Resonances and unitarity in composite Higgs models

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References

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D.E., F. Mescia and B. Yencho, PRD 88, 055002 (2013) [1307.2400]
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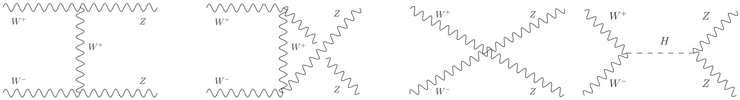
Closely related work:

R. Delgado, A. Dobado and F. Llanes, JHEP 1402 (2014) 121 [1311.5993] & 1502.04841, 1509.04725
R. Delgado, A. Dobado, M.J. Hererro and J.J, Sanz-Cillero, JHEP 1407 (2014) 149 [1404.2866]
A. Dobado, F-K. Guo, F. Llanes, 1508.03544

Motivation

We know that in the SM the Higgs boson unitarizes $W_L W_L$ scattering. Consider e.g. $W_L^+ W_L^- \to Z_L Z_L$



If any of these couplings are different from SM values, the careful balance necessary for perturbative unitarity is lost.

The first 3 diagrams are fixed by gauge invariance, but we can contemplate other Higgs-gauge boson couplings. For $s >> M_W^2$ the amplitude in the SM goes for $s \to \infty$ as

$$\frac{s}{v^2} \frac{M_H^2}{s - M_H^2} \sim \frac{M_H^2}{v^2}$$

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... but on dimensional grounds it *should* go as (cf. pion physics)

$$\frac{s}{v^2}\frac{s}{s - M_H^2} \sim \frac{s}{v^2}$$

This is what happens after *any modification* of the Higgs couplings and produces *non-unitary* amplitudes.

Adding *new effective operators* typically spoils unitarity too.

$$\mathcal{L}_{SM} \to \mathcal{L}_{SM} + \sum_i a_i \mathcal{O}_i \qquad \mathcal{O}_i \sim s^2$$

New physics may produce either type of modifications

What can the unitarity of longitudinal WW scattering tell us about anomalous couplings in EW sector?

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Parametrizing composite Higgs physics

A light "Higgs boson" with mass $M_H \sim 125$ GeV is coupled to the EW bosons according to (non-linear realization)

$$\mathcal{L}_{\rm eff} \supset -\frac{1}{2}{\rm Tr}\, W_{\mu\nu}W^{\mu\nu} - \frac{1}{4}{\rm Tr}\, B_{\mu\nu}B^{\mu\nu} + \mathcal{L}_{\rm GF} + \mathcal{L}_{\rm FP} + \sum_i \mathcal{L}_{\rm i}$$
$$+ \left[1 + 2a\left(\frac{h}{v}\right) + b\left(\frac{h}{v}\right)^2\right]\frac{v^2}{4}{\rm Tr}\, D_\mu U^\dagger D^\mu U - V(h)$$

$$U = \exp(i\,\omega\cdot\tau/v)$$
$$D_\mu U = \partial_\mu U + \frac{1}{2} i g W_\mu^i \tau^i U - \frac{1}{2} i g' B_\mu^i U \tau^3$$

and additional gauge-invariant operators are encoded in $\mathcal{L}_{\rm i}$.
Setting $a = b = 1$ (and $\mathcal{L}_{\rm i}=0$) reproduces the SM interactions

The $\mathcal{L}_i$ are a full set of C, P, and $SU(2)_L \times U(1)_Y$ gauge invariant, $d = 4$ operators that parameterize the *low-energy effects* of the *model-dependent high-energy EWSB sector* along with *a,b*. The two relevant *custodial-symmetry preserving* operators are

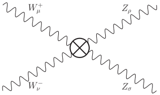
$$\mathcal{L}_4 = a_4 \left(\mathrm{Tr}\left[V_\mu V_\nu\right]\right)^2 \qquad \mathcal{L}_5 = a_5 \left(\mathrm{Tr}\left[V_\mu V^\mu\right]\right)^2 \qquad V_\mu = \left(D_\mu U\right) U^\dagger$$

The a_i could be *functions of* $\frac{h}{v}$

- For example:

		Heavy Higgs	*QCD-like technicolor*
a_4	$=$	0	$-2a_5$
a_5	$=$	$\frac{v^2}{8M_H^2}$	$\frac{N_{TC}}{96\pi^2}$

(up to logarithmic corrections)



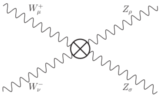
$$ig^4\left[a_4\left(g^{\mu\sigma}g^{\nu\rho} + g^{\mu\rho}g^{\nu\sigma}\right) + 2a_5 g^{\mu\nu}g^{\rho\sigma}\right]$$

There are solid indications that the "Higgs" couples to the W, Z *very similarly to the SM rules*

$$\mathcal{L}_{\rm eff} \simeq \mathcal{L}_{\rm SM} + a_4 \left(\mathrm{Tr}\left[V_\mu V_\nu\right]\right)^2 + a_5 \left(\mathrm{Tr}\left[V_\mu V^\mu\right]\right)^2$$

Then a_4 and a_5 represent anomalous 4-point couplings of the W bosons due to an extended EWSBS that however does not manifest with $O(p^2)$ couplings *noticeably different to the ones in the SM*. **Assume now that $a = b = 1$ exactly.**

These operators will lead to violations of perturbative unitarity at loop level ($\sim g^4$)



$$\sim \left(\frac{s}{v^2}\right)^2$$

Violations of unitarity are cured by the appearance of new particles or resonances

We can now use well-understood unitarization techniques to constrain these resonances and the effective couplings $\{a_i\}$

After the Higgs discovery

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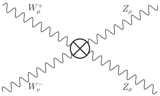

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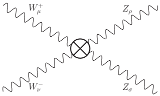

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To take home

- The Higgs unitarizes these amplitudes in SM (where $a = b = 1, \quad \{a_i\} = 0$)
- The theory is renormalizable *without* the $\{a_i\}$ if $a = b = 1$
- The $\{a_i\}$ will then be finite non-running parameters.

We would like to

- Determine how much room is left for the a_i
- Find possible additional resonances imposed by unitarity
- Should we have already seen any?
- To what extent an extended EWSBS is excluded?

Yes, there are new resonances even with relatively light masses
Their signal is very weak and not visible with limited statistics (what about a 2 TeV resonance?)
Searching for resonances is an efficient (albeit indirect) way of setting constrains on aTGC and aQGC $\Leftarrow \{a_i\}$

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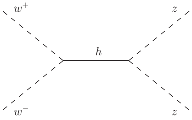
Equivalence Theorem

Most studies concerning unitarity at high energies are (understandbly) carried out using the ET

$$A(W_L^+ W_L^- \to Z_L Z_L) \to A(\omega^+ \omega^- \to \omega^0 \omega^0) + O(M_W/\sqrt{s})$$

For a light Higgs the region one needs to include tree-level Higgs exchange as well



Then one could make use of the well known chiral lagrangian techniques to derive the amplitudes and compare with experiment, including the Higss as an explicit resonance.

