

Where are the Icecube Neutrinos Beyond 2-3 PeV ?

Tom Weiler
Vanderbilt University
Nashville, TN

μ

Neutrinos carry three types of information:

- (1) Direction
- (2) Energy
- (3) Flavor

All three have interesting features in IceCube data

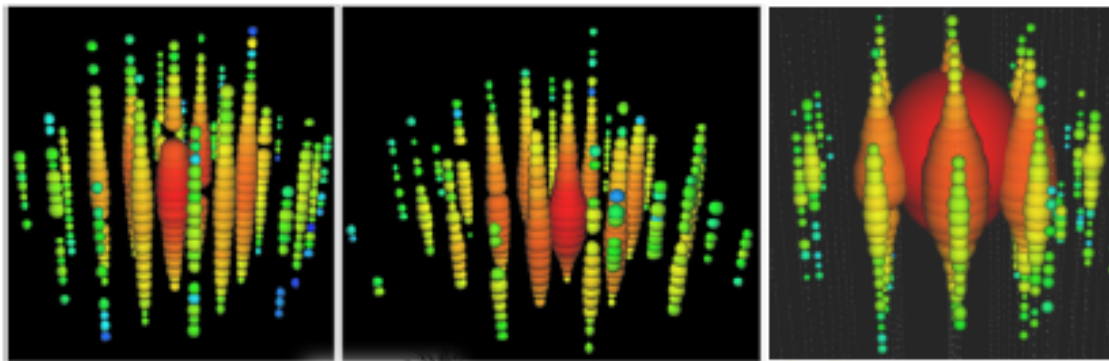
μ

Astro-Nu evidence update:

- first evidence for an extra-terrestrial flux shown at IPA2013 [*IceCube, Science 342 (2013)*]

[*IceCube, Phys.Rev.Lett. 113:101101 (2014)*]

- 3 yrs: 37 events in 988 days 5.7σ
- bkg. 8.4 ± 4.2 atm. μ and $6.6 + 5.9$ ν
- 4 yrs: 54 events $\sim 7\sigma$
- mostly ν_e CC and NC cascades



"Bert"
1.04 PeV
Aug. 2011



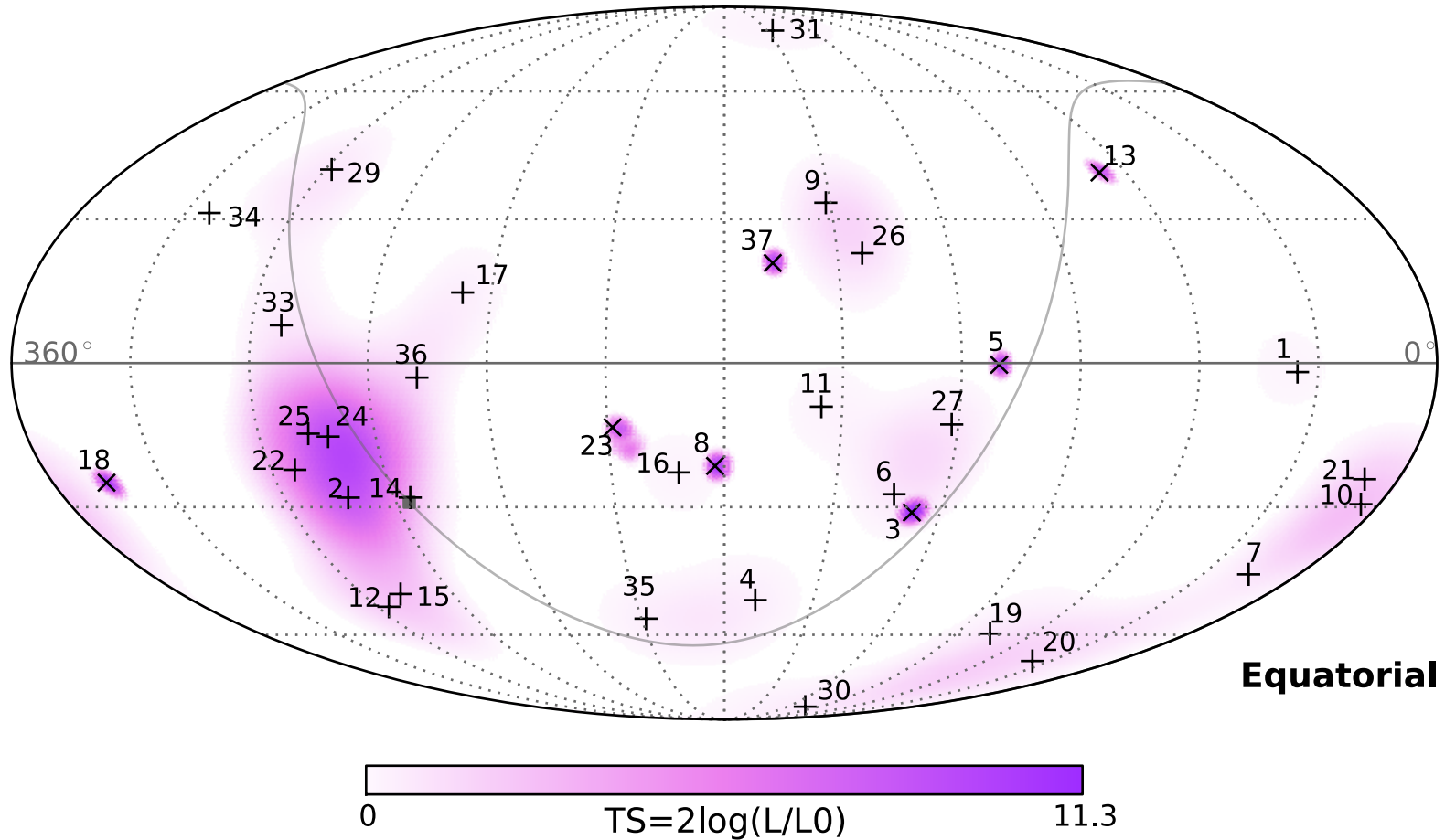
"Ernie"
1.14 PeV
Jan. 2012



"Big Bird"
2 PeV
Dec. 2012

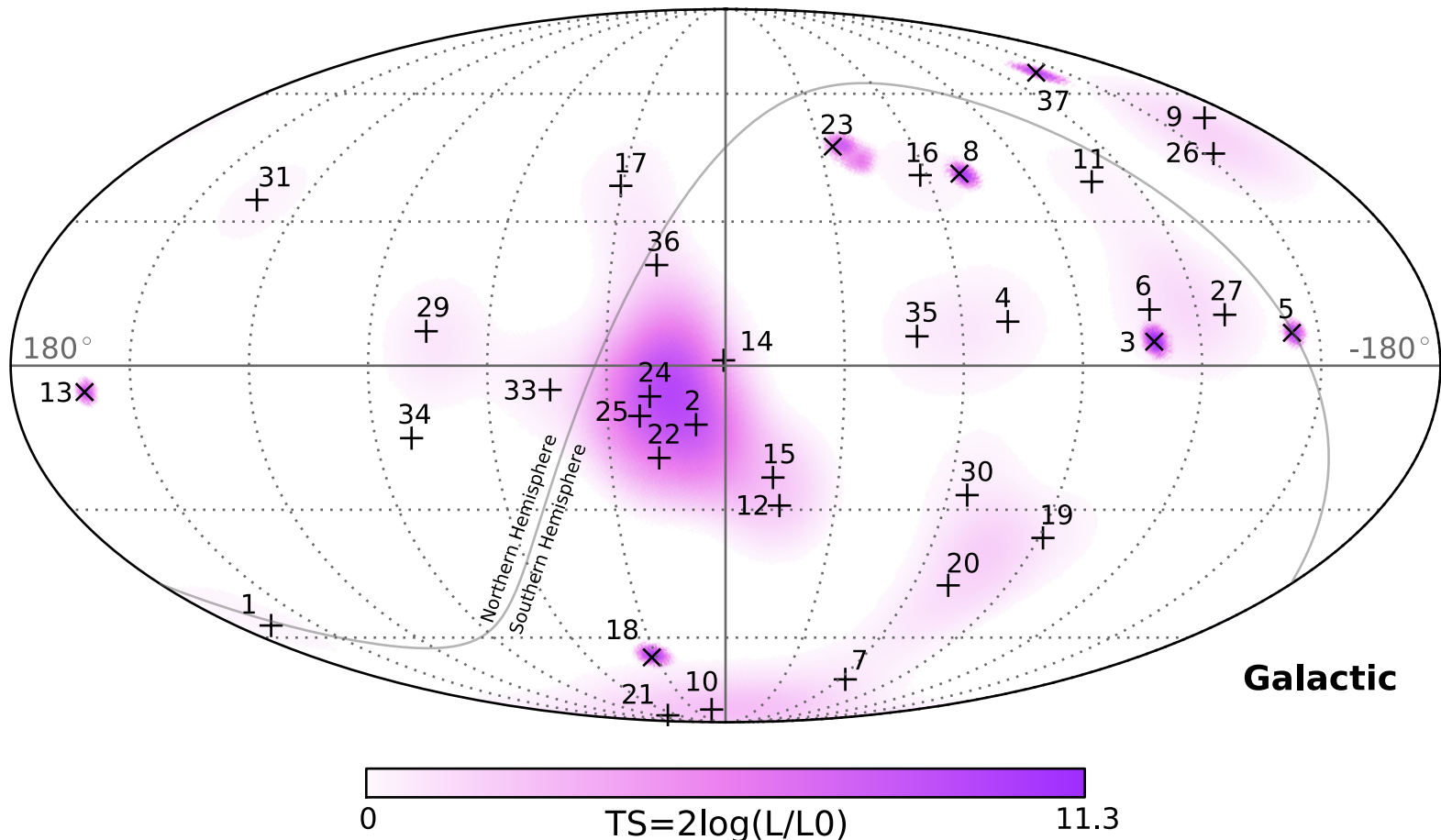
events near 1347 Dave

IceCube (Equatorial), 37 events/3 years (bkgd is $15+(2-10)$ events)

