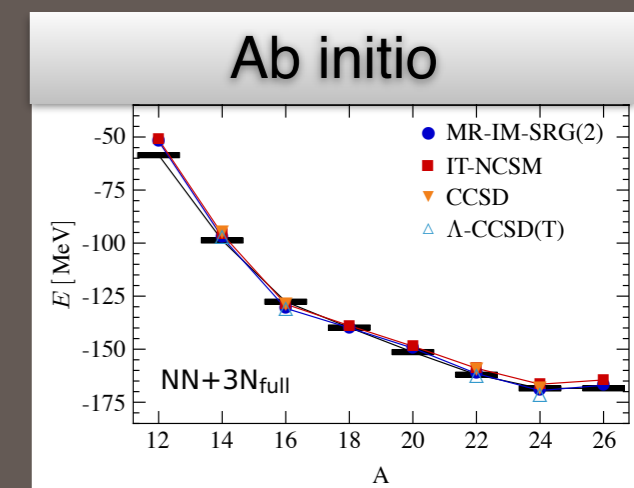
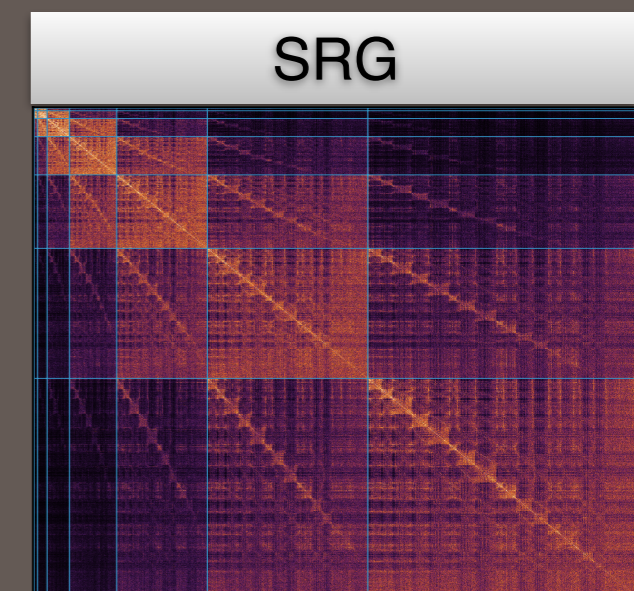
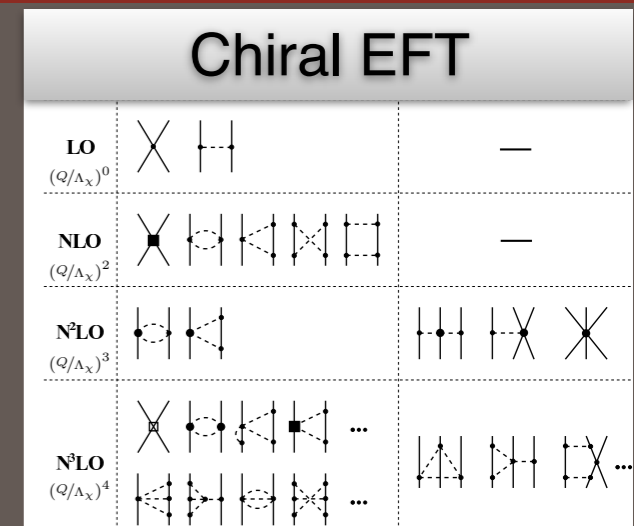
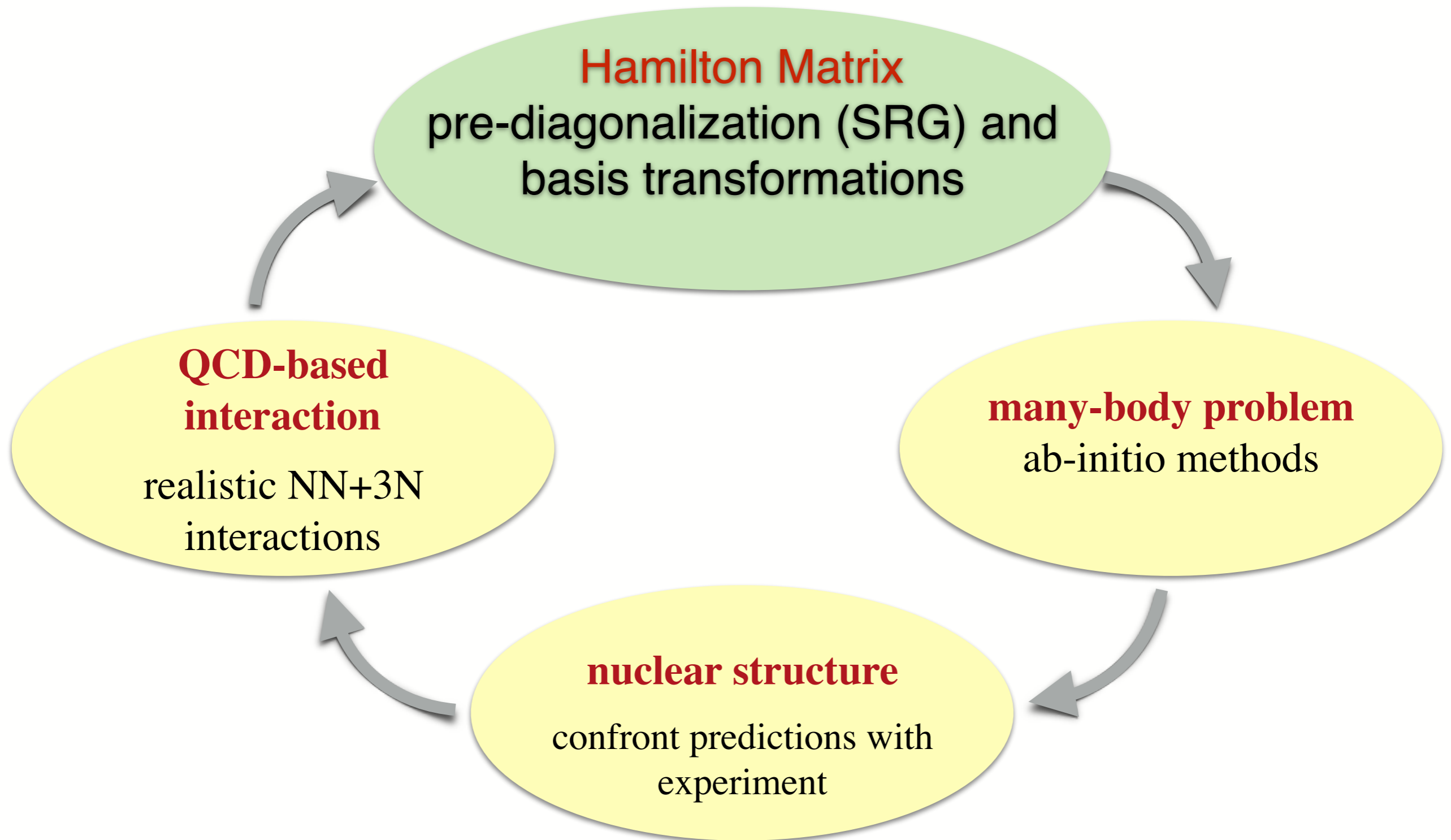


Sensitivities and Correlations from chiral NN+3N Interactions in ab initio Nuclear Structure and Reaction Calculations

INT Program INT-16-1
 May 23 2016, Seattle

Angelo Calci | TRIUMF





Chiral NN+3N Interactions

Weinberg, van Kolck, Machleidt, Entem, Meissner, Epelbaum, Krebs, Bernard,...

- **standard interaction:**

- NN @ N³LO: Entem & Machleidt, 500MeV cutoff
- 3N @ N²LO: Navrátil, local, 500MeV cutoffs & modifications

- **optimized N²LO interaction:**

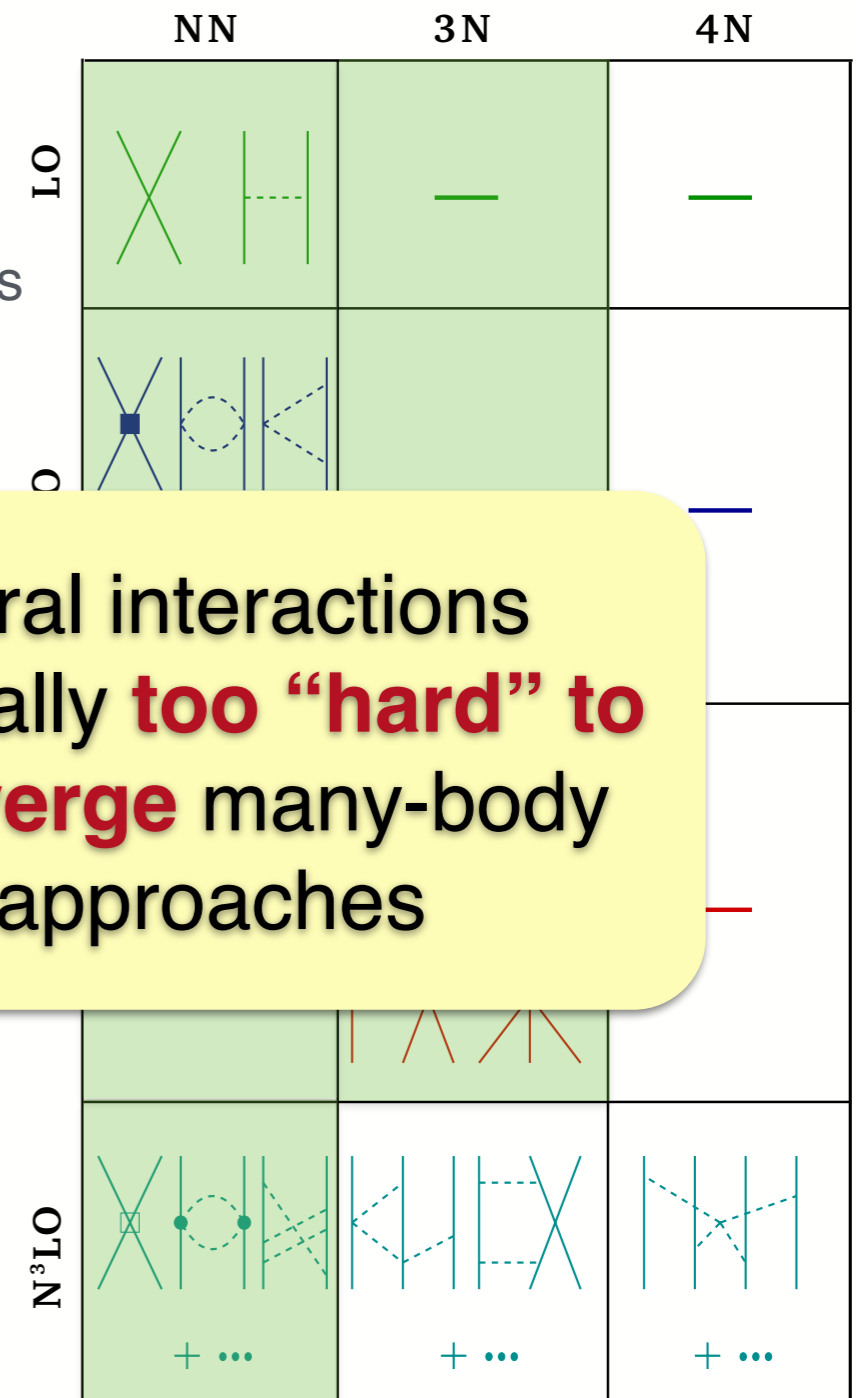
- NN: Ekström et al., 500MeV cutoff, LECs fitted by POUNDerS
- 3N: Navrátil, local, 500MeV cutoff, fit to ⁴He & T

- **Epelbaum N²LO interaction:**

- NN: Epelbaum et al., 450, . . . , 600 MeV cutoff
- 3N: Epelbaum et al., 450, . . . , 600 MeV cutoff, nonlocal

- **N²LO_{SAT} interaction:**

- NN+3N: Ekström et al., nonlocal, 450MeV cutoff, simultaneous fit to NN data and many-body observables