However for s not too large (which obviously is now an interesting region) the ET may be too crude an approximation.

We shall use as much as possible exact amplitudes

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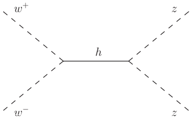

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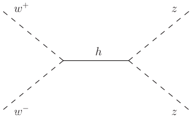

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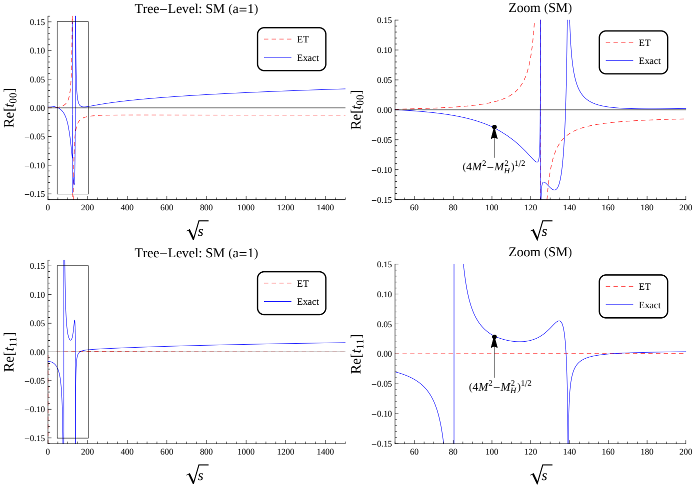
Tree-level amplitudes for $a = 1$

$a = 1$



Tree-level amplitudes for $a \neq 1$

$a > 1$

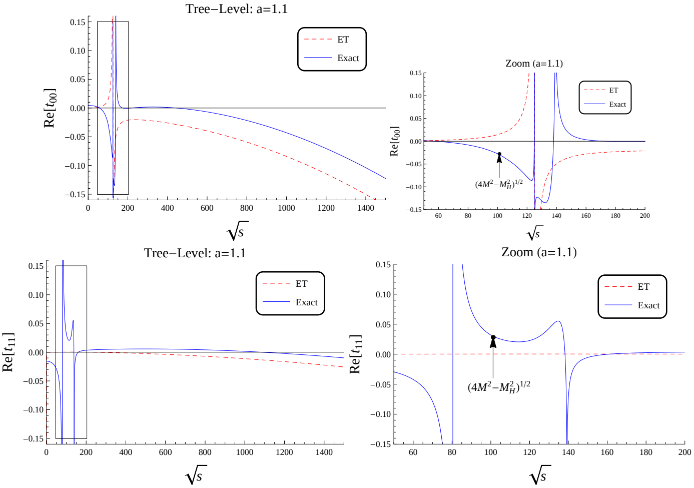


Figure : Tree-level amplitudes for $a = 1.1$

Tree-level amplitudes for $a \neq 1$

$a < 1$

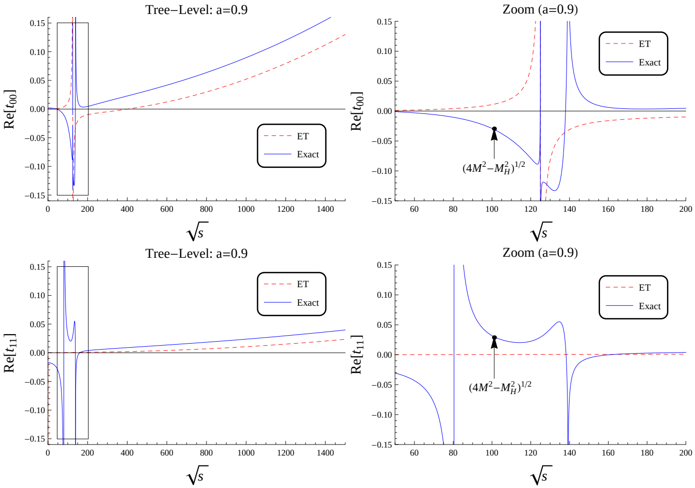
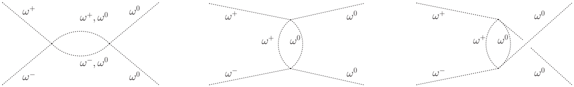


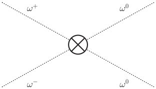
Figure : Tree-level amplitudes for $a = 0.9$

Goldstone boson amplitudes

At one-loop, the Goldstone boson scattering amplitudes *diverge*



The $\mathcal{L}_i$ serve as *counterterms* and the a_i are *scale dependent*



$$a_4 \quad \to \quad a_4(\mu) = a_4(\mu') - \frac{1}{16\pi^2}\frac{1}{12}\log(\frac{\mu^2}{\mu'^2})$$

$$a_5 \quad \to \quad a_5(\mu) = a_5(\mu') - \frac{1}{16\pi^2}\frac{1}{24}\log(\frac{\mu^2}{\mu'^2})$$

The inclusion of a dynamical Higgs brings in new diagrams and new divergences.
The above are the results for $a = 0$ and $b = 0$, *i.e. no Higgs*.

Partial wave amplitudes

We will assume $e = 0$ *(no e.m.)* i.e. the custodial limit $c_w = 1$.
The WW scattering amplitudes can then be deconstructed into amplitudes of *fixed isospin* T_I

$$\begin{aligned} T_0 &= 3A^{+-00} + A^{++++} \\ T_1 &= 2A^{+-+-} - 2A^{+-00} - A^{++++} \\ T_2 &= A^{++++} \end{aligned}$$

where $A^{+-00} = A(W_L^+ W_L^- \to Z_L Z_L)$ and all others may be expressed in terms of this amplitude through isospin and crossing symmetries

These can then be written in terms of *partial waves*

$$t_{I\,J}(s) = \frac{1}{64\pi} \int_{-1}^{1} d(\cos\theta) P_J(\cos\theta) T_I$$

which are constrained by unitarity at high energies to be $|t_{I\,J}| < 1$.
Most discussions based on unitarity are based on this simple constraint (*tree-level unitarity*)

Inverse Amplitude method

Partial wave unitarity **requires**

$$\mathrm{Im}\, t_{I\,J}(s) \;=\; \underbrace{\sigma(s)|t_{I\,J}(s)|^2}_{\substack{\text{Elastic}\\ WW \to WW}} \;+\; \underbrace{\sigma_H(s)|t_{H,I\,J}(s)|^2}_{\substack{\text{Inelastic}\\ WW \to hh}}$$

where σ and σ_H are phase space factors.