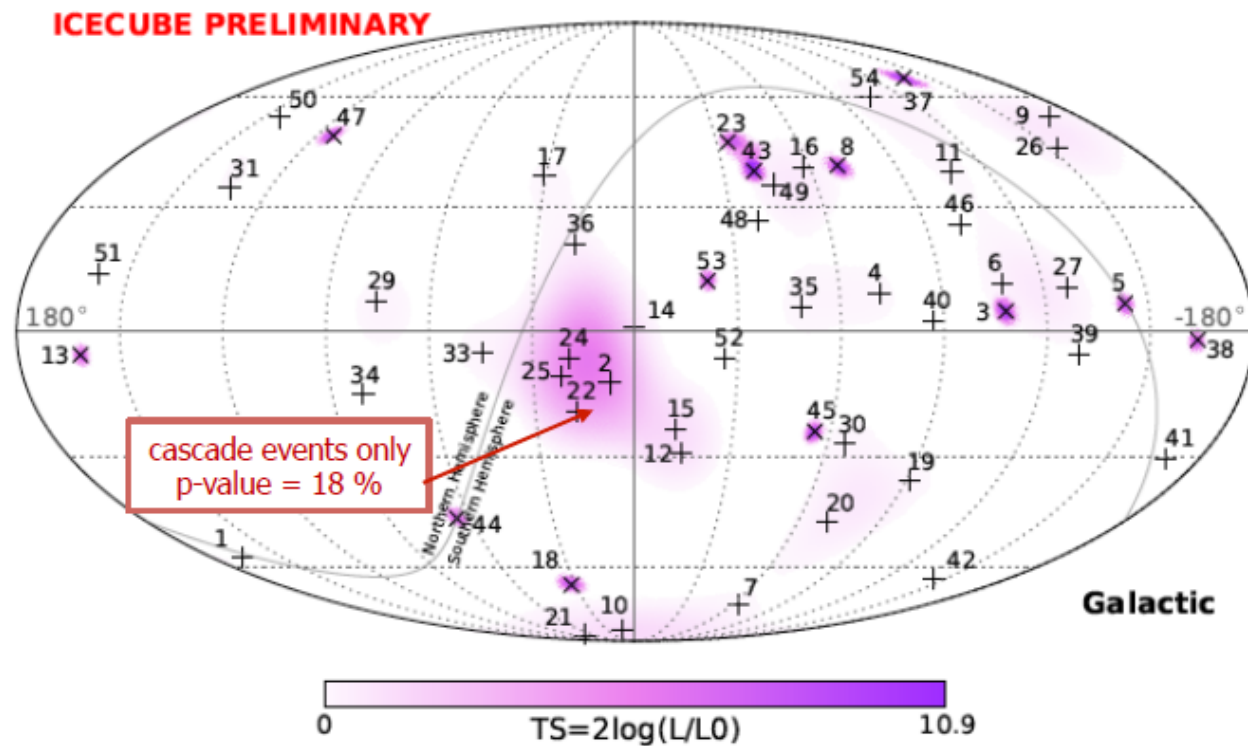


IceCube (Galactic), 37 events/3 years

Maybe Galactic Center shows a transient source (#'s 22, 24, 25) ?



4 years/ 54 events



no significant correlations – spatial or temporal

- too few events to identify sources

Energy - the “gap”, and then three events:

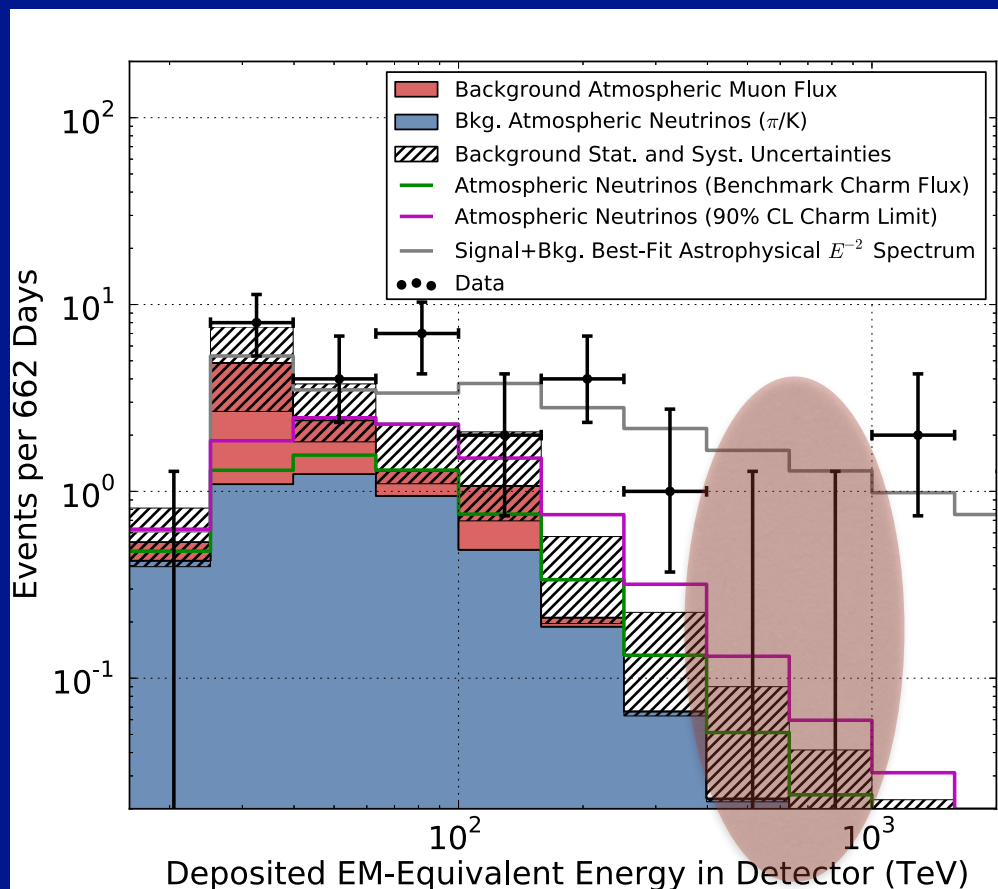


Figure 1: Distribution of the deposited energies of the observed events

Angle-Energy-Flavor Display:

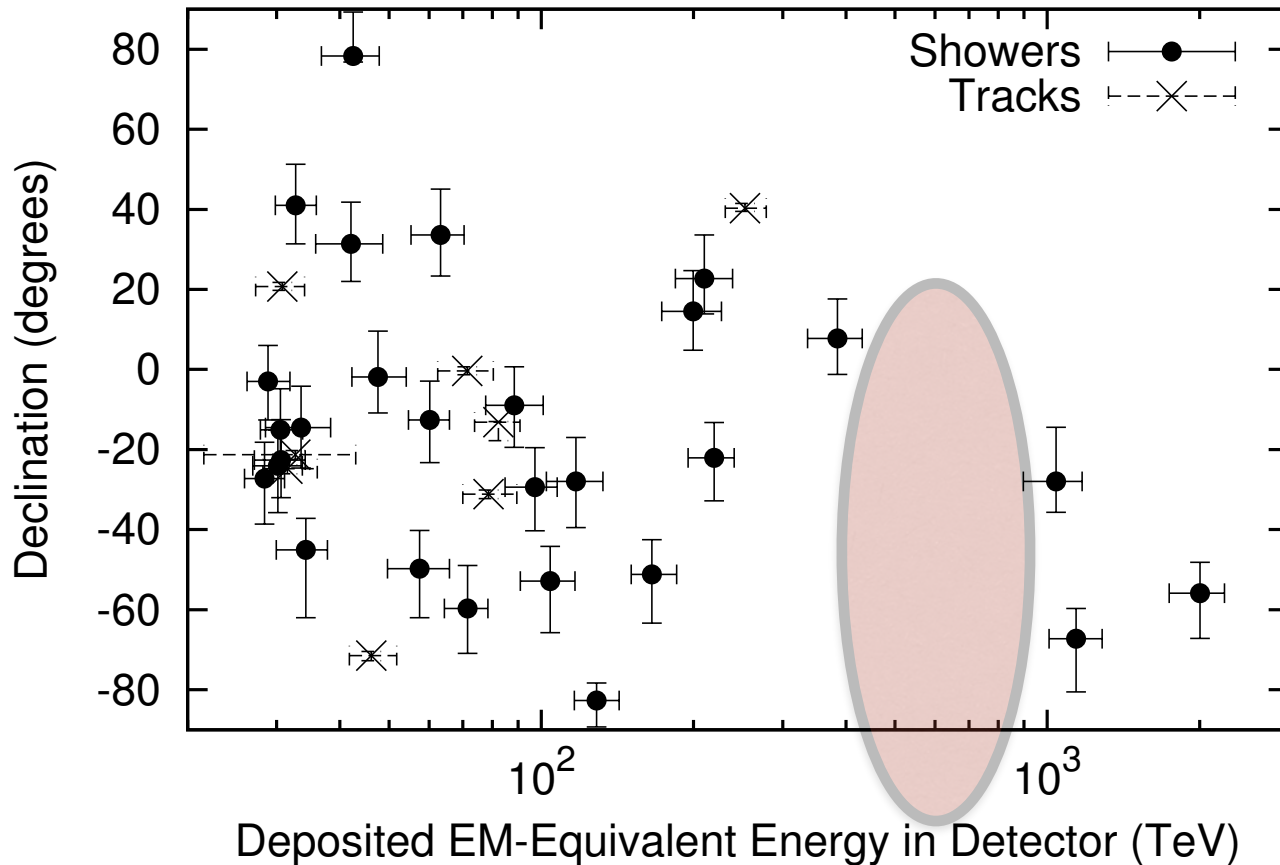
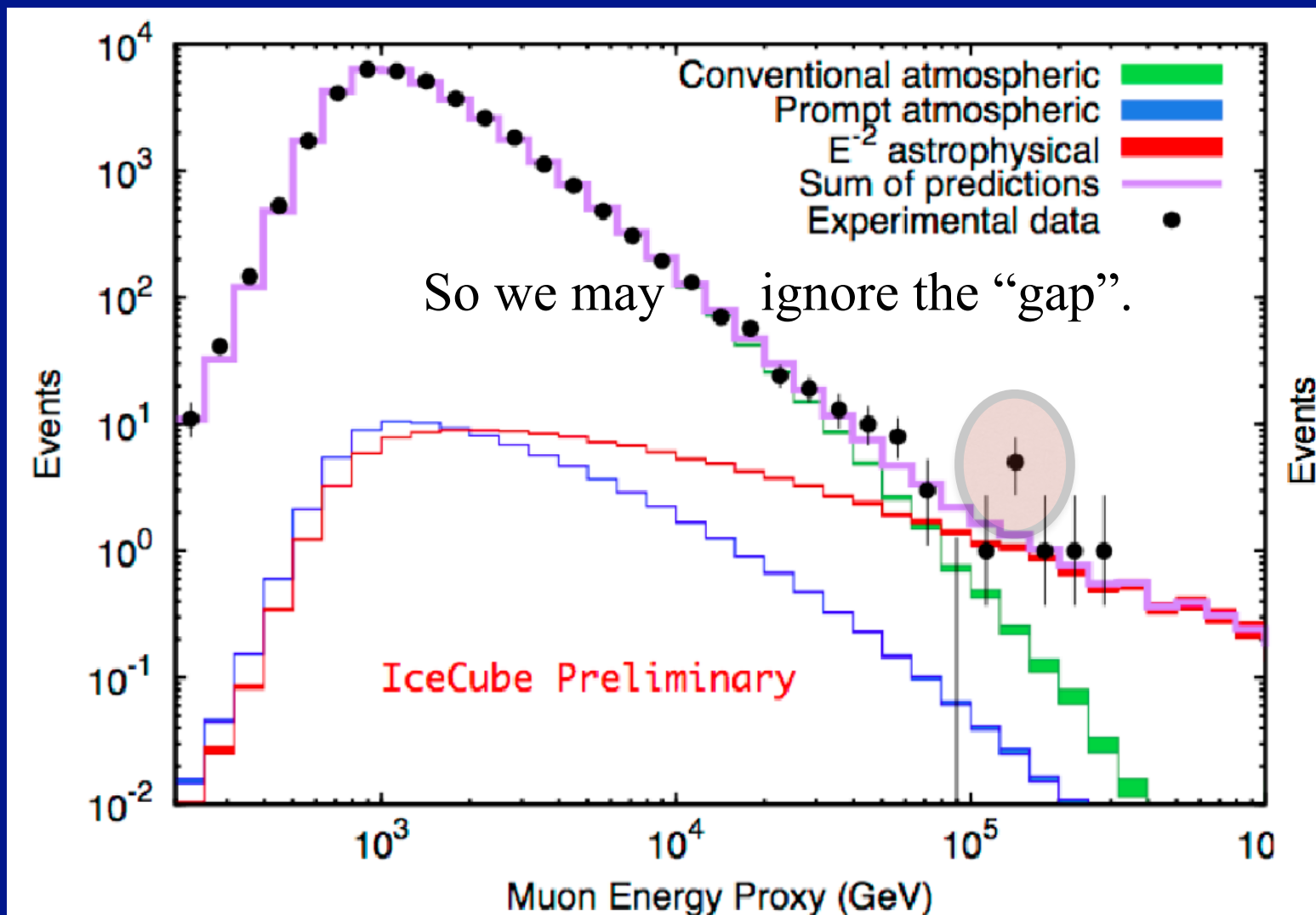
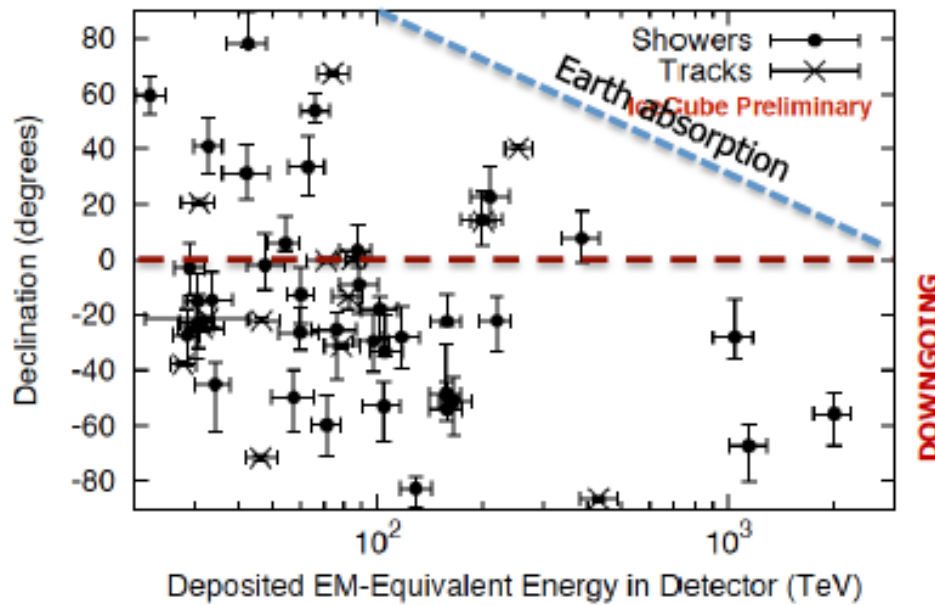


FIG. 1. Arrival angles and deposited energies of the events.

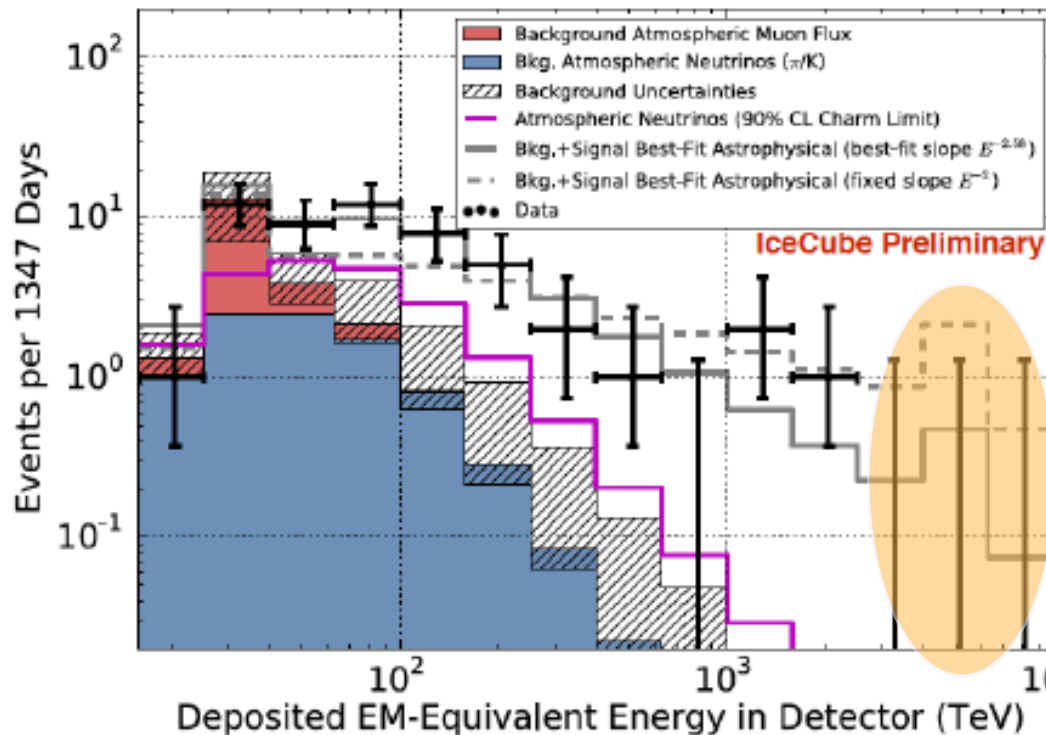
nu_mu rate vs. energy plot:



4 year
54 events



No events



and the “gap”
has narrowed

So where are the

i) continuum events
(Kistler's talk)

ii) Glashow resonance events
(Marfatia's talk)

Is there an Energy Cutoff at $\sim 2\text{PeV}$?

