

07 OCT 2015

C. Lee, LANL

"CP-Violating flavor oscillations in EWB"

based on work in papers w/ V. Cirigliano & E. Passarun

1106.0747 w/ VC & S. Tulin

0912.3523

+ M.J. Ramsey-Musolf

This morning ~~at~~

- One of Sakharov :
- 1) B
 - 2) C, CP
 - 3) arrow of time \rightarrow or ~~th eq.~~ (or CPT)

almost ~~more~~ 50 yrs (1967).

1), 2) exist in SM but not enough ~~at~~.

3) could've existed in SM (EWPT) but $m_H = 125 \text{ GeV} \Rightarrow$ no 1st order PT.

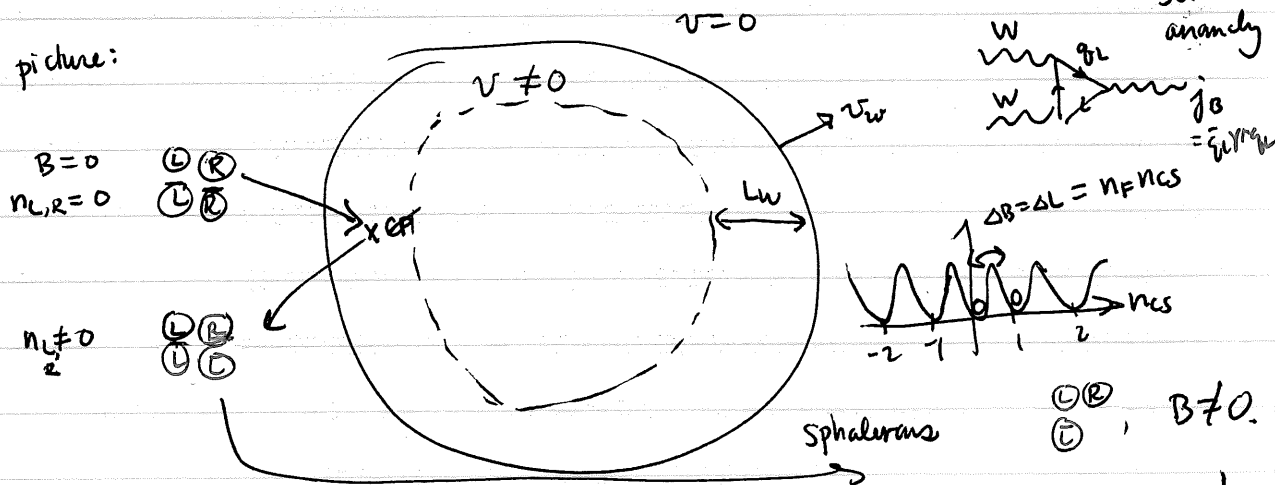
Need BSM.

I will focus on EWB, on improving the theory of quantum transport in which $\mathcal{O}(10-100)$ uncertainties existed in literature.

3 min

Construct my model of scalars exhibiting ~~some~~ the key qualitative features of EWB. (loosely based on steps like in MSSM \Rightarrow QCD.)

Big picture:



Lots of complicated physics

- 1) non-PT spheron rate
- 2) bubble wall dynamics
- 3) 1st vs. 2nd order PT
- * 4) transport theory:

$$\frac{\partial \rho_B}{\partial t} \sim \Gamma_{\text{sph}} \frac{\partial n_c}{\partial t}$$

↳ evolution eqs. in non-eq. bkgd.?

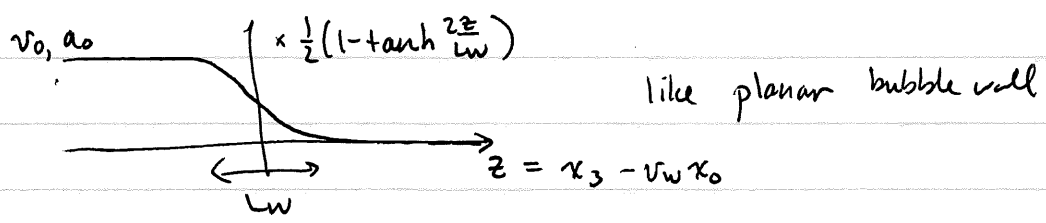
I want to give a flavor, intuitive picture for noneq. transport.

