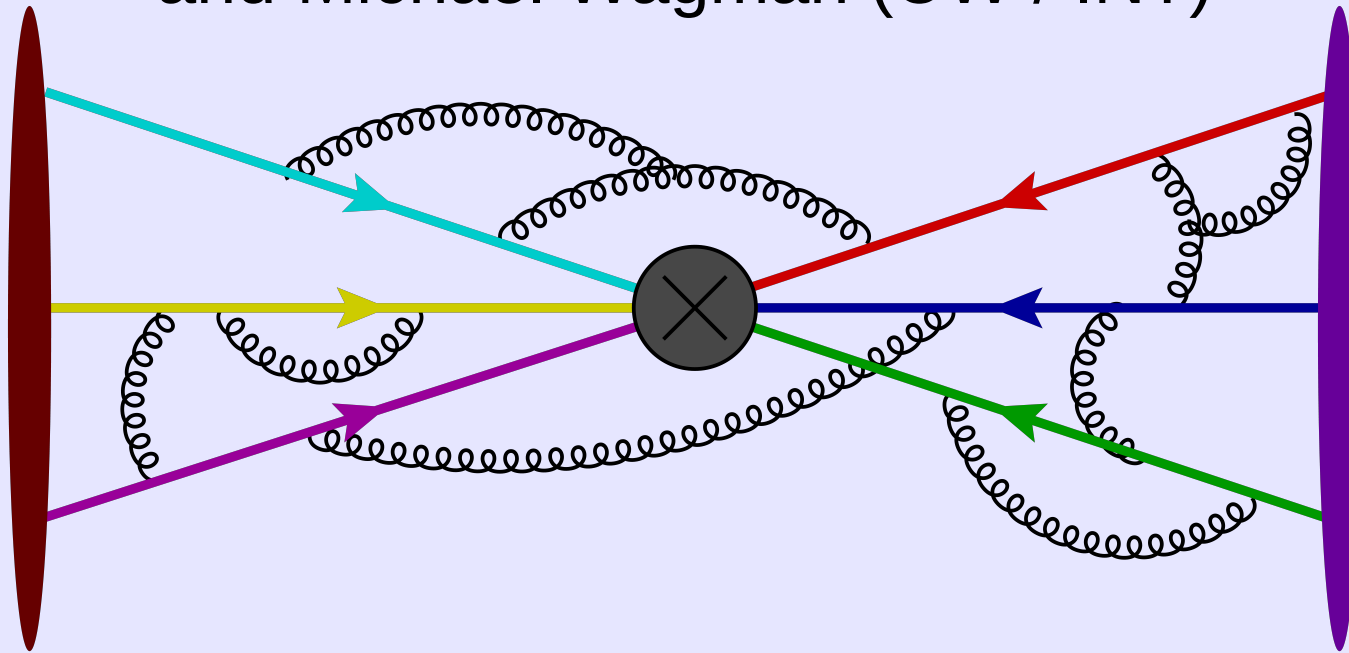


# Neutron-Antineutron Perturbative Operator Renormalization

Michael I. Buchoff (UW / INT / LLNL)  
and Michael Wagman (UW / INT)

arxiv:1506.00647  
arxiv:1502.00044



Michael Wagman, 9/24/2015 Seminar for INT Program:

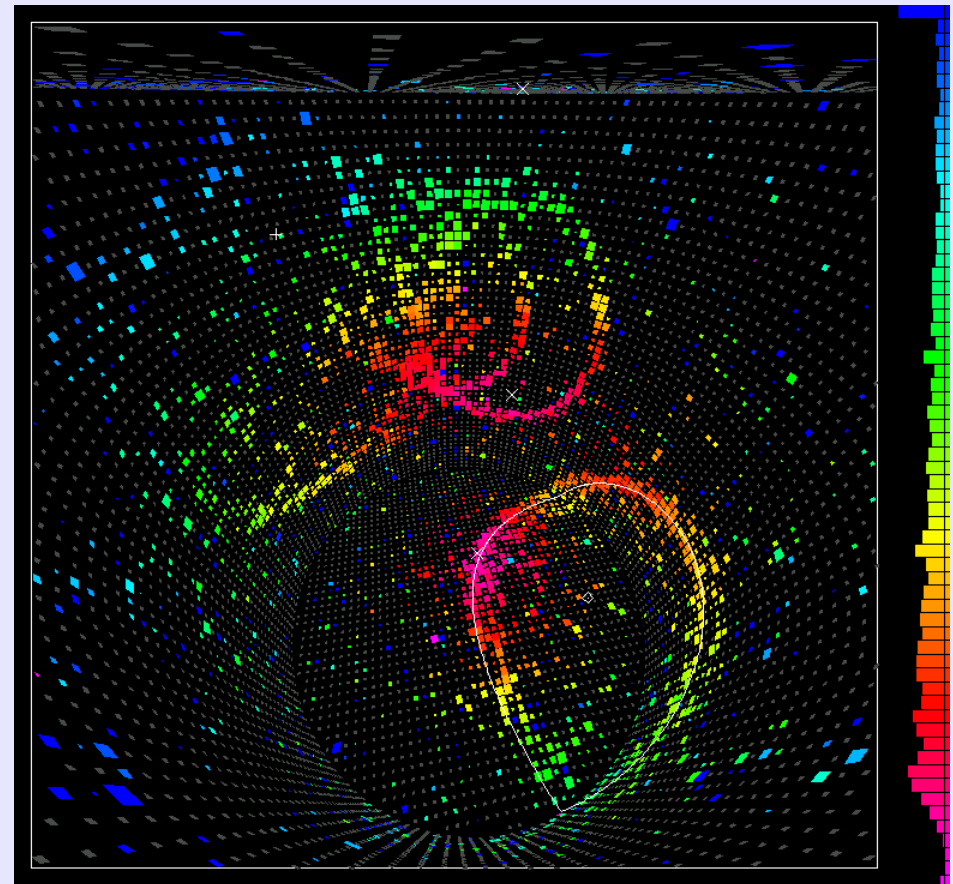
Intersections of BSM Phenomenology and QCD for New Physics Searches

# Neutron-Antineutron Transitions

Predicted by some BSM theories relevant for baryogenesis

Experimentally testable predictions require QCD matrix element calculations

Calculations with quantified uncertainty require lattice QCD and perturbative operator renormalization



