



Max-Planck-Institut für Physik
(Werner-Heisenberg-Institut)



Alexander von Humboldt
Stiftung/Foundation

Inhomogeneous thermalization and strongly coupled thermal flow via holography

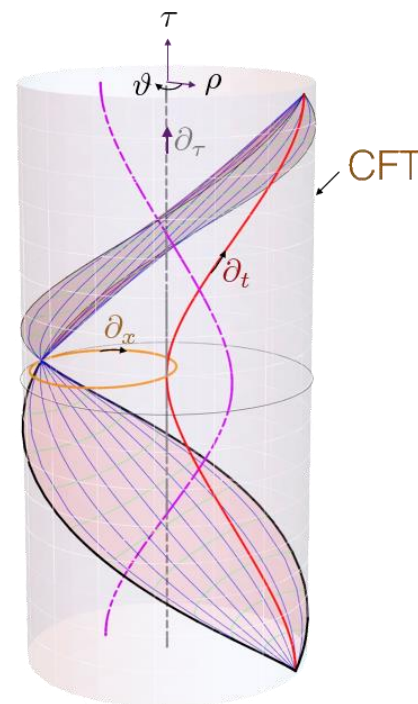
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The holographic path

Out-of-equilibrium strongly coupled field theory

- Pros of gauge/gravity approach:
 - Provides real-time analysis.
 - Allows for finite temperature setups.
 - New angle can lead to fresh new insights.
- Cons of gauge/gravity approach:
 - Caveats and limitations ($\mathcal{N}=4$ SYM)
 - Time dependence \rightarrow Solve PDEs

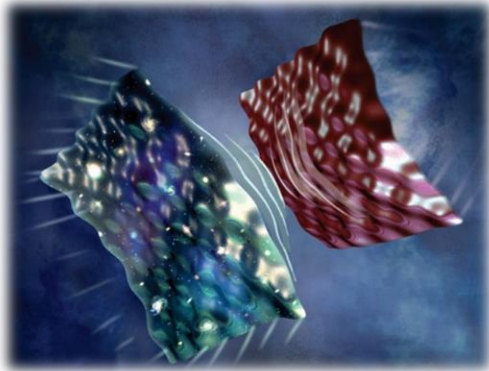


Applications

- ❖ Quark-gluon plasma thermalization [*Chesler, Yaffe, Heller, Romatschke, Mateos, vd Schee, Bantilan*]
- ❖ Turbulence in Gravity [*Lehner, Green, Yang, Zimmerman, Chesler, Adams, Liu*]
- ❖ Driven superconductors [*Rangamani, Rozali, Wong*]
- ❖ Quantum quenches [*Balasubramanian, Buchel, Myers, van Niekerk, Das*]
- ❖ Revivals [*Mas, Lopez, Serantes, da Silva*]

Outline

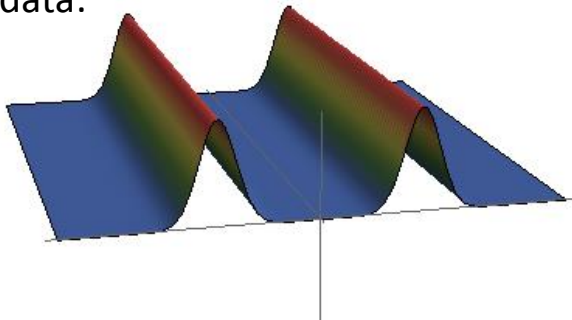
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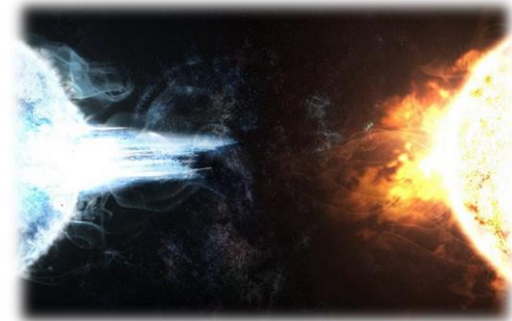
Collisions modeled as shockwaves

- Extraction of S-E tensor at boundary.
- Comparison with experiment.
- Main result: Thermalization time.

Initial data:



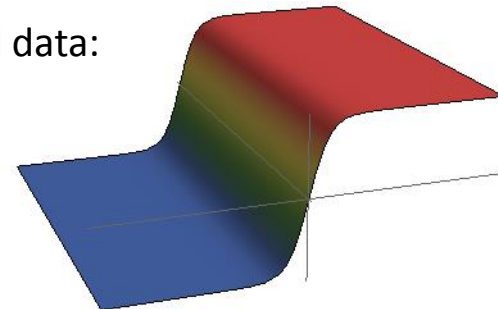
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Thermal flow modeled as dynamical horizon

- Extraction of steady state regime.
- Comparison with CFT and hydro.
- New result: Information flow.

Initial data:



Hydrodynamics approach

Conservation Law:

$$\partial_\mu T^{\mu\nu} = 0$$

- Hydro: Effective field theory of the slow evolution of expectation values of conserved currents close to equilibrium.
- Assumption: Some notion of local quasi-equilibrium.

