Beyond the Standard Model: The Low & High Energy Interface

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AMHERST CENTER FOR FUNDAMENTAL INTERACTIONS Physics at the interface: Energy, Intensity, and Cosmic frontiers University of Massachusetts Amherst

http://www.physics.umass.edu/acfi/

INT Workshop, Seattle, September 2015



Goals for this talk

- Set the context for the workshop
- Highlight (some) opportunities for low energy BSM discoveries
- Illustrate complementarity with BSM searches at the high energy frontier
- Underscore the need for on-going developments in nuclear and hadronic structure

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- Highlight (some) opportunities for low energy BSM discoveries
- Illustrate complementarity with BSM searches at the high energy frontier
- Underscore the need for on-going developments in nuclear and hadronic structure Challenges

Outline

- I. Fundamental symmetries: the BSM context
- II. LNV: $0\nu\beta\beta$ decay & the LHC
- III. CPV: EDMs, the LHC, & Baryon Asymmetry
- IV. Precision Tests (if time)
- V. Outlook

I. The BSM Context

Questions for Fundamental Physics*

- What is the origin of matter (luminous & dark) ?
- Why are neutrino masses so small ?
- Are fundamental interactions "natural"?



BSM Physics: Where Does it Live ?

BSM Physics: Where Does it Live ?



BSM Physics: Where Does it Live ?



Questions for Fundamental Physics*

- What is the origin of matter (luminous & dark) ?
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Discovering answers requires studies at three frontiers: energy, intensity, & cosmic.

*Partial List

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This talk

Low-Energy / High-Energy Interplay



Low-Energy / High-Energy Interplay



The Nuclear Physics Program

Targeted program of experiments & theory

Ature of the neutrino & search for lepton number violation

♦ Yet unseen T-violation (CP-violation)

Other key ingredients of the "New Standard Model"

Four Components **

EDM searches: BSM CPV, Origin of Matter	<i>0vββ decay searches:</i> Nature of neutrino, Lepton number violation, Origin of Matter
Electron & muon prop's &	Radioactive decays & other
interactions:	tests
SM Precision Tests, BSM	SM Precision Tests, BSM
"diagnostic" probes	"diagnostic" probes

Four Components

This talk

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Four Components

This talk

EDM searches: BSM CPV, Origin of Matter	<i>0vββ decay searches:</i> Nature of neutrino, Lepton number violation, Origin of Matter
<i>Electron & muon prop's & interactions:</i> <i>SM Precision Tests, BSM "diagnostic" probes</i>	Radioactive decays & other tests SM Precision Tests, BSM "diagnostic" probes

If time

II. LNV: $0\nu\beta\beta$ – Decay & the LHC

0vββ-Decay: LNV? Mass Term?

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.} \qquad \mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Dirac Majorana

Ονββ-Decay: LNV? Mass Term?

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana



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Majorana

Impact of observation

- Total lepton number not conserved at classical level
- New mass scale in nature, Λ
- Key ingredient for standard baryogenesis via leptogenesis



Ton Scale Experiments

$0\nu\beta\beta$ decay Experiments - Efforts Underway



Thanks: J. Wilkerson

Why Might A "Ton-Scale" Exp't See It?



Why Might A "Ton-Scale" Exp't See It?



0vββ-Decay: LNV? Mass Term?

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

$$\mathcal{C}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

"Standard" Mechanism

- Light Majorana mass generated at the conventional see-saw scale: Λ ~ 10¹² – 10¹⁵ GeV
- 3 light Majorana neutrinos mediate decay process



Why Might A "Ton-Scale" Exp't See It?



26

Why Might A "Ton-Scale" Exp't See It?

Three active light neutrinos



Interpreting the Result



Interpreting a Positive Result



Interpreting a Null Result











Neutrino Mass Hierarchy



Expected significance for rejecting wrong hierarchy hypothesis

Blennow et al, 1311.1822

Interpreting a Positive Result


Why Might A "Ton-Scale" Exp't See It?



Two parameters: Effective coupling & effective heavy particle mass

0vββ-Decay: LNV? Mass Term?

