

Strongly correlated states of trapped ultracold fermions in a $U(2)$ gauge potential



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Frontiers in quantum simulations with cold atoms
INT-Seattle, 9 April 2015

M.Burrello, **MR**, M.Roncaglia, A.Trombettoni, PRB 91, 115117 (2015)

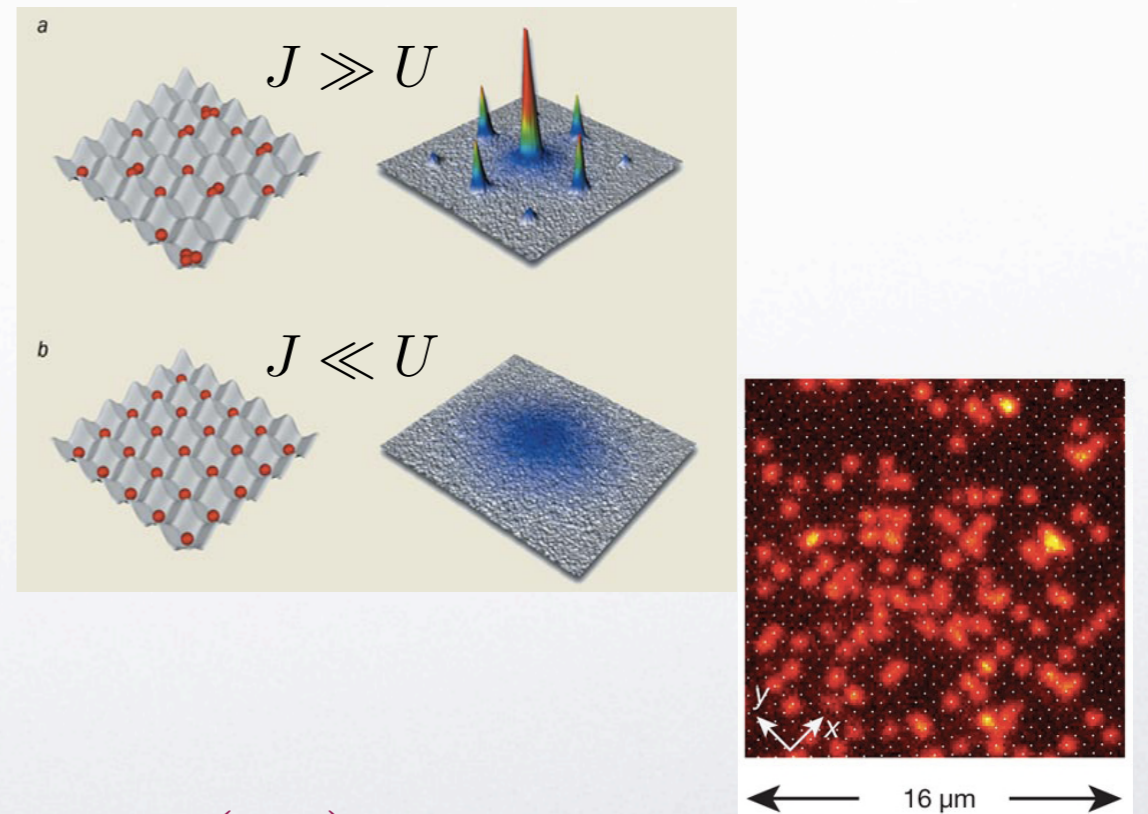
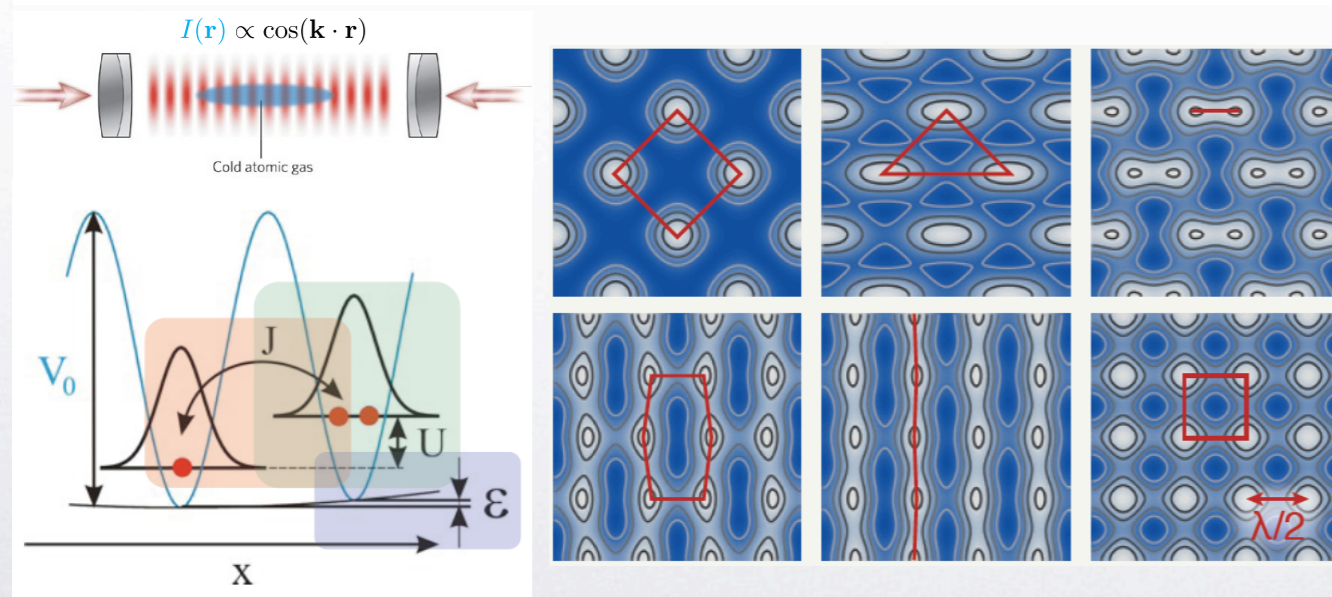
Outline

- Motivation: beyond standard
- $U(2)$ potential & deformed LL
- non-monotonic Haldane pseudopotentials
- Novel incompressible states: Haffnian?
- Entanglement spectrum
- Conclusions

Quantum Engineering

- analog implementation of solid-state systems, with added values:
 - isolated neutral quantum systems (long coherence times)
 - high tunability of microscopic parameters (also interactions!)
 - access to many microscopic observables
- **GOAL:** answering questions untreatable by classical calculations (!?)

$$H = -J \sum_{\langle i,j \rangle} b_i^\dagger b_j + \sum_i \epsilon_i \hat{n}_i + \frac{1}{2} U \sum_i \hat{n}_i (\hat{n}_i - 1)$$



M. Lewenstein, et al., Adv Phys 56, 243–379 (2007).
 I. Bloch, J. Dalibard, W. Zwerger, RMP 80, 885 (2008)

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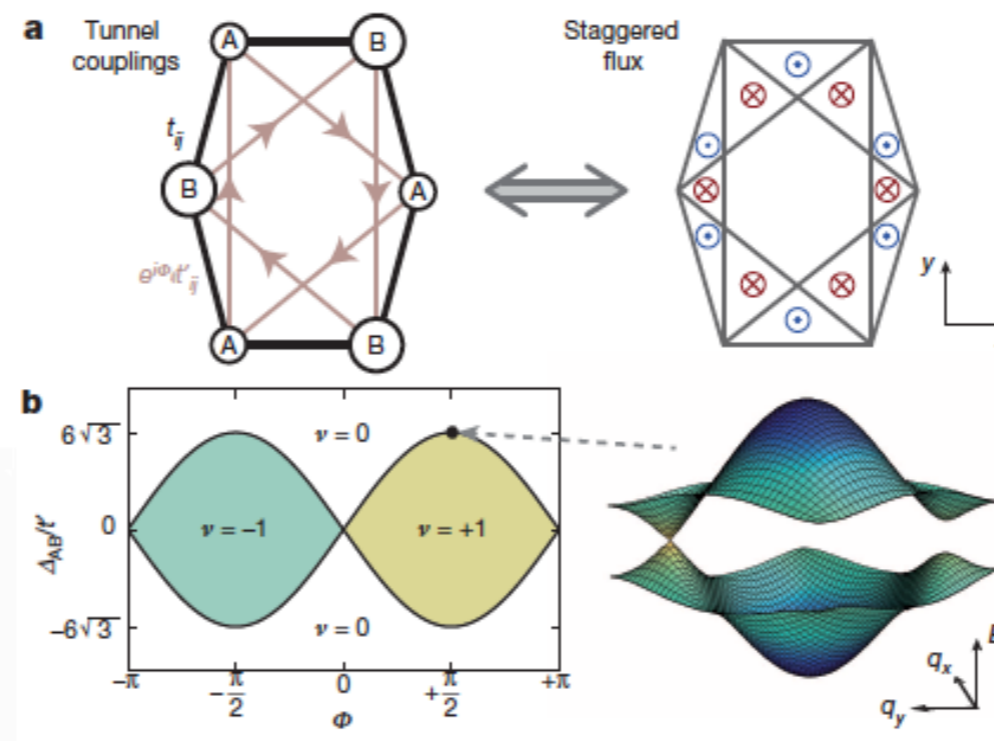
Quantum Engineering

- BUT can we go beyond emulation / simulation of existing regimes?

e.g.

Haldane model @ ETH

G. Jotzu, et al., Nature **515**, 237 (2014)



| Regime | Broken symmetry | Band structure | Berry curvature |
|---|-----------------|----------------|-----------------|
| ① $\varphi = 0^\circ; \pm 180^\circ$ $\Delta_{AB} = 0$ | — | | |
| ② $\varphi = 0^\circ; \pm 180^\circ$ $\Delta_{AB} > 0$ | IS | | |
| ③ $\varphi = 0^\circ; \pm 180^\circ$ $\Delta_{AB} < 0$ | IS | | |
| ④ $-180^\circ < \varphi < 0^\circ$ $\Delta_{AB} = 0$ | TRS | | |
| ⑤ $0^\circ < \varphi < 180^\circ$ $\Delta_{AB} = 0$ | TRS | | |

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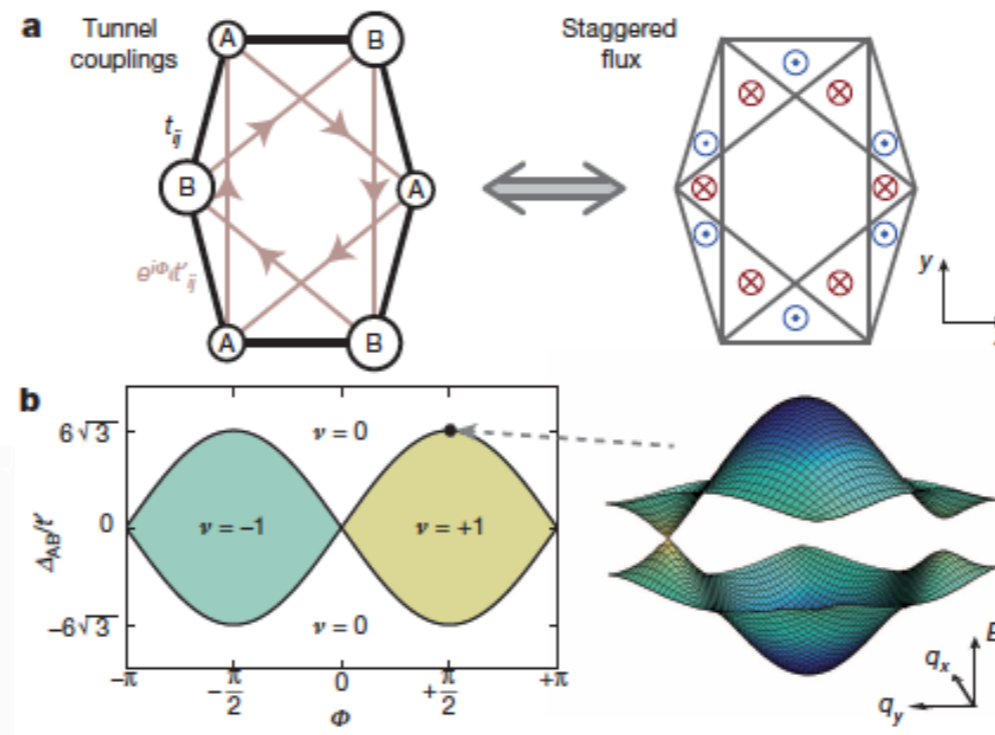
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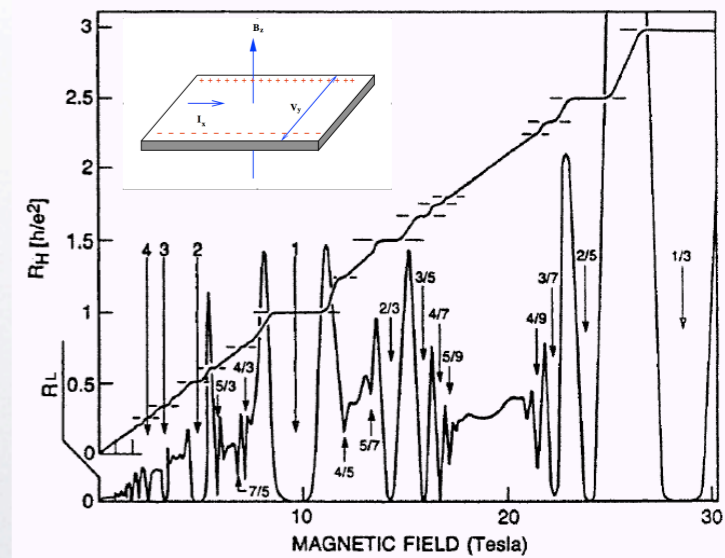
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- HERE, focus on Fractional Quantum Hall states:

- a wealth of exotic states & topological properties predicted “mathematically” for “strange” interaction & gauge potential forms
- but in semiconductors 2DEG, almost no tunability !



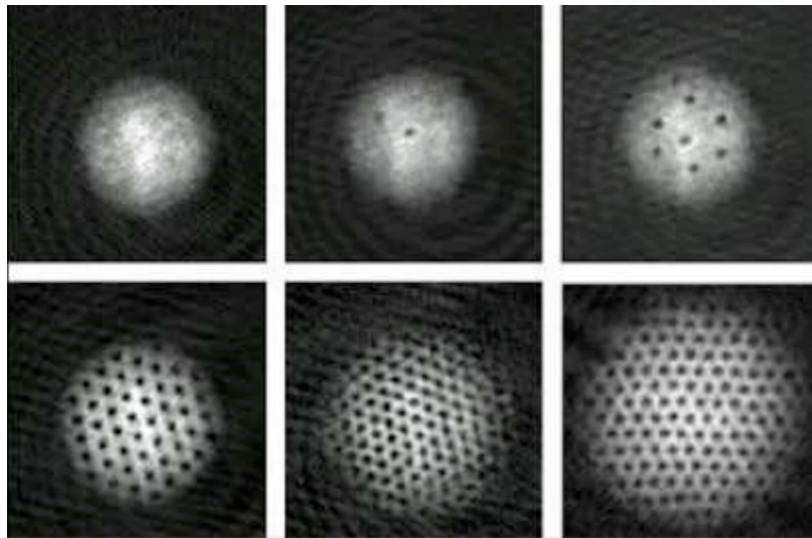
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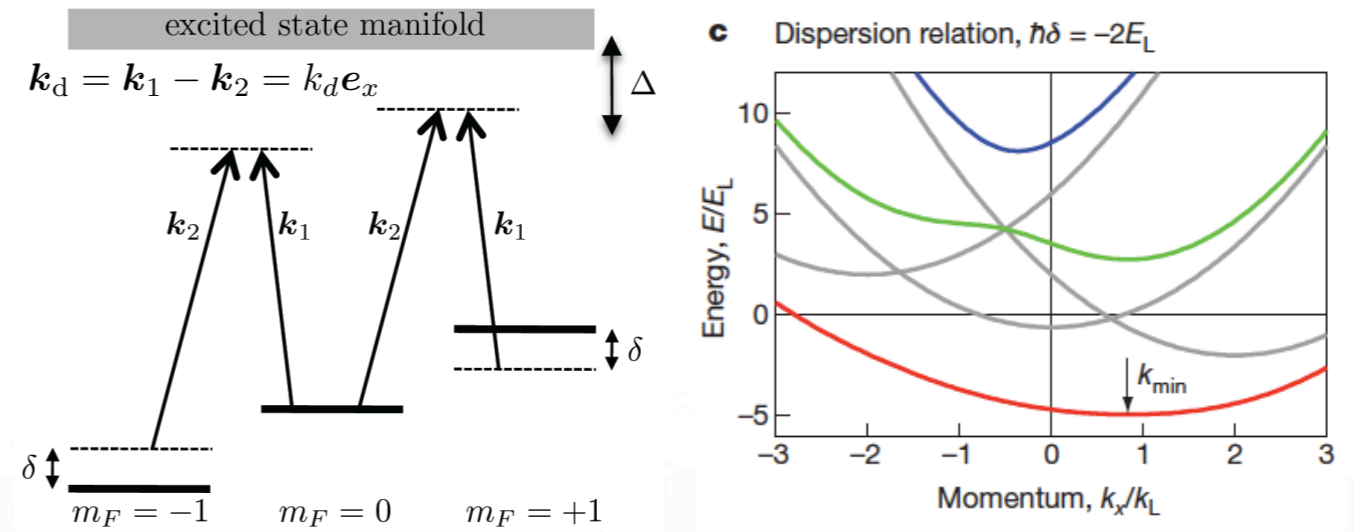
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Gauge potentials in cold atoms

