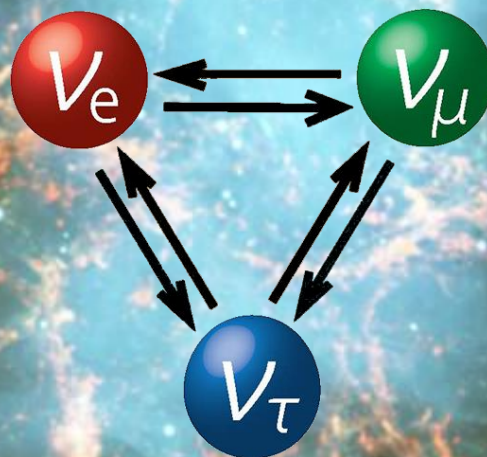


Neutrino Flavor Conversion in Supernovae



Georg Raffelt, Max-Planck-Institut für Physik, München

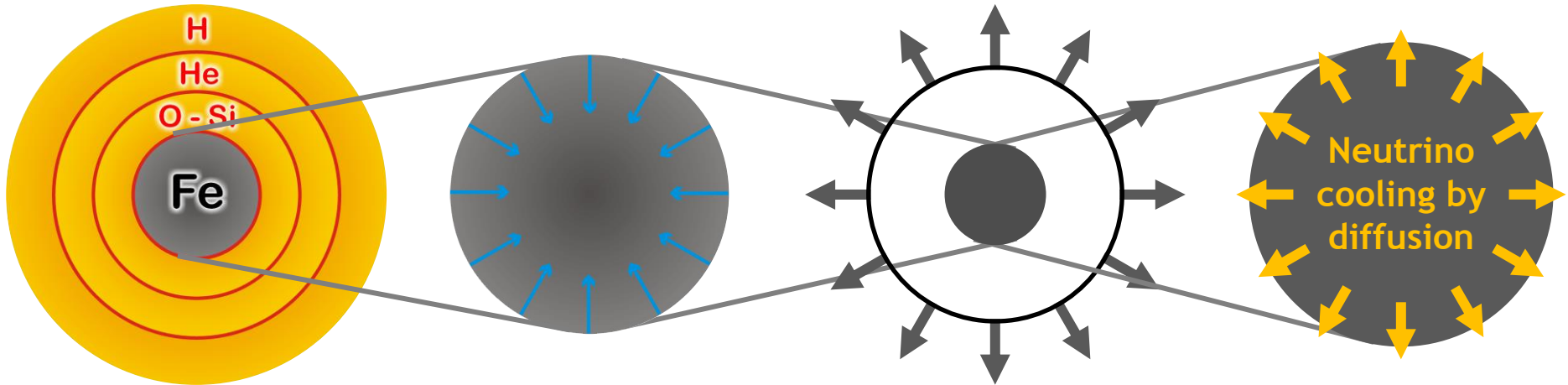
Core-Collapse Supernova Explosion

End state of a
massive star
 $M \gtrsim 6-8 M_{\odot}$

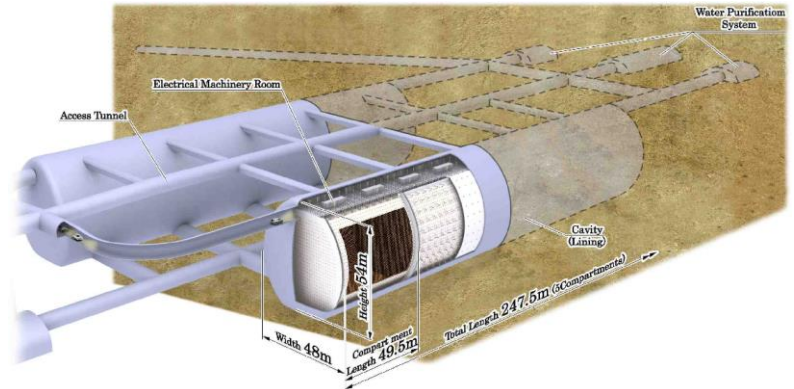
Collapse of
degenerate core

Bounce at ρ_{nuc}
Shock wave forms
explodes the star

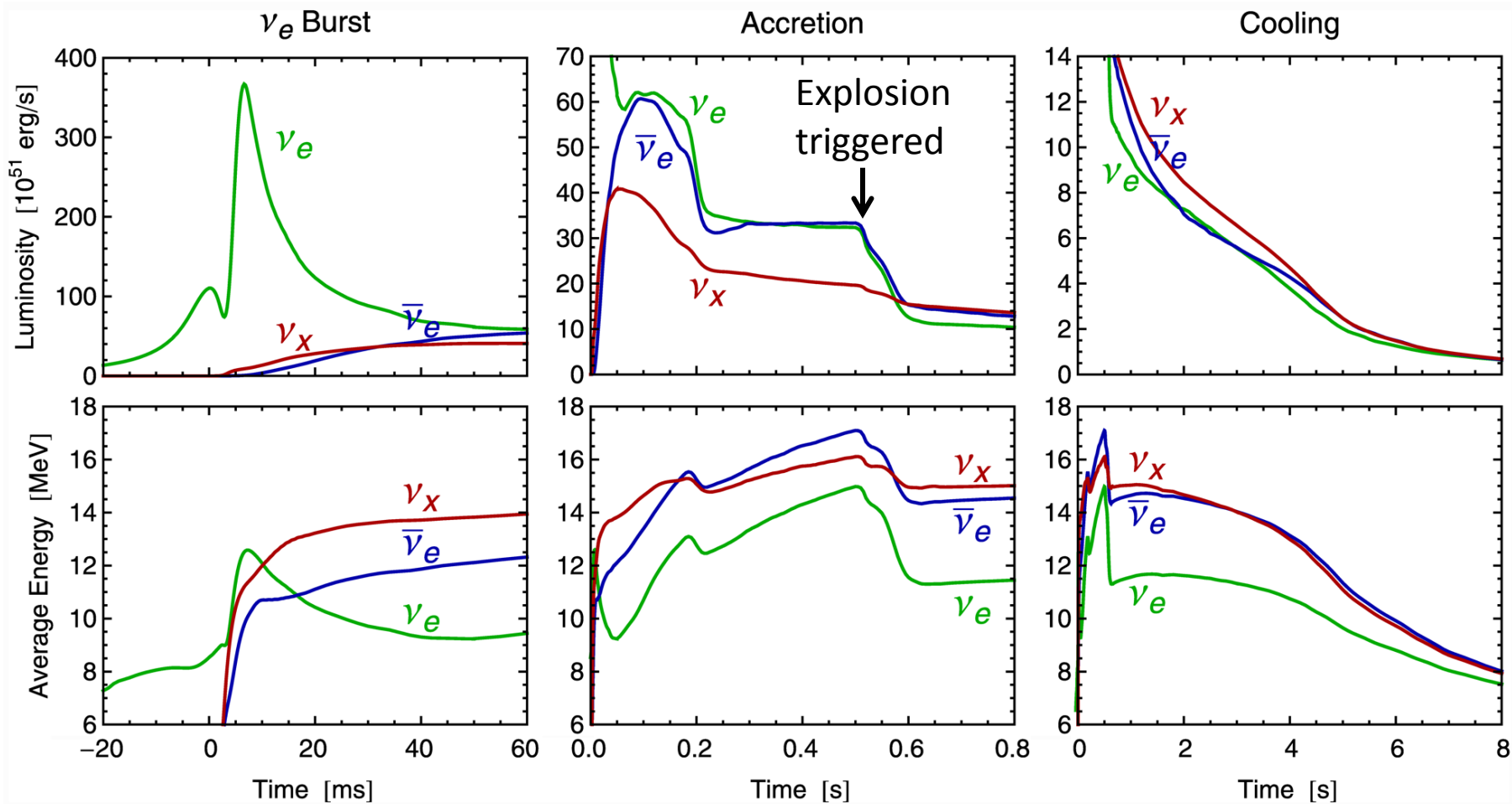
Grav. binding E
 $\sim 3 \times 10^{53}$ erg
emitted as nus
of all flavors



- Huge rate of low-E neutrinos (tens of MeV) over few seconds in large-volume detectors
- A few core-collapse SNe in our galaxy per century
- Once-in-a-lifetime opportunity



Three Phases of Neutrino Emission



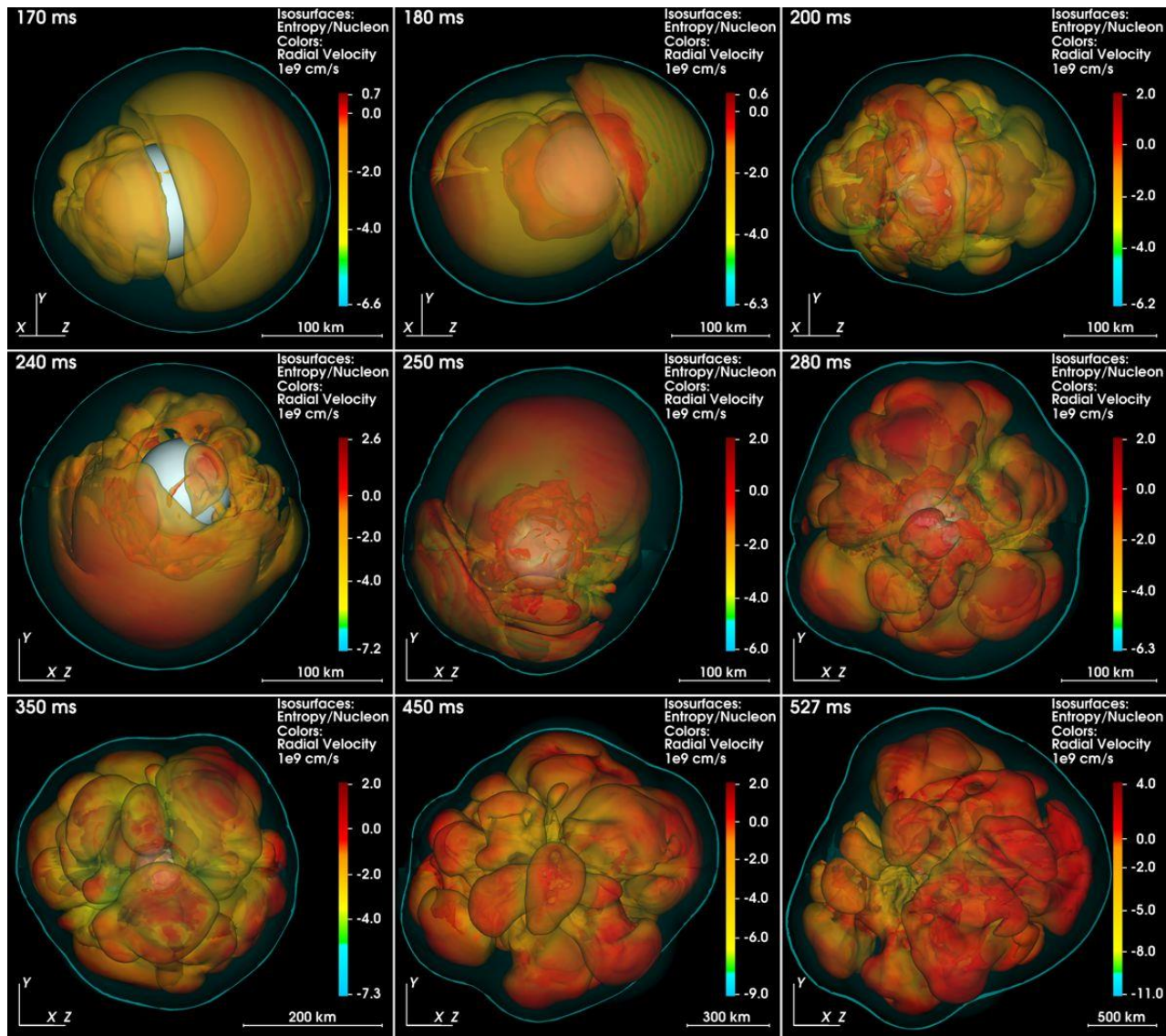
- Shock breakout
- De-leptonization of outer core layers

- Shock stalls ~ 150 km
- Neutrinos powered by infalling matter

Cooling on neutrino diffusion time scale

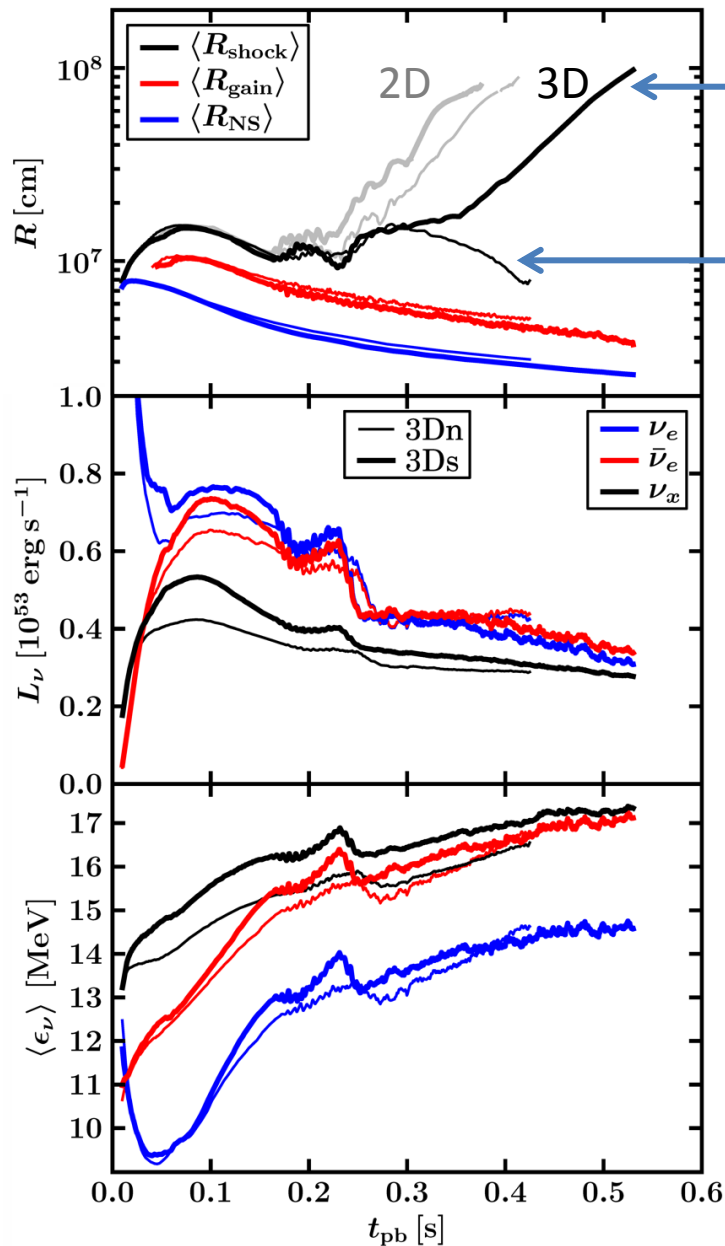
Spherically symmetric Garching model ($25 M_{\odot}$) with Boltzmann neutrino transport

Exploding 3D Garching Model (20 M_{SUN})



Melson, Janka, Bollig, Hanke, Marek & Müller, arXiv:1504.07631

Exploding 3D Garching Model (20 M_{SUN})

