

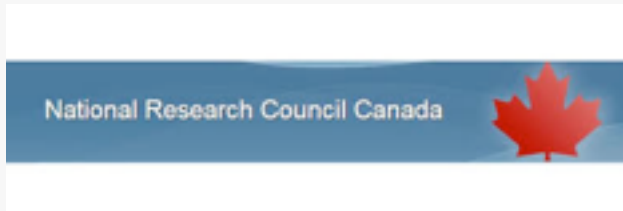
Towards ab-initio calculations of electromagnetic reactions in exotic nuclei

Sonia Bacca | Theory Department | TRIUMF
 | Physics Department | Univ. of Manitoba

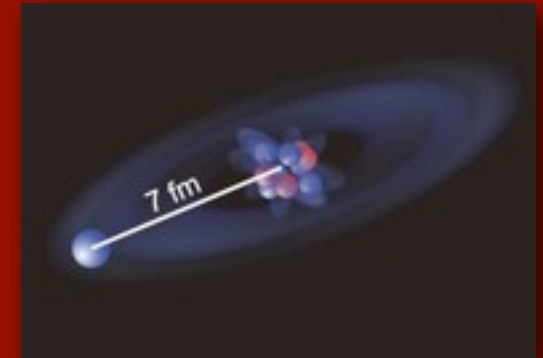
In collaboration with:

Nir Barnea, Gaute Hagen, Mirko Miorelli, Thomas Papenbrock,
 Giuseppina Orlandini, *et al.*

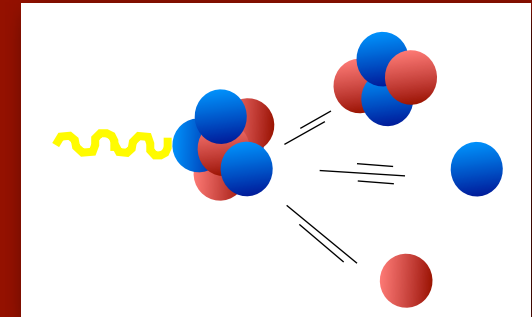
INT workshop on “Reactions and Structure of Exotic Nuclei”



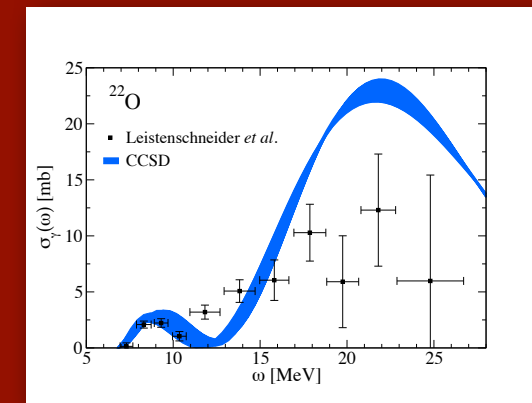
Owned and operated as a joint venture by a consortium of Canadian universities via a contribution through the National Research Council Canada
 Propriété d'un consortium d'universités canadiennes, géré en co-entreprise à partir d'une contribution administrée par le Conseil national de recherches Canada



Nuclear Halo



Electromagnetic Reactions



Pigmy Resonance

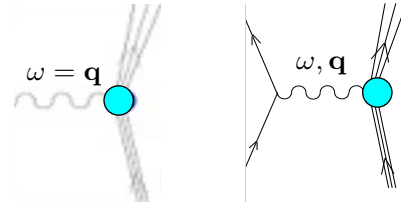
Motivations

- Electromagnetic probes (coupling constant $\ll 1$)

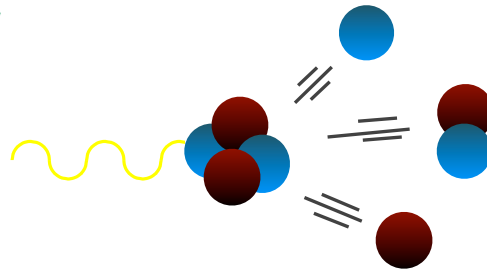
“With the electro-magnetic probe, we can immediately relate the cross section to the transition matrix element of the current operator, thus to the structure of the target itself”

[De Forest-Walecka, Ann. Phys. 1966]

$$\sigma \propto |\langle \Psi_f | J^\mu | \Psi_0 \rangle|^2$$



- For few-nucleons one can perform exact calculations both for bound and scattering states \Rightarrow test the nuclear theory on light nuclei and **extend it to heavier mass number**



- Provide important informations in other fields of physics, where nuclear physics plays a crucial role:

- Astrophysics: γ interactions with nucleonic matter, radiative capture reactions, ν interactions with nucleonic matter (vector current as em)
- Atomic physics (nuclear corrections to atomic levels, etc.)
- Particle physics (neutrino-less double beta decay, neutrino-nucleus interactions)

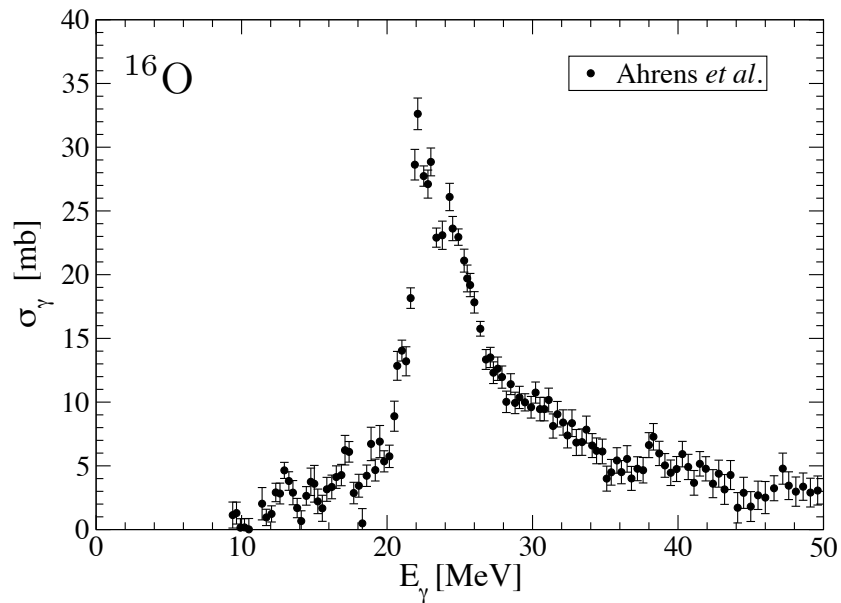


Electromagnetic Reactions

Photo-nuclear Reactions

Reactions resulting from the interaction of a photon with the nucleus

For photon energy 15-25 MeV stable nuclei across the periodic table show a wide and large peak