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

TeV LNV Mechanism

- Majorana mass generated at the TeV scale
 - Low-scale see-saw
 - Radiative m_v
- *m_{MIN}* << 0.01 eV but 0vββ-signal accessible with tonne-scale exp'ts due to heavy Majorana particle exchange



Ονββ-Decay: TeV Scale LNV

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana



TeV Scale LNV

Can it be discovered with combination of $0\nu\beta\beta$ & LHC searches ?

Simplified models

0vββ-Decay: TeV Scale LNV

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana



TeV Scale LNV

Comparing $\partial \nu \beta \beta$ & LHC sensitivities:

- LHC backgrounds
- Running effective op's to low energy
- Matching onto hadronic d.o.f.
- Long range NME contributions

Ονββ-Decay: TeV Scale LNV

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana





TeV Scale LNV

Effective operators:

$$\begin{split} \mathcal{L}_{\mathrm{LNV}}^{\mathrm{eff}} &= \frac{C_1}{\Lambda^5} \mathcal{O}_1 + \mathrm{h.c.} \\ \mathcal{O}_1 &= \bar{Q} \tau^+ d \bar{Q} \tau^+ d \bar{L} L^C \end{split}$$

$$g_{\rm eff} = C_1(\Lambda)^{1/4}$$

41

0vββ-Decay: TeV Scale LNV

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

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Ονββ-Decay: TeV Scale LNV

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$0v\beta\beta$ -Decay: TeV Scale LNV & m_v

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.} \qquad \mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Dirac Majorana

Implications for m_{v} :





Schecter-Valle: non-vanishing Majorana mass at (multi) loop level Simplified model: possible (larger) one loop Majorana mass 46

$0v\beta\beta$ -Decay: TeV Scale LNV & m_v

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.} \qquad \mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Dirac Majorana

Implications for m_{v} :



A hypothetical scenario

Low-Energy / High-Energy Interplay

TeV LNV



0vββ / LHC Interplay: Matrix Elements

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$









0vββ / LHC Interplay: Matrix Elements

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana



T. Peng, MRM, P. Winslow 1508.04444

III. CPV: EDMs, LHC, & Y_B

System	Limit (e cm)*	SM CKM CPV	BSM CPV
¹⁹⁹ Hg	3.1 x 10 ⁻²⁹	10 ⁻³³	10 ⁻²⁹
ThO	8.7 x 10 ⁻²⁹ **	10 ⁻³⁸	10 ⁻²⁸
n	3.3 x 10 ⁻²⁶	10 ⁻³¹	10 ⁻²⁶

* 95% CL ** e⁻ equivalent

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Mike Pendlebury: 1936-2015



The Guardian 9/23/15

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Mass Scale Sensitivity

System	Limit (e cm)*	SM CKM CPV	BSM CPV
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ThO	8.7 x 10 ⁻²⁹ **	10 ⁻³⁸	10 ⁻²⁸
n	3.3 x 10 ⁻²⁶	10 ⁻³¹	10 ⁻²⁶

* 95% CL ** e⁻ equivalent



Not shown: muon

Complementarity: Three Illustrations

- CPV in an extended scalar sector (2HDM): "Higgs portal CPV"
- Weak scale baryogenesis (MSSM)
- Model-independent

What is the CP Nature of the Higgs Boson ?

• Interesting possibilities if part of an extended scalar sector

Higgs Portal CPV

Inoue, R-M, Zhang: 1403.4257

CPV & 2HDM: Type I & II

 $\lambda_{6,7} = 0$ for simplicity

$$V = \frac{\lambda_1}{2} (\phi_1^{\dagger} \phi_1)^2 + \frac{\lambda_2}{2} (\phi_2^{\dagger} \phi_2)^2 + \lambda_3 (\phi_1^{\dagger} \phi_1) (\phi_2^{\dagger} \phi_2) + \lambda_4 (\phi_1^{\dagger} \phi_2) (\phi_2^{\dagger} \phi_1) + \frac{1}{2} \left[\lambda_5 (\phi_1^{\dagger} \phi_2)^2 + \text{h.c.} \right] \\ - \frac{1}{2} \left\{ m_{11}^2 (\phi_1^{\dagger} \phi_1) + \left[m_{12}^2 (\phi_1^{\dagger} \phi_2) + \text{h.c.} \right] + m_{22}^2 (\phi_2^{\dagger} \phi_2) \right\}.$$





