

Combining ab initio calculations and EFT for loosely bounded systems:



Xilin Zhang (Ohio University)

*INT “Universality in few-body systems” Program, University of
Washington, Seattle, WA, April 22, 2014*

*X. Z, K. M. Nollett and D. R. Phillips,
arXiv:1311.6822 (PRC.89.024613) ; 1401.4482 [PRC(R)]*

Outline

- Motivations

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- A toy model: spinless nucleon and core

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- Li7 capture: spins, core excitation, leading order (LO) results
- Be7 capture: nonperturbative Coulomb, LO results

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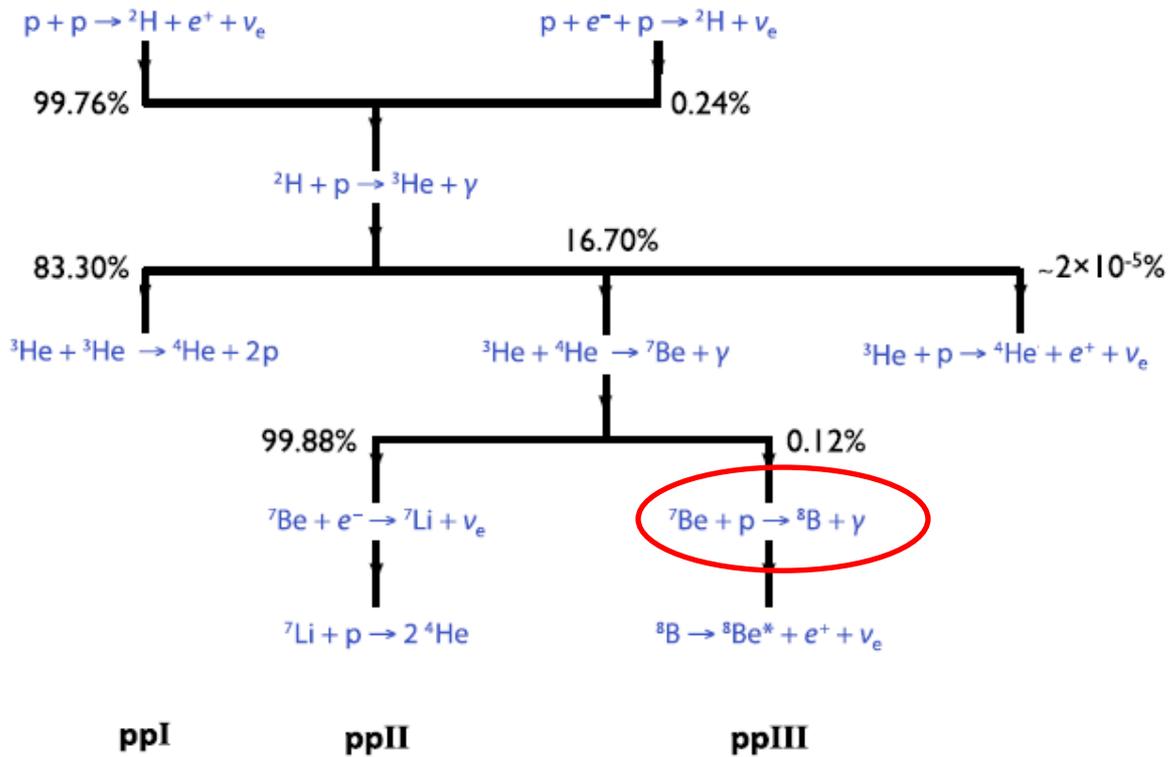
- Motivations
- A toy model: spinless nucleon and core
- $\text{Li}7$ capture: spins, core excitation, leading order (LO) results
- $\text{Be}7$ capture: nonperturbative Coulomb, LO results
- Outlook: Next-to-LO

Motivations

- Astrophysics: solar neutrino flux; solar model;...
- Neutrino mixing parameters

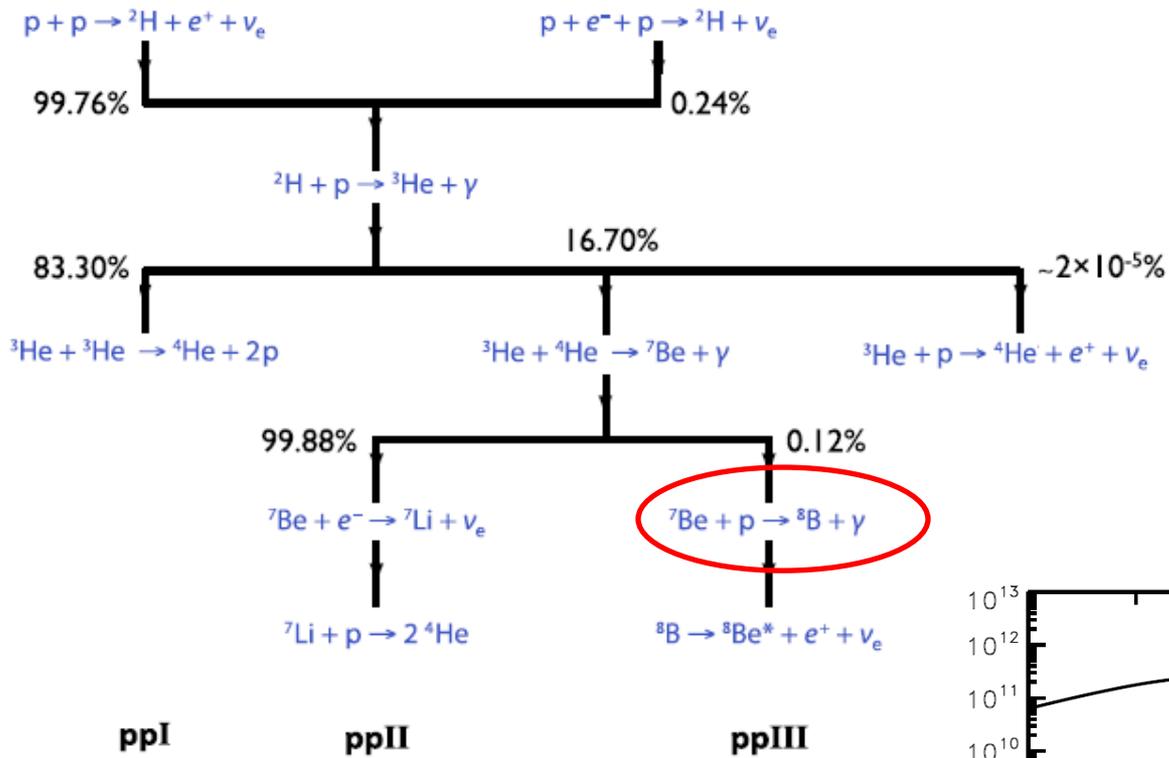
Solar
neutrino
generation

Solar neutrino generation

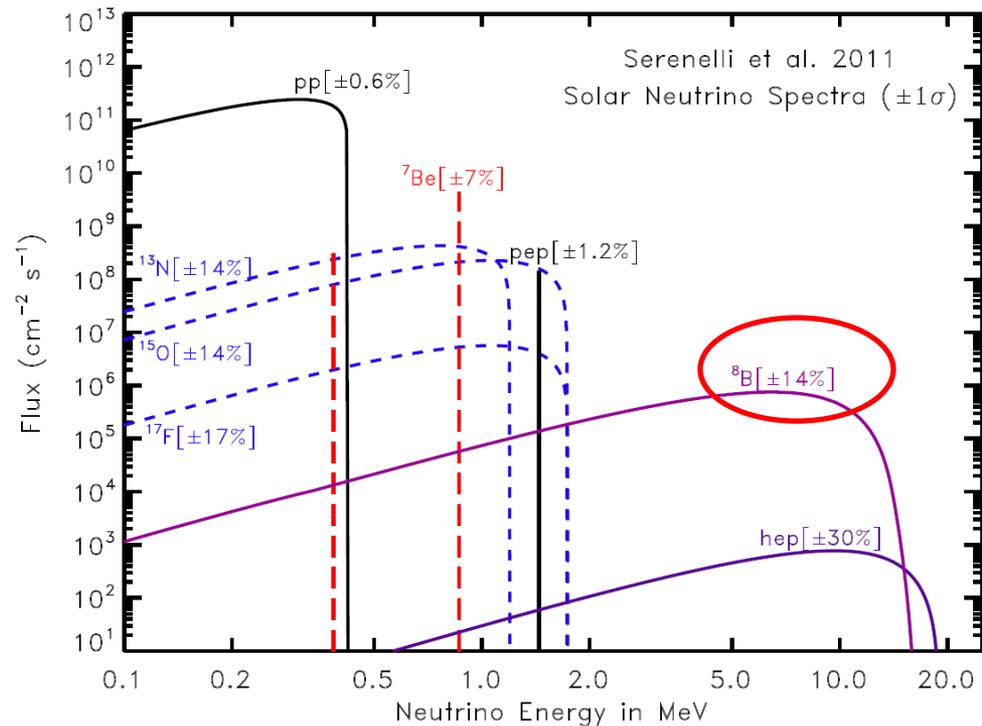


W. C. Haxton et al., arXiv:1208.5723

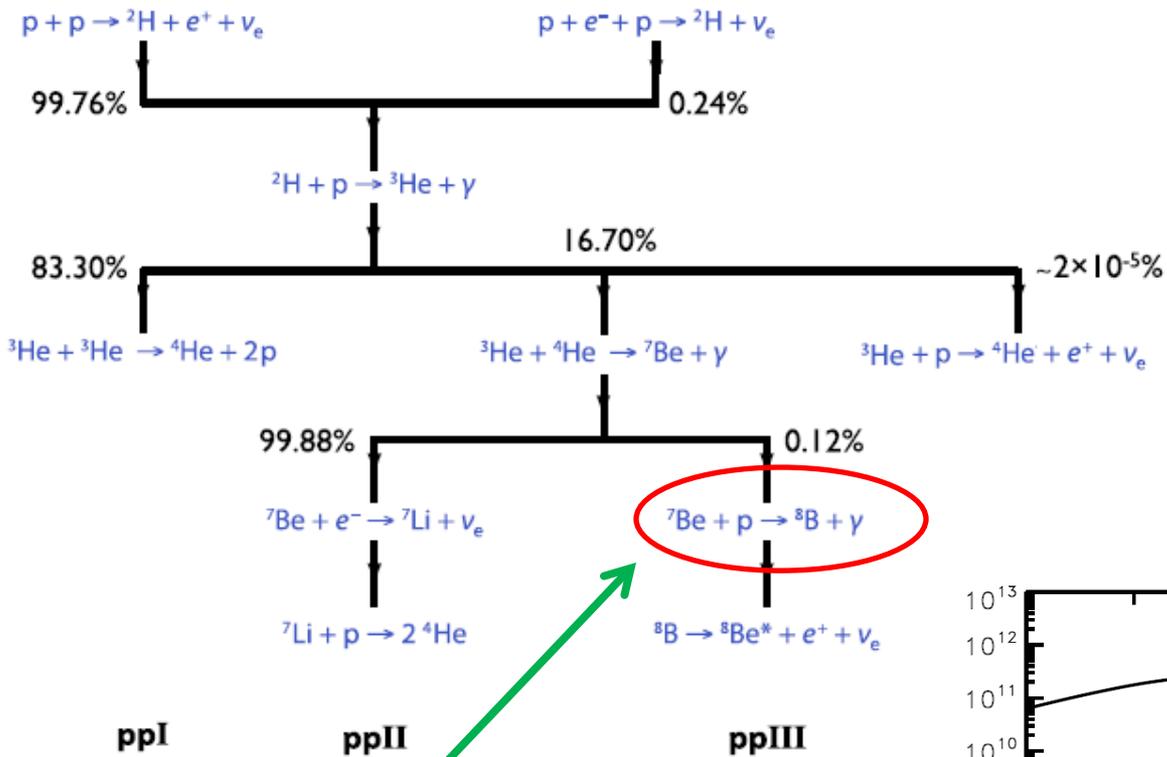
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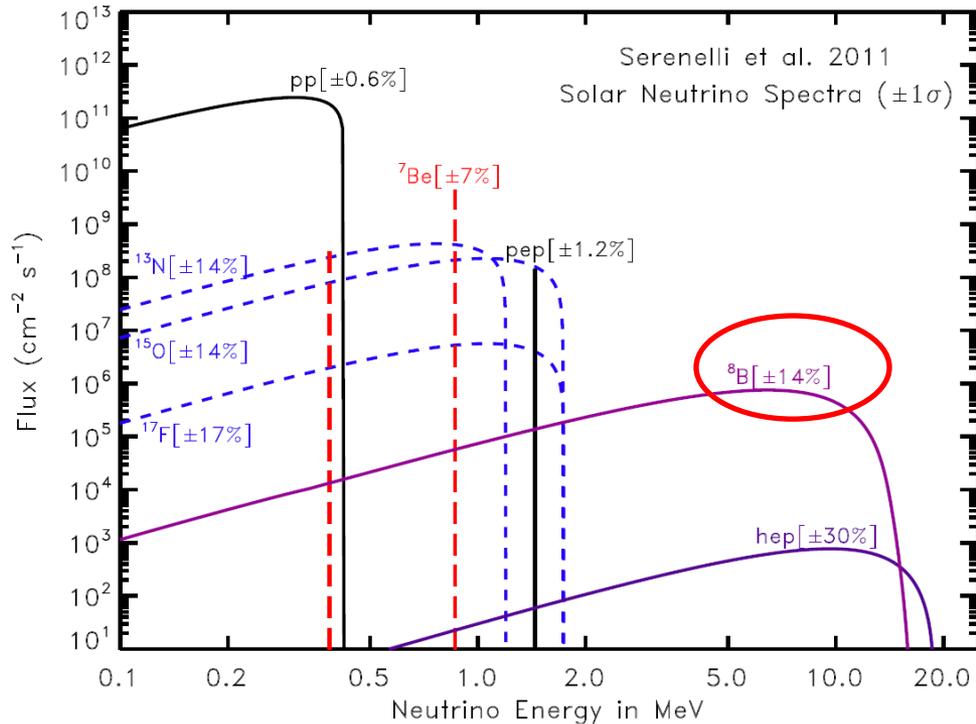
W. C. Haxton et al., arXiv:1208.5723



Solar neutrino generation



${}^7\text{Be} + p \rightarrow {}^8\text{B} + \gamma$



Not experimentally accessible

W. C. Haxton et.al., arXiv:1208.5723

SENSITIVITY OF r -PROCESS NUCLEOSYNTHESIS TO LIGHT-ELEMENT NUCLEAR REACTIONS

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National Astronomical Observatory, 2-21-1 Osawa, Mitaka, Tokyo 181-8588; and Department of Astronomy, Graduate School of Science, University of Tokyo, 7-3-1 Hongo, Bunkyo-ku, Tokyo 113-0033, Japan; sasaqui@th.nao.ac.jp

G. J. MATHEWS AND K. OTSUKI²

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MOST IMPORTANT 18 LIGHT-MASS NUCLEAR REACTIONS, ADOPTED “STANDARD” THERMONUCLEAR REACTION RATES $\lambda_i(0)$, AND UNCERTAINTIES

No.	Reaction	$N_{Av}\langle\sigma v\rangle$	$1\sigma^a$	Reference ^b
(1).....	$\alpha(\alpha n, \gamma)^9\text{Be}$	$N_{Av}^2\langle\alpha\alpha n\rangle = 2.43 \times 10^9 T_9^{-2/3} \exp[-13.490 T_9^{-1/3} - (T_9/0.15)^2](1 + 74.5 T_9)$ $+ 6.09 \times 10^5 T_9^{-3/2} \exp(-1.054/T_9)(1 - 58.80 T_9 - 1.794 \times 10^4 T_9^2)$ $+ 2.969 \times 10^6 T_9^3 - 1.535 \times 10^8 T_9^4 + 2.610 \times 10^9 T_9^5$	$\pm 35\%$	1
(2).....	$\alpha(t, \gamma)^7\text{Li}^c$	$3.032 \times 10^5 T_9^{-2/3} \exp(-8.09/T_9^{1/3})(1.0 + 0.0516 T_9^{1/3} + 0.0229 T_9^{2/3})$ $+ 8.28 \times 10^{-3} T_9 - 3.28 \times 10^{-04} T_9^{4/3} - 3.01 \times 10^{-04} T_9^{5/3}$ $+ 5.109 \times 10^5 T_{9*}^{5/6} T_9^{-3/2} \exp(-8.068/T_{9*}^{1/3})$	$\pm 30\%$	2
(3).....	$^7\text{Li}(n, \gamma)^8\text{Li}$	$4.90 \times 10^3 + 9.96 \times 10^3 T_9^{-3/2} \exp(-2.62/T_9)$	$\pm 35\%$	3

Li7 capture is used to constrain models of Be7 capture.

Motivations

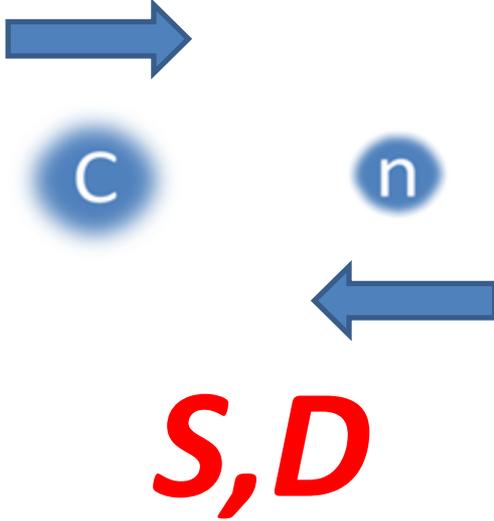
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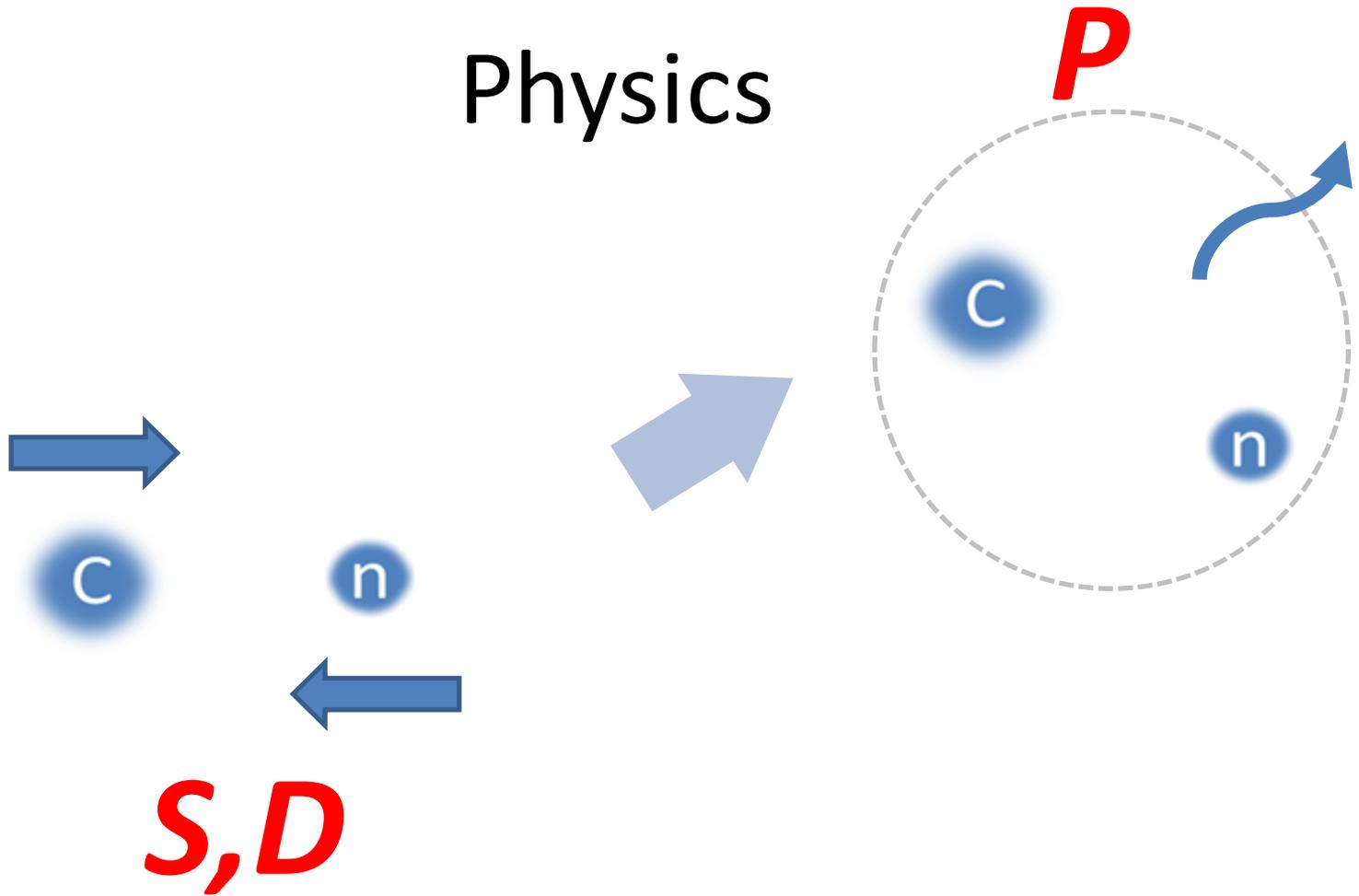
- Astrophysics: solar neutrino flux; solar model;...
- Neutrino mixing parameters
- EFT: a simple picture; systematic expansion (Lagrangian); uncertainty estimate
- Parameters: ab initio bound state information
- Ab initio reaction calculation