Given a perturbative expansion

$$t_{I\,J} \;\approx\; \underbrace{t^{(2)}_{I\,J}}_{\text{tree}} \;+\; \underbrace{t^{(4)}_{I\,J}}_{\substack{\text{one-loop}\\ +\, a_i \text{ terms}}} \;+\; \cdots$$

we can require unitarity to hold *exactly* by defining *(Note: non-coupled channels)*

$$t_{I\,J} \approx \frac{t^{(2)}_{I\,J}}{1 - t^{(4)}_{I\,J}/t^{(2)}_{I\,J}}$$

Several mild analyticity assumptions are implied.

New resonances

The unitarization of the amplitudes may result in the appearance of *new heavy resonances* associated with the high-energy theory

$$
\begin{aligned}
t_{00} &\rightarrow \text{Scalar isoscalar} \\
t_{11} &\rightarrow \text{Vector isovector} \\
t_{20} &\rightarrow \text{Scalar isotensor}
\end{aligned}
$$

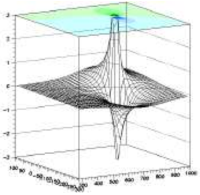
Will search for poles in $t_{I\,J}(s)$ up to $(4\pi v) \sim 3$ TeV (domain of applicability)

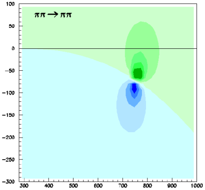
True resonances will have the phase shift pass through $+\pi/2$

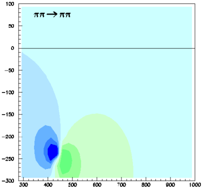
$$\delta_{IJ} = \tan^{-1}\left(\frac{\mathrm{Im}\; t_{IJ}}{\mathrm{Re}\; t_{IJ}}\right)$$

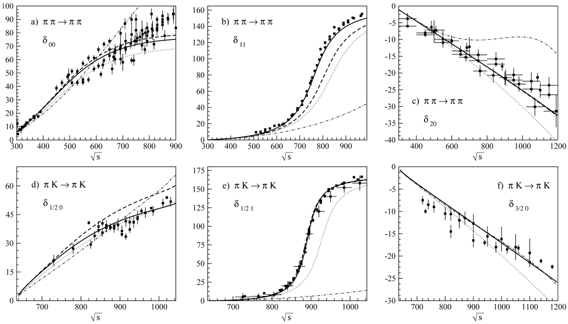
This method is known to work remarkably well in strong interactions: $\pi\pi$ scattering $\Rightarrow$ σ and ρ meson masses and widths

In hadronic physics









Truong '89, Truong,Dobado,Herrero, '90, Dobado, Pelaez,'93,'96

Integrating out a heavy higgs and recovering the resonance:

	Tree level		1-loop ($\mu = M_W$)		1-loop ($\mu = M_H$)	
MH	Mass	Width	Mass	Width	Mass	Width
500	446	28	442	27	455	30
600	521	46	511	43	538	50
700	599	70	580	63	627	81
800	680	105	647	90	726	133
900	767	157	712	121	840	201
1000	862	226	776	164	980	351
1100	968	344	838	210	1166	617
1200	1091	485	899	260	1460	1868

$$a_4 = \frac{-1}{16\pi^2}\frac{1}{12}\left(\Delta_\epsilon - \log\left(\frac{M_H^2}{\mu^2}\right) + \frac{17}{6}\right)$$

$$a_5 = \frac{M_W^2}{2g^2 M_H^2} - \frac{1}{16\pi^2}\frac{1}{24}\left(\Delta_\epsilon - \log\left(\frac{M_H^2}{\mu^2}\right) + \frac{79}{3} - \frac{27\pi}{2\sqrt{3}}\right)$$

(On-shell scheme)

See also recent work by Corbett, Eboli, Gonzalez-Garcia

Criticisms

Is this unitarization method unique?

No, it is not. Many methods exist: IAM, K-matrix approach, N/D expansions, Roy equations,...

While the quantitative results differ slightly, the gross picture does not change

For a very detailed discussion of different methods see 1502.0484 (Delgado et al.)

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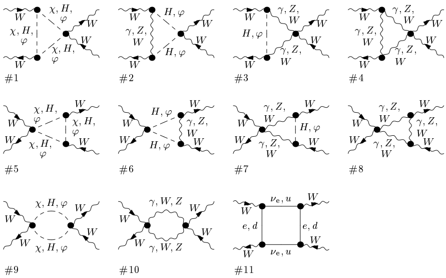
The calculation

Real problem: one-loop calculation *extremely difficult*



+ *more than 1000 other diagrams!*

Denner & Hahn (1998) [hep-ph/9711302]

Shortcut

We can take a shortcut:

$$t_{IJ}^{(4)} = \mathrm{Re}\, t_{IJ}^{(4)} + i\mathrm{Im}\, t_{IJ}^{(4)}$$

The **Optical Theorem** implies the *perturbative* relation

$$\underbrace{\mathrm{Im}\ t_{I\,J}^{(4)}(s)}_{\text{one-loop}} = \sigma(s)|\underbrace{t_{I\,J}^{(2)}}_{\text{tree}}(s)|^2 + \sigma_H(s)|t_{H,I\,J}^{(2)}(s)|^2$$

For *real part*, note that

$$\begin{aligned} \mathrm{Re}\, t_{IJ}^{(4)} &= a_i\text{-dependent terms} \quad + \quad \text{real part of loop calculation} \\ &\approx a_i\text{-dependent terms} \\ & (\text{for large } s,\ a_i) \end{aligned}$$

We approximate *real part of loop contribution* with one-loop Goldstone boson amplitudes using the Equivalence Theorem

Summary of the method

$$t_{IJ}^{(2)} \quad \rightarrow \quad \text{calculated from tree-level amplitude with } W_L$$

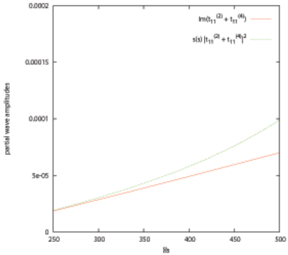
$$\text{Re}\, t_{IJ}^{(4)} \quad \rightarrow \quad \text{calculated from } a_i\text{-dependent terms with } W_L +$$