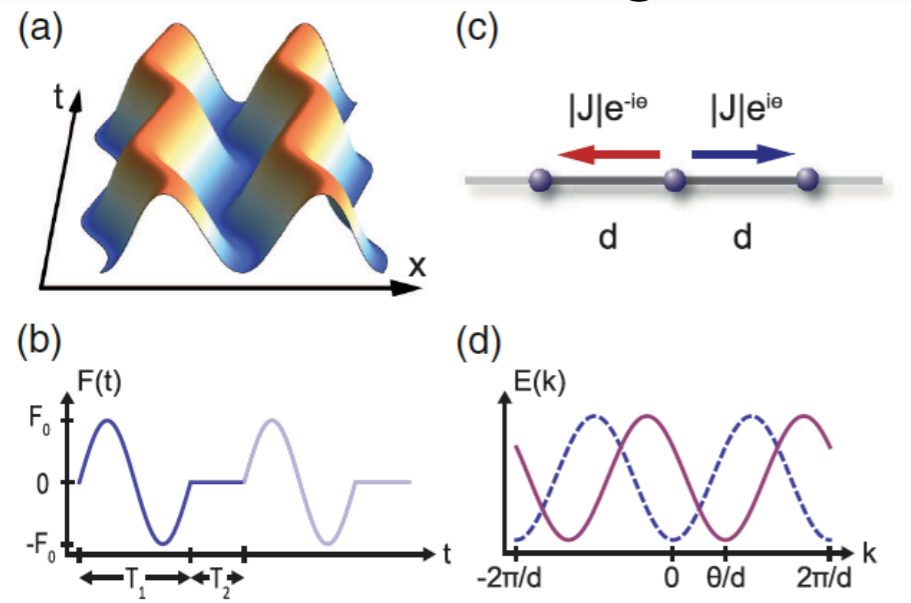
- fast rotation



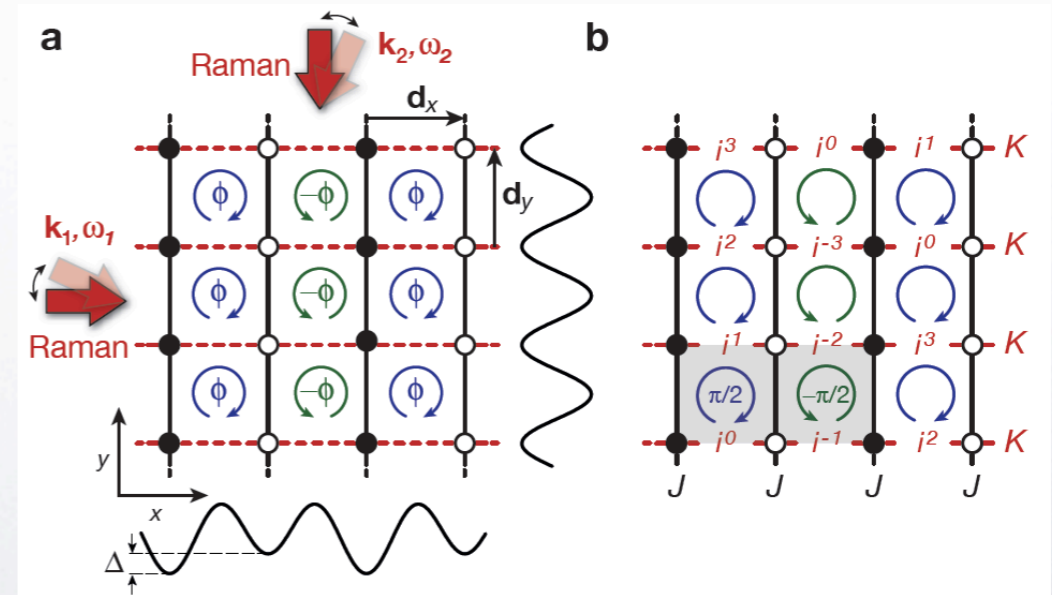
- adiabatic Berry phase



- lattice shaking



- Raman hopping



Dalibard, Gerbier, Juzeliunas, and Öhberg, RMP 83, 1523 (2011)

Goldman, Juzeliunas, Öhberg, and Spielman, Rep. Prog. Phys. 77 126401 (2014)

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Single particle Hamiltonian

$$H = (p_x \mathbb{I} + \mathcal{A}_x)^2 + (p_y \mathbb{I} + \mathcal{A}_y)^2$$

$$m = 1/2 \\ \hbar = e = c = 1$$

- U(1) magnetic field + SU(2) spin-orbit = U(2) gauge

$$\vec{\mathcal{A}} = \left(-\frac{y\mathcal{B}}{2} \mathbb{I} + q\sigma_x; \frac{x\mathcal{B}}{2} \mathbb{I} + q\sigma_y, ; 0 \right) \quad [\mathcal{A}_x, \mathcal{A}_y] \neq 0$$

- spatial components of $F^{\mu\nu} = [D^\mu, D^\nu]$ are not everything !

$$\vec{F} = \vec{\nabla} \times \vec{\mathcal{A}} + i\vec{\mathcal{A}} \times \vec{\mathcal{A}} = (\mathcal{B} \mathbb{I} - 2q^2 \sigma_z) \hat{z} \quad U(2) \neq U(1) \times U(1)$$

Single particle Hamiltonian

$$H = (p_x \mathbb{I} + \mathcal{A}_x)^2 + (p_y \mathbb{I} + \mathcal{A}_y)^2 + \frac{\omega^2 r^2}{4} \mathbb{I}$$

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- weak harmonic confinement

$$B \equiv \sqrt{\mathcal{B}^2 + \omega^2} \\ \Delta \equiv B - \mathcal{B} \sim \omega^2 / 2B \quad \longrightarrow \quad -L_z \Delta$$

Single particle Hamiltonian

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- compensating Zeeman field

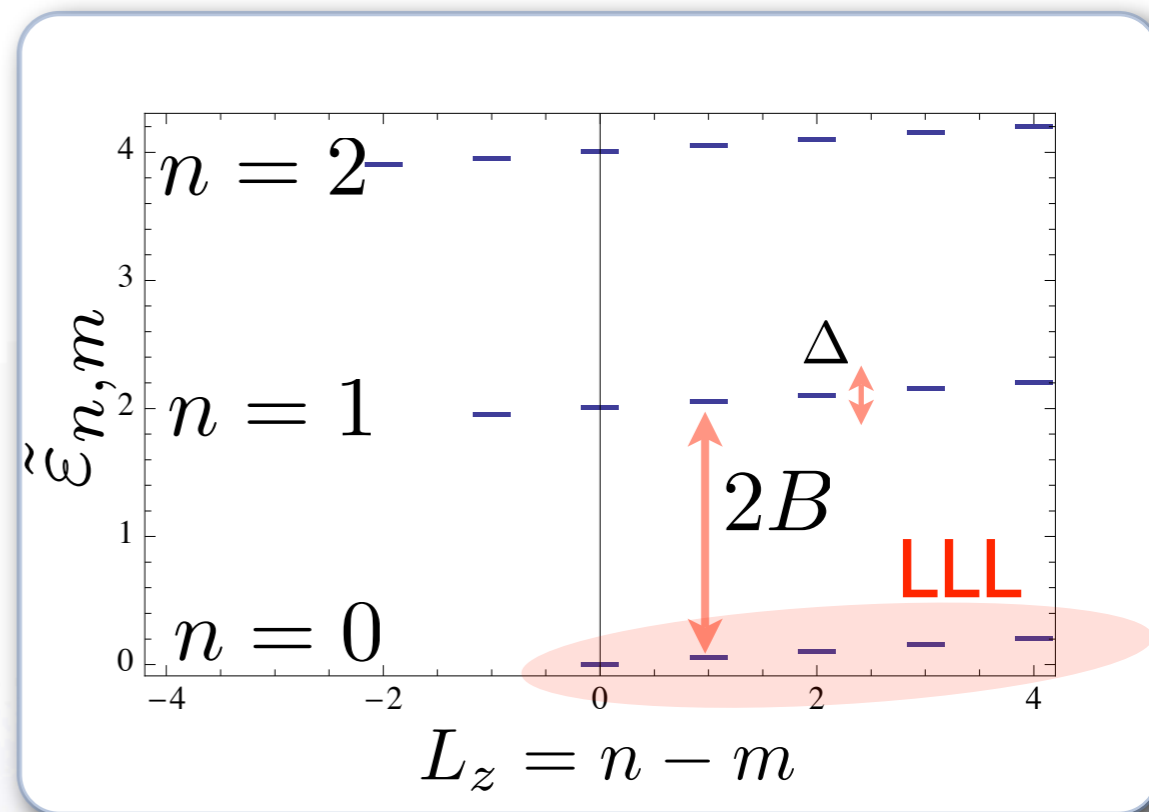
$$V_s(x, y) = -q\Delta (y\sigma_x - x\sigma_y)$$

Deformed Landau Levels

$$H = \left(p_x - \frac{yB}{2} + q\sigma_x \right)^2 + \left(p_y + \frac{xB}{2} + q\sigma_y \right)^2 - L_z \Delta$$

- $q=0$: orbital eigenstates $\psi_{n,m}$
- Lowest Landau Level (LLL) approx.

$$N\Delta \ll 2B \quad k_B T \ll 2B$$



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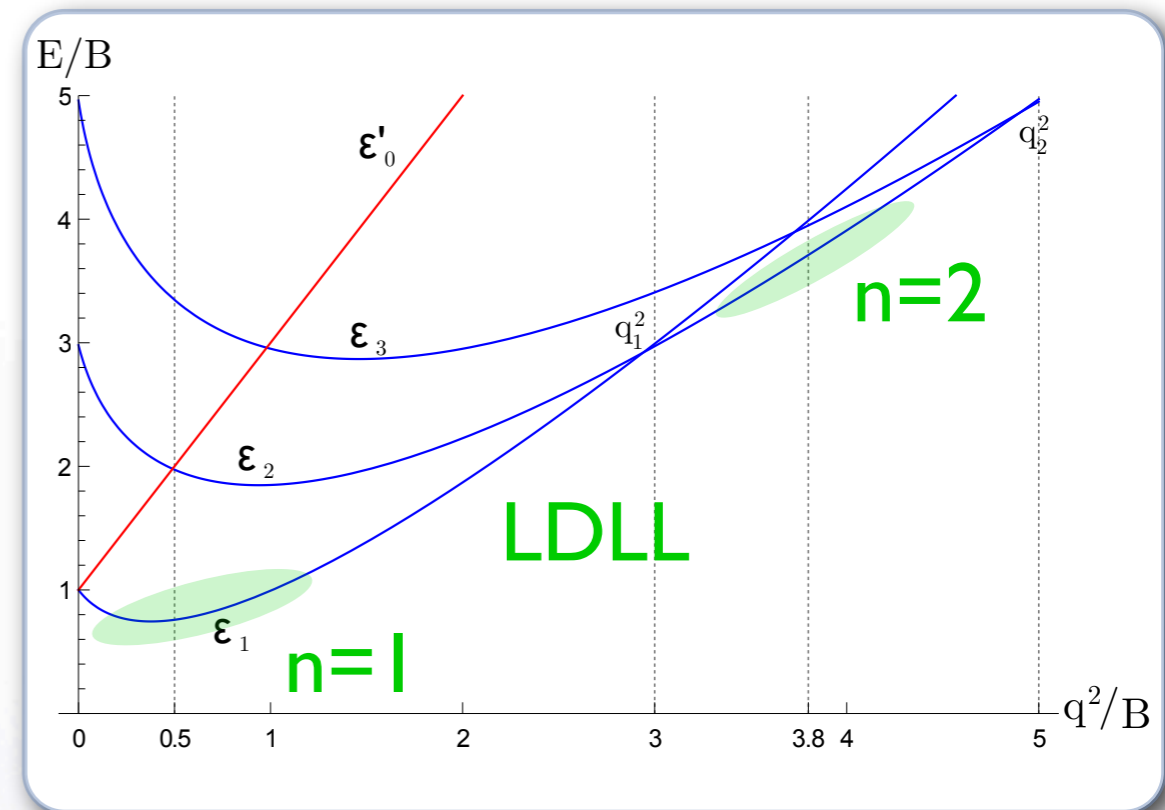
- $q \neq 0$: Jaynes-Cummings in the basis

$$\{ \psi_{n-1,m} | \uparrow \rangle, \psi_{n,m} | \downarrow \rangle \}$$

$$H_{n,m} = \tilde{\epsilon}_{n,m} \mathbb{I} + D_n \begin{pmatrix} -\cos \varphi_n & \sin \varphi_n \\ \sin \varphi_n & \cos \varphi_n \end{pmatrix}$$

- Deformed Landau Levels (DLL) $\chi_{n,m}$

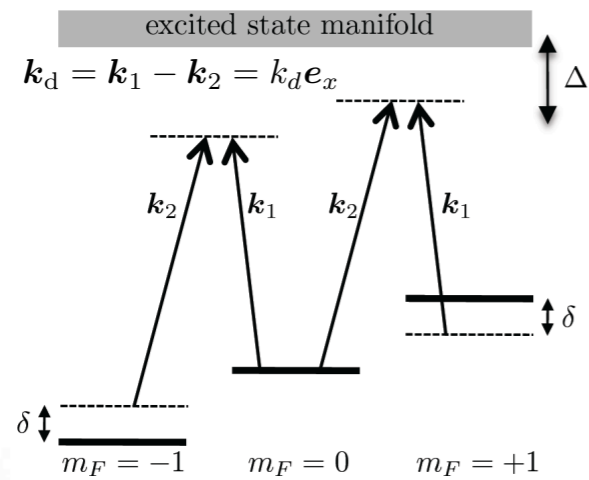
$$\epsilon_{n,m} = \tilde{\epsilon}_{n,m} - D_n$$



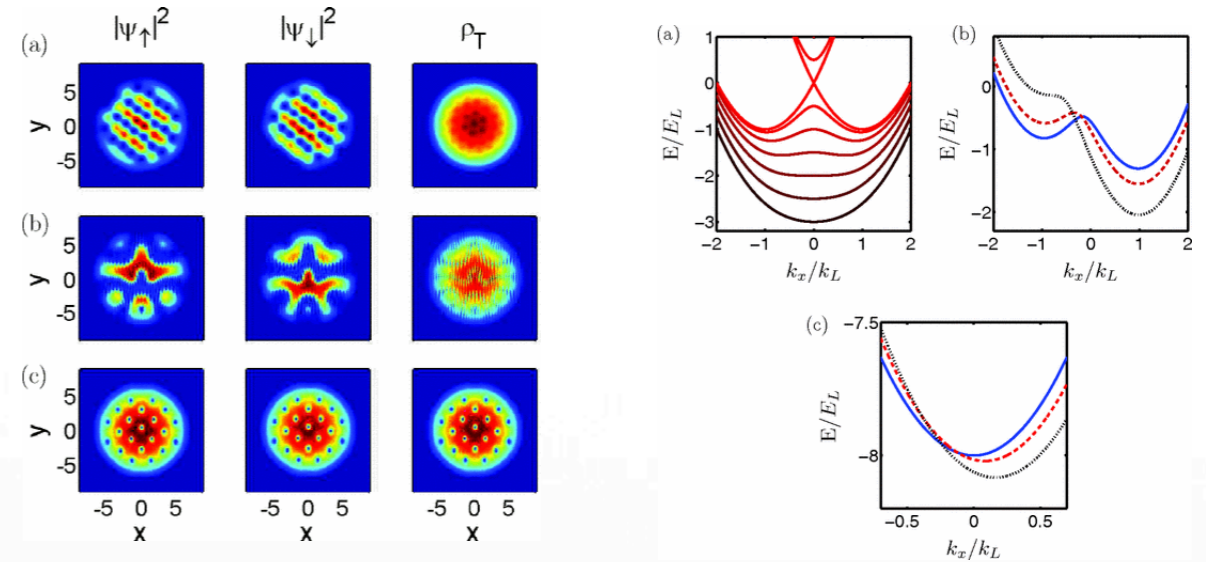
$$\sin \varphi_n = \frac{2q\sqrt{2Bn}}{\sqrt{(B-\Delta/2)^2 + 8q^2 Bn}}$$

Experimental considerations

- **U(1) magnetic field + SU(2) spin-orbit** M. Burrello, A. Trombettoni, PRA **84**, 043625 (2011)



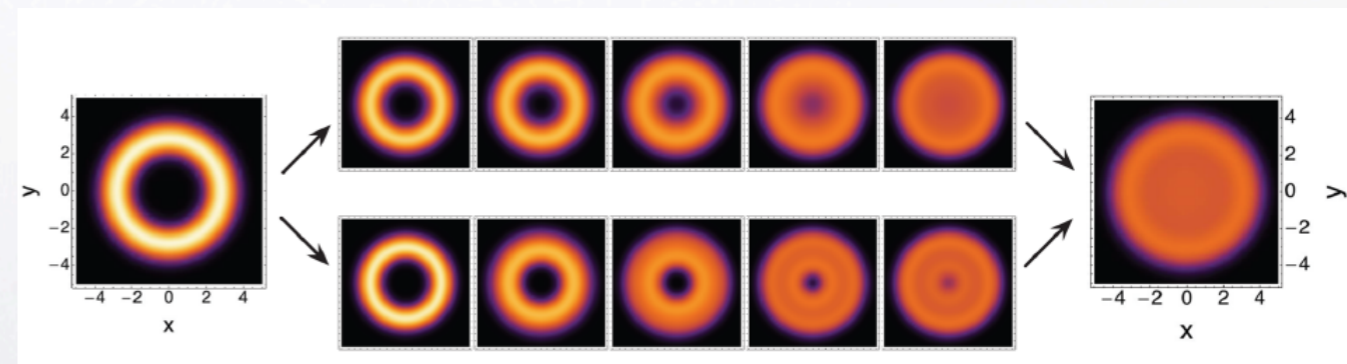
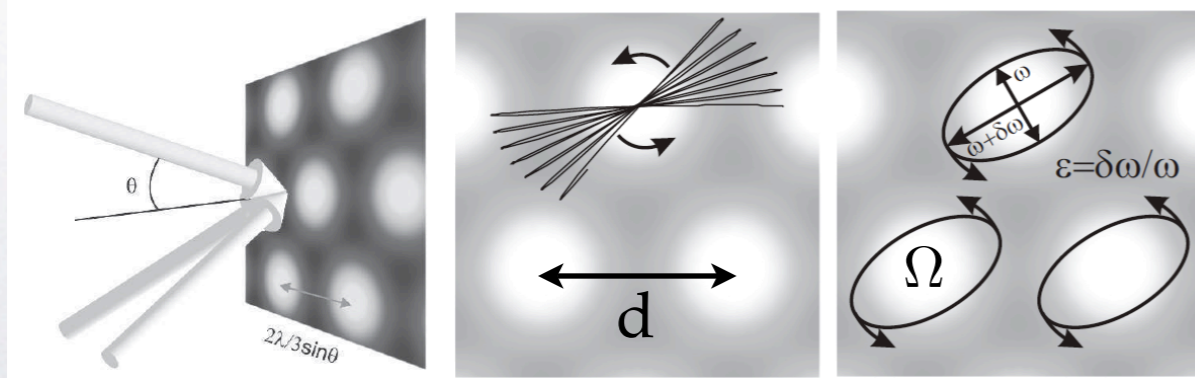
J. Ruseckas, et al., PRL **95**, 010404 (2005).