Future Reach: Higgs Portal CPV

CPV & 2HDM: Type II illustration

 $\lambda_{6.7} = 0$ for simplicity



P	re	se	nt

 $sin \alpha_b$: CPV scalar mixing

Future:	Future:	
d _n x 0.1	<i>d_n</i> x 0.01	
d _A (Hg) x 0.1	<i>d_A(Hg)</i> x 0.1	
d _{ThO} x 0.1	d _{ThO} x 0.1	
d _A (Ra)	d _A (Ra)	

Inoue, R-M, Zhang: 1403.4257

Higgs Portal CPV: EDMs & LHC

CPV & 2HDM: Type II illustration

 $\lambda_{6.7} = 0$ for simplicity



Ρ	re:	se	nt	

 $sin \alpha_b$: CPV scalar mixing

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Inoue, R-M, Zhang: 1403.4257

Higgs Portal CPV: EDMs & LHC

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Inoue, R-M, Zhang: 1403.4257

68

Higgs Portal CPV: EDMs & LHC

CPV & 2HDM: Type II illustration

 $\lambda_{6.7} = 0$ for simplicity



Present

 $sin \alpha_b$: CPV scalar mixing

Future:	Future:
d _n x 0.1	<i>d_n</i> x 0.01
d _A (Hg) x 0.1	<i>d_A(Hg)</i> x 0.7
d _{ThO} x 0.1	d _{ThO} x 0.1
d _A (Ra)	d _A (Ra)

Inoue, R-M, Zhang: 1403.4257

Low-Energy / High-Energy Interplay

Higgs Portal CPV



Had & Nuc Uncertainties

CPV & 2HDM: Type II illustration

$\lambda_{6,7} = 0$ for simplicity



Present

 $sin \alpha_b$: CPV scalar mixing

Had & Nuc Uncertainties

CPV & 2HDM: Type II illustration

$\lambda_{6,7} = 0$ for simplicity



Present

Challenge #2

 $sin \alpha_b$: CPV scalar mixing

Inoue, R-M, Zhang: 1403.4257

Was the baryon asymmetry produced during electroweak symmetry-breaking ?

- EDMs provide most powerful probe of CPV
- Phase transition → Separate talk (back up slides)

EDMs & EW Baryogenesis: MSSM



Heavy sfermions: LHC consistent & suppress 1-loop EDMs



Sub-TeV EW-inos: LHC & EWB - viable but non-universal phases





EDMs & EW Baryogenesis: MSSM



Heavy sfermions: LHC consistent & suppress 1-loop EDMs



Sub-TeV EW-inos: LHC & EWB - viable but non-universal phases



EDMs & EW Baryogenesis: MSSM



Heavy sfermions: LHC consistent & suppress 1-loop EDMs



Sub-TeV EW-inos: LHC & EWB - viable but non-universal phases



Low-Energy / High-Energy Interplay

EWB for Compressed SUSY





Model Independent: Effective Operators

$\delta_{\!f}$	fermion EDM	(3)
$oldsymbol{\widetilde{\delta}}_q$	quark CEDM	(2)
$C_{\widetilde{G}}$	3 gluon	(1)
C _{quqd}	non-leptonic	(2)
C _{lequ, ledq}	semi-leptonic	(3)
$m{C}_{arphi$ ud	induced 4f	(1)

12 total + $\overline{\theta}$

light flavors only (e,u,d)

Paramagnetic Systems: Two Sources



TI, YbF, ThO...
Paramagnetic Systems: Two Sources



Paramagnetic Systems: Two Sources



74

Paramagnetic Systems: Two Sources



IV. Precision Tests







True deviation from SM ?













Curtin et al, '14



Curtin et al, '14



Curtin et al, '14



The Hunt for a Dark Z: PVES



Dark Z: Mechanism

$$\mathcal{L} \subset -\frac{1}{4} \, \hat{B}_{\mu\nu} \, \hat{B}^{\mu\nu} - \frac{1}{4} \, \hat{Z}_{D\mu\nu} \, \hat{Z}_D^{\mu\nu} + \frac{1}{2} \, \frac{\epsilon}{\cos\theta} \, \hat{Z}_{D\mu\nu} \, \hat{B}^{\mu\nu} + \frac{1}{2} \, m_{D,0}^2 \, \hat{Z}_D^\mu \, \hat{Z}_{D\mu\nu} \, \hat{Z}_D^\mu \, \hat{Z$$

