

INT Workshop
2014

May 13,

Geometric scaling of three-body collision resonances for a ${}^6\text{Li}$ - ${}^{133}\text{Cs}$ mixture in the Efimov scenario



Funding:



MRSEC



Colin V. Parker
University of Chicago

WEAKLY-BOUND STATES OF THREE RESONANTLY-INTERACTING PARTICLES

V. N. EFIMOV

A. F. Ioffe Physico-technical Institute, USSR Academy of Sciences

Submitted February 16, 1970

Yad. Fiz. 12, 1080-1091 (November, 1970)

It is shown that if the pair forces of three identical particles are sufficiently resonant, a family of bound states of low energy is produced. The quantum numbers of all the states are the same: for spinless bosons 0^+ and for nucleons $\frac{1}{2}^+$, $T = \frac{1}{2}$. The dimension of the states is larger than the radius of the pair forces. The most favorable conditions for the appearance of a family of levels occur for three spinless neutral bosons: the conditions are less favorable for charged particles and particles with spin and isospin. The possibility of existence of such levels in a system of three particles (in the C^{12} nucleus) and of three nucleons (H^3) is considered.

Hyperspherical
equation:

$$i\hbar \frac{\partial \phi}{\partial t} = -\frac{\hbar^2}{m} \frac{\partial^2 \phi}{\partial R^2} - \frac{\hbar^2}{m} \frac{s_0^2 + 1/4}{R^2} \phi$$

\Rightarrow Scale invariant under dilation $R \rightarrow \lambda R$ and $t \rightarrow \lambda^2 t$

$\lambda = \exp(\pi/s_0) = 22.7$ for 3 identical bosons



Vitaly Efimov

From Nuclear Physics to Cold atoms

Chris Greene (Purdue)



Brett Esry (Kansas State)



VOLUME 83, NUMBER 9

PHYSICAL REVIEW LETTERS

30 AUGUST 1999

Recombination of Three Atoms in the Ultracold Limit

B. D. Esry

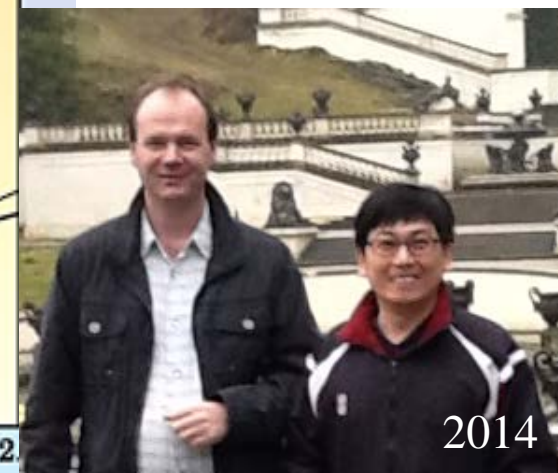
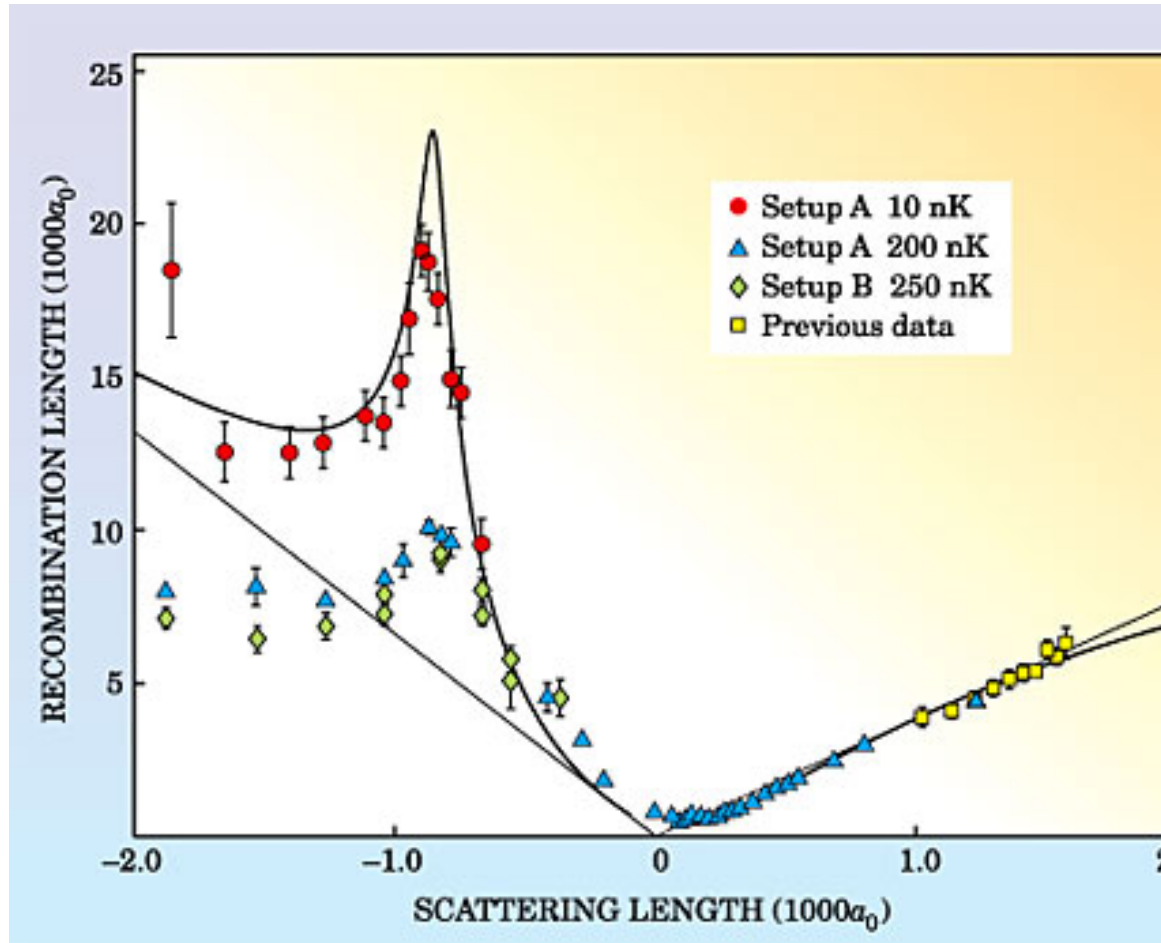
*Institute for Theoretical Atomic and Molecular Physics, Harvard-Smithsonian Center for Astrophysics,
Cambridge, Massachusetts 02138*

Chris H. Greene and James P. Burke, Jr.

