

Connecting Neutron Star Observations and the Nucleon-Nucleon Interaction with Gravitational Waves

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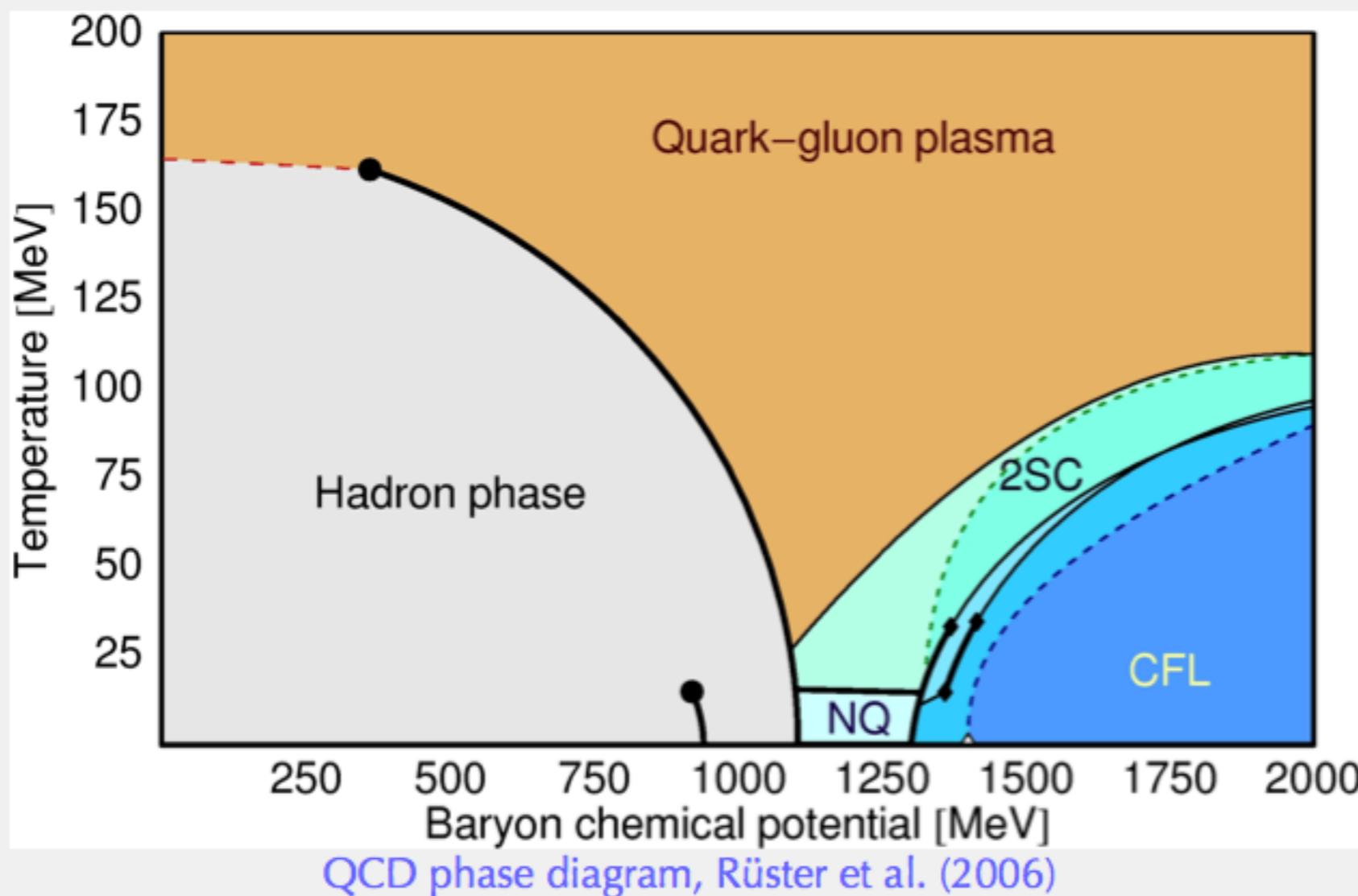
LIGO!



Outline

- Neutron star masses and radii and the EOS
- Tidal deformabilities
- Pulsar glitches and moments of inertia
- Neutron drip line

Fundamental Questions



Frontiers of nuclear science (2007):

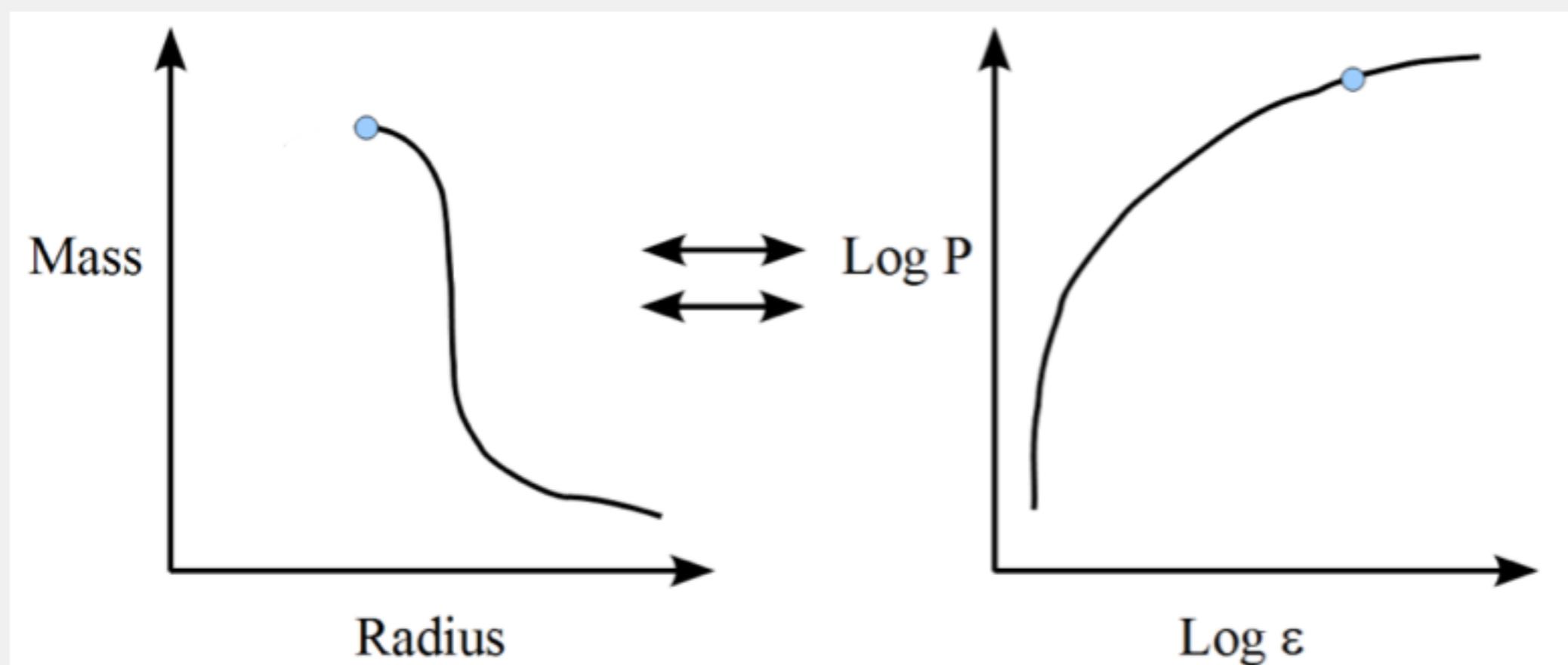
- What is the nature of neutron stars and dense matter?
- What is the nucleon-nucleon interaction?
- What is the origin of the elements?

Also,

- How do neutron stars merge?
- How do mergers generate gravitational waves?

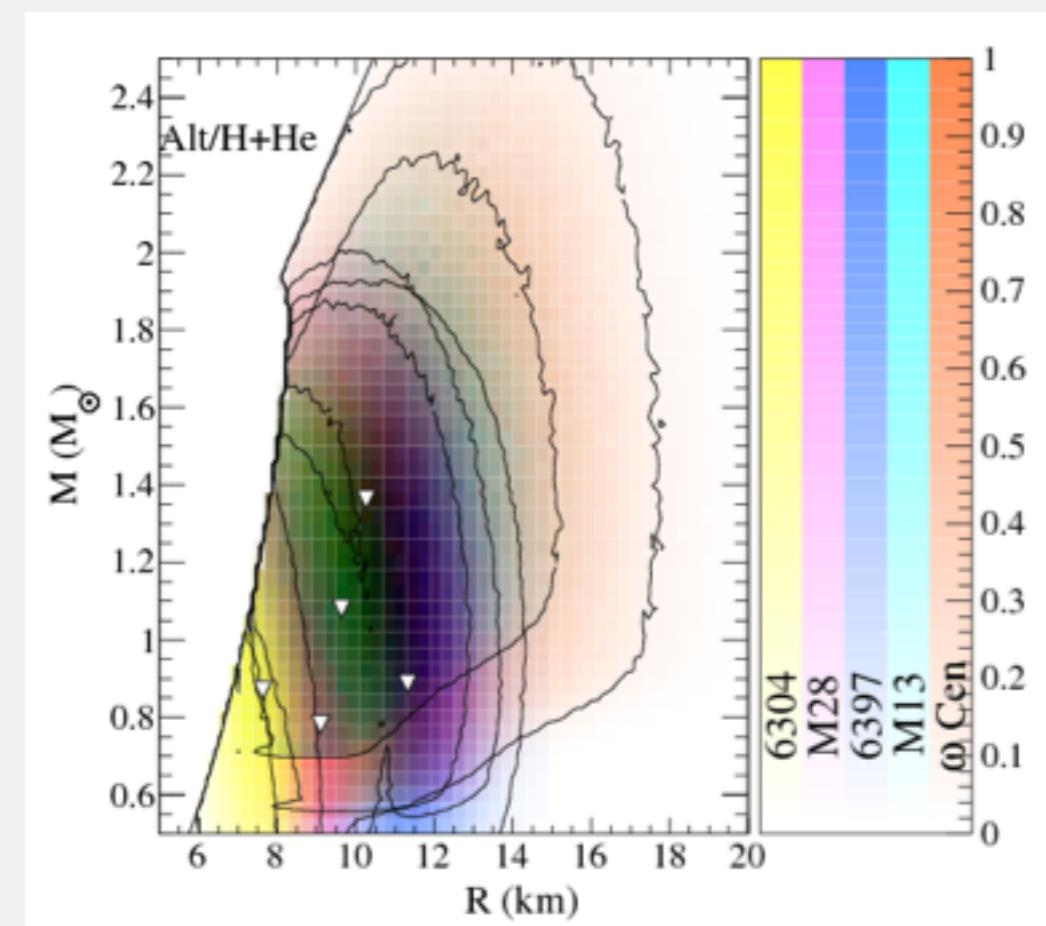
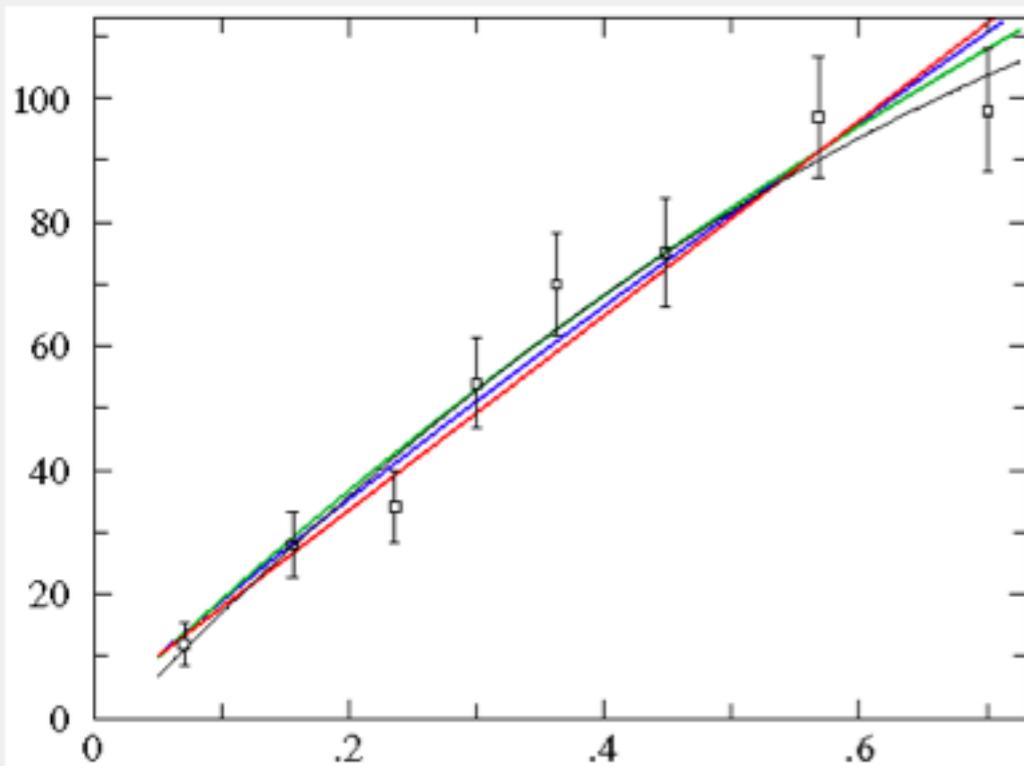
Neutron Star Masses and Radii and the EOS

- Neutron stars (to better than 10%) all lie on one universal mass-radius curve
(Largest correction is rotation - work in progress)
- Recent measurement of two $2 M_{\odot}$ neutron stars
Demorest et al. (2010), Antoniadis et al. (2013)
- As of 2007 neutron star radii constrained to 8-15 km, now 10-13 km
Lattimer and Prakash (2007); Steiner, Lattimer and Brown (2013)



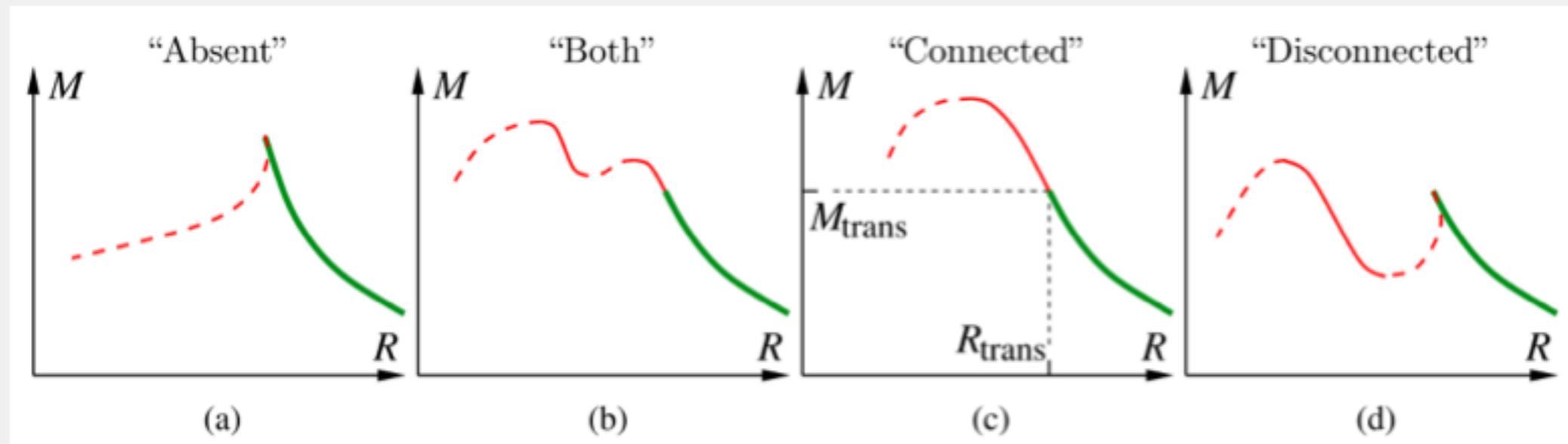
- Einstein's field equations provide a 1-1 correspondence
- Formally an underconstrained problem, but effectively over constrained if you have enough precise data (we don't yet)

χ^2 fits



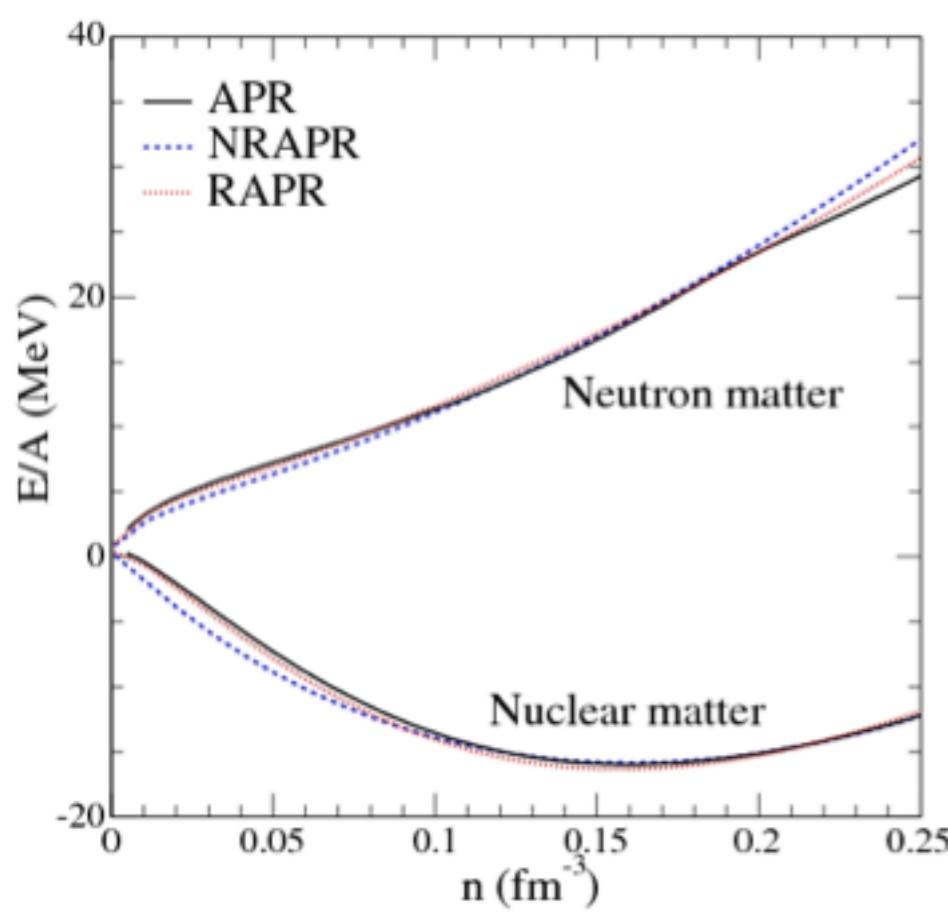
Lattimer and Steiner (2014)

- Straightforward when there's no uncertainty in the horizontal axis
- When there is: e.g. "Deming regression"
- M-R curve is not a function (central pressure is a better "independent" variable)



Alford, Han, and Prakash (2014)

EOS parameterization



Steiner et al. (2005)

Nucleonic matter:

$$\begin{aligned}\epsilon &= B + \frac{K}{18n_0^2} (n - n_0)^2 \\ &\quad + (1 - 2x)^2 (17 \text{ MeV}) \left(\frac{n}{n_0} \right)^{2/3} \\ &\quad + (1 - 2x)^2 (S - 17 \text{ MeV}) \left(\frac{n}{n_0} \right)^\gamma\end{aligned}$$

- Parameters: B, K, S , and γ
- $S(n_B) \equiv E_{\text{neut}}(n_B) - E_{\text{nuc}}(n_B)$
- $S = S(n_0); L = 3n_0 S'(n_0)$

High densities:

- Polytropes: $P(\epsilon) = K\epsilon^\Gamma$ with $\Gamma \equiv 1 + \frac{1}{n}$

