



# Electromagnetic probes of the QGP

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,Heavy Flavor and Electromagnetic Probes in Heavy Ion Collisions‘**

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# Electromagnetic probes: photons and dileptons

Feinberg (76), Shuryak (78)

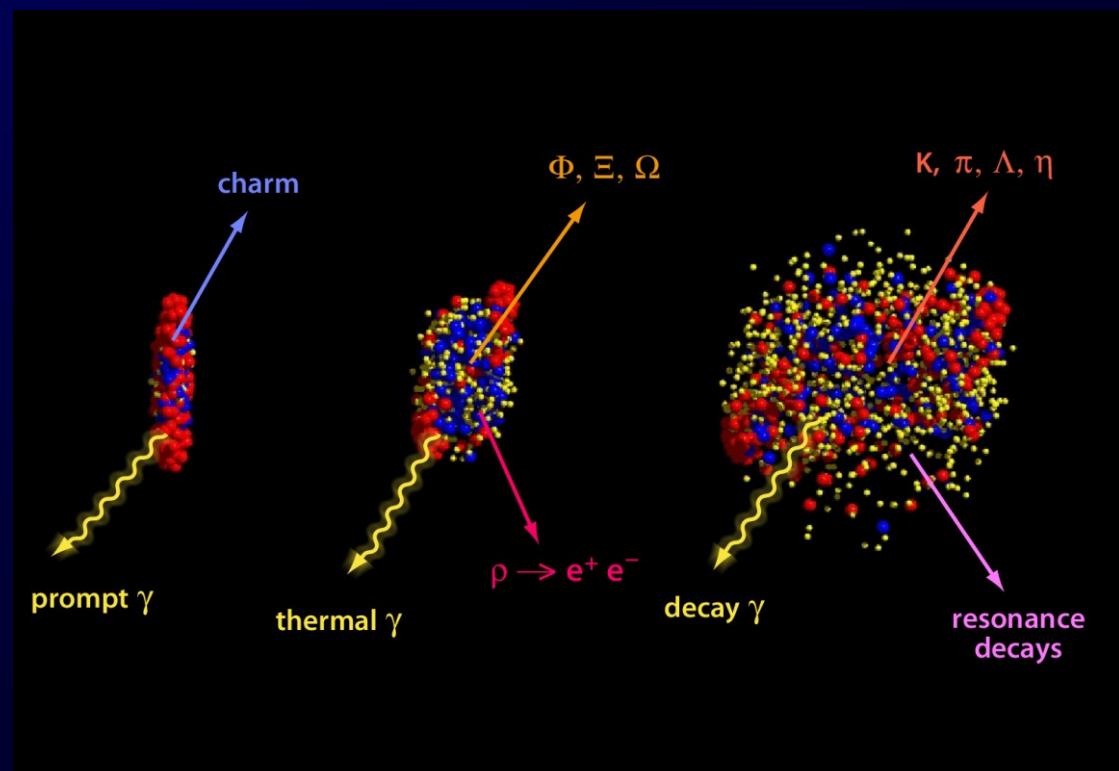
## Advantages:

- ✓ dileptons and real photons are emitted from different stages of the reaction and not effected by final-state interactions
- ✓ provide undistorted information about their production channels
- ✓ promising signal of QGP – ‘thermal’ photons and dileptons

→ Requires theoretical models which describe the dynamics of heavy-ion collisions during the whole time evolution!

## Disadvantages:

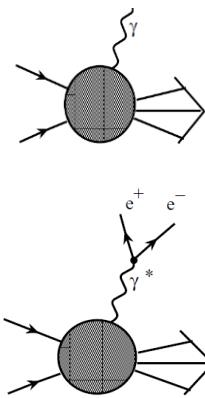
- low emission rate
- production from hadronic corona
- many production sources which cannot be individually disentangled by experimental data



# Modeling of photon/dilepton emission

## I. Emission rate from thermal field theory:

Feinberg (76), McLerran, Toimela (85),  
Weldon (90), Gale, Kapusta (91)



- **Photons:**  $q_0 \frac{d^3 R}{d^3 q} = -\frac{g_{\mu\nu}}{(2\pi)^3} \text{Im } \underline{\Pi^{\mu\nu}(q_0 = |\vec{q}|)} f(q_0, T)$
- **Dileptons:**  $E_+ E_- \frac{d^3 R}{d^3 p_+ d^3 p_-} = \frac{2e^2}{(2\pi)^6} \frac{1}{q^4} L_{\mu\nu} \text{Im } \underline{\Pi^{\mu\nu}(q_0, \vec{q})} f(q_0, T)$
- **Bose distribution:**  
 $f(q_0, T) = \frac{1}{e^{q_0/T} - 1}$
- $L_{\mu\nu}$  is the electromagnetic leptonic tensor
- $\Pi_{\mu\nu}$  is the retarded photon self energy at finite T :  $\Pi_{\mu\nu} \sim i \int d^4 x e^{ipx} \langle [J_\mu(x), J_\nu(0)] \rangle_T$

- Hadron phase: using VDM:  $\text{Im}\Pi \sim \text{Im}D^\rho$  in-medium  $\rho$ -meson spectral function from many-body approach ( cf. Rapp, Chanfrey, Wambach, NPA 617 (1997) 472 )

→ study of the in-medium properties of hadrons at high baryon density and T

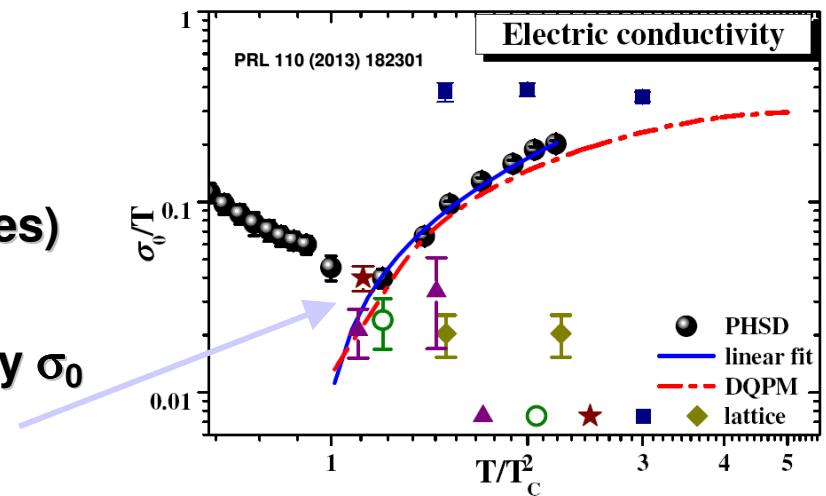
→ restoration of chiral symmetry ( $\rho$ -a<sub>1</sub>):

$\text{Im}D^\rho \sim$  chiral condensate (by Weinberg sum rules)

(cf. Hohler, Rapp, arXiv:1311.2921)

- Rates at  $q_0 \rightarrow 0$  are related to electric conductivity  $\sigma_0$
- Probe of electric properties of the QGP

$$q_0 \frac{dR}{d^4 x d^3 q} \Big|_{q_0 \rightarrow 0} = \frac{T}{4\pi^3} \sigma_0$$



PHSD plot from Cassing et al., PRL 110 (2013) 182301;  
cf. also NJL: Marty et al., PRC87 (2013) 3, 034912

# Modeling of photon/dilepton emission

## II. Emission rate from relativistic kinetic theory:

(e.g. for  $1+2 \rightarrow \gamma+3$ )

Applicable also for  
non-equilibrium  
system !

$$q_0 \frac{d^3 R}{d^3 q} = \int \frac{d^3 p_1}{2(2\pi)^3 E_1} \frac{d^3 p_2}{2(2\pi)^3 E_2} \frac{d^3 p_3}{2(2\pi)^3 E_3} (2\pi)^4 \delta^4(p_1 + p_2 - p_3 - q) \\ \times |M|^2 \frac{f(E_1)f(E_2)[1 \pm f(E_3)]}{2(2\pi)^3}$$

■  $f(E)$  - distribution function

- $M$  – invariant scattering matrix element from microscopic models
- Modeling of hadronic elementary reactions:  
Chiral models, OBE models,... (Born-type diagrams)
- Problems:
  - very limited experimental information on mm, mB elementary reactions
  - Hadrons change their properties in the hot and dense medium:  
→ from vacuum cross sections to in-medium, i.e.  
from 'T-matrix' to 'G-matrix' approaches (many-body theory)

E.g. :  $\rho$ -meson collisional broadening – important for dilepton studies!

# Production sources of photons in p+p and A+A

## □ Decay photons (in pp and AA):

$$m \rightarrow \gamma + X, m = \pi^0, \eta, \omega, \eta', a_1, \dots$$

## □ Direct photons: (inclusive(=total) – decay) – measured experimentally

### ■ hard photons:

(large  $p_T$ ,  
in pp and AA)

- prompt (pQCD; initial hard N+N scattering)

- jet fragmentation (pQCD; qq, gq bremsstrahlung)  
(in AA can be modified by parton energy loss in medium)

### ■ thermal photons:

(low  $p_T$ , in AA)

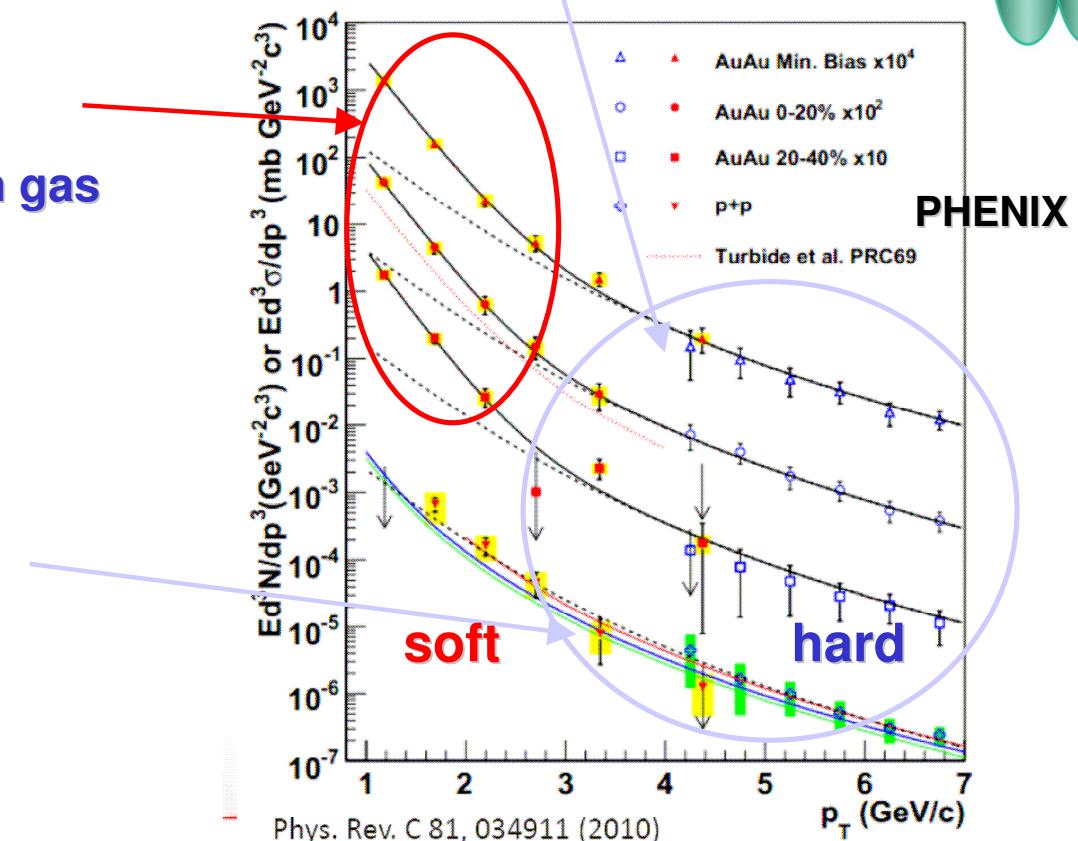
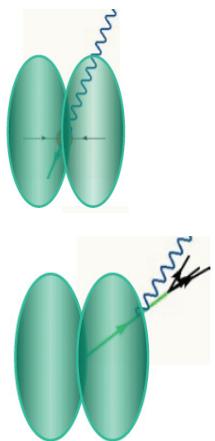
- QGP
- Hadron gas

### ■ jet- $\gamma$ -conversion in plasma

(large  $p_T$ , in AA)

### ■ jet-medium photons

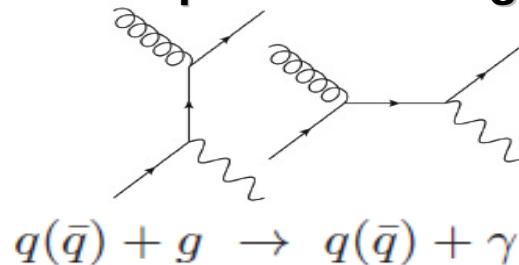
(large  $p_T$ , in AA) - scattering of  
hard partons with thermalized  
partons  $q_{\text{hard}} + g_{\text{QGP}} \rightarrow \gamma + q$ ,  
 $q_{\text{hard}} + q\bar{q}_{\text{QGP}} \rightarrow \gamma + q$



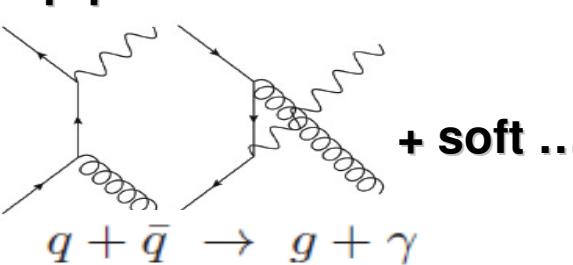
# Production sources of thermal photons

## □ Thermal QGP:

### Compton scattering



### q-qbar annihilation



HTL program (Klimov (1981), Weldon (1982),  
Braaten & Pisarski (1990); Frenkel & Taylor (1990), ...)

- Rates beyond pQCD:  
off-shell massive  $q, g$   
(used in PHSD)

O. Linnyk, JPG 38 (2011) 025105

- pQCD LO: ‘AMY’ Arnold, Moore, Yaffe, JHEP 12, 009 (2001)
- pQCD NLO: Gale, Ghiglieri (2014)

← QGP rates used in hydro !

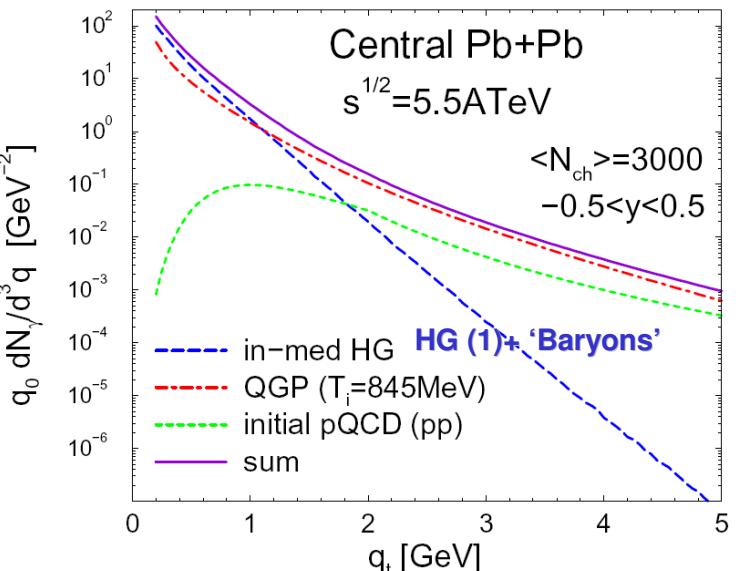
## □ Hadronic sources:

### (1) secondary mesonic interactions:

$$\pi + \pi \rightarrow \rho + \gamma, \rho + \pi \rightarrow \pi + \gamma, \pi + K \rightarrow \rho + \gamma, \dots$$

### (2) meson-meson and meson-baryon bremsstrahlung:

$$m + m \rightarrow m + m + \gamma, \quad m + B \rightarrow m + B + \gamma, \\ m = \pi, \eta, \rho, \omega, K, K^*, \dots, \quad B = p, \Delta, \dots$$



Models: chiral models, OBE, SPA ...  
Kapusta, Gale, Haglin (91), Rapp (07), ...

