# HVP contribution to the muon anomalous magnetic moment from lattice QC

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Using the Darwin (9600 core) Sandybridge/infiniband cluster at Cambridge, part of STFC's DiRAC HPC facility

Muon anomalous magnetic moment

 $\bar{\bar{S}}$ 

 $a_{\mu} =$ 2 Measure using polarised muons circulating in E and B fields. At a momentum where  $\beta \times E$  terms cancel, difference between precession and cyclotron frequencies:  $\vec{\mu} = g$ 2*m S*

 $g-2$ 

$$
\omega_a = -\frac{e}{m} a_\mu B
$$
  
BNL result:

*e*

 $a_\mu^{expt} = 11659208.9(6.3) \times 10^{-10}$ 

E989 (FNAL) will reduce exptl uncty to 1.6, starting 2017



 $\vec{\mu}\times\vec{B}$ 

### Standard Model theory expectations flavour

Contributions from QED, EW and QCD interactions. QED dominates. QCD contribs start at  $\alpha_{OED}^2$ and the rows, from top top top top to top top to both  $\mathcal{L}$ 



 $a_\mu^{QED} = 11658471.885(4) \times 10^{-10}$ 

 $a_\mu^{EW} = 15.4(2) \times 10^{-10}$ 

 $a_\mu^{E821}$  $= 11659208.9(6.3) \times 10^{-10}$ 

### Choortainty Gommatod by that from nacholite contributions Uncertainty dominated by that from hadronic contribns



### Hadronic contributions

$$
a_{\mu}^{expt} - a_{\mu}^{QED} - a_{\mu}^{EW} = 721.7(6.3) \times 10^{-10}
$$
  
=  $a_{\mu}^{HVP} + a_{\mu}^{HOHVP} + a_{\mu}^{HLBL} + a_{\mu}^{new physics}$ 

Focus on lowest order hadronic vacuum polarisation, so assume:

$$
a_{\mu}^{HLbL} = 10.5(2.6) \times 10^{-10}
$$
  

$$
a_{\mu}^{HOHVP} = -8.85(9) \times 10^{-10}
$$
  
MLO+NNLO  
<sub>Kurz et al,</sub>  
<sub>1403.6400</sub>

 $a_\mu^{HVP, no\ new\ physics} = 719.8(6.8)\times 10^{-10}$ 



### π<sup>+</sup>π: new data from BESIII; arXiv:1507.08188v2



Full analysis inc. BES data still analysis inc. BES data still at Benasque, 2015 to be done

# Lattice calculation of HVP Analytically continue to Euclidean  $q^2$ .

$$
a_{\mu}^{HVP,i} = \frac{\alpha}{\pi} \int_0^{\infty} dq^2 f(q^2) (4\pi \alpha e_i^2) \hat{\Pi}_i(q^2)
$$



Blum, hep-lat/ 0212018



Calculation with quarks Calculation required is

$$
J_\mu\left(\begin{matrix} \frac{\partial}{\partial u_1} & \cdots & \frac{\partial}{\partial u_n} \\ \frac{\partial}{\partial u_2} & \cdots & \frac{\partial}{\partial u_n} \end{matrix}\right)_\nu
$$

$$
= (q^2 g_{\mu\nu} - q_\mu q_\nu) \Pi(q^2)
$$

Fourier transform and plot out as a function of  $q^2$ 

E. Gregory, BMW, LAT15. Smeared clover action

correlation function of quark and antiquark propagators, created and destroyed by vector (photon) current



### Simpler method

For spatial vector currents at zero spatial momentum

$$
\Pi^{jj}(q^2) = q^2 \Pi(q^2) = a^4 \sum_t e^{iqt} \sum_{\vec{x}} \langle j^j(\vec{x},t)j^j(0) \rangle
$$

