

# Phases of strongly-interacting bosons on a two-leg ladder

Marie Piraud

Arnold Sommerfeld Center for Theoretical Physics, LMU, Munich

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# My collaborators

## LMU, Munich



F. Heidrich-Meisner



U. Schollwöck



I. P. McCulloch (Brisbane)

## Leibniz U., Hannover



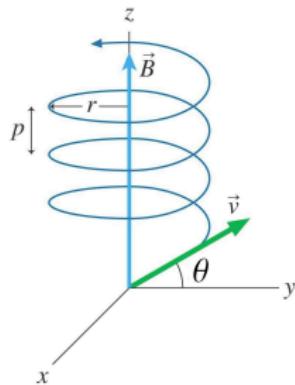
S. Greschner



T. Vekua

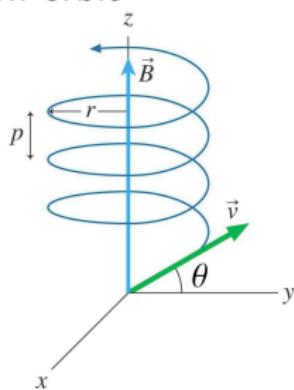
# Electron in a magnetic field

Cyclotron orbit



# Electron in a magnetic field

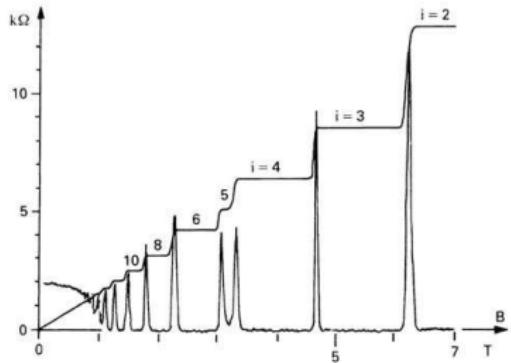
Cyclotron orbit



Landau levels (highly degenerate)

$$E_n = \hbar\omega_c(n + 1/2)$$

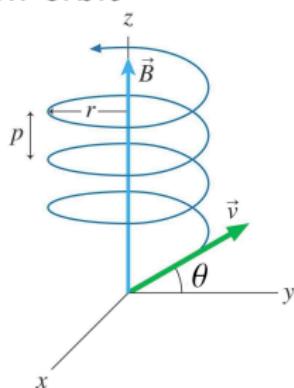
Quantum Hall Effect :  
quantization of conductance



K. von Klitzing, RMP (1986)

# Electron in a magnetic field

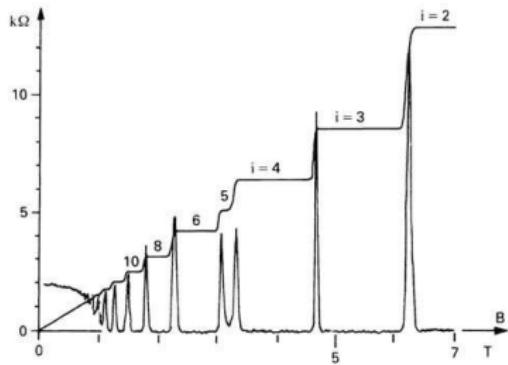
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Landau levels (highly degenerate)

$$E_n = \hbar\omega_c(n + 1/2)$$

Quantum Hall Effect :  
quantization of conductance



K. von Klitzing, RMP (1986)

Modifies transport properties & creates Topological transitions

With interactions : Fractional QHE

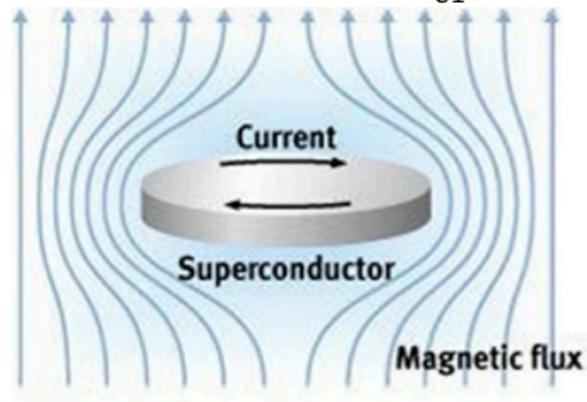
⇒ Unusual excitations, fractional charge and statistics

# Type-II Superconductor in a magnetic field

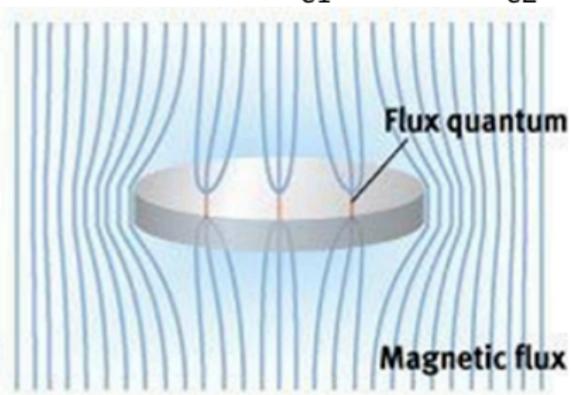
## Meissner effect

W. Meissner & R. Ochsenfeld, Naturwissenschaften (1933)

Meissner state  $H < H_{c1}$



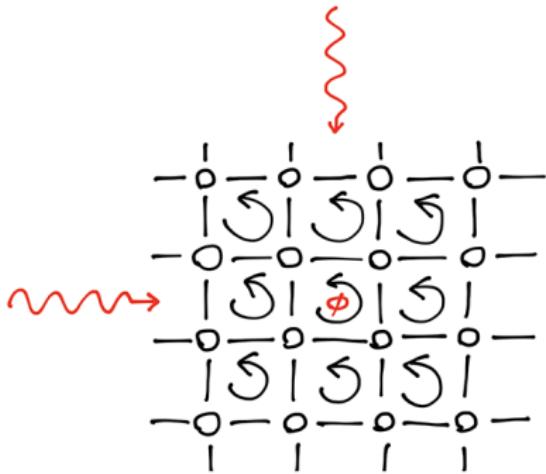
Vortex lattice  $H_{c1} < H < H_{c2}$



Source : Quantumlevitation.com

# Ultra-cold atoms in synthetic gauge fields

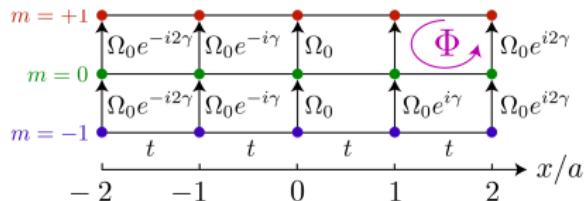
Engineer effective  $\vec{B}$  field or SO coupling on lattices :



