

Tomography of hadrons in high-energy exclusive processes

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<http://www.int.washington.edu/PROGRAMS/14-55w/>

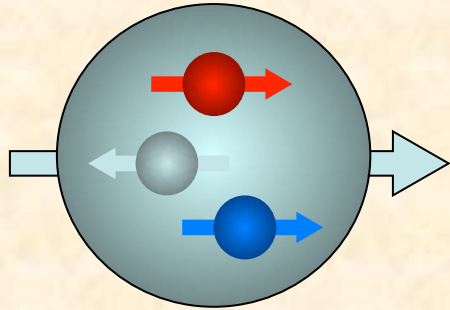
February 28, 2014

Contents

- **Introduction to internal structure of hadrons**
- **GPDs at hadron facilities**
SK, M. Strikman, K. Sudoh, PRD 80 (2009) 074003.
- **Constituent-counting rule for exotic hadrons**
H. Kawamura, SK, T. Sekihara, PRD 88 (2013) 034010.
- **GPDs and GDAs for exotic hadrons**
H. Kawamura, SK, arXiv:1312.1596, Phys. Rev. D in press (2014).
- **Comments on J-PARC project**
- **Summary**

Introduction to
Internal structure of hadrons

Nucleon Spin



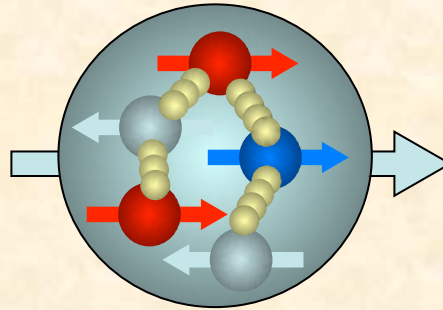
Naïve Quark Model

$$\Delta\Sigma = \Delta u_v + \Delta d_v = 1$$

Electron / muon scattering

$$\Delta\Sigma \approx 0.3$$

Almost none of nucleon spin is carried by quarks!



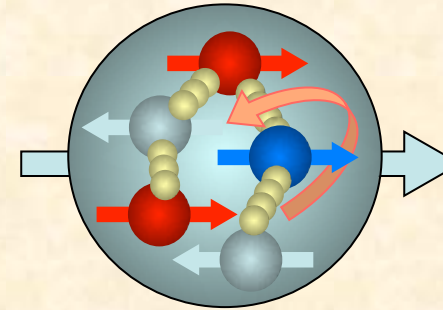
QCD

Sea-quarks and gluons?

Gluon: ΔG

Sea-quarks: Δq_{sea}

Recent data indicate ΔG is small at $x \sim 0.1$.



Orbital angular momenta ?

L_q, L_g

Future experiments

$$\text{Nucleon Spin: } \frac{1}{2} = \frac{1}{2} \underbrace{(\Delta u_v + \Delta d_v + \Delta q_{sea})}_{\Delta\Sigma} + \Delta G + L_q + L_g$$

Progress in exotic hadrons

$q\bar{q}$ Meson
 q^3 Baryon

$q^2\bar{q}^2$ Tetraquark
 $q^4\bar{q}$ Pentaquark
 q^6 Dibaryon

...
 $q^{10}\bar{q}$ e.g. Strange tribaryon

...
 gg Glueball

...

- $\Theta^+(1540)???:$ LEPS

$uudd\bar{s} ?$

Pentaquark?

- **Kaonic nuclei?**: KEK-PS, ...
 Strange tribaryons, ...

$K^- pnn, K^- ppn ?$
 $K^- pp ?$

- **X (3872), Y(3940):** Belle
 Tetraquark, $D\bar{D}$ molecule

$c\bar{c}$
 $D^0(c\bar{u})\bar{D}^0(\bar{c}u)$
 $D^+(c\bar{d})D^-(\bar{c}d) ?$

- **$D_{sJ}(2317), D_{sJ}(2460):$** BaBar, CLEO, Belle
 Tetraquark, DK molecule

$c\bar{s}$
 $D^0(c\bar{u})K^+(u\bar{s})$
 $D^+(c\bar{d})K^0(d\bar{s}) ?$

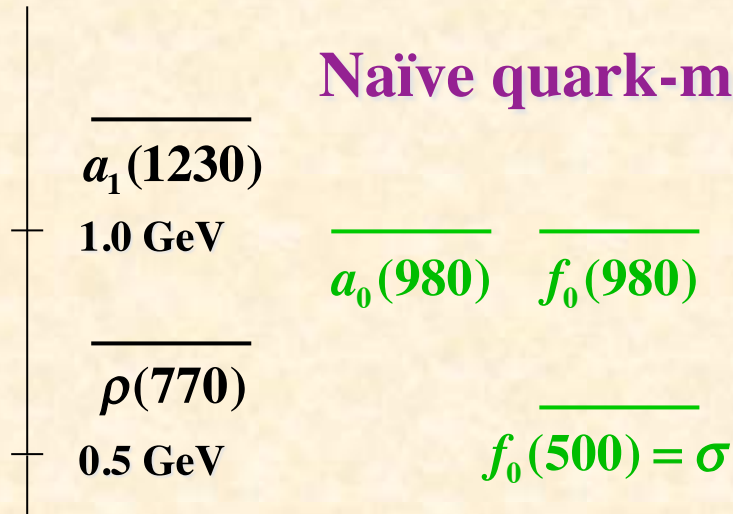
- **Z (4430):** Belle
 Tetraquark, ...

$c\bar{c}u\bar{d}, D$ molecule?

- ...

Scalar mesons $J^P=0^+$ at $M \sim 1$ GeV

Naïve quark-model



$$\sigma = f_0(500) = \frac{1}{\sqrt{2}}(u\bar{u} + d\bar{d})$$

$$f_0(980) = s\bar{s}$$

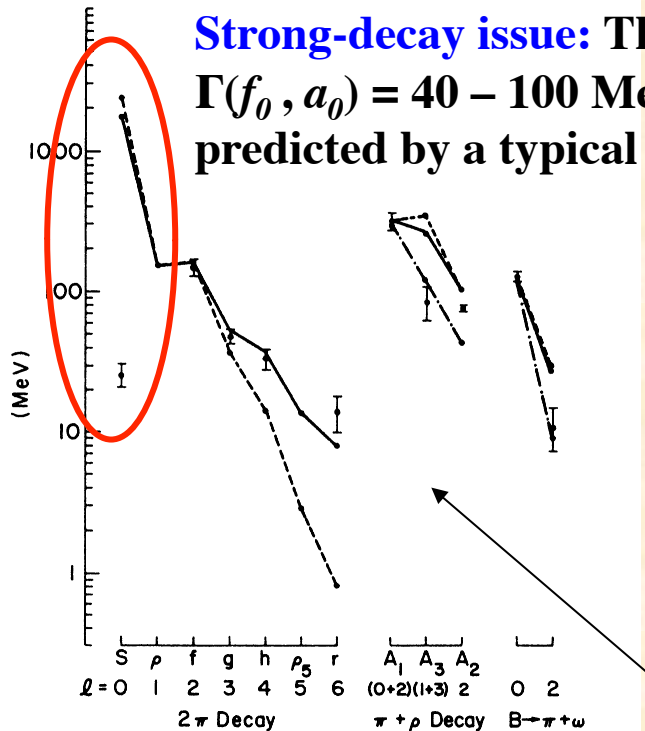
$$a_0(980) = u\bar{d}, \frac{1}{\sqrt{2}}(u\bar{u} - d\bar{d}), d\bar{u}$$

Naive model: $m(\sigma) \sim m(a_0) < m(f_0)$

↕ contradiction

Experiment: $m(\sigma) < m(a_0) \sim m(f_0)$

Strong-decay issue: The experimental widths $\Gamma(f_0, a_0) = 40 - 100$ MeV are too small to be predicted by a typical quark model.



These issues could be resolved

if f_0 (a_0) is a tetraquark ($qq\bar{q}\bar{q}$) or a $K\bar{K}$ molecule, namely an "exotic" hadron.

Radiative decay: F. E. Close, N. Isgur, and SK, Nucl. Phys. B389 (1993) 513.

SK and V. R. Pandharipande, Phys. Rev. D38 (1988) 146.

Exotic hadrons by fragmentation functions

“Favored” and “disfavored” (unfavored) fragmentation functions
 Possibility of finding exotic hadrons in high-energy processes

Hirai, SK, Oka, Sudoh,
 PRD 77 (2008) 017504.

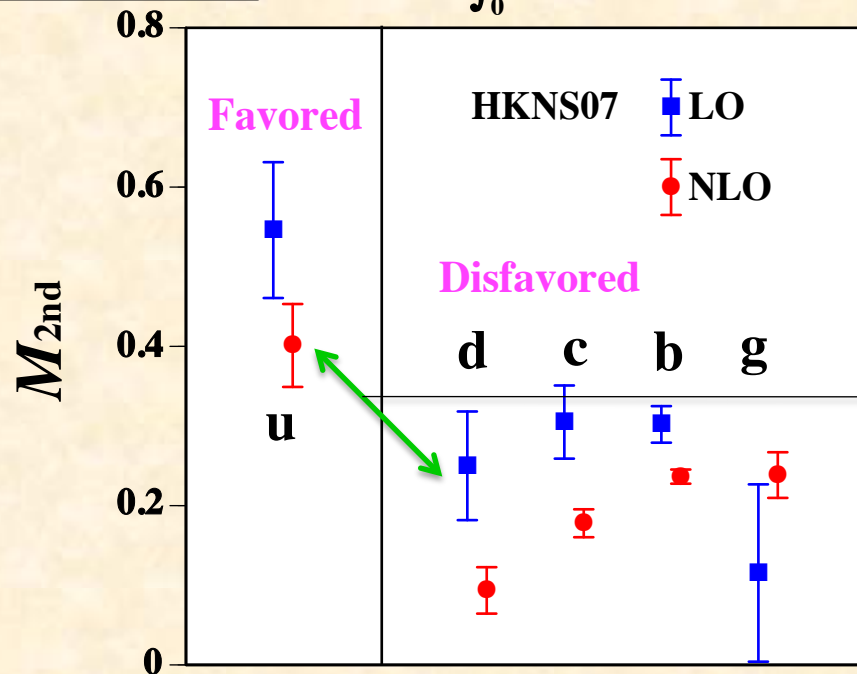
Possibilities for $f_0(980)$: $\frac{1}{\sqrt{2}}(u\bar{u} + d\bar{d})$, $s\bar{s}$, $\frac{1}{\sqrt{2}}(u\bar{u}s\bar{s} + d\bar{d}s\bar{s})$, $K\bar{K}$, or gg

e.g. if $f_0(980) = s\bar{s}$: **favored** $s, \bar{s} \rightarrow f_0$; **disfavored** $u, d, \bar{u}, \bar{d} \rightarrow f_0, \dots$

$f_0(980)$: Belle analysis is possible in principle.

Pion case

$$M_{2nd} = \int_0^1 dz z D_i^{\pi^+}(z)$$



2nd moments of
 M. Hirai, SK, T.-H. Nagai, K. Sudoh,
 PRD 75 (2007) 094009.

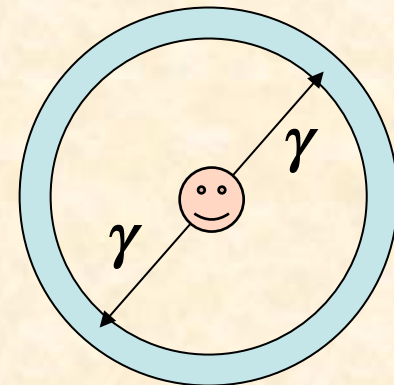
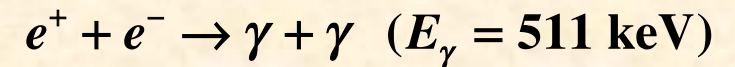
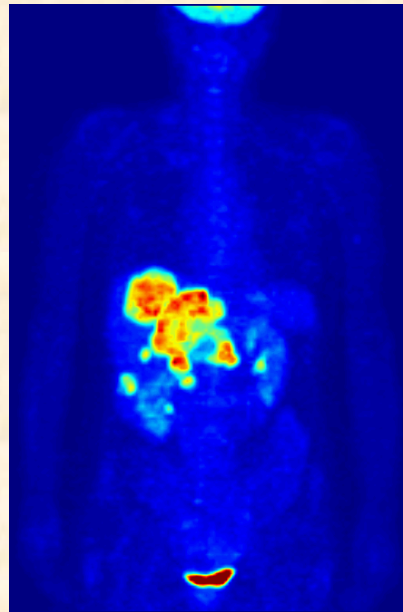
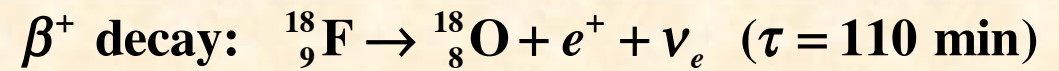
There are distinct differences between the favored and disfavored 2nd moments.
 → It could be used for exotic-hadron studies.