Department of Physics and JILA, University of Colorado, Boulder, Colorado 80309-0440

(Received 19 May 1999)

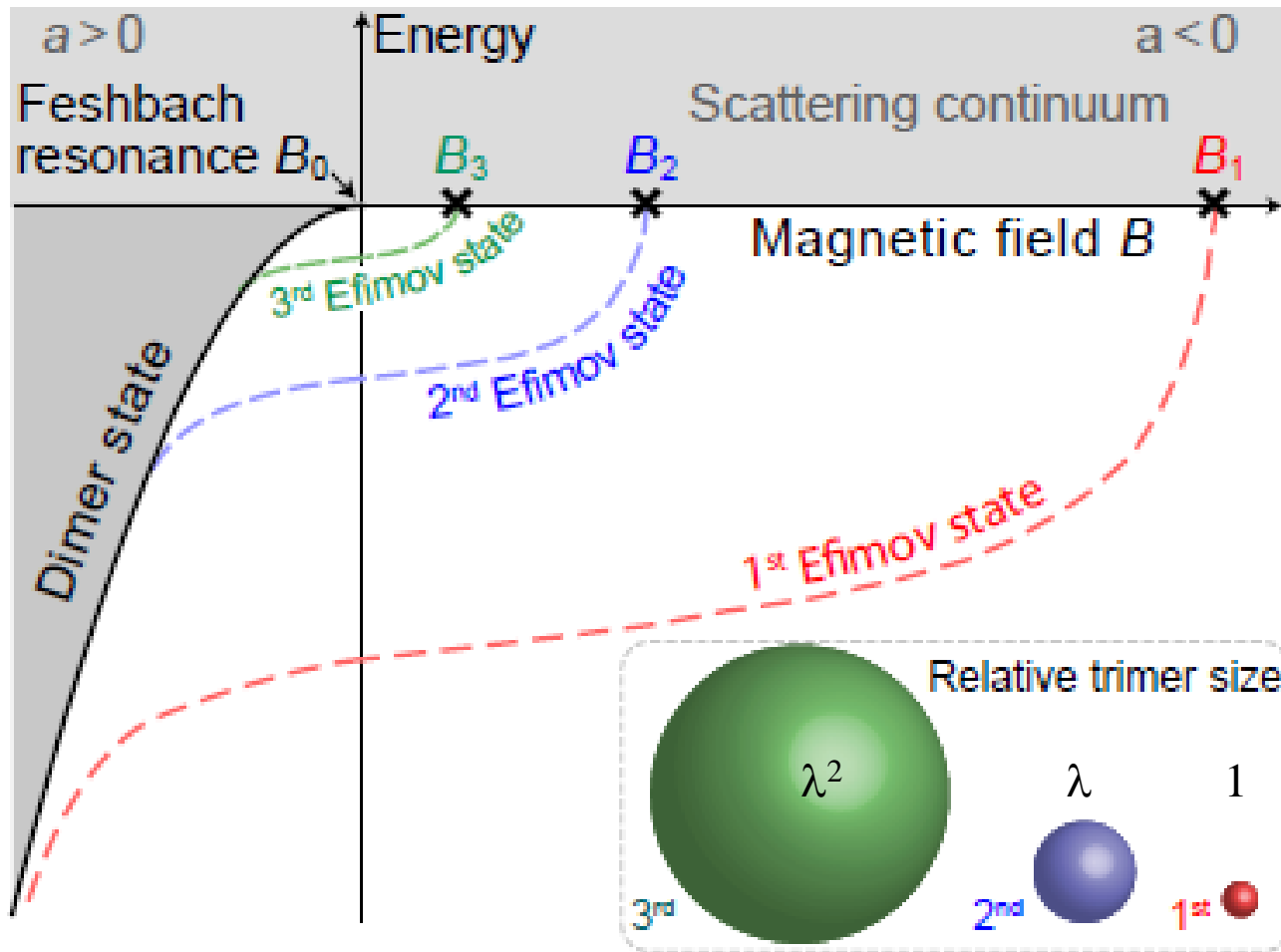
Observation of Efimov Resonance in Cesium



