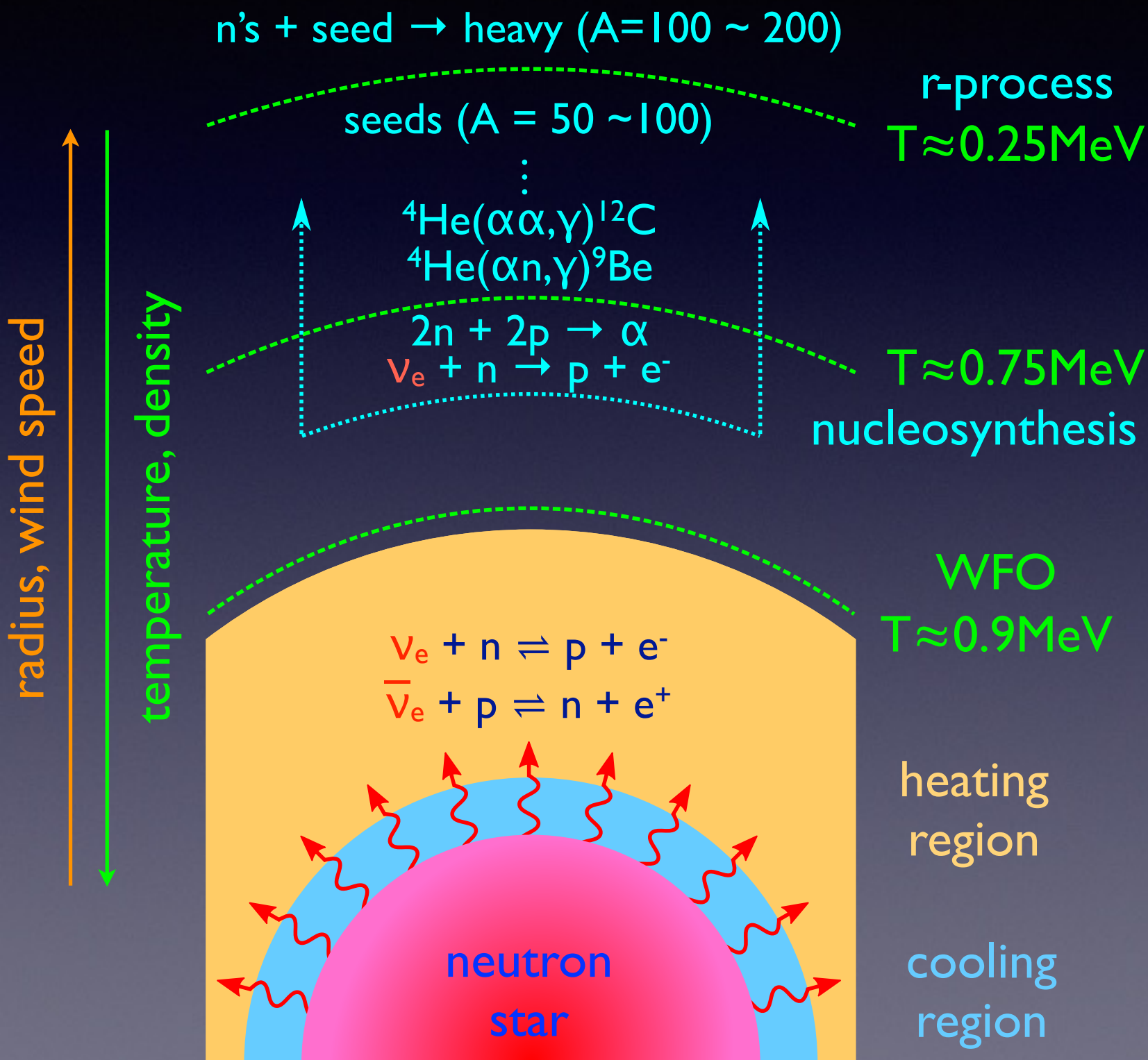


Spontaneous Symmetry Breaking in Collective Neutrino Oscillations

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Neutrinos in Supernovae



- $\sim 10^{53}$ ergs, 10^{58} neutrinos in ~ 10 seconds
- All neutrino species, 10~30 MeV
- Dominate energetics
- Influence nucleosynthesis
- Probe into SNe

Oscillations in Dense Medium

active neutrinos, coherent forward scattering only

$$\rho(t, \mathbf{x}, \mathbf{p}) = \int d^3\mathbf{x}' e^{-i\mathbf{p}\cdot\mathbf{x}'} \psi^\dagger \left(t, \mathbf{x} - \frac{1}{2}\mathbf{x}' \right) \psi \left(t, \mathbf{x} + \frac{1}{2}\mathbf{x}' \right)$$

$$(\partial_t + \hat{\mathbf{v}} \cdot \nabla) \rho = -i[H, \rho]$$

mass matrix

electron density

neutrino energy

$$H = \frac{M^2}{2E} + \sqrt{2}G_F \text{diag}[n_e, 0, 0] + H_{\nu\nu}$$

ν - ν forward scattering (self-coupling)

Neutrino Self-Coupling

mass matrix \longrightarrow
 $H = \frac{M^2}{2E} + \sqrt{2}G_F \text{diag}[n_e, 0, 0] + H_{\nu\nu}$
 neutrino energy \longleftarrow

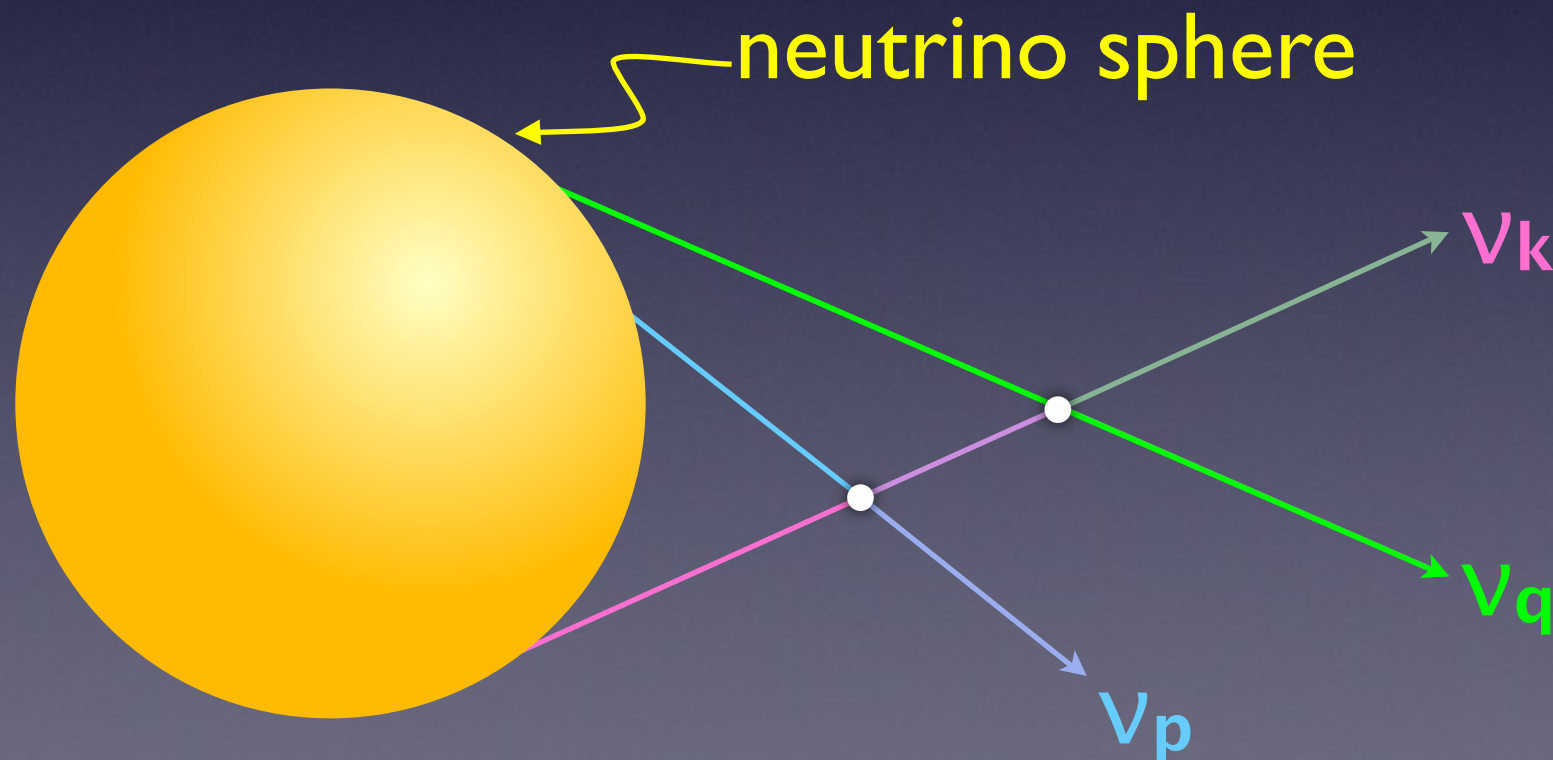
electron density \downarrow
 n_e

$H_{\nu\nu}$
 \uparrow
 ν - ν forward scattering
 (self-coupling)

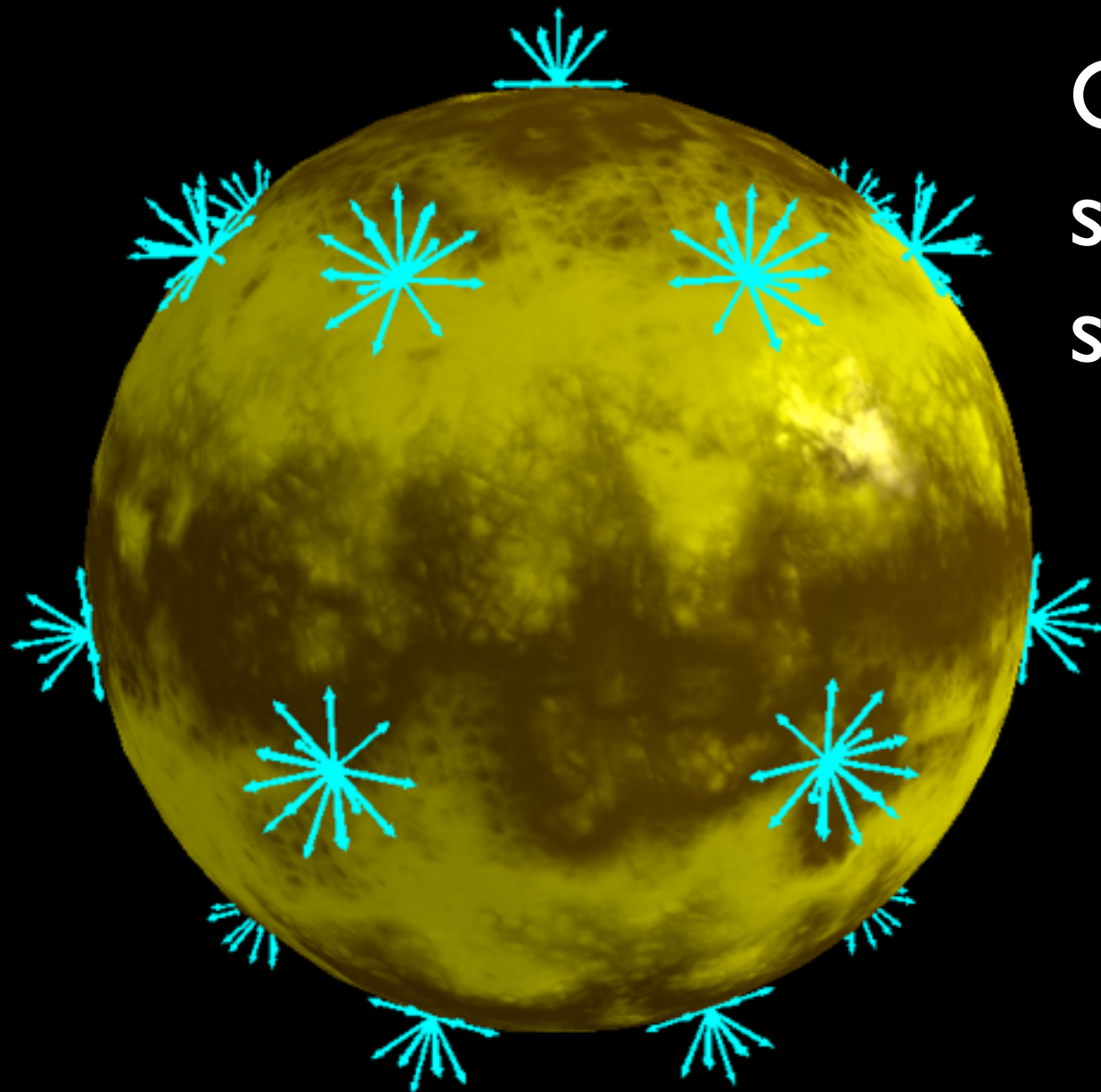
$$H_{\nu\nu} = \sqrt{2}G_F \int d^3\mathbf{p}' (1 - \hat{\mathbf{v}} \cdot \hat{\mathbf{v}}') (\rho_{\mathbf{p}'} - \bar{\rho}_{\mathbf{p}'})$$

