

Masses, r-process, neutron-stars, fission

Peter Möller

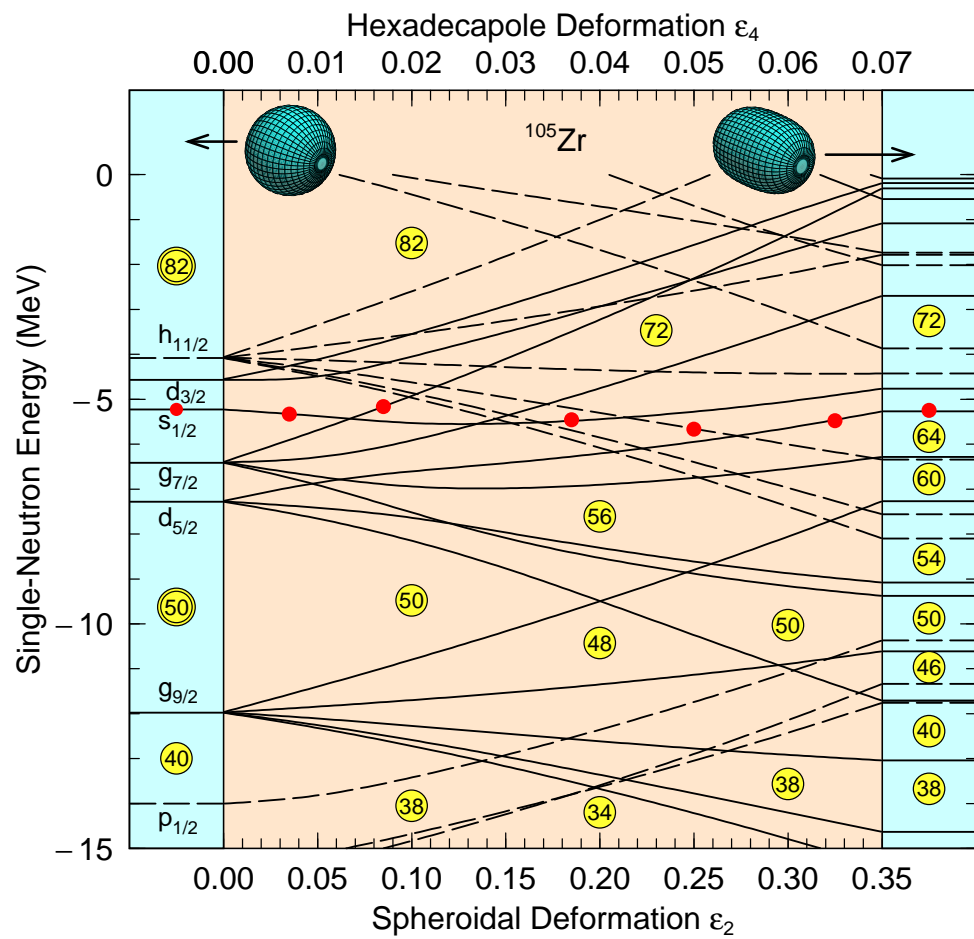
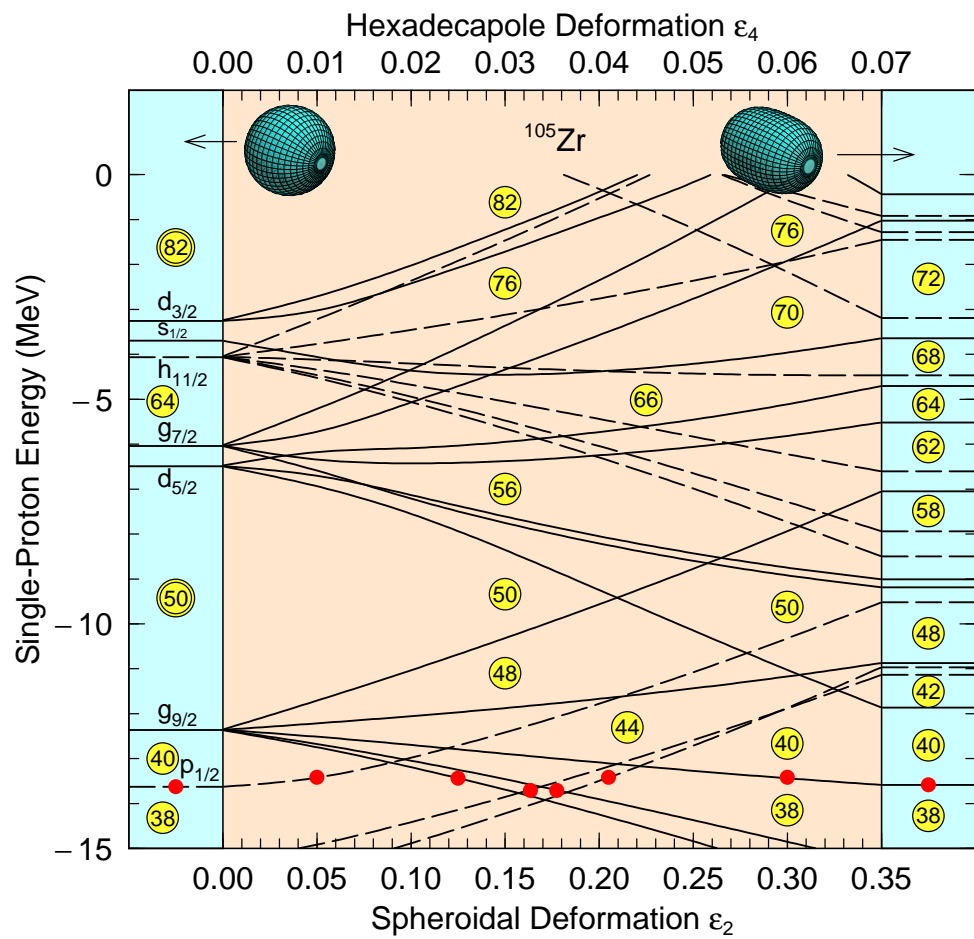
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Collaborators on this and other projects:

W. D. Myers, J. Randrup(LBL), H. Sagawa (Aizu), S. Yoshida (Hosei), T. Ichikawa(YITP), A. J. Sierk(LANL), A. Iwamoto (JAEA), S. Aberg (Lund), R. Bengtsson (Lund), S. Gupta (IIT, Ropar), and many experimental groups (e. g. K.-L. Kratz (Mainz), H. Schatz (MSU), A. Andreyev (York) . . .).

More details about masses, other projects (beta-decay,fission), associated ASCII data files, interactive access to data (type in Z, A and get specific data, contour maps) and figures are at

<http://t2.lanl.gov/nis/molleretal/>



CALCULATION OF GAMOW–TELLER β -STRENGTH FUNCTIONS IN THE RUBIDIUM REGION IN THE RPA APPROXIMATION WITH NILSSON-MODEL WAVE FUNCTIONS

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Abstract: We calculate allowed Gamow–Teller and, in a few cases, Fermi β -strength functions in a model that is applicable to studies of nuclei throughout the periodic system. For our first study we have selected a sequence of rubidium isotopes, namely ^{89}Rb – ^{99}Rb . We develop a model that uses calculated Nilsson-model wave functions, spherical or deformed, as the case may be, as the starting point for determining the wave functions of the mother and daughter nuclei in the β -decay. Pairing is treated in the BCS approximation. To account for the retardation of low-energy GT decay rates we add, as is customarily done, a simple residual interaction specific to GT decay, namely $V_{\text{GT}} = :\beta^1\text{--}\beta^{1+}:$, to the hamiltonian. This residual interaction is treated in the RPA approximation. The strength of the interaction is adjusted to get agreement between the calculated and experimental energy of the giant Gamow–Teller resonance for ^{208}Pb and ^{144}Sm . Since the present model is based on calculated wave functions and single-particle levels, studies of nuclei far from stability, where little experimental information is available, are more straightforward relative to calculations where “experimental” levels are used. The model can treat deformed nuclei employing wave functions calculated to desired accuracy, within the framework of the model, for the deformed single-particle well. The calculations show that use of single-particle parameters appropriate to the region studied and taking deformation into account is important. We find good agreement between calculated and experimental spectra over the region studied, provided an appropriate choice of single-particle parameters and deformation is made.

