GPDs in MC and DVCS

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• Exclusive channel: MC much easier...

• Background carefully substracted (model-independently)

• Detector acceptance: radiative corrections + resolution effects

DVCS process





Missing mass squared $e \ p \rightarrow e \ \gamma \ X$



Competing channels:

• π^0 electroproduction: $e \ p \to e \ p \ \pi^0 \to e \ p \ \gamma \ \gamma$

• Associated DVCS: $e \ p \rightarrow e \ N \ \pi \ \gamma$

- Non-resonant: $e \ p \rightarrow e \ N \ \pi \ \gamma$
- Resonant: $e \ p \to e \ (\Delta \text{ or } N^*)\gamma$

 $M_X^2 > (M + m_{\pi^0})^2$ $m_{\Lambda}^2 = 1.52 \,\mathrm{GeV}^2$

$$\pi^0$$
 substraction $(\pi^0 \rightarrow \gamma \gamma)$



• Symmetric decay: minimum angle in lab of 4.4° for $E_{\pi^0}^{\max} = 3.5 \,\text{GeV}$ \Rightarrow Clusters separation

• Asymmetric decay: sometimes only 1-cluster

⇒ Mistaken for DVCS event

$$\pi^0$$
 substraction $(\pi^0 \rightarrow \gamma \gamma)$



Substraction procedure:

- **(**) Compute kinematics of each *detected* π^0 (2 clusters in calorimeter).
- **2** Randomize the decay : sample $\cos \theta$ randomly between [-1,1] a big number of times (~ 5000).
- Sompute the ratio of 2-cluster/1-cluster events generated by this π^0 ($\sim 30\%$ in average).

Repeating this procedure for *each* detected π^0 provides an automatic normalization of the contamination as a function of Q^2 , t, φ , ...

Missing mass

 $\label{eq:Missing mass squared } \begin{array}{l} \text{Missing mass squared } ep \to e' \gamma X \\ M_X^2 = (e+p-e'-\gamma)^2 \qquad \qquad M_X^2 \ \text{cut} = (M+m_\pi)^2 \end{array}$



Exclusivity ensured by missing mass technique

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Associated systematic error

Varying the M_X^2 cut below the $\pi\text{-production}$ threshold gives an estimate of syst. uncertainty



Systematic error due to the missing mass cut $\lesssim 2\%$

M. Defurne

Calculations for DVCS

• M. Vanderhaeghen, et al., Phys Rev C62, 025501 (2000)

Used for data analysis so far

• I. Akushevich and A. Ilyich, Phys Rev D85, 053008 (2012)

Work in progress, only available for BH so far

Radiative corrections (RC)

Not small: $\sim \alpha \log{(Q^2/m_e^2)} \sim 10 - 20\%$

• External RC:

$$I_{\text{ext}}(E_0, \Delta E, t) = \frac{bt}{\Gamma(1+bt)} \left(\frac{\Delta E}{E_0}\right)^{bt} \left[\frac{1}{\Delta E} \left(1 - \frac{\Delta E}{E_0} + \frac{3}{4} \left(\frac{\Delta E}{E_0}\right)^2\right)\right]$$

• Internal real RC:

$$I_{\text{int}}(E_0, E, \nu) = \frac{\nu}{\Delta E} \left(\frac{\Delta E}{E_0}\right)^{\nu}, \qquad \nu = \frac{\alpha}{\pi} \left[\ln \left(\frac{Q^2}{m_e^2}\right) - 1 \right]$$

• Internal virtual RC:

Applied as a global factor depending on the kinematics: O(10%)

	E	E'	θ_e	Q^2	δ_{Vertex}	δ_{Vacuum}	$\delta_{Real,0}$	δ_S	$\delta_{Real,1}(\Delta E)$
	GeV	GeV	deg	GeV^2					
DVCS-1	5.75	3.529	15.6	1.5	-0.228	0.022	0.277	0.068	-0.295
DVCS-2	5.75	2.937	19.3	1.9	-0.236	0.022	0.286	0.069	-0.294
DVCS-3	5.75	2.344	23.8	2.3	-0.242	0.022	0.293	0.070	-0.290

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Radiative corrections

Effect in missing mass distribution



- MC takes into account the radiative lineshape
- Significant effect via the M_X^2 cut