A toy model

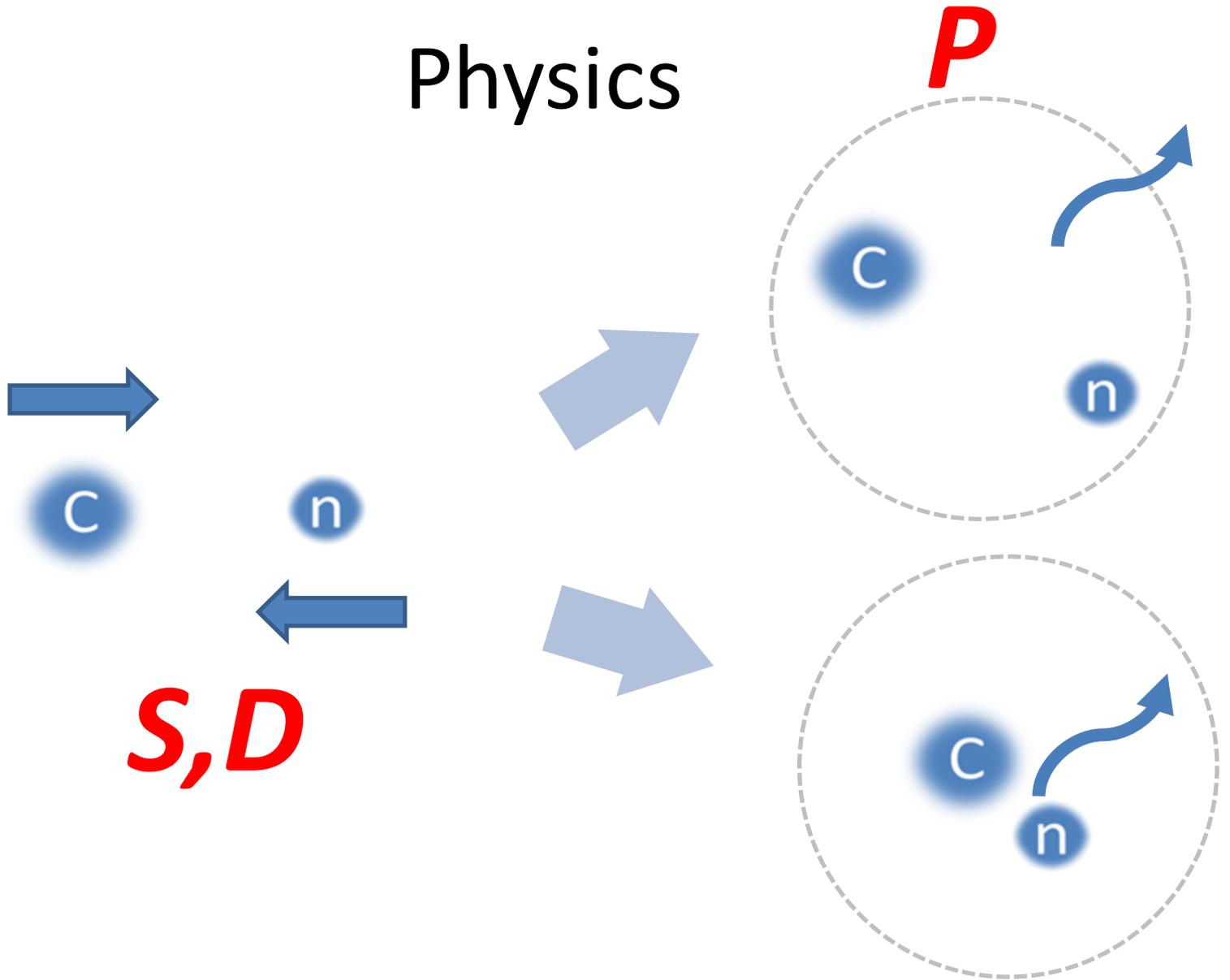
Physics



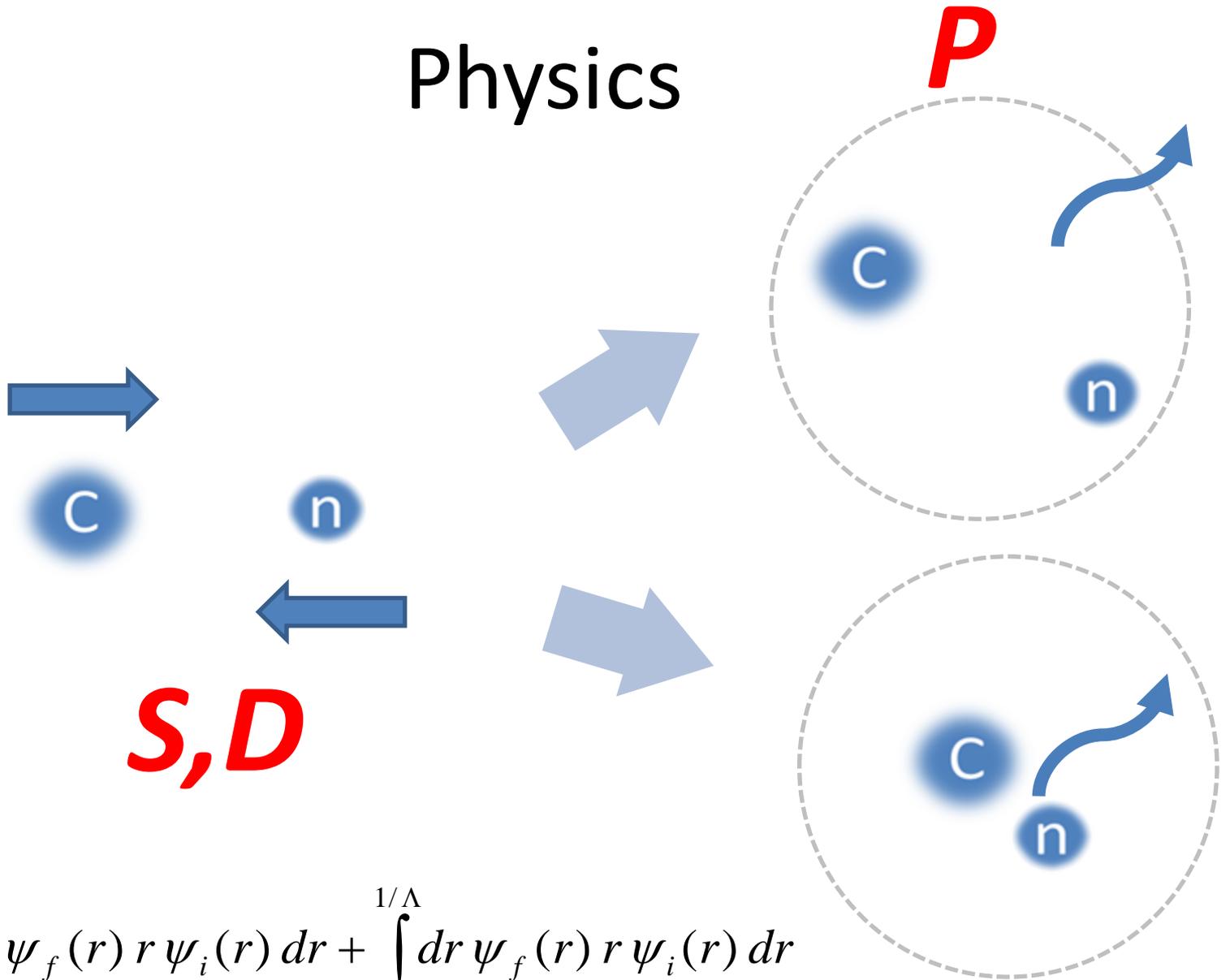
Physics



Physics

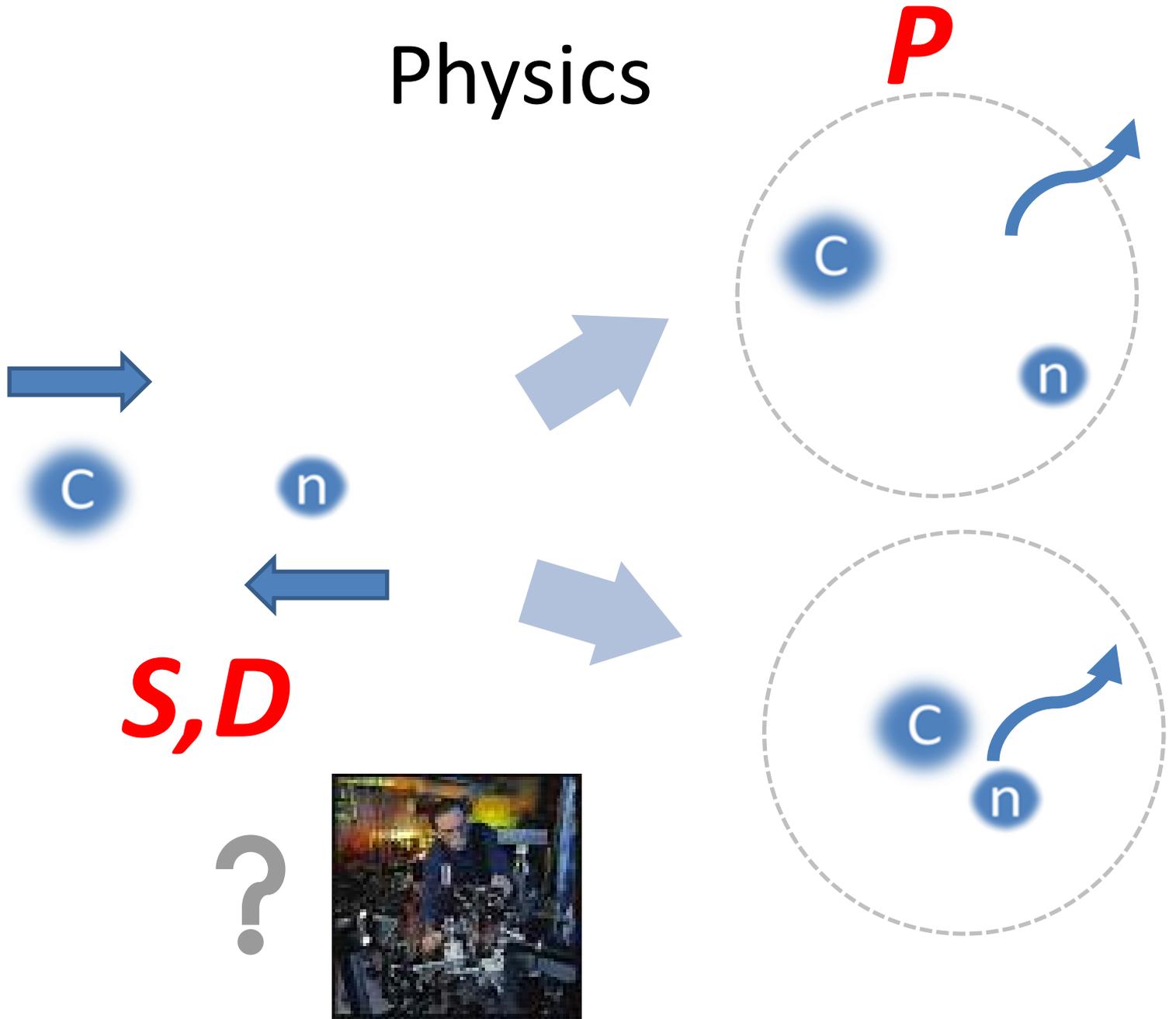


Physics

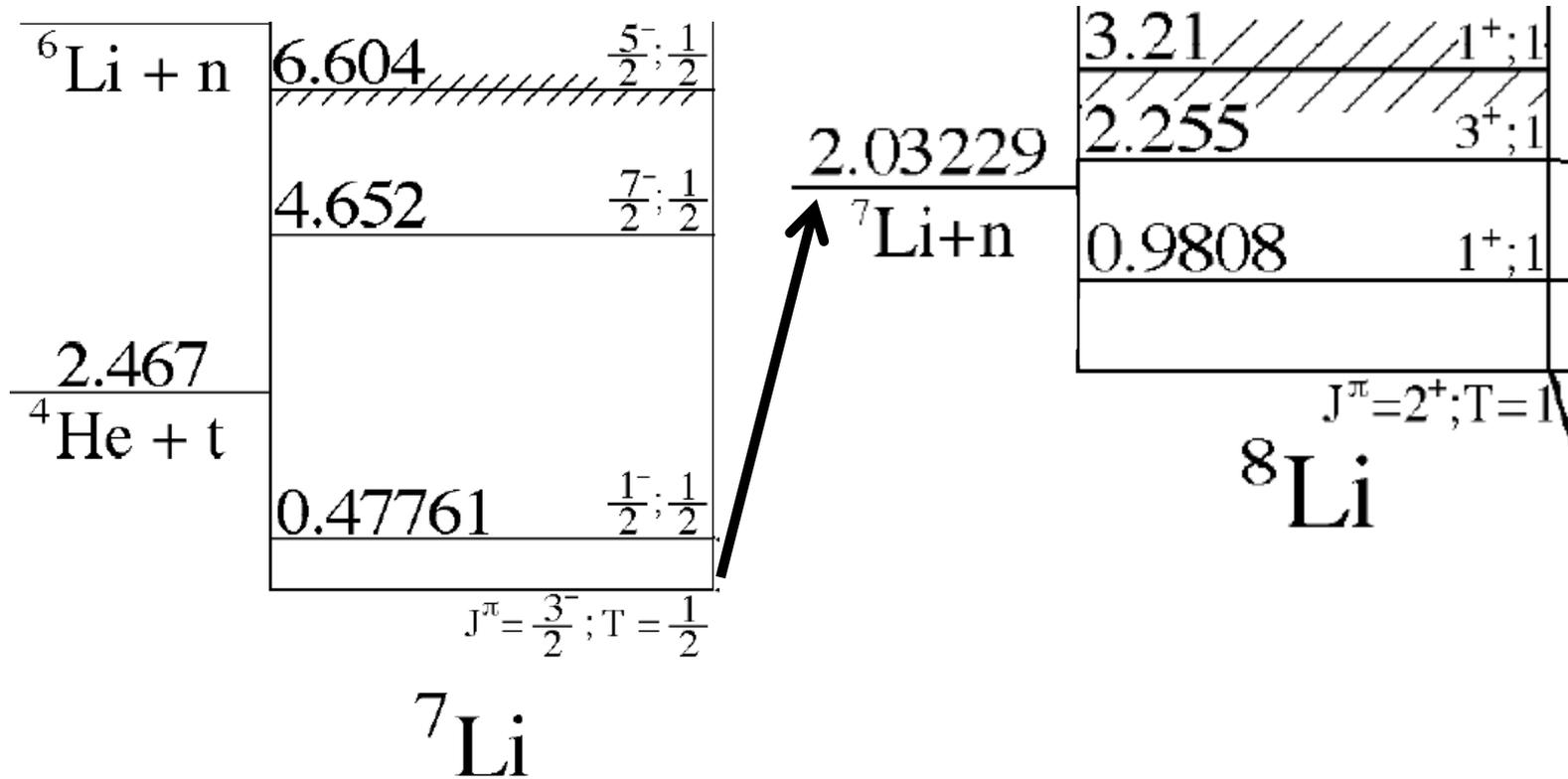


$$M \sim \int_{1/\Lambda}^{+\infty} dr \psi_f(r) r \psi_i(r) dr + \int_0^{1/\Lambda} dr \psi_f(r) r \psi_i(r) dr$$

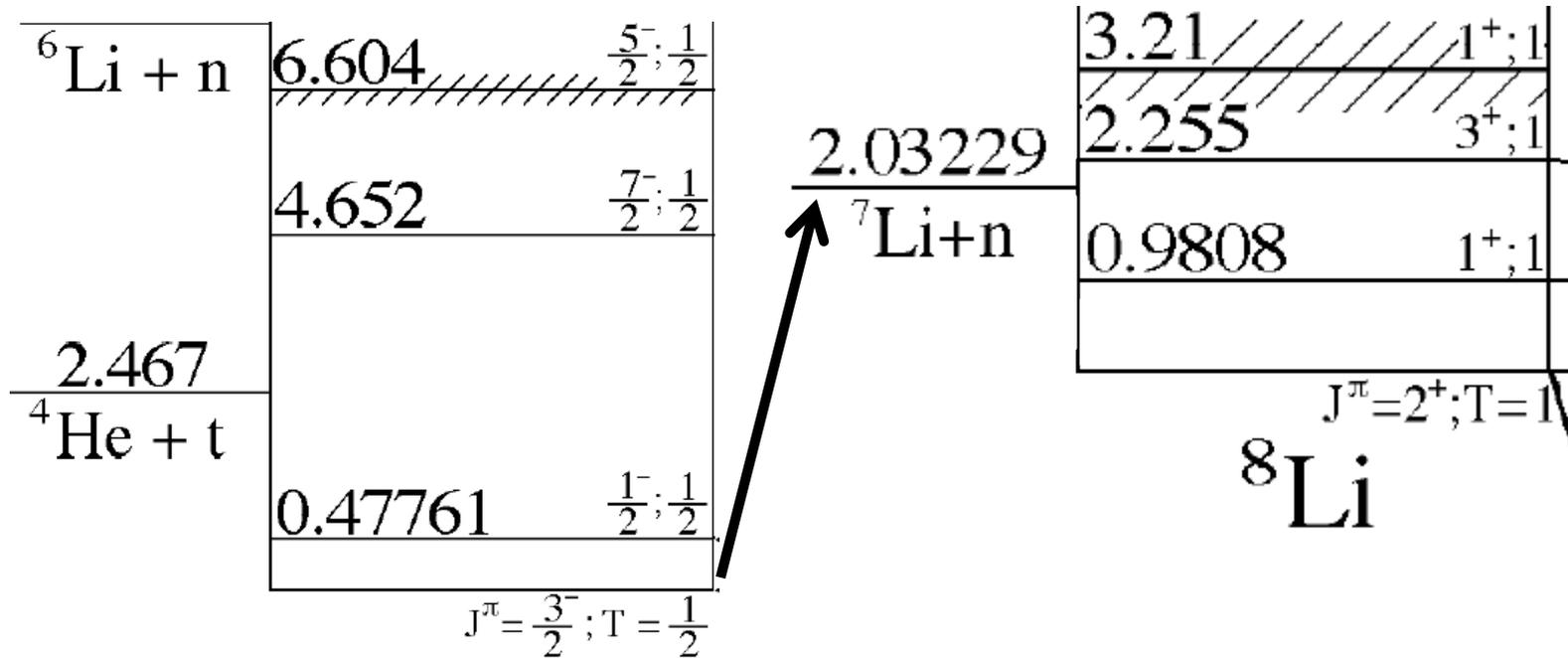
Physics



Gross features: p-wave



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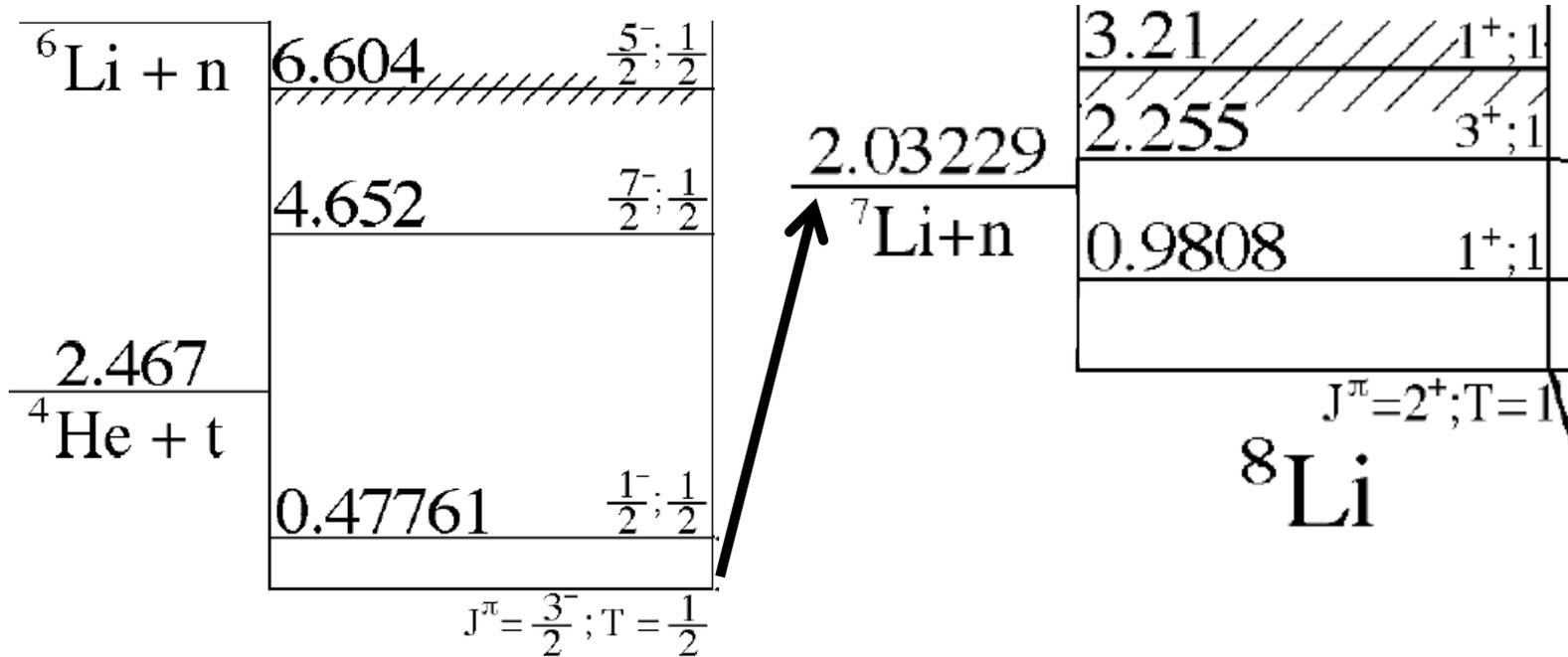
${}^7\text{Li}$

$$\Lambda \sim \sqrt{2M_{43}B_{Li7}} \sim 100 \text{ MeV}$$

$$\gamma \sim \sqrt{2M_{71}B_{Li8}} = 57.8 \text{ MeV}$$

$$\frac{\gamma}{\Lambda} \sim 0.5$$

Gross features: p-wave



${}^7\text{Li}$

*Shallow p-wave
bound state*

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Gross features: s-wave

Parameter	Channel	Value	Assigned scaling
$a_{({}^5S_2)}$	S -wave, $S = 2$	$-3.63(5)$ fm	$1/\gamma$
$a_{({}^3S_1)}$	S -wave, $S = 1$	$0.87(7)$ fm	$1/\Lambda$

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L. Koester, K. Knopf, and W. Waschkowski, Z. Phys. A 312, 81 (1983)

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Large s-wave scattering length

L. Koester, K. Knopf, and W. Waschkowski, Z. Phys. A 312, 81 (1983)

S-wave in EFT

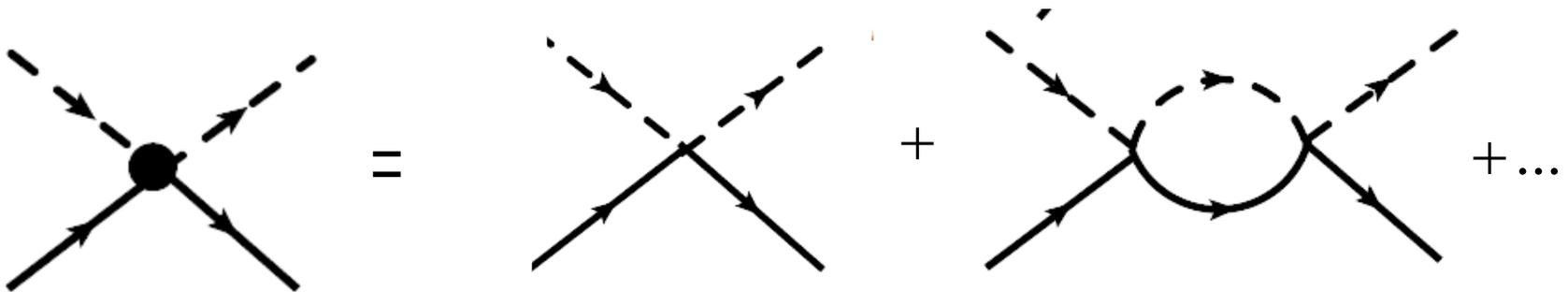
$$\mathcal{L}_0 = n^\dagger \left(i\partial_t + \frac{\nabla^2}{2M_n} \right) n + c^\dagger \left(i\partial_t + \frac{\nabla^2}{2M_c} \right) c$$

$$\mathcal{L}_S = g c^\dagger n^\dagger c n \quad g(\mu) = \frac{2\pi}{M_R \left(2\mu - \frac{1}{a_0} \right)}$$

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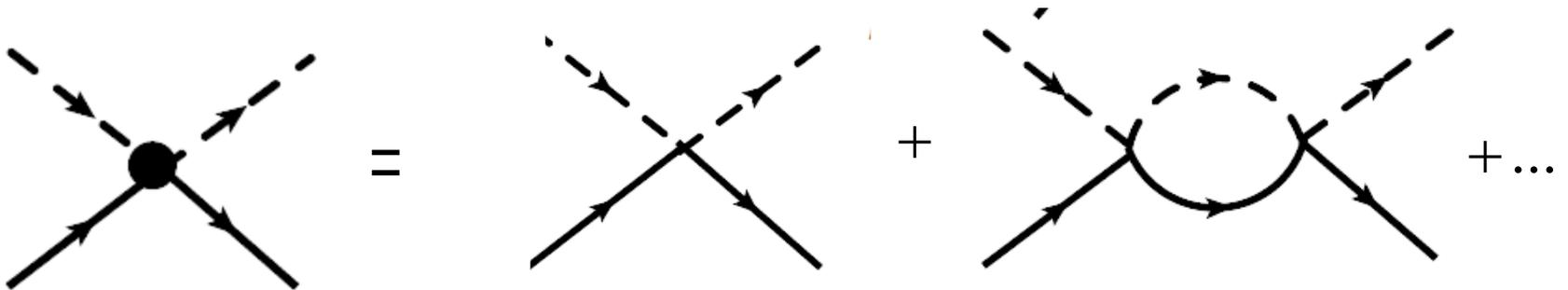
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S-wave in EFT