1. Introduction

A theoretical understanding of β -strength functions is important for the interpretation of the large amount of experimental data that is now being collected on, for instance, the giant Gamow–Teller resonance ¹⁻²⁾ and on β -decay spectra of nuclei far from stability ³⁾. More references on these subjects are found in the conference report ⁴⁾ of the 4th International Conference on Nuclei far from Stability and references quoted therein. The β -strength function must also be known for theoretical studies of phenomena where the β -strength function cannot be easily measured. Examples of such processes are the decay from the r-process line to the line of β -stability ⁵⁻⁷⁾, β -delayed particle emission ⁷⁾ and the production of transuranium elements by neutron capture and subsequent decay by β -emission, fission and neutron emission ^{8,9)}.

A variety of models have been used to calculate β -strength functions. The gross theory of β -decay ¹⁰⁾ describes, because of its statistical character, only the average

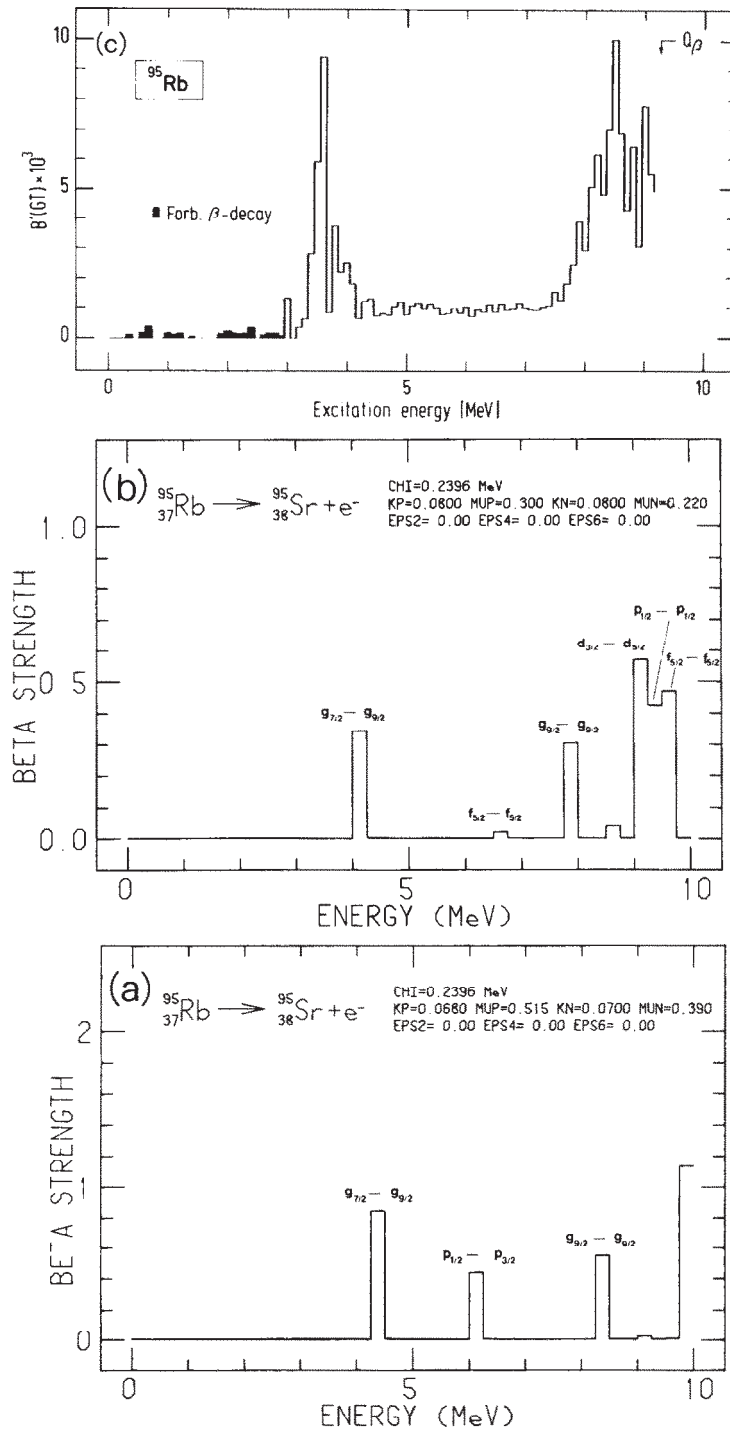


Fig. 8. Calculated β^- GT strength function for ^{95}Rb with (a) the $A=100$ set of κ and μ , and (b) the $N=60$ set of κ and μ , in units of $4f^2 \text{ MeV}^{-1}$. The peaks are labeled with the quantum numbers of the neutron (first) and proton (last) orbital involved in the transition. (c) Experimental results from ref. ³.