Similarity Renormalization Group

- Roth, Langhammer, AC et al. — Phys. Rev. Lett. 107, 072501 (2011)
- Roth, Neff, Feldmeier — Prog. Part. Nucl. Phys. 65, 50 (2010)
- Jurgenson, Navrátil, Furnstahl — Phys. Rev. Lett. 103, 082501 (2009)
- Bogner, Furnstahl, Perry — Phys. Rev. C 75 061001(R) (2007)

Similarity Renormalization Group (SRG)

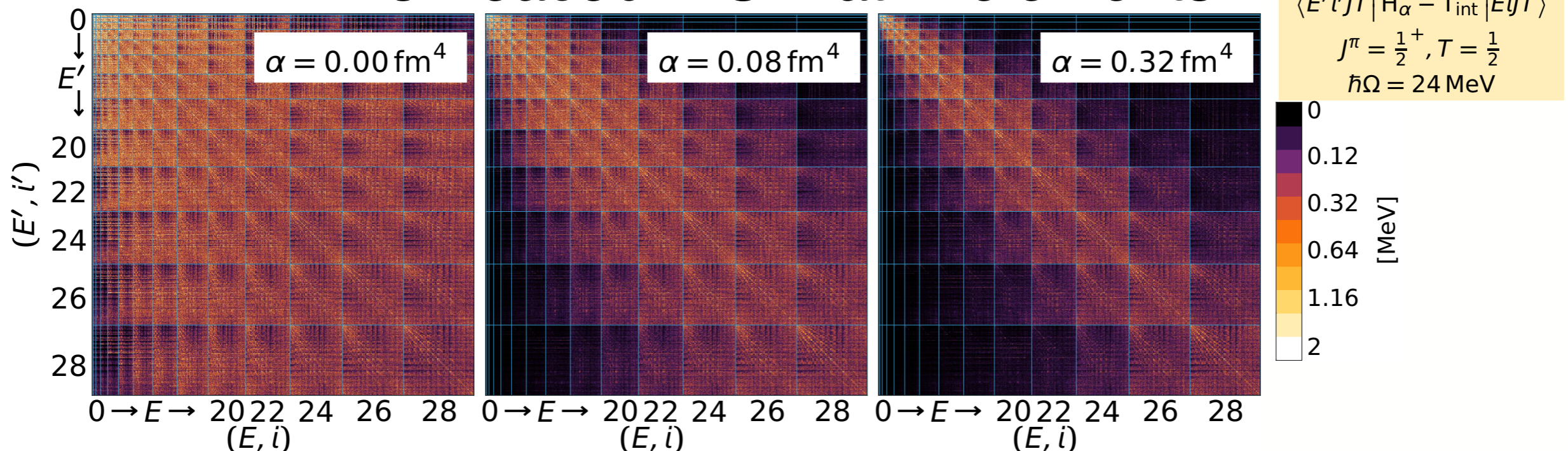
accelerate convergence by **pre-diagonalizing** the Hamiltonian with respect to the many-body basis

- unitary transformation leads to evolution equation**

$$\frac{d}{d\alpha} \tilde{H}_\alpha = [\eta_\alpha, \tilde{H}_\alpha] \quad \text{with} \quad \eta_\alpha = (2\mu)^2 [T_{\text{int}}, \tilde{H}_\alpha] = -\eta_\alpha^\dagger$$

advantages of SRG: **flexibility** and **simplicity**

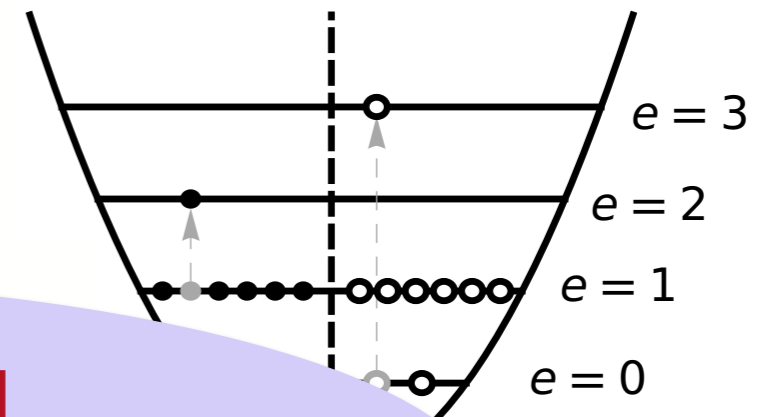
3B-Jacobi HO matrix elements



(Importance Truncated) NCSM

- solving the eigenvalue problem:

$$H |\Psi_n\rangle = E_n |\Psi_n\rangle$$



- **model space:**

spanned by
unperturbed
up to $N_{max} \hbar\Omega$

problem of NCSM

enormous increase of model space with
particle number A

Importance Truncated NCSM

- **extrapolation** of $\kappa_{min} \rightarrow 0$ recovers effect of omitted contributions
- IT-NCSM provides **same results** as full NCSM
- **expands** application **range** to larger A
- **importance truncated space** spanned by basis states with $|K_\nu| \geq \kappa_{min}$

Uncertainty Quantification of chiral Hamiltonians in Spectroscopy

Uncertainties of Chiral Interactions

in the past

- uncertainties **from many-body approach** included
- observables calculated for **single chiral Hamiltonian** (inconsistent chiral order)
- quality of chiral forces assessed by agreement with experiment

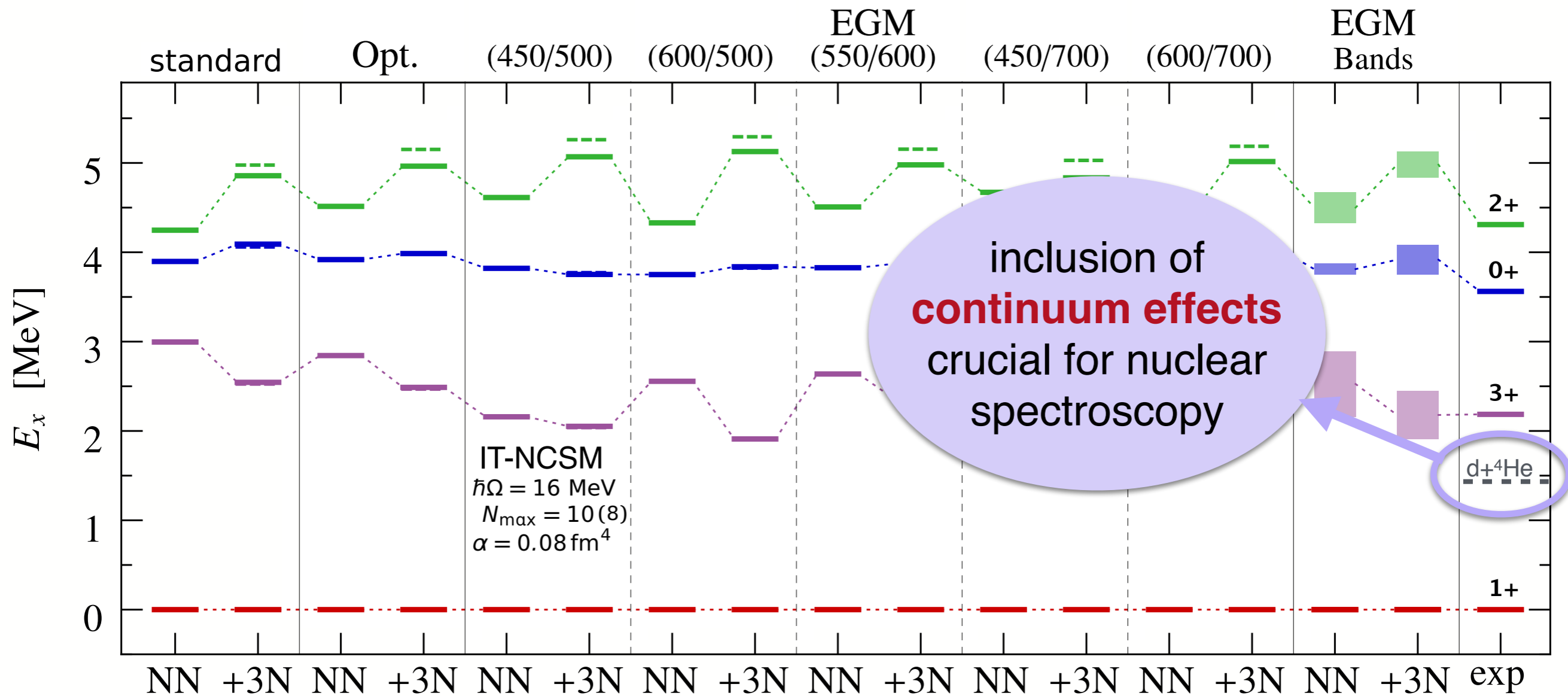
future

- start with NN+3N force at **consistent chiral order**
- use **sequence of cutoffs** and chiral **orders**
 - ➔ estimate **uncertainties** of chiral EFT and many-body approach

nuclear structure physics approaches new era

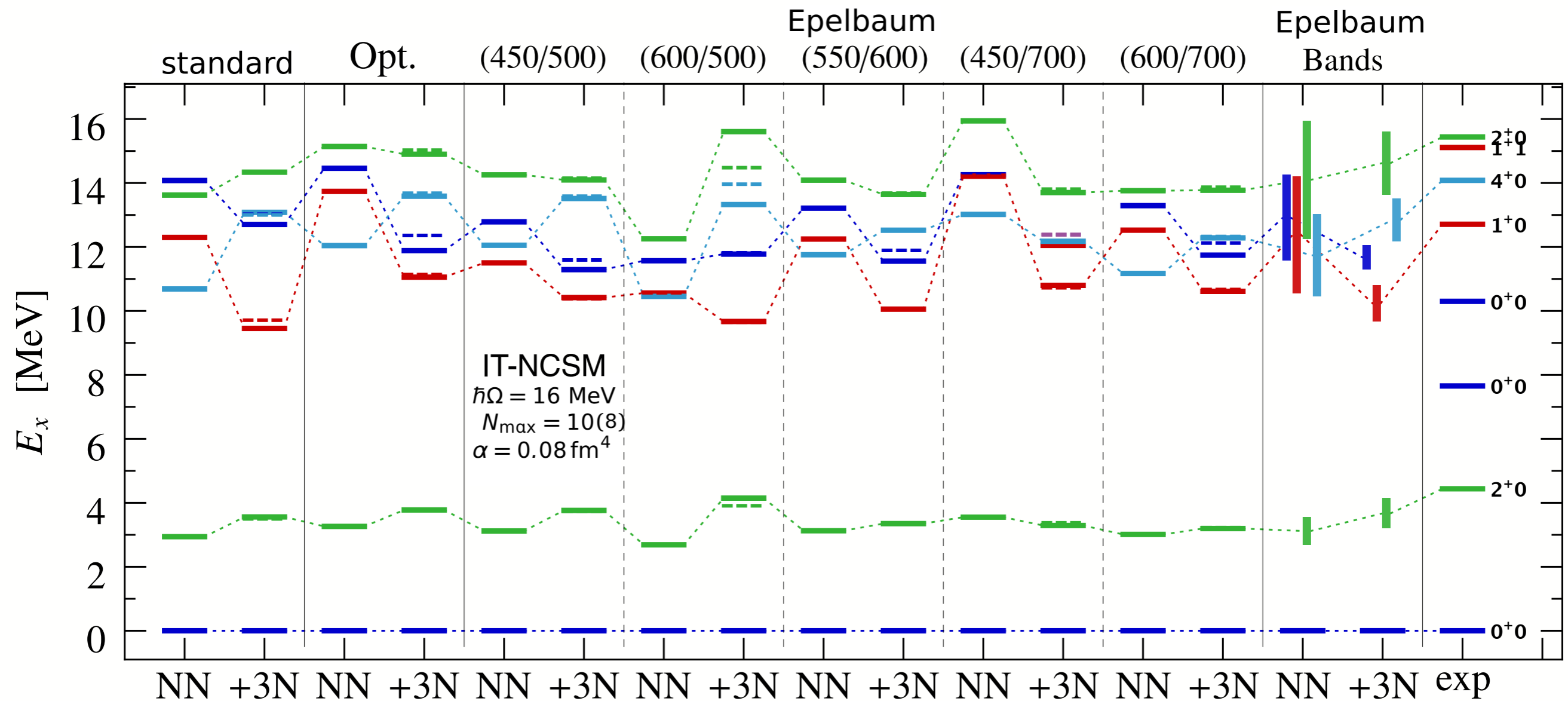
- **ongoing progress** in construction of consistent NN+3N Hamiltonians
- first step towards reliable uncertainty quantification in nuclear spectroscopy