C. Chin & E.J. Mueller, PRL Viewpoint (2013)

- Interplay interactions/ gauge fields/ dimensionality  
⇒ Promising candidate to search for new phases of matter.
- UCA permit to study **bosons** in magnetic fields
- Ladders = simplest "2D" system

In 'Synthetic dimensions'



A. Celi et al., PRL (2014)

Realizations :

M. Mancini et al., arXiv1502 :02495

B. Stuhl et al., arXiv1502 :02496

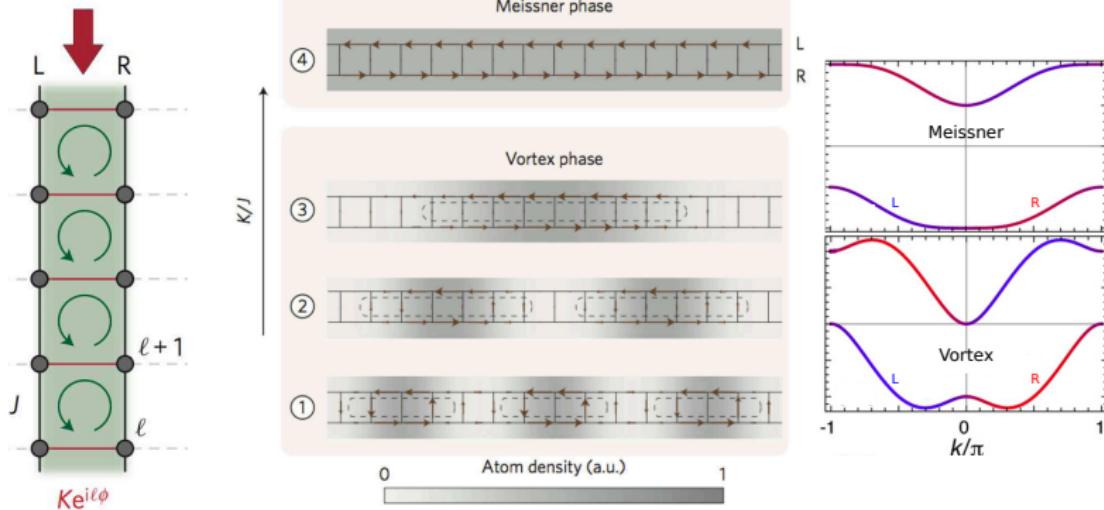
J. Dalibard & F. Gerbier, RMP (2011)

- Phases of strongly-interacting bosons on a two-leg ladder

# Weakly-interacting bosons on ladders

## Vortex/ Meissner Phase Transition (ladders with uniform $\pi/2$ magnetic flux)

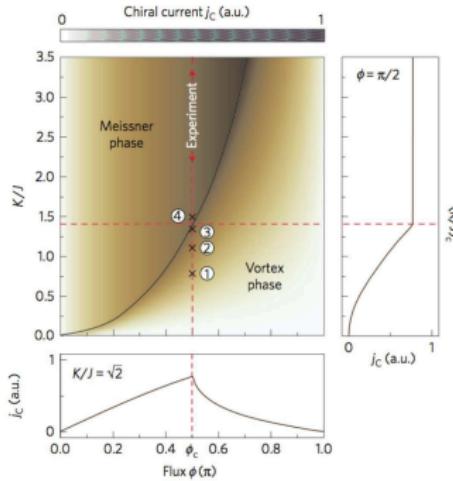
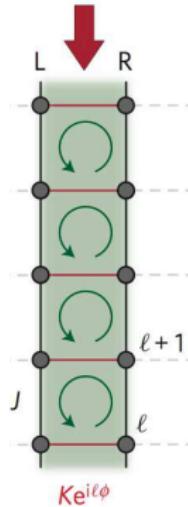
Theo : E. Orignac & T. Giamarchi, PRB (2001) , D. Hügel & B. Paredes, PRA (2014)  
Exp : M. Atala et al., Nat. Phys. (2014)



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## Interacting bosons on ladders in uniform flux

## A few theoretical results

- Field-theoretical studies E. Orignac & T. Giamarchi, PRB (2001)  
A. Tokuno & A. Georges, NJP (2014)
  - Mean-field approach R. Wei & E. J. Mueller, PRA (2014)
  - Chiral Mott Insulator A. Dhar *et al.*, PRA (2012) ; PRB (2013)
  - Mott Insulator with Meissner Currents A. Petrescu & K. Le Hur, PRL (2013)
  - Study of the momentum distribution M.-C. Cha & J.-G. Shin, PRA (2011)
  - Laughlin-1/2 states A. Petrescu & K. Le Hur, PRB (2015)

# Interacting bosons on ladders in uniform flux

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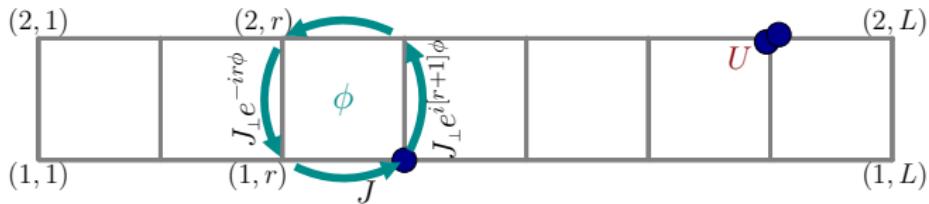
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⇒ Explore the phase diagram numerically