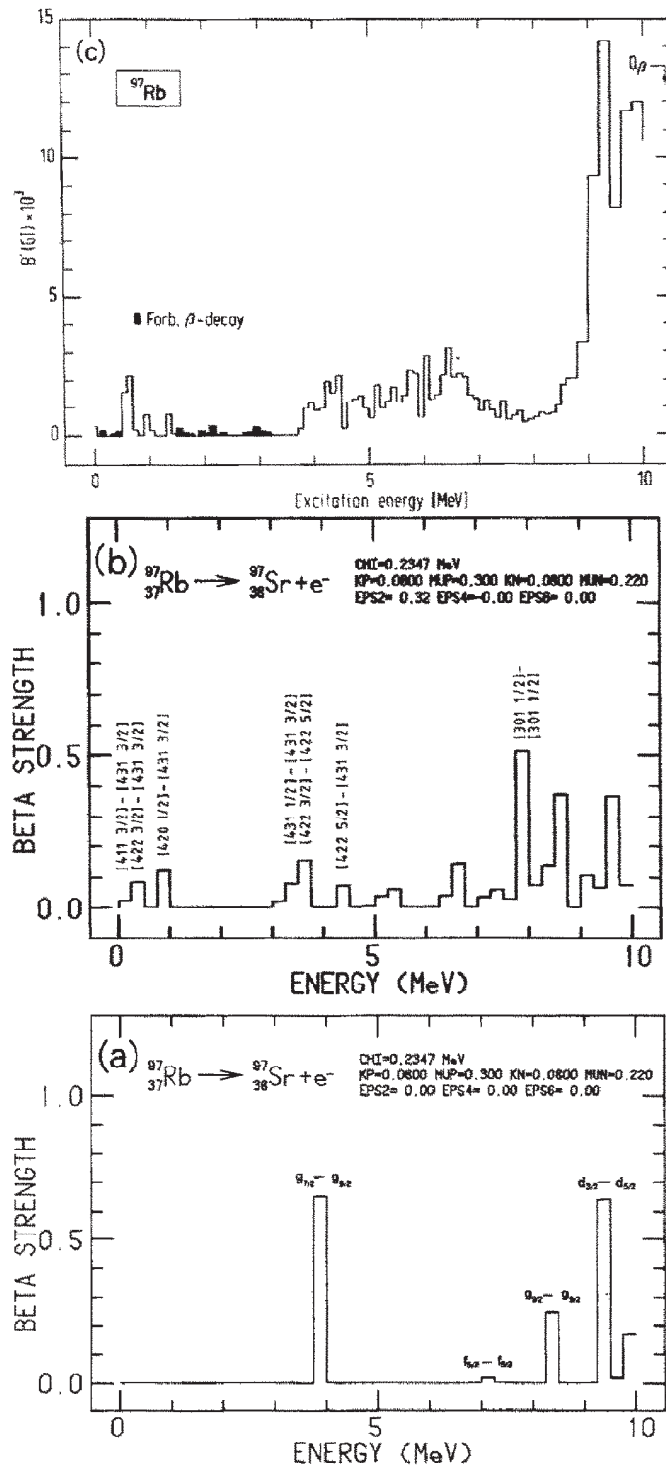


Fig. 9. Calculated β^- GT strength function for ^{97}Rb treated as (a) a spherical nucleus, and (b) a deformed nucleus, in units of $4f^2 \text{ MeV}^{-1}$. The peaks in (b) are labeled with the asymptotic quantum numbers of the neutron (first) and proton (last) orbital involved in the transition. (c) Experimental results from ref. ³⁾.

Successive FRDM enhancements

Optimization (2006)

Better search for optimum FRDM parameters.

Accuracy improvement: 0.01 MeV

New mass data base (AME2003) (2006)

Better agreement than with AME1989.