HG rates (1) used in hydro (‘TRG’ model) -  
massive Yang-Mills approach:  
Turbide, Rapp, Gale, PRC 69, 014903 (2004)

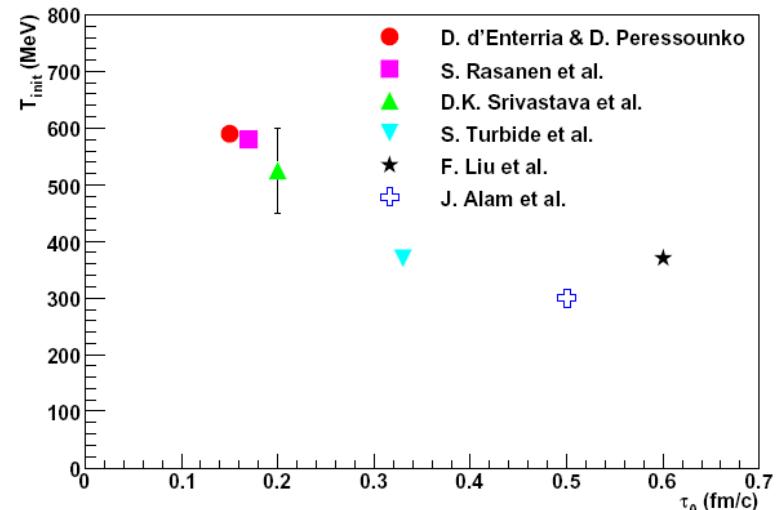
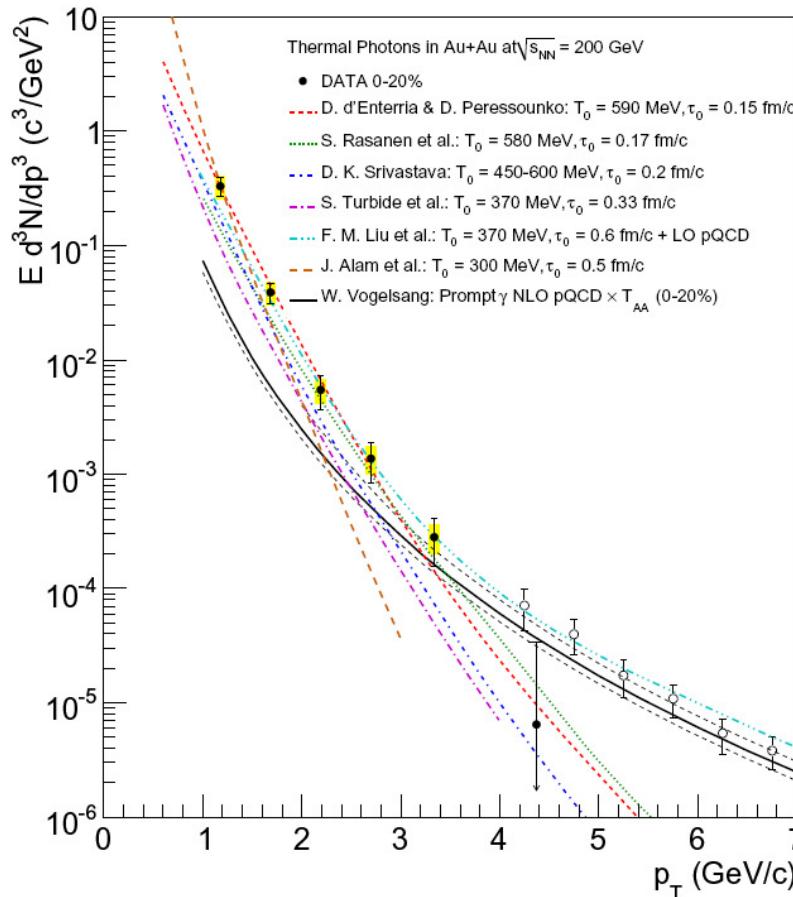
# 2010: Direct photon spectra for Au+Au at $s^{1/2}=200$ GeV

PHENIX, Phys. Rev. C81 (2010) 034911

## Variety of model predictions:

fireball, 2+1 Bjorken hydro, 3+1 ideal hydro  
with different initial conditions and EoS

Models: assume formation of a hot QGP with initial temperature  $T_{\text{init}}$  at thermalization time  $\tau_0$

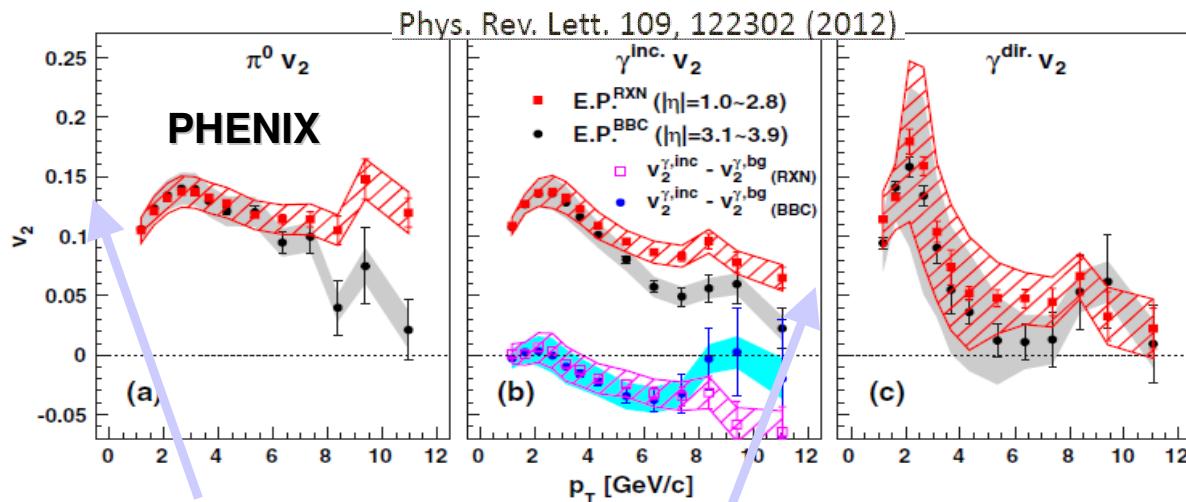


→ Huge variations in  $T_{\text{init}}$  and  $\tau_0$ !

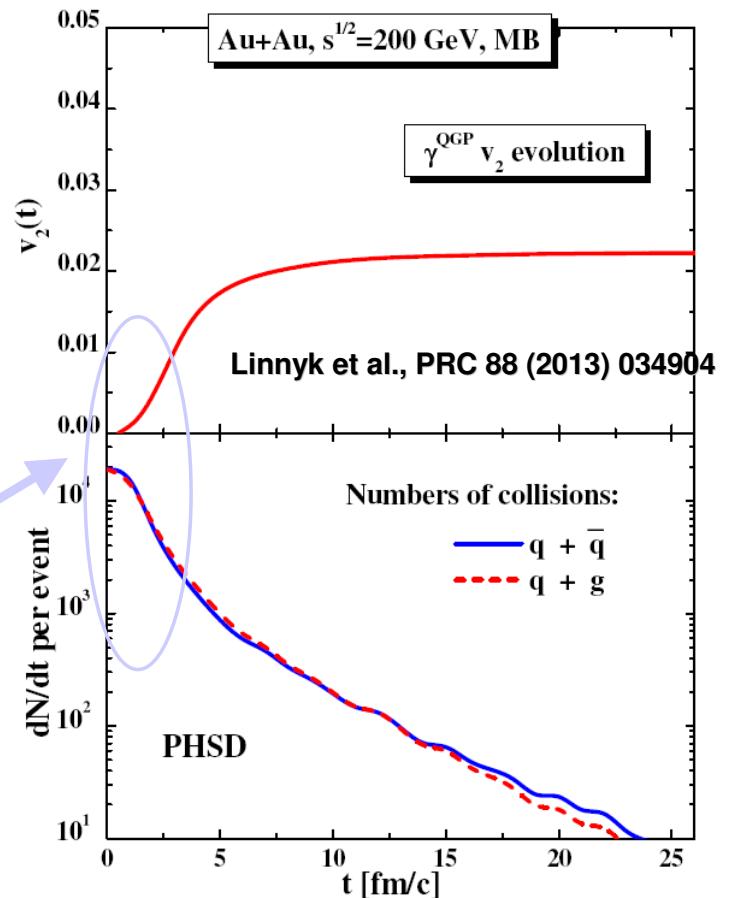
Warning: some model evolution was not fitted to the final hadron spectra!

Photon spectra show sensitivity to the dynamical evolution!

# PHENIX: Photon $v_2$ puzzle



$$\frac{dN}{d\phi} = \frac{1}{2\pi} \left( 1 + 2 \sum_{n \geq 1} v_n \cos(n(\phi - \Psi_n^{RP})) \right)$$



- PHENIX (also now ALICE):  
strong elliptic flow of photons  $v_2(\gamma^{\text{dir}}) \sim v_2(\pi)$
- Result from a variety of models:  $v_2(\gamma^{\text{dir}}) \ll v_2(\pi)$
- Problem: QGP radiation occurs at early times when flow is not yet developed → expected  $v_2(\gamma^{\text{QGP}}) \rightarrow 0$
- $v_2 = \text{weighted average } v_2 = \frac{\sum_i N^i \cdot v_2^i}{\sum_i N^i} \rightarrow \text{a large QGP contribution gives small } v_2(\gamma^{\text{QGP}})$
- NEW (QM'2014): PHENIX, ALICE experiments - large photon  $v_3$  !

Challenge for theory – to describe spectra,  $v_2$ ,  $v_3$  simultaneously !

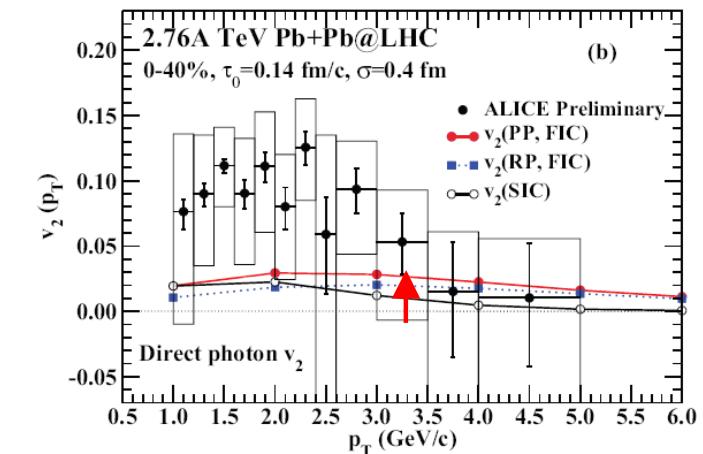
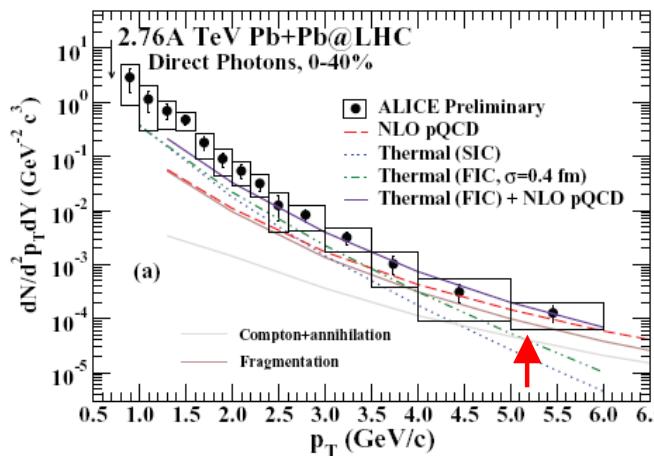
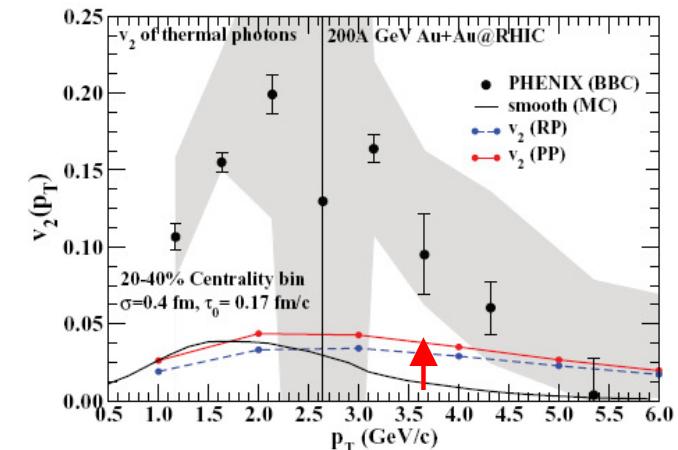
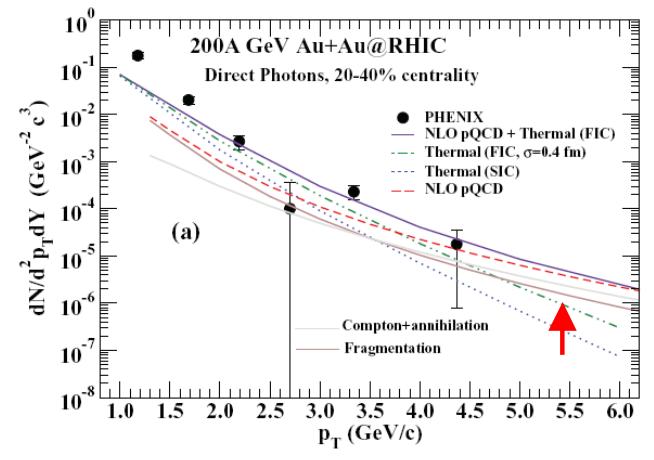
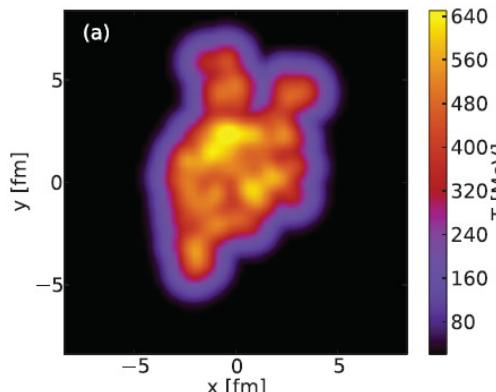
# 1. Hydro: Influence of e-b-e fluctuating initial conditions

→ From smooth Glauber initial conditions  
to event-by-event hydro with fluctuating initial conditions

Jyväskylä  
ideal hydro

- Ideal QGP and HG fluid
- Initial: „bumpy“ ebe
- MC Glauber
- EoS: IQCD

R. Chatterjee et al.,  
PRC 88, 034901 (2013)



→ Fluctuating initial conditions: slight increase at high  $p_T$  for yield and  $v_2$   
small effect, right direction!

## 2. From ideal to viscous hydro: direct photons as a QGP viscometer?

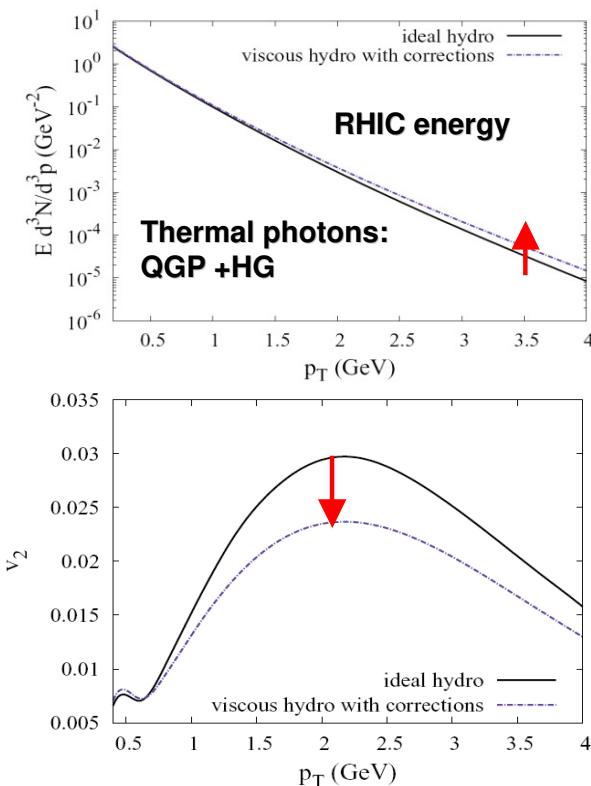
## The thermal photon emission rates with viscous corrections:

$$q \frac{dR}{d^3q}(q, T) = \Gamma_0(q, T) + \frac{\pi^{\mu\nu}}{2(e+P)} \underline{\Gamma_{\mu\nu}(q, T)},$$

**equilibrium contribution**      **first order viscous correction**

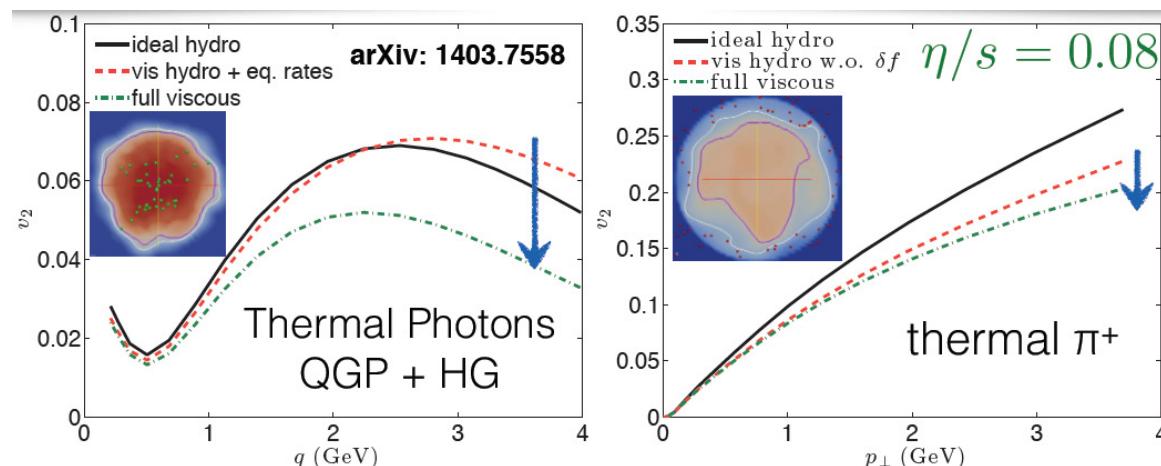
□ (3+1)D MUSIC (McGill):  
M. Dion et al., PRC84 (2011) 064901