DVCS cross section

$$\frac{d^{5}\sigma}{d^{5}\Phi} = \underbrace{\frac{d^{5}\sigma(|BH|^{2})}{d^{5}\Phi}}_{\text{Known from FF}} + \underbrace{\Gamma \eta \ \mathcal{C}^{\text{DVCS}}(\mathcal{F}, \mathcal{F}^{*})}_{|\text{DVCS}|^{2} \text{ (twist-2)}} +$$

$$(\Gamma_0^{\Re} - \cos(\phi_{\gamma\gamma})\Gamma_1^{\Re}) \Re \mathbf{e} \left[\mathcal{C}^I(\mathcal{F}) \right] + \Gamma_{0,\Delta}^{\Re} \Re \mathbf{e} \left[\mathcal{C}^I + \Delta \mathcal{C}^I \right] (\mathcal{F}) + \cos(2\phi_{\gamma\gamma})\Gamma_2^{\Re} \Re \mathbf{e} \left[\mathcal{C}^I(\mathcal{F}^{\mathrm{eff}}) \right]$$

Interference BH-DVCS

BMK-2002: Belitsky, Müller, Kirchner, Nucl. Phys. **B629**, 323 (2002)

DVCS cross section: data analysis

Challenge:

Very rapid variation with kinematics (mostly due to BH propagators)

Experimental issues:

- Integration of cross section within experimental bins
- Bin migration

Variables:
$$\mathbf{x}_v = \{k, x_B, Q^2, t, \varphi_e, \varphi, v_z\}_v$$
 $\mathbf{x}_e = \{k, x_B, Q^2, t, \varphi_e, \varphi, v_z\}_e$
Exp. bins: $\mathbf{j}_v = \{j_{x_B}, j_{Q^2}, j_t, j_{\varphi}\}_v$ $\mathbf{i}_e = \{i_{x_B}, i_{Q^2}, i_t, i_{\varphi}\}_e$

$$K(\mathbf{x}_e|\mathbf{x}_v)$$

$$\frac{d\sigma}{d\Omega}(\mathbf{x}_v) = \Gamma^{BH}(\mathbf{x}_v) + \sum_{\Lambda=1}^3 \Gamma^{\Lambda}(\mathbf{x}_v) X_{\mathbf{j}_v}^{\Lambda} \equiv \sum_{\Lambda=0}^3 \Gamma^{\Lambda}(\mathbf{x}_v) X_{\mathbf{j}_v}^{\Lambda}$$

Experimental number of counts

$$\begin{split} N(\mathbf{j}_{v}) &= \mathcal{L} \int_{\mathbf{x}_{v} \in \mathsf{Bin}(\mathbf{j}_{v})} \sum_{\Lambda=0}^{3} \Gamma^{\Lambda}(\mathbf{x}_{v}) X_{\mathbf{j}_{v}}^{\Lambda} d\mathbf{x}_{v} = \\ &= \mathcal{L} \sum_{\Lambda=0}^{3} X_{\mathbf{j}_{v}}^{\Lambda} \int_{\mathbf{x}_{v} \in \mathsf{Bin}(\mathbf{j}_{v})} \Gamma^{\Lambda}(\mathbf{x}_{v}) d\mathbf{x}_{v} \\ N(\mathbf{i}_{e}) &= \int_{\mathbf{x}_{e} \in \mathsf{Bin}(\mathbf{i}_{e})} d\mathbf{x}_{e} \sum_{\mathbf{j}_{v}} N(\mathbf{i}_{v}) K(\mathbf{x}_{e} | \mathbf{x}_{v}) = \\ &= \mathcal{L} \sum_{\mathbf{j}_{v}} \sum_{\Lambda=0}^{3} X_{\mathbf{j}_{v}}^{\Lambda} \int_{\mathbf{x}_{e} \in \mathsf{Bin}(\mathbf{i}_{e})} d\mathbf{x}_{e} \int_{\mathbf{x}_{v} \in \mathsf{Bin}(\mathbf{j}_{v})} d\mathbf{x}_{v} \Gamma^{\Lambda}(\mathbf{x}_{v}) K(\mathbf{x}_{e} | \mathbf{x}_{v}) \end{split}$$

We define a bin mapping function:

$$K_{\mathbf{i}_e,\mathbf{j}_v}^{\Lambda} = \int_{\mathbf{x}_e \in \mathsf{Bin}(\mathbf{i}_e)} \int_{\mathbf{x}_v \in \mathsf{Bin}(\mathbf{j}_v)} d\mathbf{x}_e \, d\mathbf{x}_v \, K(\mathbf{x}_e | \mathbf{x}_v) \Gamma^{\Lambda}(\mathbf{x}_v)$$