Oscillations in SN

$$H = \frac{M^2}{2E} + \sqrt{2}G_F \text{diag}[n_e, 0, 0] + H_{\nu\nu}$$



(1+3+3)D

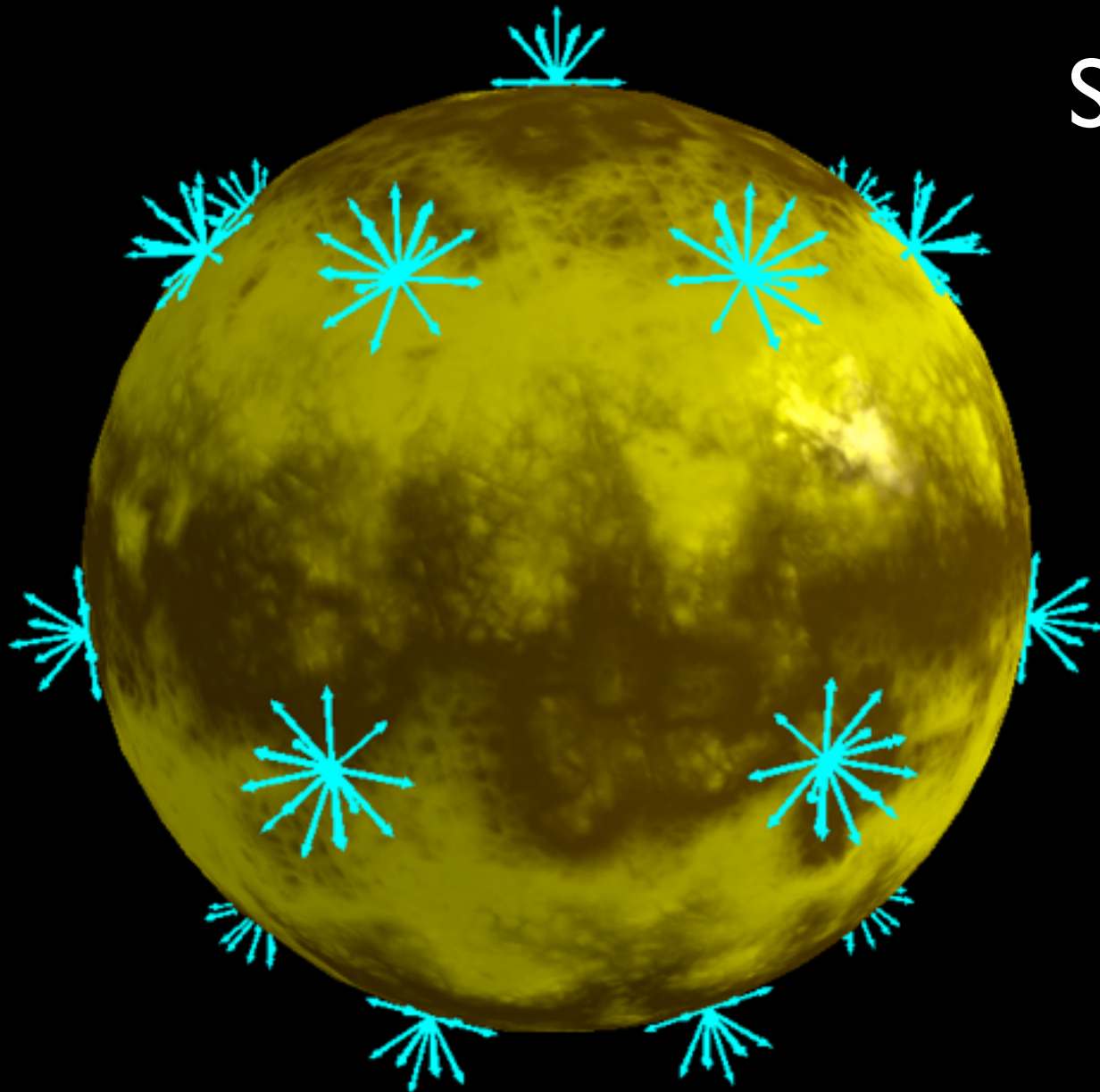


Coherent forward
scattering outside neutrino
sphere

$$\rho(t; r, \Theta, \Phi; E, \vartheta, \varphi)$$

(0+3+3)D

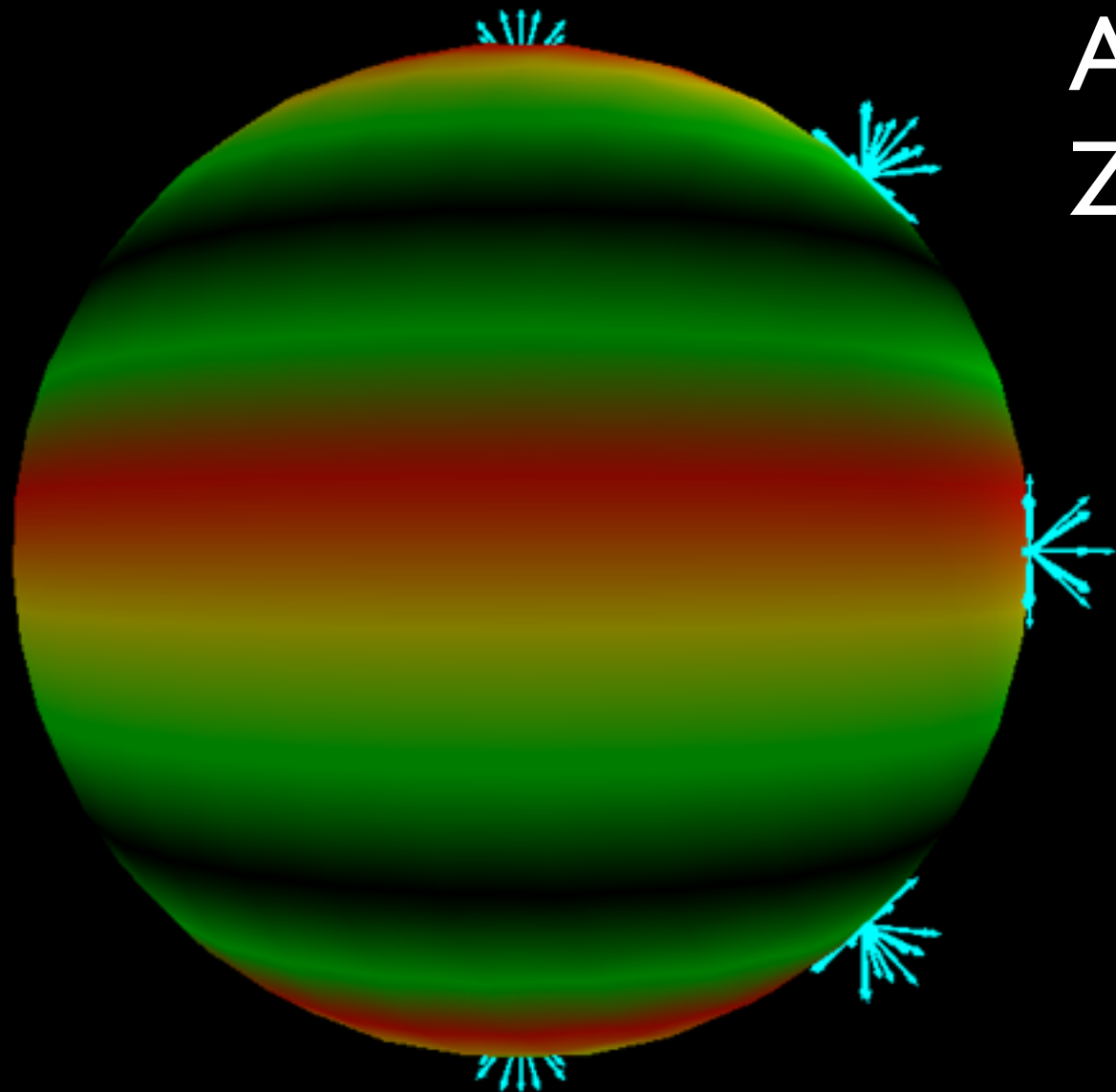
Stationary emission



$$\rho(r, \Theta, \Phi; E, \vartheta, \varphi)$$

(0+2+3)D

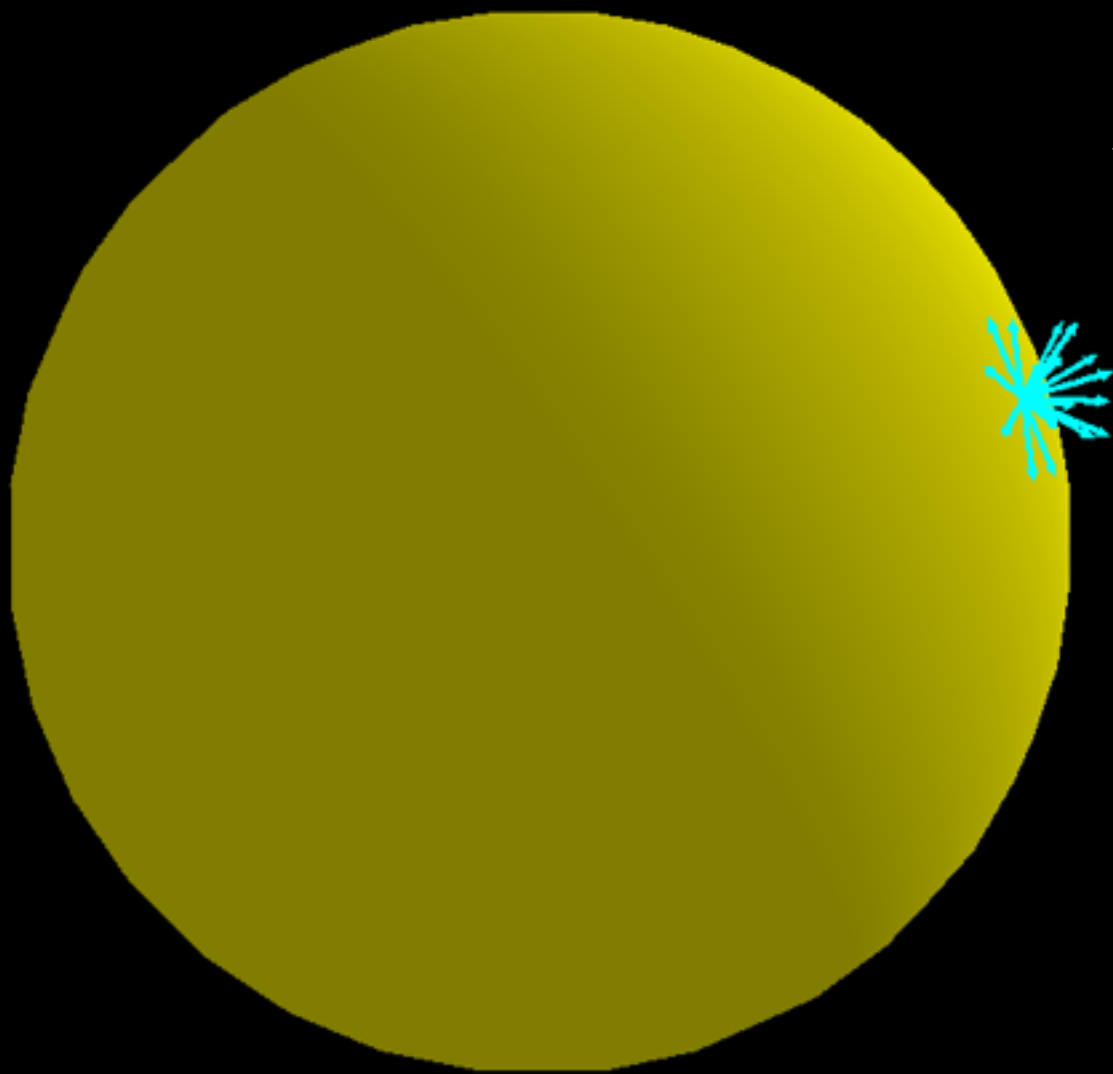
Axial symmetry around the
Z axis



$$\rho(r, \Theta; E, \vartheta, \varphi)$$

(0+1+3)D

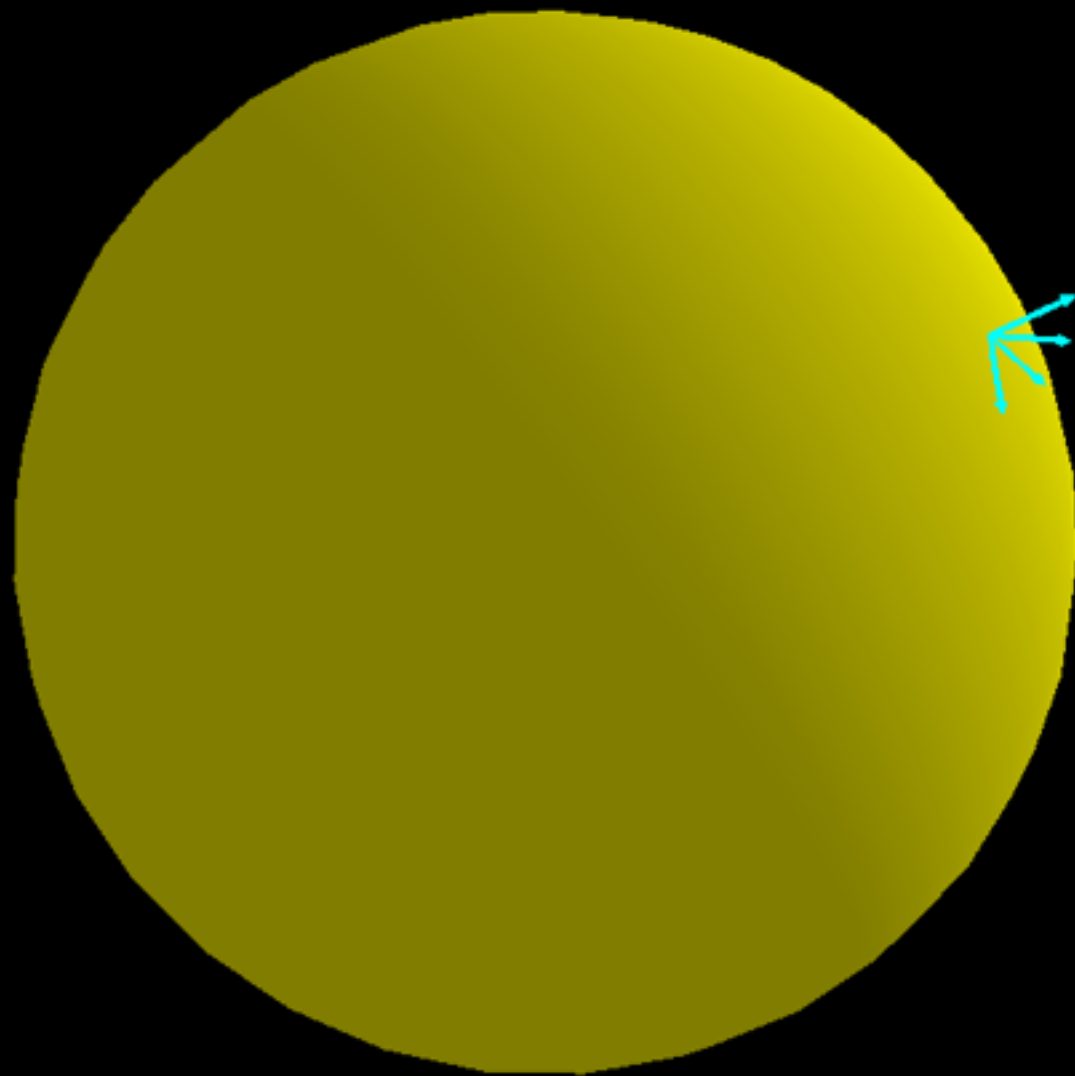
Spherical symmetry about
the center (**inconsistent?**)