# Interacting bosons on a two-leg ladder in a uniform flux : Model

$$\mathcal{H}_K = -J \sum_{\ell=1,2; r=1}^L \left( a_{\ell,r+1}^\dagger a_{\ell,r} + \text{H.c.} \right) - J_\perp \sum_{r=1}^L \left( e^{-ir\phi} a_{1,r}^\dagger a_{2,r} + \text{H.c.} \right)$$

$$\mathcal{H}_U = \frac{U}{2} \sum_{\ell=1,2; r=1}^L n_{\ell,r} (n_{\ell,r} - 1)$$



# Density Matrix Renormalization Group

Ground state and first excited states of 1D many-body systems

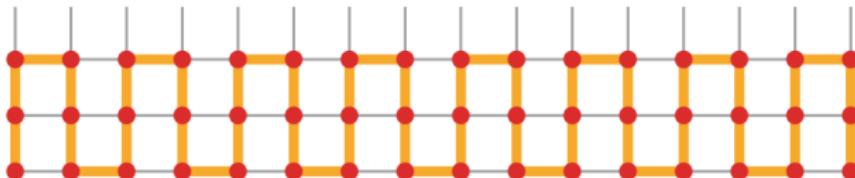
S. R. White, Phys. Rev. Lett. **69**, 2863 (1992)

U. Schollwöck, Rev. Mod. Phys. **77**, 259 (2005)

Successfully applied to (quasi-)2D systems

S. Yan *et al.*, Science (2011)

S. Depenbrock *et al.*, PRL (2012)



E. M. Stoudenmire *et al.*, Annu. Rev. Cond. Mat. Phys. (2012)

# Outline

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- Hard-Core Bosons

M. P., F. Heidrich-Meisner, I. P. McCulloch, S. Greschner, T. Vekua and U. Schollwöck, PRB **91** 140406(R) (2015)

- Phase Diagram : MI and SF, Meissner and vortex phases
- A few observables : chiral current and vortex density

- Finite interactions

S. Greschner, M. P., F. Heidrich-Meisner, I. P. McCulloch, U. Schollwöck and T. Vekua, *submitted* (2015)

- Phase Diagram : vortex lattice phases
- Chiral current : spontaneous reversal !

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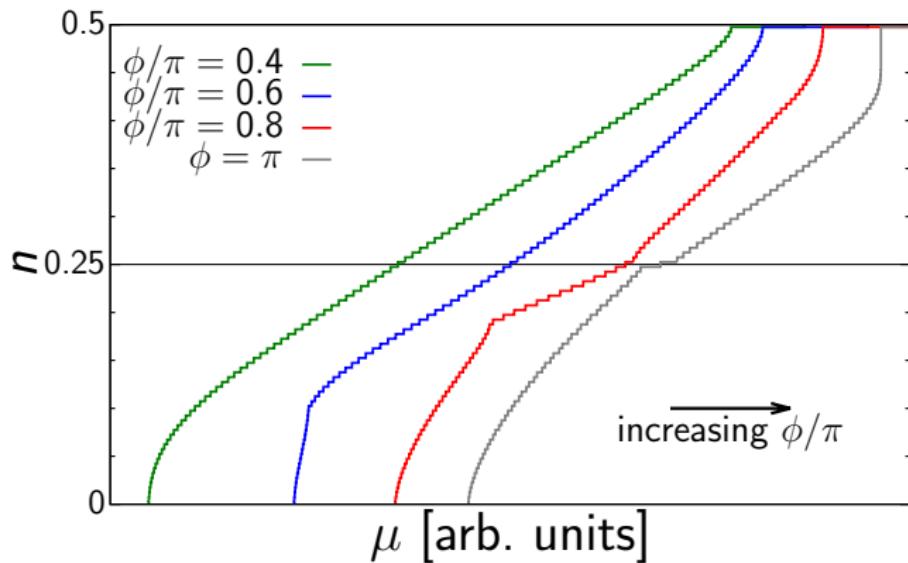
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- Phase Diagram
- Chiral current

# Equation of state

Filling  $n = N/2L$  vs. chemical potential  $\mu$



Plateau (at  $n = 0.5$ , and at  $n = 0.25$  for  $\phi = \pi$ )

⇒ incompressible state

Kink ⇒ change in the number of gapless excitations (+ use  $S_{vN}$ )

# Fractionally filled Mott insulator

## Hard-core bosons, no flux

- Strong rungs ( $J_{\perp} \gg J$ ) :  $\mathcal{H}_{\mathbf{k}} \sim -J_{\perp} \sum_r (a_{1,r}^{\dagger} a_{2,r} + \text{H.c.})$

$s_z = 0$  triplet

$$(|1,0\rangle_r + |0,1\rangle_r)/\sqrt{2}$$

$$E = -J_{\perp}$$

$s_z = \pm 1$  triplets

$$|1,1\rangle_r \text{ and } |0,0\rangle_r$$

$$E = 0$$

singlet

$$(|1,0\rangle_r - |0,1\rangle_r)/\sqrt{2}$$

$$E = J_{\perp}$$

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singlet

$$(|1,0\rangle_r - |0,1\rangle_r)/\sqrt{2}$$

$$E = J_{\perp}$$

1 boson per rung  $\Rightarrow$  'rung-triplet' phase

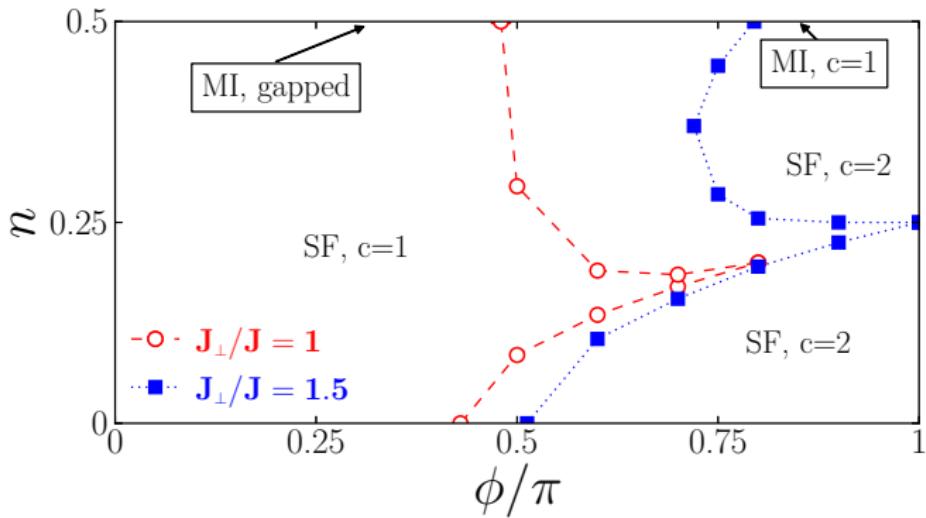


$\Rightarrow$  Mott insulator extends down to  $J_{\perp}/J = 0$  for HCBs

T. Vekua *et al.*, PRB (2003); F. Crépin *et al.*, PRB (2011).

# Phase diagram for Hard-Core Bosons

Parameter space : rung-tunneling  $J_{\perp}/J$ , flux  $\phi$ , filling  $n = N/2L$ .