Tomography

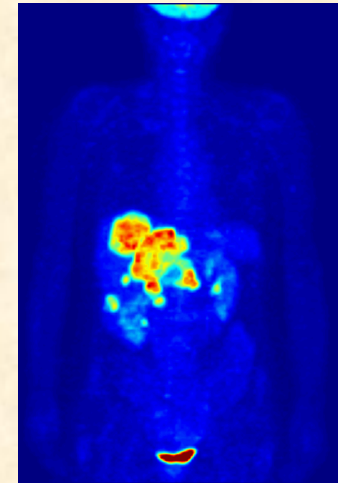
- CT (Computed Tomography)
- PET (Positron Emission Tomography)



© Jens Langner



Nucleon (hadron) tomography

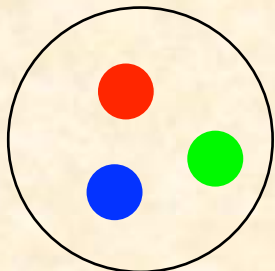


Classical density distribution

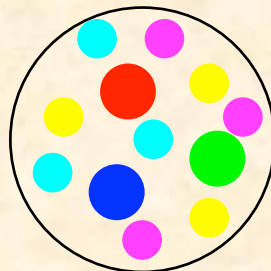
3D picture of nucleon
(Density distribution of quantum system:
Quantum tomography)

1D(Bjorken-x) picture@HERA

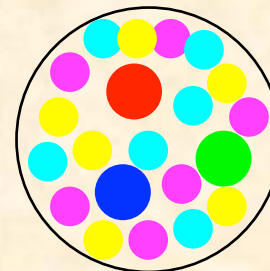
Low energy



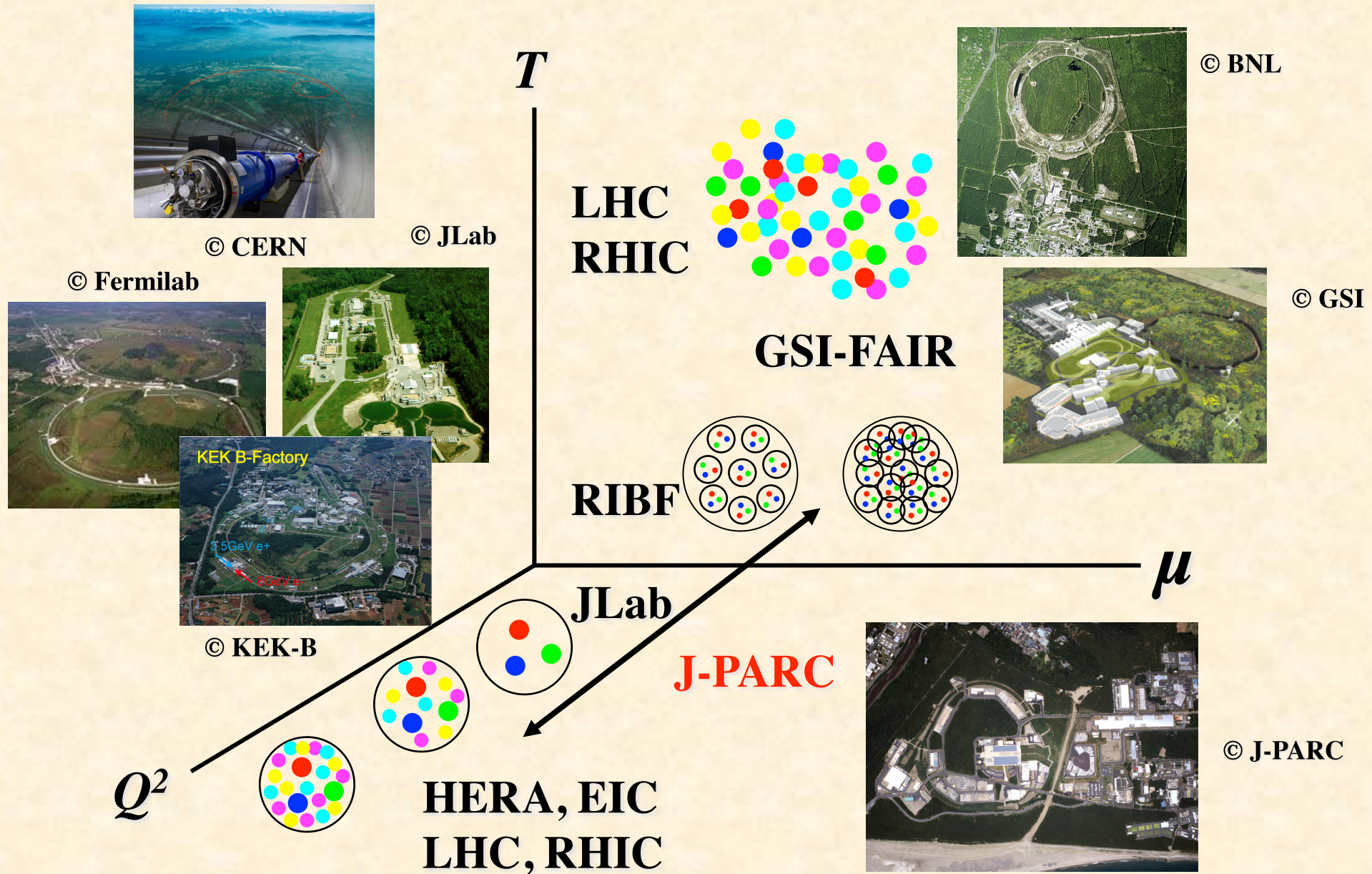
Intermediate energy



High energy



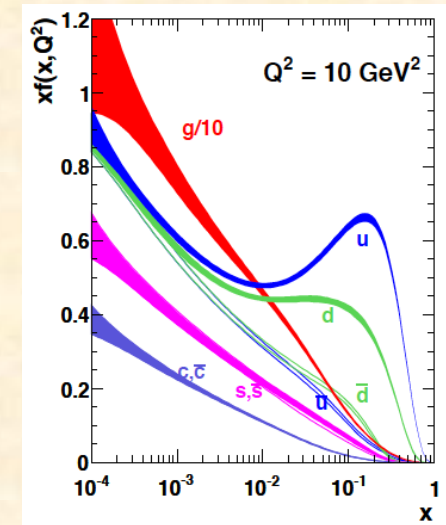
Test apparatus corresponds to “PET”



Wigner distribution and various structure functions

Wigner operator: $\hat{w}(k_+, \vec{k}_\perp, \vec{r}) \equiv \int d\xi_- d^2\xi_\perp e^{i(\xi_- k_+ - \vec{\xi}_\perp \cdot \vec{k}_\perp)} \bar{\psi}(\vec{r} - \vec{\xi} / 2) \psi(\vec{r} + \vec{\xi} / 2)$

Wigner distribution: $W(x, \vec{k}_\perp, \vec{r}) \equiv \int \frac{d^3q}{(2\pi)^3} \langle \vec{q} / 2 | \hat{w}(\vec{r}, k_+, \vec{k}_\perp) | -\vec{q} / 2 \rangle, \quad x = k_+ / p_+$



Form factor

PDF (Parton Distribution Function)

$$\int dx d^2k_\perp$$

$$\int d^2k_\perp d^3r$$

Wigner distribution $W(x, \vec{k}_\perp, \vec{r})$

3D world

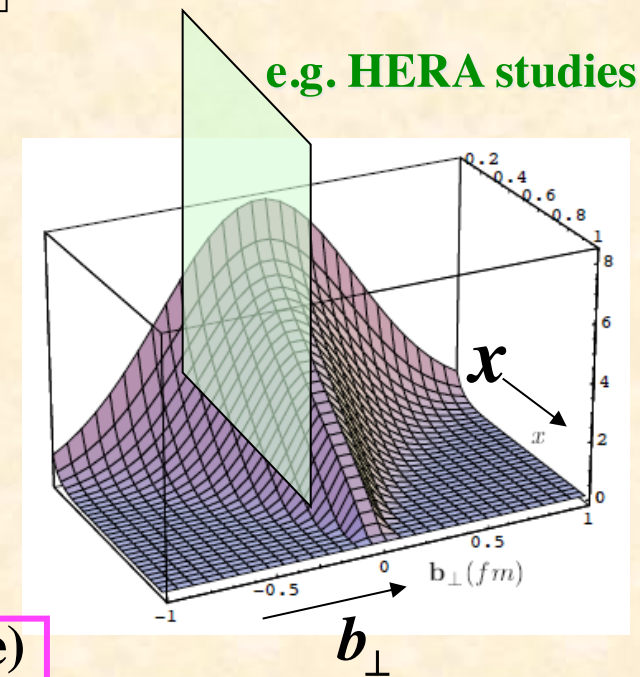


$$\int d^3r$$

TMD (Transverse Momentum Dependent) parton distribution

$$\int d^2k_\perp dz$$

GPD (Generalized Parton Distribution)



s - t crossing \rightarrow

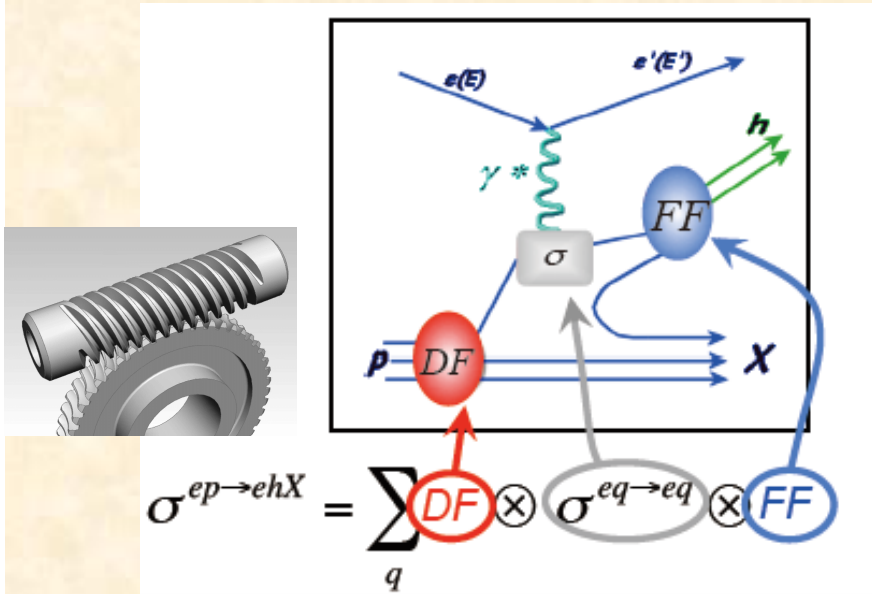
$\gamma \rightarrow h\bar{h}$

GDA (Generalized Distribution Amplitude)

TMD distributions

Contalbrigo@Jlab-PAC-2011

		quark polarisation			
nucleon polarisation		N/q	U	L	T
U	f_1 Number Density		h_1^\perp Boer-Mulders		
L	g_1 Helicity		h_{1L}^\perp Worm-gear		
T	f_{1T}^\perp Sivers	g_{1T}^\perp Worm-gear	h_1^\perp Transversity	h_{1T}^\perp Pretzelosity	



SIDIS cross section:

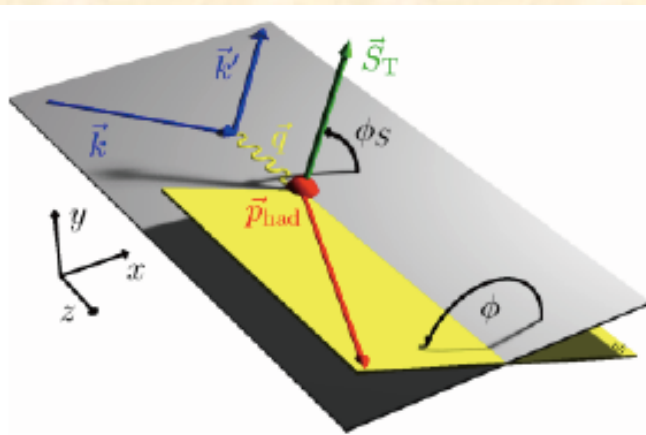
$$\frac{d^6 \sigma}{dx dy dz d\phi_S d\phi dP_{h\perp}^2} \stackrel{\text{Leading}}{\propto} S_T \left\{ \sin(\phi - \phi_S) F_{UT,T}^{\sin(\phi - \phi_S)} \right\}$$

$f_{1T}^\perp \otimes D_1$ $h_1^\perp \otimes H_1^\perp$

$$+ S_T \left\{ \varepsilon \sin(\phi + \phi_S) F_{UT}^{\sin(\phi + \phi_S)} + \varepsilon \sin(3\phi - \phi_S) F_{UT}^{\sin(3\phi - \phi_S)} \right\}$$

$g_{1T}^\perp \otimes D_1$

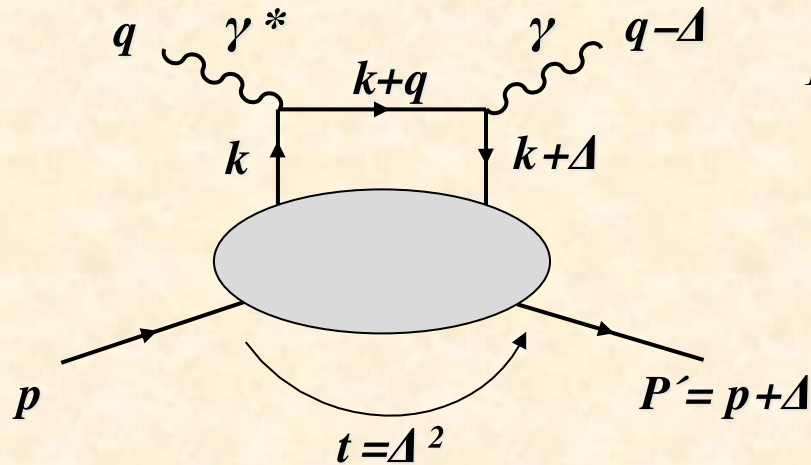
$$+ S_T \lambda_e \left\{ \sqrt{1 - \varepsilon^2} \cos(\phi - \phi_S) F_{LT}^{\cos(\phi - \phi_S)} \right\} + \dots$$



GPDs in the ERBL region at hadron facilities

SK, M. Strikman, K. Sudoh, PRD 80 (2009) 074003.