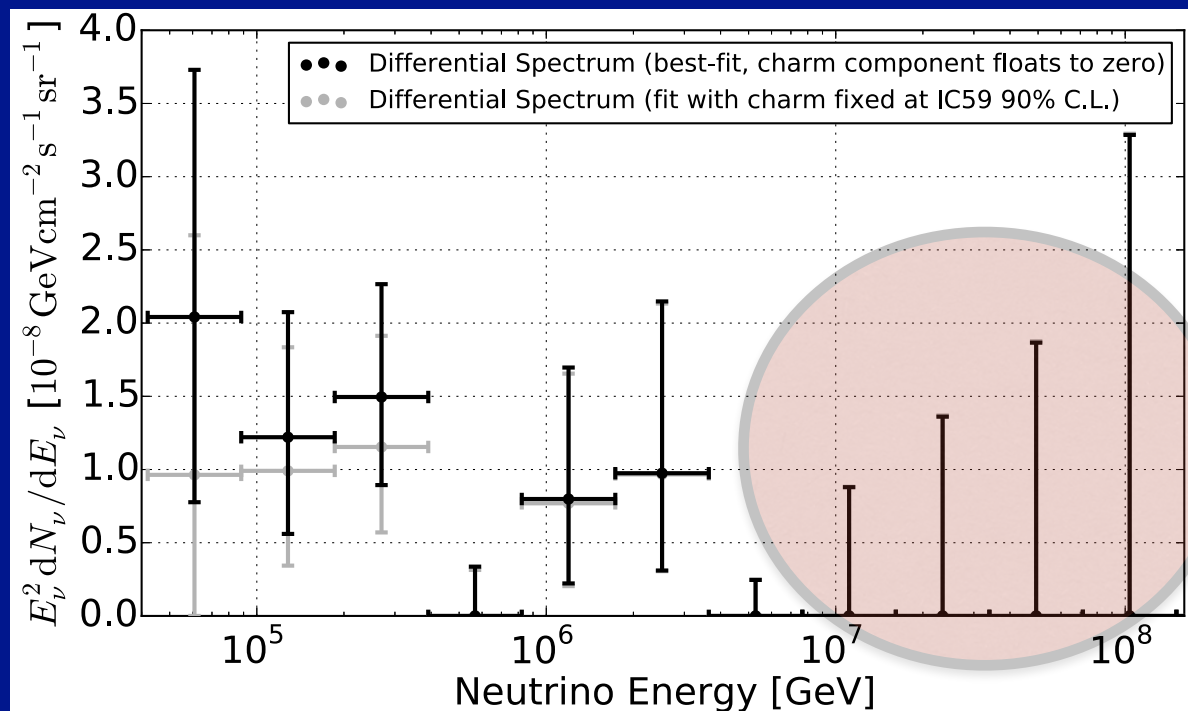


FIG. 4. Extraterrestrial neutrino flux ($\nu + \bar{\nu}$) as a function of energy. Vertical error bars indicate the $2\Delta\mathcal{L} = \pm 1$ contours of the flux in each energy bin, holding all other values, including background normalizations, fixed. These provide approximate 68% confidence ranges. An increase in the

Where are $E > 2$ PeV neutrino continuum events?

$$\left(\frac{N}{T\Omega}\right)_{\text{non-Res}} = N_{n+p} \int_{E_\nu^{\text{min}}} dE_\nu \left(\frac{dF_\nu}{dE_\nu}\right) \sigma(E_\nu)_{\nu N}$$
$$= N_{n+p} \left(\frac{6.3 \text{ PeV}}{E_\nu^{\text{min}}}\right)^{\alpha-1.40} \left(\frac{\sigma_{\nu N}^{\text{CC}}(E_\nu)}{(\alpha-1.40)} \frac{E_\nu dF_\nu}{dE_\nu}\right) \Big|_{E_\nu=6.3 \text{ PeV}},$$

i.e., 1, 0.66 (0.54), 0.52 (0.39), 0.44 (0.29), 0.38 (0.23)
for $E_\nu^{\text{min}} = 1, 2, 3, 4, 5, \text{ PeV}$

for spectral index = -2 (- 2.3) .

Glashow is a “no-show”?

$$\bar{\nu}_e + e^- \rightarrow W^-$$

$$s = M_W^2 = 2m_e E_\nu, \quad \text{so } E_R = \frac{M_W^2}{2m_e} = 6.3 \text{PeV}$$

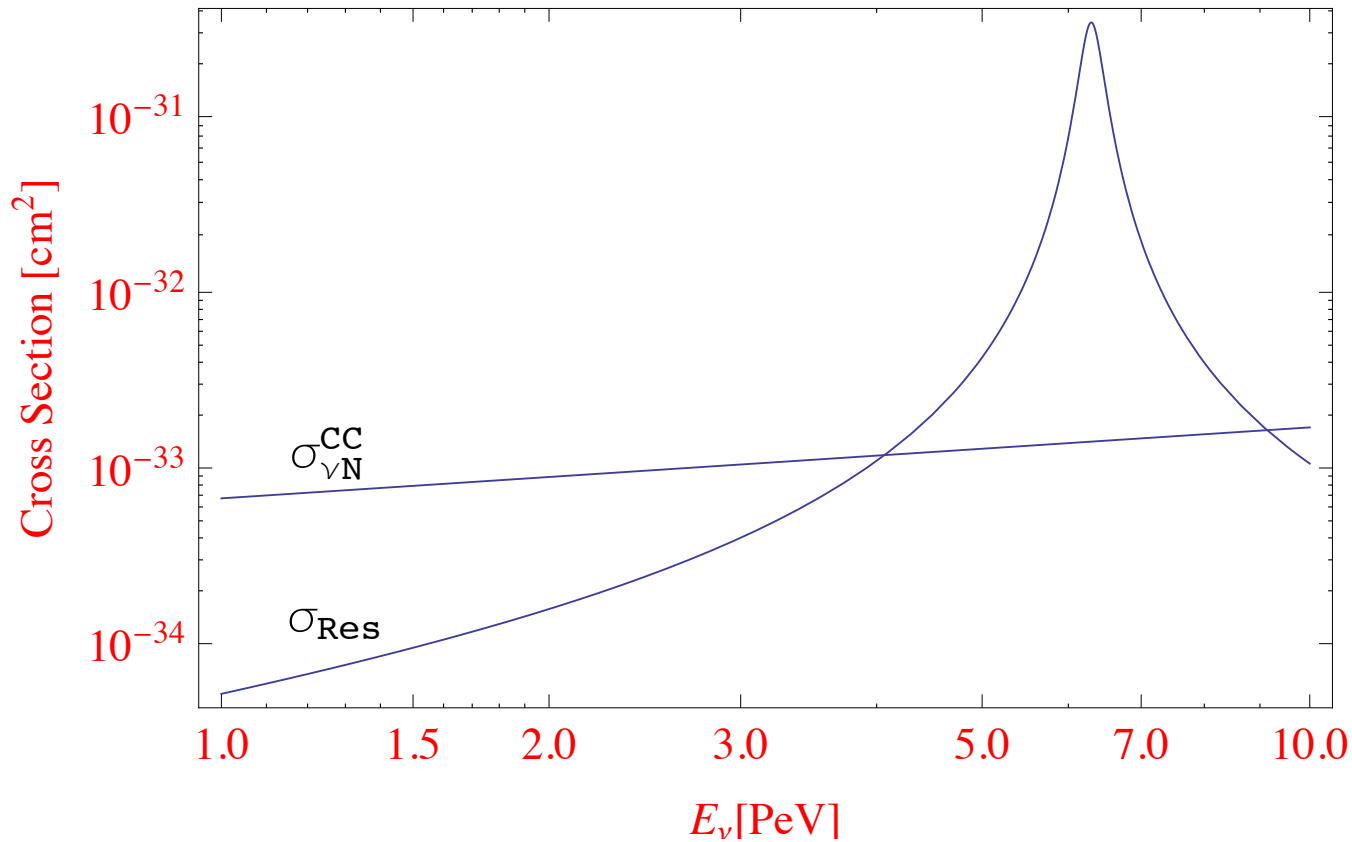


FIG. 1: Cross sections for the resonant process, $\bar{\nu}_e + e^- \rightarrow W^- \rightarrow \text{hadrons}$, and the non-resonant process, $\nu_e + N \rightarrow e^- + \text{hadrons}$, in the 1–10 PeV region.

Possible explanations:

Most probable is

(1) spectral break or spectral end at 2-3 PeV;

and next is

(2) Nature playing a statistical trick on us.