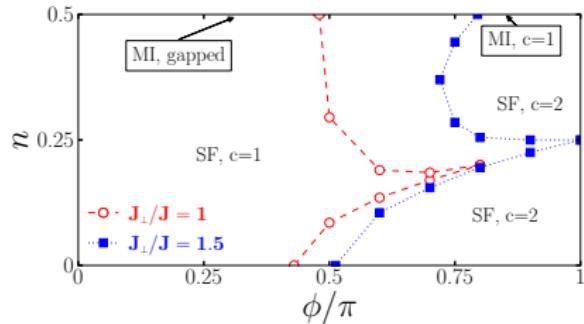


Upper half  $0.5 < n < 1$  related by particle-hole symmetry.  
Phase diagram consistent with field-theory.

# Current patterns

**Local currents**

$$j_{\ell,r}^{\parallel} = iJ \left( a_{\ell+1,r}^\dagger a_{\ell,r} - \text{H.c.} \right) \quad \text{and} \quad j_r^{\perp} = iJ_{\perp} \left( e^{-ir\phi} a_{1,r}^\dagger a_{2,r} - \text{H.c.} \right)$$



## Meissner-SF

$\phi/\pi = 0.5, J_{\perp}/J = 1.5, n=0.25$



## Vortex-SF

$\phi/\pi = 0.9, J_{\perp}/J = 1.5, n=0.1$



cf E. Orignac & T. Giamarchi, PRB (2001) , A. Tokuno & A. Georges, NJP (2014) ,  
M. Atala *et al.*, Nat. Phys. (2014)

# Current patterns

## Meissner-MI

$\phi/\pi = 0.5, J_{\perp}/J = 2, n = 0.5$

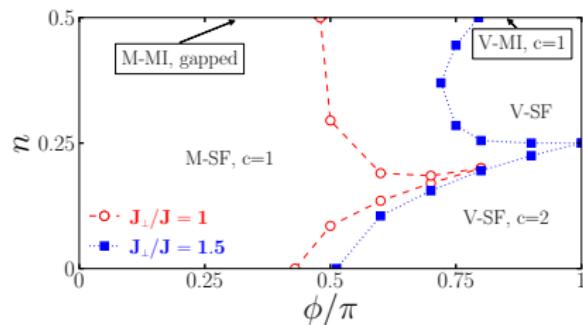


## Vortex-MI

$\phi/\pi = 0.5, J_{\perp}/J = 0.5, n = 0.5$



cf A. Petrescu & K. Le Hur, PRL (2013)



## Meissner-SF

$\phi/\pi = 0.5, J_{\perp}/J = 1.5, n=0.25$



## Vortex-SF

$\phi/\pi = 0.9, J_{\perp}/J = 1.5, n=0.1$



## Tractable limit : around $\phi \sim \pi$

---

For strong rung-tunneling  $J_{\perp} \gg J$  and  $n < 0.5$  : effective spin  $\mathcal{H}$

Pseudo-spin-1/2 on each rung  $r$  :

$$(|1,0\rangle_r + e^{ir\phi} |0,1\rangle_r)/\sqrt{2} \rightarrow |\downarrow\rangle_r \quad \text{and} \quad |0,0\rangle_r \rightarrow |\uparrow\rangle_r.$$

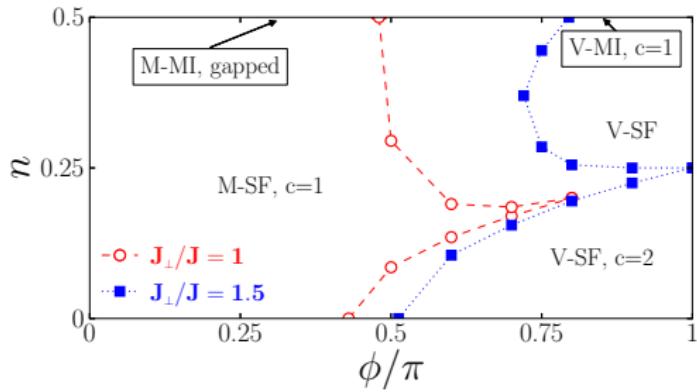
$$\mathcal{H}_{\text{eff}}(\phi = \pi) = \frac{J^2}{2|J_{\perp}|} \sum_r \left( 2S_r^z S_{r+1}^z - \left[ S_r^+ \left( \frac{1}{2} - S_{r+1}^z \right) S_{r+2}^- + \text{H.c.} \right] \right)$$

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- $(n = 0.25, \phi = \pi)$  becomes gapped  
⇒ V-SF splitted in 2 lobes (for  $J_{\perp} \gtrsim 1.3J$ )

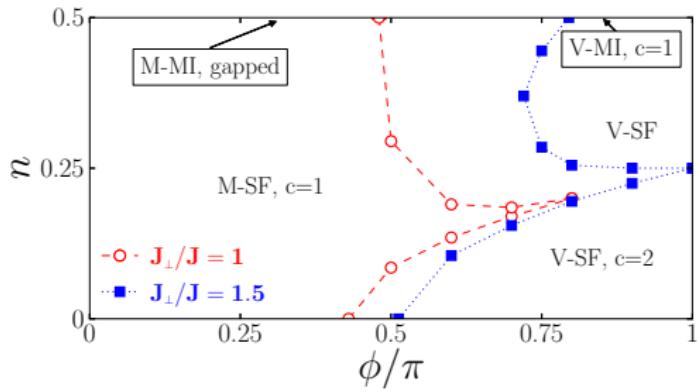


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- $(n = 0.25, \phi = \pi)$  becomes gapped  
⇒ V-SF splitted in 2 lobes (for  $J_{\perp} \gtrsim 1.3J$ )
- V-MI and upper lobe of V-SF disappears (for  $J_{\perp} \gtrsim 1.7J$ )