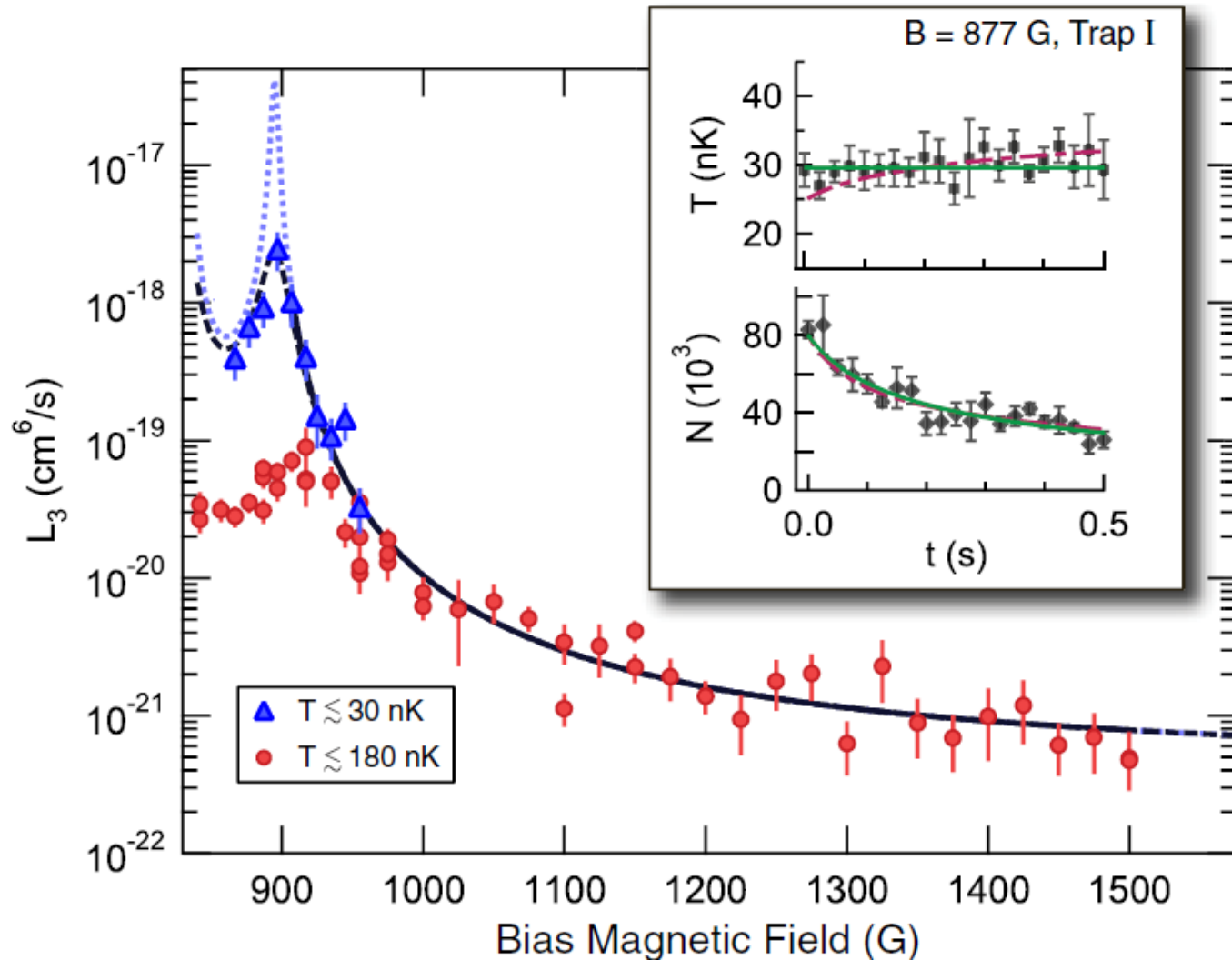
Physics Today 2006

Also found in Li6 (Penn State, Heidelberg), Li7 (Rice, Bar-Ilan), K39 (LENS), Rb85 (JILA)

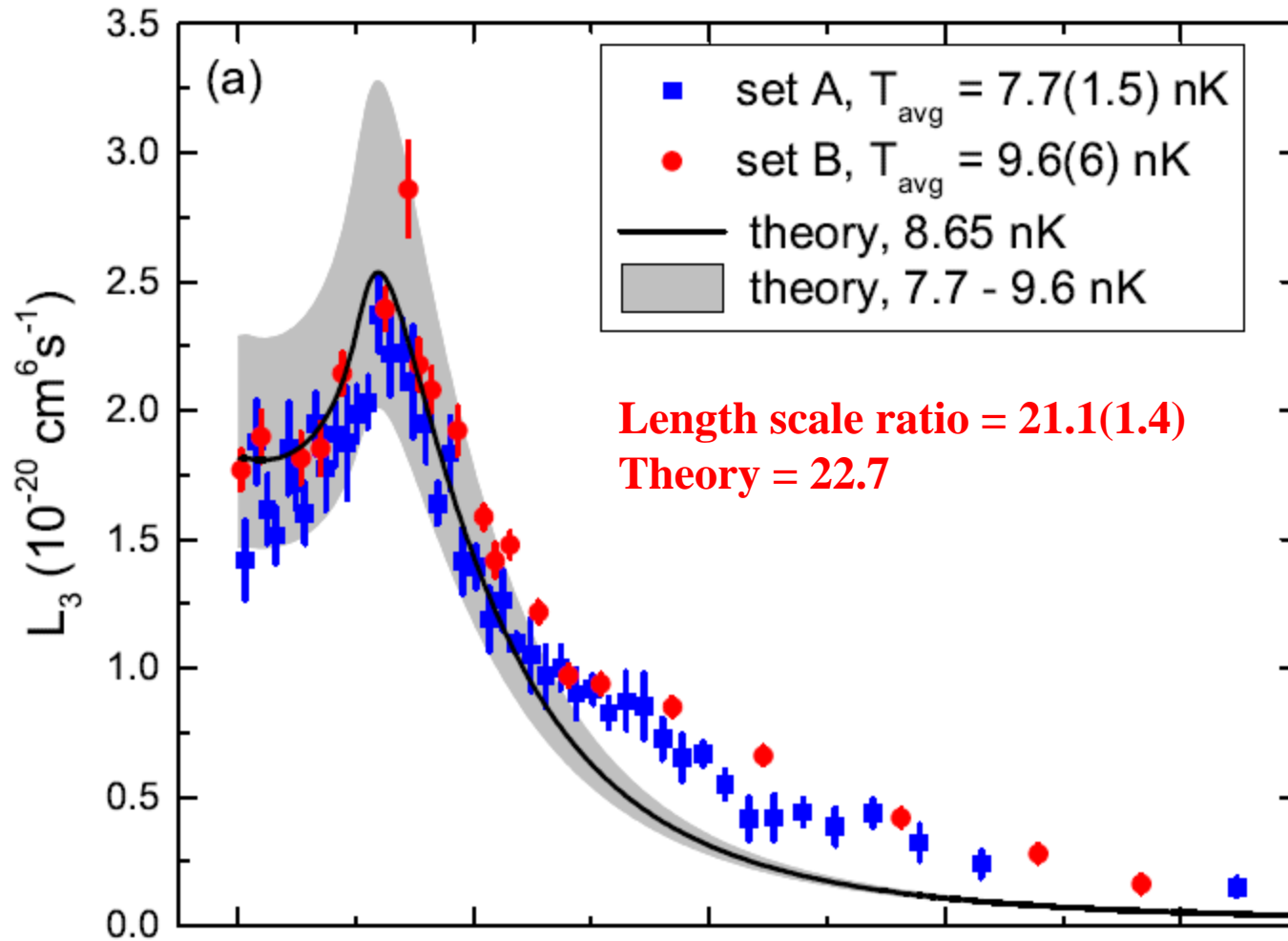
Efimov state structure near a Feshbach resonance