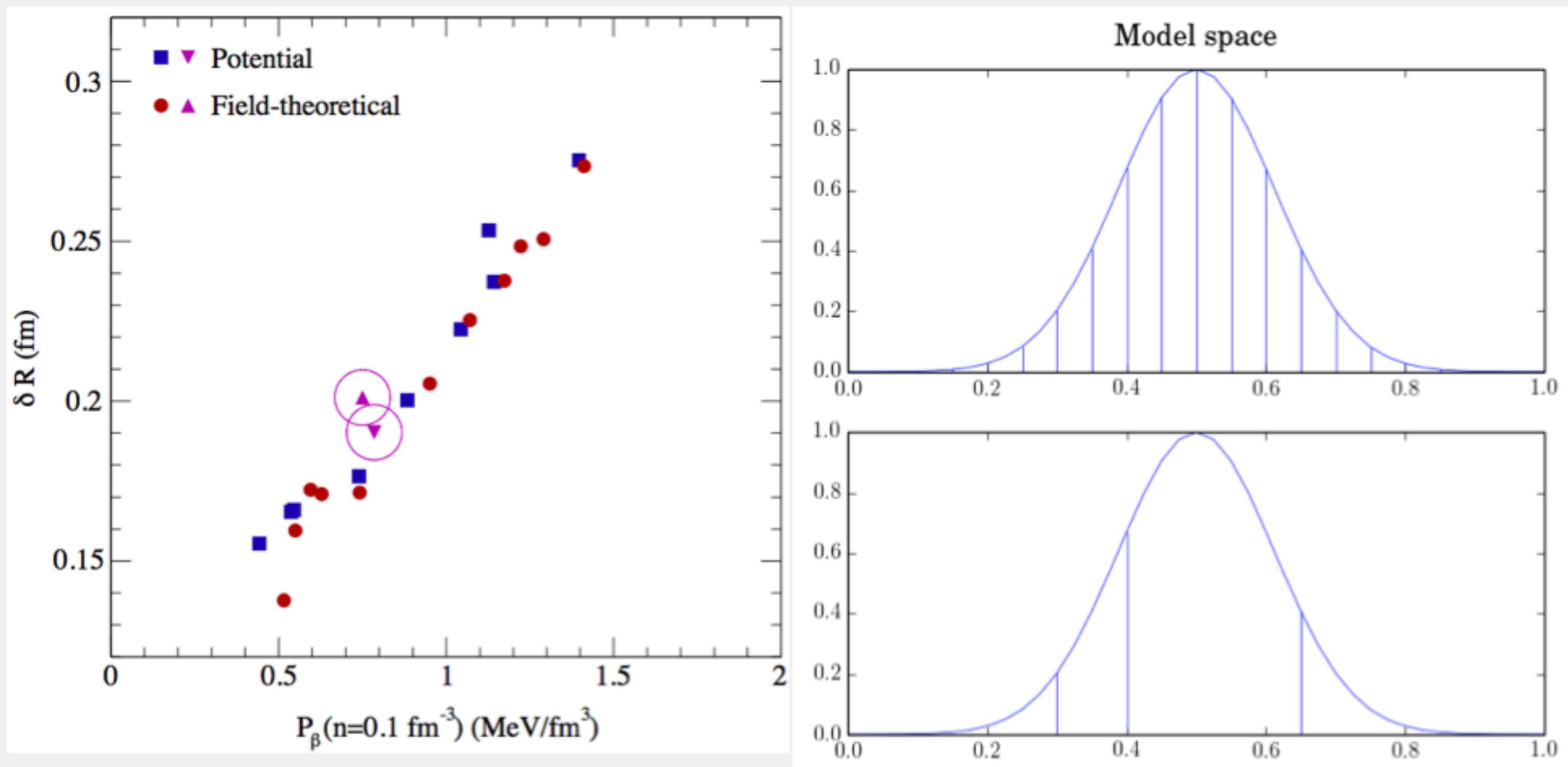
crust | ϵ_{trans} | nucleons | ϵ_1 | Polytrope 1 | ϵ_2 | Polytrope 2

- Lines: $P(\epsilon) = A_i\epsilon + P_i$

crust | ϵ_{trans} | nucleons | ϵ_1 | line 1 | ϵ_2 | line 2 | ...

Purpose of a Model

- Model comparison
Model A vs. Model B?
- Parameter estimation
However, parameters don't always have a clear physical meaning
- Predictions
Produce a probability distribution
 - Distribution over parameter space and also over model space



Radius Measurements in qLMXBs

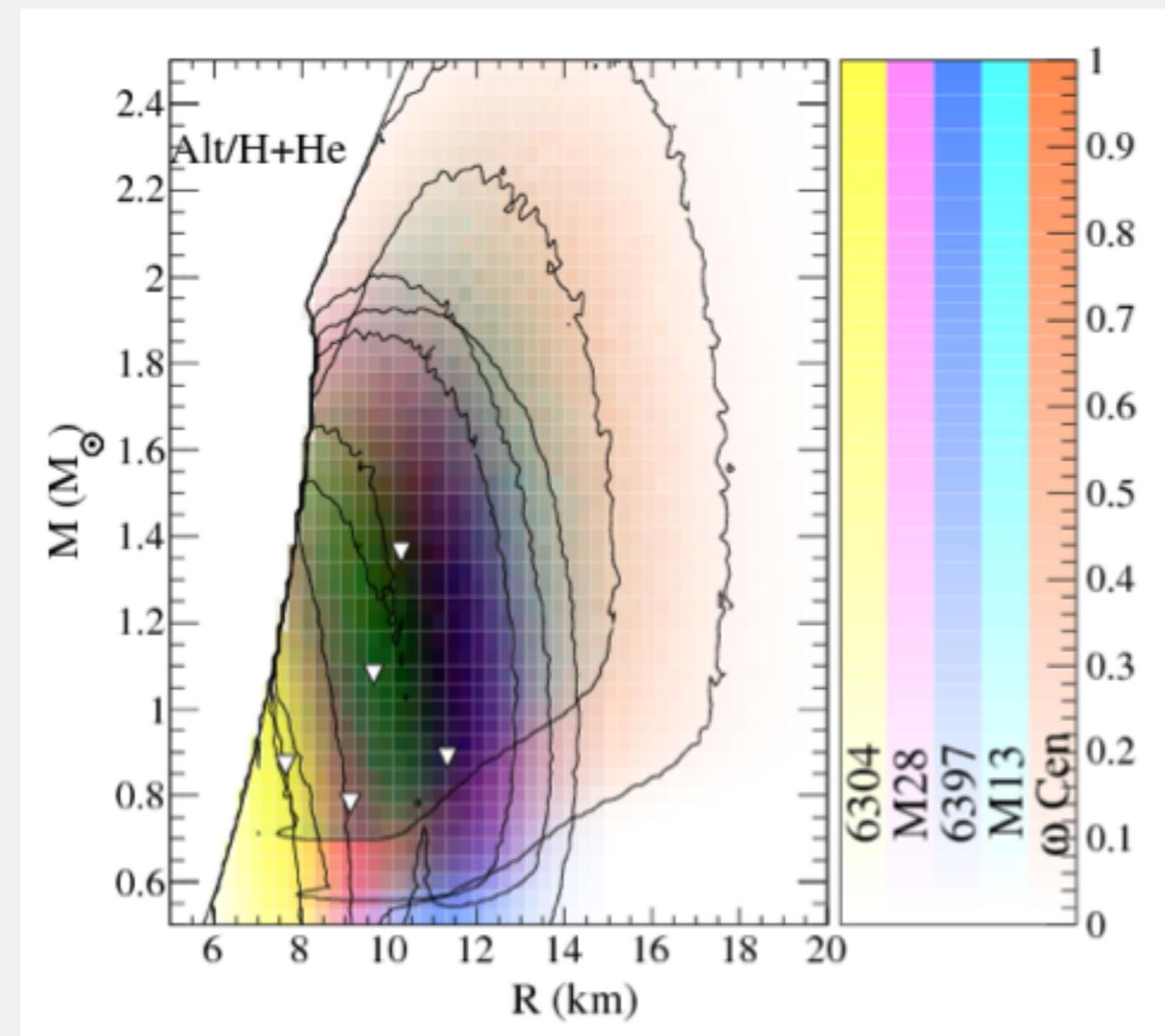
Quiescent LMXBs

- Measure flux of photons and their energy distribution
- Know distance if in a globular cluster
- Implies radius measurement

$$F \propto T_{\text{eff}}^4 \left(\frac{R_{\infty}}{D} \right)^2$$

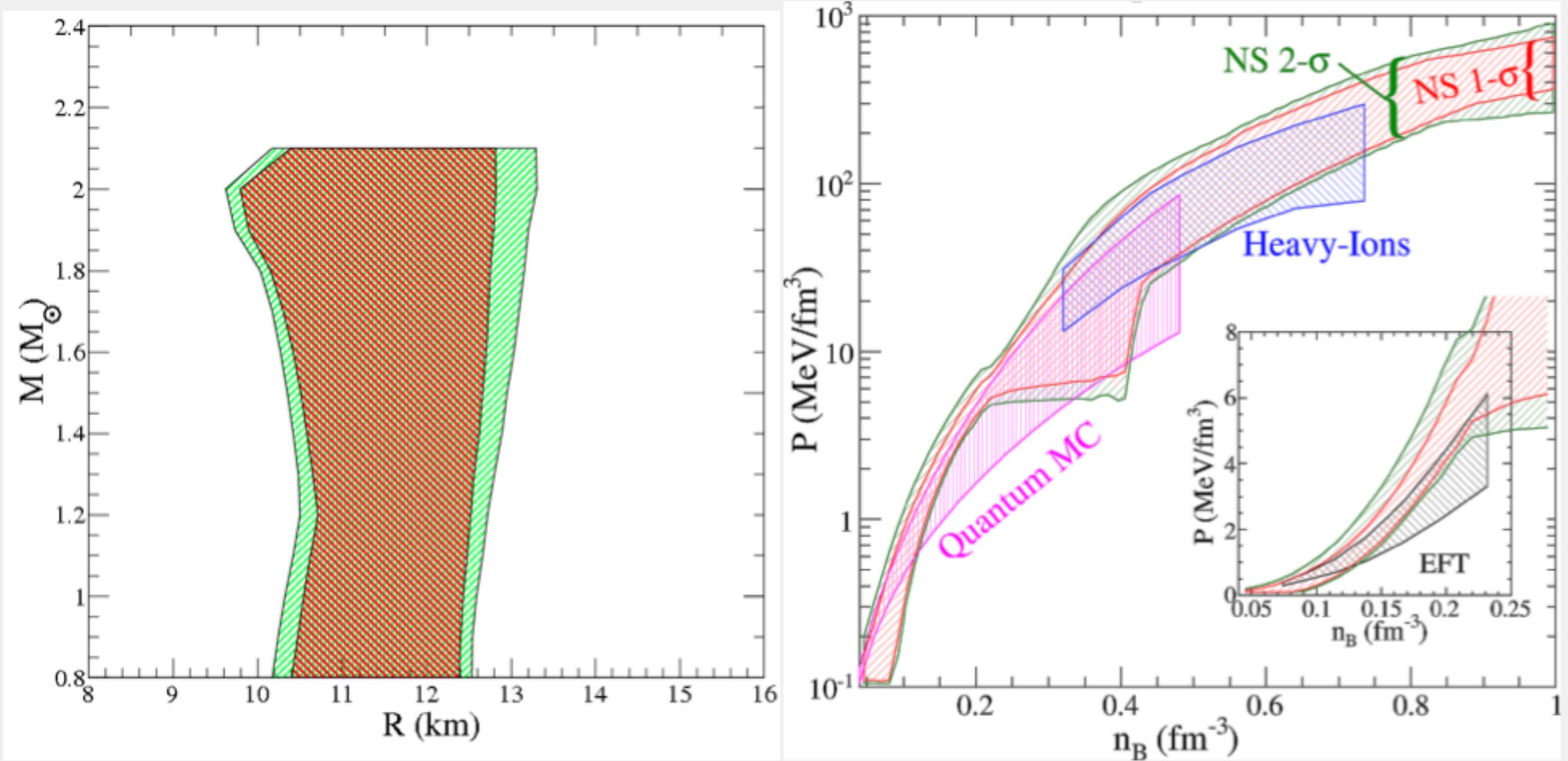
i.e. Rutledge et al. (1999)

- Many important unresolved systematics
- Need to understand X-ray absorption between source and observer
- Need information about the atmosphere, including composition
- Also need X-ray absorption and absolute flux calibration
- Inevitably give small radii for some low-mass stars



Lattimer and Steiner (2014) - Probability distributions for five neutron stars, colors added together

The M-R curve and the EOS of Dense Matter



Steiner, Lattimer, and Brown (2013); red and green outlines 68% and 95% regions

- Full Bayesian MCMC sampling of the likelihood (times prior)
- Radius of a 1.4 solar mass neutron star is 10.4 - 12.9 km
- Note the uncertainty in the EOS at a few times saturation
- These results are limited by strong systematic uncertainties
- No assumption that pressure is correlated between low and high-densities

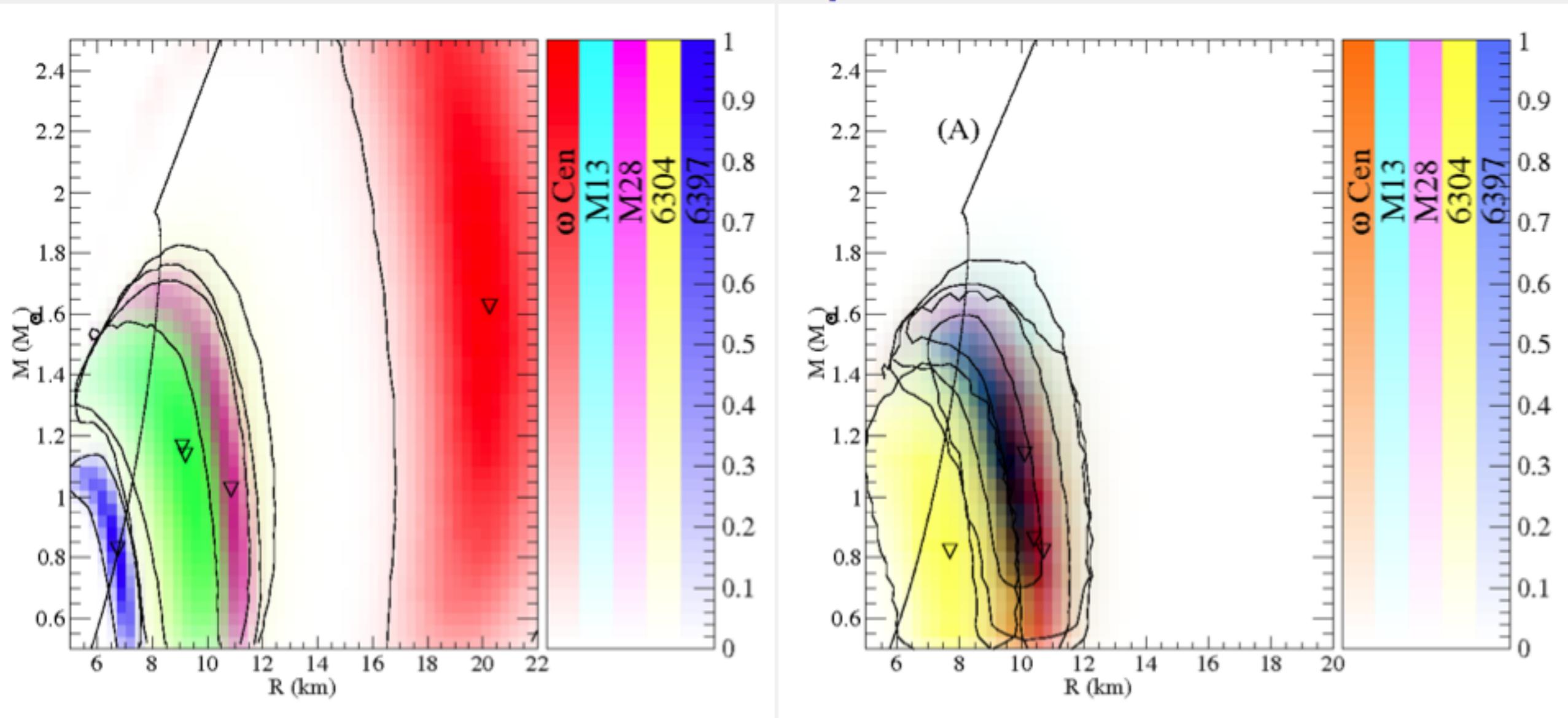
The M-R curve and the EOS of dense matter

EOS Model	Data modifications	$R_{95\%>}$	$R_{68\%>}$	$R_{68\%<}$	$R_{95\%<}$
(km)					
Variations in the EOS model					
A	-	11.18	11.49	12.07	12.33
B	-	11.23	11.53	12.17	12.45
C	-	10.63	10.88	11.45	11.83
D	-	11.44	11.69	12.27	12.54
Variations in the data interpretation					
A	I	11.82	12.07	12.62	12.89
A	II	10.42	10.58	11.09	11.61
A	III	10.74	10.93	11.46	11.72
A	IV	10.87	11.19	11.81	12.13
A	V	10.94	11.25	11.88	12.22
A	VI	11.23	11.56	12.23	12.49
Global limits		10.42	10.58	12.62	12.89

Steiner, Lattimer, and Brown (2013)

- Critical component: trying different EOS parameterizations and different interpretations of the data
- Model C allows for strong phase transitions
- Try several different models to assess systematics

As of last year...

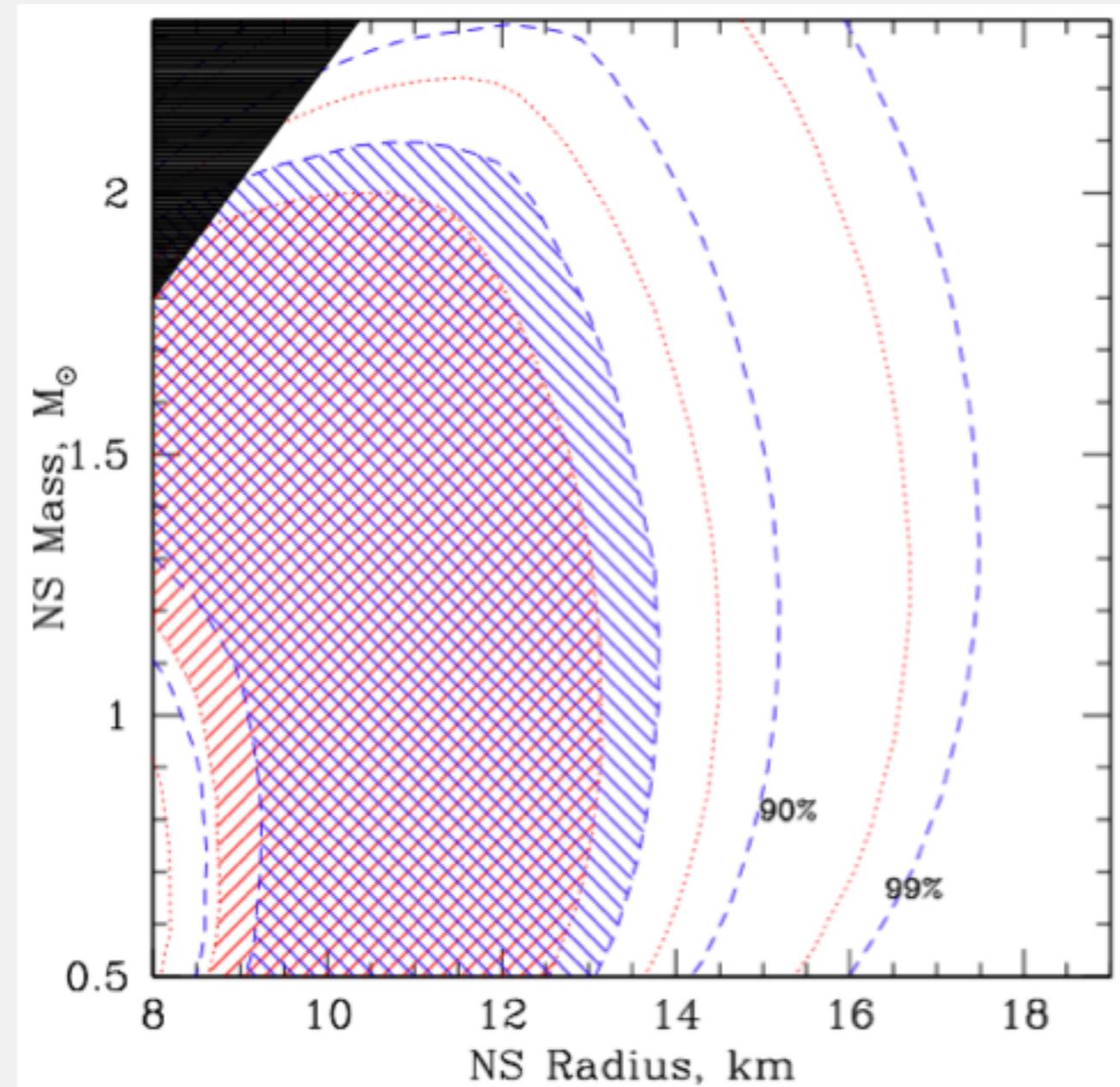


Results from Guillot et al. (2013) slightly adapted for Lattimer and Steiner (2014), relative Bayes factor of 1200 and Steiner (2014) before any assumption about the M-R curve

- R_{NS} in ω Cen : 11 km or 20 km!
- R_{NS} in NGC 6397 \sim 7 km?
- We tried different N_H values, different distances, and Helium atmospheres
- 36 separate models, millions of parameterizations