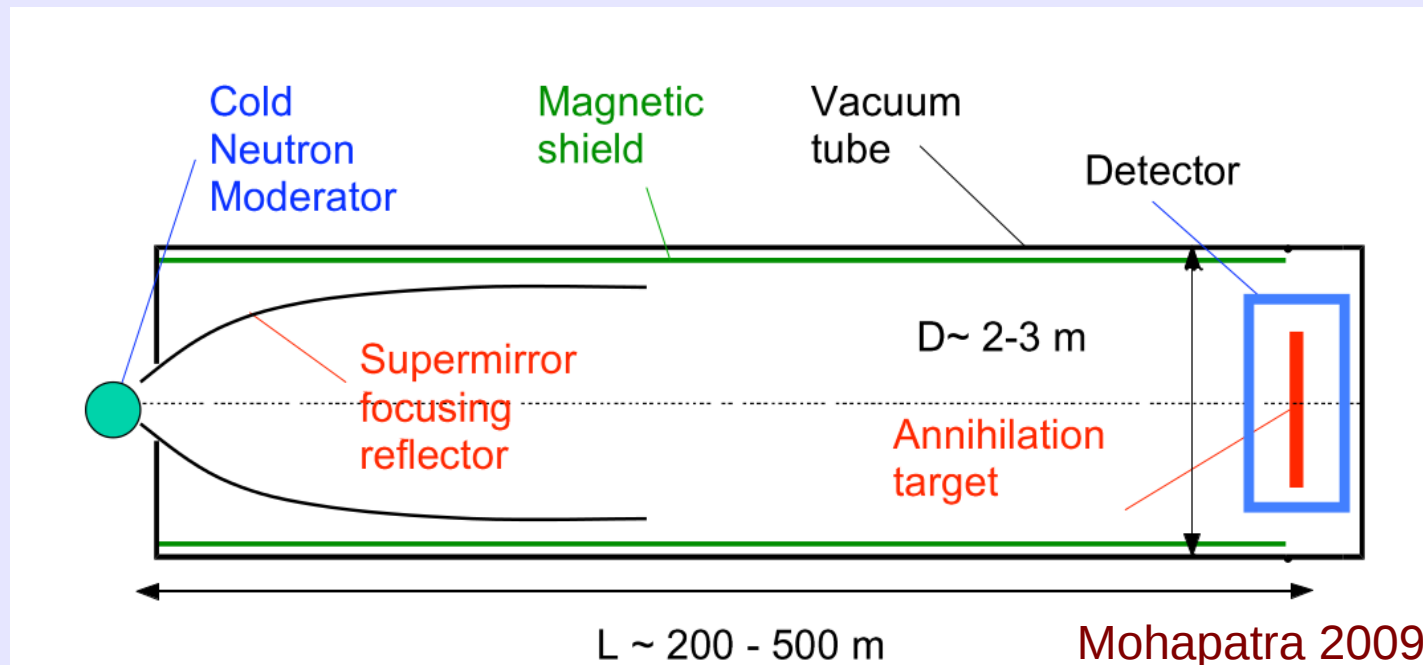
Super K, proton decay candidate event

# Neutron-Antineutron Phenomenology

Vacuum transition rate described by two-state Hamiltonian

$$\frac{1}{\tau_{n\bar{n}}} = \langle \bar{n} | H_{n\bar{n}} | n \rangle$$

Magnetic fields, nuclear interactions modify transition rate

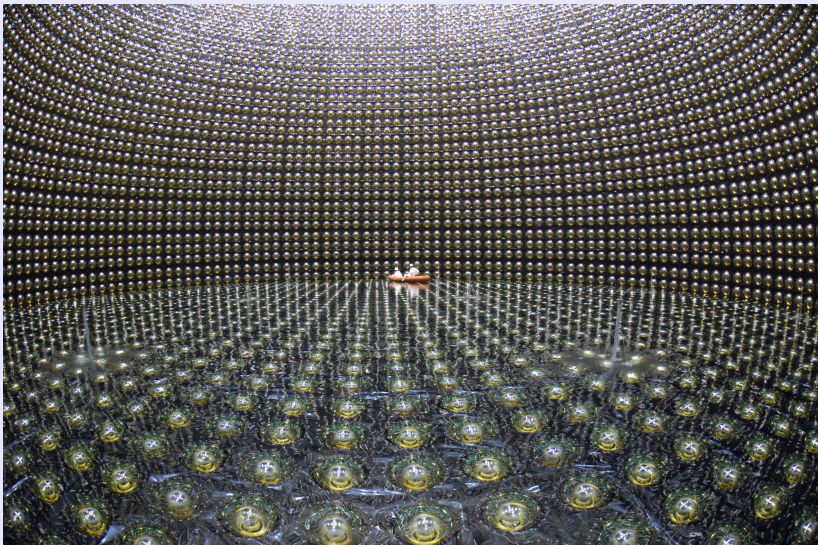




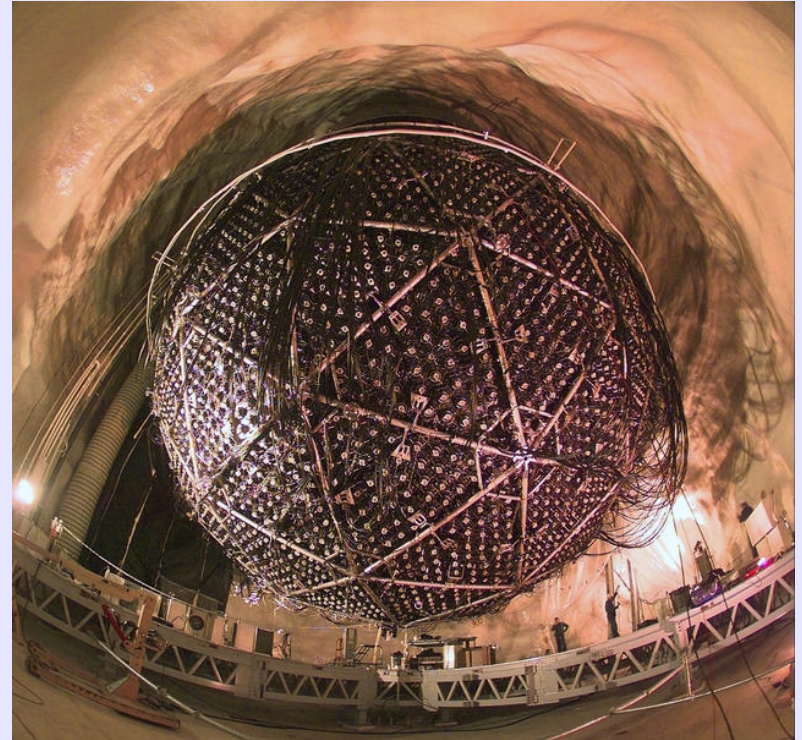
# Experimental Constraints



ILL:  $\tau_{n\bar{n}} > 2.7$  years



Super K:  $\tau_{n\bar{n}} > 11$  years



SNO:  $\tau_{n\bar{n}} > 5.7$  years  
(preliminary)



# Standard Model Results

Complete basis of gauge singlet six-quark operators,  
bag model estimates for QCD matrix elements

Rao and Shrock (1984)