${}^6\text{Li}$: Cutoff Dependence



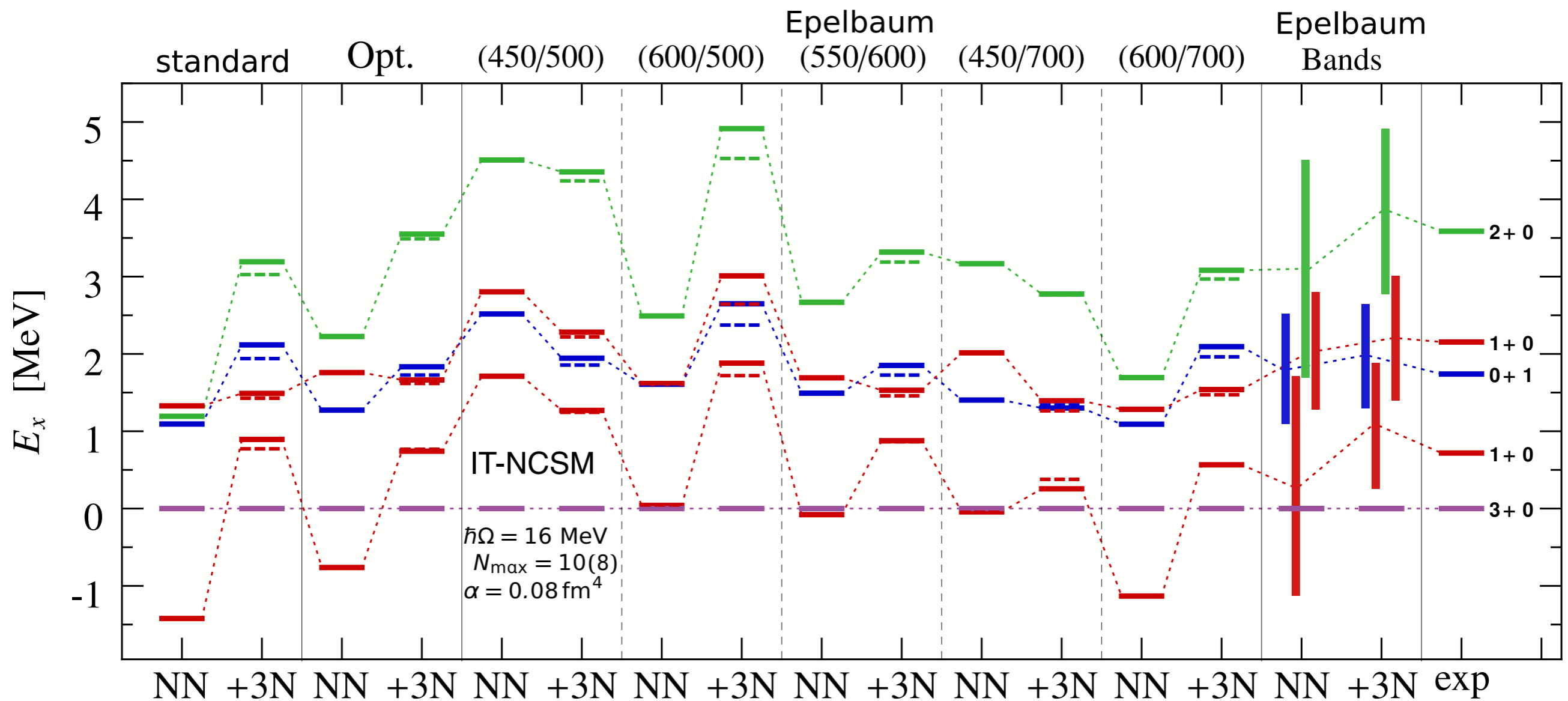
- reasonable predictions
- small cutoff dependence

^{12}C : Cutoff Dependence



- small cutoff dependence for NN+3N

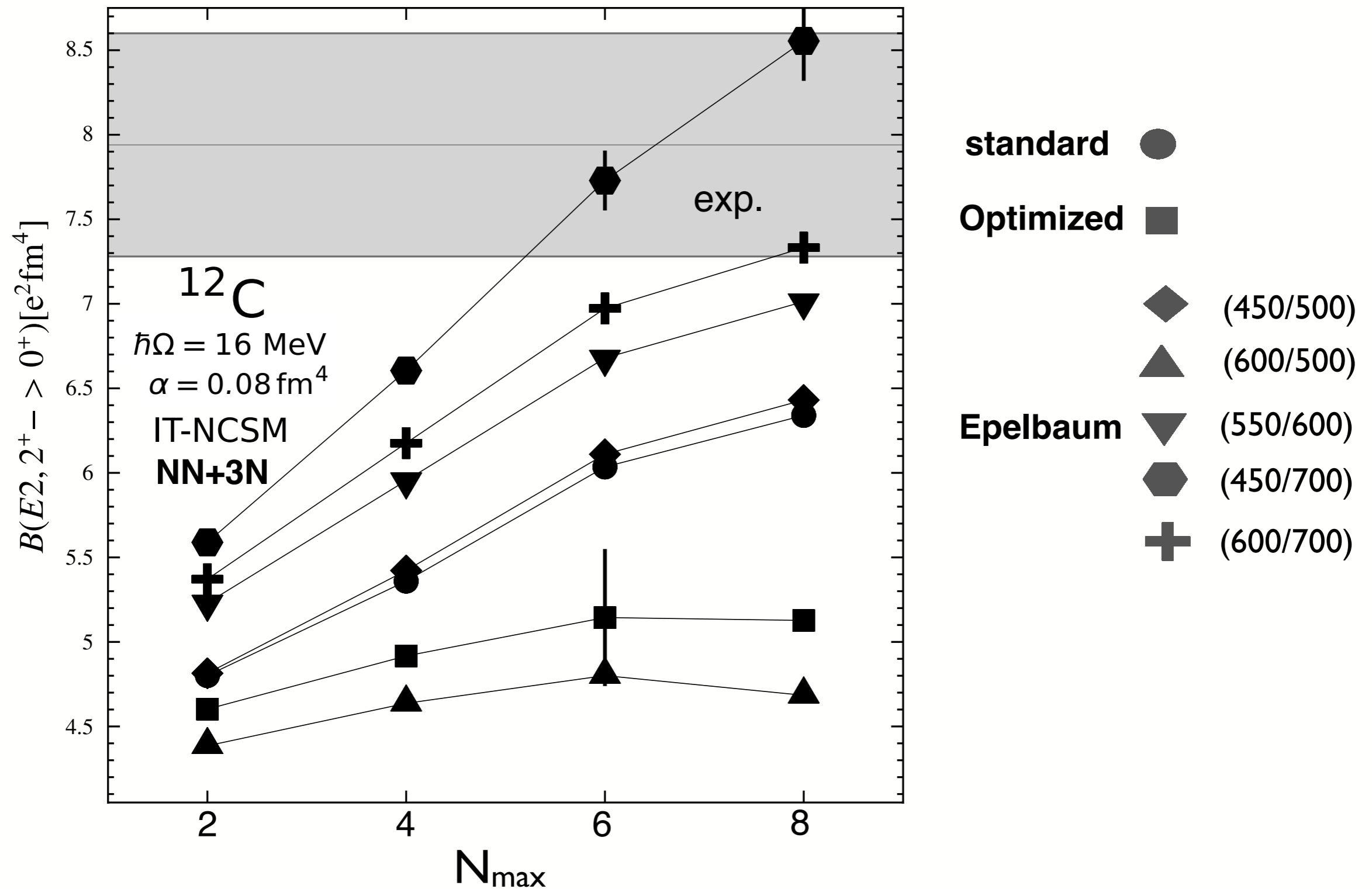
^{10}B : Cutoff Dependence



- complex system with compressed spectrum
- accurate predictions within large uncertainties

Quadrupole Observables

^{12}C : B(E2) Transition



B(E2) and Quadrupole moment

challenges in calculating quadrupole observables:

- generally **slow convergence** in HO space
- **chiral uncertainties** impact observables
- **observables** need to be **SRG evolved**

$$\langle \tilde{\Psi}_m | O | \tilde{\Psi}_n \rangle \neq \langle \Psi_m | O | \Psi_n \rangle = \langle \tilde{\Psi}_m | U^\dagger O U | \tilde{\Psi}_n \rangle$$

what we do

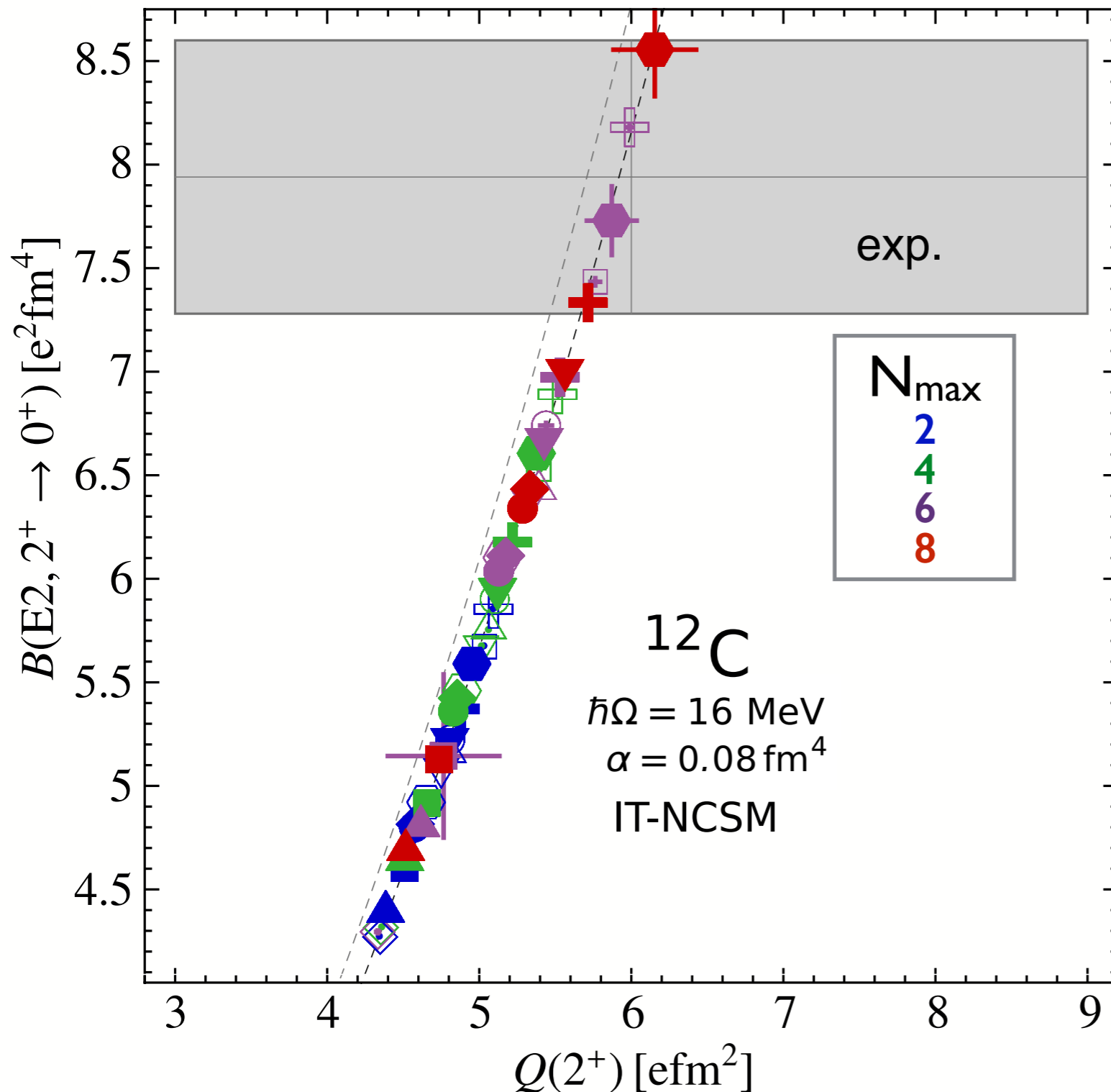
what we want

what we need to do

correlations in ab initio calculations

- can provide a solution
- information about underlying structure

^{12}C : Quadrupole Correlation



- **robust correlation** between $B(E2)$ and spectroscopic Q value
- independent of:
 - chiral interaction / cutoff
 - model space
 - SRG flow parameter