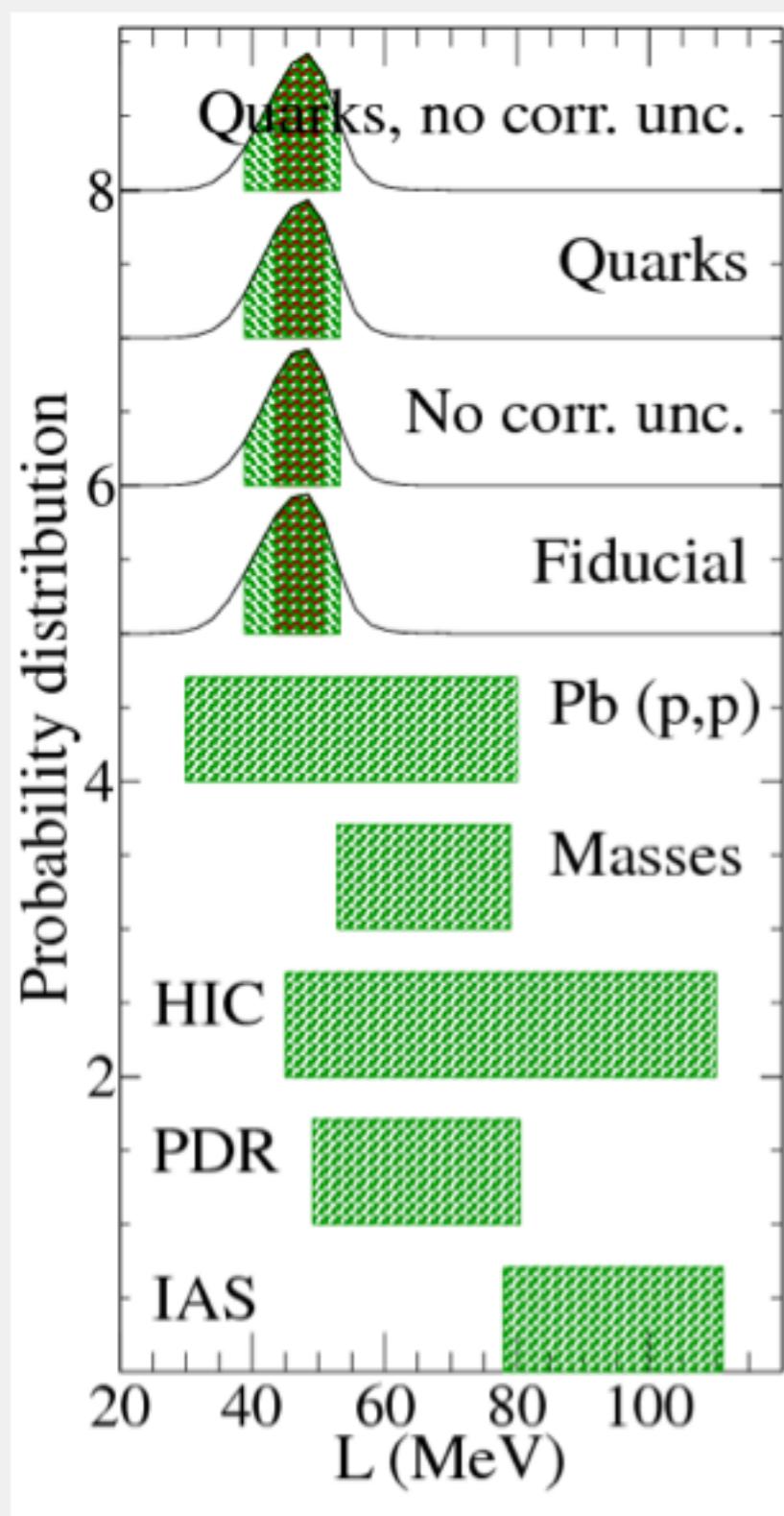
Recent Updates

- Heinke et al. (2014) confirms smaller N_H values for ω Cen with different model for the ISM and new data
- Confirmation of expectations from nuclear physics
- Radius ranges don't change that much from Steiner et al. (2013)

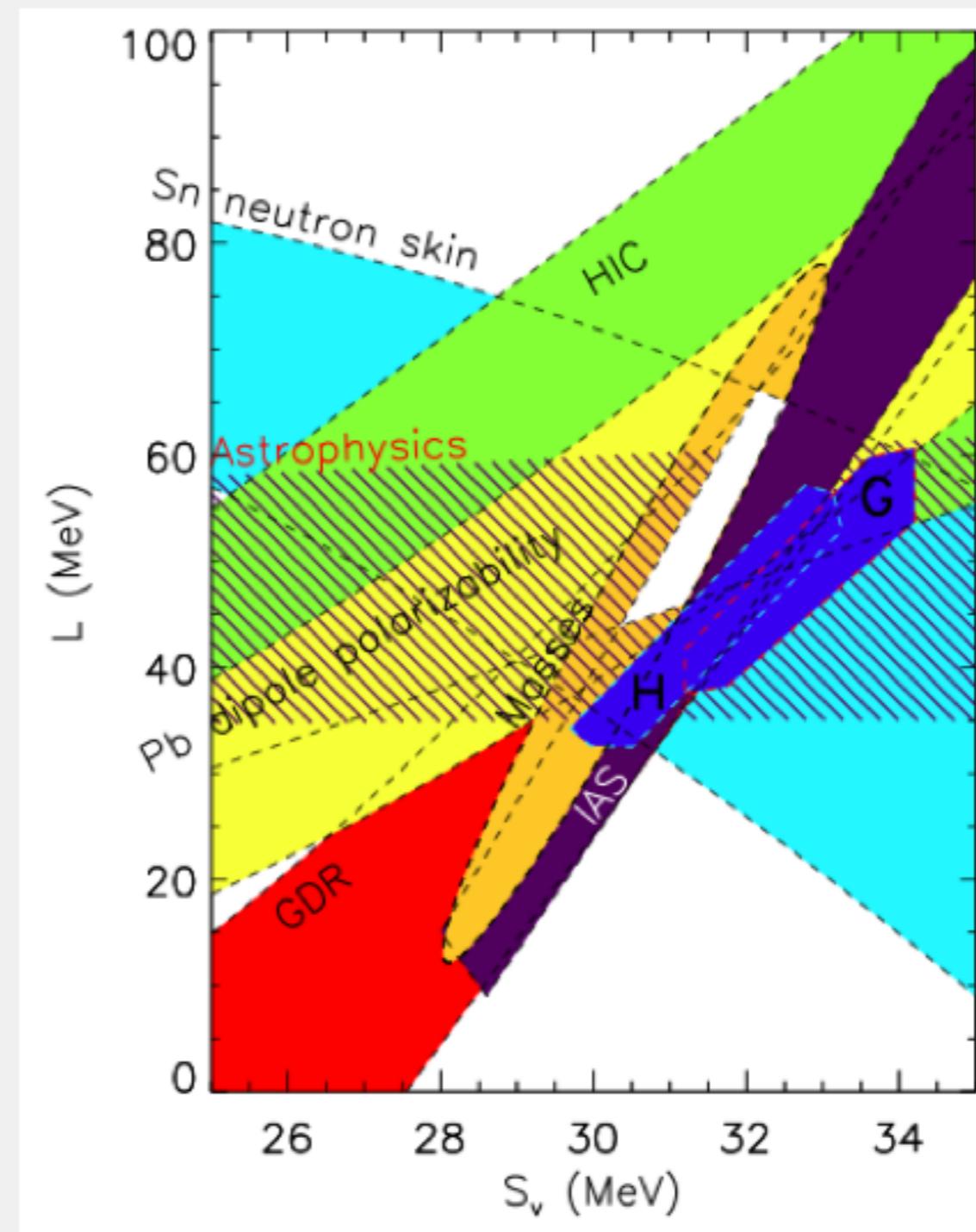


Heinke et al. (2014)

Neutron Star Constraints on L



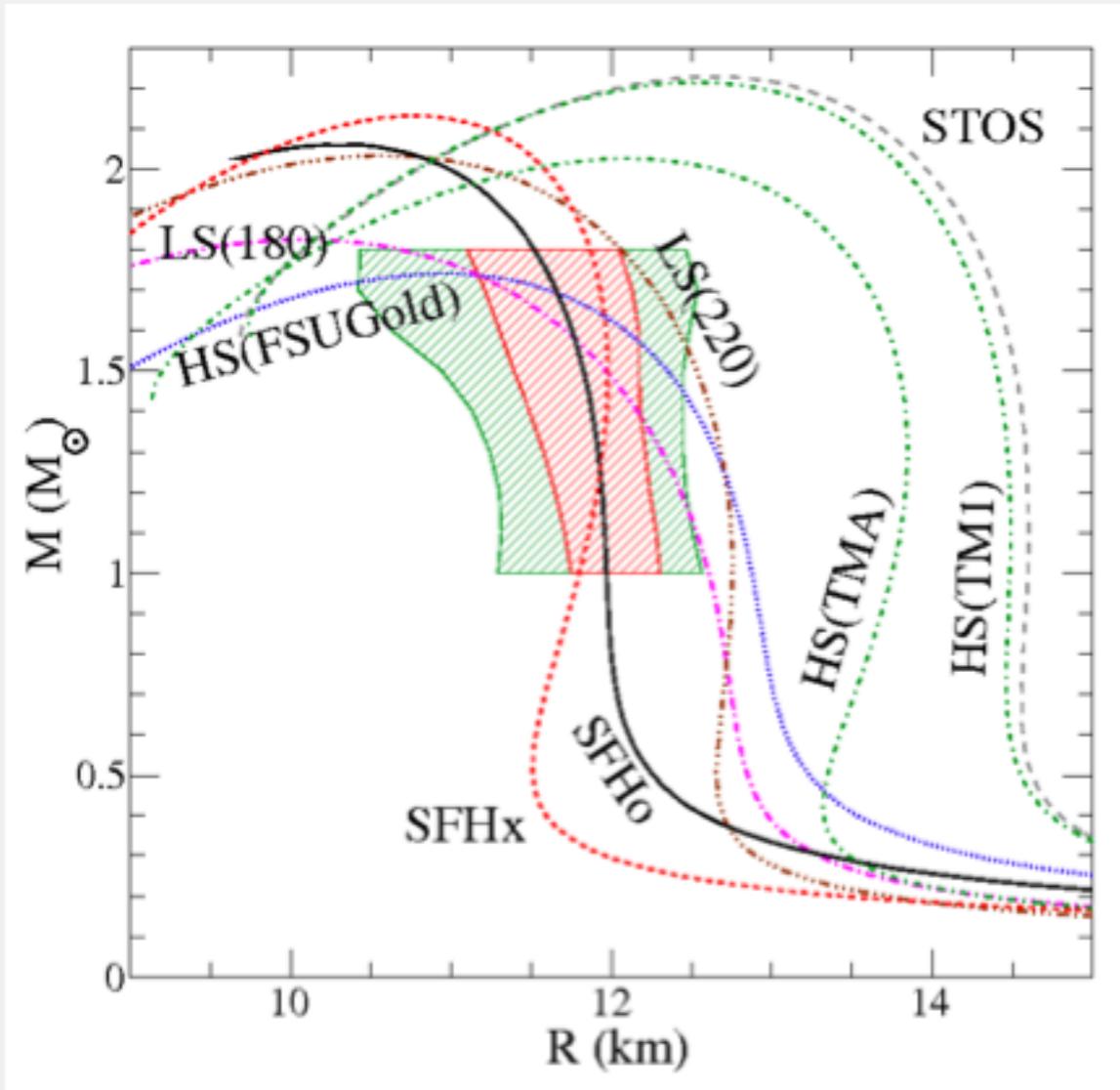
Steiner and Gandolfi (2012)
(IAS results have since come down)



Lattimer and Steiner (2014)

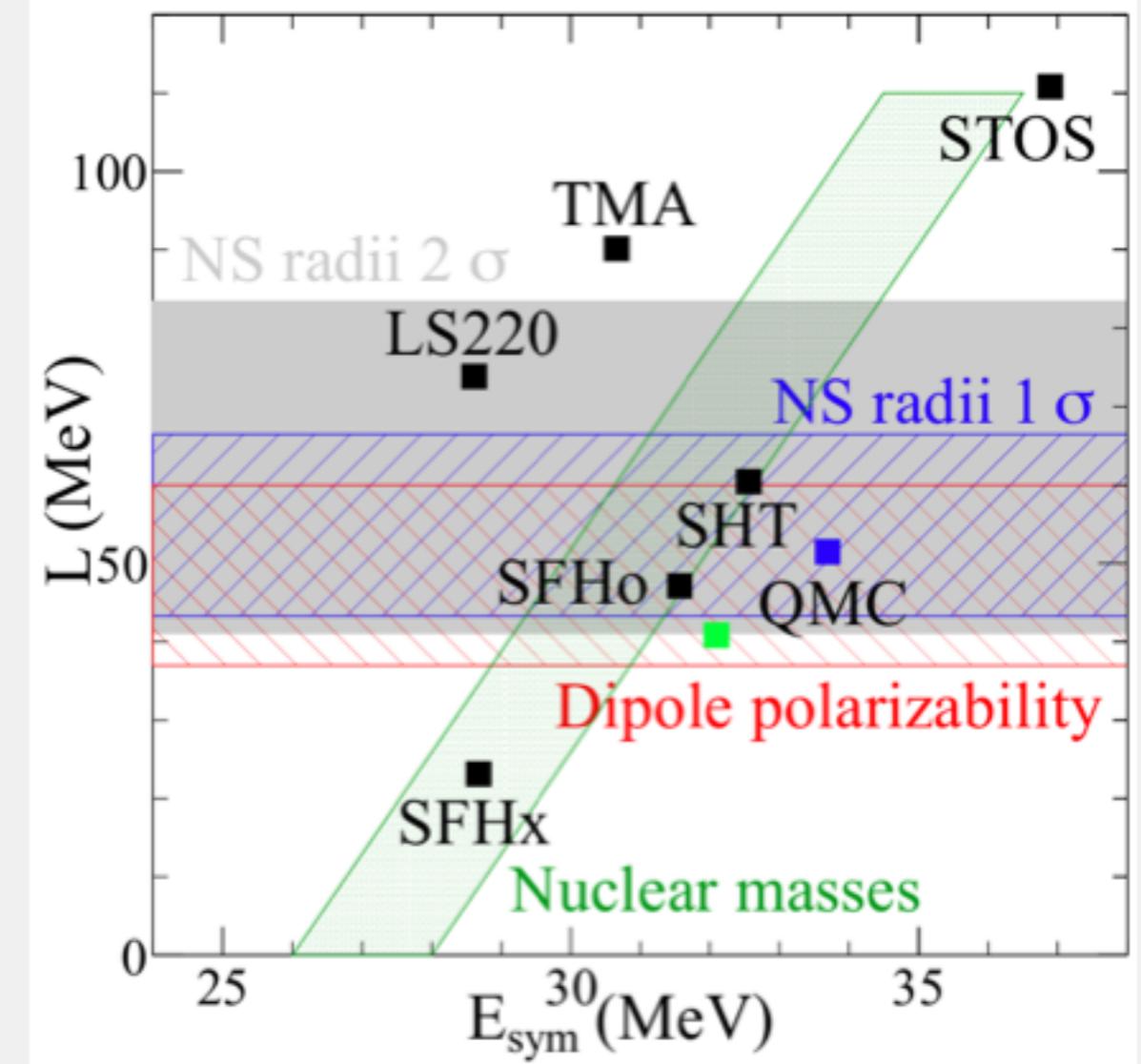
- Neutron stars strongly constrain L
- We also found $R_n - R_p < 0.2 \text{ fm}$
Confirmed by MAMI data

Supernova EOS and the Symmetry Energy



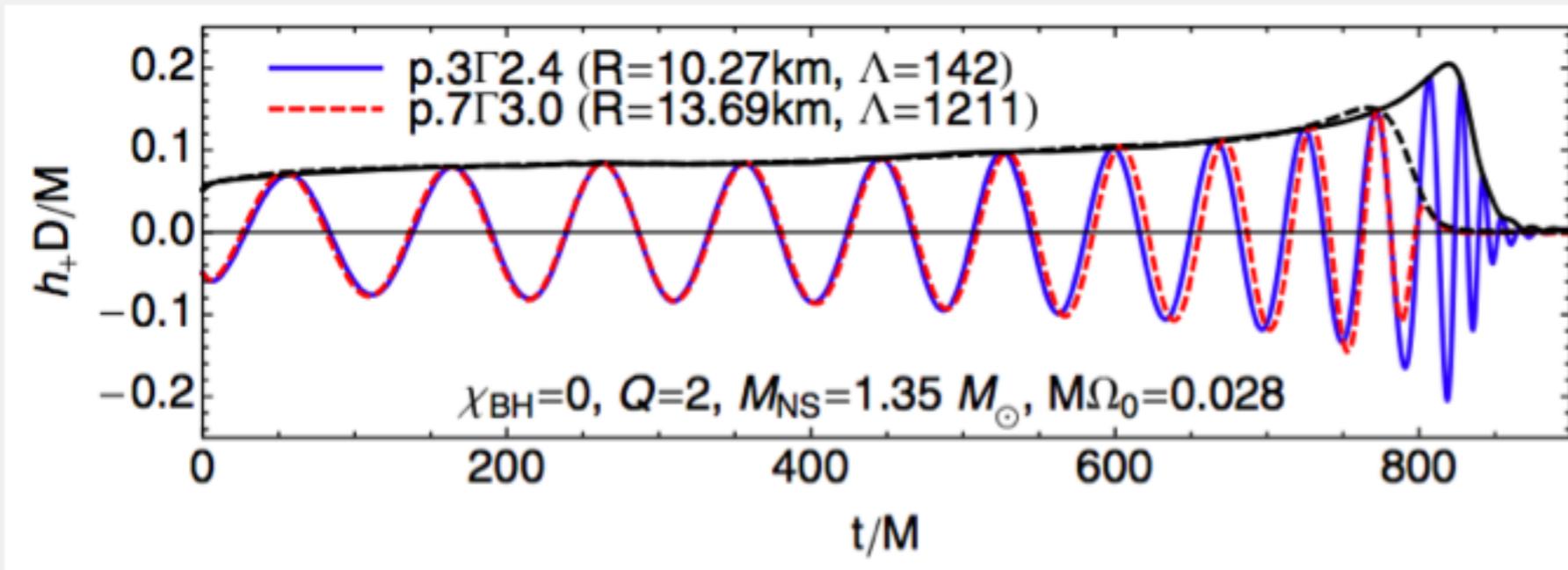
Steiner, Hempel, and Fischer (2013)

- Limited number of supernova EOSs which satisfy $M - R$ constraints and the $S - L$ correlation
- Current EOS uncertainties too small to explain explosion
- Many simulation properties are weakly correlated with the symmetry energy

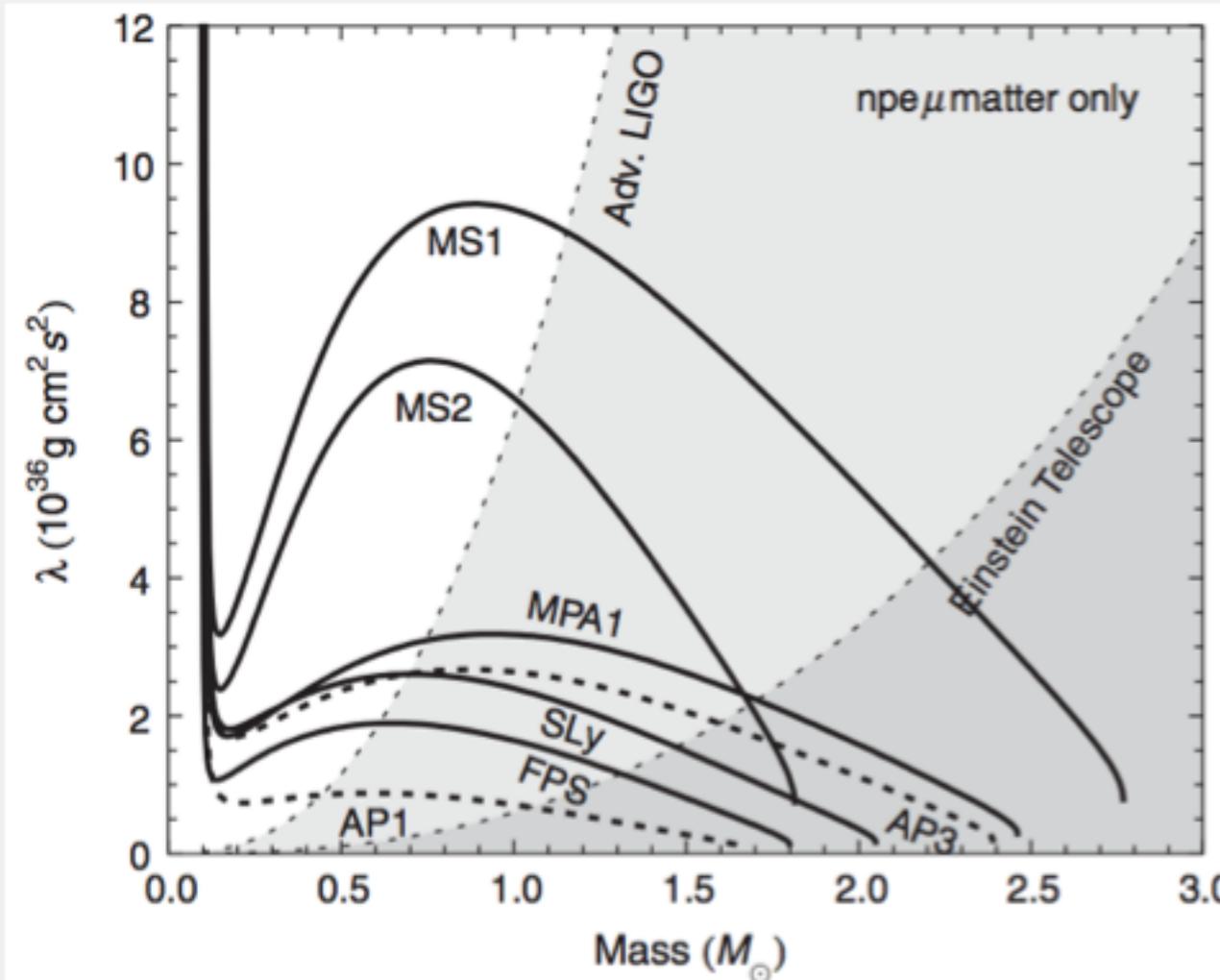


Based on Steiner, Hempel, and Fischer (2013)