Generalized Parton Distributions (GPDs)



$$P = \frac{p^+ + p'^+}{2}, \quad \Delta = p' - p$$

Bjorken variable $x = \frac{Q^2}{2p \cdot q}$

Momentum transfer squared $t = \Delta^2$

Skewness parameter $\xi = \frac{p^+ - p'^+}{p^+ + p'^+} = -\frac{\Delta^+}{2P^+}$

GPDs are defined as correlation of off-forward matrix:

$$\int \frac{dz^-}{4\pi} e^{ixP^+z^-} \langle p' | \bar{\psi}(-z/2) \gamma^+ \psi(z/2) | p \rangle \Big|_{z^+=0, \bar{z}_\perp=0} = \frac{1}{2P^+} \left[H(x, \xi, t) \bar{u}(p') \gamma^+ u(p) + E(x, \xi, t) \bar{u}(p') \frac{i\sigma^{+\alpha} \Delta_\alpha}{2M} u(p) \right]$$

Forward limit: PDFs $H(x, \xi, t) \Big|_{\xi=t=0} = f(x)$

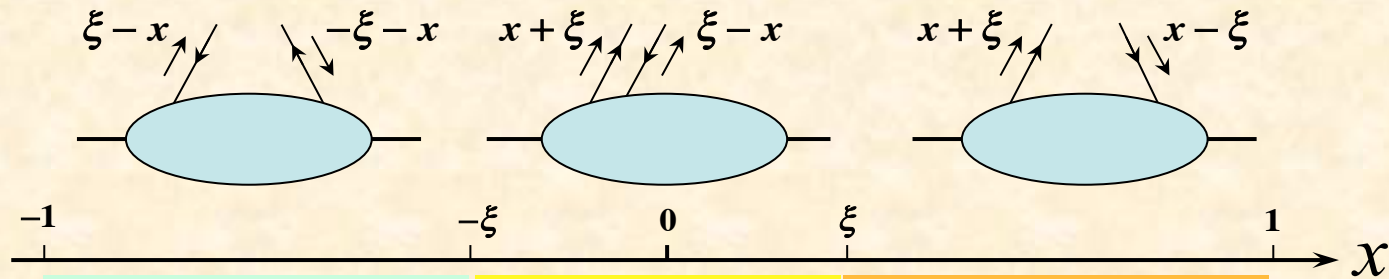
First moments: Form factors

Dirac and Pauli form factors F_1, F_2 $\int dx H(x, \xi, t) = F_1(t), \quad \int dx E(x, \xi, t) = F_2(t)$

Second moments: Angular momenta

Sum rule: $J_q = \frac{1}{2} \int dx x [H_q(x, \xi, t=0) + E_q(x, \xi, t=0)], \quad J_q = \frac{1}{2} \Delta q + L_q$

GPDs in different x regions and GPDs at hadron facilities



$-1 < x < \xi$ ($x + \xi < 0, x - \xi < 0$)

$\xi < x < 1$ ($x + \xi > 0, x - \xi > 0$)

$-\xi < x < \xi$ ($x + \xi > 0, x - \xi < 0$)

Quark distribution

Emission of quark with momentum fraction $x + \xi$
 Absorption of quark with momentum fraction $x - \xi$

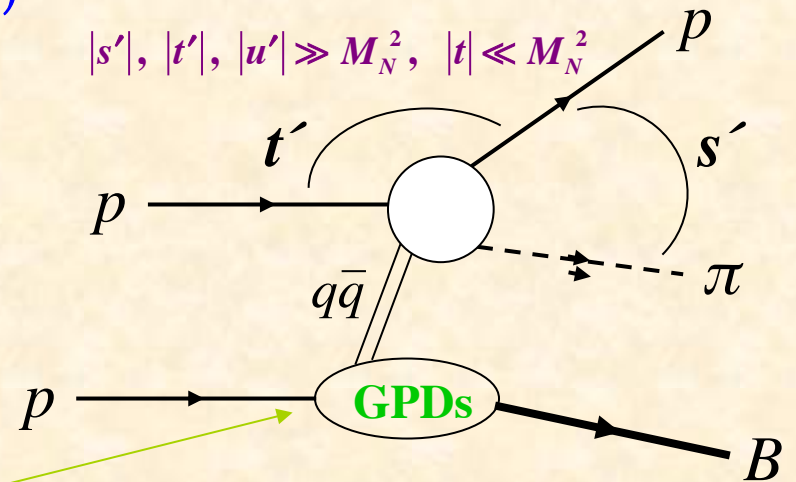
Meson-like distribution amplitude

Emission of quark with momentum fraction $x + \xi$
 Emission of antiquark with momentum fraction $\xi - x$

Antiquark distribution

Emission of antiquark with momentum fraction $\xi - x$
 Absorption of antiquark with momentum fraction $-\xi - x$

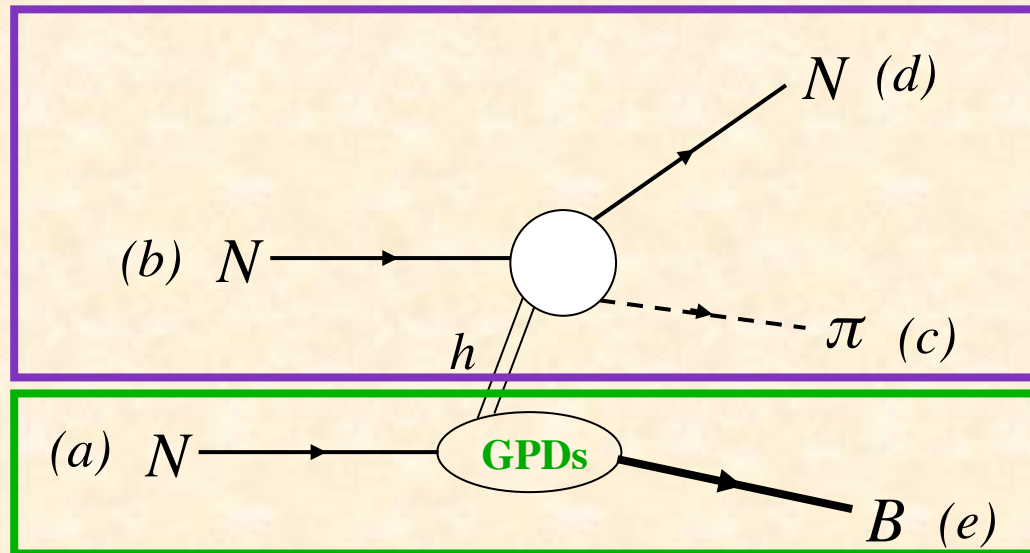
Consider a hard reaction with
 $|s'|, |t'|, |u'| \gg M_N^2, |t| \ll M_N^2$



GPDs at J-PARC: S. Kumano, M. Strikman, and K. Sudoh, PRD 80 (2009) 074003.

Efremov-Radyushkin -Brodsky-Lepage (ERBL) region

Cross section estimates



$\frac{d\sigma(s',t')}{dt'}$ so as to explain
AGS experimental data on
 $\pi + p \rightarrow \pi + p, \pi + p \rightarrow \rho + p$

This part is expressed by GPDs.

Purposes of our studies:

- (1) The ultimate purpose is to extract the GPDs in the ERBL region by measurements at hadron facilities in addition to lepton ones.
- (2) Since our work is the first one to point out the GPD studies at hadron reactions, we estimate the order of magnitude of cross sections simply by using meson-pole expressions of the GPDs.
For experimental feasibility studies.

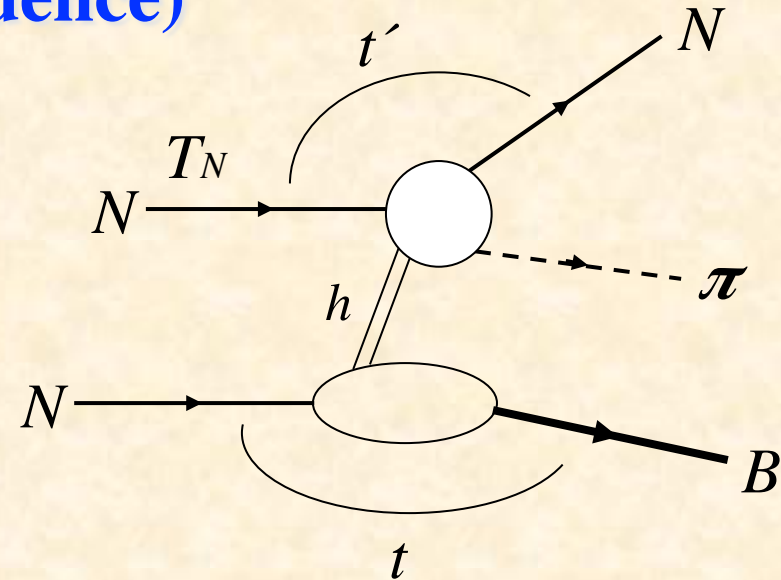
Cross section estimate (ξ dependence)

Skewness parameter: $\xi = \frac{p_N^+ - p_B^+}{p_N^+ + p_B^+}$

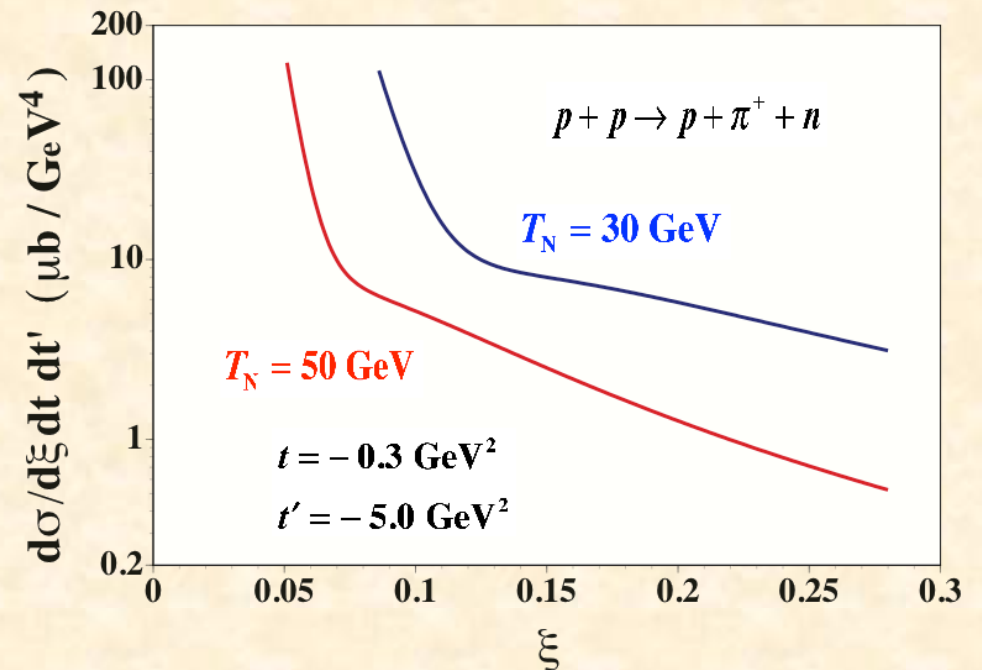
$\frac{d\sigma}{d\xi dt dt'} \left(\frac{\mu\text{b}}{\text{GeV}^2} \right)$ as a function of ξ

at fixed $T_N = 30$ (50) GeV,

$t = -0.3 \text{ GeV}^2$, $t' = -5 \text{ GeV}^2$.



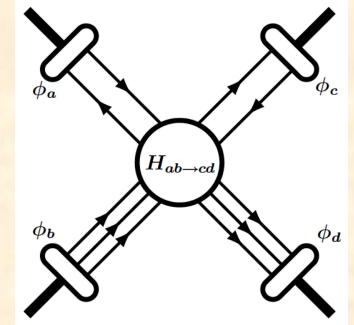
At this stage, our numerical results are for rough order of magnitude estimates on cross sections by assuming π - and ρ -like intermediate states.



Constituent-counting rule for exotic hadrons

H. Kawamura, SK, T. Sekihara, PRD 88 (2013) 034010.

Constituent-counting rule in perturbative QCD: Hard exclusive processes $a + b \rightarrow c + d$



Consider the hard exclusive hadron reaction $a + b \rightarrow c + d$

$$M_{ab \rightarrow cd} = \int d[x_a] d[x_b] d[x_c] d[x_d] \phi_c([x_c]) \phi_d([x_d]) H_M([x_a], [x_b], [x_c], [x_d], Q^2) \phi_a([x_a]) \phi_b([x_b])$$

ϕ_p = proton distribution amplitude, H_M = hard amplitude (calculated in pQCD)

Rule for estimating $M_{ab \rightarrow cd}$

(1) Feynman diagram: Draw leading and connected Feynman diagram by connecting $n / 2$ quark lines by gluons.

(2) Gluon propagators: The factor $1/P^2$ is assigned for each gluon propagator.

There are $n / 2 - 1$ gluon propagators $\sim 1/(P^2)^{n/2-1}$.

(3) Quark propagators: The factor $1/P$ is assigned for each quark propagator.

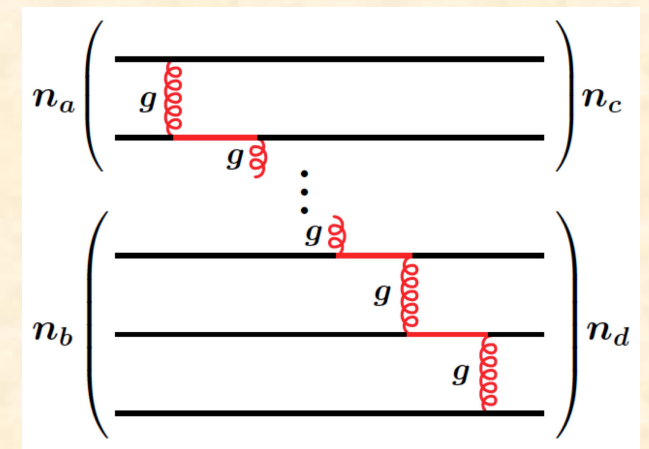
There are $n / 2 - 2$ gluon propagators $\sim 1/(P)^{n/2-2}$.