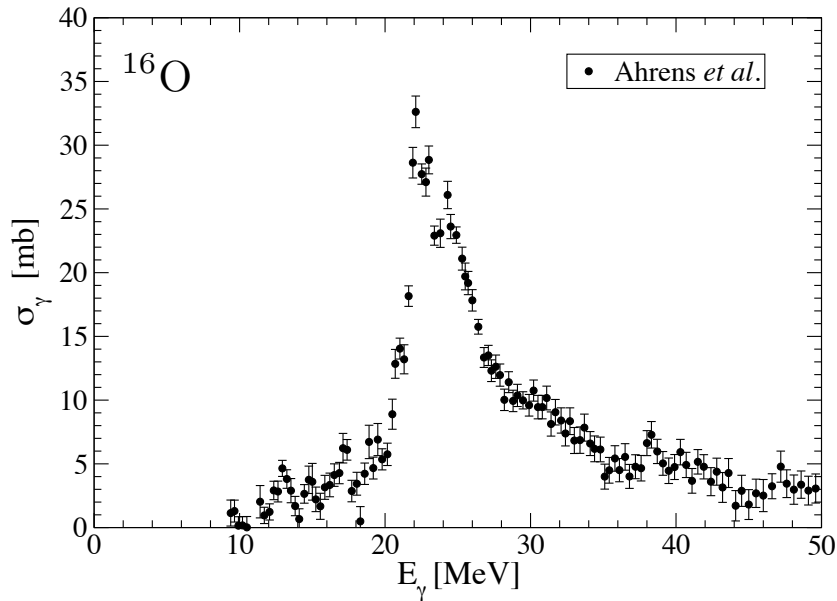


Electromagnetic Reactions

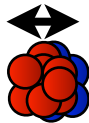
Photo-nuclear Reactions

Reactions resulting from the interaction of a photon with the nucleus

For photon energy 15-25 MeV stable nuclei across the periodic table show a wide and large peak



Giant Dipole Resonance



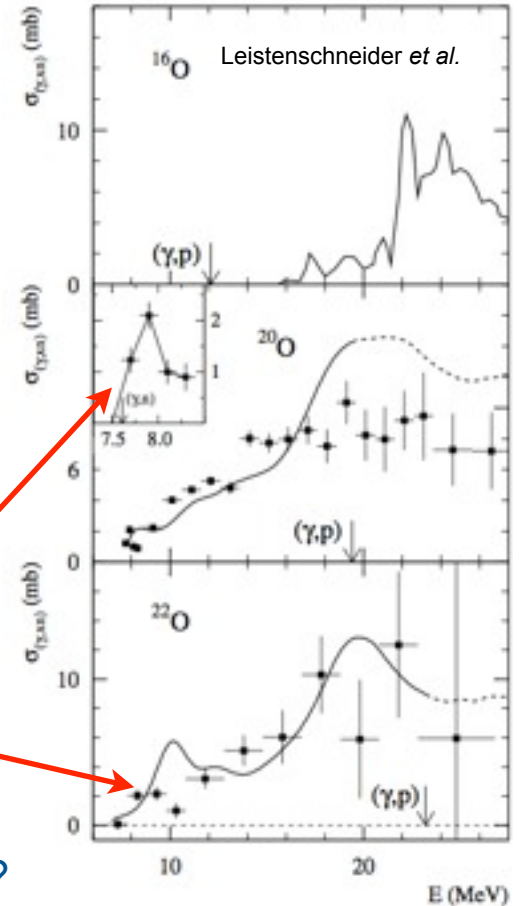
Soft Dipole Mode

- Can we give a microscopic explanation of these observations?

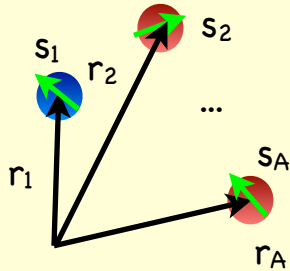
Coulomb excitations

Inelastic scattering between two charged particles. Can use unstable nuclei as projectiles.

Neutron-rich nuclei show fragmented low-lying strength



Ab-initio Theory Tools

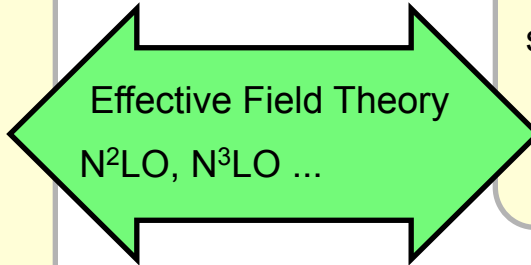
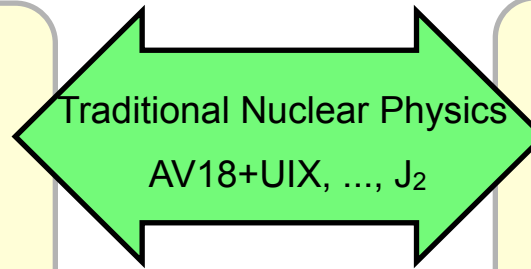


$$H|\psi_i\rangle = E_i|\psi_i\rangle$$

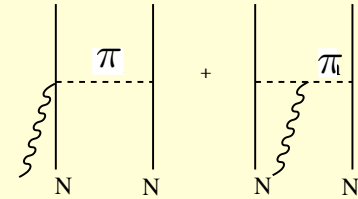
$$H = T + V_{NN} + V_{3N} + \dots$$

High precision two-nucleon potentials:
well constrained on NN phase shifts

Three nucleon forces:
less known, constraint on $A>2$ observables



$$J^\mu = J_N^\mu + J_{NN}^\mu + \dots$$



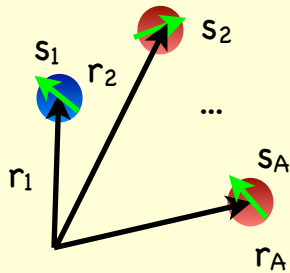
two-body currents (or MEC)
subnuclear d.o.f.

$$J^\mu \text{ consistent with } V$$

$$\nabla \cdot J = -i[V, \rho]$$

Review Paper:
Electromagnetic Reactions on Light Nuclei
S. Bacca and S. Pastore
J. Phys. G: Nucl. Part. Phys. **41** 123002 (2014).

Ab-initio Theory Tools

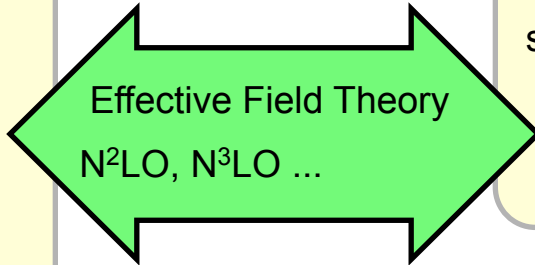
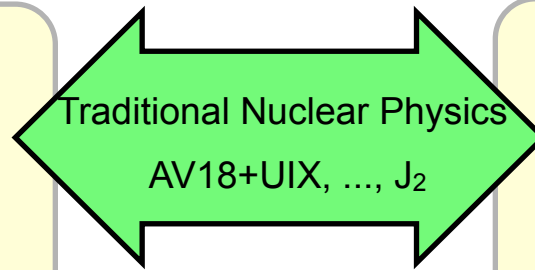


$$H|\psi_i\rangle = E_i|\psi_i\rangle$$

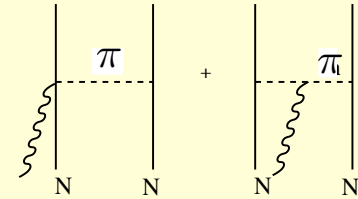
$$H = T + V_{NN} + V_{3N} + \dots$$

High precision two-nucleon potentials:
well constrained on NN phase shifts

Three nucleon forces:
less known, constraint on $A > 2$ observables



$$J^\mu = J_N^\mu + J_{NN}^\mu + \dots$$



two-body currents (or MEC)
subnuclear d.o.f.

$$J^\mu \text{ consistent with } V$$

$$\nabla \cdot J = -i[V, \rho]$$

$$\sigma \propto |\langle \Psi_f | J^\mu | \Psi_0 \rangle|^2$$

Exact Initial state &
Final state in the continuum at different energies and for different A

Review Paper:
Electromagnetic Reactions on Light Nuclei
S. Bacca and S. Pastore
J. Phys. G: Nucl. Part. Phys. **41** 123002 (2014).

