LIGO Continuous Gravitational Waves from Neutron Stars





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LIGO Scientific Collaboration and the Virgo Collaboration

(((O))) VIRGO

Binary Neutron Star Coalescence as a Fundamental Physics Laboratory, INT, UW, Seattle WA, July 3, 2014

Purpose

Neutron stars are also possible sources of detectable continuous gravitational waves.

The LIGO Scientific Collaboration and Virgo Collaboration have searched for these waves.

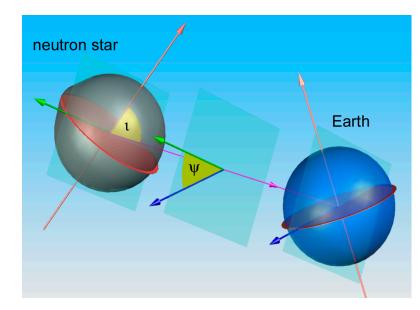
The purpose of this talk is give an overview of these searches, and discuss what can be learned from them compared to searches for gravitational-waves from binary neutron star coalescence. Recent results will be presented.

Periodic Continuous Gravitational waves

The GW signal from a triaxial pulsar can be modelled as

$$h(t) = \frac{1}{2}F_{+}(t; \psi)h_{0}(1 + \cos^{2}\iota)\cos 2\Psi(t) + F_{\times}(t; \psi)h_{0}\cos\iota\sin 2\Psi(t)$$

- •The unknown parameters are
 - h₀ amplitude of the gravitational wave signal
 - ψ polarization angle of signal; embedded in $F_{x,+}$
 - ι inclination angle of the pulsar
 - ϕ_0 initial phase of pulsar $\Phi(0)$
 - In the known pulsar searches we usually look for signals at twice the rotation frequency of the pulsars
 - For blind searches for isolated neutron stars the location in the sky and the source's frequency and its evolution are search parameters.



Generation of Continuous Gravitational Waves

Radiation generated by quadrupolar mass movements:

$$h_{\mu\nu} = \frac{2G}{rc^4} \frac{d^2}{dt^2} \left[I_{\mu\nu} \right]$$

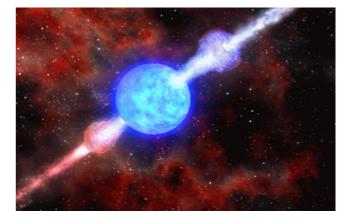


No GW from axisymmetric object rotating about symmetry axis

 $(I_{\mu\nu} = \text{quadrupole tensor}, r = \text{source distance})$

Spinning neutron star with equatorial ellipticity ε_{equat}

$$\varepsilon_{\text{equat}} = \frac{|I_{xx} - I_{yy}|}{I_{zz}}$$



gives a strain amplitude $h(f_{GW} = 2 \cdot f_{Bot})$:

Courtesy: U. Liverpool

$$h = 1.1 \times 10^{-24} \left[\frac{kpc}{r} \right] \left[\frac{f_{GW}}{kHz} \right]^2 \left[\frac{\varepsilon}{10^{-6}} \right] \left[\frac{I_{zz}}{10^{38} \text{ kg} \cdot \text{m}^2} \right]$$

Gravitational CW mechanisms

□ Equatorial ellipticity (e.g., – mm-high "bulge"):

$$h \propto \varepsilon_{\text{equat}}$$
 with $f_{GW} = 2f_{\text{rot}}$

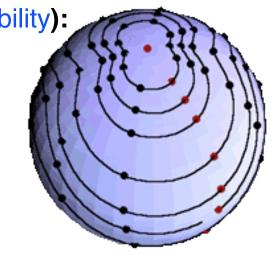
□ Poloidal ellipticity (natural) + wobble angle (precessing star):

$$h \propto \varepsilon_{poloidal} \times \theta_{wobble}$$
 with $f_{GW} = f_{rot} \pm f_{precess}$

(precession due to different L and Ω axes)

- □ Two-component (crust+superfluid) → $f_{\text{GW}} = f_{rot}$ and $2f_{\text{rot}}$
- □ r modes (rotational oscillations CFS-driven instability):
 - N. Andersson, ApJ 502 (1998) 708
 - S. Chandrasekhar PRL 24 (1970) 611
 - J. Friedman, B.F. Schutz, ApJ 221 (1978) 937

$$h \propto \alpha_{\text{r-mode}} \quad \text{with} \quad f_{GW} \cong \frac{4}{3} f_{\text{rot}}$$



Gravitational CW mechanisms

Assumption we (LSC, Virgo) have usually made to date:

Bulge is best bet for detection

→ Look for GW emission at twice the EM frequency

e.g., look for Crab Pulsar (29.7 Hz) at 59.5 Hz (troublesome frequency in North America!)

What is allowed for ε_{equat} ?

Old maximum (?) $\approx 5 \times 10^{-7}$ [$\sigma/10^{-2}$] ("ordinary" neutron star) with σ = breaking strain of crust

G. Ushomirsky, C. Cutler, L. Bildsten MNRAS 319 (2000) 902

More recent finding: $\sigma \approx 10^{-1}$ supported by detailed numerical simulation C.J. Horowitz & K. Kadau PRL 102, (2009) 191102

Recent re-evaluation: $\varepsilon_{\text{equat}} < 10^{-5}$ N.K. Johnson-McDaniel & B.J. Owen PRD 88 (2013) 044004

Gravitational CW mechanisms

Strange quark stars <u>could support</u> much higher ellipticities B.J. Owen PRL 95 (2005) 211101, Johnson-McDaniel & Owen (2013)

Maximum $\varepsilon_{\text{equat}} \approx 10^{-1}$ (!)

But what $\varepsilon_{\text{equat}}$ is <u>realistic</u>?

What could drive ε_{equat} to a high value (besides accretion)?

Millisecond pulsars have spindown-implied values lower than 10⁻⁹–10⁻⁶

New papers revisiting possible GW emission mechanisms (e.g., buried magnetic fields, accretion-driven r-modes) are also intriguing

What is the "indirect spindown limit"?