J. Radic, et al., Phys. Rev. A **84**, 063604 (2011).

- **access to LLL regime: difficult by rotation!**

N. R. Cooper, Adv Phys **57**, 539 (2008)
A. L. Fetter, Rev. Mod. Phys. **81**, 647 (2009)



Popp, et al, PRA **70**, 053612 (2004) / Gemelke et al., arXiv:1007.2677

Roncaglia, Rizzi, Dalibard, Sci. Rep. **1**, 43 (2011)

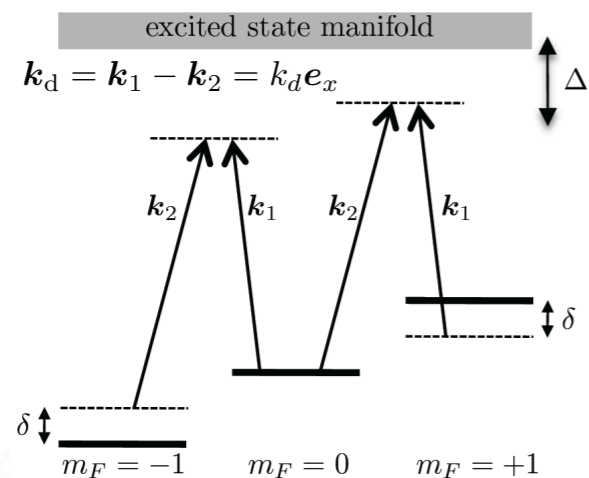
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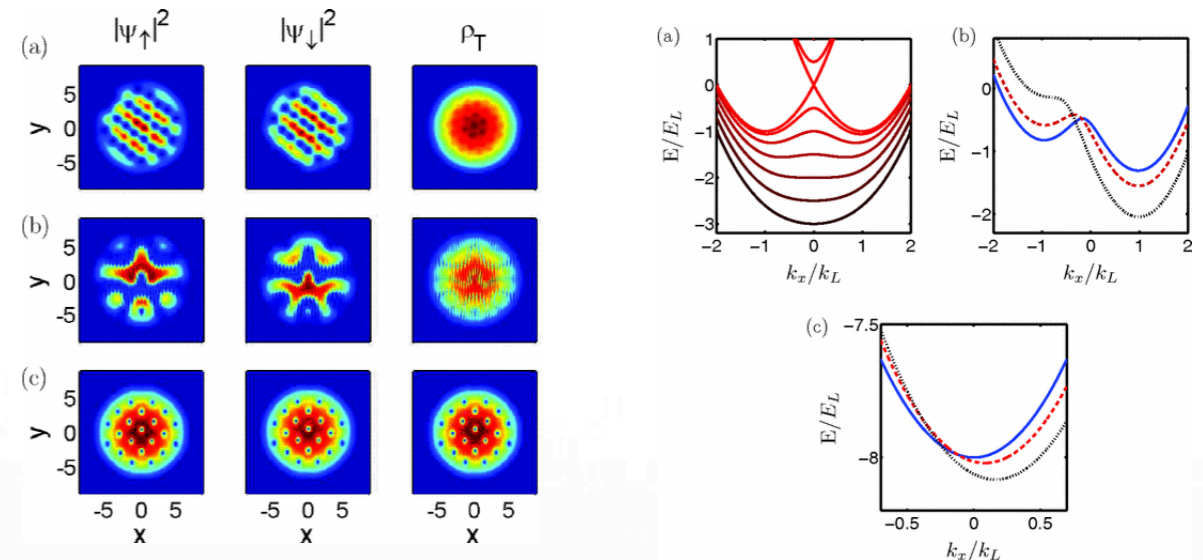
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- new hopes: synthetic optical gauge potentials & dimensions

see I. Spielman, I. Bloch, M. Lewenstein, etc.

- some rough estimates

$$\mathcal{B}\ell^2 \simeq 1 \wedge \ell \simeq 0.5\mu m \longrightarrow \mathcal{B} \simeq 4 \cdot 10^{-19} g/s$$

$$\omega \simeq 10 \div 100 \text{ Hz} \longrightarrow \Delta/\mathcal{B} \simeq 0.005 \div 0.05$$

$$\hbar/q \simeq 1\mu m \longrightarrow q^2/\hbar\mathcal{B} \simeq 0.1 \div 5$$

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Haldane pseudopotentials

- partial wave decomposition of (central) interaction potentials

$$\mathcal{H}_{\text{int}} = \sum_M \sum_{m_{\text{rel}}} \boxed{W_{m_{\text{rel}}}} \sum_{m_1, m_2} g[m_{\text{rel}}, M, m_1] g[m_{\text{rel}}, M, m_2] c_{M-m_1}^\dagger c_{m_1}^\dagger c_{m_2} c_{M-m_2}$$

- polarized electrons in Coloumb potential $W_1 \gg W_3 \gg \dots$

\implies Laughlin Ansatz ! $\nu = 1/3$ $\Psi_{1/3} = \prod_{i < j} (z_i - z_j)^3 \equiv \Theta^3$

- s-wave scattering approx. for cold bosons ! *only* $W_0 \neq 0$

$$\mathcal{H}_2 = c_2 \sum_{i < j} \delta(z_i - z_j) \quad c_2 = \sqrt{8\pi a}/\xi_z \quad \nu = 1/2 \quad \Psi_{1/2} = \prod_{i < j} (z_i - z_j)^2 \equiv \Theta^2$$

- LL filling factor $\nu \simeq \lim_{N \rightarrow \infty} \frac{N}{m_{\text{max}}}$

Haldane pseudopotentials

$$\mathcal{H}_{\text{int}} = \sum_M \sum_{m_{\text{rel}}} W_{m_{\text{rel}}} \sum_{m_1, m_2} g[m_{\text{rel}}, M, m_1] g[m_{\text{rel}}, M, m_2] c_{M-m_1}^\dagger c_{m_1}^\dagger c_{m_2} c_{M-m_2}$$

- spin 1/2 fermions within U(2) DLL

$$z = x - iy$$

$$W_{m_{\text{rel}}}^{(n)} = V_{m_{\text{rel}}}^{n-1, n-1} \cos^4 \frac{\varphi_n}{2} + V_{m_{\text{rel}}}^{n, n} \sin^4 \frac{\varphi_n}{2} + V_{m_{\text{rel}}}^{n, n-1} \frac{\sin^2 \varphi_n}{2}$$

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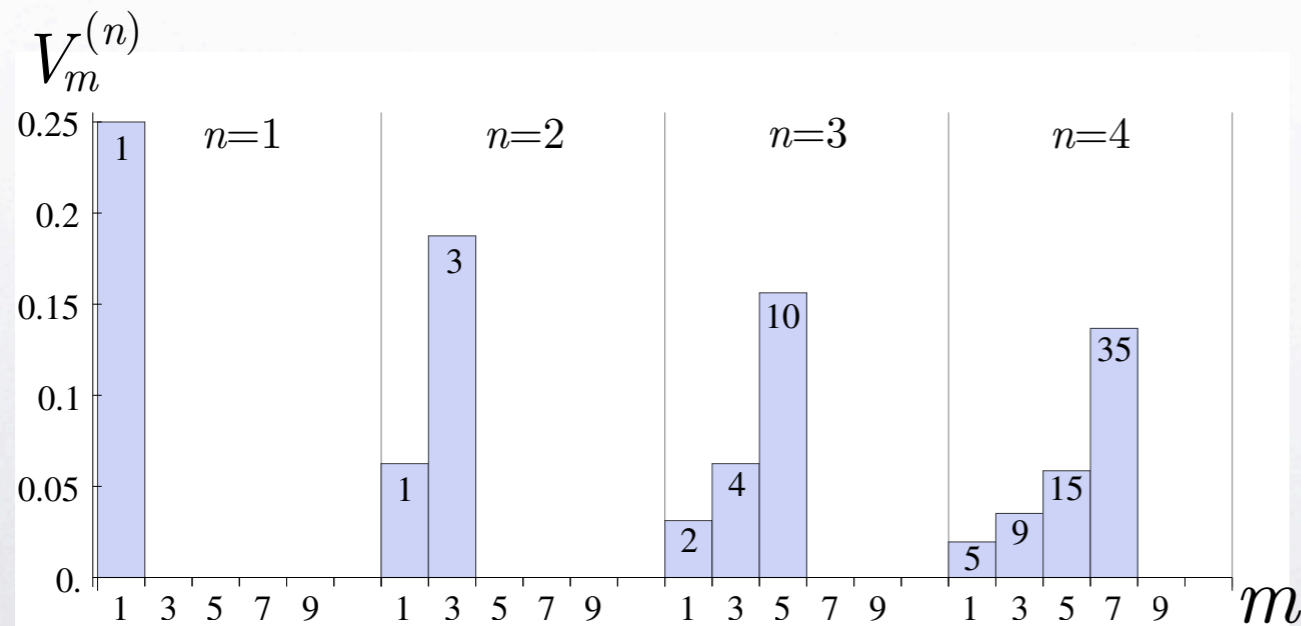
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- only interspecies contact interactions

$$\hat{V} = v \sum_{i < j} \delta(z_i - z_j) |\uparrow\downarrow\rangle\langle\uparrow\downarrow|$$

! Non-Monotonic HP !
? new FQH states ?



Haldane pseudopotentials

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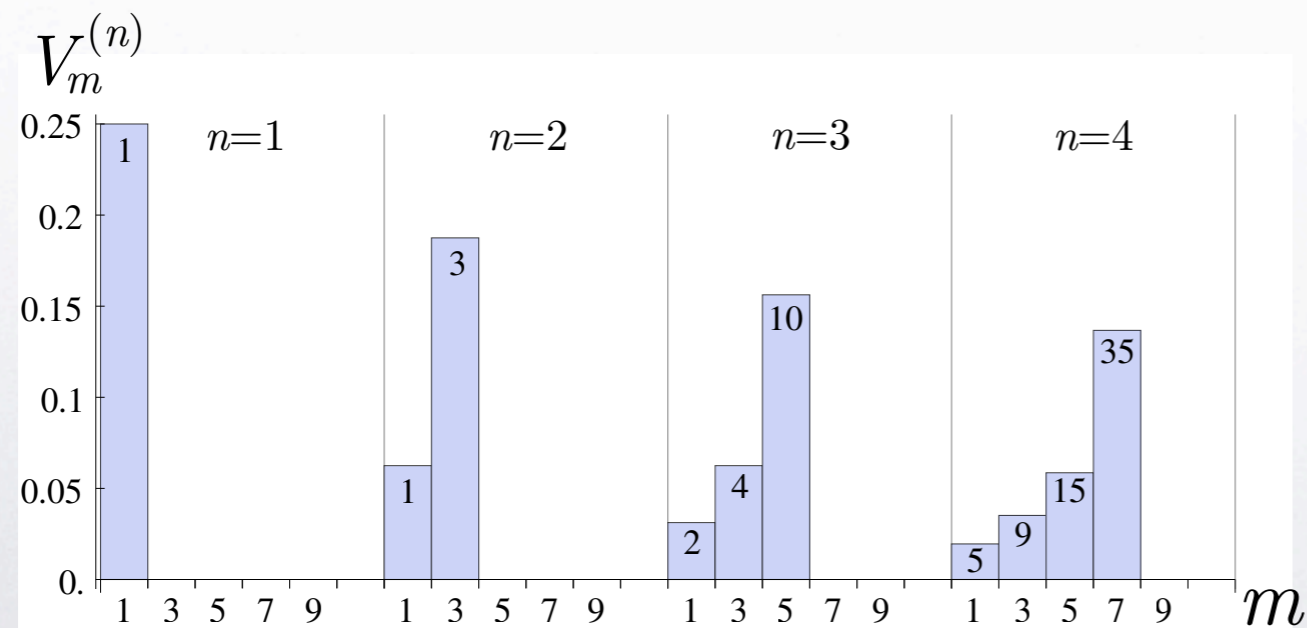
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Dipoles needed for bosons

T. Grass, et al, PRA 89, 013623 (2014)

Incompressible states

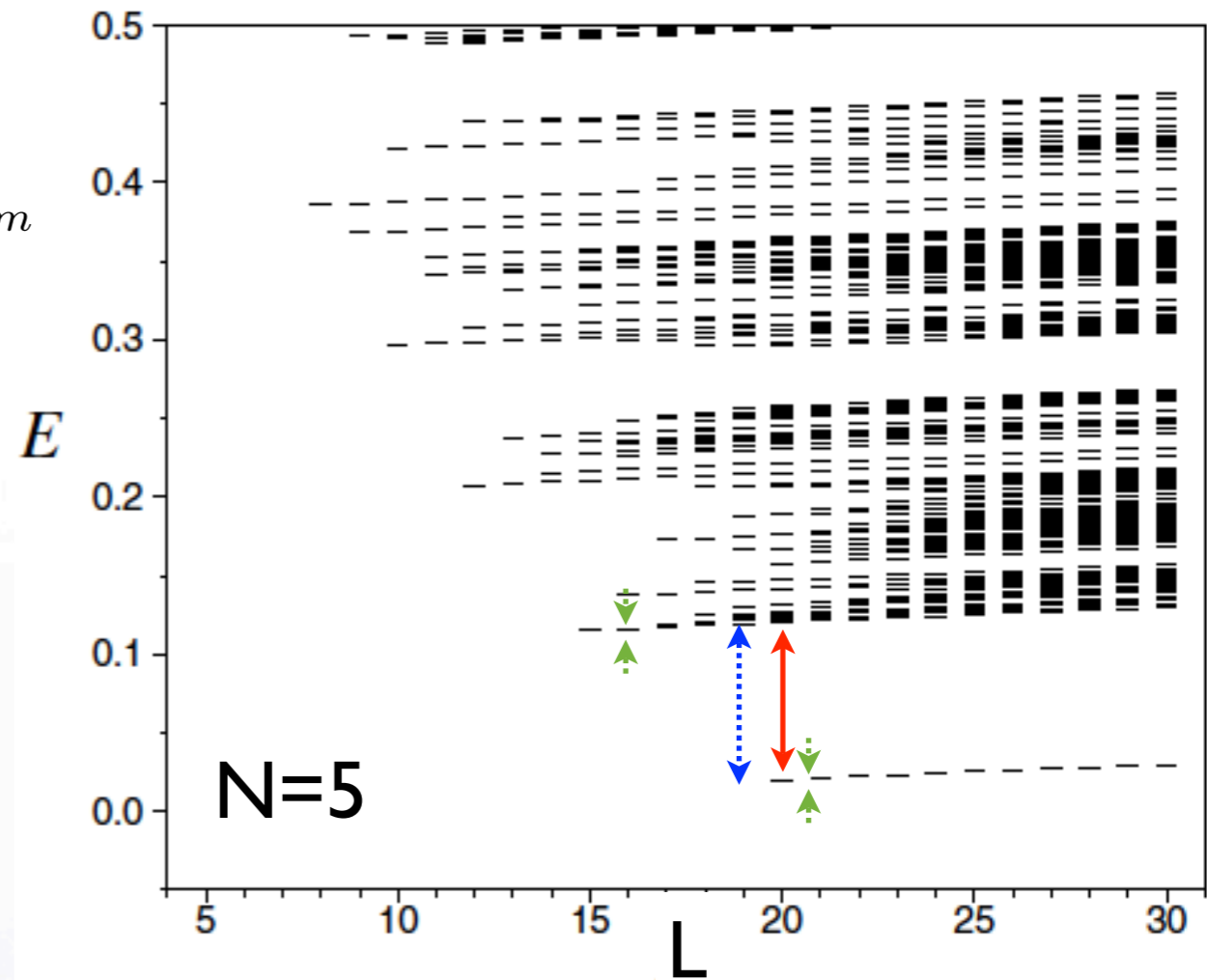
- $[\mathcal{H}, L_z] = 0$: yrast spectrum (disk)