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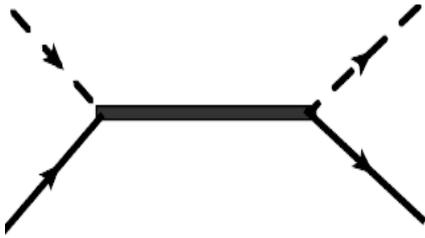
$$\mathcal{L}_S = g c^\dagger n^\dagger c n \quad g(\mu) = \frac{2\pi}{M_R \left(2\mu - \frac{1}{a_0} \right)} \quad \rightarrow$$



$$\rightarrow T = \frac{2\pi}{M_R} \frac{1}{a_0^{-1} + ik}$$

One parameter: g (or a0)

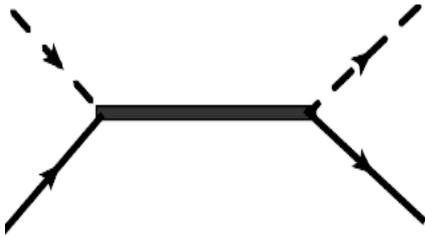
P-wave in EFT



$$\langle p' | T(E) | p \rangle = \frac{6\pi}{M_{\text{R}}} \frac{p' \cdot p}{a_1^{-1} - \frac{1}{2}r_1 k^2 + ik^3}$$

$$k^3 \cot \delta_1 = -\frac{1}{a_1} + \frac{1}{2}r_1 k^2 + \dots$$

P-wave in EFT

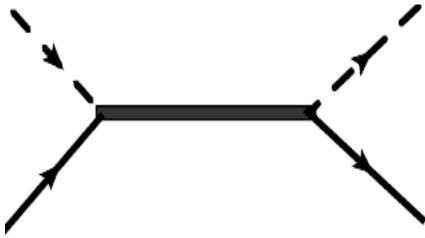


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$$a_1 \sim \frac{1}{\Lambda \gamma^2} \quad \text{and} \quad r_1 \sim \Lambda$$

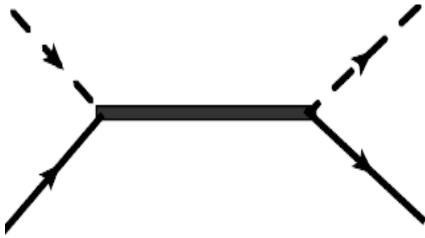


Unnatural



Natural

P-wave in EFT



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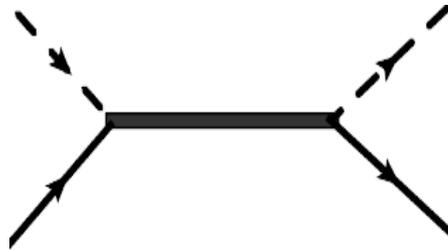
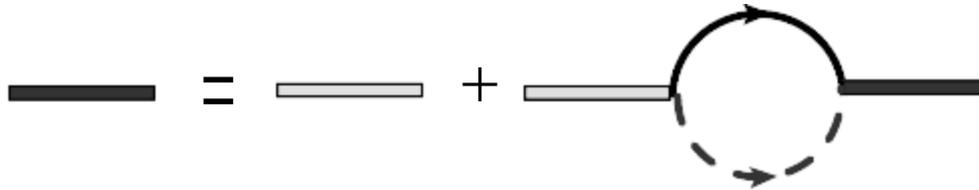
Natural

P-wave in EFT

$$\mathcal{L}_P = \pi^{\dagger i} \left(i\partial_t + \frac{\nabla^2}{2M_{\text{nc}}} + \Delta \right) \pi_i + h \pi^{\dagger i} n_i (V_n - V_c)_i c + \text{C.C.} .$$

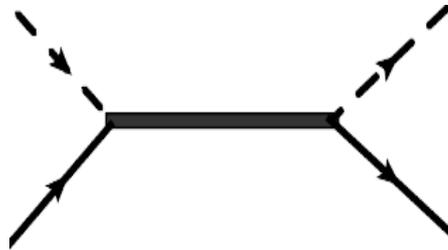
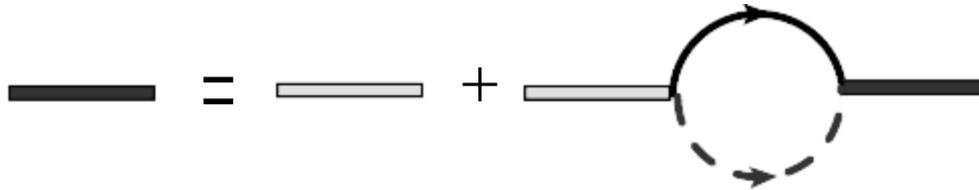
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$$\langle p' | T(E) | p \rangle = \frac{h^2}{M_R^2} (p' \cdot p) D(E, 0) = \frac{6\pi}{M_R} \frac{p' \cdot p}{a_1^{-1} - \frac{1}{2} r_1 k^2 + ik^3}$$

Two parameters: Delta and h (or a1 and r1)

P-wave in EFT

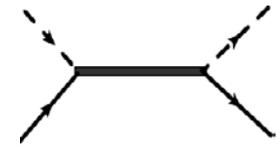
$$\langle \mathbf{r}' | \frac{1}{E - H} | \mathbf{r} \rangle = \langle \mathbf{r}' | \frac{1}{E - H_0} + \frac{1}{E - H_0} T \frac{1}{E - H_0} | \mathbf{r} \rangle$$

*Asymptotic
normalization
coefficient (ANC)*

$$\xrightarrow{E \rightarrow -B} C^2 \times \sum_j \frac{\phi_j(\mathbf{r}') \phi_j^*(\mathbf{r})}{E + B} .$$

$$\phi_j(\mathbf{r}) = \left(1 + \frac{1}{\gamma r} \right) Y_{1j}(\hat{r}) \frac{e^{-\gamma r}}{r}$$

P-wave in EFT



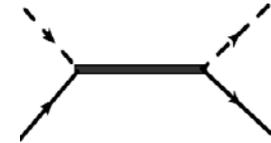
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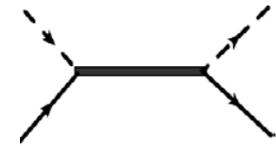
$$C = \sqrt{\frac{-2\gamma^2}{r_1 + 3\gamma}}$$

$$\frac{1}{a_1} + \frac{1}{2} r_1 \gamma^2 + \gamma^3 = 0$$

}

a_1 and r_1 (or h and Δ)

P-wave in EFT



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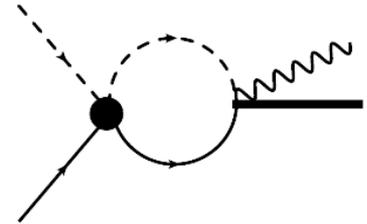
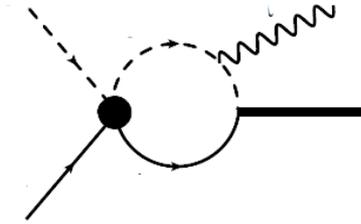
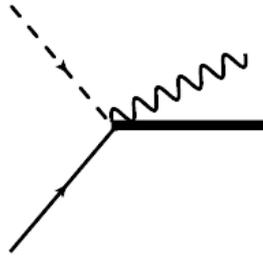
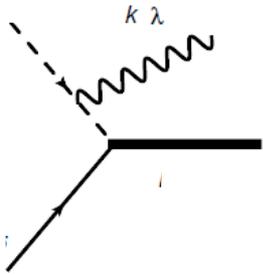
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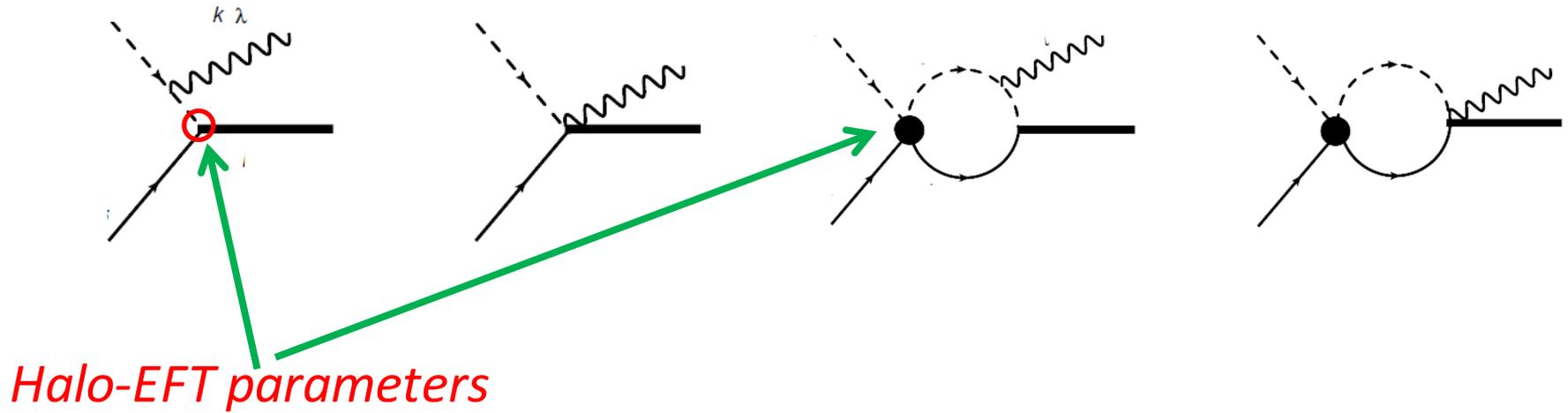
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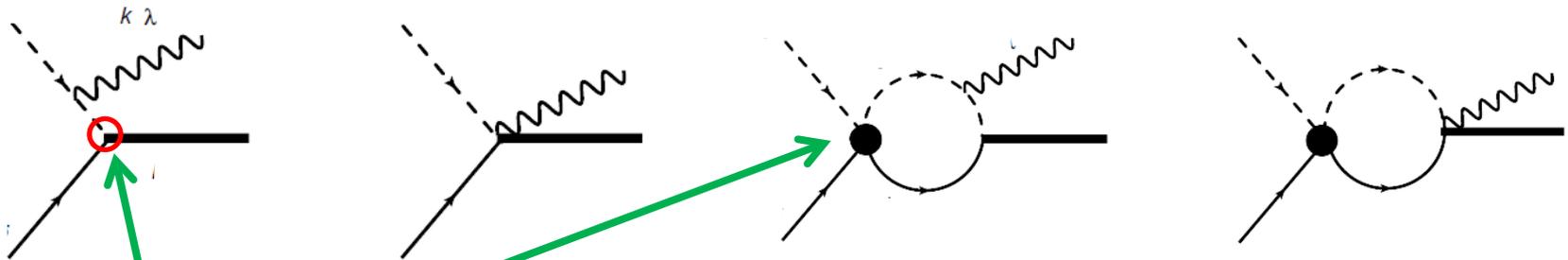
Radiative capture: LO



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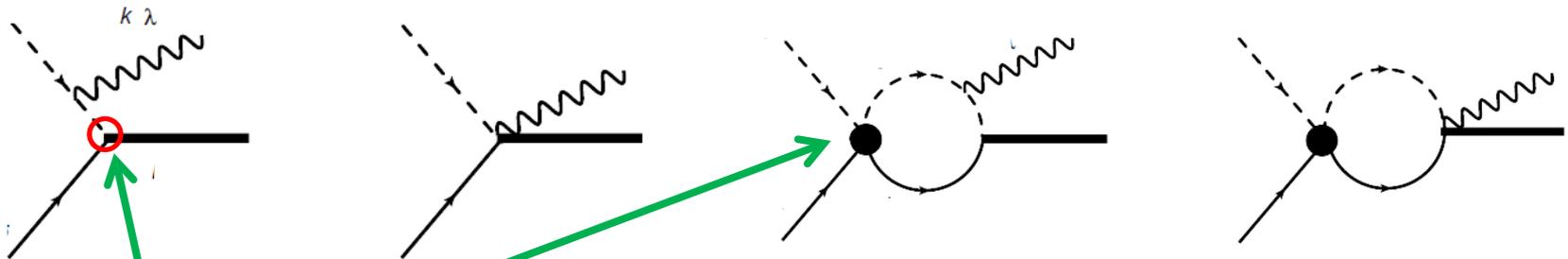


Halo-EFT parameters

$$\mathcal{M} \sim ie_c h \sqrt{Z} \left[\frac{\epsilon^*(\lambda) \cdot \mathbf{V}_c}{p_c^0 - \omega - \frac{(\mathbf{p}_c - \mathbf{k})^2}{2M_c} + i\epsilon} \left(\frac{p_c}{M_R} - \frac{\mathbf{k}}{M_c} \right)_j + (1 + X(p_c; \gamma, a_0)) \frac{\epsilon^*(\lambda)_j}{M_c} \right]$$

$$X(p_c; \gamma, a) \equiv \frac{(-)i}{a^{-1} + ip_c} \left[p_c - \frac{2}{3} i \frac{\gamma^3 - ip_c^3}{\gamma^2 + p_c^2} \right]$$

Radiative capture: LO



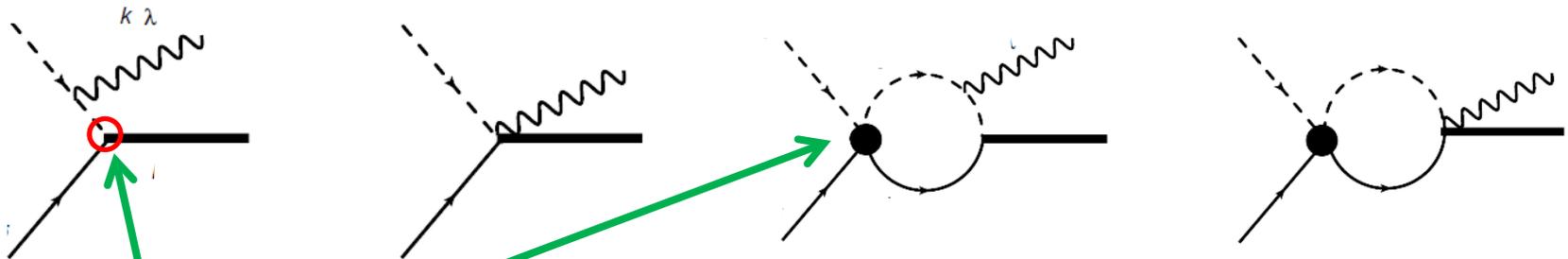
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\downarrow
 C

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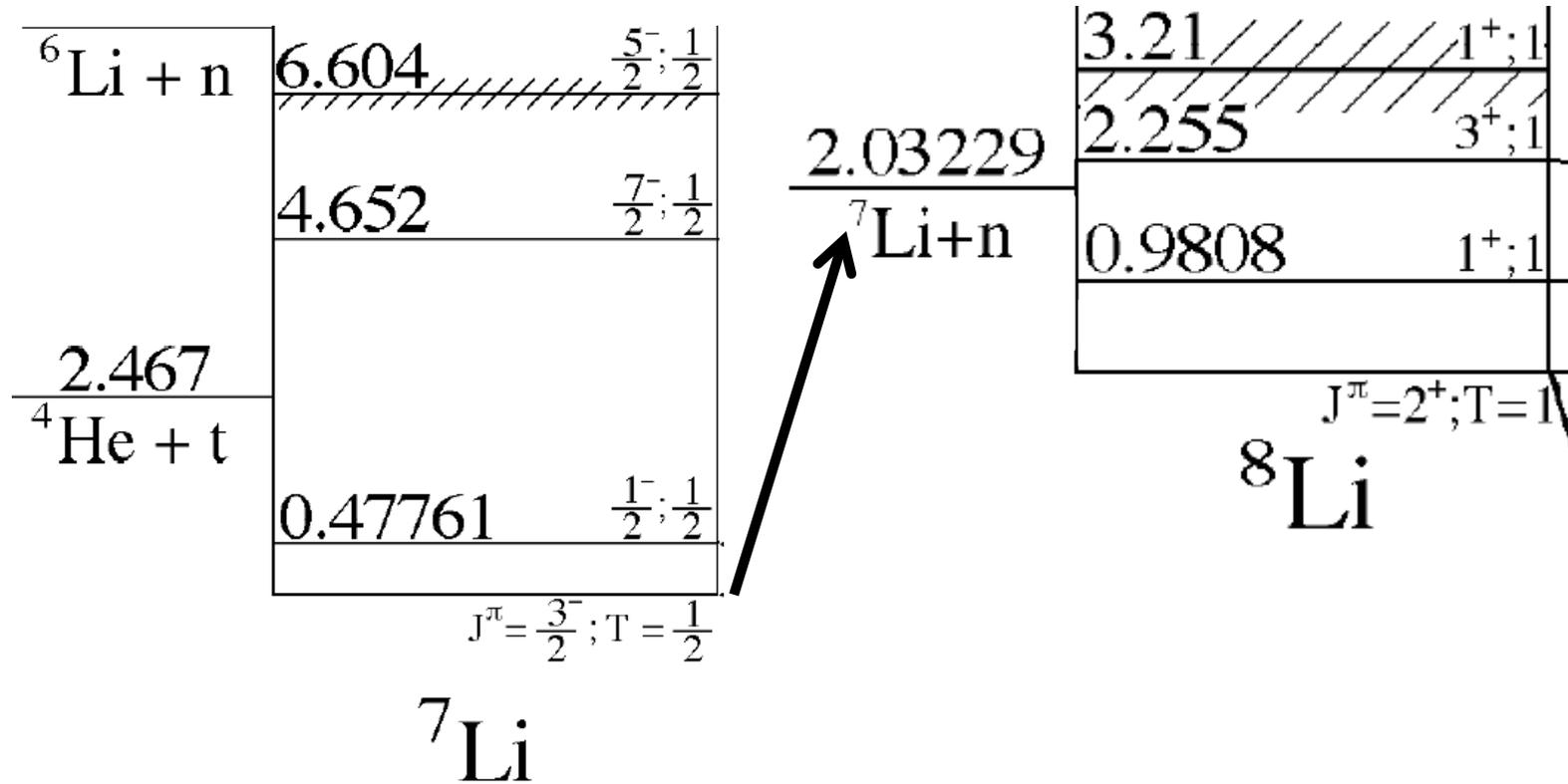
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$$a \sim \frac{1}{\gamma} \Rightarrow X \sim 1, \quad a \sim \frac{1}{\Lambda} \Rightarrow X \sim \frac{\gamma}{\Lambda}$$

