

# Contribution of Anisotropic Particle Escape to Elliptic Flow from Transport Models

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Based on *arXiv:1502.05572*:

Liang He, Terrence Edmonds, Zi-Wei Lin, Feng Liu, Denes Molnar,  
Fuqiang Wang:

*Anisotropic parton escape is the dominant source of azimuthal anisotropy*

- *in transport models (v3, much expanded from v1)*
- *from A Multi-Phase Transport (v1)*

# Outline

Current Picture of  $v_n$  Development in Heavy Ion Collisions

Method: Tracking Parton Collisions in Transport Models

Our Results

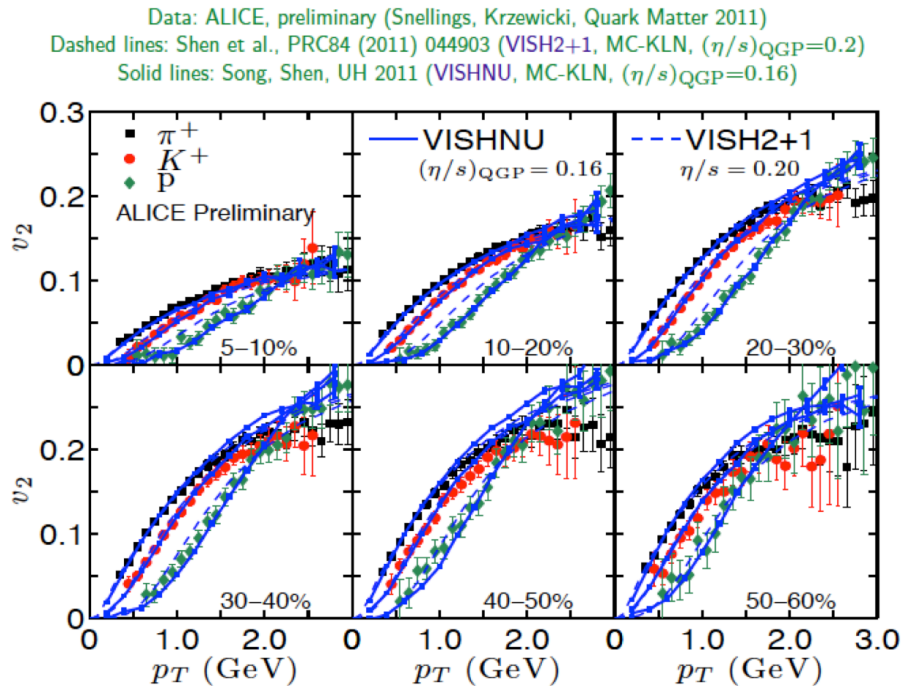
Potential Consequences

# Current Picture of $v_n$ Development

Early hydro-type collective flow in sQGP  
converts initial spatial anisotropy into final momentum-space  $v_n$

**Hydrodynamics** has been very successful  
for global observables, especially flow  $v_n$

## $v_2(p_T)$ in PbPb@LHC: ALICE vs. VISHNU



VISHNU yields correct magnitude and centrality dependence of  $v_2(p_T)$  for pions, kaons and protons!

Same  $(\eta/s)_{\text{QGP}} = 0.16$  (for MC-KLN) at RHIC and LHC!

Heinz, BES Workshop at LBNL 2014  
using viscous hydrodynamics.

**Transport model** can describe flow  $v_n$  :  
*degree of equilibration is controlled by cross section  $\sigma$*

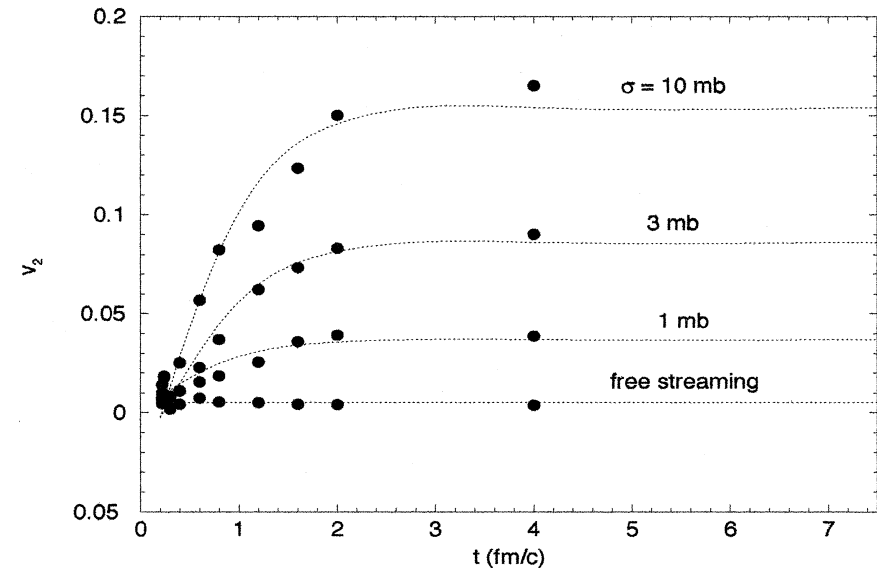
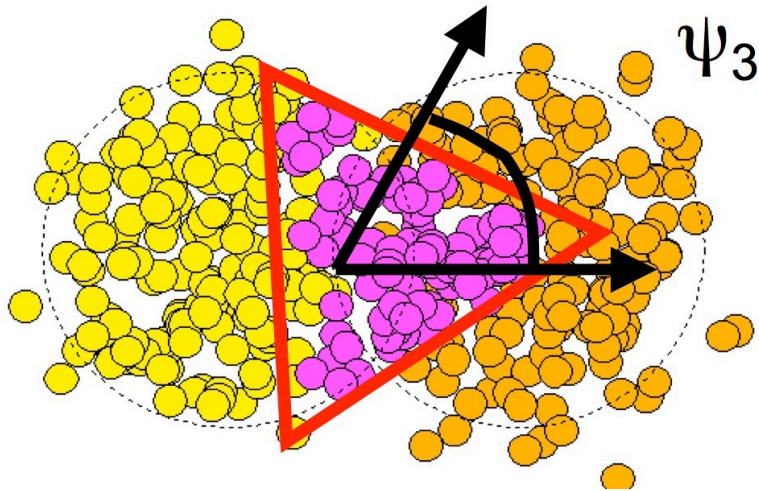


Fig. 1. Time evolution of  $v_2$  coefficient for different effective parton scattering cross sections in Au-Au collisions at  $\sqrt{s} = 200$  AGeV with impact parameter 7.5 fm. Filled circles are cascade data, and dotted lines are hyperbolic tangent fits to the data.

Zhang, Gyulassy and Ko, PLB (1999)  
using elastic parton transport.

# Current Picture of $v_n$ Development

Both **hydrodynamics** and **transport model** have been used to study  $v_n$  :

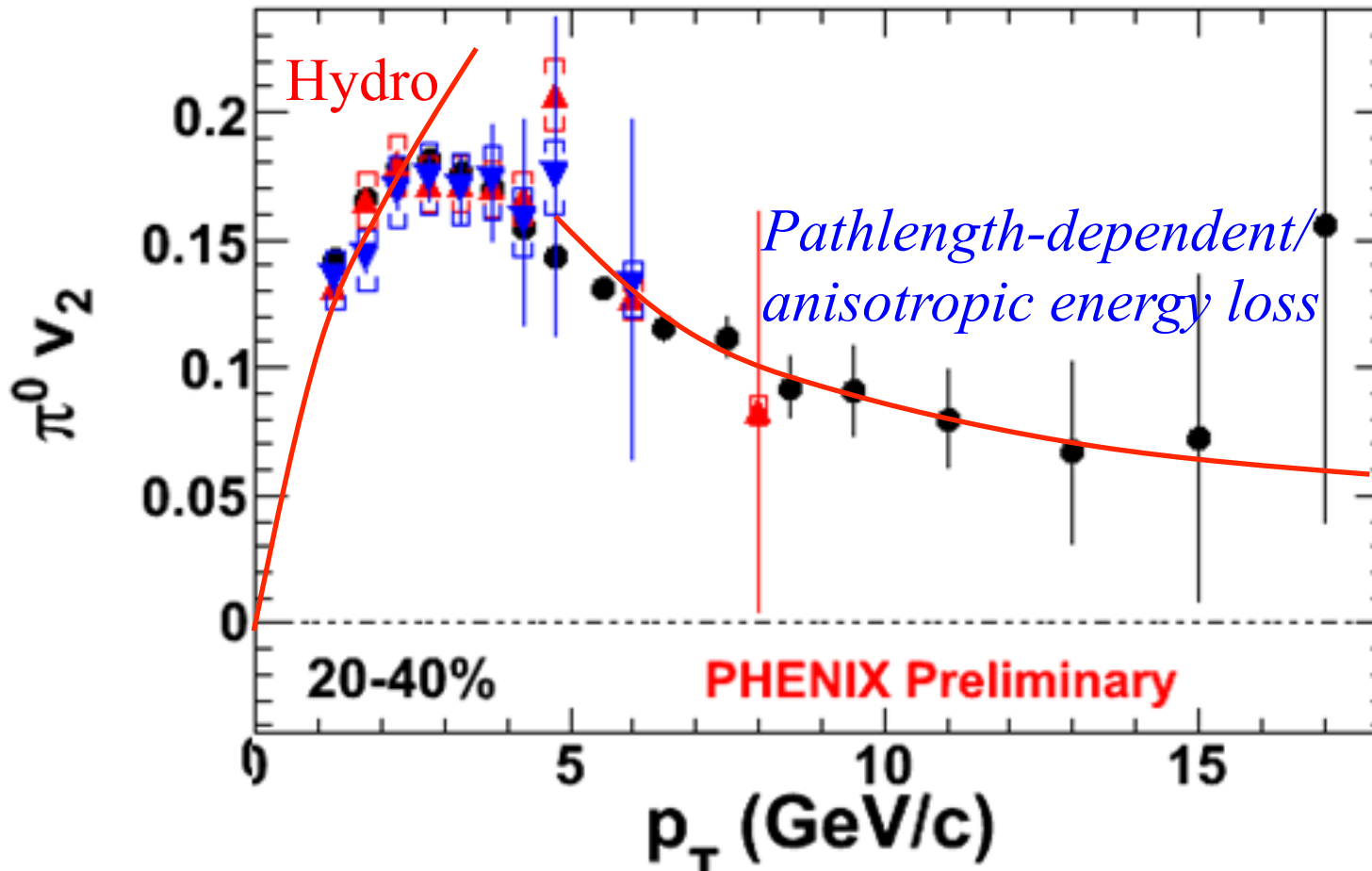


Alver and Roland, PRC (2010) discovered significant triangular flow using A Multi-Phase Transport (AMPT);  
→ intense developments of event-by-event hydrodynamics.