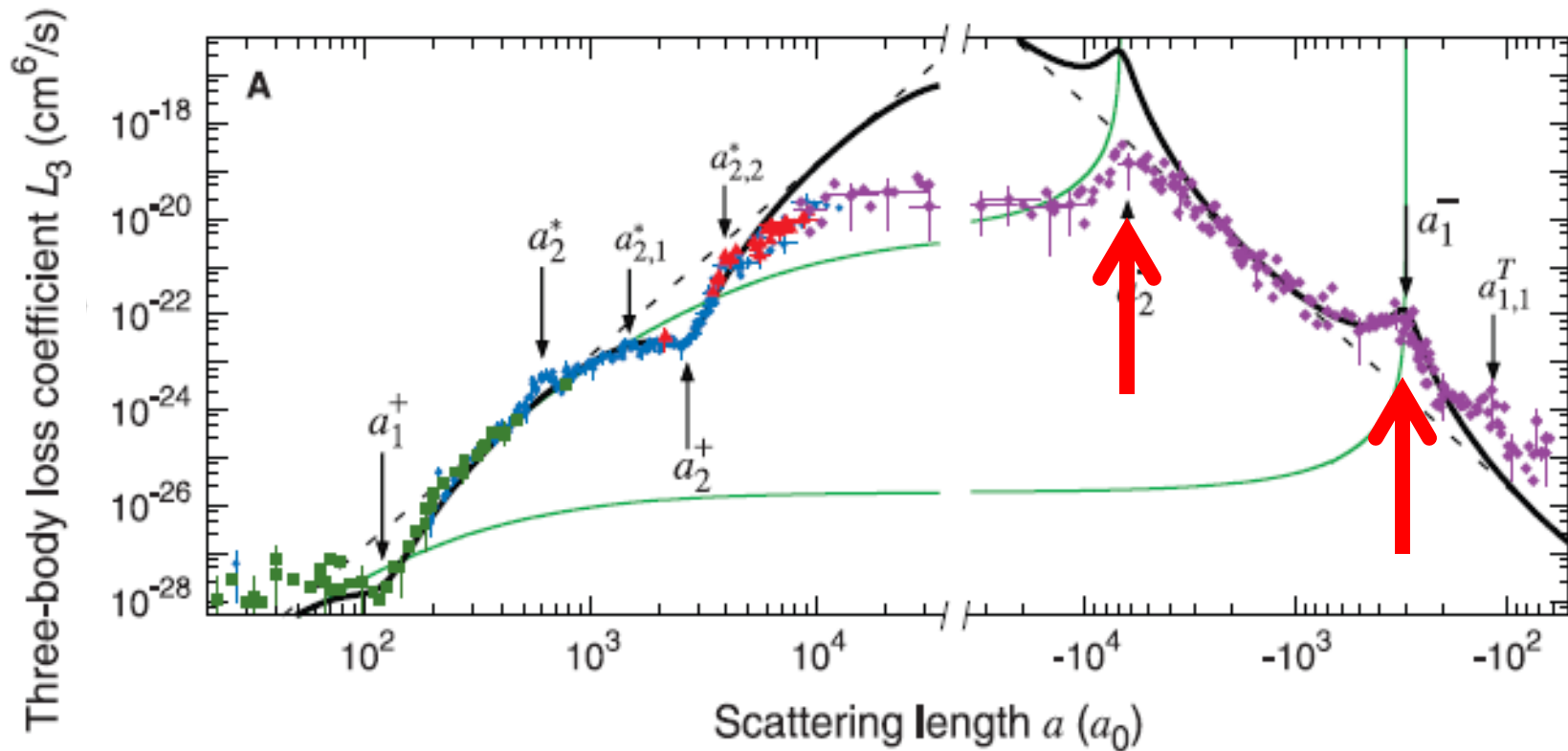
Excited Efimov state in 3-component Fermi gas



Second Efimov resonance in Cs!



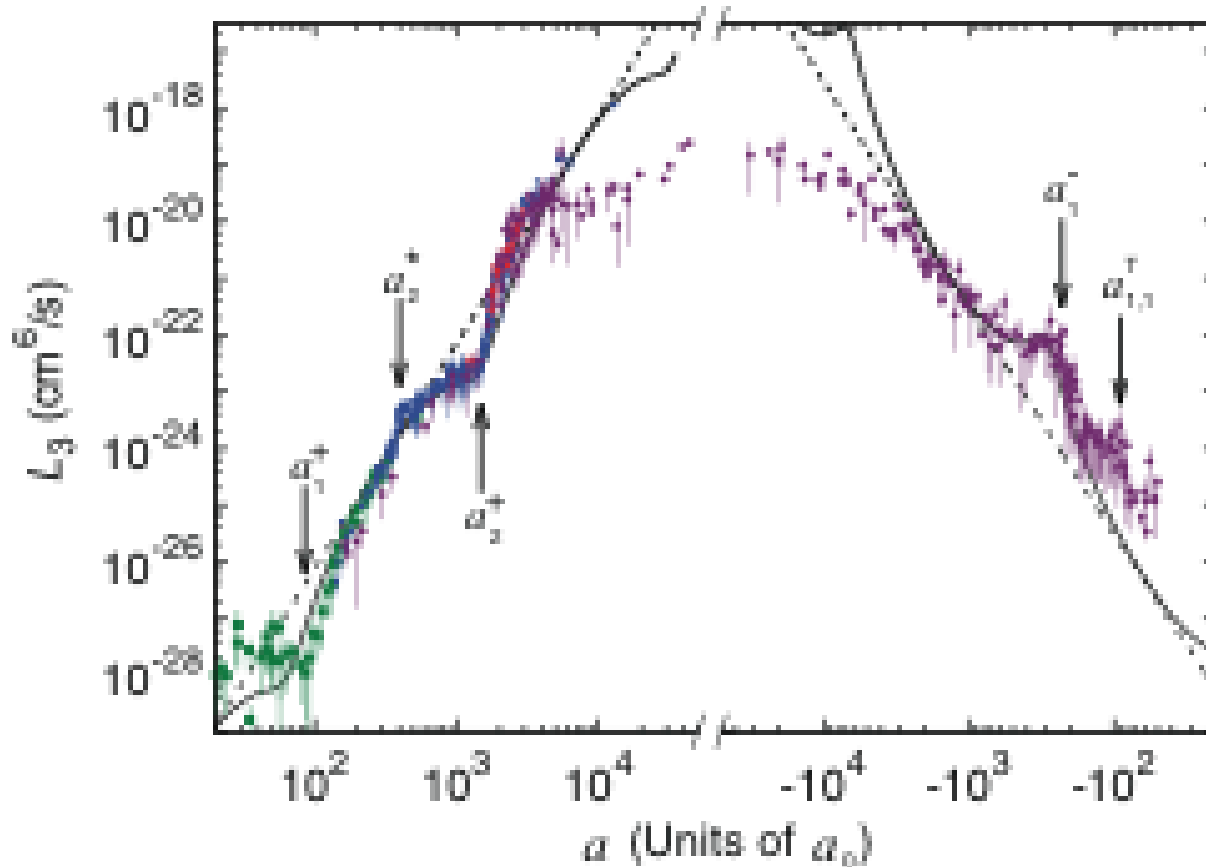
Second Efimov resonance in ${}^7\text{Li}$?



Hulet group (Rice University): *Science* 326, 1683 (2009)

New interpretation: *Physical Review A* 88, 023625 (2013).