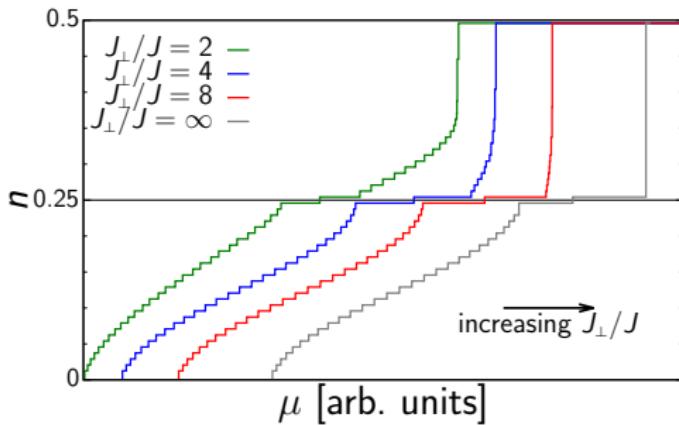
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## • Meta-magnetic transition at $\phi = \pi$ :

jump in density from  $n > 0.25$  to  $n = 0.5$  (for  $J_{\perp} \gtrsim 1.5J$ )



# Outline

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M. P., F. Heidrich-Meisner, I. P. McCulloch, S. Greschner, T. Vekua and  
U. Schollwöck, PRB **91** 140406(R) (2015)

- Phase Diagram
- A few observables

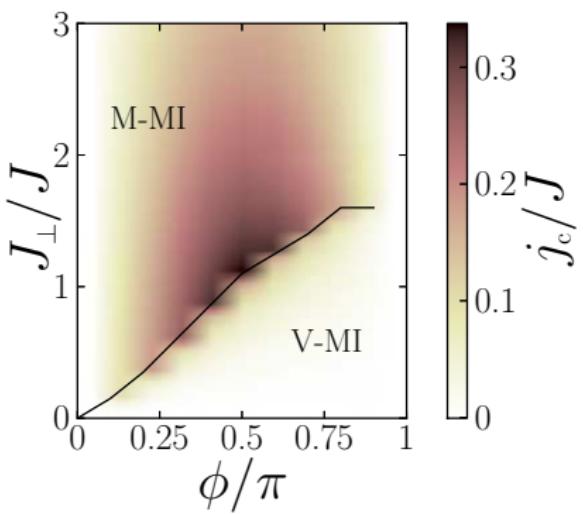
- Finite interactions

S. Greschner, M. P., F. Heidrich-Meisner, I. P. McCulloch, U. Schollwöck and  
T. Vekua, *submitted* (2015)

- Phase Diagram
- Chiral current

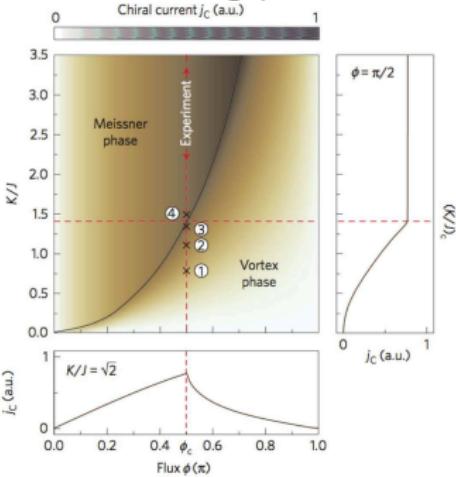
# Chiral current, HCB at $n = 0.5$

$$j_c = \frac{1}{2L} \left| \sum_r \langle j_{1,r}^{\parallel} - j_{2,r}^{\parallel} \rangle \right|$$



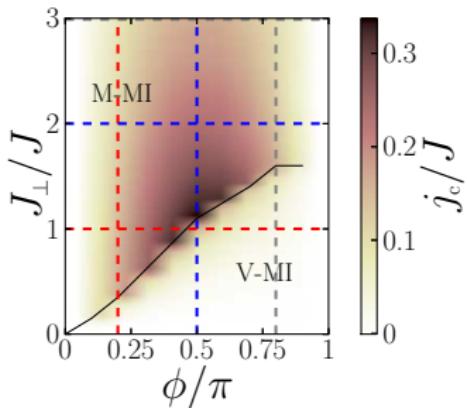
Vortex phase is disfavored

cf non-interacting phases :