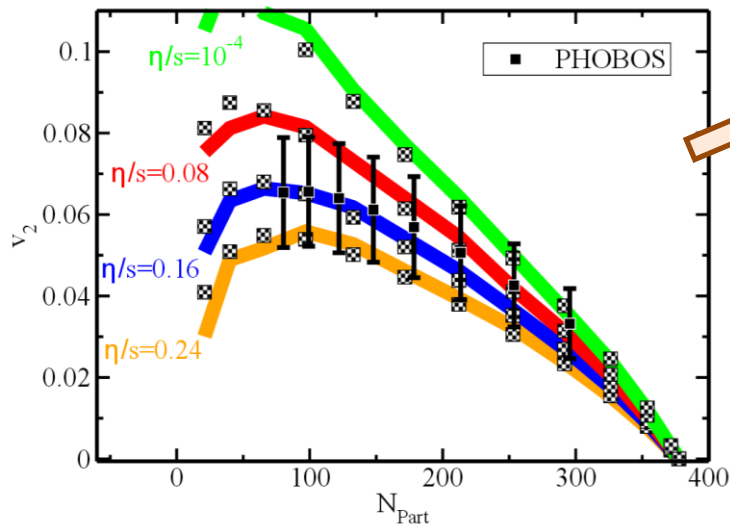
Gradient expansion:

$$T^{\mu\nu} = (\epsilon + p)u^\mu u^\nu + p g^{\mu\nu} + \Pi^{\mu\nu}$$

$\epsilon + p$: energy density
 p : pressure
 $u^\mu u^\nu$: 4-velocity
 $g^{\mu\nu}$: metric
 $\Pi^{\mu\nu}$: 1st and higher order gradient terms

Include viscosity term:

$$\Pi^{\mu\nu} = -2\eta g^{\mu\nu}$$



In HICs, the kind of physics that dominates is viscous fluid dynamics.

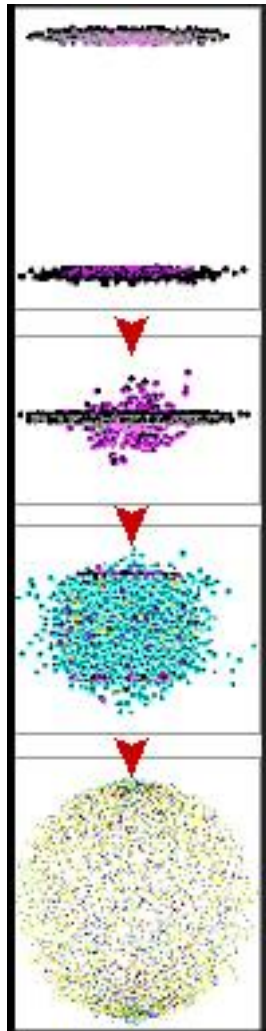
Connection to holography: $\frac{\eta}{s} = \frac{1}{4\pi} \sim 0.08$

- It should work when gradients are small.
- However, it seems to describe well HIC data!

[Luzum, Romatschke '09]

Evolution of a Heavy Ion Collision

(Picture not to scale)



Out of equilibrium

- Non-hydro degrees of freedom are dominant (at eq, they're damped).
- Gravity dual is useful description.
- Explicit and simple model: Shock wave collisions.

Hydrodynamization

- Transition to the hydrodynamics approximation.

Thermalization

- Isotropy of the diagonal components of $T^{\mu\nu}$ in local rest frame.
- Not necessary for the applicability of viscous hydro.

Hadronization

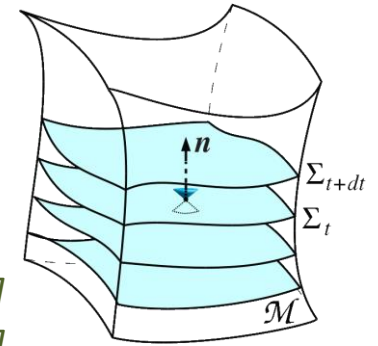
- The hadrons are back.
- Kinetic theory is applicable.

←----- $\Delta p \sim 0.7$

Gravitational formulation

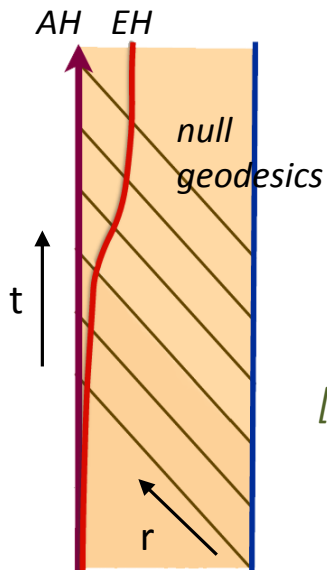
Cauchy Formulation

- ADM decomposition based on Hamiltonian formulation of GR.
- Generalized harmonic evolution scheme.
- Foliation of spacelike hypersurfaces.



[Pretorius, '05]

[Pretorius, Gubser, Bantilan '12]



Characteristic Formulation

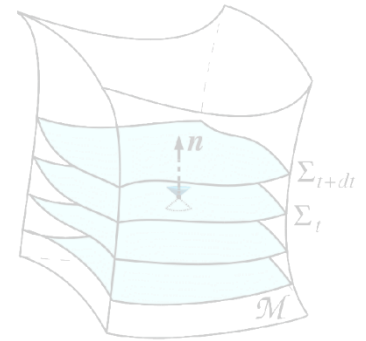
- Ingoing Eddington-Finkelstein coordinates.
- Foliation of null hypersurfaces.
- IR cutoff \rightarrow require fixed- r Apparent Horizon condition.