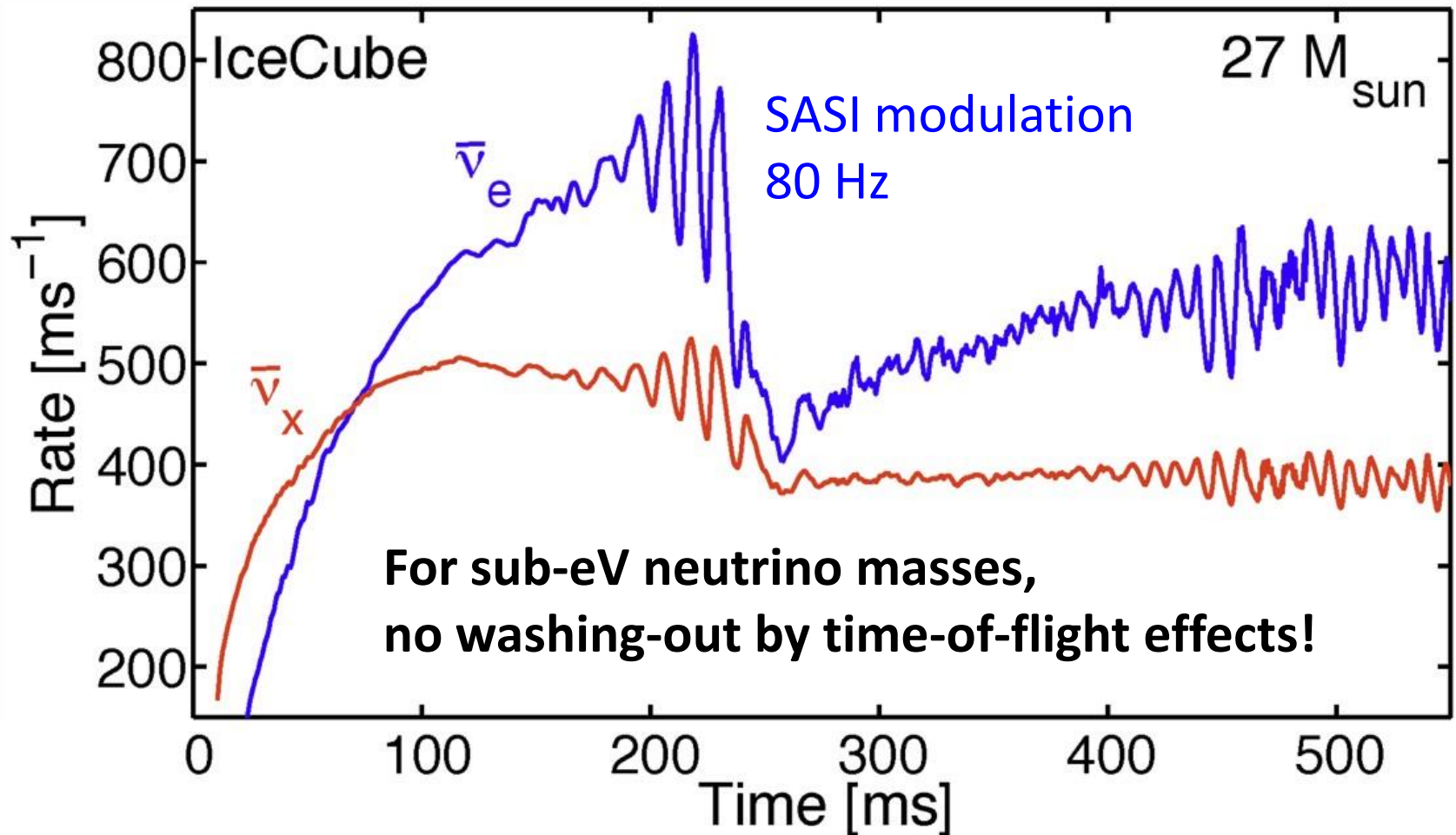


Neutrino opacity reduced (few 10%) by strange quark contribution to nucleon spin (thick lines)

“Standard” neutrino opacity (thin lines)

Melson, Janka, Bollig, Hanke, Marek & Müller,
arXiv:1504.07631

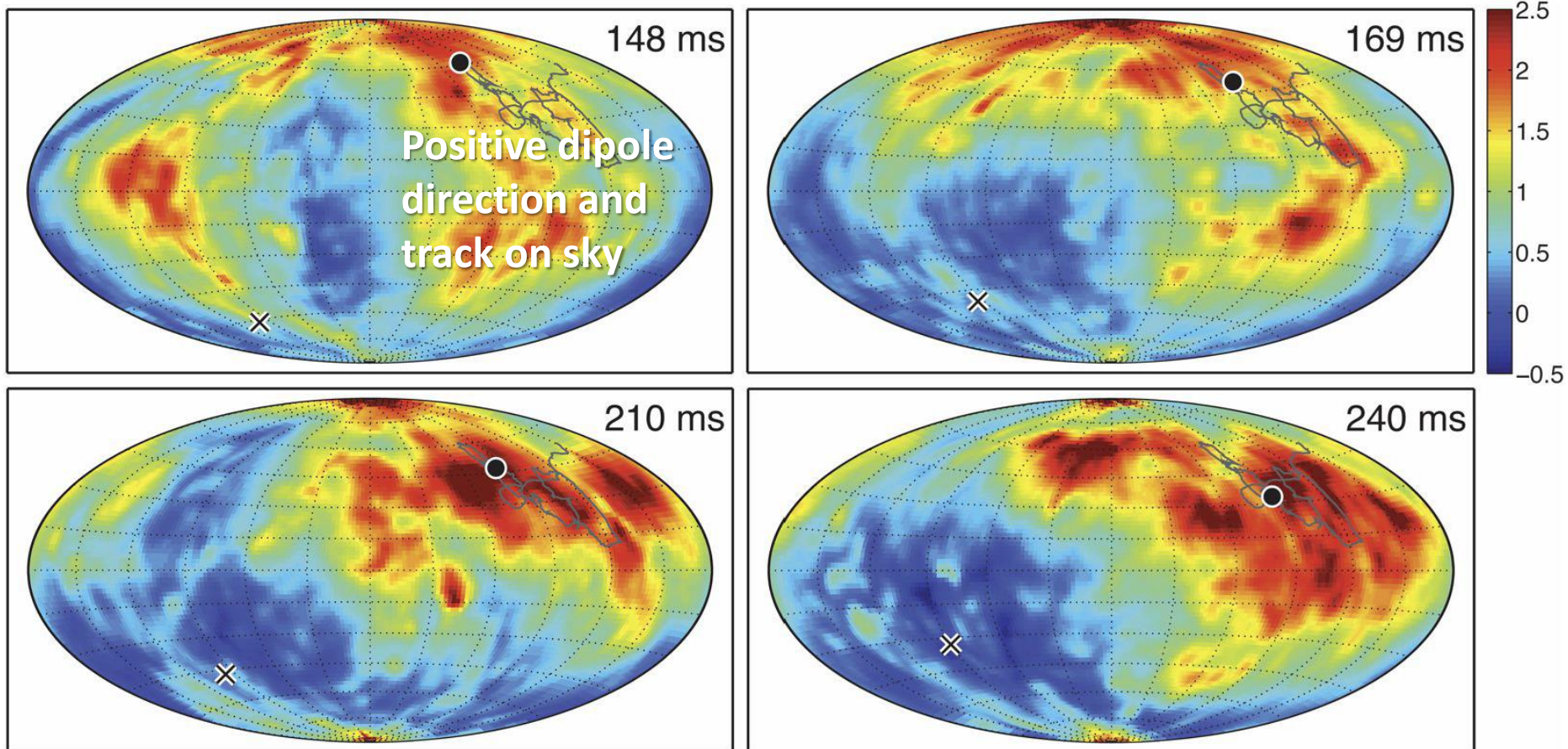
Variability seen in Neutrinos (3D Model)



Tamborra, Hanke, Müller, Janka & Raffelt, arXiv:1307.7936
See also Lund, Marek, Lunardini, Janka & Raffelt, arXiv:1006.1889

Sky Map of Lepton-Number Flux (11.2 M_{SUN} Model)

Lepton-number flux ($\nu_e - \bar{\nu}_e$) relative to 4π average
Deleptonization flux into one hemisphere, roughly dipole distribution
(LESA — Lepton Emission Self-Sustained Asymmetry)

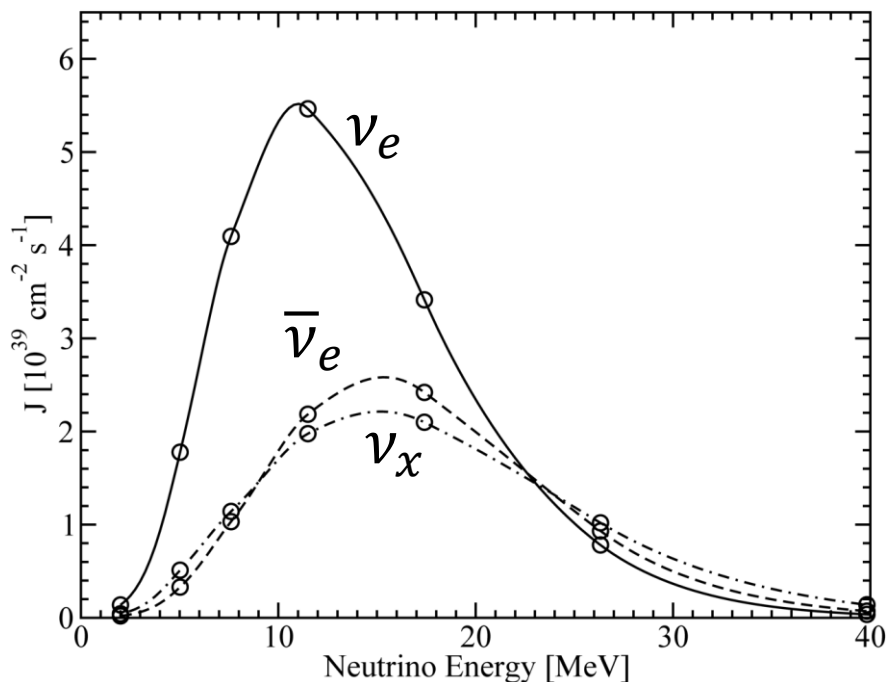


Tamborra, Hanke, Janka, Müller, Raffelt & Marek, arXiv:1402.5418

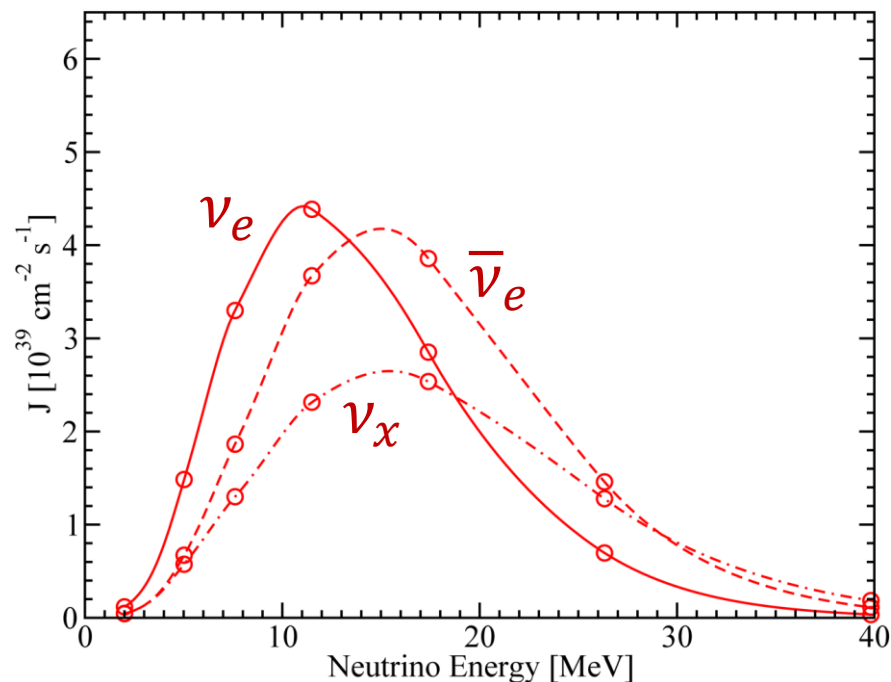
Spectra in the two Hemispheres

Neutrino flux spectra (11.2 M_{SUN} model at 210 ms) in opposite LESA directions

Direction of
maximum lepton-number flux



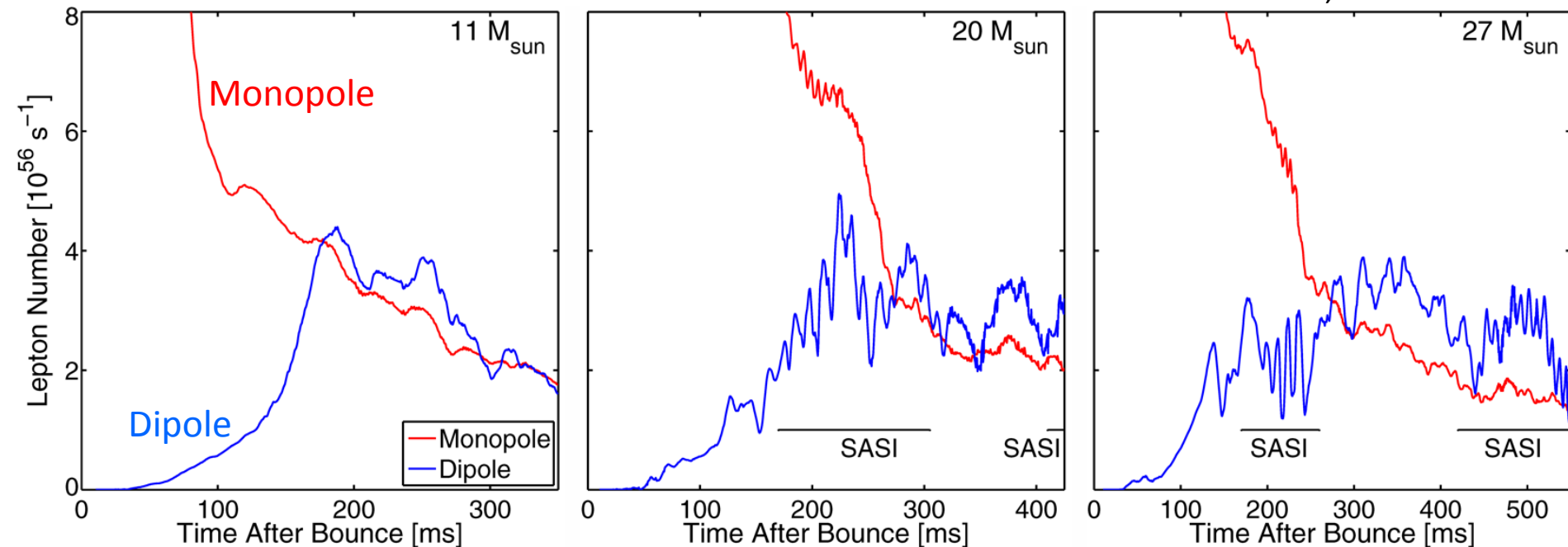
Direction of
minimum lepton-number flux



**During accretion phase, flavor-dependent fluxes
vary strongly with observer direction!**

Growth of Lepton-Number Flux Dipole

Tamborra et al., arXiv:1402.5418



- Overall lepton-number flux (monopole) depends on accretion rate, varies between models
- Maximum dipole similar for different models
- Dipole persists (and even grows) during SASI activity
- SASI and LESA dipoles uncorrelated

4000 citations

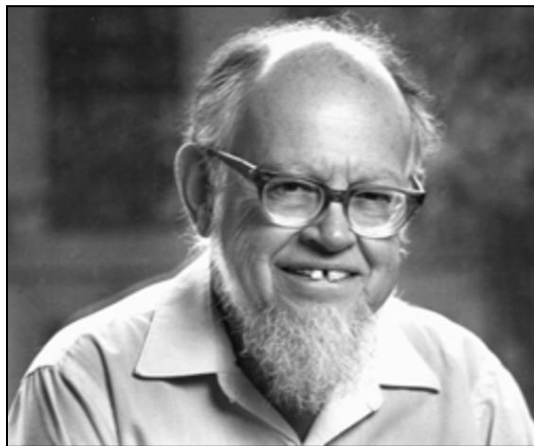
Neutrino oscillations in matter

L. Wolfenstein

Carnegie-Mellon University, Pittsburgh, Pennsylvania 15213

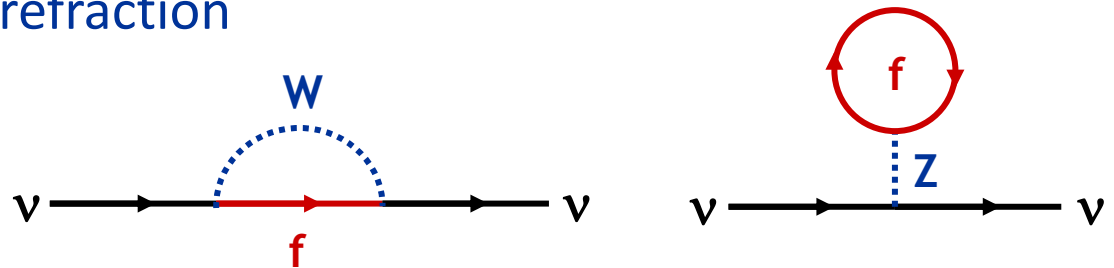
(Received 6 October 1977; revised manuscript received 5 December 1977)

The effect of coherent forward scattering must be taken into account when considering the oscillations of neutrinos traveling through matter. In particular, for the case of massless neutrinos for which vacuum oscillations cannot occur, oscillations can occur in matter if the neutral current has an off-diagonal piece connecting different neutrino types. Applications discussed are solar neutrinos and a proposed experiment involving transmission of neutrinos through 1000 km of rock.