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$$N(\mathbf{i}_{e}) = \int_{\mathbf{x}_{e}\in\mathsf{Bin}(\mathbf{i}_{e})} d\mathbf{x}_{e} \sum_{\mathbf{j}_{v}} N(\mathbf{i}_{v}) K(\mathbf{x}_{e}|\mathbf{x}_{v}) =$$
$$= \mathcal{L} \sum_{\mathbf{j}_{v}} \sum_{\Lambda=0}^{3} X_{\mathbf{j}_{v}}^{\Lambda} \int_{\mathbf{x}_{e}\in\mathsf{Bin}(\mathbf{i}_{e})} d\mathbf{x}_{e} \int_{\mathbf{x}_{v}\in\mathsf{Bin}(\mathbf{j}_{v})} d\mathbf{x}_{v} \Gamma^{\Lambda}(\mathbf{x}_{v}) K(\mathbf{x}_{e}|\mathbf{x}_{v})$$

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Fit of experimental counts: linear system

$$N^{\mathsf{MC}}(\mathbf{i}_e) = \mathcal{L} \sum_{\mathbf{j}_v, \Lambda} K^{\Lambda}_{\mathbf{i}_e, \mathbf{j}_v} X^{\Lambda}_{\mathbf{j}_v} \,.$$

$$\chi^2 = \sum_{\mathbf{i}_e} \frac{[N^{\mathsf{Exp}}(\mathbf{i}_e) - N^{\mathsf{MC}}(\mathbf{i}_e)]^2}{[\sigma^{\mathsf{Exp}}(\mathbf{i}_e)]^2}$$

$$0 = -\frac{1}{2} \frac{\partial \chi^2}{\partial X_{\mathbf{j}_v}^{\Lambda}} \bigg|_{\overline{\mathbf{X}}_{\mathbf{j}_v}} = \sum_{\mathbf{i}_e} \mathcal{L} K_{\mathbf{i}_e,\mathbf{j}_v}^{\Lambda} \frac{\mathcal{L} \sum_{\mathbf{j}'_v,\Lambda'} K_{\mathbf{i}_e,\mathbf{j}'_v}^{\Lambda'} \overline{X}_{\mathbf{j}'_v}^{\Lambda'} - N^{\mathsf{Exp}}(\mathbf{i}_e)}{[\sigma^{\mathsf{Exp}}(\mathbf{i}_e)]^2}$$

$$0 = \sum_{\mathbf{j}'_{v}, \Lambda'} \alpha^{\Lambda, \Lambda'}_{\mathbf{j}_{v}, \mathbf{j}'_{v}} \,\overline{X}^{\Lambda'}_{\mathbf{j}'_{v}} - \beta^{\Lambda}_{\mathbf{j}_{v}} \quad \forall \mathbf{j}_{v}, \Lambda$$

The linear system is defined by:

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Cross-section results

The fit parameters are:

$$\overline{X}_{\mathbf{j}_{v}}^{\Lambda} = \sum_{\mathbf{j}_{v}^{\prime}, \Lambda^{\prime}} [\alpha^{-1}]_{\mathbf{j}_{v}, \mathbf{j}_{v}^{\prime}}^{\Lambda, \Lambda^{\prime}} \beta_{\mathbf{j}_{v}^{\prime}}^{\Lambda^{\prime}}.$$

The covariance matrix of the fitted parameters is:

$$V^{\Lambda,\Lambda'}_{\mathbf{j}_{v},\mathbf{j}_{v}'} = [\alpha^{-1}]^{\Lambda,\Lambda'}_{\mathbf{j}_{v},\mathbf{j}_{v}'}.$$

$$\frac{d^5 \sigma^{\text{Exp}}(i)}{d^5 \Phi} = \frac{d^5 \sigma^{\text{Fit}}(i)}{d^5 \Phi} N_i^{\text{Exp}} / N_i^{\text{Fit}},$$

Experimental results



Belitsky, Müller, Kirchner, Nucl. Phys. B629, 323 (2002)

Results with different parametrizations

• Reanalysis with BMK-2010 and comparison

Low -t

High -t



- Stable within statistical uncertainties
- Especially at low -t where the acceptance is best known

M. Defurne

- MC for exclusive reactions is required for detector acceptance:
 - Bin migration
 - Resolution effects
 - Rapid variation of the cross section within experimental bins
- Background is usually under control is exclusivity is good enough
- Radiative corrections is necessarily part of the MC simulation
- Fast GPD and RC event generators can improve the systematic uncertainties of DVCS cross section results