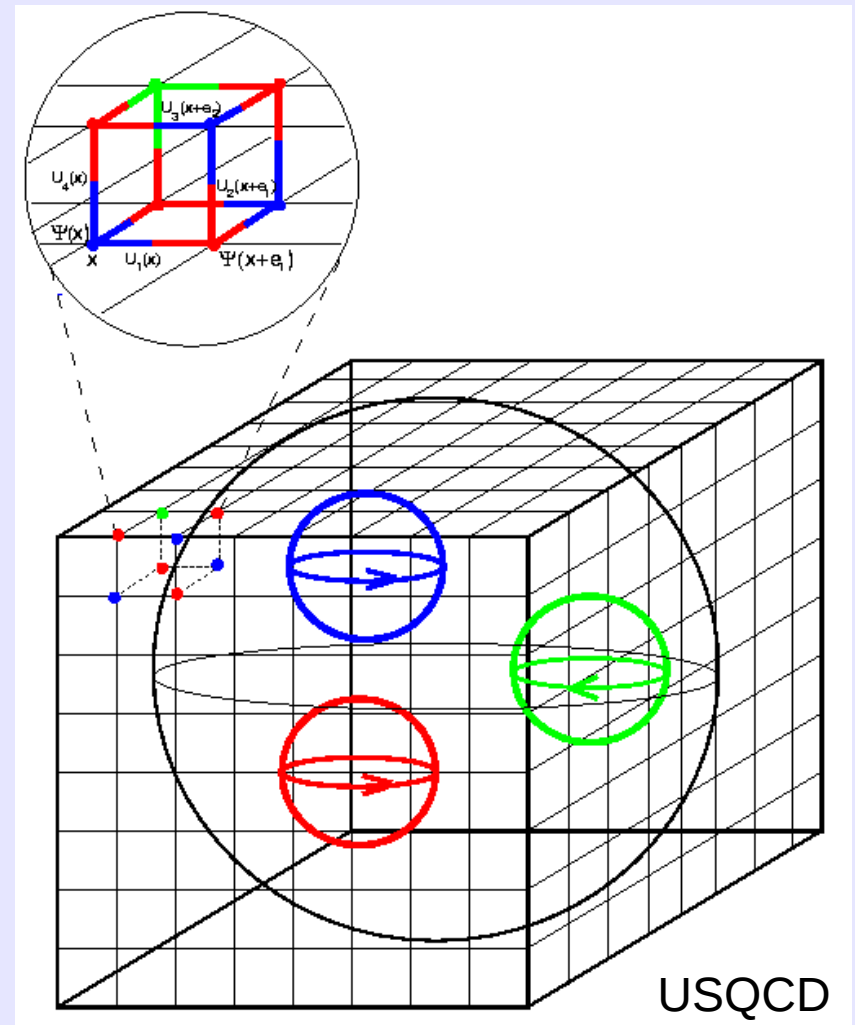
$$\mathcal{H}_{n\bar{n}} = \sum_{I=1}^{14} C_I(\mu) Q_I(\mu)$$

1-Loop QCD (and QED)  
anomalous dimensions

Caswell, Milutinovic, and Senjanovic (1983)

Preliminary lattice QCD matrix  
elements

Buchoff, Schroeder, and Wasem (2012)


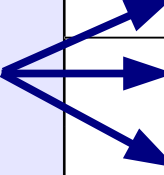
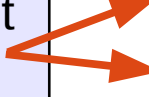


USQCD

$$\psi = \begin{pmatrix} u \\ d \end{pmatrix}$$

# The Operators

$$Q_1 = (\psi C P_R i \tau^2 \psi)(\psi C P_R i \tau^2 \psi)(\psi C P_R i \tau^2 \tau^+ \psi) T^{AAS} = \mathcal{D}_R \mathcal{D}_R \mathcal{D}_R^+ T^{AAS}$$

	Chiral Basis	Fixed-Flavor Basis	Chiral Tensor Structure	Chiral Irrep
	$Q_1$	$\mathcal{O}_{RRR}^3$	$\mathcal{D}_R \mathcal{D}_R \mathcal{D}_R^+ T^{AAS}$	$(\mathbf{1}_L, \mathbf{3}_R)$
	$Q_2$	$\mathcal{O}_{LRR}^3$	$\mathcal{D}_L \mathcal{D}_R \mathcal{D}_R^+ T^{AAS}$	$(\mathbf{1}_L, \mathbf{3}_R)$
	$Q_3$	$\mathcal{O}_{LLR}^3$	$\mathcal{D}_L \mathcal{D}_L \mathcal{D}_R^+ T^{AAS}$	$(\mathbf{1}_L, \mathbf{3}_R)$
Isospin 3 	$Q_4$	$4/5 \mathcal{O}_{RRR}^2 + 1/5 \mathcal{O}_{RRR}^1$	$\mathcal{D}_R^{33+} T^{SSS}$	$(\mathbf{1}_L, \mathbf{7}_R)$
Not SM gauge invariant 	$Q_5$	$\mathcal{O}_{RLL}^1$	$\mathcal{D}_R^- \mathcal{D}_L^{++} T^{SSS}$	$(\mathbf{5}_L, \mathbf{3}_R)$
	$Q_6$	$\mathcal{O}_{RLL}^2$	$\mathcal{D}_R^3 \mathcal{D}_L^{3+} T^{SSS}$	$(\mathbf{5}_L, \mathbf{3}_R)$
	$Q_7$	$2/3 \mathcal{O}_{LLR}^2 + 1/3 \mathcal{O}_{LLR}^1$	$\mathcal{D}_R^+ \mathcal{D}_L^{33} T^{SSS}$	$(\mathbf{5}_L, \mathbf{3}_R)$
Redundant in $D = 4$ 	$\tilde{Q}_1$	$1/3 \mathcal{O}_{RRR}^2 - 1/3 \mathcal{O}_{RRR}^1$	$\mathcal{D}_R \mathcal{D}_R \mathcal{D}_R^+ T^{SSS}$	$(\mathbf{1}_L, \mathbf{3}_R)$
	$\tilde{Q}_3$	$1/3 \mathcal{O}_{LLR}^2 - 1/3 \mathcal{O}_{LLR}^1$	$\mathcal{D}_L \mathcal{D}_L \mathcal{D}_R^+ T^{SSS}$	$(\mathbf{1}_L, \mathbf{3}_R)$