Lincoln Wolfenstein
10 Feb 1923–27 Mar 2015

Neutrinos in a medium suffer flavor-dependent refraction

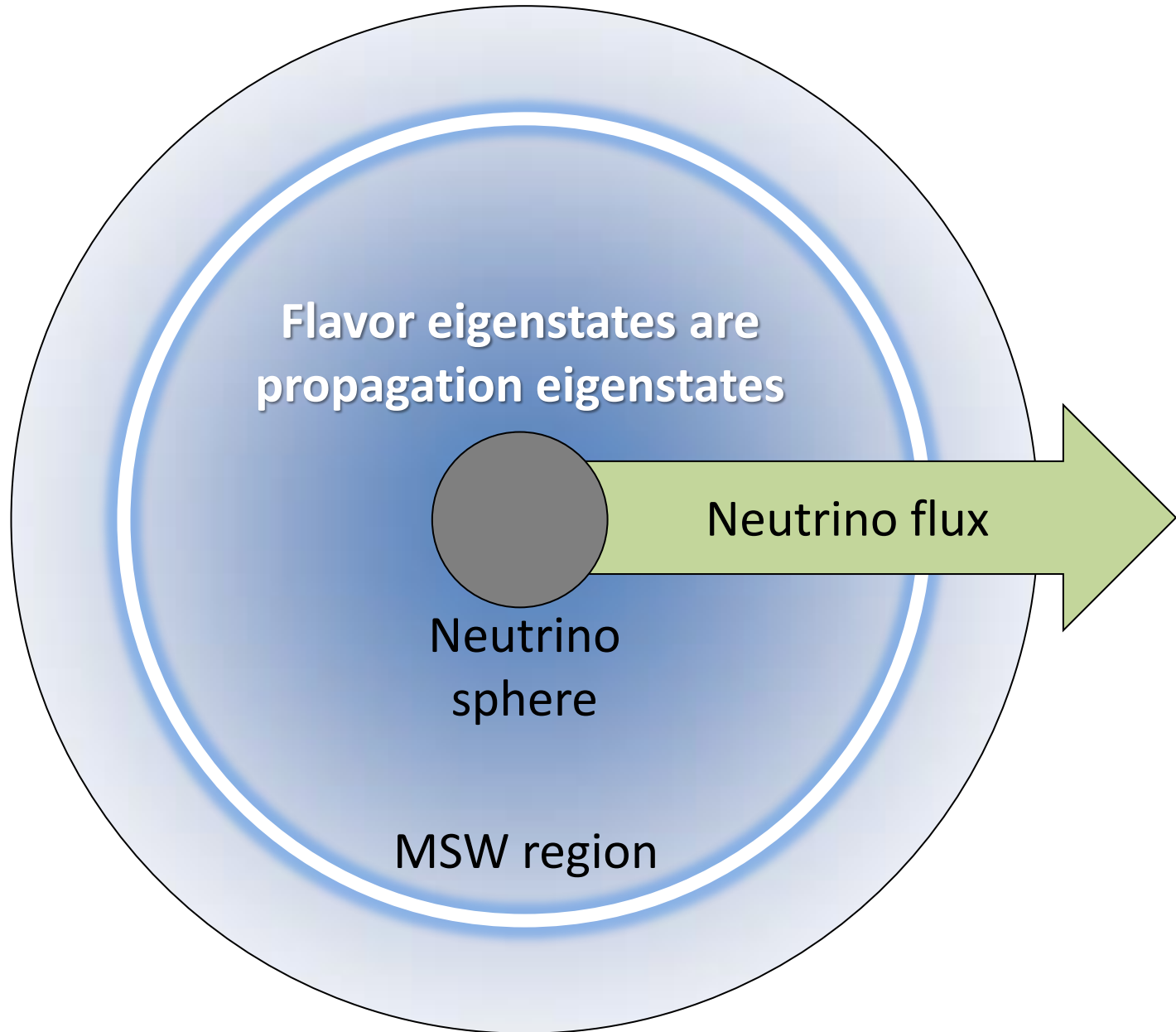


$$V_{\text{weak}} = \sqrt{2}G_F \times \begin{cases} N_e - N_n/2 & \text{for } \nu_e \\ -N_n/2 & \text{for } \nu_\mu \end{cases}$$

Typical density of Earth: 5 g/cm³

$$\Delta V_{\text{weak}} \approx 2 \times 10^{-13} \text{ eV} = 0.2 \text{ peV}$$

Flavor Oscillations in Core-Collapse Supernovae



SN Flavor Oscillations and Mass Hierarchy

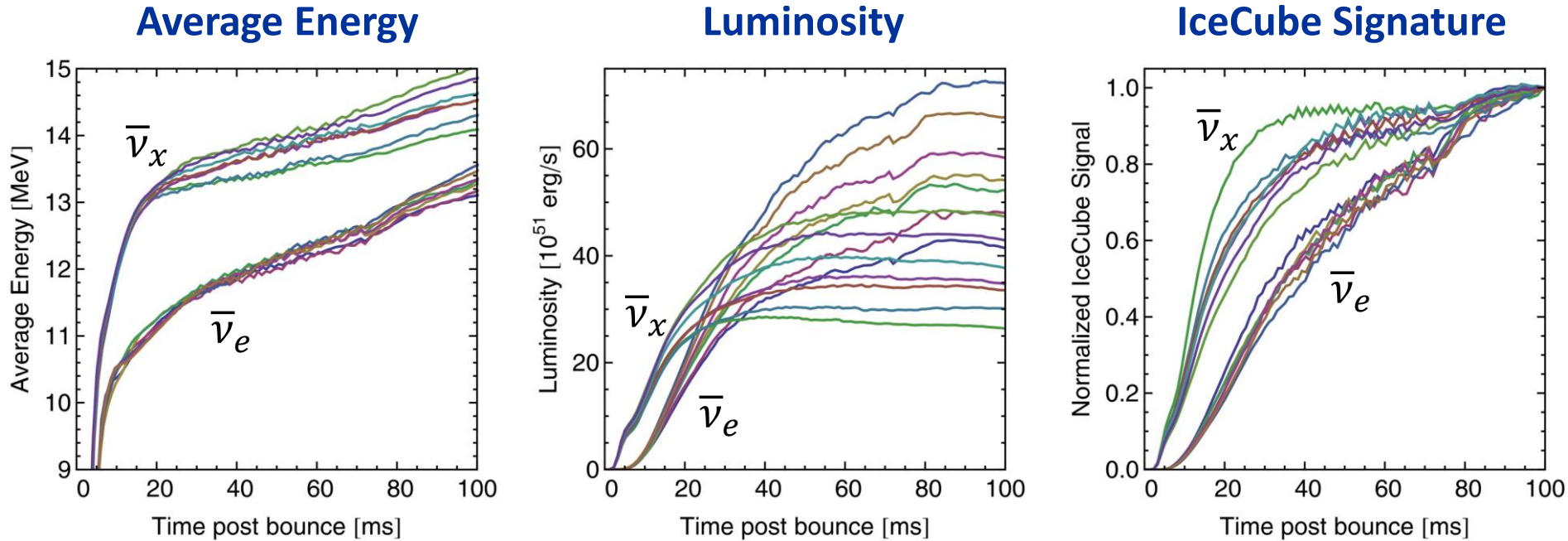
- Mixing angle Θ_{13} has been measured to be “large”
- MSW conversion in SN envelope adiabatic
- Assume that collective flavor oscillations are not important

| | Mass ordering | |
|------------------------------|----------------------------------|----------------------------------|
| | Normal (NH) | Inverted (IH) |
| ν_e survival prob. | 0 | $\sin^2 \theta_{12} \approx 0.3$ |
| $\bar{\nu}_e$ survival prob. | $\cos^2 \theta_{12} \approx 0.7$ | 0 |
| $\bar{\nu}_e$ Earth effects | Yes | No |

- When are collective oscillations important?
- How to detect signatures of hierarchy?

Early-Phase Signal in Anti-Neutrino Sector

Garching Models with $M = 12\text{--}40 M_{\odot}$



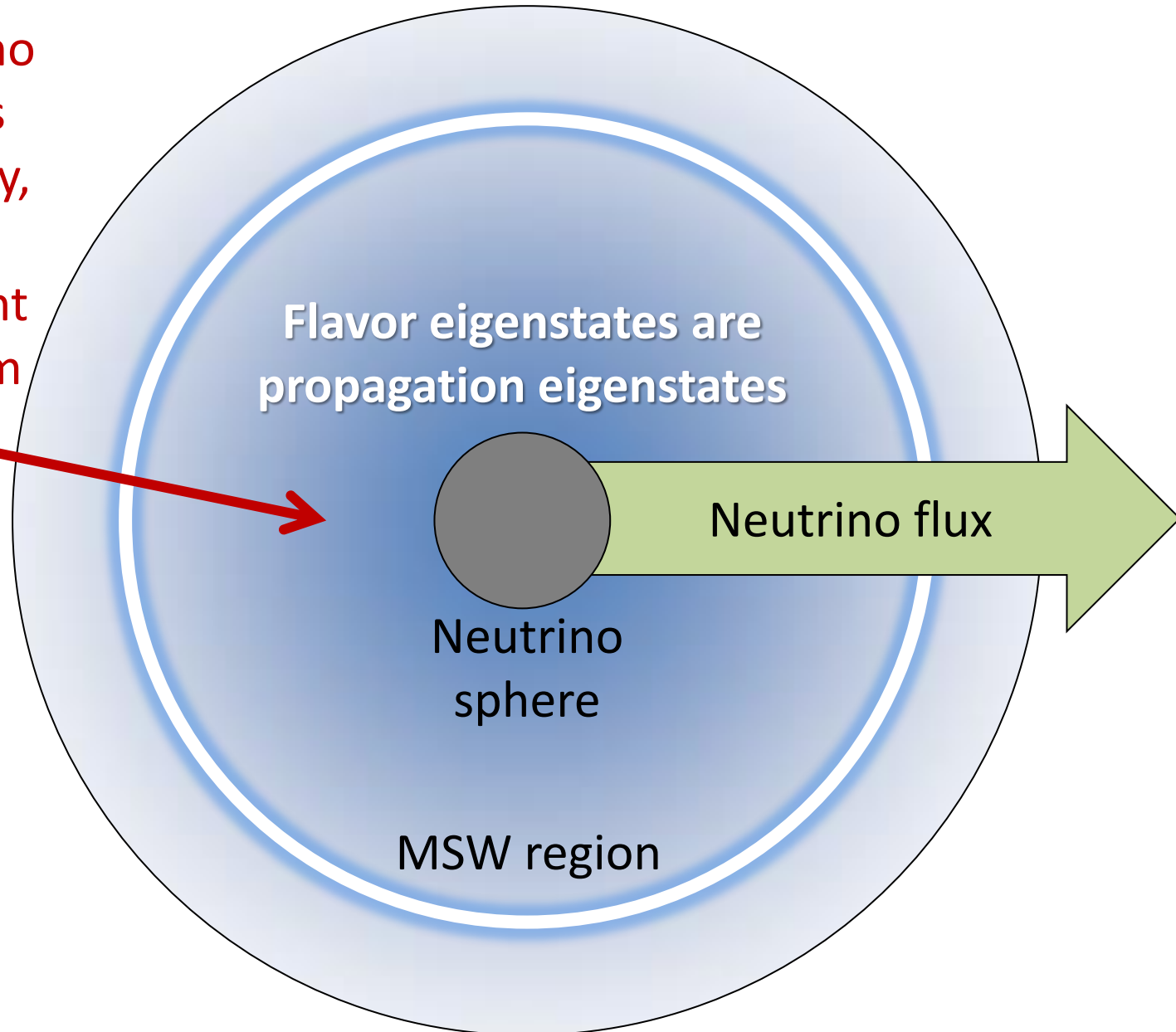
- In principle very sensitive to hierarchy, notably IceCube
- “Standard candle” to be confirmed by other than Garching models

Abbasi et al. (IceCube Collaboration) A&A 535 (2011) A109

Serpico, Chakraborty, Fischer, Hüdepohl, Janka & Mirizzi, arXiv:1111.4483

Flavor Oscillations in Core-Collapse Supernovae

Neutrino-neutrino refraction causes a flavor instability, flavor exchange between different parts of spectrum

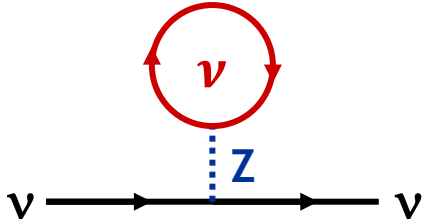


Flavor-Off-Diagonal Refractive Index

2-flavor neutrino evolution as an effective 2-level problem

$$i \frac{\partial}{\partial t} \begin{pmatrix} \nu_e \\ \nu_\mu \end{pmatrix} = H \begin{pmatrix} \nu_e \\ \nu_\mu \end{pmatrix}$$

Effective mixing Hamiltonian

$$H = \frac{M^2}{2E} + \sqrt{2}G_F \begin{pmatrix} N_e - \frac{N_n}{2} & 0 \\ 0 & -\frac{N_n}{2} \end{pmatrix} + \sqrt{2}G_F \begin{pmatrix} N_{\nu_e} & N_{\langle \nu_e | \nu_\mu \rangle} \\ N_{\langle \nu_\mu | \nu_e \rangle} & N_{\nu_\mu} \end{pmatrix}$$


The diagram shows a horizontal line representing the z-axis with arrows pointing to the right, labeled 'v' at both ends. A vertical dashed line labeled 'z' intersects this axis. Above the intersection, a red circle with a counter-clockwise arrow contains the Greek letter nu (ν).

Mass term in flavor basis: causes vacuum oscillations

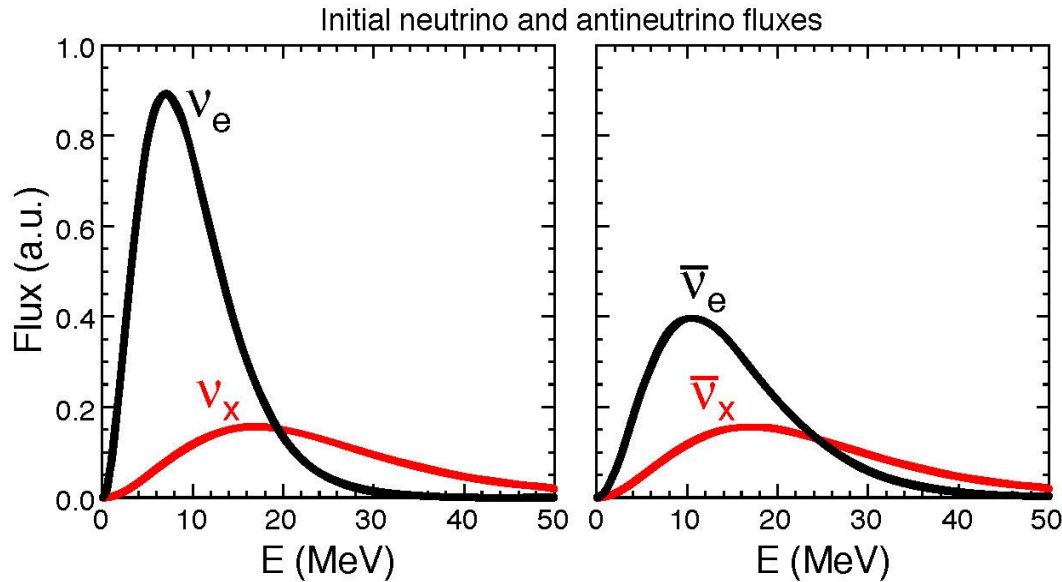
Wolfenstein's weak potential, causes MSW "resonant" conversion together with vacuum term

Flavor-off-diagonal potential, caused by flavor oscillations. (J.Pantaleone, PLB 287:128,1992)

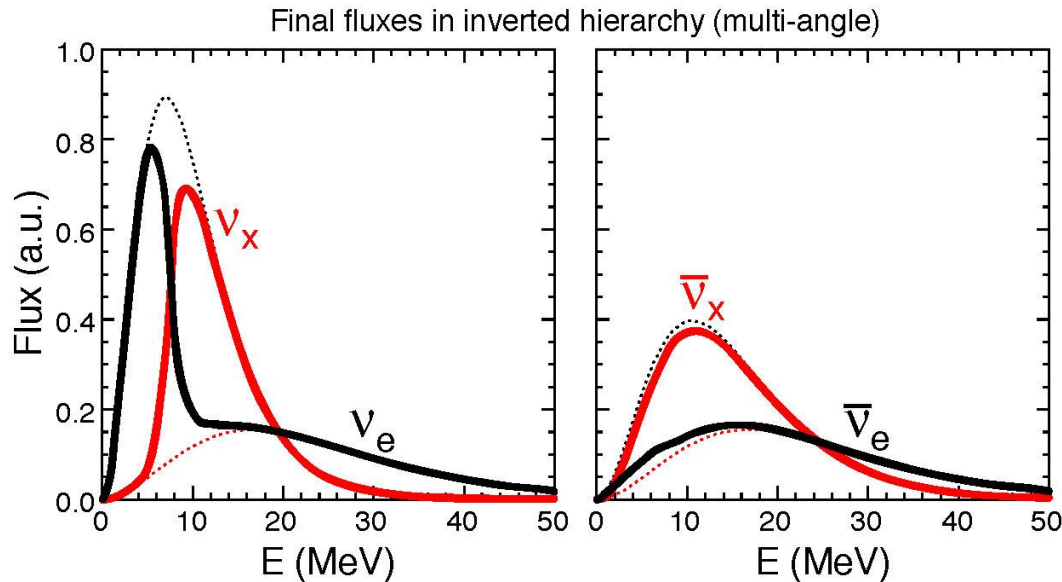
Flavor oscillations feed back on the Hamiltonian: Nonlinear effects!