Accuracy improvement: 0.04 MeV

Full 4D energy minimization (2006–2008)

Full 4D minimization($\epsilon_2, \epsilon_3, \epsilon_4, \epsilon_6$) step=0.01.

Accuracy improvement: 0.02 MeV

Axial asymmetry (2002–2006)

Also yields correct SHE gs assignments.

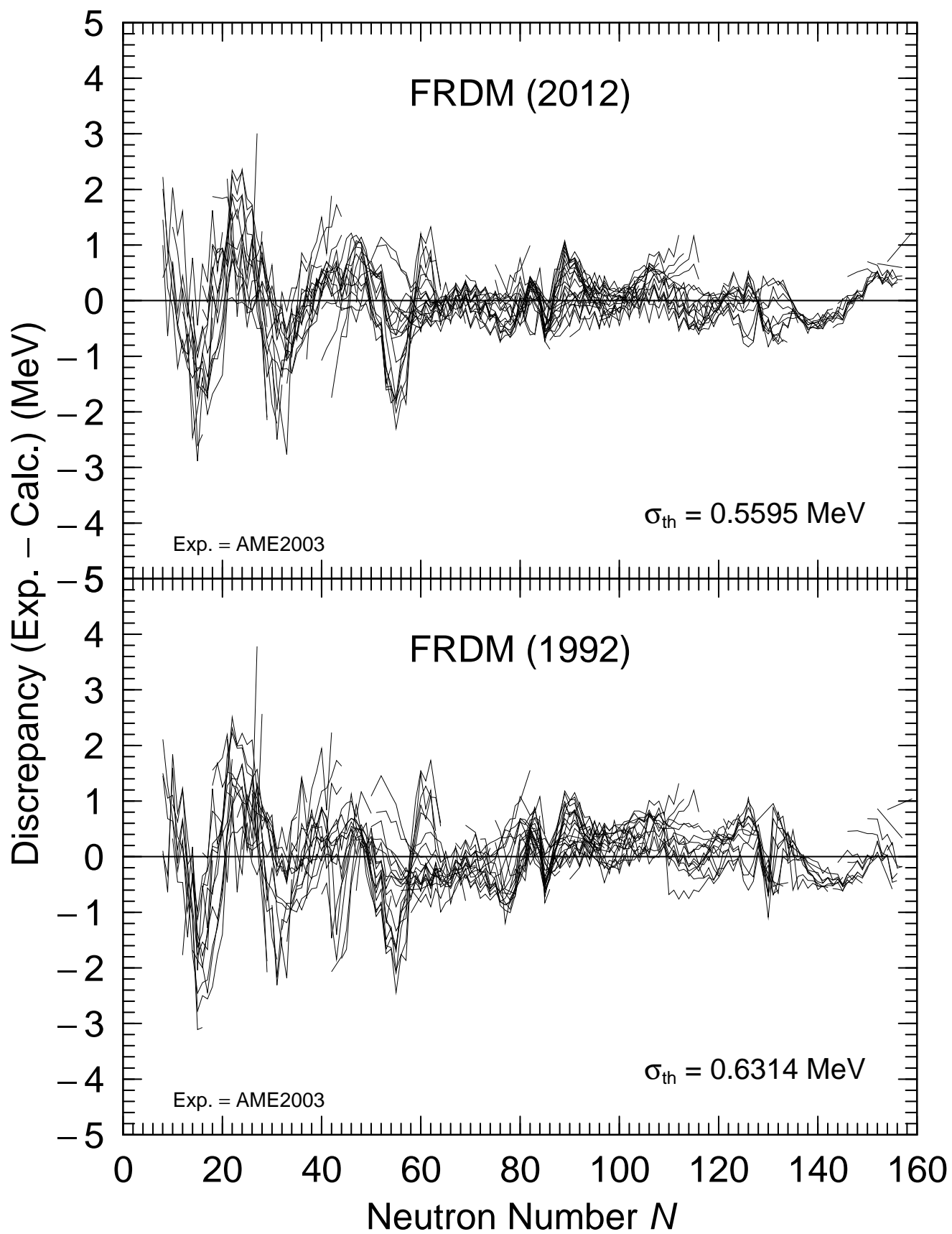