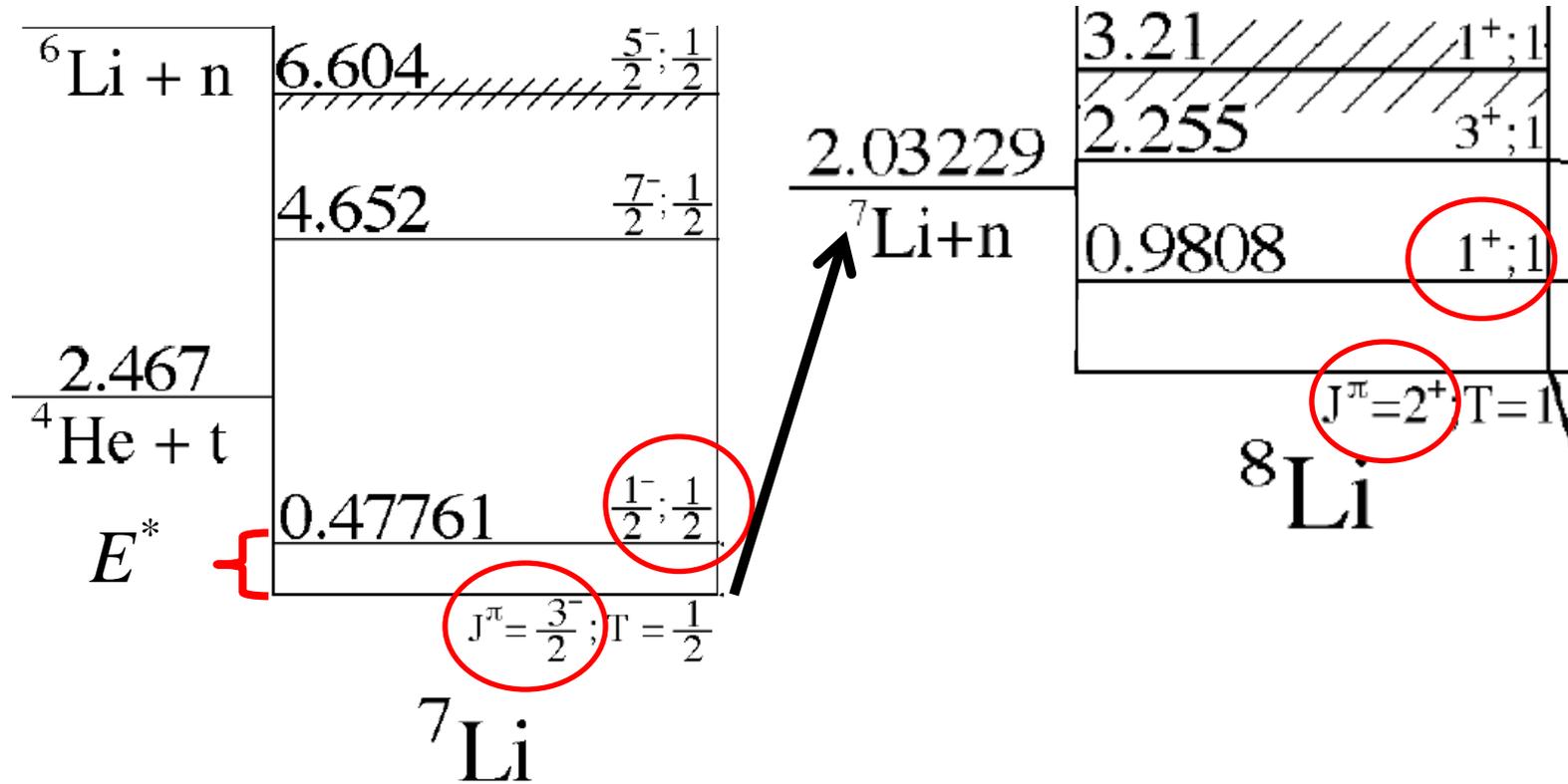


G. Rupak and R. Higa, *Phys. Rev. Lett.* 106, 222501 (2011)

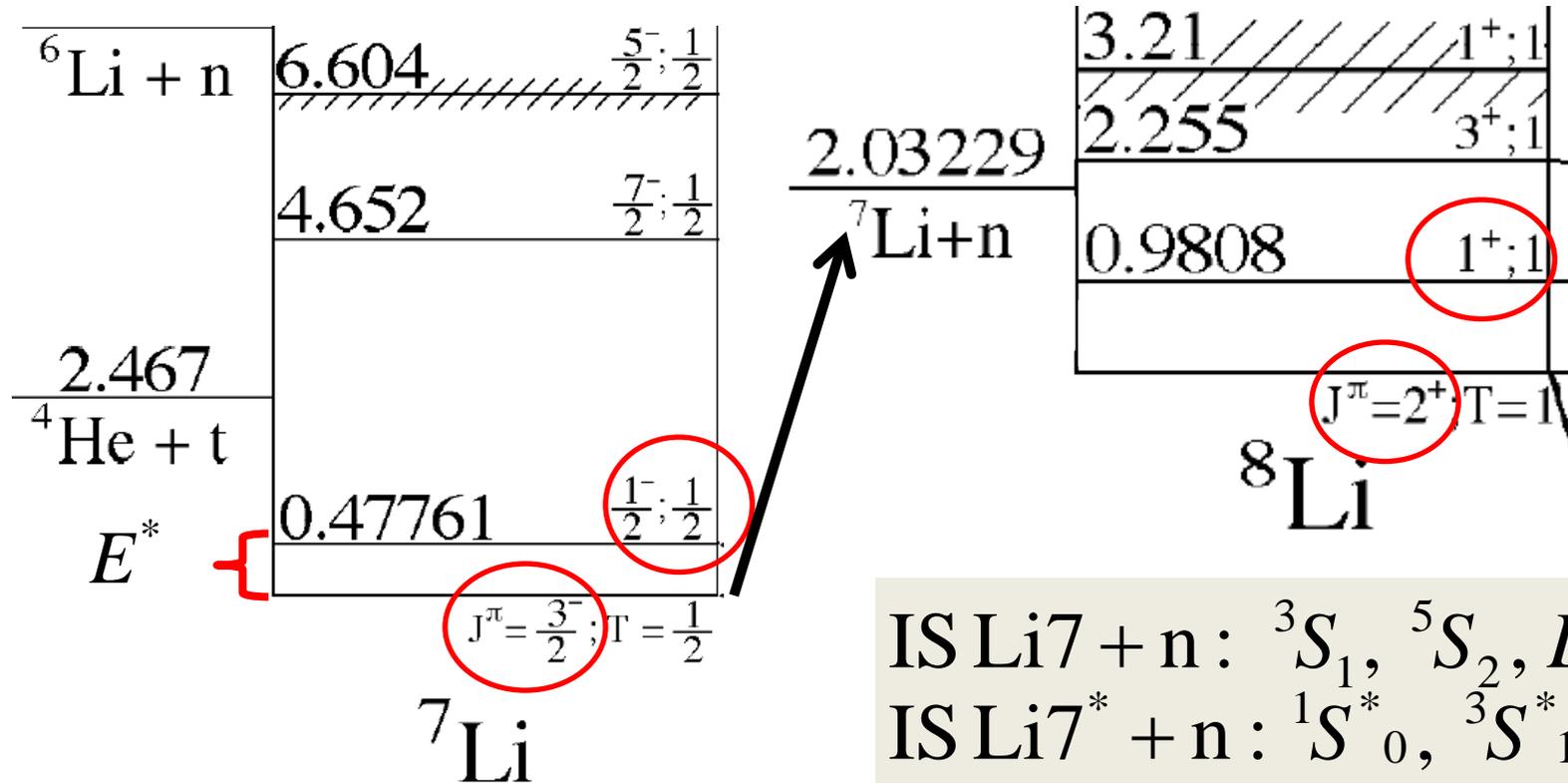
Scales, spins, core excitations



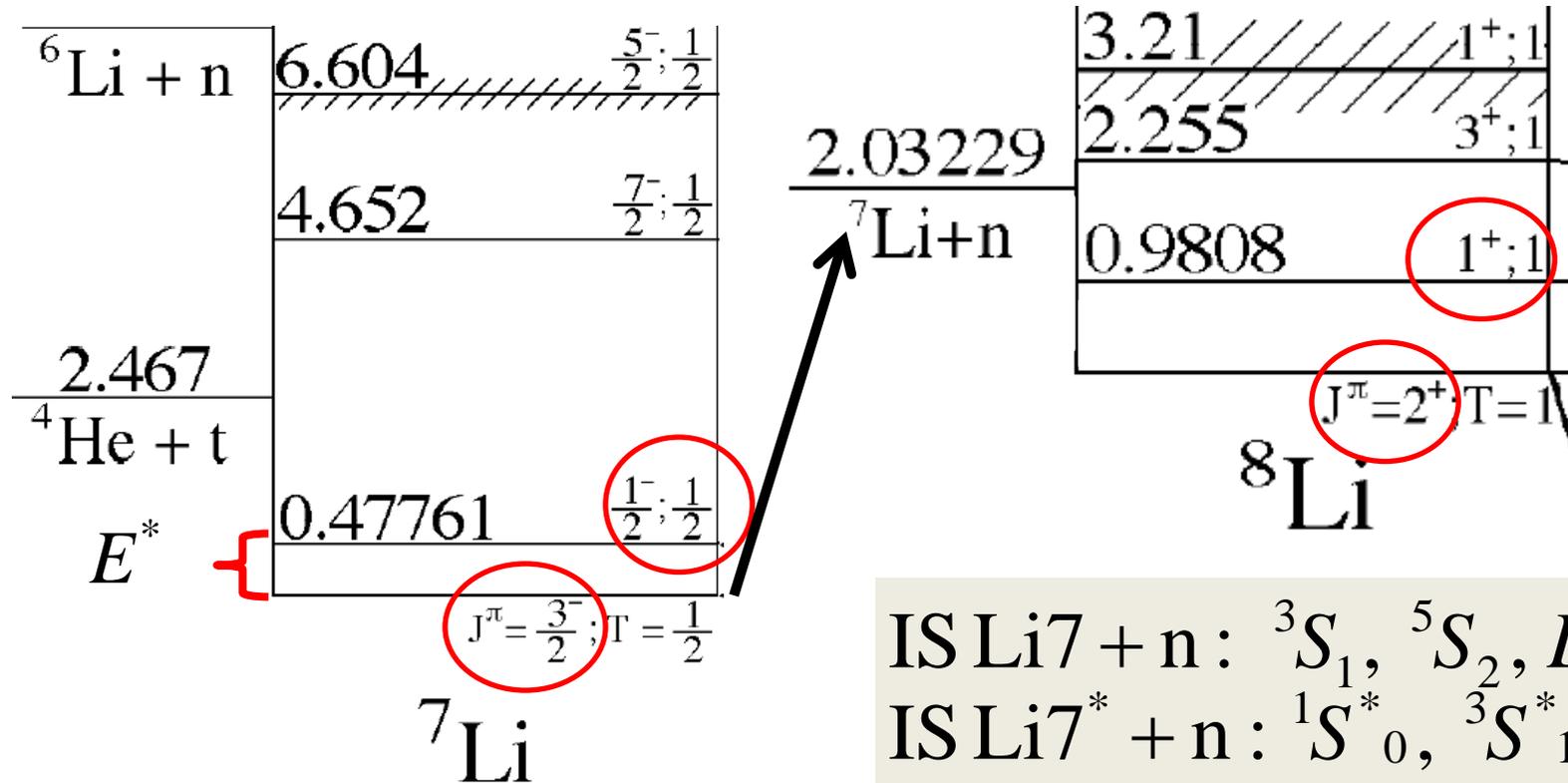
Scales, spins, core excitations



Scales, spins, core excitations

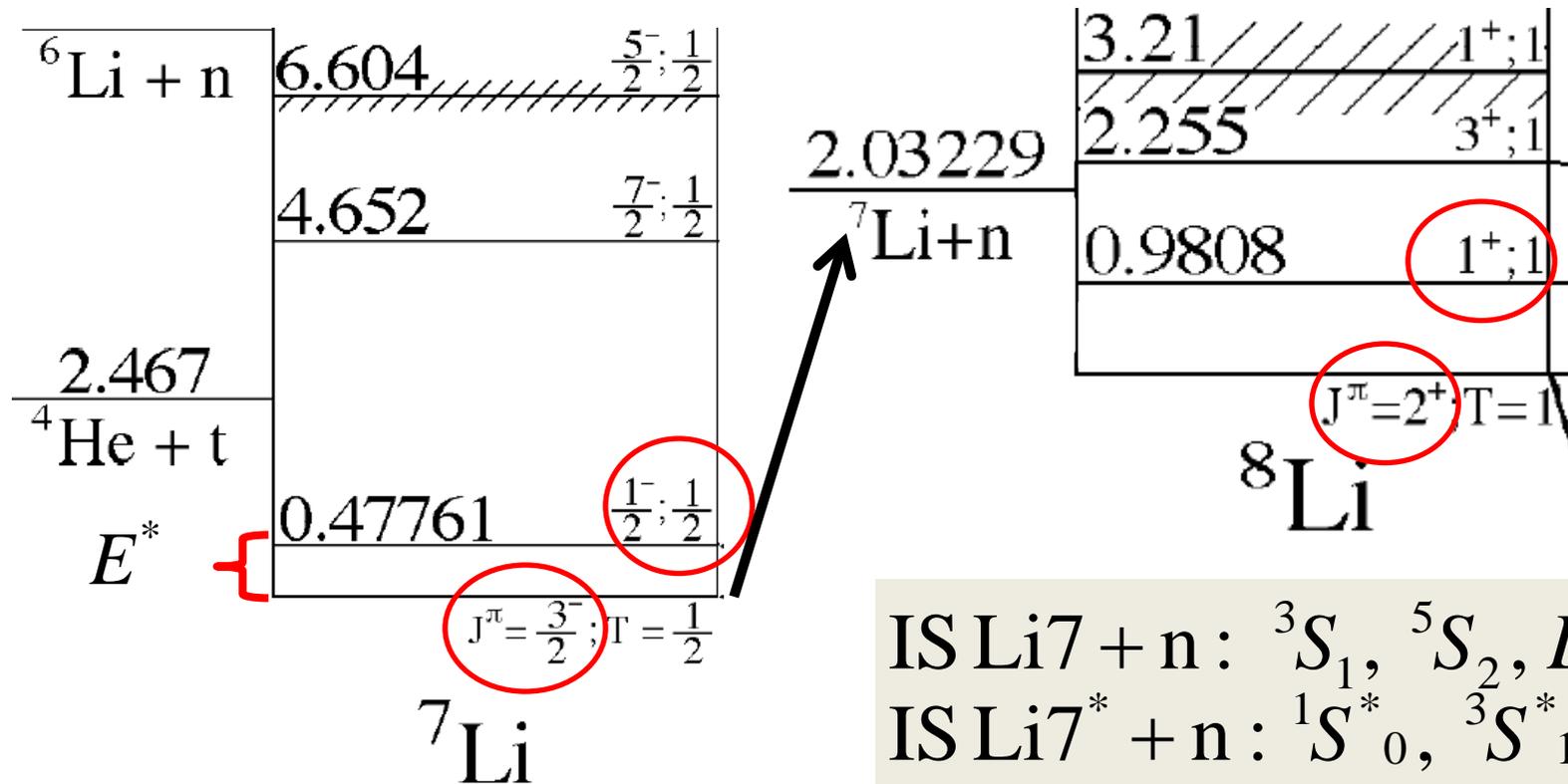


Scales, spins, core excitations



FS(2^+) $\text{Li}7 + n$: ${}^3P_2, {}^5P_2$
 FS(2^+) $\text{Li}7^* + n$: ${}^3P_2^*$

Scales, spins, core excitations



IS Li7 + n : ${}^3S_1, {}^5S_2, D$
 IS Li7* + n : ${}^1S_0^*, {}^3S_1^*$

FS(2^+) Li7 + n : ${}^3P_2, {}^5P_2$
 FS(2^+) Li7* + n : ${}^3P_2^*$

FS(1^+) Li7 + n : ${}^3P_1, {}^5P_1$
 FS(1^+) Li7* + n : ${}^1P_1^*, {}^3P_1^*$

Scales, spins, core excitations

$$\Lambda \approx 100 - 300 \text{ MeV}$$

Momentum scale	Definition	Value
γ	$\sqrt{2M_R B_{8\text{Li}}}$	57.8 MeV
γ^*	$\sqrt{2M_R (B_{8\text{Li}} + E^*)}$	65.1 MeV
γ_Δ	$\sqrt{2M_R E^*}$	30.0 MeV
$\tilde{\gamma}$	$\sqrt{2M_R B_{8\text{Li}^*}}$	41.6 MeV
$\tilde{\gamma}^*$	$\sqrt{2M_R (B_{8\text{Li}^*} + E^*)}$	51.3 MeV

Parameter	Channel	Value	Assigned scaling
$a_{({}^5S_2)}$	<i>S</i> -wave, $S = 2$	$-3.63(5) \text{ fm}$	$1/\gamma$
$a_{({}^3S_1)}$	<i>S</i> -wave, $S = 1$	$0.87(7) \text{ fm}$	$1/\Lambda$
r	<i>P</i> -wave, $J = 2$	$-1.43(2) \text{ fm}^{-1}$	Λ
\tilde{r}	<i>P</i> -wave, $J = 1$	$-1.86(6) \text{ fm}^{-1}$	Λ

EFT

$$\begin{aligned}\mathcal{L}_0 = & n^{\dagger\sigma} \left(i\partial_t + \frac{\nabla^2}{2M_n} \right) n_\sigma + c^{\dagger a} \left(i\partial_t + \frac{\nabla^2}{2M_c} \right) c_a \\ & + d^{\dagger\delta} \left(i\partial_t + \frac{\nabla^2}{2M_c} \right) d_\delta + \pi^{\dagger\alpha} \left(i\partial_t + \frac{\nabla^2}{2M_{nc}} + \Delta \right) \pi_\alpha \\ & + \tilde{\pi}^{\dagger i} \left(i\partial_t + \frac{\nabla^2}{2M_{nc}} + \tilde{\Delta} \right) \tilde{\pi}_i ,\end{aligned}$$

EFT

$$\begin{aligned}\mathcal{L}_0 &= n^{\dagger\sigma} \left(i\partial_t + \frac{\nabla^2}{2M_n} \right) n_\sigma + c^{\dagger a} \left(i\partial_t + \frac{\nabla^2}{2M_c} \right) c_a \\ &+ d^{\dagger\delta} \left(i\partial_t + \frac{\nabla^2}{2M_c} \right) d_\delta + \pi^{\dagger\alpha} \left(i\partial_t + \frac{\nabla^2}{2M_{nc}} + \Delta \right) \pi_\alpha \\ &+ \tilde{\pi}^{\dagger i} \left(i\partial_t + \frac{\nabla^2}{2M_{nc}} + \tilde{\Delta} \right) \tilde{\pi}_i ,\end{aligned}$$

$$\begin{aligned}\mathcal{L}_S &= g_{(3S_1)} c^{\dagger a'} n^{\dagger\sigma'} T_{a'\sigma'}^i T_i^{a\sigma} c_a n_\sigma \\ &+ g_{(5S_2)} c^{\dagger a'} n^{\dagger\sigma'} T_{a'\sigma'}^\alpha T_\alpha^{a\sigma} c_a n_\sigma \\ &+ g_{(3S_1^*)} d^{\dagger\delta} n^{\dagger\sigma'} T_{\delta\sigma'}^i T_i^{a\sigma} c_a n_\sigma + \text{C.C.} ,\end{aligned}$$

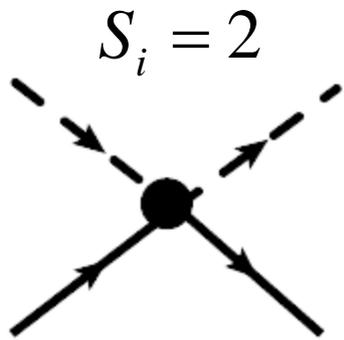
EFT

$$\begin{aligned}
 \mathcal{L}_0 = & n^{\dagger\sigma} \left(i\partial_t + \frac{\nabla^2}{2M_n} \right) n_\sigma + c^{\dagger a} \left(i\partial_t + \frac{\nabla^2}{2M_c} \right) c_a \\
 & + d^{\dagger\delta} \left(i\partial_t + \frac{\nabla^2}{2M_c} \right) d_\delta + \pi^{\dagger\alpha} \left(i\partial_t + \frac{\nabla^2}{2M_{nc}} + \Delta \right) \pi_\alpha \\
 & + \tilde{\pi}^{\dagger i} \left(i\partial_t + \frac{\nabla^2}{2M_{nc}} + \tilde{\Delta} \right) \tilde{\pi}_i ,
 \end{aligned}$$

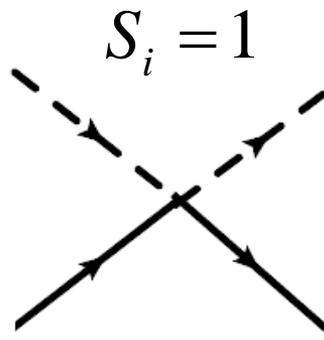
$$\begin{aligned}
 \mathcal{L}_S = & g_{(3S_1)} c^{\dagger a'} n^{\dagger\sigma'} T_{a'\sigma'}^i T_i^{a\sigma} c_a n_\sigma & \mathcal{L}_{P,gs} = & h_{(3P_2)} \pi^{\dagger\alpha} T_\alpha^{ij} T_i^{\sigma a} n_\sigma i (V_n - V_c)_j c_a \\
 & + g_{(5S_2)} c^{\dagger a'} n^{\dagger\sigma'} T_{a'\sigma'}^\alpha T_\alpha^{a\sigma} c_a n_\sigma & & + h_{(5P_2)} \pi^{\dagger\alpha} T_\alpha^{\beta j} T_\beta^{\sigma a} n_\sigma i (V_n - V_c)_j c_a \\
 & + g_{(3S_1^*)} d^{\dagger\delta} n^{\dagger\sigma'} T_{\delta\sigma'}^i T_i^{a\sigma} c_a n_\sigma + \text{C.C.} & & + h_{(3P_2^*)} \pi^{\dagger\alpha} T_\alpha^{jk} T_k^{\delta\sigma} n_\sigma i (V_n - V_{c^*})_j d_\delta + \text{C.C.} ,
 \end{aligned}$$

$$\begin{aligned}
 \mathcal{L}_{P,es} = & \tilde{h}_{(3P_1)} \tilde{\pi}^{\dagger k} T_k^{ij} T_i^{\sigma a} n_\sigma i (V_n - V_c)_j c_a \\
 & + \tilde{h}_{(5P_1)} \tilde{\pi}^{\dagger k} T_k^{\beta j} T_\beta^{\sigma a} n_\sigma i (V_n - V_c)_j c_a \\
 & + \tilde{h}_{(1P_1^*)} \tilde{\pi}^{\dagger k} T_k^{0j} T_0^{\sigma\delta} n_\sigma i (V_n - V_{c^*})_j d_\delta \\
 & + \tilde{h}_{(3P_1^*)} \tilde{\pi}^{\dagger k} T_k^{ij} T_i^{\sigma\delta} n_\sigma i (V_n - V_{c^*})_j d_\delta + \text{C.C.} .
 \end{aligned}$$

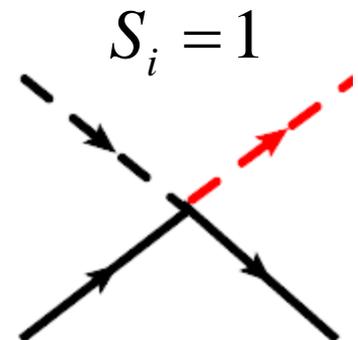
EFT



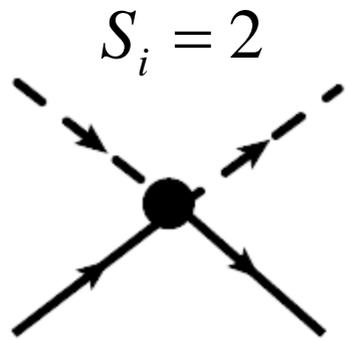
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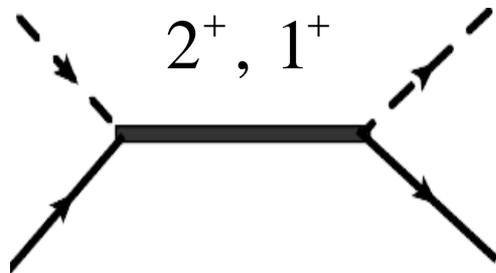
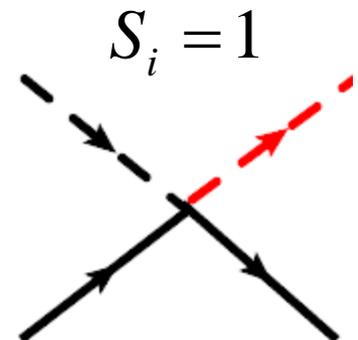
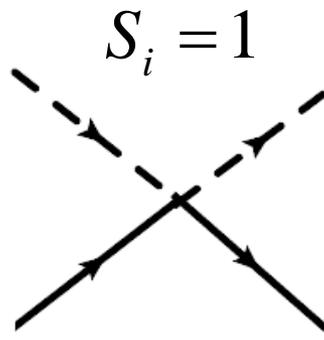
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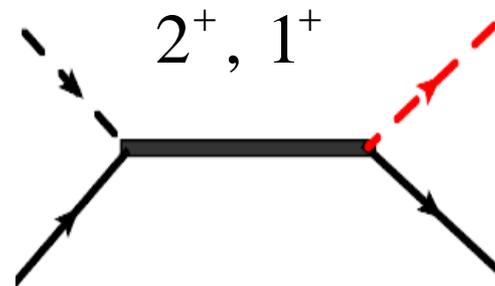
EFT



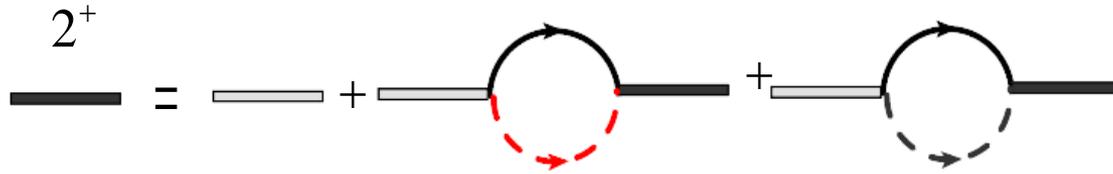
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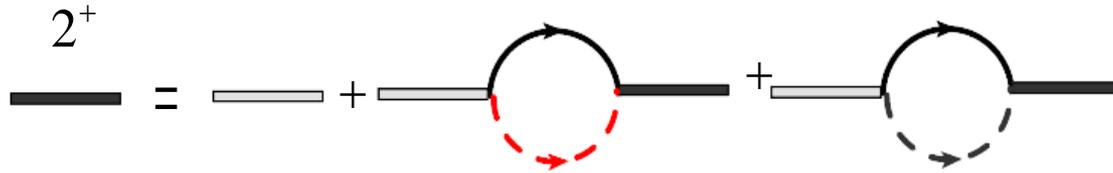
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P-wave

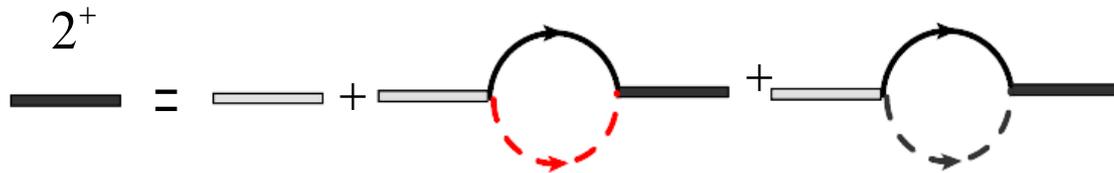


P-wave



$$\frac{C_{({}^3P_2)}^2}{h_{({}^3P_2)}^2 \gamma^2} = \frac{C_{({}^5P_2)}^2}{h_{({}^5P_2)}^2 \gamma^2} = \frac{C_{({}^3P_2^*)}^2}{h_{({}^3P_2^*)}^2 \gamma^{*2}} = \frac{Z}{3\pi}$$

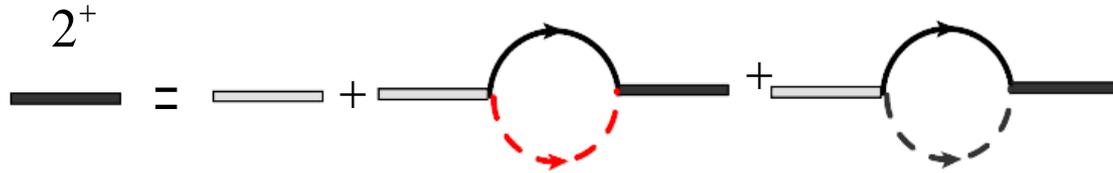
P-wave



$$\frac{C_{({}^3P_2)}^2}{h_{({}^3P_2)}^2 \gamma^2} = \frac{C_{({}^5P_2)}^2}{h_{({}^5P_2)}^2 \gamma^2} = \frac{C_{({}^3P_2^*)}^2}{h_{({}^3P_2^*)}^2 \gamma^{*2}} = \frac{\mathbf{Z}}{3\pi}$$

A function of parameters

P-wave



$$\frac{C_{({}^3P_2)}^2}{h_{({}^3P_2)}^2 \gamma^2} = \frac{C_{({}^5P_2)}^2}{h_{({}^5P_2)}^2 \gamma^2} = \frac{C_{({}^3P_2^*)}^2}{h_{({}^3P_2^*)}^2 \gamma^{*2}} = \frac{Z}{3\pi}$$