Lorentz Integral Transform Method

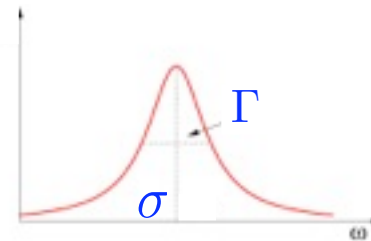
Efros, *et al.*, JPG.: Nucl.Part.Phys. **34** (2007) R459

Reduce the continuum problem to a bound-state problem



$$R(\omega) = \sum_f \left| \langle \psi_f | J^\mu | \psi_0 \rangle \right|^2 \delta(E_f - E_0 - \omega)$$

$$L(\sigma, \Gamma) = \int d\omega \frac{R(\omega)}{(\omega - \sigma)^2 + \Gamma^2} = \langle \tilde{\psi} | \tilde{\psi} \rangle < \infty$$



where $|\tilde{\psi}\rangle$ is obtained solving

$$(H - E_0 - \sigma + i\Gamma) |\tilde{\Psi}\rangle = J^\mu |\Psi_0\rangle$$

- Due to imaginary part Γ the solution $|\tilde{\psi}\rangle$ is unique
- Since $\langle \tilde{\psi} | \tilde{\psi} \rangle$ is finite, $|\tilde{\psi}\rangle$ has bound state asymptotic behaviour



$$L(\sigma, \Gamma) \xrightarrow{\text{inversion}} R(\omega)$$

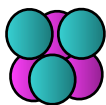
The exact final state interaction (FSI) is included in the continuum rigorously!

Solved for A=3,4,6,7 with **hyper-spherical harmonics expansions** and for A=4 with **NCSM**

Giant Dipole Resonance in $A=6$ with Hyperspherical Harmonics

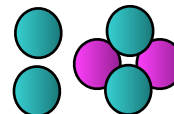
$$\sigma_\gamma = \frac{4\pi^2\alpha}{3} \omega R^{E1}(\omega)$$