Less probable are

(3) Particle resonances, which now fail the filled-in “gap”;

But what if

(4) sigma gets stronger (need 10^4 factor) ? or

(5) Lorentz invariance fails, signified by

- an E_{nu}, or

- a Gamma_{nu} cutoff ?!



- **Galactic:** (full or partial contribution)
 - diffuse or unidentified Galactic γ -ray emission [Fox, Kashiyama & Meszaros'13]
[MA & Murase'13; Neronov, Semikoz & Tchernin'13; Neronov & Semikoz'14; Guo, Hu & Tian'14]
 - extended Galactic emission [Su, Slatjer & Finkbeiner'11; Crocker & Aharonian'11]
[Lunardini & Razzaque'12; MA & Murase'13; Razzaque'13; Lunardini *et al.*'13]
[Taylor, Gabici & Aharonian'14]
 - heavy dark matter decay [Feldstein *et al.*'13; Esmaili & Serpico '13; Bai, Lu & Salvado'13]
- **Extragalactic:**
 - association with sources of UHE CRs [Kistler, Stanev & Yuksel'13]
[Katz, Waxman, Thompson & Loeb'13; Fang, Fujii, Linden & Olinto'14]
 - active galactic nuclei (AGN) [Stecker'91,'13; Kalashev, Kusenko & Essey'13]
[Murase, Inoue & Dermer'14; Kimura, Murase & Toma'14; Kalashev, Semikoz & Tkachev'14]
 - gamma-ray bursts (GRB) [Murase & Ioka'13]
 - starburst galaxies [Loeb & Waxman'06; He *et al.*'13; Yoast-Hull, Gallagher, Zweibel & Everett'13]
[Murase, MA & Lacki'13; Anchordoqui *et al.*'14; Chang & Wang'14]
 - hypernovae in star-forming galaxies [Liu *et al.*'13]
 - galaxy clusters/groups [Murase, MA & Lacki'13; Zandanel *et al.*'14]
 - ...

Slide from M. Ahlers, NeuTel 2015

Mass-Scales and Energy Cutoff in terms of Boost Factor

Anchordoqui, Barger, Goldberg, Learned, Marfatia, Pakvasa, Paul, TJW (2014)

$$\Gamma_\nu = \left(\frac{E_\nu}{2 \text{ PeV}} \right) \left(\frac{0.05 \text{ eV}}{m_\nu} \right) \times 0.4 \cdot 10^{17}$$

No other massive particle can probe Γ 's this high !

whereas

$$\frac{M_{\text{Planck}}}{v_{\text{weak}}} = \frac{1.2 \times 10^{28} \text{ eV}}{247 \text{ GeV}} = 0.5 \times 10^{17} ;$$

Suggests

$$\Gamma_\nu^{\text{max}} = M_{\text{Planck}}/M_{\text{weak}}$$

Learned & TJW (2014)

and a possible connection
to Gravity/spacetime foam.

Neutrino Energy Maximum:

$$\begin{aligned}
 E_{\nu}^{\max} &= \frac{m_{\nu} M_{\text{Planck}}}{M_{\text{weak}}} \\
 &= 2.5 \left(\frac{m_{\nu}}{0.05 \text{ eV}} \right) \left(\frac{M_{\text{Planck}}}{1.2 \times 10^{28} \text{ eV}} \right) \left(\frac{247 \text{ GeV}}{v_{\text{weak}}} \right) \text{ PeV} .
 \end{aligned}$$

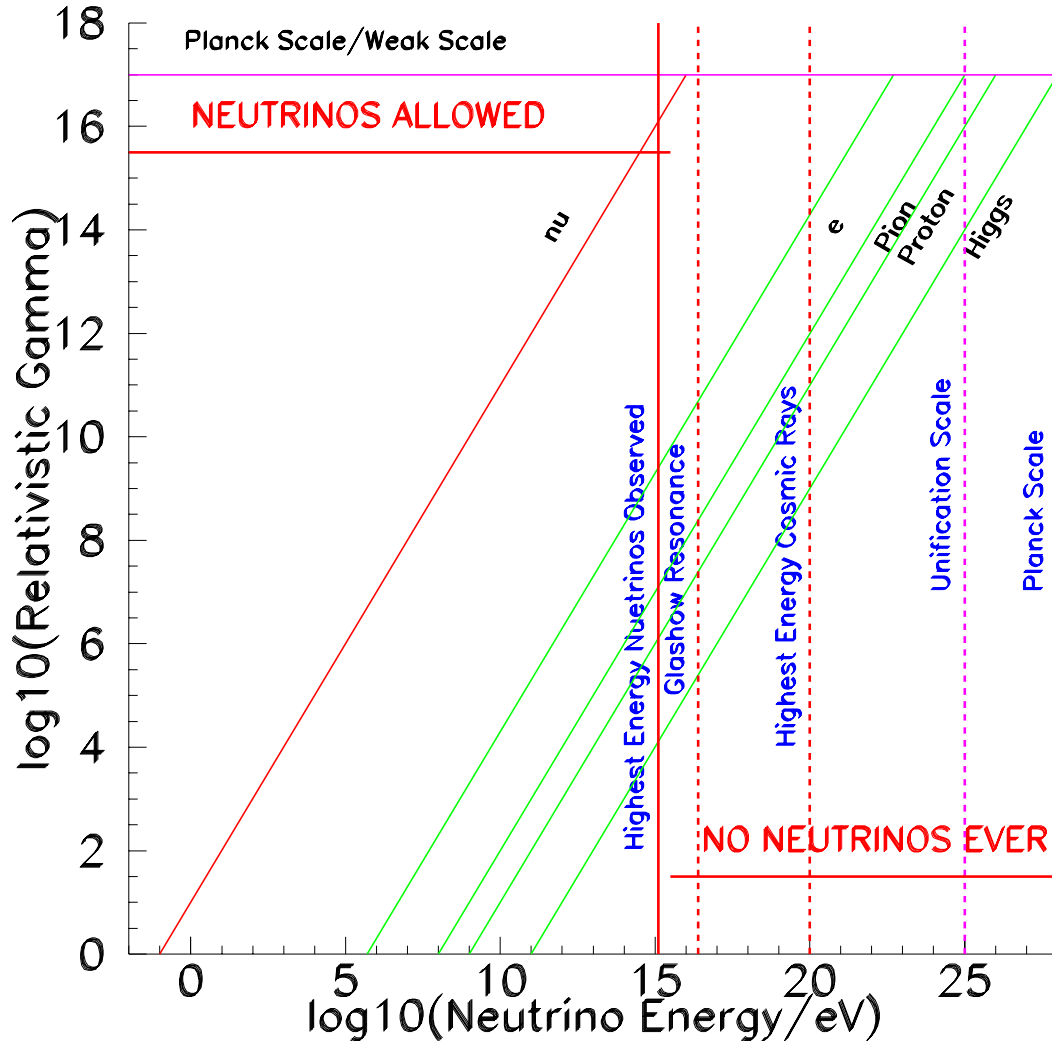
In what frame?

Nature provides THE preferred frame, the Cosmic Rest Frame.

So E_{ν}^{\max} can be written as $u_{\beta}^{\text{CRF}} (p_{\nu}^{\max})^{\beta}$, where $u_{\beta}^{\text{CRF}} = (1, \vec{0})$.

And $(p_{\nu}^{\max})^{\beta}$ transforms as usual four-vector.

The End of the Neutrino Spectrum



Reasons (excuses):

- 1- LI is “emergent”
low-energy symmetry;
- 2- Weak int’n “size” is Higgs
vev fluctuation, contracted
by Lorentz to Planck size:
 - a) spacetime foam
 - b) strong gravity/geometry
 - c) extra dim’ns open up
 - d) Planck scale LIV
 - e) ...make your own model
(like Bj);
doesn’t matter,
as it is an Xp1l issue!

Learned & TJW (2014)