M. Atala et al., Nat. Phys. (2014)

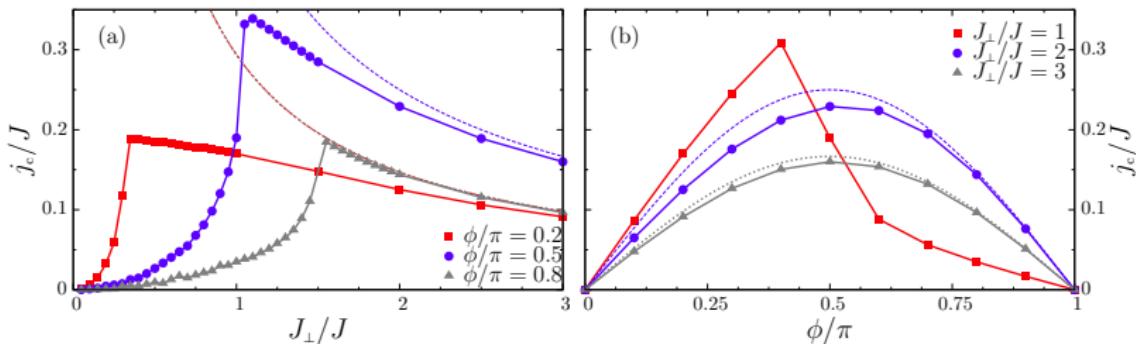
# Chiral current, HCB at $n = 0.5$



At  $J_{\perp} \ll J$  and small  $\phi$  :  
 $j_c \sim (J_{\perp}/\phi)^2$

At  $J_{\perp} \gg J$  :  
 $j_c = \frac{J^2}{J_{\perp}} \sin(\phi)$

no saturation for HCB in gapped phases



# Vortex density, HCB at $n = 0.5$

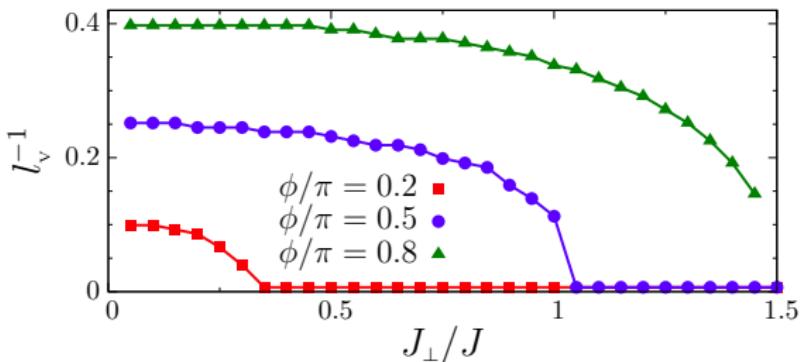
(a1)  $\phi/\pi = 0.5, J_{\perp}/J = 0.05$



(a2)  $\phi/\pi = 0.5, J_{\perp}/J = 0.5$



(a3)  $\phi/\pi = 0.5, J_{\perp}/J = 2$



Inverse spatial extension of the vortices :  $\rho_V = l_V^{-1}$  (obtained from the Fourier transform of the rung current  $\langle j_r^{\perp} \rangle$ ) .

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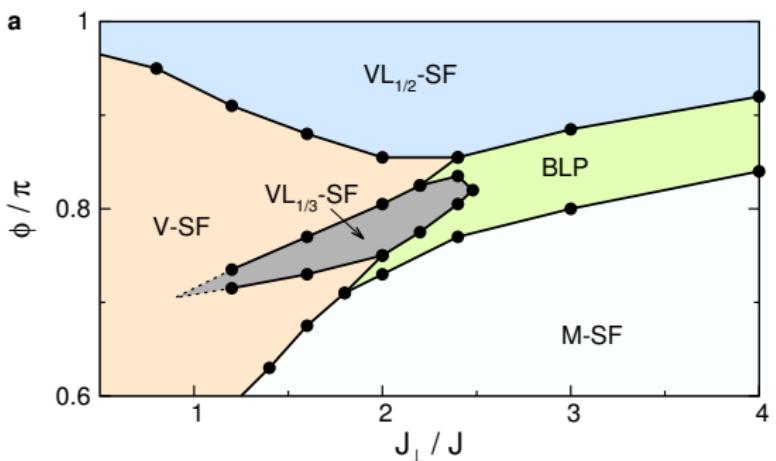
- Phase Diagram
- A few observables

- Finite interactions

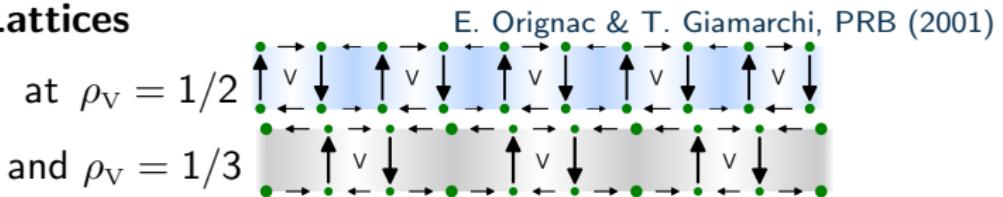
S. Greschner, M. P., F. Heidrich-Meisner, I. P. McCulloch, U. Schollwöck and  
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- Chiral current

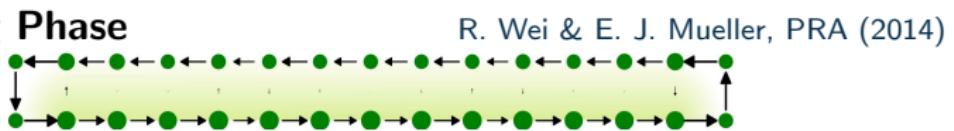
# Phase Diagram, at $U = 2J$ and $n = 0.8$