Transport at large-enough cross section will approach hydrodynamics.

**It is generally believed:** for low- $P_T$  in high-energy heavy ion collisions, the mechanism of  $v_n$  development from transport model (*via particle interactions*) is in principle the same as viscous **hydrodynamics** (*via pressure gradient*).

# Current Picture of $v_n$ Development



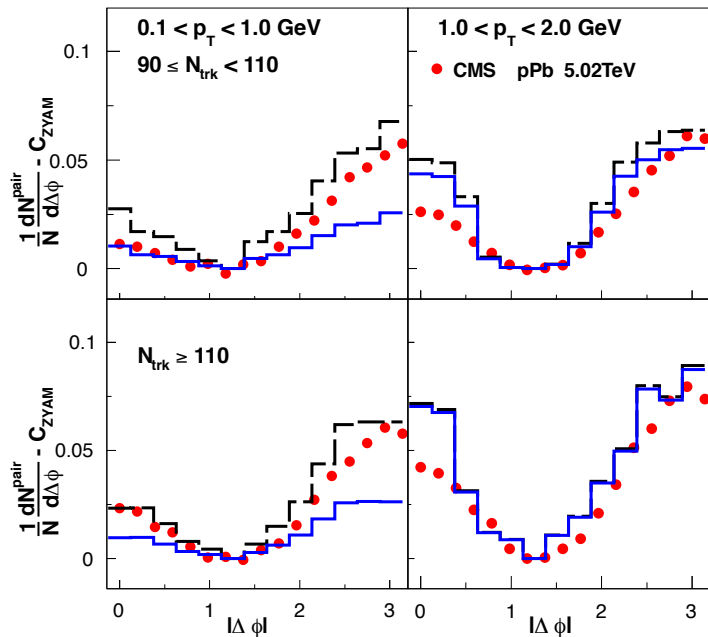
## A different paradigm for high- $P_T$

It is generally believed:

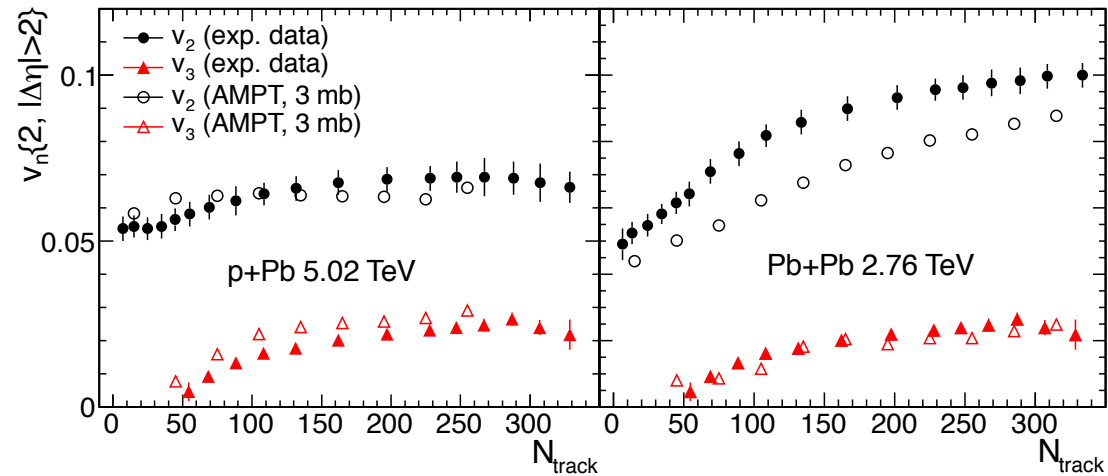
high- $P_T$  observables cannot be described by **hydrodynamics**,  
one needs particle **transport** (*plus energy loss, fragmentation, etc*)

# Current Picture of $v_n$ Development

Small systems: both **hydrodynamics** and **transport** can describe flow



Bozek and Broniowski, PLB (2013)  
using e-by-e viscous hydrodynamics.



Bzdak and Ma, PRL (2014)  
using A Multi-Phase Transport (AMPT).

**Puzzle for small systems** such as p+Pb or d+Au:

Mean free path may be comparable to the system size;

is **hydrodynamics** still applicable to such small systems?

# Tracking Parton Collisions in Transport Models

## Method:

Study  $v_n$  development

by tracking the complete collision history of each parton,

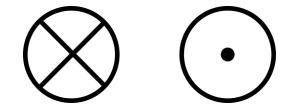
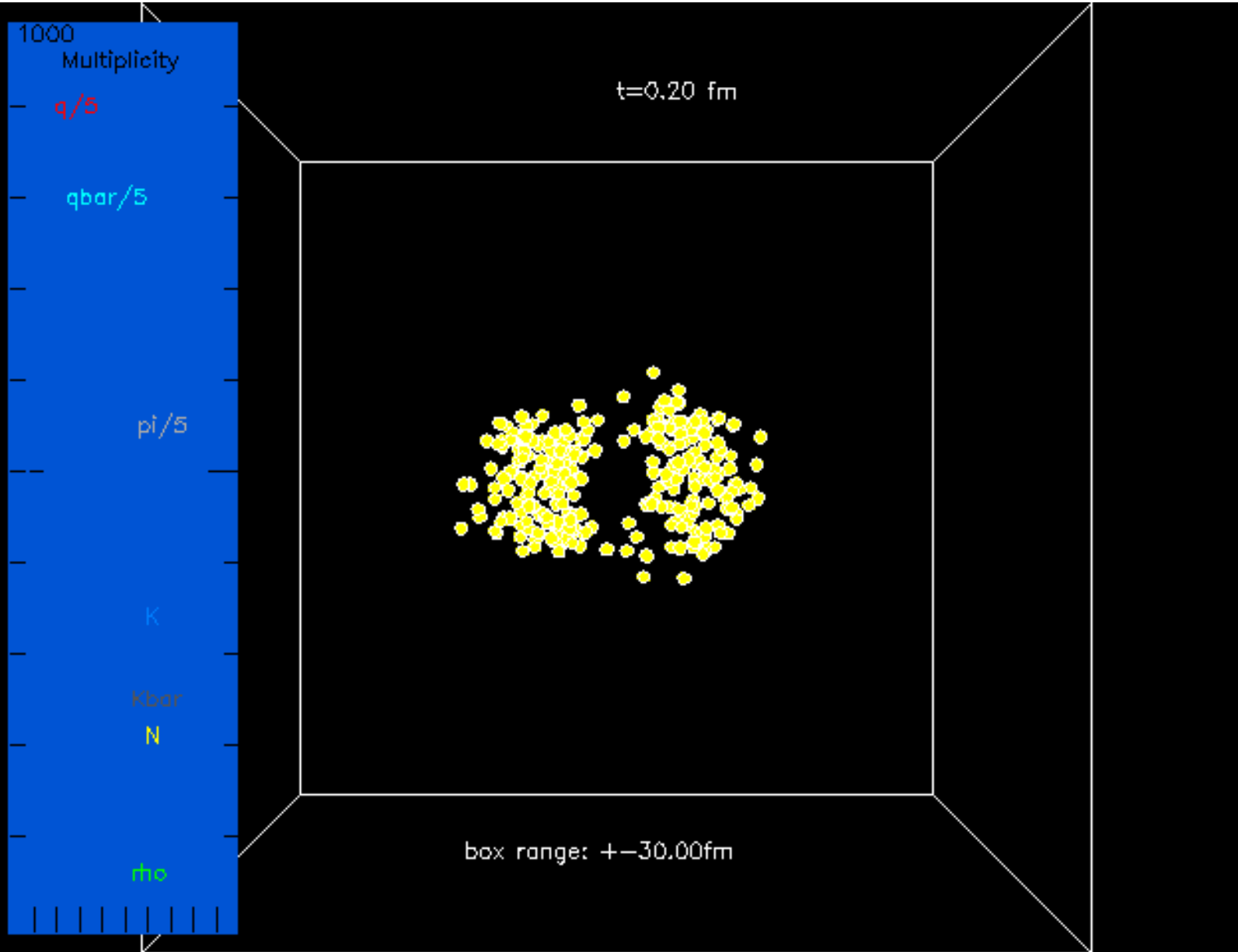
**including**

- 3 parton populations  
(*e.g. freezeout partons, active partons & all partons*)
- $v_n$  versus **Ncoll** (*number of collisions suffered by a parton*)
- $v_n$  versus time

*Most results shown here are obtained with AMPT (string melting version); some obtained with MPC (elastic version of the parton cascade).*

*We only study the parton stage here.*

# A $b=10$ fm Au+Au event at 200A GeV from String Melting AMPT



Beam axes

*60fm-long box;*

*shows only  
formed particles.*

*1<sup>st</sup> frame:  
right after the  
primary collision,  
only spectator  
nucleons are  
formed.*

