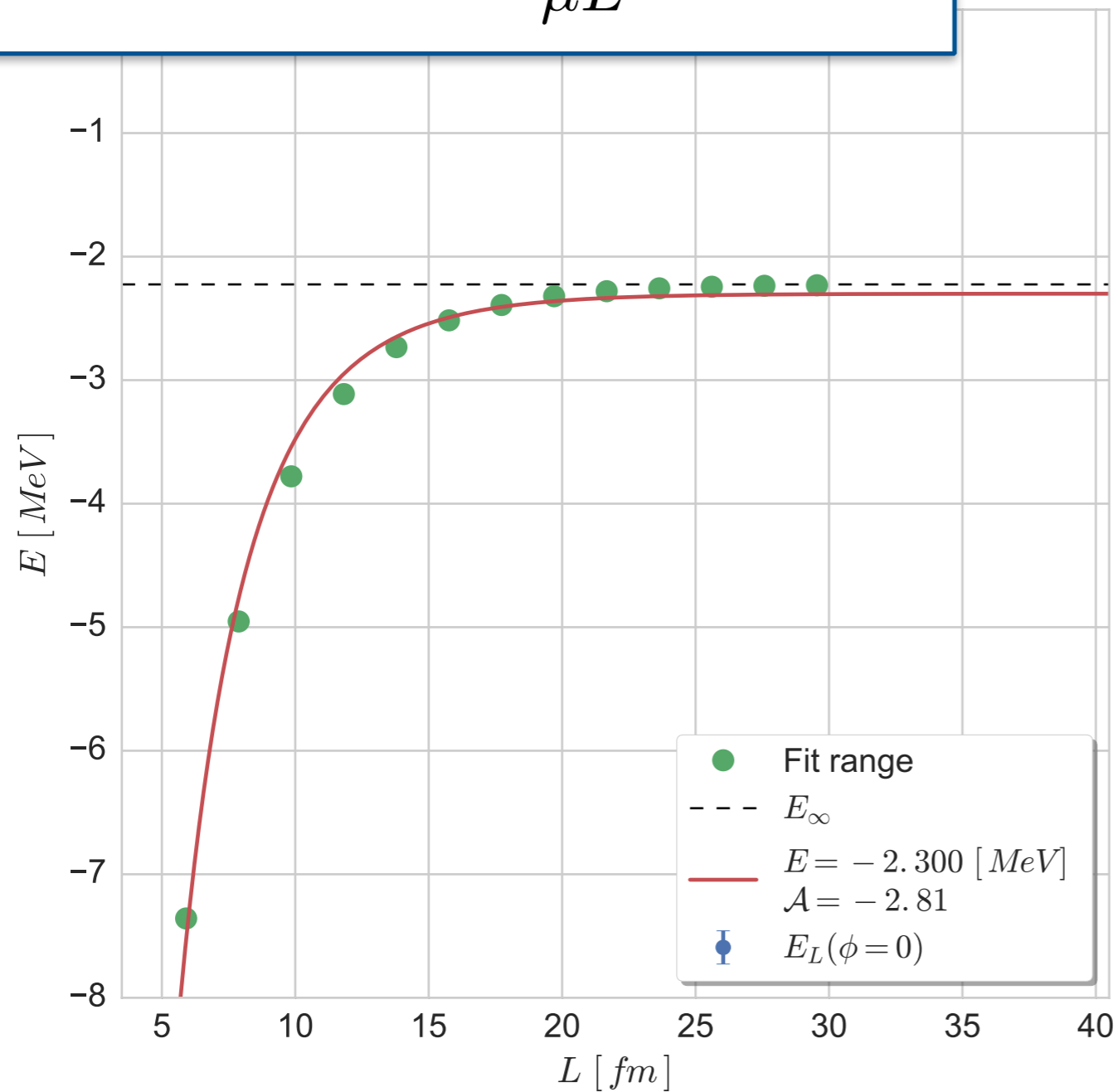
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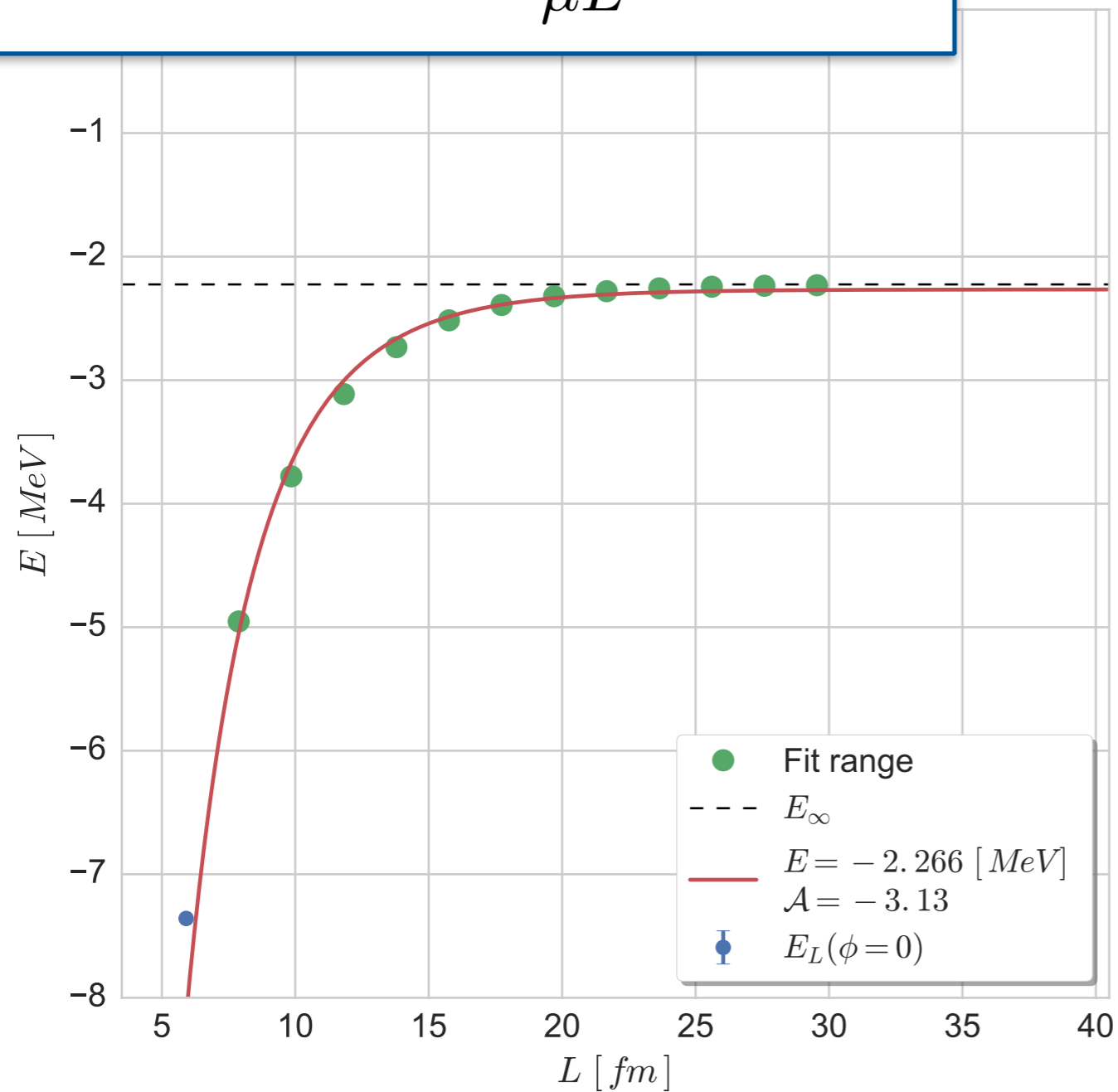
- ▶ **Goal:**
Reliable results for small boxes
- ▶ **Data uncertainties**
Errors:
Numerical errors
- ▶ **Problem:**
Convergence?

$$E_L(E, \mathcal{A}) \stackrel{!}{=} A_0 \frac{e^{-\kappa(E)L}}{\mu L} \mathcal{A} + E$$



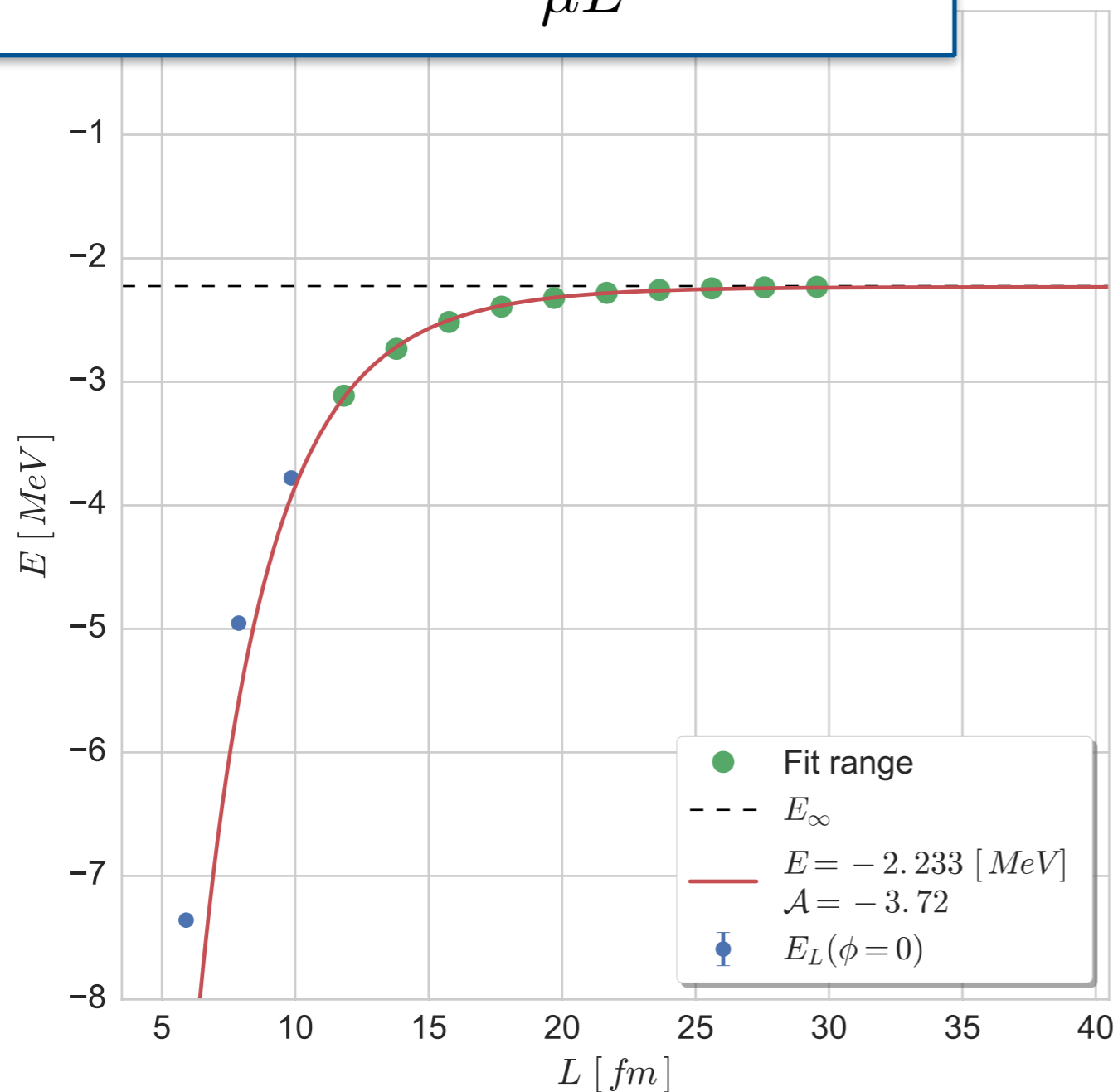
- ▶ **Goal:**
Reliable results for small boxes
- ▶ **Data uncertainties**
Errors:
Numerical errors
- ▶ **Problem:**
Results for small boxes affect fitting

$$E_L(E, \mathcal{A}) \stackrel{!}{=} A_0 \frac{e^{-\kappa(E)L}}{\mu L} \mathcal{A} + E$$