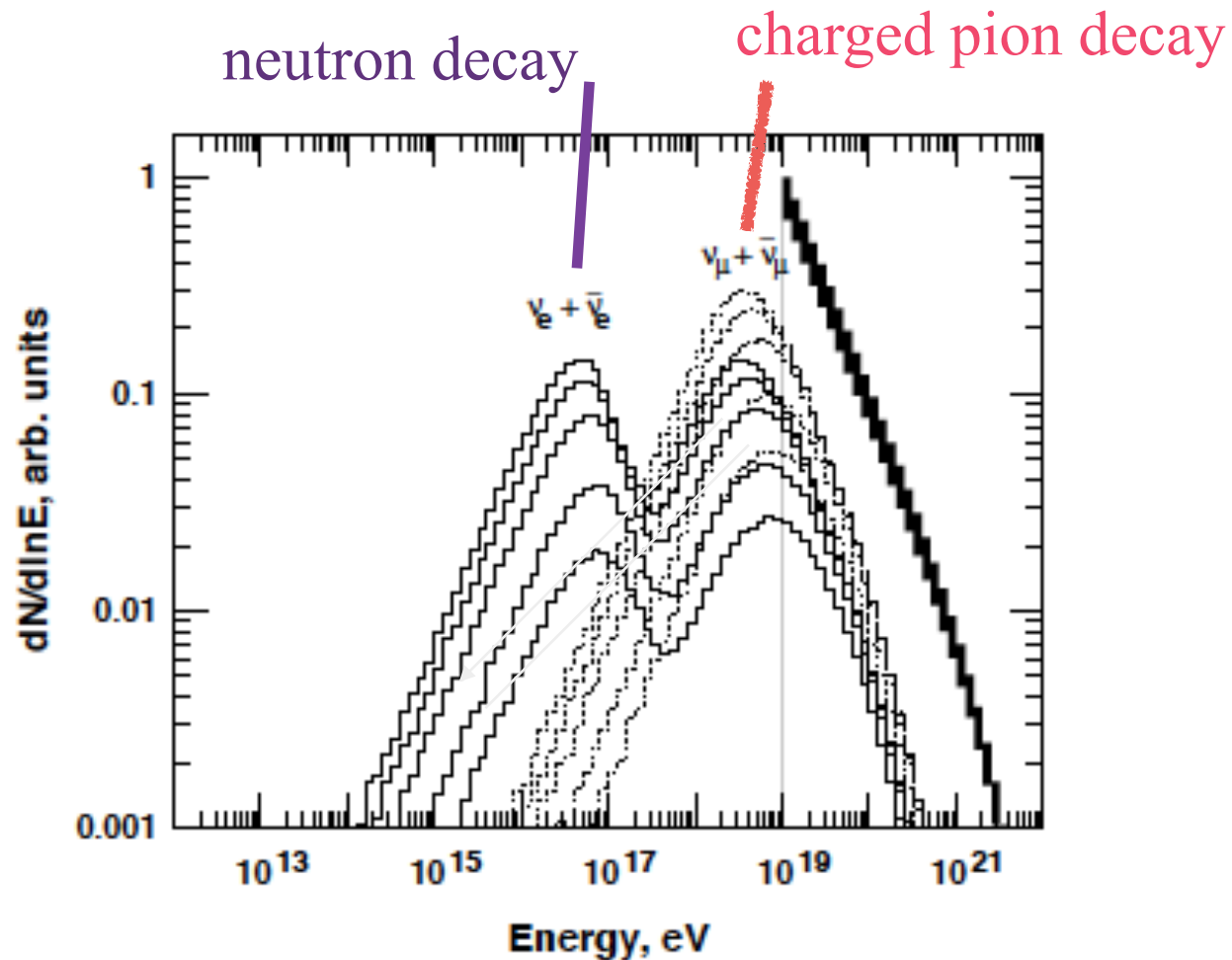
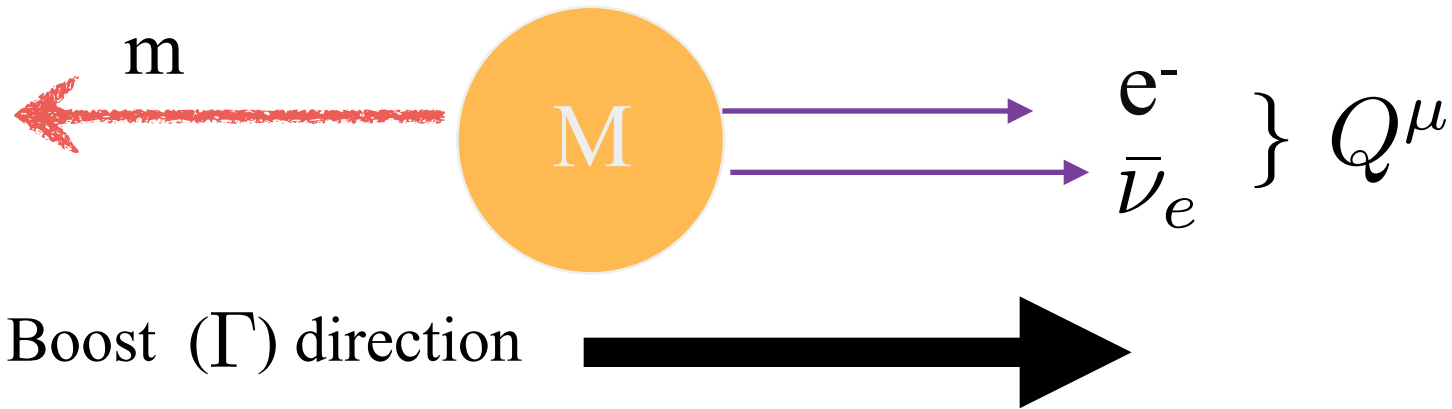


FIG. 2. Neutrino fluxes produced during the propagation of protons over 10, 20, 50, 100, and 200 Mpc (from bottom up) in a 1 nG random magnetic field. The heavy histogram shows the proton injection spectrum defined in Eq. (1).

Not only are many PeV neutrinos (and electrons) forbidden, but also the neutron, muon, and charged pion are stabilized by $E(\text{lepton})^{\text{max}}$



$$Q_0^{\text{LAB}} = \Gamma \bar{Q}_0 - \sqrt{\Gamma^2 - 1} \bar{Q} \approx \Gamma(\bar{Q}_0 - \bar{Q}) = \Gamma \frac{Q^2}{\bar{Q}_0 + \bar{Q}} > \frac{\Gamma Q^2}{2\bar{Q}_0}$$

$$\text{So, } E_M^{\text{LAB}} = M\Gamma < \left(\frac{2\bar{Q}_0 M}{Q^2} \right) Q_0^{\text{LAB}}$$

implies absolute stability if $E_M > \frac{M^2 - m^2}{m_e^2} (E_e + E_\nu)^{\text{max}}$

Anchordoqui, Barger, Goldberg, Learned, Marfatia, Pakvasa, Paul, TJW (2014)

Sufficient (but not necessary) conditions for absolute stability:

$$\frac{E_M^{\text{stability}}}{(E_e + E_\nu)}$$

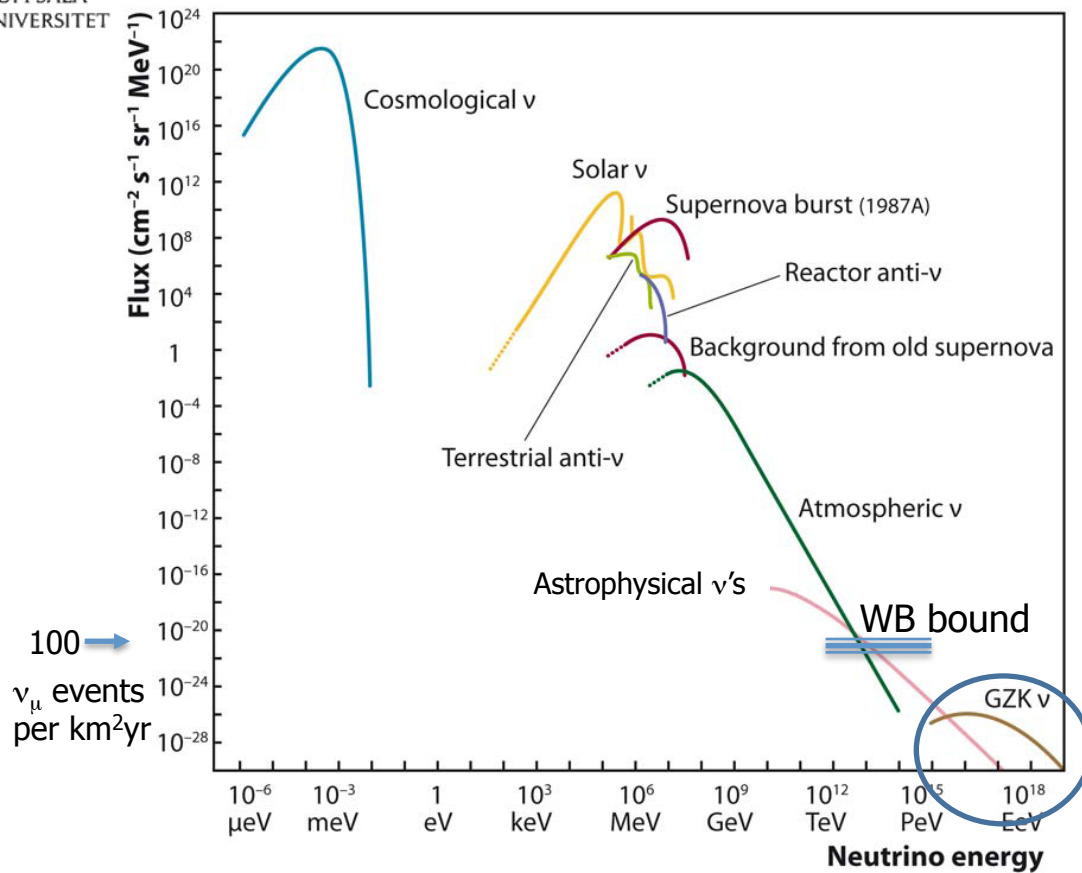
$n \rightarrow p e^- \bar{\nu}_e$	$\frac{m_n^2 - m_p^2}{m_e^2}$	9.7×10^3
$\mu \rightarrow \nu_\mu e \nu_e$	$\frac{m_\mu^2}{m_e^2}$	4.3×10^4
$\pi^\pm \rightarrow \gamma e^\pm \nu_e$	$\frac{m_{\pi^\pm}^2}{m_e^2}$	7.5×10^4
$\pi^\pm \rightarrow \pi^0 e^\pm \nu_e$	$\frac{m_{\pi^\pm}^2 - m_{\pi^0}^2}{m_e^2}$	4.5×10^3
$\tau \rightarrow \nu_\tau e \nu_e$	$\frac{m_\tau^2}{m_e^2}$	1.2×10^7

Denton, Marfatia
& TJW (2015)

Cosmogenic's another "no-show" ?