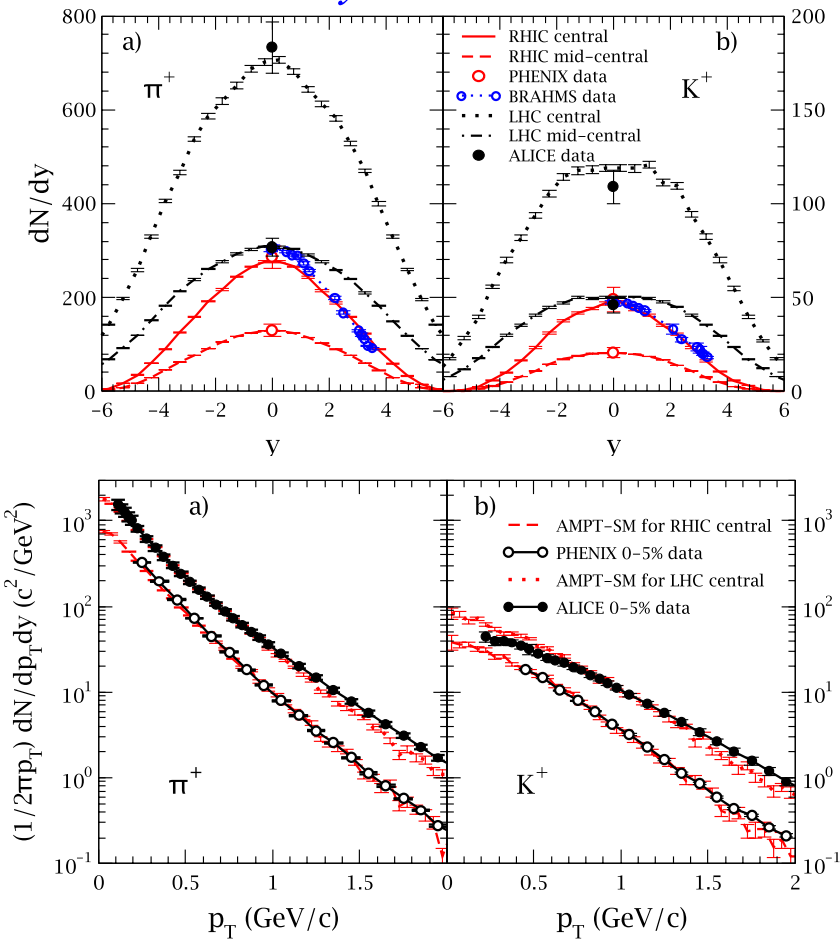
Particle # vs time



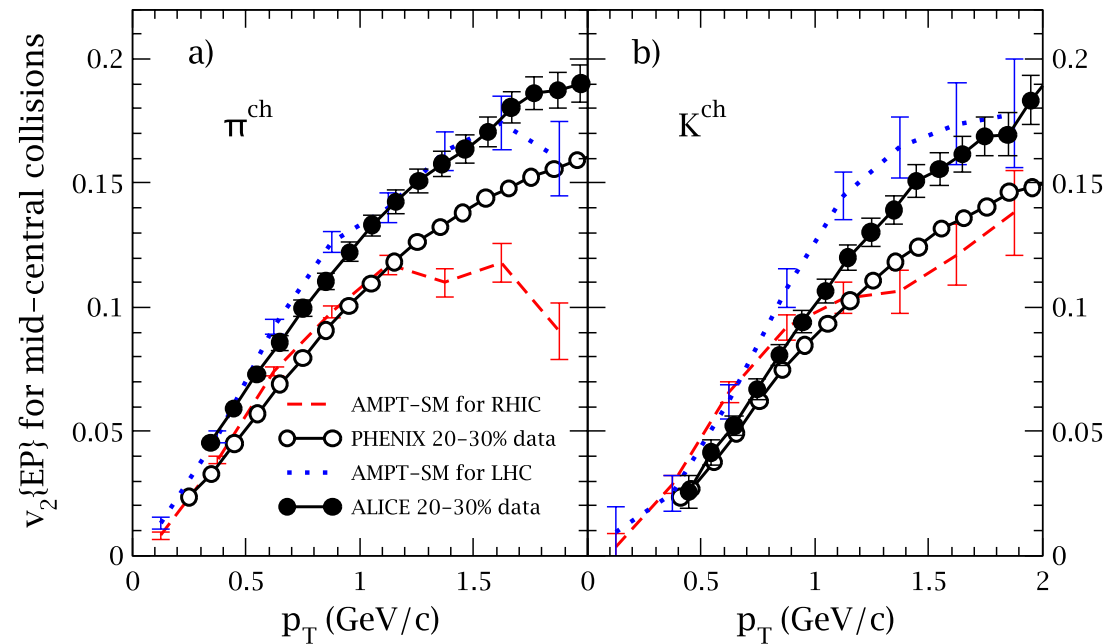
# Constraining Parameters of the String Melting AMPT

Same parameters for Au+Au as in [ZWL PRC \(2014\)](#),  
 which described low-pt ( $<2\text{GeV}/c$ )  $\pi$  & K data on  $dN/dy$ ,  $p_T$  spectra &  $v_2$   
 in central & mid-central events of 200A GeV Au+Au.

$dN/dy$  of  $\pi$  & K



$v_2$  of  $\pi$  & K (AuAu@200A GeV  $b=7.3\text{fm}$ )

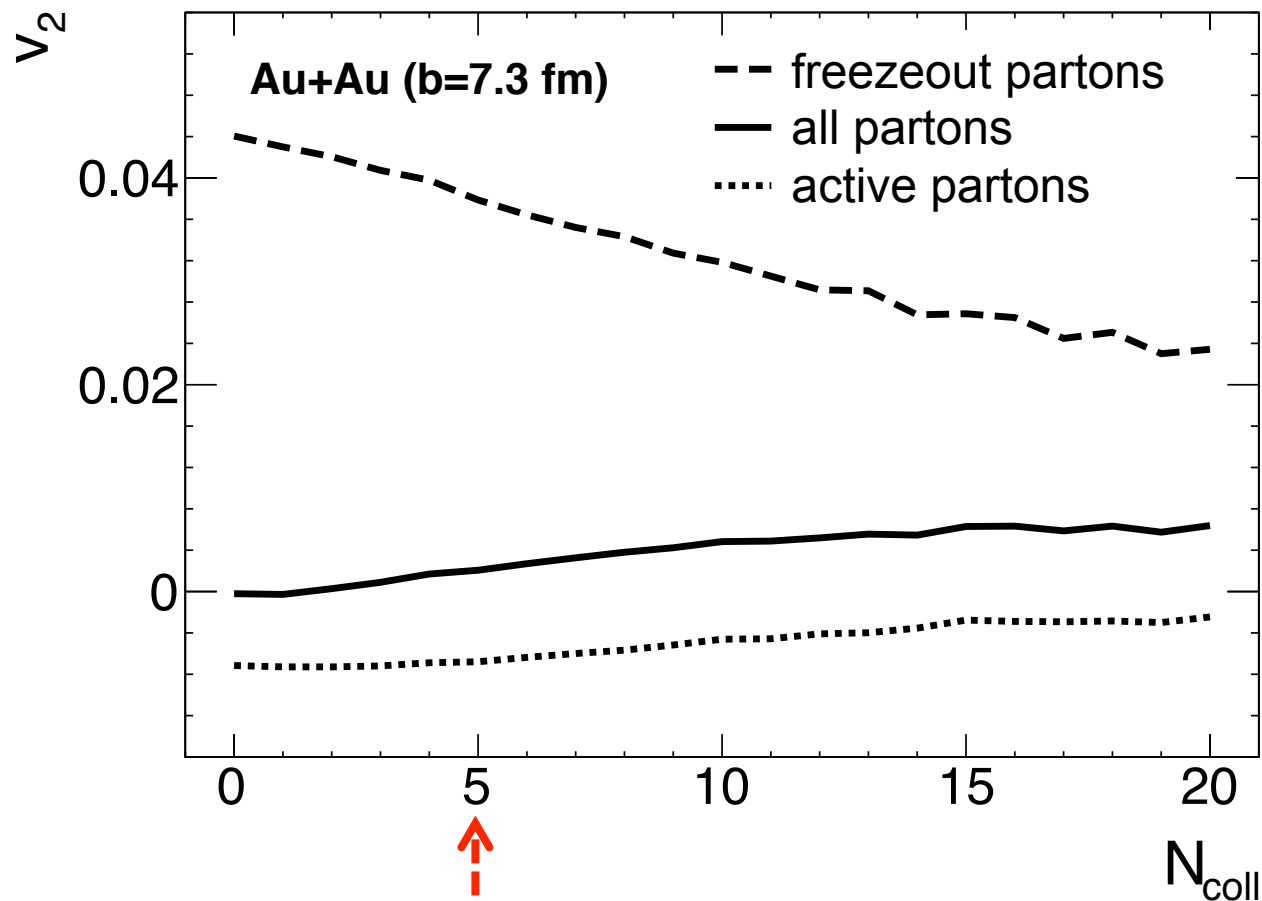


Parton cross section  $\sigma = 3\text{mb}$

$p_T$  spectra of  $\pi$  & K (in central collisions)

# Results: $v_2$ versus collision # of each parton

**N<sub>coll</sub>**: *number of collisions suffered by a parton*



3 parton populations at a given  $N_{\text{coll}}$ :

*freezeout partons:* freeze out after exactly  $N_{\text{coll}}$  collisions;

*active partons:* will collide again, freeze out after  $>N_{\text{coll}}$  collisions;

*all partons:* sum of the above two populations

(i.e. all formed partons that have survived  $N_{\text{coll}}$  collisions).

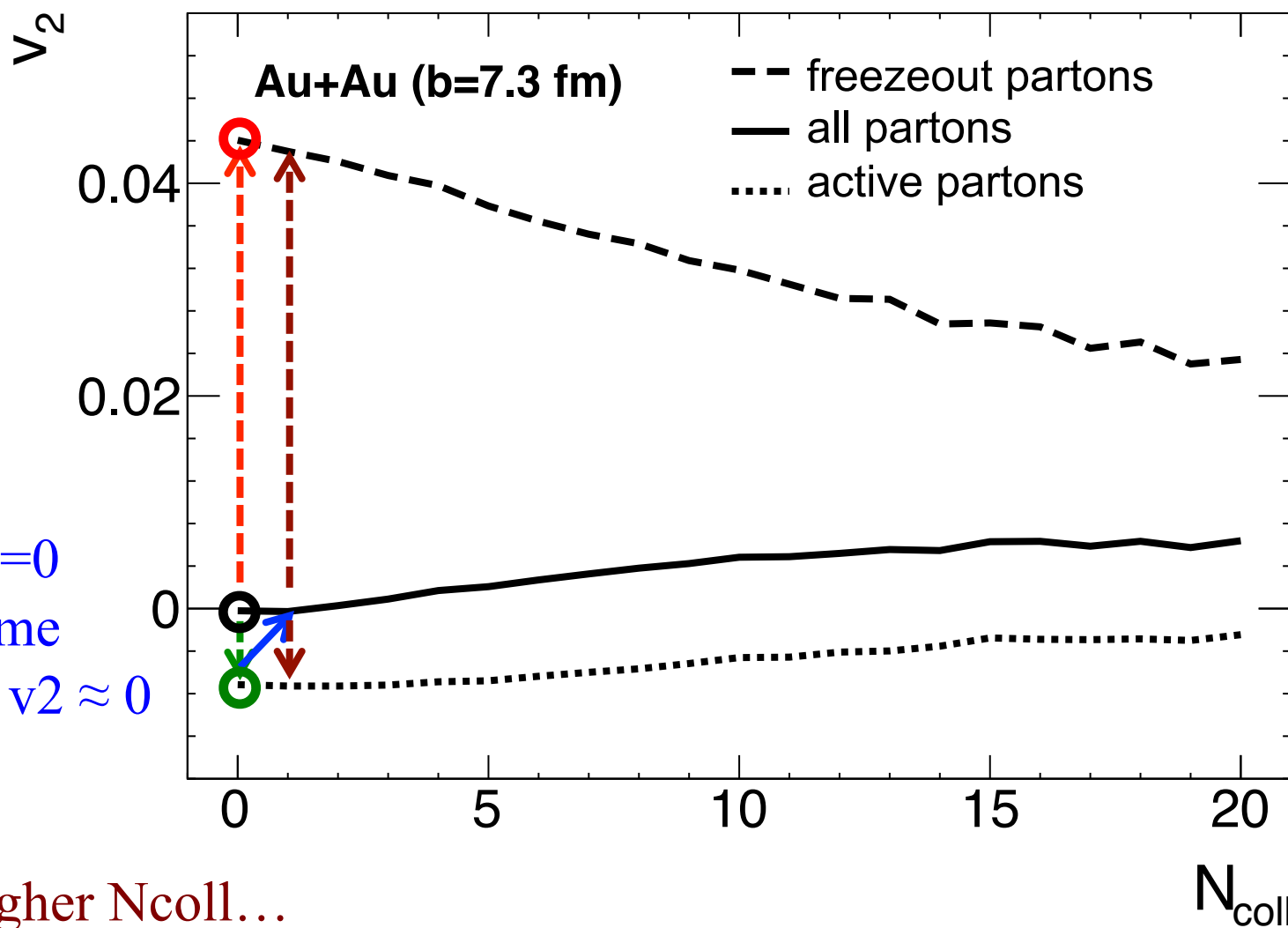
# Results: $v_2$ versus collision # of each parton