[Chesler, Yaffe '08]

Gravitational formulation

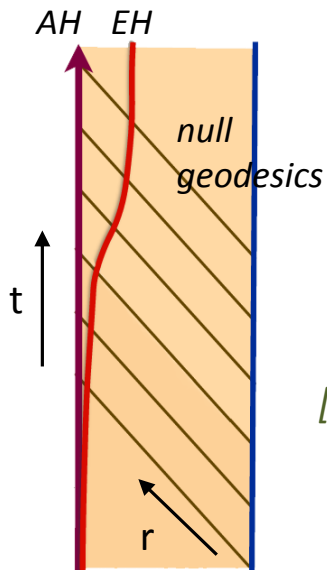
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Characteristic Formulation

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[Chesler, Yaffe '08]

$$\partial_{rt}g = \mathcal{F}[S, g] \Rightarrow \partial_t \mathcal{G} = \mathcal{F}, \quad \partial_r g = \mathcal{G}$$

$$\text{with conditions } \mathcal{C}_{ri}[S] = 0, \quad \mathcal{C}_{ti}[S] = 0$$

Gauge freedom:
coord. reparametrization

$$r \rightarrow r + \xi(x^i) \Rightarrow \text{AH fixed}$$

Bianchi identities imply they are
boundary constraints: $\nabla_r \mathcal{C}_{ri} = 0$

Outcome overview

- Plasma thermalizes very quickly: Hydro applicable within $\tau < 1/T$. *[Chesler, Yaffe '11]*
- Linearizing far-from-eq. State around final state:
Surprisingly accurate description. *[Heller, Mateos, van der Schee, Trancanelli '12]*
- Fully dynamical simulation of a HIC:
Holography + Hydro + cascade *[Pratt, Romatschke, van der Schee '13]*
- Simulation of BH collisions in AdS *[Bantilan, Romatschke '14]*
- Successful simulation of 4+1 evolution for off-center collisions *[Chesler, Yaffe '15]*

Ansatz equations

Null holographic coord. \rightarrow

$$ds^2 = 2 dt dr - A dt^2 + 2F_i dt dx^i + \Sigma^2 \tilde{g}_{ij} dx^i dx^j$$

Required initial data \rightarrow

Determinant of spatial part \rightarrow

Boundary EOM:

$$\partial_t a_4 = -\frac{4}{3} \partial_y f_4$$

$$\partial_t f_4 = -\frac{1}{4} \partial_y a_4 - 2 \partial_y b_4$$

Introducing ripples

$$ds^2 = 2 dt dr - A dt^2 + 2dt (F dy + G dx_1) + \Sigma^2 \left[e^{C-2B} \cosh D dy^2 + e^{B-C} \cosh D dx_1^2 + 2e^{-B/2} \sinh D dy dx_1 + e^B dx_2^2 \right]$$

Simplification from 3+1 to 2+1 by:

$$h(t, r, y, x_1) \rightarrow h_0(t, r, y) + \epsilon e^{ikx_1} \delta h(t, r, y)$$

Alternative: Use polar coordinates.

Assuming boost invariance and breaking rotational symmetry,

$$x_1, x_2 \rightarrow \rho, \phi$$

$$t = \tau \cosh \eta$$

$$y = \tau \sinh \eta$$

$$h(\tau, r, \rho, \phi) \rightarrow h_0(\tau, r, \rho) + \epsilon e^{im\phi} \delta h(\tau, r, \rho)$$



Initial Condition

Gaussian shocks

- Two separated shocks with finite thickness and energy density, moving toward each other:

$$ds^2 = \frac{L^2}{z^2} \left[-dt^2 + dz^2 + dy^2 + d\vec{x}^2 + \Phi(t, y, \vec{x}, z) (dy - dt)^2 \right]$$

Exact analytic solution, in Fefferman-Graham coordinates.

$\Phi(t, y, \vec{x}, z = 0)$ is the bdry. energy density, arbitrary if

$$\left(\partial_z^2 - \frac{3}{z} \partial_z + \nabla_{\vec{x}}^2 \right) \Phi = 0$$

- Transform from FG to EF coord. → Infalling geodesic congruence → Read initial data:

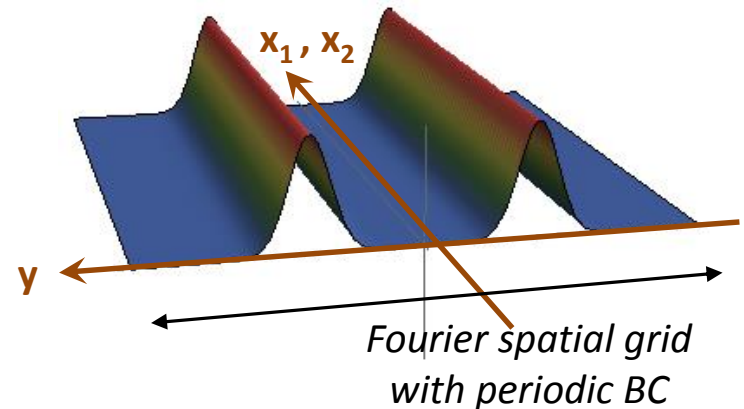
$$B(0, r, y), a_4(0, y), f_4(0, y)$$

- Generalize for perturbations. Freedom of choice, simple example: $\delta\mathcal{E} = \epsilon e^{ik x_1} \mathcal{E}$

- Need to add a (small) regulator energy density.

choose profile:

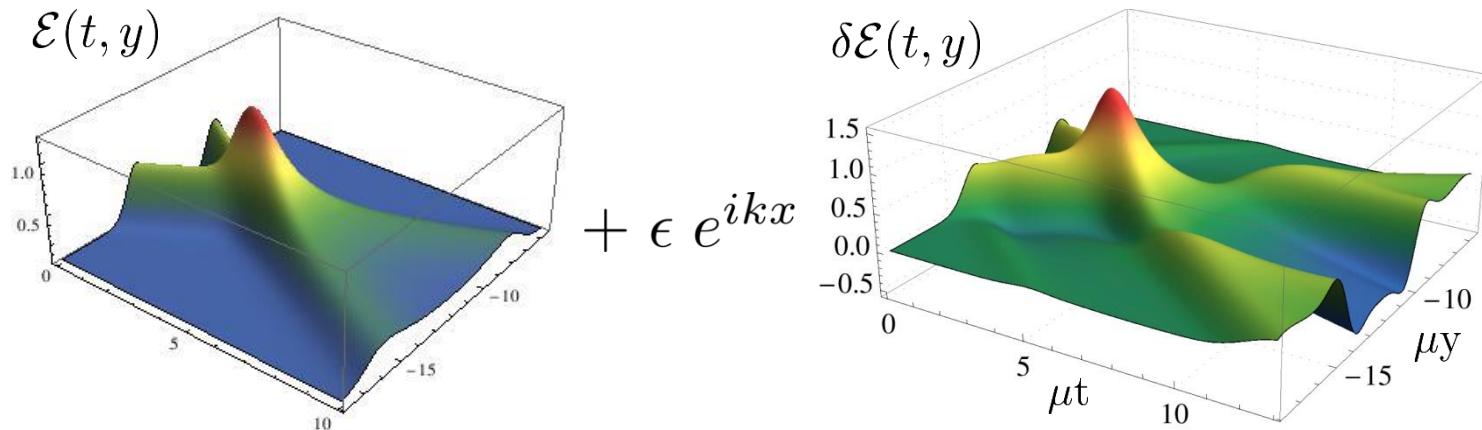
$$\Phi(t, y) = \frac{\mu^3}{2\pi w^2} e^{-(t \pm y)^2 / 2w^2}$$



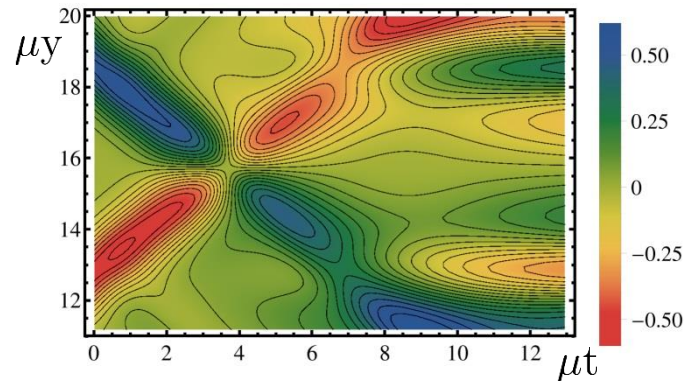
Stress-Energy Tensor results

- The inhomogeneities are evolved on top of the dynamical background.

Energy density:

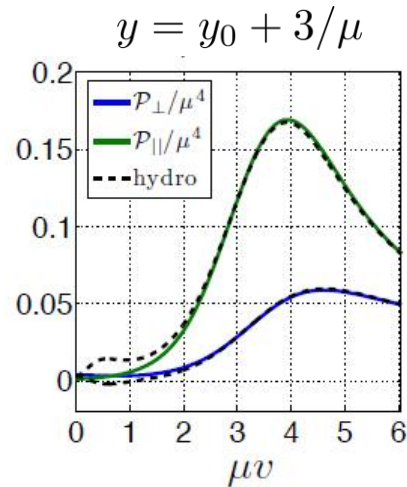
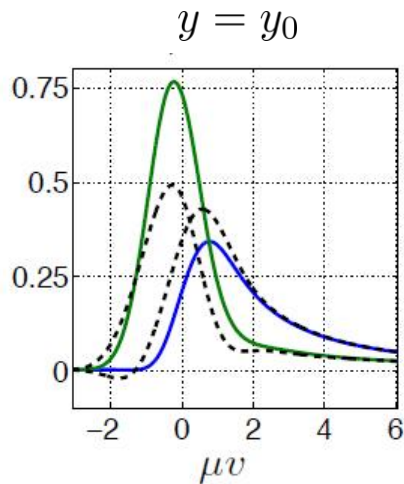


Longitudinal energy flux:



[Fernández '15]

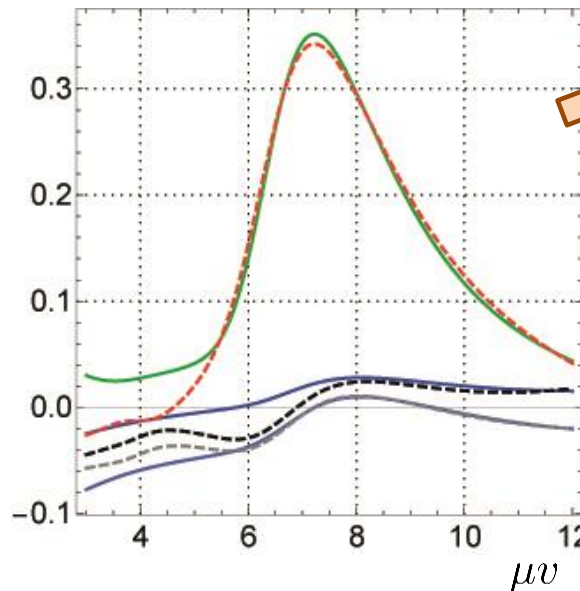
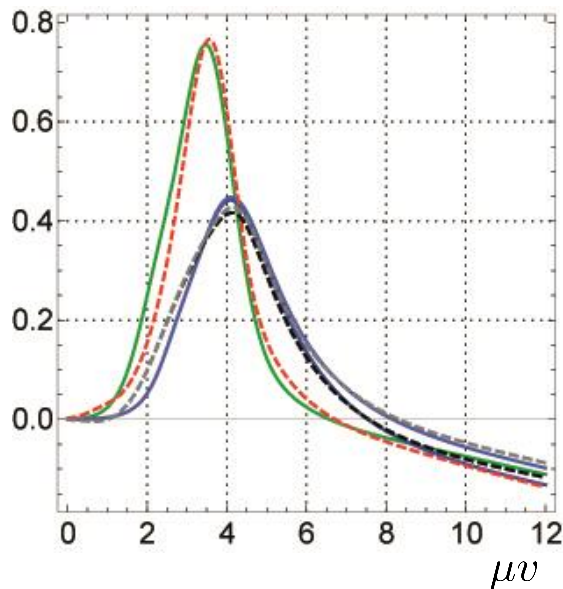
Comparing with hydrodynamics