$$\rho(r; E, \vartheta, \varphi)$$

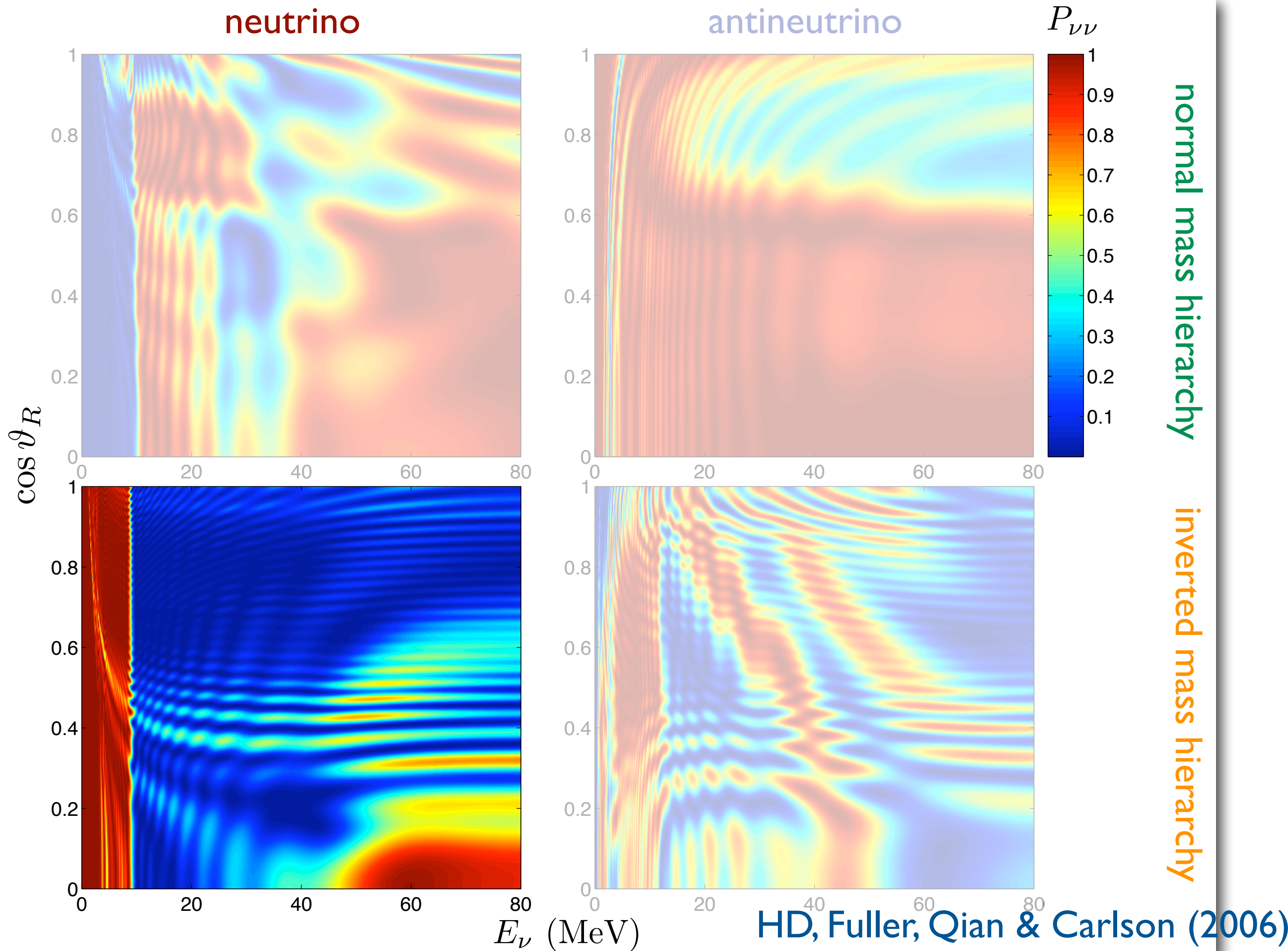
(0+1+2)D

Multi-Angle/Bulb Model

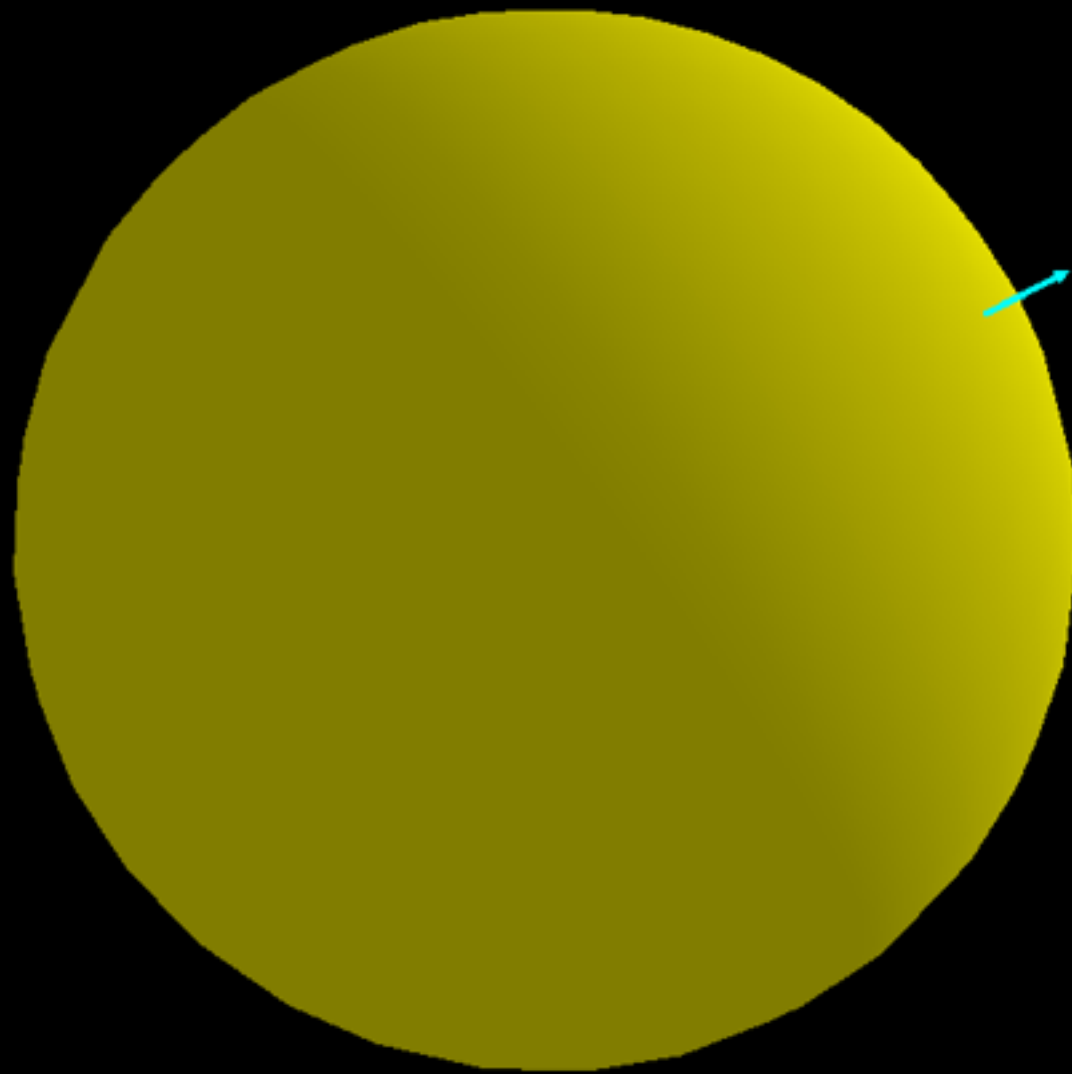


Azimuthal symmetry around
any radial direction

$$\rho(r; E, \vartheta)$$



(0+1+1)D Single-Angle Model

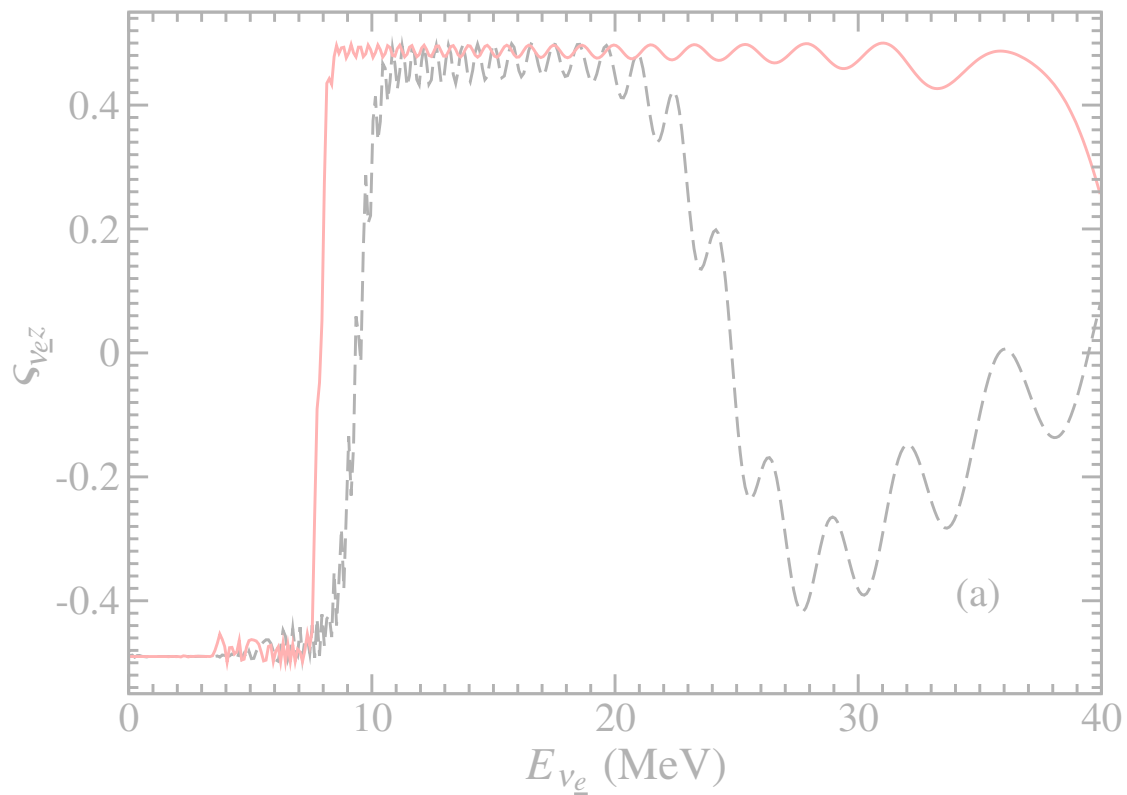