- viscous QGP and HG fluid
  - Initial: ‚bumpy‘ ebe from IP-Glasma
  - EoS: IQCD



□ (2+1)D VISH2+1 (Ohio State) :  
C. Shen et al., arXiv:1308.2111, arXiv:1403.7558

- viscous QGP and HG fluid
  - Initial: ‘bumpy’ ebe from MC Glauber /KLN
  - EoS: IQCD



## → Effect of shear viscosity:

- \* small enhancement of the photon yield
  - \* suppression of photon  $v_2$
  - \* effect on  $v_2$  for photons is stronger than for hadrons

### 3. Influence of Glasma initial conditions with initial flow

#### □ (3+1)D MUSIC - 2014:

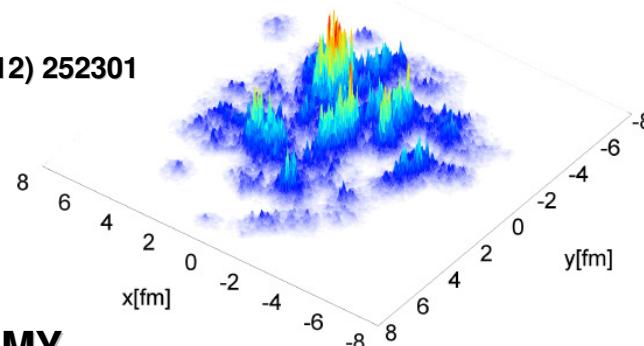
J-F. Paquet et al. (2014)

- viscous QGP and HG fluid ( $\eta/s=0.22$ )

- Initial: 'bumpy' ebe from IP-Glasma → generate initial flow due to fluctuations of IC

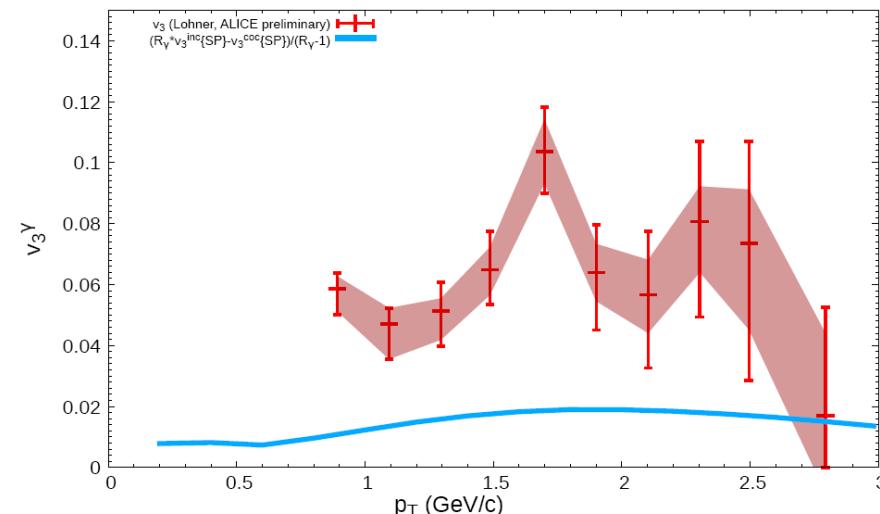
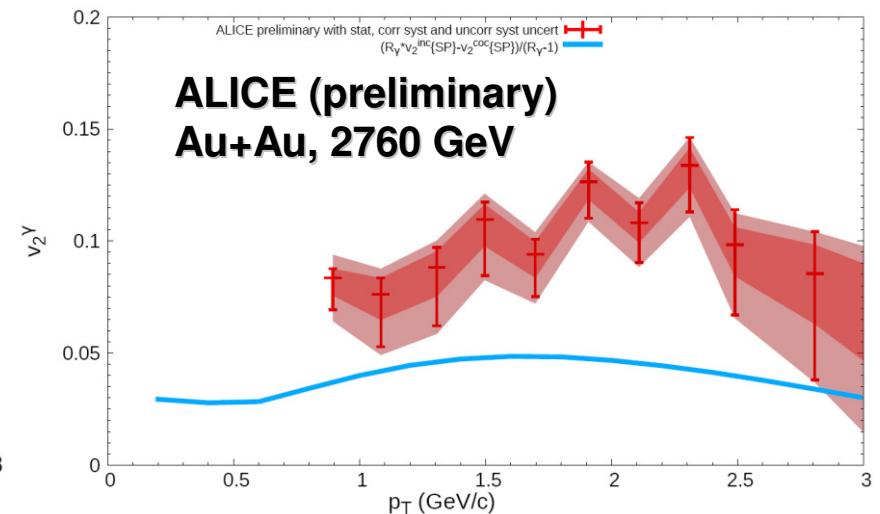
IP-Glasma:

Schenke et al., PRL108 (2012) 252301



- EoS: IQCD
- QGP photon rate: AMY
- HG photon rate: TGR for meson gas with viscous corrections + Rapp spectral function for p-mesons to account for the baryonic contributions

- MUSIC with IC-Glasma describes hadronic flow  $v_n$  systematics at RHIC & LHC, however, missing  $v_2, v_3$  of photons!



→ ,Bumpy' ebe from IP-Glasma - small effect

## 4. Hydro with pre-equilibrium flow

□ **Initial' flow:** rapid increase in bulk  $v_2$  in fireball model

van Hees, Gale, Rapp, PRC84 (2011) 054906

□ **pre-equilibrium flow in (2+1)D VISH2+1 - 2014:**

C. Shen et al., arXiv:1308.2111, arXiv:1403.7558; Talk by C. Shen @ QM'2014

- viscous QGP and HG fluid ( $\eta/s=0.18$ )
- Initial: 'bumpy' ebe from MC Glauber /KLN
- EoS: IQCD
- QGP photon rate: AMY
- HG photon rate: TGR for meson gas with viscous corrections

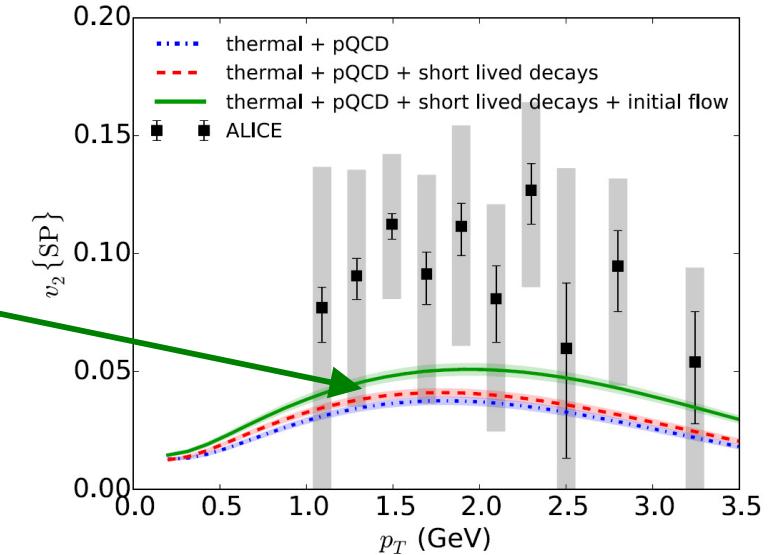
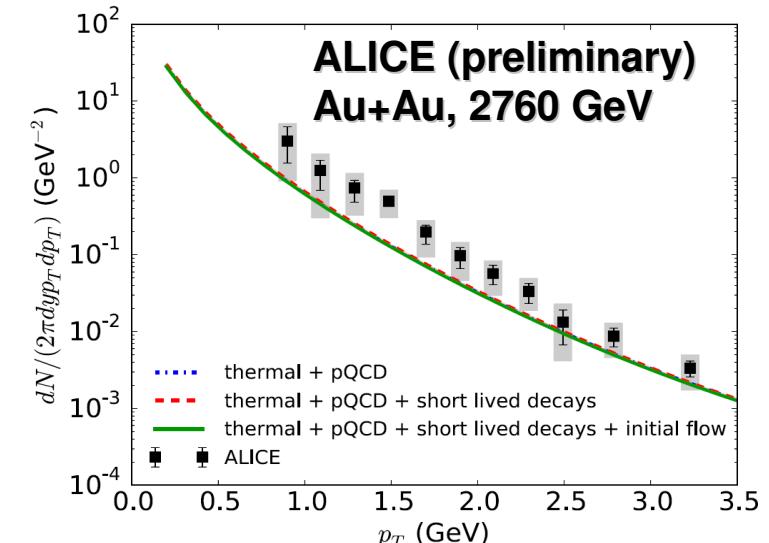
• Generation of **pre-equilibrium flow**:  
 using **free-streaming model** to evolve the partons  
 right after the collisions to  $0.6 \text{ fm/c}$   
 + Landau matching to switch to viscous hydro

→ quick development of momentum anisotropy  
 with saturation near  $T_c$



→ **Pre-equilibrium flow:**

- small effect on photon spectra
- slight **increase of  $v_2$**



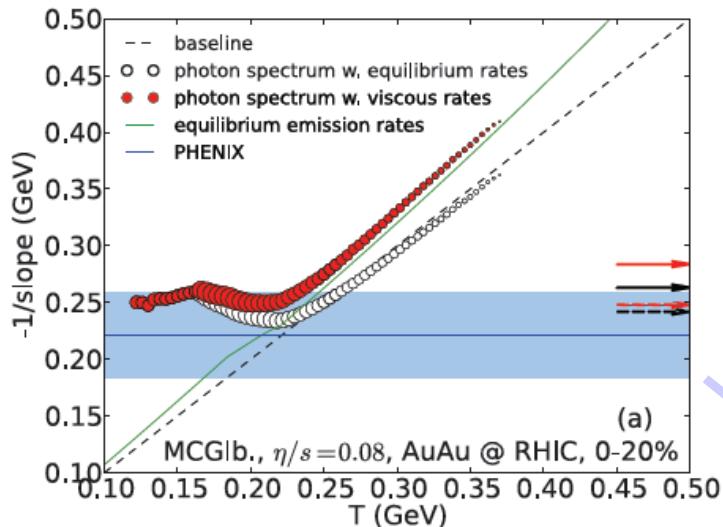
Warning: results can be considered as  
**upper limit** for the pre-equilibrium flow effect!

# Are thermal photons a QGP thermometer?



## □ (2+1)d viscous hydro VISH2+1 (Ohio)

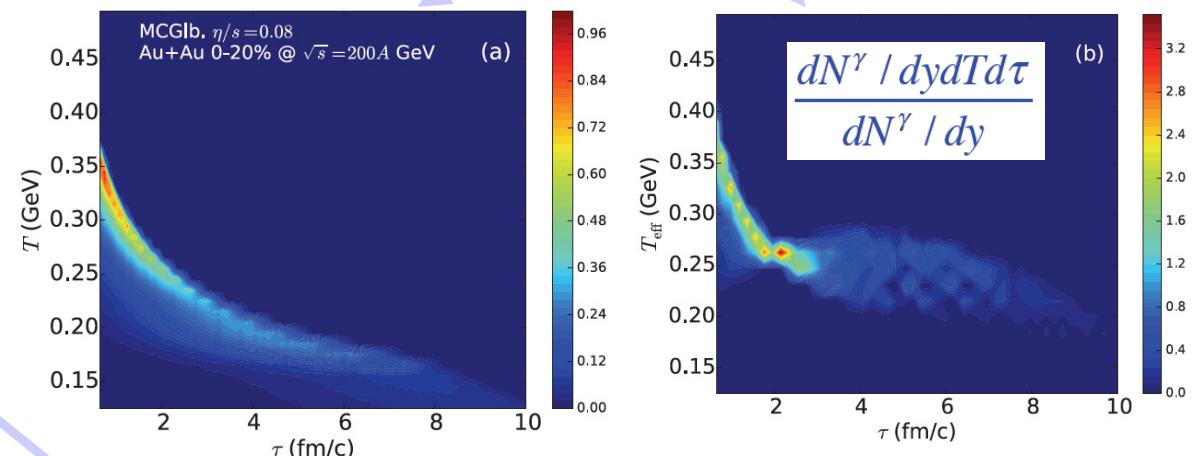
- Time evolution of the effective temperature
- $T_{\text{eff}} = -1/\text{slope}$  vs. local fluid cell temperature  $T$



Range of photon emission	Fraction of total photon yield	
	AuAu@RHIC 0–20% centr.	PbPb@LHC 0–40% centr.
$T = 120\text{--}165 \text{ MeV}$	17%	15%
$T = 165\text{--}250 \text{ MeV}$	62%	53%
$T > 250 \text{ MeV}$	21%	32%
$\tau = 0.6\text{--}2.0 \text{ fm}/c$	28.5%	26%
$\tau > 2.0 \text{ fm}/c$	71.5%	74%

C. Shen et al., PRC89 (2014) 044910; arXiv:1308.2440

- Contour plots of differential photon yield vs. time and temperature  $T$  and  $T_{\text{eff}}$ :



## Exp. Data:

- RHIC:  $T_{\text{eff}} = 221 + 19 + 19 \text{ MeV}$
- LHC:  $T_{\text{eff}} = 304 + 51 \text{ MeV}$

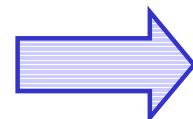
- Measured  $T_{\text{eff}} >$  'true'  $T$  
$$T_{\text{eff}} = \sqrt{\frac{1+v}{1-v}} T$$
- 'blue shift' due to the radial flow!
- only ~1/3 at LHC and ~1/4 at RHIC of total photons come from hot QCD ( $T > 250 \text{ MeV}$ )

## ❑ Further improvements of hydro models ?

- Bulk viscosity
- Modeling of initial pre-equilibrium effects
- ...



- Non-equilibrium dynamics ?
- Missing strength related to hadronic stage?

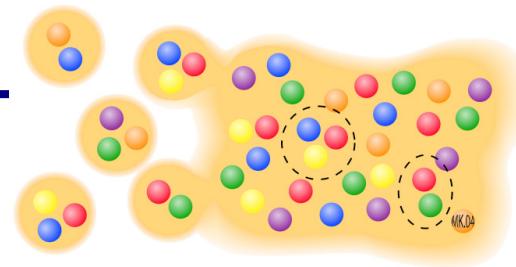


From hydro to non-equilibrium  
microscopic transport models :

use PHSD as a 'laboratory' for that



# From hadrons to partons



In order to study the **phase transition** from hadronic to partonic matter – **Quark-Gluon-Plasma** – we need a consistent non-equilibrium (transport) model with

- explicit **parton-parton interactions** (i.e. between quarks and gluons) beyond strings!
- explicit **phase transition** from hadronic to partonic degrees of freedom
- IQCD EoS for partonic phase

**Transport theory:** off-shell Kadanoff-Baym equations for the Green-functions  $S_{<h}(x,p)$  in phase-space representation for the **partonic** and **hadronic phase** (applicable for strongly interacting system!)