 $V_0(H,S) = -\mu^2 |H|^2 + \lambda |H|^4 - \mu_S^2 |S|^2 + \lambda_S |S|^4 + \kappa |S|^2 |H|^2$

Dark Z: Mechanism

$$\mathcal{L} \subset -\frac{1}{4} \,\hat{B}_{\mu\nu} \,\hat{B}^{\mu\nu} - \frac{1}{4} \,\hat{Z}_{D\mu\nu} \,\hat{Z}_D^{\mu\nu} + \frac{1}{2} \,\frac{\epsilon}{\cos\theta} \,\hat{Z}_{D\mu\nu} \,\hat{B}^{\mu\nu} + \frac{1}{2} \,m_{D,0}^2 \,\hat{Z}_D^\mu \,\hat{Z}_{D\mu}$$

Kinetic Mixing

Mass Mixing

$$V_0(H,S) = -\mu^2 |H|^2 + \lambda |H|^4 - \mu_S^2 |S|^2 + \lambda_S |S|^4 + \kappa |S|^2 |H|^2$$

Higgs Mixing

Dark Z: Mechanism



V. Outlook

- Tests of fundamental symmetries & neutrino properties provide powerful windows into key open questions in fundamental physics
- There exists a rich interplay with BSM searches at the high energy frontier & both frontiers are essential
- Exciting opportunities for discovery and insight lie at the frontier interface
- Fully realizing them poses new challenges for hadronic & nuclear structure theory

Stay Tuned !

Back Up Slides



$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.} \qquad \mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Dirac Majorana

Our analysis:

- Include backgrounds
- Incorporate QCD running
- Include long-distance contributions to nuclear matrix elements

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.} \qquad \mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Dirac Majorana

Backgrounds:

- Charge flip
- Jet faking electron

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

$$\mathcal{C}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

Backgrounds:

- Charge flip
- Jet faking electron



e⁺ transfers most of p_T to conversion e⁻; Z / γ^* + jets \rightarrow apparent e⁻ e⁻ jj event

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

Backgrounds:

- Charge flip
- Jet faking electron



 e^+ transfers most of p_T to conversion e^- ; b's not tagged \rightarrow apparent $e^- e^-$ jj event 97

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.} \qquad \mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Dirac Majorana

Backgrounds: Bin in η and apply charge flip prob



98



$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.} \qquad \mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Dirac Majorana

Backgrounds: Cuts

- H_T
- MET
- M_{//}

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Dirac

Majorana

Backgrounds: Cuts

$\sigma({ m fb})$	Signal	Backgrounds									$\frac{\mathrm{s}}{\sqrt{\mathrm{s}+\mathrm{B}}}\left(\sqrt{\mathrm{fb}}\right)$
		Diboson			Charge Flip		Jet Fake				
		W^-W^-+2j	W^-Z+2j	ZZ+2j	$Z/\gamma^*\!+\!2\mathrm{j}$	$t\overline{t}$	$t\overline{t}$	\overline{t} +3j	W^-+3j	4j	
Before Cuts	0.142	0.541	6.682	0.628	903.16	68.2	6.7	0.45	15.09	362.352	0.0038
Signal Selection	0.091	0.358	4.66	0.435	721.7	28.9	2.37	0.22	11.73	72.03	0.0031
$H_T(\text{jets}) > 650 \text{ GeV}$	0.054	0.04	0.187	0.015	5.6	0.266	0.025	0.0003	0.102	0.027	0.0213
$m_{\ell_1\ell_2} > 130 \text{ GeV}$	0.039	0.029	0.105	0.008	0.163	0.127	0.024	$3x10^{-4}$	0.101	0.027	0.0493
$E_T < 40 \text{ GeV}$	0.036	0.005	0.036	0.007	0.126	0.014	0.005	$3x10^{-5}$	0.03	0.017	0.0684
$(\eta_{j_{1,2}} - \eta_{\ell_{1,2}})_{max} < 2.2$	0.033	0.003	0.022	0.005	0.093	0.009	0.004	$2x10^{-5}$	0.019	0.011	0.0738

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.} \qquad \mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Dirac Majorana