Trajectory independent
neutrino flavor evolution

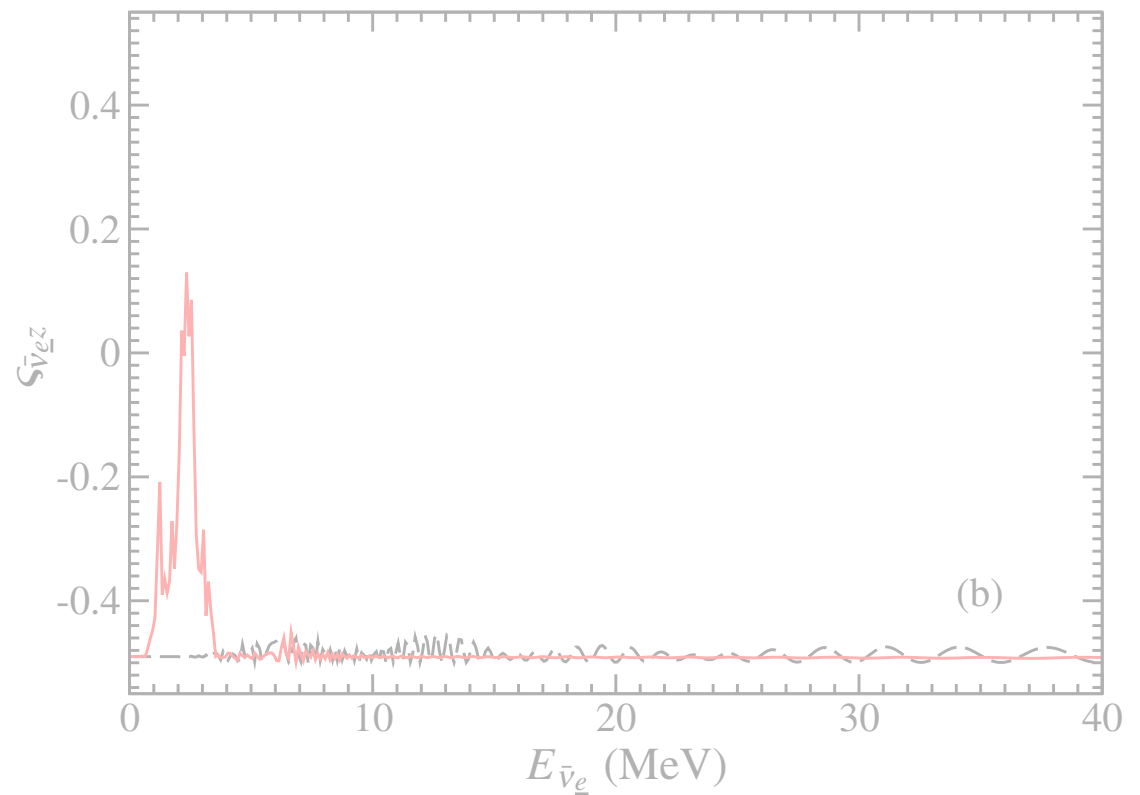
$$\rho(r; E)$$

Equivalent to an
homogeneous and isotropic
neutrino gas evolving with
time.

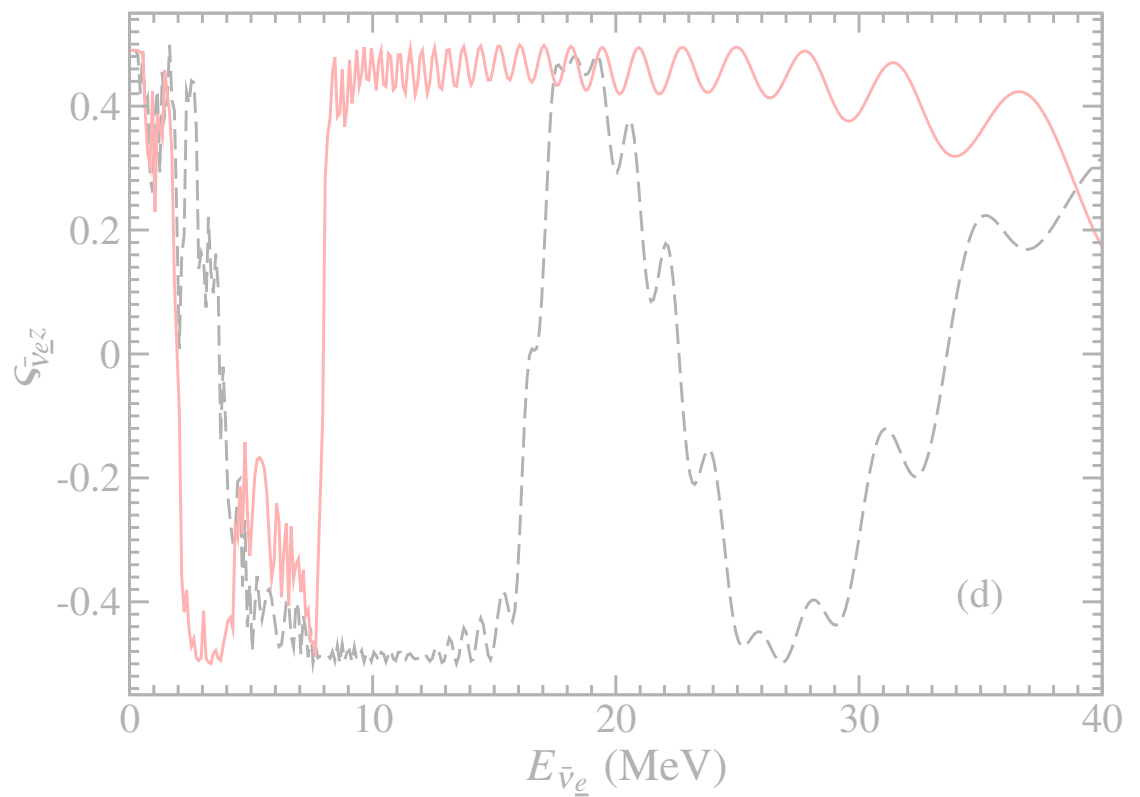
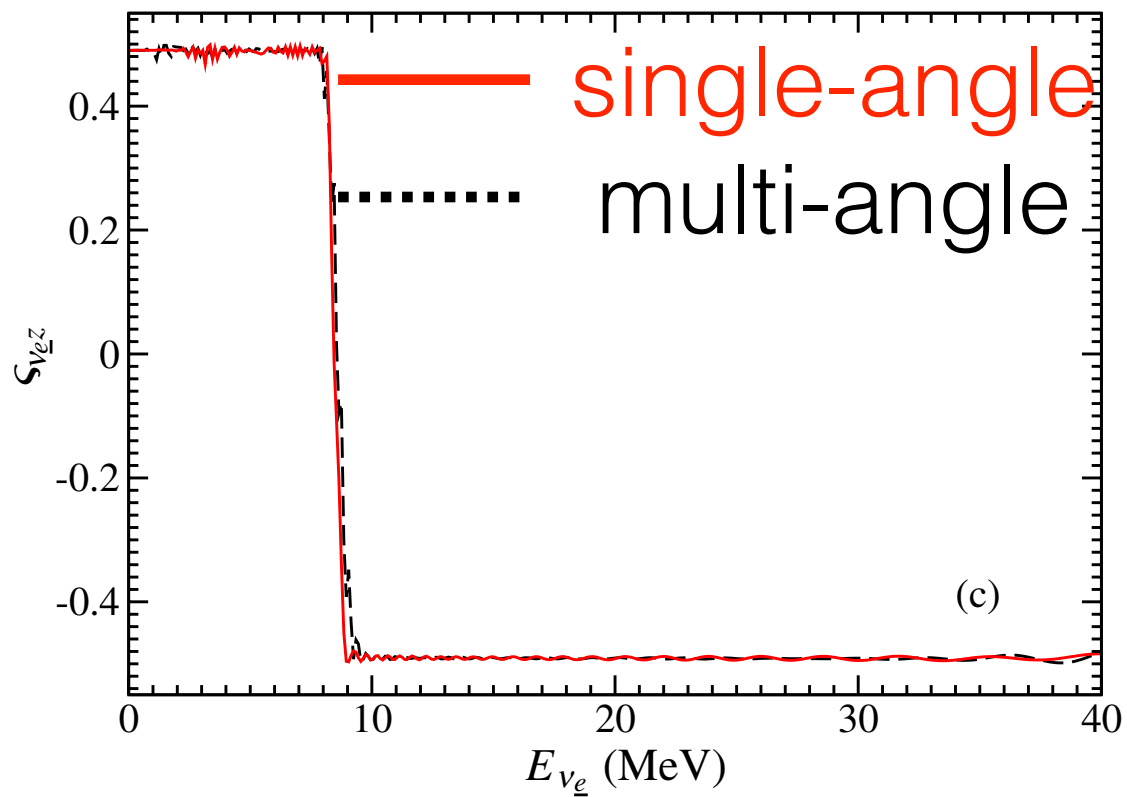
neutrino



