

#### Universal Aspects of Dipolar Scattering Christopher Ticknor

## May 15 at INT



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- Universal Dipolar Scattering

   Theory of a long range scattering potential

   2D Universal and tilted dipolar scattering
  - Interplay of interaction and geometry

• Universal 3-body dipolar recombination (not included)  $L_3 \propto D^4$ ,  $a^2 D^2$ 









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- But dipoles are different
- s=3 $\sigma_l \propto E^0$

- Long range
- Anisotropic
- Strong

 $V_{dd} = \frac{\vec{d} \cdot \vec{d} - 3(\hat{R} \cdot \vec{d})^2}{\mathbf{p}^3}$ 

 $V_{dd} \sim \langle d \rangle^2 \frac{1 - 3\cos^2(\theta)}{r^3}$ 



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## **Dipolar Scattering**



- The system can be magnetic or electric.
- Dipolar length scale

$$D = \langle d \rangle^2 m_r / \hbar^2$$

Energy scale

$$E_D = \langle d \rangle^2 / D^3 = \hbar^6 / \langle d \rangle^4 m_r^3$$

Strong interactions, many exotic theories.









Equation of Motion

$$\left(-\frac{\nabla^2}{2m_r} + \langle \mu \rangle^2 \frac{1 - 3\cos^2(\theta)}{R^3}\right) \psi = E \psi$$

$$y = R/D$$
  

$$R\psi = \sum_{l} Y_{lm} F_{l}(R)$$

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## **Dipolar Scattering**



Equation of Motion → Universal!

• Except for short range  $y_0 = R_0 / D \ll 1$ 

$$\left(-\frac{d^{2}}{2 d y^{2}}+\frac{l(l+1)}{2 y^{2}}\right)F_{l}+\sum_{l'}\frac{C_{ll'}}{y^{3}}F_{l'}=\frac{E}{E_{D}}F_{l}$$

$$C_{ll'} = \langle lm | 1 - 3\cos^2(\theta) | l'm \rangle$$

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#### Semi-classical solution from Eikonal approximation



Ticknor PRL 100, 133202 (2008) UNCLASSIFIED





Born Approximation: **all** dipolar coupled partial waves have an energy independent cross section!

$$\sigma_{ll'} \propto |\langle F_l | 1/R^3 | F_{l'} \rangle \langle lm | 1 - 3\cos^2(\theta) | l'm \rangle|^2$$

$$\sigma_e = 1.117 D^2$$
  
$$\sigma_o = 3.351 D^2$$

Bohn, Cavagnero, and Ticknor NJP 11 055039 (2009)

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## Experimental observation from Innsbruck.

Threshold universality



From K. Aikawa et al. PRL 112 010404 (2013)





#### **2D dipolar scattering with a tilt** Christopher Ticknor

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 $V_{dd}$ 

Strong

Long range

$$\begin{array}{l} R \to \rho \\ \hat{d} \to \cos\left(\alpha\right) \hat{z} + \sin\left(\alpha\right) \hat{x} \end{array}$$

 $=\frac{\vec{d}\cdot\vec{d}-3(\hat{R}\cdot\vec{d})^2}{\mathbf{p}^3}$ 

$$V_{dd} = \langle d \rangle^2 \frac{1 - 3\cos(\theta)^2 \sin^2(\alpha)}{\rho^3}$$

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## 2D dipolar scattering





- Universal
  - Repulsive diagonal potential for m=0
  - Diagonal in partial wave
  - Short range not really important





**2D Universal Dipolar Scattering** 

Semi-classical regime



 $\sigma_{sc} = \frac{4}{k} \sqrt{\pi Dk}$ 

#### Ticknor, PRA 80 052702 (2009)

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## Interplay of geometry and interaction

NATIONAL EST.1943  $\hat{d} \cdot \hat{\rho} \neq 0$ ZAttractive XRepulsive  $V_{dd} = \langle d \rangle^2 \frac{1 - 3\cos(\theta)^2 \sin^2(\alpha)}{\sigma^3}$ UNCLASSIFIED

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#### Breakdown of superfluidity in a BEC

NATIONAL LABORATOR - EST.1943 -

#### An amazing BEC experiment.



PRL 104, 160401 (2010)

Selected for a Viewpoint in Physics PHYSICAL REVIEW LETTERS

week ending 23 APRIL 2010

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#### **Observation of Vortex Dipoles in an Oblate Bose-Einstein Condensate**

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#### Interplay of geometry and interaction

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— EST.1943 —

- 2D dipolar gas
   with tilted polarization <sup>∞</sup>/<sub>\*</sub>
- Anisotropic superfluid character!