## Parton-Hadron-String-Dynamics (PHSD)

QGP phase described by

### Dynamical QuasiParticle Model (DQPM)

W. Cassing, E. Bratkovskaya, PRC 78 (2008) 034919;  
NPA831 (2009) 215;

W. Cassing, EPJ ST 168 (2009) 3

A. Peshier, W. Cassing, PRL 94 (2005) 172301;  
Cassing, NPA 791 (2007) 365: NPA 793 (2007)

# Dynamical QuasiParticle Model (DQPM) - Basic ideas:

DQPM describes QCD properties in terms of 'resummed' single-particle Green's functions – in the sense of a two-particle irreducible (2PI) approach:

$$\text{Gluon propagator: } \Delta^{-1} = P^2 - \Pi$$

$$\text{gluon self-energy: } \Pi = M_g^2 - i2\Gamma_g\omega$$

$$\text{Quark propagator: } S_q^{-1} = P^2 - \Sigma_q$$

$$\text{quark self-energy: } \Sigma_q = M_q^2 - i2\Gamma_q\omega$$

- the resummed properties are specified by complex self-energies which depend on temperature:
  - the real part of self-energies ( $\Sigma_q, \Pi$ ) describes a dynamically generated mass ( $M_q, M_g$ );
  - the imaginary part describes the interaction width of partons ( $\Gamma_q, \Gamma_g$ )
- space-like part of energy-momentum tensor  $T_{\mu\nu}$  defines the potential energy density and the mean-field potential (1PI) for quarks and gluons ( $U_q, U_g$ )
- 2PI framework guarantees a consistent description of the system in- and out-of equilibrium on the basis of Kadanoff-Baym equations

# The Dynamical QuasiParticle Model (DQPM)

- Basic idea: interacting quasi-particles: massive quarks and gluons ( $g, q, \bar{q}$ ) with Lorentzian spectral functions :

$(i = q, \bar{q}, g)$

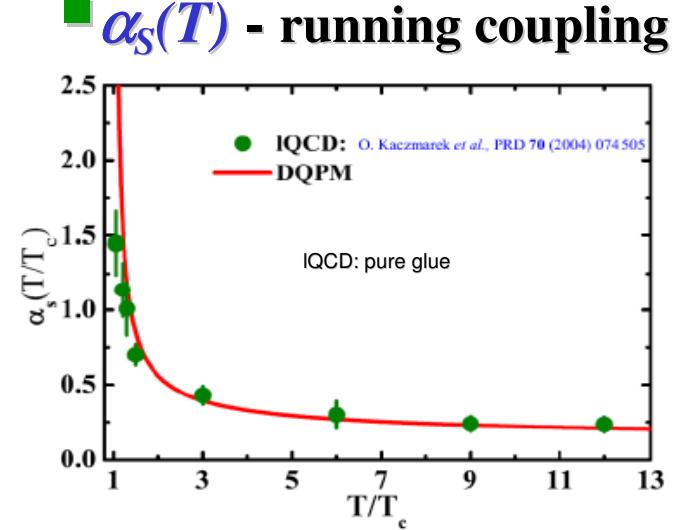
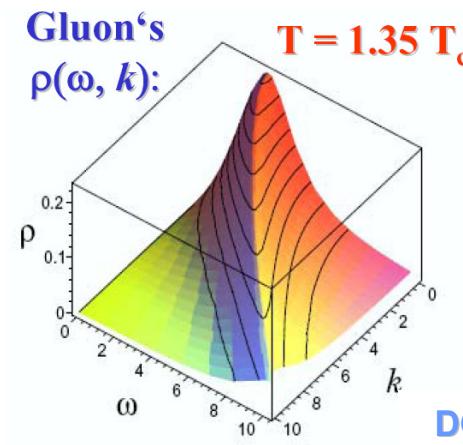
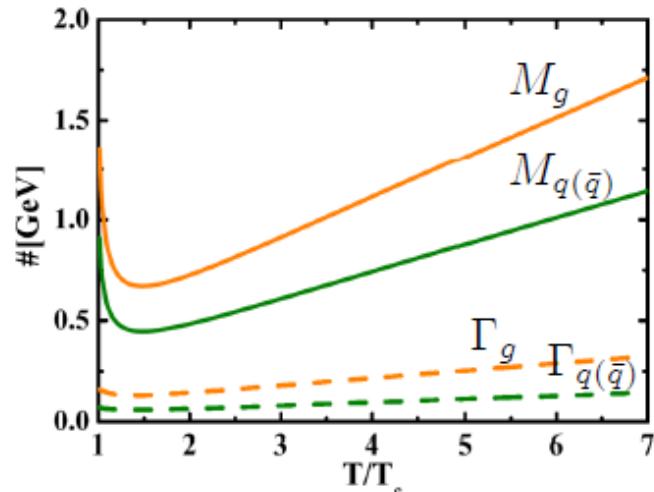
$$\rho_i(\omega, T) = \frac{4\omega\Gamma_i(T)}{\left(\omega^2 - \vec{p}^2 - M_i^2(T)\right)^2 + 4\omega^2\Gamma_i^2(T)}$$

- Modeling of the quark/gluon masses and widths  $\rightarrow$  HTL limit at high T

$$M_i(T) \propto \alpha_s(T) f_{HTL}(T), \quad \Gamma_i(T) \propto \alpha_s(T) f_{HTL}(T)$$

- fit to lattice (lQCD) results (e.g. entropy density) with 3 parameters

→ Quasi-particle properties:  
large width and mass for gluons and quarks



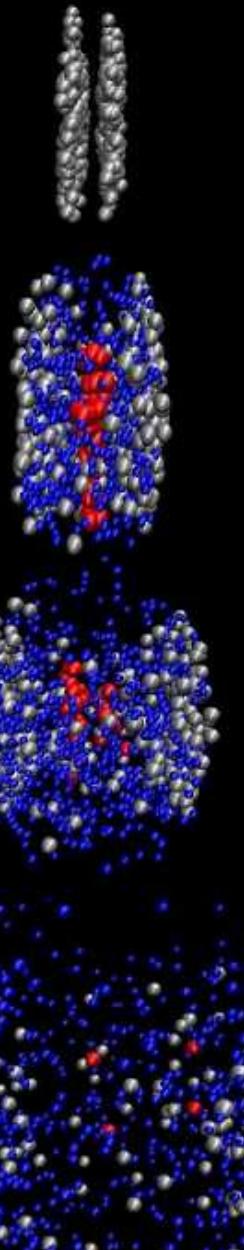
$$T_c = 158 \text{ MeV}$$

$$\varepsilon_c = 0.5 \text{ GeV/fm}^3$$

DQPM: Peshier, Cassing, PRL 94 (2005) 172301;  
Cassing, NPA 791 (2007) 365; NPA 793 (2007)

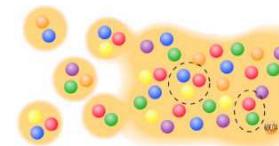


# Parton Hadron String Dynamics

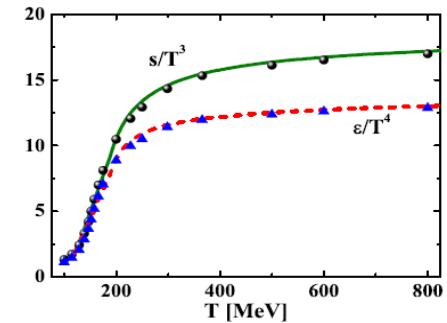


## I. From hadrons to QGP:

- Initial A+A collisions:
  - string formation in primary NN collisions
  - strings decay to pre-hadrons ( $B$  - baryons,  $m$  – mesons)
- Formation of QGP stage by dissolution of pre-hadrons into massive colored quarks + mean-field energy based on the **Dynamical Quasi-Particle Model (DQPM)** which defines **quark spectral functions**, masses  $M_q(\varepsilon)$  and widths  $\Gamma_q(\varepsilon)$  + **mean-field potential  $U_q$**  at given  $\varepsilon$  – local energy density (related by lQCD EoS to  $T$  - temperature in the local cell)

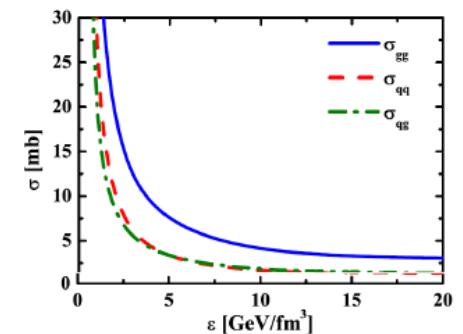


**QGP phase:**  
 $\varepsilon > \varepsilon_{\text{critical}}$



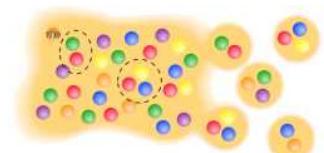
## II. Partonic phase - QGP:

- quarks and gluons (= ,dynamical quasiparticles‘) with off-shell spectral functions (width, mass) defined by the DQPM
- in **self-generated mean-field potential** for quarks and gluons  $U_q$ ,  $U_g$
- **EoS of partonic phase**: ‘crossover‘ from lattice **QCD** (fitted by DQPM)
- **(quasi-) elastic and inelastic** parton-parton interactions: using the effective cross sections from the DQPM



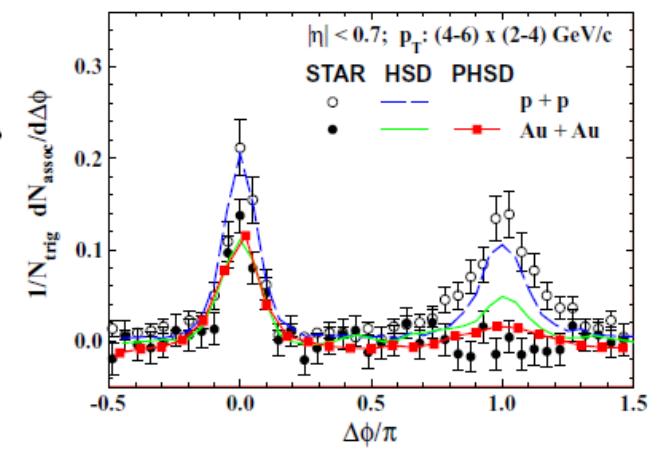
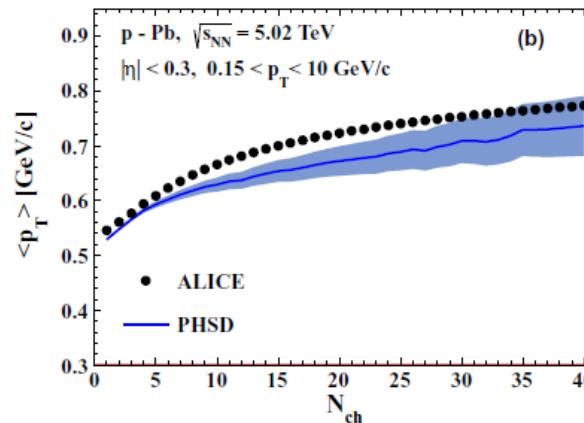
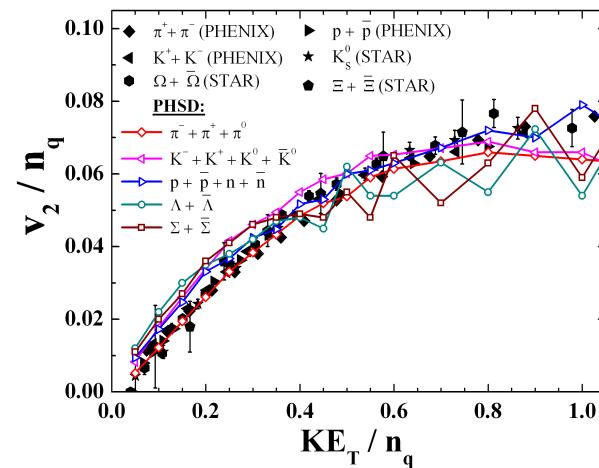
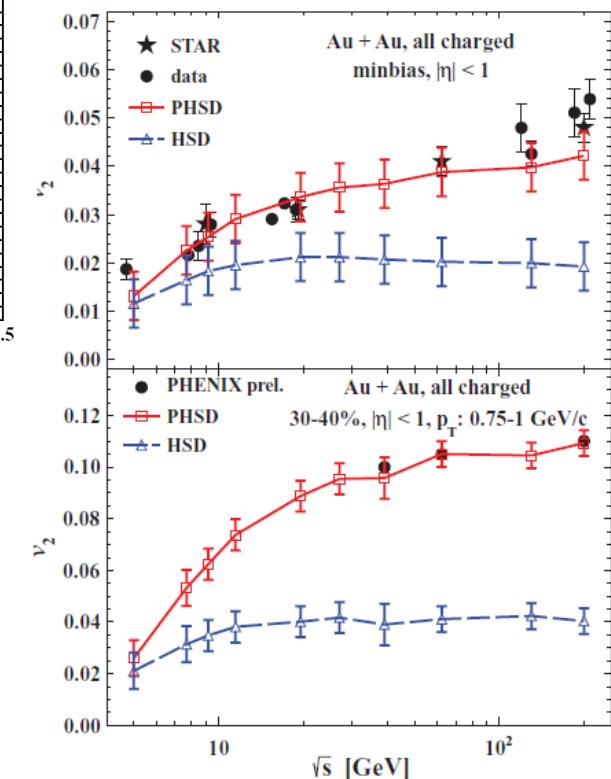
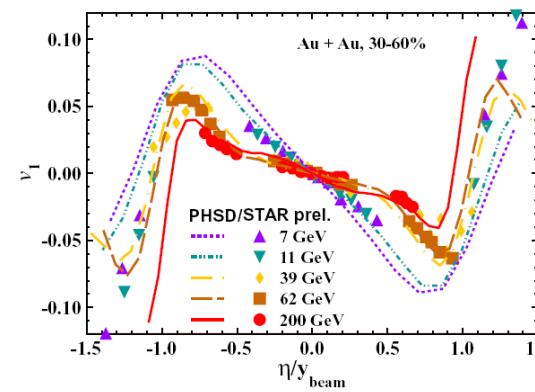
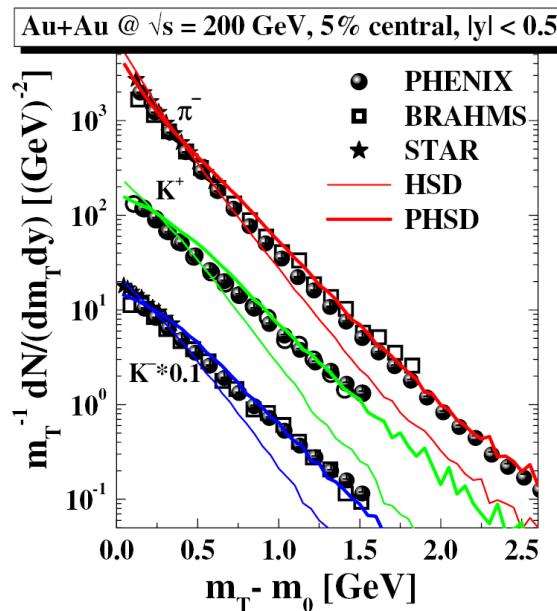
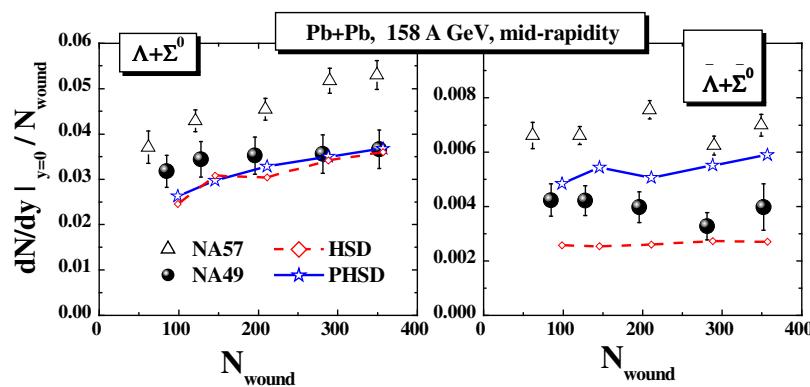
## III. Hadronization: based on DQPM

- **massive, off-shell (anti-)quarks** with broad spectral functions hadronize to off-shell mesons and baryons or color neutral excited states - ‘strings‘ (strings act as ,doorway states‘ for hadrons)



## IV. Hadronic phase: hadron-string interactions – off-shell HSD

# PHSD for HIC (highlights)

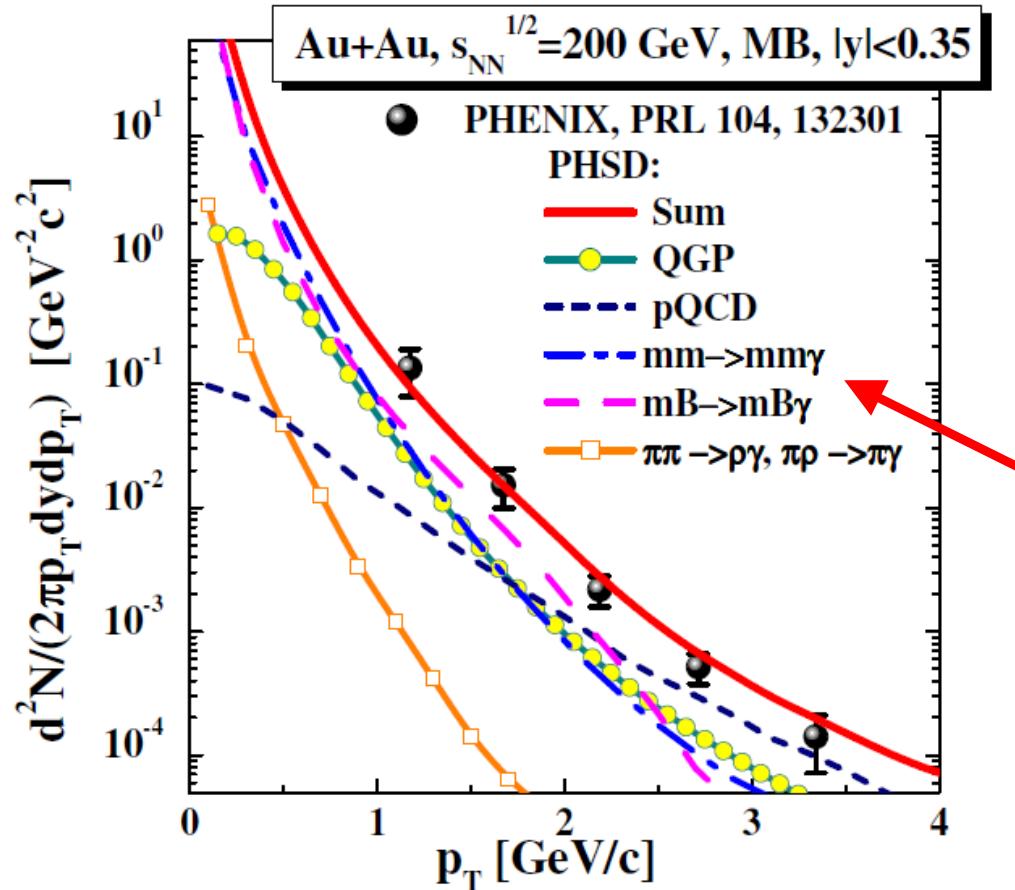


▪ PHSD provides a consistent description of HIC

# PHSD: photon spectra at RHIC: QGP vs. HG ?

Linnik et al., PRC88 (2013) 034904;  
PRC 89 (2014) 034908

- Direct photon spectrum (min. bias)



## PHSD:

- **QGP** gives up to ~50% of direct photon yield below 2 GeV/c

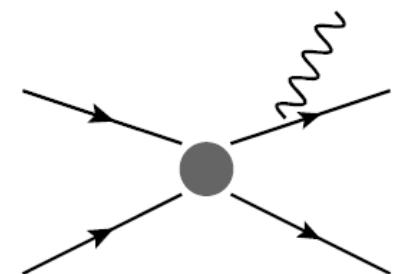
! sizeable contribution of hadronic sources  
 – meson-meson (mm) and meson-Baryon (mB) bremsstrahlung

$$m+m \rightarrow m+m+\gamma,$$

$$m+B \rightarrow m+B+\gamma,$$

$$m=\pi, \eta, \rho, \omega, K, K^*, \dots$$

$$B=p$$



!!! mm and mB bremsstrahlung channels can not be subtracted experimentally !