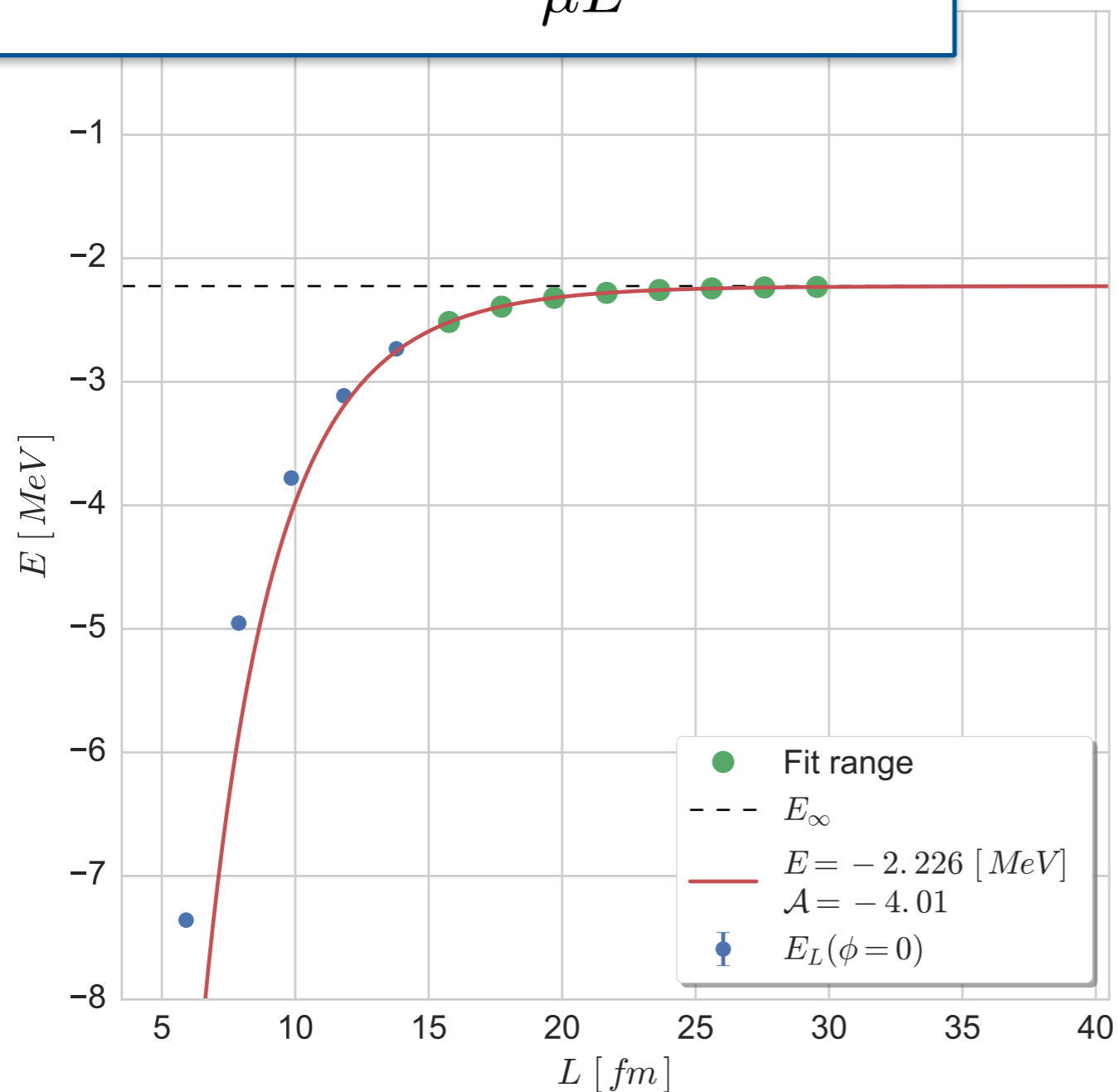
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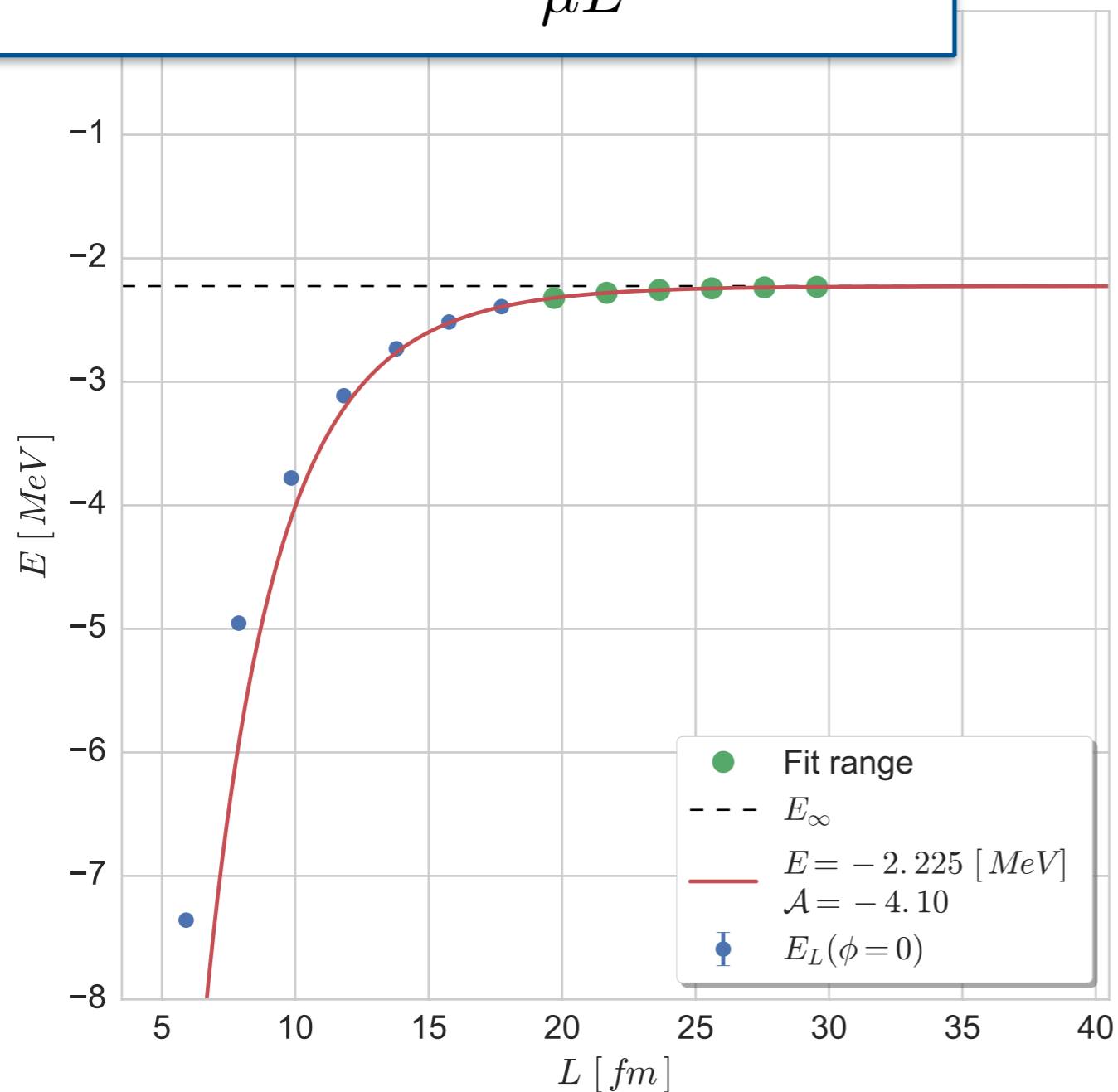
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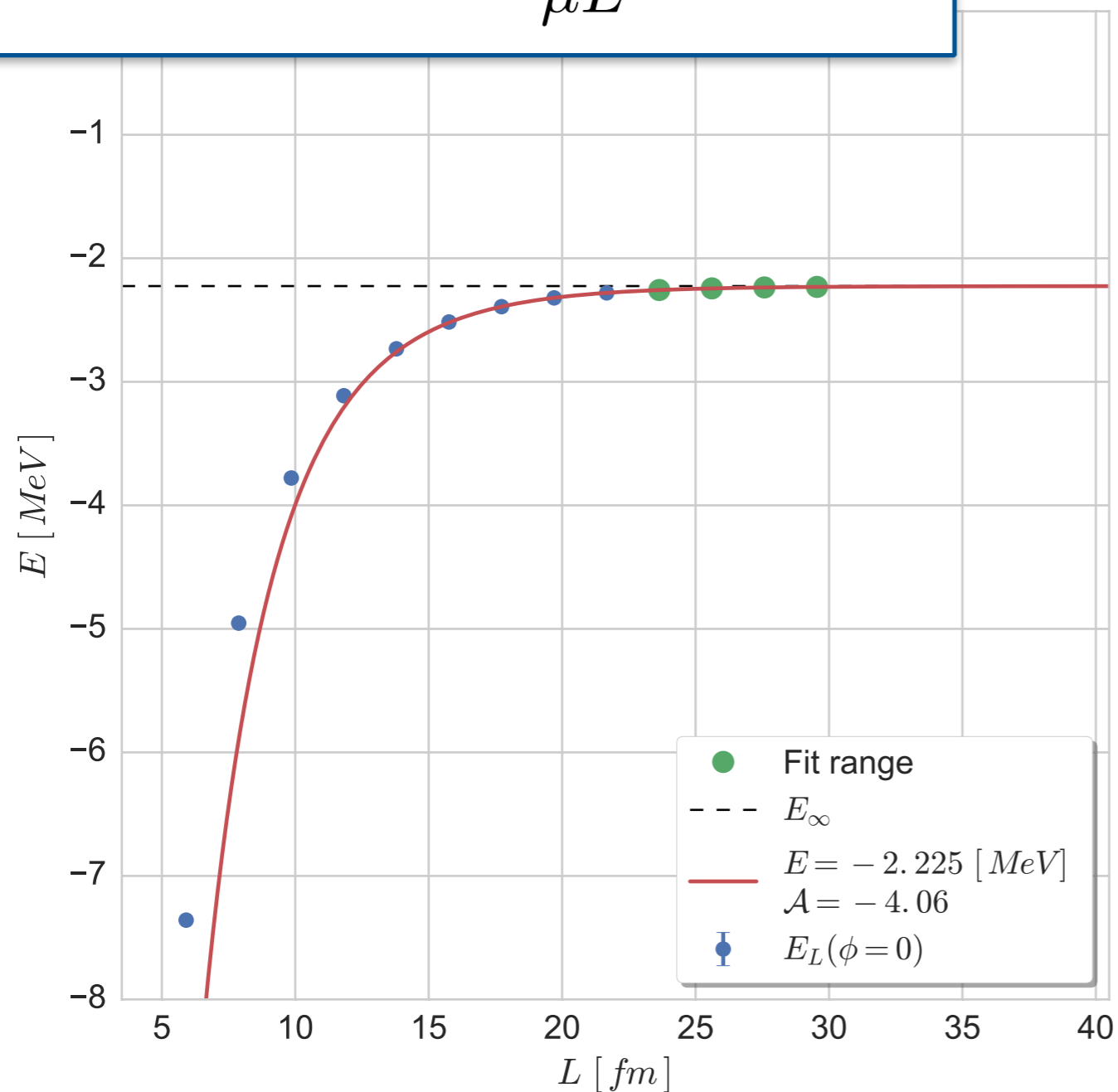
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- ▶ **Goal:**
Reliable results for small boxes
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Errors:
Numerical errors
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Results for small boxes affect fitting

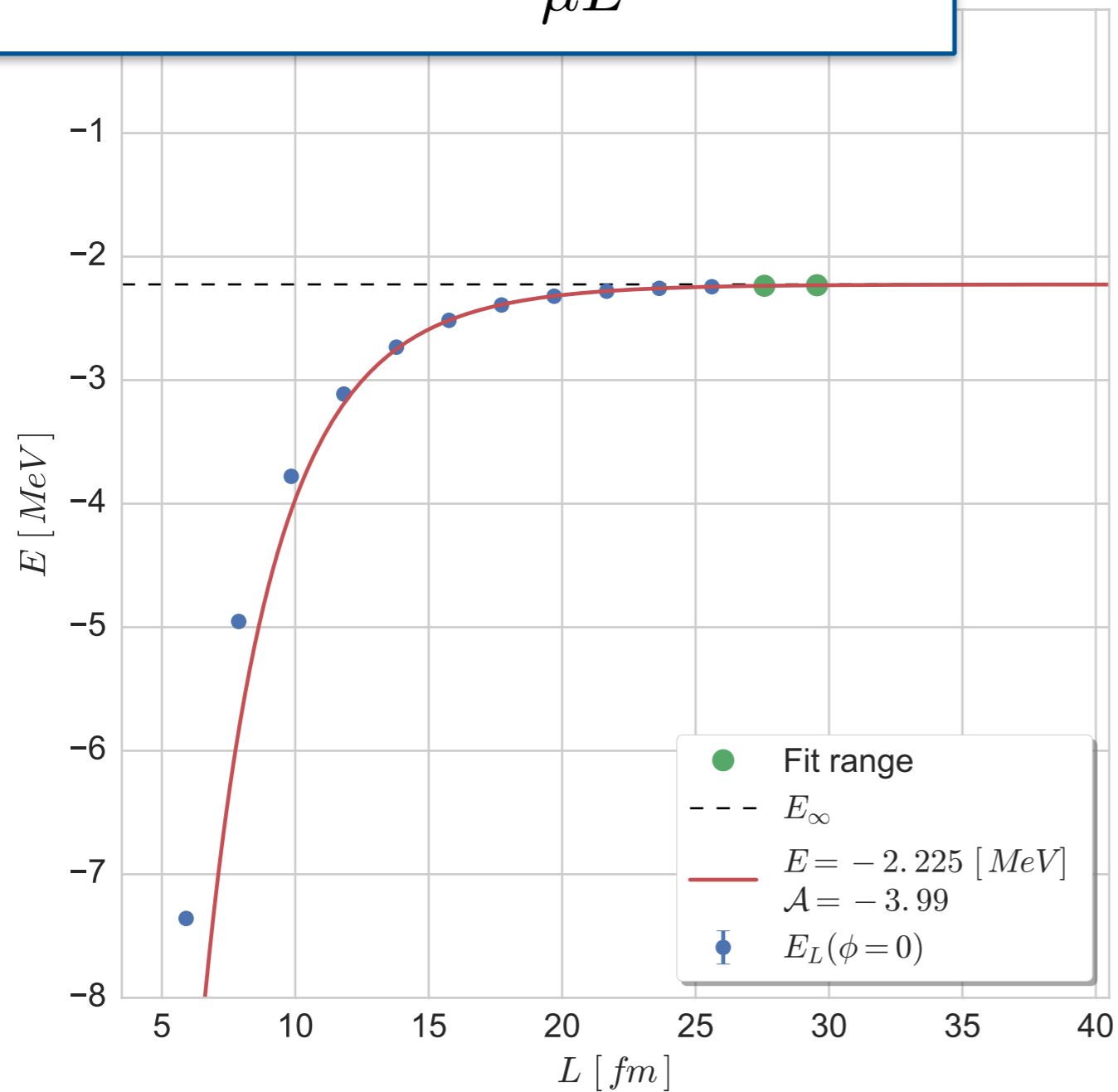
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- ▶ **Goal:**
Reliable results for small boxes
- ▶ **Data uncertainties**
Errors:
Numerical errors
- ▶ **Problem:**
Results for small boxes affect fitting

$$E_L(E, \mathcal{A}) \stackrel{!}{=} A_0 \frac{e^{-\kappa(E)L}}{\mu L} \mathcal{A} + E$$

$$E = -2.225 \text{ [MeV]}$$
$$\mathcal{A} = -3.99$$



- ▶ **Goal:**
Reliable results for small boxes
- ▶ **Data uncertainties Errors:**
Numerical errors
- ▶ **Problem:**
Results for small boxes affect fitting
- ▶ **Solution:**
'Re-weight' data points

Error Analysis

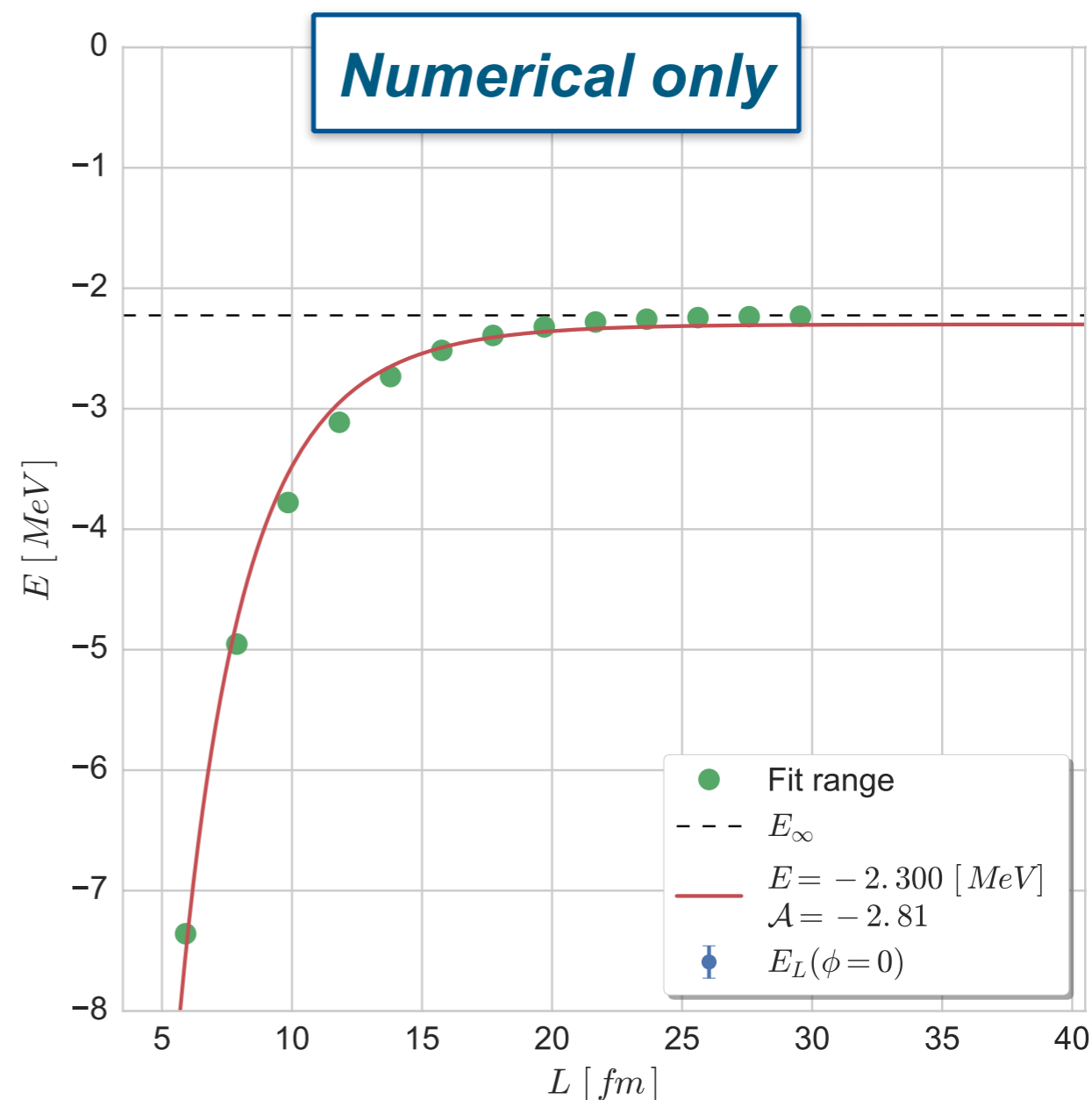
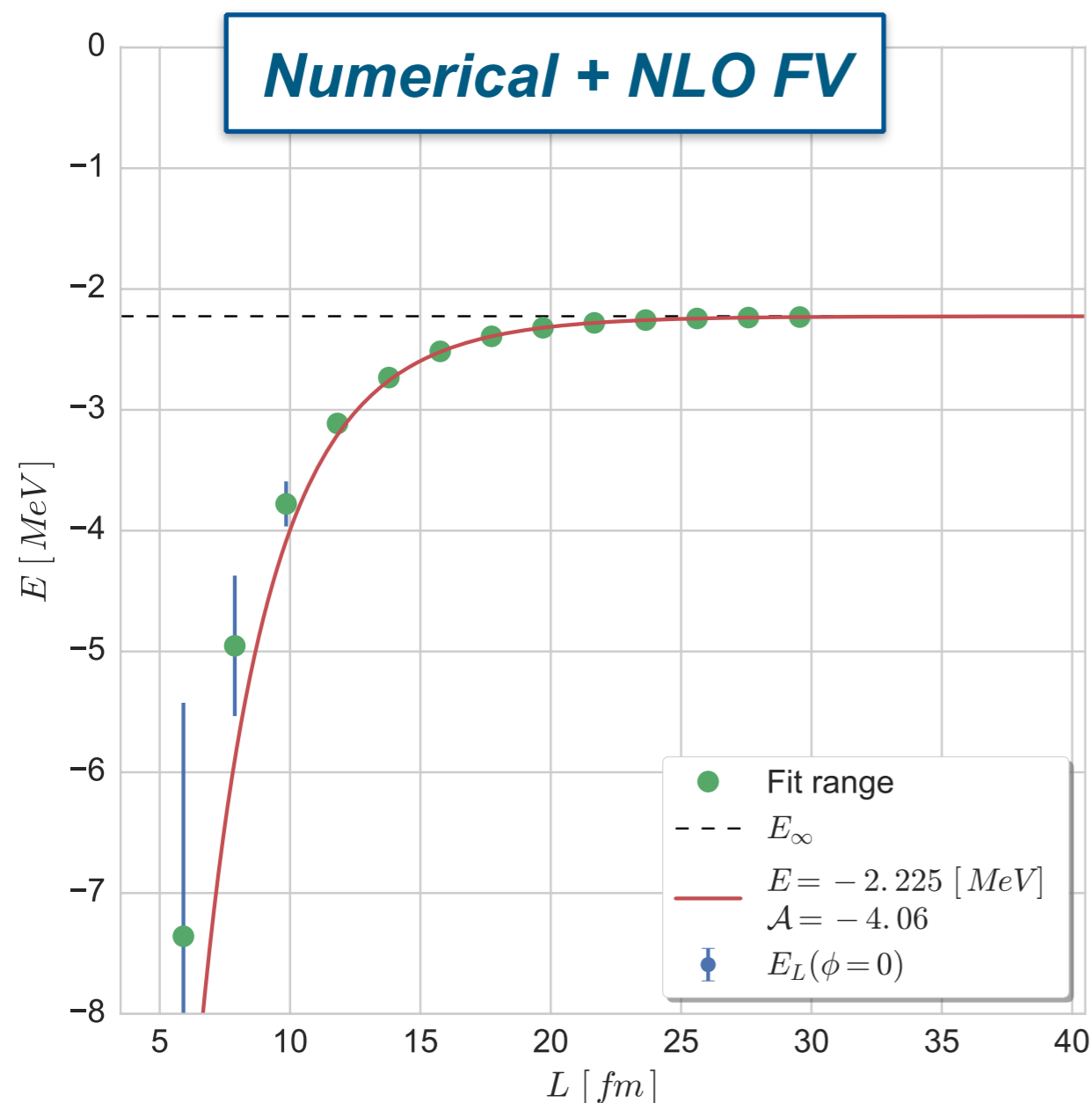
$$\Delta E_L^{(NLO)} \simeq A_0 \frac{e^{-\sqrt{2}\kappa L}}{\mu L}$$