B(E2) and Quadrupole moment

simple rotational model:

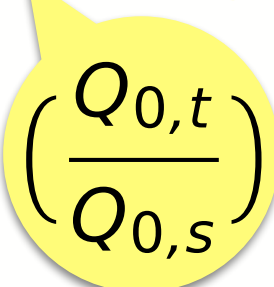
- observables connected via intrinsic quadrupole moment Q_0

$$Q(J) = \frac{3K^2 - J(J+1)}{(J+1)(2J+3)} Q_{0,s}$$

$$B(E2, J_i \rightarrow J_f) = \frac{5}{16\pi} Q_{0,t}^2 \left(\begin{matrix} J_i & 2 & J_f \\ K & 0 & K \end{matrix} \right)$$

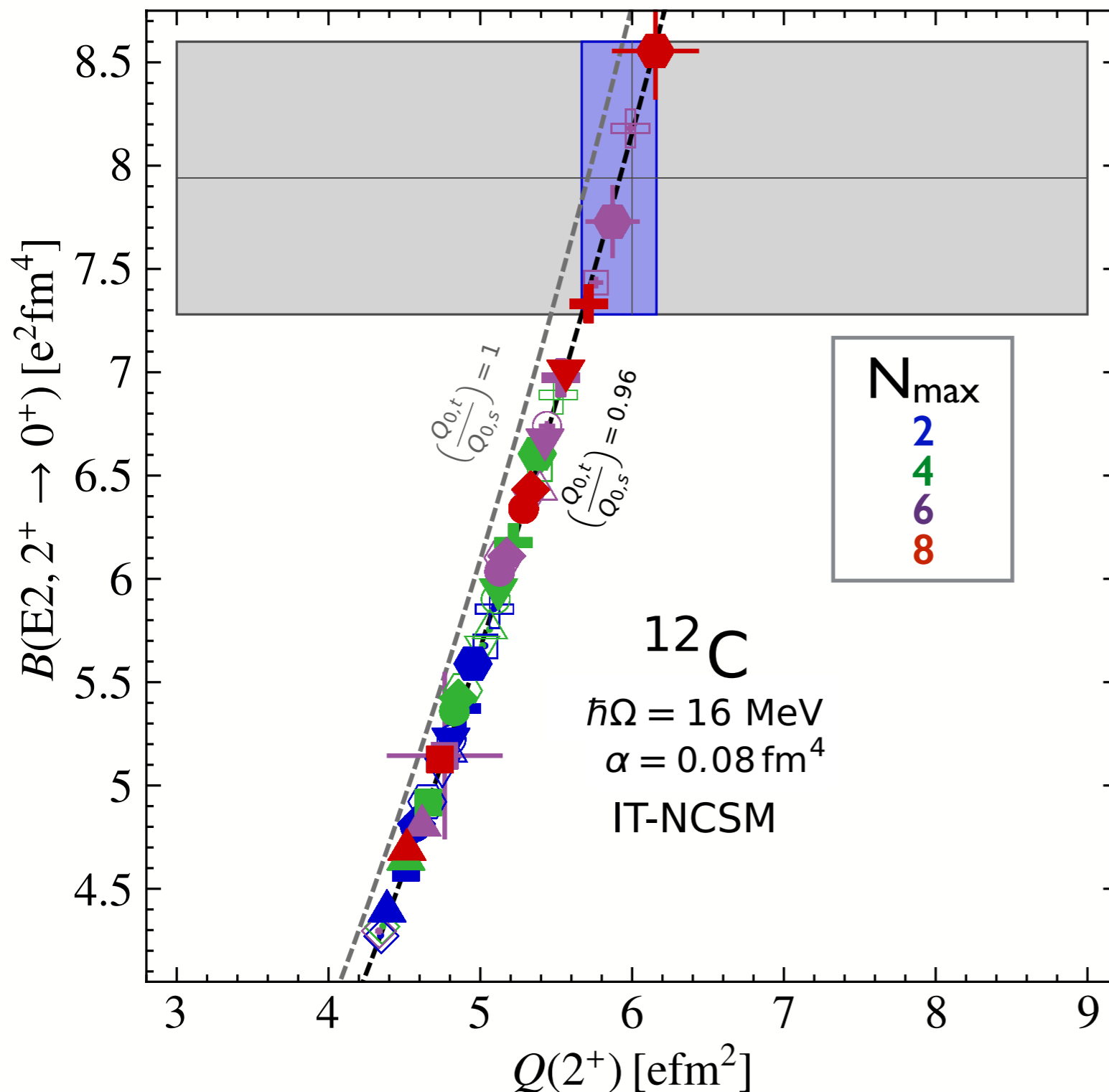


$$B(E2, J_i \rightarrow J_f) = \frac{5}{16\pi} \frac{((J+1)(2J+3))^2}{(3K^2 - J(J+1))^2} \left(\begin{matrix} J_i & 2 & J_f \\ K & 0 & K \end{matrix} \right) \left(\frac{Q_{0,t}}{Q_{0,s}} \right)^2 Q(J)^2$$

only parameter


quadratic relation between B(E2) and Quadrupole moment

^{12}C : Quadrupole Correlation

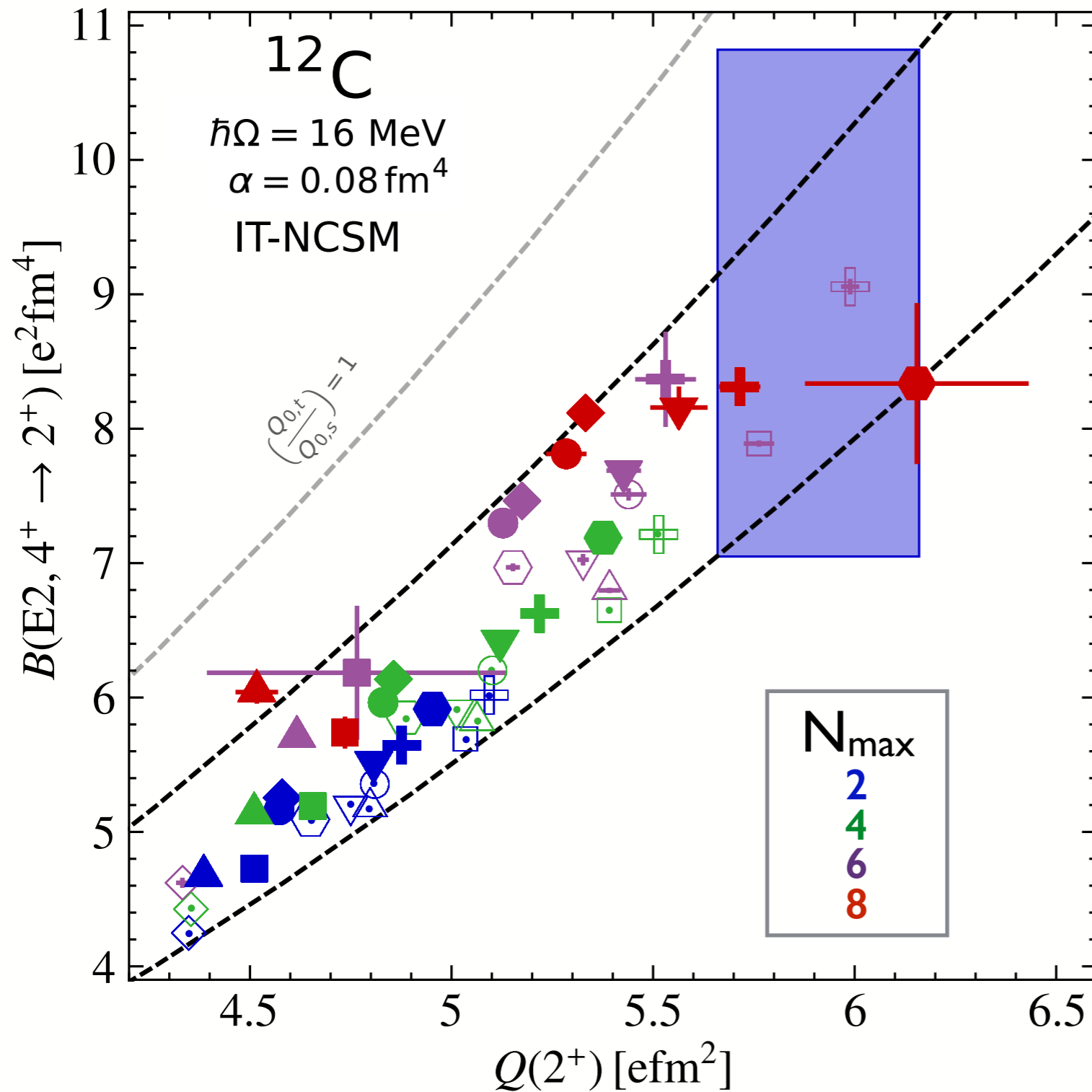


- **robust correlation** between $B(E2)$ and spectroscopic Q value
- independent of:
 - chiral interaction / cutoff
 - model space
 - SRG flow parameter

can combine correlation curve with measured $B(E2)$ to **predict Q** precisely

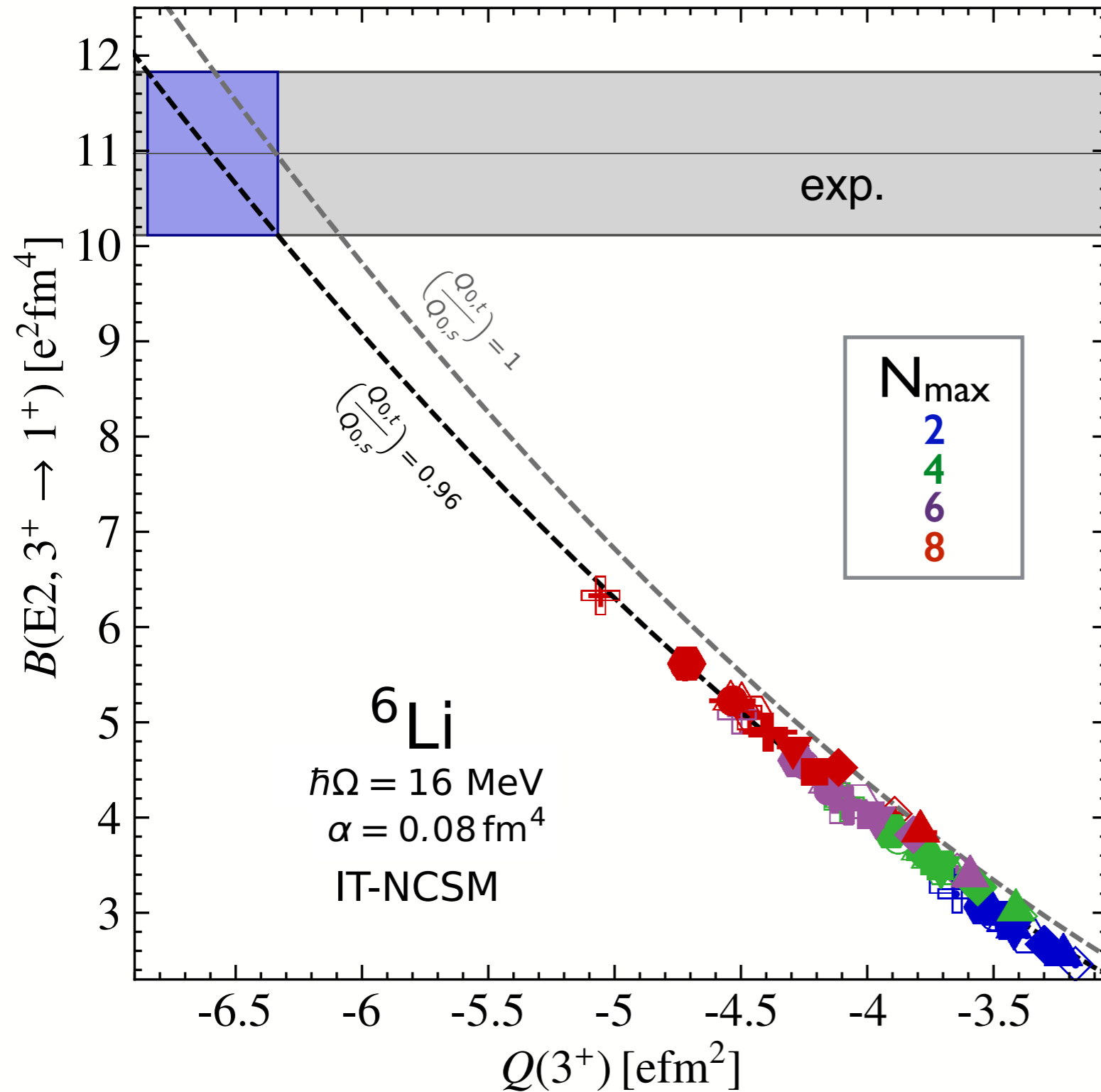
	std	Opt.	Epelbaum				
NN	□	○	◇	△	▽	⬡	+
NN+3N	■	●	◆	▲	▼	⬤	+