(4) External quarks: The factor \sqrt{P} is assigned for each external quark.

There are n gluon propagators $\sim (\sqrt{P})^n$.

$$M_{ab \rightarrow cd} \sim \frac{1}{(P^2)^{n/2-1}} \frac{1}{(P)^{n/2-2}} (\sqrt{P})^n = \frac{(P)^{n/2}}{(P)^{n-2} (P)^{n/2-2}} = \frac{1}{(P)^{n-4}} \sim \frac{1}{s^{n/2-2}}$$

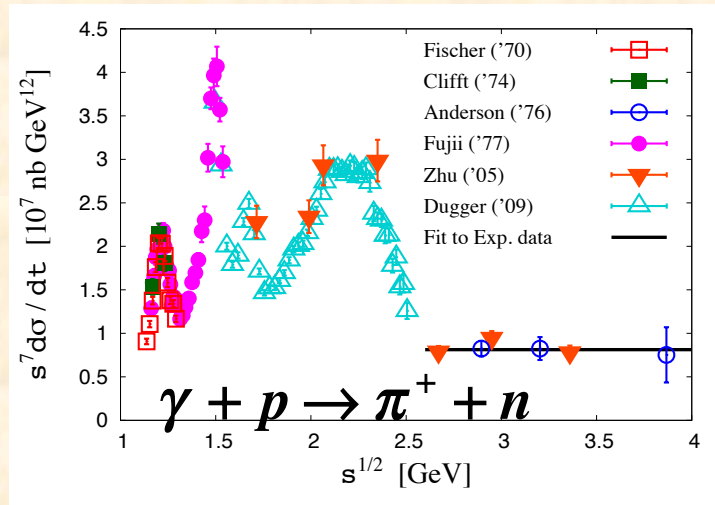
Cross section: $\frac{d\sigma_{ab \rightarrow cd}}{dt} \simeq \frac{1}{16\pi^2} \sum_{\text{spol}} |M_{ab \rightarrow cd}|^2 \sim \frac{1}{s^{n-2}}$



Constituent-counting rule, Transition from hadron degrees of freedom to quark-gluon ones

Typical current situation

- Transition from hadron d.o.f to quark d.o.f.
- (Looks like) Constituent-counting scaling



BNL: C. White et al., PRD 49 (1994) 58.

No.	Interaction	Cross section		$n-2$ ($\frac{d\sigma}{dt} \sim 1/s^{n-2}$)
		E838	E755	
1	$\pi^+ p \rightarrow p\pi^+$	132 ± 10	4.6 ± 0.3	6.7 ± 0.2
2	$\pi^- p \rightarrow p\pi^-$	73 ± 5	1.7 ± 0.2	7.5 ± 0.3
3	$K^+ p \rightarrow pK^+$	219 ± 30	3.4 ± 1.4	$8.3^{+0.6}_{-1.0}$
4	$K^- p \rightarrow pK^-$	18 ± 6	0.9 ± 0.9	≥ 3.9
5	$\pi^+ p \rightarrow p\rho^+$	214 ± 30	3.4 ± 0.7	8.3 ± 0.5
6	$\pi^- p \rightarrow p\rho^-$	99 ± 13	1.3 ± 0.6	8.7 ± 1.0
13	$\pi^+ p \rightarrow \pi^+ \Delta^+$	45 ± 10	2.0 ± 0.6	6.2 ± 0.8
15	$\pi^- p \rightarrow \pi^+ \Delta^-$	24 ± 5	≤ 0.12	≥ 10.1
17	$pp \rightarrow pp$	3300 ± 40	48 ± 5	9.1 ± 0.2
18	$\bar{p}p \rightarrow \bar{p}p$	75 ± 8	≤ 2.1	≥ 7.5

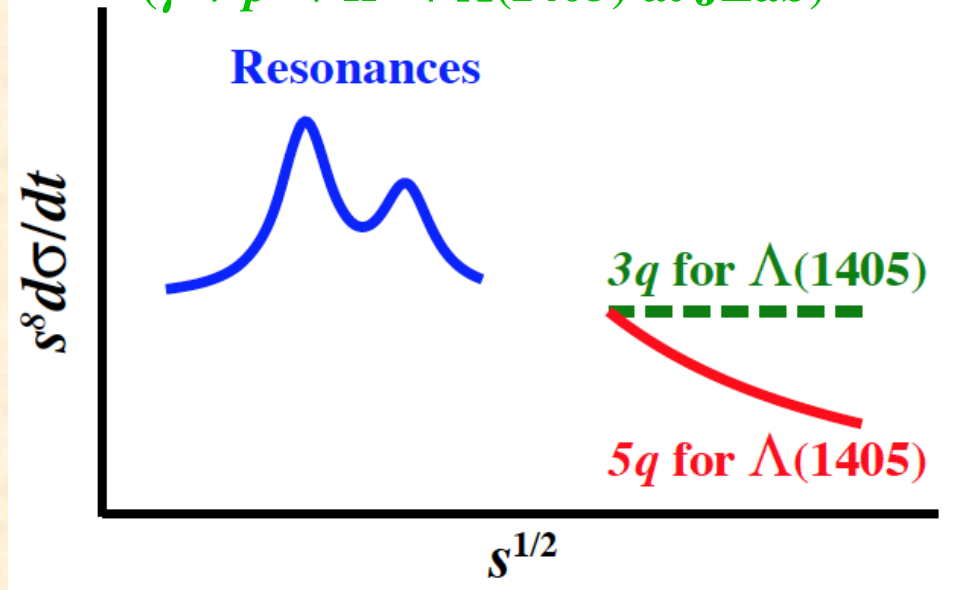
Our idea

- Transition from hadron d.o.f to quark d.o.f. for exotic-hadron production
- **Internal structure of exotic hadrons by constituent-counting scaling**

Exotic hadron production

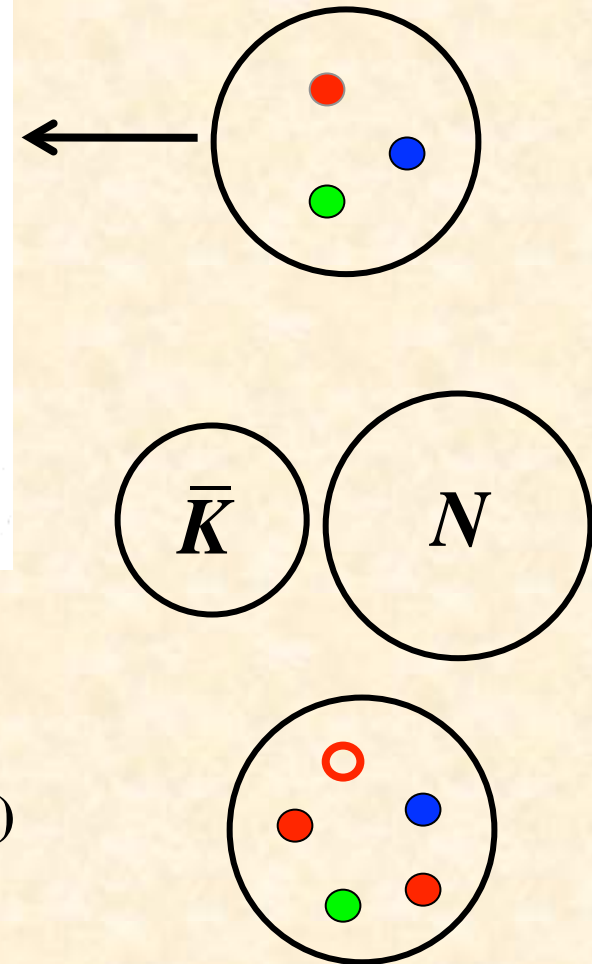
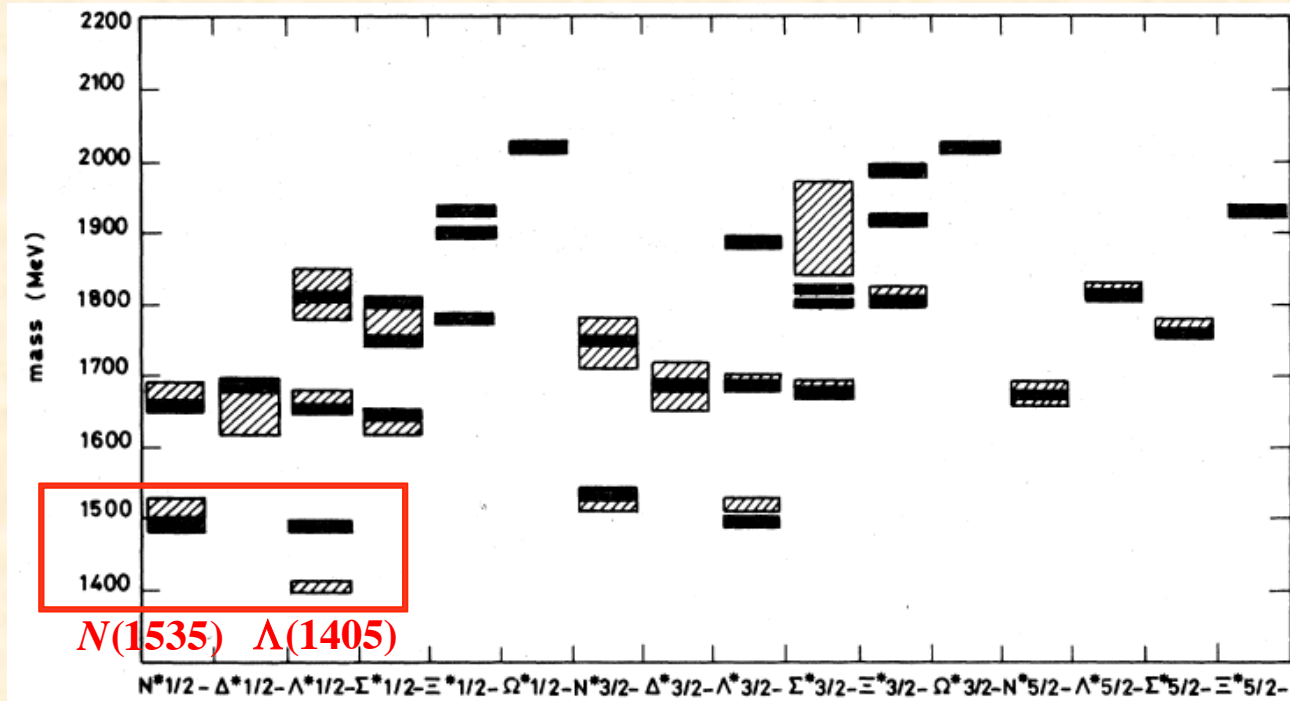
$$\pi^0 + p \rightarrow K^0 + \Lambda(1405)$$

$$(\gamma + p \rightarrow K^+ + \Lambda(1405) \text{ at JLab})$$



$\Lambda(1405)$: exotic hadron?

Negative-parity baryons
N. Isgur and G. Karl,
PRD 18 (1978) 4187.

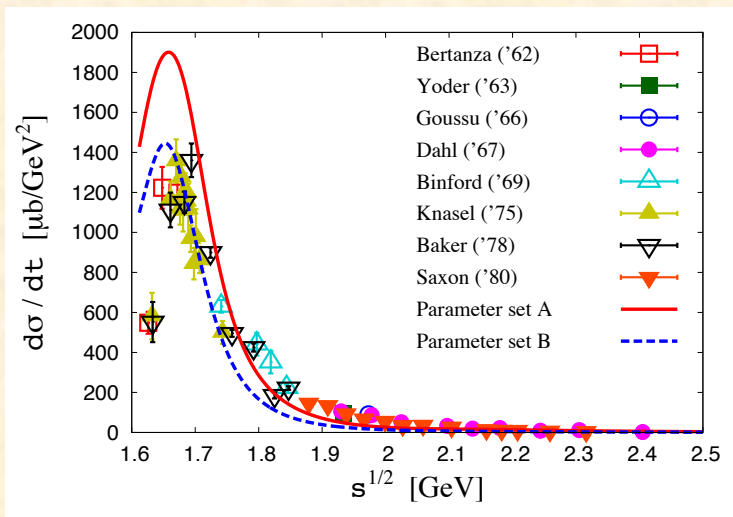
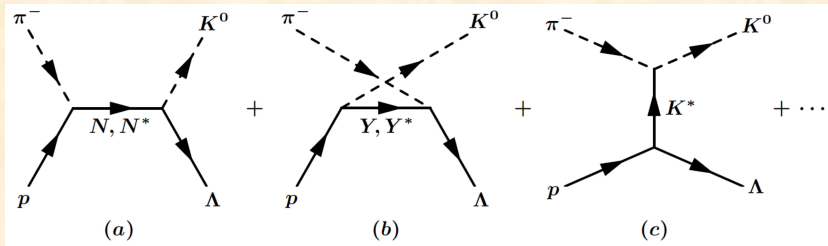


Most spectra agree with the ones by a $3q$ -picture

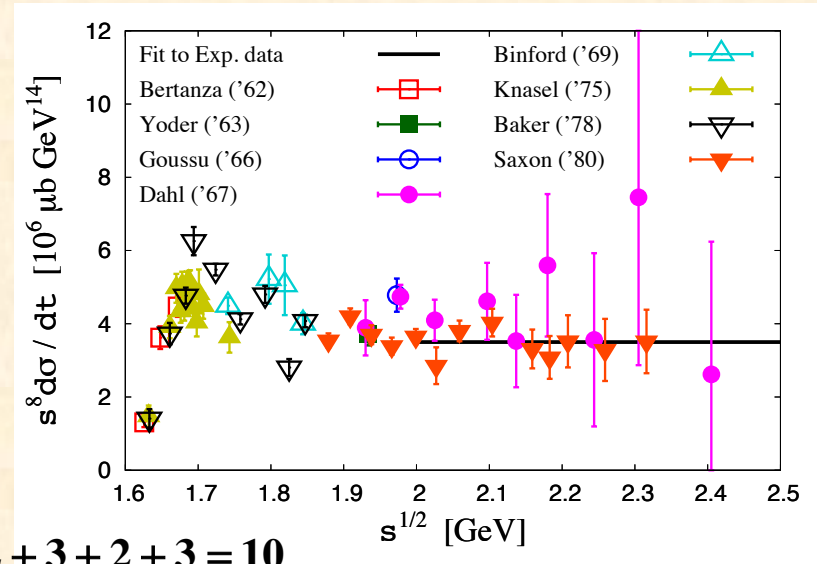
- Only $\Lambda(1405)$ deviates from the measurement.
- Difficult to understand the small mass of $\Lambda(1405)$ in comparison with $N(1535)$.
→ $\bar{K}N$ molecule or penta-quark ($qqqq\bar{q}$)?