Neutron Star Tidal Deformabilities



Lackey et al. (2014)

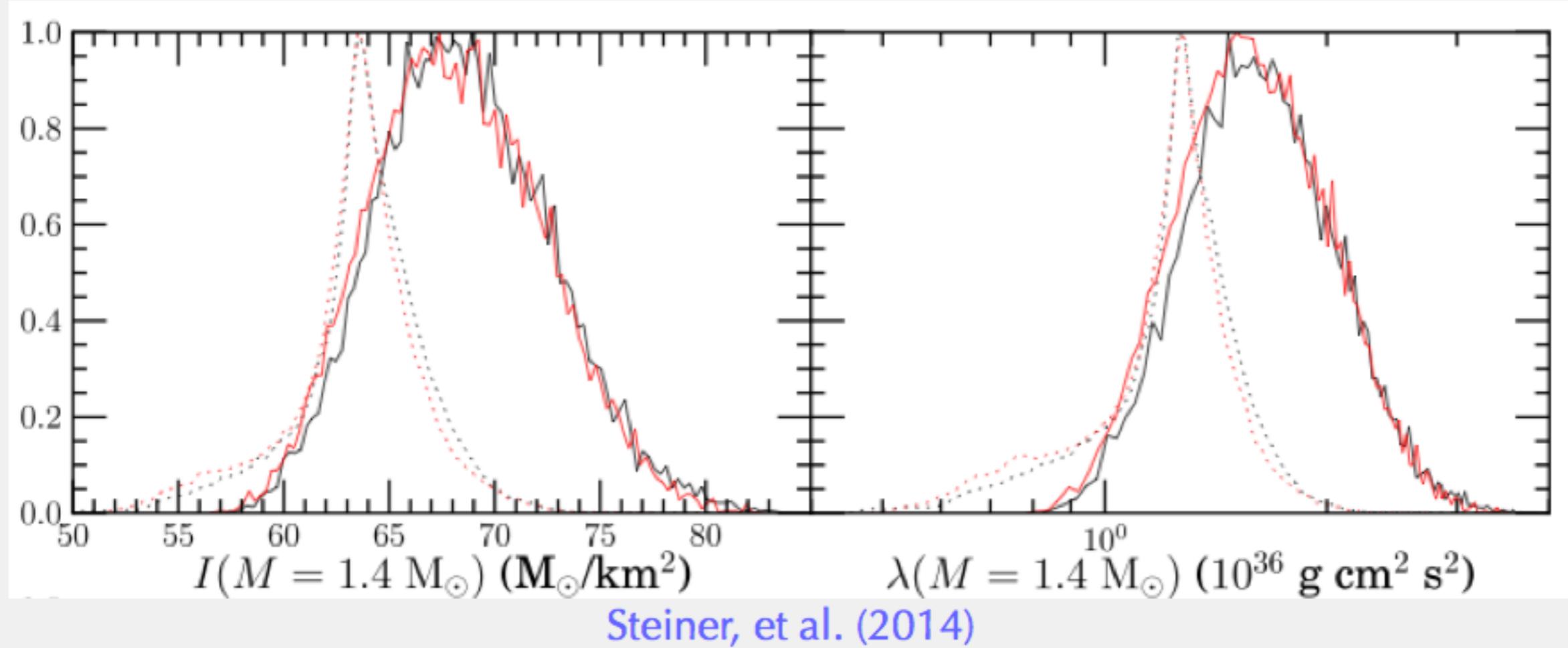


- Gravitational wave signal from an NS merger measures tidal deformability λ
- Point masses early on; deformation near 400 Hz
- Easier to detect larger tidal deformations

Hinderer et al. (2010)

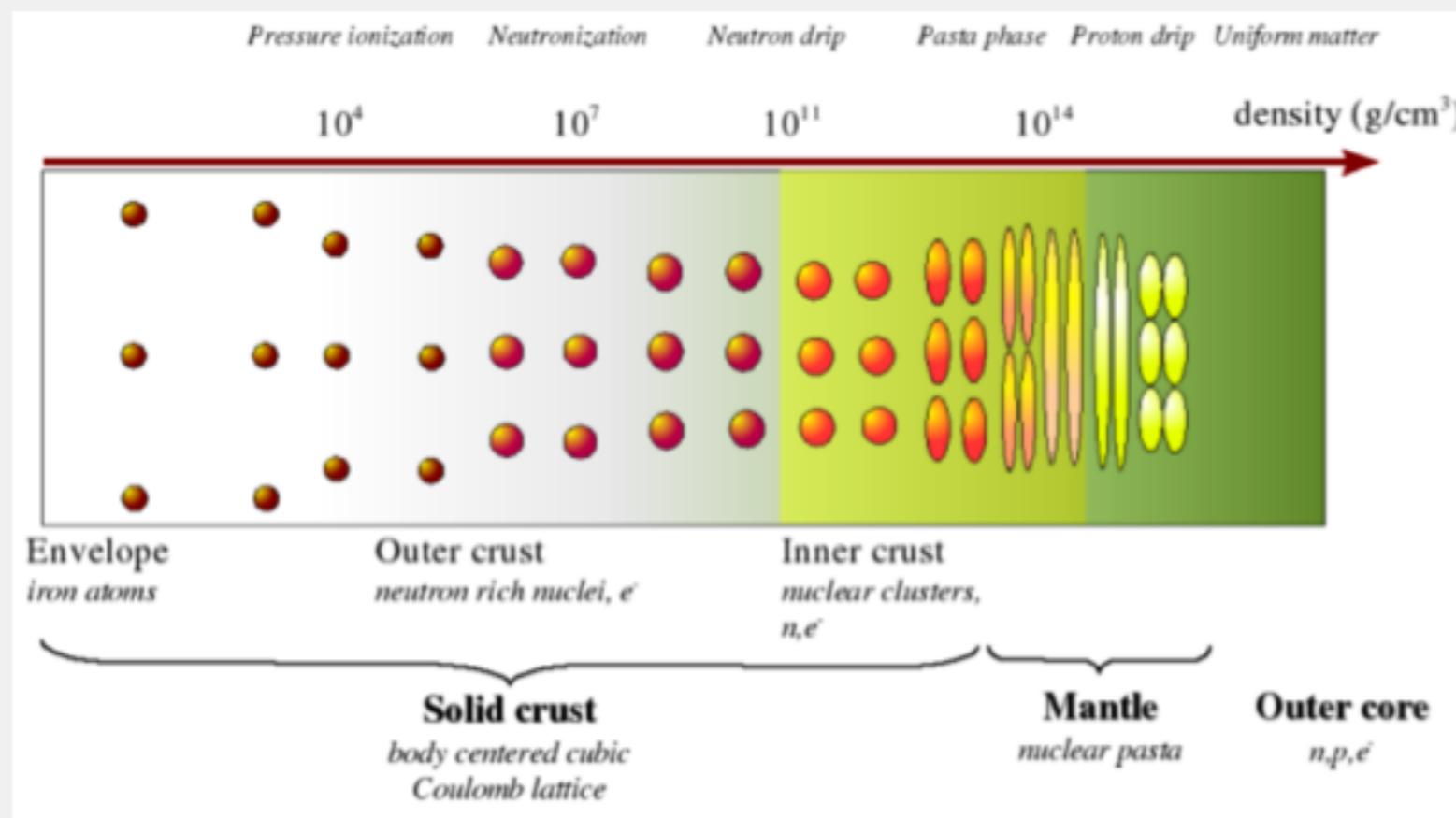
Tidal Deformabilities

- Current neutron star mass and radius observations suggest tidal deformabilities are small



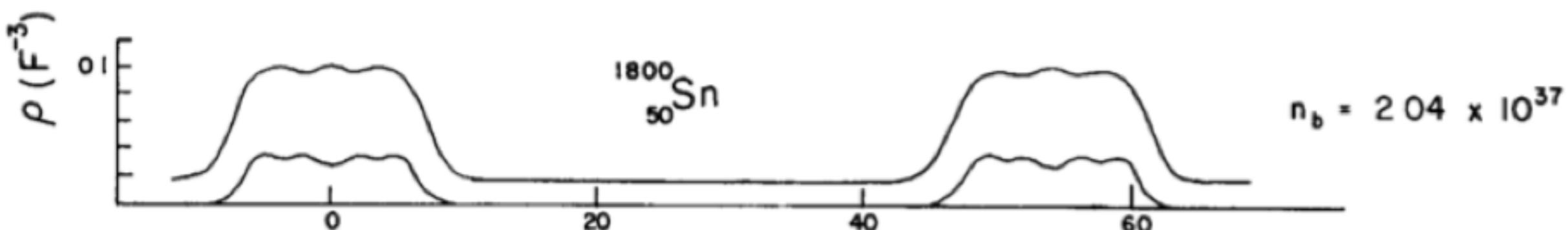
- LIGO sensitivity improved by
 - multiple measurements
[del Pozzo, et al.](#)
 - and by directly simulating the EOS
[Lackey et al.](#)
 - Using information beyond 400 Hz
[Damour and Nagar](#)
- For large enough signal-to-noise, LIGO may have smaller systematics

Structure of Matter in the Neutron Star Crust



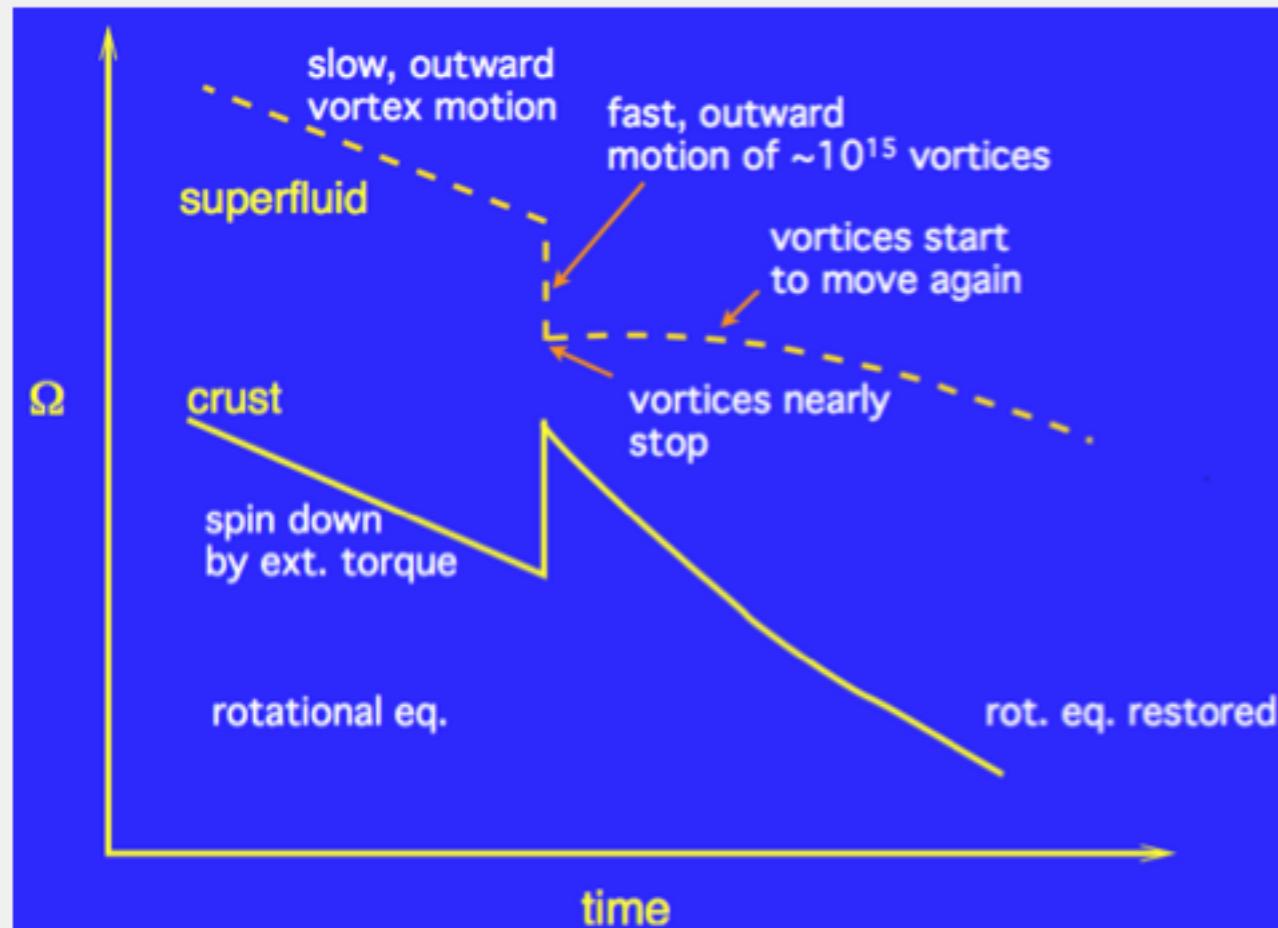
Picture from N. Chamel

- Neutron-rich nuclei
- Sea of superfluid neutrons
- Crust-core transition



Negele and Vautherin (1973!)

Pulsar Glitch Mechanism



Picture from B. Link

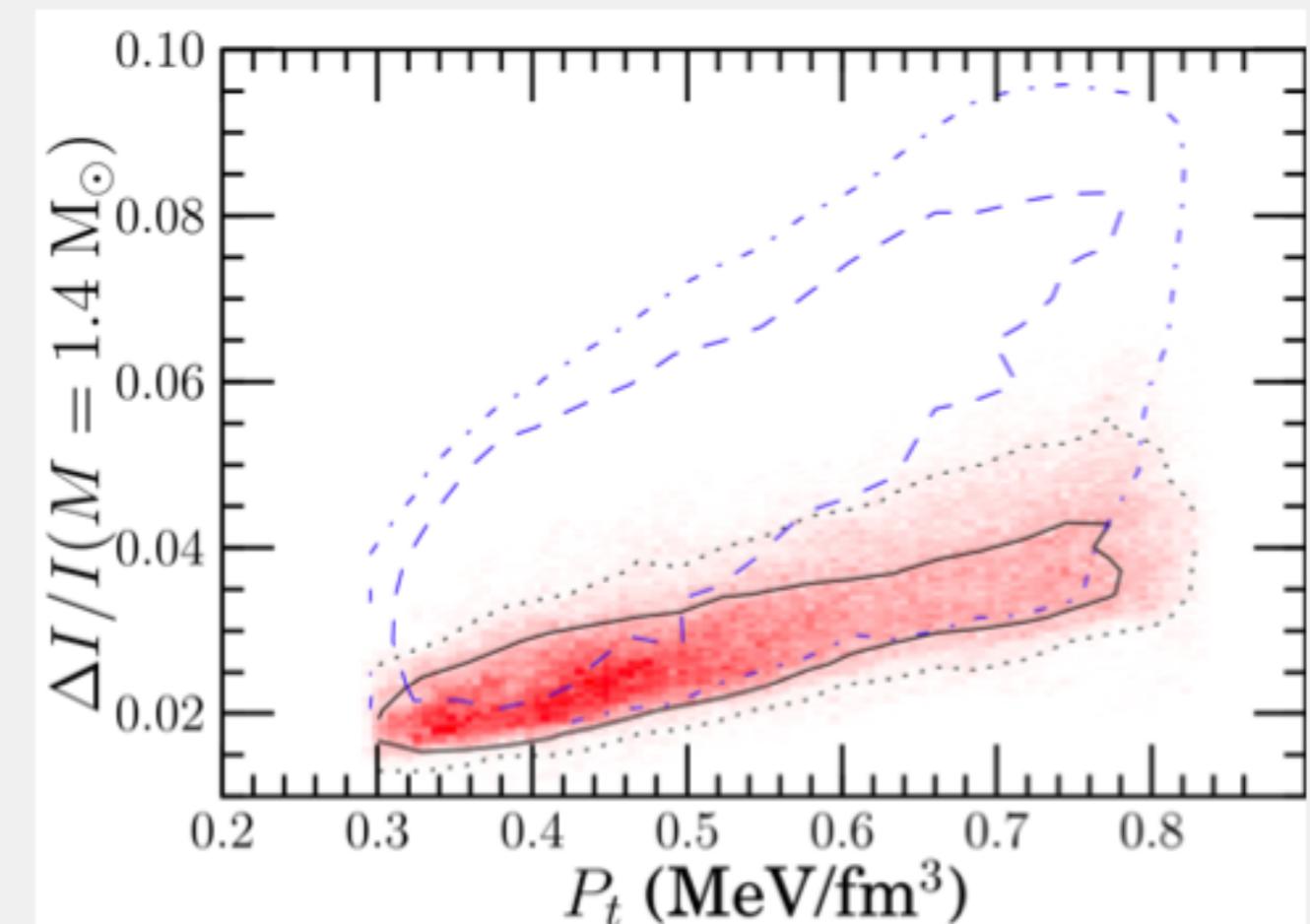
- Superfluid component, decoupled from rotation at the surface
- Natural to associate the superfluid component with the superfluid neutrons in the crust
- What is the mechanism for the sudden change?

- Superfluid vortices pinned to the lattice
- Neutron star spins down, vortices bend creating tension, eventually they must shift lattice sites
- Quasi-free neutrons are entrained with the lattice

Chamel 2012, Chamel et al. 2013

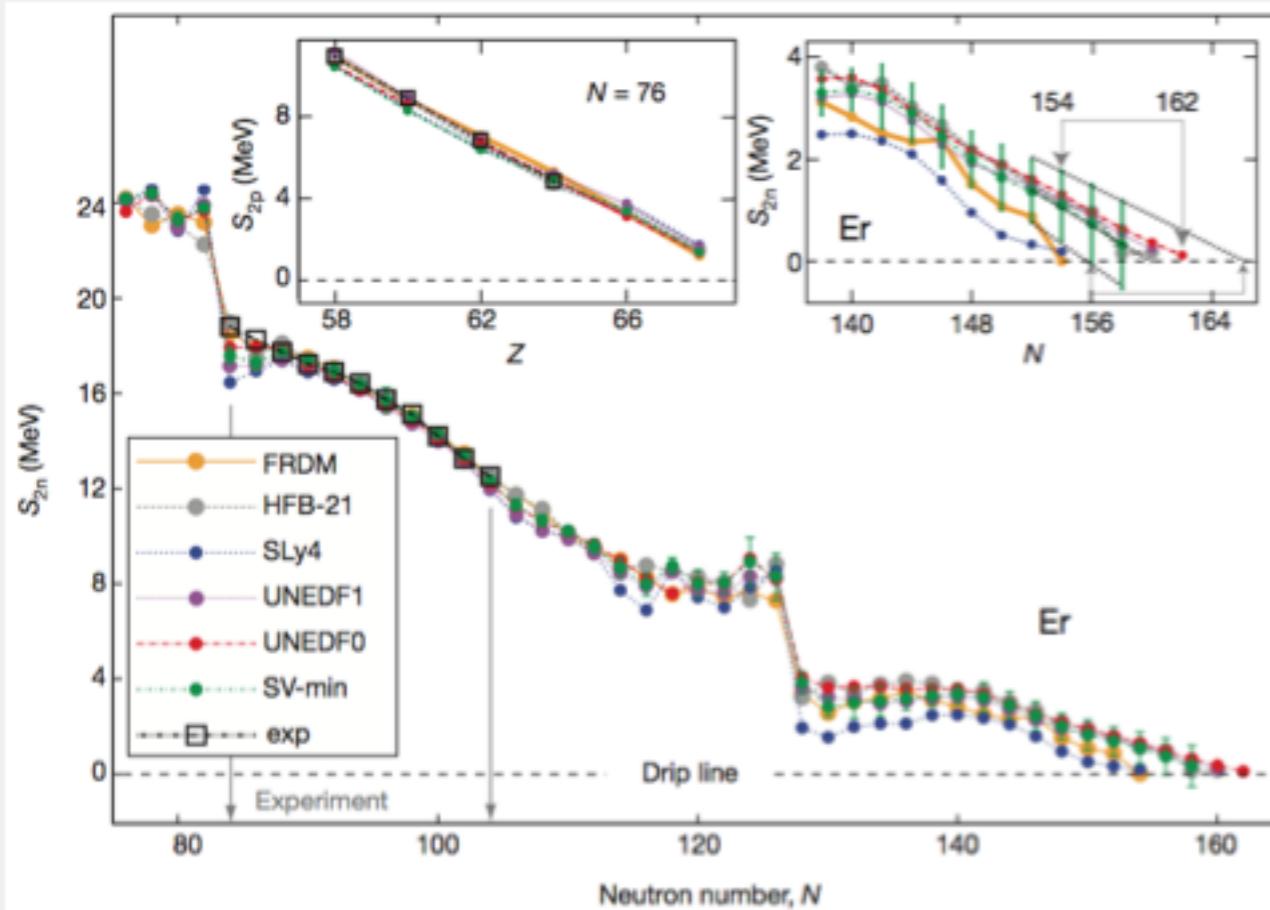
Is There Enough Superfluid in the Crust?