- Compare with the pressures that would follow from the viscous hydrodynamic const. relations.

Mid-rapidity

- Dramatic rise in the pressures.
- Very anisotropic and out-of-eq.
- Hydro holds almost from the beginning.



Isotropization time and hydrodynamization time are completely different.

[Fernández '15]

Outlook: QGP chiral anomaly

In a Heavy Ion Collision:

- Strong magnetic fields are produced by the charged “spectators” ($B \sim 5 - 15 m_\pi^2/e$)
- Anomalous electric currents → Observable effects in hadron production.

In relativistic chiral plasma, the anomaly is present:

$$\partial_\mu J^\mu = a_1 F \wedge F + a_3 \text{Tr} G \wedge G$$

→ Anomalous electric current: $\vec{J} = \sigma_B \vec{B}$

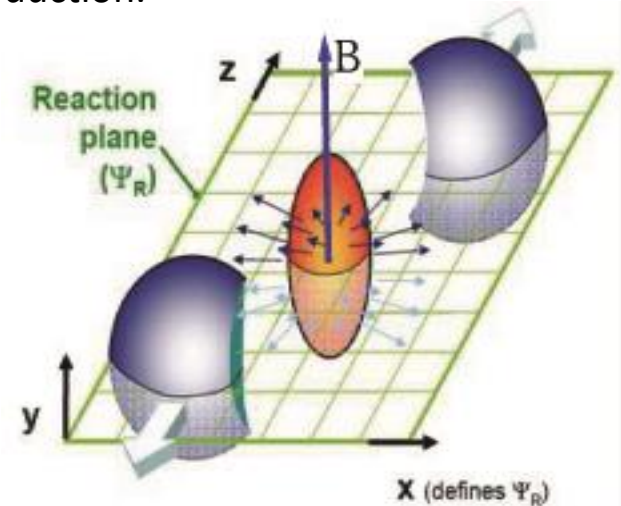
Holomodel ingredients

- Massive $A_\mu \rightarrow$ Emulate dynamical gluons for $\partial_\mu J^\mu \neq 0$.
- Scalar field $\theta \rightarrow$ Recover gauge invariance.
- CS term \rightarrow Non-dynamical part of anomaly.

❖ Stückelberg-Chern-Simons theory:

$$S = \int d^5x \sqrt{-g} \left[R + \Lambda - \frac{1}{4e} F^2 - \frac{m^2}{2} (A_\mu - \partial_\mu \theta)^2 + \frac{\kappa}{3} \epsilon^{\mu\alpha\beta\gamma\delta} (A_\mu - \partial_\mu \theta) F_{\alpha\beta} F_{\gamma\delta} \right]$$

...evolution of axial charge?



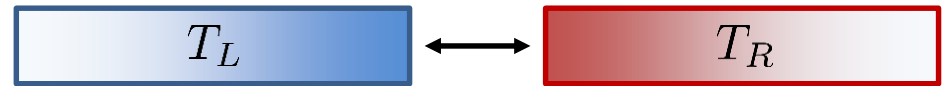
[Rebhan, Schmitt, Stricker '09]
[Gürsoy, Kharzeev, Rajagopal '14]

Universal regime of thermal transport

[Bernard, Doyon '12]

Thermal quench in 1+1

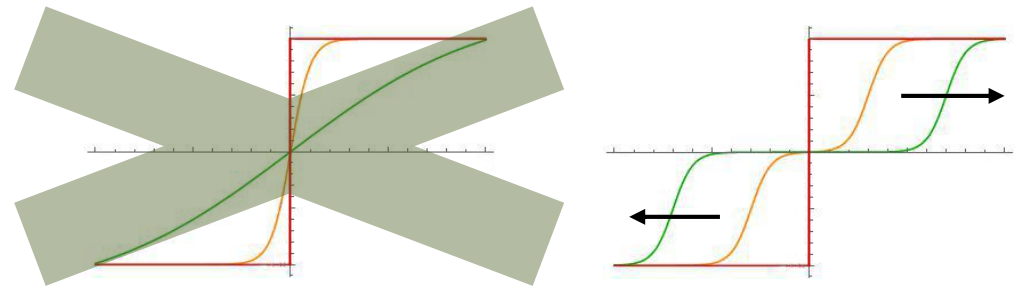
Two exact copies initially at equilibrium, independently thermalized.