Ordinary-hadron production $\pi^- + p \rightarrow K^0 + \Lambda$ as a reference

At low energies



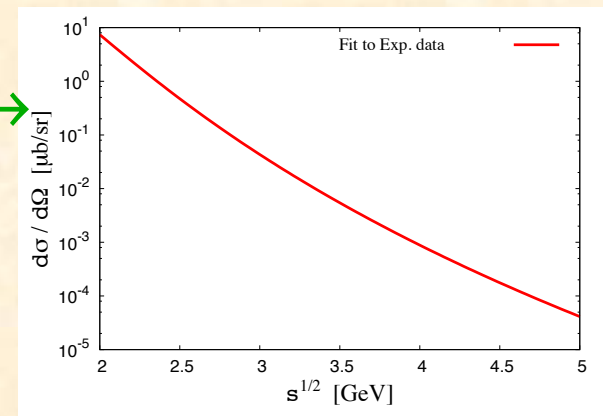
From low to higher energies



$$n = 2 + 3 + 2 + 3 = 10$$

$$\frac{d\sigma_{ab \rightarrow cd}}{dt} = \frac{\text{const}}{s^{n-2}}, \quad n = 10.1 \pm 0.6, \text{ encouraging!}$$

Our prediction at high energies →



Exotic-hadron production $\pi^- + p \rightarrow K^0 + \Lambda(1405)$

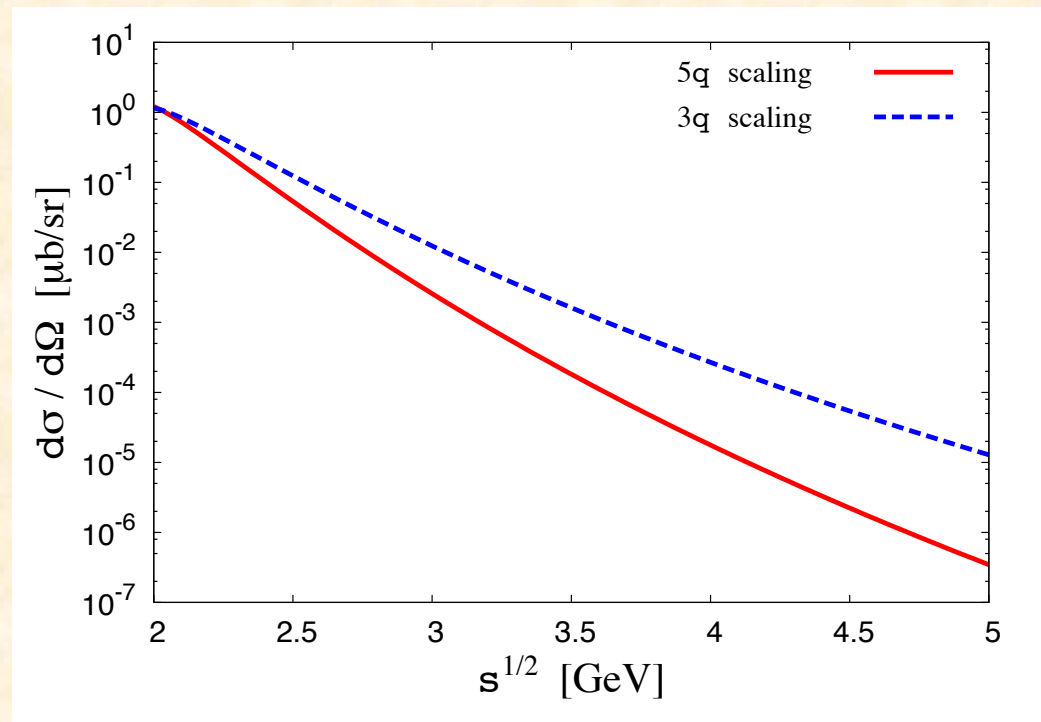
Theoretical and experimental situation
is no as good as the one for the ground Λ .

$$n = 2 + 3 + 2 + 3 = 10 \text{ if } \Lambda(1405) = \text{three-quark state}$$

$$= 2 + 3 + 2 + 5 = 12 \text{ if } \Lambda(1405) = \text{five-quark state} \\ \text{(including } \bar{K}N \text{ molecule)}$$

$$\frac{d\sigma_{ab \rightarrow cd}}{dt} = \frac{\text{const}}{s^{n-2}}, \quad n = 10 \text{ or } 12$$

Our prediction at high energies



GPDs and GDAs for exotic hadrons

H. Kawamura and SK

arXiv:1312.1596 (Phys. Rev. D in press)

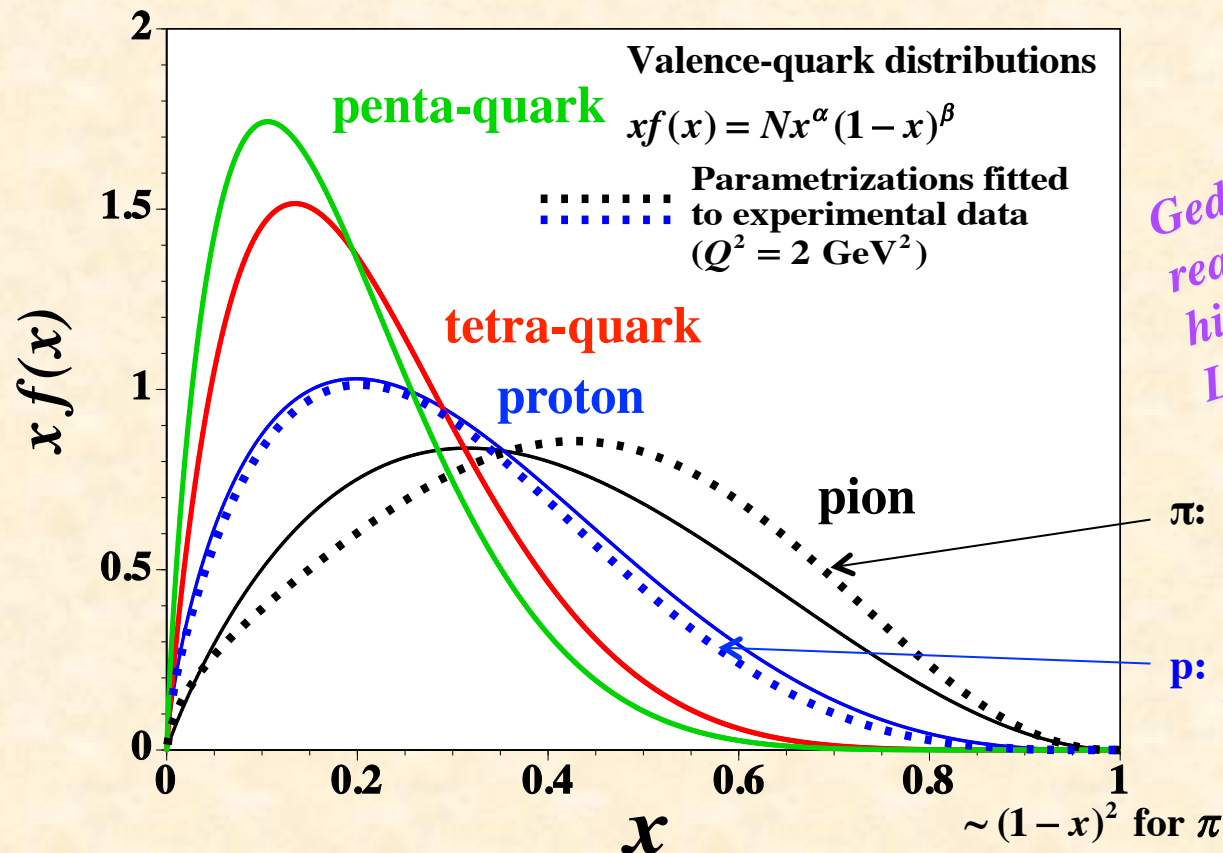
Simple function of GPDs

$$H_q^h(x,t) = f(x)F(t,x)$$

M. Guidal, M.V. Polyakov,
A.V. Radyushkin, M. Vanderhaeghen,
PRD 72, 054013 (2005).

Longitudinal-momentum distribution (PDF) for valence quarks: $f(x) = q_v(x) = c_n x^{\alpha_n} (1-x)^{\beta_n}$

- Valence-quark number sum rule (charge and baryon numbers): $\int_0^1 dx f(x) = n$
- Constituent counting rule at $x \rightarrow 1$: $\beta_n = 2n - 3 + 2\Delta S$ (n = number of constituents)
- Momentum carried by quarks $\langle x \rangle_q \simeq \int_0^1 dx x f(x)$



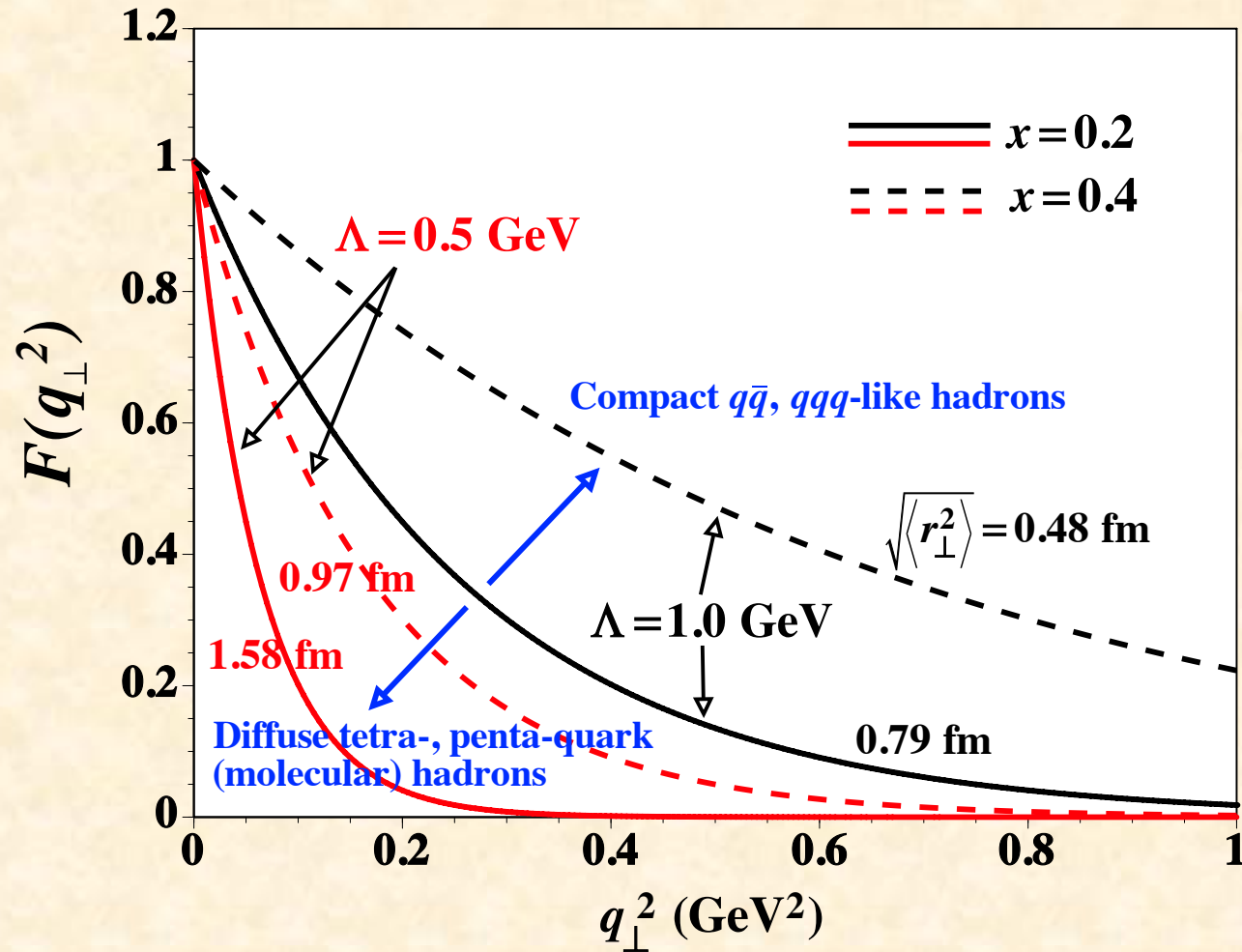
*Gedankenexperiment, but ...
read our paper for studying exotics in
high-energy processes at KEK-B,
Linear Collider, ...*

π : M. Aicher, A. Schafer, W. Vogelsang,
PRL 105 (2010) 252003.

p : A. D. Martin, R. G. Roberts,
W. J. Stirling, PLB 636, 259 (2006)

Two-dimensional form factor

$$H_q^h(x,t) = f(x)F(t,x), \quad F(t,x) = e^{(1-x)t/(x\Lambda^2)}, \quad \langle r_\perp^2 \rangle = \frac{4(1-x)}{x\Lambda^2}$$



GPDs for exotic hadrons !?