- We require 1.6% of I to explain glitches in Vela
[Link, Epstein, and Lattimer \(1999\)](#)
- Entrainment: 75-85% of otherwise superfluid neutrons 'connected' to the lattice
[N. Chamel \(2012\)](#)
- Current M and R observations suggest there is not enough I in the crust
[See Andersson et al. \(2012\)](#)
- Unless the systematics force much larger neutron star radii and P_t is large

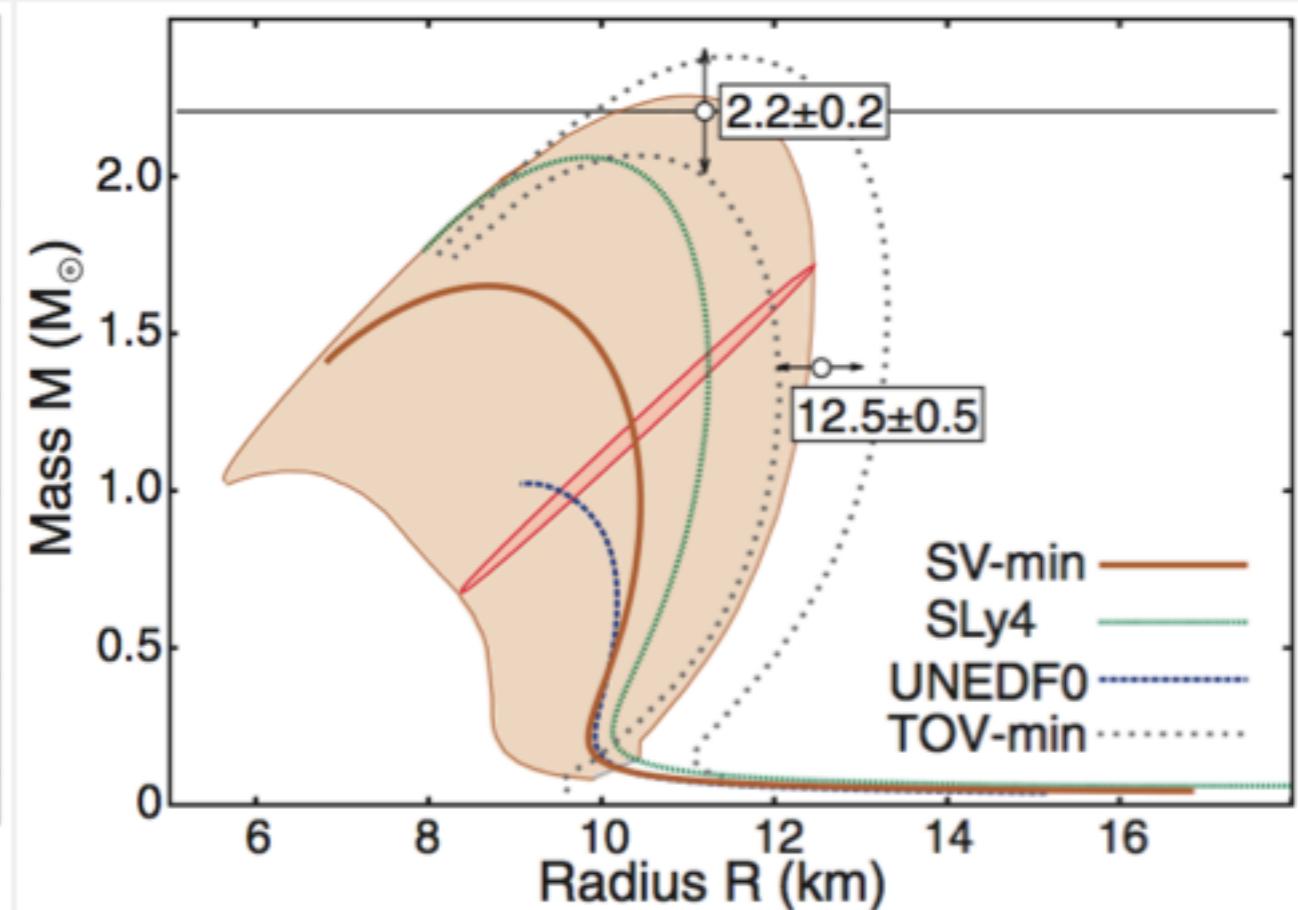


[Steiner et al. \(2014\)](#); black and red are with M & R observations, blue contours are with $I = 70 \text{ M}_\odot \text{ km}^2$

Neutron Star Radii and the Neutron Drip Line



Erler et al. (2012)



Erler et al. (2013)

- It might be possible to constrain the neutron drip line from neutron radii
- But the connection between these observables has not been fully explored
- In particular:
 - The isovector part of the Skyrme functional is incomplete
Sheikh et al. (2014)
 - Current work assumes low- and high-density matter are correlated

Summary

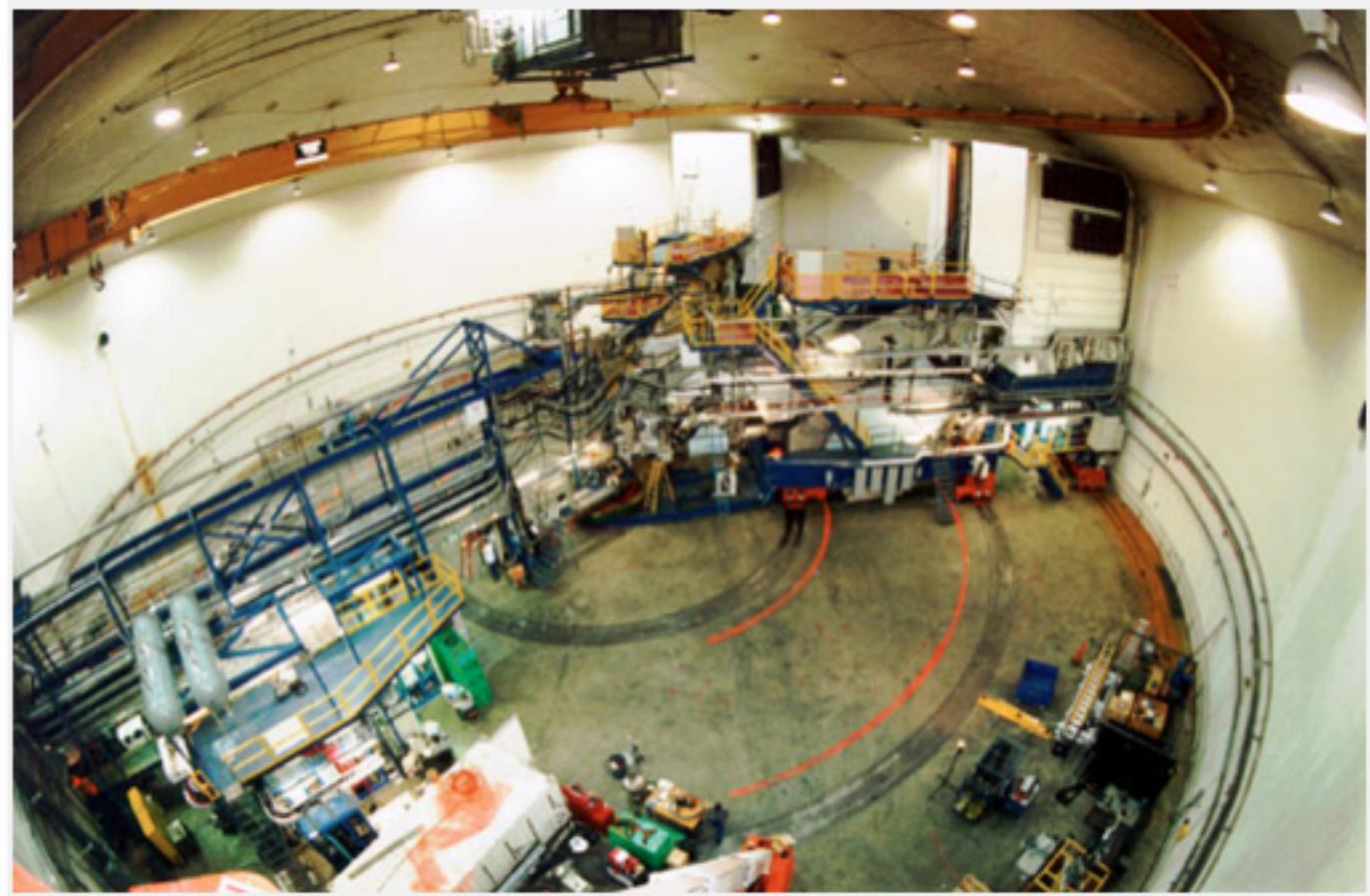
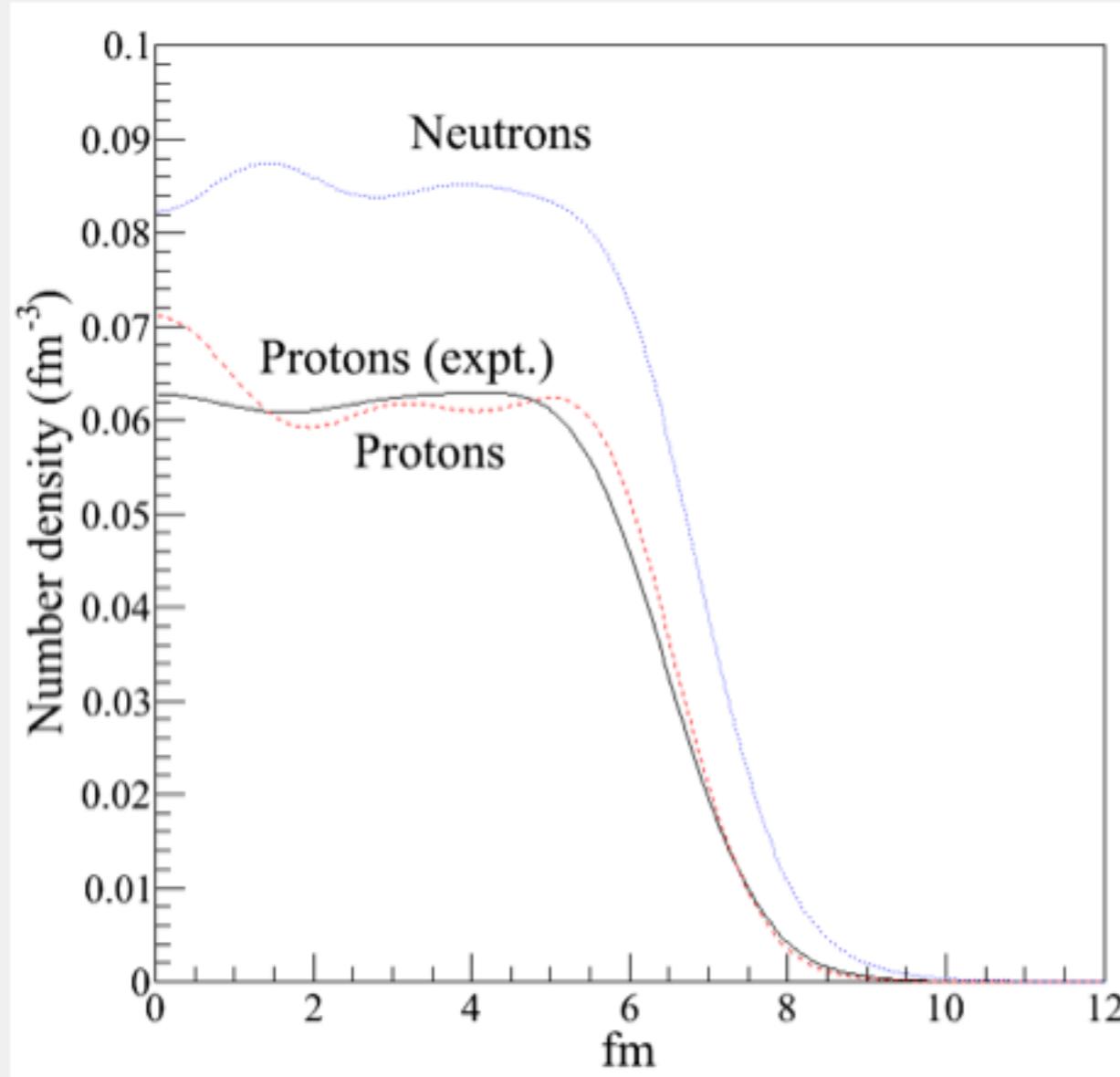
- Currently available neutron star mass and radius observations constrain the universal neutron star $M - R$ curve
 - Neutron star radii are likely between 10.4 and 12.9 km
 - We now have constraints on the EOS
 - $60 < I < 75 M_{\odot} \text{ km}^2$
 - $1 < \lambda < 3 \times 10^{36} g \text{ cm}^2 \text{ s}^2$
- Constrain the nucleon-nucleon interaction and QCD.
 - $(41) 43 \text{ MeV} < L < 67 (83) \text{ MeV}$
- Current observations imply there is not enough I to explain glitches
- Potential for further connections and new observations

The Neutron Skin Thickness of Lead

- Lead-208: 82 protons, 126 neutrons

$$R_n^2 \equiv \int r^2 n_n(r) d^3r \quad R_p^2 \equiv \int r^2 n_p(r) d^3r$$

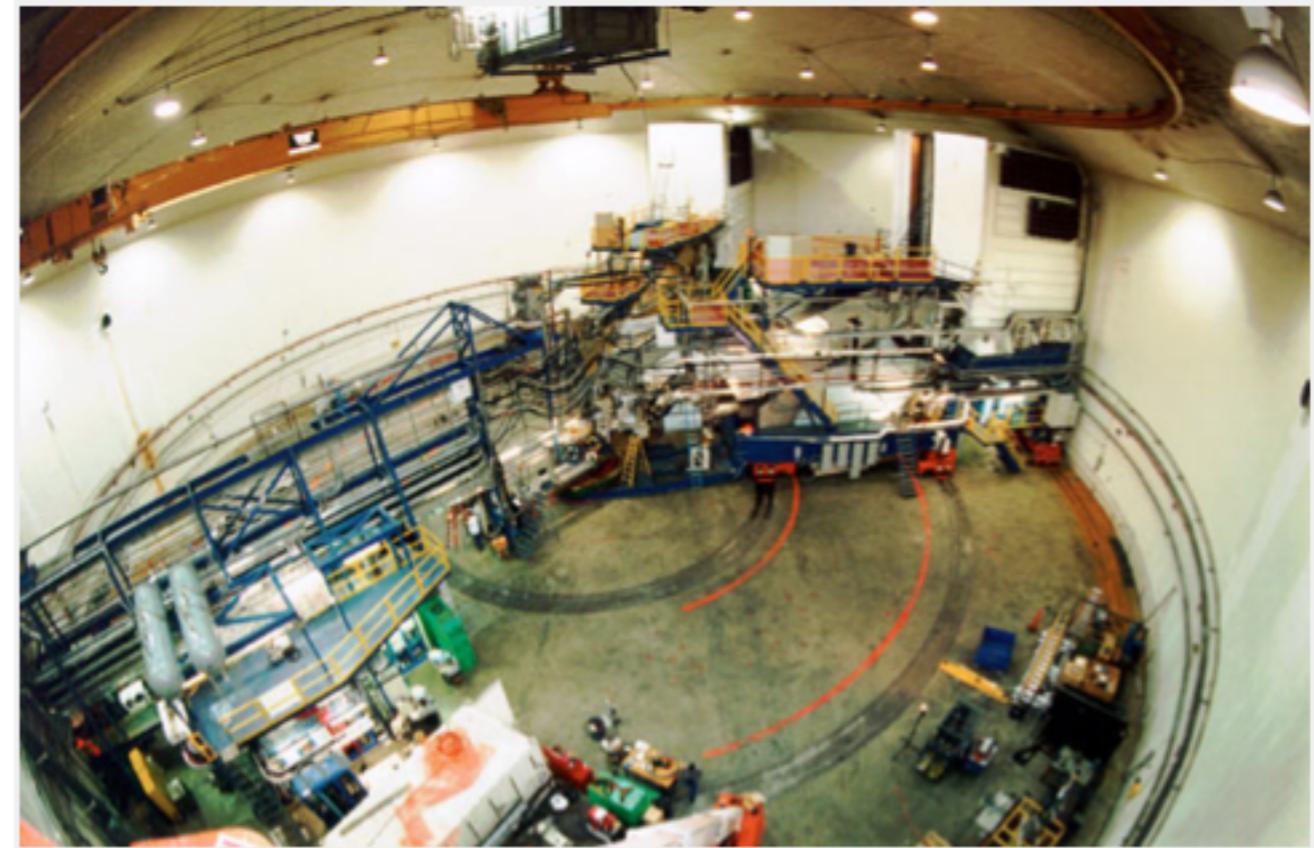
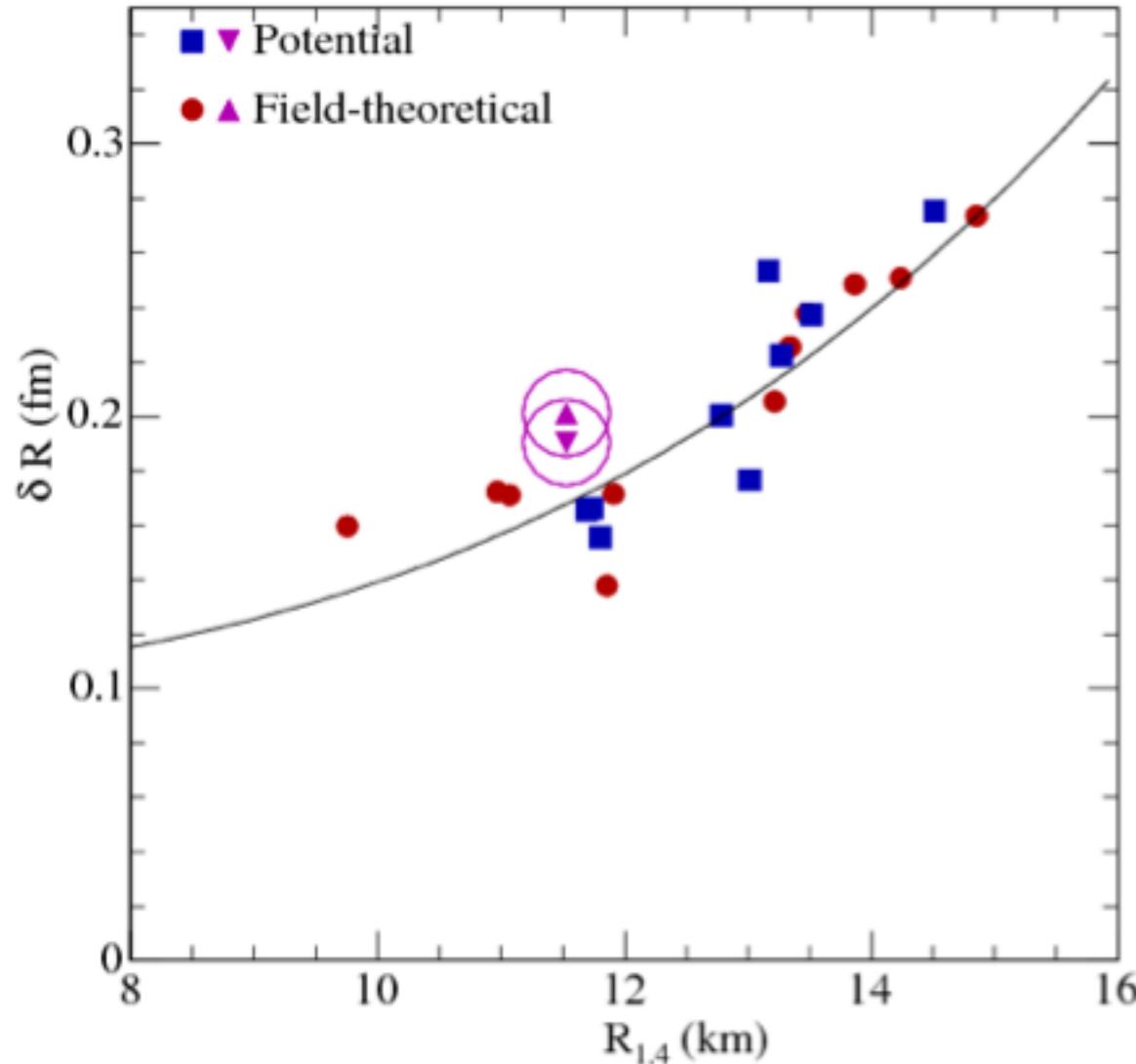
- Neutron radii are hard to measure, use parity-violating electron scattering
- Weak charge of neutron \gg weak charge of proton, i.e.
 $| -1 | \gg 1 - 4 \sin^2 \theta_W$