It is useful to define the "indirect spindown limit" for a known pulsar, under the assumption that it is a "gravitar", i.e., a star spinning down due to gravitational wave energy loss

Unrealistic for known stars, but serves as a useful benchmark

Equating "measured" rotational energy loss (from measured period increase and reasonable moment of inertia) to GW emission gives:

$$h_{SD} = 2.5 \times 10^{-25} \left[\frac{kpc}{d} \right] \sqrt{\left[\frac{1kHz}{f_{GW}} \right] \left[\frac{-df_{GW} / dt}{10^{-10} Hz / s} \right] \left[\frac{I}{10^{45} g \cdot cm^2} \right]}$$

Example:

Crab
$$\rightarrow$$
 h_{SD} = **1.4** × **10**⁻²⁴

 $(d=2 \text{ kpc}, f_{GW} = 59.5 \text{ Hz}, df_{GW}/dt = -7.4 \times 10^{-10} \text{ Hz/s})$



Finding a completely unknown CW Source

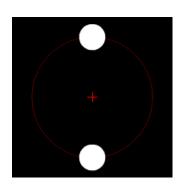
Serious technical difficulty: Doppler frequency shifts

- Frequency modulation from earth's rotation (v/c ~ 10-6)
- Frequency modulation from earth's orbital motion (v/c ~ 10⁻⁴)
- → Coherent integration of 1 year gives frequency resolution of 30 nHz
- → 1 kHz source spread over 6 million bins in ordinary FFT!

Additional, related complications:

Daily amplitude modulation of antenna pattern

Spin-down of source



Orbital motion of sources in binary systems

Finding a completely unknown CW Source

Modulations / drifts complicate analysis enormously:

- Simple Fourier transform inadequate
- Every sky direction requires different demodulation

Computational scaling:

```
Single coherence time – Sensitivity improves as (T_{coherence})^{1/2} but cost scales with (T_{coherence})^{6+}
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- → Restricts T_{coherence} < 1-2 days for all-sky search
- → Exploit <u>coincidence</u> among different spans

Alternative:

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Semi-coherent stacking of spectra (e.g., T_{coherence} = 30 \text{ min}) \rightarrow Sensitivity improves only as (N_{stack})^{1/4}
```

→ All-sky survey at full sensitivity = Formidable challenge Impossible?

Frequency Modulation

$$f(t) \cong \left(1 + \frac{\vec{v}(t)}{c} \cdot \hat{n}\right) \left[f_0 + f_1(t - t_0) + \dots\right]$$

Relativistic corrections can be included in the actual code.

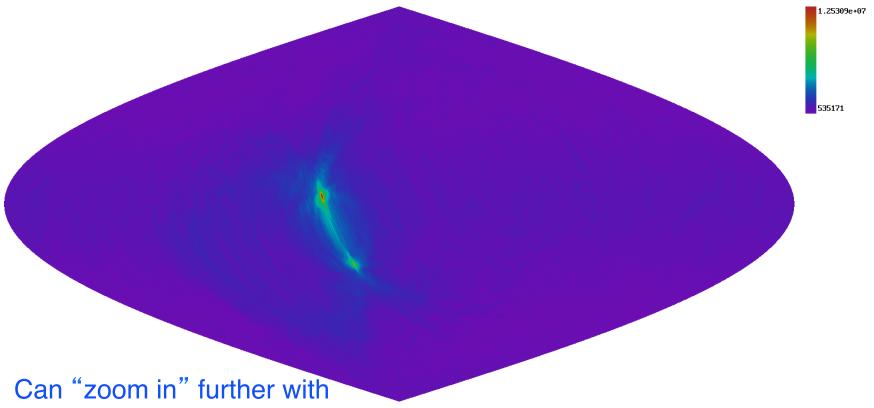
$$\dot{f}(t) \cong \left(\frac{\vec{a}(t)}{c} \cdot \hat{n}\right) \left[f_0 + f_1(t - t_0)\right] + \left(1 + \frac{\vec{v}(t)}{c} \cdot \hat{n}\right) f_1 + \dots$$

$$S = \left(\frac{\vec{a}_{orb}(t)}{c} \cdot \hat{n}\right) f_0 + f_1$$

 $S = \left(\frac{\vec{a}_{orb}(t)}{c} \cdot \hat{n}\right) f_0 + f_1$ For analysis < 1 yr sky points with small S have small doppler variation making them harder to distinguish GWs from instrument lines at these points.

But three substantial benefits from modulations:

- Reality of signal confirmed by need for corrections
- Corrections give precise direction of source
- Single interferometer can make definitive discovery



follow-up algorithms once we lock on to source

[V. Dergachev, PRD 85 (2012) 062003M. Shaltev & R. Prix, PRD 87 (2013) 084057

Sky map of strain power for signal injection (semi-coherent search)

Targeted (matched-filter) algorithm applied to <u>195</u> known pulsars over LIGO S5/S6 and Virgo VSR2/VSR4 data

Lowest (best) upper limit on strain:

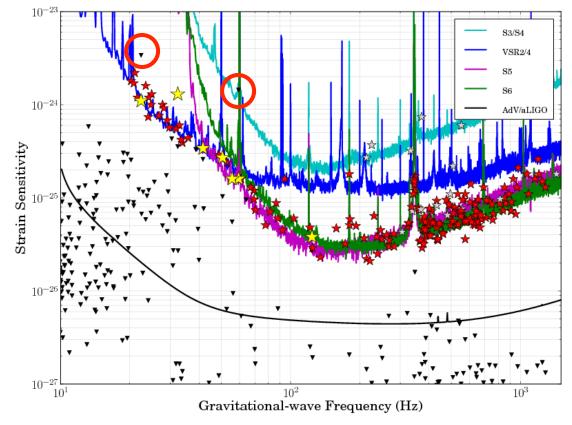
$$h_0 < 2.1 \times 10^{-26}$$

Lowest (best) upper limit on ellipticity:

$$\varepsilon < 6.7 \times 10^{-8}$$

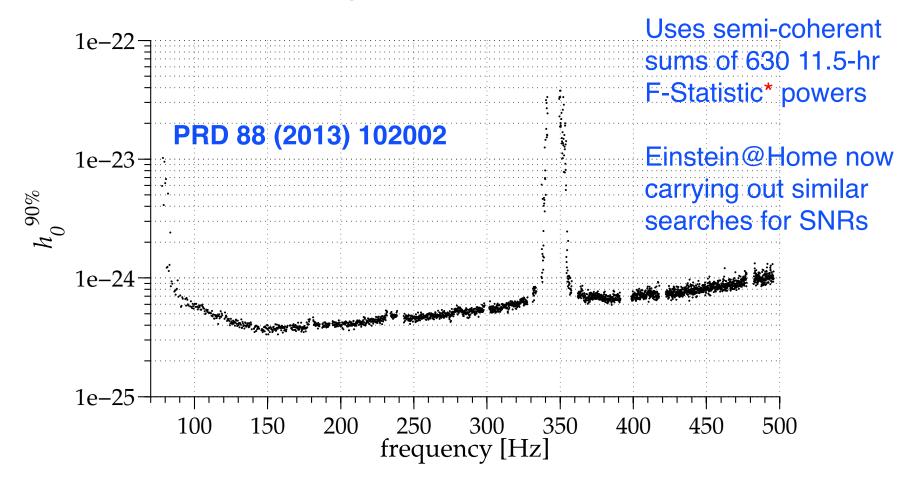
Crab limit at 1% of total energy loss

Vela limit at 10% of total energy loss

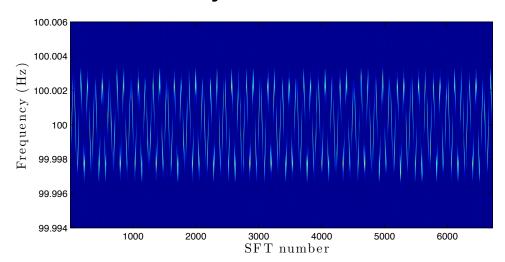


J. Aasi et al. 2014 ApJ 785 119.

Directed-search algorithm applied to the galactic center using LIGO S5 data (knowing direction improves sensitivity)



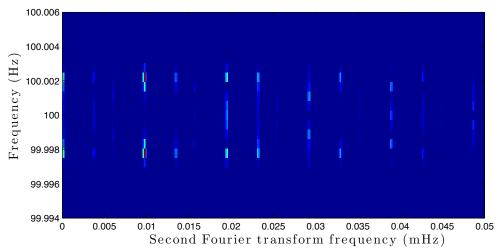
First all-sky search for unknown binary CW sources



Uses TwoSpect* algorithm:

Sample spectrogram (30-minute FFTs) for

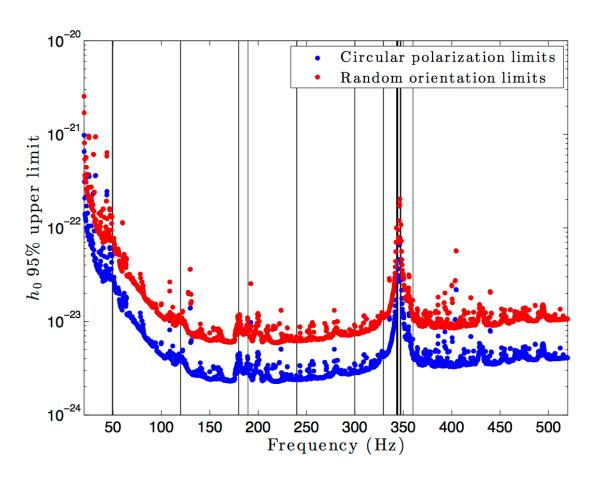
simulated strong signal (Earth's motion already demodulated)



Result of Fourier transforming each row of spectrogram

→ Concentrates power in orbital harmonics

First all-sky search for unknown binary CW sources



TwoSpect results for unknown spinning neutron stars in binary systems:

The blue dots in this plot show the upper limits on the circularly polarized gravitational wave strain amplitude.

The red dots show the upper limit on the randomly polarized gravitational wave strain amplitude.

arXiv:1405.7904

Summary

No discoveries yet, but...

Still examining data we have taken
 (computationally bound – E@H: 1 Petaflop, 100K volunteers)

- Major upgrade of LIGO & Virgo under way now
 - Advanced LIGO & Virgo
 - Improves range more than an order of magnitude
 - Moore's Law will help too...

Electromagnetic observations (radio, x-ray, γ -ray) of nearby neutron stars helpful now – and later

Extra Slides

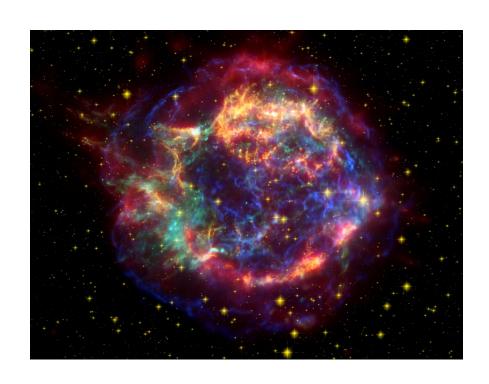
Not all known sources have measured timing

Compact central object in the Cassiopeia A supernova remnant

Birth observed in 1681 – One of the youngest neutron stars known

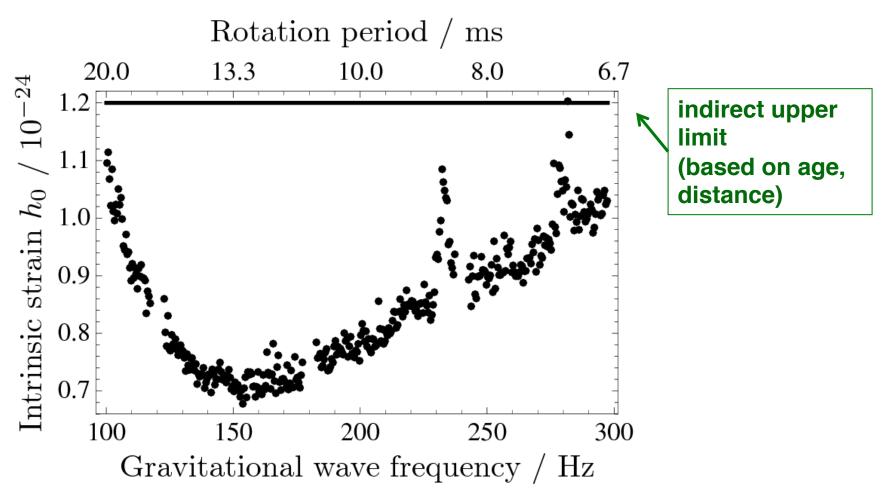
Star is observed in X-rays, but no pulsations observed