Chiral symmetry provides a basis with no operator mixing

# Perturbative Renormalization

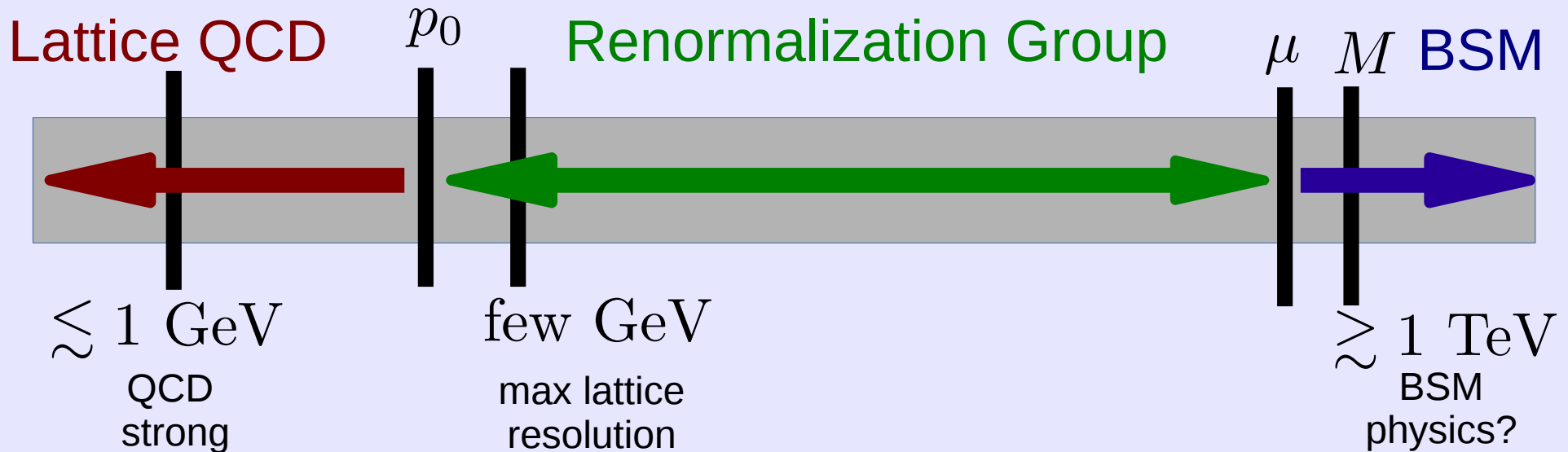
$$\mathcal{H}_{n\bar{n}} = \sum_I C_I^{\overline{MS}}(\mu) Q_I^{\overline{MS}}(\mu) = \sum_I C_I^{\overline{MS}}(\mu) U_I(\mu, p_0) Q_I^{RI}(p_0)$$

$$U_I(\mu, p_0) = \left( \frac{\alpha_s(p_0)}{\alpha_s(\mu)} \right)^{-\gamma_I^{(0)}/2\beta_0} \left[ 1 - r_I^{(0)} \frac{\alpha_s(p_0)}{4\pi} + \left( \frac{\beta_1 \gamma_I^{(0)}}{2\beta_0^2} - \frac{\gamma_I^{(1)}}{2\beta_0} \right) \frac{\alpha_s(p_0) - \alpha_s(\mu)}{4\pi} + O(\alpha_s^2) \right]$$

One-loop matching

Two-loop running,  
needed for  $\alpha_s(p_0)$  accuracy

Negligible





# Regularization-Independent Renormalization

RI-MOM scheme: apply momentum subtraction to vertex functions  $\Lambda_I$  at lattice matching scale  $p_0$

Martinelli et al (1995)

$$\begin{aligned}
 [\Lambda_I]_{ijklmn}^{\alpha\beta\gamma\delta\eta\zeta}(p) &= \frac{1}{5} \left\langle Q_I(0) \bar{u}_i^\alpha(p) \bar{u}_j^\beta(p) \bar{d}_k^\gamma(p) \bar{d}_l^\delta(-p) \bar{d}_m^\eta(-p) \bar{d}_n^\zeta(-p) \right\rangle \Big|_{amp} \\
 &+ \frac{3}{5} \left\langle Q_I(0) \bar{u}_i^\alpha(p) \bar{u}_j^\beta(-p) \bar{d}_k^\gamma(p) \bar{d}_l^\delta(p) \bar{d}_m^\eta(-p) \bar{d}_n^\zeta(-p) \right\rangle \Big|_{amp} \\
 &+ \frac{1}{5} \left\langle Q_I(0) \bar{u}_i^\alpha(-p) \bar{u}_j^\beta(-p) \bar{d}_k^\gamma(p) \bar{d}_l^\delta(p) \bar{d}_m^\eta(p) \bar{d}_n^\zeta(-p) \right\rangle \Big|_{amp}
 \end{aligned}$$

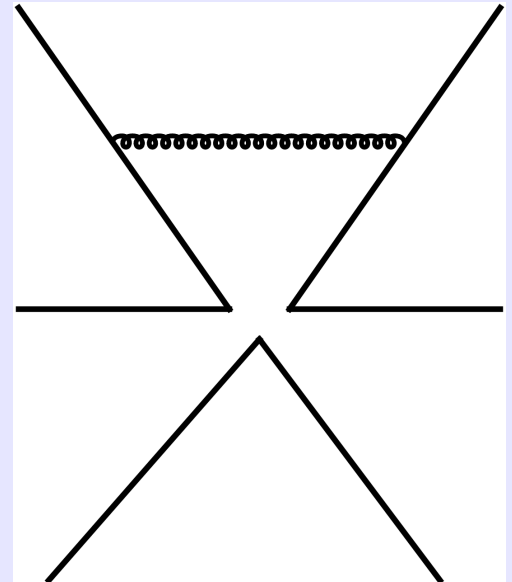
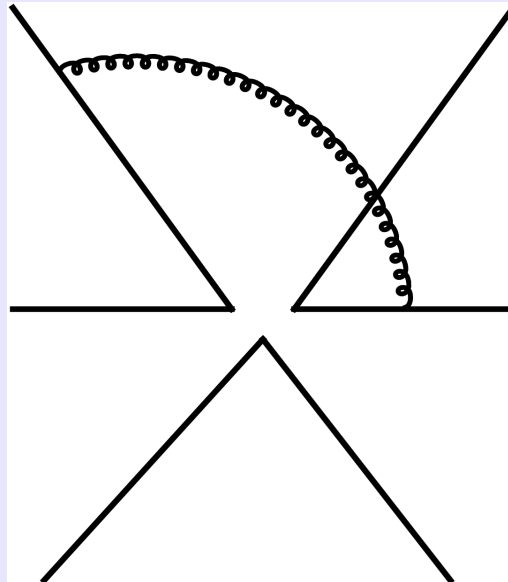
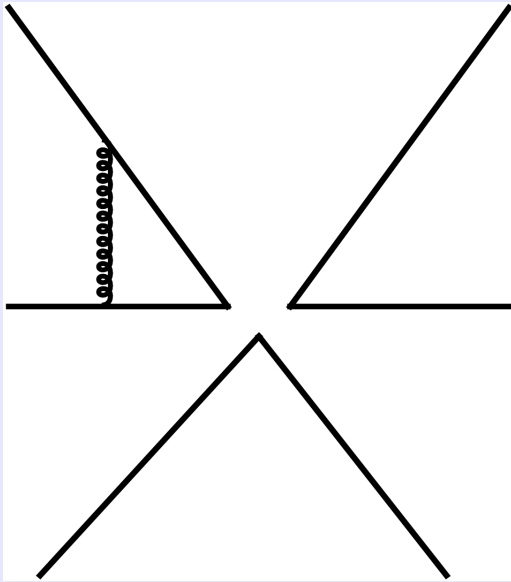
Syritsyn (2015)