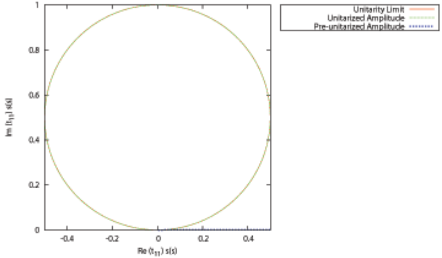
$$\text{real part of one-loop Goldstone boson scattering}$$

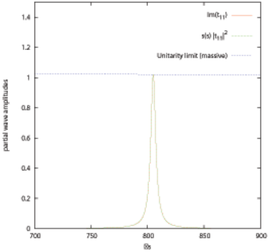
$$\text{Im}\, t_{IJ}^{(4)} \quad \rightarrow \quad \begin{cases} \sigma(s)|t_{IJ}^{(2)}|^2 + \sigma_H(s)|t_{H,IJ}^{(2)}|^2 & \text{if } I = 0 \\ \sigma(s)|t_{IJ}^{(2)}|^2 & \text{otherwise} \end{cases}$$

In addition, we neglect coupled channels (justified in as much as $a^2 - b$ is zero or very small in all cases).

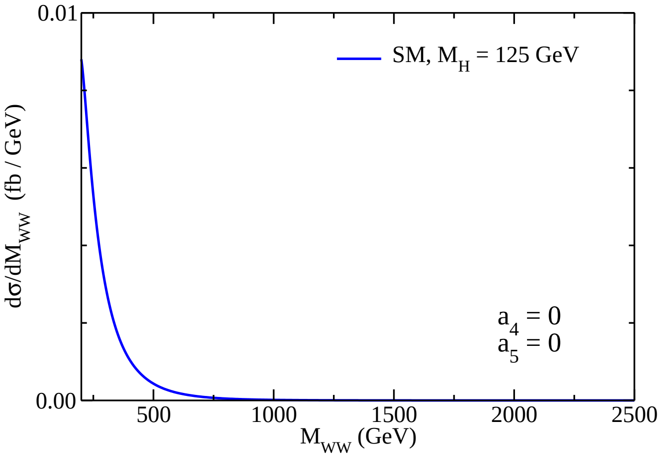
(Results from the Madrid group further justify this assumption)







Vector channel ($a_4 = 0.008, \quad a_5 = 0$)

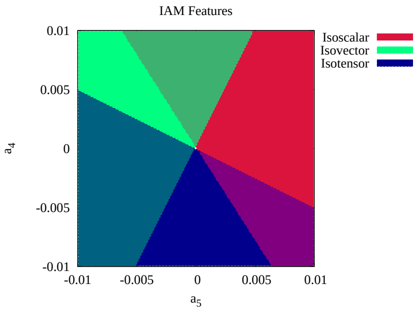


For SM ($a_4 = a_5 = 0$) there are *no additional resonances*.

Are there EWSB resonances?

Must search for poles in the second Riemann sheet — the phase shift must go through $+\pi/2$ at the resonance.

Are there any *physically acceptable* resonances?



The blue-shaded area leads to *acausal resonances*. These values for a_4 and a_5 are *unphysical* — they cannot be realized in any effective theory with a meaningful UV completion.

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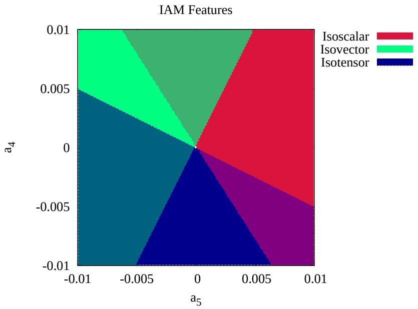

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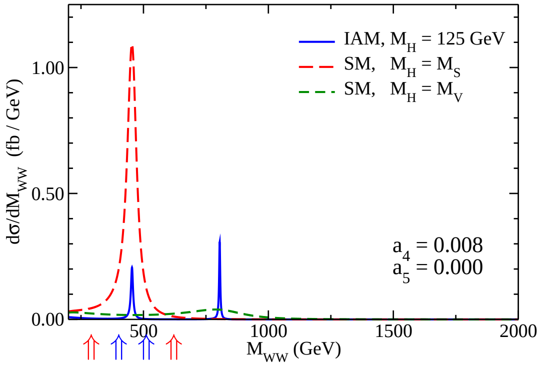
Are these resonances detectable?

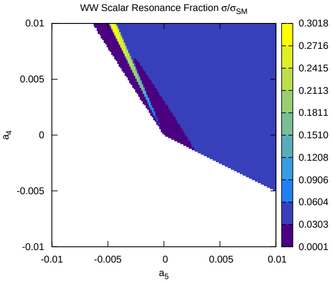


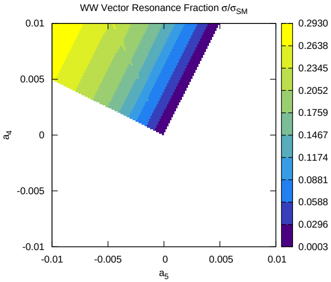
Signal of IAM scalar/vector vs. SM Higgs of *same mass*

The large contribution that the SM Higgs represents leaves little room for additional resonances.

Note: only in $WW \to WW$ or $WW \to ZZ$ channels!

Hard to detect!





Scalar

$R \sim 0.0001 - 0.30$

Typically < 0.1

Vector

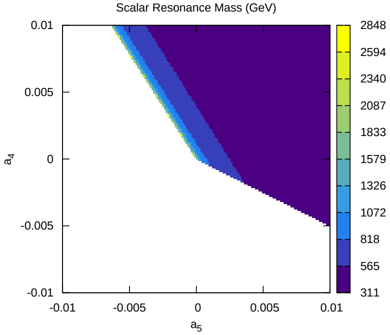
$R \sim 0.0003 - 0.3$

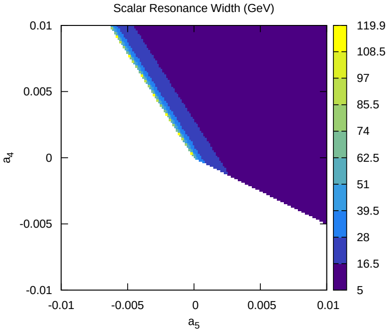
Typically ~ 0.2

They *could still be there*, but would give a small signal.