## Vortex Lattices



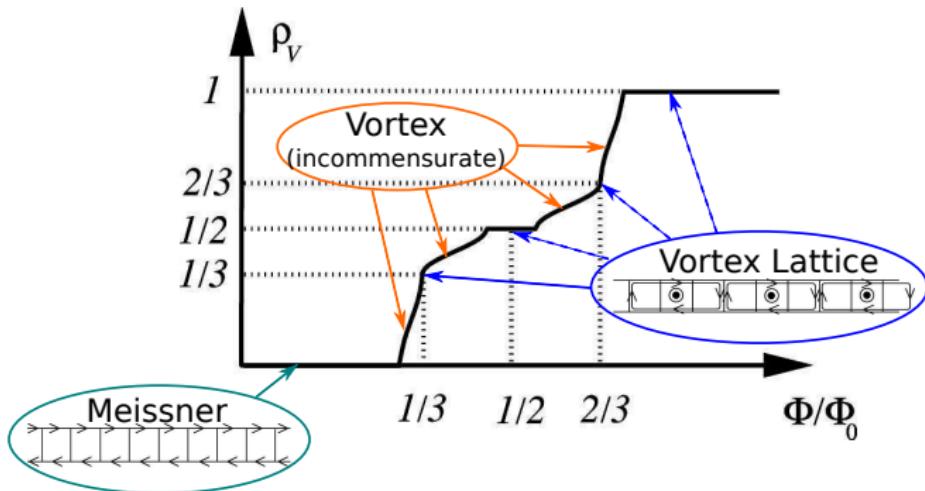
## Biased Leg Phase



# Vortex Lattices

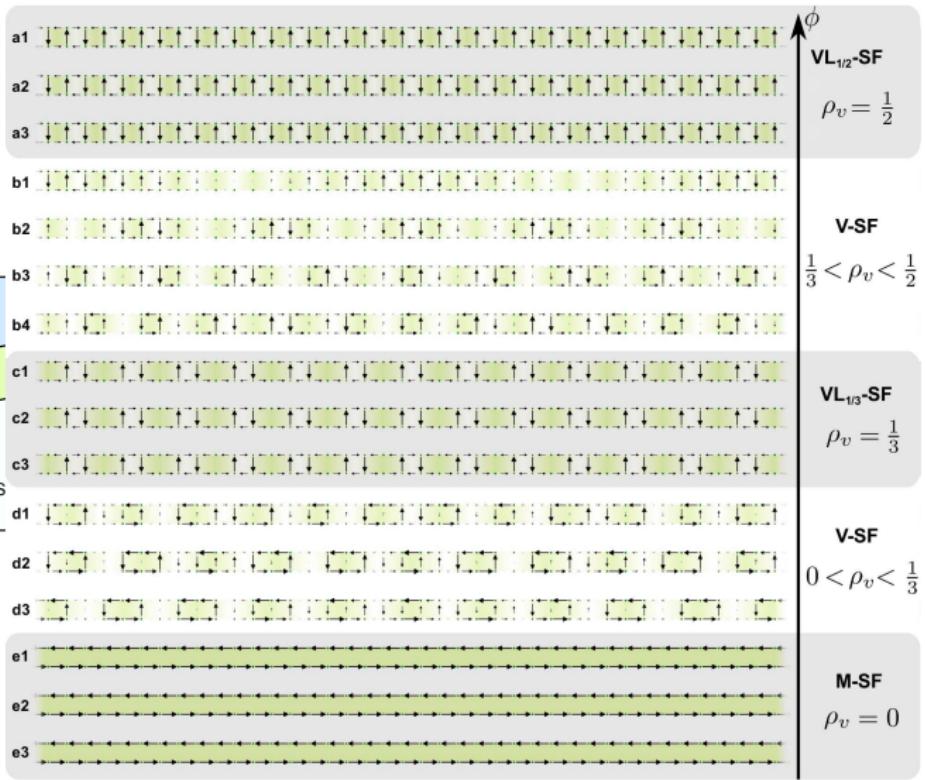
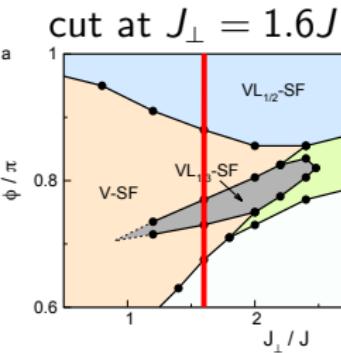
## Commensurate/ incommensurate transitions

E. Orignac & T. Giamarchi, PRB (2001)

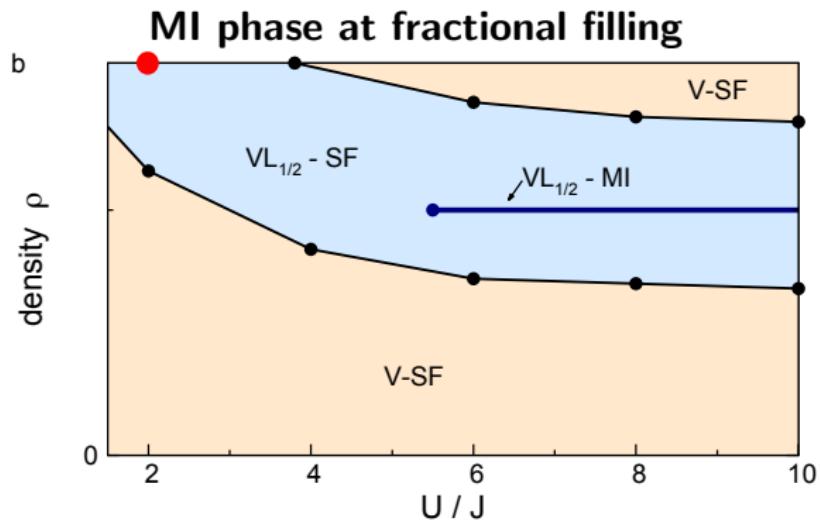
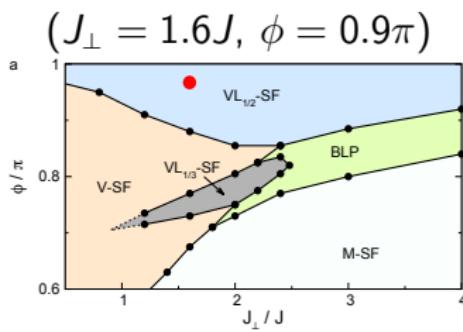


The size of the plateaux reduces with interactions

# Phase Diagram, at $U = 2J$ and $n = 0.8$



# Phase Diagram : Mott Insulating phases



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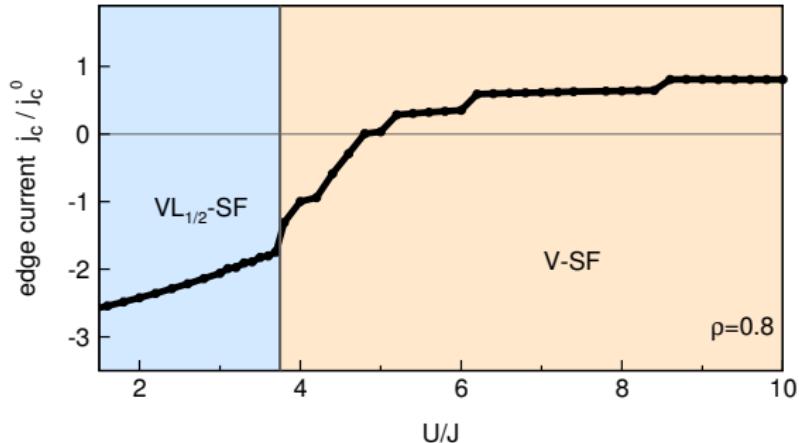
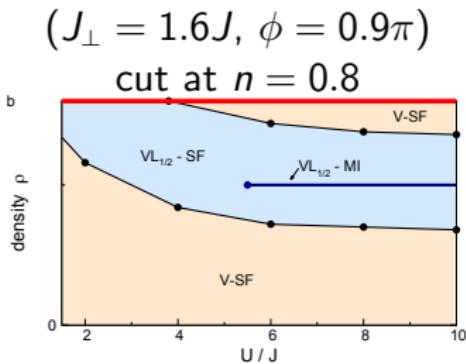
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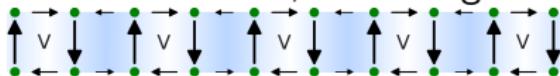
S. Greschner, M. P., F. Heidrich-Meisner, I. P. McCulloch, U. Schollwöck and  
T. Vekua, *submitted* (2015)