Time-moments of lattice current-current correlators

$$
G_{2n} \equiv a^4 \sum_{t} \sum_{\vec{x}} t^{2n} Z_V^2 \langle j^j(\vec{x}, t) j^j(0) \rangle
$$
  
=  $(-1)^n \left. \frac{\partial^{2n}}{\partial q^{2n}} q^2 \hat{\Pi}(q^2) \right|_{q^2=0}$   
 $J \left( \oint_Q \right)$ 

 $\hat{\Pi}(q^2) = \sum$  $\infty$ *j*=1  $q^{2j}\Pi_j$  with  $\Pi_j = (-1)^{j+1} \frac{G_{2j+2}}{(2j+2)^j}$  $(2j+2)!$ with

#### Allows us to reconstruct  $\Pi(q^2)$  and integrate  $\hat{\Pi}(q^2)$

Use Pade approximants (ratio of m/n polynomials) rather than Taylor expansion for better large  $q^2$  behaviour.

Test Pade approximants in similar scenarios (1-loop quark vacuum polarisation, with noise added)



Improved precision allows higher order Pade - we use [2,2]

### CHARM contribution

HPQCD 1004.4285, 1208.2855

# Part of the set of calculations that gave  $m_c, M(J/\psi) - M(\eta_c), \Gamma(J/\Psi \to e^+e^-), \Gamma(J/\psi \to \eta_c \gamma)$

Used HISQ valence quarks on MILC 2+1 asqtad configs.  $Z_v$  from contnm QCD pert. th.

Extrapolation to physical point allows us to compare directly to moments from e+e- expt. in charm region

$$
a_{\mu}^{HVP,c} = 14.4(4) \times 10^{-10}
$$
  
HPQCD 1403.1778



### BOTTOM contribution

HPQCD 1110.6887, 1309.5797, 1408.5768

### Part of the set of calculations that gave

 $m_b, M(\Upsilon) - M(\eta_b), M(\Upsilon') - M(\eta'_b), \Gamma(\Upsilon \to e^+e^-), \Gamma(\Upsilon' \to e^+e^-)$ 



#### STRANGE contribution polarization. Results are shown for 400 di↵erent simulations,



### Check mass and decay constant of  $\phi$  from these correlators against expt





## LIGHT contribution  $m_u = m_d$

HISQ valence quarks on MILC  $2+1+1$  HISQ configs. Use Z<sub>v</sub> from s calc.  $\mathcal{L}_{\text{v}}$  from s calc.

Multiple a (use  $w_0$ ),  $m_1$  (inc. phys.), volumes (at ml/ms=0.1). New ingredient since correlators much noisier. Use:  $\sum_{i=1}^{n}$ 

$$
G(t) = \begin{cases} G_{\text{data}}(t) & \text{for } t \le t^* \longleftarrow \text{from Monte Carlo} \\ G_{\text{fit}}(t) & \text{for } t > t^* \longleftarrow \text{from multi-exponential fit} \end{cases}
$$

$$
t^* = 1.5 \text{fm} \frac{6}{4} m_{\rho} \text{ so } 70\% \text{ of result from G}_{\text{data}}
$$

- $\left\{ \begin{array}{ll} 0.014 & \text{if } 0.014 \leq \theta \leq 0.014 \leq \theta \leq$ understand  $\rho$  on lattice, inc. finite-volume from  $\pi\pi$ . • 80% of result comes from  $\rho$  meson pole, so need to
	- $\cdot$  10% from  $\pi\pi$  sensitive to finite-volume and m  $\epsilon_0 \pi\pi$ taste-issues for staggered quarks). • 10% from  $\pi\pi$ , sensitive to finite-volume and  $m_{\pi}$  (so  $\pi$

One approach is to correct Taylor coefficients  $\int\limits_j^{latt} (\pi\pi)) \Bigg[$  $\mathsf T$ 

 $\hat{\Pi}^{latt}_i$ 

Remove lattice  $\pi\pi$ using effective theory of  $\rho, \pi, \gamma$ inc. staggered quark effects and finite vol. Jegerlehner +Szafron, 1101.2872