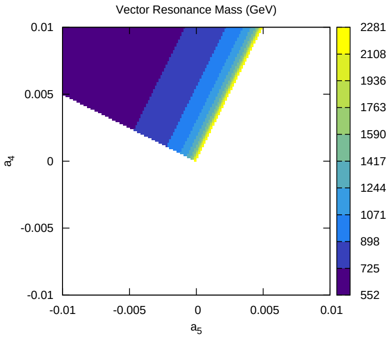
N.B.: Notice the enhancement region

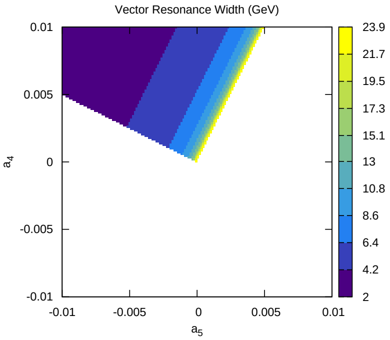
Scalar properties





- $M_S \sim 300 - 3000$ GeV
- $\Gamma_S \sim 5 - 120$ GeV

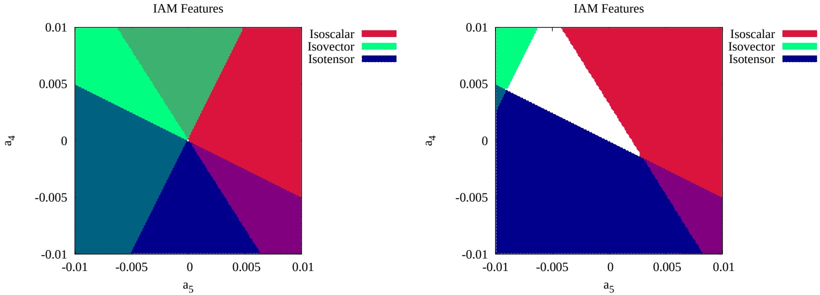




- $M_V \sim 550 - 2300$ GeV
- $\Gamma_V \sim 2 - 24$ GeV

Bounds on the anomalous QGC a_4 and a_5

Allowed regions for the anomalous couplings a_i (in white) if no resonance is found below ...



< 3 TeV < 600 GeV

Current limits: $a_4 \in [0.094, 0.10]$; $a_5 \in [0.23, 0.26]$

What if the hWW couplings are not exactly the SM ones?

Nothing prevents us from carrying out the same programme

$$\mathcal{L}_{\rm eff} = -\frac{1}{2}{\rm Tr} W_{\mu\nu}W^{\mu\nu} - \frac{1}{4}{\rm Tr} B_{\mu\nu}B^{\mu\nu} + \sum_i \mathcal{L}_{\rm i} + \mathcal{L}_{\rm GF} + \mathcal{L}_{\rm FP}$$
$$+ \left[1 + 2a\left(\frac{h}{v}\right) + b(\frac{h}{v})^2\right] \frac{v^2}{4}{\rm Tr} D_\mu U^\dagger D^\mu U + \frac{1}{2}(\partial_\mu h)^2 - \frac{1}{2}M_H^2 h^2$$
$$-d_3(\lambda v)h^3 - d_4\frac{1}{4}h^4$$

This *effective theory* is *non-renormalizable* and the a_i will be required to absorb the divergences

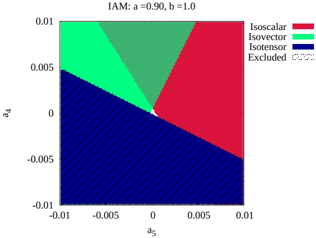
$$\delta a_4 = \Delta_\epsilon \frac{1}{(4\pi)^2}\frac{-1}{12}(1-a^2)^2$$

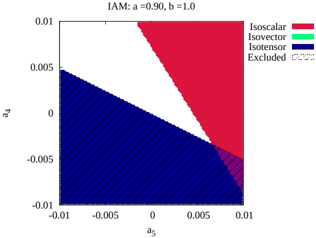
$$\delta a_5 = \Delta_\epsilon \frac{1}{(4\pi)^2}\frac{-1}{24}\left[(1-a^2)^2 + \frac{3}{2}((1-a^2)-(1-b))^2\right]$$

We have set $d_3 = d_4 = 1$ for simplicity.

Looking for resonances when $a \neq 1$

Exclusion zones and bounds on a_i for $a = 0.9$ ($b = a^2$)



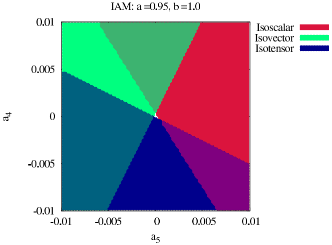


No resonances < 3 TeV

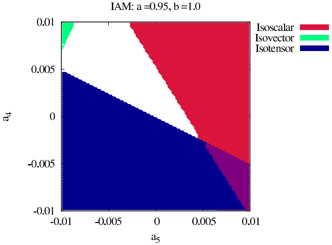
$M_R > 600$ GeV

The allowed regions for the anomalous couplings are slightly larger.

Bounds on a_i for $a = 0.95$ ($b = a^2$)



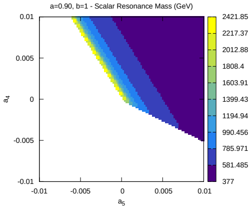
No resonances

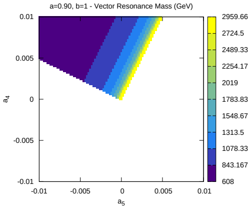


$M_R < 600$ GeV

The characteristics of the resonances tend smoothly to the $a = 1$ case (hWW coupling as in the SM).

Masses of scalar and vector resonances for $a = 0.9$



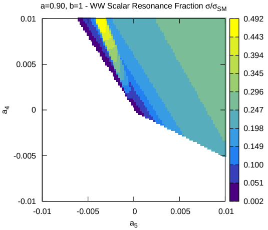


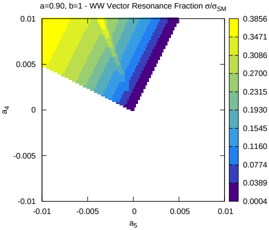
Resonances tend to be slightly heavier and broader than for $a = 1$

The parameter b is only marginally visible in the widths.

There are constraints on vector masses from S, T, U parameter constraints in some models. *e.g. Pich, Rosell, Sanz-Cillero, 2013*.

Visibility of the resonances for $a = 0.9$





Visibility of resonances for $a < 1$

Like for $a = 1$ the signal is always much lower than the one for a Higgs of the same mass.

For $a = 1$ typically $\sigma_{\text{resonance}}/\sigma_{\text{Higgs}} < 0.1$,
For $a < 1$ typically $\sigma_{\text{resonance}}/\sigma_{\text{Higgs}} \simeq 0.2$.