#### Ticknor et al. PRL 106 065301





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# Interplay of geometry and interaction Los Alamos

Binding Energies as function of a/D and  $\alpha$ 



#### Ticknor PRA 84 032702

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$$f(\theta_i, \theta_f) = \frac{e^{i\pi/4}}{\sqrt{2\pi k}} \sum_{mm'} e^{-im\theta_i} T_{mm'} e^{im'\theta_f}$$

$$\sigma(\theta_i) = \int d\theta_f |f(\theta_i, \theta_f)|^2$$

$$\sigma = \frac{1}{k} \sum_{mm'} |T_{mm'}|^2$$

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0.16

0.14

0.18

 $\rho_0/D$ 

0.2

0.22

0.24

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#### Interplay of geometry and interaction



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## Conclusions

## Universal dipolar scattering in 3D.

- Many partial waves always contribute.
- Dipolar scattering in 2D is universal.
- If polarization is tilted then:
  - Tune anisotropy of system.

Money from LANL and LDRD.



#### Interplay of geometry and interaction OS lamos . NATIONAL LABORATORY EST.1943 l 0.2 $\left|f\left(\theta_{i},\theta_{f}\right)\right|^{2}$ 0.1 0.3 ${\mathcal X}$ -0.2 -0.1 0.1 0.2 -0.1 0.2 -0.2 0. 0.15 -0.3 -0.1 0.2 0.3 -0.2 0.1 0.10 0.05 -0.1 0.15 -0.15 -0.05 0.05 0.10 -0.10 -0.2 -0.05

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-0.10



-0.3

#### Interplay of geometry and interaction



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 Studied formation of 2 body molecule from 3 body dipolar system.



Ticknor and Rittenhouse, PRL 105 013201 (2010)

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Also see: Wang and Greene PRA 85 022704 (2012)

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#### Example wave functions.







Now use Fermi Golden Rule (FGR)

$$\langle 3-body \ plane \ wave | V_{dd} | molecule + third \rangle$$

$$V_{dd} \propto \sum_{k m_1 m_2} \frac{R_{<}^k}{R_{>}^{3+k}} [Y_{k, m_1} \otimes Y_{k+2, m_2}]_{20}$$





#### Interplay of geometry and interaction





#### Interplay of geometry and interaction NATION EST. 1943



 $\alpha/\pi = 0.315$ 



## Interplay of geometry and interaction

EST. 1943

 $\hat{d}\cdot\hat{\rho}\neq 0$ 

$$\langle m|1 - 3(\hat{d} \cdot \hat{\rho})^2 | m' \rangle$$
$$= \left(1 - \frac{3}{2}\sin^2(\alpha)\right) \delta_{mm'} - \frac{3}{4}\sin^2(\alpha) \delta_{mm'\pm 2}$$



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- Semiclassical Universality first predicted numerically
- True of Fermions and Bosons and distinguishable particles.



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# Interplay of geometry and interaction $\hat{d} \cdot \hat{\rho} \neq 0$

$$k \sigma_{m \to m'} = \frac{4 (Dk)^2}{(m^2 - 1/4)^2} (1 - \frac{3}{4} \sin(\alpha)^2)^2$$

$$k \sigma_{m' \to m+2} = \frac{4 (Dk)^2}{(m-1/2)(m+3/2)} (\frac{3}{4} \sin(\alpha)^2)^2$$

$$k \sigma_{\mp 1 \to \pm 1} = \frac{4(Dk)^2}{(m^2 - 1/4)^2} (\frac{3}{4} \sin(\alpha)^2)^2$$



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