$$E1 = \sum_i^Z (z_i - Z_{cm})$$



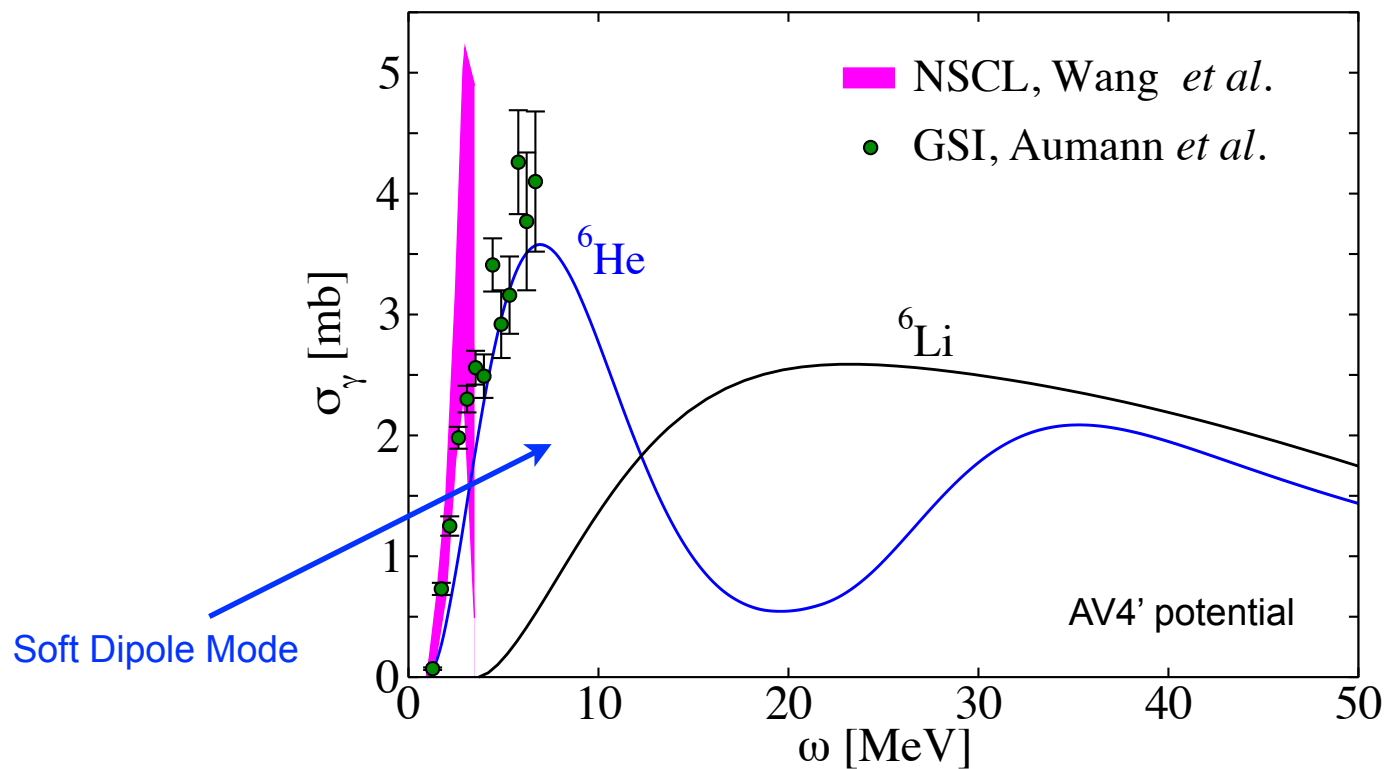
${}^6\text{Li}$

stable



${}^6\text{He}$

unstable $T_{1/2} = 806$ ms

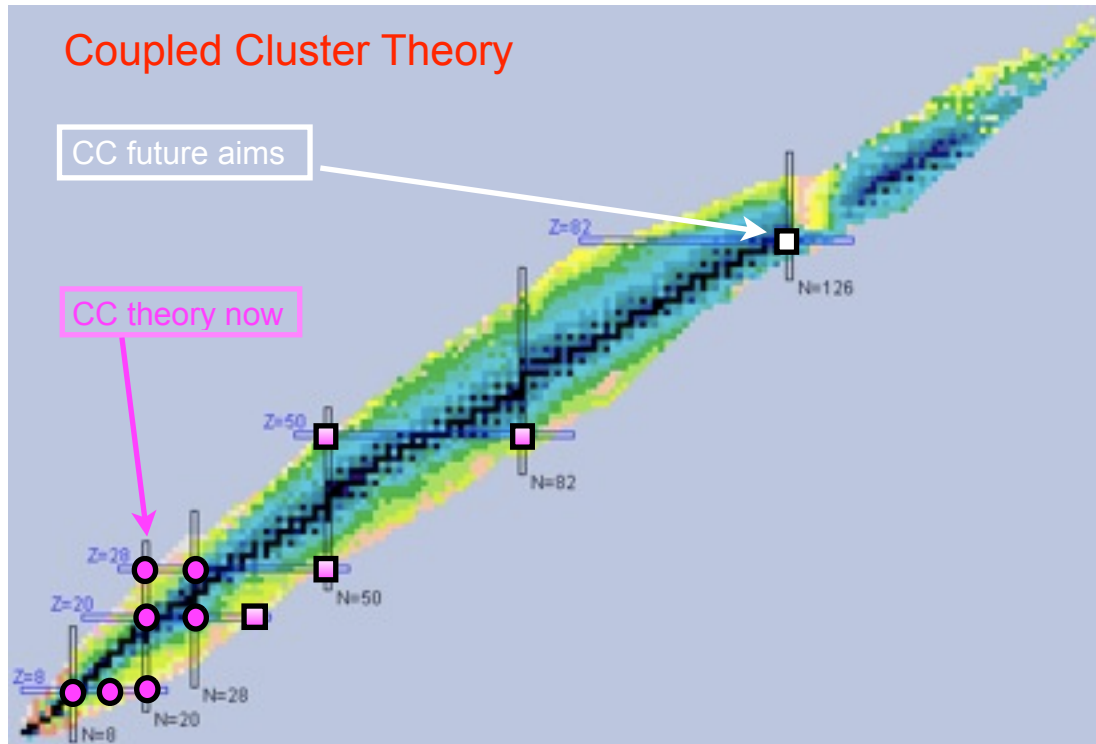


S.B. *et al.* PRL **89** 052502/PRC **69** 052502 (2004)

Extension to medium-mass nuclei

Develop new many-body methods that can extend the frontiers to heavier and neutron nuclei

Coupled Cluster Theory



- CC is optimal for closed shell nuclei ($\pm 1, \pm 2$)

Uses particle coordinates

$$|\psi_0(\vec{r}_1, \vec{r}_2, \dots, \vec{r}_A)\rangle = e^T |\phi_0(\vec{r}_1, \vec{r}_2, \dots, \vec{r}_A)\rangle$$

→ reference SD with any sp states

$$T = \sum T_{(A)} \quad \text{cluster expansion}$$

$$T_1 = \sum_{ia} t_i^a a_a^\dagger a_i$$

$$T_2 = \frac{1}{4} \sum_{ij,ab} t_{ij}^{ab} a_a^\dagger a_b^\dagger a_j a_i \quad \dots$$

T_1 T_2 T_3



CCSD
CCSDT

For the ground state energy

$$E_0 = \langle \phi_0 | e^{-T} H e^T | \phi_0 \rangle \quad \bar{H} = e^{-T} H e^T \quad \text{similarity transformed Hamiltonian}$$

$$0 = \langle \phi_i^a | e^{-T} H e^T | \phi_0 \rangle$$

$$0 = \langle \phi_{ij}^{ab} | e^{-T} H e^T | \phi_0 \rangle$$

Leads to CCSD equations for the t-amplitudes

Model space truncation $N \leq N_{max}$

Computational load $n_o^2 n_u^4$

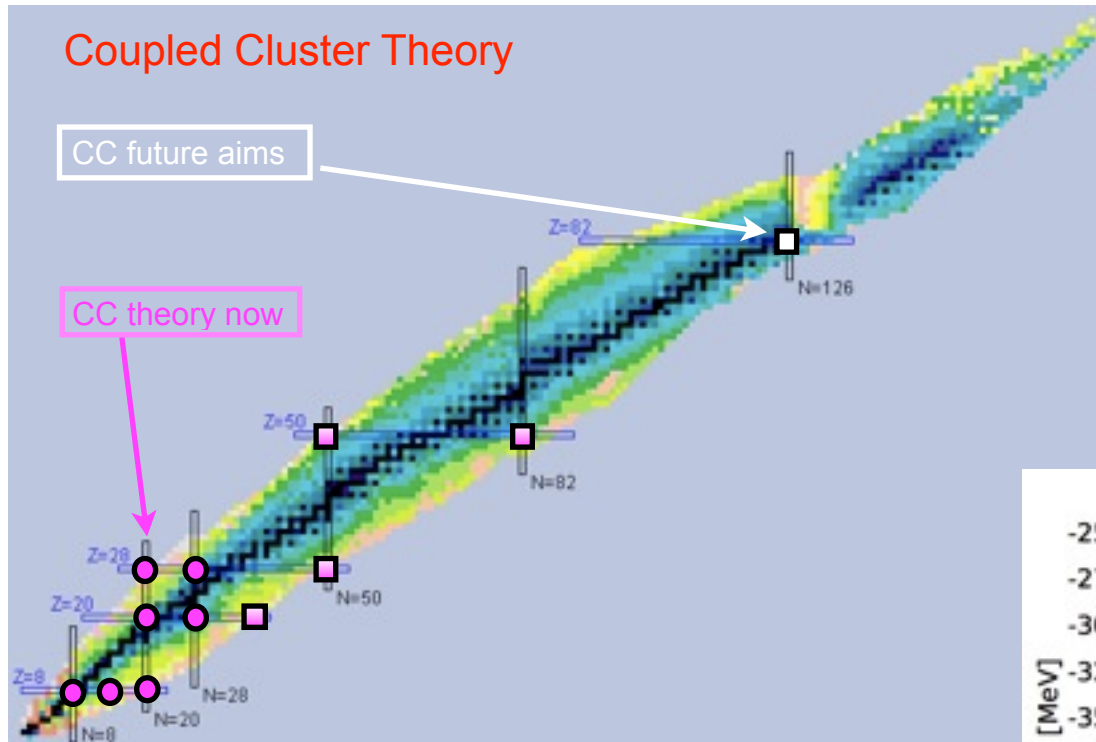
Extension to medium-mass nuclei

Develop new many-body methods that can extend the frontiers to heavier and neutron nuclei

Coupled Cluster Theory

CC future aims

CC theory now



- CC is optimal for closed shell nuclei ($\pm 1, \pm 2$)

Uses particle coordinates

$$|\psi_0(\vec{r}_1, \vec{r}_2, \dots, \vec{r}_A)\rangle = e^T |\phi_0(\vec{r}_1, \vec{r}_2, \dots, \vec{r}_A)\rangle$$

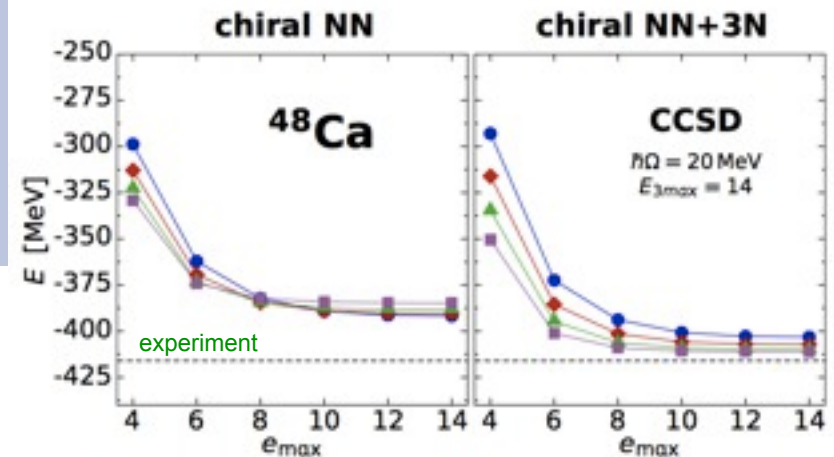
→ reference SD with any sp states

CC is a very mature theory for g.s., see e.g.