**At  $N_{\text{coll}}=0$ :**

all partons:  $v_2=0$  by symmetry;

escaped/freezeout:  $v_2 \approx 4.5\%$ ;

Active:  $v_2 < 0$ .



**At  $N_{\text{coll}}=1$ :**

active partons at  $N_{\text{coll}}=0$

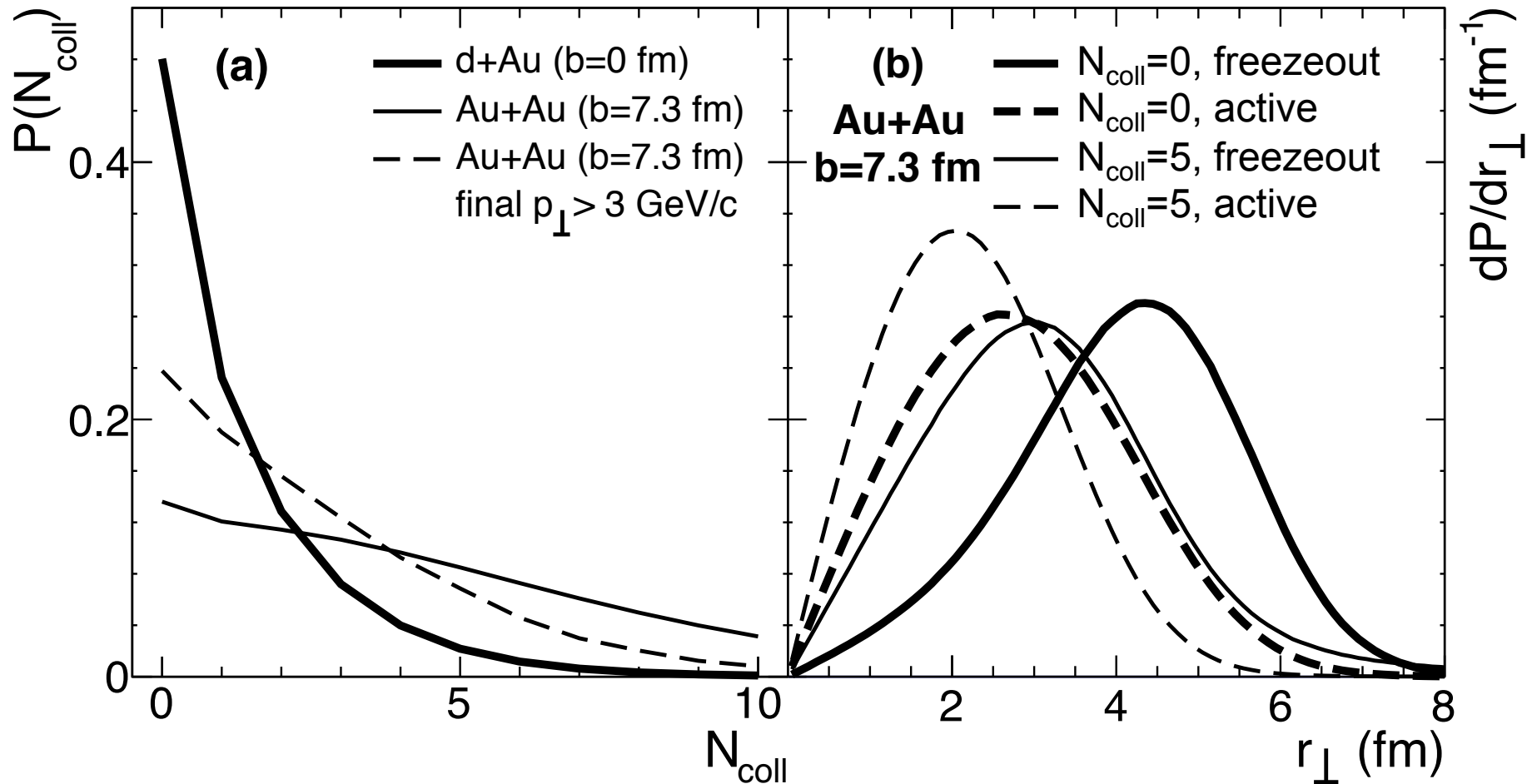
collide once each, become

all partons at  $N_{\text{coll}}=1$ :  $v_2 \approx 0$

this repeats itself to higher  $N_{\text{coll}}$ ...

$N_{\text{coll}}$

# Results: freezeout vs collision # of a parton



$\langle N_{\text{coll}} \rangle = 4.6$  for Au+Au  
 $= 1.2$  for d+Au.

Freezeout in the outer region  
 (~surface emission, but not from a sharp surface);

freezeout region moves in with  $N_{\text{coll}}$ .

# Results: anisotropic particle escape

**At  $N_{\text{coll}}=0$ :**

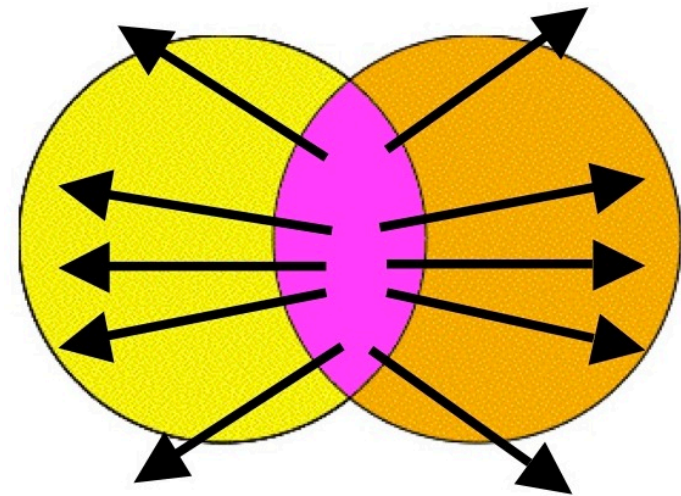
Escaped:  $v_2 \approx 4.5\%$ ,

purely due to

**anisotropic escape probability**

(response to geometrical shape only,  
no effect from collective flow)

In simplified picture of elliptic flow



**At  $N_{\text{coll}} \geq 1$ :**

Escaped:  $v_2 > 0$

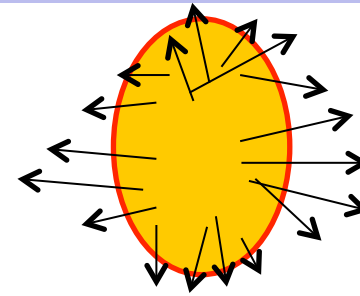
due to

the above **anisotropic escape probability**

**modified by collective flow** of all active partons

# Results: anisotropic particle escape

Final  $v_2$  is generated by interactions, which generate anisotropic collective flow and freezeout/escape from an anisotropic shape:



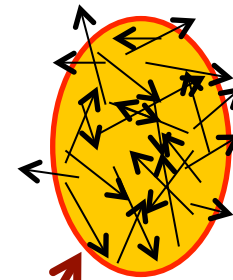
$f(\mathbf{x}, \mathbf{p}, t)$

Let's view  $v_2$  as coming from 2 separate but compounding sources:

## 1) anisotropic escape probability

*we define this term as:*

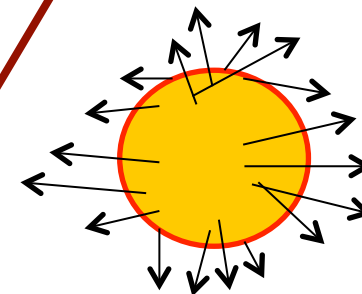
effect due to spatial anisotropy only  
if there were 0 collective flow



$f(\mathbf{x}, \text{iso } \mathbf{p}, t)$

## 2) collective flow of all active partons

effect from anisotropic collective flow only  
if there were no spatial anisotropy