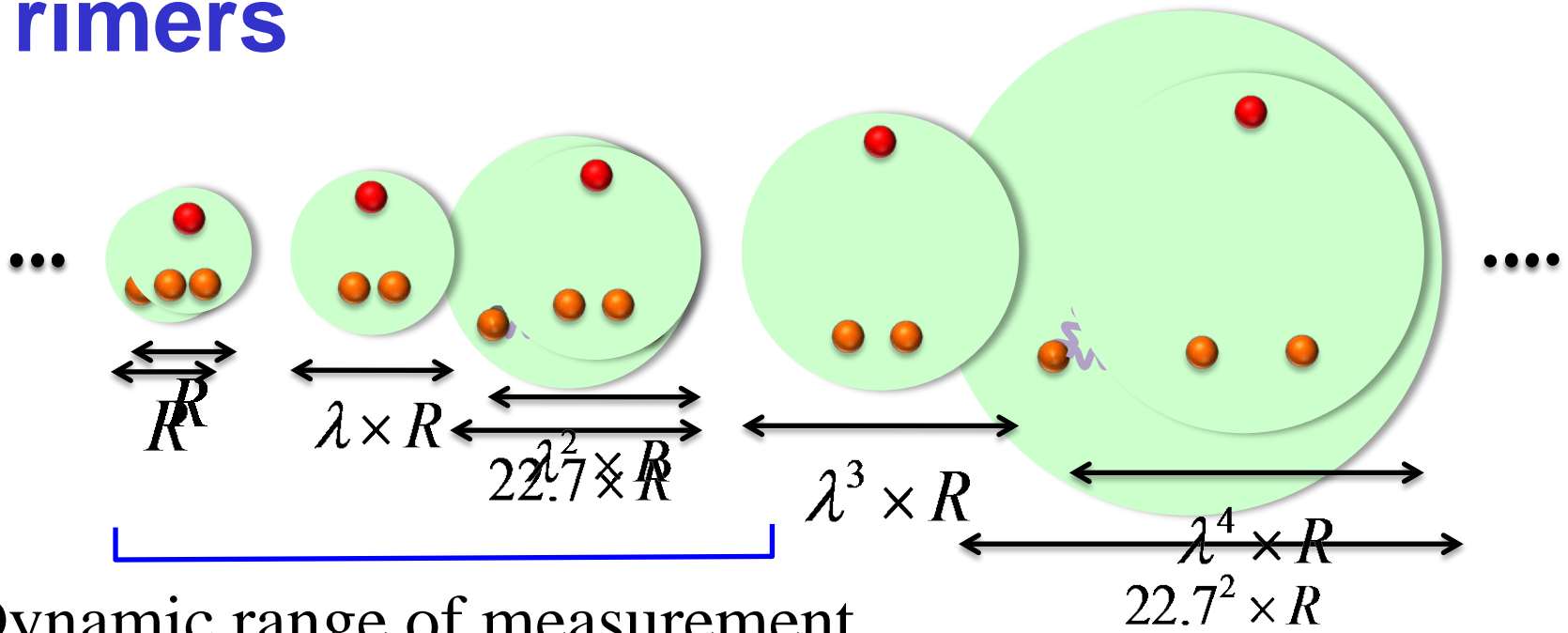
Second Efimov resonance in ${}^7\text{Li}$?



Hulet group (Rice University): *Science* 326, 1683 (2009)

New interpretation: *Physical Review A* 88, 023625 (2013).

Homonuclear vs. Heteronuclear Trimers

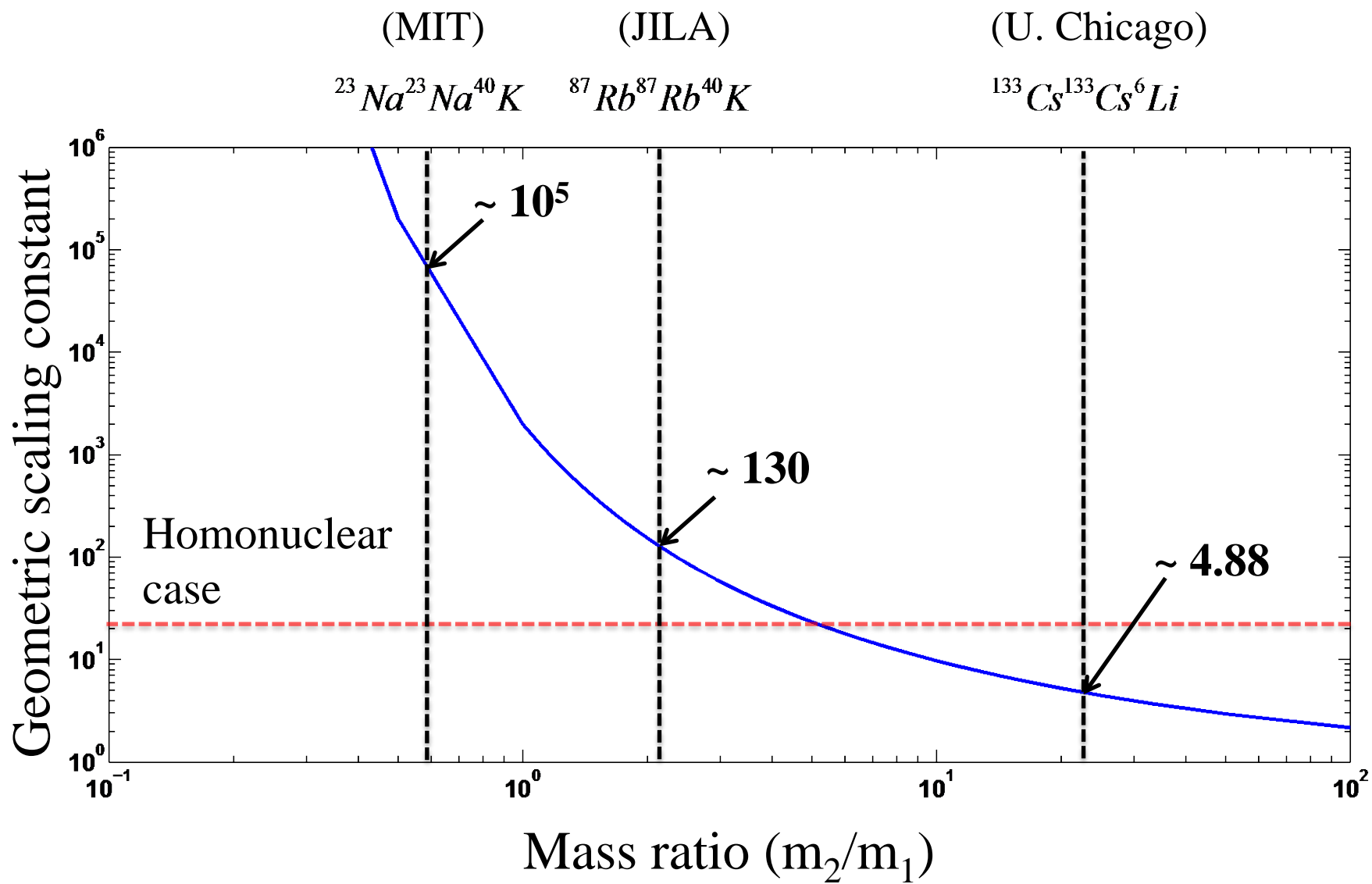


Dynamic range of measurement

λ : Determined by

- (1) **Mass ratio of constituent atoms**
- (2) Number of resonant pair interactions
- (3) Symmetry of particle statistics