Low energy: N







$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

Low energy: QCD Running

$$\mathcal{O}_{1} = (\bar{u}_{L}d_{R})(\bar{u}_{L}d_{R})(\bar{e}_{L}e_{R}^{c}),$$

$$\mathcal{O}_{2} = (\bar{u}_{L}\sigma^{\mu\nu}d_{R})(\bar{u}_{L}\sigma_{\mu\nu}d_{R})(\bar{e}_{L}e_{R}^{c}),$$

$$\mathcal{O}_{3} = (\bar{u}_{L}t^{a}d_{R})(\bar{u}_{L}t^{a}d_{R})(\bar{e}_{L}e_{R}^{c}),$$

$$\mathcal{O}_{4} = (\bar{u}_{L}t^{a}\sigma^{\mu\nu}d_{R})(\bar{u}_{L}t^{a}\sigma_{\mu\nu}d_{R})(\bar{e}_{L}e_{R}^{c}).$$

$$\gamma^{ij} = -\frac{\alpha_s}{2\pi} \begin{pmatrix} 8 & 0 & 0 & 1\\ 0 & -8/3 & 48 & 0\\ 0 & 2/9 & -1 & 5/12\\ 32/3 & 0 & 20 & 19/3 \end{pmatrix}$$

$$\mathcal{L}_{\text{eff}} = \sum_{j} \frac{C_j(\mu)}{\Lambda^5} \mathcal{O}_j(\mu) + h.c.,$$

$$\mu \frac{d}{d\mu} C = \gamma^T C$$

104

$$\mathcal{L}_{\text{mass}} = y \bar{L} \tilde{H} \nu_R + \text{h.c.}$$

Dirac

$$\mathcal{L}_{\text{mass}} = \frac{y}{\Lambda} \bar{L}^c H H^T L + \text{h.c.}$$

Majorana

Low energy: QCD Running

 $\begin{aligned} \mathcal{O}_1 &= (\bar{u}_L d_R) (\bar{u}_L d_R) (\bar{e}_L e_R^c), \\ \mathcal{O}_2 &= (\bar{u}_L \sigma^{\mu\nu} d_R) (\bar{u}_L \sigma_{\mu\nu} d_R) (\bar{e}_L e_R^c), \\ \mathcal{O}_3 &= (\bar{u}_L t^a d_R) (\bar{u}_L t^a d_R) (\bar{e}_L e_R^c), \\ \mathcal{O}_4 &= (\bar{u}_L t^a \sigma^{\mu\nu} d_R) (\bar{u}_L t^a \sigma_{\mu\nu} d_R) (\bar{e}_L e_R^c). \end{aligned}$

Assuming $C_k = 1$ at $\mu = 5$ GeV \rightarrow Effective DBD amplitude for O_1 substantially weaker for given LHC constraints





Low energy: Nuclear Matrix Elements: Long Range Effects



Exploit Chiral Symmetry & EFT ideas



Low energy: Nuclear Matrix Elements: Long Range Effects



Our work

Helo et al

Exploit Chiral Symmetry & EFT ideas

Why Might A "Ton-Scale" Exp't See It?


Why Might A "Ton-Scale" Exp't See It?



Lightest neutrino mass (eV) ightarrow



EW Phase Transition: St'd Model



Lattice: Endpoint

Lattice	Authors	$M_{\rm h}^C~({ m GeV})$
4D Isotropic	[76]	80 ± 7
4D Anisotropic	[74]	72.4 ± 1.7
3D Isotropic	[72]	72.3 ± 0.7
3D Isotropic	[70]	72.4 ± 0.9

S'td Model: 1st order EWPT requires light Higgs

EW Phase Transition: MSSM







112

EW Phase Transition: MSSM





Katz, Perelstein, R-M, Winslow 1509.02934 113

EW Phase Transition: Higgs Portal





- Renormalizable
- φ : singlet or charged under SU(2)_L x U(1)_Y
- Generic features of full theory (NMSSM, GUTS...)
- More robust vacuum stability
- Novel patterns of SSB

Precision Tests

Electron Scattering



16

Electron Scattering



Electron Scattering

Search for additional neutral weak force that is inaccessible to the Large Hadron Collider





MOLLER: PV ee

SoLID & EIC: PV eD