Conservation eqs & tracelessness:

$$\partial_x \langle T^{xx} \rangle = -\partial_t \langle T^{tx} \rangle = 0$$

$$\langle T^{xx} \rangle = \langle T^{tt} \rangle$$



$$\left. \begin{aligned} \langle T^{tx} \rangle &= F(x-t) - F(x+t) \\ \langle T^{tt} \rangle &= F(x-t) + F(x+t) \end{aligned} \right\} \begin{array}{l} \text{shock waves emanating from interface,} \\ \text{converge to non-equilibrium } \textit{Steady State}. \end{array}$$

Long time limit

$$\langle T^{tx} \rangle = cg(T_L - T_R)$$

Energy flow

$$\langle J_E \rangle \neq 0$$

(Lorentz-boosted thermal distribution)

Thermal transport in $d > 1$

[Bhaseen, Doyon, Lucas, Schalm '13]

Holographic dual

- In 1+1, holographic dual is unique: Boosted BTZ black hole.
- In $d+1$, assume ctant. homogeneous heat flow as well:

$$\langle T^{\mu\nu} \rangle = a_d T^{d+1} (\eta^{\mu\nu} + (d+1)u^\mu u^\nu) \quad \rightarrow$$

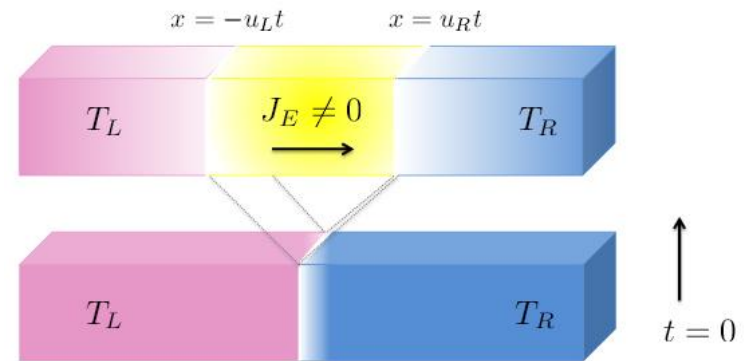
also unique non-singular solution:
Boosted black brane.

Generalization to any d:

- Effective dimension reduction to 1+1.
- *Dissipation*: Energy can be exchanged among the various constituents.
- Linear response regime:

$$|T_L - T_R| \ll T_L + T_R$$

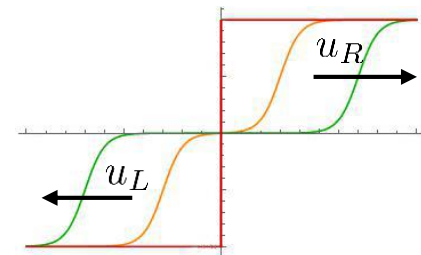
\rightarrow Hydro eqs. explicitly solvable.



$\langle \vec{J}_E \rangle \neq 0$ even if systems asymptotically far apart.

❖ Heat transport dominated by diffusion instead of flow?

\rightarrow $\left\{ \begin{array}{l} \text{width} \propto \sqrt{t} \\ \text{distance} \propto t \end{array} \right.$



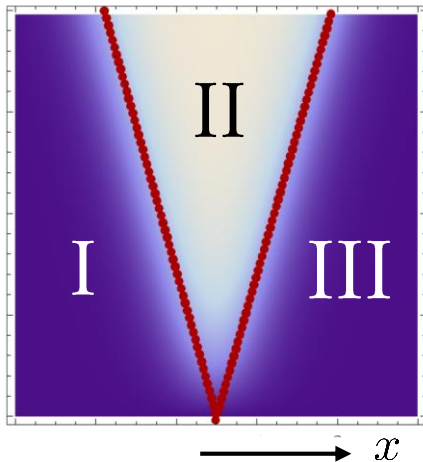
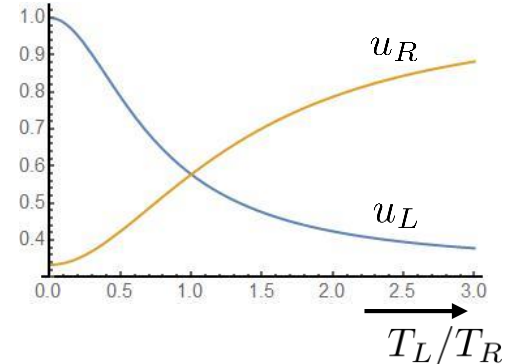
Hydrodynamics approach

➤ Generic d prediction:

Final steady state characterized by

$$\langle T^{tx} \rangle = a_d \frac{T_L^{d+1} - T_R^{d+1}}{u_L + u_R} \quad \text{where} \quad a_d \sim \frac{L^d}{G_N} \quad \text{and}$$

➤ Shockwave velocities: $u_L u_R = \frac{1}{d} = c_s^2$



Hydrodynamical evolution of 3 regions

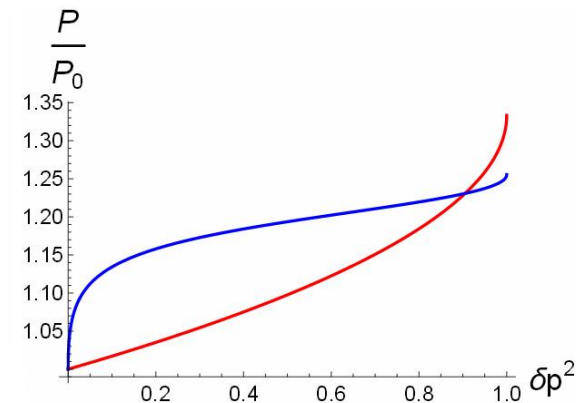
Match solutions \rightarrow asymptotics of the central region?