The situation for $a < 1$ is not radically different from $a = 1$

Resonances (particularly in the vector channel) are slightly more difficult to appear

They tend to be slightly heavier and broader

They give a slightly larger experimental signal

This situation changes drastically for $a > 1$

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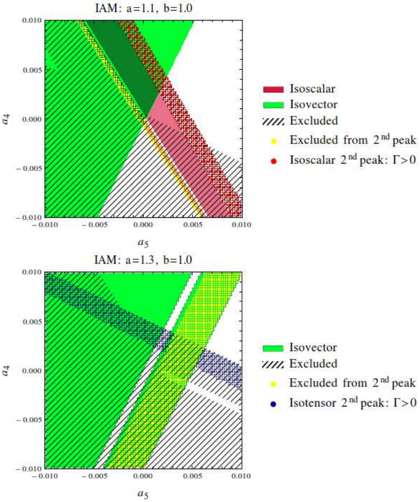
Spectrum of resonances for $a > 1$



'Something' happens when $a > 1$... *(Falkowski, Rychkov, Urbano [2012]; Espriu, Mescia [2014]; Bellazzini, Martucci, Torre [2014])*

Sum rule

Assuming the strict ET for all s

$$1 - a^2 = \frac{v^2}{6\pi}\int_0^\infty \frac{ds}{s}(2\sigma_{I=0}(s)^{tot} + 3\sigma_{I=1}(s)^{tot} - 5\sigma_{I=2}(s)^{tot}),$$

(Falkowski, Rychkov and Urbano, 2012)

However, we have seen that the analytic structure of the real amplitudes is more complex. Then

- LHS is modified to $3 - a^2 + \mathcal{O}(g^2)$
- The integral along the $|s| \to \infty$ does not necessarily vanish. *(Bellazzini, Martucci and Torre, 2014)*
- LHS gets renormalized while RHS does not...

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The difference between $1 - a^2$ and $3 - a^2$ can be traced back to the inclusion of W exchange in the t-channel.

If the propagatiog degrees of freedom remain *unchanged* all the way to $s = \infty$ (big 'if' !) the W t-channel contributes to the exterior circuit and gives $c_\infty = 2$ and restores the $1 - a^2$ on the LHS of the sum rule obtained when W propagation is ignored.
A trivial consequence of Cauchy's theorem

However, there is no guarantee that the propagating degrees of freedom remain *unchanged*. In fact in strongly interacting theories one might expect $c_\infty = 0$

Yet, as we have seen there is a substantial difference between $a > 1$ and $a < 1$ in the spectrum of dynamical resonances.

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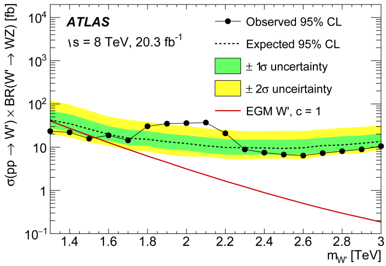
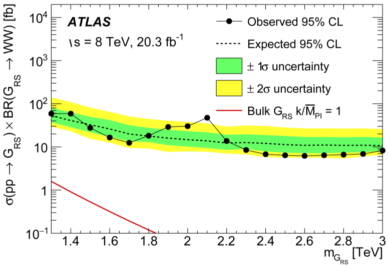
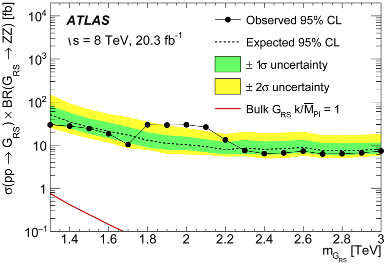
Conclusions I

- Unitarity is a powerful constraint on scattering amplitudes. Its validity is well tested.
- Even in the presence of a light Higgs, it can help constrain anomalous couplings by helping predict heavier resonances.
- An extended EWSBS would typically have such resonances even in the presence of a light 'Higgs'
- However their properties are radically different from the 'standard lore'
- Current LHC Higgs search results do not yet probe the IAM resonances: at least $10\times$ statistics is required
- . . . .

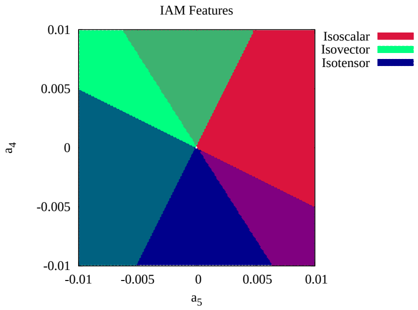
A resonance at 2 TeV?

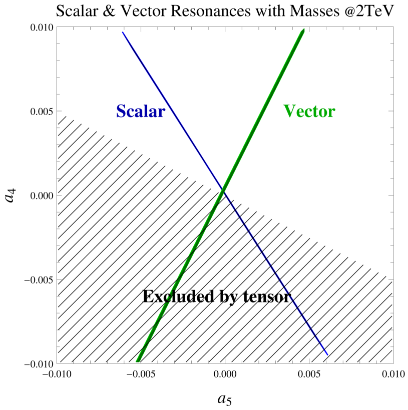




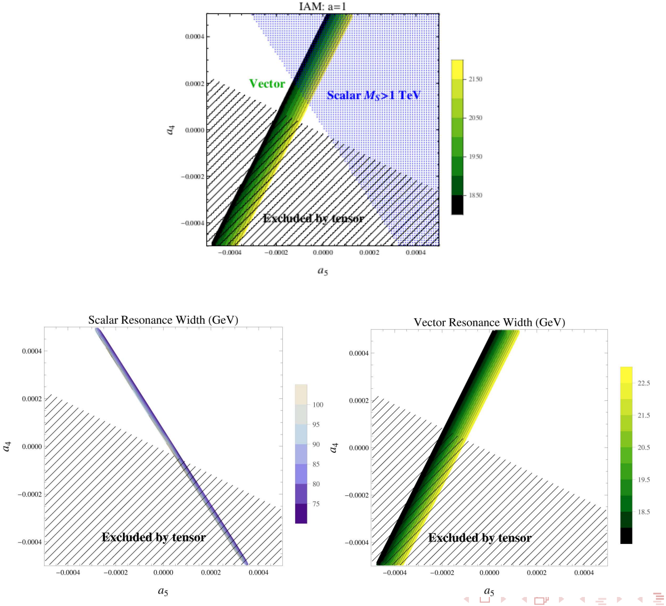


Redution of parameter space after resonance in 1800 GeV ¡ M_R ¡ 2200 GeV



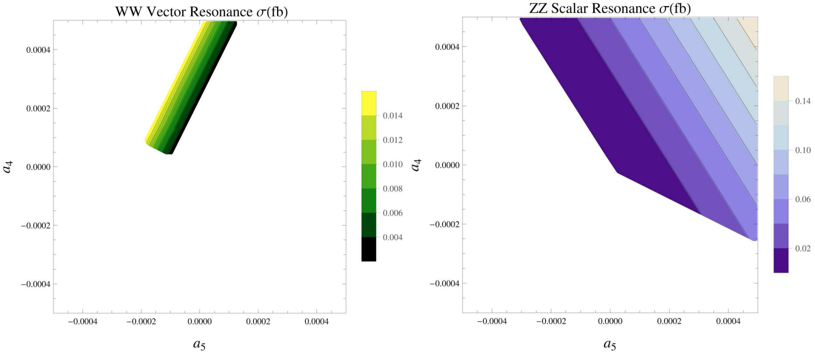


Blow-up and widths



Cross-sections

ATLAS results suggest cross-sections around *10 fb*.
Channel identification is dubious.