4 parameters: 3 h + 1 Delta,
or 3 C + gamma



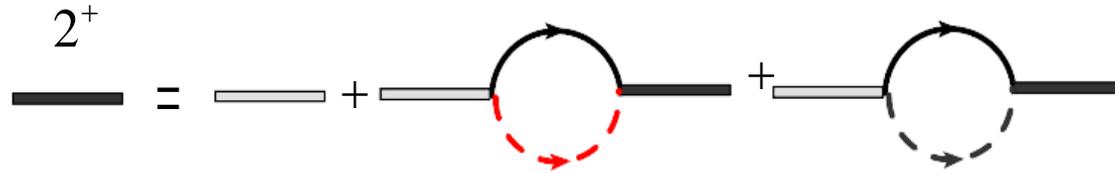
	$C_{({}^3P_2)}$	$C_{({}^5P_2)}$	$C_{({}^3P_2^*)}$
Nollett	-0.283(12)	-0.591(12)	-0.384(6)
	-0.284(23)	-0.593(23)	



K. M. Nollett and R. B. Wiringa, *PRC* 83, 041001 (2011)

L. Trache, et.al., *Phys. Rev. C* 67, 062801(R) (2003)

P-wave



$$\frac{C_{({}^3P_2)}^2}{h_{({}^3P_2)}^2 \gamma^2} = \frac{C_{({}^5P_2)}^2}{h_{({}^5P_2)}^2 \gamma^2} = \frac{C_{({}^3P_2^*)}^2}{h_{({}^3P_2^*)}^2 \gamma^{*2}} = \frac{Z}{3\pi}$$

4 parameters: 3 h + 1 Delta,
or 3 C + gamma



5 parameters



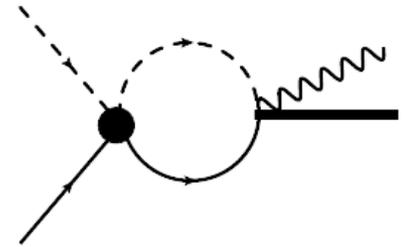
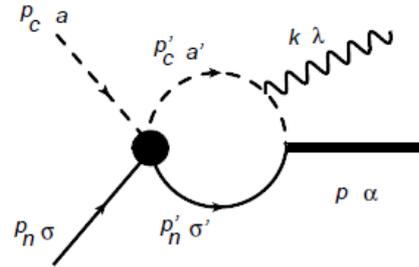
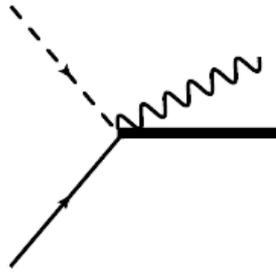
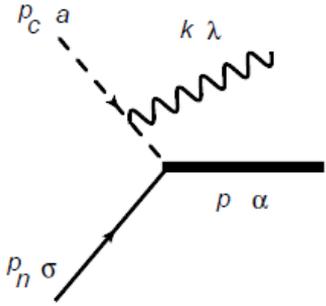
	$C_{({}^3P_2)}$	$C_{({}^5P_2)}$	$C_{({}^3P_2^*)}$	$\tilde{C}_{({}^3P_1)}$	$\tilde{C}_{({}^5P_1)}$	$\tilde{C}_{({}^1P_1^*)}$	$\tilde{C}_{({}^3P_1^*)}$
Nollett	-0.283(12)	-0.591(12)	-0.384(6)	0.220(6)	0.197(5)	-0.195(3)	-0.214(3)
	-0.284(23)	-0.593(23)		0.187(16)	0.217(13)		



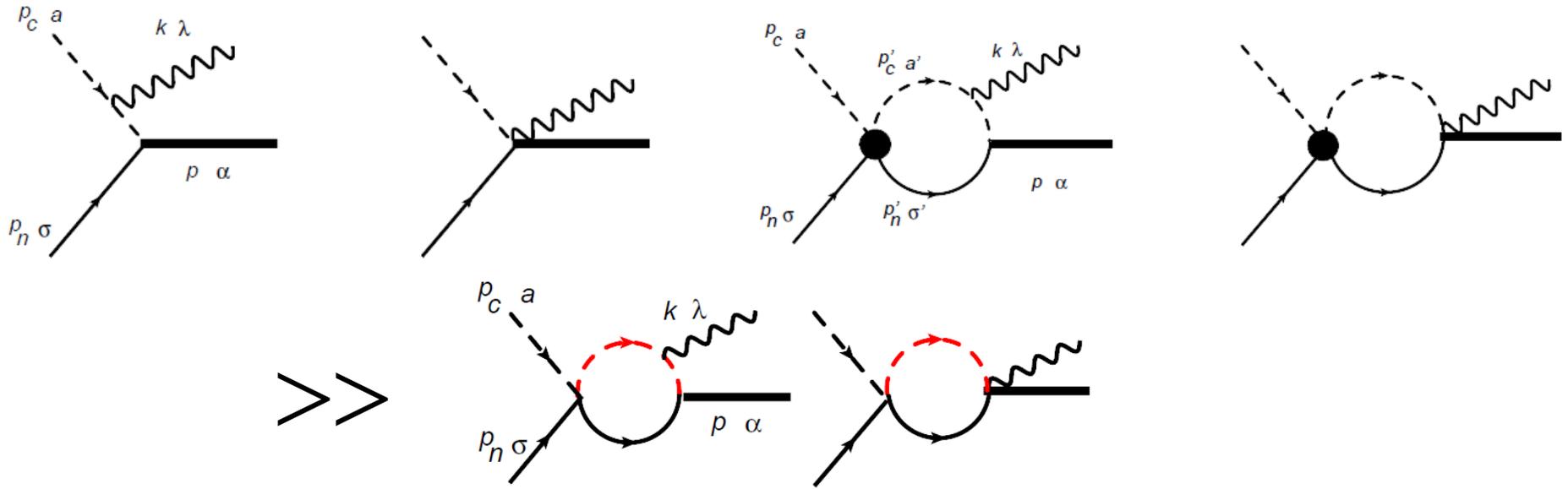
K. M. Nollett and R. B. Wiringa, *PRC* 83, 041001 (2011)

L. Trache, et.al., *Phys. Rev. C* 67, 062801(R) (2003)

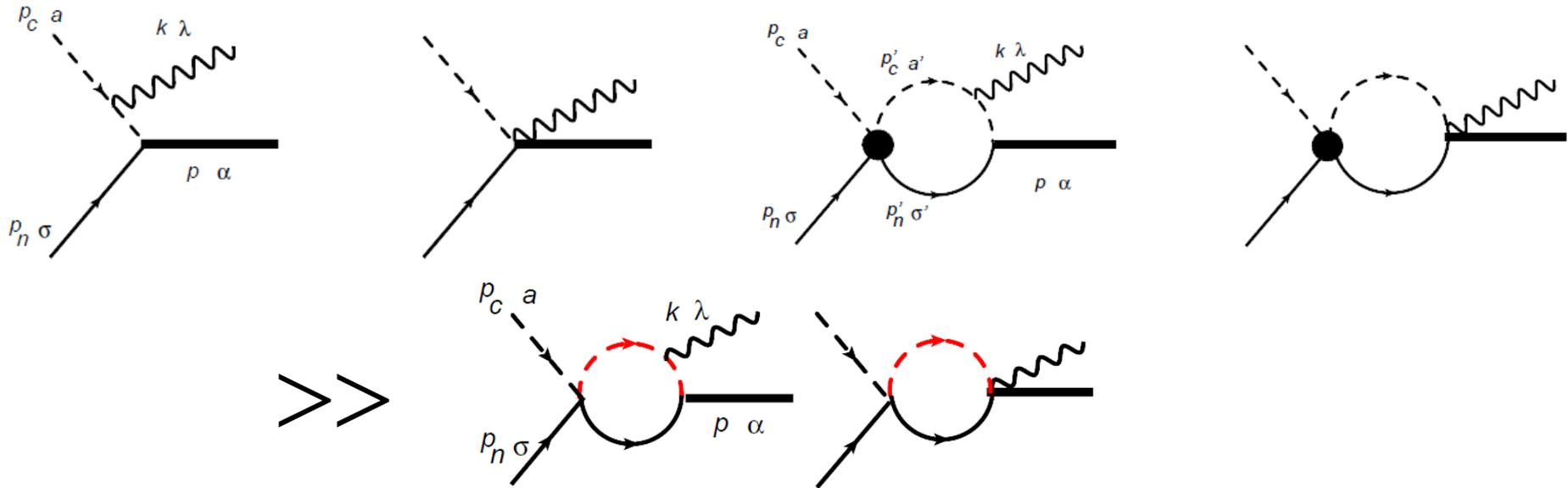
Radiative captures: LO



Radiative captures: LO



Radiative captures: LO



Initial total spin $S_i=2$

$$\mathcal{M} = ie_c h({}^3P_2) \sqrt{8Z^{LO} M_n M_c M_n} \left(T_\beta^{\sigma a} T_\alpha^{\beta j} \right) \left[\frac{\epsilon^*(\lambda) \cdot V_c}{p_c^0 - \omega - \frac{(p_c - k)^2}{2M_c} + i\epsilon} \left(\frac{p_c}{M_R} - \frac{k}{M_c} \right)_j + (1 + X(p_c; \gamma, a({}^3S_2))) \frac{\epsilon^*(\lambda)_j}{M_c} \right]$$

$$X(p_c; \gamma, a) \equiv \frac{(-)i}{a^{-1} + ip_c} \left[p_c - \frac{2}{3} i \frac{\gamma^3 - ip_c^3}{\gamma^2 + p_c^2} \right] \quad a \sim \frac{1}{\gamma} \Rightarrow X \sim 1, \quad a \sim \frac{1}{\Lambda} \Rightarrow X \sim \frac{\gamma}{\Lambda}$$

$$\begin{aligned}
\sum_{\sigma,a}^{\alpha,\lambda} |\mathcal{M}|^2 &= \frac{5}{3} 64\pi\alpha Z_c^2 \frac{3\pi}{\gamma^2} \frac{M_n^2}{M_R} \left(C_{({}^5P_2)}^{\text{LO}} \right)^2 \left[|1 + X(p_c; \gamma, a({}^5S_2))|^2 - \frac{2p_c^2 \sin^2 \theta}{p_c^2 + \gamma^2} \left(\frac{\gamma^2}{p_c^2 + \gamma^2} + \text{Re} \{ X(p_c; \gamma, a({}^5S_2)) \} \right) \right] \\
&+ \frac{5}{3} 64\pi\alpha Z_c^2 \frac{3\pi}{\gamma^2} \frac{M_n^2}{M_R} \left(C_{({}^3P_2)}^{\text{LO}} \right)^2 \left[1 - \frac{p_c^2 \sin^2 \theta}{p_c^2 + \gamma^2} \frac{2\gamma^2}{p_c^2 + \gamma^2} \right]
\end{aligned}$$

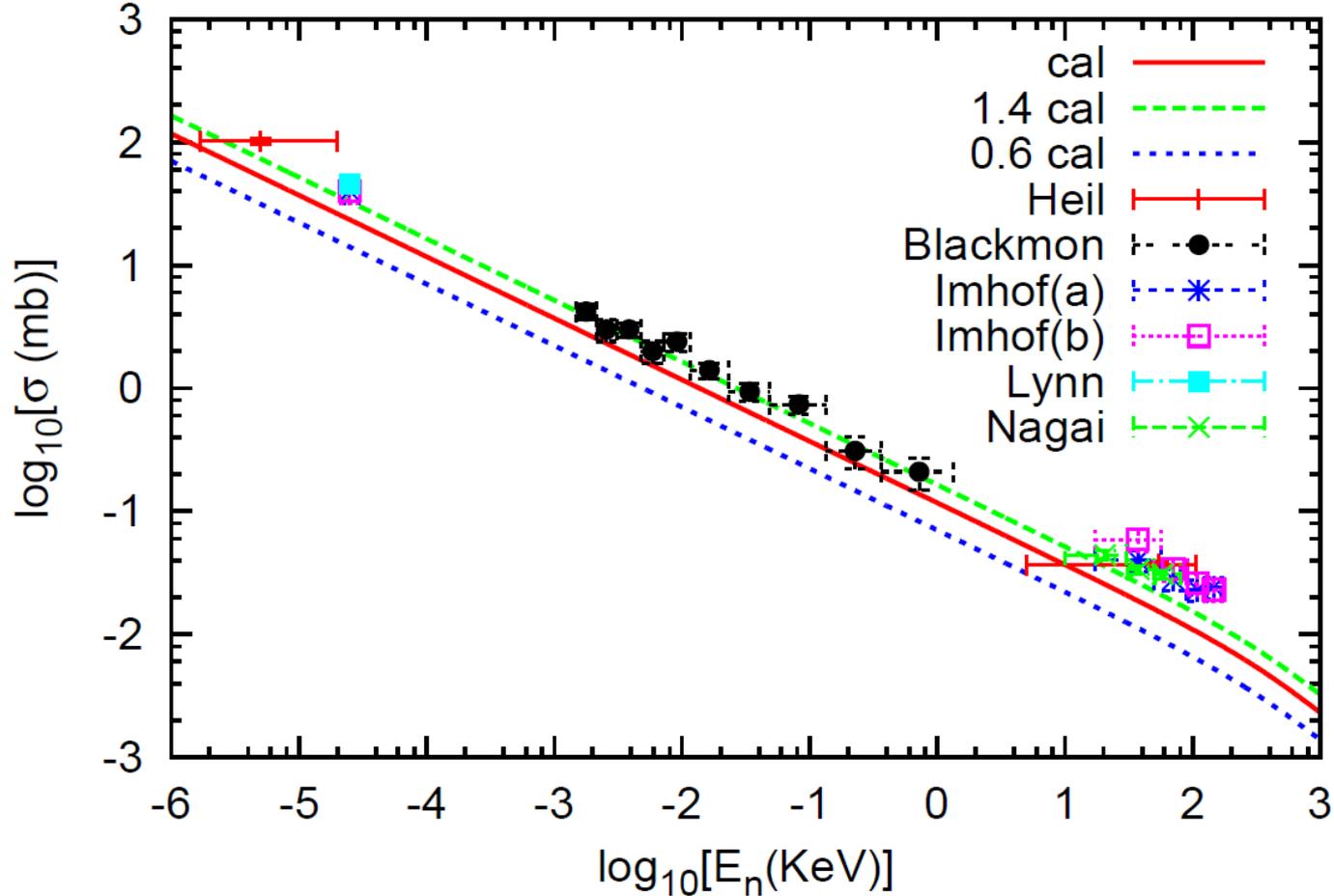
$$X(p_c; \gamma, a) \equiv \frac{(-)i}{a^{-1} + ip_c} \left[p_c - \frac{2}{3} i \frac{\gamma^3 - ip_c^3}{\gamma^2 + p_c^2} \right]$$

$$\sum_{\sigma,a}^{\alpha,\lambda} |\mathcal{M}|^2 = \frac{5}{3} 64\pi\alpha Z_c^2 \frac{3\pi}{\gamma^2} \frac{M_n^2}{M_R} \left(C_{({}^5P_2)}^{\text{LO}} \right)^2 \left[|1 + X(p_c; \gamma, a({}^5S_2))|^2 - \frac{2p_c^2 \sin^2 \theta}{p_c^2 + \gamma^2} \left(\frac{\gamma^2}{p_c^2 + \gamma^2} + \text{Re} \{ X(p_c; \gamma, a({}^5S_2)) \} \right) \right] \\ + \frac{5}{3} 64\pi\alpha Z_c^2 \frac{3\pi}{\gamma^2} \frac{M_n^2}{M_R} \left(C_{({}^3P_2)}^{\text{LO}} \right)^2 \left[1 - \frac{p_c^2 \sin^2 \theta}{p_c^2 + \gamma^2} \frac{2\gamma^2}{p_c^2 + \gamma^2} \right]$$

$$X(p_c; \gamma, a) \equiv \frac{(-)i}{a^{-1} + ip_c} \left[p_c - \frac{2}{3} i \frac{\gamma^3 - ip_c^3}{\gamma^2 + p_c^2} \right]$$

$$\sum_{i,f} |\mathcal{M}|^2 = 64\pi\alpha Z_c^2 \frac{3\pi}{\tilde{\gamma}^2} \frac{M_n^2}{M_R} \left\{ \left(\tilde{C}_{({}^3P_1)}^{\text{LO}} \right)^2 \left[1 - \frac{p_c^2 \sin^2 \theta}{p_c^2 + \tilde{\gamma}^2} \left(\frac{2\tilde{\gamma}^2}{p_c^2 + \tilde{\gamma}^2} \right) \right] \right. \\ \left. + \left(\tilde{C}_{({}^5P_1)}^{\text{LO}} \right)^2 \left[|1 + X(p_c; \tilde{\gamma}, a({}^5S_2))|^2 - \frac{2p_c^2 \sin^2 \theta}{p_c^2 + \tilde{\gamma}^2} \left(\frac{\tilde{\gamma}^2}{p_c^2 + \tilde{\gamma}^2} + \text{Re} \{ X(p_c; \tilde{\gamma}, a({}^5S_2)) \} \right) \right] \right\}$$

LO results on $\text{Li7}(n,\text{gamma})\text{Li8}(\text{Li8}^*)$



N. K. Timofeyuk *et.al.*, *PRL* 91, 232501 (2003); D. Howell *et.al.*, *PRC* 88, 025804 (2013);
D. Gul'ko *et.al.*, *SJNP* 6, 477 (1968); E. Lynn *et.al.*, *PRC* 44, 764 (1991);
Y. Nagai *et. al.*, *PRC* 71, 055803 (2005); J. C. Blackmon *et. al.*, *PRC* 54, 383 (1996); J. E. Lynn *et. al.*, *PRC*
44, 764 (1991); M. Heil *et.al.*, *Astro. J.* 507, 997 (1998); W. L. Imhof *et.al.*, *PR* 114, 1037 (1959).

LO results on $\text{Li7}(n,\gamma)\text{Li8}(\text{Li8}^*)$

$$\frac{\sigma[(S_i = 1) \rightarrow 2^+]}{\sigma[(S_i = 2) \rightarrow 2^+]} = \frac{\left(C_{(3P_2)}^{\text{LO}}\right)^2}{\left(C_{(5P_2)}^{\text{LO}}\right)^2 \left(1 - \frac{2}{3}\gamma a_{(5S_2)}\right)^2}$$



$$\frac{\sigma[(S_i = 2) \rightarrow 2^+]}{\sigma(\rightarrow 2^+)} = 0.93(2) \quad [> 0.86]$$

A. D. Gul'ko, S. S. Trostin, and A. Hudoklin, *Sov. J. Nucl. Phys.* 6, 477 (1968);
J. E. Lynn, E. T. Jurney, and S. Raman, *Phys. Rev. C* 44, 764 (1991);
Y. Nagai et. al., *Phys. Rev. C* 71, 055803 (2005).

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 $\frac{\sigma[(S_i = 2) \rightarrow 2^+]}{\sigma(\rightarrow 2^+)} = 0.93(2) \quad [> 0.86]$

$$\frac{\sigma[(S_i = 2) \rightarrow 1^+]}{\sigma(\rightarrow 1^+)} = 0.65(6) \quad \text{or} \quad 0.75(7).$$

A. D. Gul'ko, S. S. Trostin, and A. Hudoklin, *Sov. J. Nucl. Phys.* 6, 477 (1968);
J. E. Lynn, E. T. Journey, and S. Raman, *Phys. Rev. C* 44, 764 (1991);
Y. Nagai et. al., *Phys. Rev. C* 71, 055803 (2005).

LO results on Li7(n,gamma)Li8(Li8*)

$$\frac{\sigma[(S_i = 1) \rightarrow 2^+]}{\sigma[(S_i = 2) \rightarrow 2^+]} = \frac{\left(C_{(3P_2)}^{\text{LO}}\right)^2}{\left(C_{(5P_2)}^{\text{LO}}\right)^2 \left(1 - \frac{2}{3}\gamma a_{(5S_2)}\right)^2}$$

→ $\frac{\sigma[(S_i = 2) \rightarrow 2^+]}{\sigma(\rightarrow 2^+)} = 0.93(2) \quad [> 0.86]$

$$\frac{\sigma[(S_i = 2) \rightarrow 1^+]}{\sigma(\rightarrow 1^+)} = 0.65(6) \quad \text{or} \quad 0.75(7),$$

$$\frac{\sigma(\rightarrow 1^+)}{\sigma(\rightarrow 2^+)} = \frac{3 \left(\tilde{C}_{(3P_1)}^{\text{LO}}\right)^2 + \left(\tilde{C}_{(5P_1)}^{\text{LO}}\right)^2 \left|1 - \frac{2}{3}a_{(5S_2)}\tilde{\gamma}\right|^2}{5 \left(C_{(3P_2)}^{\text{LO}}\right)^2 + \left(C_{(5P_2)}^{\text{LO}}\right)^2 \left|1 - \frac{2}{3}a_{(5S_2)}\gamma\right|^2}$$

→ $\frac{\sigma(\rightarrow 2^+)}{\sigma} = 0.88(4) \quad [0.89(1)]$

A. D. Gul'ko, S. S. Trostin, and A. Hudoklin, *Sov. J. Nucl. Phys.* 6, 477 (1968);
 J. E. Lynn, E. T. Jurney, and S. Raman, *Phys. Rev. C* 44, 764 (1991);
 Y. Nagai et. al., *Phys. Rev. C* 71, 055803 (2005).