antineutrino



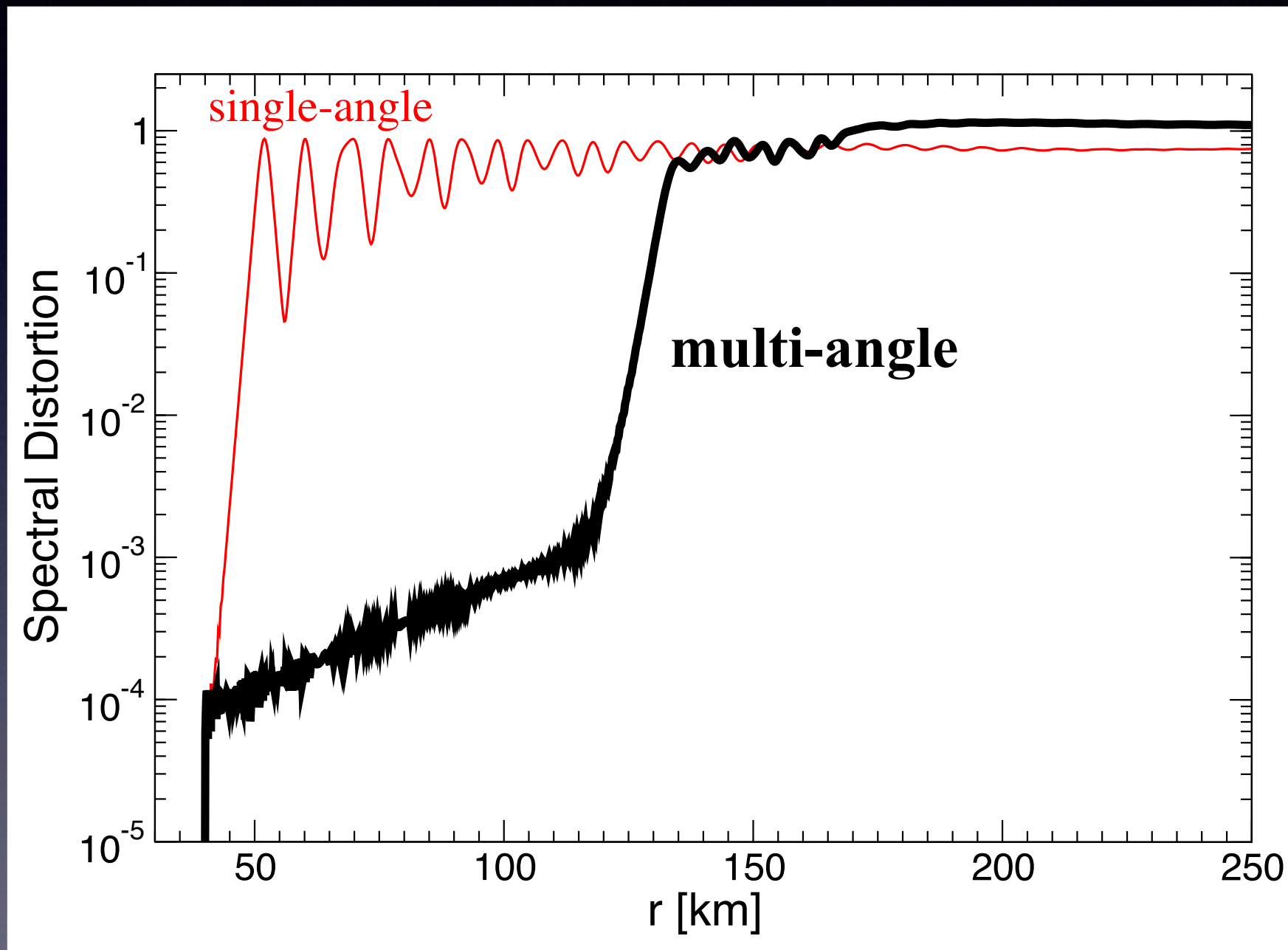
normal mass hierarchy



inverted mass hierarchy

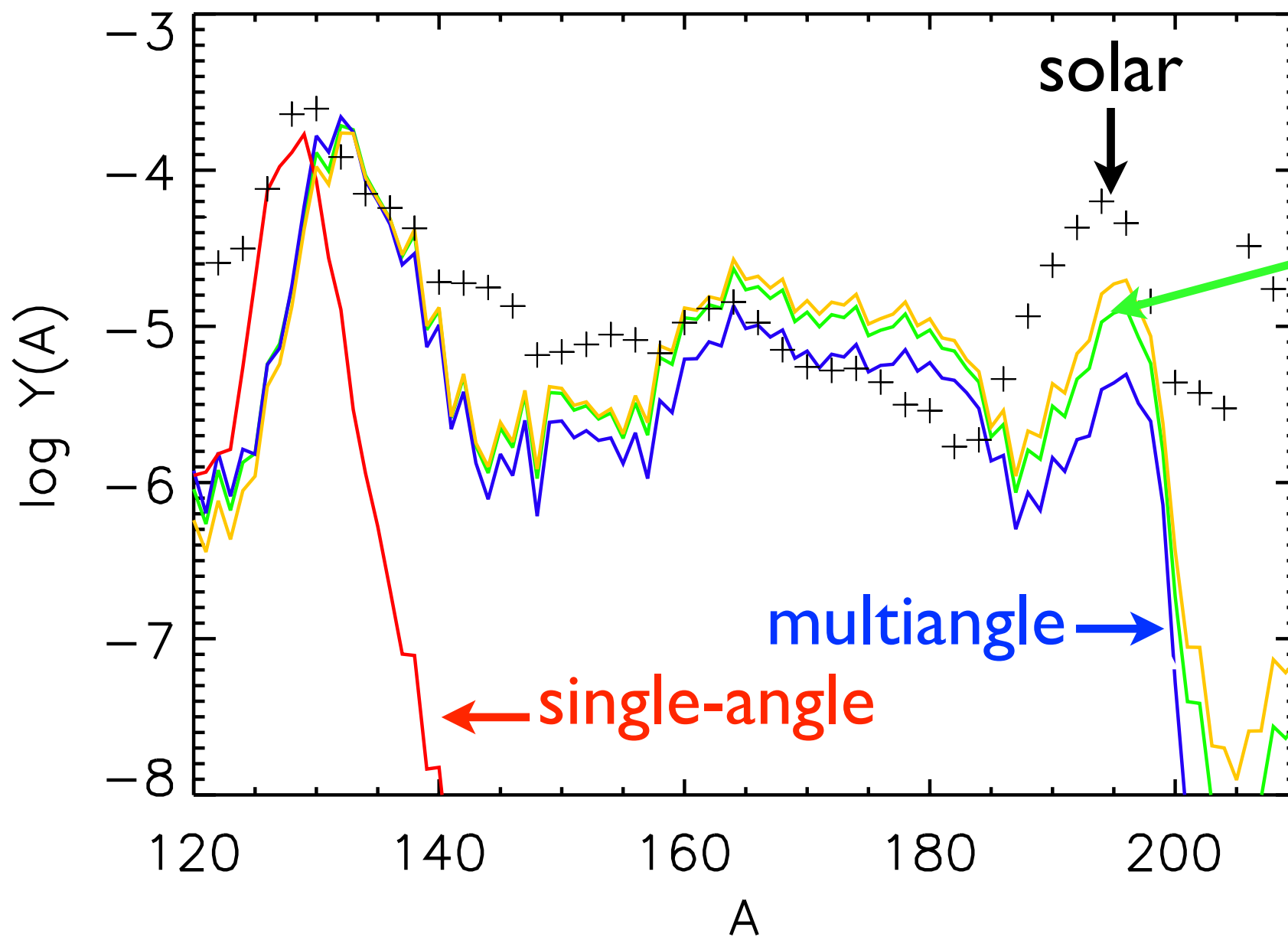
HD, Fuller, Qian & Carlson (2006)

Multiangangle Suppression



HD & Friedland (2010)

Nucleosynthesis



no osc.

Duan, Friedland,
McLaughlin & Surman
(2011)

Conclusions

- Dimensionality of the supernova model matters in collective neutrino oscillations.

Line Model

Neutrinos emitted in single energy and two beams.

Two active flavors, “no matter & small mixing angle”

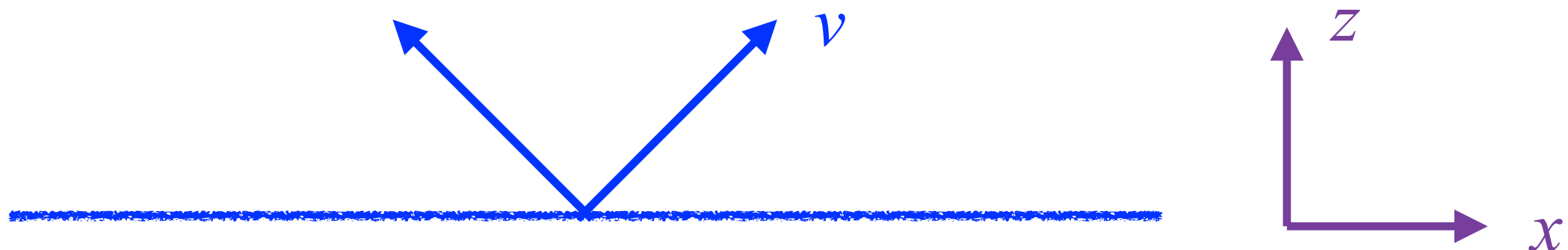
time independent, x translation symmetry, left-right symmetry

$$i\hat{\mathbf{v}} \cdot \nabla \rho = [-(\omega/2)\eta\sigma_3 + \mathbf{H}_{\nu\nu}, \rho]$$

$$i\hat{\mathbf{v}} \cdot \nabla \bar{\rho} = [-(-\omega/2)\eta\sigma_3 + \mathbf{H}_{\nu\nu}, \bar{\rho}]$$

$$\omega = \frac{|\Delta m^2|}{2E}$$

$$\eta = \begin{cases} +1 & \text{NH} \\ -1 & \text{IH} \end{cases}$$



Line Model

Neutrinos emitted in single energy and two beams.

Two active flavors, “no matter & small mixing angle”

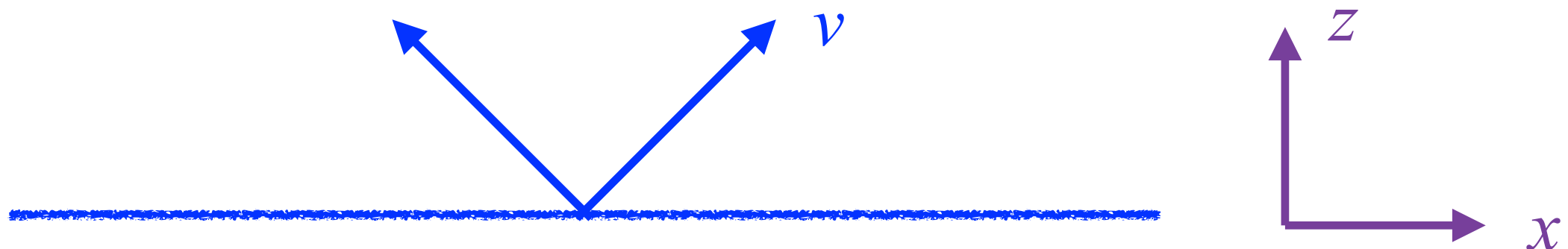
time independent, x translation symmetry, left-right symmetry

$$i\hat{\mathbf{v}} \cdot \nabla \rho = [-(\omega/2)\eta\sigma_3 + \cancel{H_{\nu\nu}}, \rho]$$

$$i\hat{\mathbf{v}} \cdot \nabla \bar{\rho} = [-(\omega/2)\eta\sigma_3 + \cancel{H_{\nu\nu}}, \bar{\rho}]$$

$$\rho(z) = e^{i(\omega/2)(z/v_z)\sigma_3} \rho(0) e^{-i(\omega/2)(z/v_z)\sigma_3}$$

$$\bar{\rho}(z) = e^{i(\omega/2)(z/v_z)\sigma_3} \bar{\rho}(0) e^{-i(\omega/2)(z/v_z)\sigma_3}$$



Line Model

Neutrinos emitted in single energy and two beams.