Mass Dependence



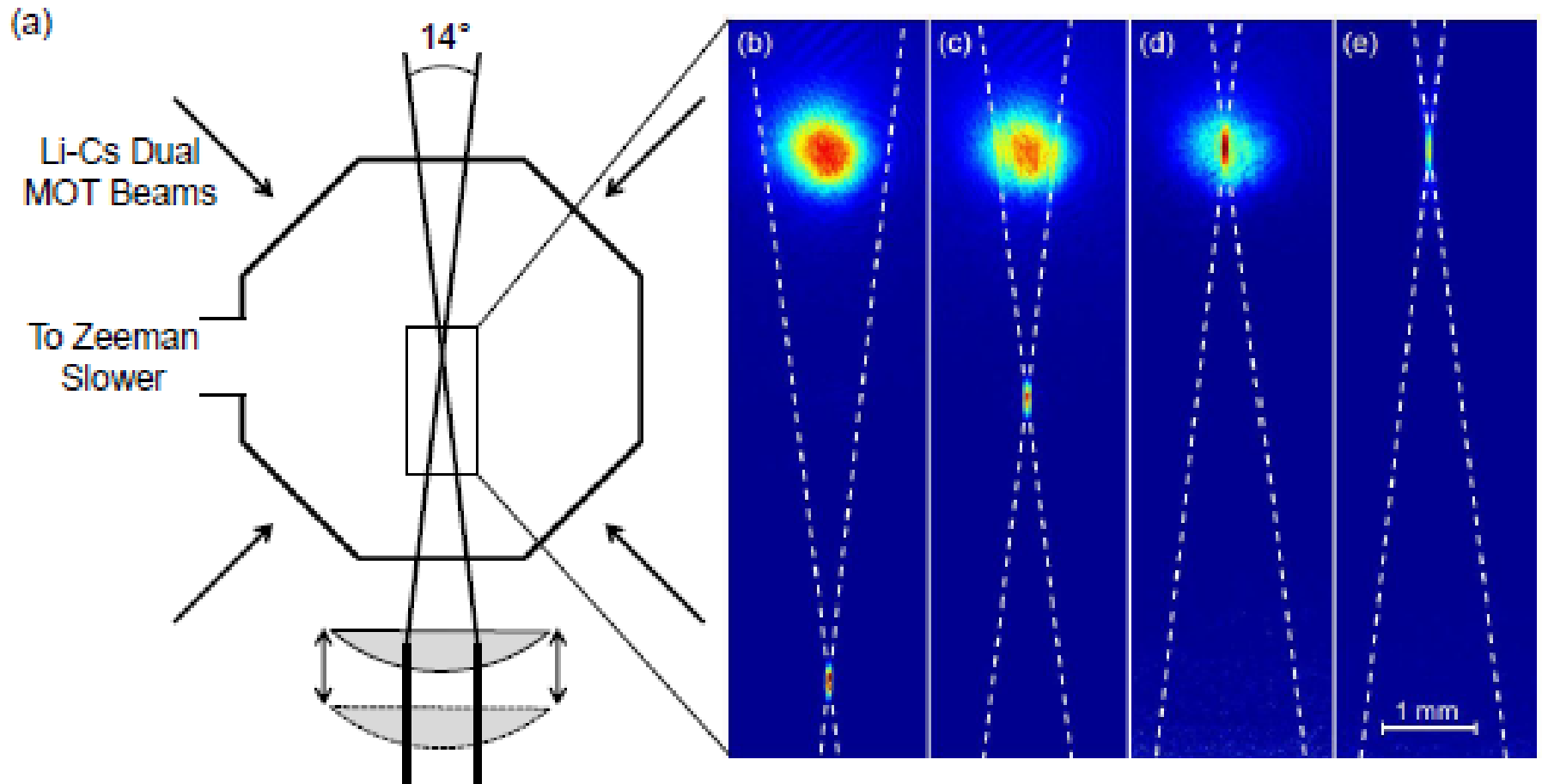
Efimov states in atomic mixtures

2 heavy Bosons + 1 light atom

$B-F$	e^{π/s_0}	Two features		Three features	
		$ a_{\min} $	$E_{\max}(\text{nK})$	$ a_{\min} $	$E_{\max}(\text{nK})$
$^{133}\text{Cs}-^6\text{Li}$	4.877	3×10^3	1500	2×10^4	60.0
$^{87}\text{Rb}-^6\text{Li}$	6.856	8×10^3	230	6×10^4	5.00
$^{23}\text{Na}-^6\text{Li}$	36.28	9×10^5	$\ll 0.1$	3×10^7	$\ll 0.1$
$^7\text{Li}-^6\text{Li}$	$> 10^2$	$\gg 10^8$	$\ll 0.1$	$\gg 10^8$	$\ll 0.1$
$^{133}\text{Cs}-^{40}\text{K}$	47.02	2×10^6	$\ll 0.1$	9×10^7	$\ll 0.1$
$^{87}\text{Rb}-^{40}\text{K}$	$> 10^2$	$\gg 10^8$	$\ll 0.1$	$\gg 10^8$	$\ll 0.1$
$^{23}\text{Na}-^{40}\text{K}$	$> 10^2$	$\gg 10^8$	$\ll 0.1$	$\gg 10^8$	$\ll 0.1$
$^7\text{Li}-^{40}\text{K}$	$> 10^2$	$\gg 10^8$	$\ll 0.1$	$\gg 10^8$	$\ll 0.1$

J.P. D'Incao and B.D. Esry, PRA (2006)