Spectral Split

Initial
fluxes at
neutrino
sphere



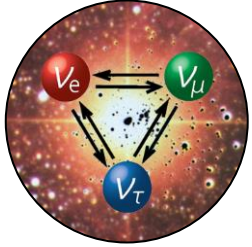
After
collective
trans-
formation



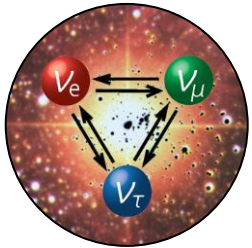
Figures from
Fogli, Lisi,
Marrone & Mirizzi,
arXiv:0707.1998

Explanations in
Raffelt & Smirnov
arXiv:0705.1830
and 0709.4641
Duan, Fuller,
Carlson & Qian
arXiv:0706.4293
and 0707.0290

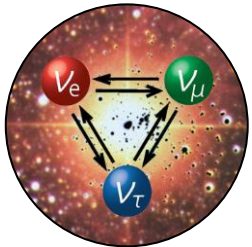
Self-Induced Flavor Conversion



No net flavor conversion of ensemble
(in contrast to MSW conversion)

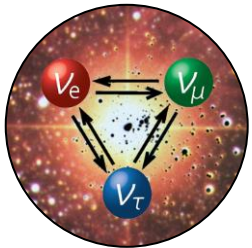


Flavor content exchanged
between different momentum modes
(or nus and anti-nus changing together)



Interacting neutrino system: Coupled oscillators

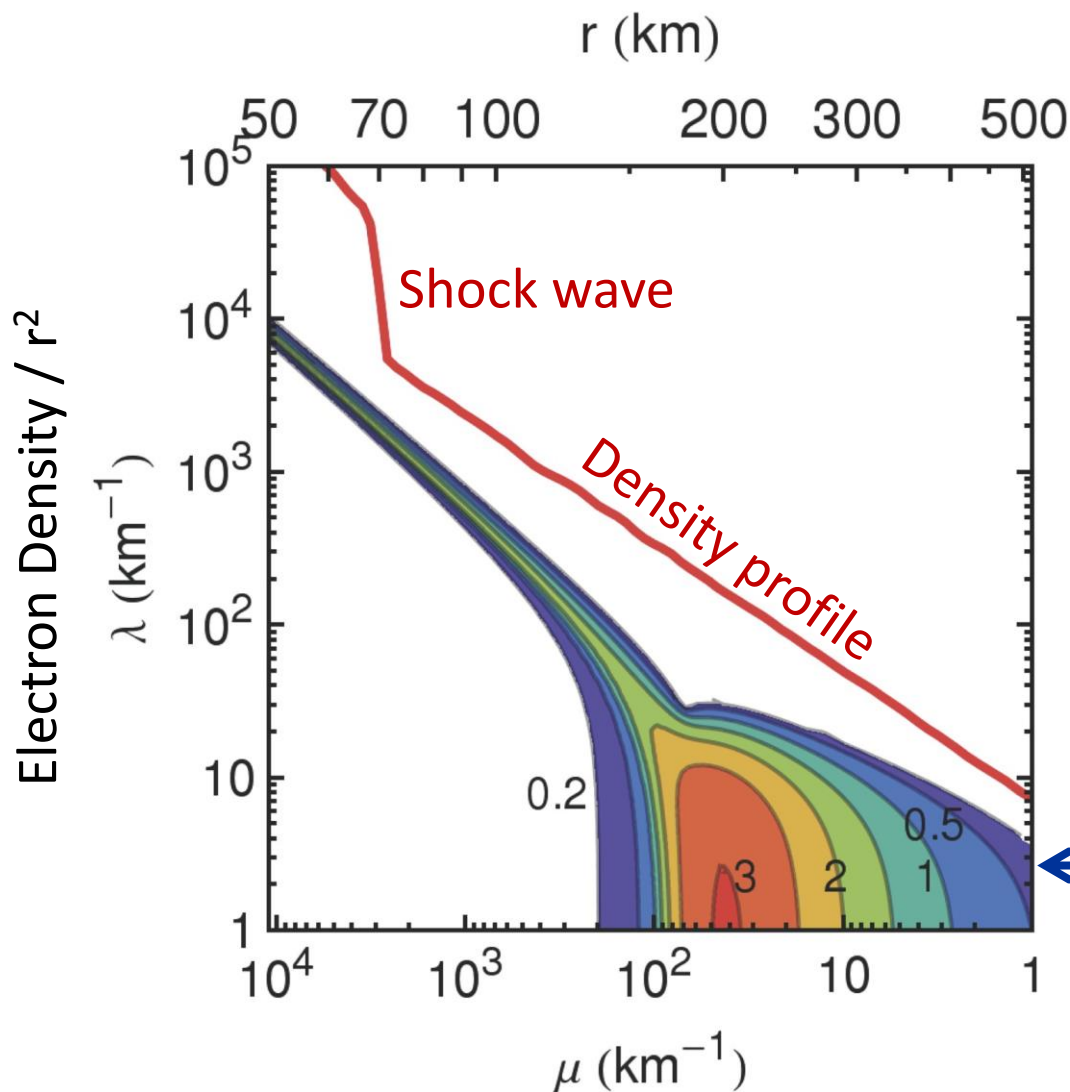
- Collective harmonic oscillation modes
- Exponential run-away modes



Instability required to get started

- Exponentially growing off-diagonals in density matrix
- Linearized stability analysis to find growing modes

Multi-Angle Multi-Energy Stability Analysis



The studied $15 M_{\odot}$ accretion-phase models (Garching) are stable against collective flavor conversion (2-flavor, inverted hierarchy)

Contours of growth rate κ [km^{-1}]

Sarikas, Raffelt, Hüdepohl & Janka, arXiv:1109.3601

Multi-Angle Matter Effect

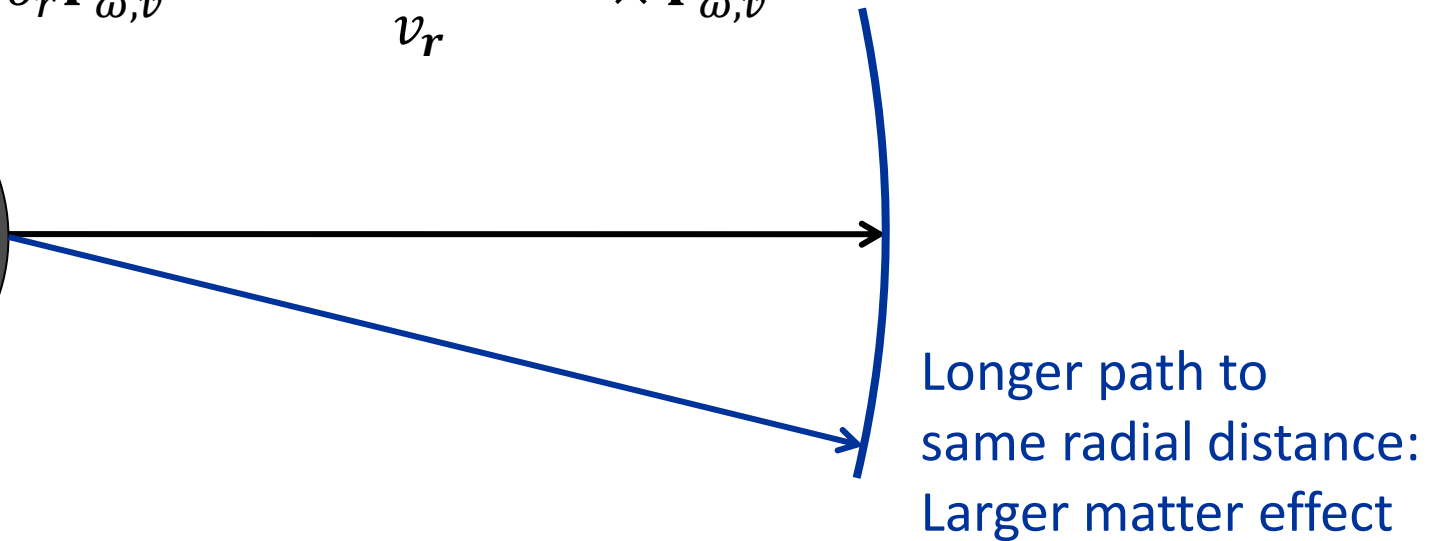
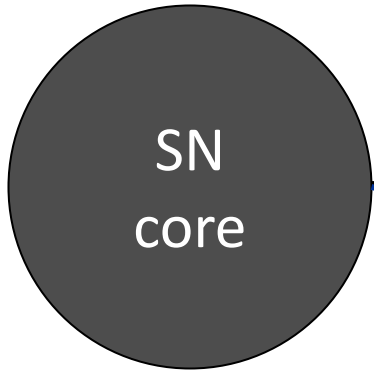
Liouville form of oscillation equation

$$\dot{\mathbf{P}}_{\omega, \mathbf{v}} + (\mathbf{v} \cdot \nabla_r) \mathbf{P}_{\omega, \mathbf{v}} = (\omega \mathbf{B} + \lambda \mathbf{L} + \mu \mathbf{P}) \times \mathbf{P}_{\omega, \mathbf{v}}$$

Drops out for stationary solutions

$$\begin{array}{cc} \uparrow & \uparrow \quad \uparrow \\ \sqrt{2} G_F N_e & \sqrt{2} G_F N_\nu \end{array}$$

$$\partial_r \mathbf{P}_{\omega, \mathbf{v}} = \frac{\omega \mathbf{B} + \lambda \mathbf{L} + \mu \mathbf{P}}{v_r} \times \mathbf{P}_{\omega, \mathbf{v}}$$

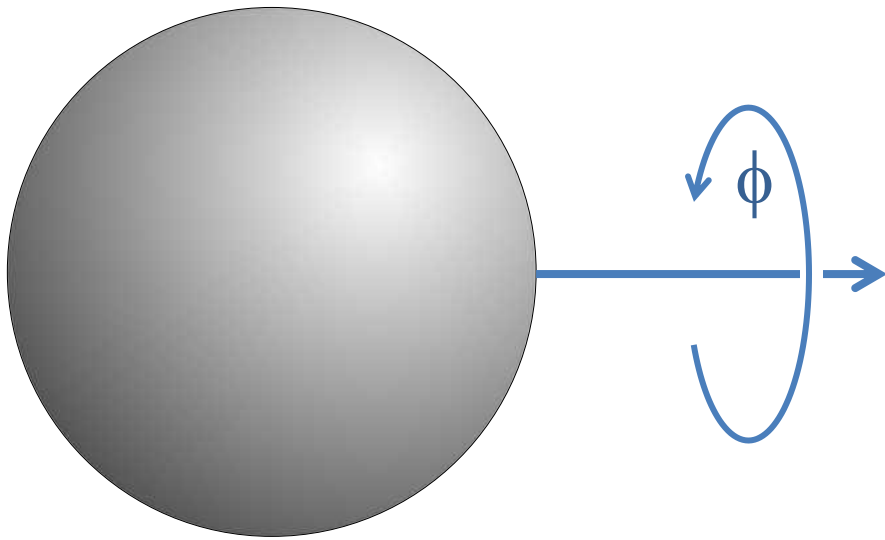


Self-induced conversion suppressed for $N_e \gtrsim N_\nu$

Esteban-Pretel, Mirizzi, Pastor, Tomàs, Raffelt, Serpico & Sigl, arXiv:0807.0659

Symmetry Breaking in Collective Flavor Oscillations

Assume globally spherically symmetric neutrino emission from SN core
→ Axial symmetry in chosen direction



Self-induced neutrino flavor conversion in both hierarchies (unless suppressed by multi-angle matter effect)

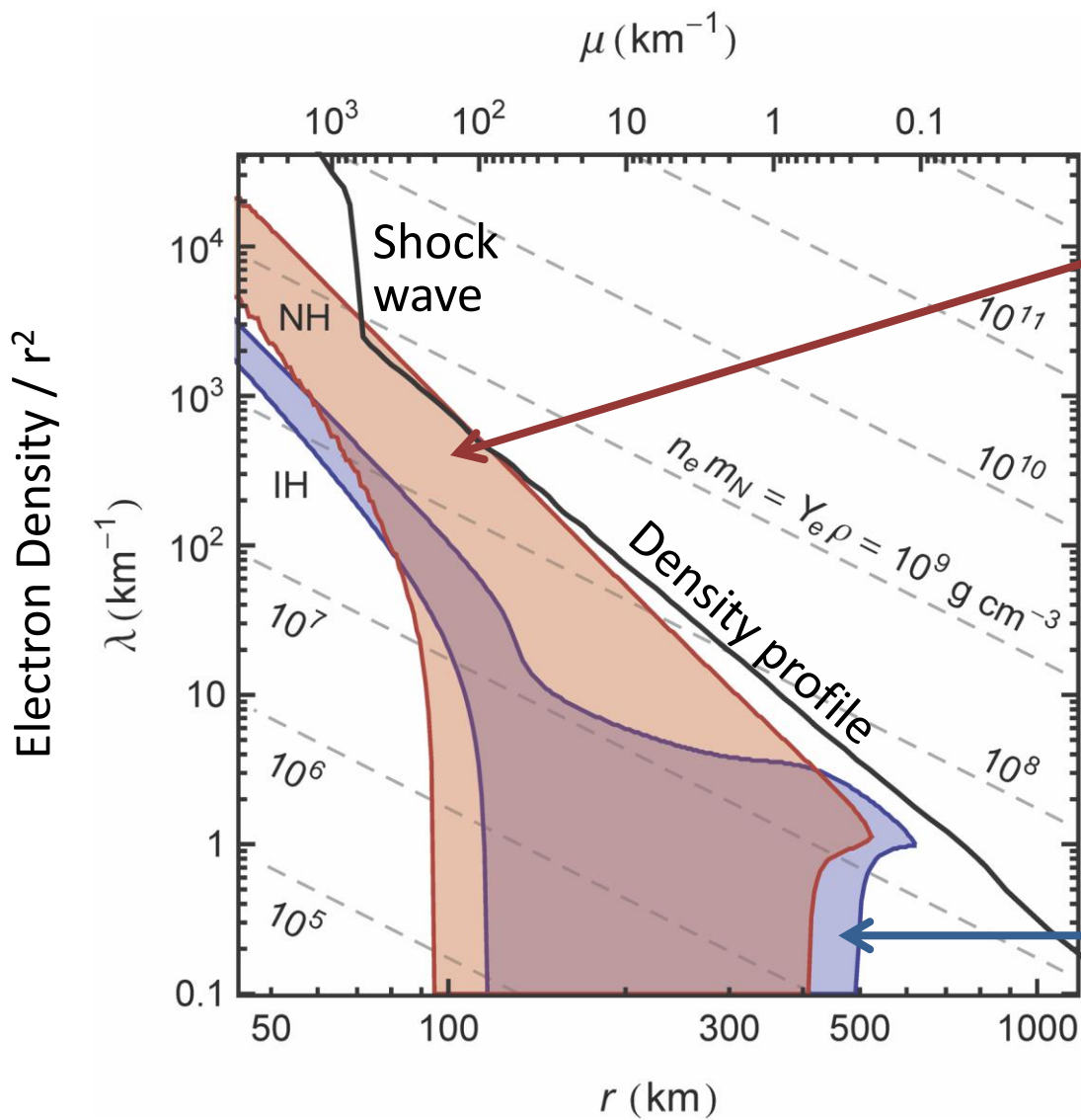
- **Axially symmetric solution:**
Conversion for inverted hierarchy
(usual result)
- **Spontaneous breaking of axial symmetry:**
Dipole solution ($\propto \cos \phi$ or $\sin \phi$)
Conversion for normal hierarchy
(Was missed by enforcing axial symmetry because of axially symmetric emission)