Requires a broad band search over accessible band



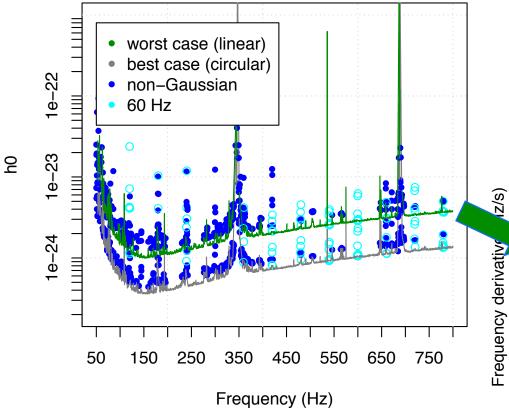
Cassiopeia A

Search for Cassiopeia A – Young age (~300 years) requires search over 2nd derivative



Ap. J. 722 (2010) 1504

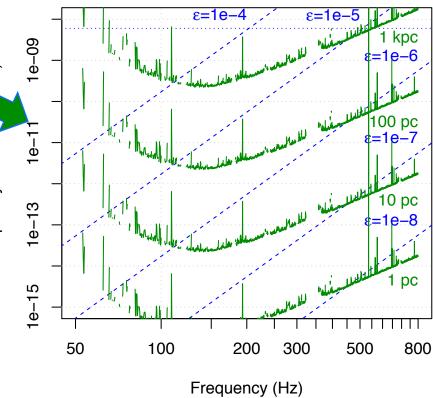




PRD 85 (2012) 022001

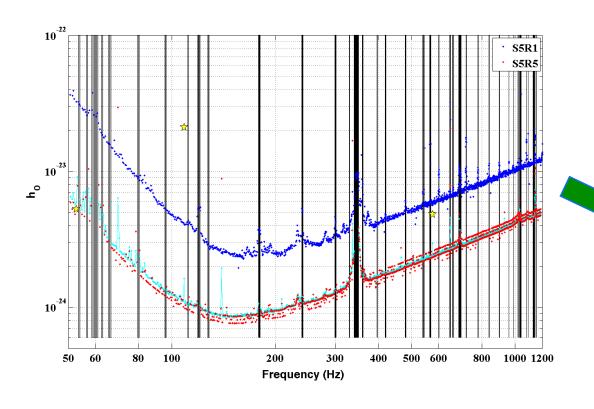
Semi-coherent stacks of 30-minute, demodulated power spectra ("PowerFlux")

Astrophysical reach:



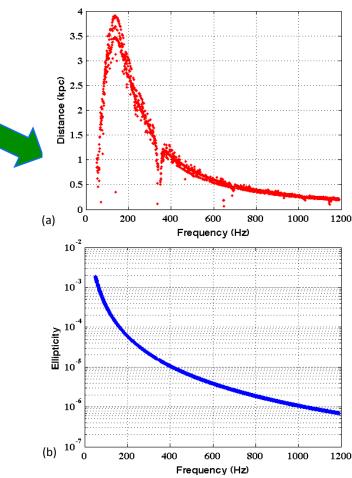
S5 all-sky results:

Einstein@Home semi-coherent sums of 121 25-hour F-Statistic powers (2 interferometers)



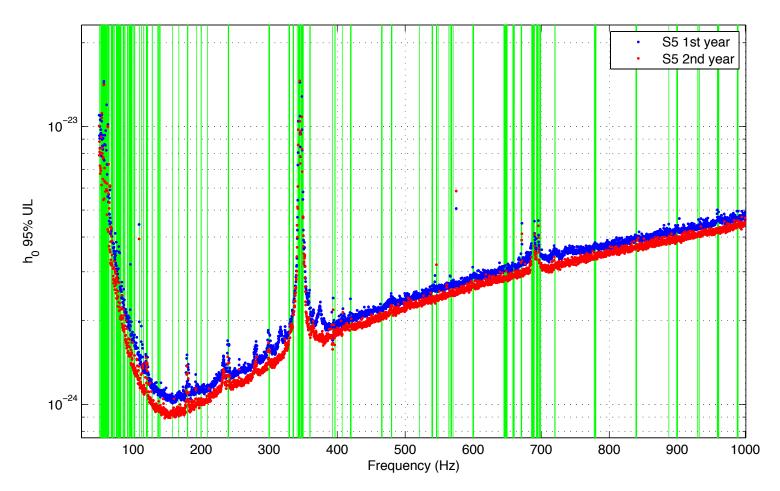
PRD 87 (2013) 042001





S5 all-sky results:

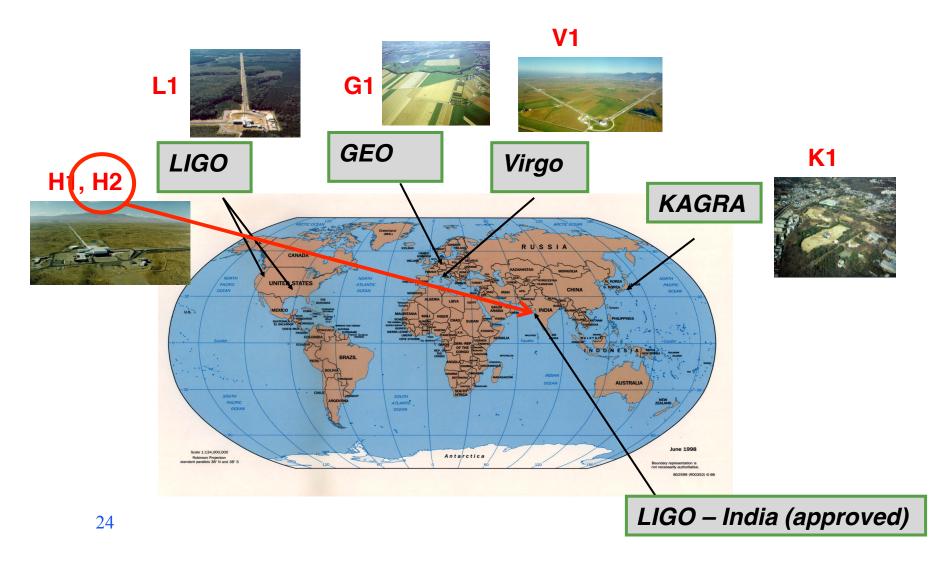
Hough-transform search based on ~68K 30-minute demodulated spectra (3 interferometers)



The Global Interferometer Network

The three (two) LIGO, Virgo and GEO interferometers are part of a Global Network.