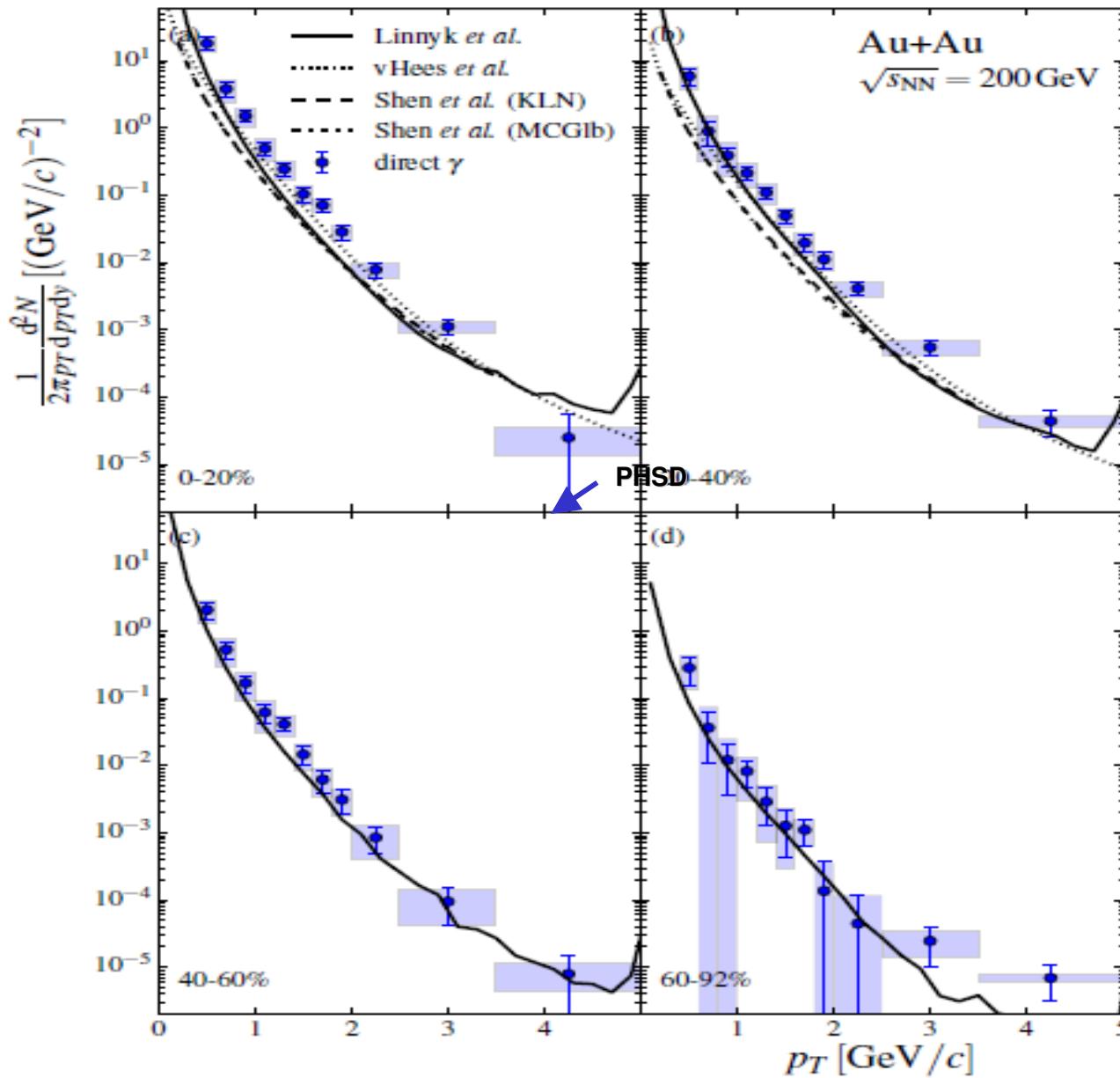
The slope parameter $T_{eff}$ (in MeV)			
PHSD		PHENIX [38]	
QGP	hadrons	Total	
$260 \pm 20$	$200 \pm 20$	$220 \pm 20$	$233 \pm 14 \pm 19$

# Photon $p_T$ spectra at RHIC for different centralities

from talk by S. Mizuno at QM'2014

PHENIX data - arXiv:1405.3940

PHSD predictions:  
O. Linnyk et al, Phys. Rev. C 89 (2014) 034908



□ PHSD approximately reproduces the centrality dependence

□ mm and mB bremsstrahlung is dominant at peripheral collisions

!!! Warning:  
large uncertainties in the Bremsstrahlung channels in the present PHSD results !

# Bremsstrahlung – trivial ‘background’?

## □ Uncertainties in the Bremsstrahlung channels in the present PHSD results :

1) based on the Soft-Photon-Approximation (SPA) (factorization = strong x EM)

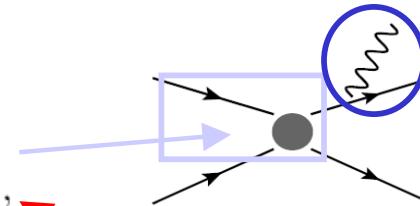
### □ Soft Photon Approximation (SPA):

$$m_1 + m_2 \rightarrow m_1 + m_2 + \gamma$$

C. Gale, J. Kapusta, Phys. Rev. C 35 (1987) 2107

$$q_0 \frac{d^3\sigma^\gamma}{d^3q} = \frac{\alpha}{4\pi} \frac{\bar{\sigma}(s)}{q_0^2}$$

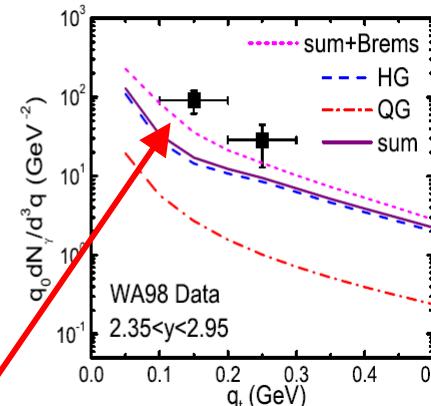
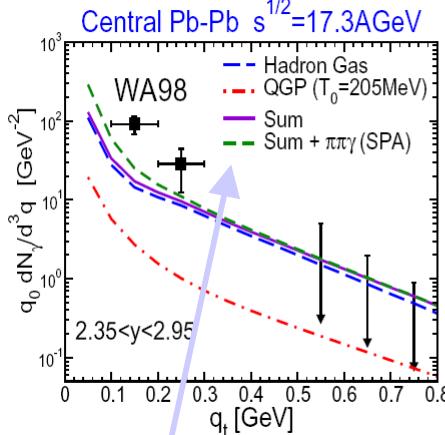
$$\bar{\sigma}(s) = \frac{s - (M_1 + M_2)^2}{2M_1^2} \sigma(s),$$



2) no experimental constraint on  $m+m$  and  $m+B$  differential elastic cross sections

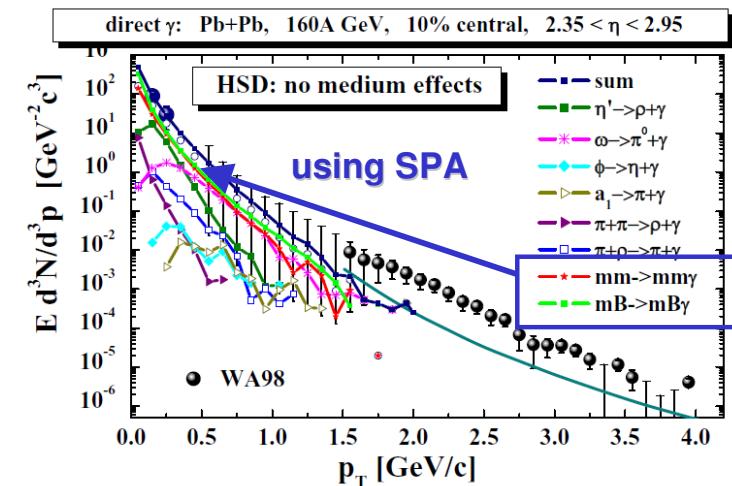
## □ Bremsstrahlung: seen at SPS - WA98

Fireball model: Liu, Rapp, Nucl. Phys. A 96 (2007) 101



- effective chiral model for  $\pi\pi \rightarrow \pi\pi\gamma$ ,  $\pi K \rightarrow \pi K\gamma$   
bremsstrahlung gives larger contribution than SPA

HSD: E. B., Kiselev, Sharkov, PR C78 (2008) 034905



→ mm and mB Bremsstrahlung seems to be an important source of soft photons!  
More work has to be done to have it under control!

# Centrality dependence of the 'thermal' photon yield

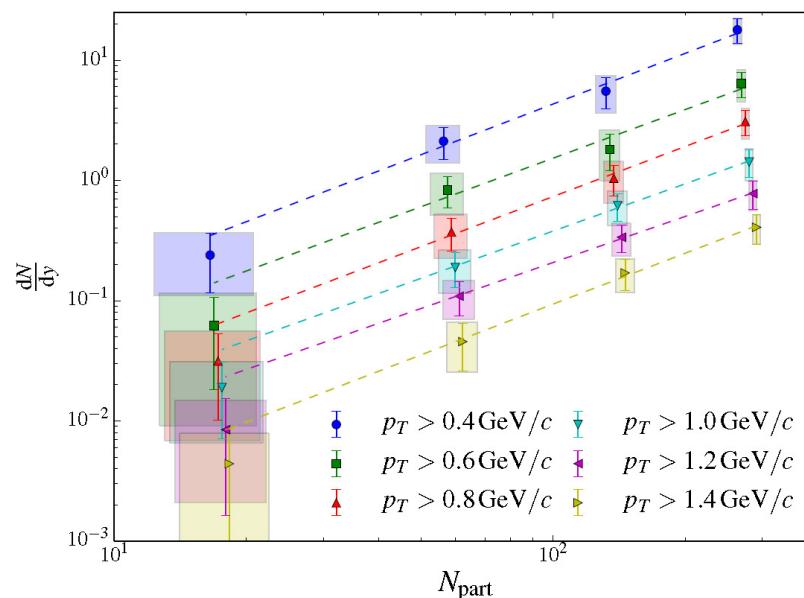
O. Linnyk et al, Phys. Rev. C 89 (2014) 034908

**PHENIX** (arXiv:1405.3940):

scaling of **thermal** photon yield vs centrality:

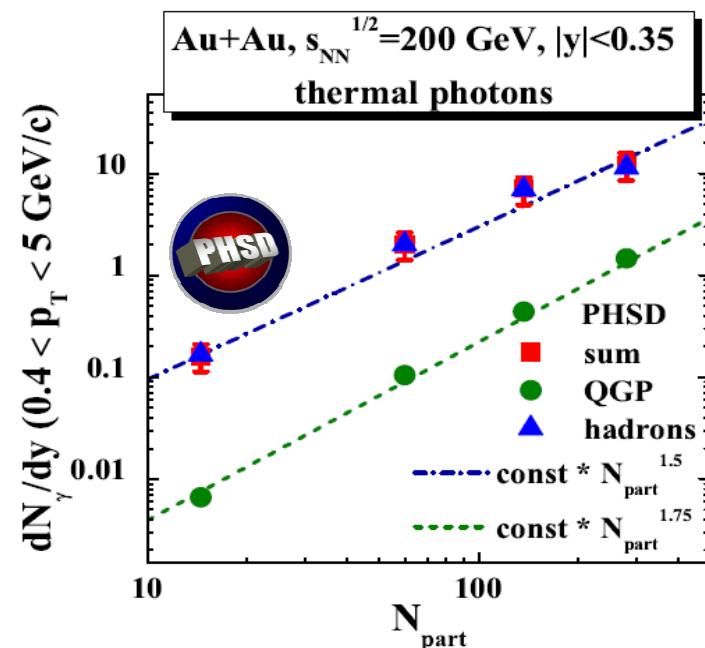
$$dN/dy \sim N_{\text{part}}^{\alpha} \text{ with } \alpha \sim 1.48 \pm 0.08$$

('Thermal' photon yield = direct photons - pQCD)



**PHSD predictions:**

- **Hadronic channels** scale as  $\sim N_{\text{part}}^{1.5}$
- **Partonic channels** scale as  $\sim N_{\text{part}}^{1.75}$



□ **PHSD:** scaling of the thermal photon yield with  $N_{\text{part}}^{\alpha}$  with  $\alpha \sim 1.5$

□ similar results from **viscous hydro**:

$$(2+1)\text{d VISH2+1: } \alpha(\text{HG}) \sim 1.46, \alpha(\text{QGP}) \sim 2, \alpha(\text{total}) \sim 1.7$$

→ What do we learn?

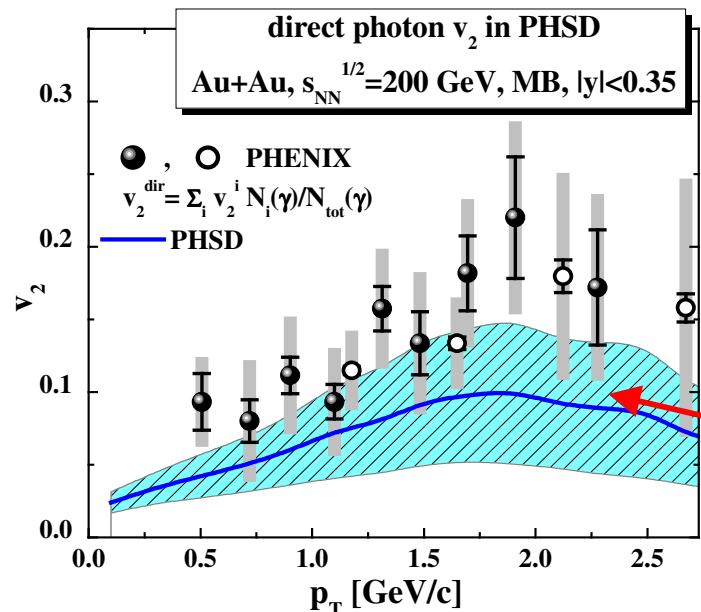
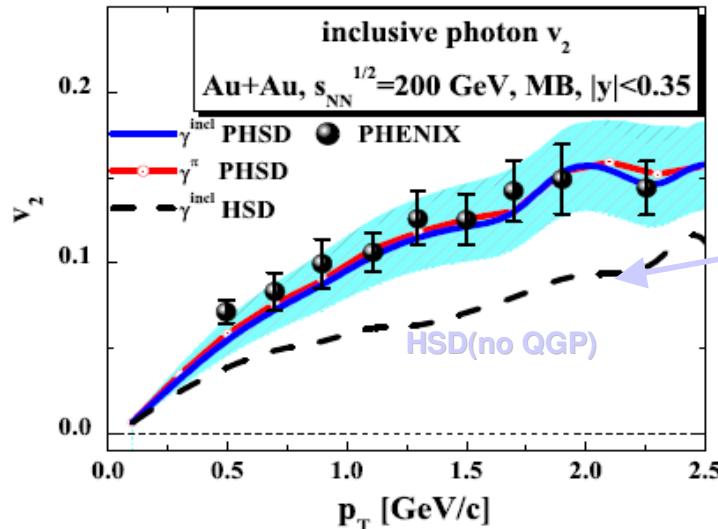
Indications for a dominant **hadronic origin of thermal photon production!**?

# Are the direct photons a barometer of the QGP?



- Do we see the **QGP pressure** in  $v_2(\gamma)$  if the photon production is **dominated by hadronic sources**?