$$T^{\mu\nu} = P(d u^\mu u^\nu + \eta^{\mu\nu}) + \pi^{\mu\nu} + \mathcal{O}(\partial^2) \quad \text{with} \quad p = \frac{\epsilon}{d-1}$$

looking for sols. of the form:

$$u^\mu = (\cosh[\theta(t, x)], \sinh[\theta(t, x)], \vec{0})$$

Two possible Steady State configurations:
 Thermodynamic branch, second branch



Late time prediction
is matched

[Bhaseen, Doyon, Lucas, Schalm '13]
 [Chang, Karch, Yarom '13]

Holographic approach

[Amado, Yarom '15]

Energy transport: Lorentz-boosted thermal distribution

→ Gravity dual: Boosted black brane

$$ds^2 = \frac{L^2}{z^2} \left[\frac{dz^2}{f(z)} - f(z) (\cosh \theta dt - \sinh \theta dx)^2 + (\cosh \theta dx - \sinh \theta dt)^2 + dx_{\perp}^2 \right]$$

Numerical details

➤ Full out-of-equilibrium evolution of Einstein's equations.

- Ansatz for Characteristic Evolution (Eddington-Finkelstein coord.):

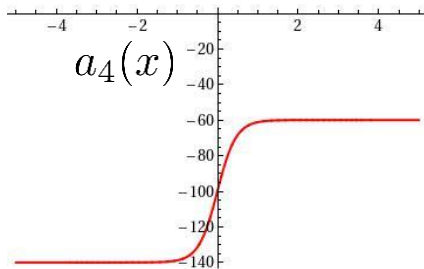
$$ds^2 = 2dt (dr - A dt - F dx) + \Sigma^2 (e^B dx_{\perp}^2 + e^{-B} dx^2)$$

- Initial data:

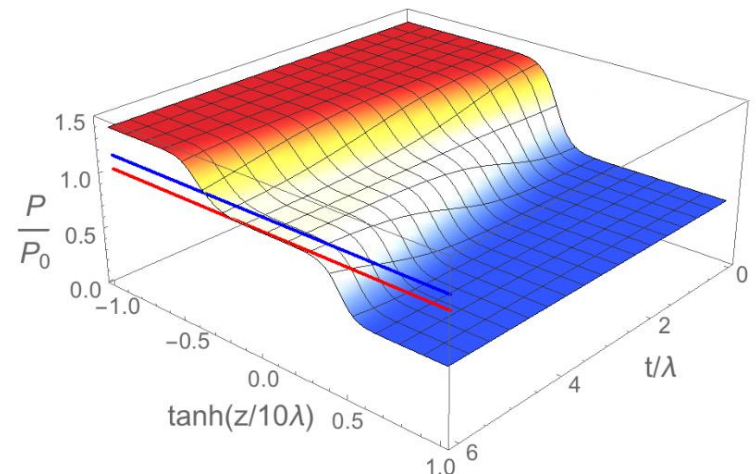
$$B = 0,$$

$$f_4 = 0,$$

and



➔ evolution



Check
via
holography

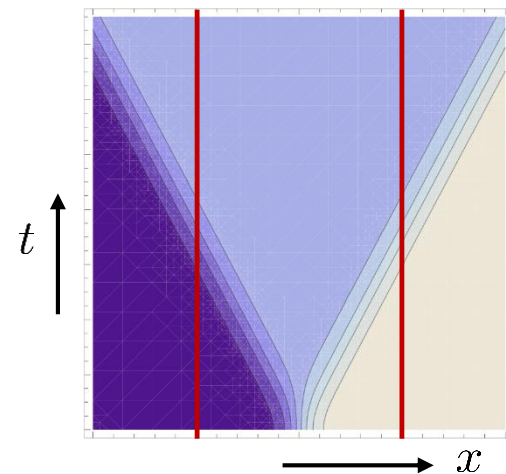
Information flow

How does information get exchanged between the systems which are isolated at $t=0$?

- Entanglement entropy

→ Holographic measure: $S_E = \frac{\text{Area}(\gamma_A)}{4G_N^{d+2}}$

- Regularization: Subtract result at $T=0$: $S_{EE}^{div} \sim \frac{1}{z_{uv}}$



- Spatial induced metric on surface:

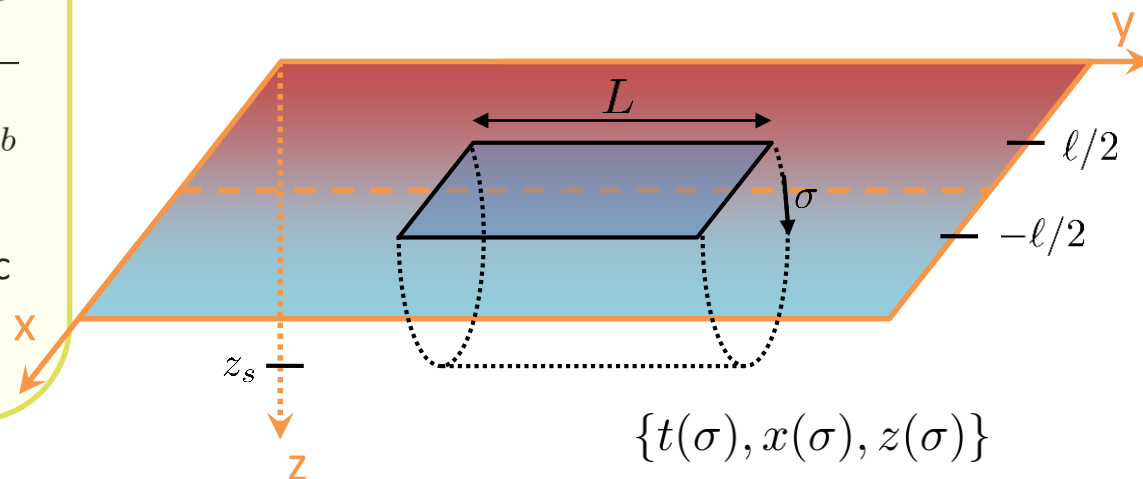
$$ds^2 = h_{\sigma\sigma}d\sigma^2 + g_{yy}dy^2$$