Two active flavors, “no matter & small mixing angle”

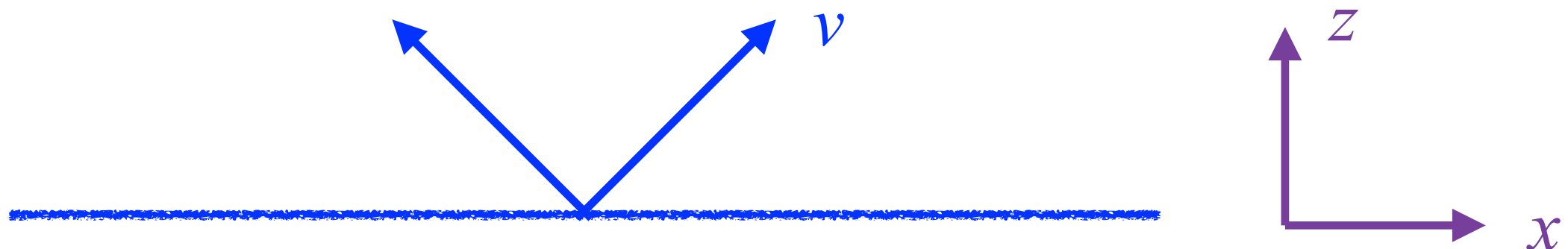
time independent, x translation symmetry, left-right symmetry

$$i\hat{\mathbf{v}} \cdot \nabla \rho = [-(\omega/2)\eta\sigma_3 + \mathbf{H}_{\nu\nu}, \rho]$$

$$i\hat{\mathbf{v}} \cdot \nabla \bar{\rho} = [-(\omega/2)\eta\sigma_3 + \mathbf{H}_{\nu\nu}, \bar{\rho}]$$

$$\rho(z) = e^{i(\Omega/2)(z/v_z)\sigma_3} \rho(0) e^{-i(\Omega/2)(z/v_z)\sigma_3}$$

$$\bar{\rho}(z) = e^{i(\Omega/2)(z/v_z)\sigma_3} \bar{\rho}(0) e^{-i(\Omega/2)(z/v_z)\sigma_3}$$



Line Model

Electron flavor neutrinos and antineutrinos

$$\rho \propto \begin{bmatrix} 1 & \epsilon \\ \epsilon^* & 0 \end{bmatrix}$$

$$\bar{\rho} \propto \begin{bmatrix} 1 & \bar{\epsilon} \\ \bar{\epsilon}^* & 0 \end{bmatrix}$$

$$i\partial_z \begin{bmatrix} \epsilon \\ \bar{\epsilon} \end{bmatrix} = v_z^{-1} \begin{bmatrix} -\eta\omega - \alpha\mu & \alpha\mu \\ -\mu & \eta\omega + \mu \end{bmatrix} \begin{bmatrix} \epsilon \\ \bar{\epsilon} \end{bmatrix}$$

$$\alpha = n_{\bar{\nu}}/n_{\nu}$$

$$\mu \propto G_F n_{\nu}$$

Ω can be complex in IH and $\frac{2\omega}{(1 + \sqrt{\alpha})^2} < \mu < \frac{2\omega}{(1 - \sqrt{\alpha})^2}$

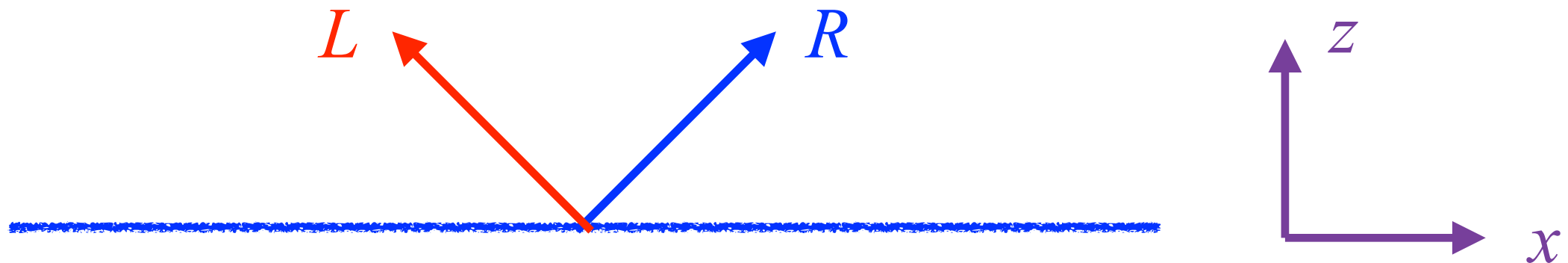
Linear (flavor) stability analysis (Banerjee et al, 2011)

Line Model

time independent, x translation symmetry, ~~left-right symmetry~~

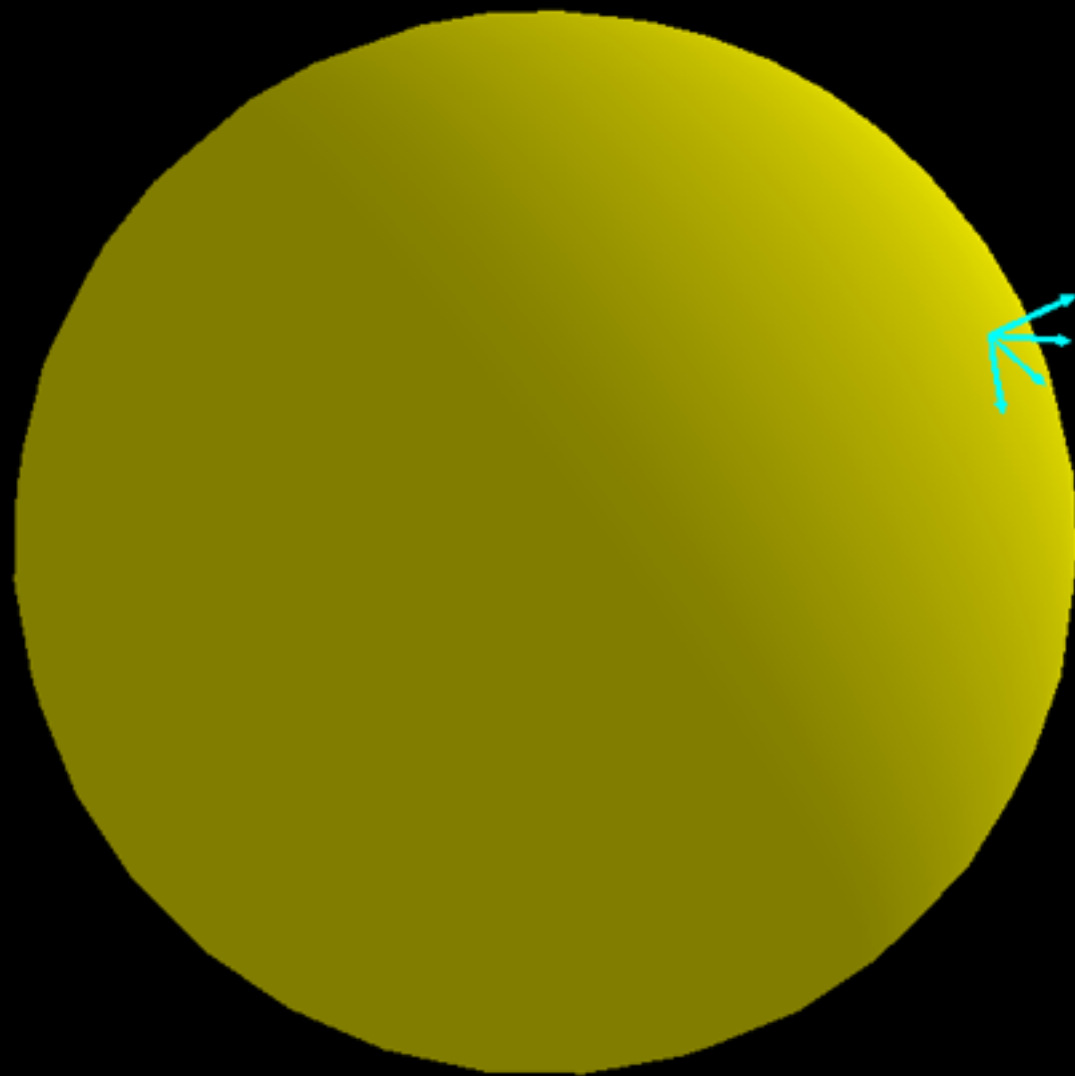
$$i\partial_z \begin{bmatrix} \epsilon_+ \\ \bar{\epsilon}_+ \\ \epsilon_- \\ \bar{\epsilon}_- \end{bmatrix} = \begin{bmatrix} \Lambda_+ & \\ & \Lambda_- \end{bmatrix} \begin{bmatrix} \epsilon_+ \\ \bar{\epsilon}_+ \\ \epsilon_- \\ \bar{\epsilon}_- \end{bmatrix} \quad \epsilon_{\pm} = \frac{\epsilon_L \pm \epsilon_R}{2}$$

$\Omega_{+/-}$ can be complex in IH/NH



(0+1+2)D

Multi-Angle/Bulb Model

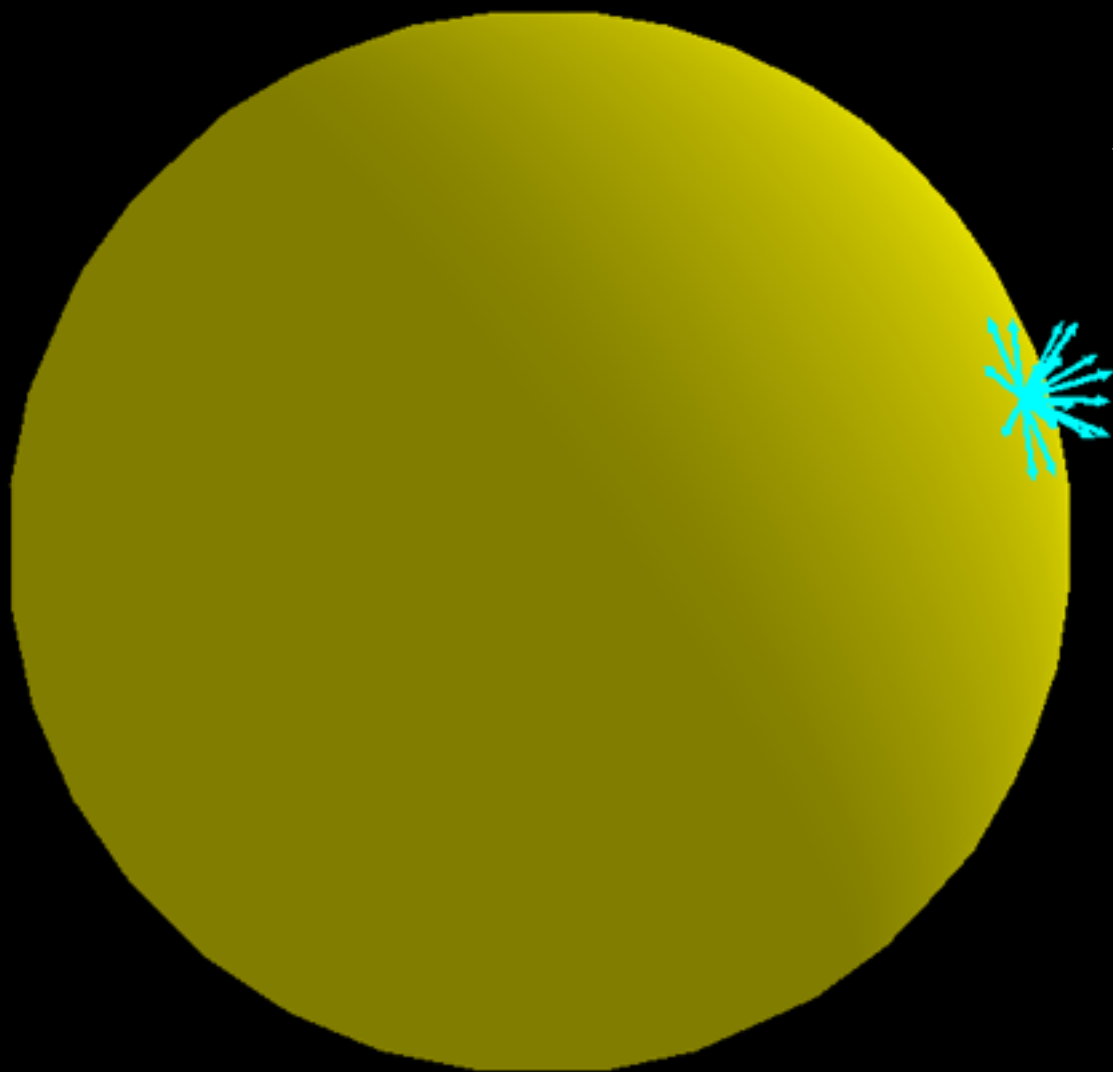


Azimuthal symmetry around
any radial direction

$$\rho(r; E, \vartheta)$$

(0+1+3)D

Spherical symmetry about the center



$$\rho(r; E, \vartheta, \varphi)$$

z

L, R

Raffelt et al (2013)
Mirizzi (2013)