RI-MOM condition applicable to lattice and dim reg operators

$\mathcal{P}_I$  projects out contribution from tree-level  $\Lambda_I^{(0)}$   $\longrightarrow [\mathcal{P}_I \Lambda_J(p_0, \mu = p_0)] = \delta_{IJ}$

# One-Loop Matching

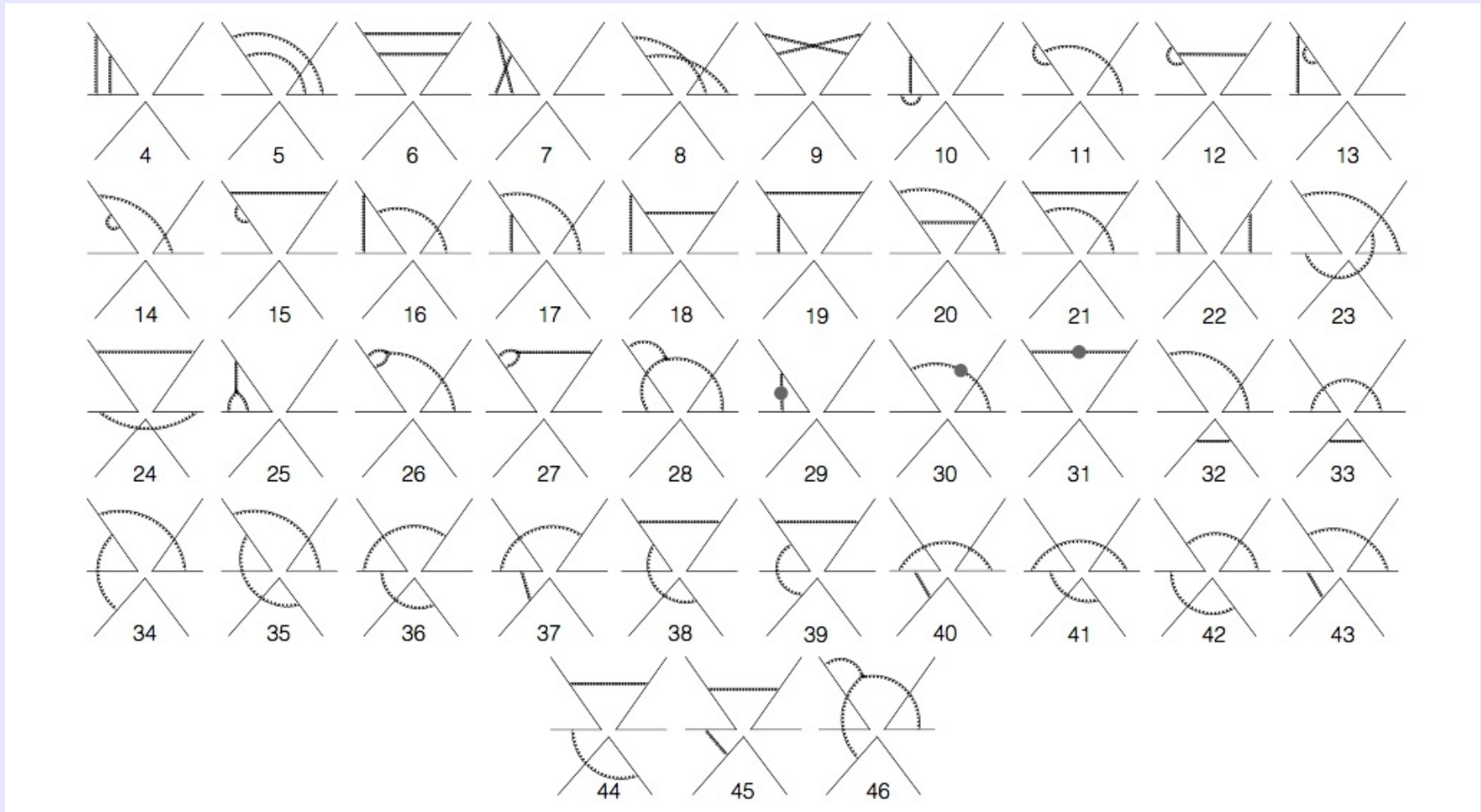
15 one-loop Feynman diagrams contribute



One quark pair is not involved in the interaction. Identical diagrams contribute to four-quark weak matrix elements and proton decay

# Two-Loop Running

Need  $\ln \mu^2$  coefficients of 320 two-loop Feynman diagrams



Covers all diagram topologies needed for two-loop running of any operator built from spin singlet quark pairs



# Evanescent Operators

Operator basis incomplete in dim reg, Fierz identities broken

$$Q_1 = (\psi C P_R i \tau^2 \psi) (\psi C P_R i \tau^2 \psi) (\psi C P_R i \tau^2 \tau^+ \psi) T^{AAS}$$

Vanishes in  $D = 4$

$$E_1^a \equiv (\psi C P_R \sigma_{\mu\nu} i \tau_2 \psi) (\psi C P_R \sigma_{\mu\nu} i \tau_2 \psi) (\psi C P_R i \tau_2 \tau^+ \psi) T^{SSS} - 12Q_1$$

Complete basis for one-loop counterterms found by adding evanescent operators


Buras and Weisz (1990), Dugan and Grinstein (1991), Herrlich and Nierste (1995), ...

$$\begin{aligned} \frac{1}{\epsilon} (\psi C P_R \sigma_{\mu\nu} i \tau_2 \psi) (\psi C P_R \sigma_{\mu\nu} i \tau_2 \psi) (\psi C P_R i \tau_2 \tau^+) T^{SSS} + (\overline{\text{MS}} \text{ counterterm}) \\ = \frac{1}{\epsilon} 12Q_1 + \frac{1}{\epsilon} E_1^a \end{aligned}$$

# Evanescent Operators

Complication: evanescent counterterms have different color structures than corresponding two-loop subdiagrams

$$(\psi C P_R \sigma_{\mu\nu} i\tau_2 \psi)(\psi C P_R \sigma_{\mu\nu} i\tau_2 \psi)(\psi C P_R i\tau_2 \tau_+ \psi) T^{SSS} \equiv 12Q_1 + E_1^a$$