B min

Consider the model $\Phi = \begin{pmatrix} \phi_L \\ \phi_R \end{pmatrix}$

$$\mathcal{L} = \partial_\mu \Phi^\dagger \partial^\mu \Phi - \Phi^\dagger M^2(x) \Phi + \text{int}$$

where $M^2(x) = \begin{pmatrix} m_L^2 & v(x) e^{i\alpha(x)} \\ v(x) e^{-i\alpha(x)} & m_R^2 \end{pmatrix}$ (like $\tilde{t}_{L,R}$ interaction w/ Higgs vev)



$\mathcal{L}_{\text{int}} = -\frac{1}{2} \Lambda^2 \Phi^\dagger y \Phi$ $y = \begin{pmatrix} y_L & 0 \\ 0 & y_R \end{pmatrix}$ → defines "flavor" or "interaction" basis

↓
real scalar
acting as thermal bath

(try gauge boson)

B min

Mass basis: $\phi = \begin{pmatrix} \phi_1 \\ \phi_2 \end{pmatrix} = U^\dagger \Phi$

$$U^\dagger = \begin{pmatrix} \cos\theta & -\sin\theta e^{-i\alpha} \\ \sin\theta e^{i\alpha} & \cos\theta \end{pmatrix}(x)$$

$$\tan 2\theta = \frac{2v(x)}{m_L^2 - m_R^2}$$

$$m^2 = U^\dagger M^2 U = \begin{pmatrix} m_1^2 & 0 \\ 0 & m_2^2 \end{pmatrix} \quad m_{1,2}^2 = \frac{1}{2} (m_L^2 + m_R^2) \pm \frac{1}{2} \sqrt{(m_L^2 - m_R^2)^2 + 4v^2}$$

$$\rightarrow -\frac{1}{2} A^2 \phi^\dagger \Upsilon \phi \quad \Upsilon = U^\dagger \gamma U = \begin{pmatrix} \gamma_{11} & \gamma_{12} \\ \gamma_{21} & \gamma_{22} \end{pmatrix}$$

$$\mathcal{L} = \partial_\mu \phi^\dagger \partial^\mu \phi - \phi^\dagger m^2 \phi + \text{int} \\ - \phi^\dagger \Sigma^\mu \partial_\mu \phi + \partial_\mu \phi^\dagger \Sigma^\mu \phi - \phi^\dagger \Sigma^\mu \Sigma_\mu \phi$$

$$\Sigma_\mu = U^\dagger \partial_\mu U \\ = \begin{pmatrix} 0 & -e^{ia} \\ e^{ia} & 0 \end{pmatrix} \partial_\mu \theta + \begin{pmatrix} i \sin 2\theta & \frac{1}{2} \sin 2\theta e^{-ia} \\ \frac{1}{2} \sin 2\theta e^{ia} & -i \sin 2\theta \end{pmatrix} \partial_\mu a.$$

17 min

Derive transport eq. for particle densities from correlation functions

$$\tilde{G} = \begin{pmatrix} G^t & -G^< \\ G^> & -G^{\bar{t}} \end{pmatrix} (x, y) = \begin{pmatrix} \langle T \phi(x) \phi^\dagger(y) \rangle & -\langle \phi^\dagger(y) \phi(x) \rangle \\ \langle \phi(x) \phi^\dagger(y) \rangle & -\langle \bar{T} \phi(x) \phi^\dagger(y) \rangle \end{pmatrix}$$

Schwinger-Dyson:

$$\tilde{G}(x, y) = \tilde{G}^{(0)}(x, y) + \int d^4 z d^4 w \tilde{G}^0(x, z) \tilde{\Pi}(z, w) \tilde{G}(w, y)$$

Eqs of motion for G^z :

$$[\partial_x^2 + m^2(x) + 2\Sigma_\mu(x) \partial_x^\mu + \Sigma_\mu \Sigma^\mu(x) + \partial_x^\mu \Sigma_\mu(x)] G^z(x, y) = -i \int d^4 z [\tilde{\Pi}(x, z) \tilde{G}(z, y)]^z$$

sim. for y

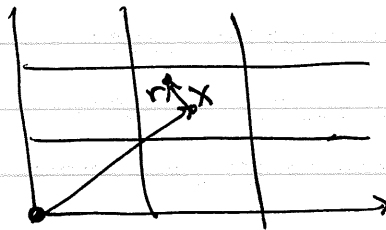
(Rest of this talk I am just going to unpack this eq.)

want distributions in momentum space:

21 min

Wigner transform

$$G(k; X) \equiv \int d^4 r e^{ik \cdot r} G(X + \frac{r}{2}, X - \frac{r}{2})$$



Solve for k distribution in each X box.