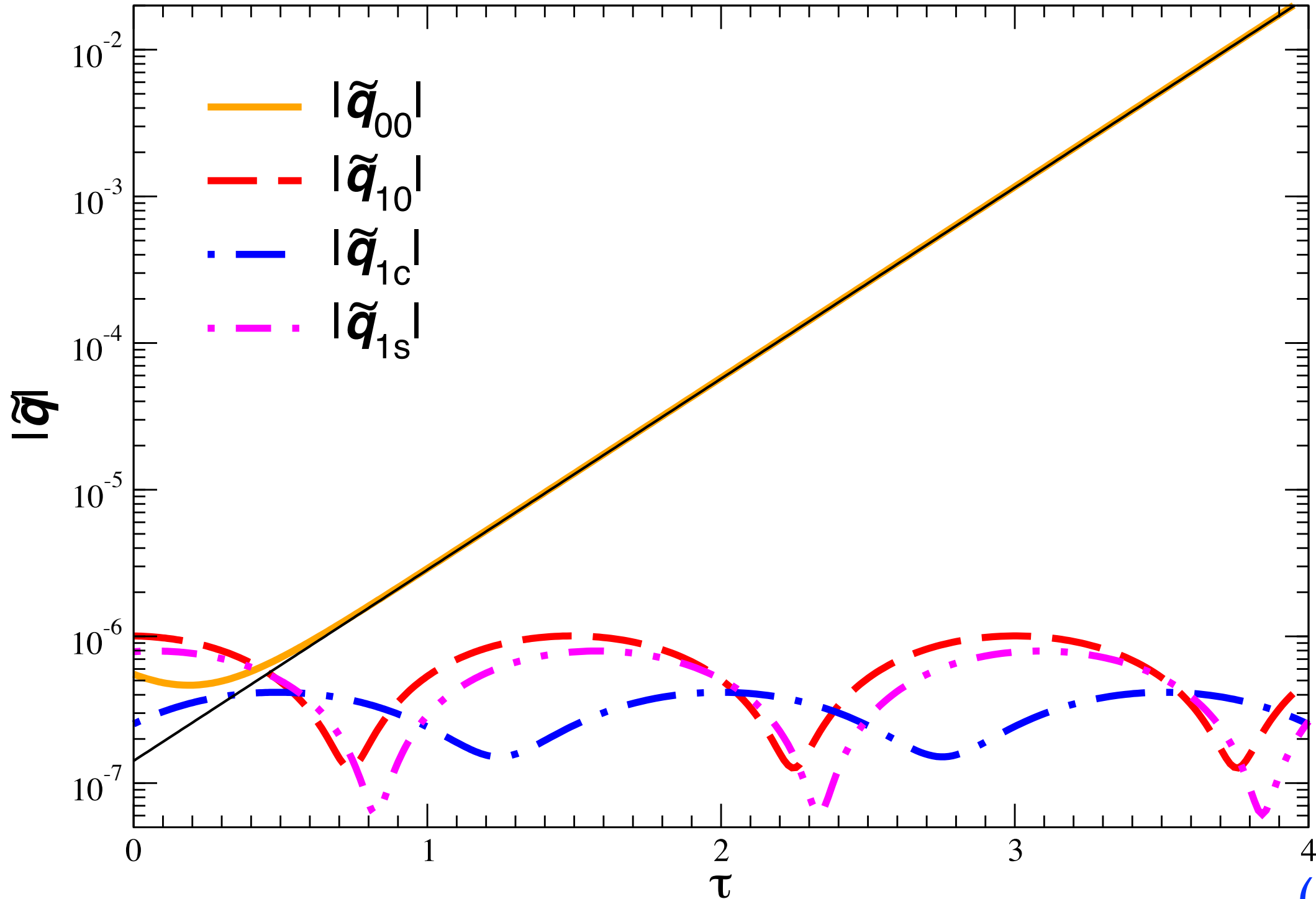
Homogeneous Gas

$$(1 - \hat{\mathbf{v}} \cdot \hat{\mathbf{v}}')\mu = 4\pi\mu \left[Y_{0,0}(\mathbf{p})Y_{0,0}^*(\mathbf{p}') - \frac{1}{3} \sum_{m=0,\pm 1} Y_{1,m}(\mathbf{p})Y_{1,m}^*(\mathbf{p}') \right]$$

- Multipole modes are decoupled in the linear regime
- $l=0$: $\mu_{\text{eff}} = \mu$, unstable in IH
- $l=1$: $\mu_{\text{eff}} = -\mu/3$ unstable in NH (breaking isotropy)

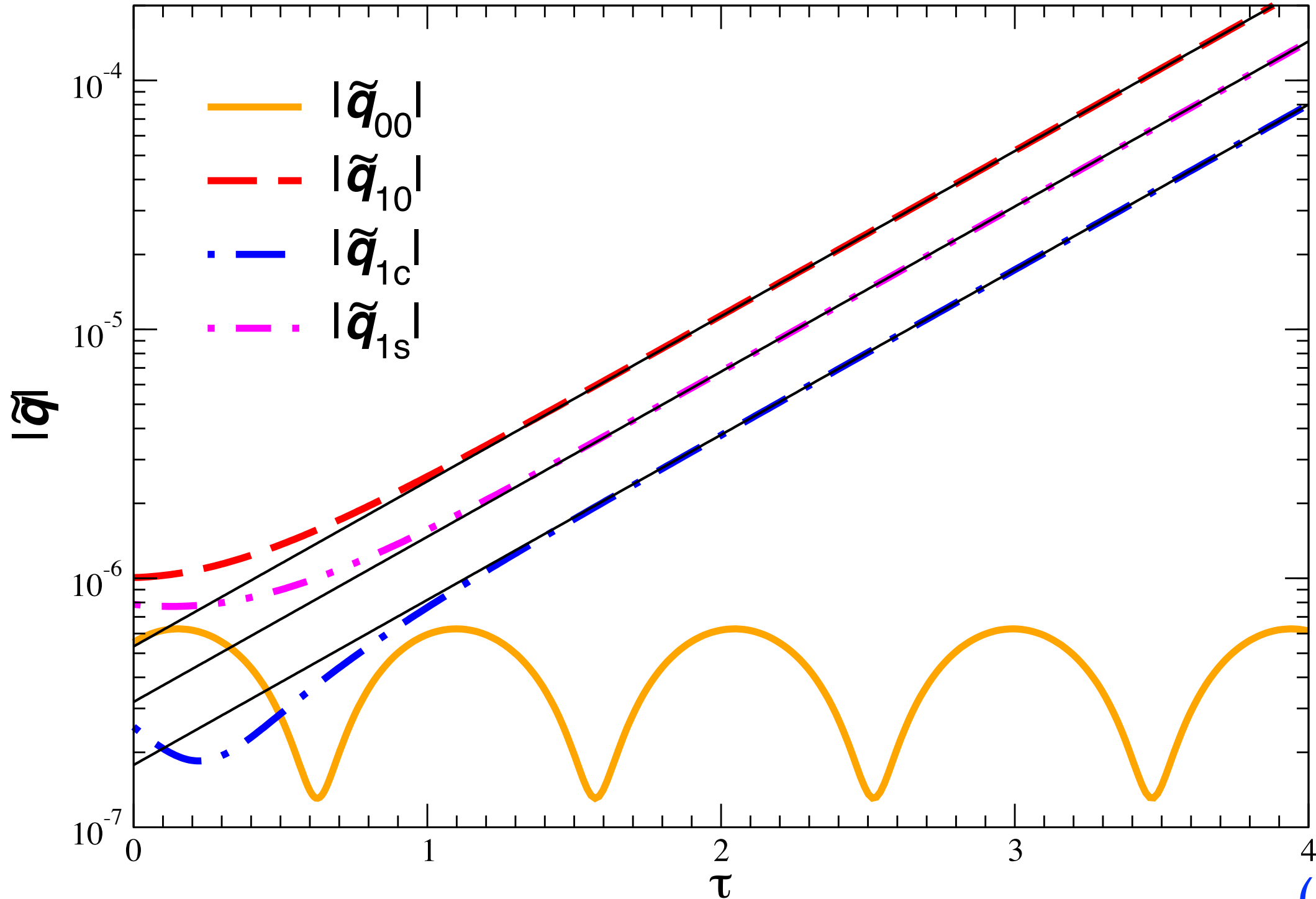
(HD, 2013)

Inverted Hierarchy



(HD, 2013)

Normal Hierarchy



(HD, 2013)

Conclusions

- Dimensionality of the supernova model matters in collective neutrino oscillations.
- Spontaneous breaking of **directional symmetry in momentum space** \Rightarrow Collective neutrino oscillations in both mass hierarchies

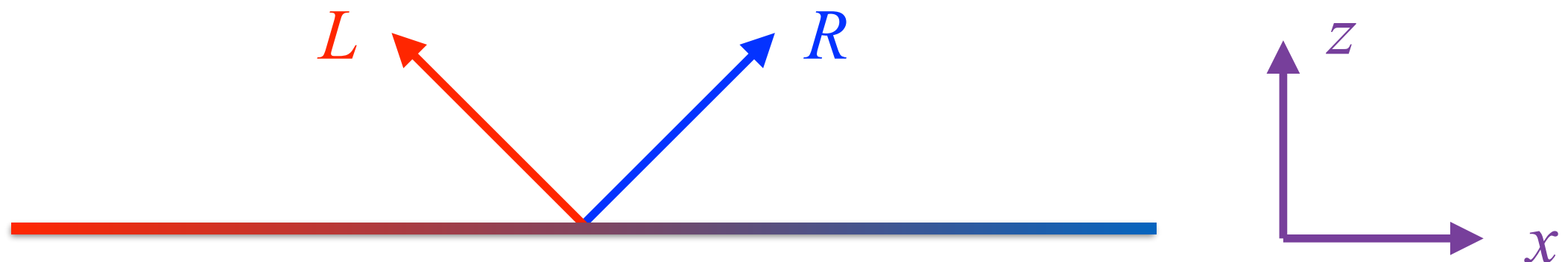
Line Model

time independent, ~~x translation symmetry, left-right symmetry~~

Periodic condition

$$\epsilon_m^\pm = \frac{1}{L} \int_0^L \left[\frac{\epsilon_L(x) \pm \epsilon_R(x)}{2} \right] e^{-i(2m\pi)(x/L)} dx$$