G. Raffelt, S. Sarikas & D. de Sousa Seixas

Axial symmetry breaking in self-induced flavor conversion of SN neutrino fluxes

PRL 111 (2013) 091101 [arXiv:1305.7140]

Instability Footprints

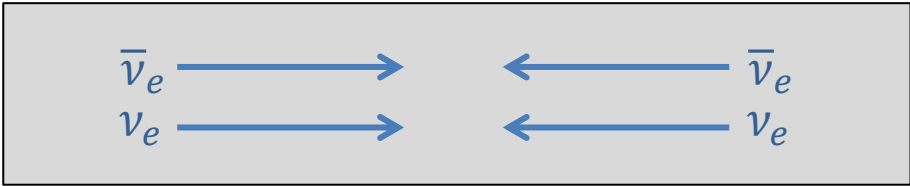
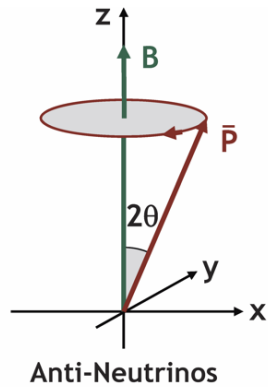
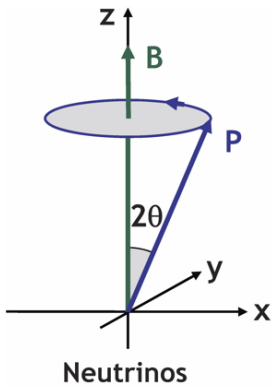


Axial-symmetry breaking (MAA) instability (normal ordering NH) is "more dangerous" to trigger self-induced flavor conversion

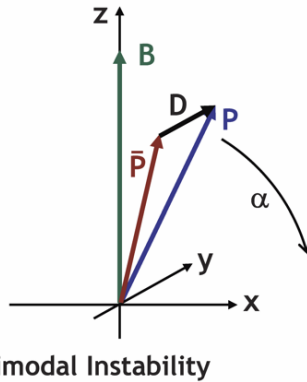
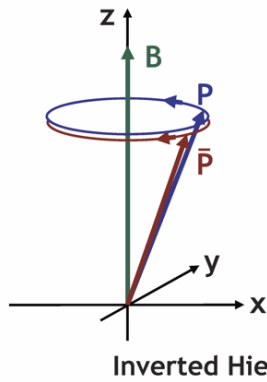
Traditional "bimodal" instability (inverted mass ordering IH)

Raffelt, Sarikas & Seixas, arXiv:1305.7140

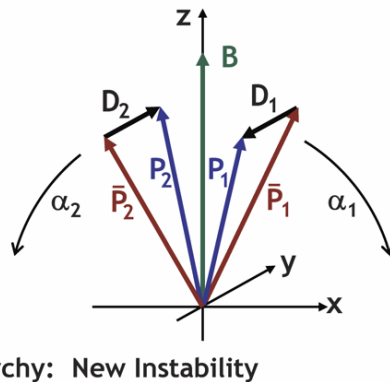
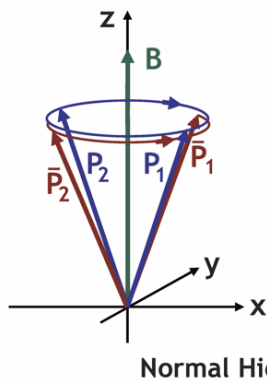
Colliding Beam Model



Raffelt & Seixas, arXiv:1307.7625



Left- and right-moving neutrinos
behave symmetrically
Instability for inverted mass ordering (IH)



Left-right symmetry breaking:
- Anti-symmetric mode for
normal mass ordering (NH)
- Corresponds to axial symmetry breaking
in SN case (MAA instability)

Symmetry Assumptions

Neutrino transport and flavor oscillations: 7D problem

$$(\partial_t + \vec{v} \cdot \vec{\nabla}_x + \vec{F} \cdot \vec{\nabla}_p) \rho(t, \vec{x}, \vec{p}) = -i [H(t, \vec{x}, \vec{p}), \rho(t, \vec{x}, \vec{p})] + \mathcal{C}[\rho(t, \vec{x}, \vec{p})]$$

↑
Ignore external forces
(e.g. no grav. deflection)

↑
Includes vacuum, matter,
nu-nu refraction

↑
Ignore collision term:
Free streaming

- **Homogeneous, isotropic system evolving in time (“early universe”)**
or 1D homogeneous evolving in time (“colliding beams”)

$$\partial_t \rho(t, E) = -i [H(t, E), \rho(t, E)]$$

- **Stationary, spherically symmetric, evolving with radius (“supernova”)**

$$v_r \partial_r \rho(r, E, \theta) = -i [H(r, E, \theta), \rho(r, E, \theta)]$$

↑

↑
Zenith angle of nu momentum \vec{p}

Radial velocity depends on θ , leads to multi-angle matter effect

- Ordinary differential equations in “time” or “radius” with maximal symmetries
- Misses dominant solutions (spontaneous symmetry breaking)

Spatial Symmetry Breaking (SSB)

Oscillation equation with explicit transport term

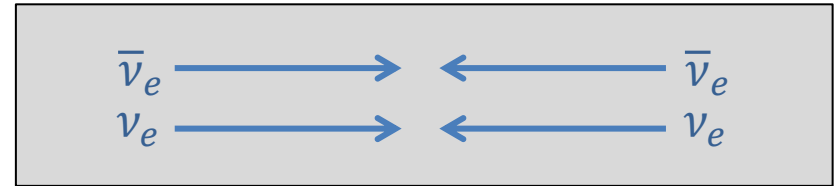
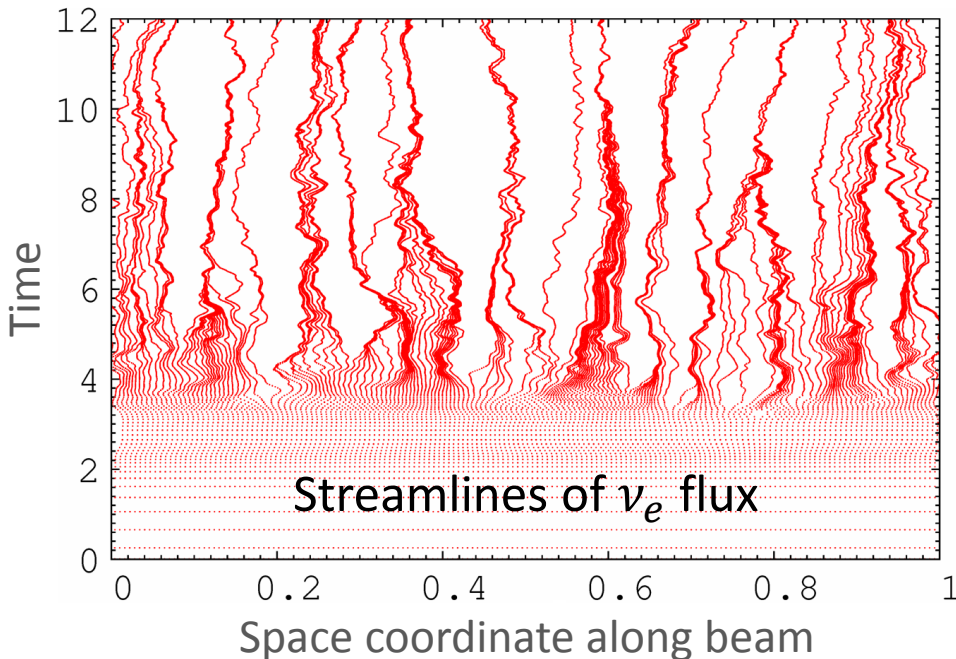
$$i \underbrace{(\partial_t + \vec{v} \cdot \vec{\nabla}_x)} \rho(t, \vec{x}, \vec{p}) = [H(t, \vec{x}, \vec{p}), \rho(t, \vec{x}, \vec{p})]$$

Without flavor oscillations: free streaming

Spatial Fourier transform (plane wave expansion)

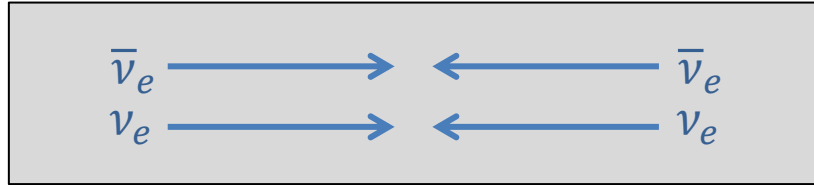
$$(i\partial_t + \vec{v} \cdot \vec{k}) \rho(t, \vec{k}, \vec{p}) = \int d^3\vec{x} e^{-i\vec{k}\cdot\vec{x}} [H(t, \vec{x}, \vec{p}), \rho(t, \vec{x}, \vec{p})]$$

Interaction term couples different Fourier modes



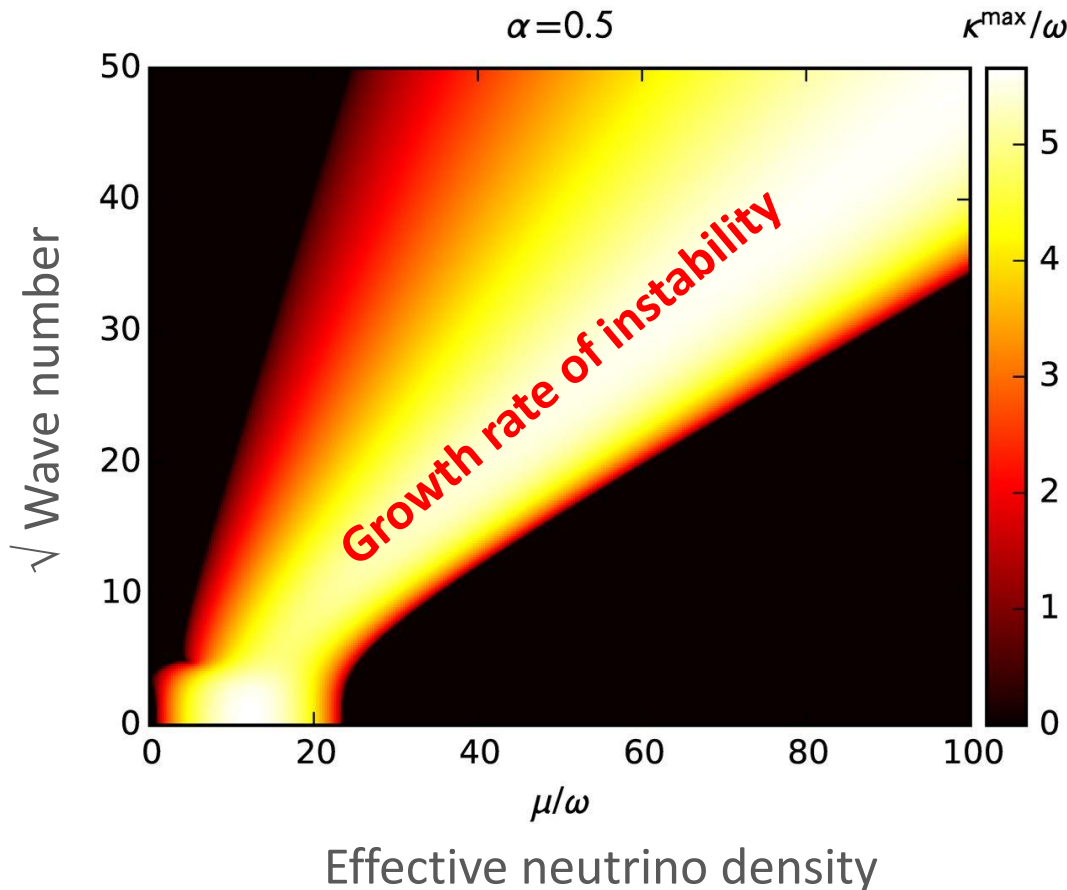
Mirizzi, Mangano & Saviano
arXiv:1503.03485

Spatial Symmetry Breaking (SSB)



Linearized stability analysis for colliding-beam model

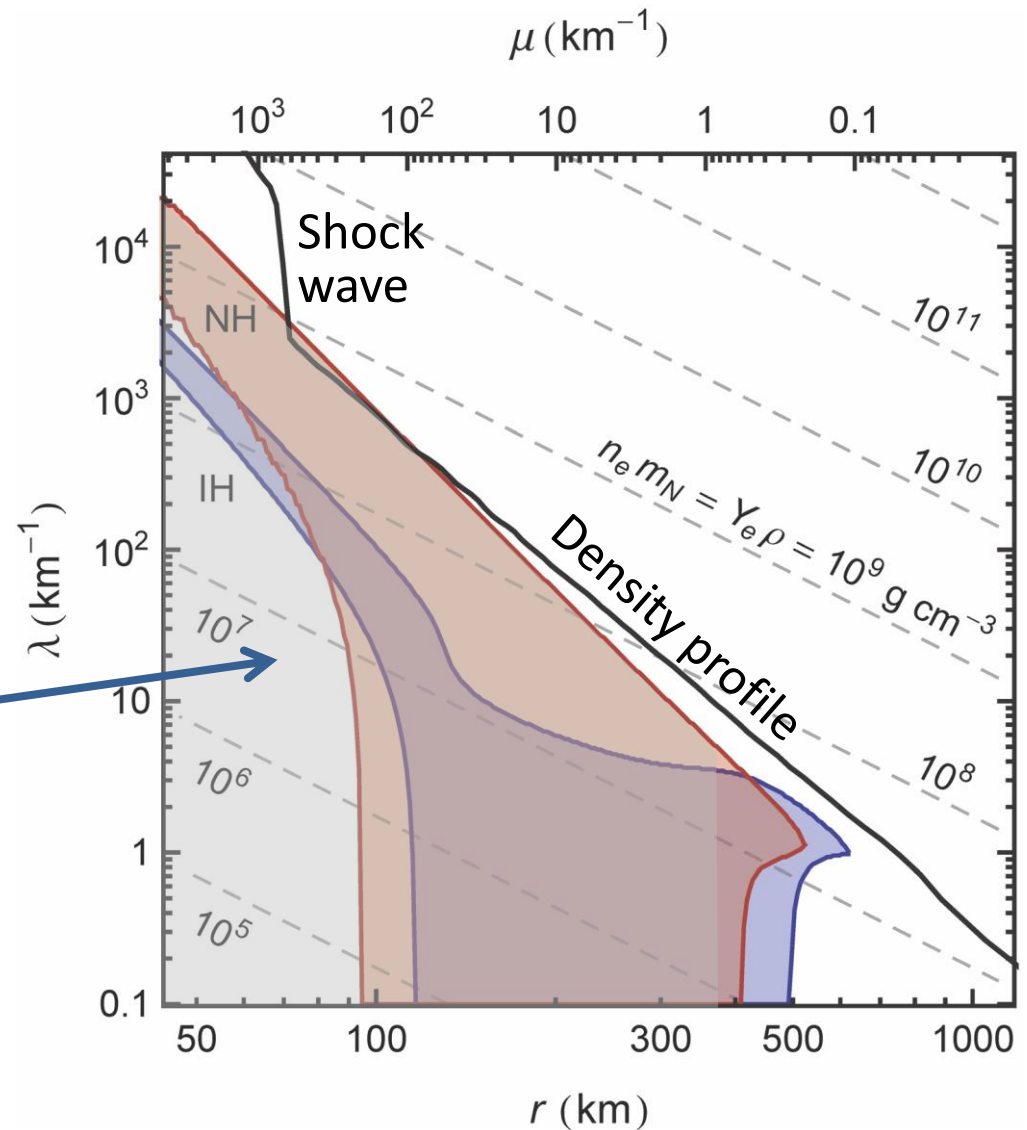
Duan & Shalgar, arXiv:1412.7097



- Instability footprint shifted to larger neutrino density μ for larger wave number k
- For any neutrino density, unstable for some k -range
- No flavor-stable conditions exist for homogeneous neutrino gas (no "sleeping top" regime)

Small-Scale Instabilities

- Small-scale modes “fill in” the stability footprint for large neutrino density
- Largest-scale mode is “most dangerous” to cross SN density profile



Chakraborty, Hansen, Izaguirre & Raffelt, Work in progress (2015)

Space-Time-Dependent Problem in Supernova

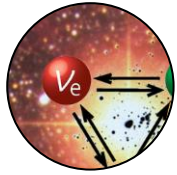
- Neutrino momentum distribution not limited to “outward” direction
- Important “halo” flux even at large distance
- Large 3D effects

→ Inhomogeneous, anisotropic, non-stationary problem

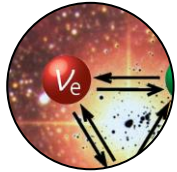
Really no self-induced flavor conversion below shock-wave or even below ν -sphere?