Wigner transforms obey complicated eq with infinite number of derivatives in k and X.

Need power counting scheme.

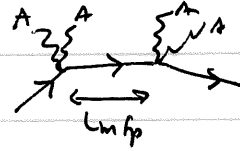
25 min

Expand in 3 ϵ 's:

$$E_w \sim \frac{L_{int}}{L_w}$$

Where $L_{int} \sim \frac{1}{|k|} \sim \frac{1}{T}$ muff

$$E_{coll} \sim \frac{L_{int}}{L_{mp}}$$



$$E_{coll} \sim y_{L,R} \text{ (pert. th.)}$$

$$E_{osc} \sim \frac{L_{int}}{L_{osc}}$$

$$L_{osc} \sim \frac{1}{\omega_1 - \omega_2} \sim \frac{|k|}{\Delta m^2}$$

27 min

$$O(\epsilon^0): \quad (k^2 - \bar{m}^2) G^z(k; x) = 0 \quad \bar{m}^2 = \frac{m_1^2 + m_2^2}{2}$$

"constraint eq"

$$\Rightarrow G^z(k; x) = 2\pi \delta(k^2 - \bar{m}^2) F^z(k; x)$$

↳ so far arbitrary

(actually $m_1^2 = m_1^2$
 $m_2^2 = m_2^2$)

$$F^> = \theta(k^0) [\mathbb{1} + f(k; x)] + \theta(-k^0) \bar{f}(k; x)$$

$$F^< = \theta(k^0) f(k; x) + \theta(-k^0) [\mathbb{1} + \bar{f}(-k; x)]$$

(commutation rels.)

On eq., $f = n_B = \frac{1}{e^{kx} - 1}$.

30 min

$O(\epsilon^1)$: "kinetic eqs." tell us how f 's evolve.

We get:

$$\left. \begin{aligned} (u \cdot \partial_x + \vec{F} \cdot \vec{\nabla}_k) f(k; x) &= -[i\omega_E + u \cdot \Sigma, f_m] + \mathcal{C}[f, \bar{f}] \\ (u \cdot \partial_x + \vec{F} \cdot \vec{\nabla}_k) \bar{f} &= +[i\omega_E - u \cdot \Sigma, \bar{f}_m] + \mathcal{C}[\bar{f}, f] \end{aligned} \right\}$$

where $u^\mu = (1, \frac{k}{\omega_k})$ $\vec{F} = -\vec{\nabla}_k \omega_E$

planar
→

$$v_{rel} \partial_z f = -[i\omega_E + v_{rel} \Sigma_z, f] + \mathcal{C}[f, \bar{f}]$$

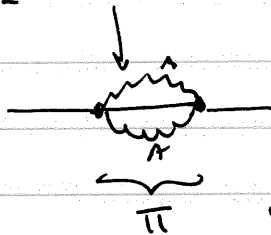
$$v_{rel} \partial_z \bar{f} = [i\omega_E - v_{rel} \Sigma_z, \bar{f}] + \mathcal{C}[\bar{f}, f]$$

where $v_{rel} = \frac{k_z}{\omega_k} - v_w$.

33 min

Collision term:

$$\mathcal{C}[f, \bar{f}] = \int_0^\infty \frac{dk^0}{2\pi} \frac{1}{2} \left(\{ \pi^>(k, x); G^<(k, x) \} - \{ \pi^<(k, x); G^>(k, x) \} \right)$$



$\phi A \rightarrow \phi A$
or $\phi \phi^\dagger \leftrightarrow A A$ scattering pair annihilation/creation.

not going to write more explicitly here.