Hagen *et al.* PRL **101**, 092502 (2008), PRC **82**, 03433 (2010)
PRL **108**, 242501 (2012), PRL **109**, 032502 (2012)

What about electro-weak reactions?

LIT+CC can possibly extend calculations of inelastic reactions into medium-mass nuclei!



R. Roth *et al.*, Phys. Rev. Lett. **109**, 052501 (2012)

LIT with Coupled Cluster Theory

New theoretical method aimed at extending *ab-initio* calculations towards medium-mass

S.B. *et al.*, PRL **111**, 122502 (2013)

$$(H - z^*)|\tilde{\Psi}\rangle = J^\mu|\psi_0\rangle$$

with $z = E_0 + \sigma + i\Gamma$

$$L(\sigma, \Gamma) = \langle \tilde{\Psi} | \tilde{\Psi} \rangle$$



$$(\bar{H} - z^*)|\tilde{\Psi}_R(z^*)\rangle = \bar{\Theta}|\Phi_0\rangle$$

$$\bar{H} = e^{-T} H e^T$$

$$\bar{\Theta} = e^{-T} \Theta e^T$$

$$L(\sigma, \Gamma) = \langle \tilde{\Psi}_L | \tilde{\Psi}_R \rangle = -\frac{1}{2\pi} \Im \left\{ \langle \bar{0}_L | \bar{\Theta}^\dagger \left[|\tilde{\Psi}_R(z^*)\rangle - |\tilde{\Psi}_R(z)\rangle \right] \right\}$$

with $|\tilde{\Psi}_R(z^*)\rangle = \hat{R}(z^*)|\Phi_0\rangle$

Formulation based on the solution of an **Equation of Motion with source** No approximations done so far!

Present implementation in the CCSD scheme

$$T = T_1 + T_2$$

$$\hat{R} = \hat{R}_0 + \hat{R}_1 + \hat{R}_2$$

Chiral Effective Field Theory

Weinberg, van Kolck, Kaplan, Savage, Wise,
Epelbaum, Meissner, Nogga, Machleidt,...

Construct an effective Lagrangian which respects the chiral symmetry and is written as an expansion in powers of $\frac{Q}{\Lambda_b}$

$$\mathcal{L} = \sum_k c_k \left(\frac{Q}{\Lambda_b} \right)^k$$

Details of short distance physics not resolved, but captured in low energy coupling constants, fit to experiment once

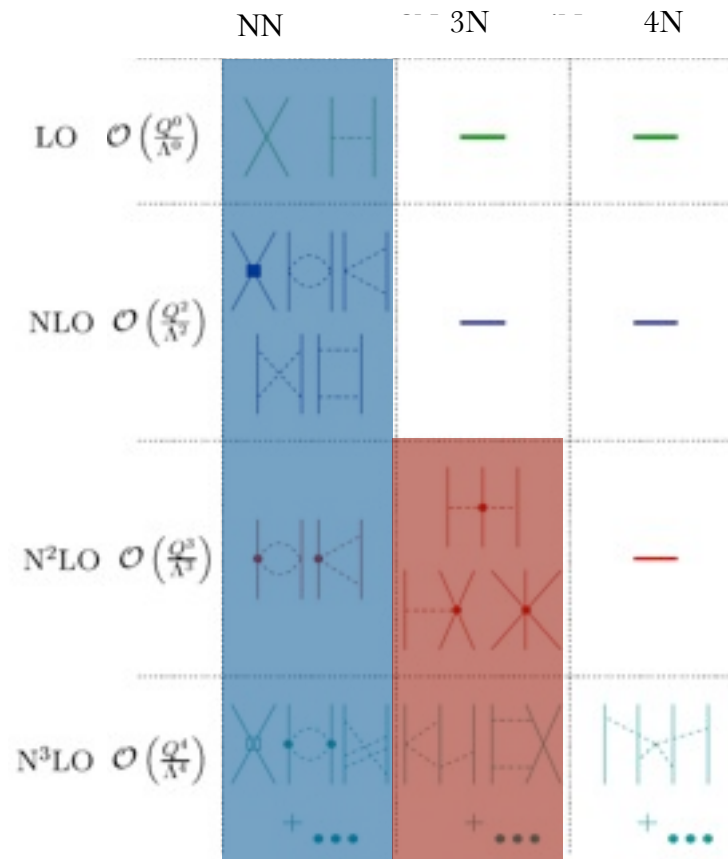
Systematic expansion of the potential in many-body terms

Traditional Paradigm:

calibrate NN on NN scattering data
calibrate 3N on 3N data

Tests and Results from NN only

$$H(\Lambda) = T + V_{NN}(\Lambda) + V_{3N}(\Lambda) + V_{4N}(\Lambda) + \dots$$



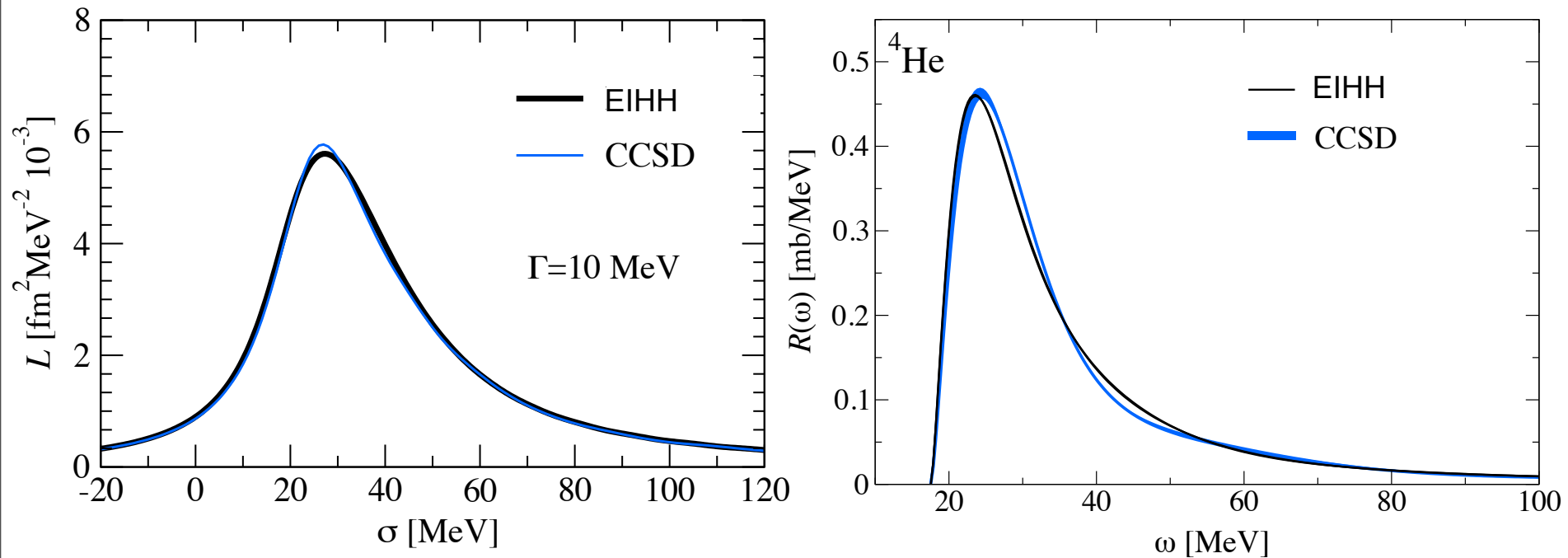
$$V_{NN} > V_{3N} > V_{4N}$$

LIT with Coupled Cluster Theory

New theoretical method aimed at extending *ab-initio* calculations towards medium mass

Validation for ${}^4\text{He}$

➔ Comparison of CCSD with exact hyperspherical harmonics (EIHH) with NN forces at $N^3\text{LO}$



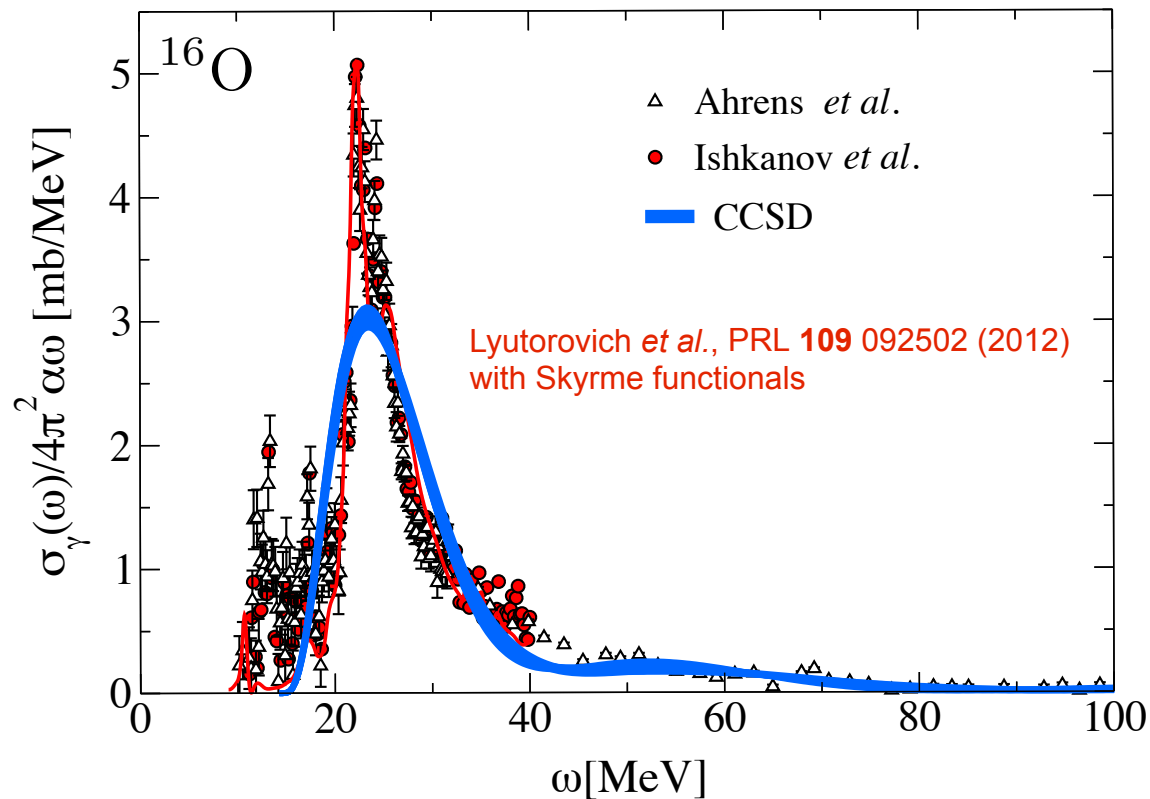
The comparison with exact theory is very good!

LIT with Coupled Cluster Theory

New theoretical method aimed at extending *ab-initio* calculations towards medium mass

S.B. *et al.*, PRL **111**, 122502 (2013)

Extension to Dipole Response Function in ^{16}O with NN forces derived from χEFT (N^3LO)