$$E_L(L, \phi) - E_\infty \simeq \Delta E_L^{(LO)}(L, \phi) + \Delta E_L^{(NLO)}(L, \phi)$$

$$=: \Delta E_L^{(LO)}(L, \phi) + \delta(\Delta E_L^{(LO)}(L, \phi))$$

$$E = -2.225 \text{ [MeV]}$$

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Error Analysis

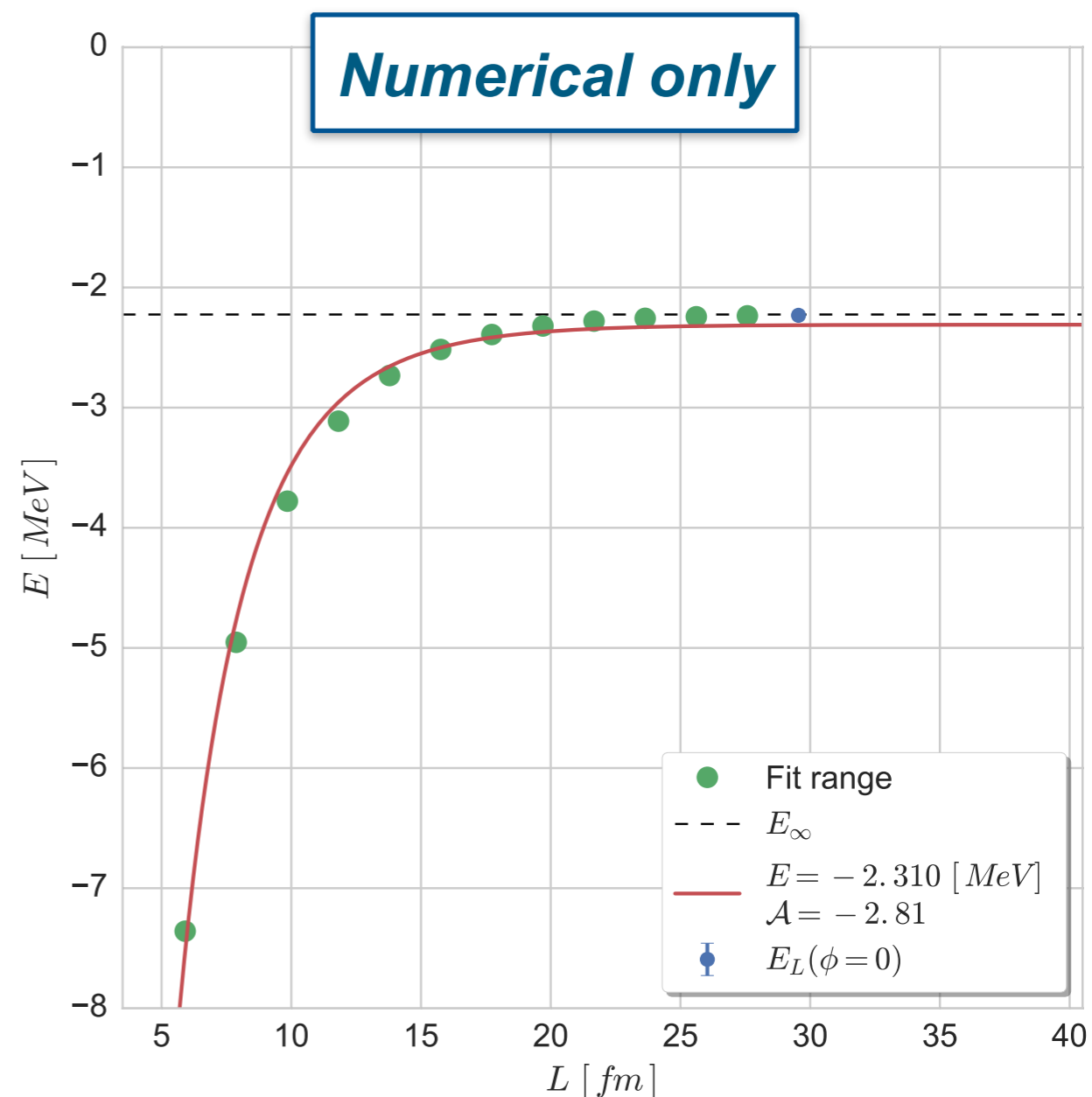
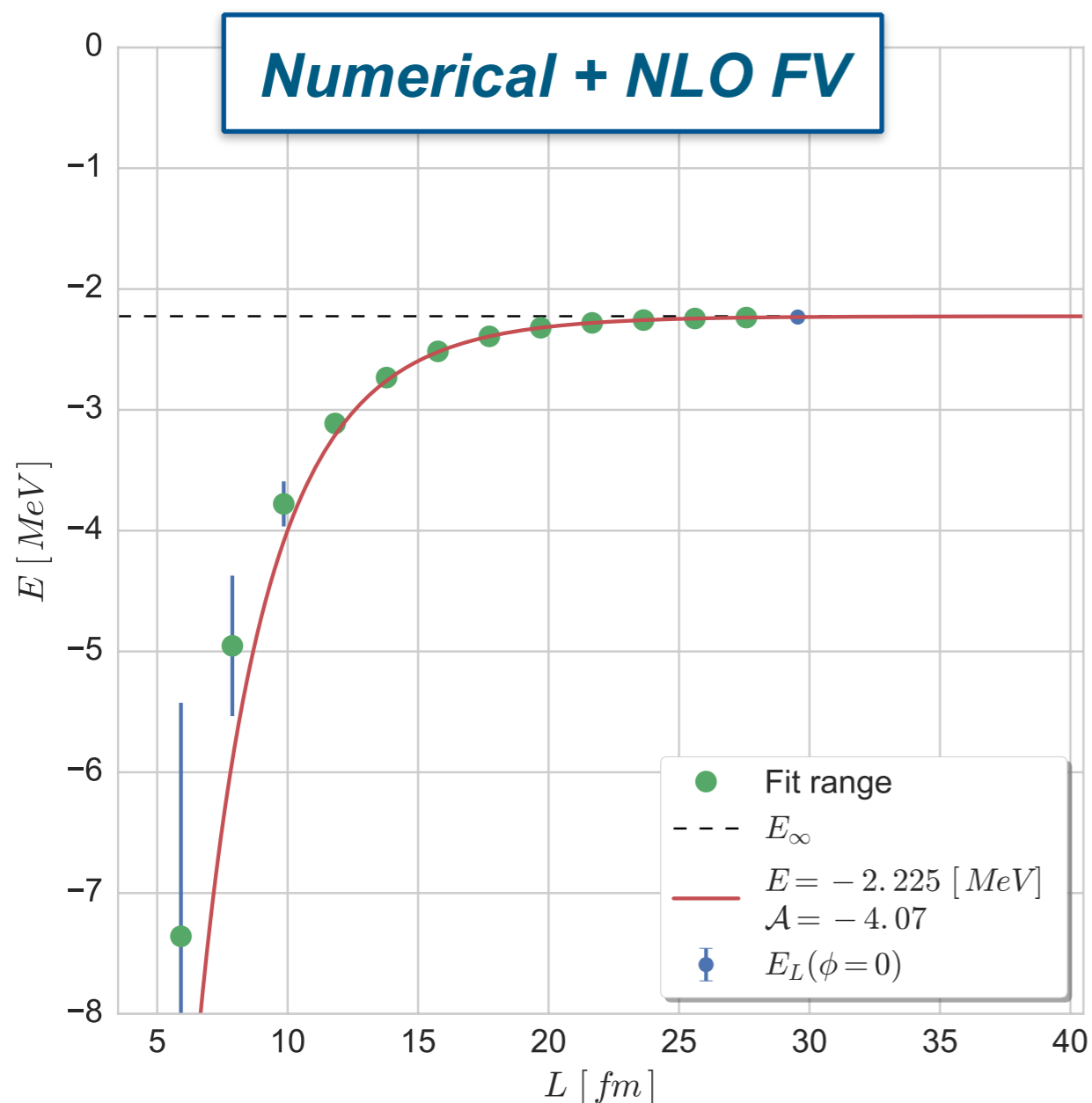
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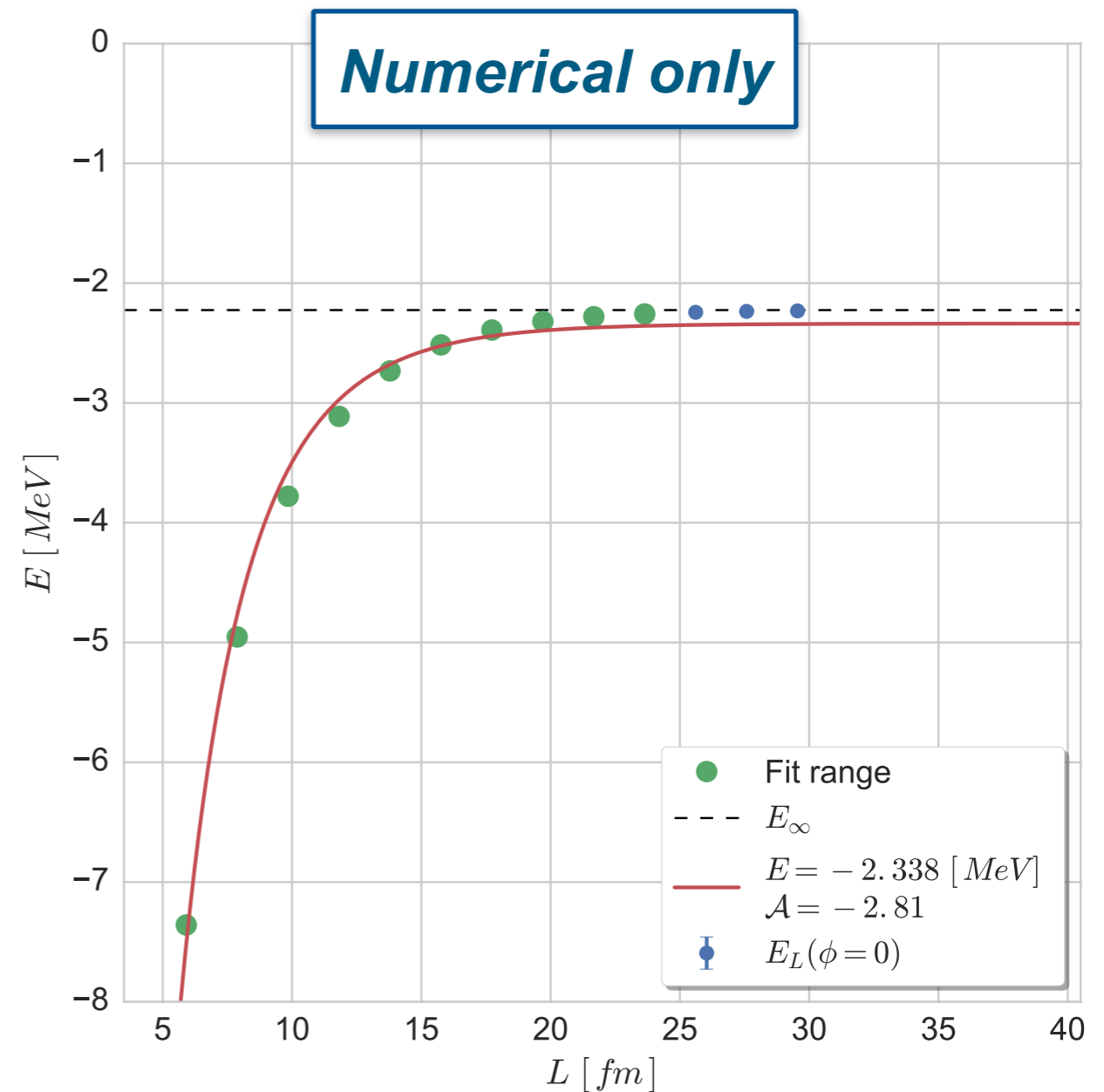
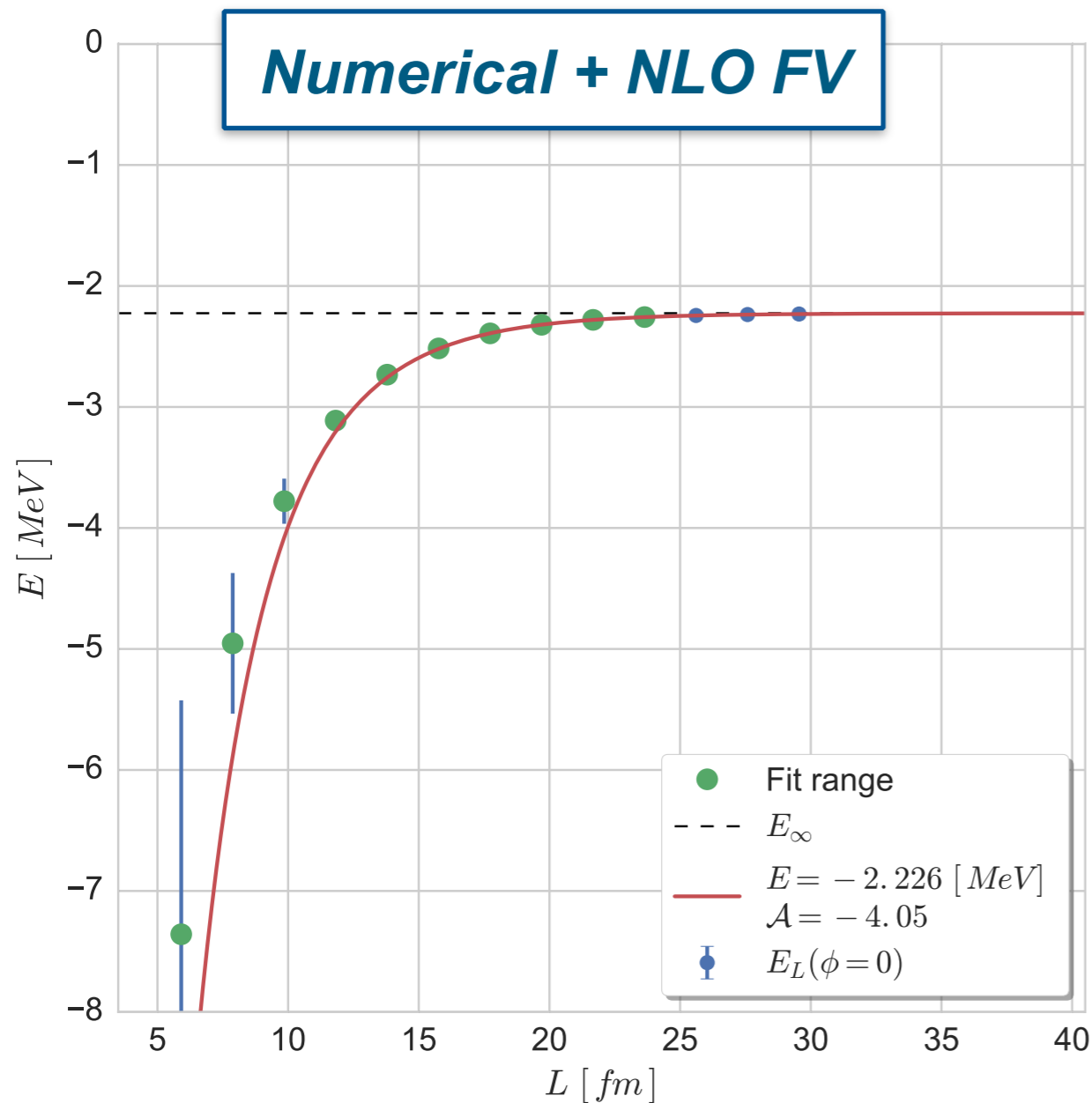
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Error Analysis