^{12}C : Quadrupole Correlation



- **less robust correlation** between $B(E2)$ and spectroscopic Q value
- can be used to predict $B(E2, 4^+ \rightarrow 2^+)$ rather precise

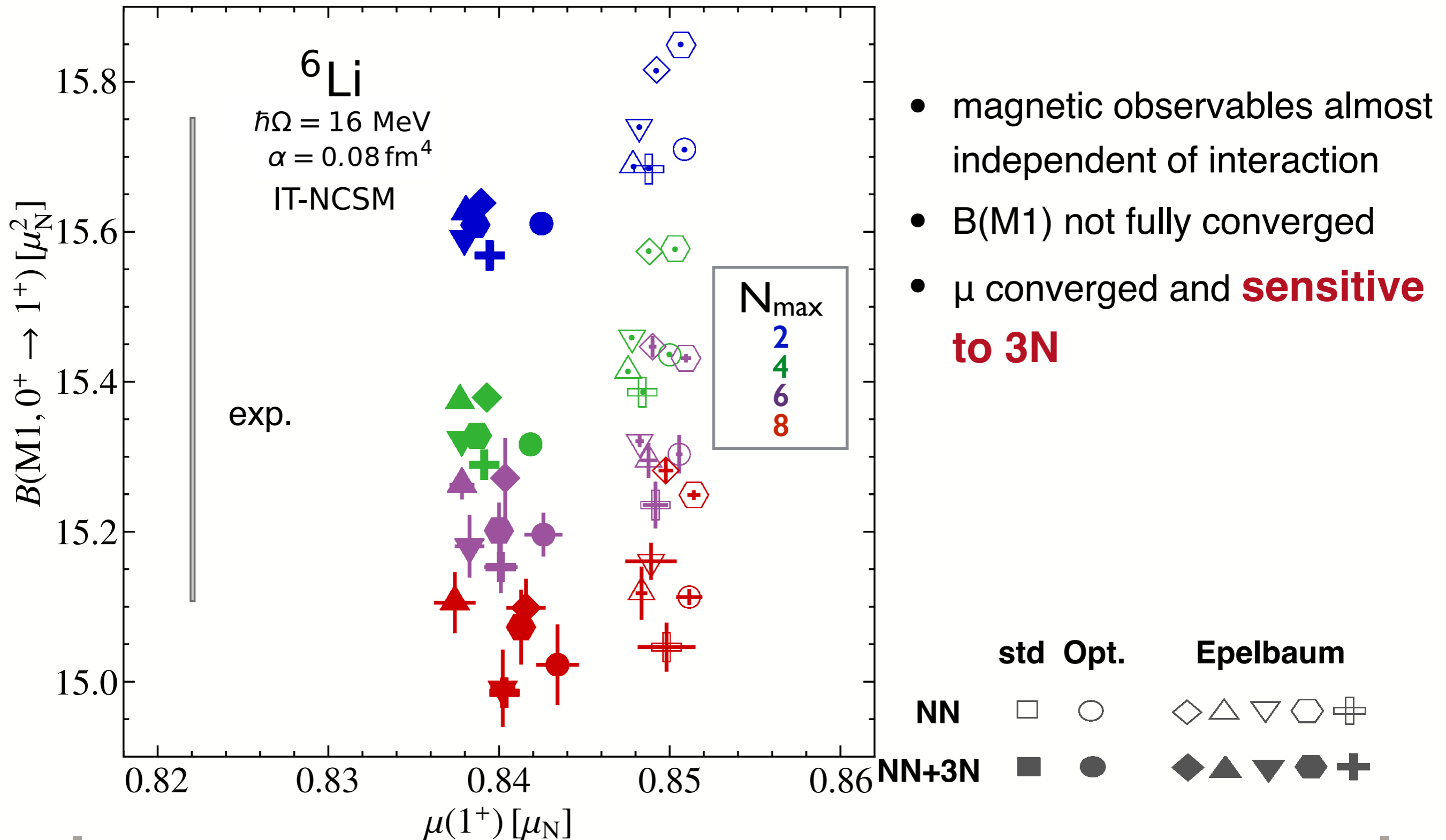
${}^6\text{Li}$: Quadrupole Correlation



- **robust correlation** between $B(E2)$ and spectroscopic Q value
- independent of:
 - chiral interaction / cutoff
 - model space
 - SRG flow parameter

	std	Opt.	Epelbaum				
NN	□	○	◇	△	▽	◻	+
NN+3N	■	●	◆	▲	▼	◼	+

${}^6\text{Li}$: Dipol Observables



NCSM with Continuum

with
P. Navrátil, R. Roth,
J. Langhammer, S. Quaglioni, G. Hupin

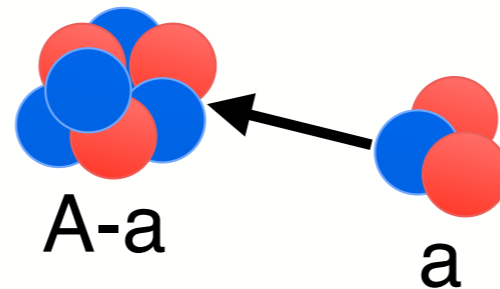
NCSM with Continuum

Baroni, Navrátil, Quaglioni

Phys. Rev. Lett. 110, 022505 (2013)

ab-initio description of nuclei

bound states &
spectroscopy



resonances &
scattering states

(IT-)NCSM

ab initio description of
nuclear clusters

RGM


describing relative
motion of clusters

NCSMC
continuum effects in spectroscopy

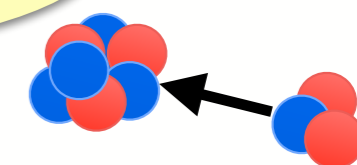
NCSM with Continuum (NCSMC)

- representing $H|\psi^{J\pi T}\rangle = E|\psi^{J\pi T}\rangle$ using the **over-complete basis**

$$|\psi^{J\pi T}\rangle = \sum_{\lambda} c_{\lambda} |\Psi_{A E_{\lambda} J^{\pi T}}\rangle + \sum_{\nu} \int dr r^2 \frac{\chi_{\nu}(r)}{r} |\xi_{\nu r}^{J\pi T}\rangle$$



expansion in A-body
(IT-)NCSM eigenstates



relative motion of clusters
NCSM/RGM expansion

- leads to NCSMC equation

$$\begin{pmatrix} H_{NCSM} & h \\ h & \mathcal{H} \end{pmatrix} \begin{pmatrix} c \\ \chi(r)/r \end{pmatrix} = E \begin{pmatrix} \mathbb{1} & g \\ g & \mathbb{1} \end{pmatrix} \begin{pmatrix} c \\ \chi(r)/r \end{pmatrix}$$

- with 3N contributions in

H_{NCSM}

covered by
(IT-)NCSM

h

given by
 $\langle \Psi_{A E_{\lambda} J^{\pi T}} | H | \xi_{\nu r}^{J\pi T} \rangle$

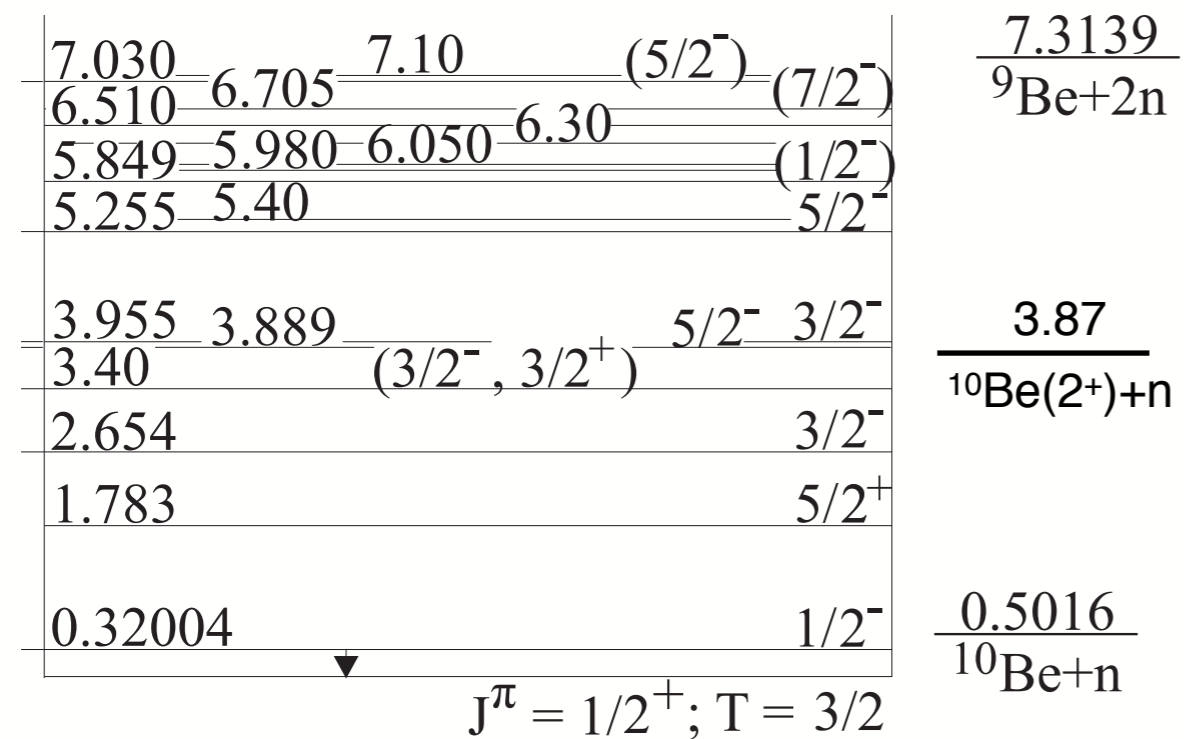
\mathcal{H}

contains NCSM/RGM
Hamiltonian kernel

Neutron-rich halo Nucleus ^{11}Be

- Halo states**

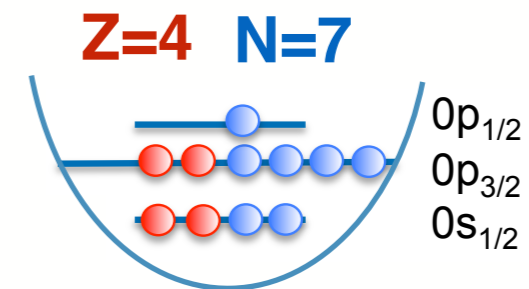
weakly bound $J=1/2$ states dominated by $n-^{10}\text{Be}$