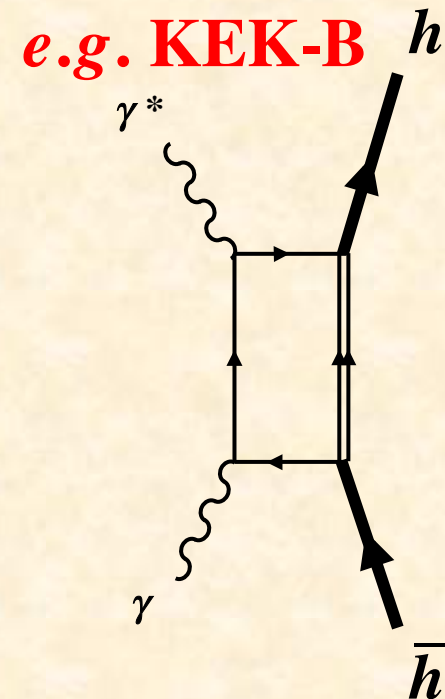
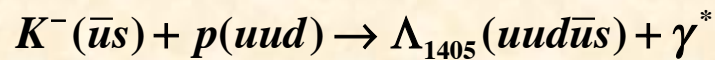
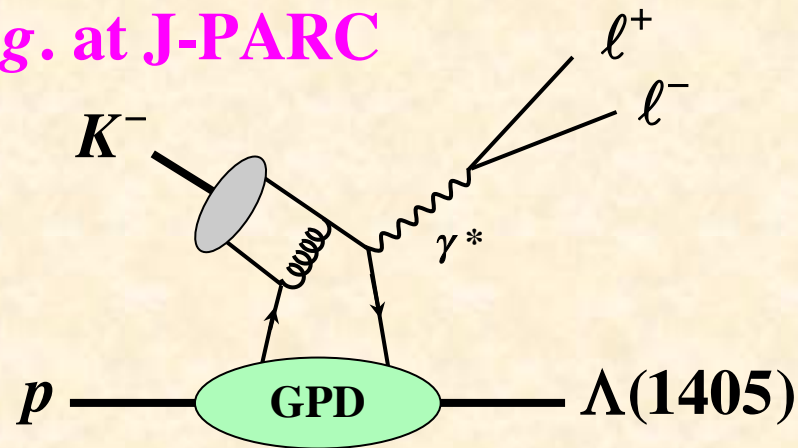
Because stable targets do not exist for exotic hadrons, it is not possible to measure their GPDs in a usual way.

→ Transition GPDs

or

→ $s \leftrightarrow t$ crossed quantity = GDAs at KEK-B, Linear Collider

e.g. at J-PARC



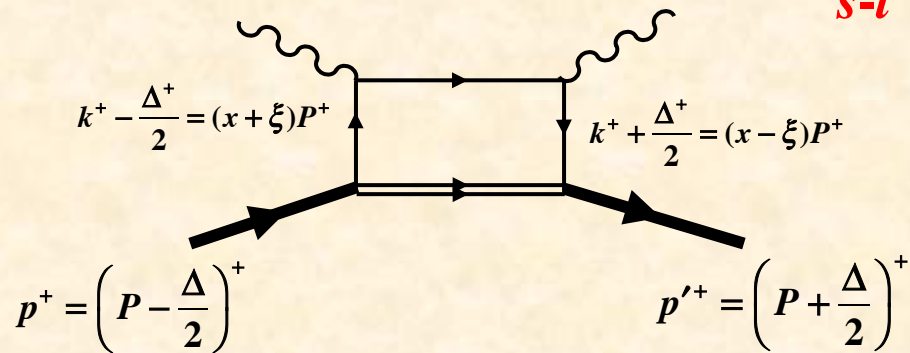
GPD $H_q^h(x, \xi, t)$ and GDA $\Phi_q^{h\bar{h}}(z, \zeta, W^2)$

$$\text{GPD: } H_q(x, \xi, t) = \int \frac{dy^-}{4\pi} e^{ixP^+y^-} \langle h(p') | \bar{\psi}(-y/2) \gamma^+ \psi(y/2) | h(p) \rangle \Big|_{y^+=0, \vec{y}_\perp=0}, \quad P^+ = \frac{(p+p')^+}{2}$$

$$\text{GDA: } \Phi_q(z, \zeta, s) = \int \frac{dy^-}{2\pi} e^{izP^+y^-} \langle h(p) \bar{h}(p') | \bar{\psi}(-y/2) \gamma^+ \psi(y/2) | 0 \rangle \Big|_{y^+=0, \vec{y}_\perp=0}$$

$$\text{DA: } \Phi_q^h(z, \zeta, s) = \int \frac{dy^-}{2\pi} e^{izP^+y^-} \langle h(p) | \bar{\psi}(-y/2) \gamma^+ \gamma_5 \psi(y/2) | 0 \rangle \Big|_{y^+=0, \vec{y}_\perp=0}$$

$H_q^h(x, \xi, t)$



$$P = \frac{p+p'}{2}, \quad \Delta = p' - p$$

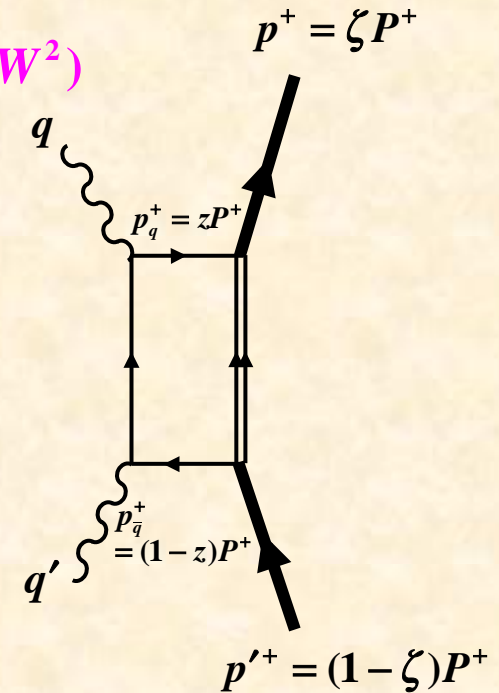
Bjorken variable: $x = \frac{Q^2}{2p \cdot q}$

Momentum transfer squared: $t = \Delta^2$

Skewness parameter: $\xi = \frac{p^+ - p'^+}{p^+ + p'^+} = -\frac{\Delta^+}{2P^+}$

\longleftrightarrow
s-t crossing

$\Phi_q^{h\bar{h}}(z, \zeta, W^2)$



Bjorken variable for γ^* : $z = \frac{Q^2}{2q \cdot q'}$

Light-cone momentum ratio for h in $h\bar{h}$: $\zeta = \frac{p^+}{P^+} = \frac{1 + \beta \cos \theta}{2}$

Invariant mass of $h\bar{h}$: $W^2 = (p+p')^2$

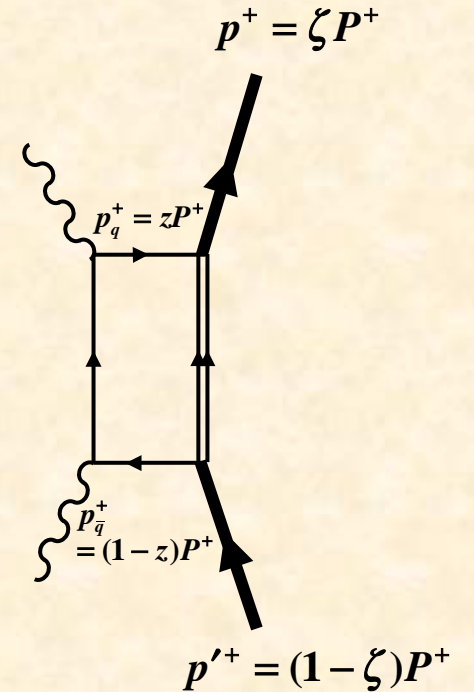
Functional form of GDAs $\Phi_q^{h\bar{h}}(z, \zeta, W^2)$

Kinematical variables

$$z = \frac{k^+}{P^+}, \quad \zeta = \frac{1}{2}(1 + \beta \cos \theta), \quad \beta = \sqrt{1 - \frac{4m_\pi^2}{W^2}}, \quad W^2 = (p + p')^2$$

Sum rule

$$\int_0^1 dz (2z - 1) \Phi_q^{h\bar{h}(I=0)}(z, \zeta, W^2) = -2M_{2(q)}^h \zeta (1 - \zeta) F_{h(q)}(W^2)$$



Exotic signatures should appear in the functional form (e.g. form factor).

GDA could be expressed by

$$\Phi_q^{h\bar{h}(I=0)}(z, \zeta, W^2) = N_h z^\alpha (1 - z)^\beta (2z - 1) \zeta (1 - \zeta) F_{h(q)}(s)$$

$$N_h = - \frac{2M_{2(q)}^h}{B(\alpha + 1, \beta + 1)} \frac{(\alpha + \beta + 3)(\alpha + \beta + 2)}{(\alpha + 2)(\alpha + 1) + (\beta + 2)(\beta + 1) - 2(\alpha + 1)(\beta + 1)}$$

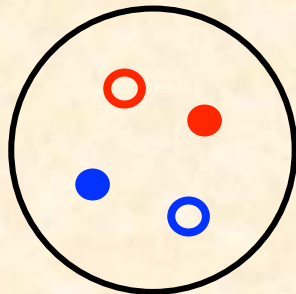
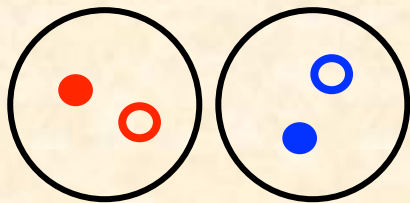
Cross section: form factor dependence

$$\Phi_q^{h\bar{h}(I=0)}(z, \zeta, W^2) = N_h z^\alpha (1-z)^\beta (2z-1)\zeta(1-\zeta) F_h(W^2)$$

Ordinal $q\bar{q}$



**Molecule $K\bar{K}$
or tetra-quark $qq\bar{q}\bar{q}$**



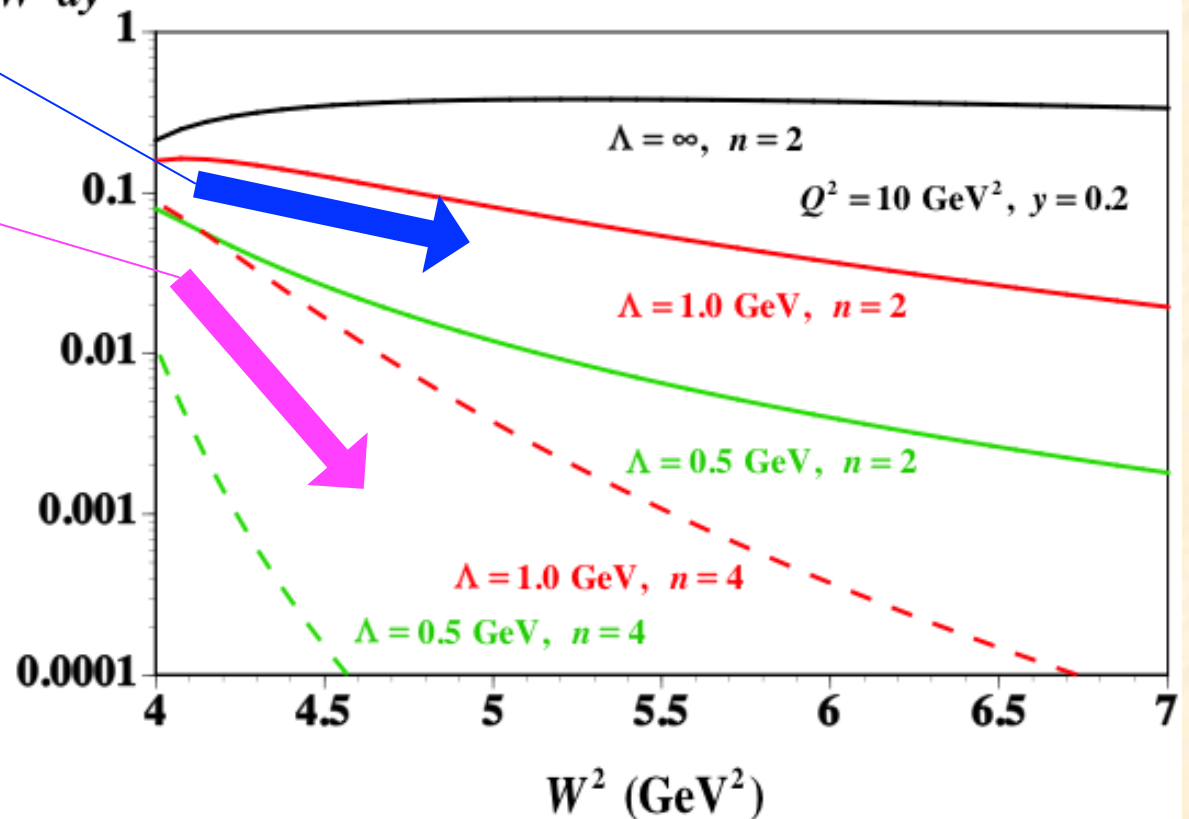
$$F_h(W^2) = \frac{1}{[1 + (W^2 - 4m_h^2) / \Lambda^2]^{n-1}}$$