Accuracy improvement: 0.01 MeV

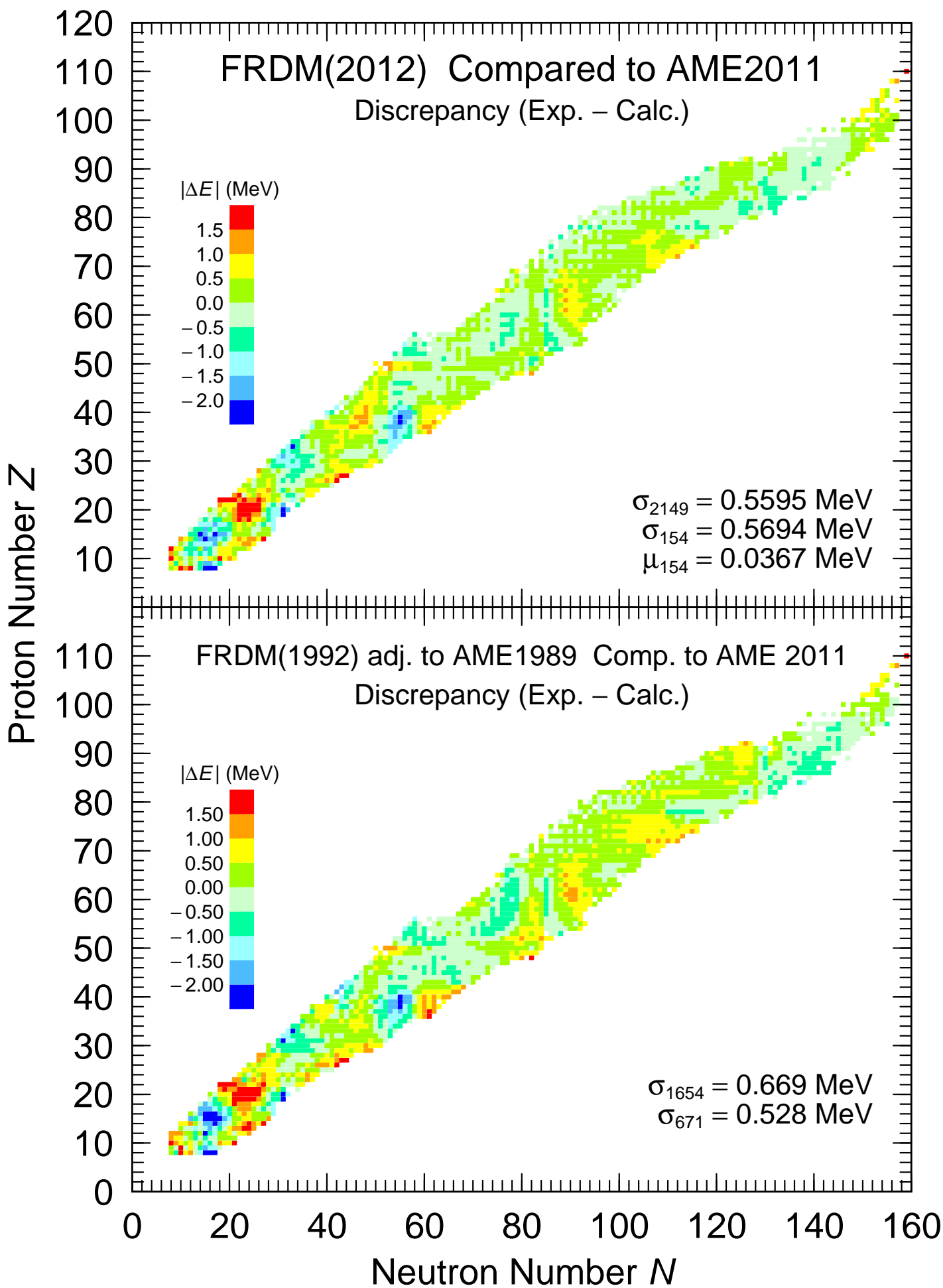
***L* variation (2009–2011)**

Accuracy improvement: 0.02 MeV

Improved gs correlation energies (2012)

Accuracy improvement: 0.01 MeV





Trap Data from Haettner et al. (PRL 106 (2011) 122501)