$T^{AAS}$  color structure  


Non-local divergences don't cancel diagram-by-diagram

Solution: change of evanescent basis

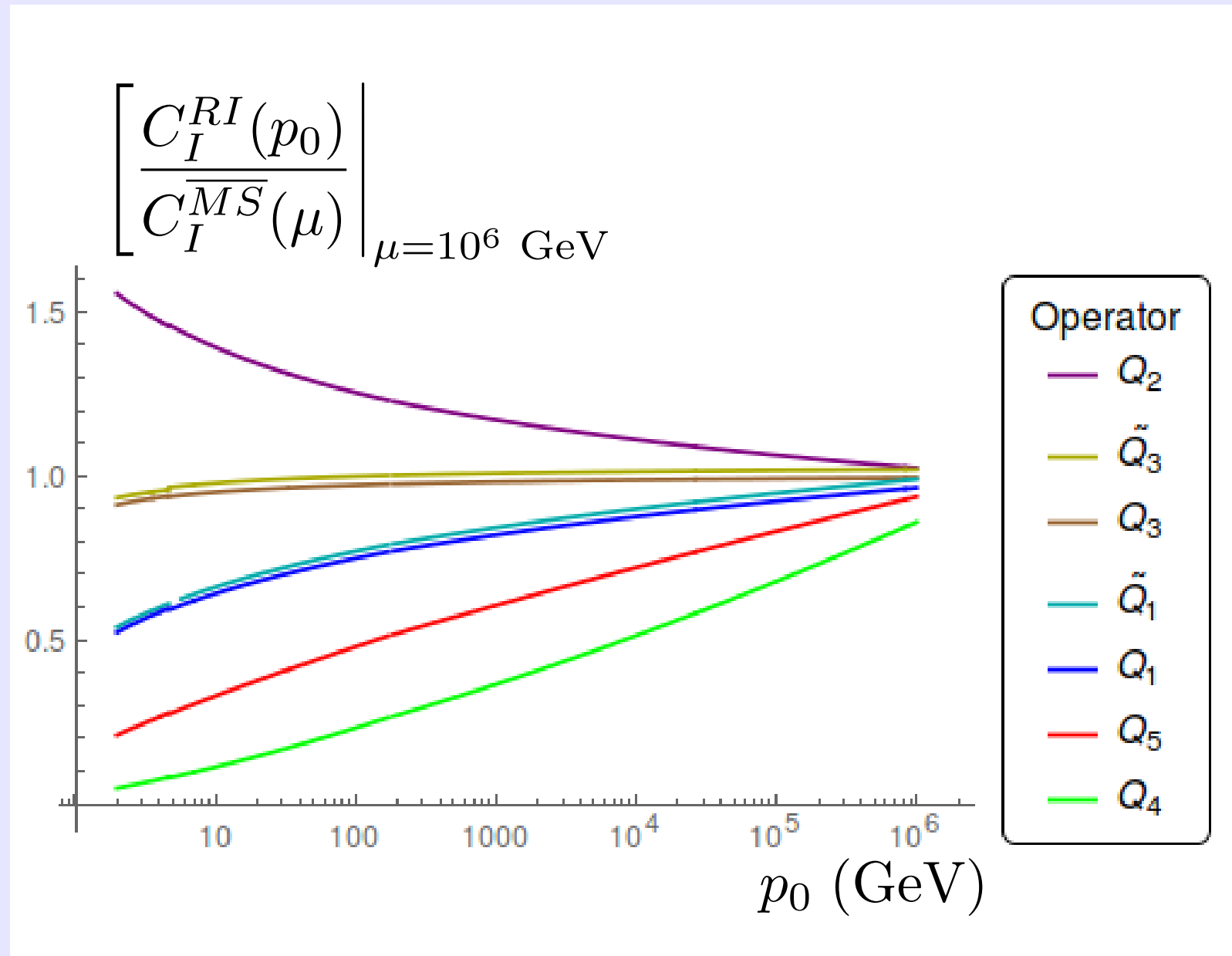
$$\begin{aligned}
 & (\psi C P_R \sigma_{\mu\nu} i\tau_2 \psi)(\psi C P_R \sigma_{\mu\nu} i\tau_2 \psi)(\psi C P_R i\tau_2 \tau_+ \psi) T^{SSS} \\
 & \equiv 8(\psi^\alpha [C P_R]^{\alpha\delta} i\tau_2 \psi^\beta)(\psi^\gamma [C P_R]^{\gamma\beta} i\tau_2 \psi^\delta)(\psi^\eta [C P_R]^{\eta\zeta} i\tau_2 \tau_+ \psi^\zeta) T^{SSS} + E_1^{a'} \\
 & = 12Q_1 - 6(Q_1 - \tilde{Q}_1) + E_1^{a'}
 \end{aligned}$$

# Results

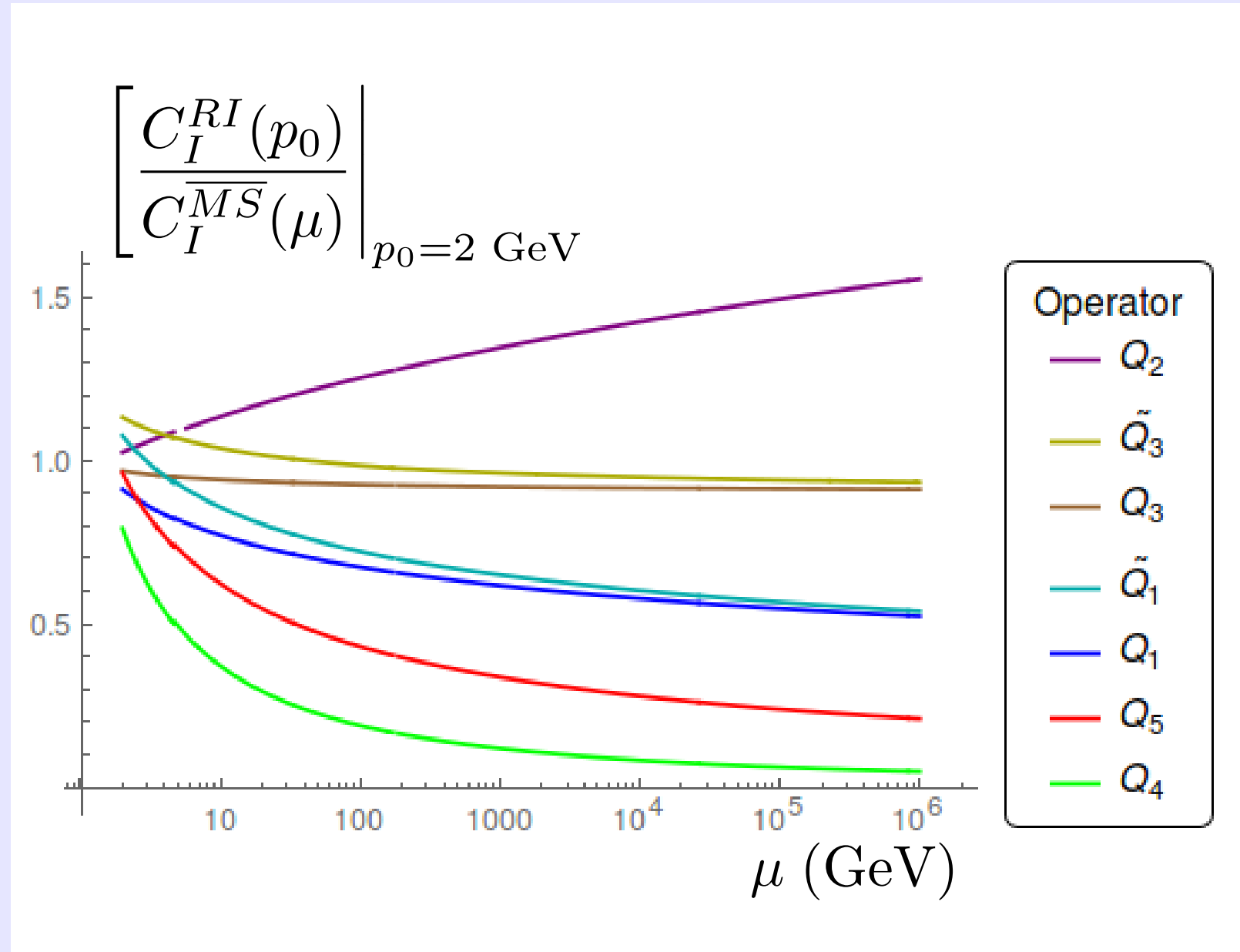
Operator	$\gamma^{(0)}$	$\gamma^{(1)}$	$r^{(0)}$
$Q_1$	4	$335/3 - 34N_f/9$	$101/30 + 8/15 \ln 2$
$Q_2$	-4	$91/3 - 26N_f/9$	$-31/6 + 88/15 \ln 2$
$Q_3$	0	$64 - 10N_f/3$	$-9/10 + 16/5 \ln 2$
$Q_4$	24	$229 - 46N_f/3$	$177/10 - 64/5 \ln 2$
$Q_5$	12	$238 - 14N_f$	$49/10 - 24/5 \ln 2$
$\tilde{Q}_1$	4	$797/3 - 118N_f/9$	$-109/30 + 8/15 \ln 2$
$\tilde{Q}_3$	0	$218 - 38N_f/3$	$-79/10 + 16/5 \ln 2$