- Investigations to date are simplified case studies
- May not represent real SNe

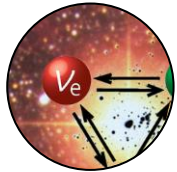
Status of Collective Flavor Conversion



Self-induced flavor conversion is an instability in flavor space of the interacting neutrino ensemble



Space-time dependent phenomenon (not simply stationary or homogeneous)



Solutions do not respect symmetries of initial system
Instabilities can occur on all scales

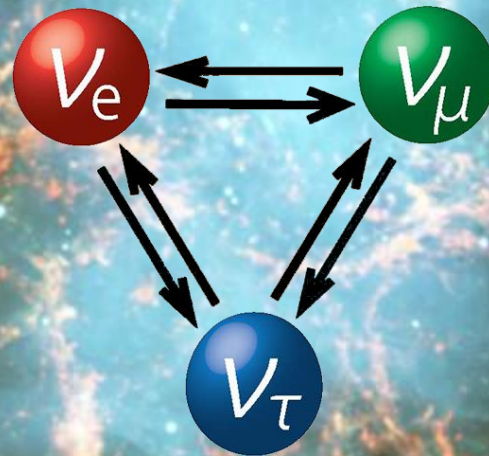


**Essentially back
to the drawing board ...**

Literature on Spatial Symmetry Breaking (SSB)

1. Axial symmetry breaking in self-induced flavor conversion of SN neutrino fluxes
Raffelt, Sarikas & Seixas, PRL 111 (2013) 091101
2. Neutrino flavor pendulum in both mass hierarchies
Raffelt & Seixas, PRD 88 (2013) 045031
3. Chaotic flavor evolution in an interacting neutrino gas
Hansen & Hannestad, PRD 90 (2014) 025009
4. Damping the neutrino flavor pendulum by breaking homogeneity
Mangano, Mirizzi & Saviano, PRD 89 (2014) 073017
5. Spontaneous breaking of spatial symmetries in collective neutrino oscillations
Duan & Shalgar, arXiv:1412.7097
6. Self-induced flavor instabilities of a dense neutrino stream in a 2D model
Mirizzi, Mangano & Saviano, arXiv:1503.03485
7. Self-induced flavor-conversion anisotropy of supernova neutrinos
Chakraborty, Hansen, Izaguirre & Raffelt, work in progress (2015)

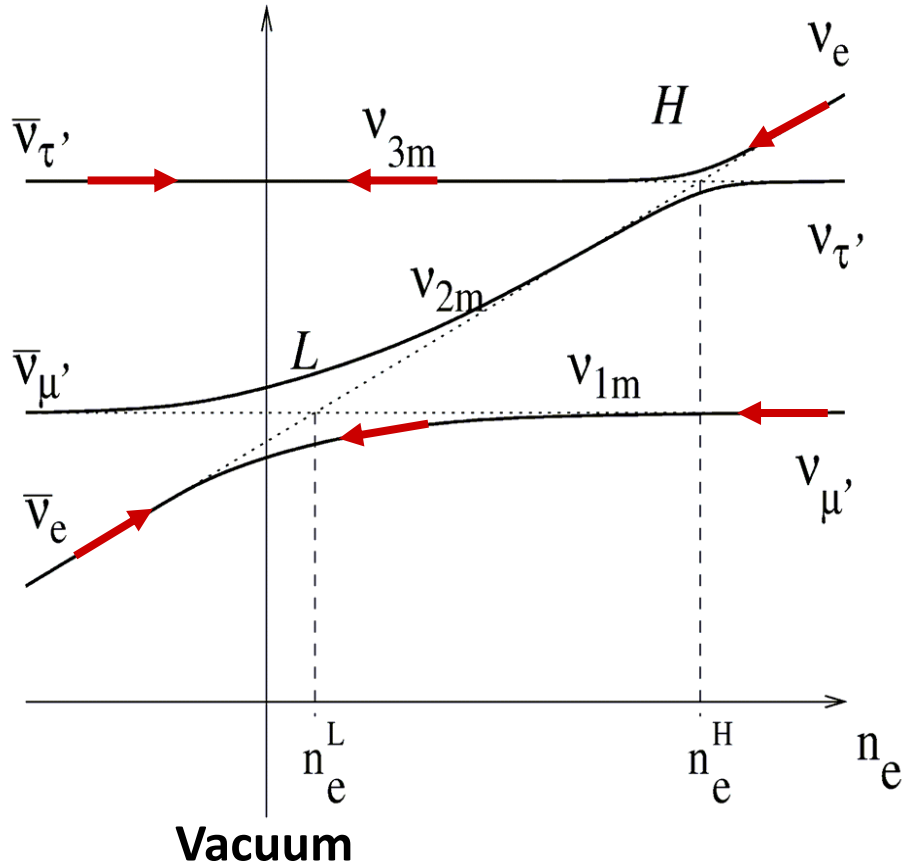
More theory progress is needed to understand
flavor conversion of supernova neutrinos!



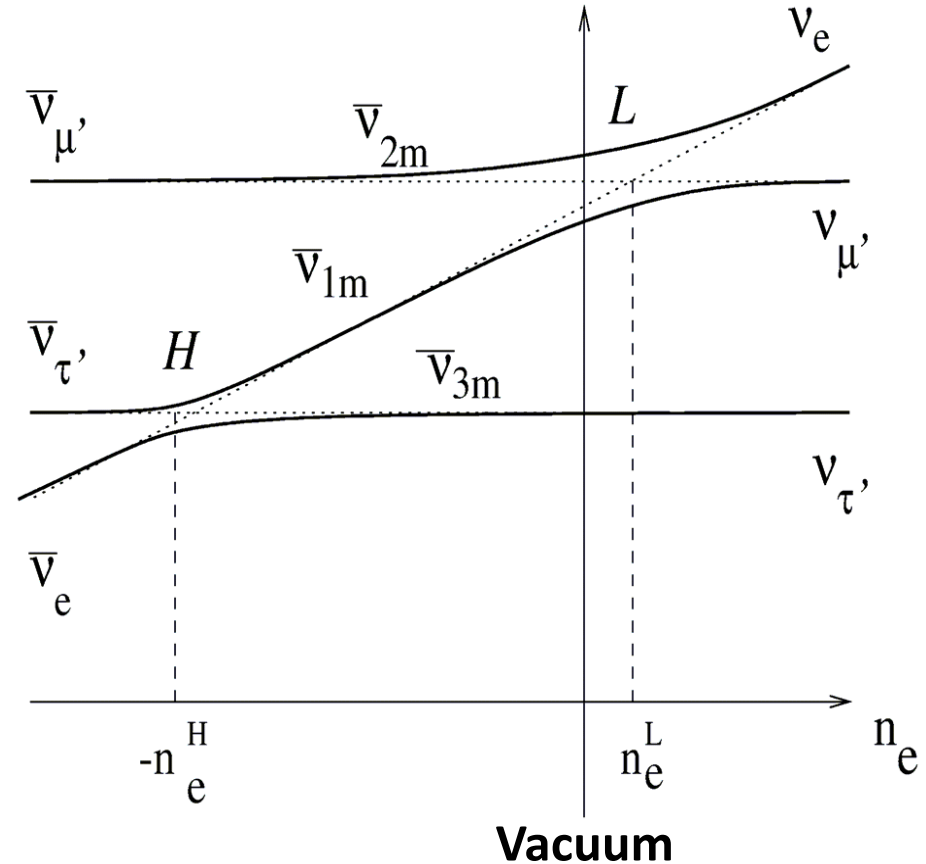
Backup

Three-Flavor Eigenvalue Diagram

Normal mass hierarchy



Inverted mass hierarchy



Dighe & Smirnov, Identifying the neutrino mass spectrum from a supernova neutrino burst, astro-ph/9907423

Three Ways to Describe Flavor Oscillations

Schrödinger equation in terms of “flavor spinor”

$$i\partial_t \begin{pmatrix} \nu_e \\ \nu_\mu \end{pmatrix} = H \begin{pmatrix} \nu_e \\ \nu_\mu \end{pmatrix} = \frac{\Delta m^2}{2E} \begin{pmatrix} \cos 2\theta & \sin 2\theta \\ \sin 2\theta & -\cos 2\theta \end{pmatrix} \begin{pmatrix} \nu_e \\ \nu_\mu \end{pmatrix}$$

Neutrino flavor density matrix

$$\rho = \begin{pmatrix} \langle \nu_e | \nu_e \rangle & \langle \nu_e | \nu_\mu \rangle \\ \langle \nu_\mu | \nu_e \rangle & \langle \nu_\mu | \nu_\mu \rangle \end{pmatrix}$$

Equivalent commutator form of Schrödinger equation

$$i\partial_t \rho = [H, \rho]$$

Expand 2×2 Hermitean matrices in terms of Pauli matrices

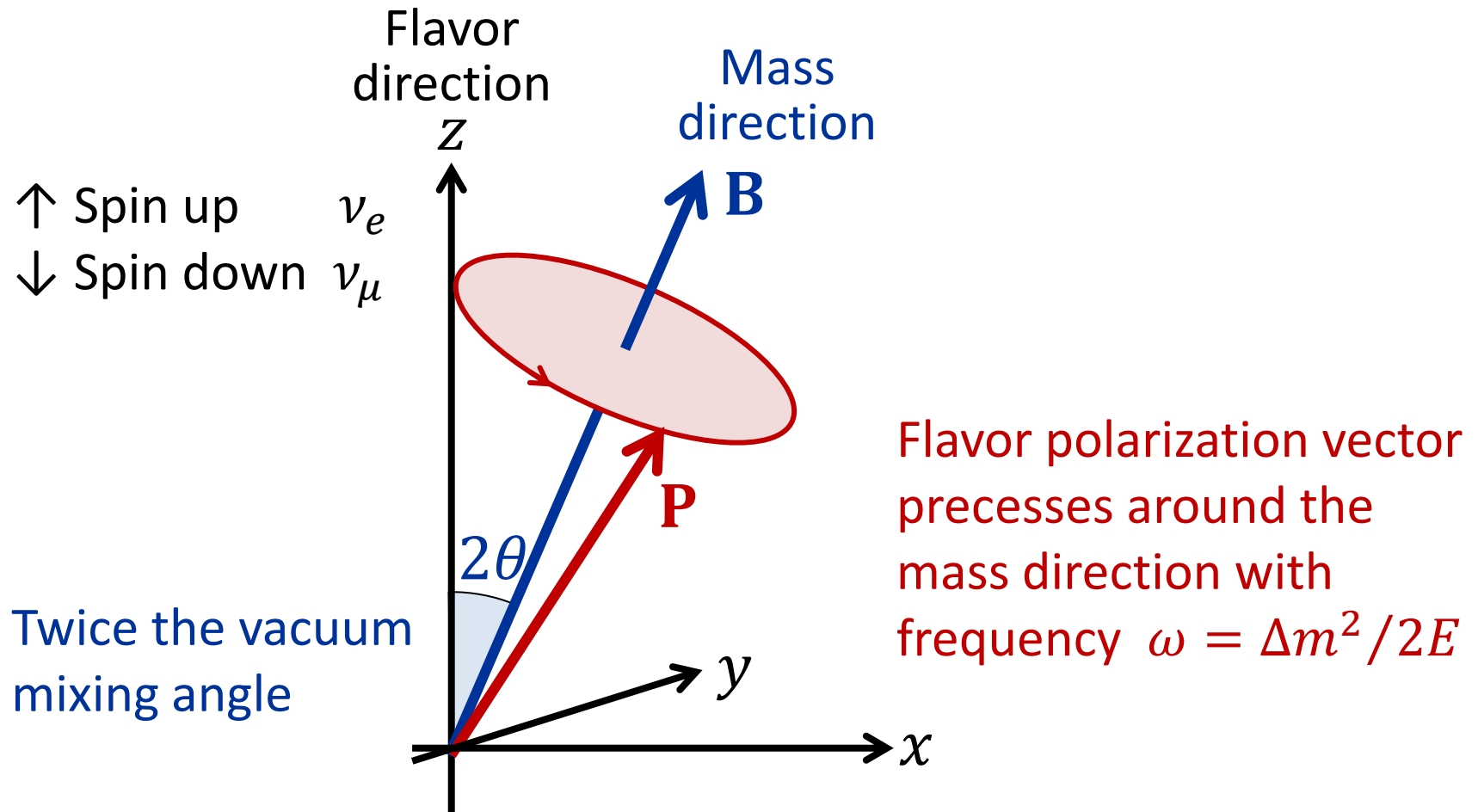
$$\rho = \frac{1}{2} [\text{Tr}(\rho) + \mathbf{P} \cdot \boldsymbol{\sigma}] \quad \text{and} \quad H = \frac{\Delta m^2}{2E} \mathbf{B} \cdot \boldsymbol{\sigma} \quad \text{with} \quad \mathbf{B} = (\sin 2\theta, 0, \cos 2\theta)$$

Equivalent spin-precession form of equation of motion

$$\dot{\mathbf{P}} = \omega \mathbf{B} \times \mathbf{P} \quad \text{with} \quad \omega = \frac{\Delta m^2}{2E}$$

\mathbf{P} is “polarization vector” or “Bloch vector” or “flavor isospin vector”

Flavor Oscillation as Spin Precession

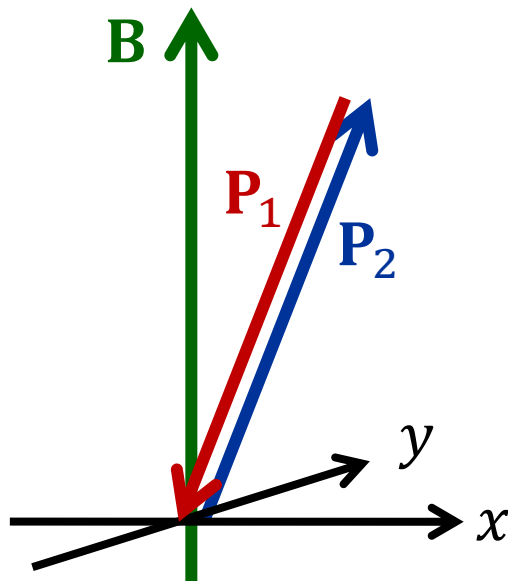


Instability in Flavor Space

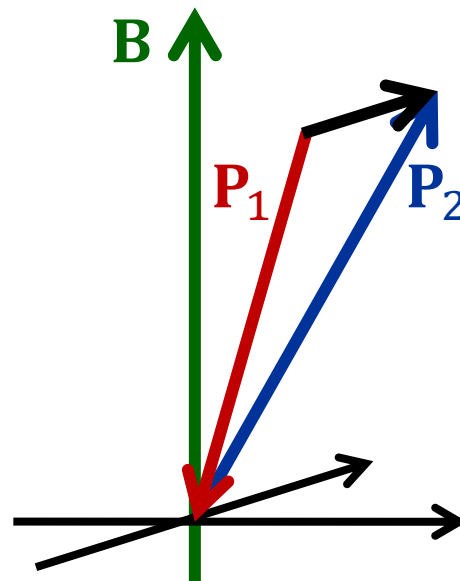
Two-mode example in co-rotating frame, initially $\mathbf{P}_1 = \downarrow$, $\mathbf{P}_2 = \uparrow$ (flavor basis)

$$\dot{\mathbf{P}}_1 = [-\omega \mathbf{B} + \mu (\mathbf{P}_1 + \mathbf{P}_2)] \times \mathbf{P}_1$$

$$\dot{\mathbf{P}}_2 = [+ \omega \mathbf{B} + \mu \underbrace{(\mathbf{P}_1 + \mathbf{P}_2)}_{0 \text{ initially}}] \times \mathbf{P}_2$$



- Initially aligned in flavor direction and $\mathbf{P} = 0$
- Free precession $\pm \omega$