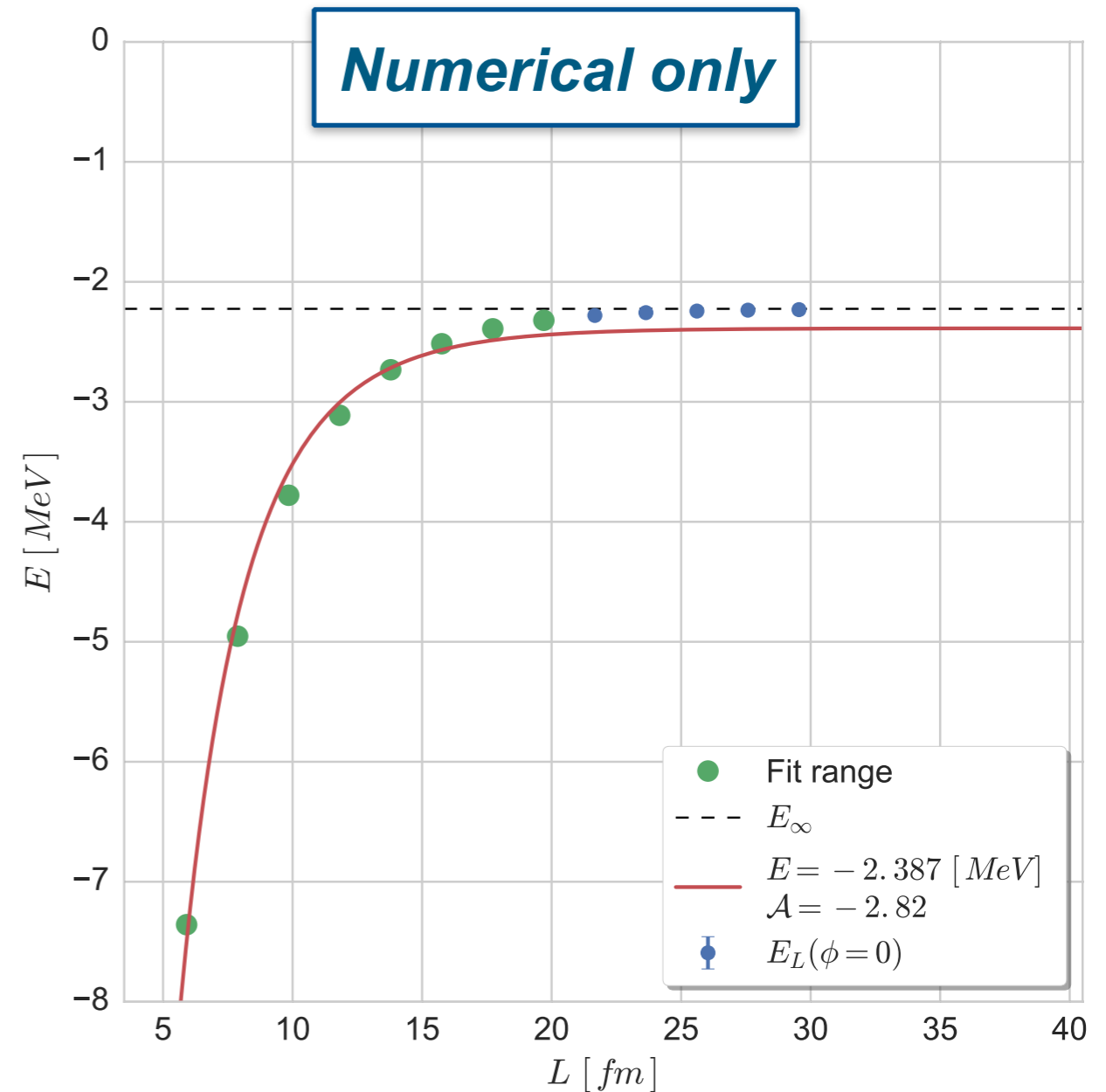
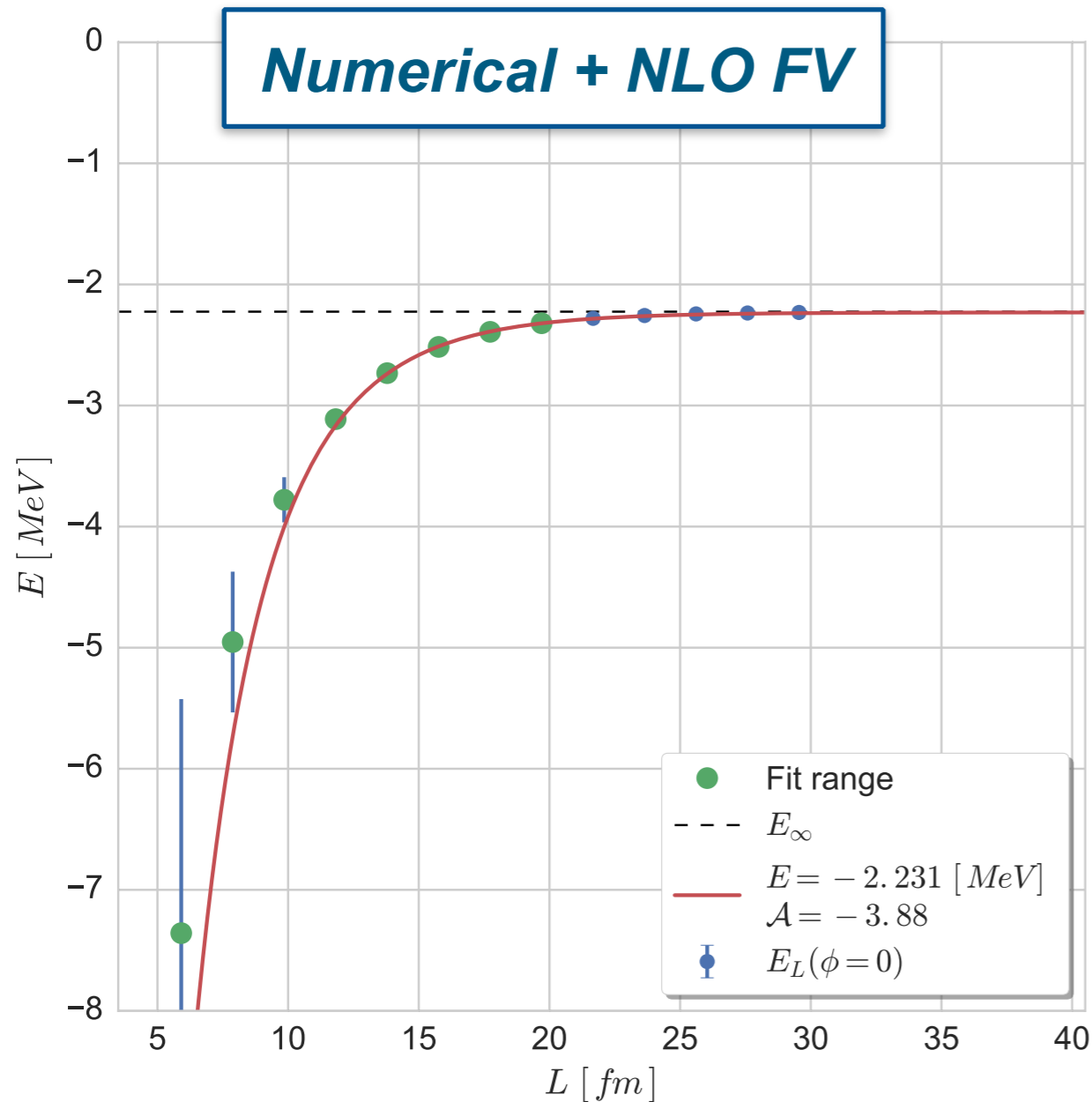
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Error Analysis

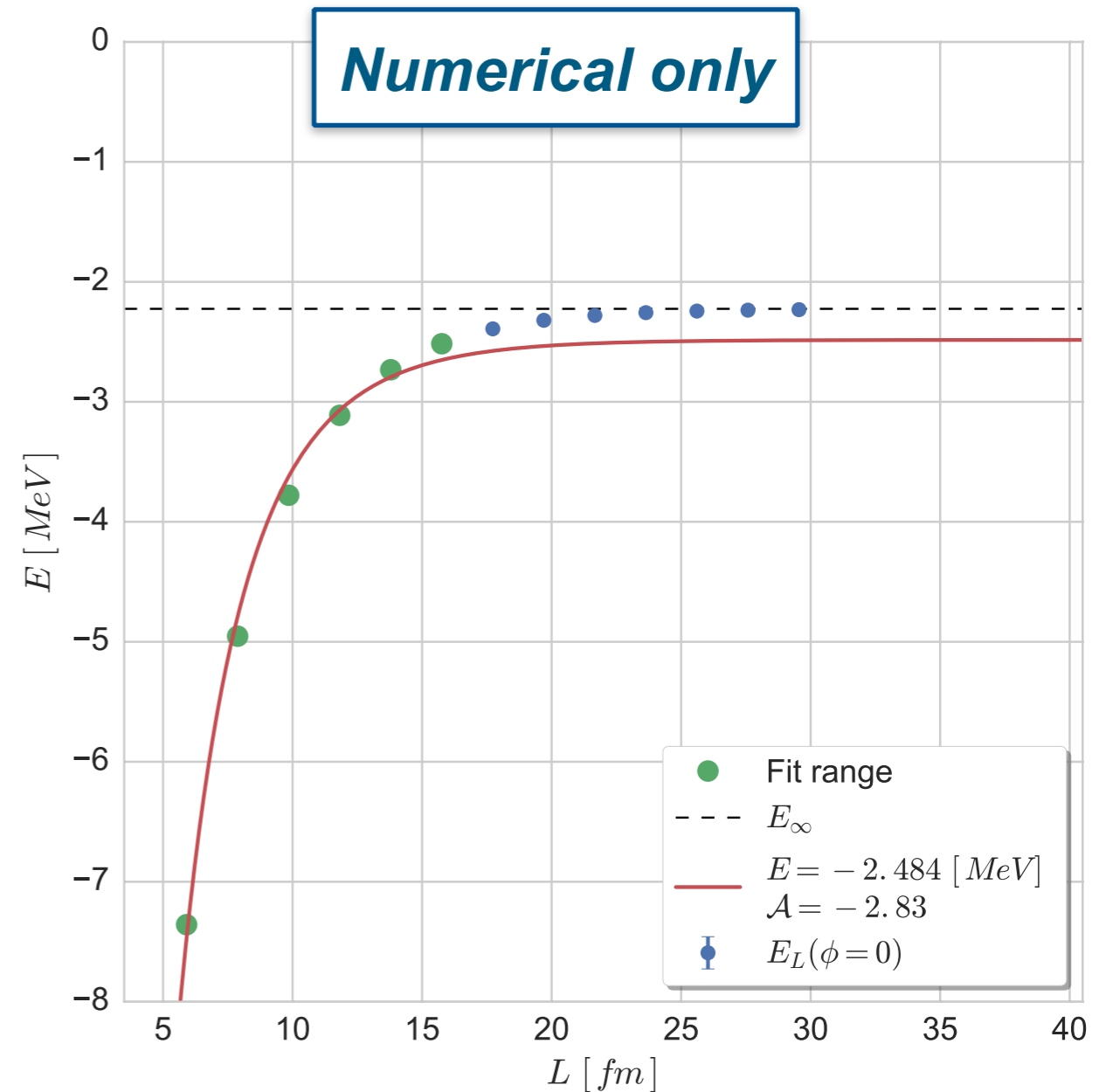
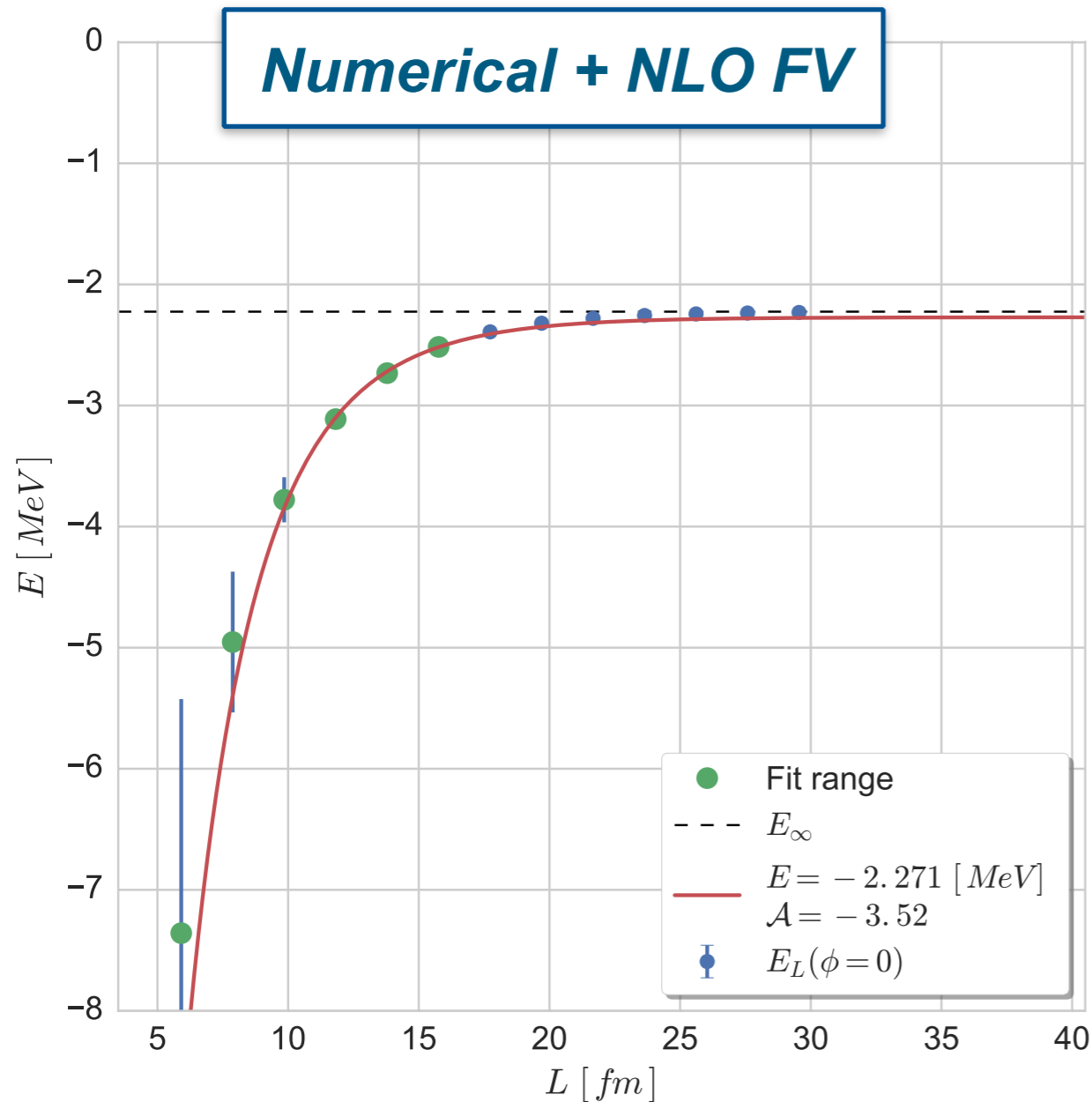
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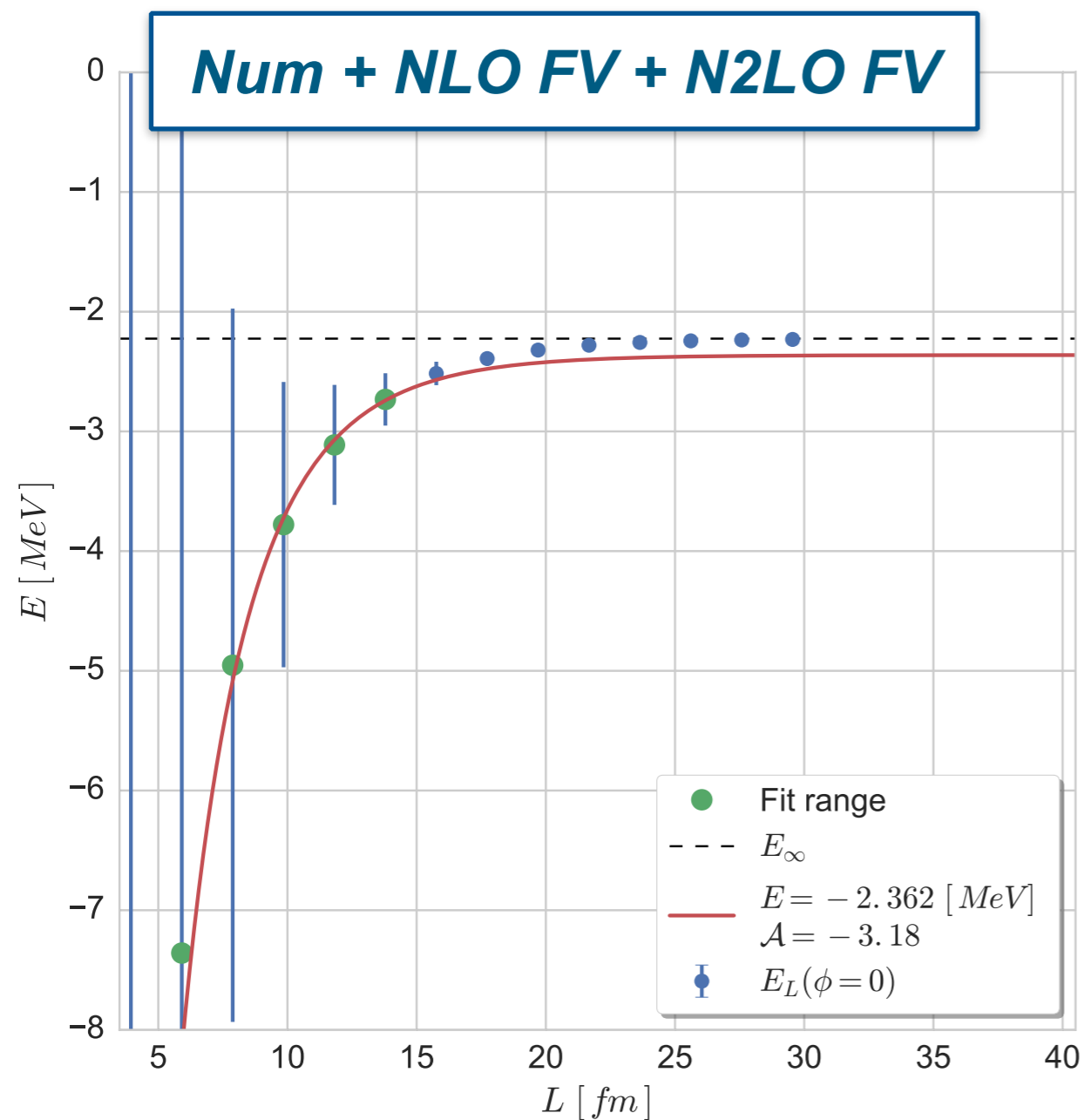
$$\mathcal{A} = -3.99$$