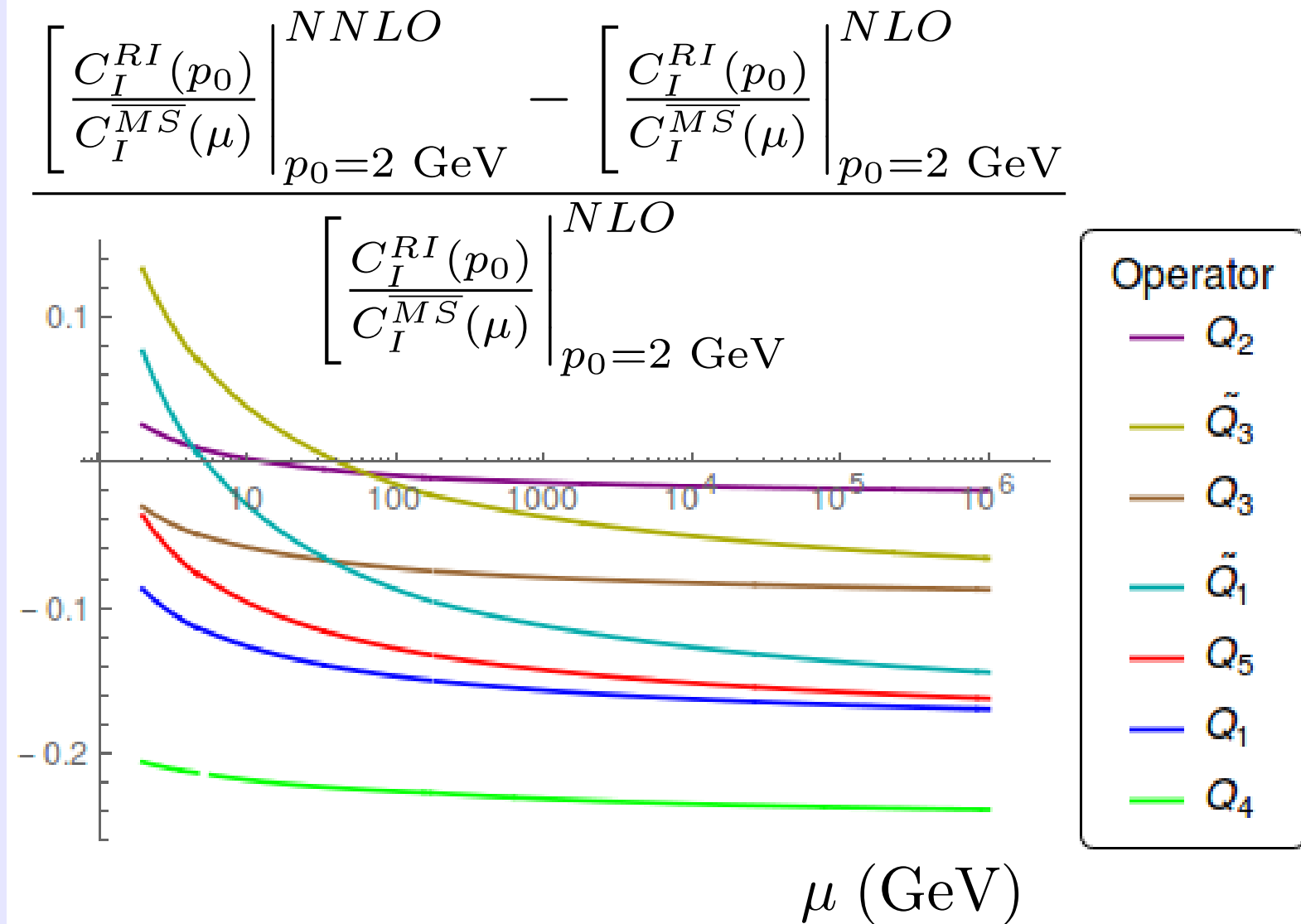
# Perturbative Renormalization Results – Theory View