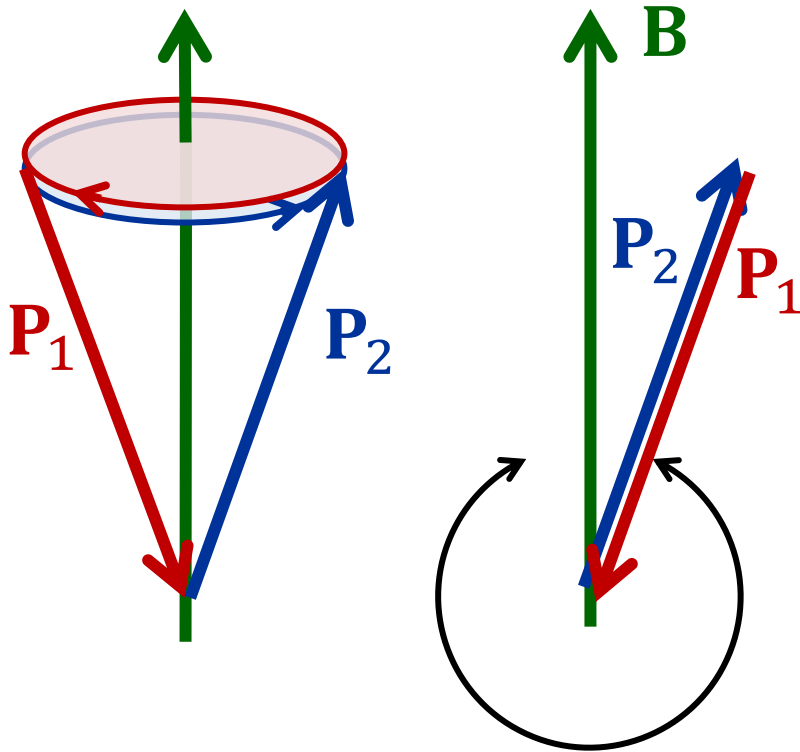
- After a short time, transverse \mathbf{P} develops by free precession

$$\mathbf{P} = \mathbf{P}_1 + \mathbf{P}_2$$

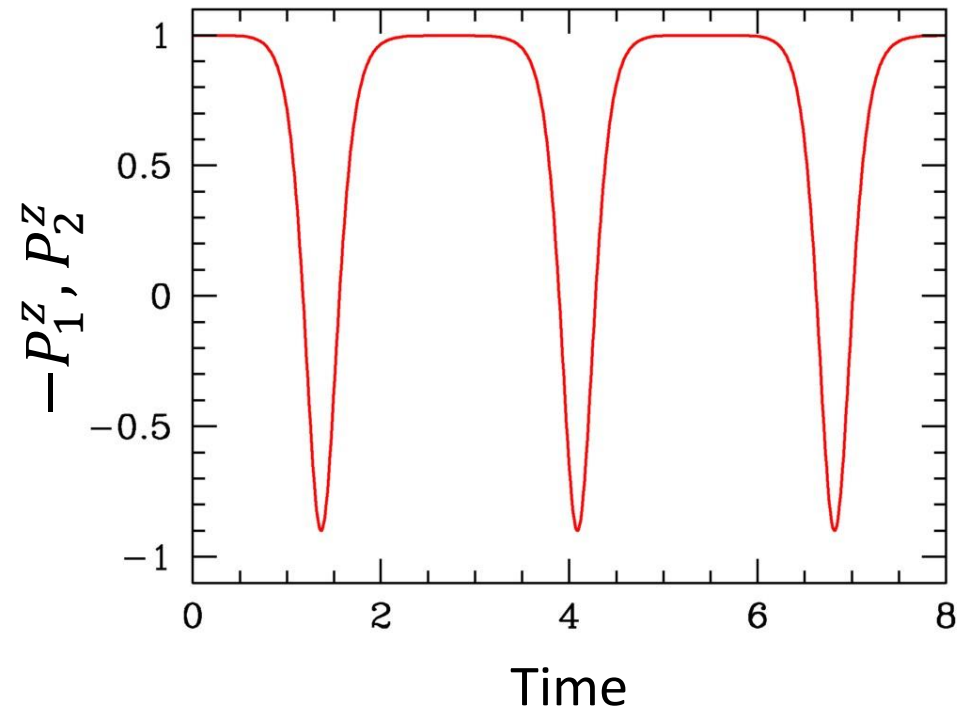
Matter effect transverse to mass and flavor directions
Both \mathbf{P}_1 and \mathbf{P}_2 tilt around \mathbf{P} if μ is large

Two Spins with Opposite Initial Orientation

No interaction ($\mu = 0$)
Free precession in
opposite directions

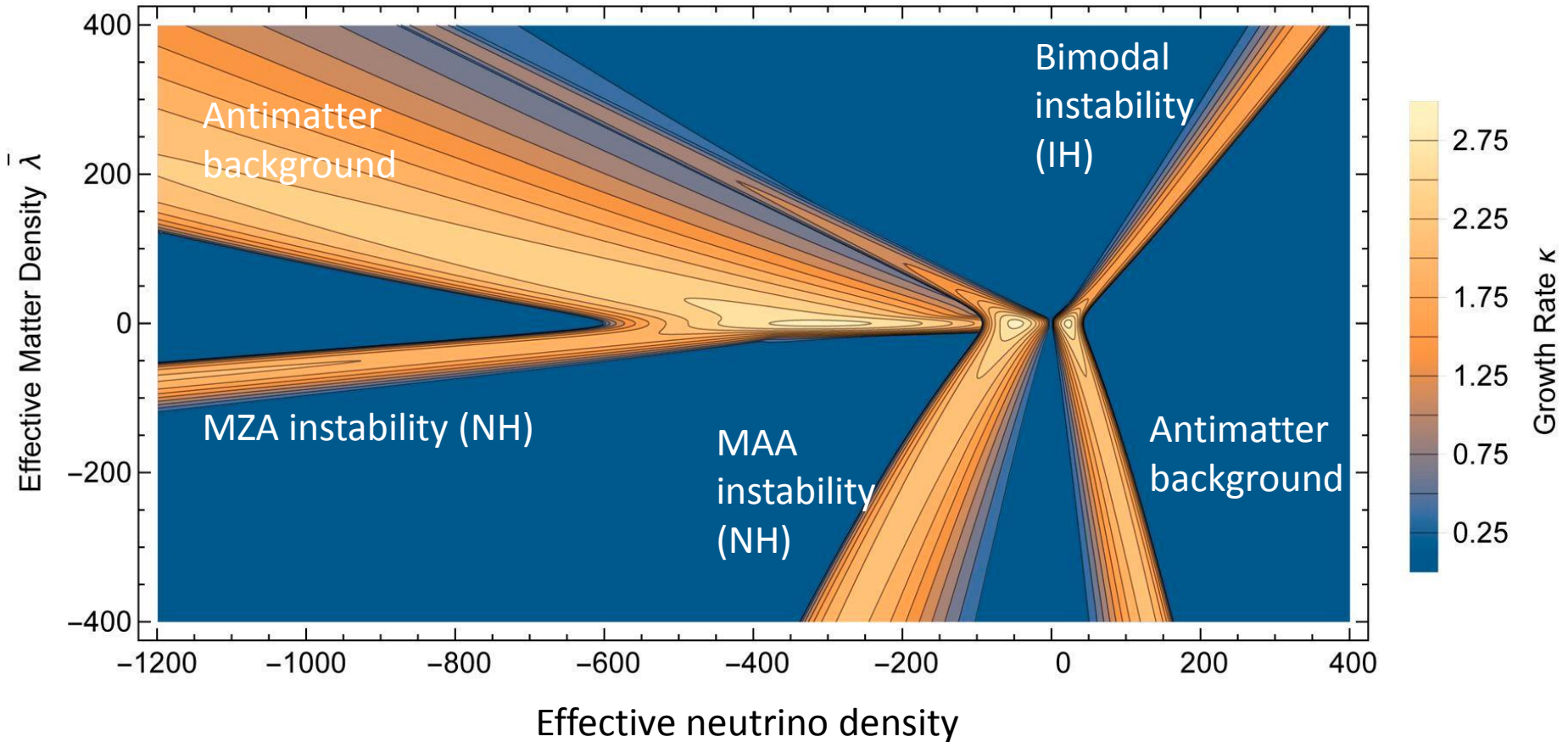


Even for very small mixing angle,
large-amplitude flavor oscillations



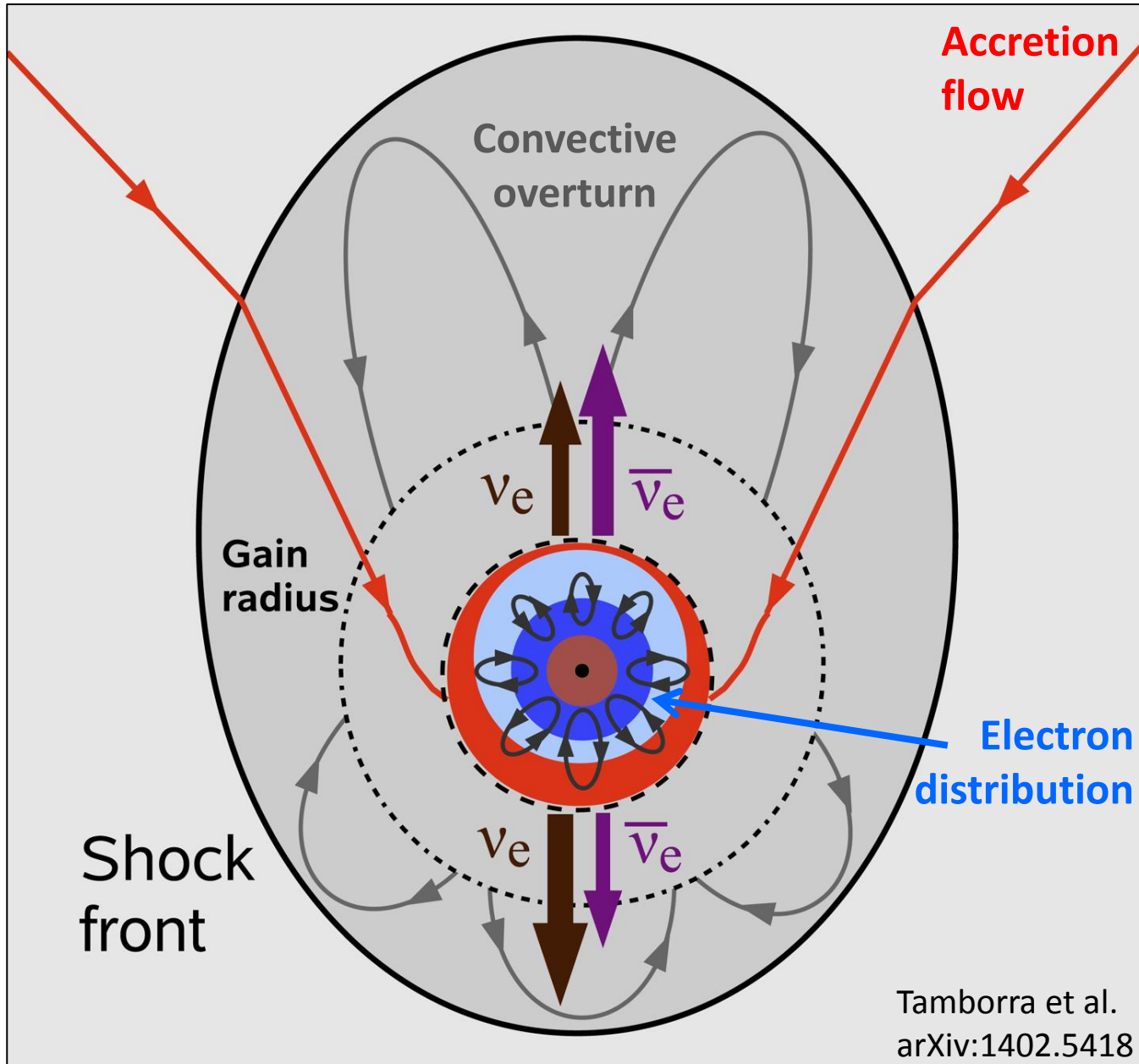
Strong interaction
($\mu \rightarrow \infty$)
Pendular motion

Multi-Angle Matter Effect in Supernovae



Chakraborty, Hansen, Izaguirre & Raffelt, Work in progress (2015)

Schematic Theory of LESA

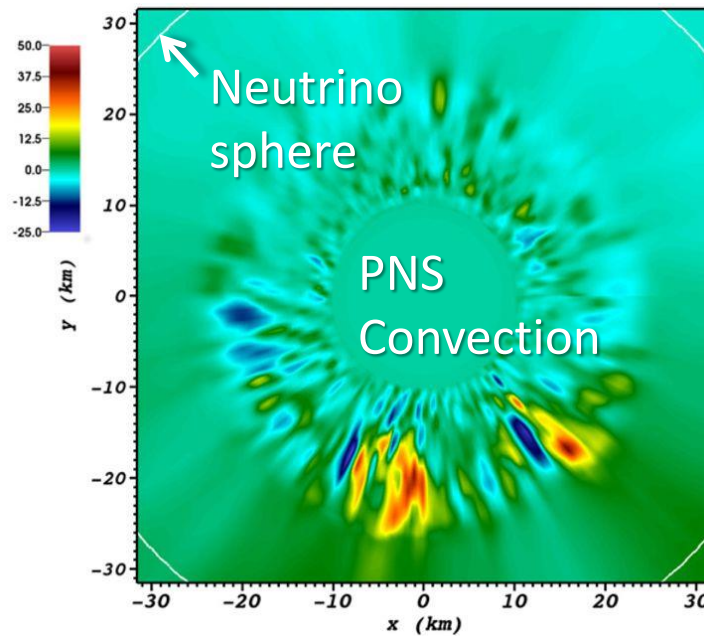
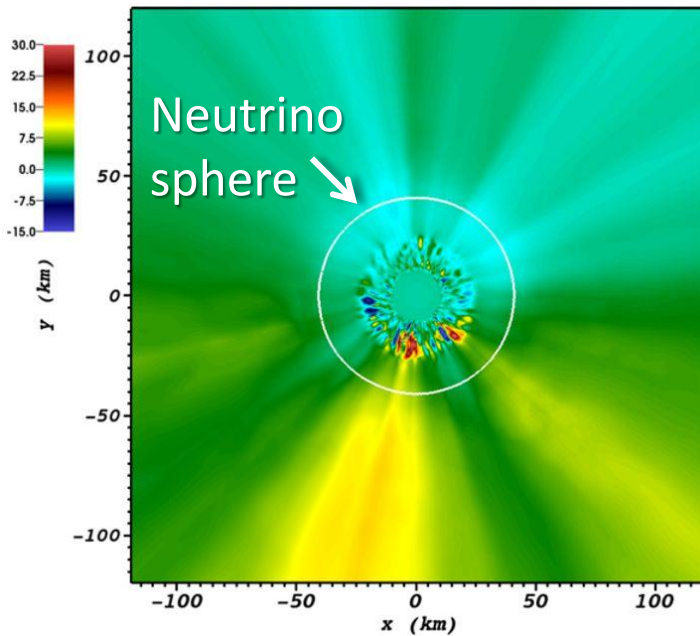


Feedback loop consists of asymmetries in

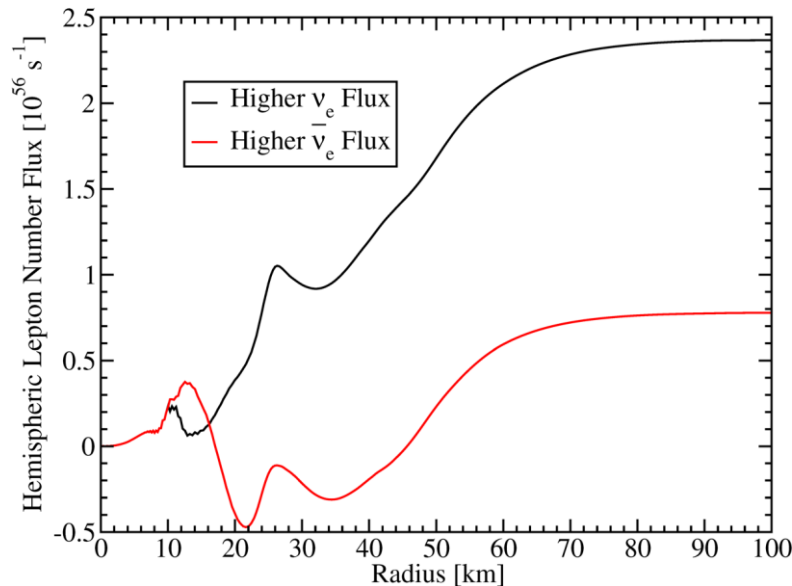
- accretion rate
- lepton-number flux
- neutrino heating rate
- dipole deformation of shock front