$$\mathcal{E}(L) = E(L) + L \Delta \quad L_z = Nn - \sum m a_m^\dagger a_m$$

- incompressibility gap

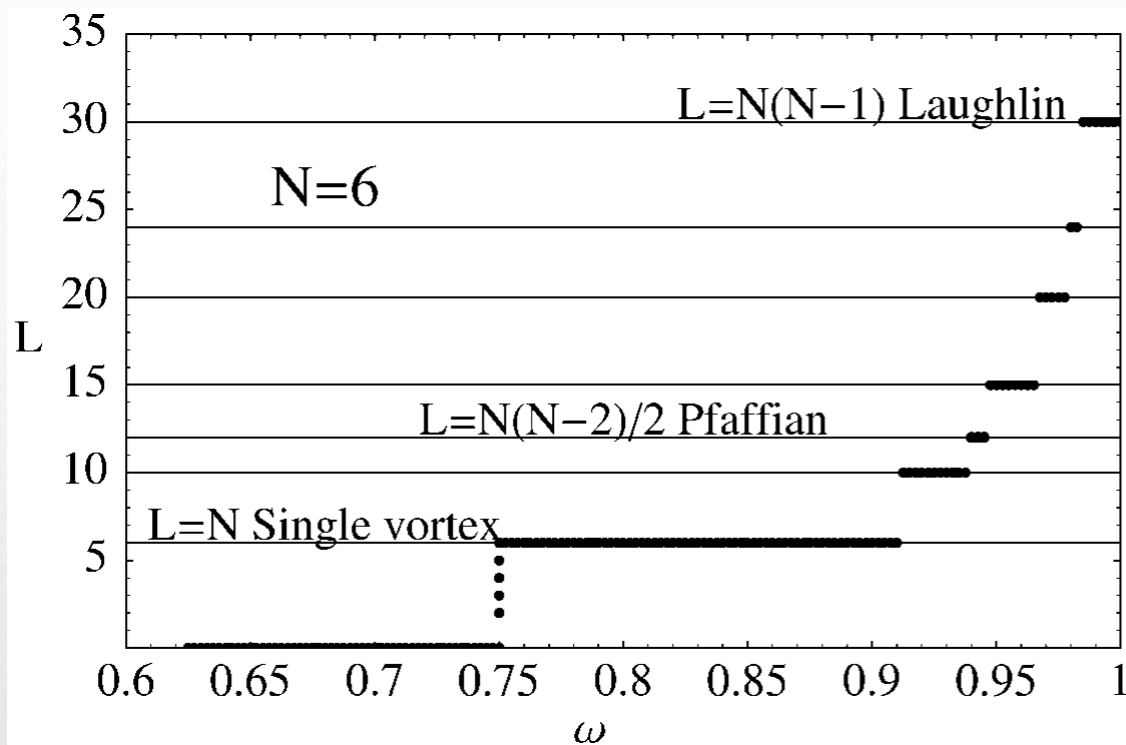
$$\mathcal{D}(L) = \min[\delta E(L), E(L) - E(L-1)]$$

decides over stabilized states



Paredes, et al., PRL **87**, 010402 (2001)

- quasi-hole excitation $\Delta_{\text{qh}} \simeq \delta$



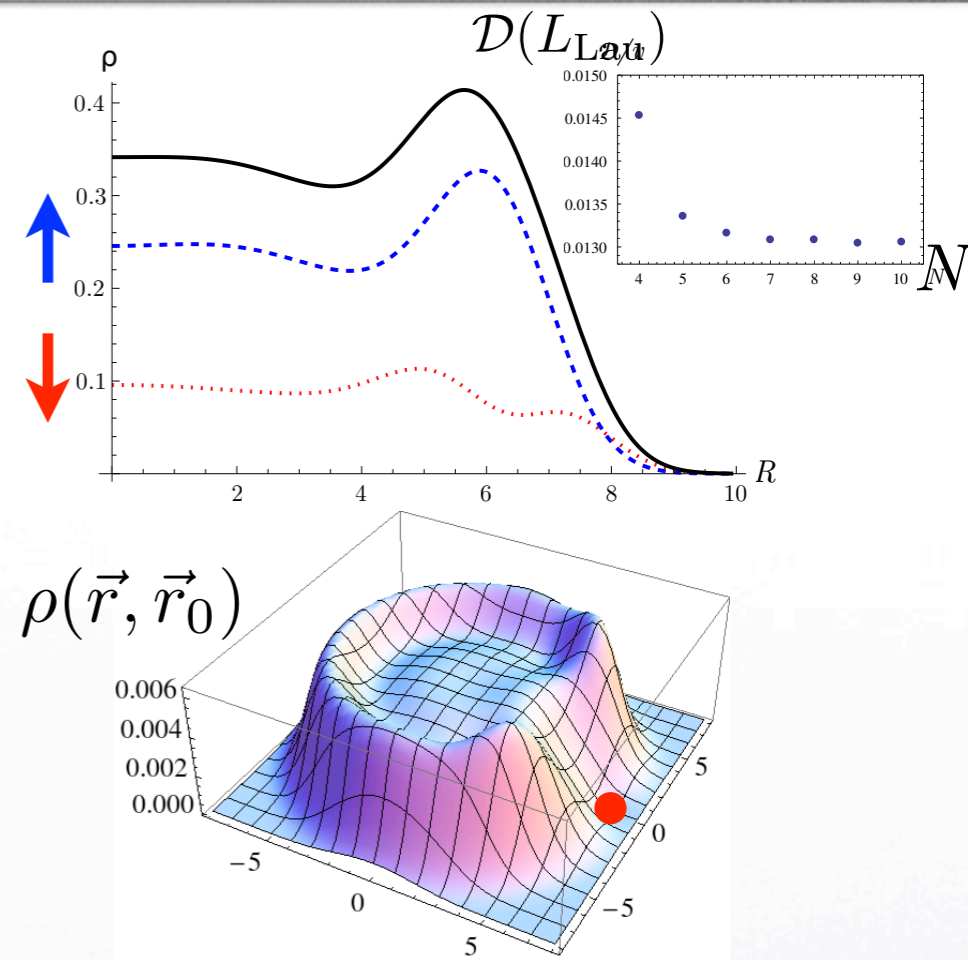
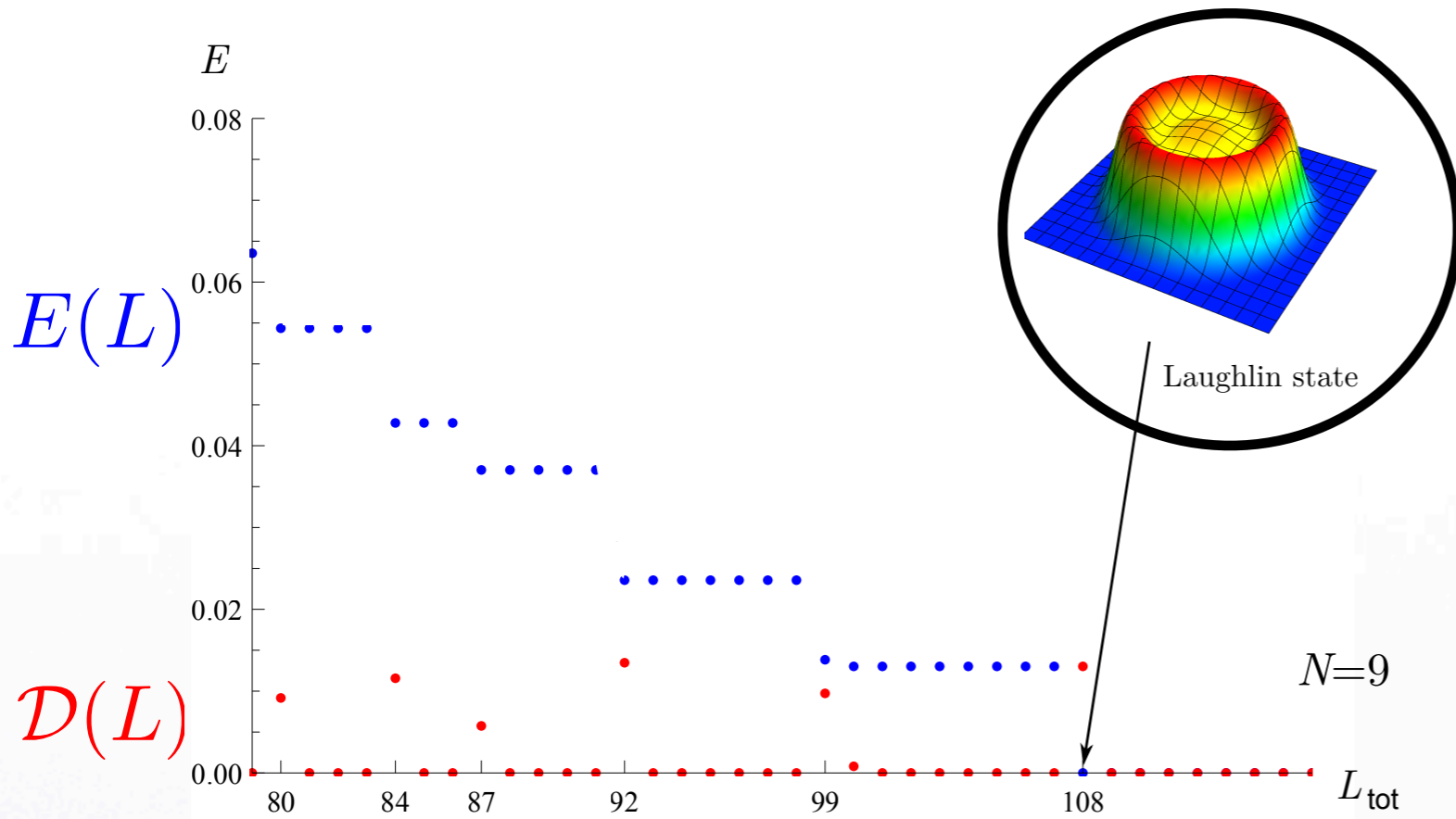
Wilkin, Gunn, PRL **84**, 6 (2000)

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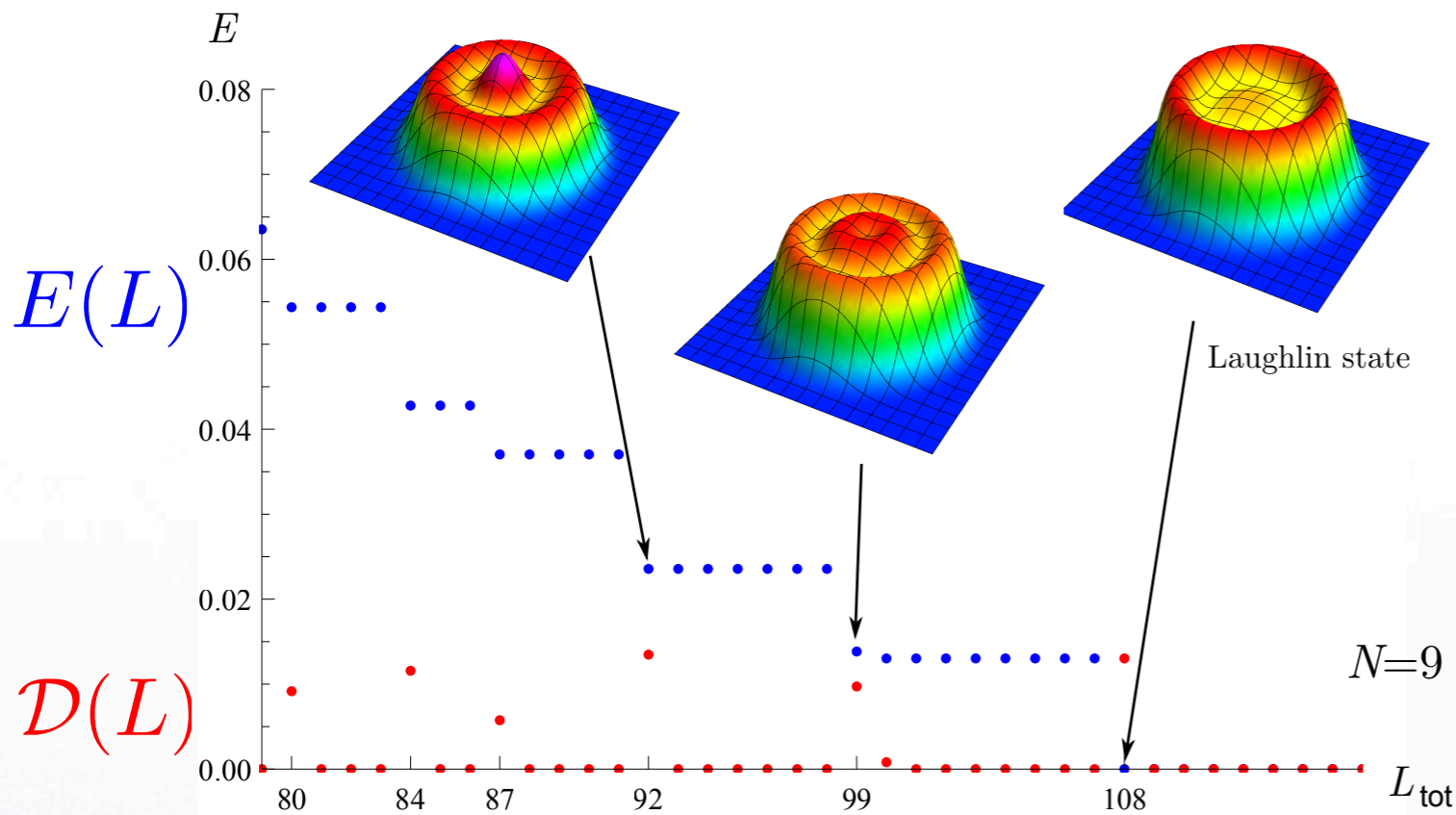
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1st DLL: Laughlin



- Laughlin ansatz $\nu = 1/3$ $\Psi_{1/3} = \prod_{i < j} (z_i - z_j)^3 \equiv \Theta^3$
- incompressibility gap $E_{\text{qp}} \simeq \mathcal{D}(L_{\text{Lau}}) \approx 0.013\nu$ $L_{\text{Lau}} \equiv 3N(N-1)/2$