- It is considered as isospin mirror of Li^7 capture on the nucleon level
- From EFT/core+proton picture, they are quite different due to strong Coulomb effect



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- From EFT/core+proton picture, they are quite different due to strong Coulomb effect

*E. Ryberg, C. Forssén, H.-W. Hammer and L. Platter,
PRC 89, 014325 (2014)*

Nonperturbative Coulomb effect

$$k_C \equiv Q_c Q_n \alpha_{EM} M_R \quad \eta \equiv \frac{k_C}{k} \sim 1 \quad \textit{Sommerfeld para.}$$

Nonperturbative Coulomb effect

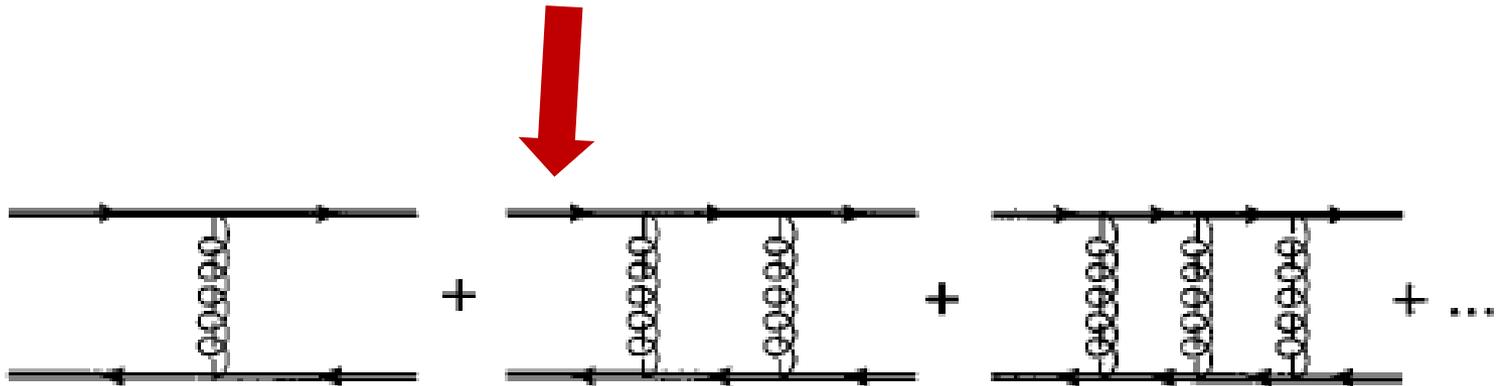
$$k_C \equiv Q_c Q_n \alpha_{EM} M_R \quad \eta \equiv \frac{k_C}{k} \sim 1 \quad \text{Sommerfeld para.}$$

$$\frac{1}{E - H_0 - V_c - V_s} = \frac{1}{E - H_0 - V_c} + \frac{1}{E - H_0 - V_c} V_s \frac{1}{E - H_0 - V_c} + \dots$$

Nonperturbative Coulomb effect

$$k_C \equiv Q_c Q_n \alpha_{EM} M_R \quad \eta \equiv \frac{k_C}{k} \sim 1 \quad \text{Sommerfeld para.}$$

$$\frac{1}{E - H_0 - V_c - V_s} = \frac{1}{E - H_0 - V_c} + \frac{1}{E - H_0 - V_c} V_s \frac{1}{E - H_0 - V_c} + \dots$$



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$$\chi_{\mathbf{k}}^{(\pm)}(\mathbf{r}) = e^{-\frac{\pi}{2}\eta} e^{i\mathbf{k}\cdot\mathbf{r}} \Gamma(1 \pm i\eta) M(\mp i\eta, 1; \pm ikr - ikr)$$

Kummer function

$$\chi_{\mathbf{k}}^{(\pm)*}(r=0) \chi_{\mathbf{k}}^{(\pm)}(r=0) = \frac{2\pi\eta}{e^{2\pi\eta} - 1} = C_{\eta,0}^2$$

$$\chi_{\mathbf{k}}^{(\mp)*}(r=0) \chi_{\mathbf{k}}^{(\pm)}(r=0) = C_{\eta,0}^2 e^{\pm 2i\sigma_0} \quad \text{Coulomb barrier, and phase}$$

$$C_{\eta,l} = \frac{2^l e^{-\frac{\pi}{2}\eta} |\Gamma(l+1+i\eta)|}{\Gamma(2l+2)} \quad e^{2i\sigma_l} \equiv \frac{\Gamma(l+1+i\eta)}{\Gamma(l+1-i\eta)}$$

Effective range expansion in EFT

$$\begin{aligned}\langle \chi_{\mathbf{p}'}^{(-)} | T_{cs}(E) | \chi_{\mathbf{p}}^{(+)} \rangle &= (-) \frac{2\pi}{M_{\text{R}}} \frac{\chi_{\mathbf{p}'}^{(-)*}(0) \chi_{\mathbf{p}}^{(+)}(0)}{-a_0^{-1} - 2k_C H(\eta)} \\ &\rightarrow (-) \frac{2\pi}{M_{\text{R}}} \frac{C_{\eta,0}^2 e^{2i\sigma_0}}{-a_0^{-1} - 2k_C H(\eta)}\end{aligned}$$

$$C_{\eta,0}^2 k(\cot \delta_0 - i) = -\frac{1}{a_0} + \dots - 2k_C H(\eta) \quad H(\eta) = \psi(i\eta) + 1/(2i\eta) - \ln(i\eta)$$

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X. Kong and F. Ravndal, NPA 665, 137 (2000).

R. Higa, H. -W. Hammer and U. van Kolck, NPA 809, 171 (2008).

Effective range expansion in EFT

*X. Kong and F. Ravndal,
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$$\begin{aligned}\langle \chi_{\mathbf{p}'}^{(-)} | T_{cs}(E) | \chi_{\mathbf{p}}^{(+)} \rangle &= (-) \frac{2\pi}{M_R} \frac{\chi_{\mathbf{p}'}^{(-)*}(0) \chi_{\mathbf{p}}^{(+)}(0)}{-a_0^{-1} - 2k_C H(\eta)} \\ &\rightarrow (-) \frac{2\pi}{M_R} \frac{C_{\eta,0}^2 e^{2i\sigma_0}}{-a_0^{-1} - 2k_C H(\eta)}\end{aligned}$$

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$$\begin{aligned}\langle \chi_{\mathbf{p}'}^{(-)} | T_{cs}(E) | \chi_{\mathbf{p}}^{(+)} \rangle &= (-) \frac{6\pi}{M_R} \frac{\partial \chi_{\mathbf{p}'}^{(-)*}(0) \partial \chi_{\mathbf{p}}^{(+)}(0)}{-\frac{1}{a_1} + \frac{r_1}{2} k^2 - k^2(1 + \eta^2) 2k_C H(\eta)} \\ &\rightarrow (-) \frac{6\pi}{M_R} \frac{k^2 C_{\eta,1}^2 e^{2i\sigma_1}}{-\frac{1}{a_1} + \frac{r_1}{2} k^2 - k^2(1 + \eta^2) 2k_C H(\eta)}\end{aligned}$$

$$C_{\eta,1}^2 k^3 (\cot \delta_1 - i) = -\frac{1}{a_1} + \frac{r_1}{2} k^2 + \dots - k^2(1 + \eta^2) 2k_C H(\eta)$$

Effective range expansion in EFT

$$\begin{aligned} \langle \chi_{\mathbf{p}'}^{(-)} | T_{cs}(E) | \chi_{\mathbf{p}}^{(+)} \rangle &= (-) \frac{2\pi}{M_R} \frac{\chi_{\mathbf{p}'}^{(-)*}(0) \chi_{\mathbf{p}}^{(+)}(0)}{-a_0^{-1} - 2k_C H(\eta)} \\ &\rightarrow (-) \frac{2\pi}{M_R} \frac{C_{\eta,0}^2 e^{2i\sigma_0}}{-a_0^{-1} - 2k_C H(\eta)} \end{aligned}$$

*One parameter:
g (or a0)*

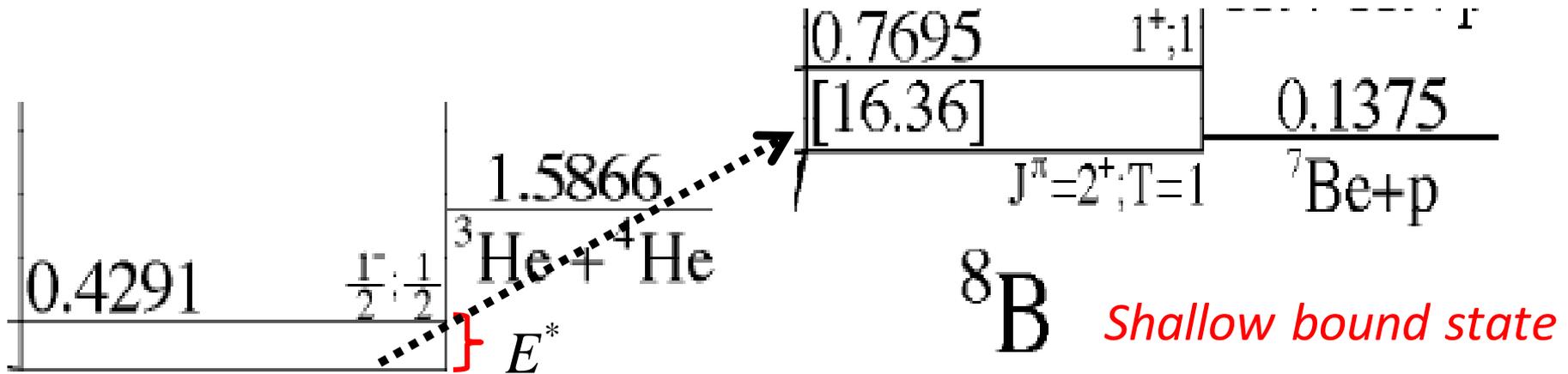
$$C_{\eta,0}^2 k (\cot \delta_0 - i) = -\frac{1}{a_0} + \dots - 2k_C H(\eta) \quad H(\eta) = \psi(i\eta) + 1/(2i\eta) - \ln(i\eta)$$

$$\begin{aligned} \langle \chi_{\mathbf{p}'}^{(-)} | T_{cs}(E) | \chi_{\mathbf{p}}^{(+)} \rangle &= (-) \frac{6\pi}{M_R} \frac{\partial \chi_{\mathbf{p}'}^{(-)*}(0) \partial \chi_{\mathbf{p}}^{(+)}(0)}{-\frac{1}{a_1} + \frac{r_1}{2} k^2 - k^2(1 + \eta^2) 2k_C H(\eta)} \\ &\rightarrow (-) \frac{6\pi}{M_R} \frac{k^2 C_{\eta,1}^2 e^{2i\sigma_1}}{-\frac{1}{a_1} + \frac{r_1}{2} k^2 - k^2(1 + \eta^2) 2k_C H(\eta)} \end{aligned}$$

$$C_{\eta,1}^2 k^3 (\cot \delta_1 - i) = -\frac{1}{a_1} + \frac{r_1}{2} k^2 + \dots - k^2(1 + \eta^2) 2k_C H(\eta)$$

Two parameters: Delta and h (or a1 and r1)

Scales, spins, core excitations



${}^8\text{B}$ *Shallow bound state*

$[-0.24]$ $J^\pi = \frac{3^-}{2}; T = \frac{1}{2}$

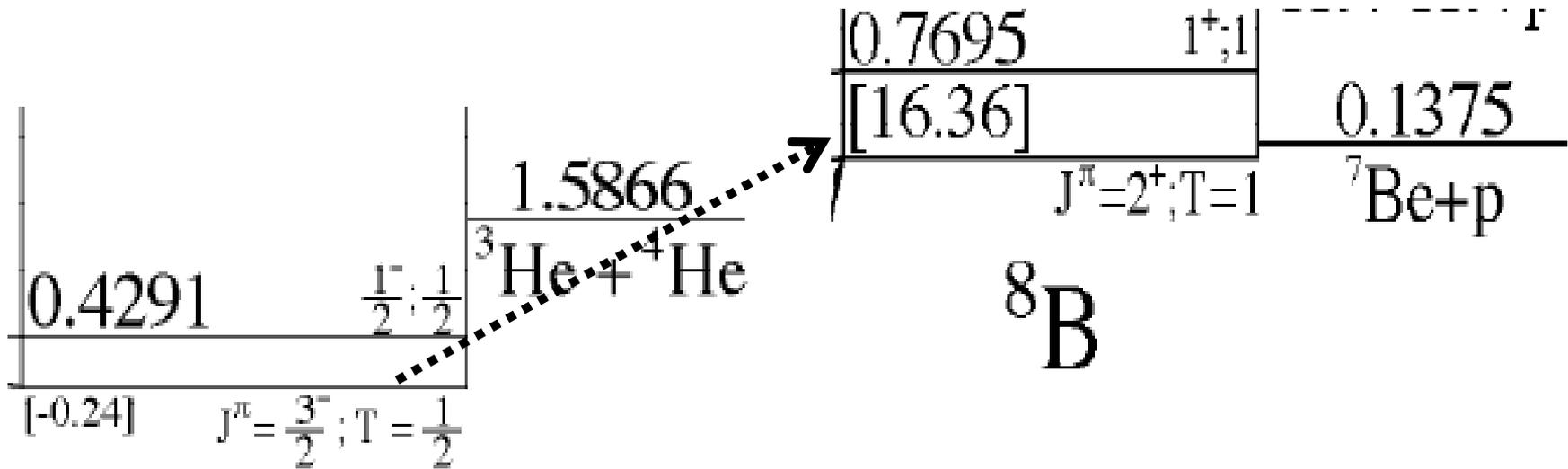
${}^7\text{Be}$

$$\frac{p_c}{\Lambda}, \frac{\gamma}{\Lambda}, \frac{\tilde{\gamma}}{\Lambda} \sim 0.2$$

$$\eta = \frac{k_c}{k} \sim 1$$

Momentum scale	Definition	Value
$k_C \sim \gamma$	$Q_c Q_n \alpha_{EM} M_R$	24.02 MeV
γ	$\sqrt{2M_R B_{8B}}$	15.04 MeV
Λ	$\sqrt{2M'_R B_{7Be}}$	70 MeV
$\gamma^* \sim \gamma$	$\sqrt{2M_R (B_{8B} + E^*)}$	30.53 MeV
$\gamma_\Delta \sim \gamma$	$\sqrt{2M_R E^*}$	26.57 MeV
$a_{3S_1}, a_{5S_2} \sim 1/\gamma$	scattering lengths	Varies
$r_0 \sim 1/\Lambda$	$l = 0$ effective ranges	Varies
$a_1 \sim \gamma^{-2} \Lambda^{-1}$	scattering volume	1054.1 fm ³
$r_1 \sim \Lambda$	$l = 1$ effective "range"	-0.34 fm ⁻¹

Repeat

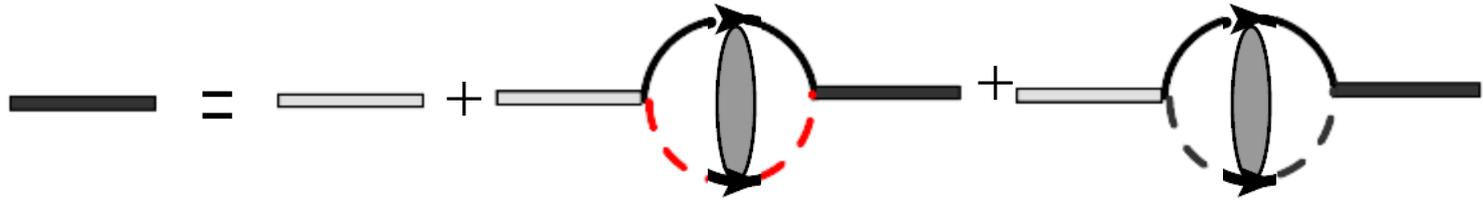


${}^7\text{Be}$

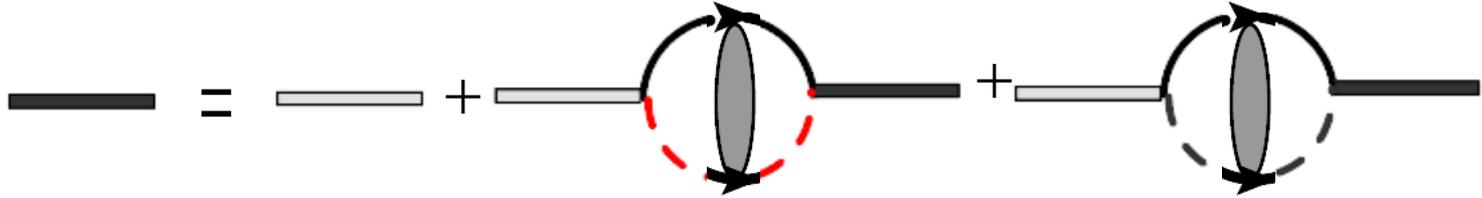
IS $\text{Be}7 + \text{p}$: ${}^3S_1, {}^5S_2, D$
 IS $\text{Be}7^* + \text{p}$: ${}^1S_0^*, {}^3S_1^*$

FS(2^+) $\text{Be}7 + \text{p}$: ${}^3P_2, {}^5P_2$
 FS(2^+) $\text{Be}7^* + \text{p}$: ${}^3P_2^*$

P-wave



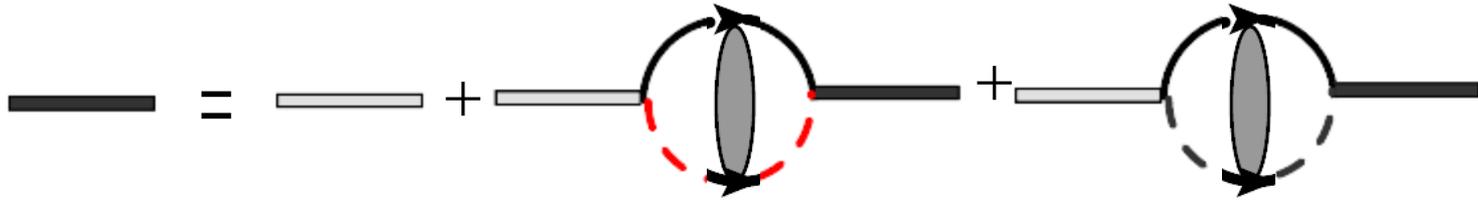
P-wave



$$\frac{C_Y^2}{h_Y^2 \gamma^2 \Gamma^2 (2 + k_C / \gamma)} = \frac{C_{({}^3P_2^*)}^2}{h_{({}^3P_2^*)}^2 \gamma^{*2} \Gamma^2 (2 + k_C / \gamma^*)} = \frac{Z}{3\pi}$$

$$Y = {}^3P_2 \text{ and } {}^5P_2$$

P-wave

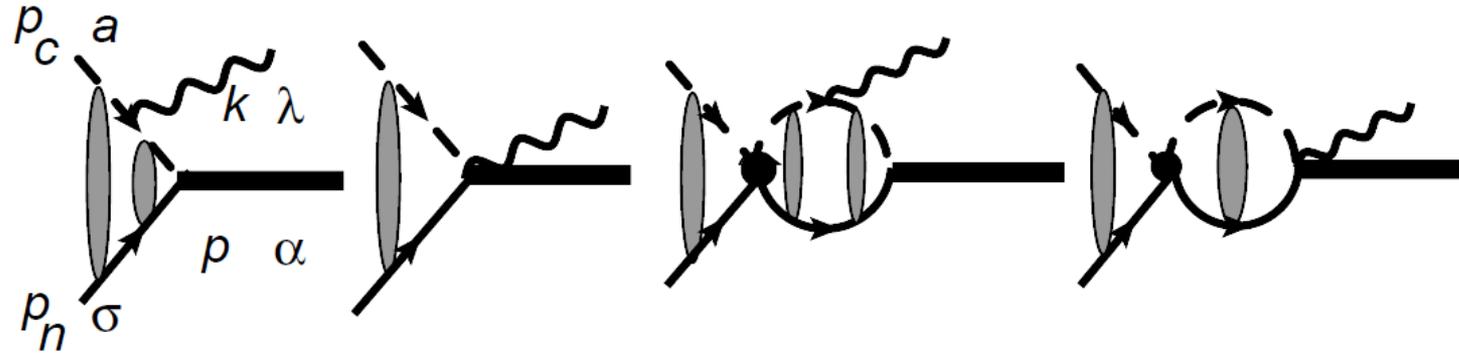


$$\frac{C_Y^2}{h_Y^2 \gamma^2 \Gamma^2 (2 + k_C / \gamma)} = \frac{C_{({}^3P_2^*)}^2}{h_{({}^3P_2^*)}^2 \gamma^{*2} \Gamma^2 (2 + k_C / \gamma^*)} = \frac{Z}{3\pi}$$

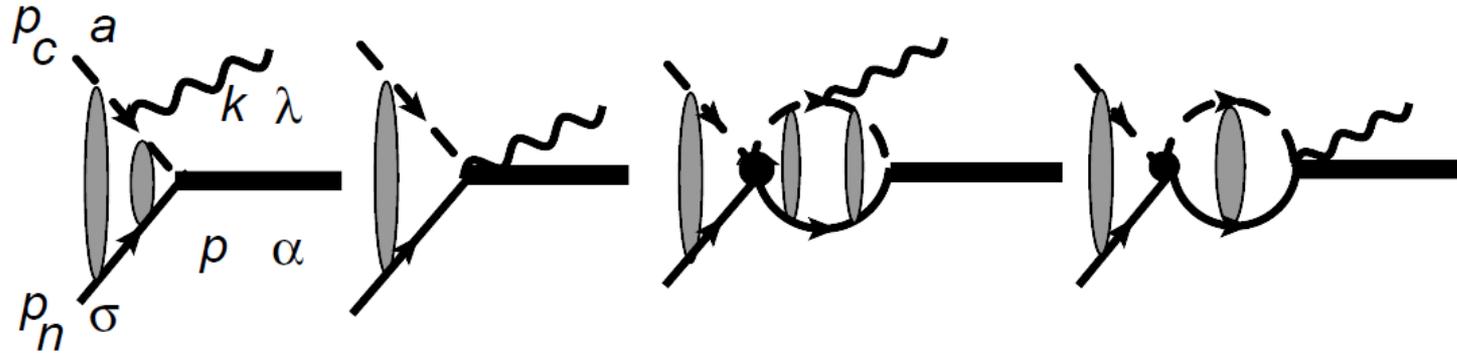
$$Y = {}^3P_2 \text{ and } {}^5P_2$$

4 parameters: 3 h + 1 Delta, or 3 C + gamma

Radiative captures: LO



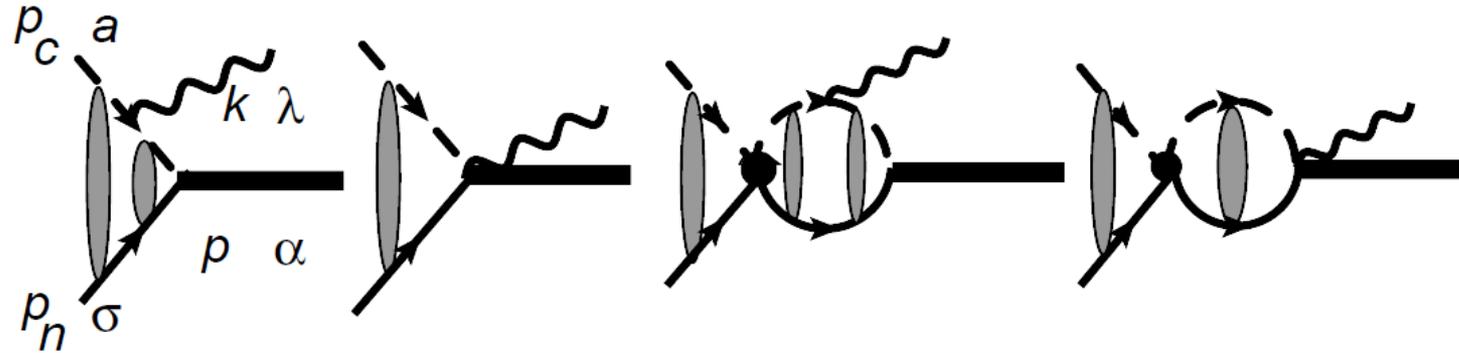
Radiative captures: LO



$$S(E) = \frac{e^{2\pi\eta}}{e^{2\pi\eta} - 1} \frac{Z_{eff}^2}{M_R^2} \frac{\pi}{24} \omega k_C (\gamma^2 + k^2)^2 \frac{5}{3} \times$$

$$\left[C_{({}^3P_2)}^{LO}{}^2 (|S({}^3S_1)|^2 + 2|D|^2) + C_{({}^5P_2)}^{LO}{}^2 (|S({}^5S_2)|^2 + 2|D|^2) \right]$$

Radiative captures: LO



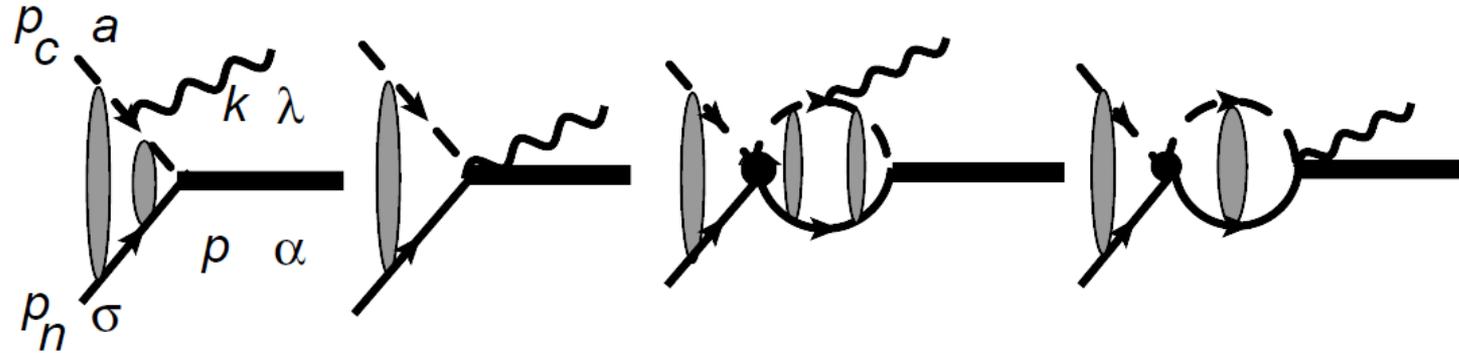
$$S(E) = \frac{e^{2\pi\eta}}{e^{2\pi\eta} - 1} \frac{Z_{eff}^2}{M_R^2} \frac{\pi}{24} \omega k_C (\gamma^2 + k^2)^2 \frac{5}{3} \times$$

$$\left[C_{({}^3P_2)}^{LO}{}^2 (|S({}^3S_1)|^2 + 2|D|^2) + C_{({}^5P_2)}^{LO}{}^2 (|S({}^5S_2)|^2 + 2|D|^2) \right]$$

$$S(X) \equiv \int_0^{+\infty} dr W_{-\eta_B, \frac{3}{2}}(2\gamma r) r \left[\frac{C_{\eta,0} G_0(k, r)}{-a_{(X)}^{-1} - 2k_C H(\eta)} + \frac{F_0(k, r)}{C_{\eta,0} k} \frac{-a_{(X)}^{-1} - 2k_C \text{Re}[H(\eta)]}{-a_{(X)}^{-1} - 2k_C H(\eta)} \right]$$

$$D \equiv \int_0^{+\infty} dr W_{-\eta_B, \frac{3}{2}}(2\gamma r) r \frac{F_2(k, r)}{C_{\eta,0} k}$$