^{11}Be

- parity inversion**

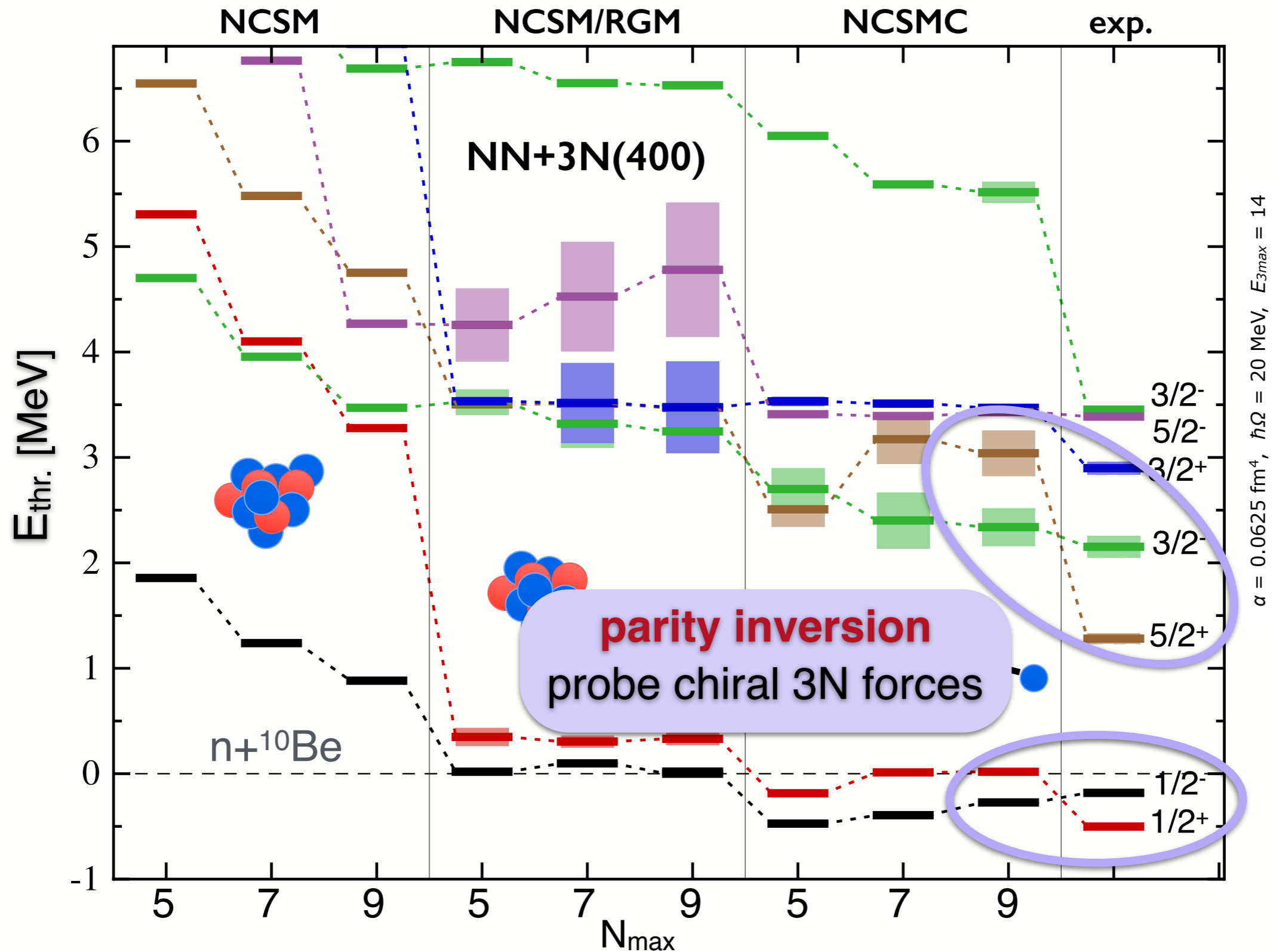
shell model predicts g.s. to be $J^\pi=1/2^-$



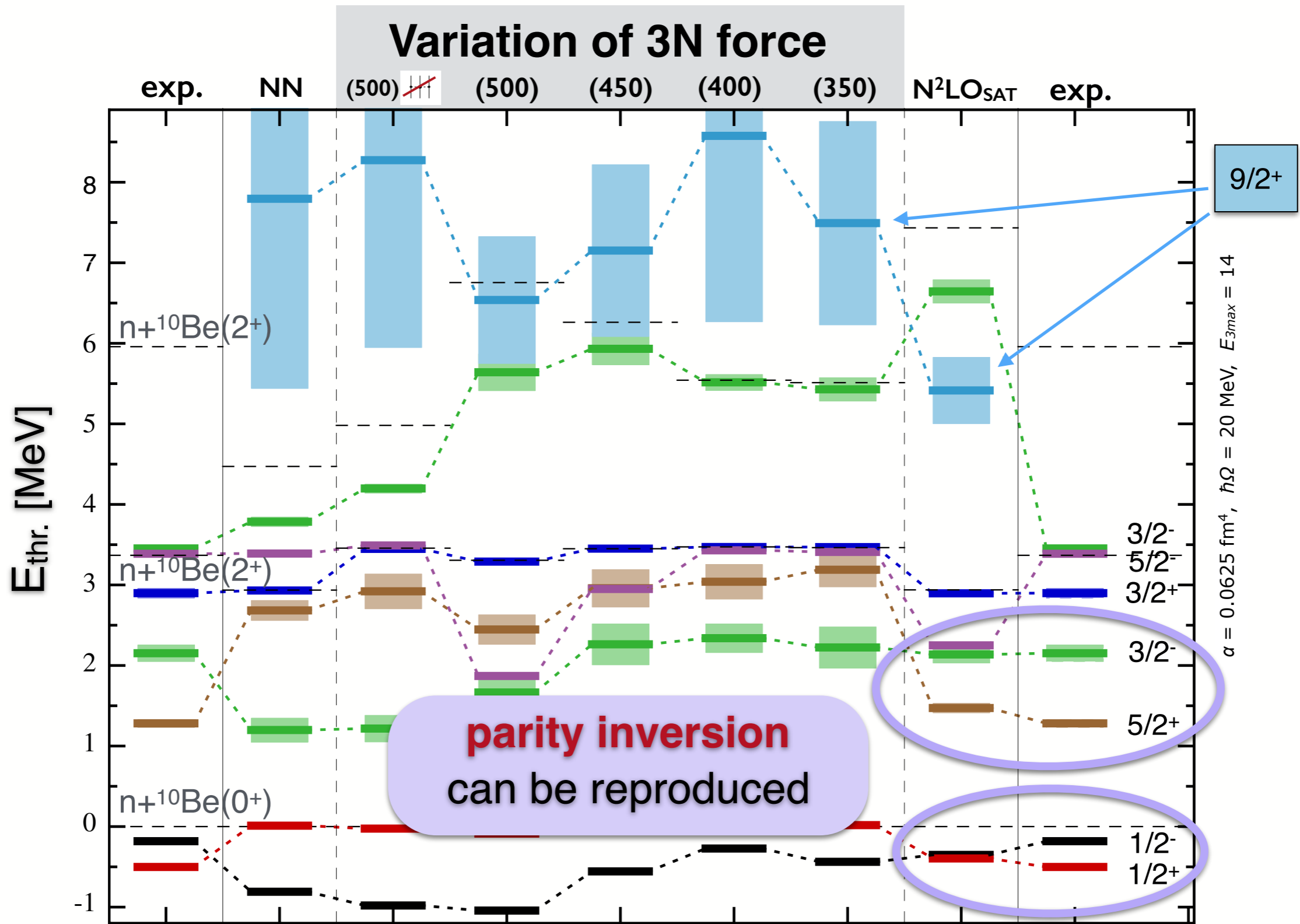
ab initio NCSMC

- include $n-^{10}\text{Be}$ continuum ($0^+, 2^+, 2^+$ states of ^{10}Be)
- include 6 negative and 3 positive parity states of ^{11}Be

^{11}Be excitation spectrum

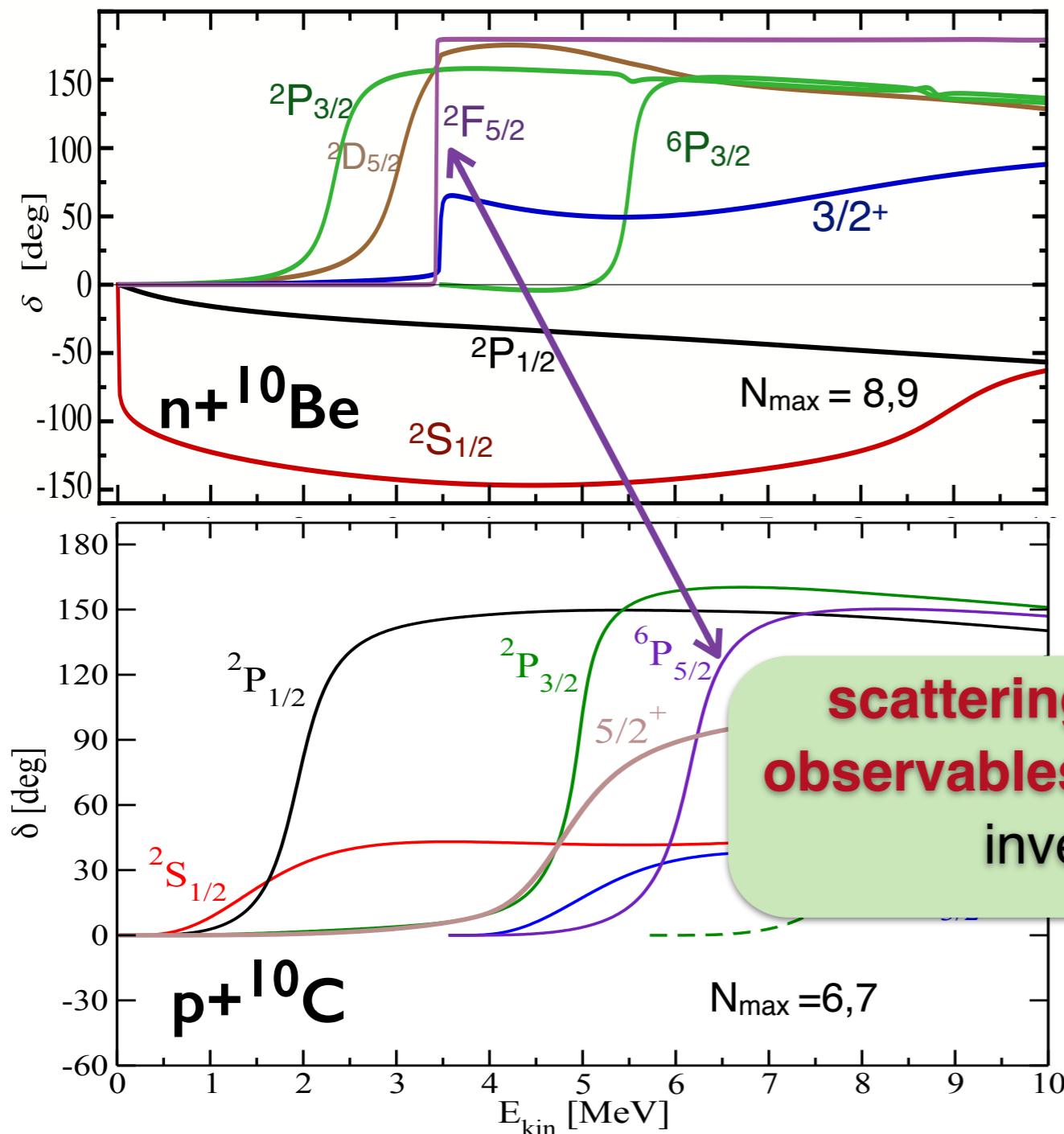


^{11}Be excitation spectrum



Mirror nuclei: ^{11}Be and ^{11}N

standard NN+3N(400)