PHSD: Linnyk et al.,  
PRC88 (2013) 034904;  
PRC 89 (2014) 034908



1)  $v_2(\gamma^{incl}) = v_2(\pi^0)$  - **inclusive photons mainly come from  $\pi^0$  decays**

▪ HSD (without QGP) underestimates  **$v_2$  of hadrons** and inclusive photons by a factor of 2, whereas the PHSD model with QGP is consistent with exp. data

→ The **QGP causes the strong elliptic flow of photons indirectly**, by enhancing the  $v_2$  of final hadrons due to the partonic interactions

**Direct photons (inclusive(=total) – decay):**

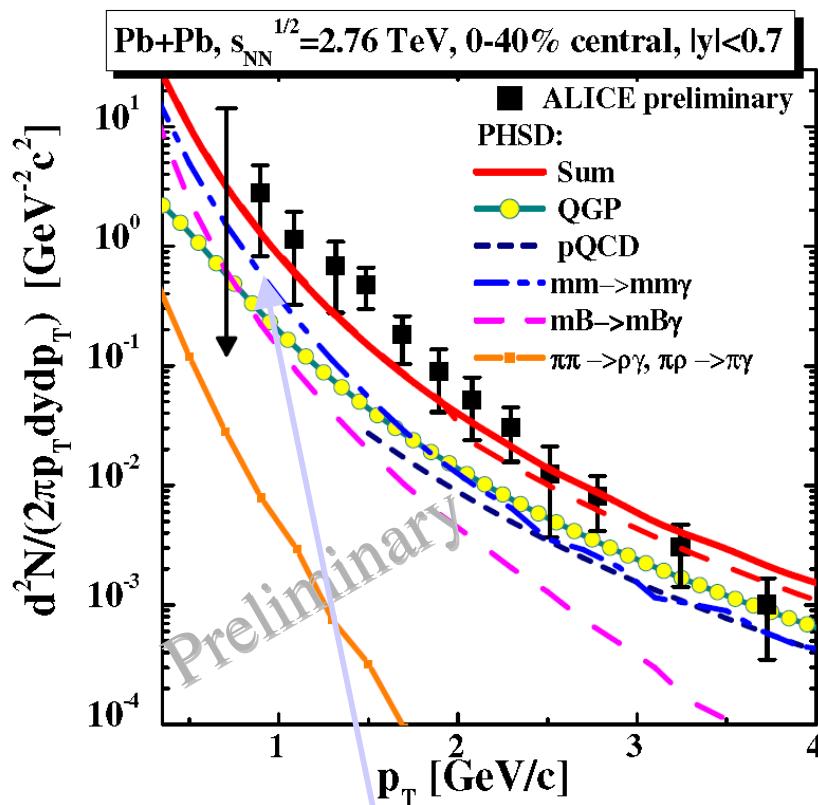
2)  $v_2(\gamma^{dir})$  of direct photons in PHSD underestimates the PHENIX data :

$v_2(\gamma^{QGP})$  is very small, but QGP contribution is up to 50% of total yield → lowering flow

→ **PHSD:  $v_2(\gamma^{dir})$  comes from mm and mB bremsstrahlung !**

# Photons from PHSD at LHC

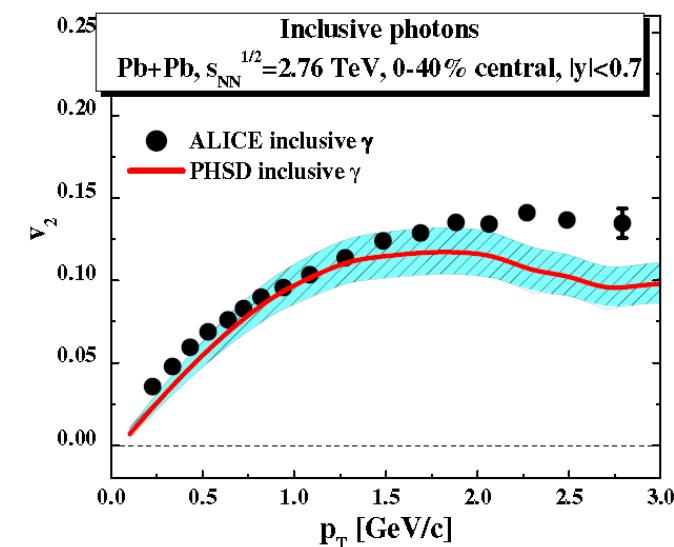
PHSD- preliminary: Olena Linnyk



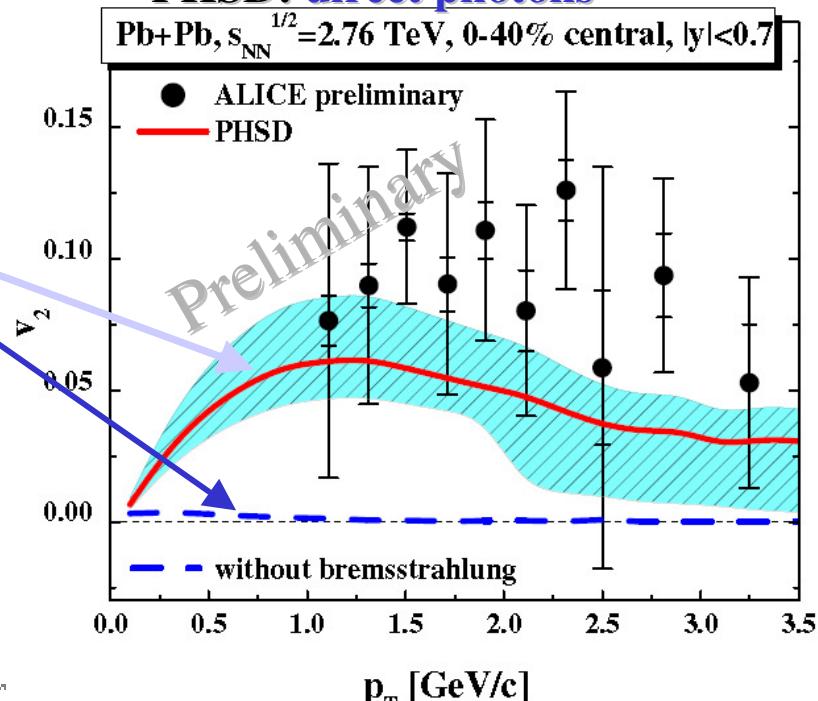
- ❑ Is the considerable elliptic flow of direct photons at the LHC also of **hadronic origin** as for RHIC?!
- ❑ The photon elliptic flow at LHC is lower than at RHIC due to **a larger relative QGP contribution / longer QGP phase.**

→ LHC (similar to RHIC):  
hadronic photons dominate spectra and  $v_2$

## PHSD: $v_2$ of inclusive photons



## PHSD: direct photons



# Towards the solution of the $v_2$ puzzle



- Is hadronic bremsstrahlung a „solution“?

## Other scenarios:

- Early-time magnetic field effects ?

(Basar, Kharzeev, Skokov, PRL109 (2012) 202303; Basar, Kharzeev, Shuryak, arXiv:1402.2286)

„... a novel photon production mechanism stemming from the **conformal anomaly of QCD-QED** and the **existence of strong (electro)magnetic fields** in heavy ion collisions.“

Exp. checks:  $v_3$ , centrality dependence of photon yield (PHENIX: arXiv:1405.3940)

- Glasma effects ?

(L. McLerran, B. Schenke, arXiv: 1403.7462)

„... Photon distributions from the Glasma are **steeper** than those computed in the Thermalized Quark Gluon Plasma (TQGP). Both the **delayed equilibration of the Glasma** and a possible anisotropy in the pressure lead to a slower expansion and mean times of photon emission of fixed energy are increased.“

- Pseudo-Critical Enhancement of thermal photons near  $T_c$  ?

(H. van Hees, M. He, R. Rapp, arXiv:1404.2846)

- non-perturbative effects?

semi-QGP - cf. talk by S. Lin at QM'2014

- ???



# ... shining in the darkness

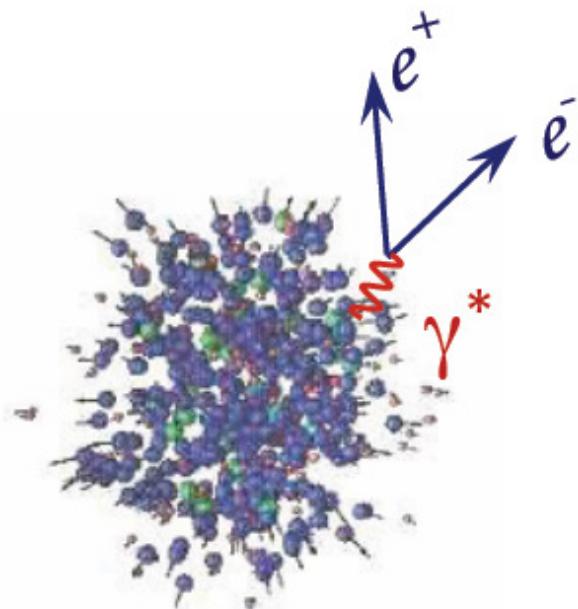
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## Some messages from the ‘photon adventure’:

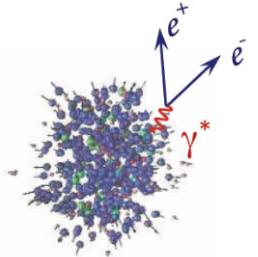
- **Can thermal QGP accelerate photons enough?!**
- **The role of hadronic sources (like bremsstrahlung) were underestimated?**
- **The importance of initial phases of the reaction:**  
Large photon  $v_2$  requires the development of pre-equilibrium / initial flow ?!
- **New sources of photon emission? Why not seen in dileptons?!**
- **The photons provide a critical test for the theoretical models:**  
models constructed to reproduce the ‚hadronic world‘ fail to explain the photon experimental data!
- ➔ **Additional impuls for the development of dynamical models:**  
e.g. from ideal to ebe viscous hydro, EoS, realistic non-equilibrium dynamics ...

**Photons – one of the most sensitive probes for the dynamics of HIC!**

# Dileptons: from SPS to LHC



# Dilepton sources

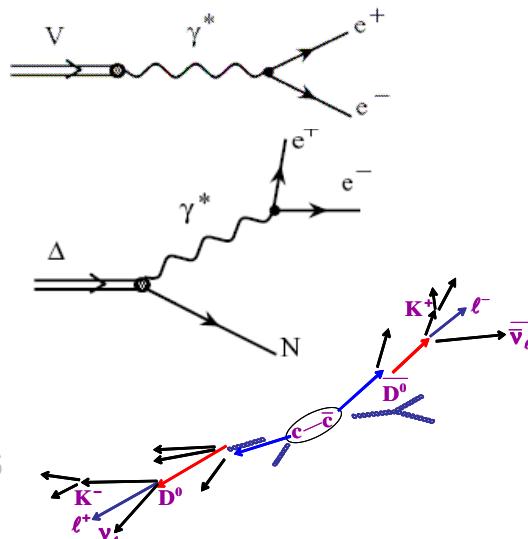


## □ from the QGP via partonic ( $q, \bar{q}, g$ ) interactions:

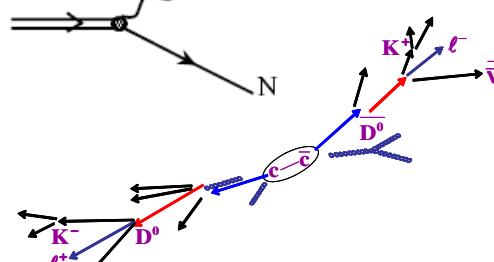


## □ from hadronic sources:

- direct decay of vector mesons ( $\rho, \omega, \phi, J/\Psi, \Psi'$ )

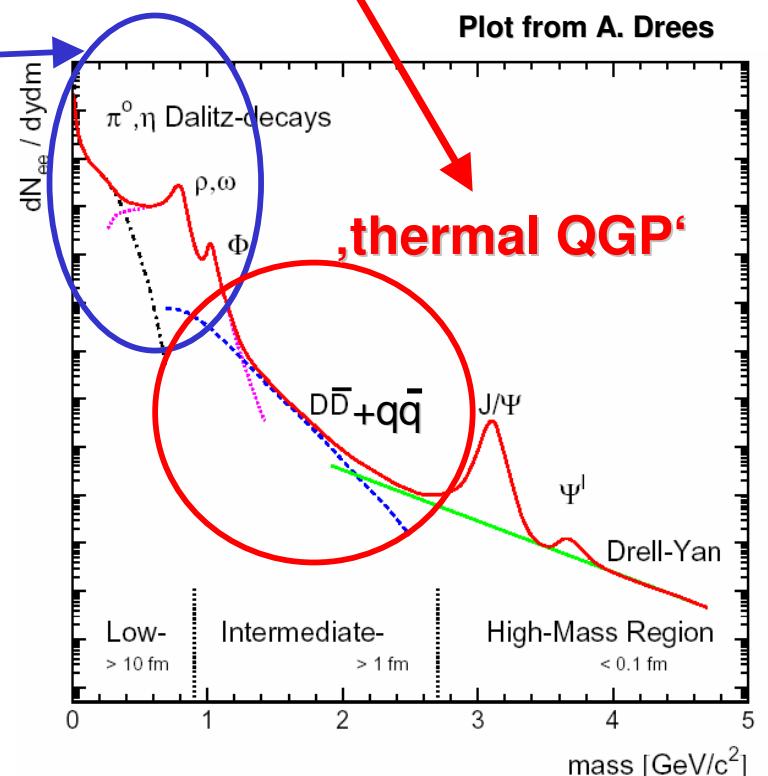


- Dalitz decay of mesons and baryons ( $\pi^0, \eta, \Delta, \dots$ )

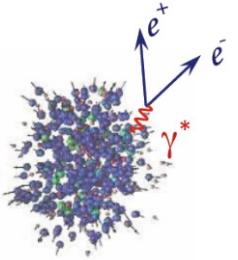


- correlated D+Dbar pairs

- radiation from multi-meson reactions ( $\pi + \pi, \pi + \rho, \pi + \omega, \rho + \rho, \pi + a_1$ ) - ,4 $\pi$