A Translatable Crossed Dipole Trap

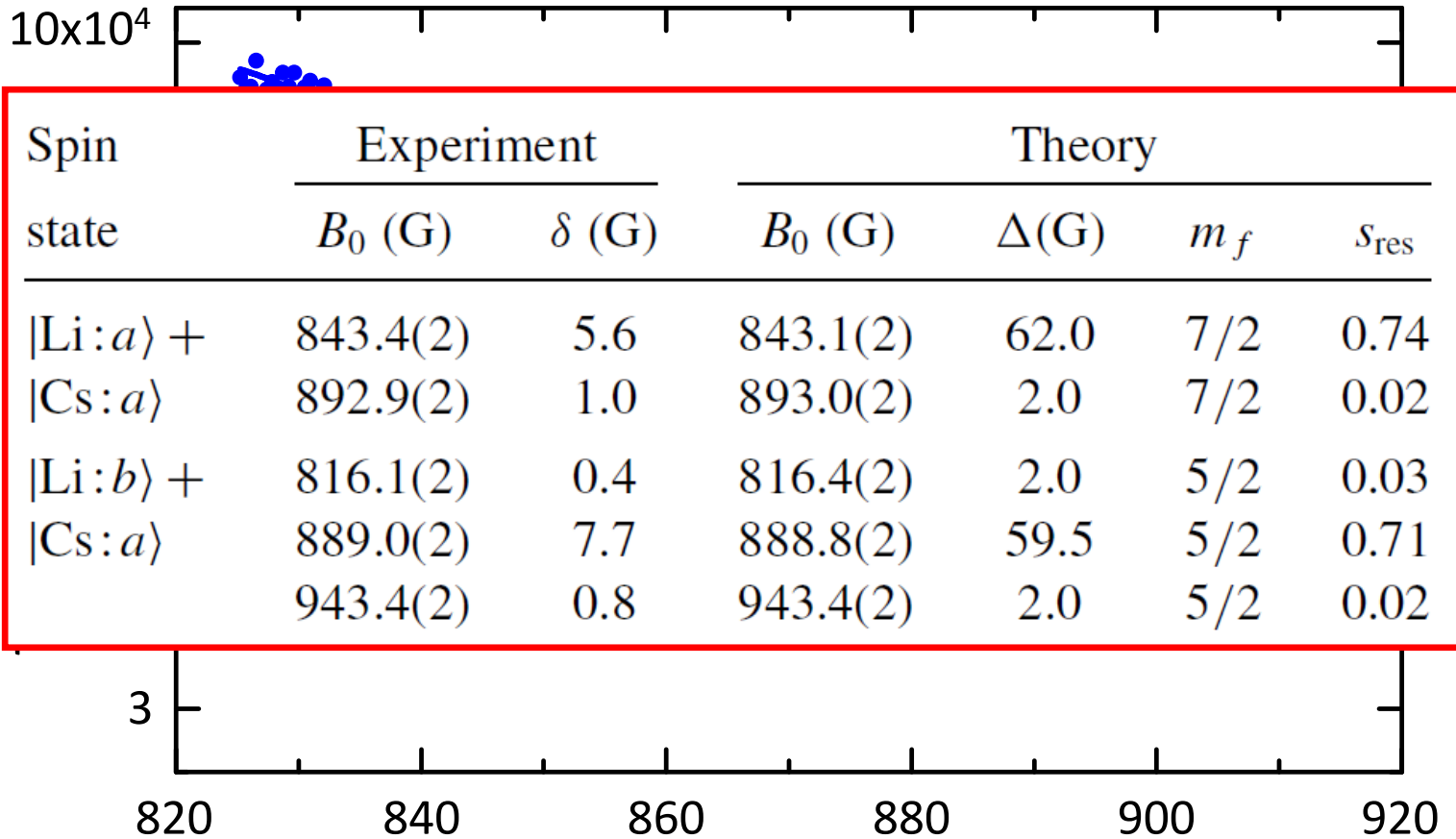


Li : $T = 3 \mu\text{K}$; $N_{\text{Li}} = \sim 10^5$

Cs: $T = 16 \mu\text{K}$; $N_{\text{Cs}} = \sim 10^5$

Trap Loss of Li-a State + Cs-a State Mixture

Lithium number



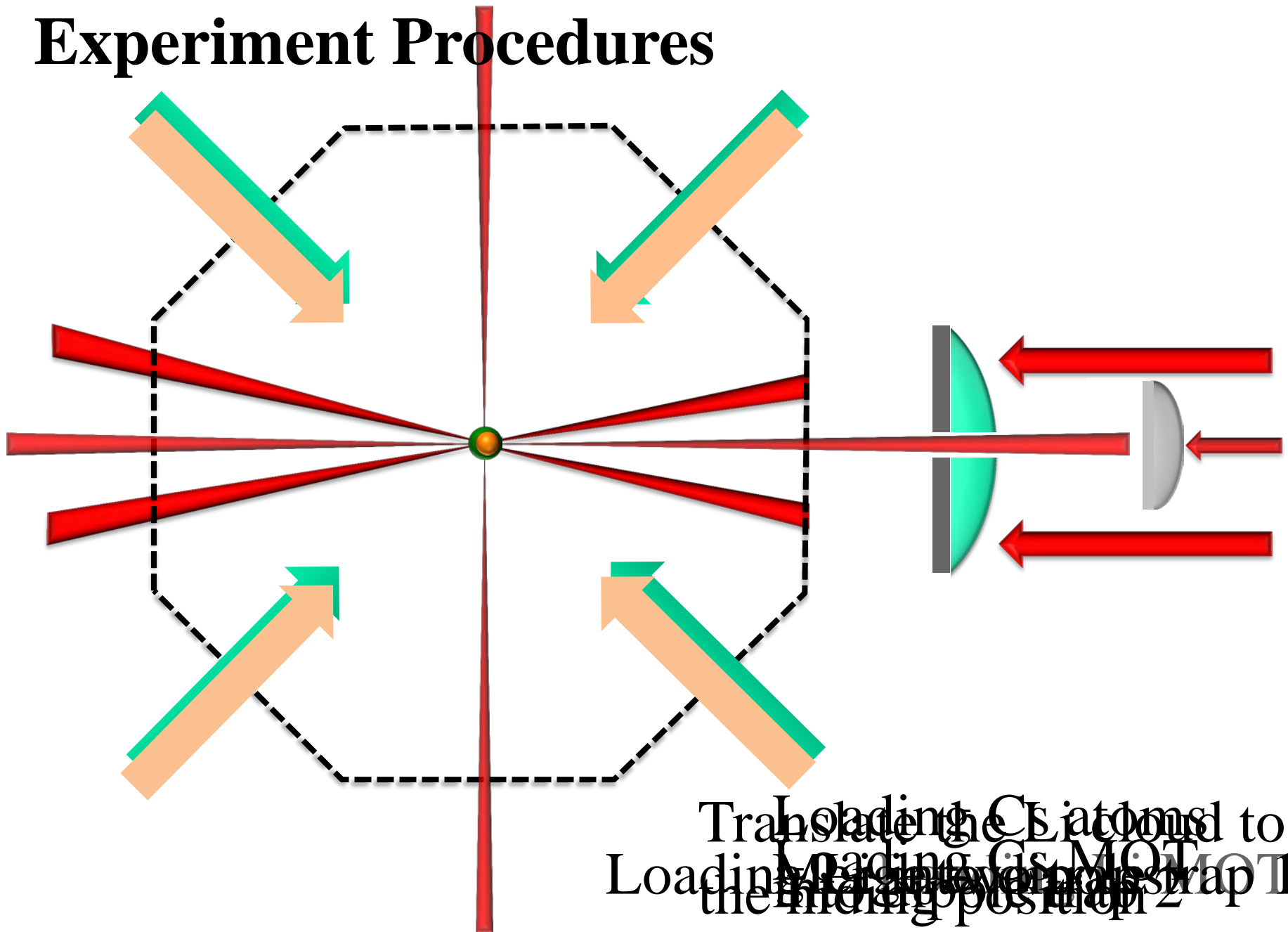
Tempera

Magnetic field B (G)