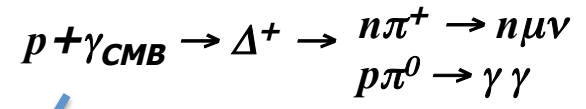


UPPSALA
UNIVERSITET



"guaranteed" flux

Greisen-Zatsepin-Kuzmin cosmogenic ν 's

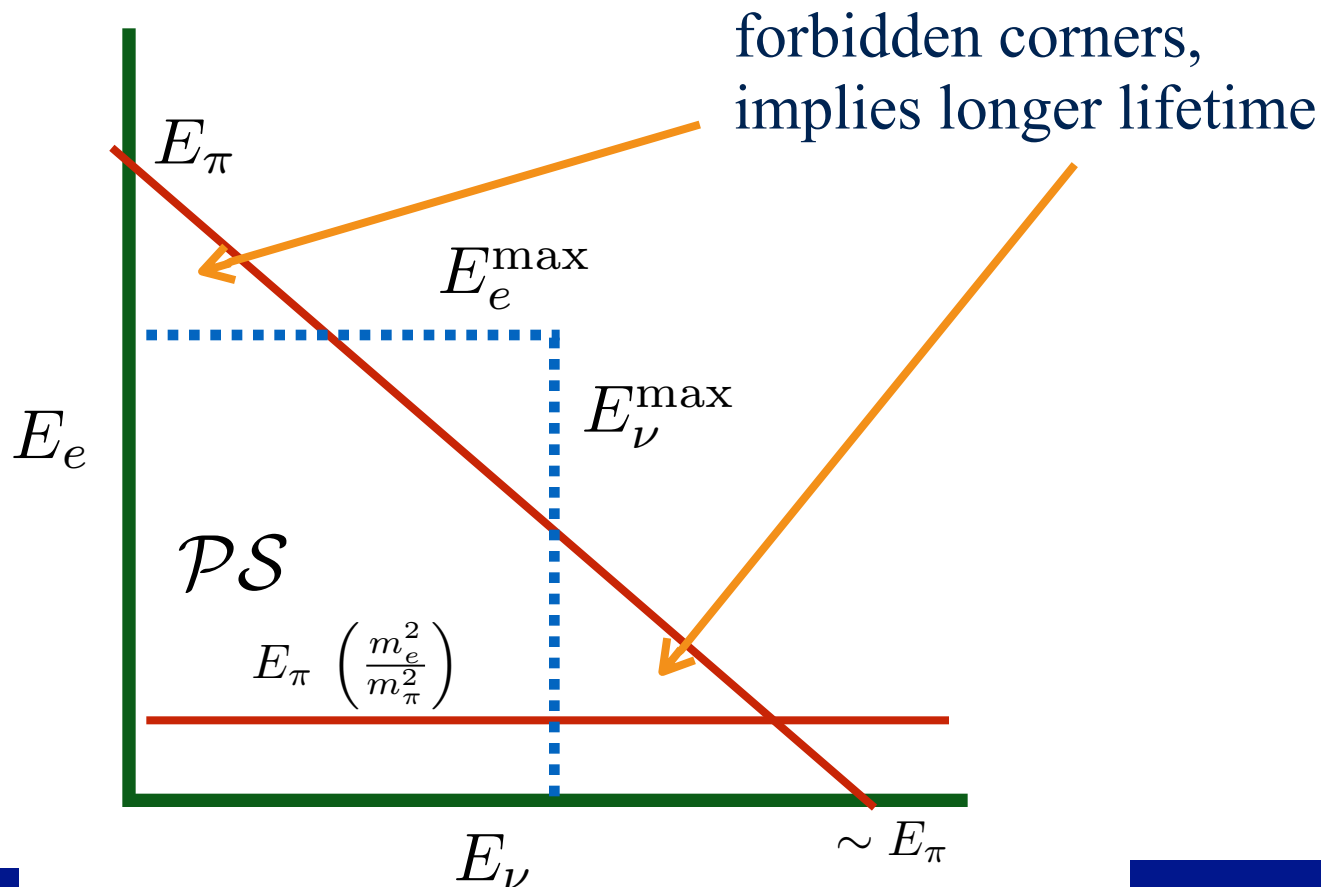


- atmospheric ν 's dominate < 100 TeV
- astrophysical ν 's (perhaps) > 100 TeV

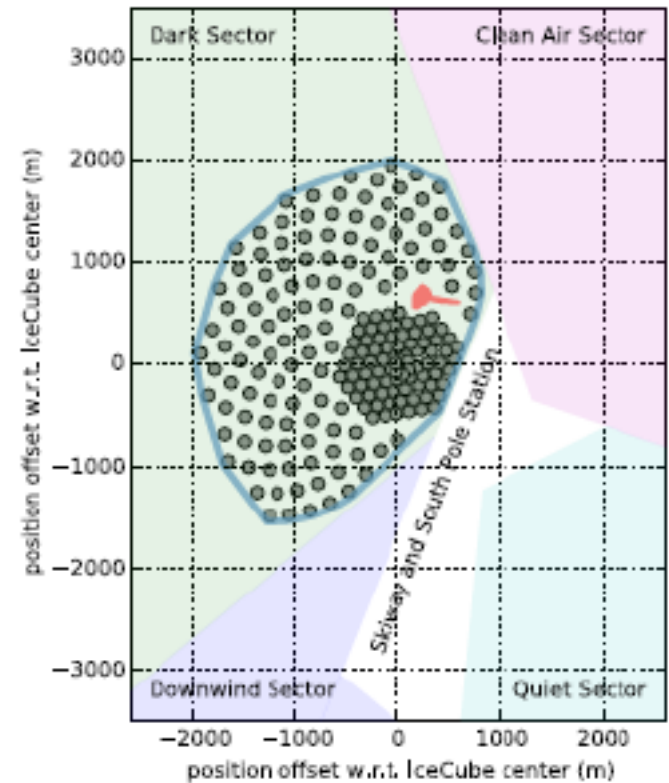
Necessary and Sufficient much harder (in progress)

3-body phase space is

$$dLIPS^{(3)} = \frac{1}{32\pi^3} dE_1 dE_2$$



May receive help from
IceCube Gen-2,
to answer questions:

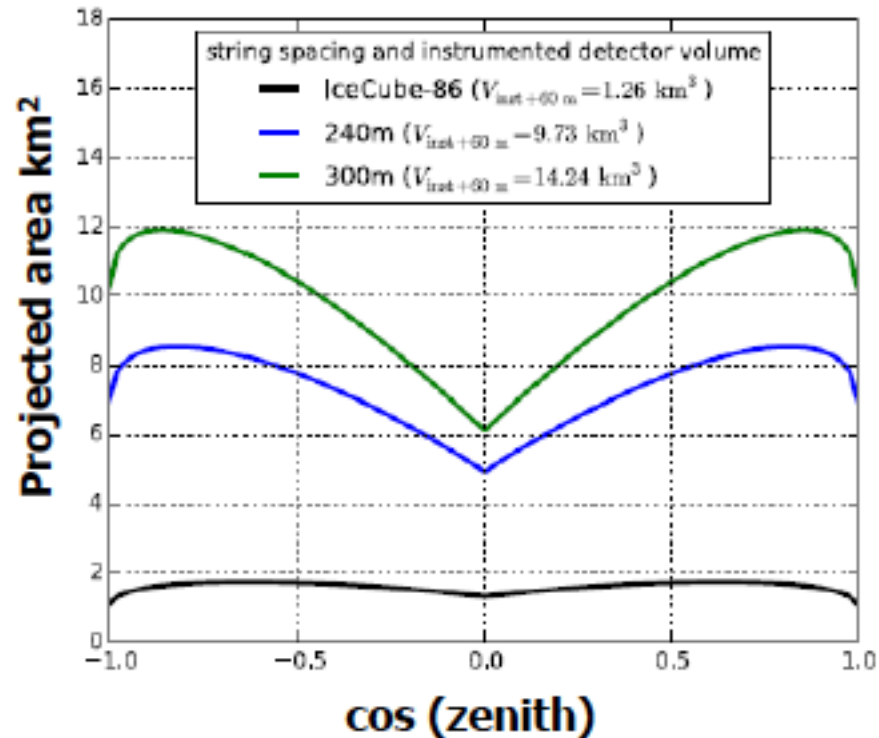


STRAWMAN DETECTOR

- 120 additional strings
- length 1.3 km
- average spacing 240 m
- volume 9.7 km^3

Expect 5-10 increase in effective area,
=> 5-10 increase in EVENT RATE:

INCREASE IN VOLUME AND PROJECTED AREA



In Summary:

Multi-PeV and Glashow resonance rates not yet in danger, but worth watching.

Glashow resonance can reveal ν source dynamics on other side of Universe (Danny M's talk).