- **Estimate uncertainty:**
compute new correlated distributions within error bars for propagation

$$E = -2.225 \text{ [MeV]}$$

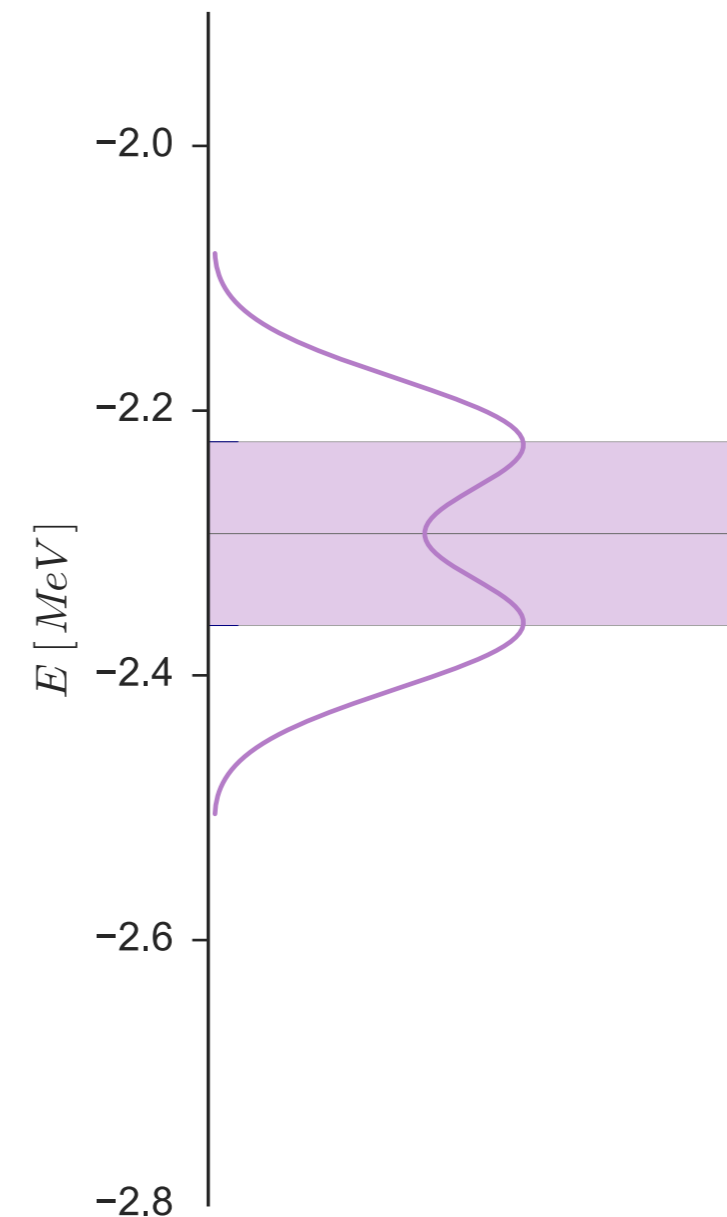
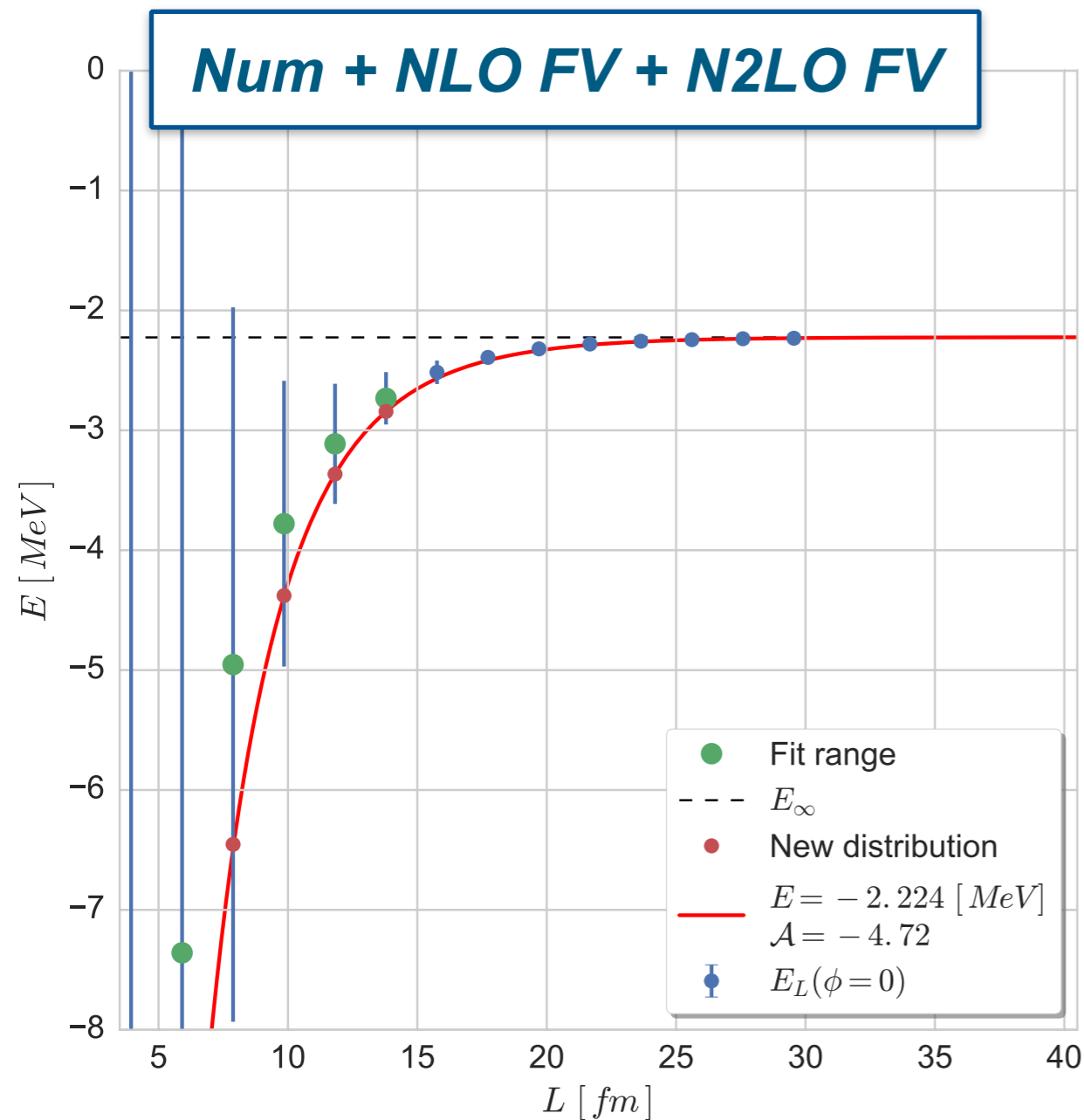
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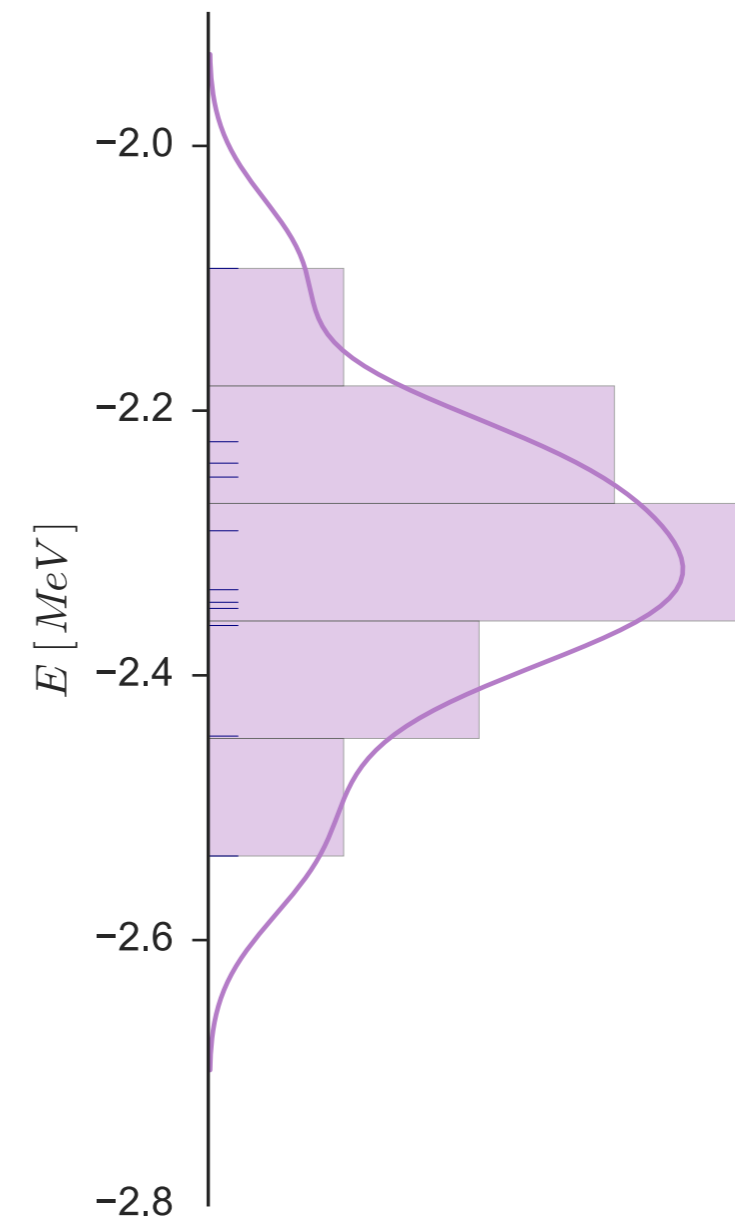
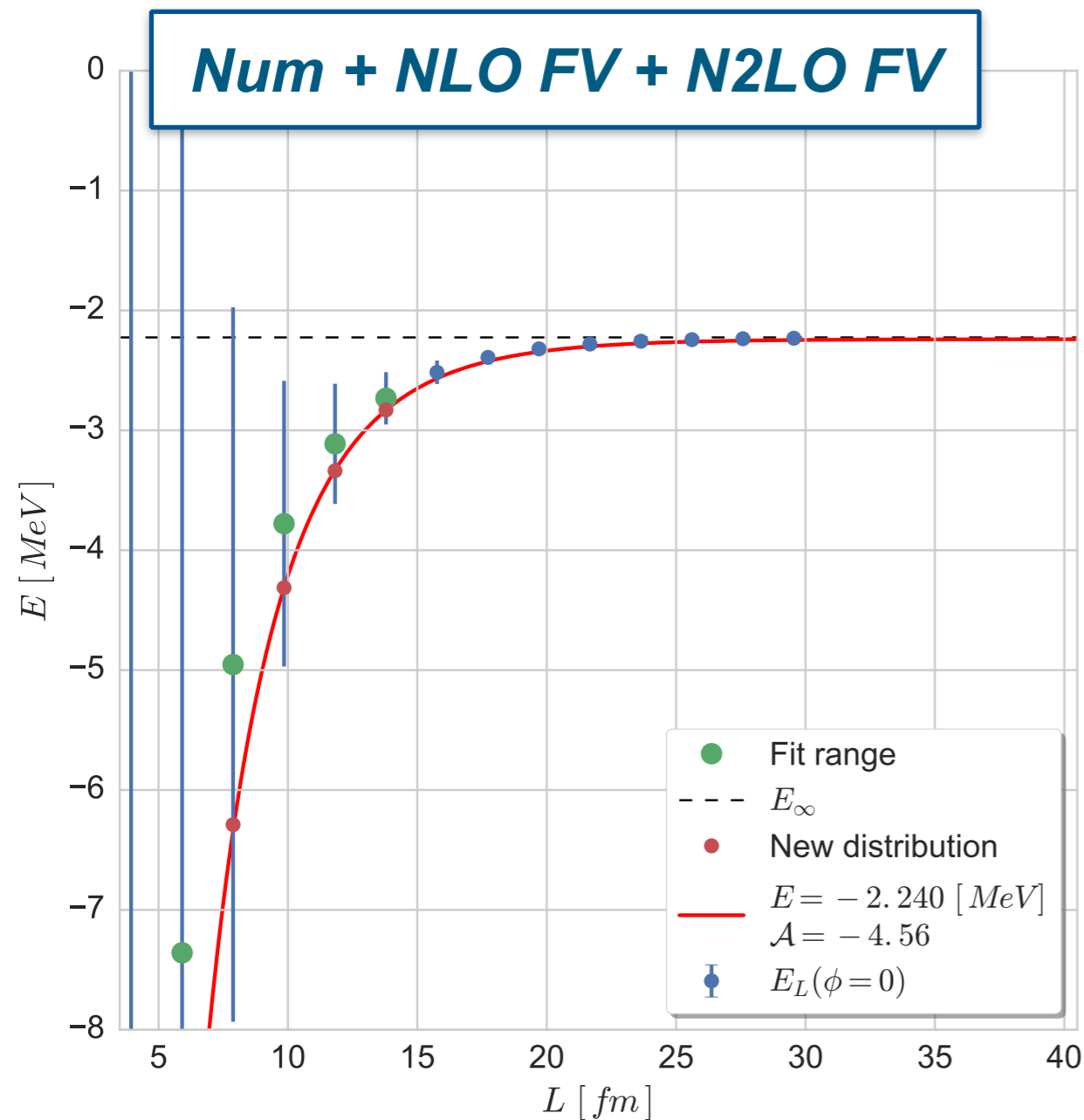
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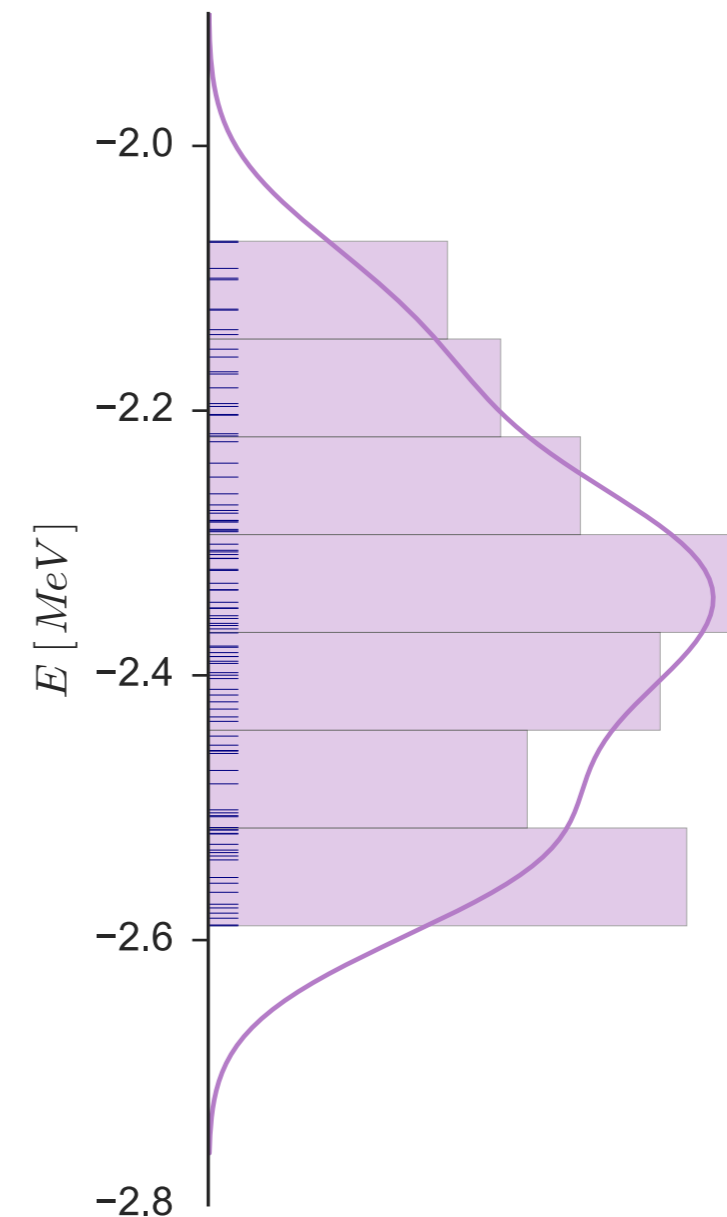
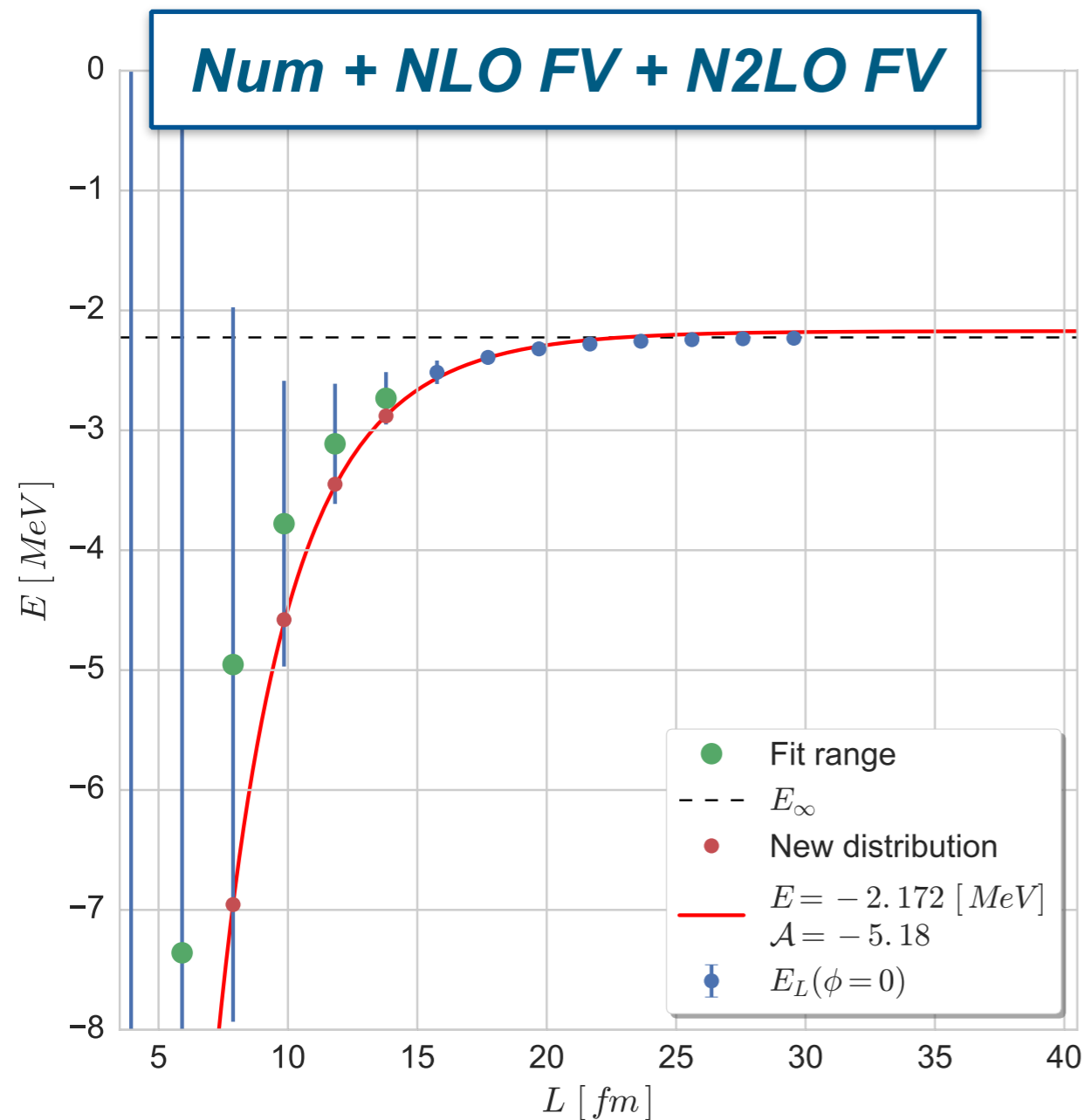
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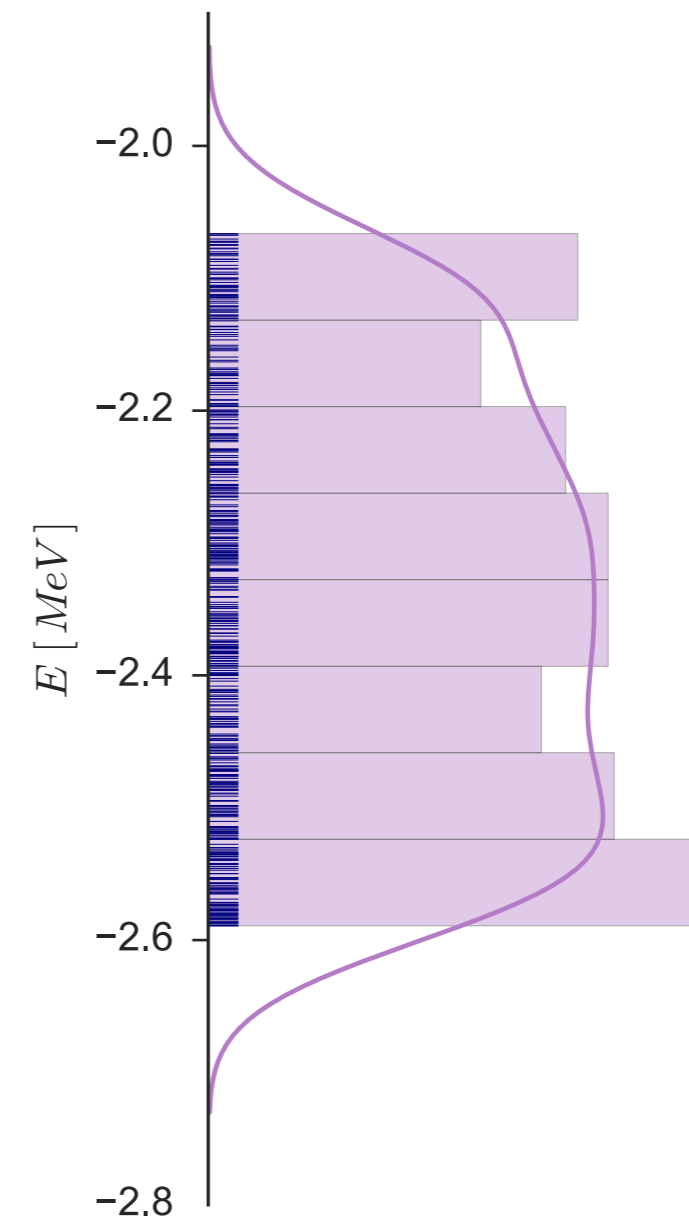
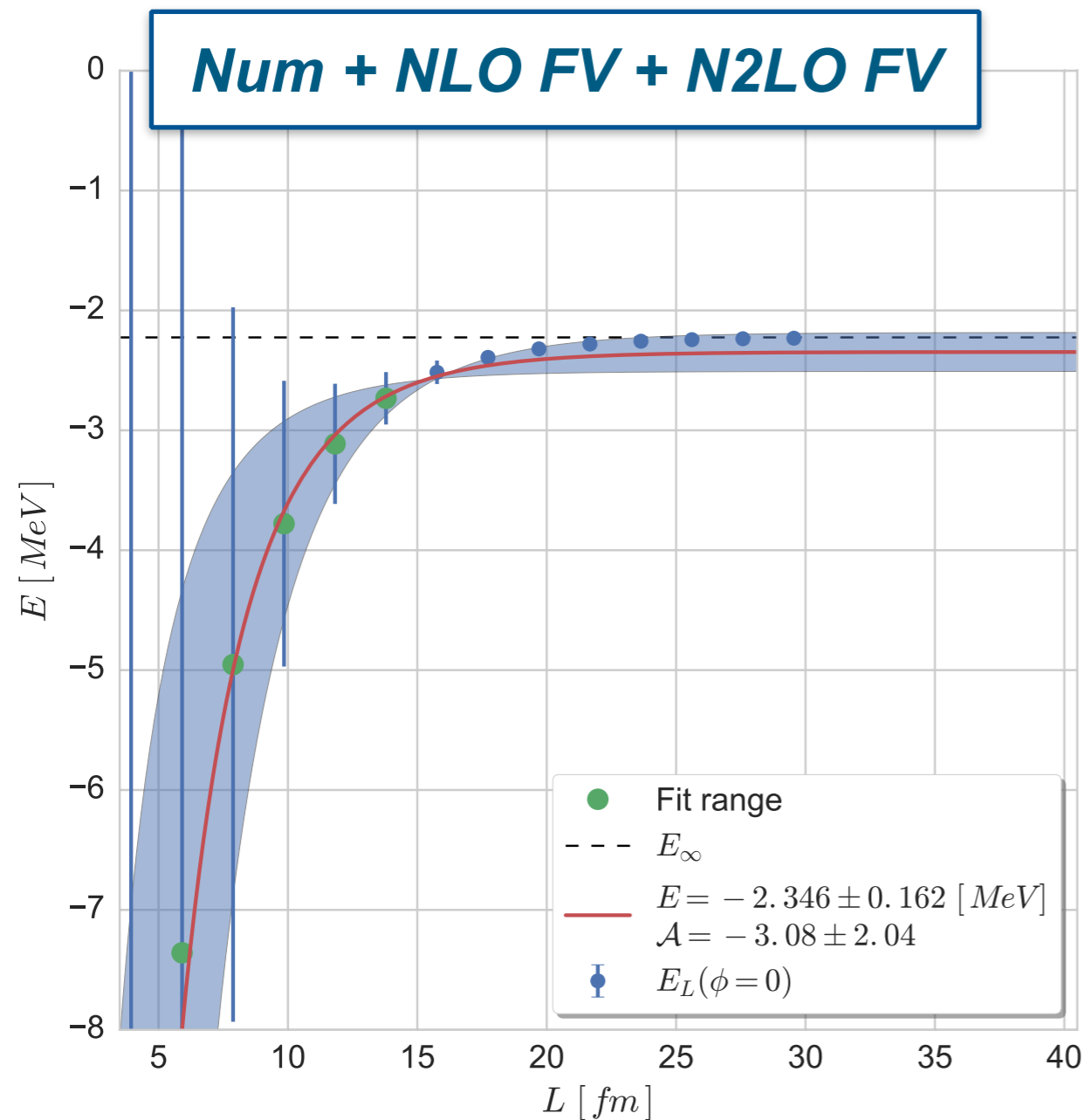
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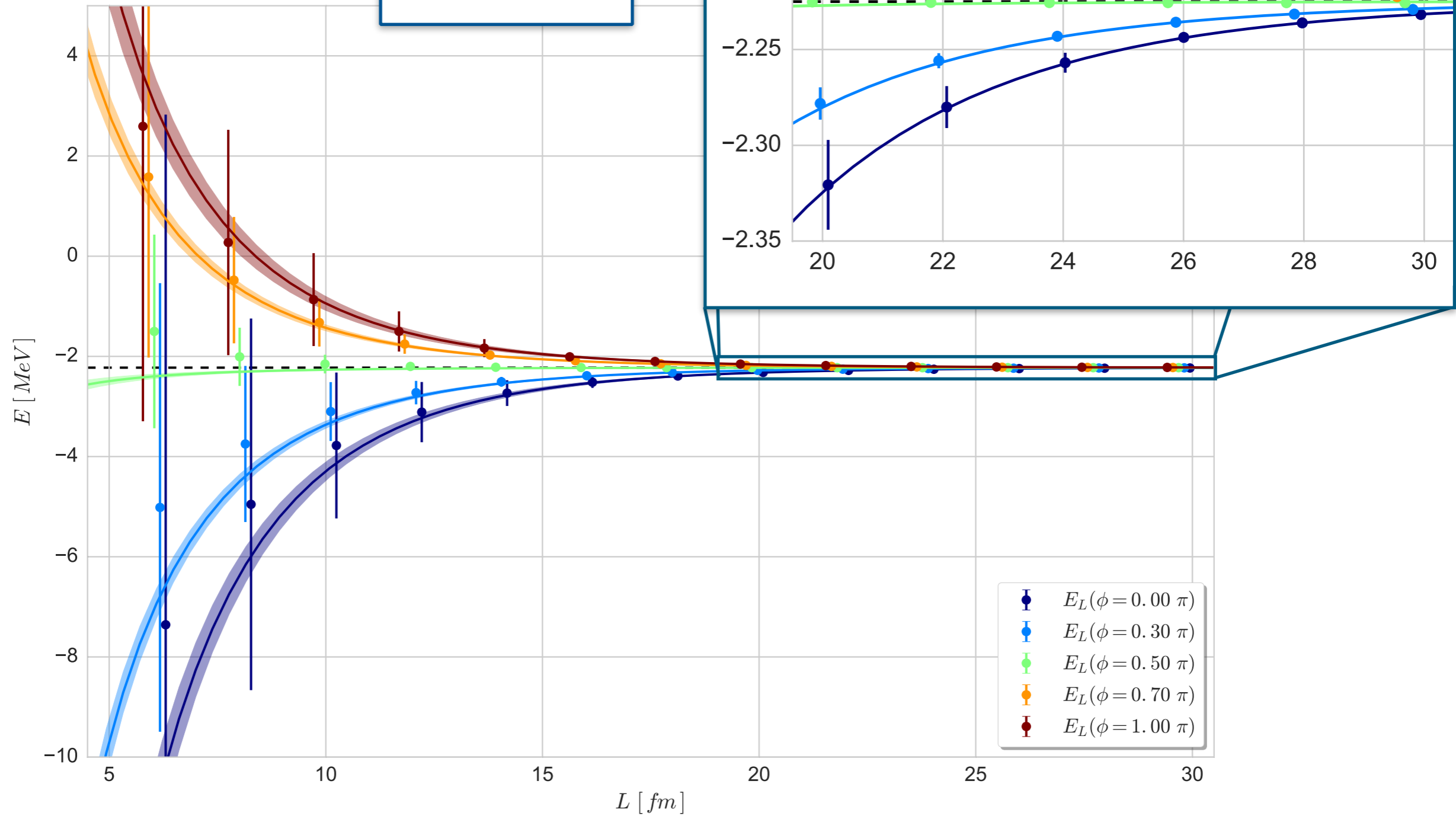