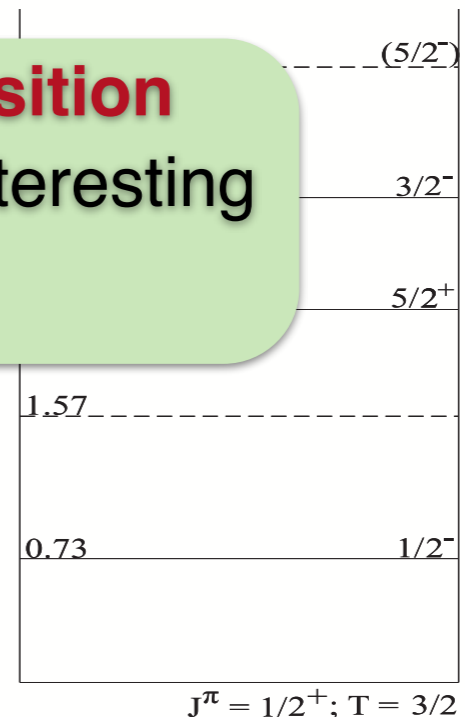


$\alpha = 0.0625 \text{ fm}^4, \hbar\Omega = 20 \text{ MeV}, E_{3\text{max}} = 14$

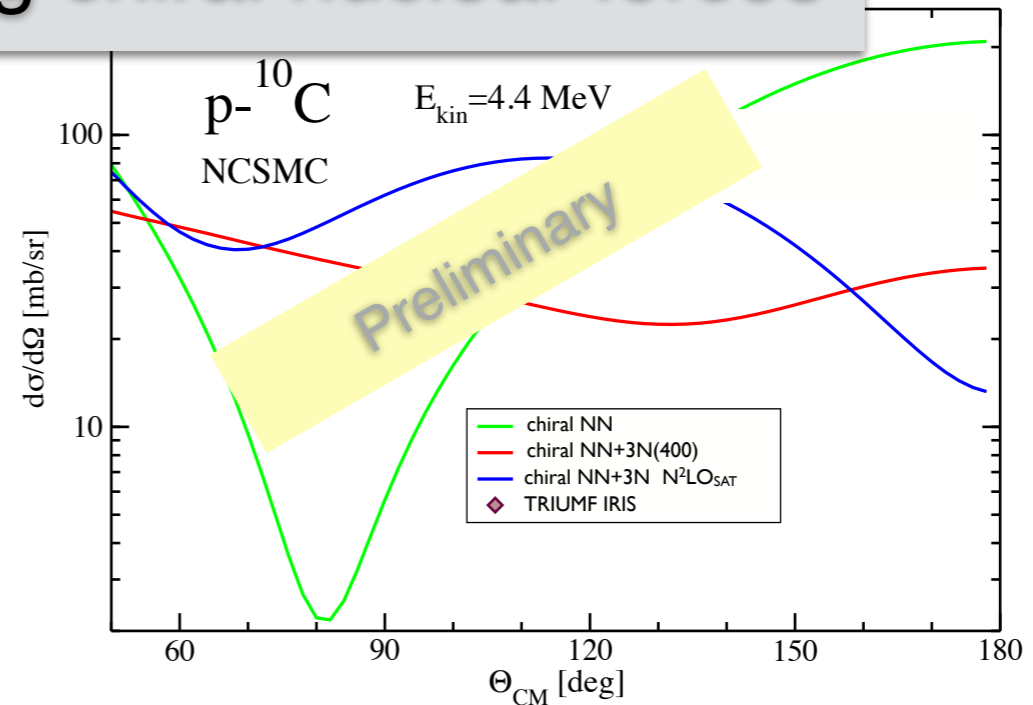
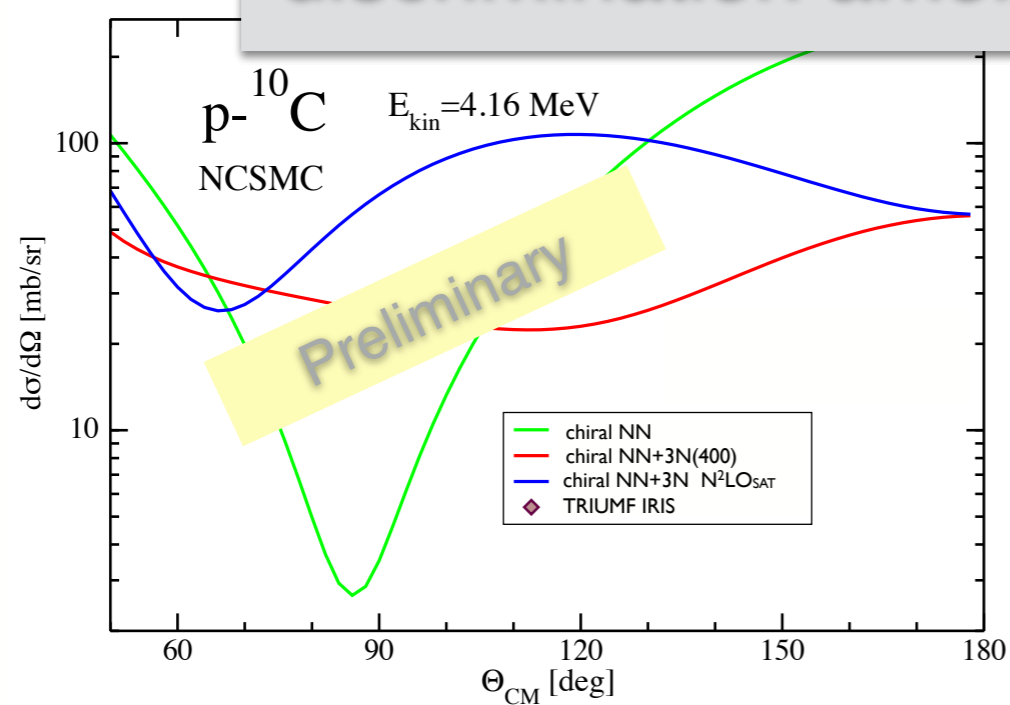
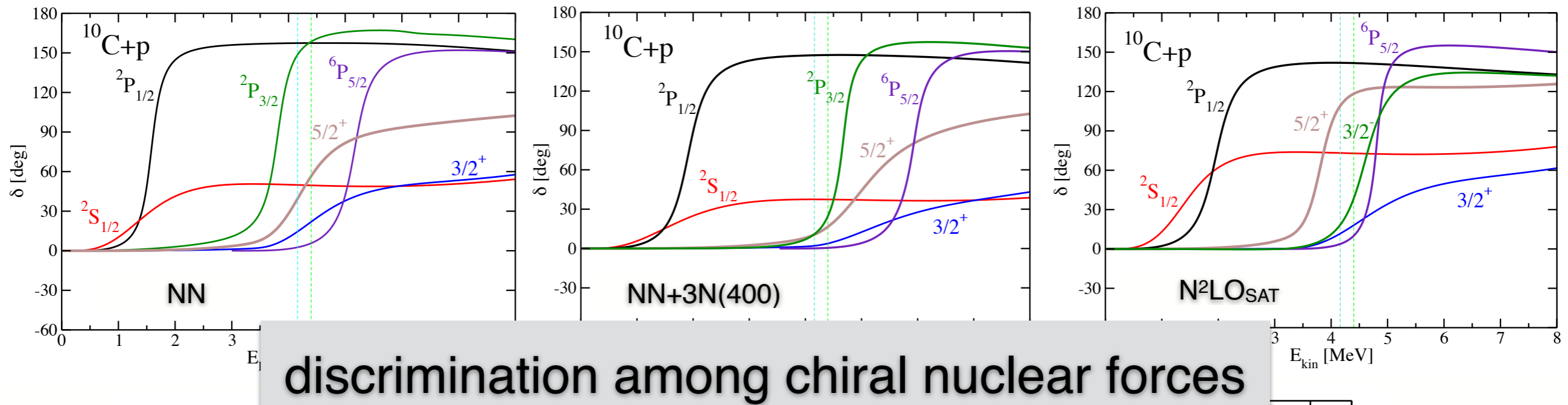
7.030	6.705	7.10	(5/2 ⁻)	(7/2 ⁻)	$\frac{7.3139}{^9\text{Be}+2\text{n}}$
6.510	5.980	6.050	6.30	(1/2 ⁻)	
5.849	5.40			5/2 ⁻	
5.255					
3.955	3.889		(3/2 ⁻ , 3/2 ⁺)	5/2 ⁻ 3/2 ⁻	$\frac{3.87}{^{10}\text{Be}(2^+)+\text{n}}$
3.40				3/2 ⁻	
2.654				5/2 ⁺	
1.782				1/2 ⁻	$\frac{0.5016}{^{10}\text{Be}+\text{n}}$

ab initio calculations predict 3/2⁺-state

scattering and transition observables enable interesting investigations



$p+^{10}\text{C}$ Scattering: Structure of ^{11}N resonances



IRIS collaboration:
A. Kumar, R. Kanungo, A. Sanetullaev *et al.*

A. Calci, P. Navratil, G. Hupin, S. Quaglioni, R. Roth *et al* with IRIS collaboration, in preparation

NCSMC with approximated 3N forces

with
P. Navrátil, R. Roth, E. Gebrerufael

NCSM with Continuum (NCSMC)

- representing $H|\psi^{J\pi T}\rangle = E|\psi^{J\pi T}\rangle$ using the **over-complete basis**

$$|\psi^{J\pi T}\rangle = \sum_{\lambda} c_{\lambda} |\Psi_{A E_{\lambda} J^{\pi T}}\rangle + \sum_{\nu} \int dr r^2 \frac{\chi_{\nu}(r)}{r} |\xi_{\nu r}^{J\pi T}\rangle$$

expansion in A-body
(IT-)NCSM eigenstates

relative
NCSM

- leads to NCSMC equation

$$\begin{pmatrix} H_{NCSM} & h \\ h & \mathcal{H} \end{pmatrix} \begin{pmatrix} c \\ \chi(r)/r \end{pmatrix} = E \begin{pmatrix} 1 \\ g \end{pmatrix}$$