1st DLL: Laughlin



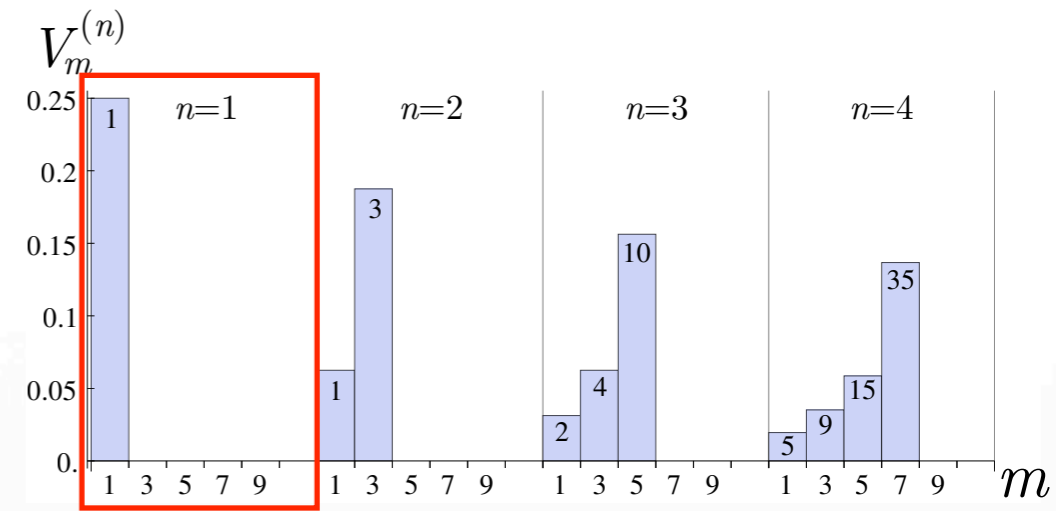
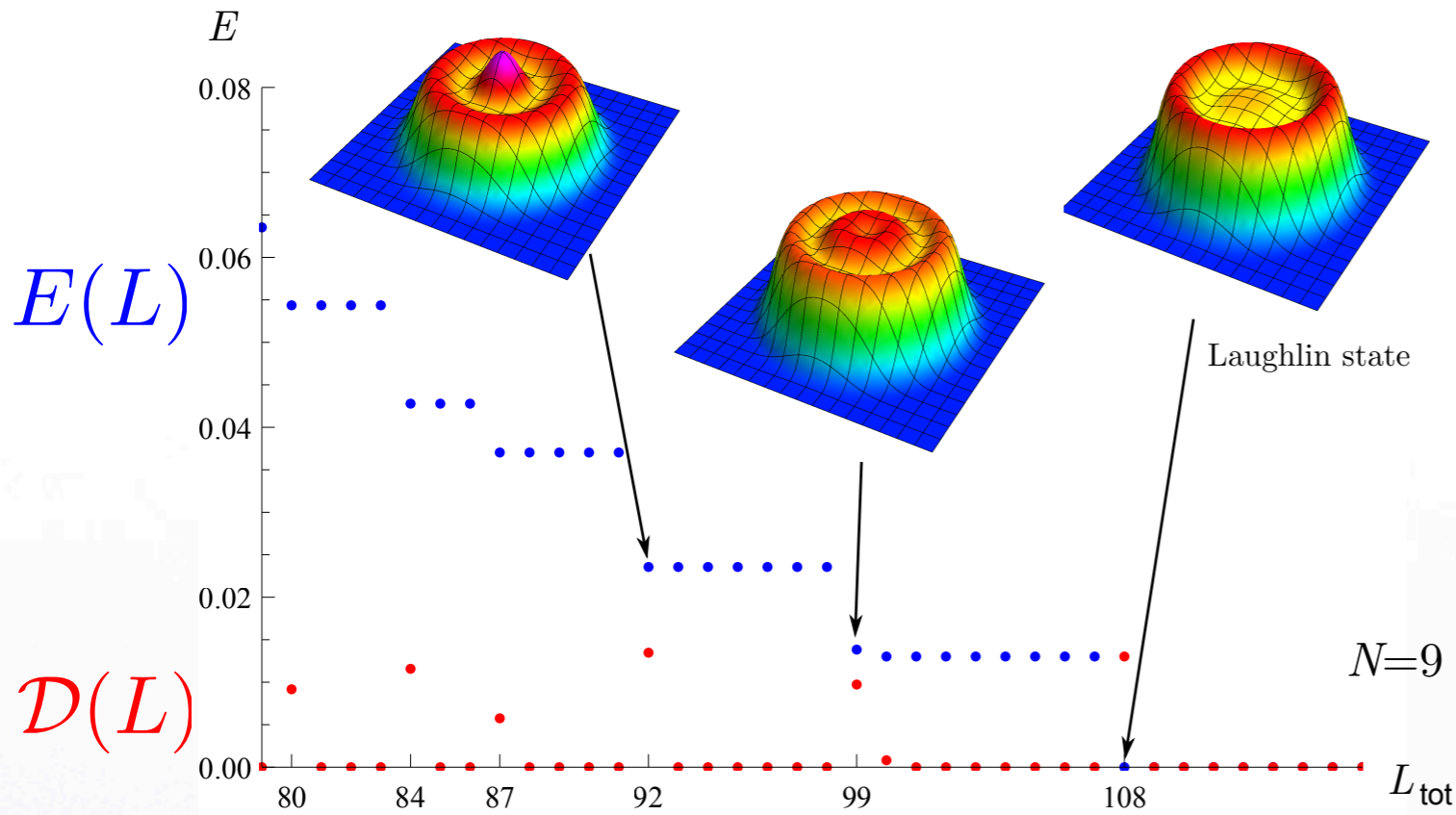
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- plateau lengths $(N, N-2, N-4, \dots) = \text{CF theory}$ A. Cappelli, et al, PRB **58**, 16291 (1998)
G. Dev, J. K. Jain, PRB **45**, 1223 (1992)

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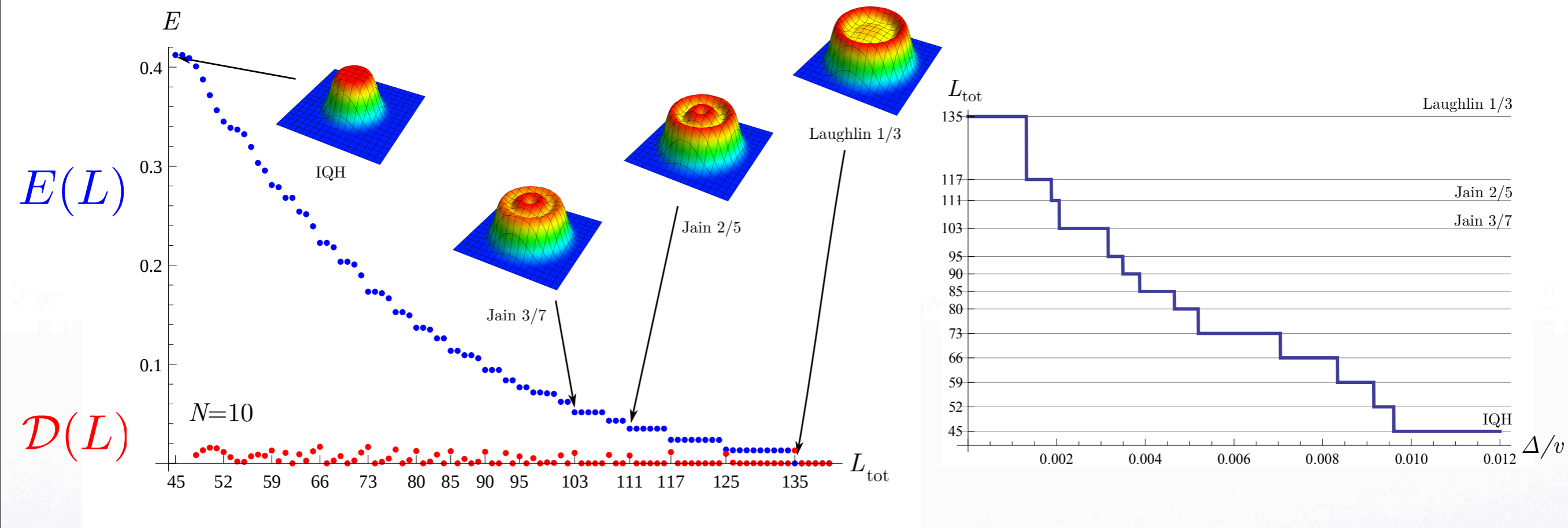
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1st DLL: Jain states



- numerics with small particle numbers \implies labelling complicated

e.g. $N=10$ $L=111$ Lau + 3qp or Jain 2/5 ?

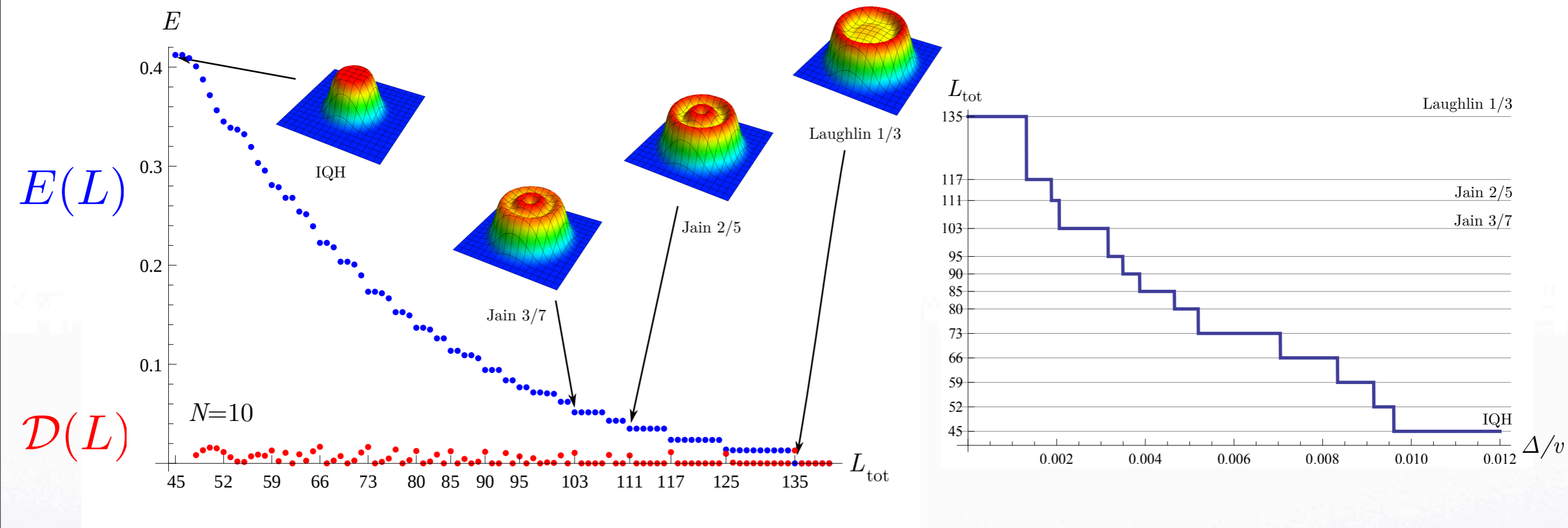
U. Zülicke, J.J. Palacios, A.H. MacDonald, PRB **67**, 045303 (2003)

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M.Burrello, **MR**,
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PRB **91**, 115117 (2015)

1st DLL: Jain states



- numerics with small particle numbers \implies labelling complicated

e.g. $N=10$ $L=111$ $Lau + 3qp$ or **Jain 2/5** ?

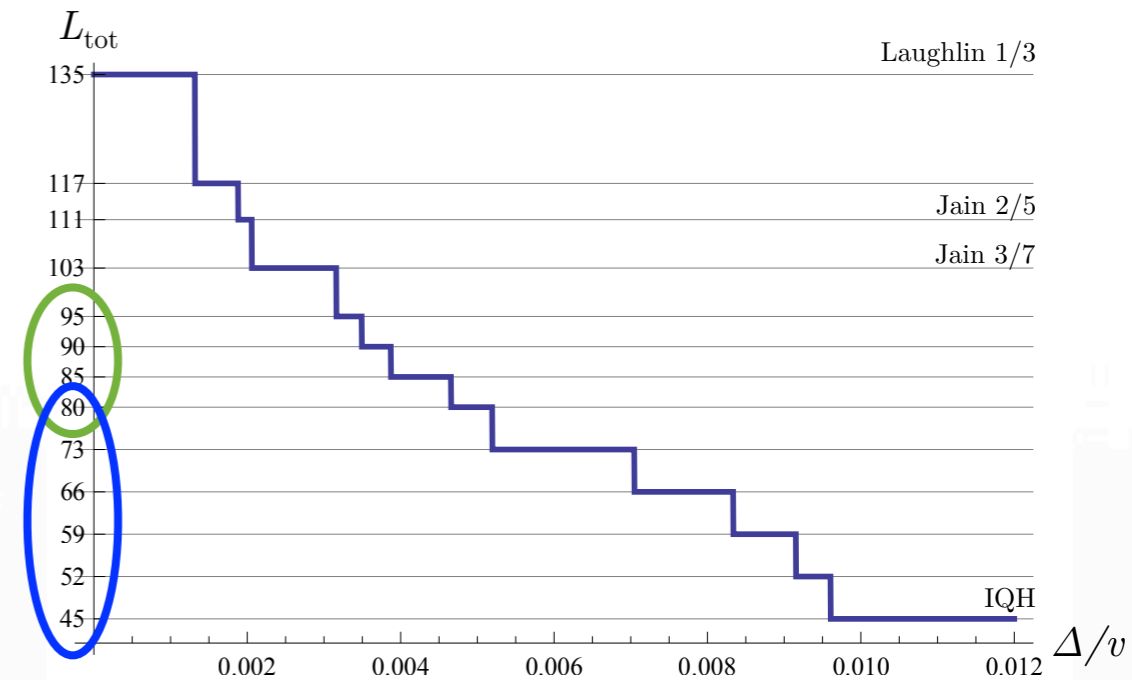
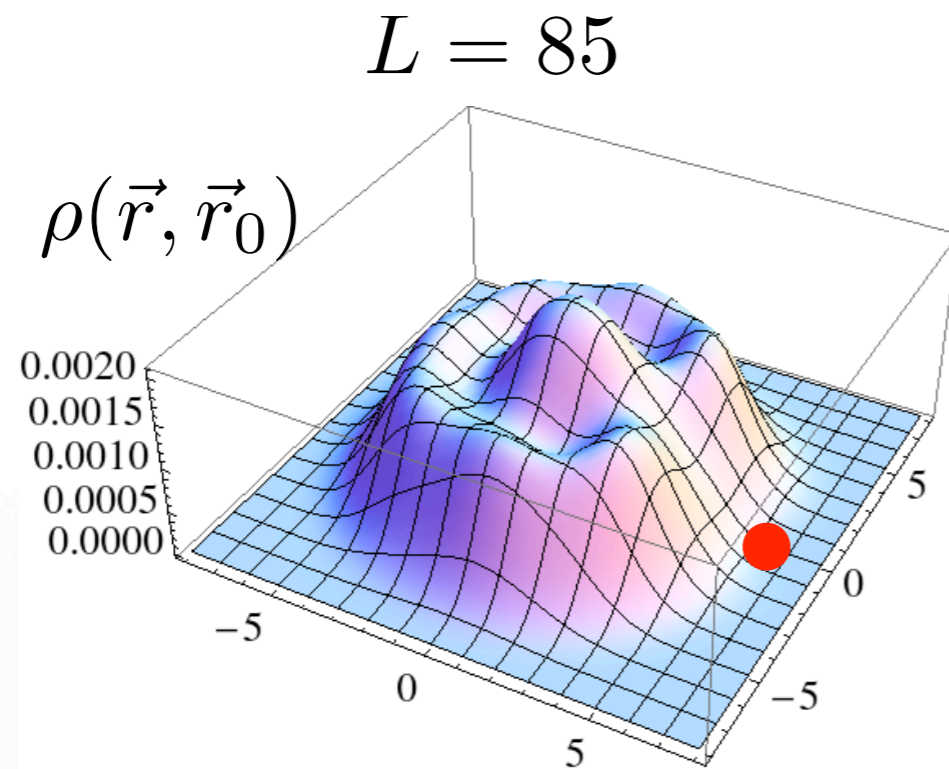
U. Zülicke, J.J. Palacios, A.H. MacDonald, PRB **67**, 045303 (2003)

Strongly correlated states
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1st DLL: other stable states



- regular pattern of ground states: \sim pseudo-crystals (atoms+fluxes)

NO Pfaffian $\Psi_{\text{Pf}} \propto \text{Pf} \left[\frac{1}{z_i - z_j} \right] \prod_{i < j} (z_i - z_j)^2$ at $L = N(N - 1) - N/2 = 85$

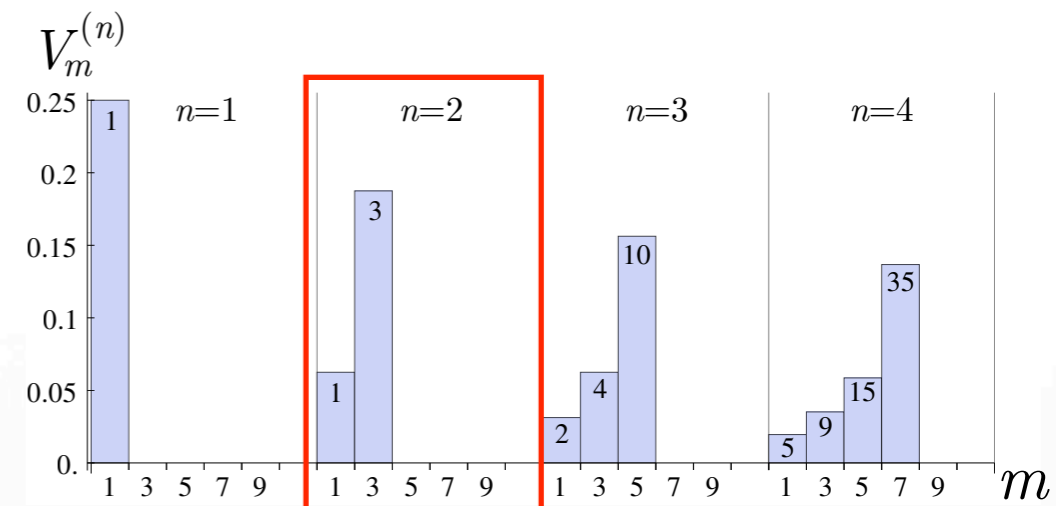
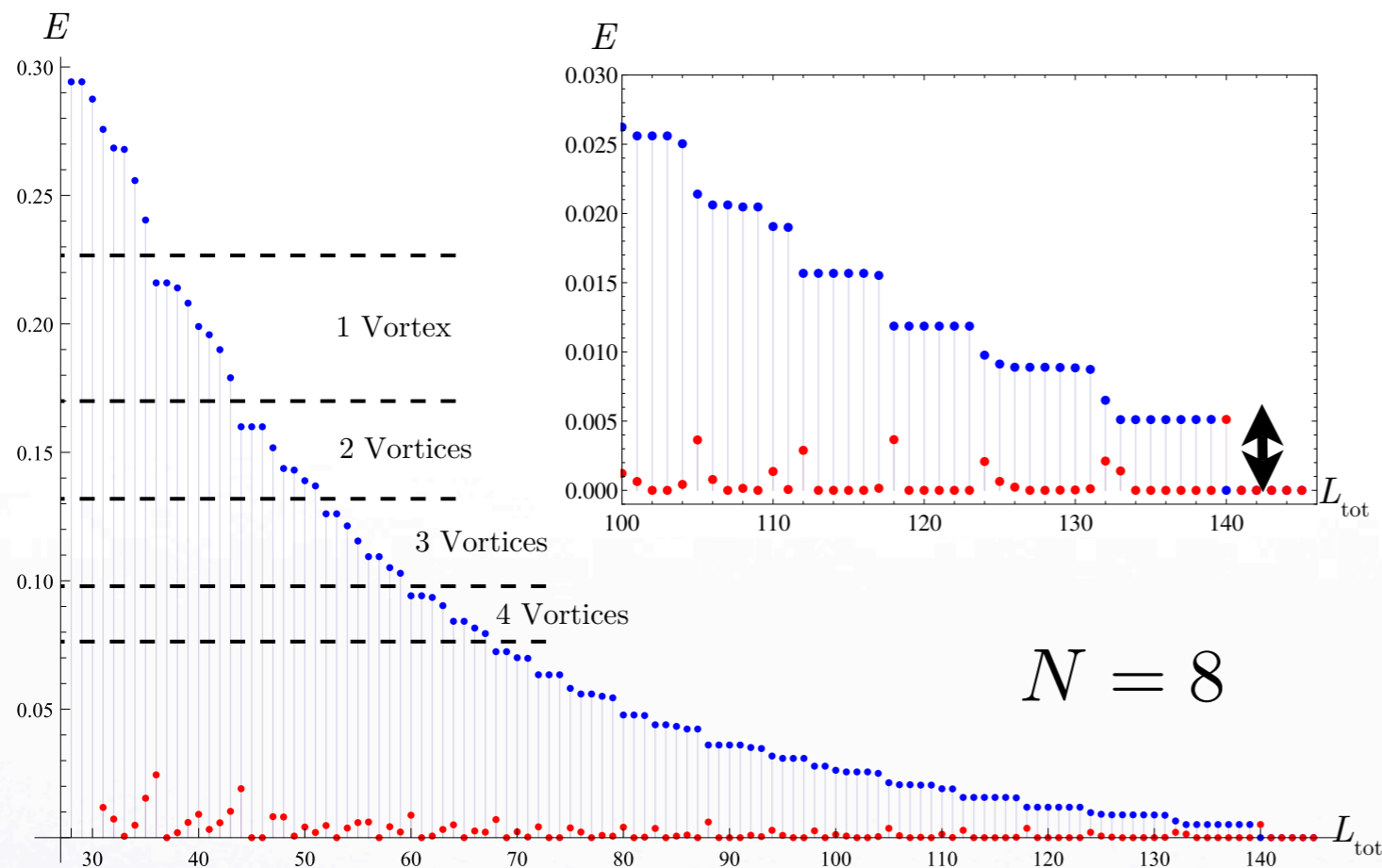
Outline