 $\prod$ 

 $\hat{\Pi}_{\dot{z}}^{latt}$ 

 $^{\mathit{latt}}_{j} \to (\Pi$ 

 $\hat{\Pi}^{latt}_i$ 

 $\frac{latt}{j}-\Pi$ 

Rescale using exptl elaborating on ETMC : 1308.4327. Reduces lattice systematics from light quark mass effects  $m_{\rho}$ 

 $m_\rho^{2j,latt}$ 

2*j,expt*

*m*

 $\overline{\rho}$ 

Restore  $\pi\pi$ from continuum effective theory

 $+ \Pi$ 

 $\hat{\prod}^{cont}_{i}$ 

 $\frac{cont}{j}(\pi\pi)$ 

 $\pi\pi$  contribution distorted at physical point using staggered quarks on these coarse lattices. Important to inc. other masses. But note: need 7fm lattice to reduce finite vol effects below 1% for contnm  $\pi\pi$ 





Future: improve statistics at physical point, finer lattices



Focus has been on<sup>t</sup> stochastic methods.<sup>+</sup>Using same source+  $\overline{\hat{\lambda}}$ for l and s helps  $\longrightarrow \longrightarrow \longrightarrow$ current with charge  $1/3$  (so  $e^2$  factor is  $1/5$  of connected) Guelpers, Mainz, LAT14 *Godiscus* (*x*<sup>0</sup>) = *ZV x*<sup>0</sup> *x*<sup>0</sup> *y su* et *lactul is*  $\frac{1}{2}$ <br> $\frac{2}{3}$ 





### HadSpec results e.g. Hadspec, 1309.2608

Use instead many  $(\sim 150)$  source vectors (eigenvectors of gauge-covariant Laplacian) for both conn. and disc. correlators to obtain good signal.



PRELIMINARY

Fitting and normalising to connected light, gives HVP disc. contribn of  $\sim -0.2\%$ anisotropic Hadspec+HPQCD,in prep.

clover action

Simple (but conservative) argument on size of disc. pieces l-l disc.pieces provide key difference between  $\omega$  and  $\rho$ 

$$
2D_{ll} = -\frac{f_{\rho}^2 m_{\rho}}{2} e^{-m_{\rho}t} + \frac{f_{\omega}^2 m_{\omega}}{2} e^{-m_{\omega}t}
$$

$$
\frac{\hat{\Pi}_{j,disc}}{\hat{\Pi}_{j,conn}} = \frac{1}{2} \left[ \frac{m_{\rho}^{2j+2} f_{\omega}^2}{m_{\omega}^{2j+2} f_{\rho}^2} - 1 \right]
$$

We do not have accurate information on decay constants because of width of  $\rho$ , mixing of  $\omega$  etc Taking  $f_{\rho} = 0.21(1) \,\text{GeV}, f_{\omega} = 0.20(1) \,\text{GeV}$ 

Disc. contribin reduced by factor of 5 from electric charge

 $\rightarrow$  HVP : disc-ll/conn-ll = -1.5(1.5) %

Adding contributions to (g-2)/2



CONCLUSIONS: Lattice - continuum comparison



A lot of progress but lattice uncty (all from u/d) still too big. Need to calc. QED,  $m_u/m_d$  effects (~1% and positive?) and disc. (negative). More calculations underway (Mainz, BMW, RBC/UKQCD …)

### Backup Slides

### Precision electroweak Higgs bounds

 $\text{sees HVP through}$   $\alpha_{QED}$  but  $\text{Hagiwara et a}$ sensitivity to range of exptl data is different Hagiwara et al, 1105.3149

Gfitter,1107.0975



Keep an eye on the 'big' picture whilst doing this …..



few MeV uncertainties in many cases

Keep an eye on the 'big'picture whilst doing this …..