Radiative captures: LO



$$S(E) = \frac{e^{2\pi\eta}}{e^{2\pi\eta} - 1} \frac{Z_{eff}^2}{M_R^2} \frac{\pi}{24} \omega k_C (\gamma^2 + k^2)^2 \frac{5}{3} \times$$

$$\left[C_{(3P_2)}^{LO}{}^2 (|S(^3S_1)|^2 + 2|D|^2) + C_{(5P_2)}^{LO}{}^2 (|S(^5S_2)|^2 + 2|D|^2) \right]$$

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*Coulomb
wavefunc.*

$F \rightarrow j$
 $G \rightarrow n$
 $W \rightarrow h$

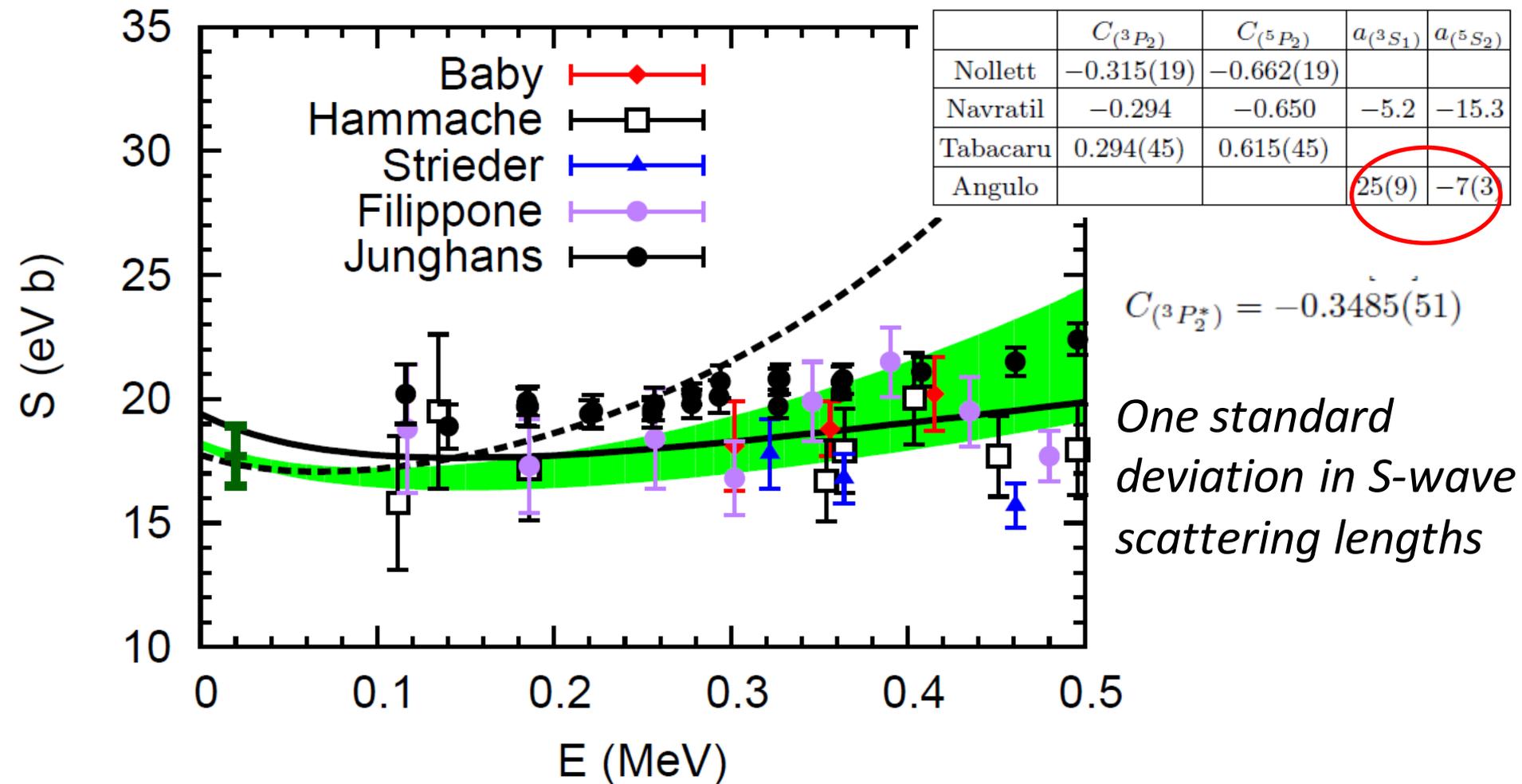
LO results on Be7(p,gamma)B8

	$C_{(3P_2)}$	$C_{(5P_2)}$	$a_{(3S_1)}$	$a_{(5S_2)}$
Nollett	-0.315(19)	-0.662(19)		
Navratil	-0.294	-0.650	-5.2	-15.3
Tabacaru	0.294(45)	0.615(45)		
Angulo			25(9)	-7(3)

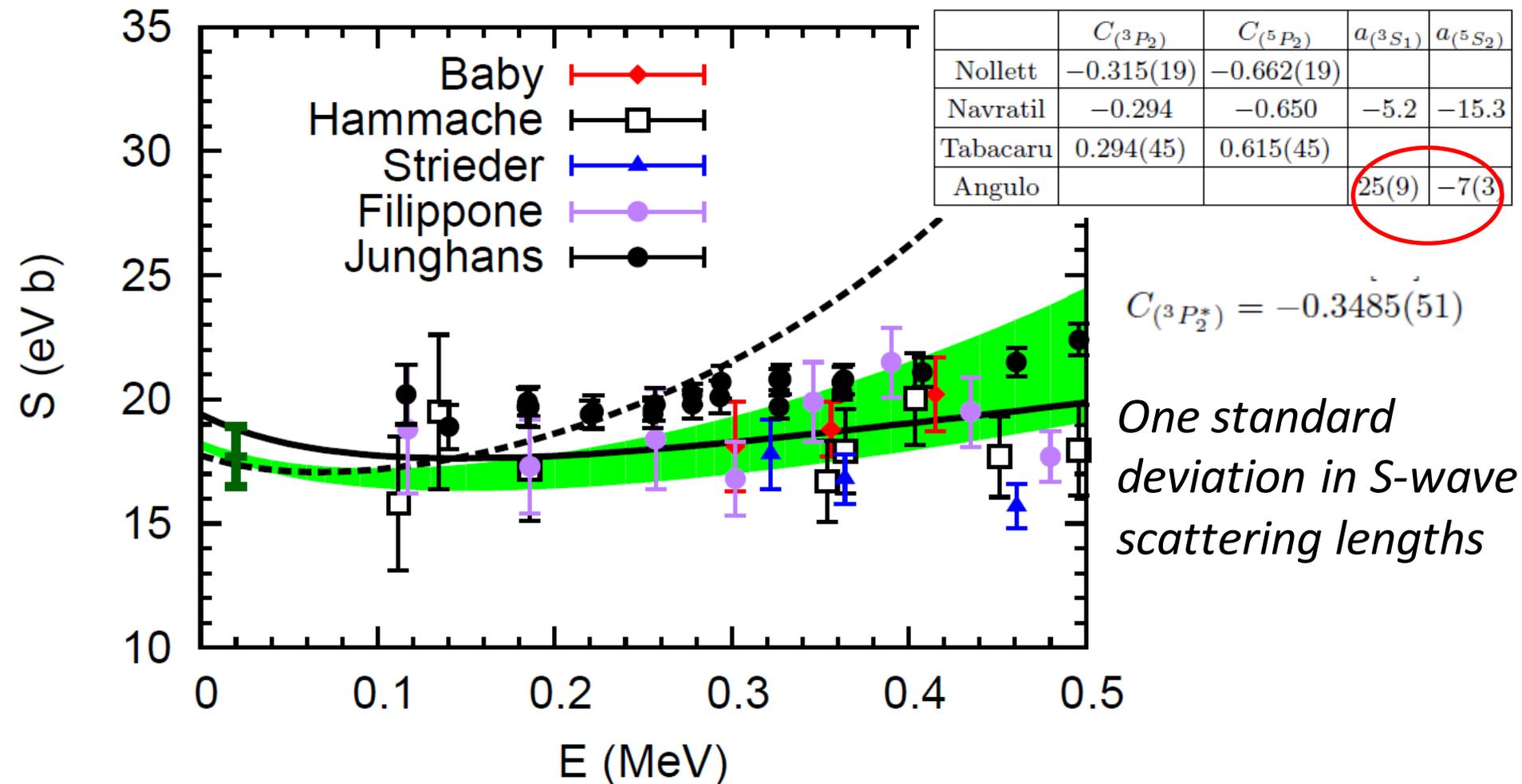
$$C_{(3P_2^*)} = -0.3485(51)$$

P. Navratil, R. Roth and S. Quaglioni, *Phys. Lett. B* 704, 379 (2011);
C. Angulo *et. al.*, *Nucl. Phys. A* 716, 211 (2003);
G. Tabacaru, *et. al.*, *Phys. Rev. C* 73, 025808 (2006)

LO results on Be7(p,gamma)B8

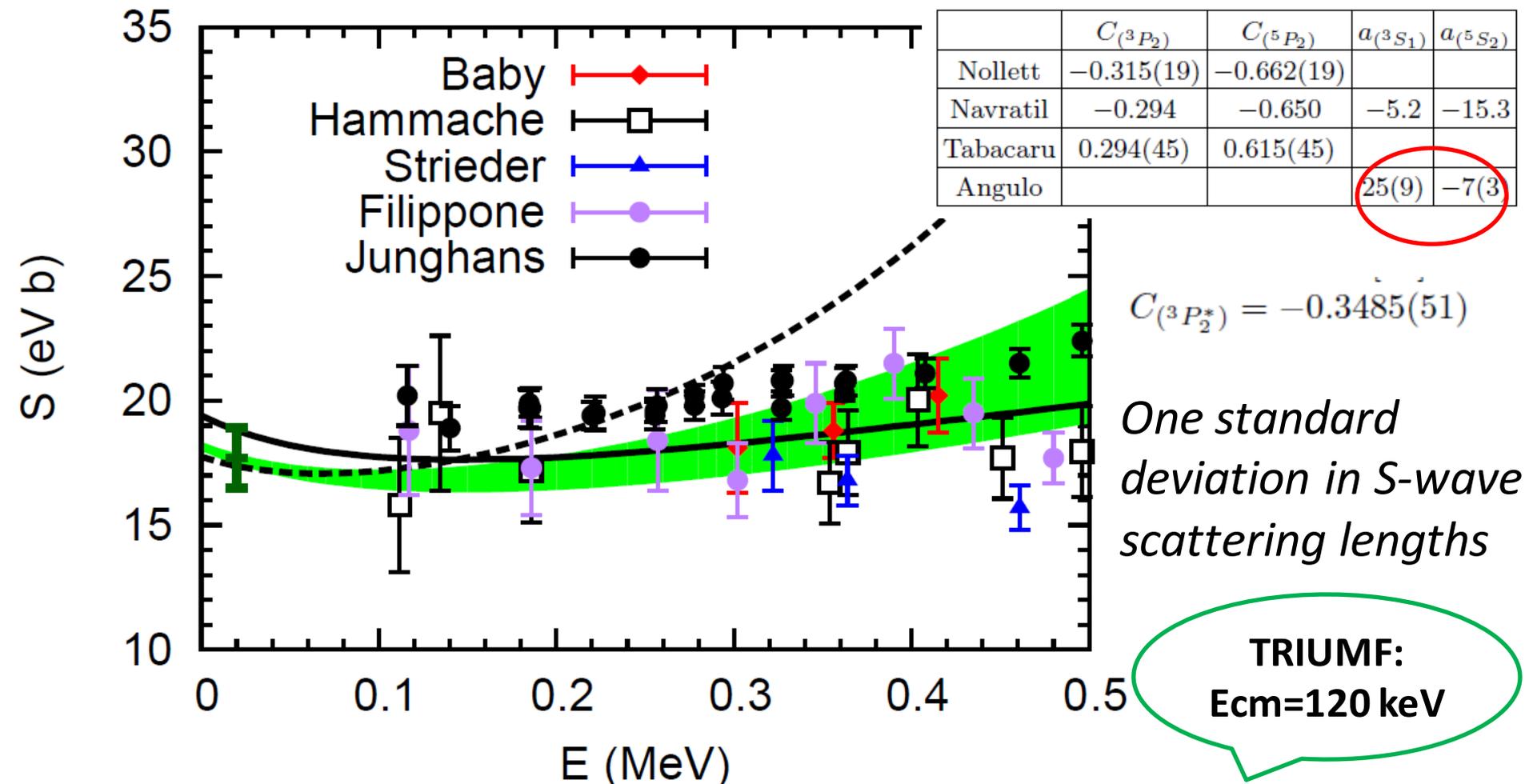


LO results on Be7(p,gamma)B8



Need better measurement of S-wave scattering lengths and/or effective ranges to extrapolate data to zero energy

LO results on Be7(p,gamma)B8



Need better measurement of S-wave scattering lengths and/or effective ranges to extrapolate data to zero energy

LO results on Be7(p,gamma)B8

$$S(E) = S(0)(1 + d_1 E + d_2 E^2) \quad \text{Fit to } 0 < E < 50 \text{ keV}$$

- L. T. Baby, *et. al.*, [ISOLDE Collaboration], *Phys. Rev.Lett.* 90, 022501 (2003);
F. Hammache, *et. al.*, *Phys. Rev. Lett.* 86, 3985 (2001);
F. Strieder, *et. al.*, *Nucl. Phys. A* 696, 219 (2001);
B. W. Filippone, *et. al.*, *Phys. Rev. C* 28, 2222 (1983);
A. R. Junghans, *et. al.*, *Phys. Rev. C* 68, 065803 (2003);
A. R. Junghans, *et. al.*, *Phys. Rev. C* 81, 012801 (2010).

LO results on Be7(p,gamma)B8

$$S(E) = S(0)(1 + d_1 E + d_2 E^2) \quad \text{Fit to } 0 < E < 50 \text{ keV}$$

	$S(0)$ (eV b)	$S_{(^3S_1)}(0)$	d_1 (MeV ⁻¹)	d_2 (MeV ⁻²)
No+A	18.2 ± 1.2	3.1 ± 0.4	-1.62	10.3
Na	17.8	3.0	-1.26	10.8
T+A	15.7 ± 2.7	2.7 ± 0.8	-1.62	10.3
	20.8 ± 1.6		-1.5 ± 0.1	6.5 ± 2.0

 E. G. Adelberger, et al., *Rev. Mod. Phys.* 83, 195 (2011)

L. T. Baby, et al., [ISOLDE Collaboration], *Phys. Rev.Lett.* 90, 022501 (2003);

F. Hammache, et al., *Phys. Rev. Lett.* 86, 3985 (2001);

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B. W. Filippone, et al., *Phys. Rev. C* 28, 2222 (1983);

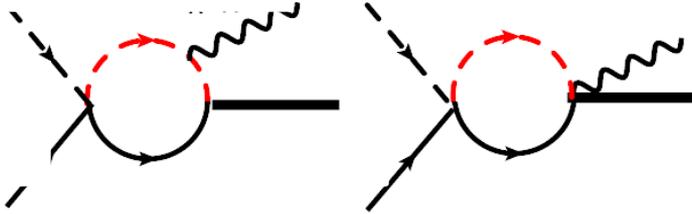
A. R. Junghans, et al., *Phys. Rev. C* 68, 065803 (2003);

A. R. Junghans, et al., *Phys. Rev. C* 81, 012801 (2010).

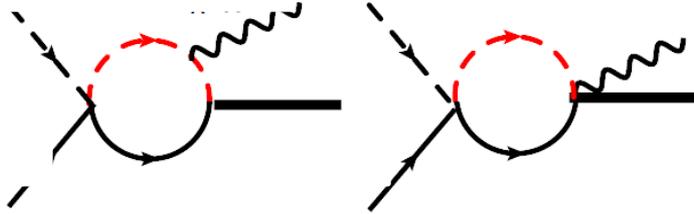
Summary

- EFT + ab initio works as expected at LO
- LO needs s-wave scattering length, p-wave ANCs, and binding momentum
- The p-wave is a coupled-channel problem
- For Be7 capture, improving s-wave measurement is important for extrapolating data to stellar energies.

Outlook: NLO

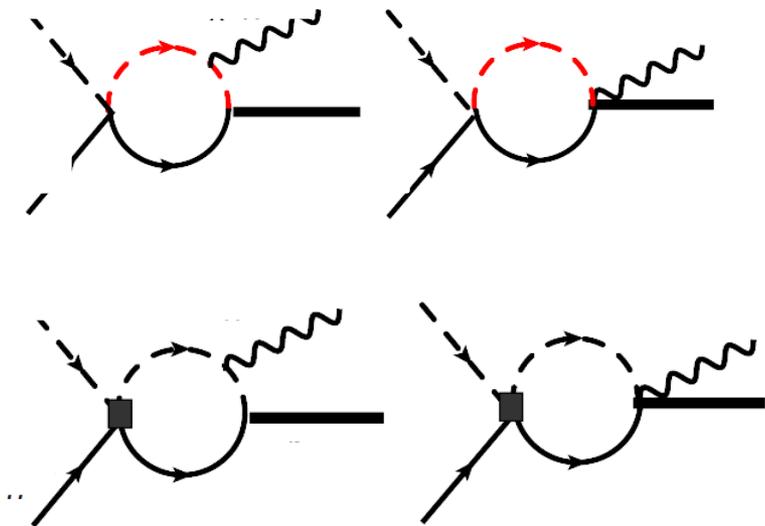


Outlook: NLO



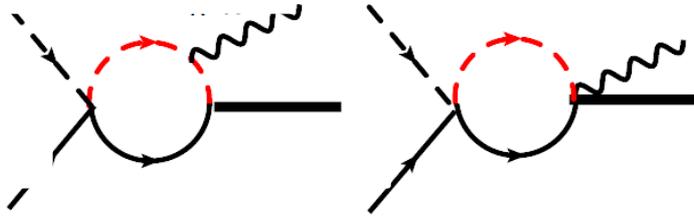
*Existence of
threshold*

Outlook: NLO

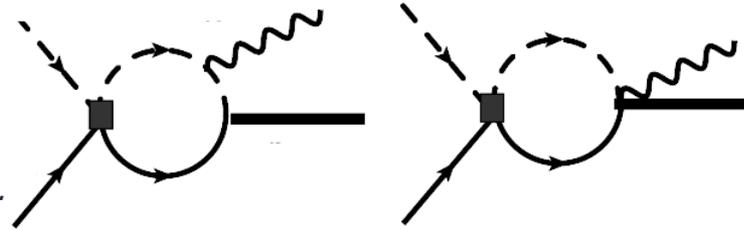


*Existence of
threshold*

Outlook: NLO

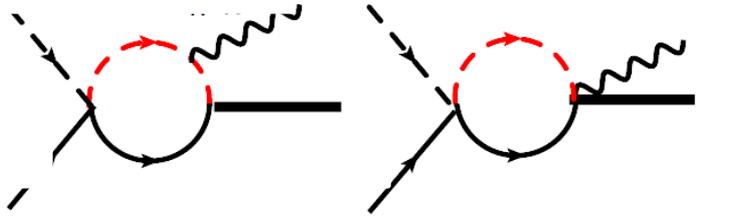


*Existence of
threshold*

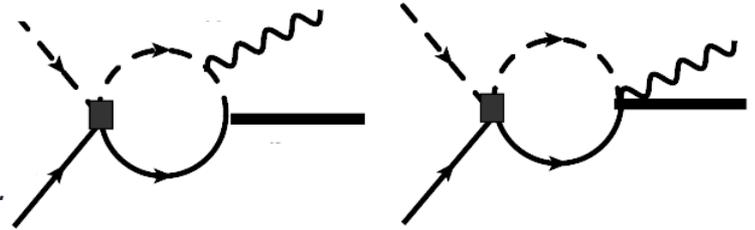


*Improve the initial
state multiple
scattering*

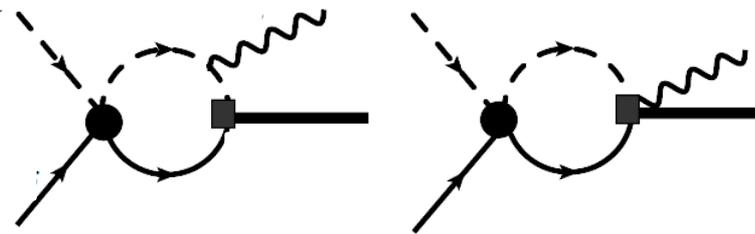
Outlook: NLO



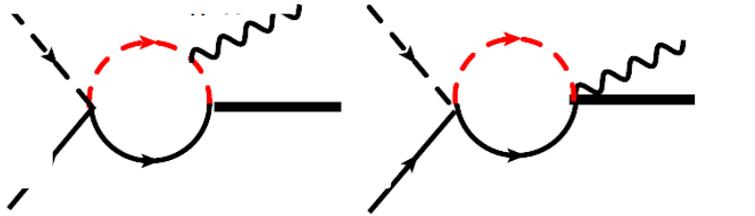
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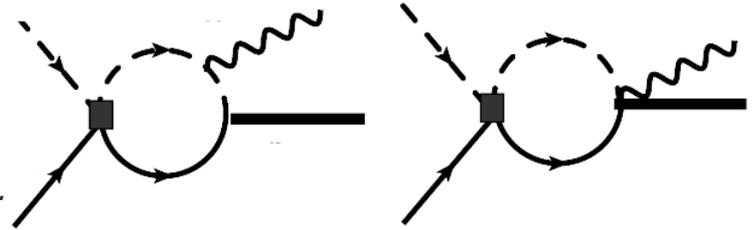
*Improve the initial
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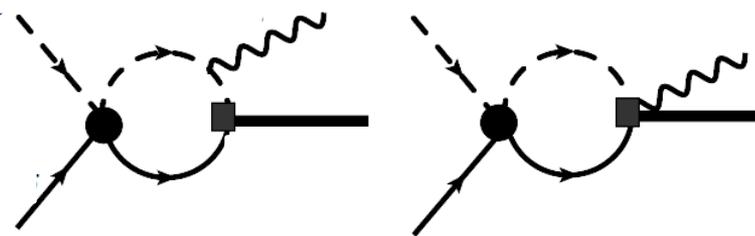
Outlook: NLO



*Existence of
threshold*

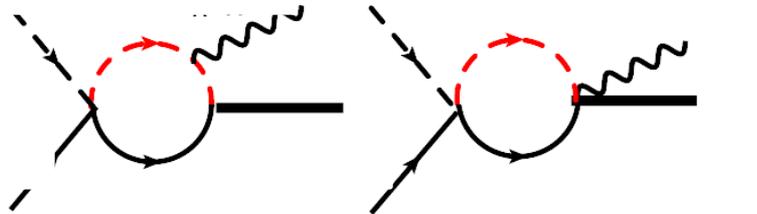


*Improve the initial
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scattering*