Tamborra et al.
arXiv:1402.5418

LESA Dipole and PNS Convection



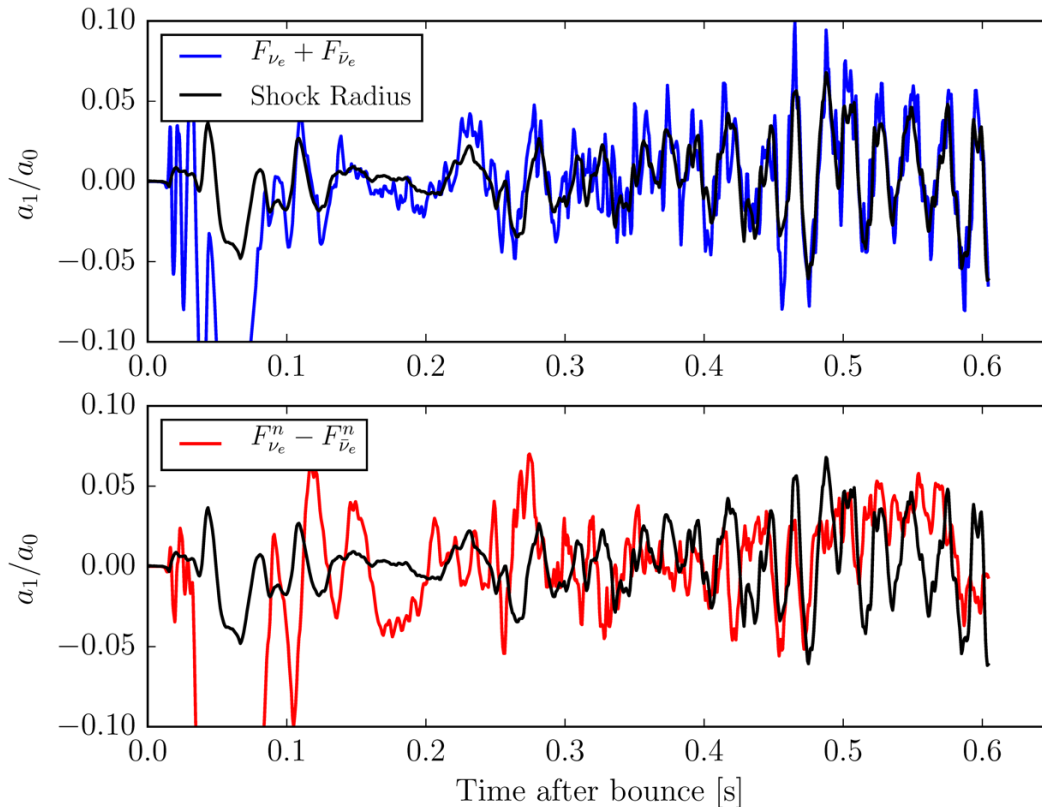
Color-coded
lepton-number flux
along radial rays
(11.2 M_{SUN} model
at 210 ms)



**Lepton flux dipole builds up mostly
below the neutrino sphere
in a region of strong convection
in the proto-neutron star (PNS)**

Is the LESA Phenomenon Real?

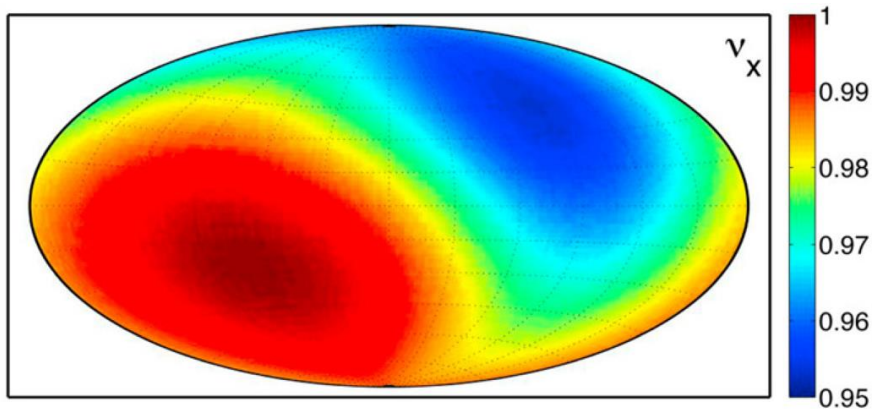
- Couch & O'Connor (2014) also find LESA in their 3D models
- Dolence, Burrows & Zhang (arXiv:1403.6115), 2D models: No LESA dipole at all



Red curve:
Lepton-number dipole $\times 5$
No evidence for beyond-noise
dipole evolution
(Fig.11 of arXiv:1403.6115)

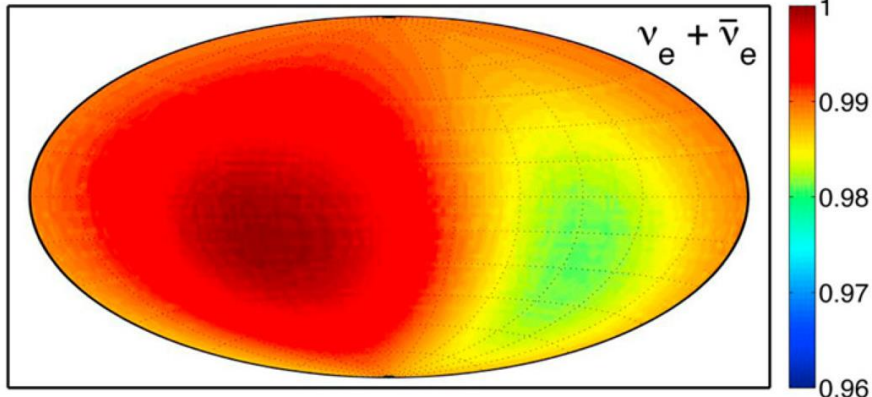
Different method of neutrino radiative transfer, different interaction rates, and many other physics differences — needs to be understood

Sky Distribution of Number Fluxes (11.2 M_{SUN})

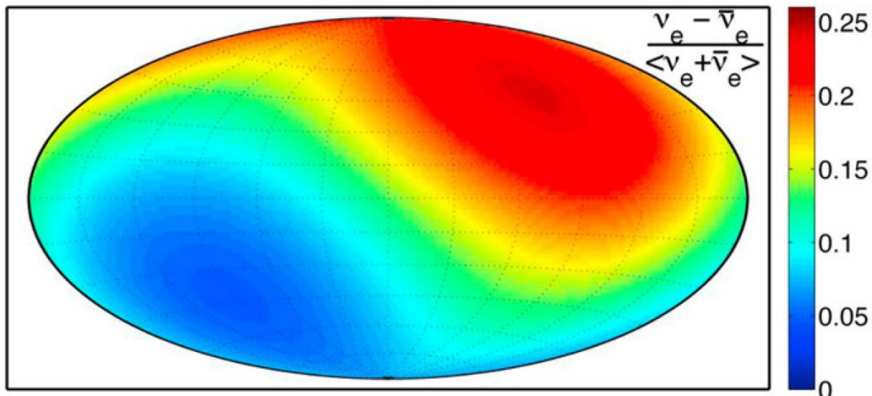


Neutrino number flux distribution
for 11.2 M_{SUN} model
integrated over 150–250 ms

Heavy-flavor neutrino fluxes (ν_x)
nearly isotropic

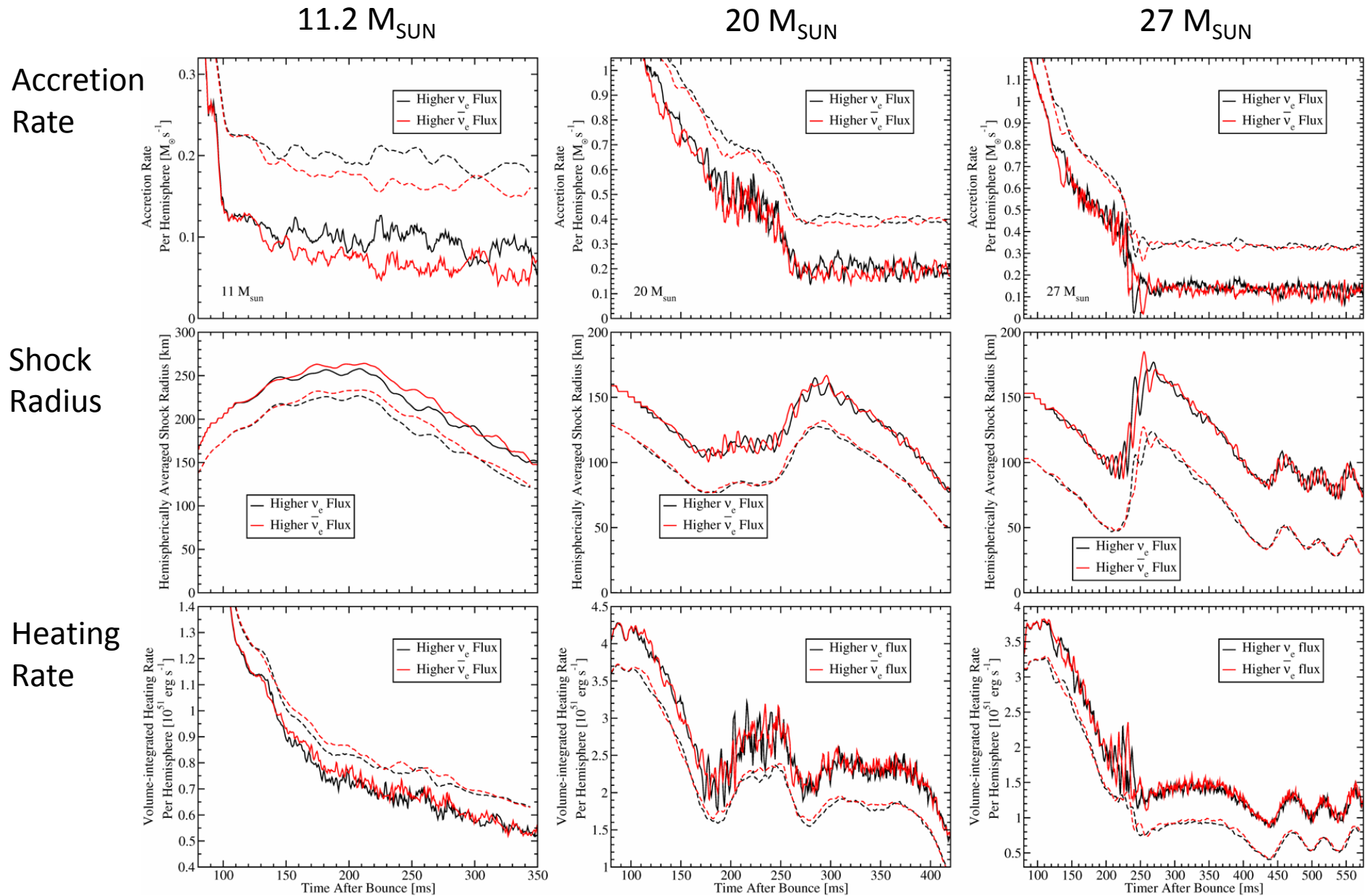


Flux of $\nu_e + \bar{\nu}_e$ nearly isotropic



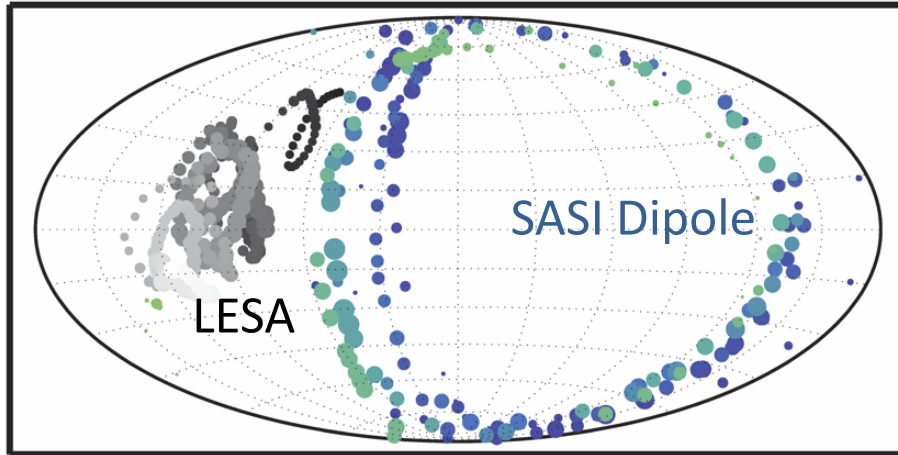
Lepton-number flux ($\nu_e - \bar{\nu}_e$)
has strong dipole distribution

Asymmetries of Elements Relevant for LESA

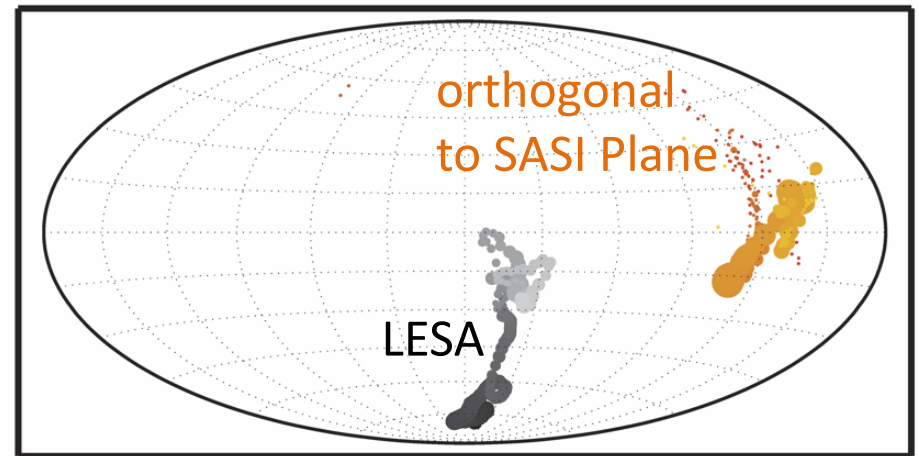
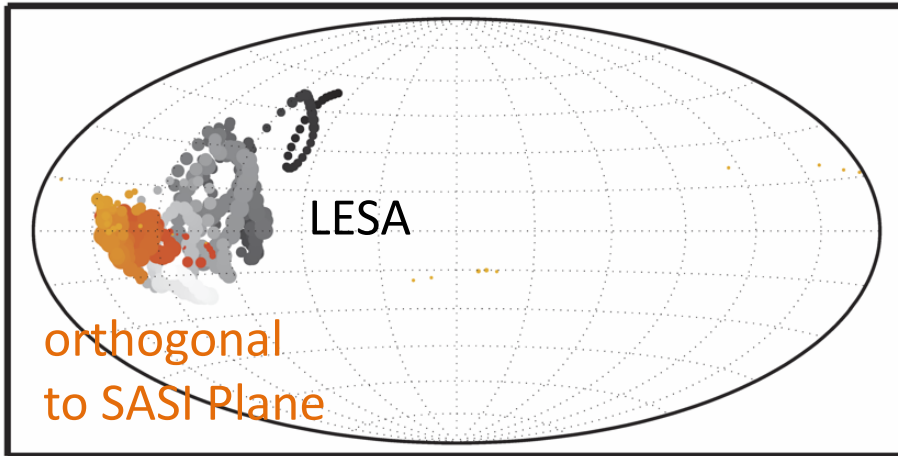
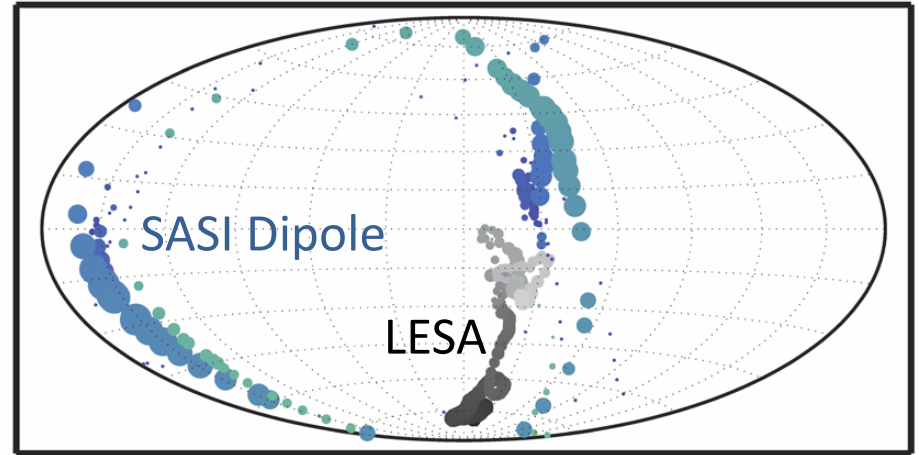


LESA vs. SASI Dipole Motions

20 M_{sun} , [170,300] ms



27 M_{sun} , [170,260] ms



No apparent directional correlation between SASI and LESA