Two-Body Results

Numerical Results:

Two-body level

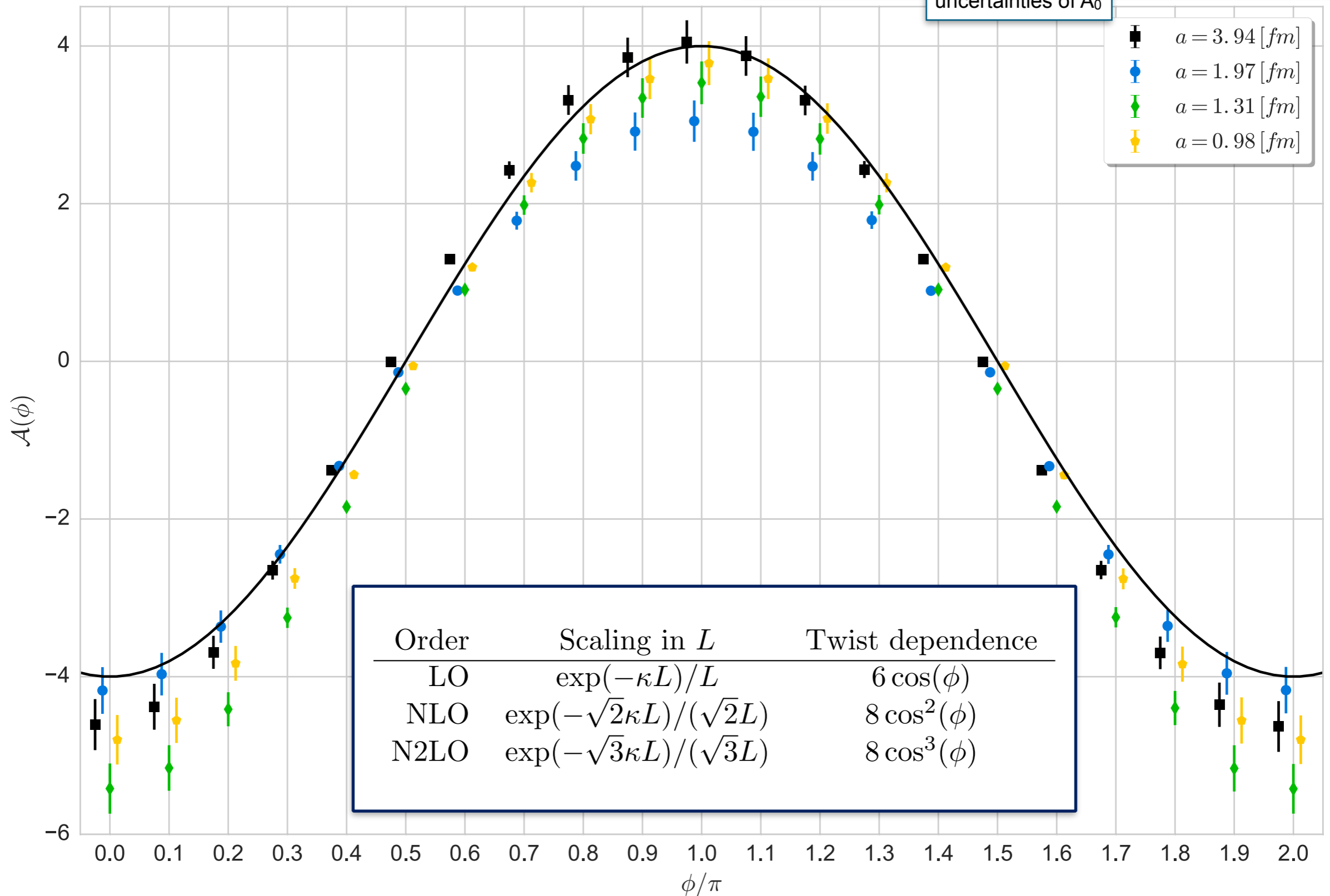
$a = 1.97 \text{ fm}$



Numerical Analysis

$$E_L(E, \mathcal{A}) \stackrel{!}{=} A_0 \frac{e^{-\kappa(E)L}}{\mu L} \mathcal{A} + E$$

Not including uncertainties of A_0



The Three-Body Case

3-Body — What are the twists?

“Spectrum of three-body bound states in a finite volume”,
U. G. Meißner, G. Ríos, A. Rusetsky
[arXiv:1412.4969],
Phys.Rev.Lett. **114**

$$\Delta E_L^{(LO)}(L, \{\vec{\phi}_i = \vec{0}\}) = \sum_{i=1}^3 \sum_{(\vec{n}_i, \vec{n}_j, \vec{n}_k) \in M_i} v(\vec{n}_i, \vec{n}_j, \vec{n}_k)$$

Set which minimizes relative hyper radius

$$v(\vec{n}_i, \vec{n}_j, \vec{n}_k) := \int d^3\vec{x}_i d^3\vec{y}_i \psi_\infty^*(\vec{x}_i, \vec{y}_i) V_i(x_i) \psi_\infty(\vec{x}_i - (\vec{n}_j + \vec{n}_k)L, \vec{y}_i + \frac{1}{\sqrt{3}}(\vec{n}_j + \vec{n}_k - 2\vec{n}_i)L)$$

For twisted boundaries

Jacobi coordinate shift expressed in one particle coordinates

$$\Delta E_L^{(LO)}(L, \{\vec{\phi}_i\}) = \sum_{i=1}^3 \sum_{(\vec{n}_i, \vec{n}_j, \vec{n}_k) \in M_i} v(\vec{n}_i, \vec{n}_j, \vec{n}_k) e^{-i \sum_{l=1}^3 \vec{\phi}_l \cdot \vec{n}_l}$$

Result at unitary limit

$$\Delta E_L^{(LO)}(L, \{\vec{\phi}_i\}) = \frac{\mathcal{N}_{PB}^{(LO)}}{9} \frac{\exp\left(-\frac{2}{\sqrt{3}}\kappa L\right)}{(\kappa L)^{3/2}} \sum_{i,j=1}^3 \cos(\vec{e}_j \cdot \vec{\phi}_i)$$

3-Body — What are the twists?