- Motivation: beyond standard
- $U(2)$ potential & deformed LL
- non-monotonic Haldane pseudopotentials
- Novel incompressible states: Haffnian?
- Entanglement spectrum
- Conclusions

Outline

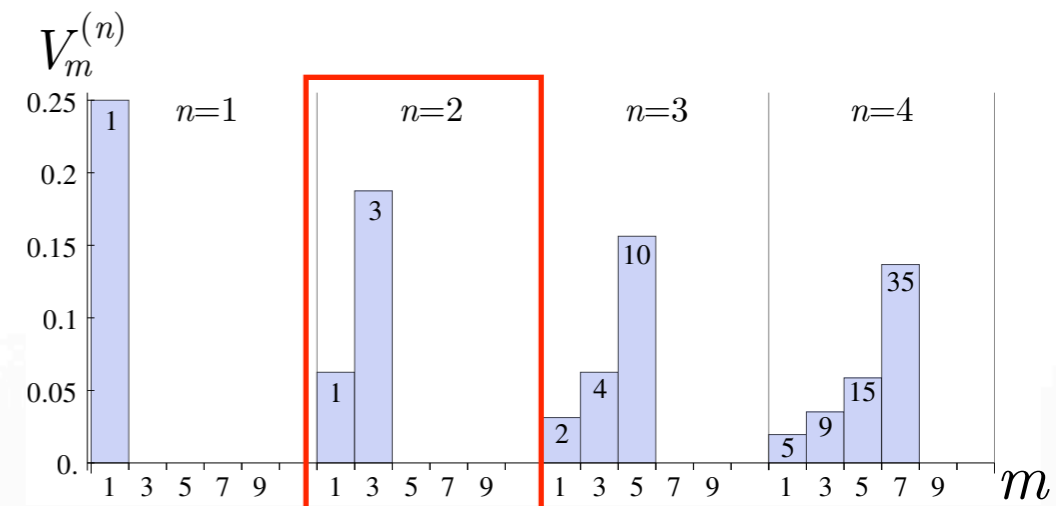
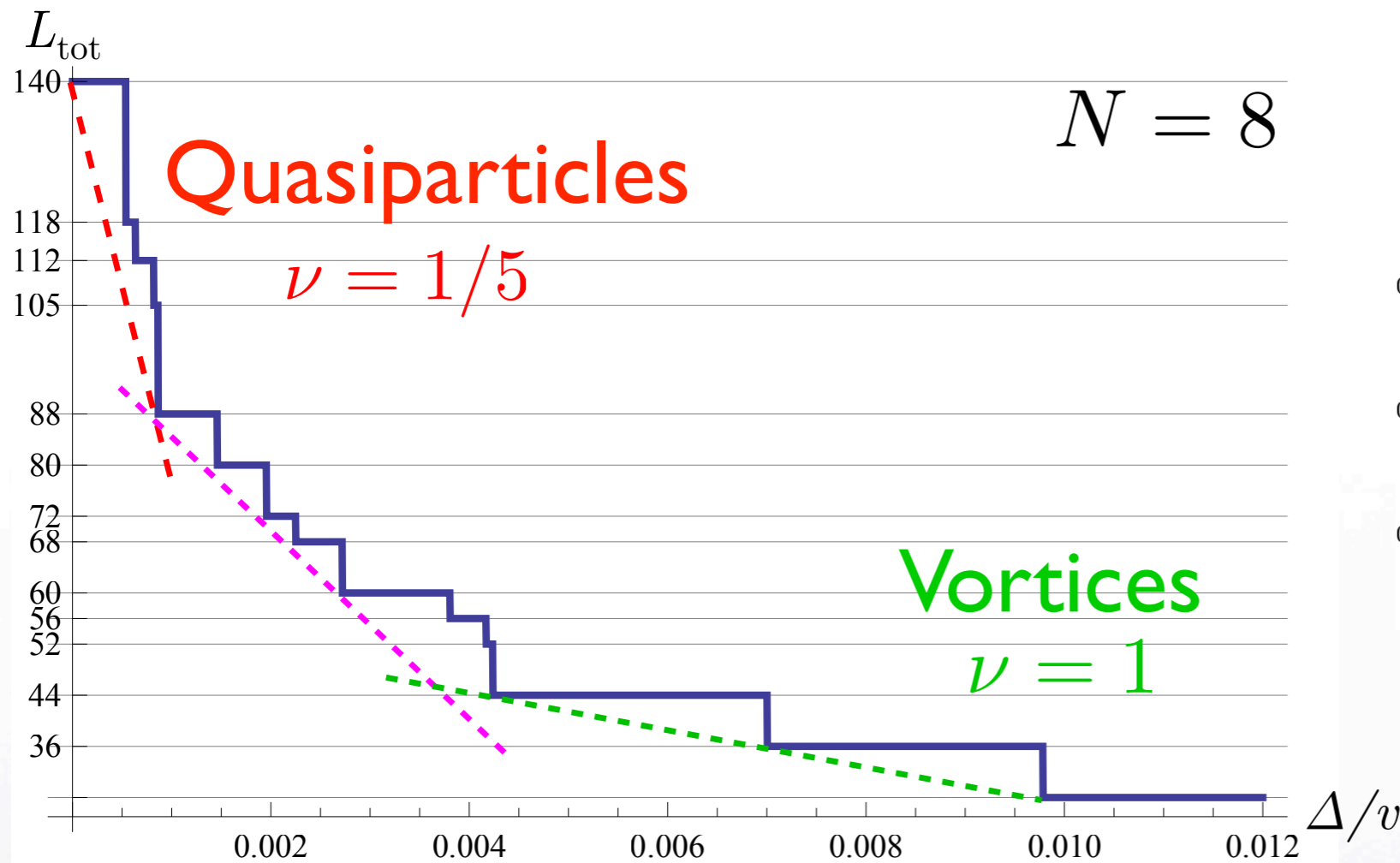
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2nd DLL: Laughlin & quasiparticles



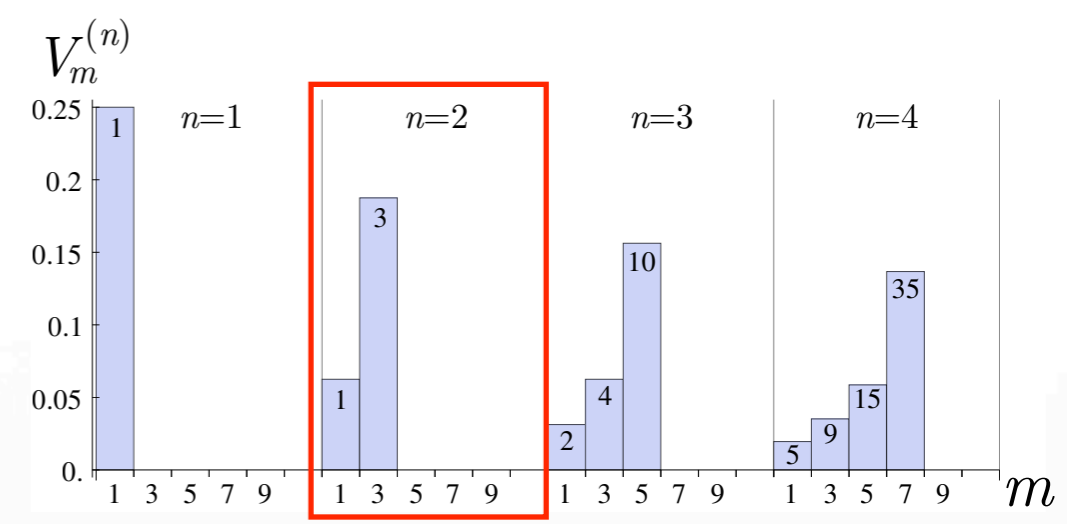
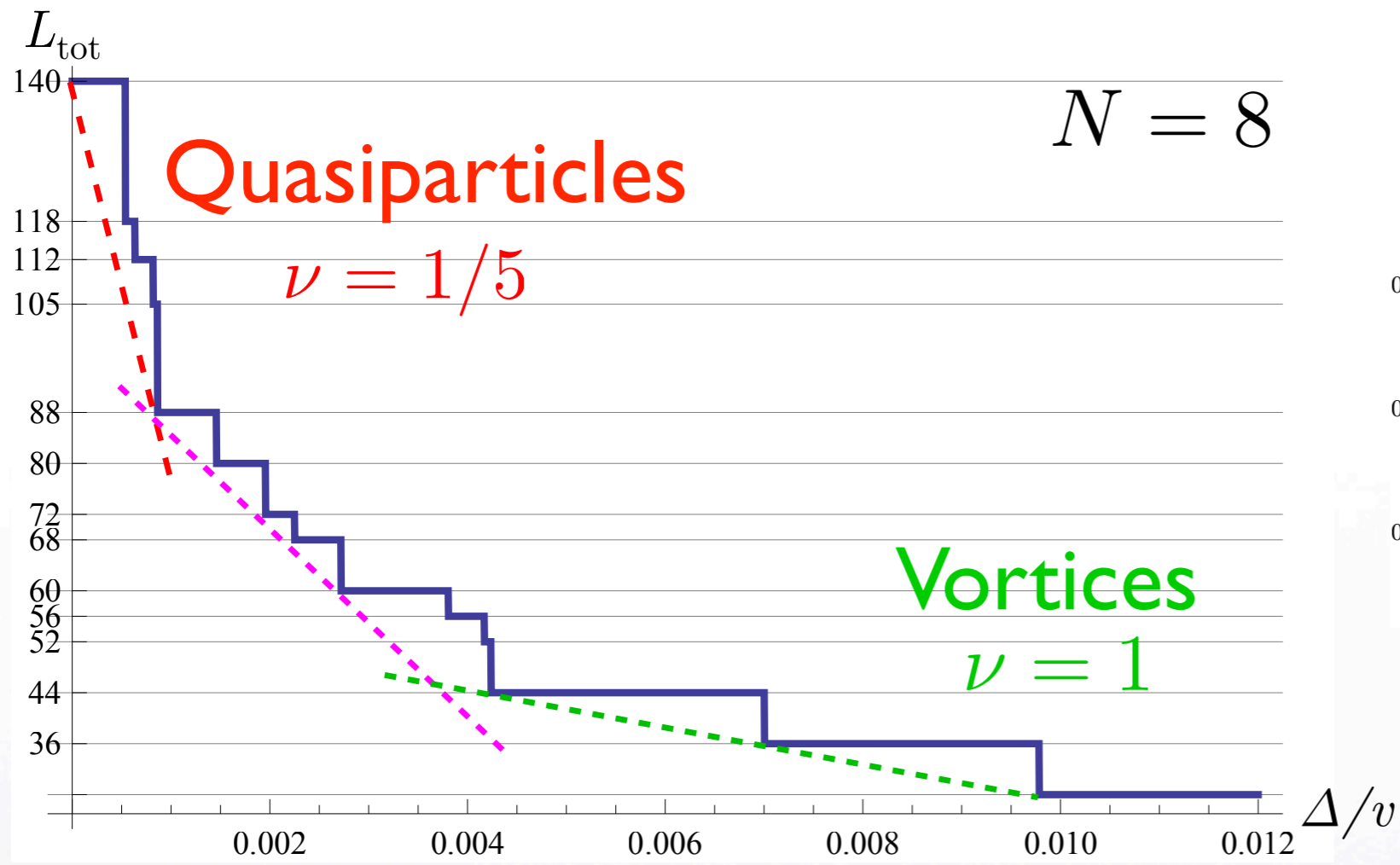
- Laughlin ansatz $\nu = 1/5$ $\Psi_{1/5} = \prod_{i < j} (z_i - z_j)^5 \equiv \Theta^5$
- incompressibility gap $E_{\text{qp}}^{(1/5)} \simeq \mathcal{D}(L_{\text{Lau}}) \approx 0.005v \simeq 0.3E_{\text{qp}}^{(1/3)}$
- smoothed transitions due to W_3 ...

2nd DLL: three regimes

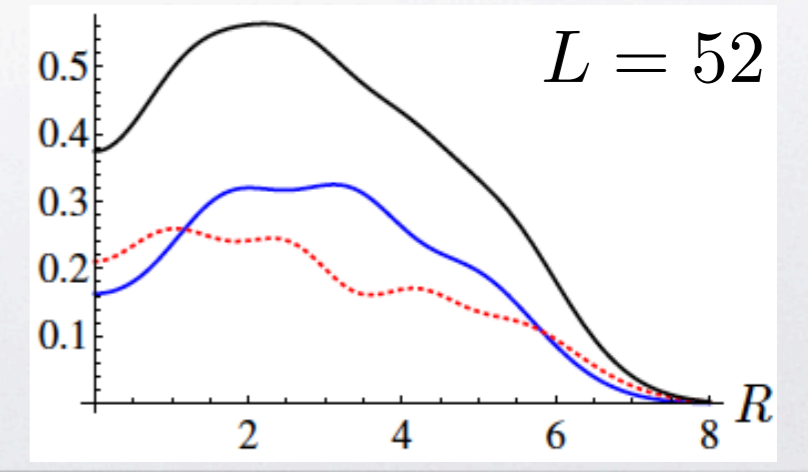
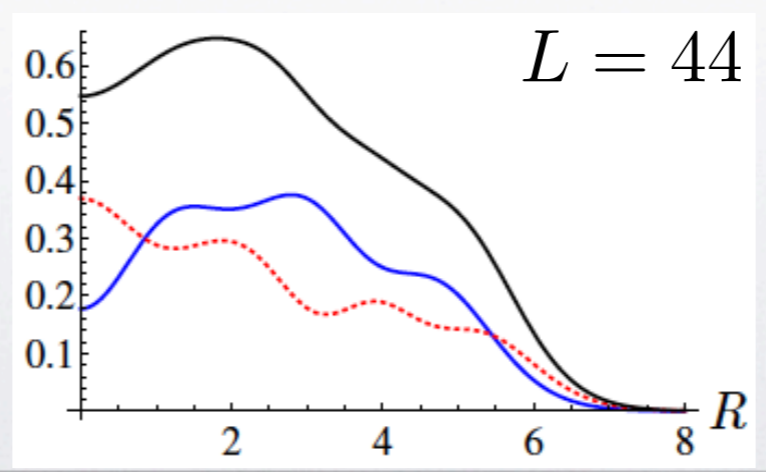
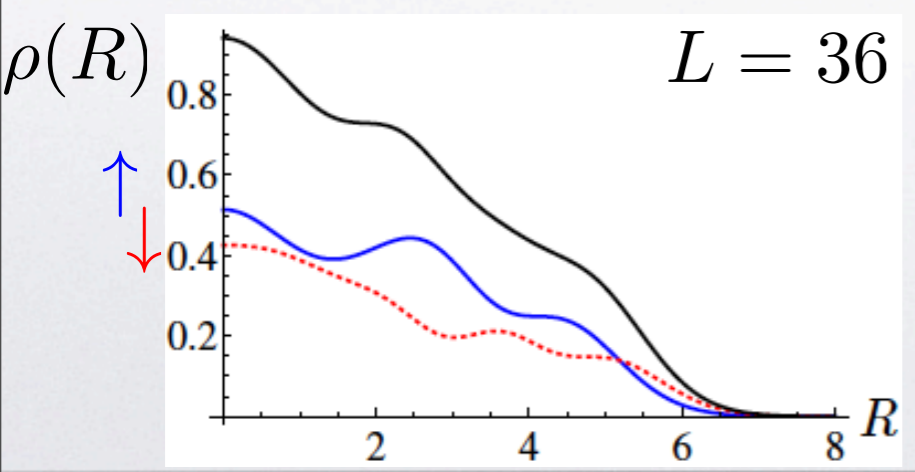