Jefferson Lab's Hall A
Measured $R_n - R_p = 0.33 \pm 0.16$ fm

The Neutron Skin Thickness of Lead

- The quantity $\delta R \equiv R_n - R_p$ is related to L as are neutron star radii



Jefferson Lab's Hall A: Measuring R_n

Steiner, Prakash, Lattimer, and Ellis (2005),
based on Horowitz and Piekarewicz (2001)

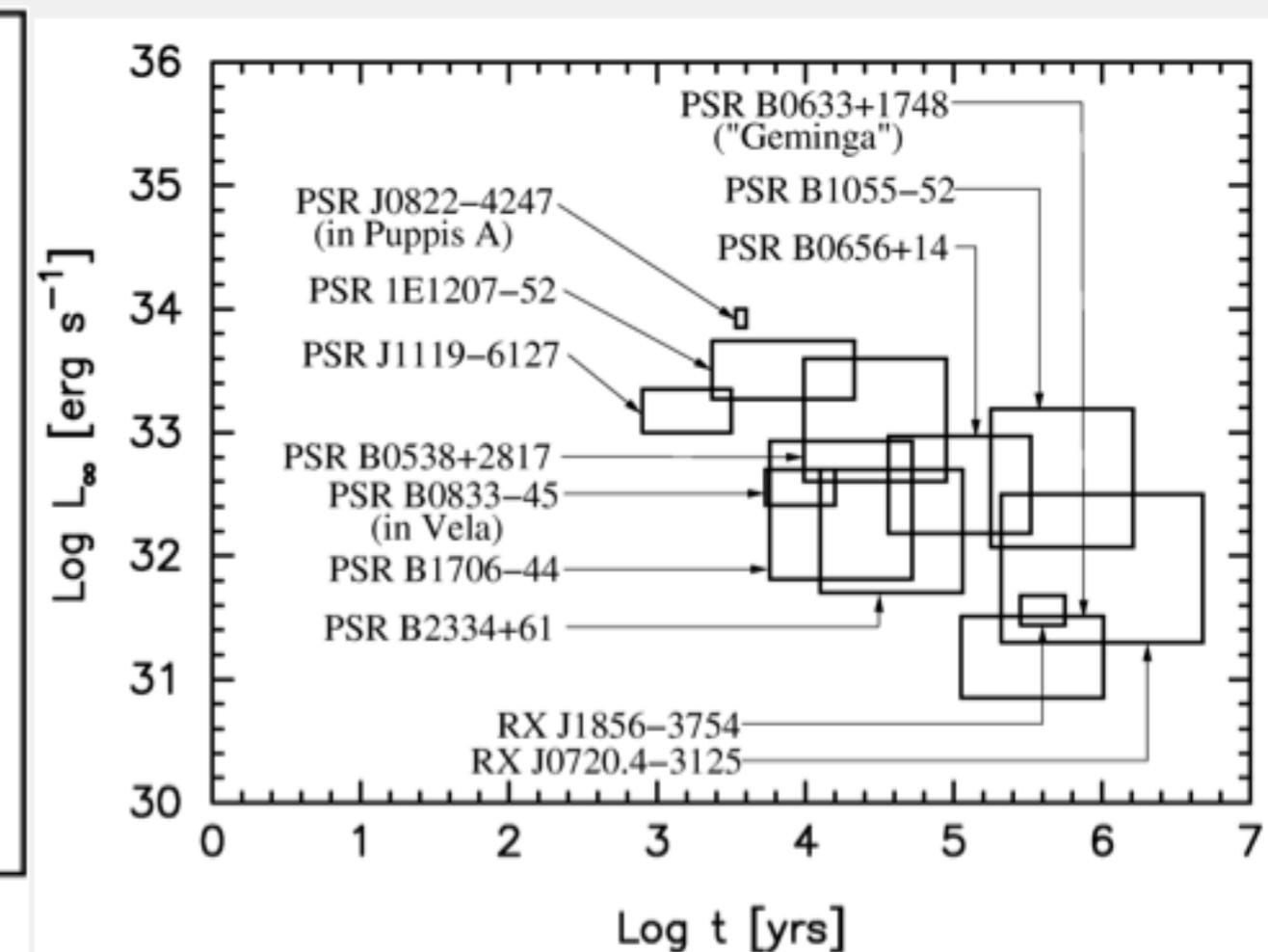
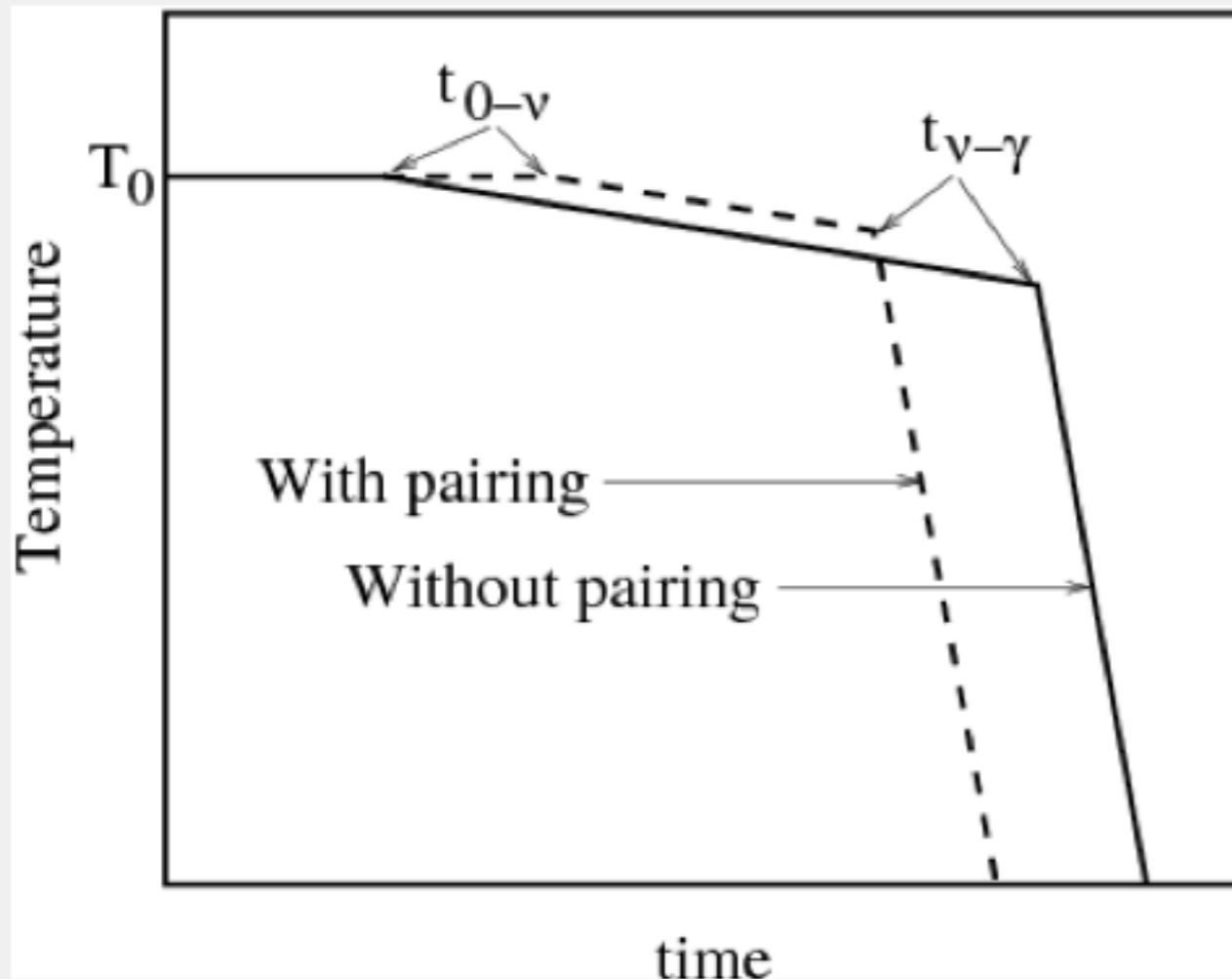
- We find $\delta R < 0.2$ fm (68%) from neutron star observations

Thermal Emission from Isolated Neutron Stars

- After ~ 10 years, the star is isothermal \Rightarrow one temperature = T

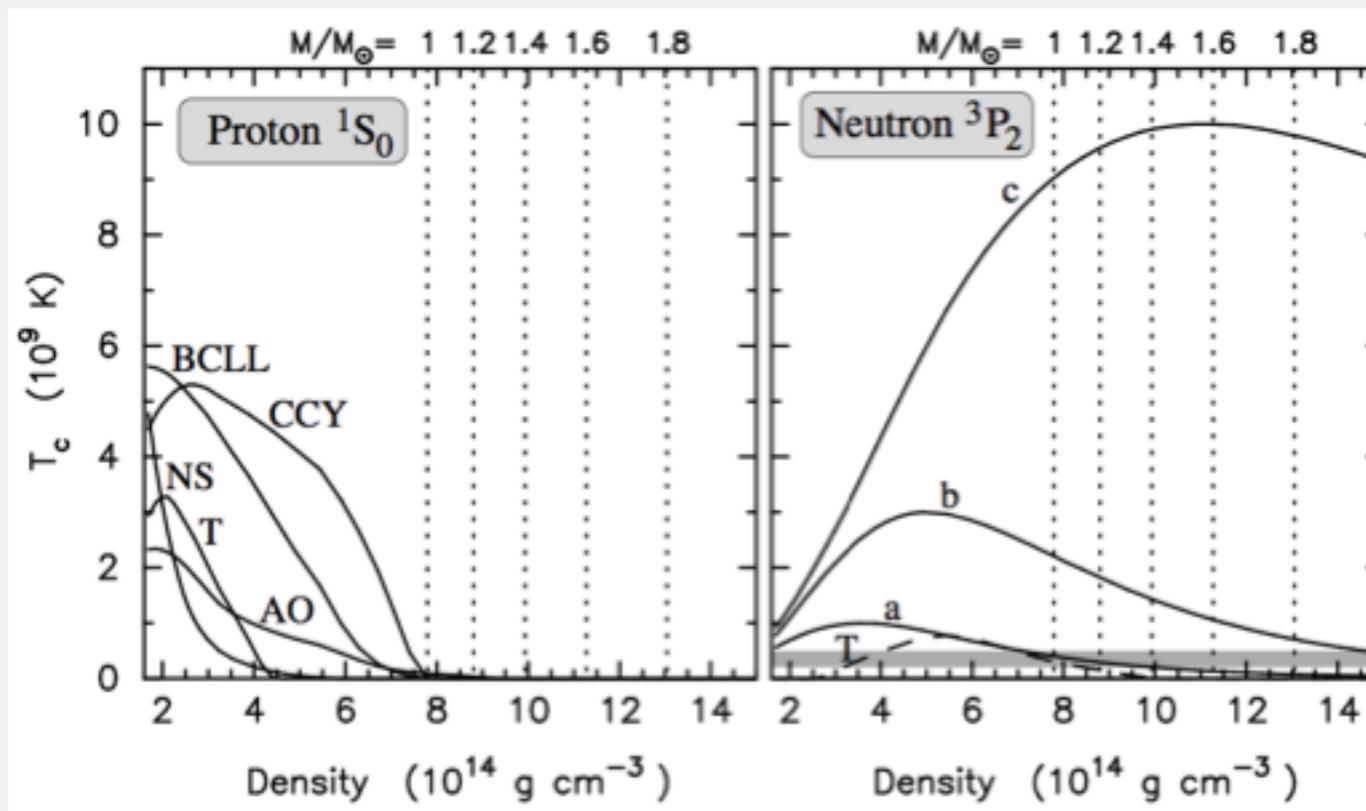
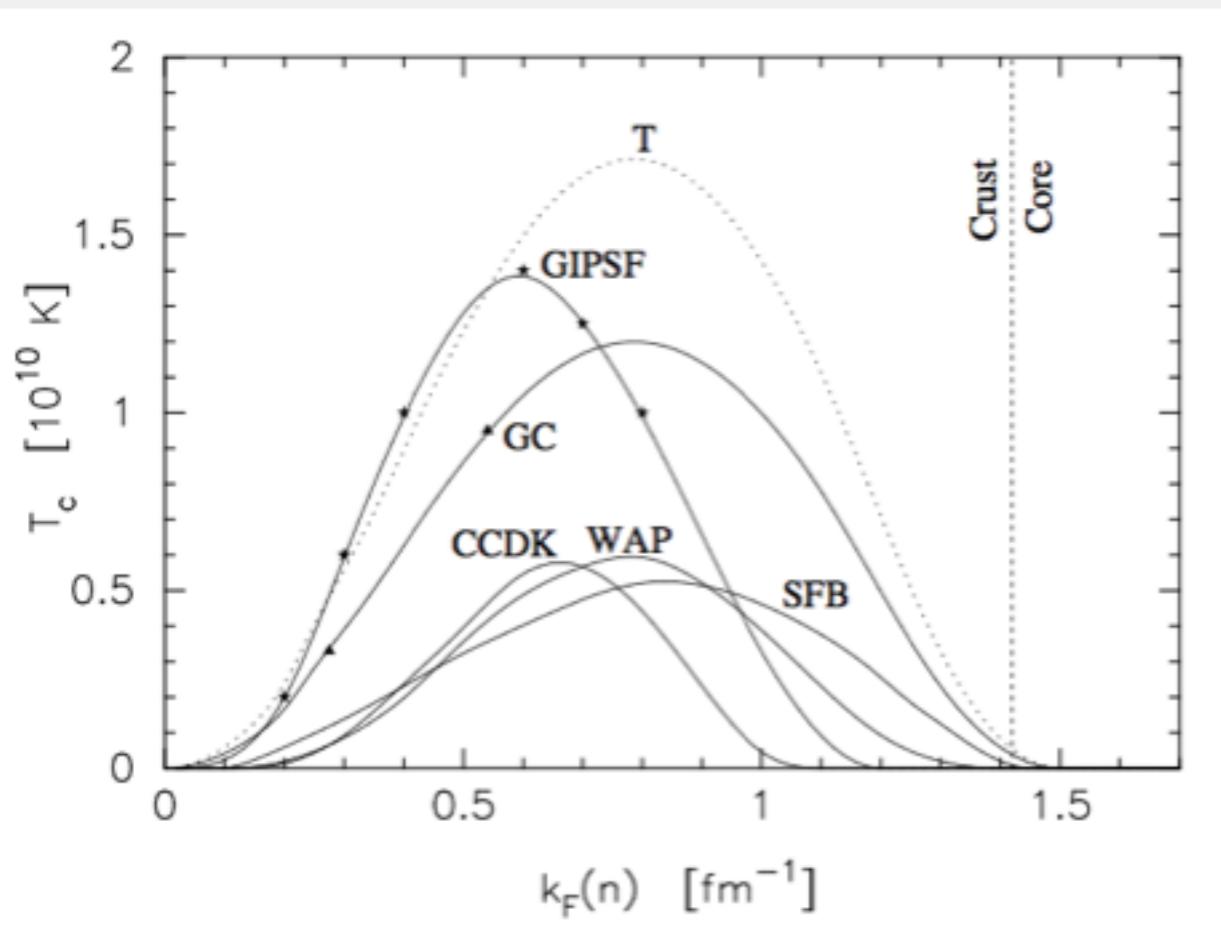
$$C_V \frac{dT}{dt} = L_\nu + L_\gamma, \quad L_\gamma \sim T^{2+4\alpha}, \quad L_\nu \sim T^8 \text{ (Modified Urca)}, \quad C_V \sim CT$$

- Age assumed from spin-down age or associated with a supernova remnant



Neutron Star Superfluidity

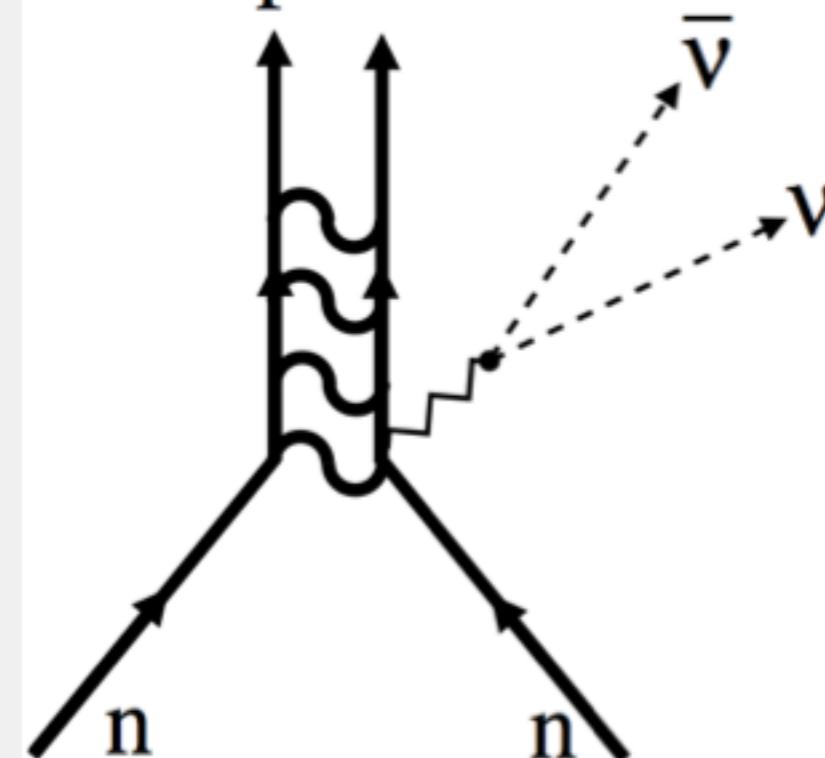
(See our review at 1302.6626)



- 1S_0 gap increases with increasing density, but drops off at higher densities because of n-n repulsion
- Superfluidity can block cooling processes
- ...but it opens up new ways of cooling