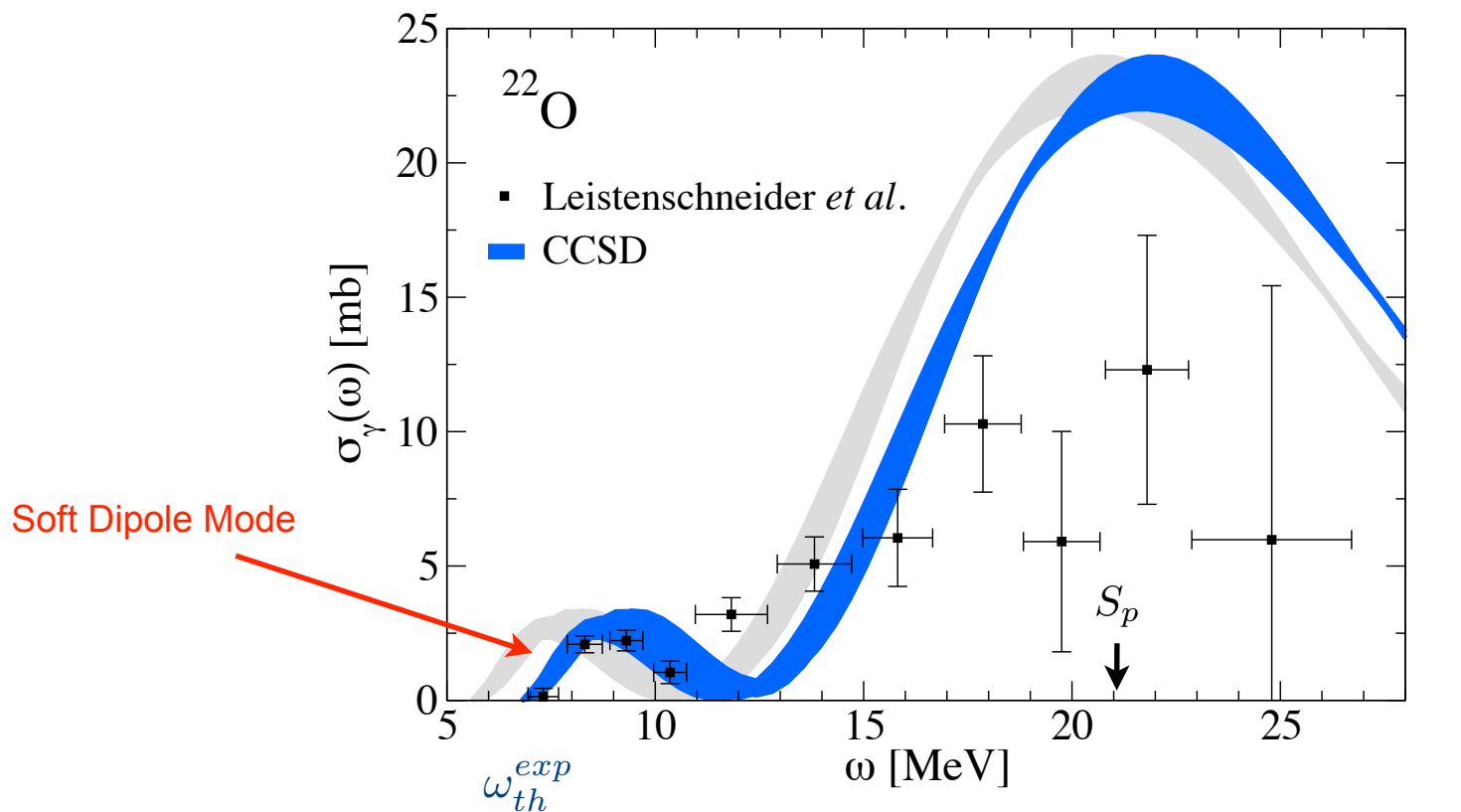


The position of the GDR in ^{16}O is described from first principles for the first time!

Addressing neutron-rich nuclei

^{22}O with NN forces derived from χEFT (N^3LO)

S.B. *et al.*, PRC **90**, 064619 (2014)

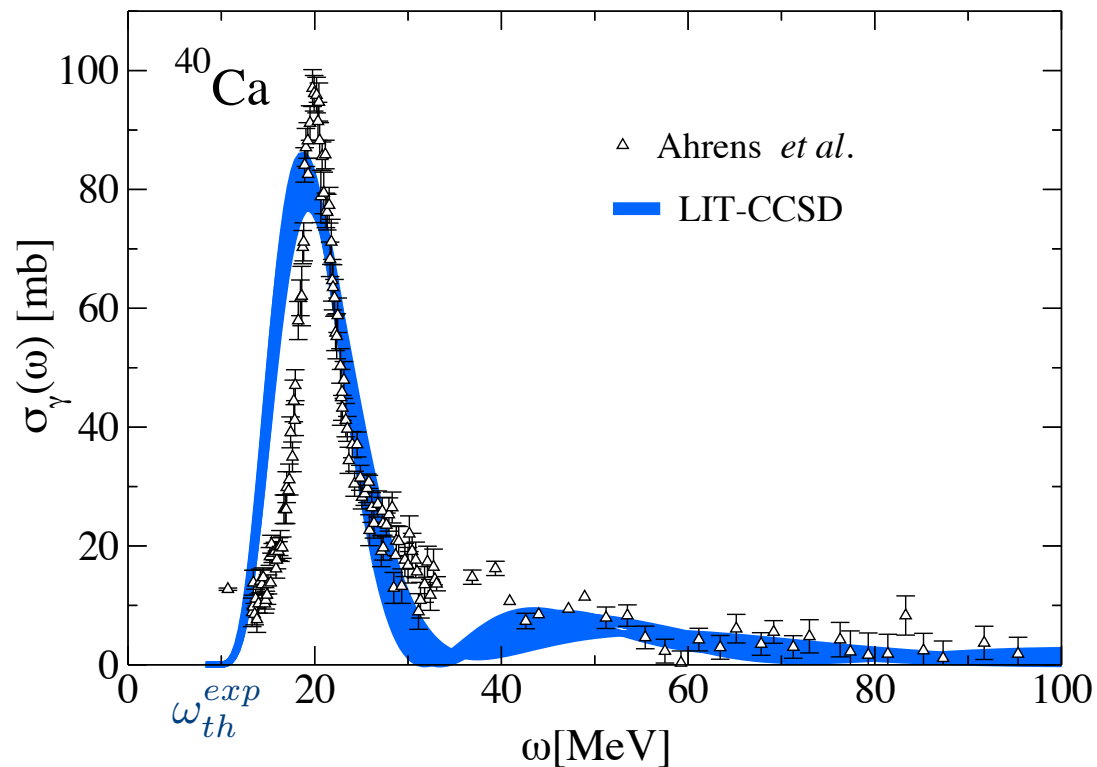


➔ Soft dipole mode emerges from a first principle calculation

Addressing heavier nuclei

^{40}Ca with NN forces derived from χEFT (N^3LO)

S.B. *et al.*, PRC **90**, 064619 (2014)