- Laughlin ansatz $\nu = 1/5$ $\Psi_{1/5} = \prod_{i < j} (z_i - z_j)^5 \equiv \Theta^5$

2nd DLL: three regimes



- skyrmionic spin texture around vortex cores

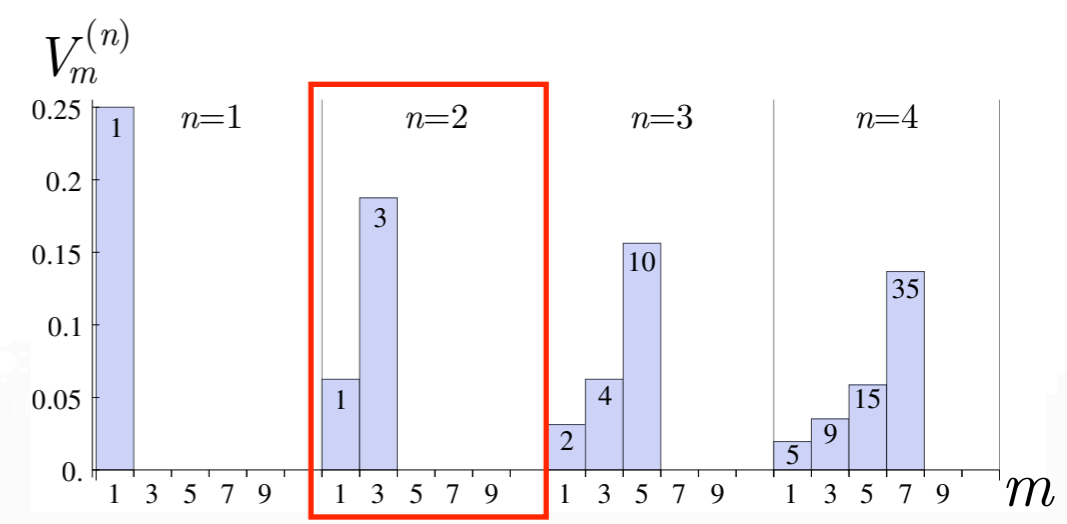
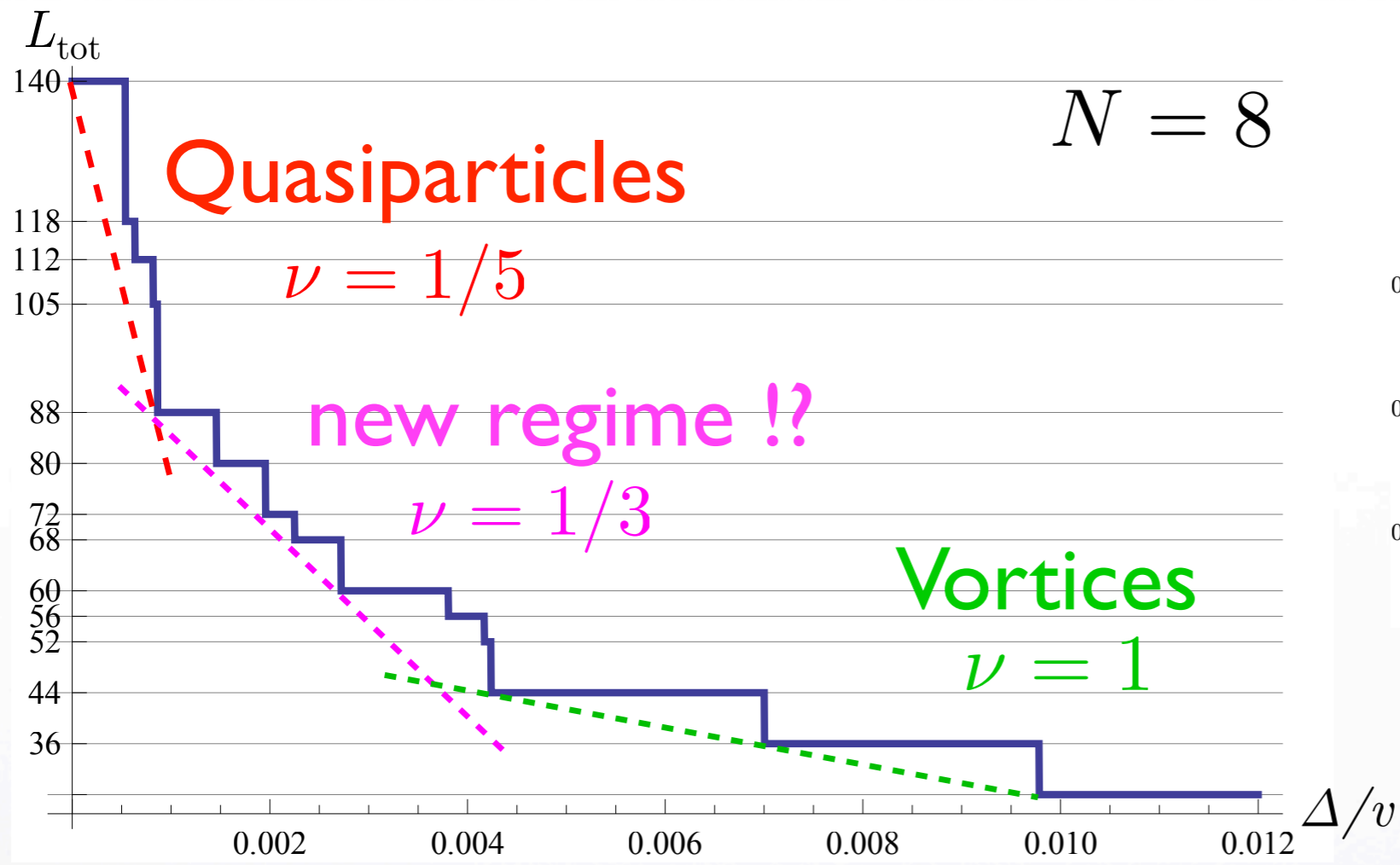


Strongly correlated states of trapped ultracold fermions in a U(2) gauge potential

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2nd DLL: three regimes



• NO Laughlin

~~$\nu = 1/3 \quad \Psi_{1/3} = \prod_{i < j} (z_i - z_j)^3 \equiv \Theta^3$~~

• Haffnian !?

$L_{\text{Hf}} = 76$

$\nu = 1/3 \quad \Psi_{\text{Hf}} \sim \mathcal{S} \left(\frac{1}{(z_1 - z_2)^2 \dots (z_{N-1} - z_N)^2} \right) \prod_{i,j} (z_i - z_j)^3$

X.-G. Wen, Nucl. Phys. B **419**, 455 (1994) // D. Green, arXiv:cond-mat/0202455

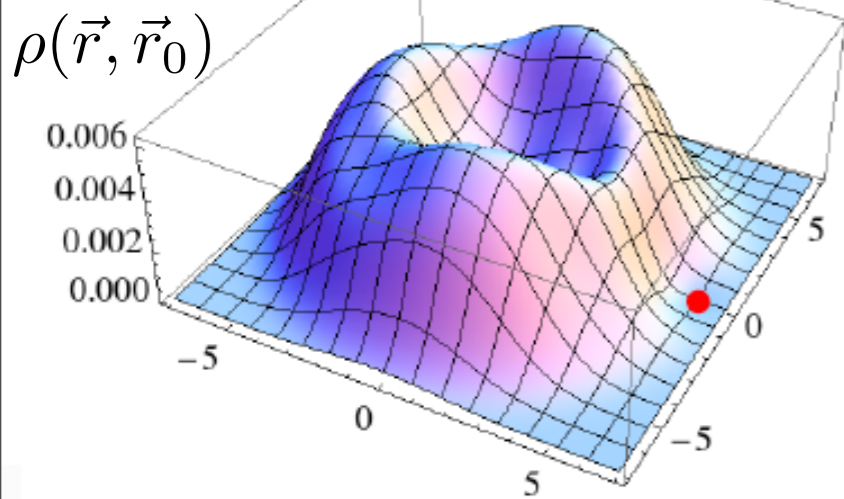
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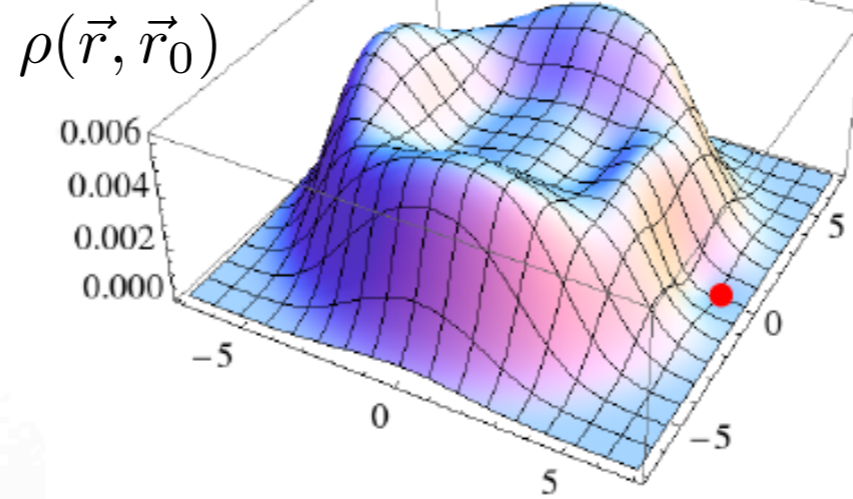
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2nd DLL: Haffnian !?

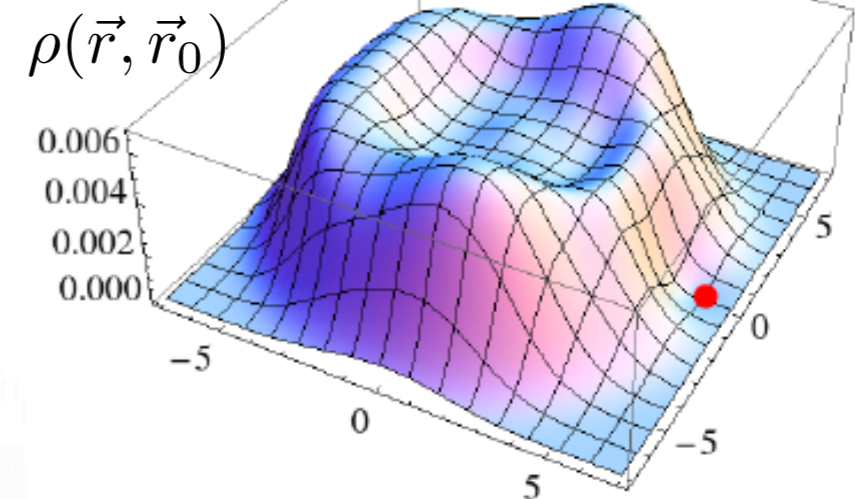
$L = 72$



$L = 76$



$L = 80$



• **Haffnian !?**

$L_{\text{Hf}} = 76$

$$\nu = 1/3 \quad \Psi_{\text{Hf}} \sim \mathcal{S} \left(\frac{1}{(z_1 - z_2)^2 \dots (z_{N-1} - z_N)^2} \right) \prod_{i,j} (z_i - z_j)^3$$

X.-G. Wen, Nucl. Phys. B **419**, 455 (1994) // D. Green, arXiv:cond-mat/0202455

- effective (d-wave) pairing $\Leftrightarrow N/2 = 4$ peaks
- stabilized states (by Δ) are rather $L_{\text{Hf}} \pm N/2$
- three-body interactions needed ...

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Entanglement spectrum: intro

- quantum information approach: (robust & intrinsic)

$$S = A \cup \bar{A} \quad \rho_A = \text{Tr}_{\bar{A}} |\psi\rangle\langle\psi| = \sum_l \lambda_l |l\rangle\langle l| \quad \text{ES}_l \equiv -\ln \lambda_l$$

H. Li and F.D.M. Haldane, PRL **101**, 010504 (2008)

- orbital partition \sim real space partition $\rho_A^{(O)} = \left(\sum_{\{\vec{n}''\}_{\bar{A}}} \Psi_{\vec{n} \otimes \vec{n}''}^* \Psi_{\vec{n}' \otimes \vec{n}''} \right) |\vec{n}\rangle_A \langle \vec{n}'|$

$\text{ES}^{(O)} \sim$ info. on edge excitations

X.-L. Qi, H. Katsura, A.W. Ludwig, PRL **108**, 196402 (2012)

- particle partition irrespective of position $\rho_A^{(P)} = \left(\langle\psi| \prod_{j \leq N_A} a_{m_j}^\dagger \prod_{k \leq N_A} a_{m'_k} |\psi\rangle \right) |\vec{m}\rangle \langle \vec{m}'|$

$\text{ES}^{(P)} \sim$ bulk & excitation properties

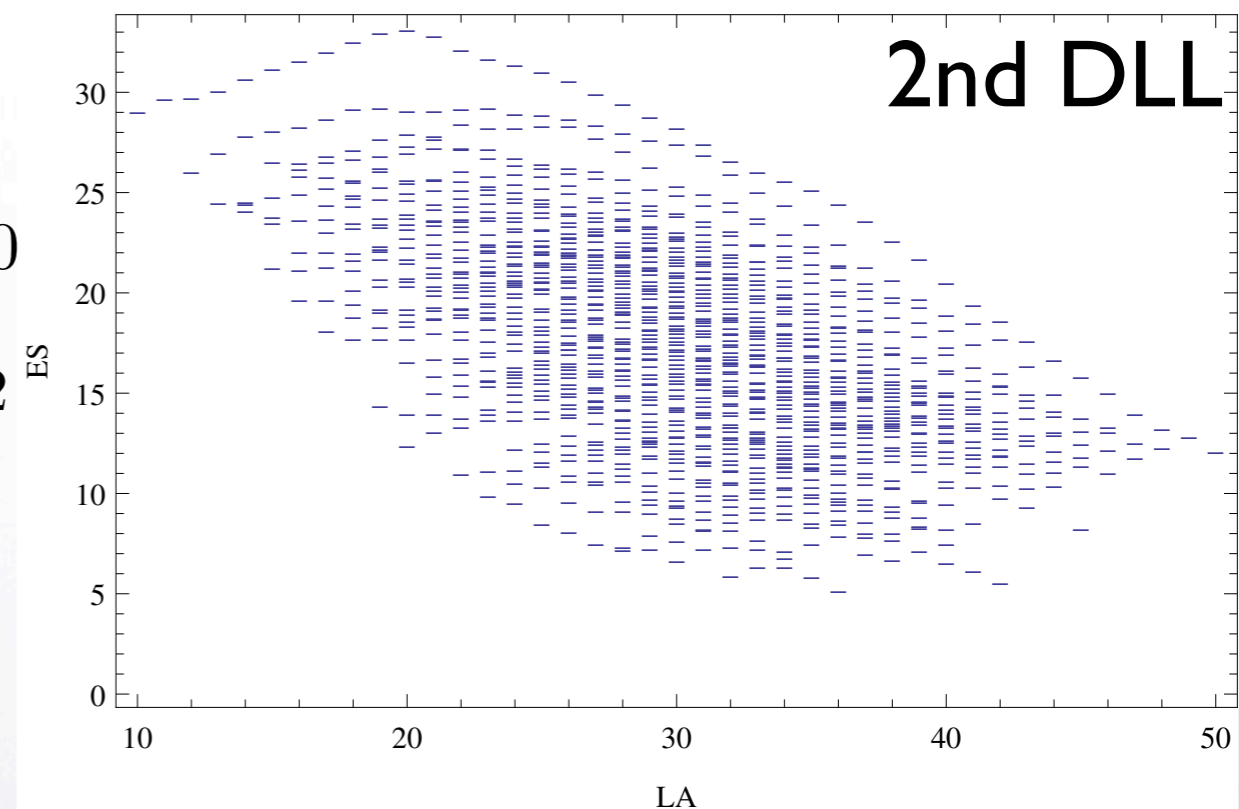
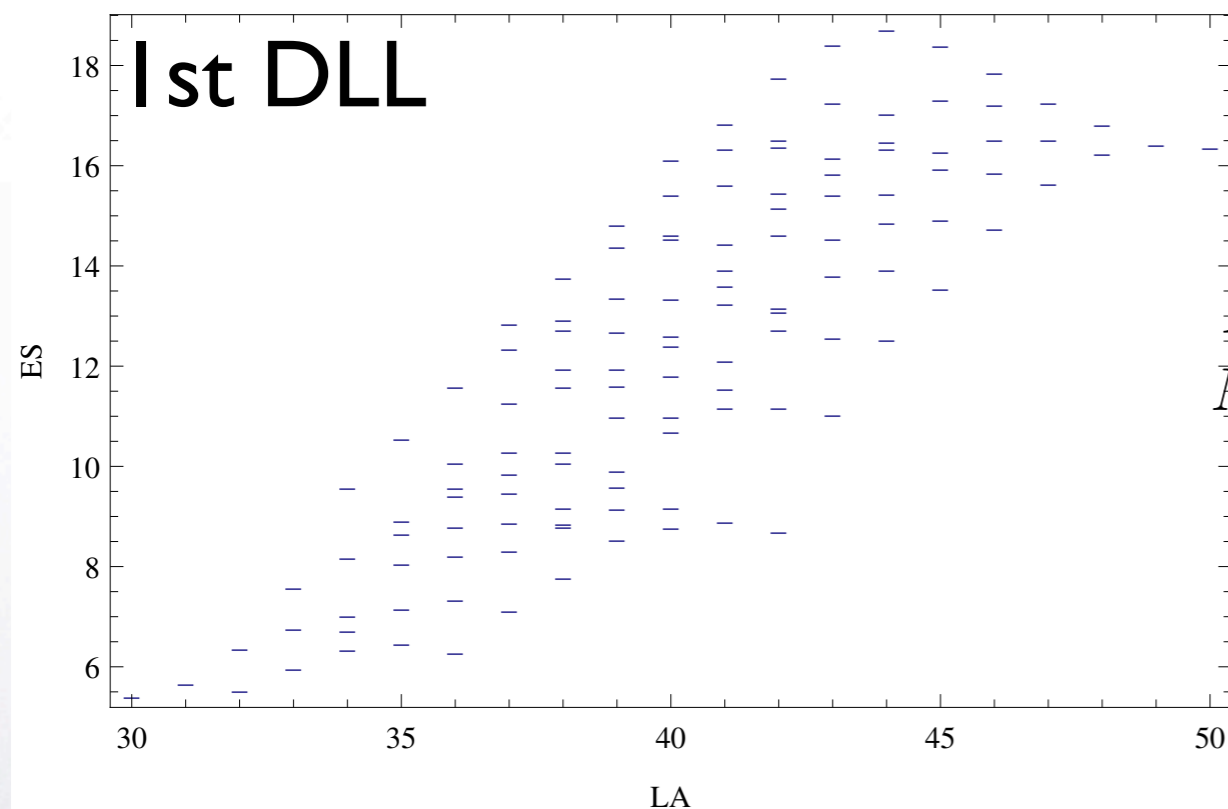
A. Sterdyniak, N. Regnault and B. A. Bernevig, PRL **106**, 100405 (2011)