35 min

Make sense of physical meaning of kin eqs:

$$f = \begin{pmatrix} f_{11} & f_{12} \\ f_{21} & f_{22} \end{pmatrix}$$

$$= p_0 \mathbb{1} + \vec{p} \cdot \vec{\sigma}$$

$$\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$\sigma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

$$\sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$\bar{f}$$

$$= \bar{p}_0 \mathbb{1} + \vec{\bar{p}} \cdot \vec{\sigma}$$

(Bloch decomposition)

$$\text{write } \frac{1}{v_{rel}} \mathcal{C} = -D_0 \mathbb{1} - \vec{\sigma} \cdot \vec{D}$$

kin eqs \Rightarrow

$$\partial_z p_0 = -D_0 [p, \bar{p}] \quad \partial_z \bar{p}_0 = -D_0 [\bar{p}, p]$$

$$\partial_z \vec{p} = (\vec{B}_0 + \vec{B}_E) \times \vec{p} - \vec{D} [p, \bar{p}]$$

$$\partial_z \vec{\bar{p}} = -(\vec{B}_0 - \vec{B}_E) \times \vec{\bar{p}} - \vec{D} [\bar{p}, p]$$

$$\vec{B}_0 = \begin{pmatrix} 0 \\ 0 \\ \frac{\omega_1 - \omega_2}{v_{rel}} \end{pmatrix}$$

$$\vec{B}_E = \begin{pmatrix} 2 \sin \alpha \theta' + \sin 2\theta \cos \alpha \alpha' \\ -2 \cos \alpha \theta' + \sin 2\theta \sin \alpha \alpha' \\ 2 \sin^2 \theta \alpha' \end{pmatrix}$$

like spin precession in mag. field

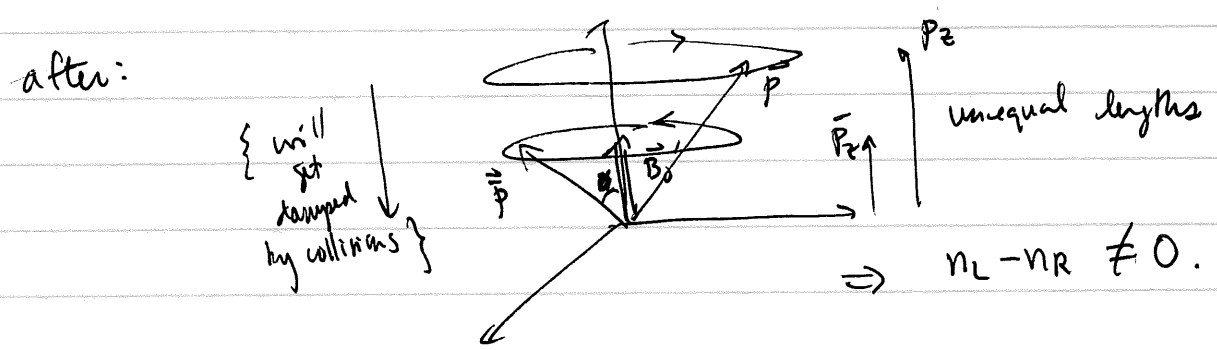
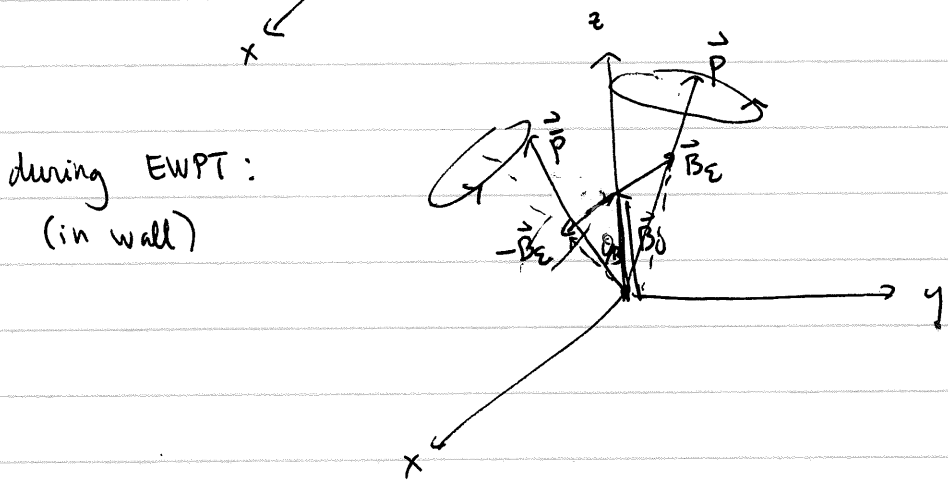
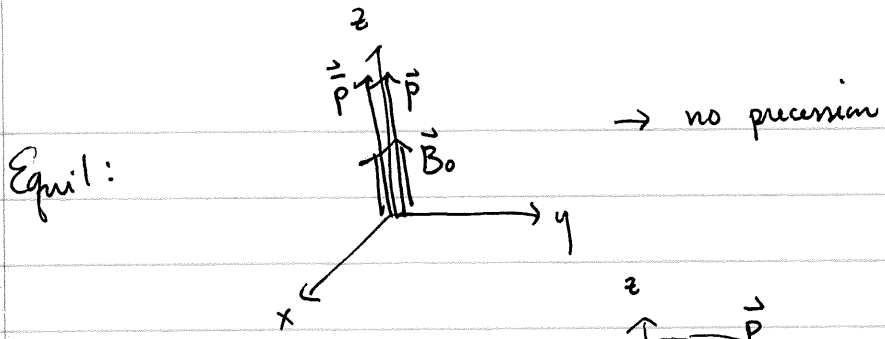
here