WW channel can be mediated by $I=1,2$. ZZ channel can be mediated by $I=0,2$.
Only resonant contribution to the cross-section is shown.

Comments and conclusions II

Cross sections are at least one order order of magnitude too small.
However:

- Near-degeneracy of scalar and vector gives also enhancement (also helps to explain the signal in *all* 3 channels
- Departures from $a = 1$ enhances the scalar signal.
  - Ex: going from $a = 1$ to $a = 0.95$ doubles the cross-section
  - Current experimental limit (90 % CL): $a \in [0.67, 1.33]$
  - In some models $a < 1$ implies breaking of custodial symmetry $\Rightarrow$ scalar isotriplets can exist, coupling proportionally to the custodial breaking parameter.
- Direct coupling of the resonances to quarks? (Drell-Yan)

All this is under current investigation!

THANK YOU!

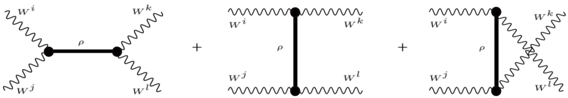
Back-up slides

Using form factors

$$\langle 0|V_\mu|\omega(q)\omega(q')\rangle = iF_V^+(q+q')_\mu + iF_V^-(q-q')_\mu$$

CVC $\Rightarrow$ $F_V^+ = 0$. Unitarity implies

$$\mathrm{Im}F_V = F_V t^* \quad t = A(\omega\omega \to \omega\omega) \quad K_V \equiv (s - M_V)^2 F_V$$

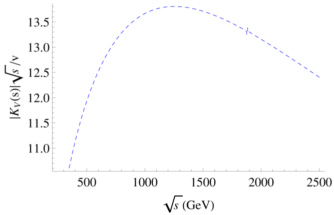


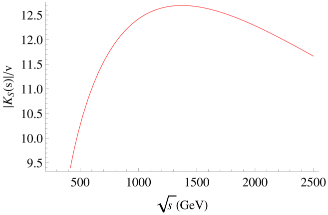
Then (symbolically)

$$t = K_V \frac{1}{s - M_V^2} K_V^* = \frac{t^{(2)}}{1 - t^{(4)}/t^{(2)}}$$

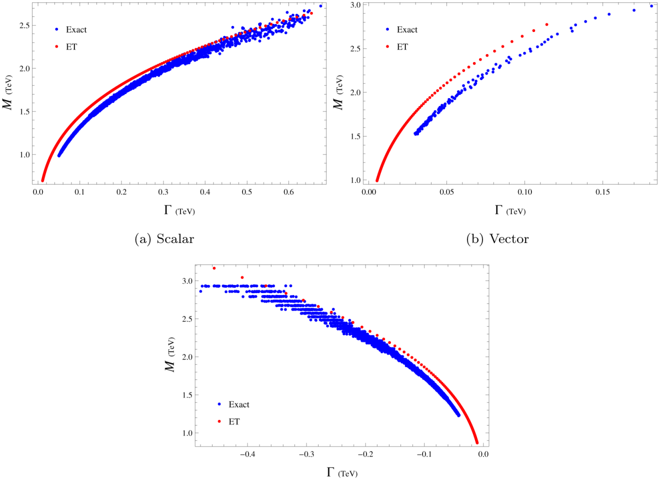
(Note: for I=1 u and t channels cancel each other)

Form factors



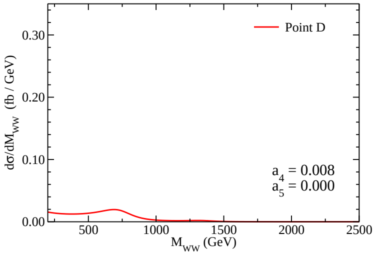


Comparing ET and 'exact' calculations

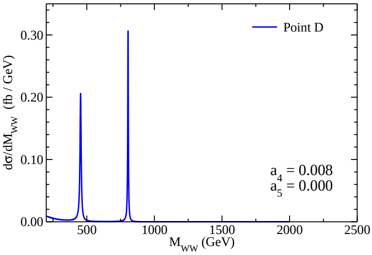


(a) Scalar (b) Vector

(c) Tensor



No light Higgs



With light Higgs

Here $\sqrt{s} = 8$ TeV

Compare before/after for same point (ex: Point D $a_4 = 0.008$, $a_5 = 0.000$)

- Different continuum
- Masses have changed positions
- Widths are *much narrower*

Resonance signal

We can estimate how observable these signals are by comparing to a heavy SM Higgs of the same mass $\rightarrow$ *look at LHC Higgs search data*

For a resonance of mass M_R and width Γ_R, let

$$\sigma^{peak} \equiv \int_{M_R-2\Gamma_R}^{M_R+2\Gamma_R} \left[dM_{WW} \times \frac{d\sigma}{dM_{WW}} \right]$$

$$\sigma_{SM}^{peak} \equiv \int_{M_H-2\Gamma_H}^{M_H+2\Gamma_H} \left[dM_{WW} \times \frac{d\sigma_{SM}}{dM_{WW}} \right]$$

Then for a heavy Higgs with $M_H \rightarrow M_R$ and $\Gamma_H(M_H \rightarrow M_R)$

$$R \equiv \left(\frac{\sigma^{peak}}{\sigma_{SM}^{peak}} \right)$$

compares the strength of the resonance regions of the same mass.