Constituent-counting rule

$n = 2$: ordinary meson

$n = 4$: molecule or tetra-quark

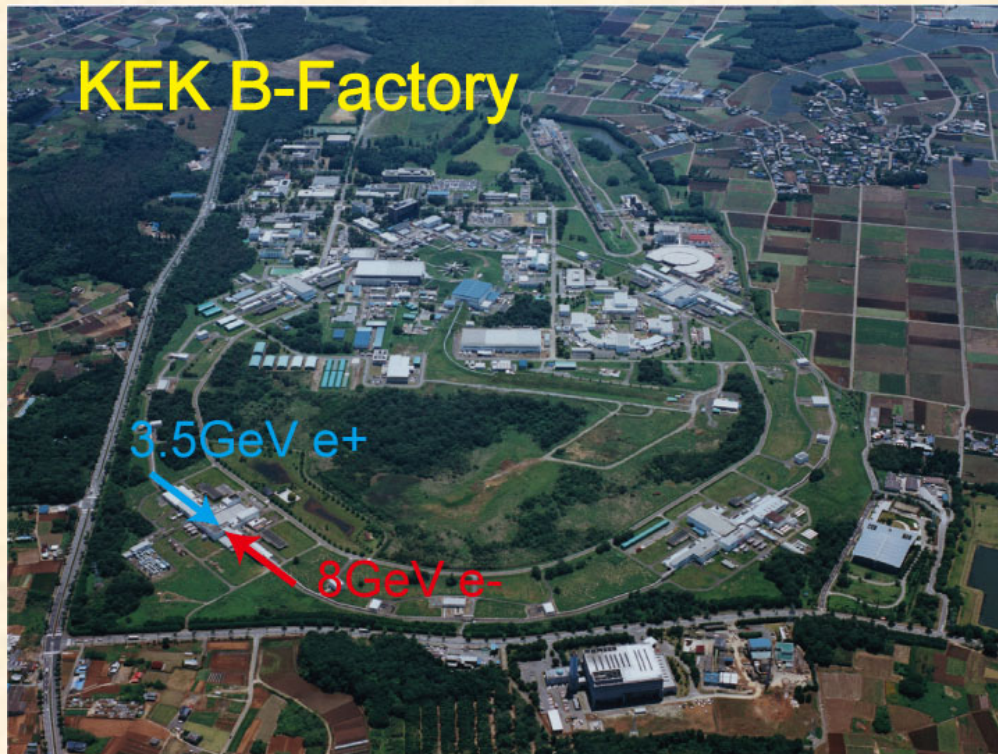
$$\frac{d\sigma_{ee \rightarrow eeMM}}{dQ^2 dW^2 dy} \text{ (fb / GeV}^4\text{)}$$



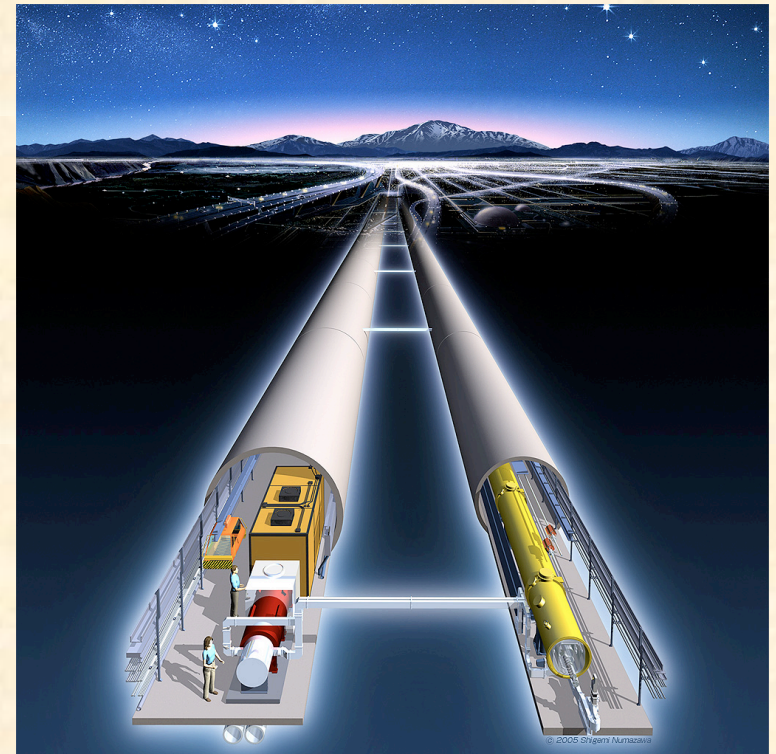
Experimental studies of GDAs in future

$\gamma\gamma \rightarrow h\bar{h}$ for internal structure of exotic hadron candidate h

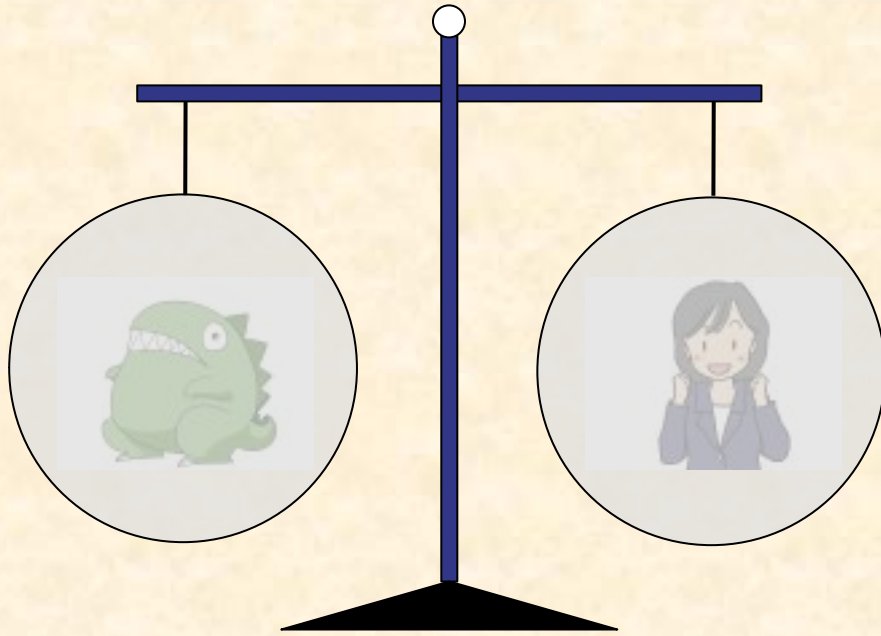
KEK B-factory



Linear Collider ?



Search for exotic hadrons ...



It is difficult to determine whether or not a hadron is exotic by low-energy observables, masses, decay widths, ...

(Already, history of a half century)



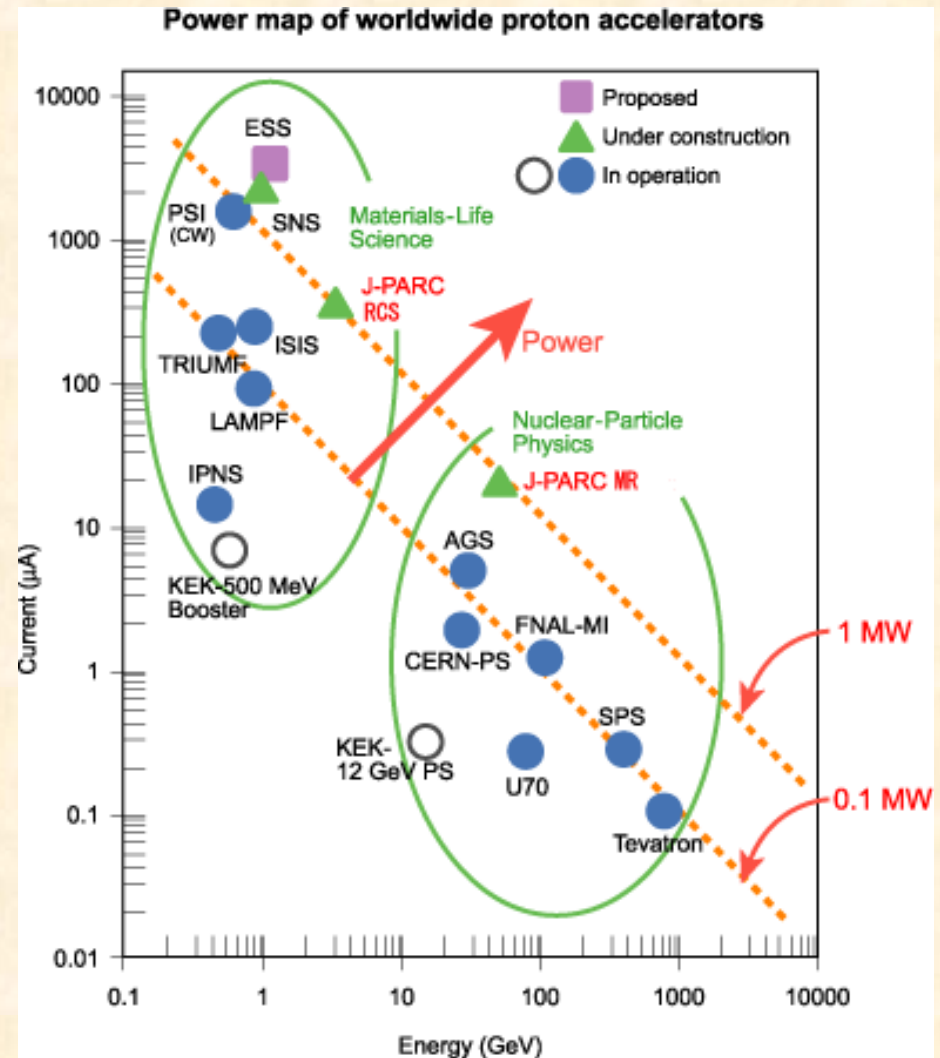
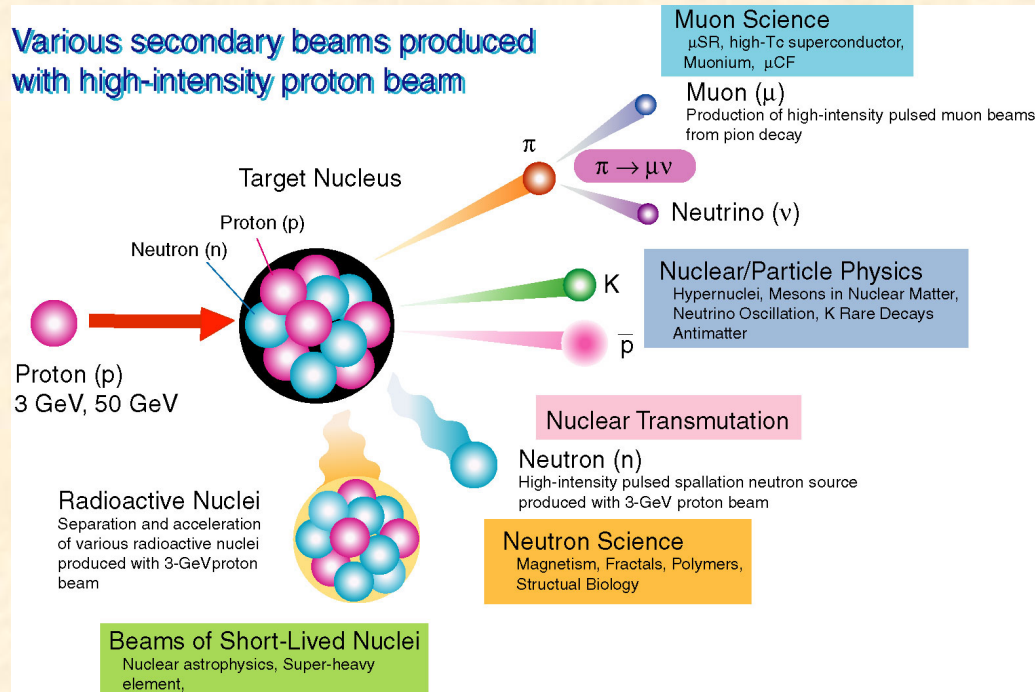
By the tomography, we may determine



**Comments on
J-PARC project**

High-Intensity Frontier of Proton Accelerator

High-intensity proton beam
 → High-intensity secondary beams
 (Neutrino, Kaon, Pion, Neutron ...)