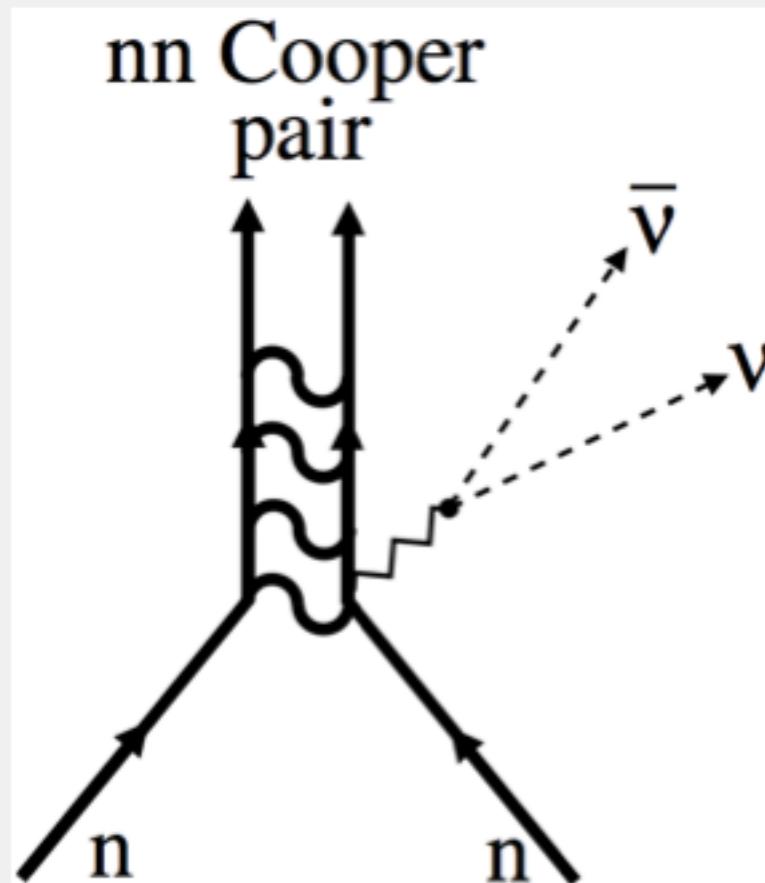
Steiner and Reddy (2009)

nn Cooper pair

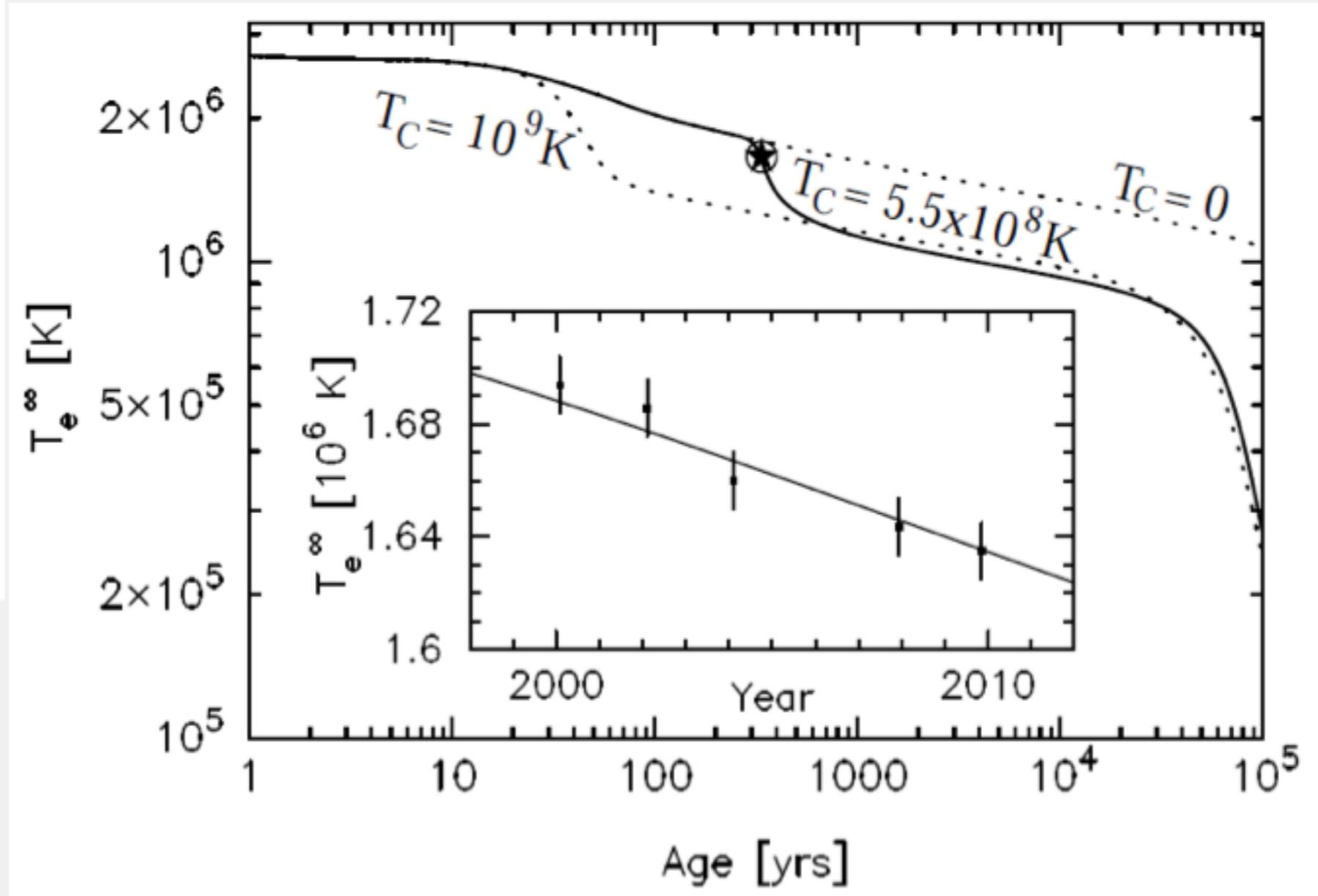


Detecting Neutron Star Superfluidity

- The large slope is only well reproduced by the neutron triplet superfluid transition and associated emissivity
- Cas A requires a very particular triplet gap $\Delta(T = 0) \propto T_C$



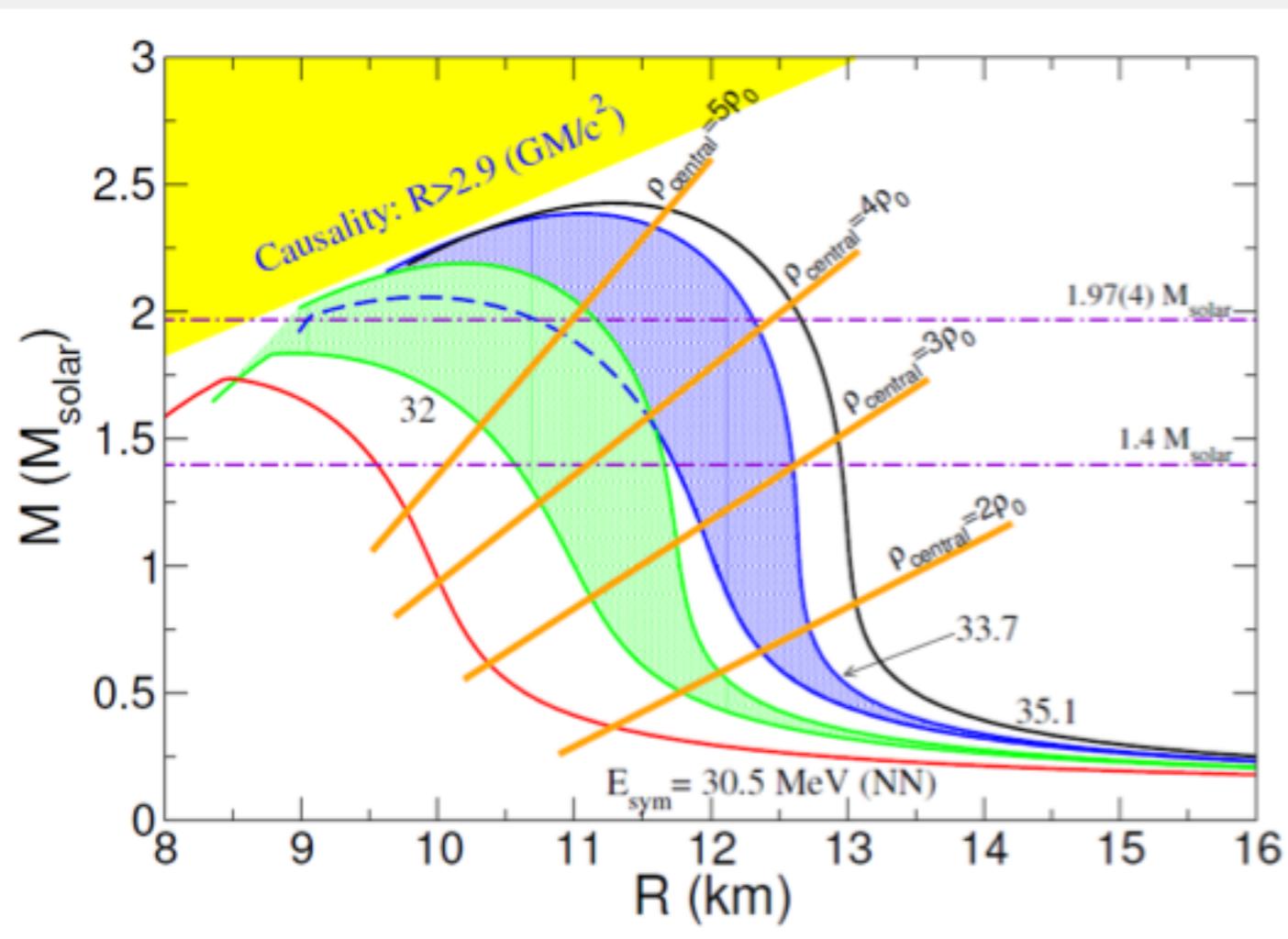
- If you form a Cooper pair, you gain energy



PRE X-ray bursts

- van Paradijs et al. pioneer the idea, it's rarely used until Özel writes several papers starting in 2007ish, getting small radii
- We demonstrate that photosphere radii are large at touchdown, add qLMXB data, use some nuclear physics, and get ~ 11 km radii.
- Suleimanov gets larger radii (14 km) for a long burst in XTE J1701, and claims other PRE X-ray data is poisoned by accretion ([Suleimanov et al. 2011](#))
- Yet the larger radius is somewhat inconsistent with qLMXB radii ([Steiner et al. 2012](#))
- Becomes clear that there may be (at least) two types of PRE X-ray bursts, which have different properties. Also some variation in normalization between bursts.
([Work by G. Zhang](#))
- Güver et al. do a systematic analysis of several sources and show that the fit of XTE J1701 is poor, but good for other sources ([Güver et al. 2012a and 2012b](#))
- Guillot et al. revisit qLMXB measurements, still find small radii ([Guillot et al. 2013](#))
- Lattimer and I re-revisit them, still finding smaller radii, but larger uncertainties
([Lattimer et al. 2013](#))
- Work with Suleimanov finds XTE J1701 is complicated by a boundary layer (possibly explaining the poor fit?) ([Retvinsev et al. 2013](#))
- Status: Larger (~ 14 km) radii are not preferred and result in poorer fits, unless you presume something has gone terribly wrong in qLMXBs. Nevertheless, PRE X-ray bursts are not well-understood and much work remains.

Connection to Nuclear Three-Body Forces



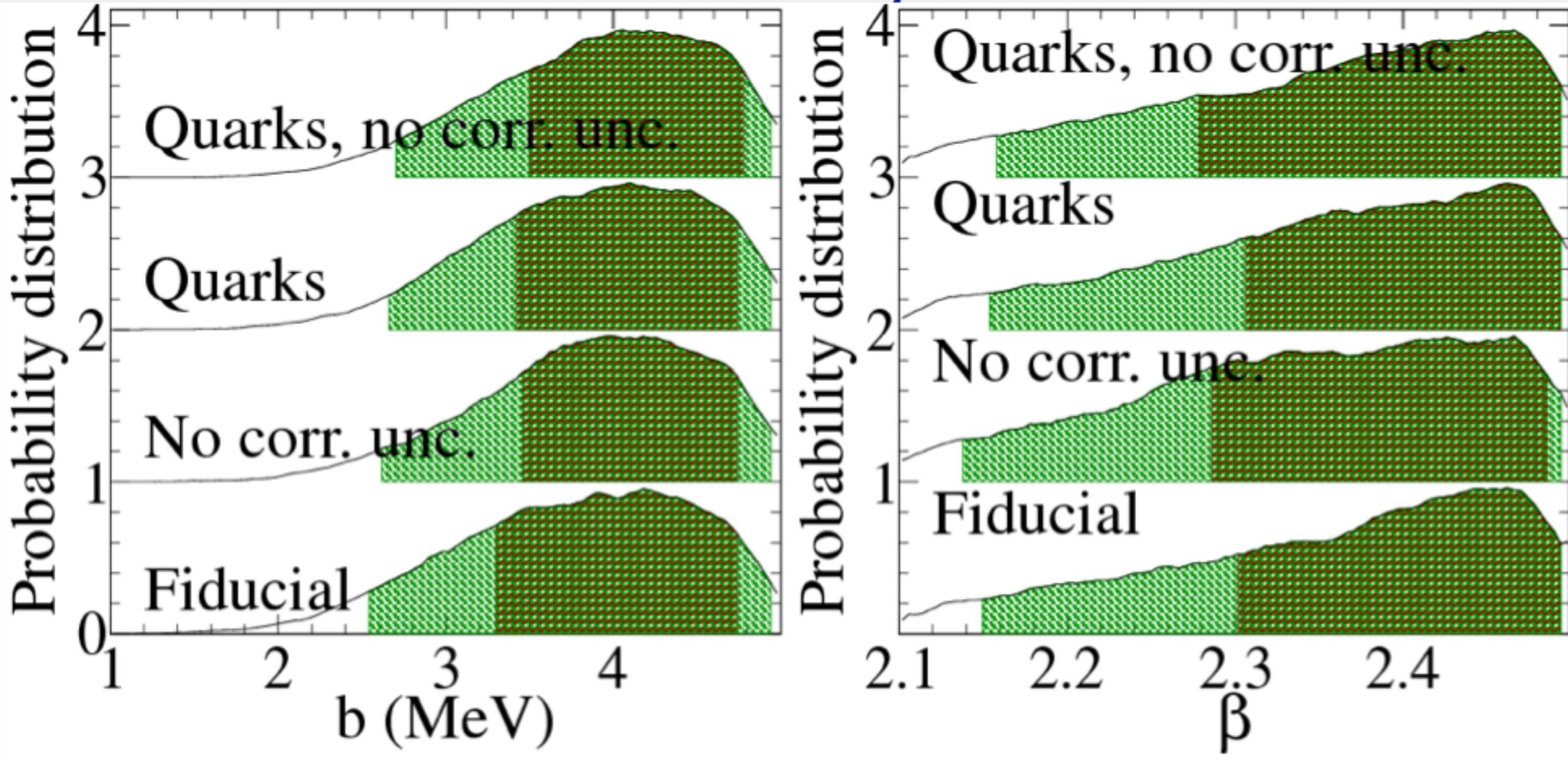
Colored regions denote
different three-body forces

$$E_{\text{neut}} = a \left(\frac{n}{n_0} \right)^\alpha + b \left(\frac{n}{n_0} \right)^\beta$$

Gandolfi, Carlson, and Reddy (2012)

- Three-nucleon interactions are important for nuclei and neutron star radii
- Quantum Monte Carlo (AFDMC) + Stellar structure
- How do neutron star observations constrain b and β ?

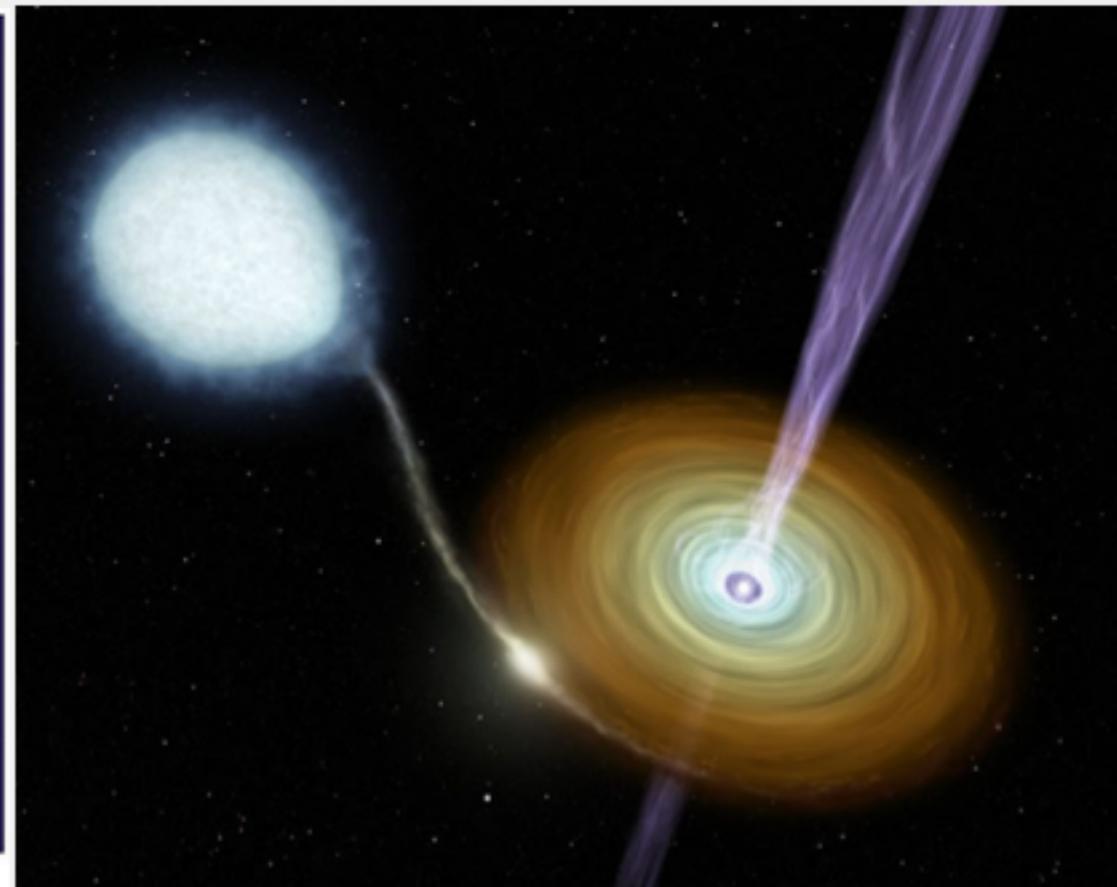
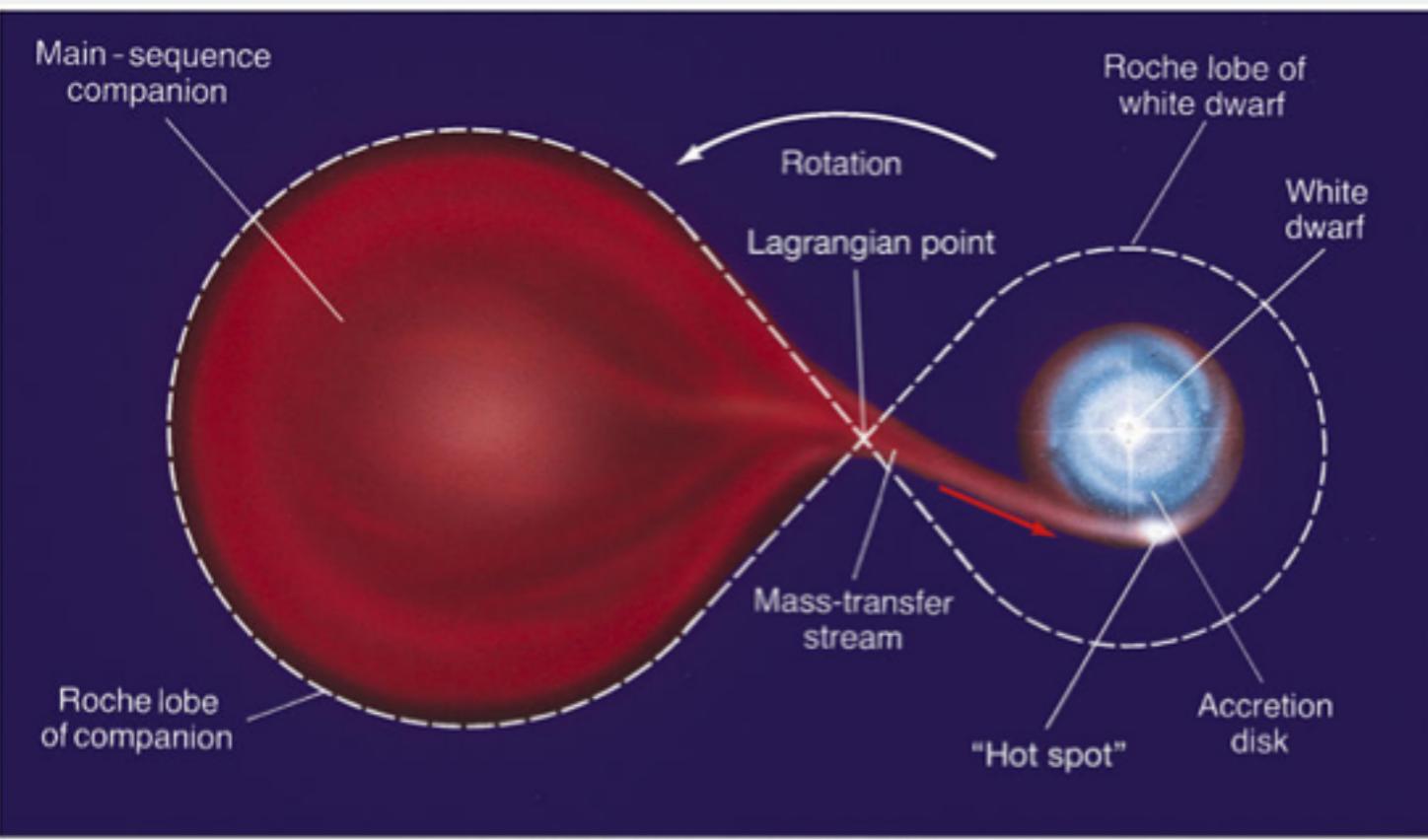
Constraints on Three-Body Force Parameters