- with 3N contributions in

H_{NCSM}

covered by
(IT-)NCSM

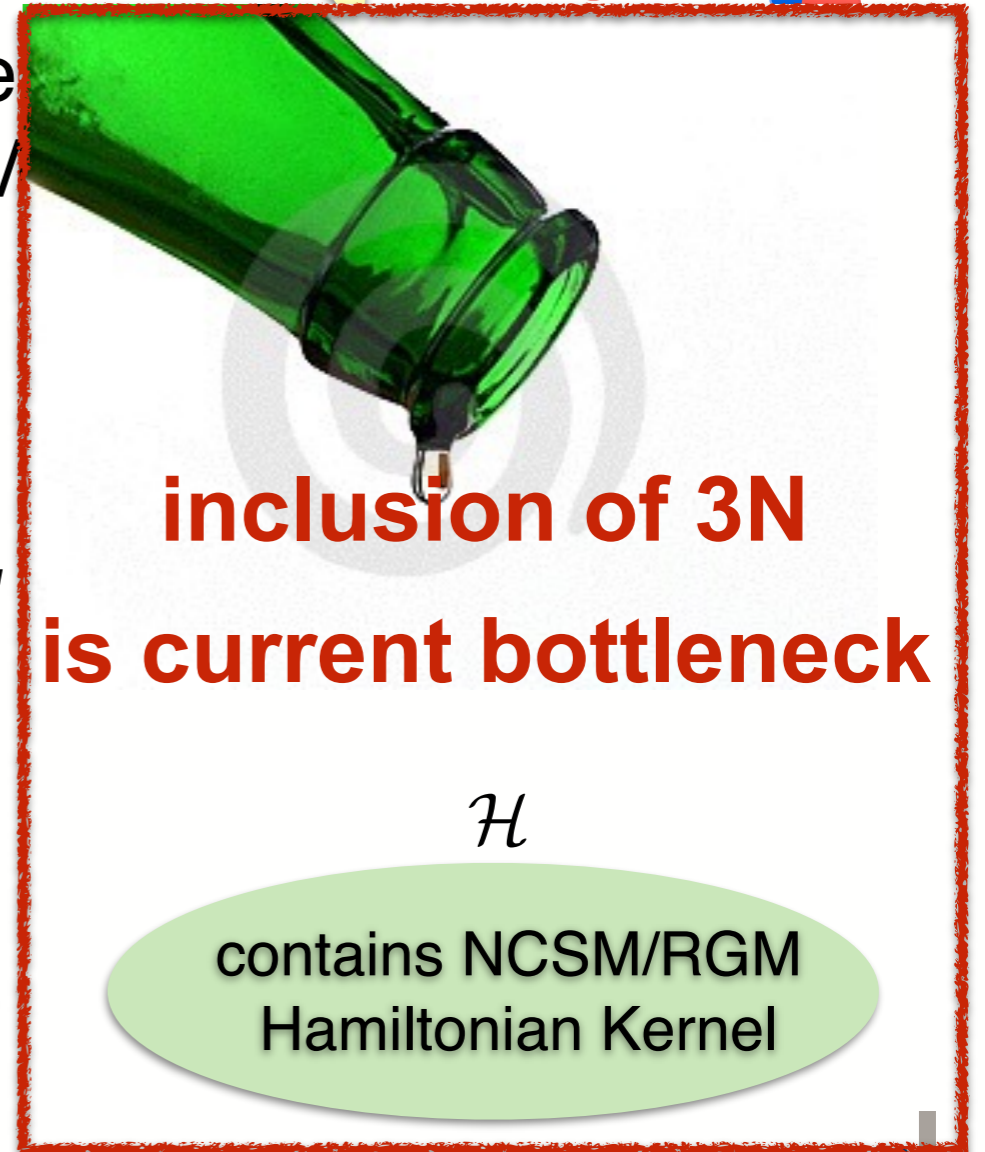
h

given by
 $\langle \Psi_{A E_{\lambda} J^{\pi T}} | H | \xi_{\nu r}^{J\pi T} \rangle$

\mathcal{H}

contains NCSM/RGM
Hamiltonian Kernel

**inclusion of 3N
is current bottleneck**



Normal-ordering (NO) approximation

- standard tool to **reduce particle rank**
- generally NO can be considered as basis transformation

$$V_{3N} \approx \underbrace{\tilde{V}_{0N} + \tilde{V}_{1N} + \tilde{V}_{2N}}_{\text{contain information of reference state and initial 3N force}} + \cancel{\tilde{V}_{3N}}$$

contain information of reference state and initial 3N force

see talk by
Ragnar Stroberg
on Thursday

- interested in direct description of **open-shell systems**
 - multi-reference normal ordering (MR-NO)
 - generalization of wicks theorem [Kutzelnigg, Mukherjee]

NCSM/RGM kernels with MR-NO contributions

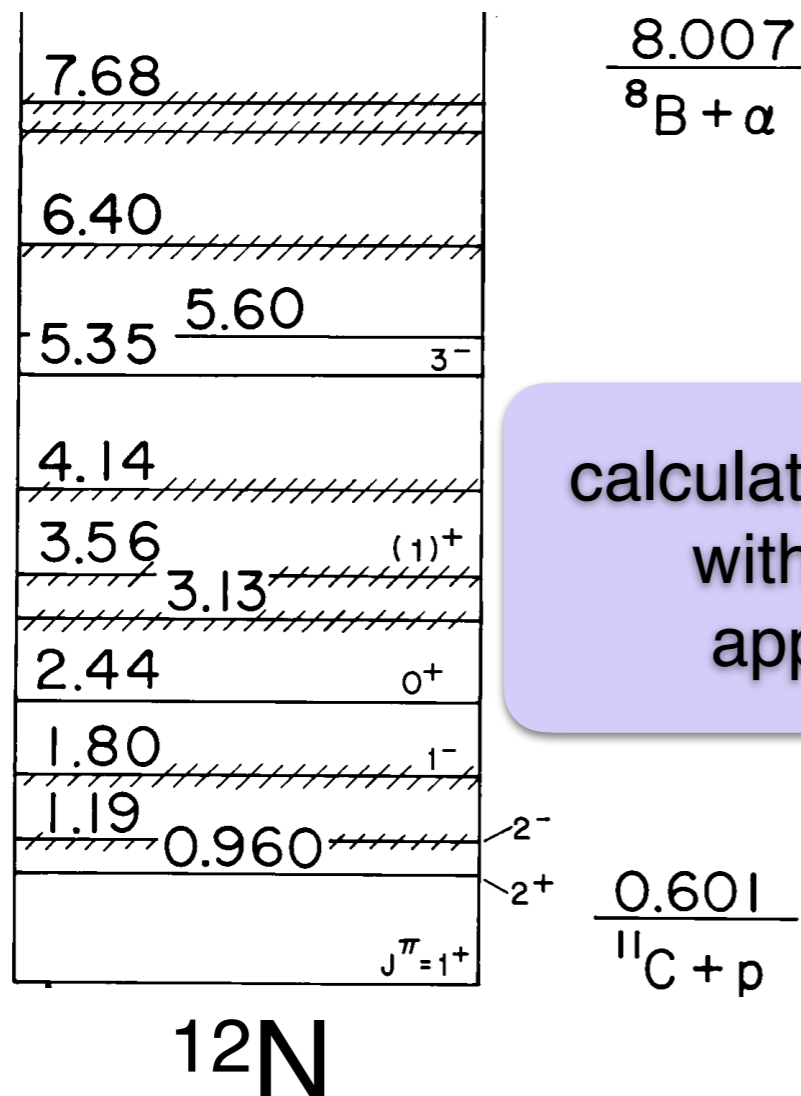
- reduces computational costs tremendously
- impressively accurate approximation

First application: ^{12}N

- ideal candidate**

weakly bound $J=1^+$ state
dominated by $p\text{-}^{11}\text{C}$

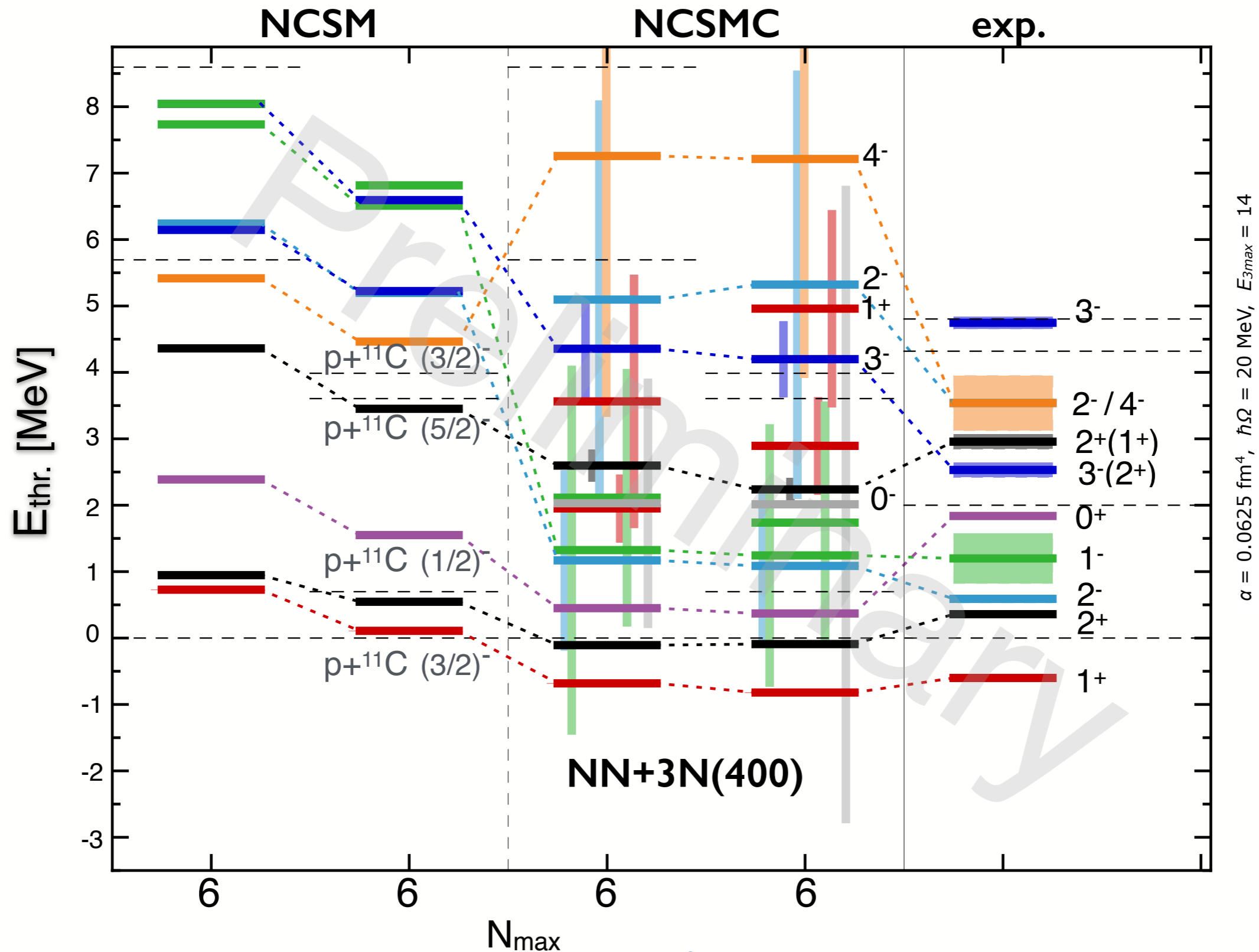
- some low lying resonances not measured precisely
- $^{11}\text{C}(p,\gamma)^{12}\text{N}$ can bypass triple-alpha process
- planned experiment at TUDA facility at TRIUMF



ab initio NCSMC

- include $p\text{-}^{11}\text{C}$ continuum ($1/2^-, 5/2^-, 3/2^-$ states of ^{11}C)
- include 4 negative and 6 positive parity states of ^{12}N
- MR-NO with respect to $N_{\text{max}}=0$ eigenstate of ^{12}N

^{12}N with continuum effects



- exciting **progress** in construction of **chiral forces**
 - **N²LO_{opt/sat/sim/sep}** different strategies to fit LECs
 - **local interaction** up to N²LO for quantum Monte Carlo
Gezerlis, Tews, Epelbaum et al. Phys. Rev. C 90, 054323 (2014)
 - **minimally non-local NN** up to N³LO include Δ in TPE component
Piarulli, Girlanda, Marcucci et al. Phys. Rev. C 91, 024003 (2015)
 - ...

LENPIC Collaboration

- improved NN up to N⁴LO
Epelbaum et al. EPJ A 51, 5(2015); PRL 115, 122301(2015)
- **3N** up to N³LO
 - allow to vary cutoff and chiral order to quantify uncertainty

self-contained framework to
employ **present and future chiral NN+3N
+4N interactions** in a variety of many-body
methods

- **p-shell spectra** provide powerful testbed for chiral potentials
 - constrain and benchmark nuclear forces
- inclusion of **continuum effects** and 3N forces crucial
 - **normal-ordering technique** enable heavier targets and **complex projectiles**

Thank you! Merci!

● thanks to my collaborators

- P. Navrátil, R. Stroberg, J. Holt,
J. Dohet-Eraly
TRIUMF Vancouver, Canada

- R. Roth, K. Hebel, E. Gebrerufael
Institut für Kernphysik, TU Darmstadt

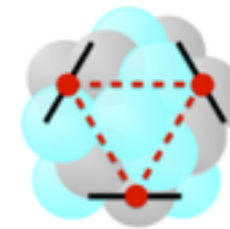
- S. Binder
University of Tennessee, Knoxville

- J. Vary, P. Maris
Iowa State University, USA

- H. Hergert
MSU, USA



LENPIC



- S. Quaglioni,
C. Romero-Redondo
LLNL Livermore, USA
- G. Hupin
Université Paris-Sud, France



Deutsche
Forschungsgemeinschaft
DFG



Exzellente Forschung für
Hessens Zukunft



COMPUTING TIME