Aerial photograph



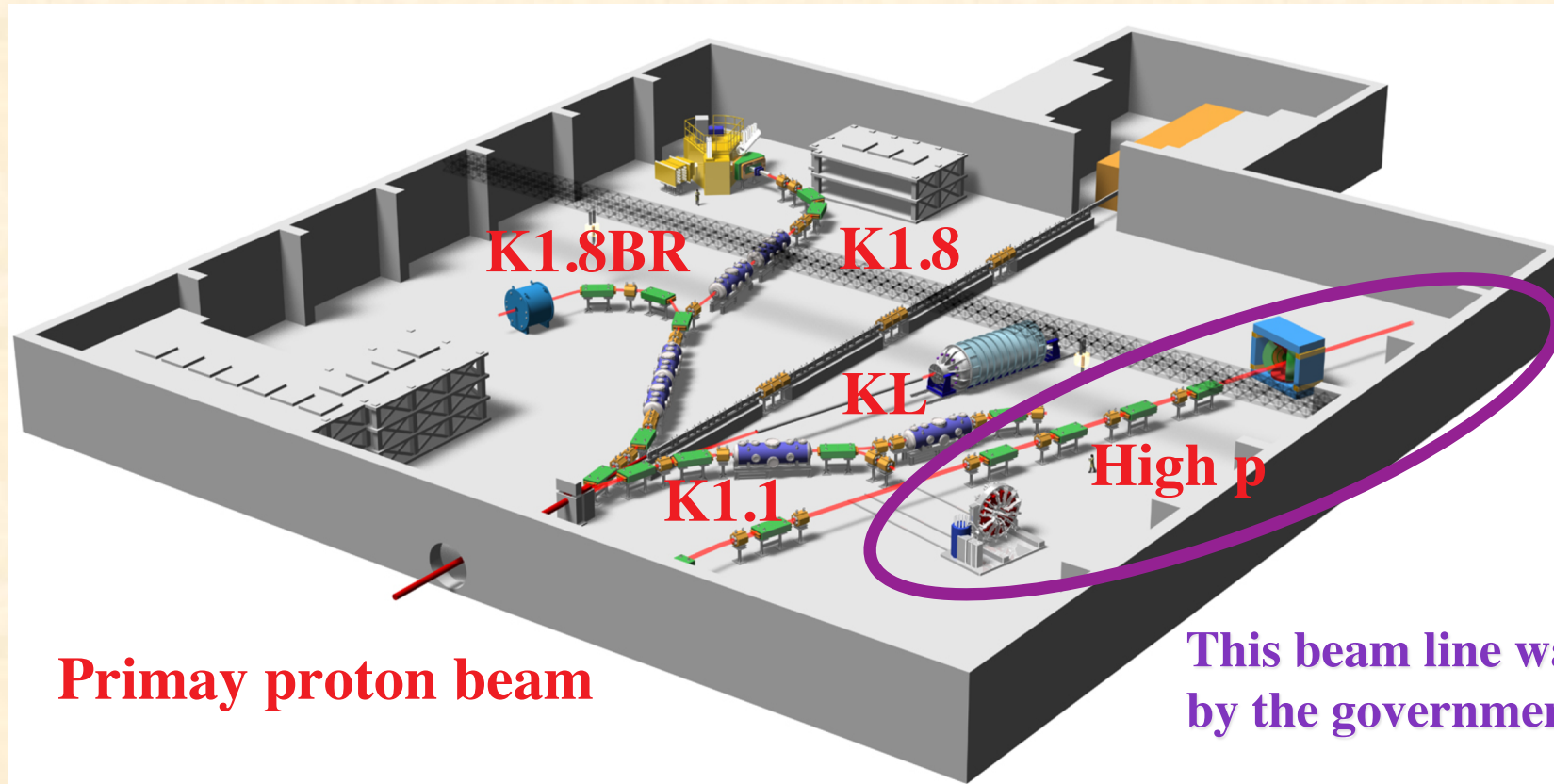
**KEK Tokai campus,
J-PARC branch
KEK theory center**

**Neutrino
facility**

**Hadron
facility**

Hadron facility

Recent workshop on high-momentum beamline physics,
January 15 - 18, 2013, KEK,
<http://www-conf.kek.jp/hadron1/j-parc-hm-2013/>



This beam line was approved
by the government in 2013.

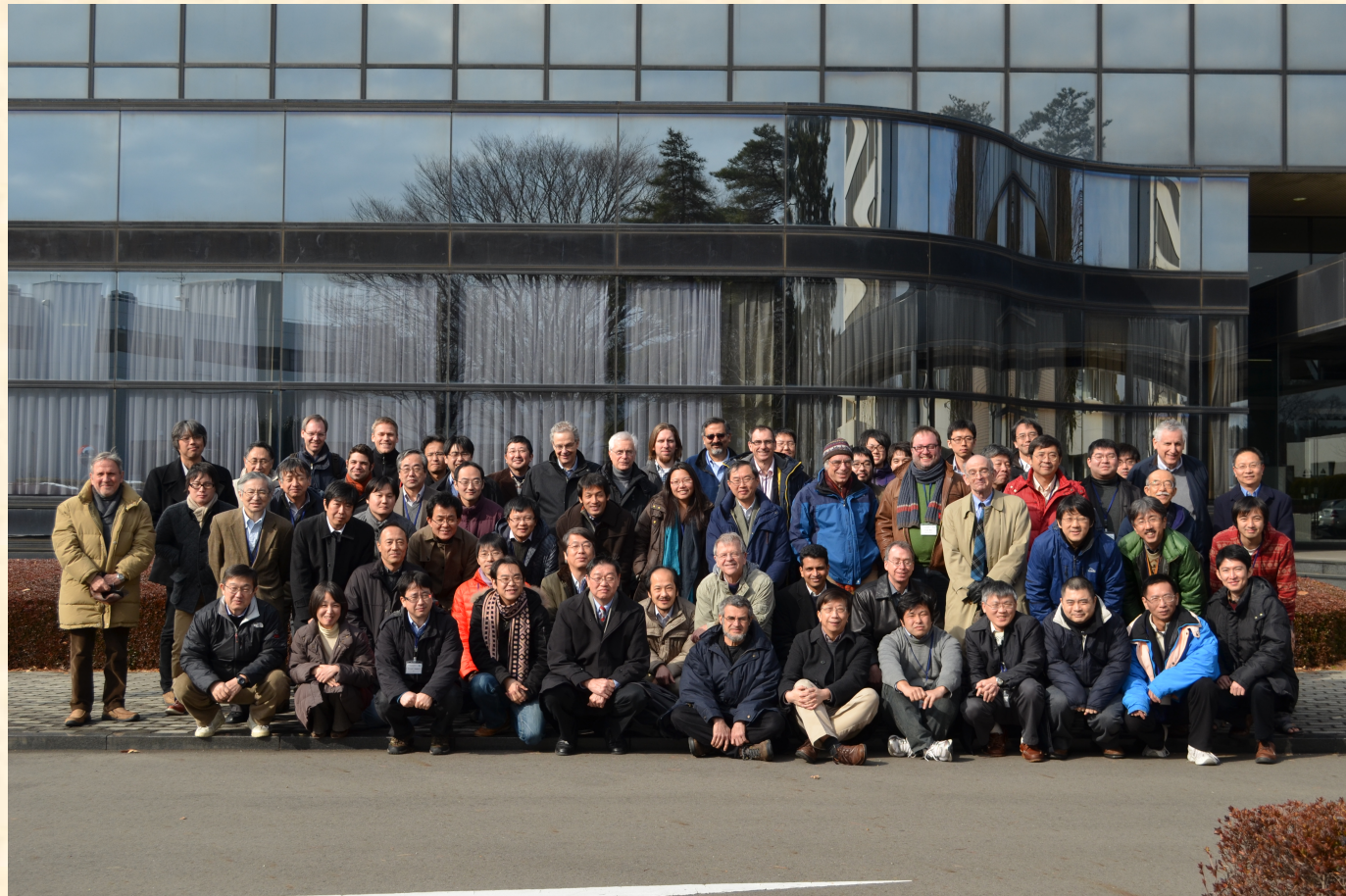
You may propose your experiments!

- Proton beam up to 30 GeV
- Unseparated hadron (pion, ...) beam up to 15~20 GeV

Hadron physics with high-momentum hadron beams at J-PARC in 2013

<http://www-conf.kek.jp/hadron1/j-parc-hm-2013/>

88 participants (~100 including non-registered ones)



Theory activities at J-PARC

J-PARC Branch, KEK Theory Center

Institute of Particle and Nuclear Studies, KEK
203-1, Shirakata, Tokai, Ibaraki, 319-1106, Japan
<http://j-parc-th.kek.jp>

4 permanent KEK staffs (A. Dote, K. Itakura, S. Kumano, O. Morimatsu)

+ 1 research fellow (T. Marruyama)

+ 5 visiting staffs (T. Harada, S. Hirenzaki, A. Hosaka, K. Saito, K. Tanaka)

Hyper-nuclear physics

Charm physics

Structure functions

Hadron masses in medium

Neutrino-nuclear interactions

If you are interested in organizing a workshop
or joining activities, please inform us.

Proposals on high-energy hadron physics

http://j-parc.jp/researcher/Hadron/en/Proposal_e.html

J. C. Peng, S. Sawada *et al.*

Proposal

Measurement of High-Mass Dimuon Production at the
50-GeV Proton Synchrotron

S. Choi *et al.*

Letter of Intent to J-PARC PAC
for

Study of Parton Distribution Function of
Mesons via Drell-Yan Process at J-PARC
at High-p beamline

Y. Goto *et al.*

Proposal

Polarized Proton Acceleration at J-PARC

**The high-momentum had not been approved financially until 2013,
so these proposals were deferred.**

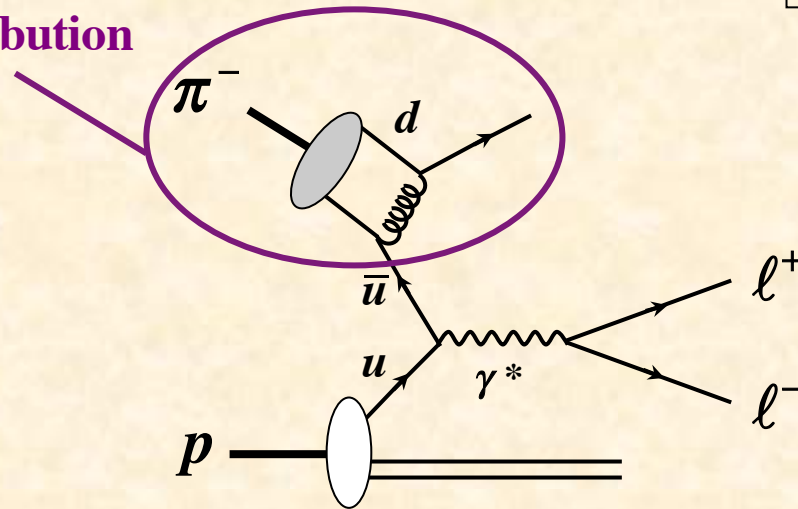
W.-C. Chang, J.-C. Peng, S. Sawada *et al.*,
possible J-PARC experiment?

**New LoI / proposal
under consideration!**

Toward a new proposal

W.-C. Chang, J.-C. Peng, S. Sawada *et al.*,
possible J-PARC experiment?

pion distribution



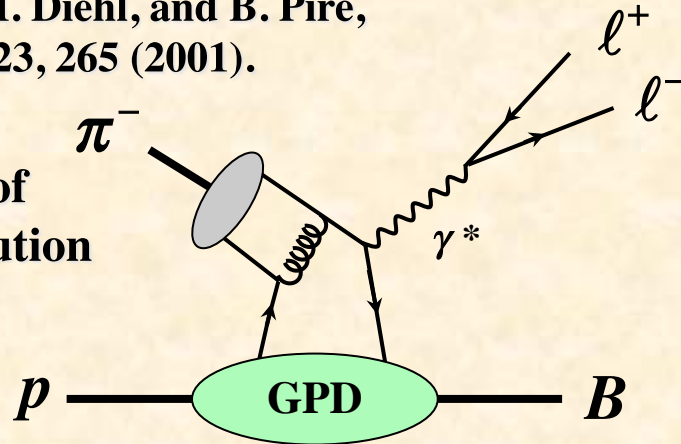
$$\pi^- (\bar{u}s) + p(uud) \rightarrow l^+ l^- + X$$

A. Brandenburg, S. J. Brodsky,
V. V. Khoze, and D. Müller,
Phys. Rev. Lett. 73 (1994) 939.

Investigation of
• Pion distribution amplitude

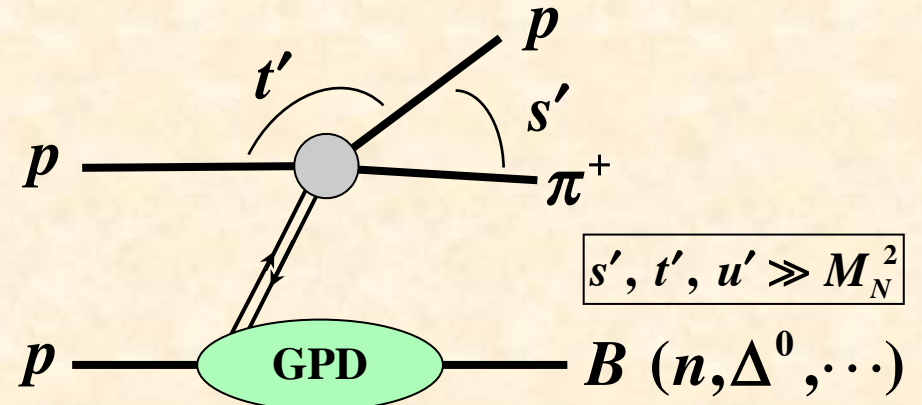
E. R. Berger, M. Diehl, and B. Pire,
Phys. Lett. B 523, 265 (2001).

Investigation of
• Pion distribution amplitude
• GPDs



$$\pi^- (\bar{u}d) + p(uud) \rightarrow B(udd) + \gamma^* (\rightarrow l^+ l^-)$$

SK, M. Strikman, K. Sudoh,
PRD 80 (2009) 074003



$$s', t', u' \gg M_N^2$$

$$B (n, \Delta^0, \dots)$$

In progress for LoI / proposal

**Wen-Chen Chang (Academia Sinica),
Jen-Chieh Peng (U. Illinois), ...**

See the slides of J-PARC workshops in 2014:

<http://research.kek.jp/people/kumanos/conf/conf14.html>

<http://j-parc-th.kek.jp/collabo/2014/02-13/hadron-sf-2014-02-13.html>

**Physics to be investigated in the high momentum beam line of
hadron hall at J-PARC**

Wen-Chen Chang

Institute of Physics, Academia Sinica, Taipei 11529, Taiwan

Hiroyuki Kawamura and Shunzo Kumano

KEK Theory Center, Institute of Particle and Nuclear Studies,

High Energy Accelerator Research Organization (KEK)

Jen-Chieh Peng

Department of Physics, University of Illinois at

Urbana-Champaign, Urbana, Illinois 61801, USA

Shin'ya Sawada

Institute of Particle and Nuclear Studies,

High Energy Accelerator Research Organization (KEK)

Summary

GPDs at hadron facilities

GPDs can be investigated by not only DVCS at lepton facilities but also exclusive reactions at hadron facilities.

Constituent-counting rule for exotic hadrons

High energies = Quark and gluon degrees of freedom

It is appropriate to use high-energy processes for determination of internal configurations for exotic-hadron candidates.

GPDs and GDAs for exotic hadrons

3D structure of hadrons can be investigated by GDAs ($s \leftrightarrow t$).

Related experimental projects

J-PARC, COMPASS, JLab, KEK-B, Liner Collider, ...

J-PARC: Propose your experiments, Propose joint workshops, ...

The End

The End