Steiner and Gandolfi (2012)

- Values of a and α are unconstrained, but constraints on b and β
- Left and right plot boundaries exhaust expected range
- Neutron star radii are indirectly constraining nuclear three-body forces
- Limitation: if hyperons (or other strong phase transition) happens at a sufficiently low density, then results are modified

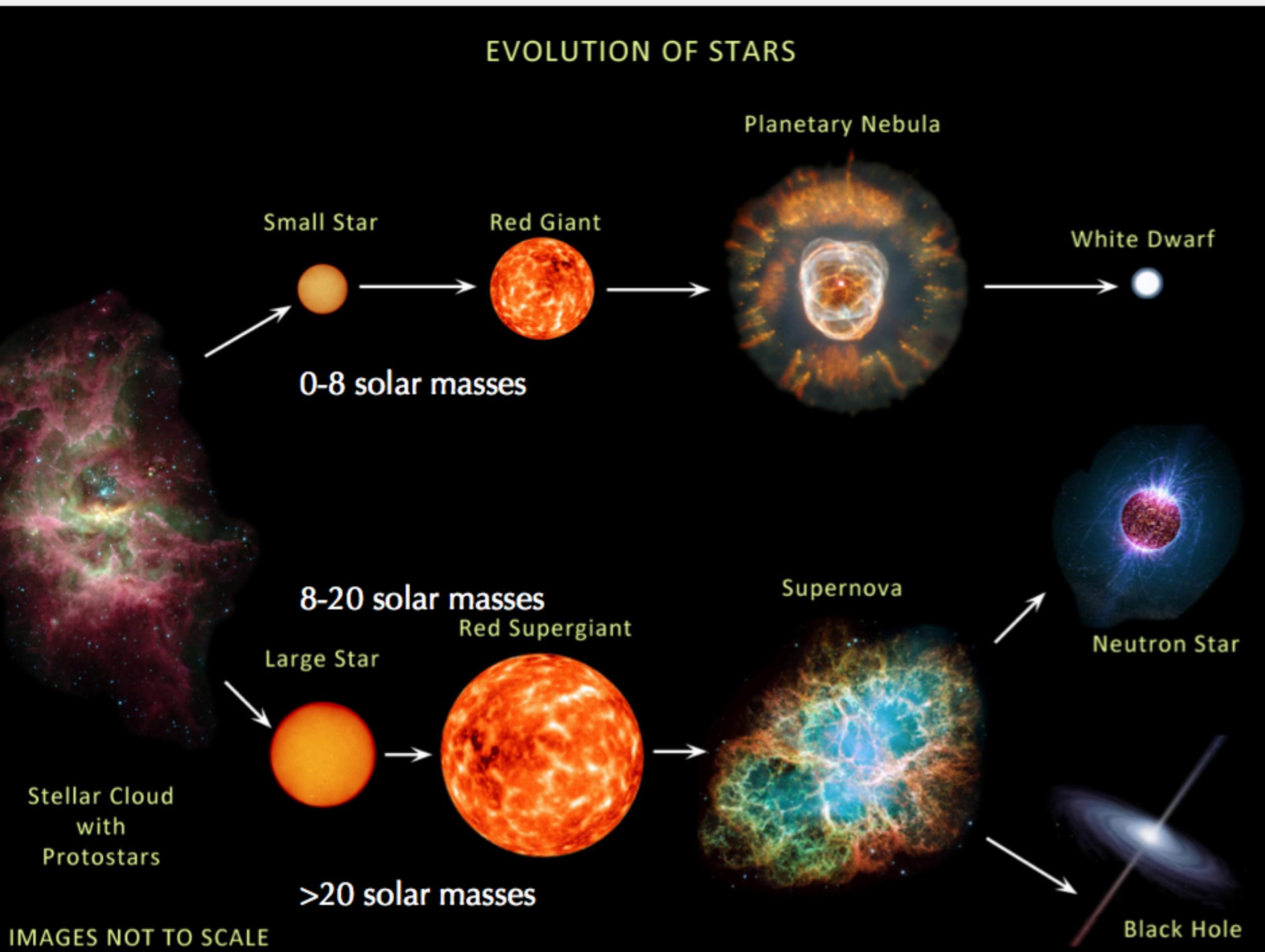
Accreting Neutron Stars: LMXBs



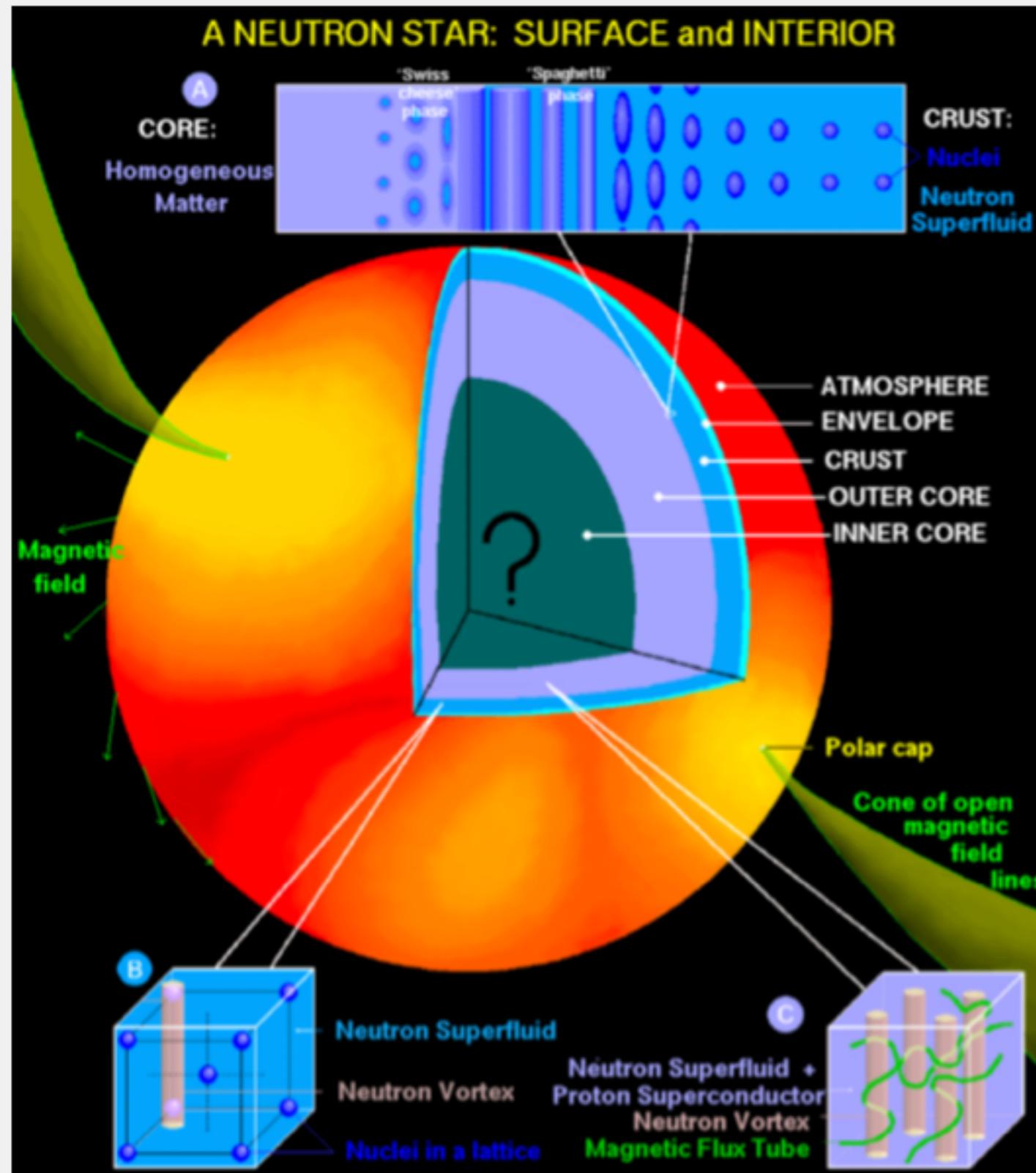
- Most stars have companions: neutron stars can have main-sequence, "normal star", companions
- Accretion heats the crust and is episodic
- At high enough density, H and He are unstable to thermonuclear explosions, i.e. X-ray bursts

Stellar Evolution

EVOLUTION OF STARS



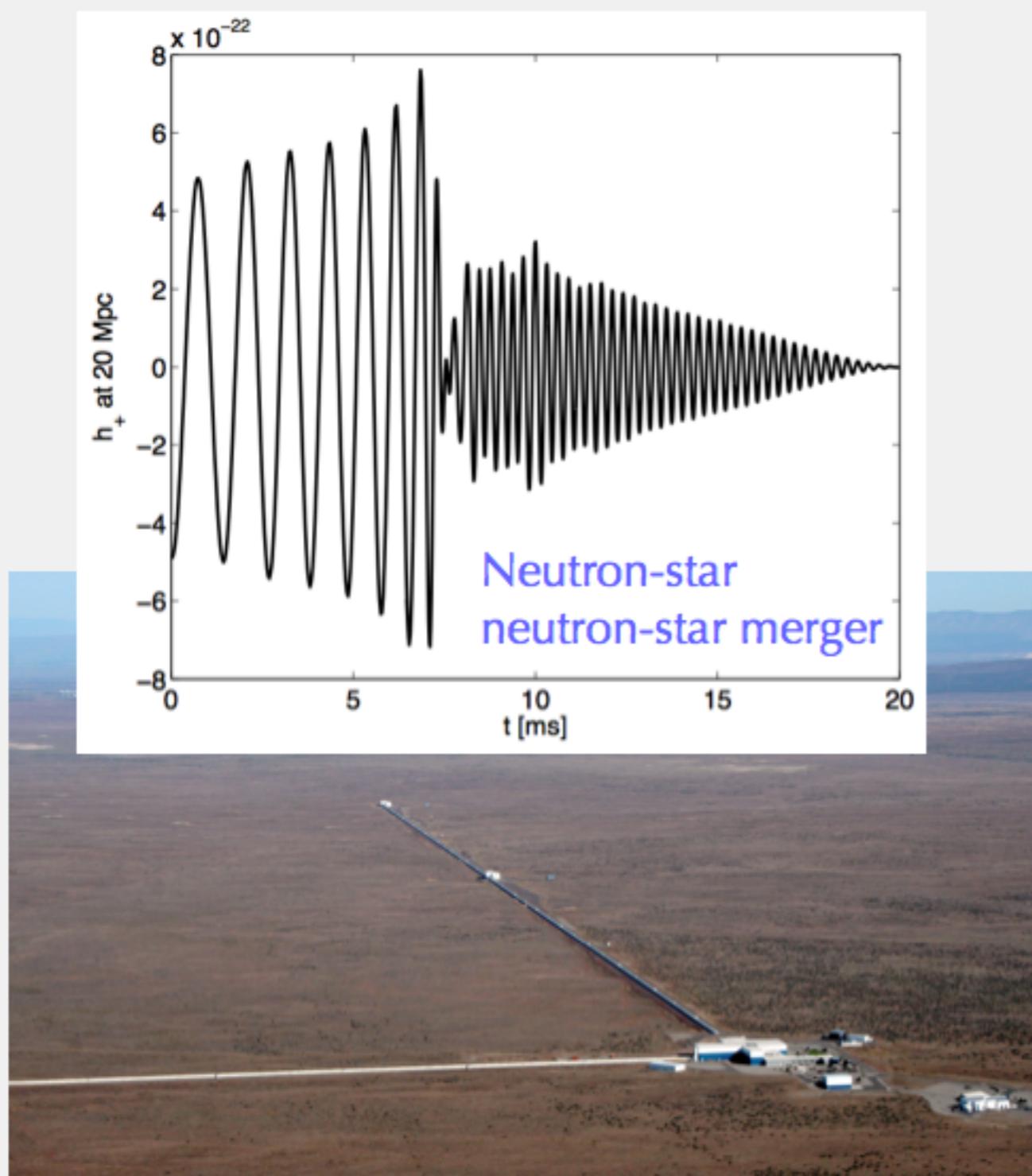
Neutron Star Composition



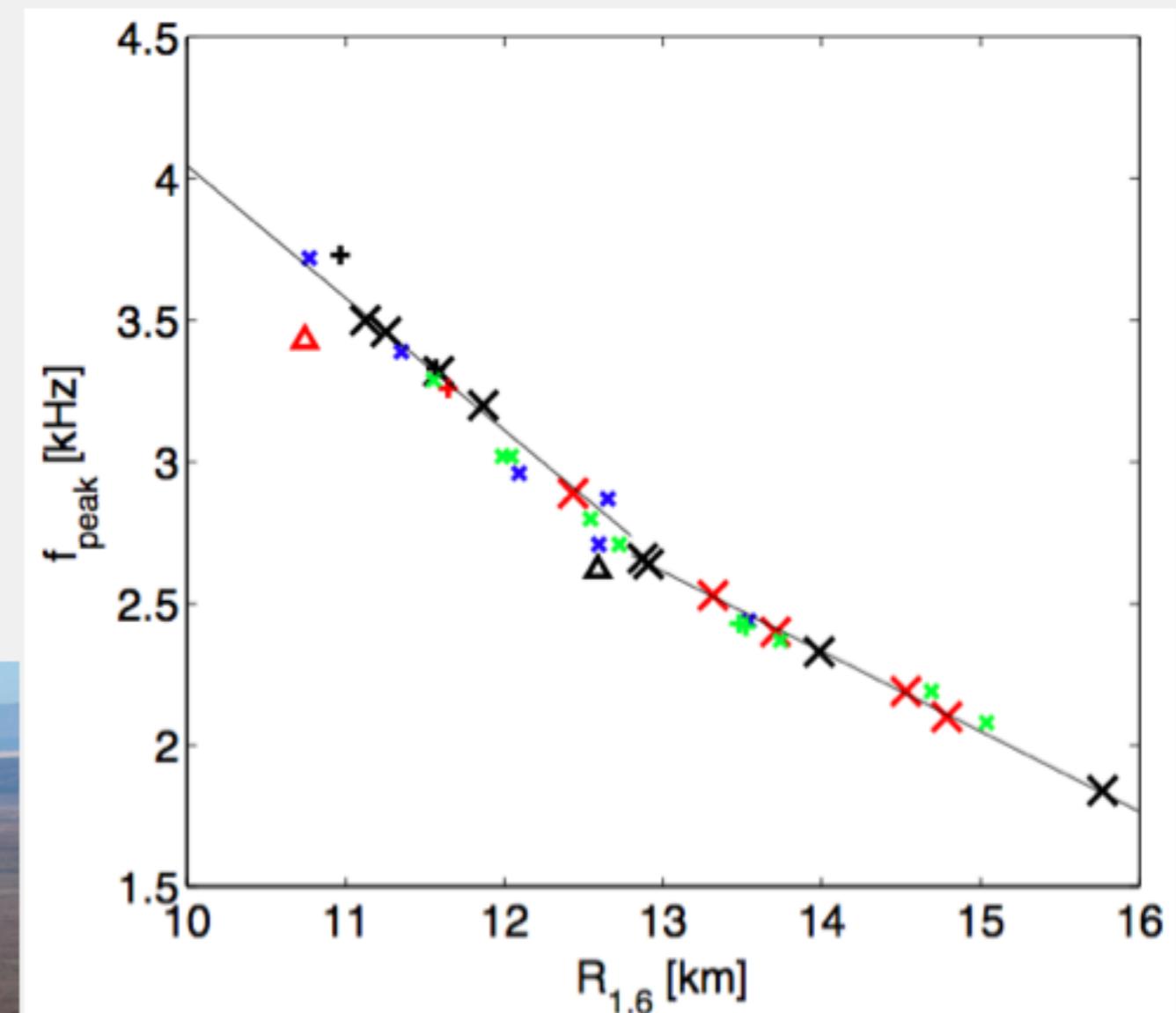
- Outer crust: of neutron-rich nuclei
- Inner crust: neutron-rich nuclei embedded in a sea of quasi-free superfluid neutrons
- Outer core: fluid of neutrons, protons, and electrons
- Inner core: hyperons, Bose condensates, deconfined quark matter
- **What are the correct degrees of freedom for the effective field theory which describes dense matter?**

Figure by Dany Page

Neutron Star Radii and LIGO Observations



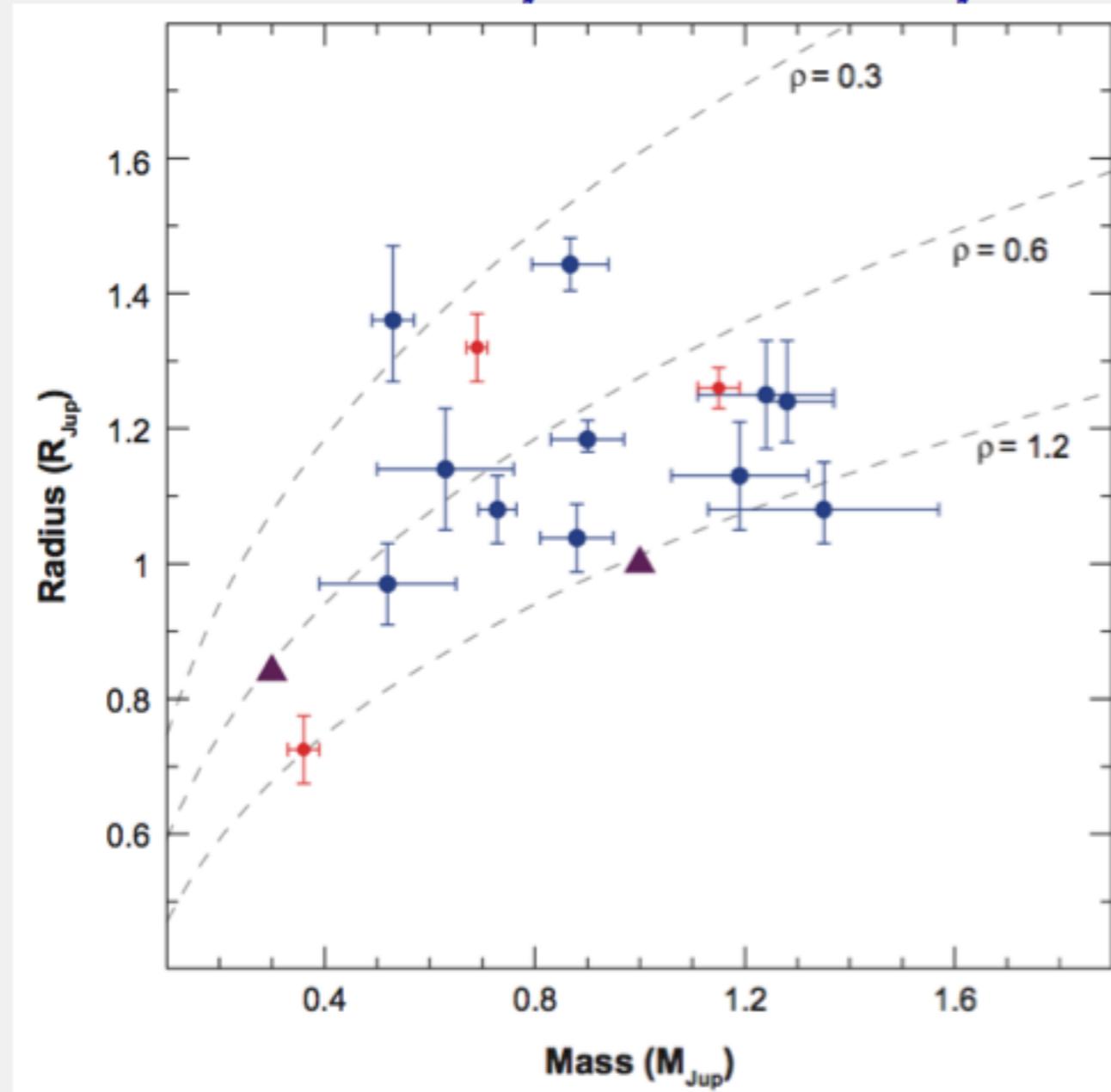
Laser Interferometer Gravitational
Wave Observatory Hanford



Bauswein, Janka, Hebeler, and Schwenk (2012)

- h_+ is the GW amplitude
- First detection in $\sim 2\text{-}3$ years
- Neutron star radii determine f_{peak}

Planetary Diversity



Udry et al. (2007)

- Varying composition, thus varying radius for a fixed mass