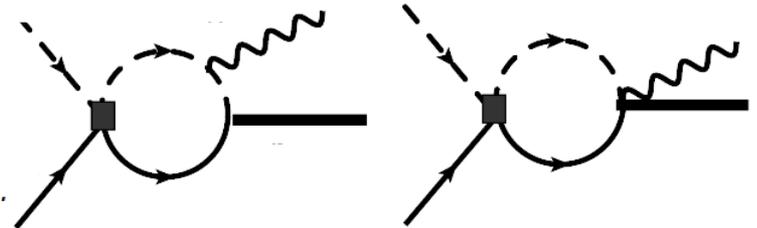


*Improve the final
state interaction*

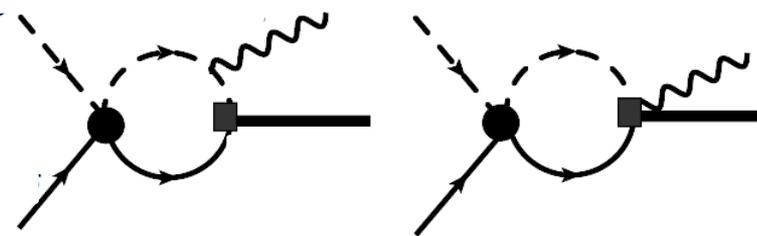
Outlook: NLO



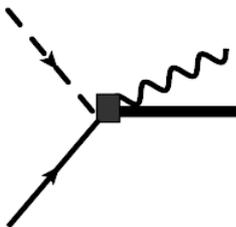
*Existence of
threshold*



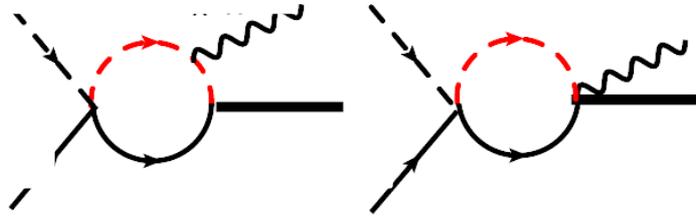
*Improve the initial
state multiple
scattering*



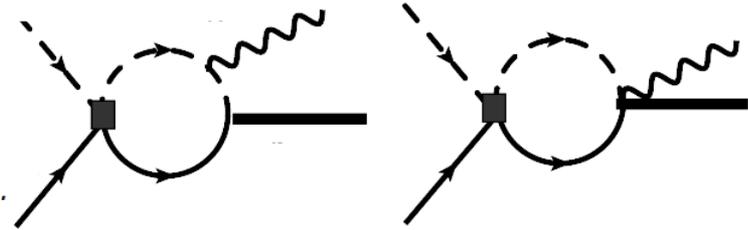
*Improve the final
state interaction*



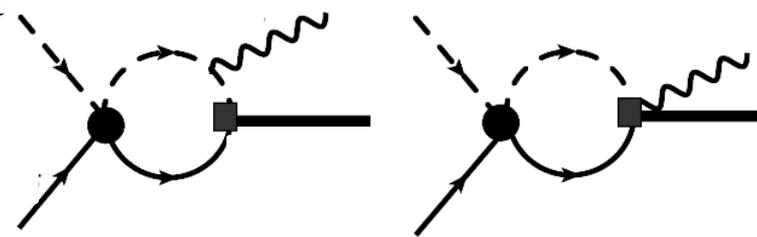
Outlook: NLO



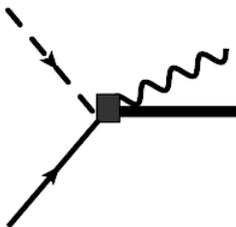
*Existence of
threshold*



*Improve the initial
state multiple
scattering*

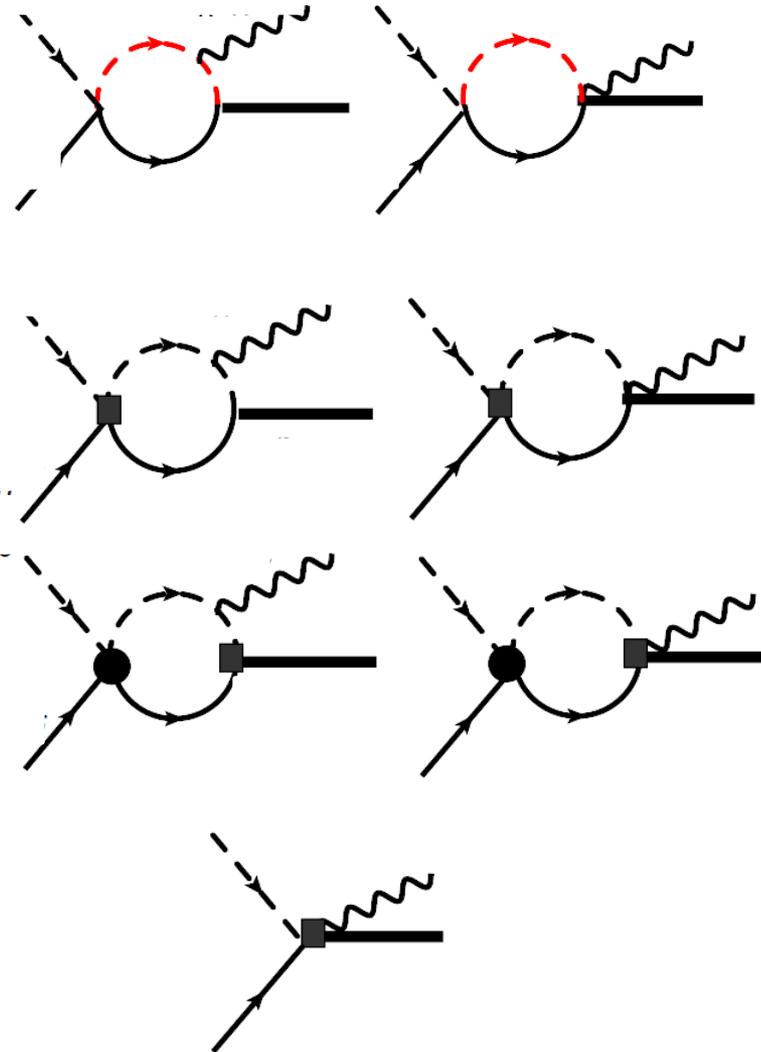


*Improve the final
state interaction*



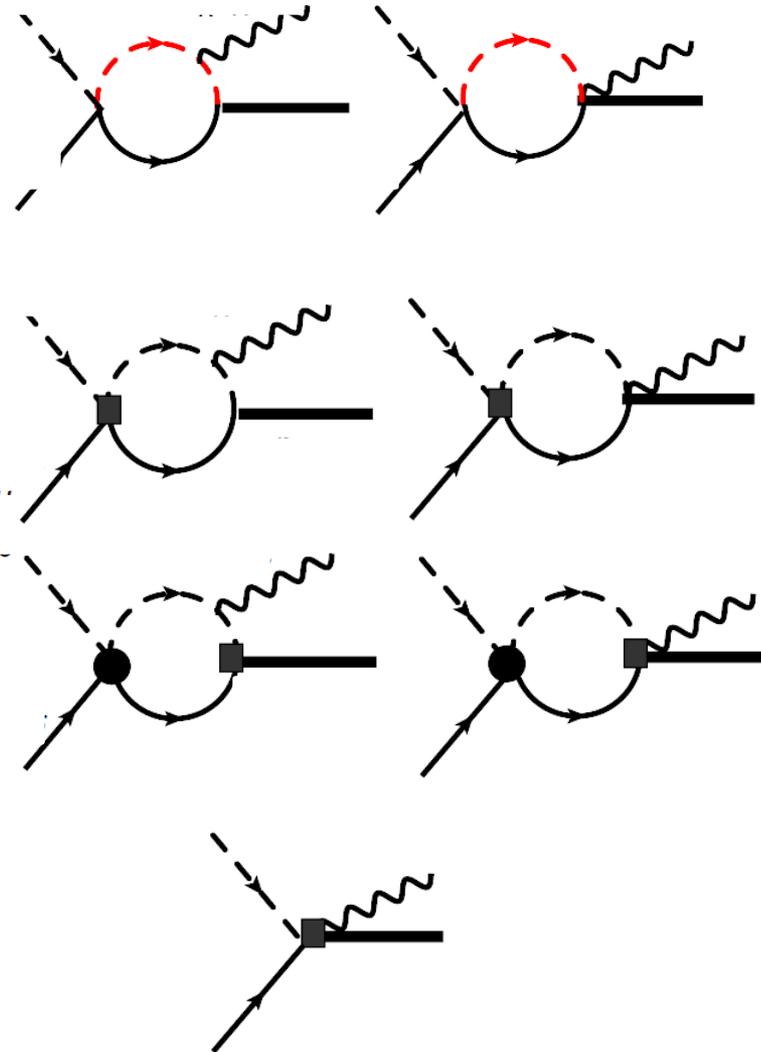
*Improve the short
distance contribution*

Outlook: NLO



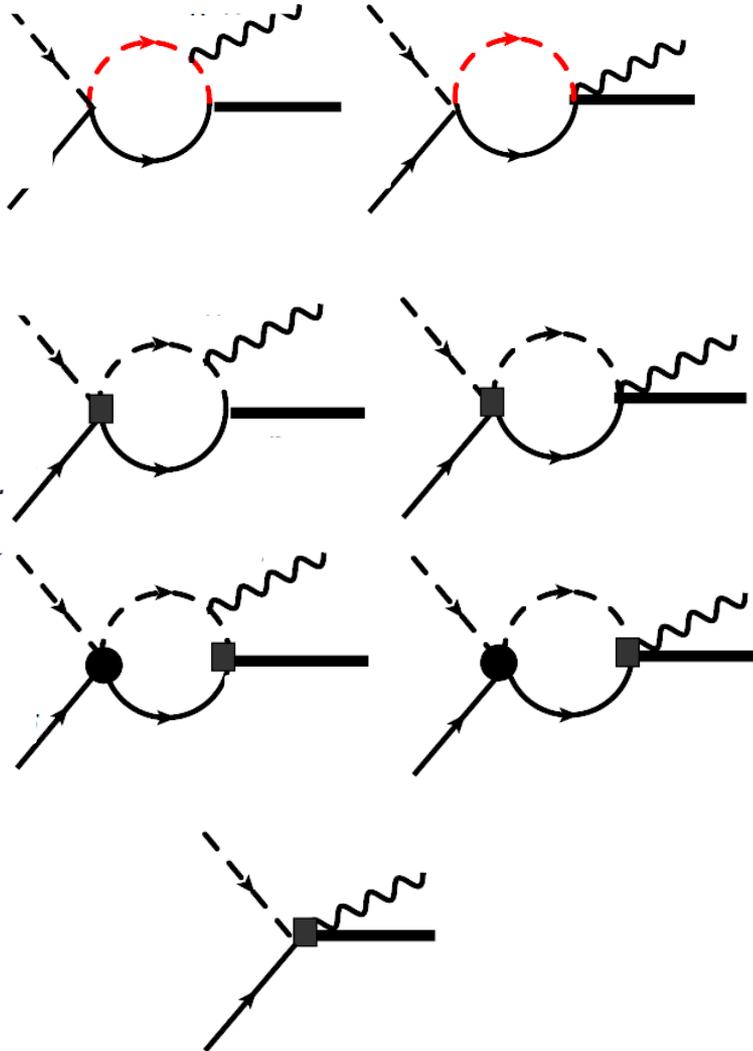
- Need to fix higher order couplings, i.e., need more “observables”.

Outlook: NLO



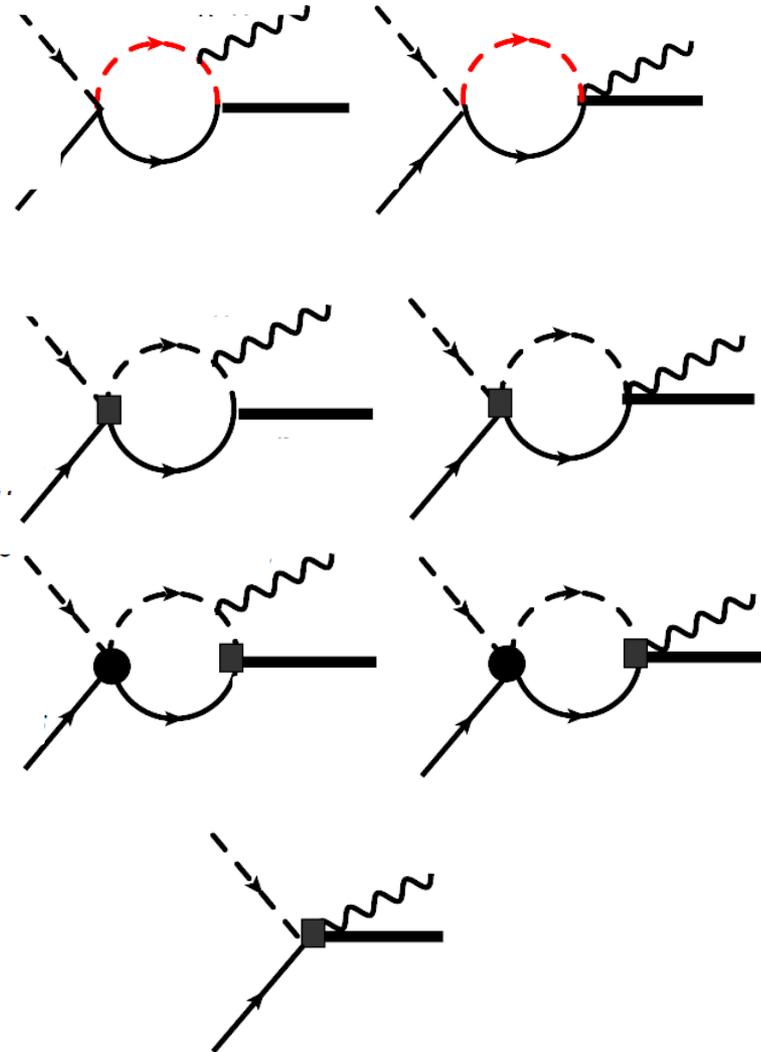
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- Extract from *direct* ab initio calculations (short distance)?

Outlook: NLO



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- Change the boundary conditions and the background fields?

Outlook: NLO



- Need to fix higher order couplings, i.e., need more “observables”.
- Extract from *direct* ab initio calculations (short distance)?
- Change the boundary conditions and the background fields?
- Use data directly?

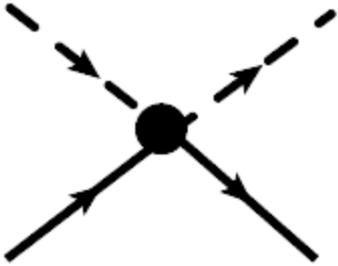
backup

- Capture cross section

$$\sigma(E) = \frac{S(E)}{E} \exp[-2\pi\eta(E)],$$

- 20 keV \sim fb
- 1MeV \sim mb

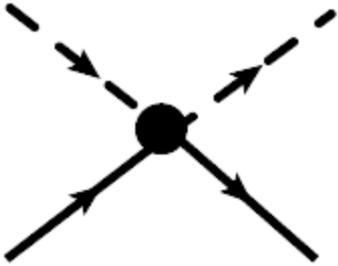
S-wave in EFT



Effective range expansion:

$$T = \frac{2\pi}{M_R} \frac{1}{-k \cot \delta_0 + ik}$$
$$-k \cot \delta_0 = \frac{1}{a_0} - \frac{1}{2} r_0 k^2 + \dots$$

S-wave in EFT



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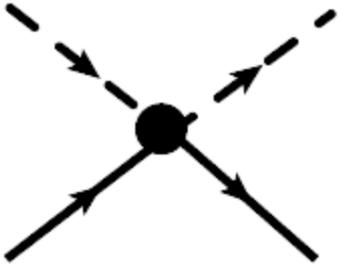
Effective range expansion:

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$$T = \frac{2\pi}{M_{\text{R}}} a_0 \left[1 - ia_0 k + \left(\frac{a_0 r_0}{2} - a_0^2 \right) k^2 + \dots \right]$$

Natural
 $a_0 \sim \frac{1}{\Lambda}$ and $r_0 \sim \frac{1}{\Lambda}$

S-wave in EFT



$$T = \frac{2\pi}{M_R} \frac{1}{-k \cot \delta_0 + ik}$$

Effective range expansion: $-k \cot \delta_0 = \frac{1}{a_0} - \frac{1}{2}r_0k^2 + \dots$

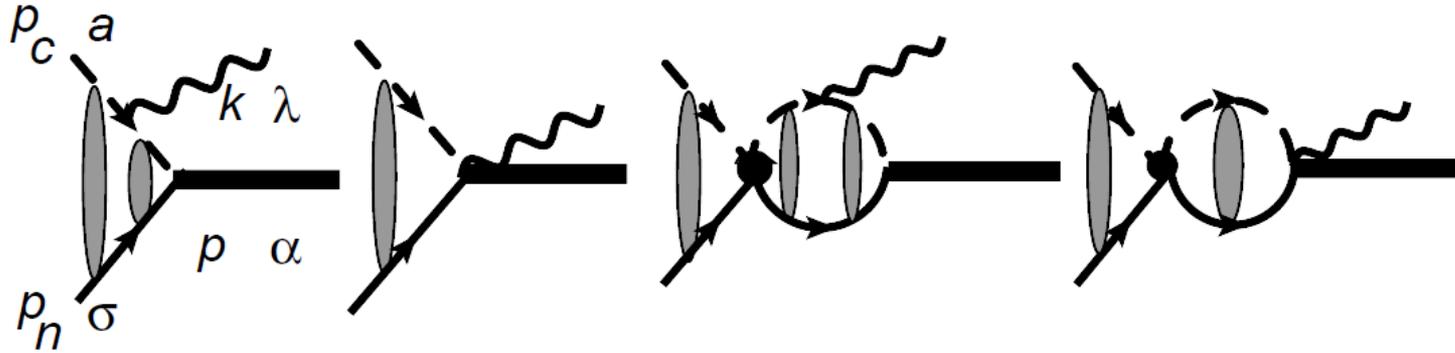
$$T = \frac{2\pi}{M_R} a_0 \left[1 - ia_0k + \left(\frac{a_0r_0}{2} - a_0^2 \right) k^2 + \dots \right] \quad \text{Natural}$$

$a_0 \sim \frac{1}{\Lambda}$ and $r_0 \sim \frac{1}{\Lambda}$

$$T = \frac{2\pi}{M_R} \frac{1}{a_0^{-1} + ik} \left(1 + \frac{r_0k^2}{2} \frac{1}{a_0^{-1} + ik} + \dots \right) \quad \text{Unnatural}$$

$a_0 \sim \frac{1}{\gamma}$ but $r_0 \sim \frac{1}{\Lambda}$

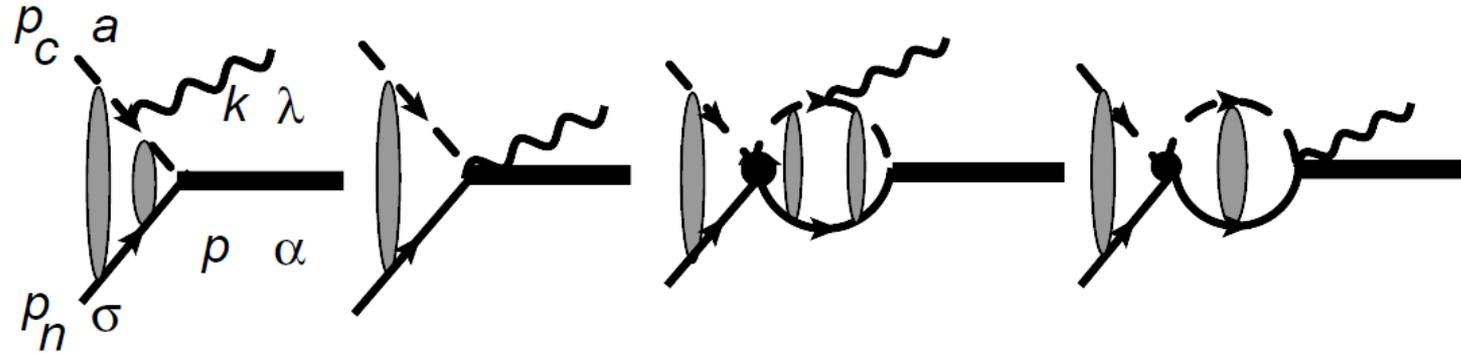
Radiative captures: LO



$$\langle \pi^\alpha | L_{EM} | \chi_{\mathbf{p}}^{(+), \delta, a} \rangle \equiv T_i^{\delta a} T_\alpha^{ij} \mathcal{M}_j \quad \text{Initial total spin } S_i=1$$

$$\mathcal{M}_j = (-i) C_{\eta,0} C_{({}^3P_2)}^{\text{LO}} \frac{Z_{\text{eff}}}{M_R} \frac{2\pi}{\sqrt{3}} (\gamma^2 + k^2) \left[e^{i\sigma_0} \epsilon_j^* Y_{00}(\hat{\mathbf{p}}) \mathcal{S}({}^3S_1) + e^{i\sigma_2} \epsilon_k^* \sqrt{2} T_j^{ka} Y_{2a}(\hat{\mathbf{p}}) \mathcal{D} \right]$$

Radiative captures: LO



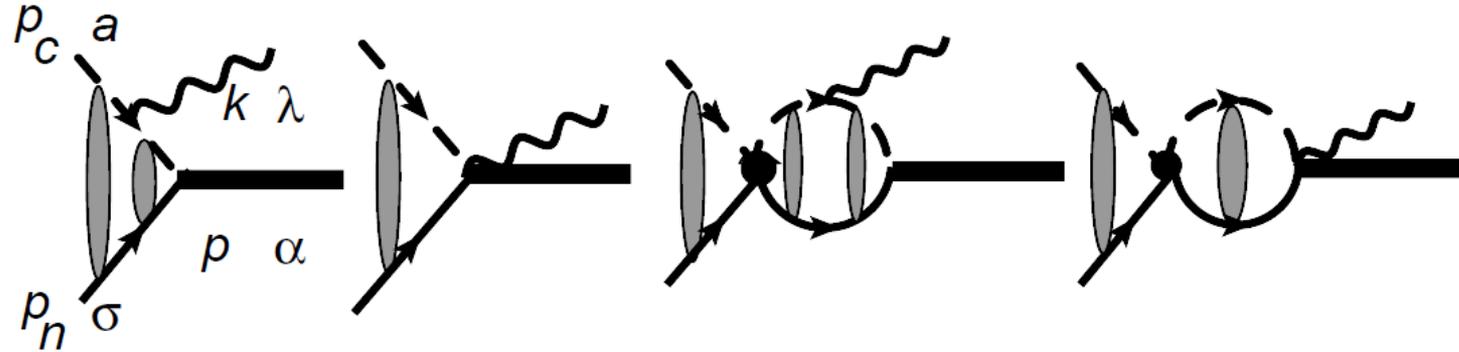
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$$\mathcal{S}(X) \equiv \int_0^{+\infty} dr W_{-\eta_B, \frac{3}{2}}(2\gamma r) r \left[\frac{C_{\eta,0} G_0(k, r)}{-a_{(X)}^{-1} - 2k_C H(\eta)} + \frac{F_0(k, r)}{C_{\eta,0} k} \frac{-a_{(X)}^{-1} - 2k_C \text{Re}[H(\eta)]}{-a_{(X)}^{-1} - 2k_C H(\eta)} \right]$$

$$\mathcal{D} \equiv \int_0^{+\infty} dr W_{-\eta_B, \frac{3}{2}}(2\gamma r) r \frac{F_2(k, r)}{C_{\eta,0} k}$$

Radiative captures: LO



$$\langle \pi^\alpha | L_{EM} | \chi_{\mathbf{p}}^{(+)} , \delta, a \rangle \equiv T_i^{\delta a} T_\alpha^{ij} \mathcal{M}_j \quad \text{Initial total spin } S_i=1$$

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*Coulomb
wavefunc.*

$F \rightarrow j$
 $G \rightarrow n$
 $W \rightarrow h$