# Perturbative Renormalization Results – Experimental View



# Relative Size of NNLO Corrections

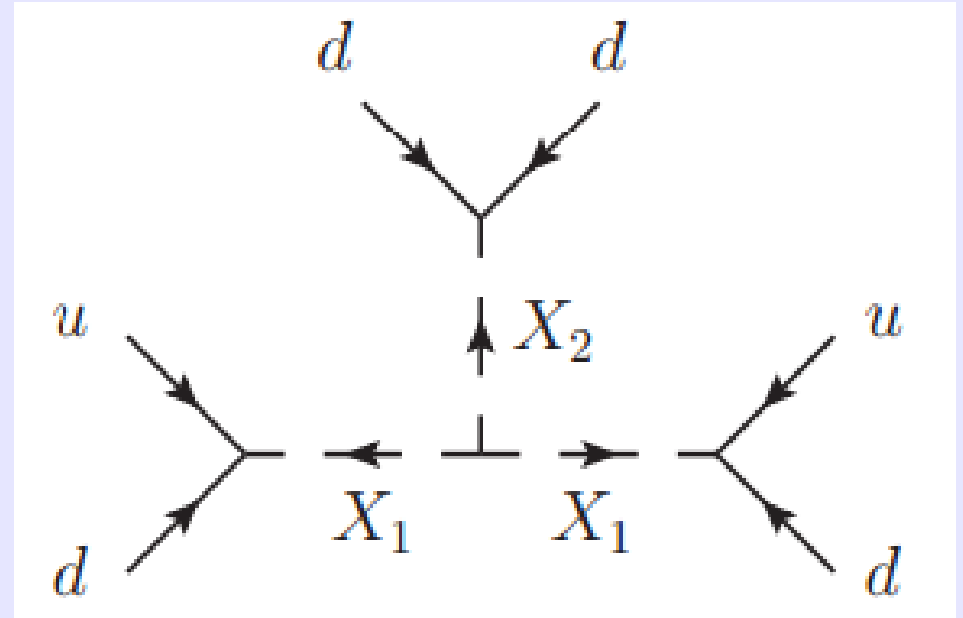




# Example BSM Application

Tree-level BSM matching for a simplified model

Arnold, Fornal, and Wise (2013)



Assuming all dimensionless couplings =1 at  $\mu = M$

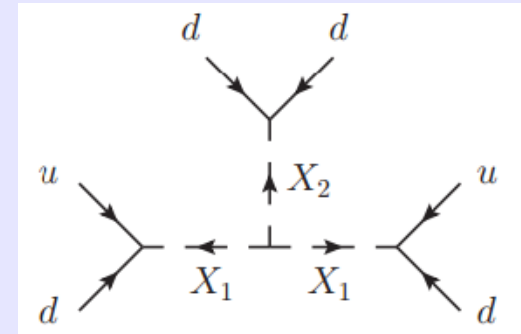
$$\mathcal{H}_{\text{eff}} = \frac{1}{16M^5} \left[ Q_4^{\overline{MS}}(M) + \frac{3}{5} \tilde{Q}_1^{\overline{MS}}(M) \right]$$

$$= \frac{1}{16M^5} \left[ U_4(M, p_0) Q_4^{RI}(p_0) + \frac{3}{5} \tilde{U}_1(M, p_0) \tilde{Q}_1^{RI}(p_0) \right]$$

# Example BSM Application

LO error:  $O\left(\frac{\alpha_s(p_0)}{\pi} \ln(M^2/p_0^2)\right) \sim 2.46$

Uncontrolled perturbative errors  
without NLO 1-loop running



Arnold, Fornal, and Wise (2013)

NLO error:  $O\left(\frac{\alpha_s(p_0)}{\pi}\right) \sim 0.09$        $O\left(\frac{\alpha_s(p_0)^2}{\pi^2} \ln(M^2/p_0^2)\right) \sim 0.23$

Assuming  $M = 500$  TeV

$$1/\tau_{n\bar{n}} = \left( \underbrace{6.74}_{\text{LO}} \quad \underbrace{-2.44}_{\text{NLO Running}} \quad \underbrace{\pm 0.16}_{\text{Lattice Statistical}} \right) \times 10^{-34} \text{ GeV}$$

Super K

$$1/\tau_{n\bar{n}} < 2 \times 10^{-33} \text{ GeV} \quad M > \left( \underbrace{402}_{\text{LO}} \quad \underbrace{-34}_{\text{NLO}} \quad \underbrace{\pm 3}_{\text{Lattice Statistical}} \right) \text{ TeV}$$

Preliminary lattice QCD matrix elements: [Buchoff, Schroeder, and Wasem \(2012\)](#)

See [Sergey Syritsyn Lattice 2015](#) talk and

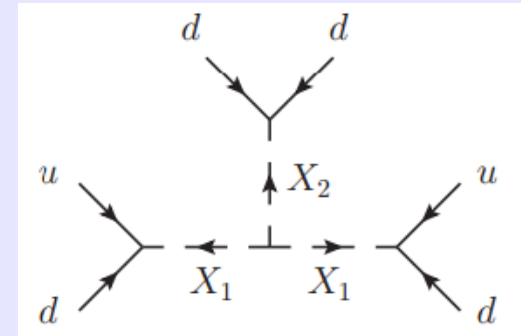
[Michael Buchoff INT Workshop](#) talk next week

for updated lattice results by [Buchoff, Schroeder, Syritsyn, and Wasem](#)

# Example BSM Application

NNLO error:  $O\left(\frac{\alpha_s(p_0)^3}{\pi^3} \ln(M^2/p_0^2)\right) \sim 0.02$

$O\left(\frac{\alpha_s(p_0)^2}{\pi^2}\right) \sim 0.01$       NLO Convergence Rate  $\sim 0.14^2 = 0.02$



Arnold, Fornal, and Wise (2013)

$$1/\tau_{n\bar{n}} = \left( \underbrace{6.74}_{\text{LO}} \quad \underbrace{-2.44}_{\text{NLO Running}} \quad \underbrace{+0.33}_{\text{NNLO Matching}} \quad \underbrace{-0.94}_{\text{NNLO Running}} \quad \underbrace{\pm 0.16}_{\text{Lattice Statistical}} \right) \times 10^{-34} \text{ GeV}$$

$$= (3.68 \pm 0.16) \times 10^{-34} \text{ GeV}$$

$$M > \left( \underbrace{402}_{\text{LO}} \quad \underbrace{-34}_{\text{NLO}} \quad \underbrace{+5}_{\text{NNLO matching}} \quad \underbrace{-16}_{\text{NNLO running}} \quad \underbrace{\pm 3}_{\text{Lattice Statistical}} \right) \text{ TeV}$$

$$M > 357 \pm 3 \text{ TeV}$$

Bound reaches  $M > 500 \text{ TeV}$  with factor of 5.4 tighter lifetime bound

Bounds with fully quantified uncertainties can be made once lattice systematics are included [Buchhoff, Schroeder, Syritsyn and Wasem](#)

# Summary

Color and flavor structure matter when determining bounds for particular BSM models:

At  $M = 500$  TeV, perturbative renormalization multiplies lattice results by (1.54) for  $Q_2$  vs (0.05) for  $Q_4$

2-loop running and 1-loop matching correct 1-loop running result of Caswell *et al.* by  $\lesssim 26\%$

Assuming same rate of convergence, remaining perturbative uncertainty  $\lesssim 7\%$



# Summary