$p_0 \sim$ total occupation #

$p_3 \sim$ L-R asymmetry

$p_{12} \sim$ off-diagonal correlations in flavor space

42 min



largest effect when $\Theta_B \sim 1$ $L_w \sim L_{osc}$
"non-adiabatic"

previous lit assumed $\Theta_B \ll 1$, not always true,
mixed layer of flow-osc. source.

50 min.

extra:

$$\vec{B}_E \times \vec{p} \rightarrow \begin{pmatrix} B_x \\ B_y \\ 0 \end{pmatrix} \times \begin{pmatrix} 0 \\ 0 \\ n_0(\omega_1) - n_0(\omega_2) \end{pmatrix} + \begin{pmatrix} 0 \\ 0 \\ B_z \end{pmatrix} \times \begin{pmatrix} p_x \\ p_y \\ 0 \end{pmatrix}$$

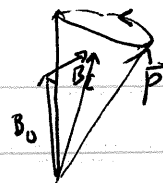
$\sim \epsilon \quad \sim \epsilon$

For $\Theta_B \sim 1 \rightarrow O(100)$ enhancement of $n_L - n_R$.

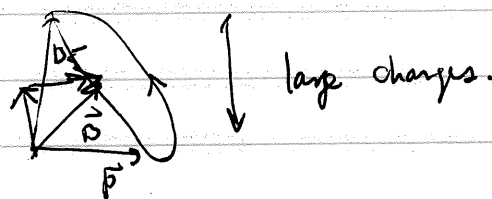
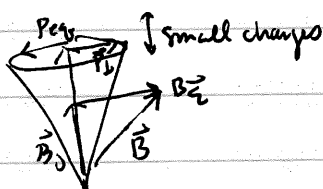
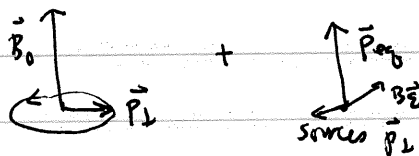
Under study:

- simplify w/ diffusion eq.
- realistic phenomen.

exact: $v_{\text{rel}} \partial_z \vec{p} = (\vec{B}_0 + \vec{B}_\Sigma) \times \vec{p}$



approx $\vec{B}_0 \times \begin{pmatrix} p_x \\ p_y \\ 0 \end{pmatrix} + \vec{B}_\Sigma \times \begin{pmatrix} 0 \\ v_B(\omega_1) - v_B(\omega_2) \\ 0 \end{pmatrix}$



$$\begin{matrix} (\vec{B}_0 + \vec{B}_\Sigma) \times \vec{p} \\ \downarrow \quad \downarrow \quad \downarrow \\ e \quad e \quad e \end{matrix}$$

Conclusions /

future:

- new possibly dominant source for \mathcal{L}^{P} in EWB
- simplification to diffusion eqs. (retaining flav. osc.)
- realistic models (incl. fermions)

55 min.