- Action:

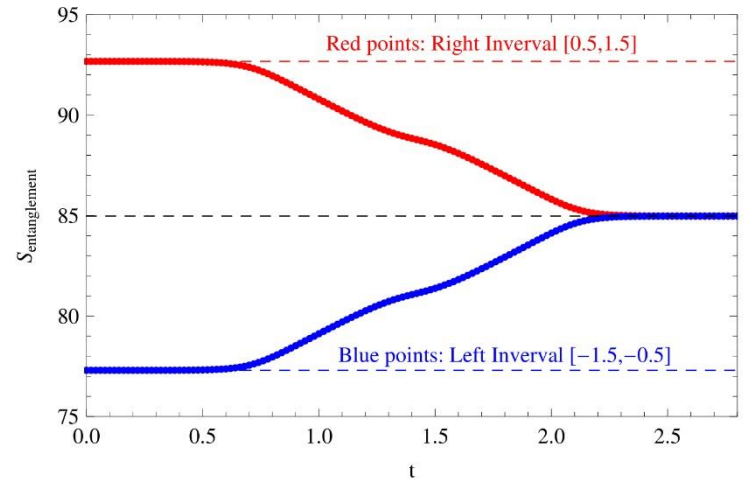
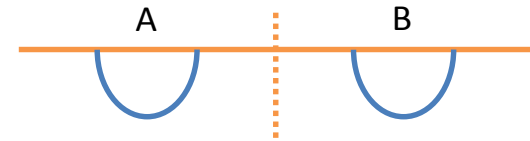
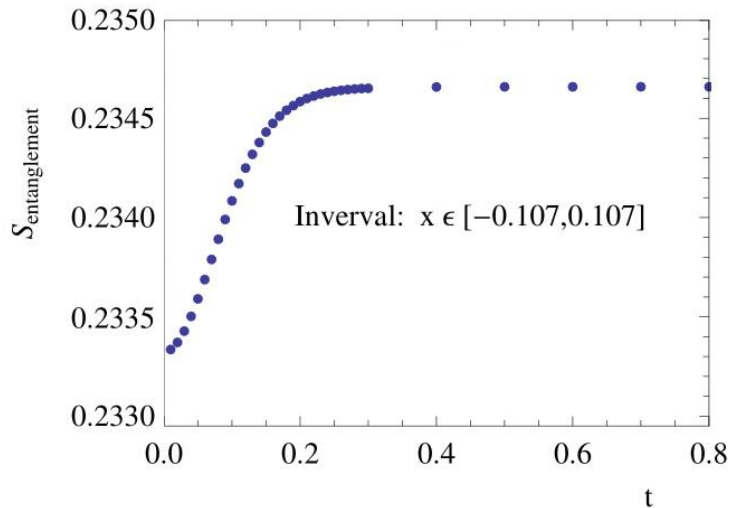
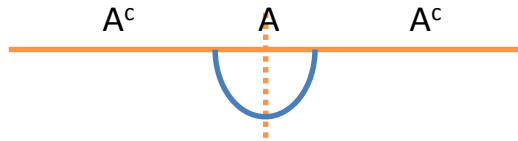
$$\mathcal{I} = L \int ds \sqrt{\frac{\partial x^a}{\partial \sigma} \frac{\partial x^b}{\partial \sigma} \tilde{g}_{ab}}$$

- Geodesic eqs. for effective metric

$$\tilde{g}_{ab} = g_{yy}g_{ab}$$



Entanglement Entropies

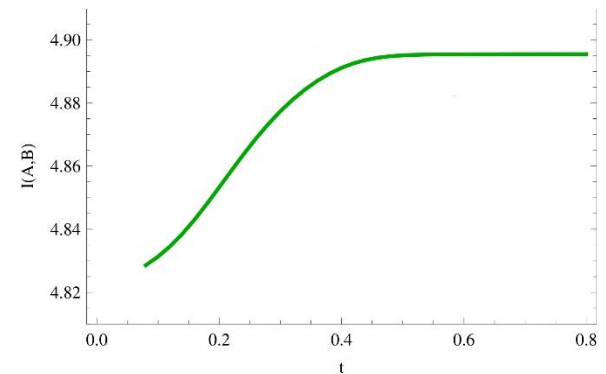


Mutual Information

What do we learn about A by looking at B?

$$I(A, B) = S(A) + S(B) - S(A + B)$$

- Shockwaves transport information.
- M.I. grows as shocks pass through the region.



[Erdmenger, Fernández, Flory, Megías, Straub '15]

Outlook

- QGP anomalies

- a. *Time evolution of the axial charge*
- b. *Axial asymmetry in hadronic distribution?*

- Azimuthal perturbations with radial flow

- a. *Comparison with elliptic flow*
- b. *Connection to hydrodynamics*

- Thermal flow analysis

- a. *Generalize: $2+1 \rightarrow 3+1$, what changes?*
- b. *Asymmetry – Intervals are not reached simultaneously*

Thank you!