Multiple signal detections will increase detection confidence and provide better precision on source locations and wave polarizations



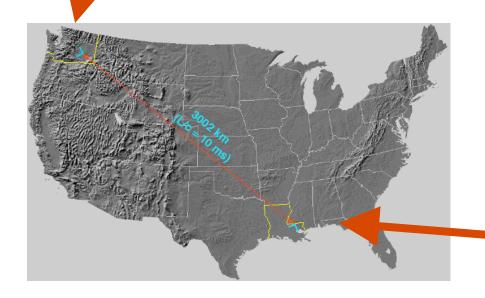
LIGO Observatories

Hanford



Observation of nearly simultaneous signals 3000 km apart rules out terrestrial artifacts

Livingston





Virgo

Have begun collaborating with Virgo colleagues (Italy/France)

Took data in coincidence for last ~4 months of latest science run

Data exchange and joint analysis underway

Will coordinate closely on detector upgrades and future data taking



3-km Michelson Interferometer just outside Pisa, Italy



GEO600

Work closely with the GEO600 Experiment (Germany / UK / Spain)

- Arrange coincidence data runs when commissioning schedules permit
- GEO members are full members of the LIGO Scientific Collaboration
- Data exchange and strong collaboration in analysis now routine
- Major partners in proposed Advanced LIGO upgrade



600-meter Michelson Interferometer just outside Hannover, Germany

LIGO Detector Facilities



Vacuum System

- •Stainless-steel tubes (1.24 m diameter, ~10⁻⁸ torr)
- Gate valves for optics isolation
- Protected by concrete enclosure



LIGO Detector Facilities

LASER

- □ Infrared (1064 nm, 10-W) Nd-YAG laser from Lightwave (now commercial product!)
- Elaborate intensity & frequency stabilization system, including feedback from main interferometer

Optics

- □ Fused silica (high-Q, low-absorption, 1 nm surface rms, 25-cm diameter)
- Suspended by single steel wire
- □ Actuation of alignment / position via magnets & coils



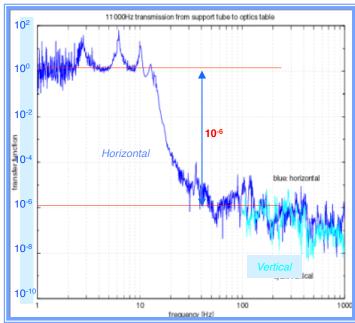


LIGO Detector Facilities

Seismic Isolation

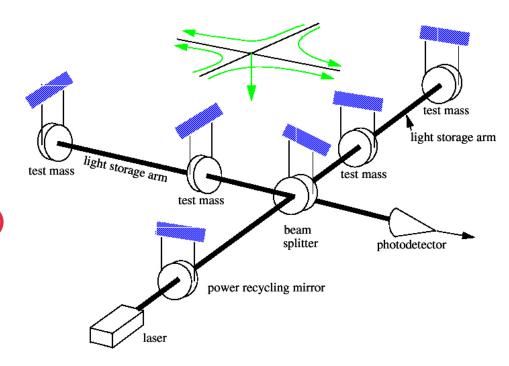
- □ Multi-stage (mass & springs) optical table support gives 10⁶ suppression
- □ Pendulum suspension gives additional 1 / f ² suppression above ~1 Hz



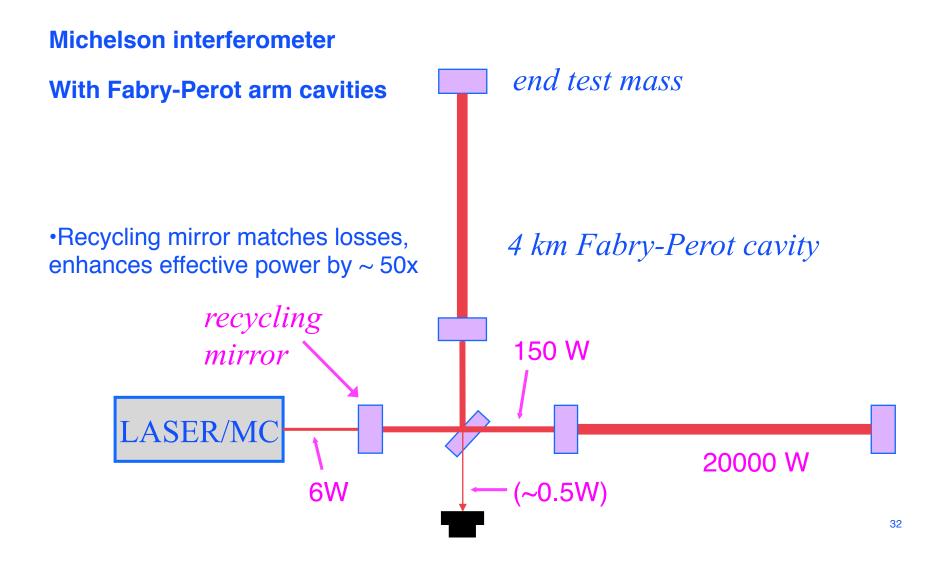


Gravitational Wave Detection

- □ Suspended Interferometers (IFO's)
 - Suspended mirrors in "free-fall"
 - Michelson IFO is "natural" GW detector
 - Broad-band response (~20 Hz to few kHz)
 - → Waveform information (e.g., chirp reconstruction)



LIGO Interferometer Optical Scheme

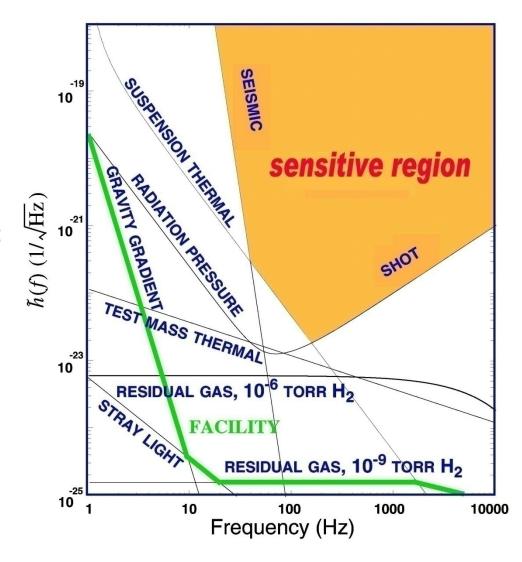


What Limits the Sensitivity of the Interferometers?