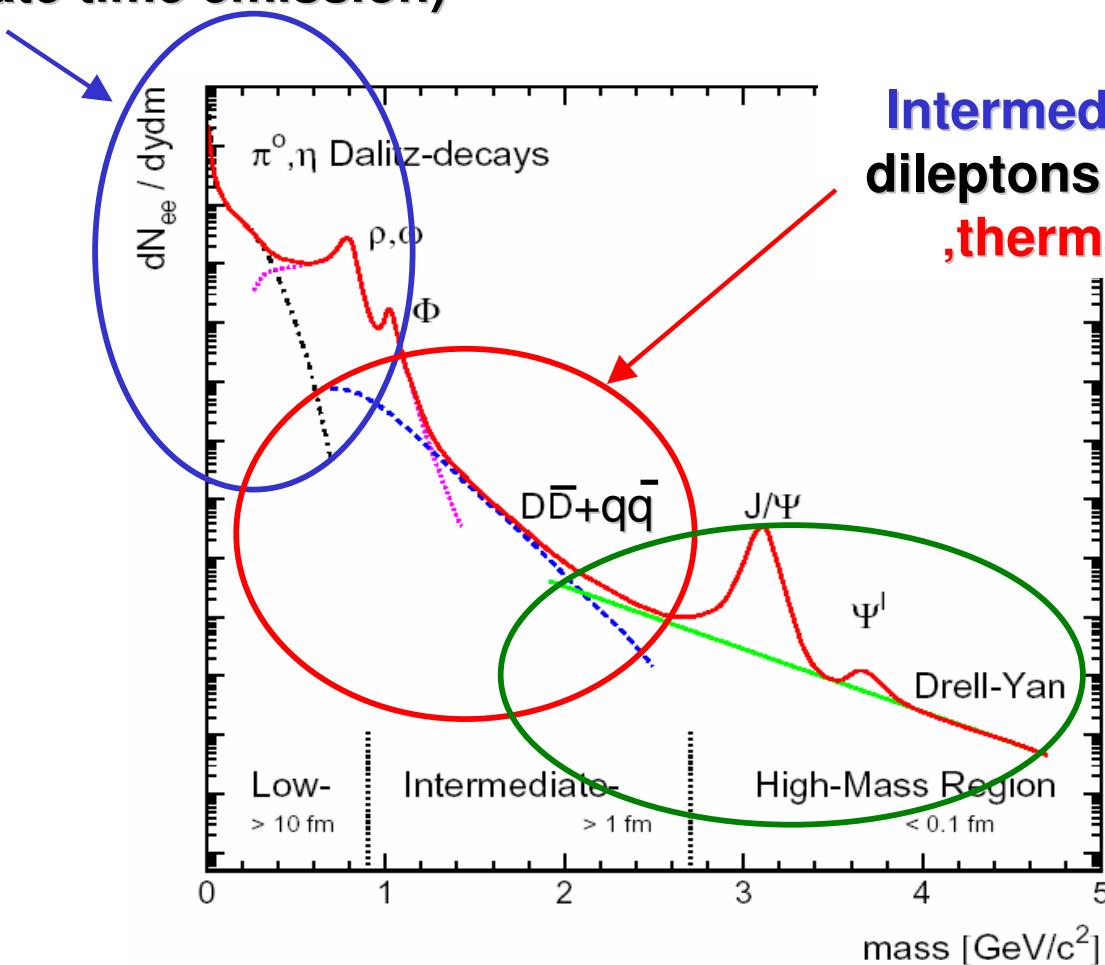


# Physics with dilepton



**Low mass dileptons**  
-probe of **hadronic  
in-medium effects**  
(late time emission)

**Advantage of dileptons:** additional „degree of freedom“ ( $M$ ) allows to disentangle various sources

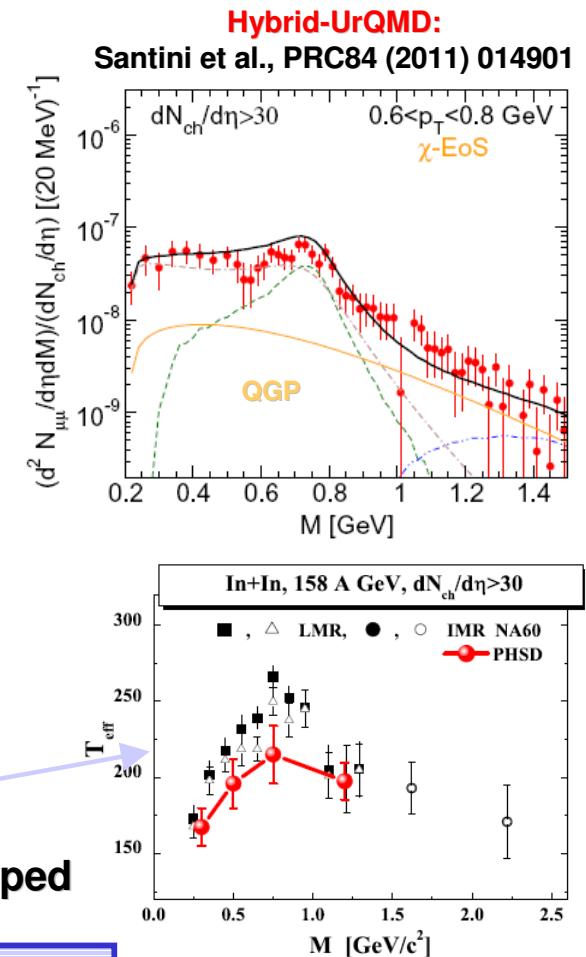
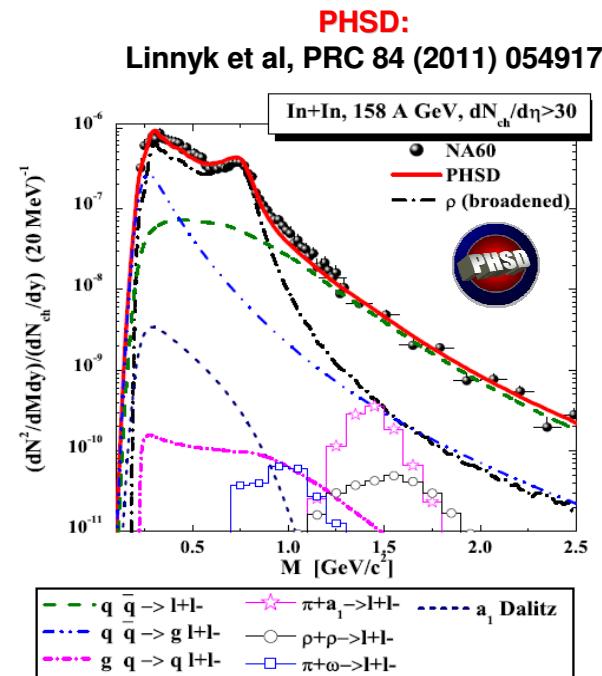
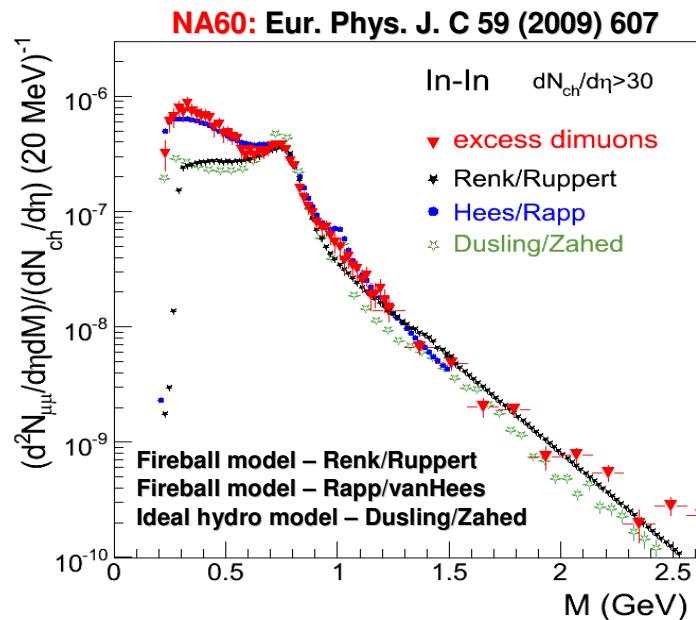


**Intermediate mass  
dileptons – probe of  
,thermal QGP‘**

**High-mass dileptons**  
– probe of  
**pQGP and  
hard probes**  
(early time emission)

# Lessons from SPS: NA60

## Dilepton invariant mass spectra:

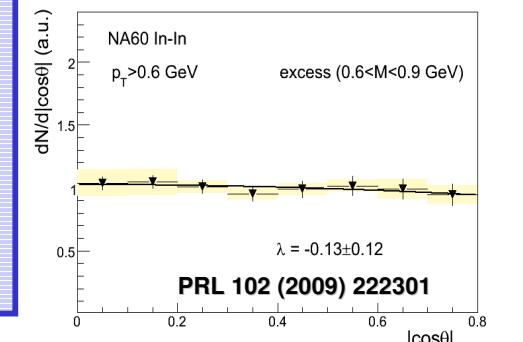


## Inverse slope parameter $T_{eff}$ :

spectrum from QGP is softer than from hadronic phase since the QGP emission occurs dominantly before the collective radial flow has developed

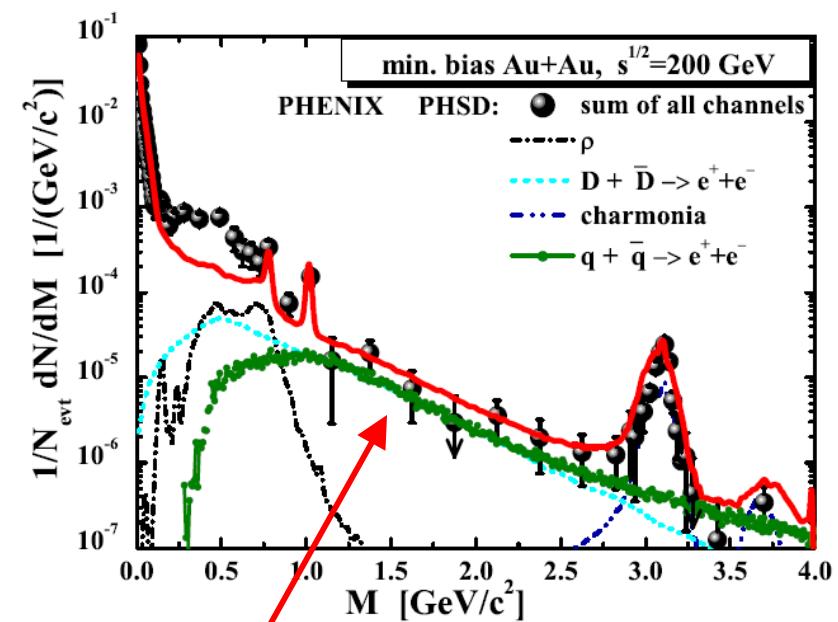
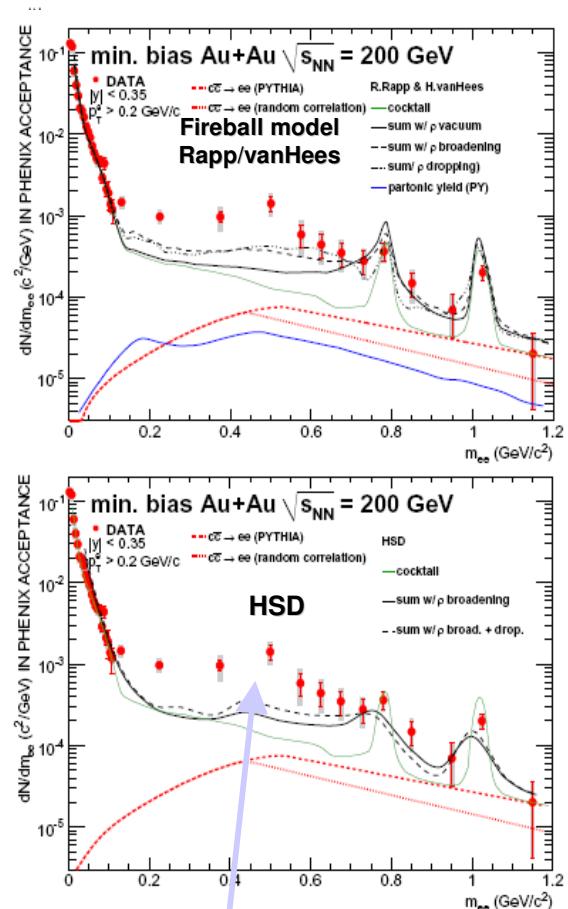
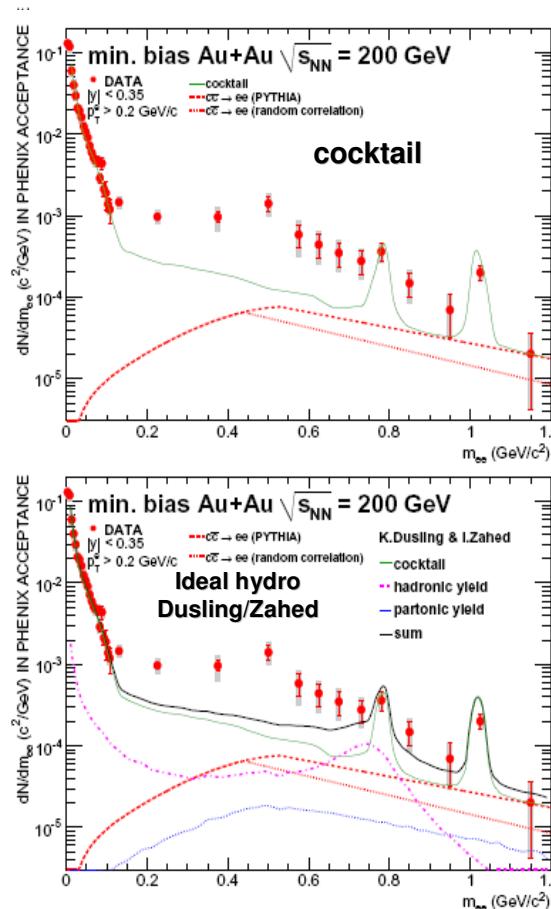
## Message from SPS: (based on NA60 and CERES data)

- 1) Low mass spectra - evidence for the **in-medium broadening of p-mesons**
- 2) Intermediate mass spectra above 1 GeV - dominated by **partonic radiation**
- 3) The rise and fall of  $T_{eff}$  – evidence for the thermal **QGP radiation**
- 4) Isotropic angular distribution – indication for a **thermal origin of dimuons**

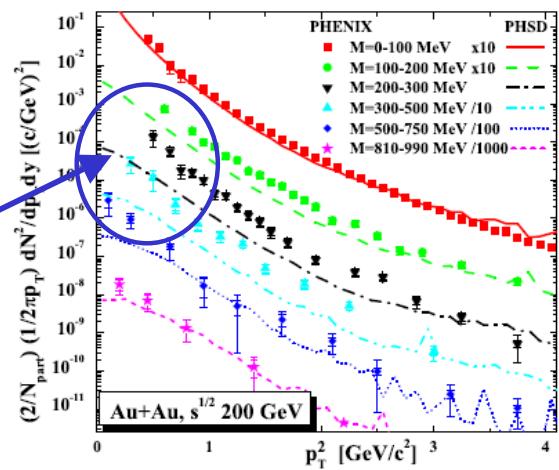


# Dileptons at RHIC: PHENIX

PHENIX: PRC81 (2010) 034911



Linnyk et al., PRC 85 (2012) 024910

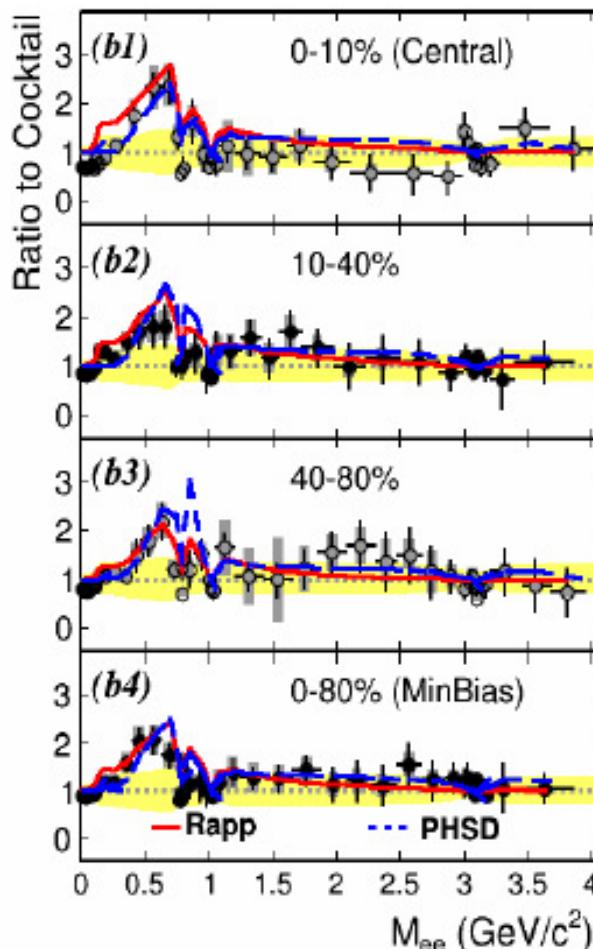
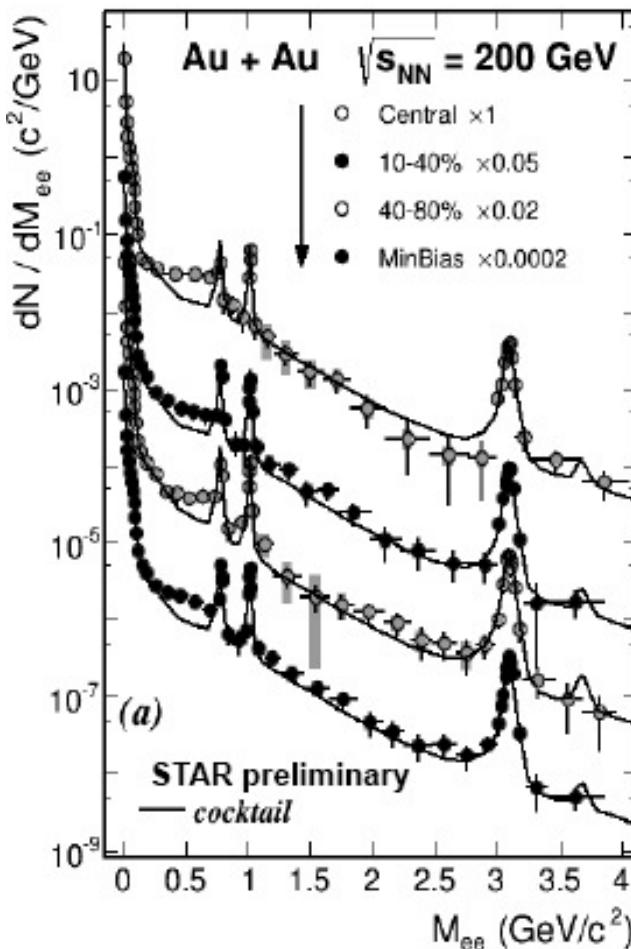


## Message:

- Models provide a good description of pp data and peripheral Au+Au data, however, fail in describing the excess for central collisions even with in-medium scenarios for the vector meson spectral function
- The ‘missing source’(?) is located at low  $p_T$
- Intermediate mass spectra – dominant QGP contribution

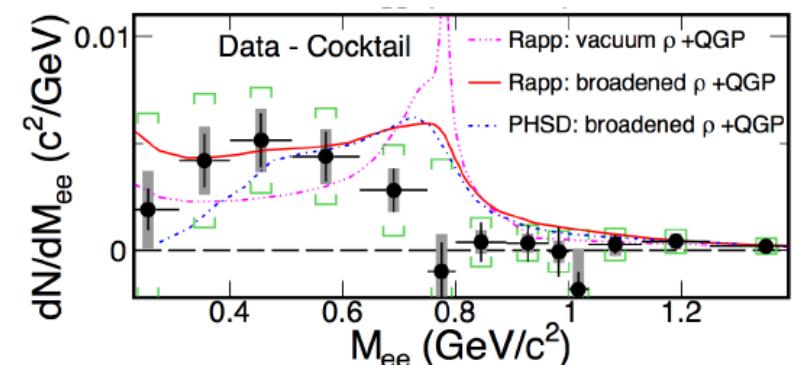
# Dileptons at RHIC: STAR data vs model predictions