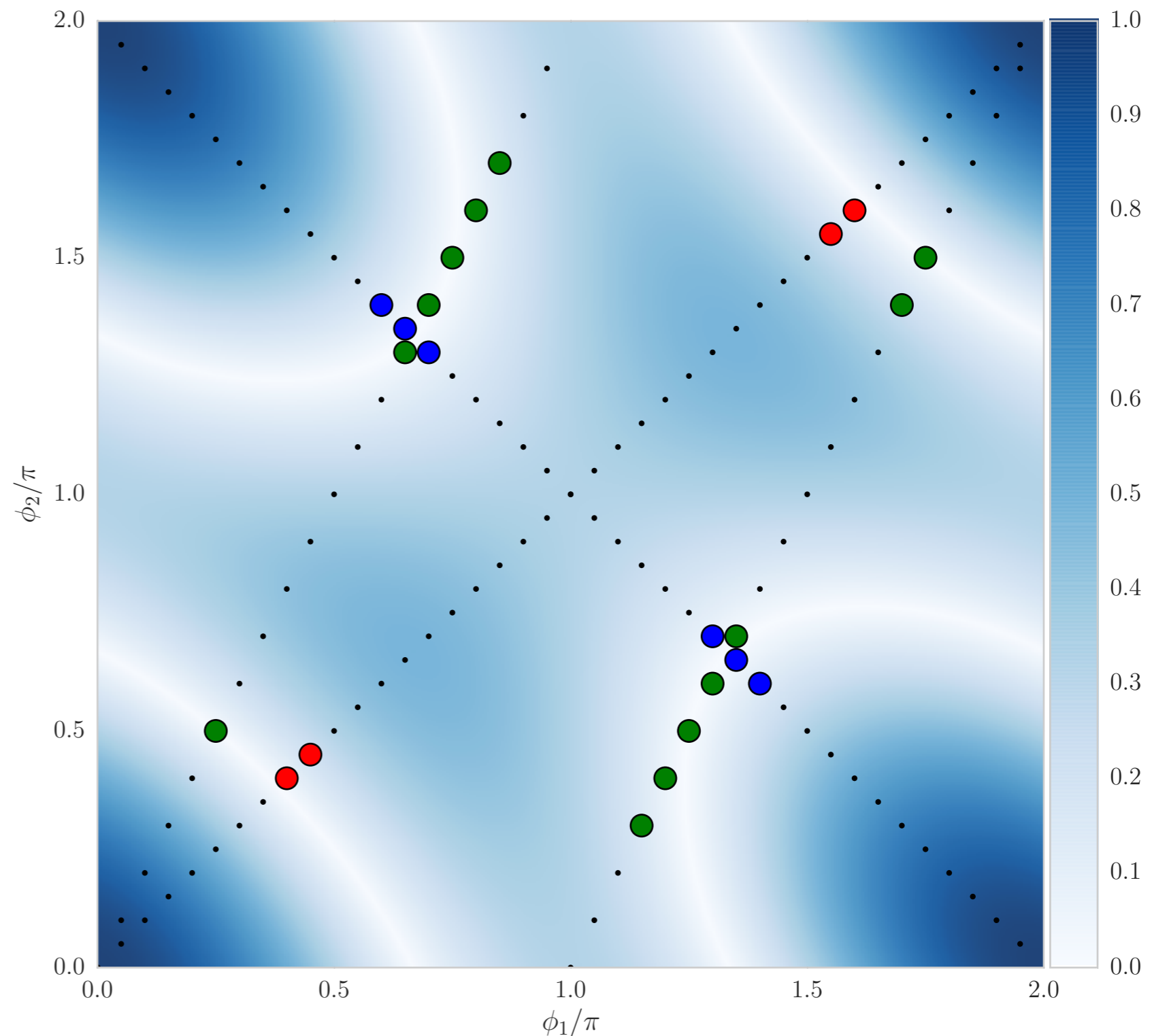
Relative LO Error Amplitude

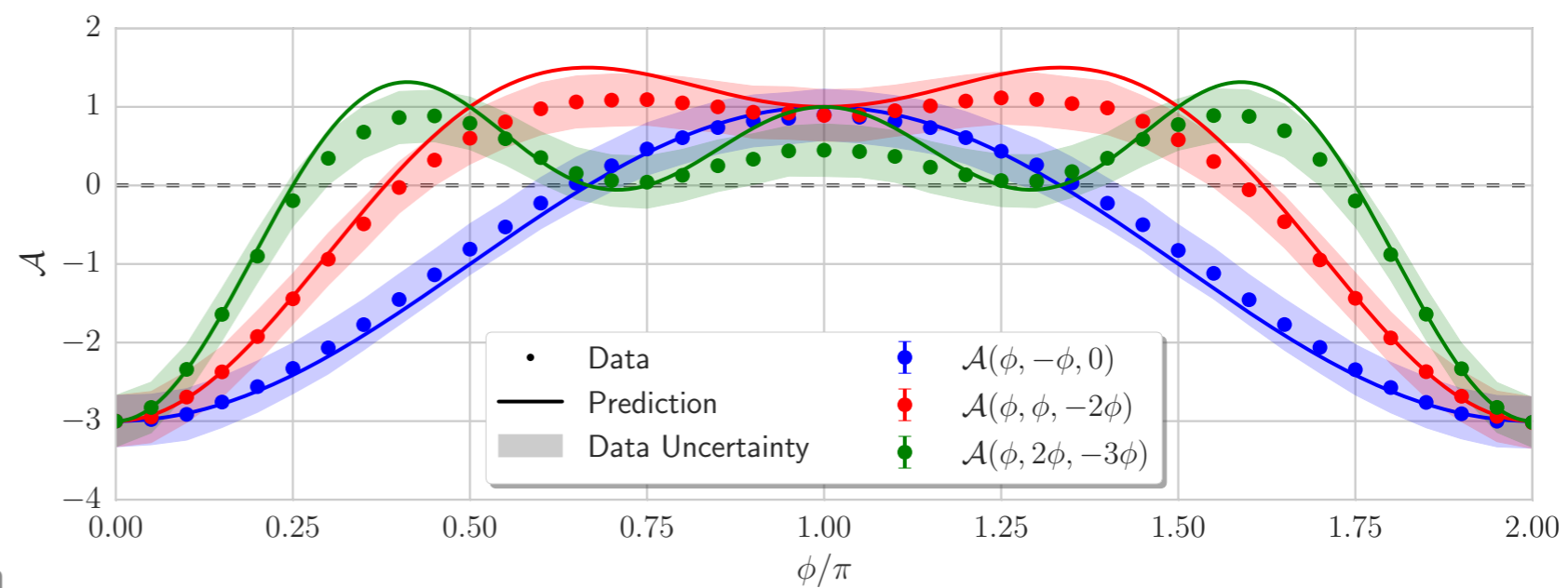
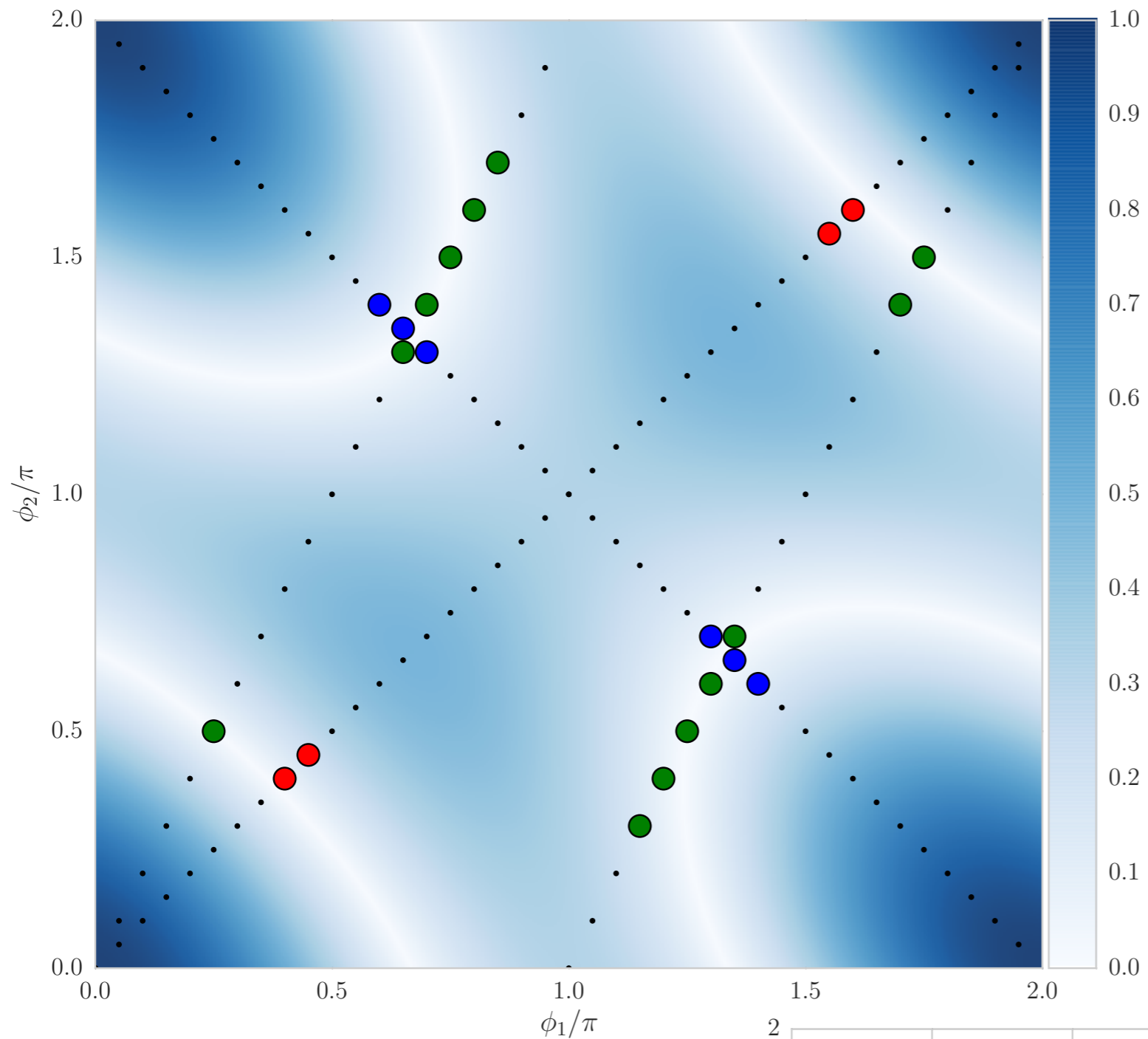
CMS constraint

$$\mathcal{A}(\phi_1, \phi_2) = \cos(\phi_1) + \cos(\phi_2) + \cos(\phi_1 + \phi_2)$$

Numerical Findings

- *Probed twists*
- *'iPB' consistent twists*

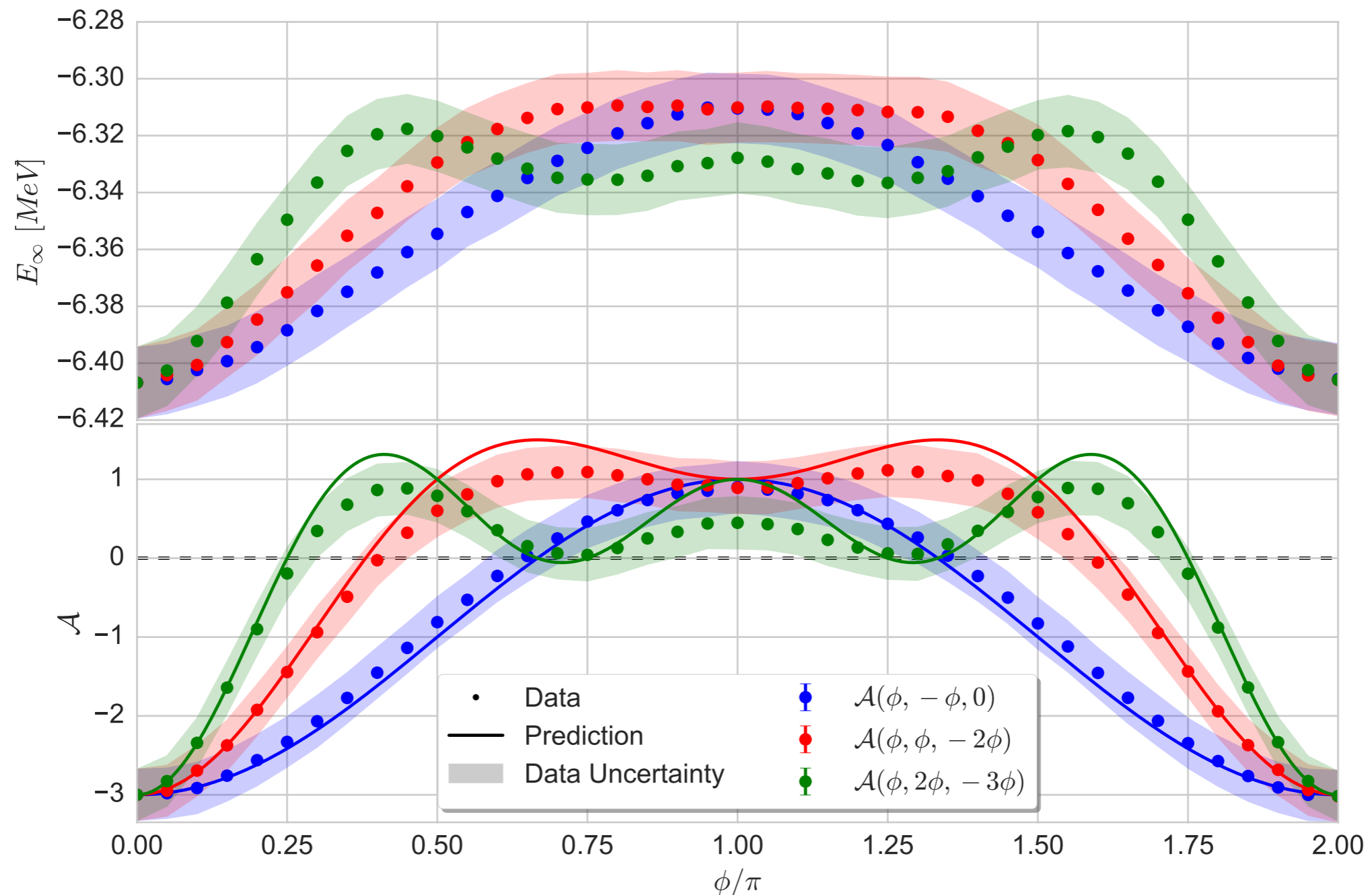




3-Body results

$$\Delta E_L^{(LO)}(L, \phi_i) = A(\phi_i) \frac{e^{-\frac{2}{\sqrt{3}}\kappa L}}{(\kappa L)^{3/2}}$$

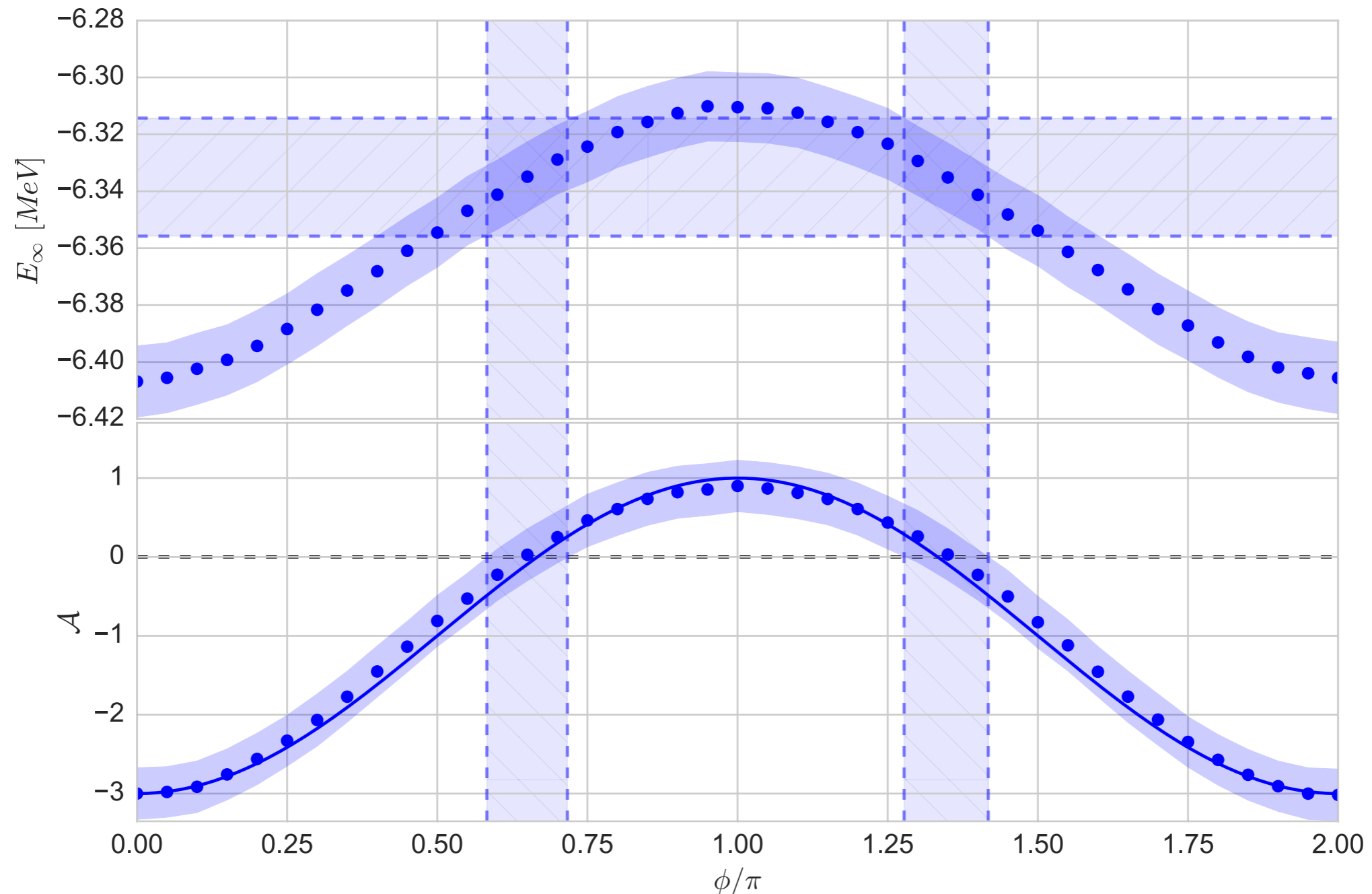
$$\mathcal{A}(\phi_i) := 3 \frac{A(\phi_i)}{A_{\max}} = \sum_{i=1}^3 \cos(\phi_i)$$