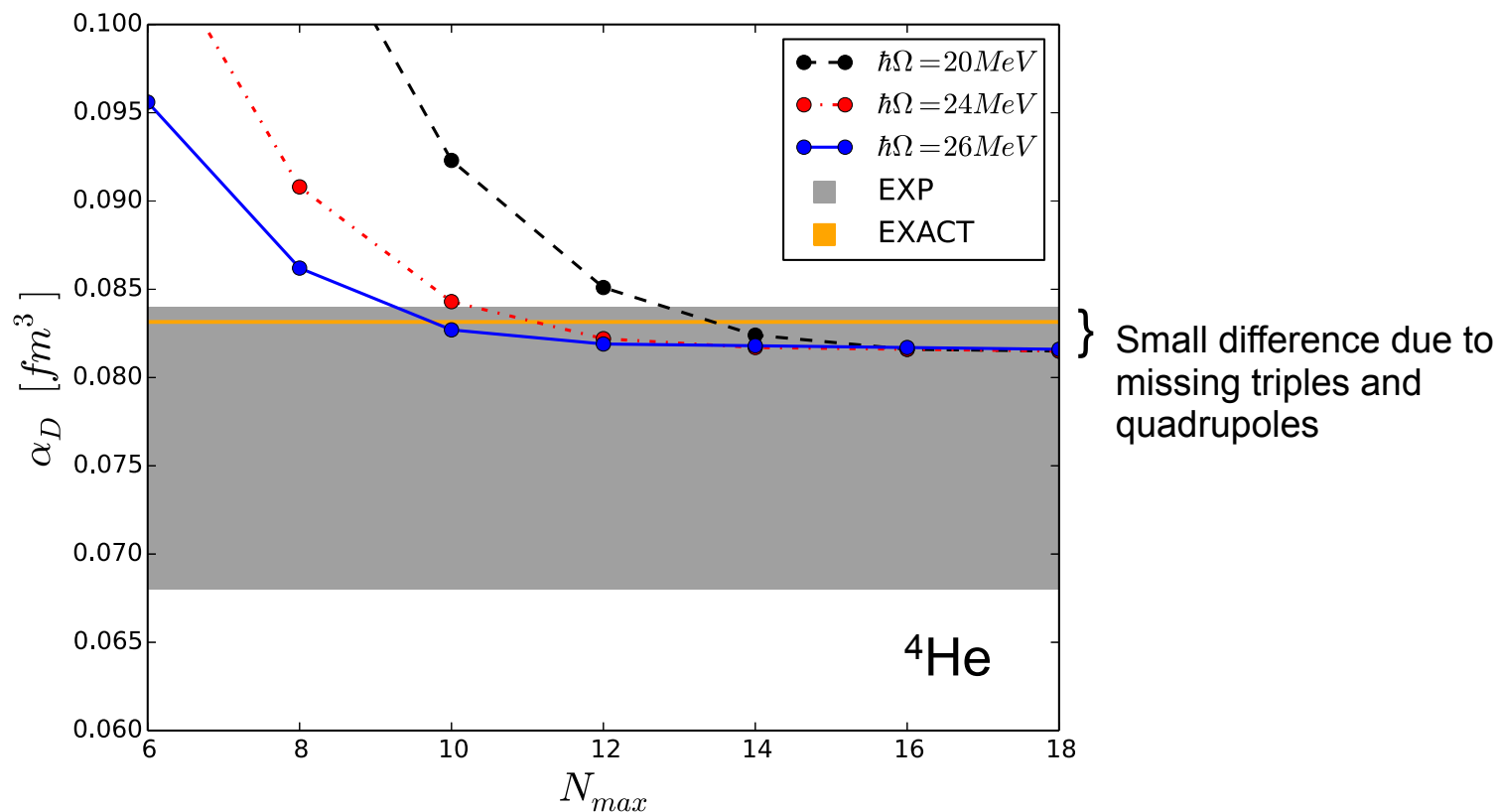
Electric Dipole Polarizability

$$\alpha_D = 2\alpha \int_{\omega_{th}}^{\infty} d\omega \frac{R(\omega)}{\omega}$$

Electric Dipole Polarizability

M. Miorelli *et al.*, in preparation (2015)

Validation for ${}^4\text{He}$ with NN(N³LO)

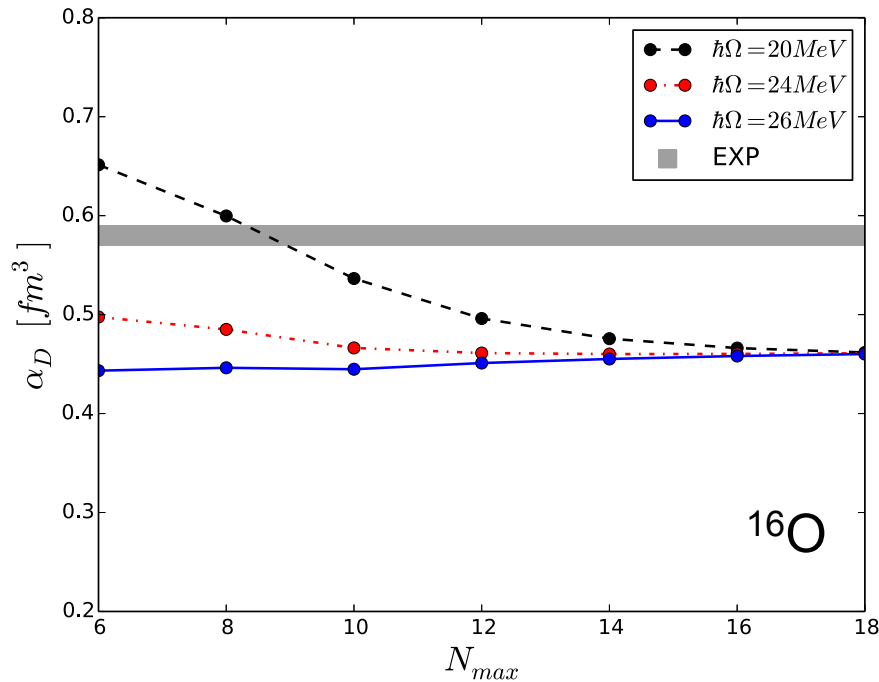


The comparison with experiment is very good

Electric Dipole Polarizability

M. Miorelli *et al.*, in preparation (2015)

Medium-mass nuclei with NN(N³LO)

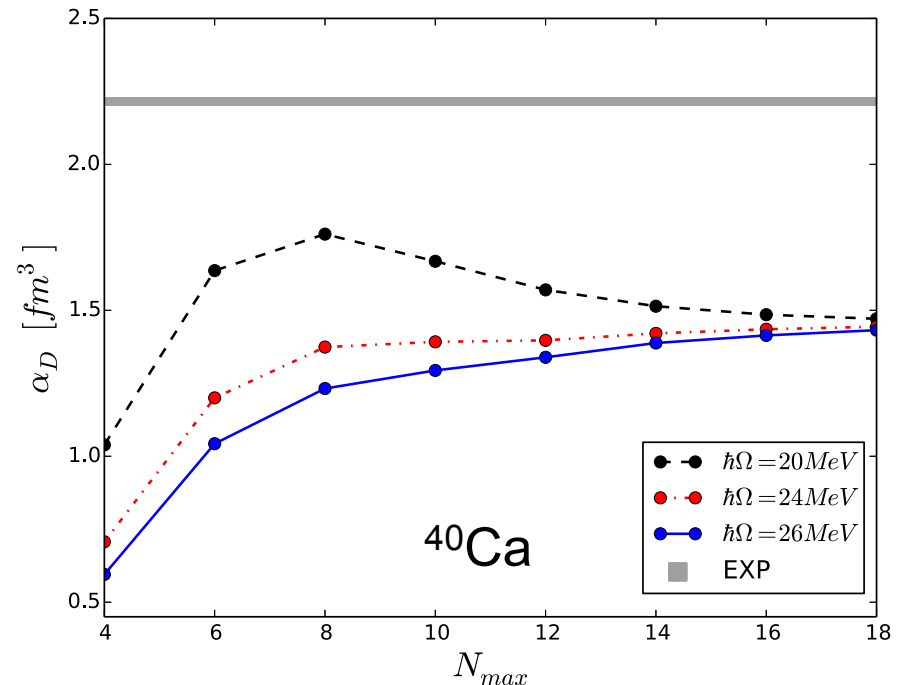


$$\alpha_D = 0.46 \text{ fm}^3$$

$$\alpha_D^{exp} = 0.585(9) \text{ fm}^3$$

$$R_{ch} = 2.3 \text{ fm}$$

$$R_{ch}^{exp} = 2.6991(52) \text{ fm}$$



$$\alpha_D = 1.47 \text{ fm}^3$$

$$\alpha_D^{exp} = 2.23(3) \text{ fm}^3$$

$$R_{ch} = 3.05 \text{ fm}$$

$$R_{ch}^{exp} = 3.4776(19) \text{ fm}$$

The present Hamiltonian underestimates both radii and electric dipole polarizabilities

Conclusions and Outlook

- Electromagnetic break up reactions are very rich observables to test our understanding of nuclear forces
- Extending first principles calculations to medium mass nuclei is possible and very exciting: more applications/impact on experiments in the future

Perspectives

- Other multipole excitation: electric quadrupole or monopole, magnetic dipole in medium mass nuclei
- Add triples excitations: \hat{R}_3, T_3
- Extend first principle calculations to the weak sector: Gamow-Teller transitions
See talk by Gaute Hagen on Wed

Thank you!