For $a = b = 1$ these results reproduce the SM prediction, i.e. no counterterms (renormalizable theory)

$$\delta a_4 = 0, \qquad \delta a_5 = 0$$

For $a = b = 0$ one gets the 'no Higgs' results (EChL)

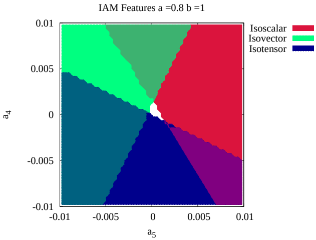
$$\delta a_4 = \Delta_\epsilon \frac{1}{(4\pi)^2} \frac{-1}{12}, \qquad \delta a_5 = \Delta_\epsilon \frac{1}{(4\pi)^2} \frac{-1}{24}$$

They should bring 'natural' finite contributions from NP:

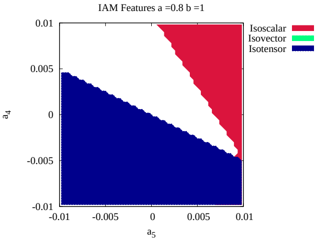
$$a_4|_{\rm finite} \simeq \frac{1}{(4\pi)^2} \frac{-1}{12} (1 - a^2)^2 \log \frac{v^2}{f^2}$$

$$a_5|_{\rm finite} \simeq \frac{1}{(4\pi)^2} \frac{-1}{24} \left[(1 - a^2)^2 + \frac{3}{2} ((1 - a^2) - (1 - b))^2 \right] \log \frac{v^2}{f^2},$$

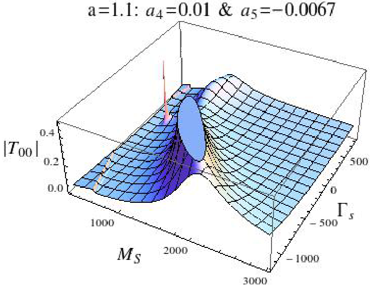
Bounds on a_i for $a = 0.8$ ($b = a^2$)

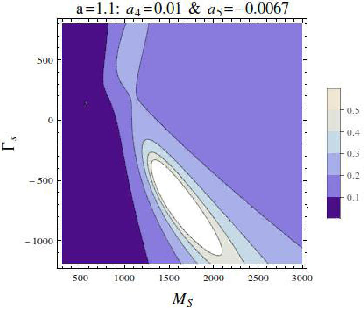


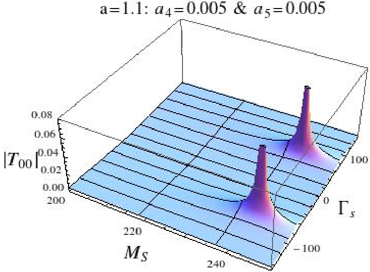
All Resonances



$M_R < 600$ GeV


a=1.1: a_4=0.01 & a_5=−0.0067
$|T_{00}|$
0.4
0.2
0.0
1000
2000
3000
M_S
500
0
−500
−1000
Γ_s


a=1.1: a_4=0.01 & a_5=−0.0067
500
0
−500
−1000
Γ_s
500
1000
1500
2000
2500
3000
M_S
0.5
0.4
0.3
0.2
0.1


a=1.1: a_4=0.005 & a_5=0.005
0.08
0.06
0.04
0.02
0.00
$|T_{00}|$
200
220
240
M_S
100
0
−100
Γ_s

$$\mathcal{L}_{QGC} = e^2\left[g_1^{\gamma\gamma} A^\mu A^\nu W_\mu^- W_\nu^+ - g_2^{\gamma\gamma} A^\mu A_\mu W^{-\nu} W_\nu^+\right]$$
$$+ e^2 \frac{c_\mathrm{w}}{s_\mathrm{w}}\left[g_1^{\gamma Z} A^\mu Z^\nu \left(W_\mu^- W_\nu^+ + W_\mu^+ W_\nu^-\right) - 2 g_2^{\gamma Z} A^\mu Z_\mu W^{-\nu} W_\nu^+\right]$$
$$+ e^2 \frac{c_\mathrm{w}^2}{s_\mathrm{w}^2}\left[g_1^{ZZ} Z^\mu Z^\nu W_\mu^- W_\nu^+ - g_2^{ZZ} Z^\mu Z_\mu W^{-\nu} W_\nu^+\right]$$
$$+ \frac{e^2}{2 s_\mathrm{w}^2}\left[g_1^{WW} W^{-\mu} W^{+\nu} W_\mu^- W_\nu^+ - g_2^{WW}\left(W^{-\mu} W_\mu^+\right)^2\right] + \frac{e^2}{4 s_\mathrm{w}^2 c_\mathrm{w}^4} h^{ZZ} (Z^\mu Z_\mu)^2$$

SM values: $g_1^{\gamma,Z} = \kappa^{\gamma,Z} = 1$, $\lambda^{\gamma,Z} = 0$ and $\delta_Z = \frac{\beta_1 + g'^2 \alpha_1}{c_\mathrm{w}^2 - s_\mathrm{w}^2}$ $\quad g_{1/2}^{VV'} = 1$, $h^{ZZ} = 0$

$$\Delta g_1^\gamma = 0 \qquad \Delta\kappa^\gamma = g^2(\alpha_2 - \alpha_1) + g^2\alpha_3 + g^2(\alpha_9 - \alpha_8)$$
$$\Delta g_1^Z = \delta_Z + \frac{g^2}{c_\mathrm{w}^2}\alpha_3 \qquad \Delta\kappa^Z = \delta_Z - g'^2(\alpha_2 - \alpha_1) + g^2\alpha_3 + g^2(\alpha_9 - \alpha_8)$$

$$\Delta g_1^{\gamma\gamma} = \Delta g_2^{\gamma\gamma} = 0 \qquad \Delta g_2^{ZZ} = 2\Delta g_1^{\gamma Z} - \frac{g^2}{c_\mathrm{w}^4}(\alpha_5 + \alpha_7)$$
$$\Delta g_1^{\gamma Z} = \Delta g_2^{\gamma Z} = \delta_Z + \frac{g^2}{c_\mathrm{w}^2}\alpha_3 \qquad \Delta g_1^{WW} = 2 c_\mathrm{w}^2 \Delta g_1^{\gamma Z} + 2g^2(\alpha_9 - \alpha_8) + g^2\alpha_4$$
$$\Delta g_1^{ZZ} = 2\Delta g_1^{\gamma Z} + \frac{g^2}{c_\mathrm{w}^4}(\alpha_4 + \alpha_6) \qquad \Delta g_2^{WW} = 2 c_\mathrm{w}^2 \Delta g_1^{\gamma Z} + 2g^2(\alpha_9 - \alpha_8) - g^2(\alpha_4 + 2\alpha_5)$$
$$h^{ZZ} = g^2\left[\alpha_4 + \alpha_5 + 2(\alpha_6 + \alpha_7 + \alpha_{10})\right]$$