- Seismic noise & vibration limit at low frequencies
- Atomic vibrations (Thermal Noise) inside components limit at mid frequencies
- Quantum nature of light (Shot Noise) limits at high frequencies
- Myriad details of the lasers, electronics, etc., can make problems above these levels

Best design sensitivity:

~ 3 x 10⁻²³ Hz^{-1/2} @ 150 Hz



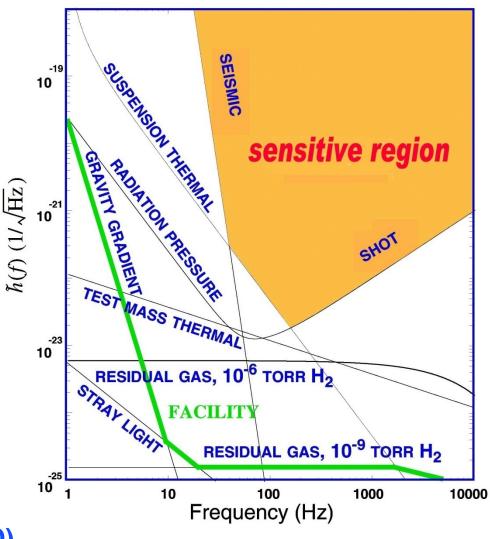
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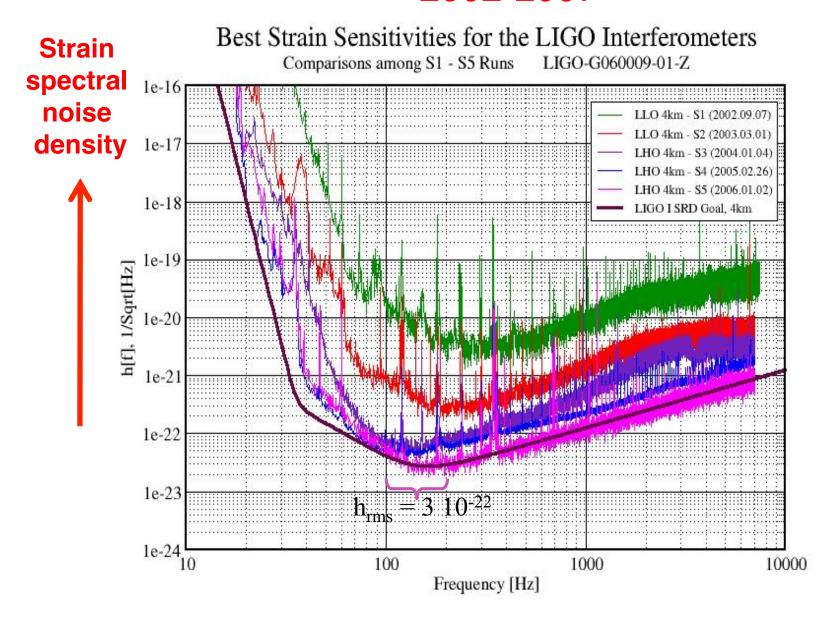
achieved Best design sensitivity:

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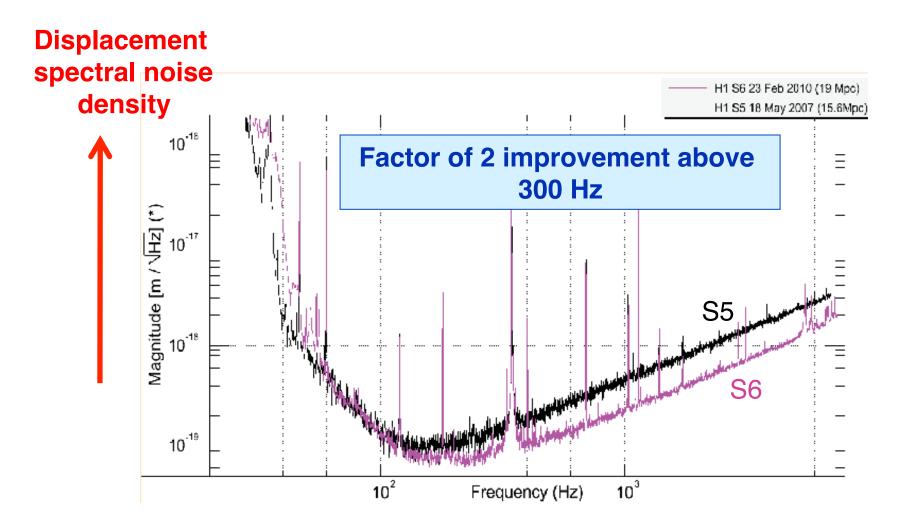
< 2 x 10⁻²³ (enhanced LIGO)



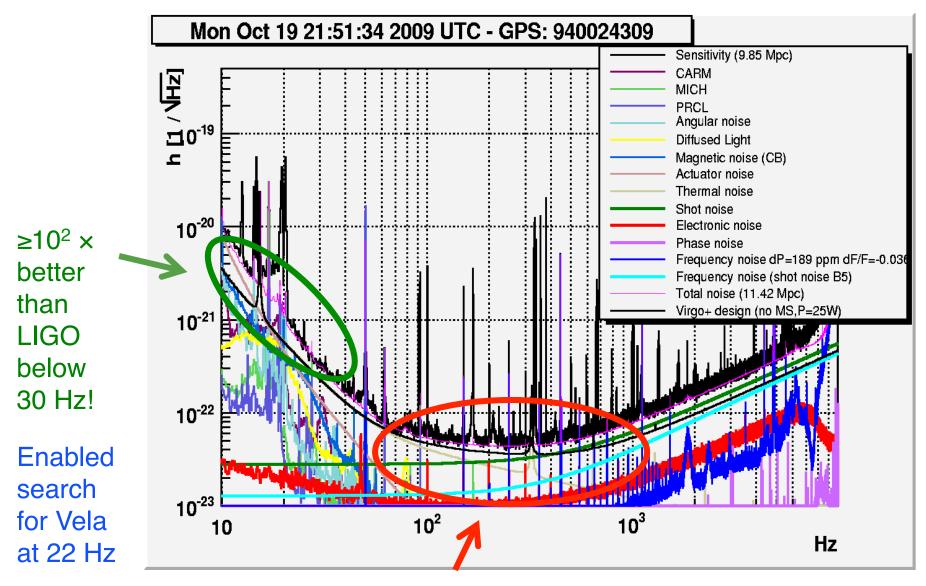
LIGO S1 → S5 Sensitivities ("Initial LIGO") 2002-2007