- Phase Diagram
- **Chiral current**

# Chiral current : Swimming against the tide !

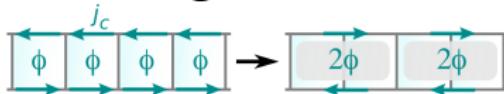


In and near the  $VL_{1/2}$ , the current circulates in the clock-wise direction, on average :

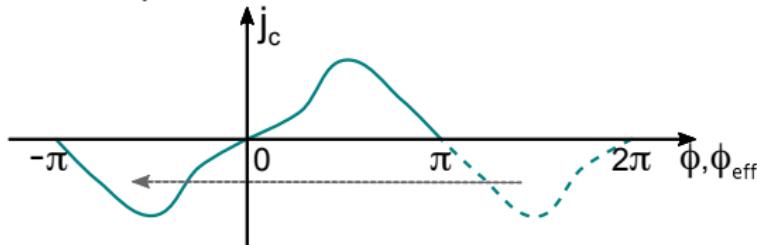


# Chiral current : Swimming against the tide !

- In the VL<sub>1/2</sub> phase : spontaneous symmetry breaking  
⇒ **doubling of the unit cell**



- $j_c$  is odd and  $2\pi$ -periodic with the flux



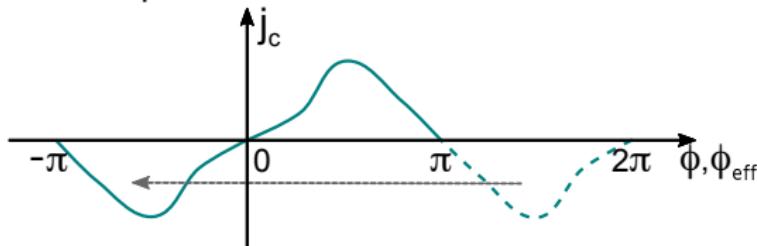
- If we have  $\odot \mathcal{B}$  on  $\square$  with  $\phi \in [\pi/2, \pi]$ ;  
the effective flux on  $\square\square$  ( $2\phi$ ) corresponds to  $\otimes \mathcal{B}$ .  
⇒ **spontaneous current reversal**

# Chiral current : Swimming against the tide !

- In the VL<sub>1/2</sub> phase : spontaneous symmetry breaking  
⇒ **doubling of the unit cell**



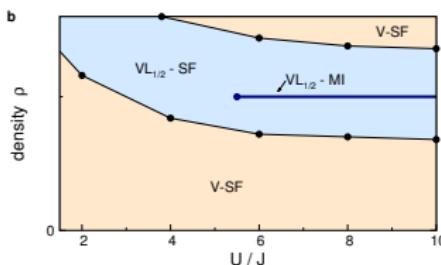
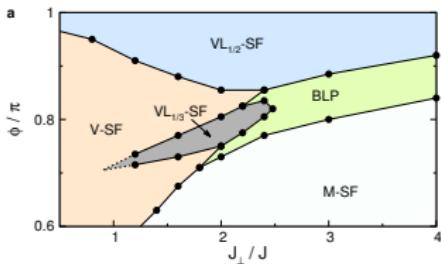
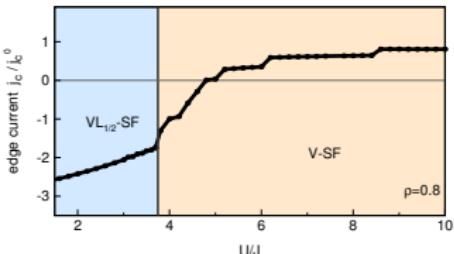
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⇒ **spontaneous current reversal**

⇒ Genuine many-body effect

# Current reversal : Remarks

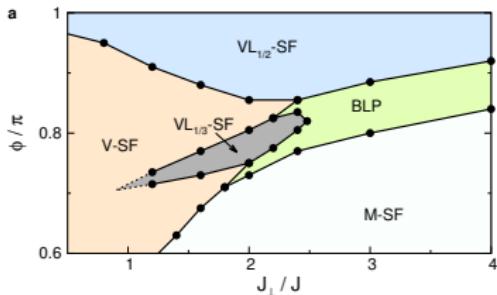


- The amplitude of the reversal is  $\sim 3j_c^0$
- The  $VL_{1/3}$  does not lead to a current reversal (because  $\phi > 2\pi/3$ )
- The current reversal is robust against the presence of a mass gap
- Study at  $n \gg 1$  : the current reversal occurs at arbitrary small  $U/J$
- It is stable against a finite number of defects, and at finite  $T$

# Conclusion

## Strongly-interacting bosons on a ladder with uniform flux :

Phase diagram and observables (DMRG and field theory)



### Plethora of phases :

- Mott insulating and superfluid, Meissner and Vortex phases
- Vortex lattices ( $\rho_v = 1/2$  and  $1/3$ )
- Biased Leg Phase

Vortex lattice phases,  
Spontaneous current reversal

M. P., F. Heidrich-Meisner, I. P. McCulloch, S. Greschner, T. Vekua and

U. Schollwöck, PRB **91** 140406(R) (2015)

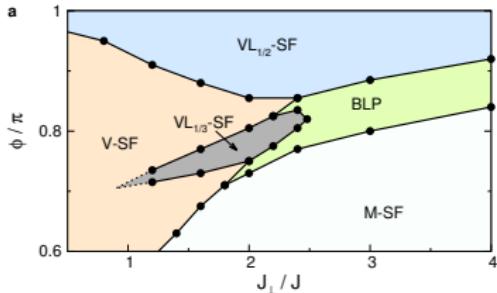
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**Thank you for your attention !**