## Centrality dependence of dilepton yield



(Talk by P. Huck at QM'2014)

## Excess in low mass region, min. bias



### Models:

- **Fireball model – R. Rapp**
- **PHSD**

### Low masses:

collisional broadening of  $\rho$

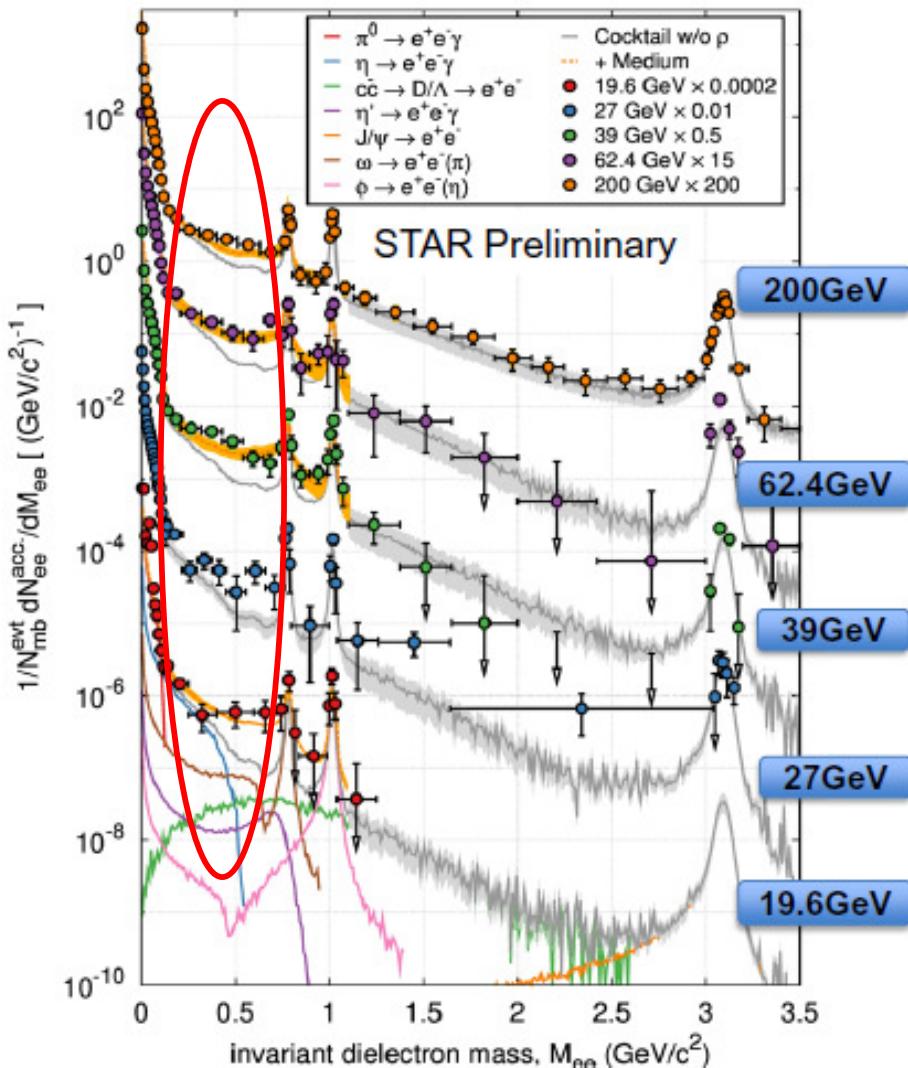
### Intermediate masses:

QGP dominant

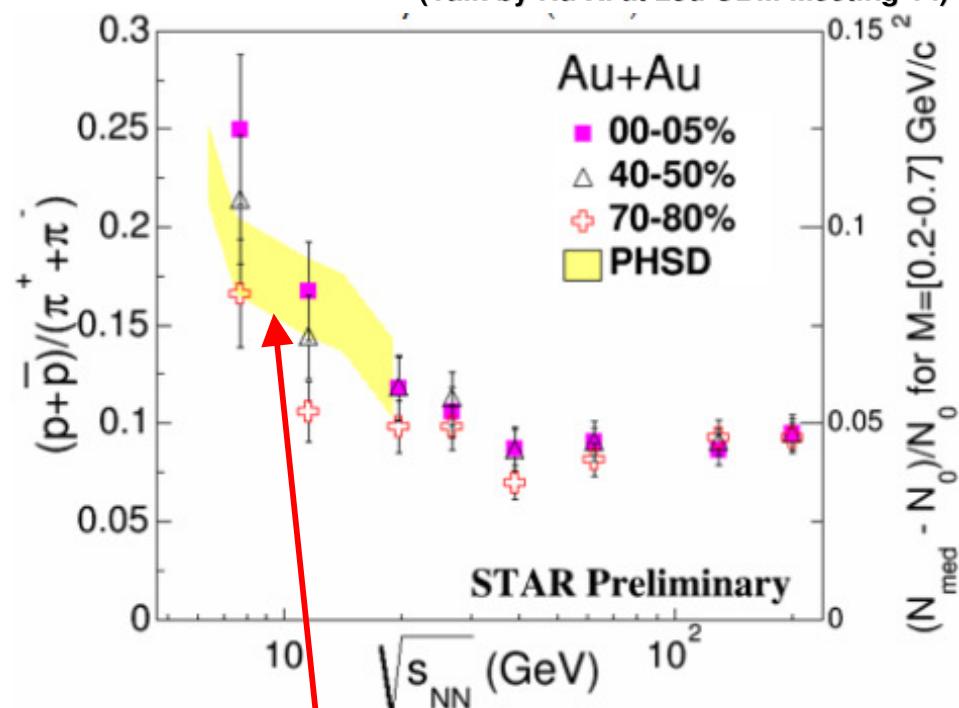
**Message:** STAR data are described by models within a **collisional broadening scenario** for the vector meson spectral function + **QGP**

# Dileptons from RHIC BES: STAR

(Talk by Nu Xu at QM'2014)



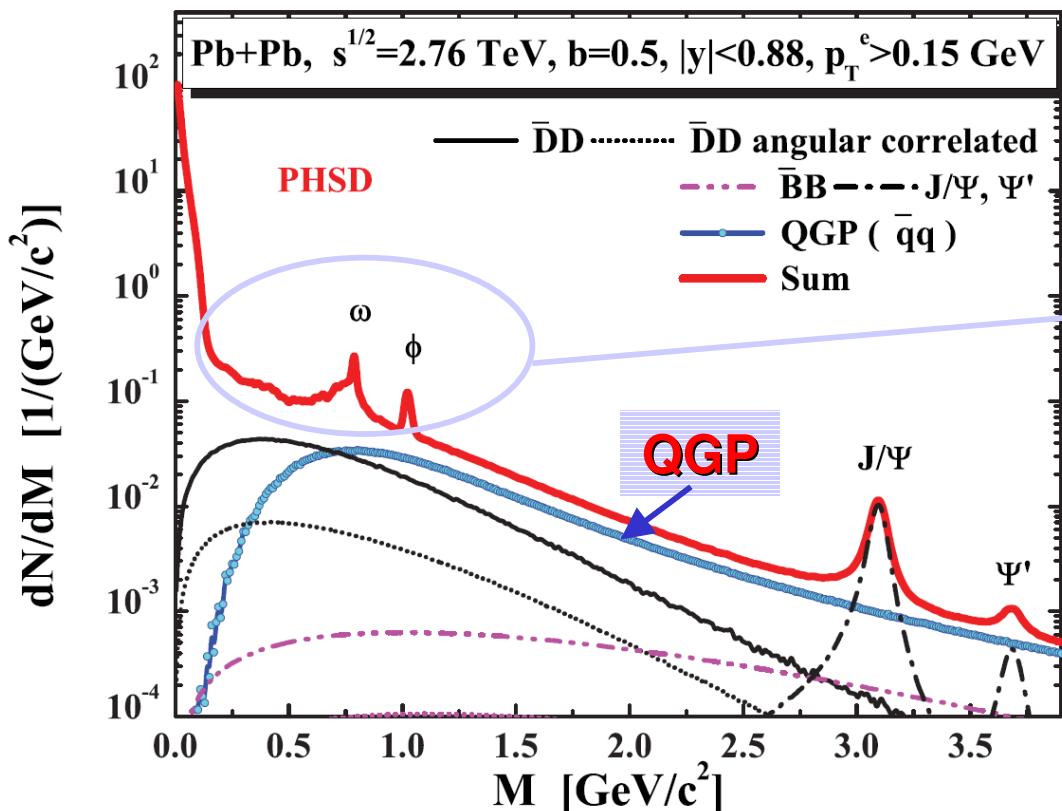
(Talk by Nu Xi at 23d CBM Meeting'14)



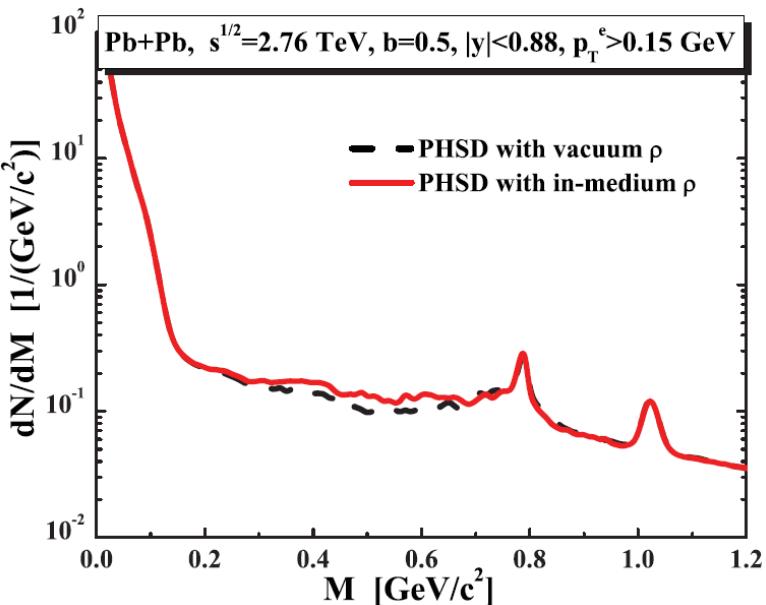
## Message:

- BES-STAR data show a constant low mass excess (scaled with  $N(\pi^0)$ ) within the measured energy range
  - PHSD model: excess increasing with decreasing energy due to a longer  $\rho$ -propagation in the high baryon density phase
- Good perspectives for future experiments –  
CBM(FAIR) / MPD(NICA)

# Dileptons at LHC



O. Linnyk, W. Cassing, J. Manninen, E.B., P.B. Gossiaux, J. Aichelin, T. Song, C.-M. Ko,  
Phys.Rev. C87 (2013) 014905; arXiv:1208.1279



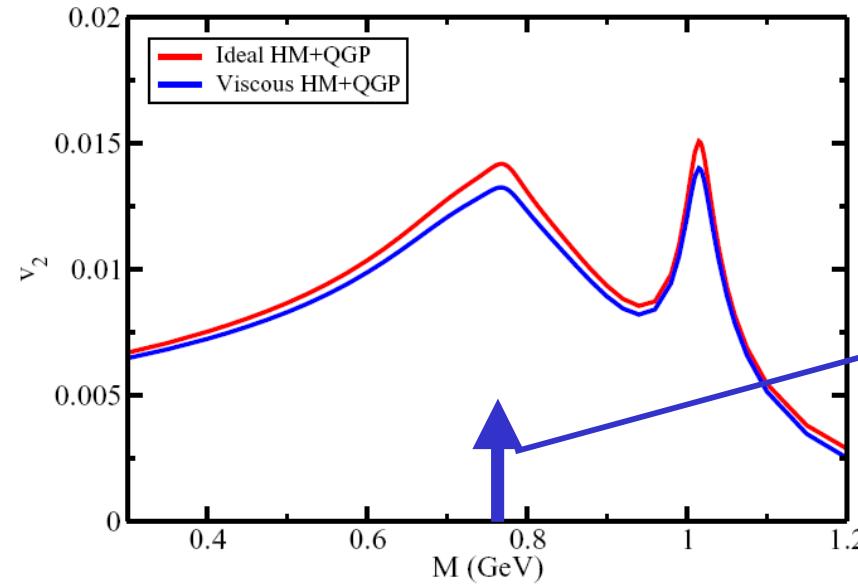
## Message:

- low masses - hadronic sources: in-medium effects for  $\rho$  mesons are small
- intermediate masses: QGP + D/Dbar
  - charm ‘background’ is smaller than thermal QGP yield
  - **QGP( $\bar{q}q$ ) dominates at  $M>1.2 \text{ GeV} \rightarrow$  clean signal of QGP at LHC!**

# Perspectives with dileptons: $v_n$

Vujanovic, Young, Schenke, Rapp, Jeon, Gale, PRC 89 (2014) 034904

(3+1)d MUSIC: Au+Au, RHIC, 10% central

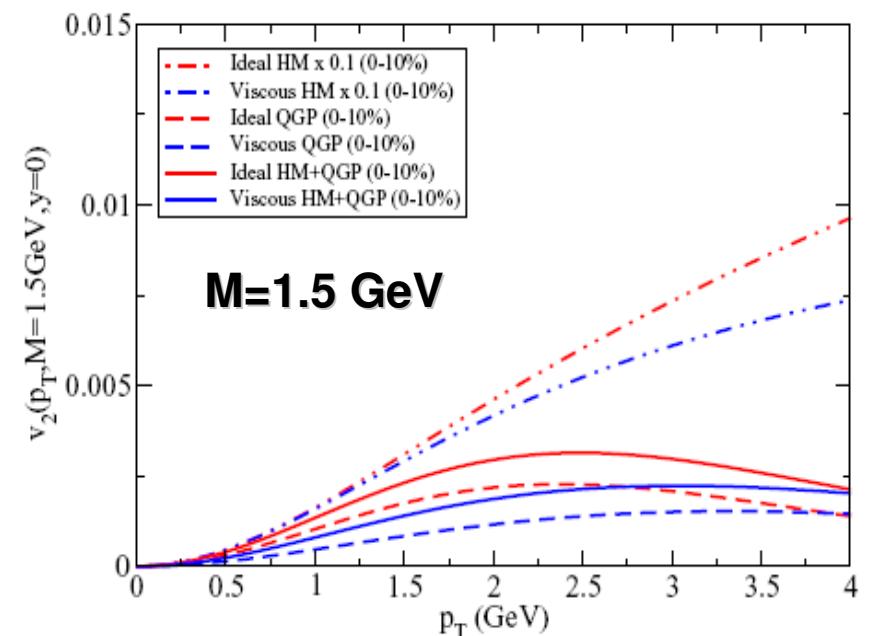
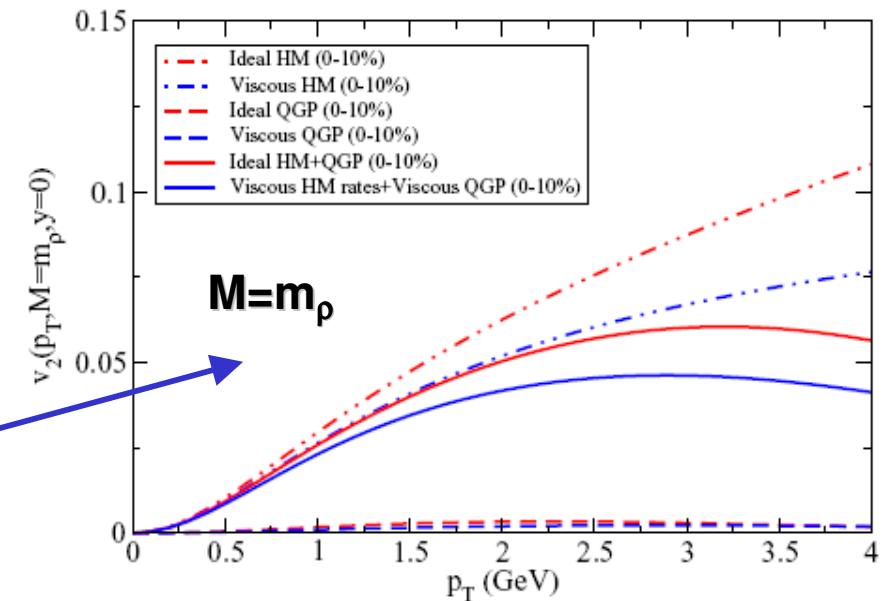


$v_2$  (similar for  $v_3$ ):

sensitive to the EoS and  $\eta/s$

sensitive to the sources

Dileptons: advantages compared to photons – extra degree of freedom  $M$  allows to disentangle the sources!



# Messages from dilepton data

## □ Low dilepton masses:

- Dilepton spectra show sizeable changes due to the in-medium effects
  - modification of the properties of vector mesons (as collisional broadening) - which are observed experimentally
- In-medium effects can be observed at all energies from SIS to LHC

## □ Intermediate dilepton masses:

- The QGP ( $\bar{q}q$ ) dominates for  $M > 1.2 \text{ GeV}$
- Fraction of QGP grows with increasing energy; at the LHC it is dominant

## Outlook:

- \* experimental energy scan
- \* experimental measurements of dilepton's higher flow harmonics  $v_n$