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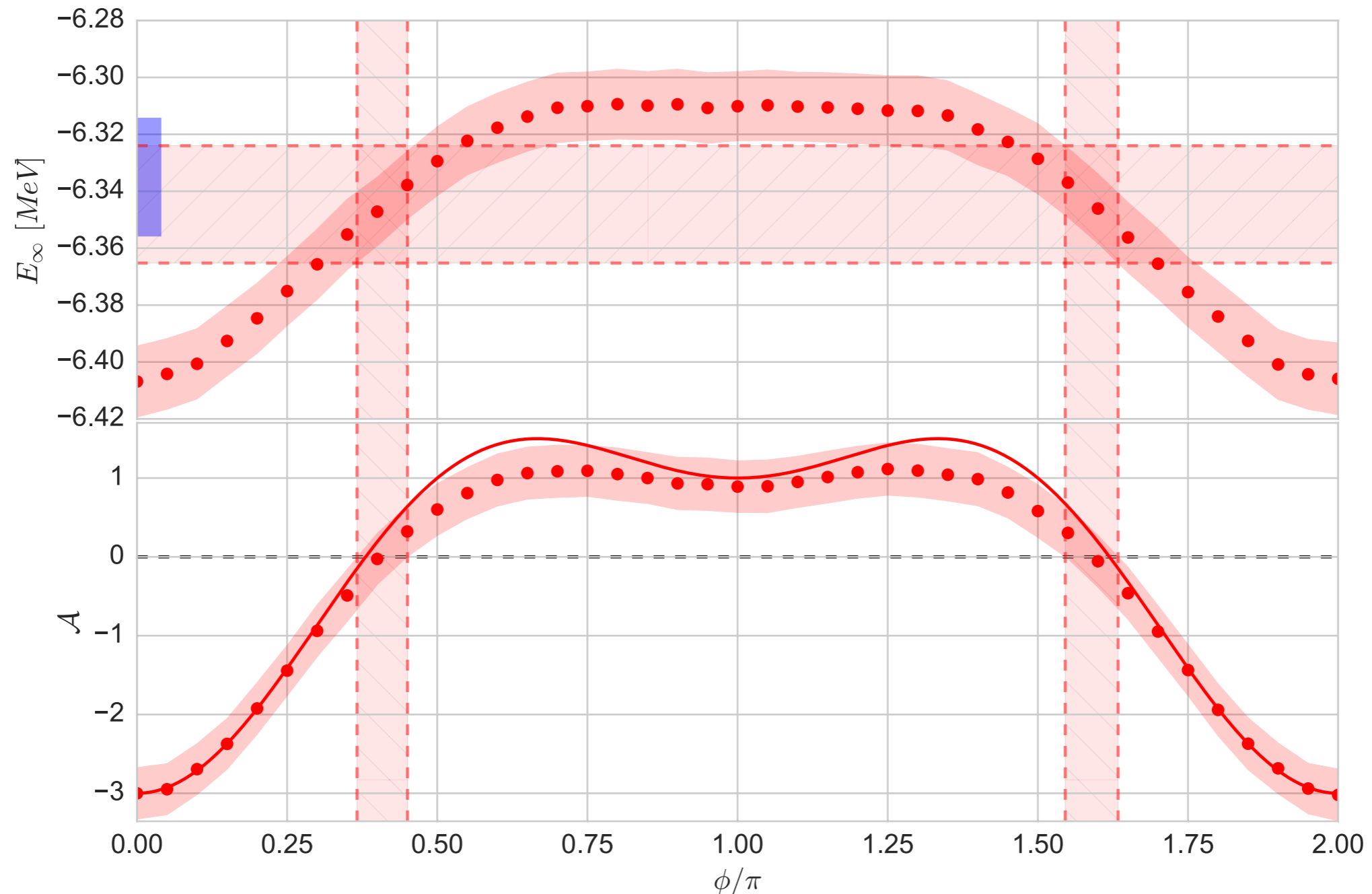
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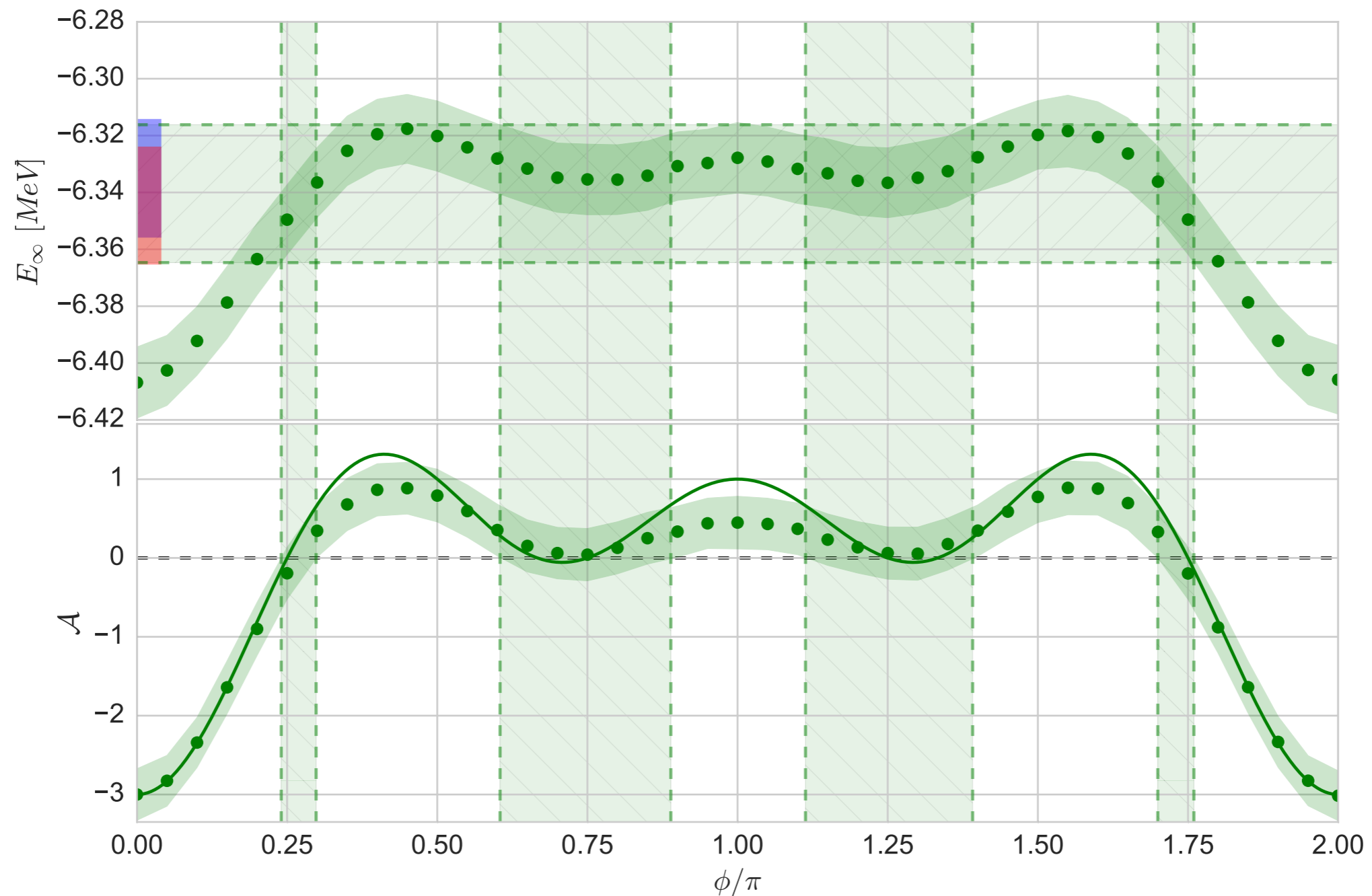
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3-Body results

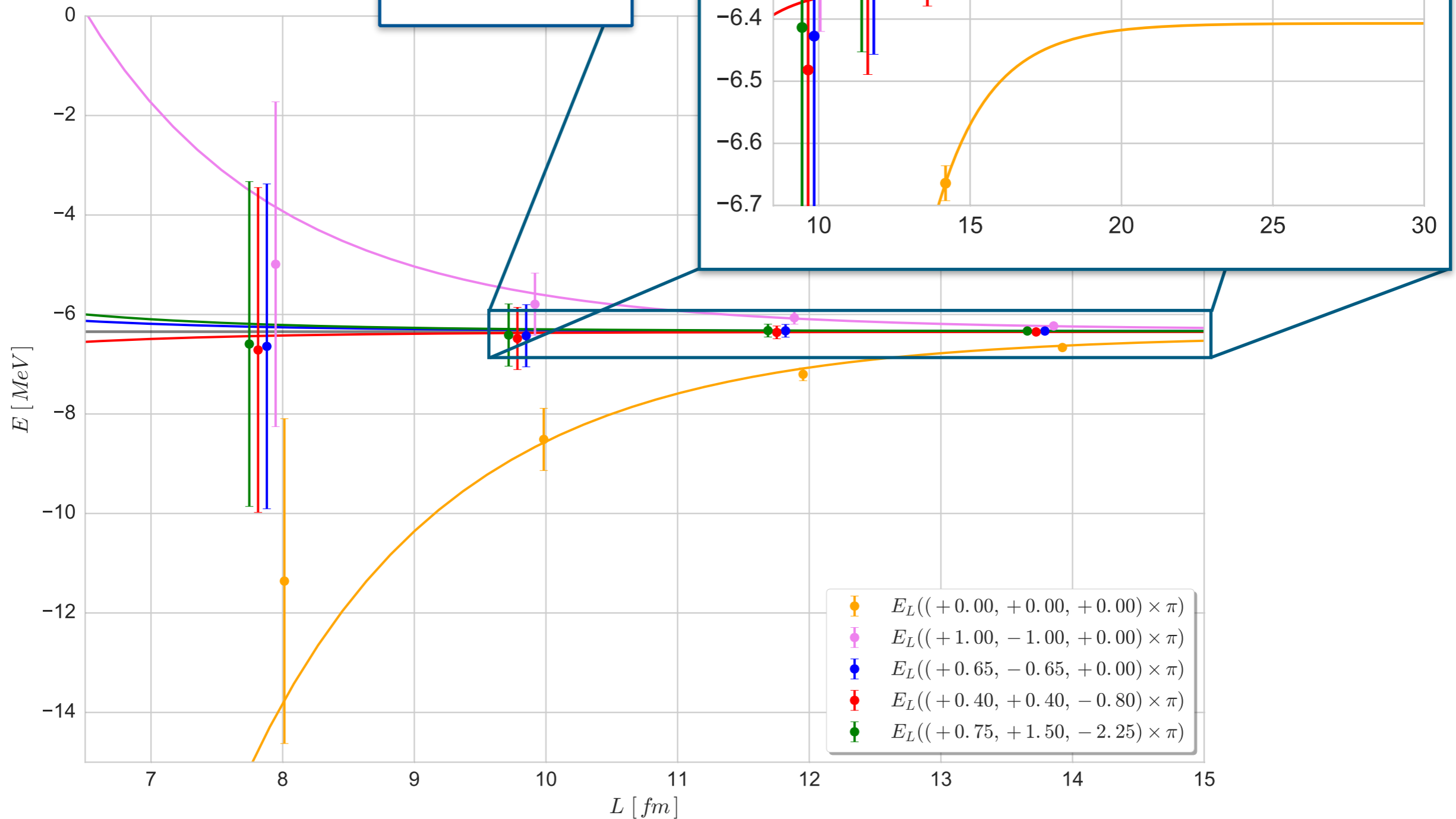
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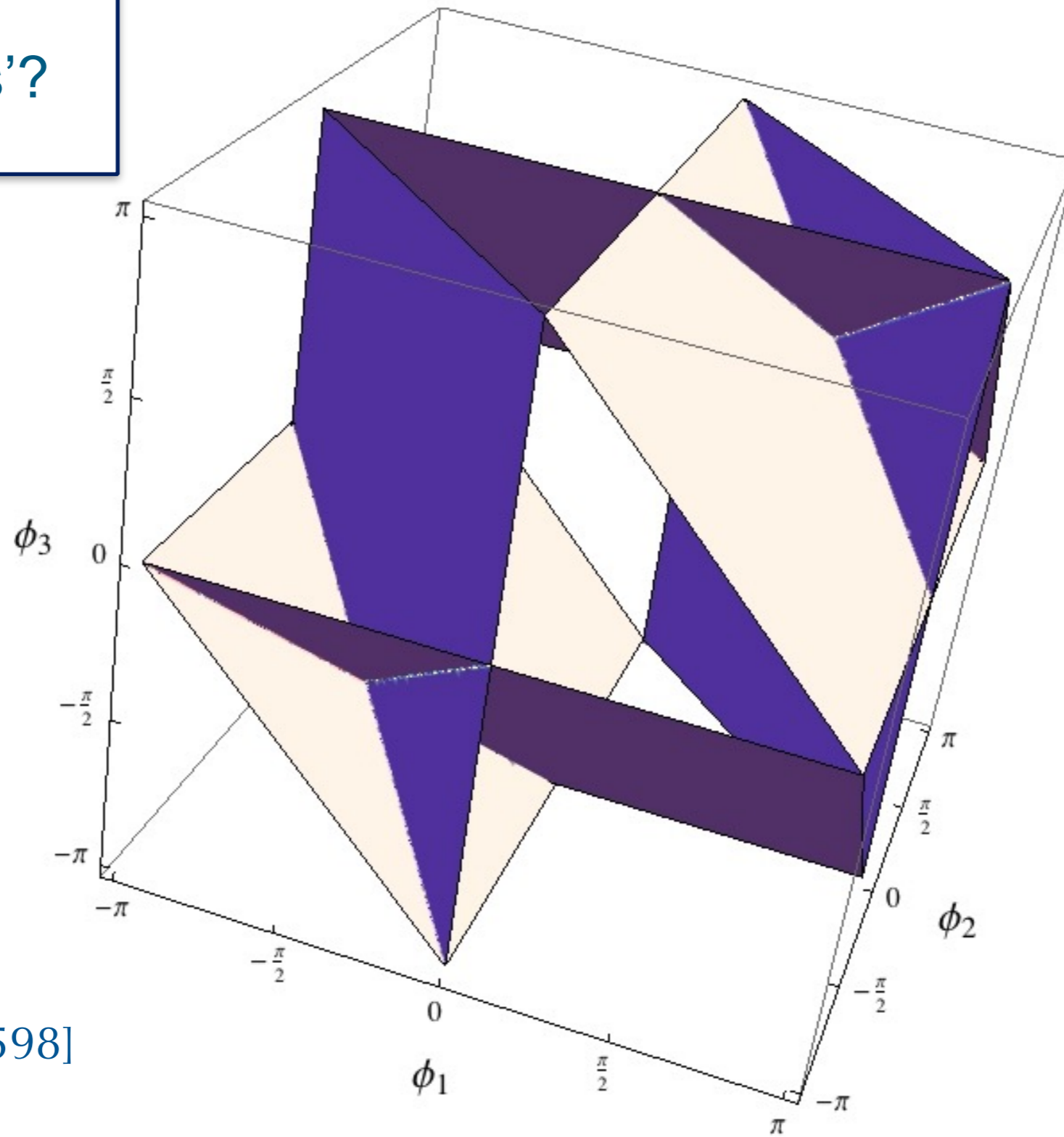
Numerical Results:

Three-body level



- ▶ **Twists are easy to implement**
 - *Multiply off diagonal terms by phase*
 - *Claim: Twist do not increase sign oscillations (Future)*
- ▶ **Twisted boundaries for reduced FV effects**
 - *Twist averaging vs iPBs*
 - *One can extract infinite volume results without knowing the functional dependence*
- ▶ **Effects of twists on non-bound states?**
 - *Mapping out the phase shifts?*
- ▶ **iPBs for larger systems (more Nuclei)? (Future)**
 - *Can one find the twist dependence for larger systems? (Seems like yes)*
- ▶ **iPBs for other partial waves? (Future)**
 - *Two-body results already known*
- ▶ **Can one relate this to relativistic systems?**
- ▶ **Are there time like twists?**
 - *Reduce the lattice in time direction and simultaneously increase precision*

4-Body 'iPBs'?



[arXiv:1511.06598]

Thank you for your attention