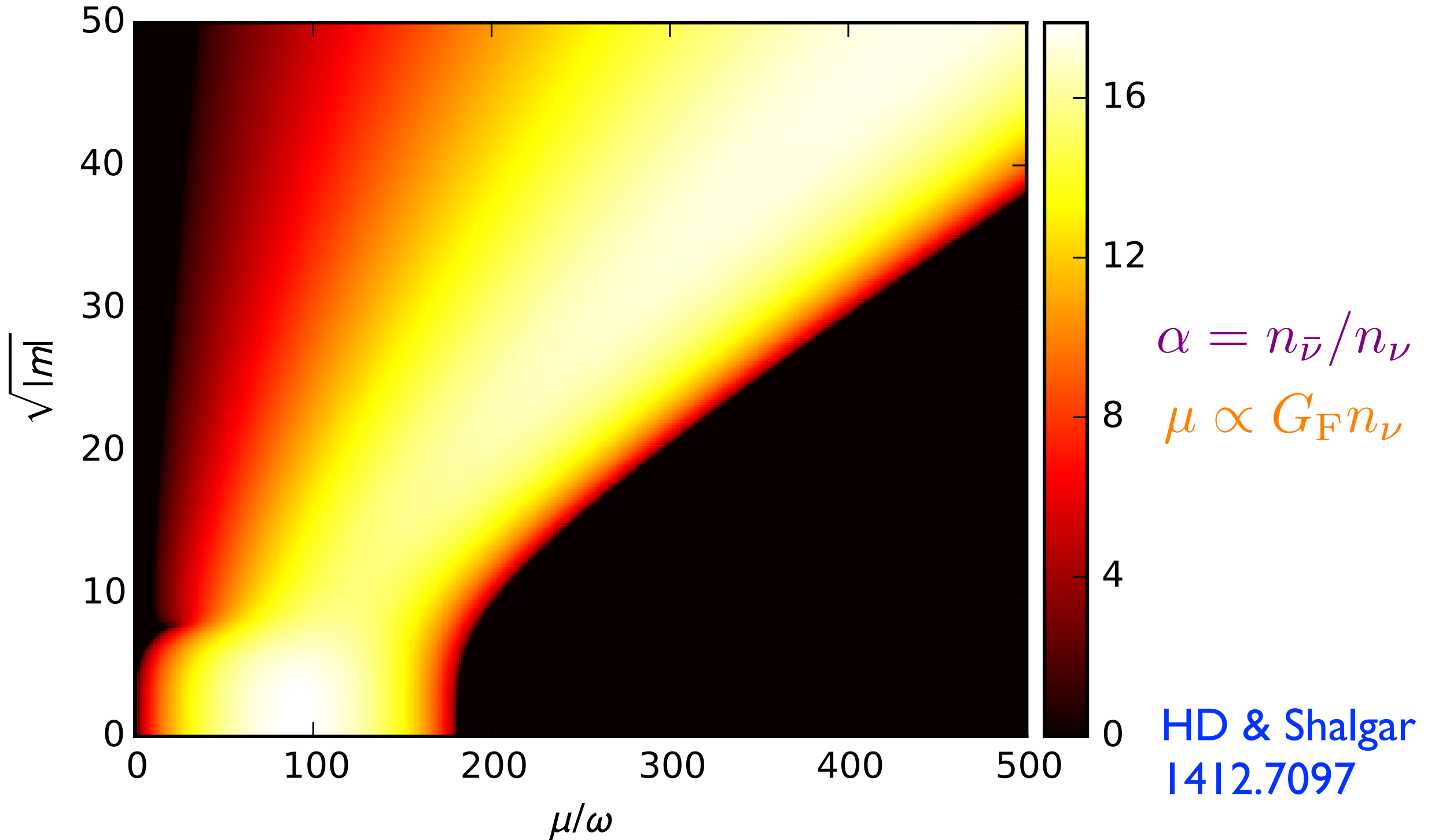
$$i\partial_z \begin{bmatrix} \epsilon_m^+ \\ \bar{\epsilon}_m^+ \\ \epsilon_m^- \\ \bar{\epsilon}_m^- \end{bmatrix} = \Lambda_m \cdot \begin{bmatrix} \epsilon_m^+ \\ \bar{\epsilon}_m^+ \\ \epsilon_m^- \\ \bar{\epsilon}_m^- \end{bmatrix}$$



A Toy Model

$\alpha = 0.8$

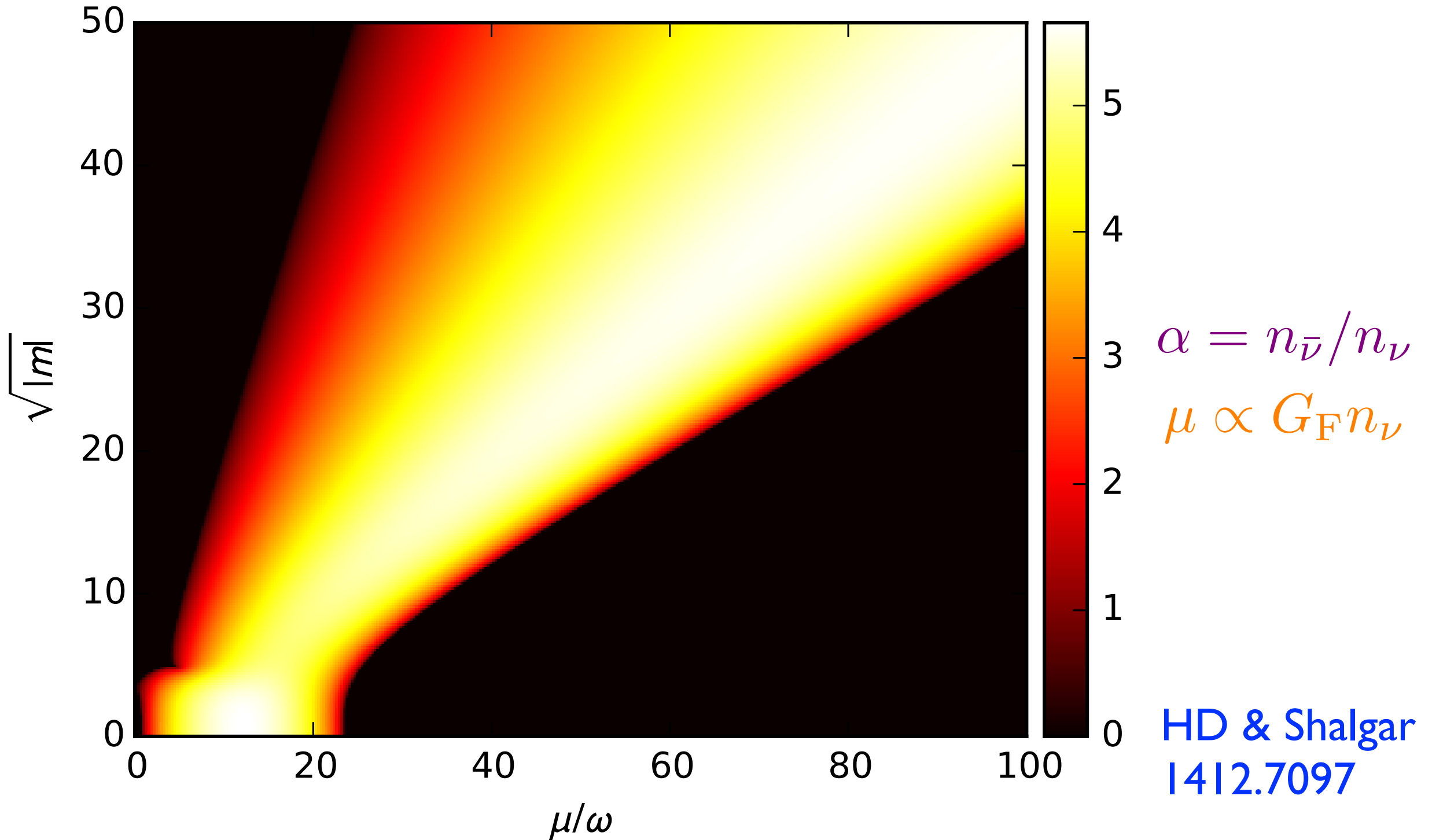
κ^{\max}/ω



A Toy Model

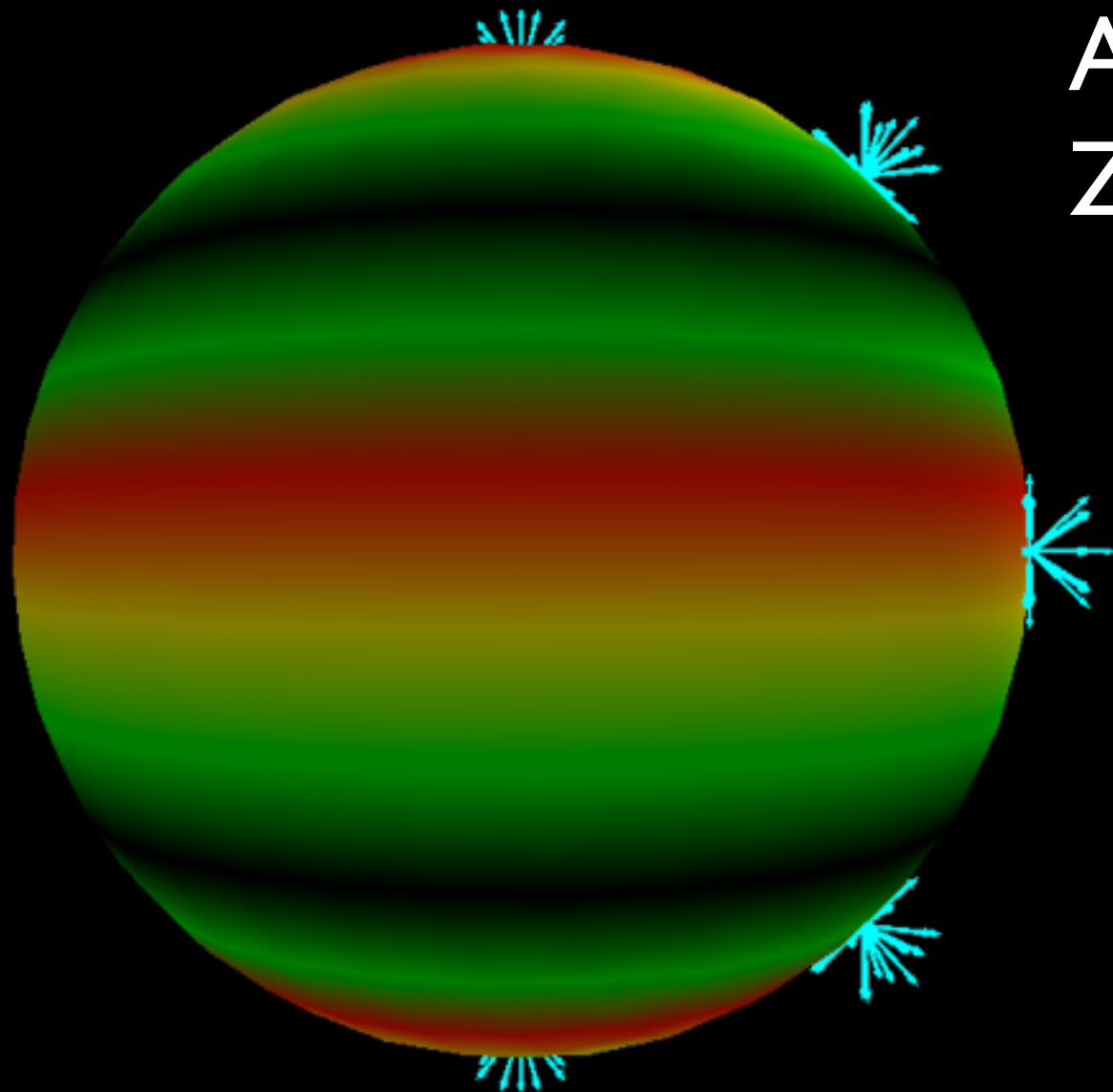
$\alpha = 0.5$

κ^{\max}/ω



(0+2+2)D

Axial symmetry around the
Z axis



$$\rho(r, \Theta; E, \cancel{\theta}, \varphi)$$

z x L, R

Conclusions

- Dimensionality of the supernova model matters in collective neutrino oscillations.
- Spontaneous breaking of **directional symmetry in momentum space** \Rightarrow Collective neutrino oscillations in both mass hierarchies
- Spontaneous breaking of **spatial symmetry in position space** \Rightarrow Collective neutrino oscillations at larger density
- Stay tuned