ES: Laughlin 1/3

- orbital partition \sim real space partition $\rho_A^{(O)} = \left(\sum_{\{\vec{n}''\}_{\bar{A}}} \Psi_{\vec{n} \otimes \vec{n}''}^* \Psi_{\vec{n}' \otimes \vec{n}''} \right) |\vec{n}\rangle_A \langle \vec{n}'|$

$$N_A = \sum_{m \leq M_A} n_m \quad L_A = \sum_{m \leq M_A} m n_m$$

$$L_{1/3} = 3N(N - 1) = 135$$



- counting Laughlin edge modes
1, 1, 2, 3, 5, 7, 11, ...

- no clear ES gap ...

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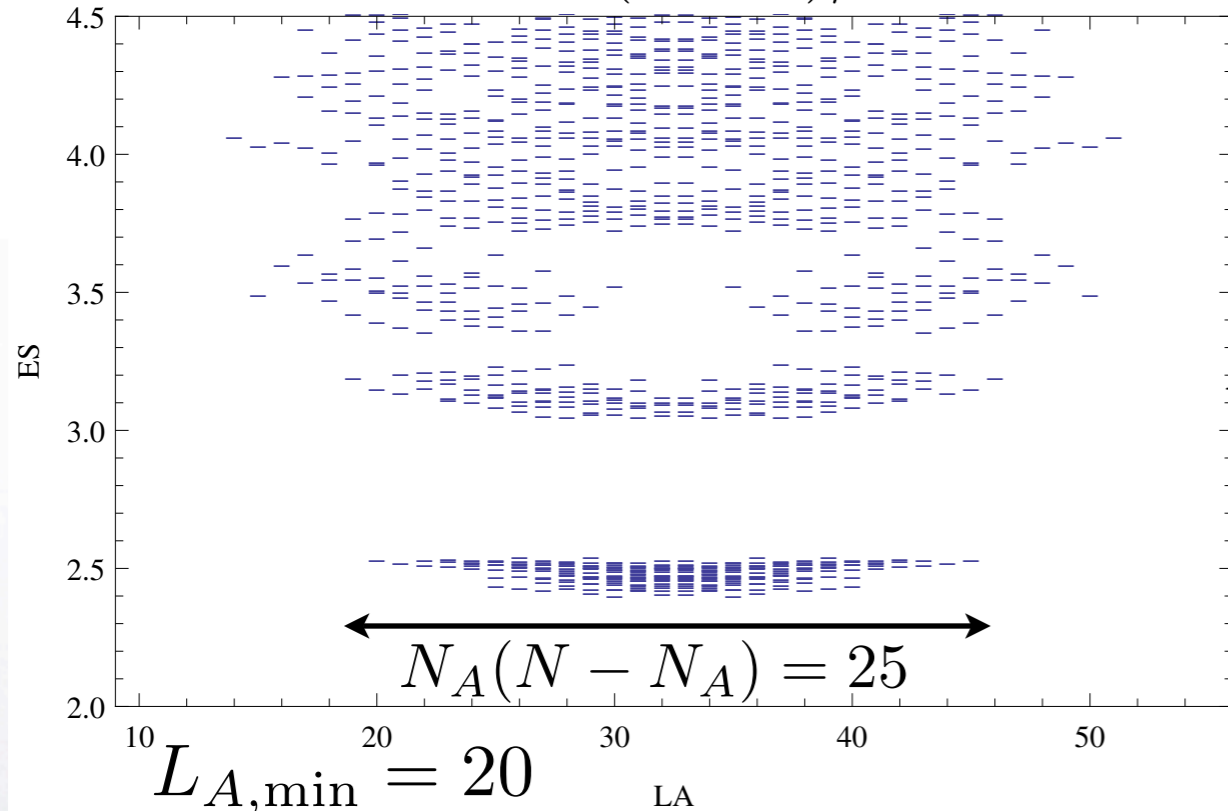
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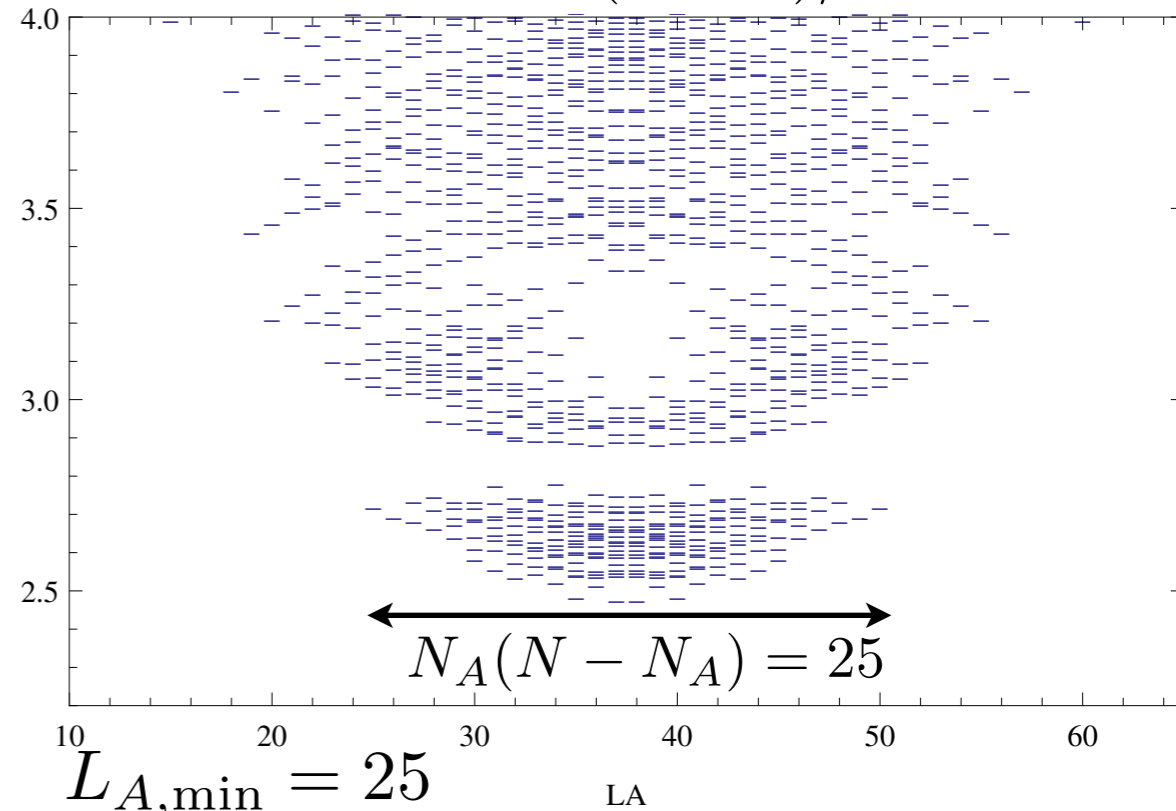
ES: vortices

- particle partition irrespective of position $\rho_A^{(P)} = \left(\langle \psi | \prod_{j \leq N_A} a_{m_j}^\dagger \prod_{k \leq N_A} a_{m'_k} | \psi \rangle \right) |\vec{m}\rangle \langle \vec{m}'|$

$$L = 65 = N(N - 1)/2 + 2N$$



$$L = 75 = N(N - 1)/2 + 3N$$



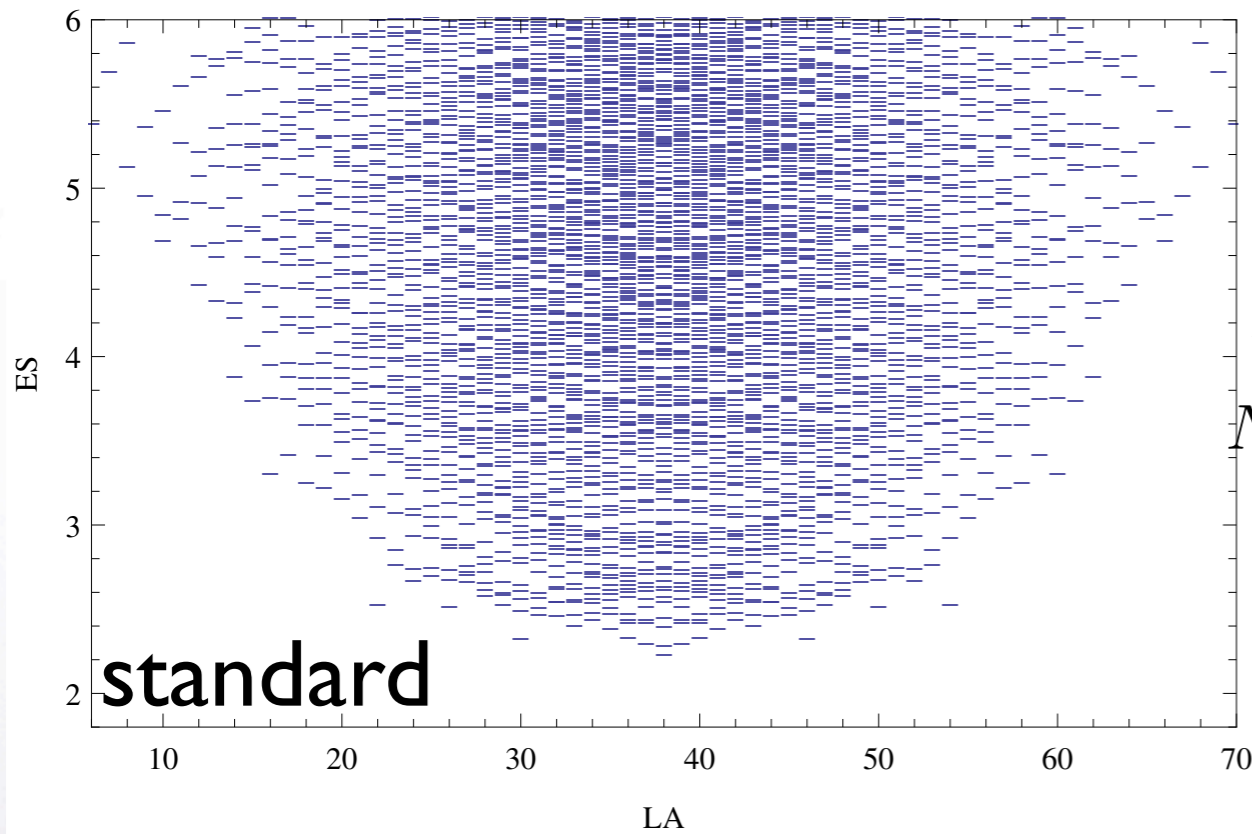
- consistent with central vortex ansatz

$$\Psi_{Vp} = \prod_i z_i^p \prod_{i < j} (z_i - z_j) \longrightarrow L_{A,\min} = N_A(N_A - 1)/2 + pN_A$$

ES: Haffnian candidate

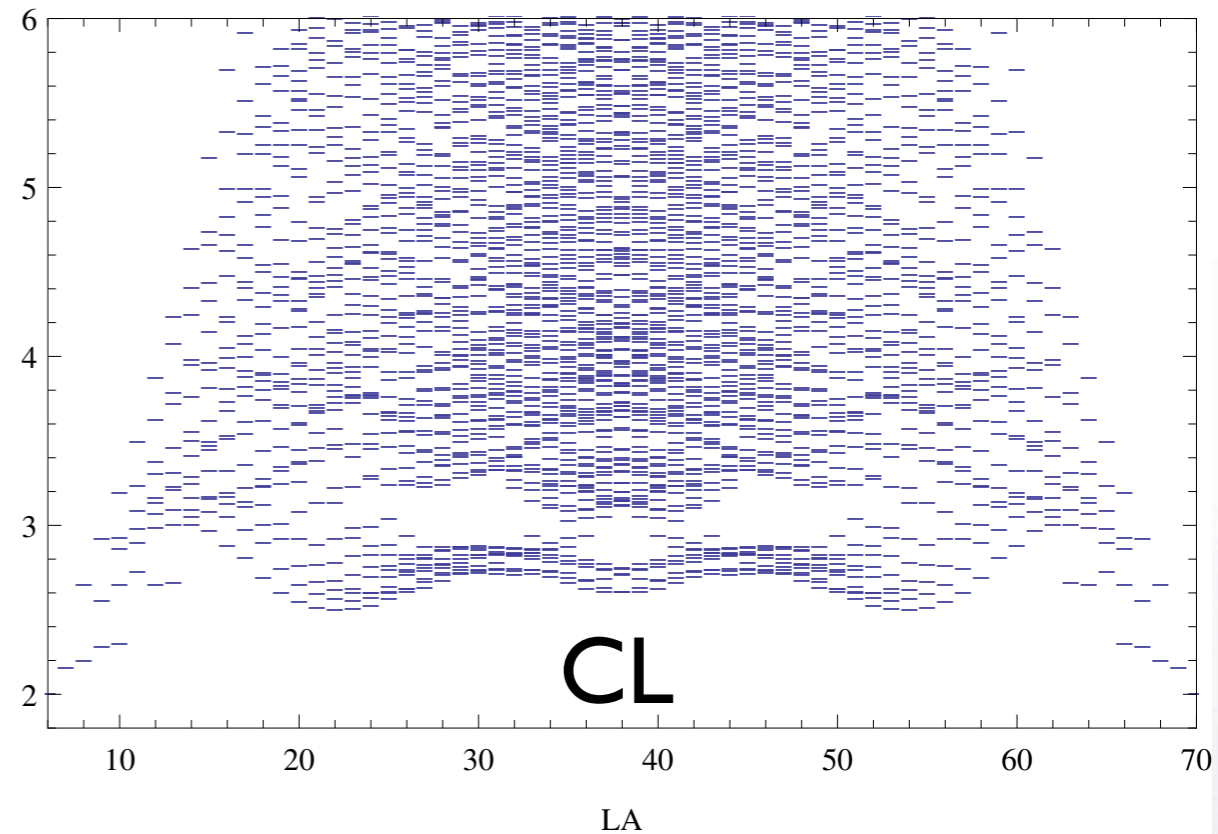
- particle partition irrespective of position $\rho_A^{(P)} = \left(\langle \psi | \prod_{j \leq N_A} a_{m_j}^\dagger \prod_{k \leq N_A} a_{m'_k} | \psi \rangle \right) |\vec{m}\rangle \langle \vec{m}'|$

$$L = 3N(N - 1) - N = 76$$



$$N = 8$$

$$N_A = 4$$



- conformal limit to cancel geom. factors $\psi_{\vec{m}}^{(CL)} \equiv \psi_{\vec{m}} \cdot \prod_{j \leq N} \sqrt{m_j!}$

R. Thomale, A Sterdyniak, N. Regnault, B.A. Bernevig, PRL **104**, 180502 (2010)

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Conclusions & Outlook

- U(2) potential:
 - i) deformed LL (spin textures...)
 - ii) non-monotonic HP from s-wave only
- Novel incompressible states: Haffnian? d-wave pairing?
- Entanglement spectrum -- theoretical detector
- other LL deforming potentials? absence of Zeeman comp.?
- degeneracy points between LL's ?
- three-body terms? dissipation-induced?

*M.Roncaglia, MR, J.I. Cirac,
PRL 104, 096803 (2010)*

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M. Roncaglia
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A. Trombettoni
SISSA-Trieste

Burrello, Rizzi, Roncaglia, Trombettoni, PRB 91, 115117 (2015)



JOHANNES GUTENBERG
UNIVERSITÄT MAINZ



all of you
for attention !

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