$f(\text{iso } \mathbf{x}, \mathbf{p}, t)$

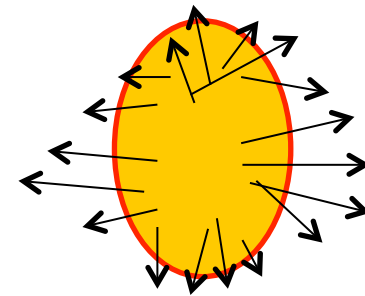
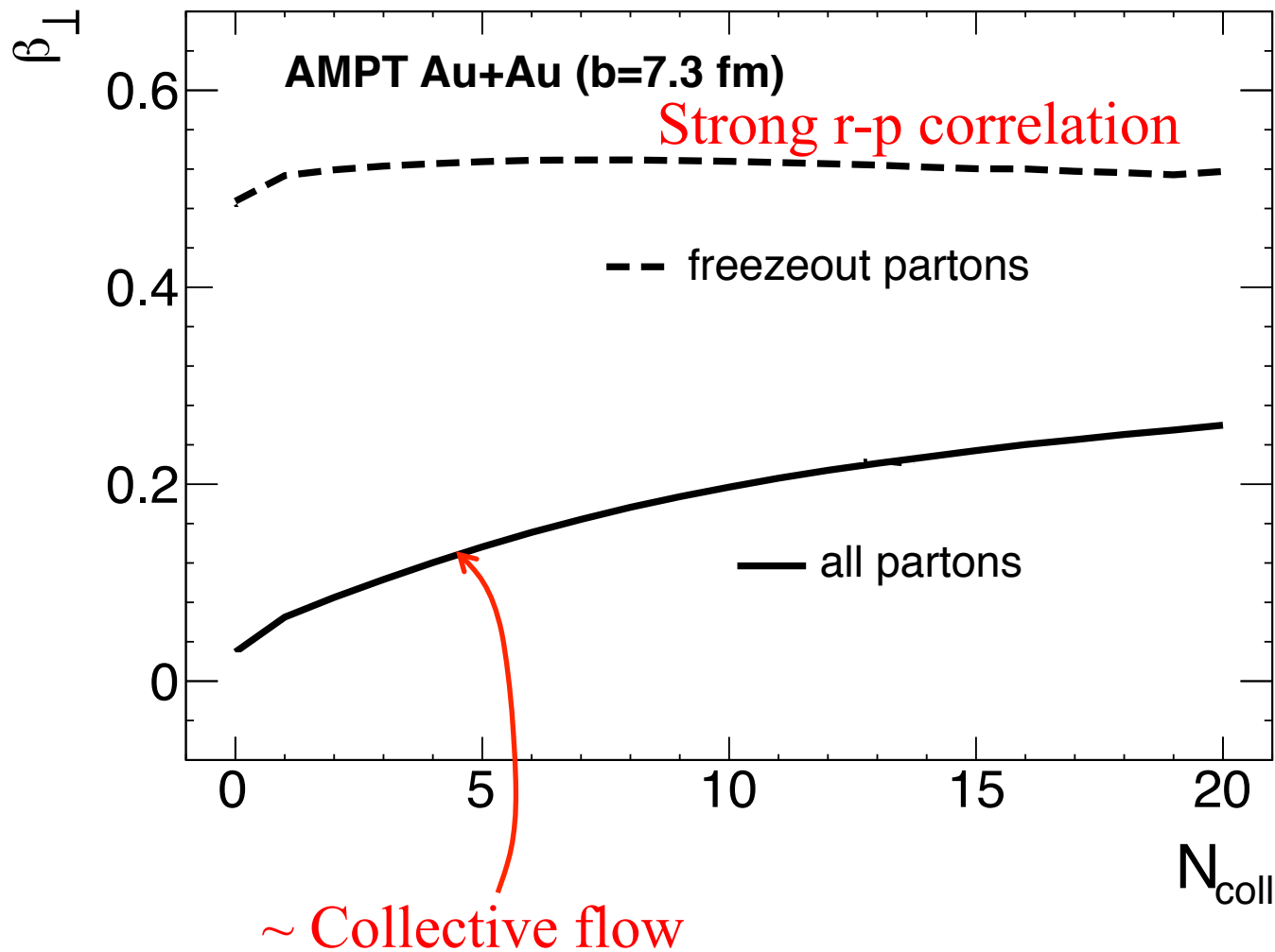
They are coupled in the actual evolution, so we design a **Random Test** to estimate 1):

$\phi$  is randomized (after each scattering) to destroy collective flow

# Results: space-momentum correlation vs collective flow

$$\beta_{\perp} = \left\langle \frac{\vec{r}_{\perp} \cdot \vec{p}}{r_{\perp} p} \right\rangle$$

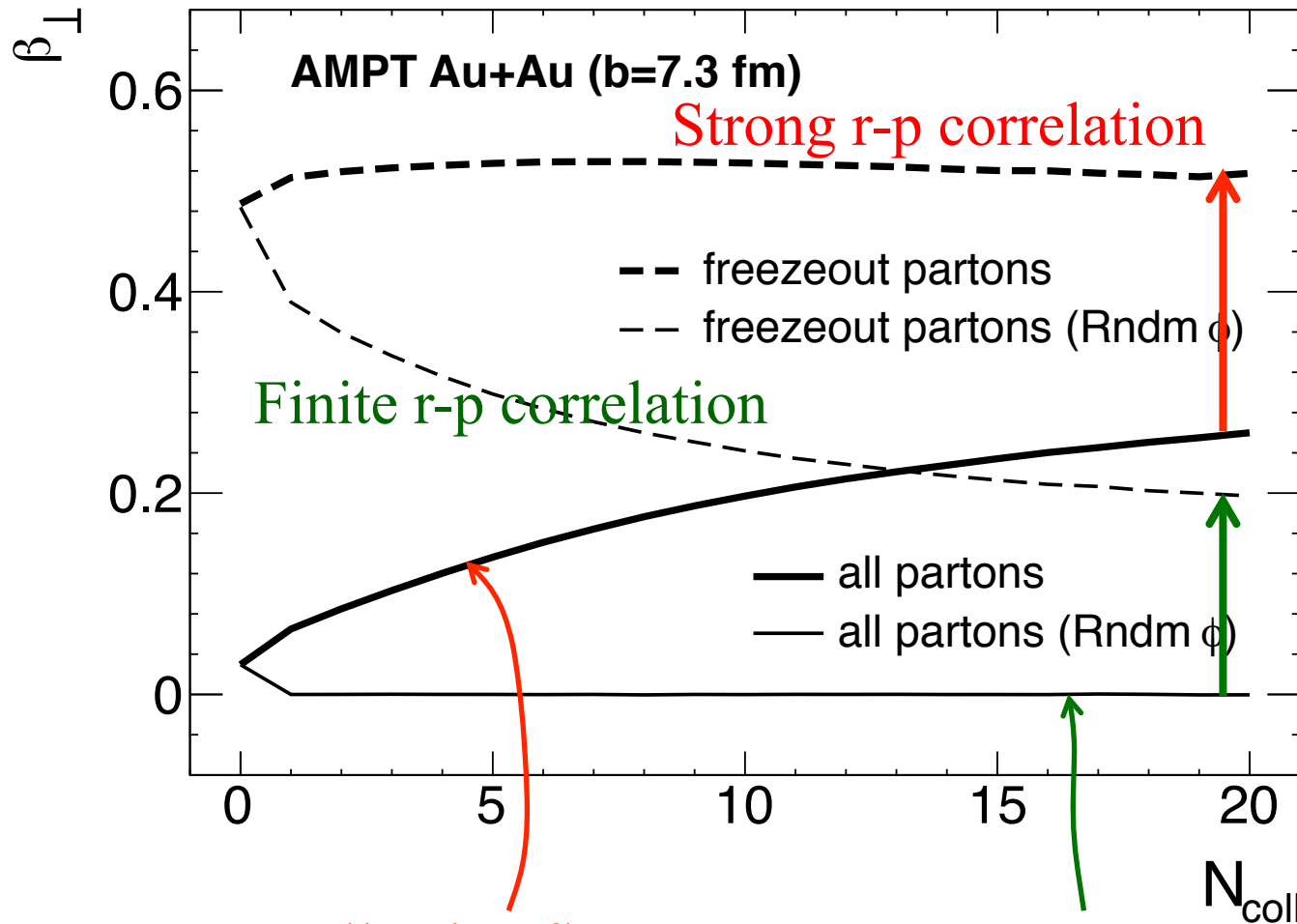
reflects space-momentum correlation,  
~ transverse flow velocity



# Results: space-momentum correlation vs collective flow

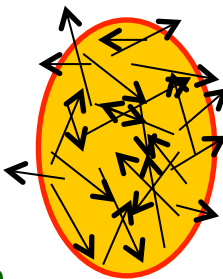
$$\beta_{\perp} = \left\langle \frac{\vec{r}_{\perp} \cdot \vec{p}}{r_{\perp} p} \right\rangle$$

reflects space-momentum correlation,  
 $\sim$  transverse flow velocity



Addition increase  
 from escape mechanism

r-p correlation  
 purely from  
 escape mechanism,  
*not from collective flow*



Collective flow

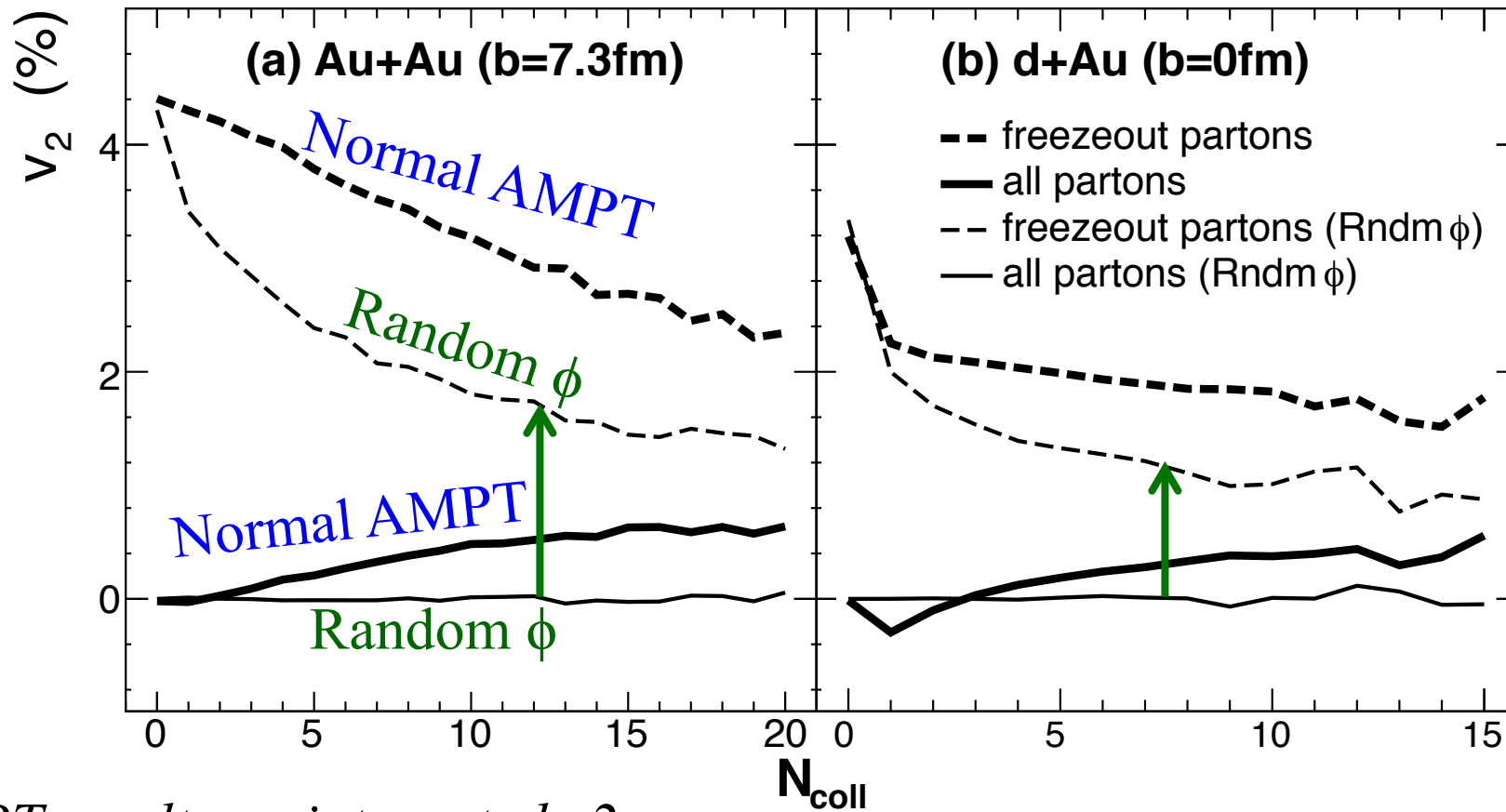
Random Test:  
 collective flow is destroyed