"Enhanced LIGO" (July 2009 – Oct 2010)



Virgo sensitivity in VSR2 (part of LIGO S6)



"Locking" the Inteferometer

Sensing gravitational waves requires sustained resonance in the Fabry-Perot arms and in the recycling cavity

- → Need to maintain half-integer # of laser wavelengths between mirrors
- → Feedback control servo uses error signals from imposed RF sidebands
- → Four primary coupled degrees of freedom to control
- → Highly non-linear system with 5-6 orders of magnitude in light intensity

Also need to control mirror rotation ("pitch" & "yaw")

→ Ten more DOF's (but less coupled)

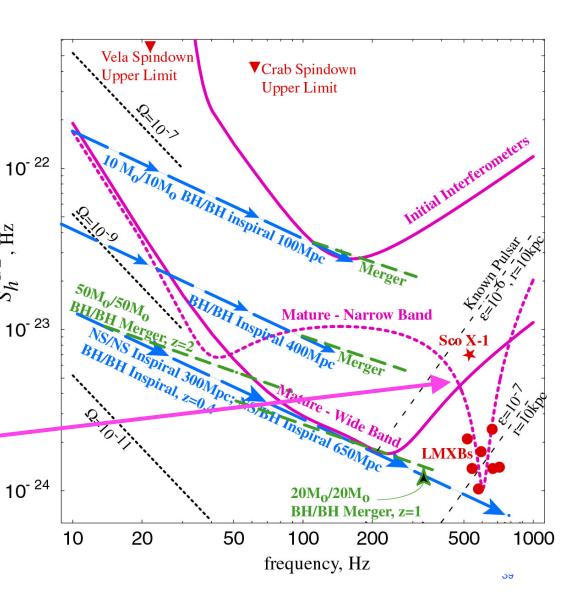
And need to stabilize laser (intensity & frequency), keep the beam pointed, damp out seismic noise, correct for tides, etc.,...

Advanced LIGO

Sampling of source strengths vis a vis Initial LIGO and Advanced LIGO

Lower h_{rms} and wider bandwidth both important

"Signal recycling" offers potential for tuning shape of noise curve to improve sensitivity in target band (e.g., known pulsar cluster) 10⁻²⁴



Advanced LIGO

Increased laser power:

10 W → 180 W

Improved shot noise (high freq)



Higher-Q test mass:

Fused silica with better optical coatings

Lower internal thermal noise in bandwidth

Increased test mass:

 $10 \text{ kg} \rightarrow 40 \text{ kg}$

Compensates increased radiation pressure noise

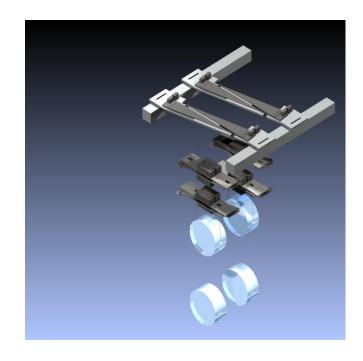
Advanced LIGO

Detector Improvements:

New suspensions:

Single → Quadruple pendulum

Lower suspensions thermal noise in bandwidth

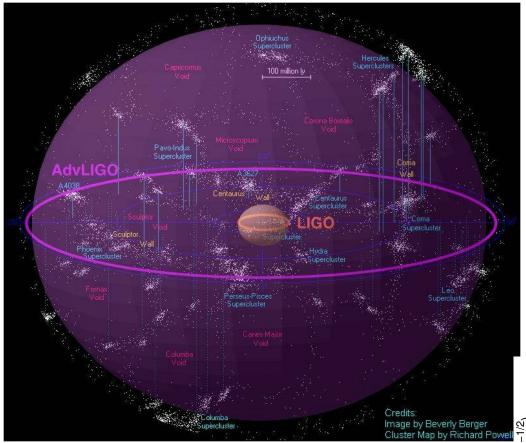




Improved seismic isolation:

Passive → Active

Lowers seismic "wall" to ~10 Hz



(Range x ~10 → Volume x ~1000)

But that sensitivity will not be achieved instantly...

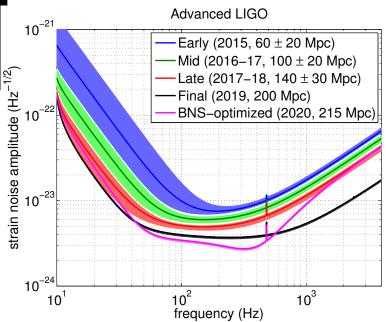
Advanced LIGO

Neutron Star Binaries:

Average range ~ 200 Mpc

Most likely rate ~ 40/year

The science from the first 3 hours of Advanced LIGO should be comparable to 1 year of initial LIGO



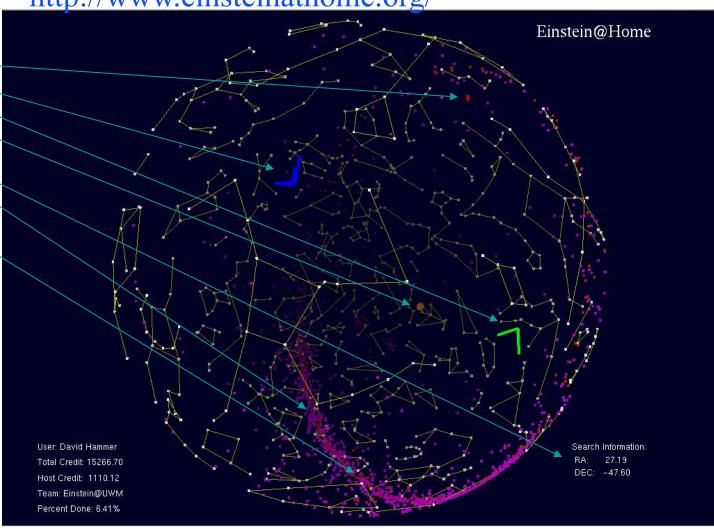
arXiv: 1304.0670



http://www.einsteinathome.org/

- □ GEO-600 Hannover _
- LIGO Hanford
- □ LIGO Livingston
- Current search point
- Current search coordinates
- □ Known pulsars
- Known supernovae remnants

Your computer can help too!





Sources

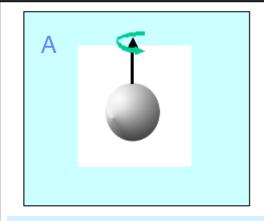
Search methods can detect any type of periodic source.