If events above a few PeV do not arrive, then either/or

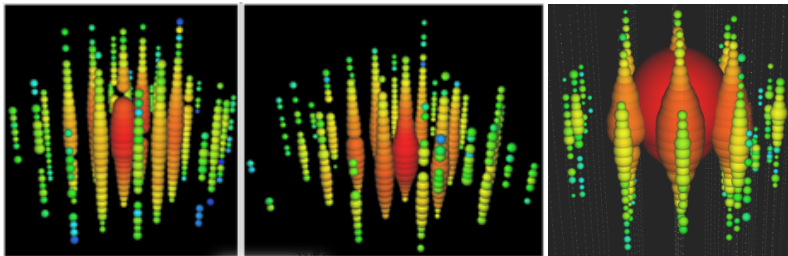
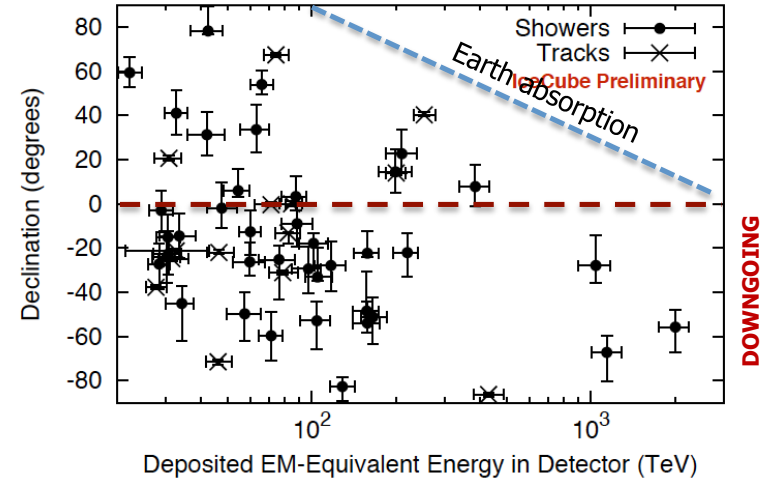
- (a) Nature cuts off/breaks the sources
($E_p \sim 20 E_\nu$ for pion chain);
- (b) $\sigma_{\nu N}$ is anomalously large at \sim PeV (by $\gtrsim 10^4$);
- (c) new fundamental physics at scale $\Gamma_\nu \sim 10^{17}$.

Extra Slides

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[*IceCube, Phys.Rev.Lett. 113:101101 (2014)*]

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1.04 PeV
Aug. 2011

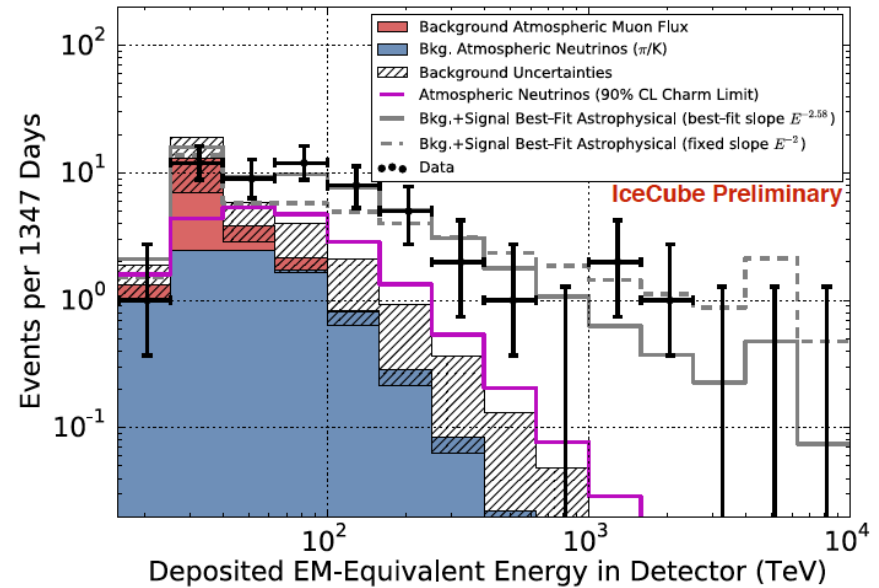


"Ernie"
1.14 PeV
Jan. 2012



"Big Bird"
2 PeV
Dec. 2012

- zenith distribution consistent with **isotropic** astrophysical flux



Glashow Resonance - Formulas:

$$\left(\frac{N}{T\Omega}\right)_{\text{Res}} = \frac{N_p}{2m_e} (\pi M_W \Gamma_W) \sigma_{\text{Res}}^{\text{peak}} \left. \frac{dF_{\bar{\nu}_e}}{dE_{\bar{\nu}_e}} \right|_{E_{\bar{\nu}_e}=6.3\text{PeV}},$$

$$\sigma_{\text{Res}}^{\text{peak}} = \frac{24\pi \text{B}(W^- \rightarrow \bar{\nu}_e e^-) \text{B}(W^- \rightarrow \text{had})}{M_W^2} = 3.4 \times 10^{-31} \text{cm}^2.$$

The “Resonometer” of Cosmic Nu Source Models: Barger, Fu, Learned, Marfatia, Pakvasa, TJW, arXiv'd today

TABLE I: Neutrino flavor ratios at source, component of $\bar{\nu}_e$ in total neutrino flux at Earth after mixing and decohering, and consequent relative strength of Glashow resonance, for six astrophysical models. (Neutrinos and antineutrinos are shown separately, when they differ.)

	Source flavor ratio		Earthly flavor ratio		$\bar{\nu}_e$ fraction in flux (\mathcal{R})
$pp \rightarrow \pi^\pm$ pairs	(1:2:0)		(1:1:1)		18/108 = 0.17
w/ damped μ^\pm	(0:1:0)		(4:7:7)		12/108 = 0.11
$p\gamma \rightarrow \pi^+$ only	(1:1:0)	(0:1:0)	(14:11:11)	(4:7:7)	8/108 = 0.074
w/ damped μ^+	(0:1:0)	(0:0:0)	(4:7:7)	(0:0:0)	0
charm decay	(1:1:0)		(14:11:11)		21/108 = 0.19
neutron decay	(0:0:0)	(1:0:0)	(0:0:0)	(5:2:2)	60/108 = 0.56

(Kaons change little)

Glashow event rates vs. continuum:

TABLE II: Ratio of resonant event rate around the 6.3 PeV peak to non-resonant event rate above $E_\nu^{\min} = 1, 2, 3, 4, 5$ PeV. The single power-law spectral index α is taken to be 2.0 and 2.3 for the non-parenthetic and parenthetic values, respectively. As an example, the single power-law extrapolation from the three events observed just above 1 PeV predicts a mean number of observed resonance events around 6.3 PeV equal to the first numerical column times 3.

E_ν^{\min} (PeV)	1	2	3	4	5
$pp \rightarrow \pi^\pm$ pairs	0.33 (0.29)	0.50 (0.53)	0.64 (0.77)	0.76 (1.0)	0.87 (1.2)
damped μ^\pm	0.22 (0.18)	0.33 (0.34)	0.42 (0.50)	0.49 (0.64)	0.56 (0.79)
$p\gamma \rightarrow \pi^+$ only	0.14 (0.12)	0.22 (0.23)	0.28 (0.33)	0.33 (0.43)	0.38 (0.53)
damped μ^+	0 (0)	0 (0)	0 (0)	0 (0)	0 (0)
charm decay	0.37 (0.32)	0.56 (0.60)	0.72 (0.86)	0.85 (1.1)	0.98 (1.4)
neutron decay	1.1 (0.94)	1.7 (1.8)	2.1 (2.5)	2.5 (3.3)	2.9 (4.0)

Glashow Resonance - rates vs. continuum

The numbers of expected resonant events is greatly reduced from the ratio of resonant to non-resonant cross sections by the additional factors. The cross section ratio at 6.3 PeV is 240: see Fig. 1. The $\frac{\Gamma_W}{M_W}$ ratio is 1/38. The α -dependent factor $\left[(\alpha - 1.40) \left(\frac{1 \text{ PeV}}{6.3 \text{ PeV}} \right)^{\alpha - 1.40} \right]$ yields about 0.2 for both α 's of interest, 2.0 and 2.3. The end result is about $2\mathcal{R}$ for the ratio of resonant events to non-resonant events above 1 PeV.

Since three down-going shower events have been observed at IceCube in the 1-2 PeV region, the expected number of Glashow events is found by multiplying the first numerical column of Table II by three. These expected resonant event numbers are 1.0 (0.9), 0.7 (0.5), 0.4 (0.4), 0 (0), 1.1 (1.0), and 3.3 (2.8), for the six models, and for $\alpha = 2.0$ (2.3), respectively. Since no 6.3 PeV events are observed,¹ all models remain viable except perhaps the final one, where neutron decay to pure $\bar{\nu}_e$ predicts some resonance events at Earth. In terms of Poisson statistics, when $\langle N \rangle$ events are expected, the probability that none are observed is $P(0|\langle N \rangle) = e^{-\langle N \rangle}$; thus, the six models yield Poissonian occurrence probabilities of 37%, 50%, 67%, 100%, 33%, and 4% for $\alpha = 2$, and slightly larger probabilities for $\alpha = 2.3$.