# Results: contribution of escape mechanism to final $v_2$

$v_2$  from **Random Test**:

purely from escape mechanism, *at 0 collective flow*



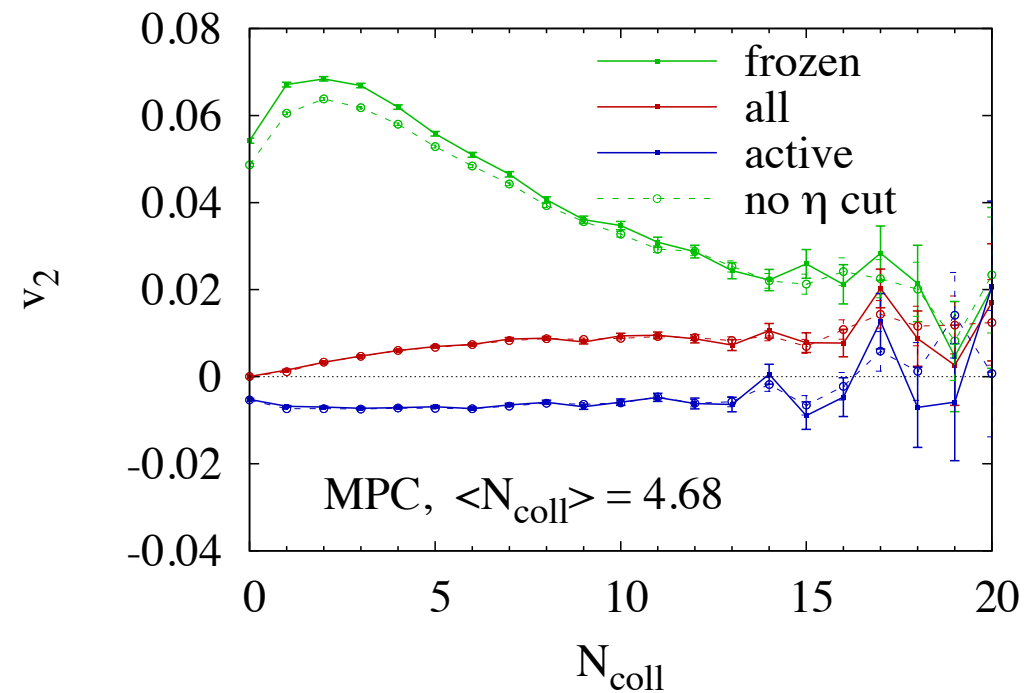
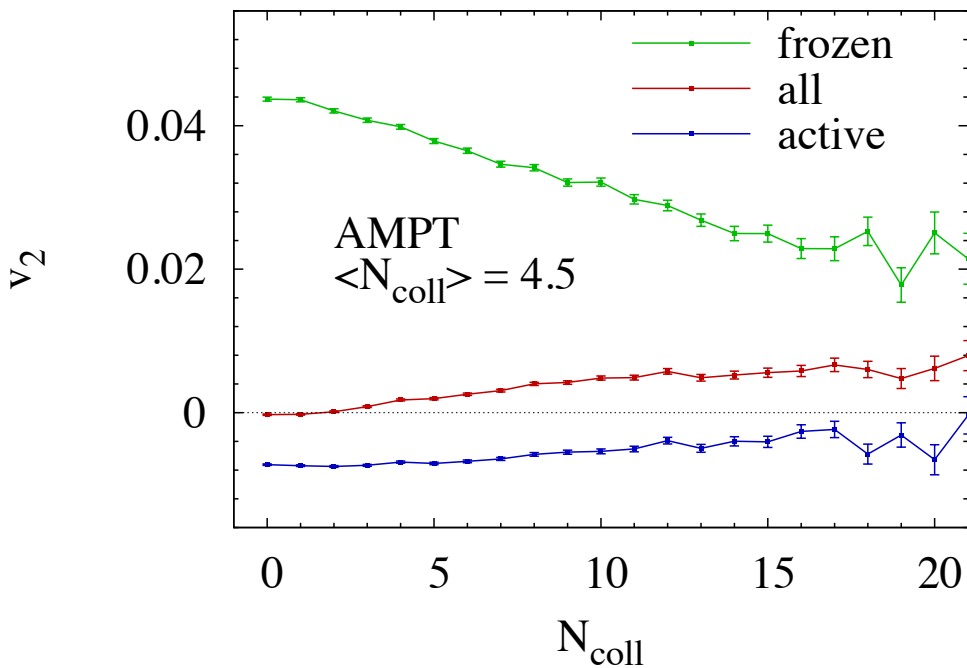
*AMPT results on integrated  $v_2$ :*

	Normal	Random $\phi$	Ratio:	
Au+Au	3.9%	2.7%	69%	
d+Au	2.7%	2.5%	93%	17/25

*~contribution from pure escape*

# Results: this is a general feature of transport models

MPC gives essentially the same results as AMPT at similar  $\langle N_{\text{coll}} \rangle$  despite differences in *parton initial condition (number, density profile,  $P_T$  spectrum),  $d\sigma/dt$ , formation time, parton-subdivision.*

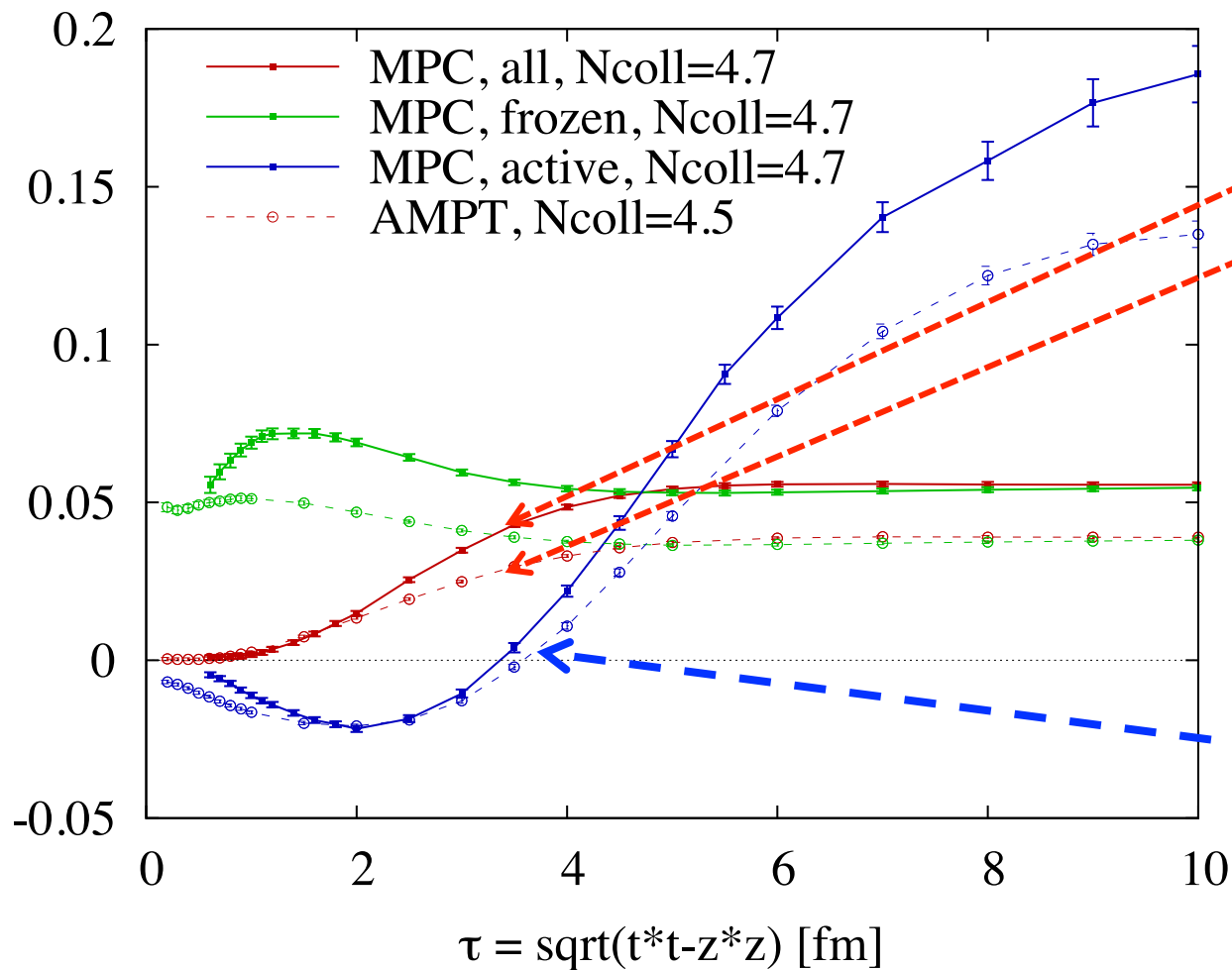


# Results: this is a general feature of transport models

MPC gives essentially the same results as AMPT at similar  $\langle N_{\text{coll}} \rangle$