Upper limits are set on gravitational-wave amplitude, h₀, of rotating triaxial ellipsoid.

Credits:

A. image by Jolien Creighton; LIGO Lab Document G030163-03-Z.

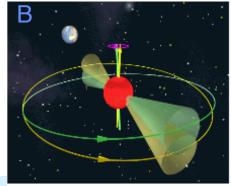
- B. image by M. Kramer; Press Release PR0003, University of Manchester - Jodrell Bank Observatory, 2 August 2000.
- C. image by Dana Berry/NASA; NASA News Release posted July 2, 2003 on Spaceflight Now.
- D. image from a simulation by Chad Hanna and Benjamin Owen; B. J. Owen's research page, Penn State University.



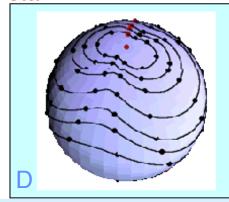
Mountain on neutron



Accreting neutron star



Precessing neutron star

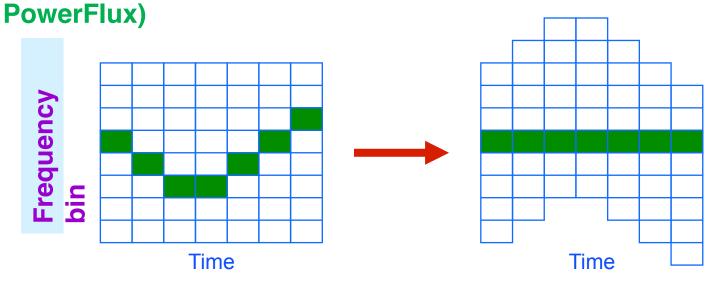


Oscillating neutron star

Searching for continuous waves

Several approaches tried or in development:

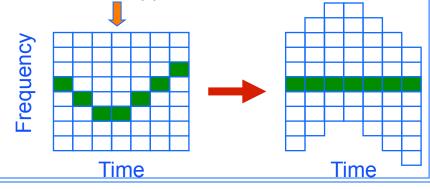
 Summed powers from many short (30-minute) FFTs with skydependent corrections for Doppler frequency shifts → "Semicoherent" (StackSlide, Hough transform,



 Push up close to longest coherence time allowed by computing resources (~1 day) and look for coincidences among outliers in different data stretches (Einstein@Home)

Methods

- Semicoherent Methods
 - StackSlide: add the power
 - Hough: add weighted 1 or 0
 - PowerFlux: add weighted power Track Doppler shift and df/dt



- **Coherent Methods**
 - Bayesian Param. **Estimation**
 - Maximum Likelihood &

$$P(x | h) = \frac{1}{\sqrt{2\pi}\sigma_1} e^{\frac{-(x_1 - h_1)}{2\sigma_1^2}} \frac{1}{\sqrt{2\pi}\sigma_2} e^{\frac{-(x_2 - h_2)^2}{2\sigma_2^2}} \dots$$

$$P(h \mid x) = P(h)P(x \mid h)/P(x) \Rightarrow$$
Time Domain

$$\chi^{2} = \sum_{j} \frac{(x_{j} - h_{j})^{2}}{\sigma_{j}^{2}} \Rightarrow \left(\sum_{j} \frac{x_{j} h_{j}}{\sigma_{j}^{2}} - \frac{1}{2} \sum_{j} \frac{h_{j} h_{j}}{\sigma_{j}^{2}}\right)$$

- Frequency Domain $\Rightarrow (\log \Lambda)_{\max} \Rightarrow F$
- •Weights depend on both noise and antenna patterns:
- Methods can include multi-detector data and coincidence
- •Hierarchical Methods: combine the above to maximize sensitivity.

What is the "indirect spindown limit"?

If a star's age is known (e.g., historical SNR), but its spin is unknown, one can still define an <u>indirect</u> spindown upper limit by assuming gravitar behavior has dominated its lifetime:

$$\tau = \frac{f}{4 \left(\frac{df}{dt} \right)}$$

And substitute into h_{SD} to obtain [K. Wette, B. Owen,... CQG 25 (2008) 235011]

$$h_{ISD} = 2.2 \times 10^{-24} \left[\frac{kpc}{d} \right] \sqrt{\left[\frac{1000 \, yr}{\tau} \right] \left[\frac{I}{10^{45} \, g \cdot cm^2} \right]}$$

Example:

Cassiopeia A
$$\rightarrow$$
 h_{ISD} = **1.2** × **10**⁻²⁴ (d=3.4 kpc, **T**=328 yr)

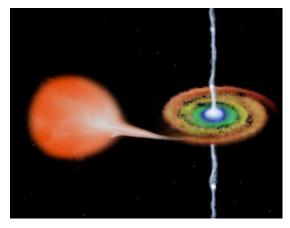
What is the "X-ray flux limit"?

For an LMXB, equating accretion rate torque (inferred from X-ray luminosity) to gravitational wave angular momentum loss (steady state) gives: [R.V. Wagoner ApJ 278 (1984) 345; J. Papaloizou & J.E. Pringle MNRAS 184 (1978) 501; L. Bildsten ApJ 501 (1998) L89]

$$h_{X-ray} \approx 5 \times 10^{-27} \sqrt{\left[\frac{600 Hz}{f_{sig}}\right] \left[\frac{F_x}{10^{-8} erg \cdot cm^{-2} \cdot s^{-1}}\right]}$$

Example: Scorpius X-1

→ $h_{X-ray} \approx 3 \times 10^{-26} [600 \text{ Hz} / f_{sig}]^{1/2}$ ($F_x = 2.5 \times 10^{-7} \text{ erg} \cdot \text{cm}^{-2} \cdot \text{s}^{-1}$)



Courtesy: McGill U.

S1:

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S2:

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S3-S4:

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All-sky search for periodic gravitational waves in LIGO S4 data – PRD 77 (2008) 022001

The Einstein@Home search for periodic gravitational waves in LIGO S4 data – PRD 79 (2009) 022001

Upper limit map of a background of gravitational waves – PRD 76 (2007) 082003 (Cross-correlation – Sco X-1)

S5:

Beating the spin-down limit on gravitational wave emission from the Crab pulsar – APJL 683 (2008) 45

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First search for gravitational waves from the youngest known neutron star – APJ 722 (2010) 1504

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S5:

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