Time-dependence of 3 parton populations:

*frozen partons, active partons, & all partons*

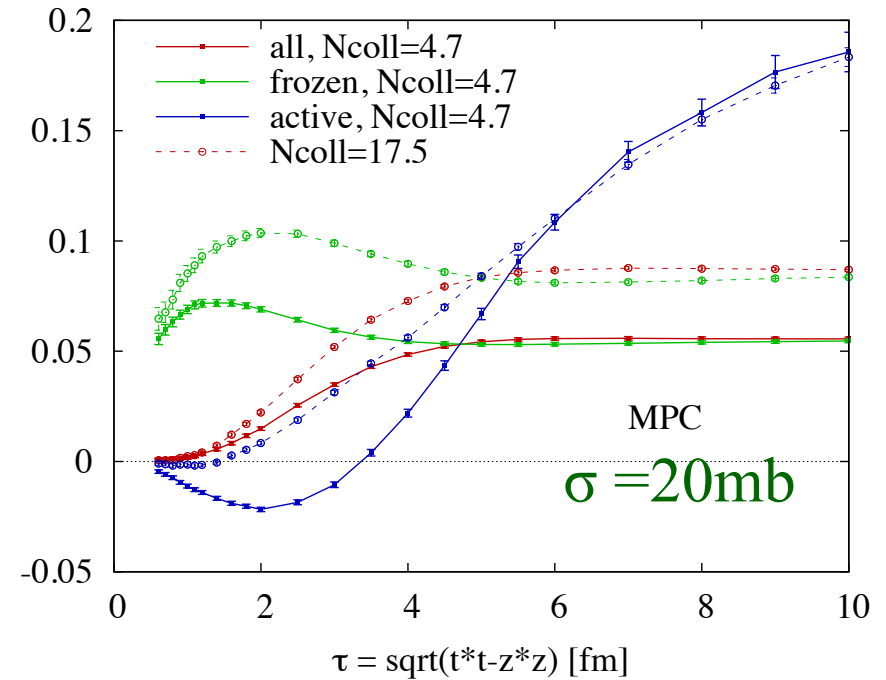
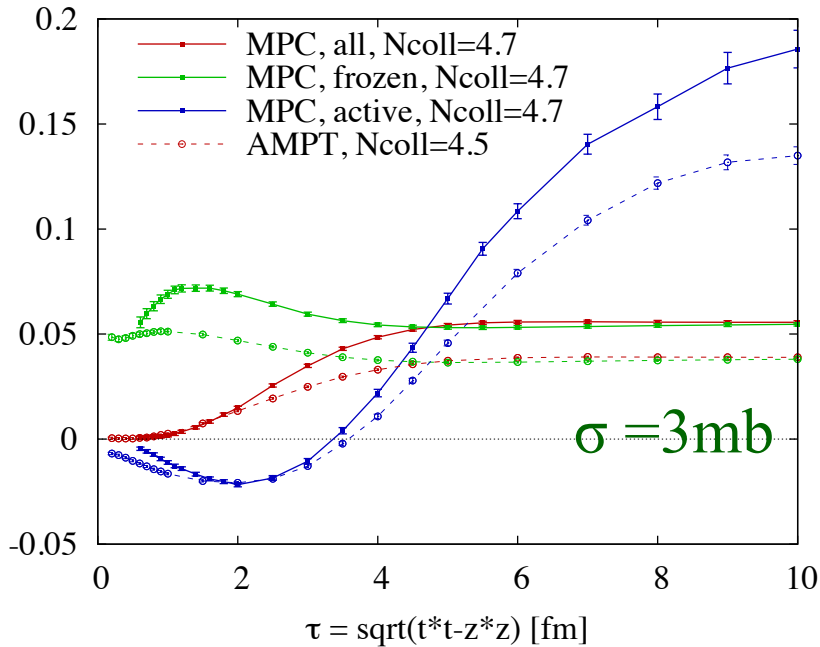


*$v_2$  of all partons  
gradually builds up  
within  $\sim 4$  fm/c*

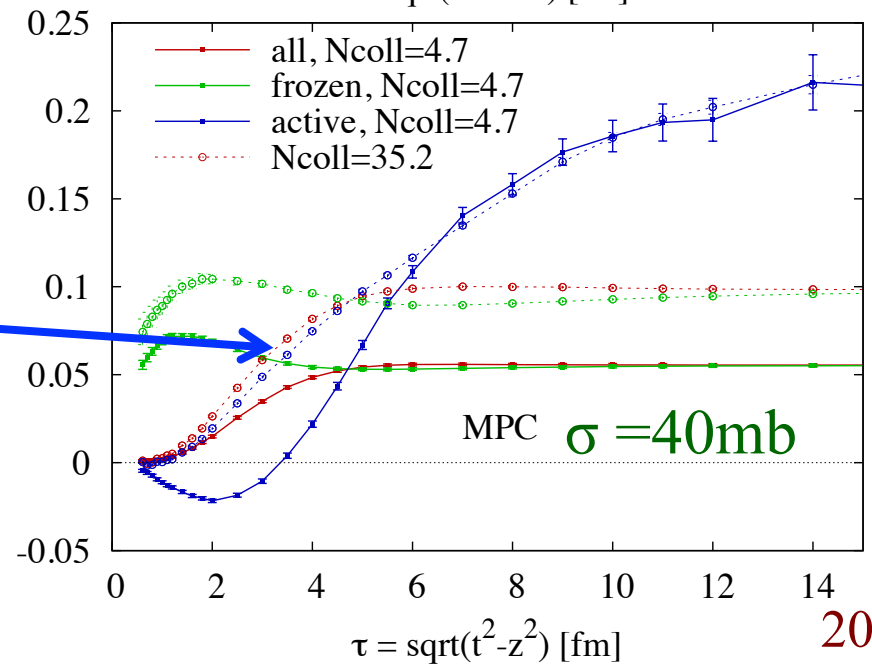
*$v_2$  of active partons  
is small or negative  
during most of  $v_2$  build-up*

# Results: Anisotropic Particle Escape versus Hydrodynamic Flow

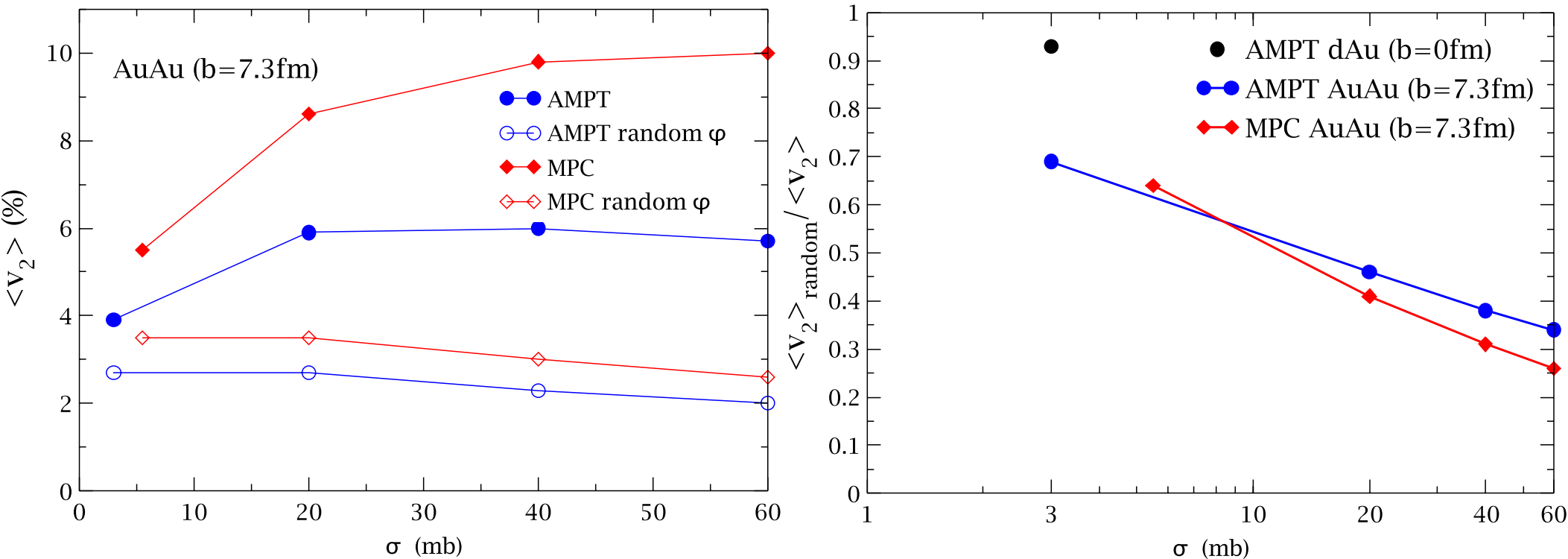
When will **hydrodynamic flow** dominate?



At very high  $\langle N_{\text{coll}} \rangle$  or  $\sigma$  :  
 $v_2$  of active partons  
 follows the total  $v_2$  closely  
 during the early  $v_2$  build-up



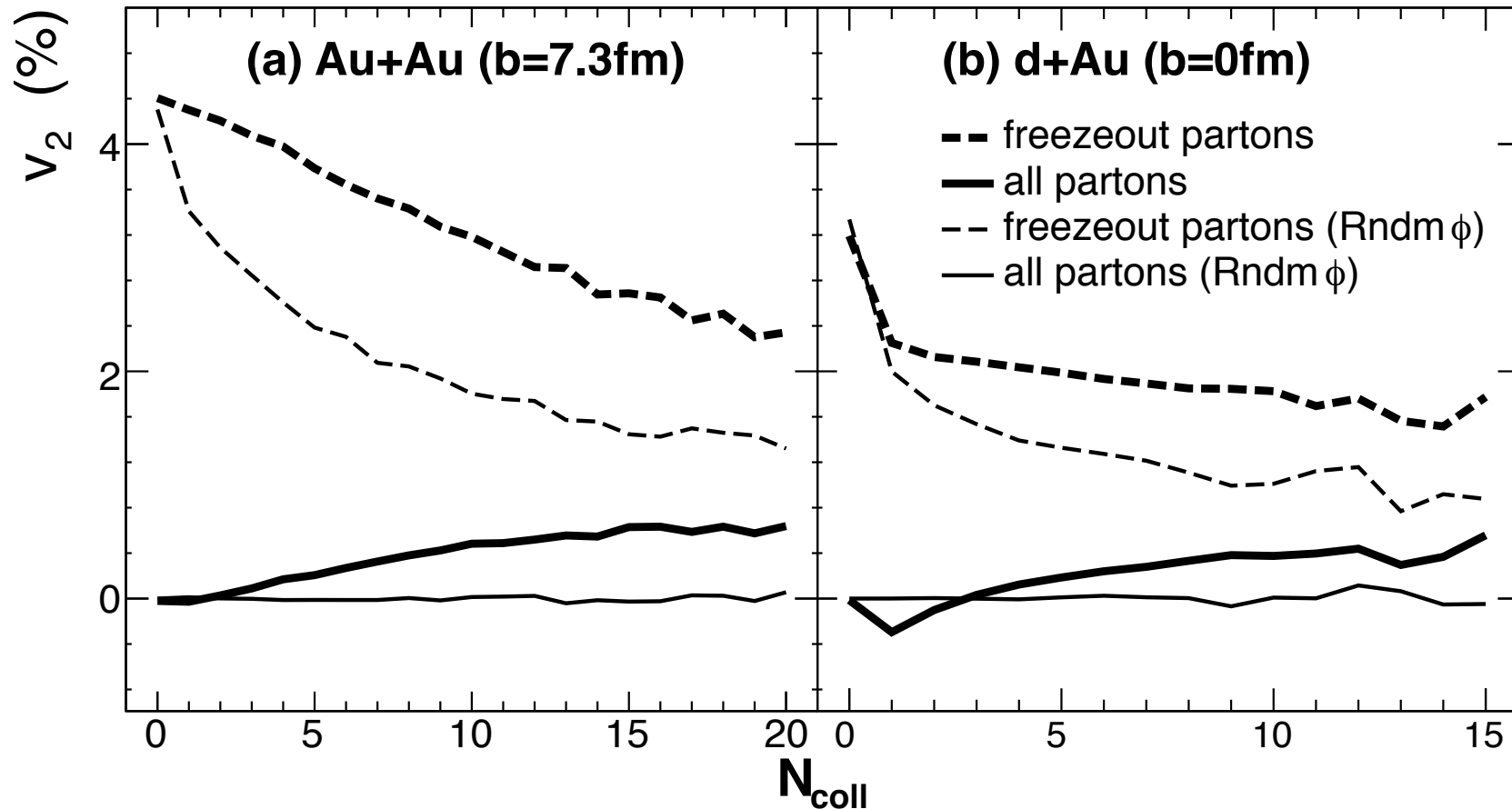
# Results: Anisotropic Particle Escape versus Hydrodynamic Flow



Escape mechanism is dominant  
for small system  $v_2$   
& maybe even for semi-central AuAu at RHIC

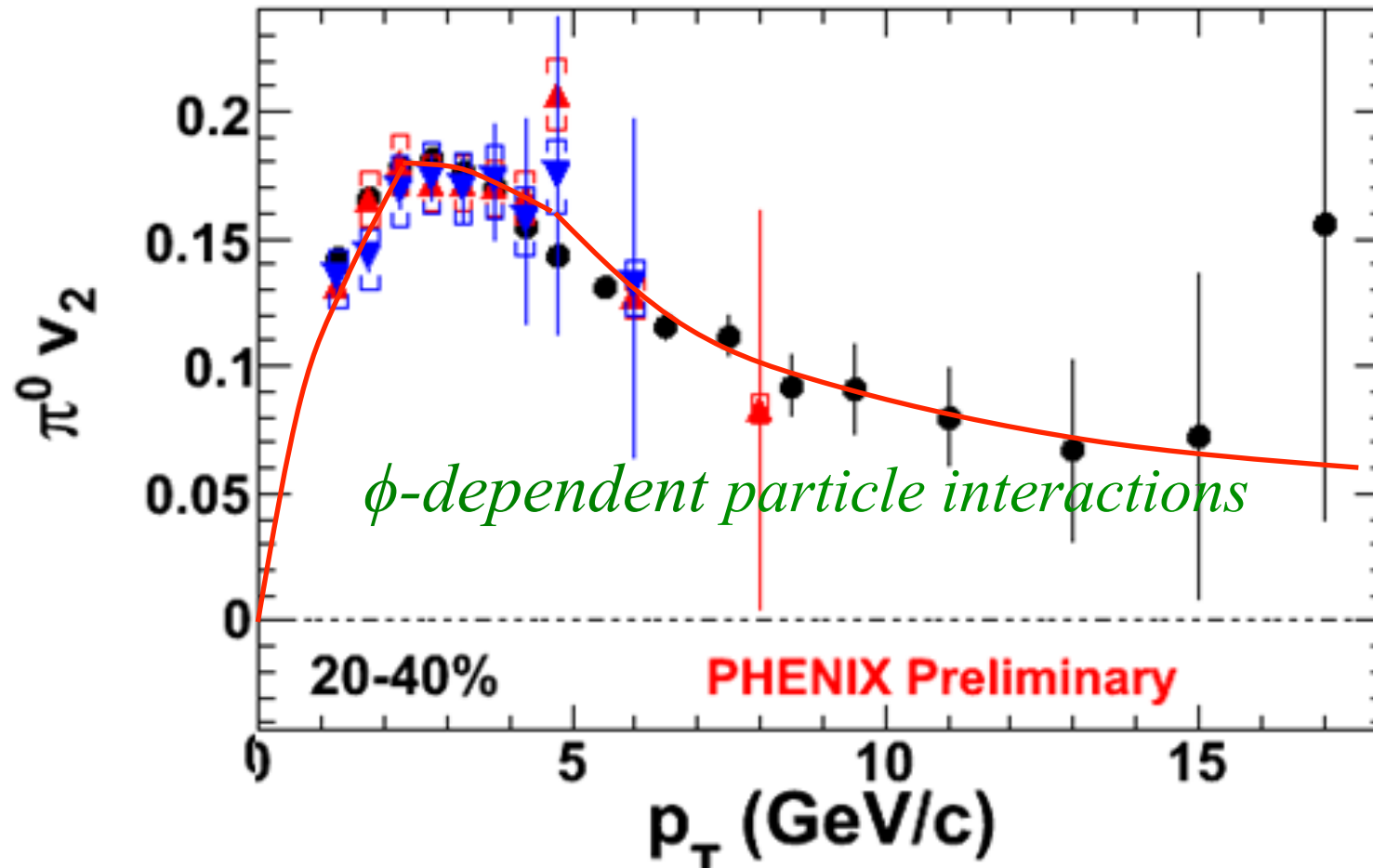
Hydrodynamic collective flow  
will dominate at very high  $\sigma$  or  $\langle N_{\text{coll}} \rangle$

# Potential Consequences



- **Explains similar anisotropic flows**  
**observed in small systems and in large systems:**  
since both are dominated by the escape mechanism  
( $\phi$ -dependent interactions & escape probability)

# Potential Consequences



- **Main reason for  $v_n$  at low  $P_T$  and high  $P_T$  are qualitatively the same** since both are dominated by the escape mechanism ( *$\phi$ -dependent particle interactions including energy-loss*)

# Potential Consequences

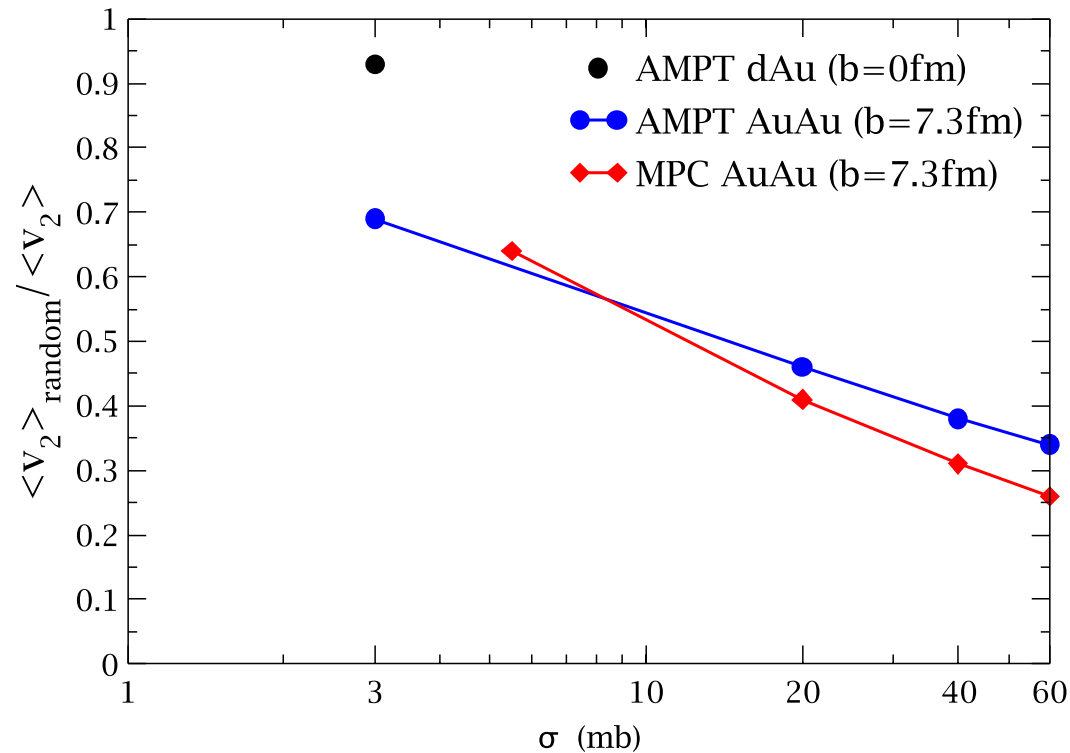
## Hydrodynamics

needs to include the escape mechanism  
(negative part of Cooper-Frye?  
continuous emission?)

will affect sQGP properties extracted  
by comparing  $v_n$  with hydrodynamics

Hydrodynamics can dominate  $v_n$   
at very high interaction strength  
or collision number  $\langle N_{\text{coll}} \rangle$ :

is it the case for heavy ion collisions?





# Anisotropic Particle Escape versus **Hydrodynamic Flow**

**Hydrodynamics** describe vn data well.

**AMPT/transport** describes vn data well.

Are they essentially the same?

How can we differentiate them with experimental observables?

Can they be improved to converge toward each other?