Tung et al., PRA 87 010702 (2013)

Repp et al., PRA 87 010701 (2013)

Experiment Procedures

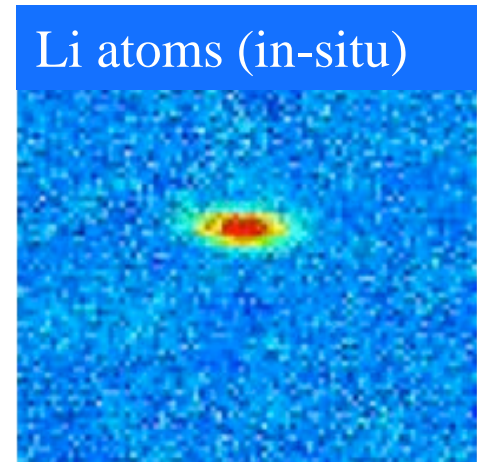
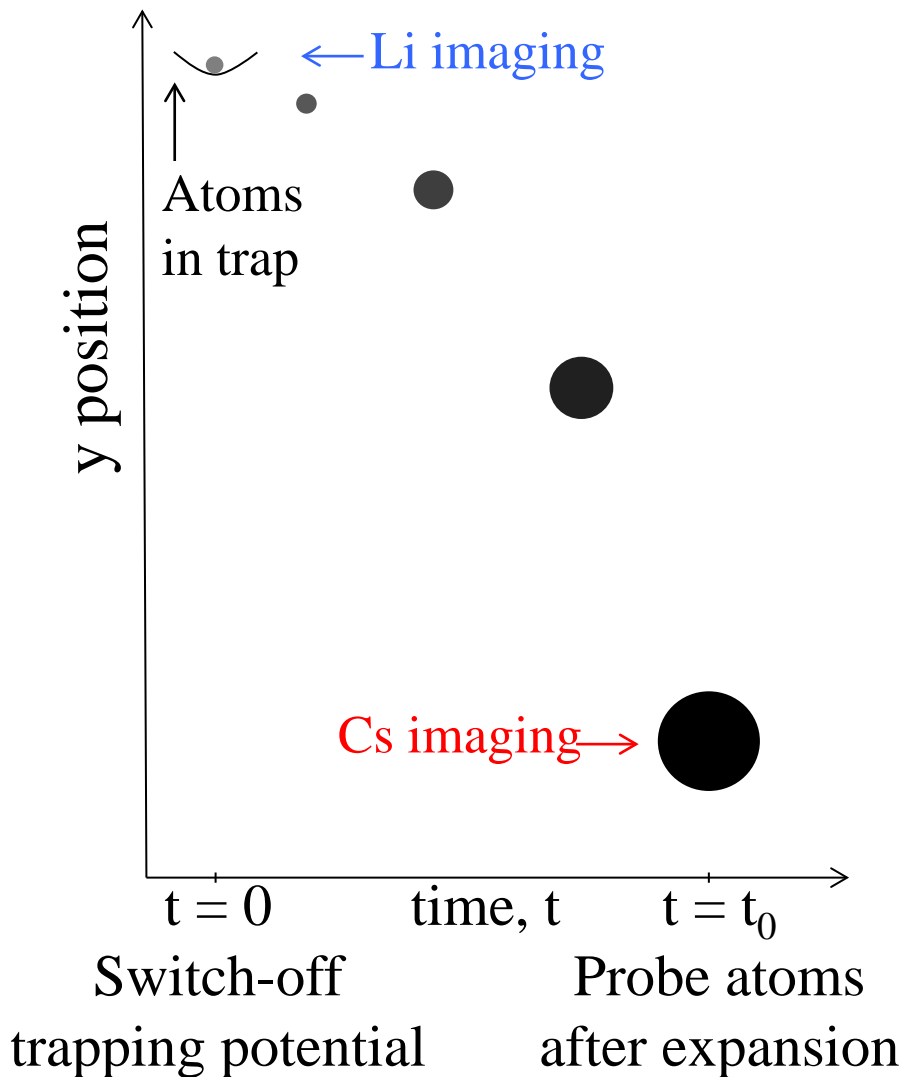


Translating the Cs atoms to Loading position
Loading Cs MOT the initial position

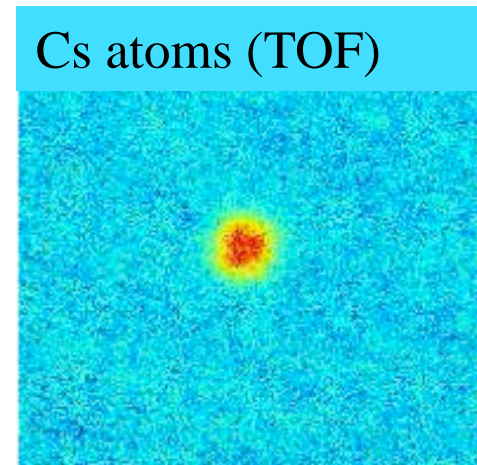
Dual Resonant Absorption Imaging

The time of flight method

Resonant absorption imaging



30×10^3

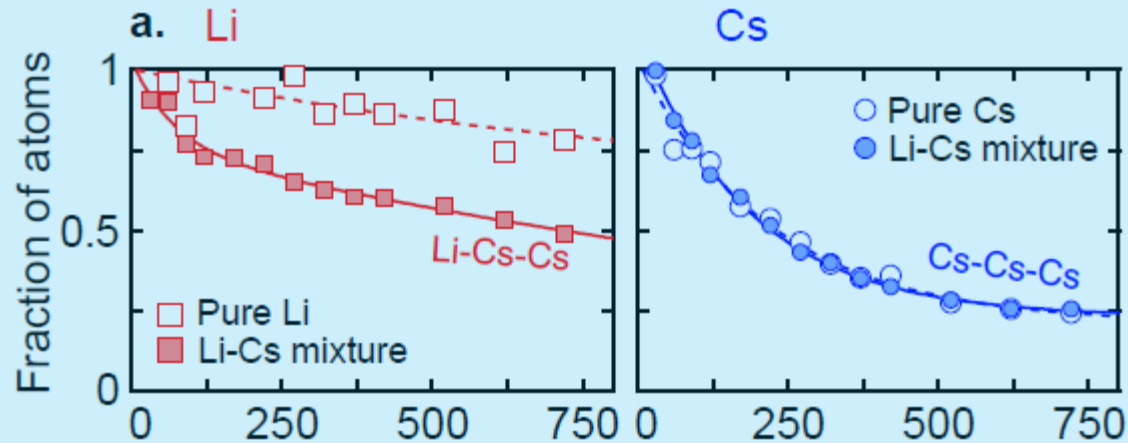


15 ms

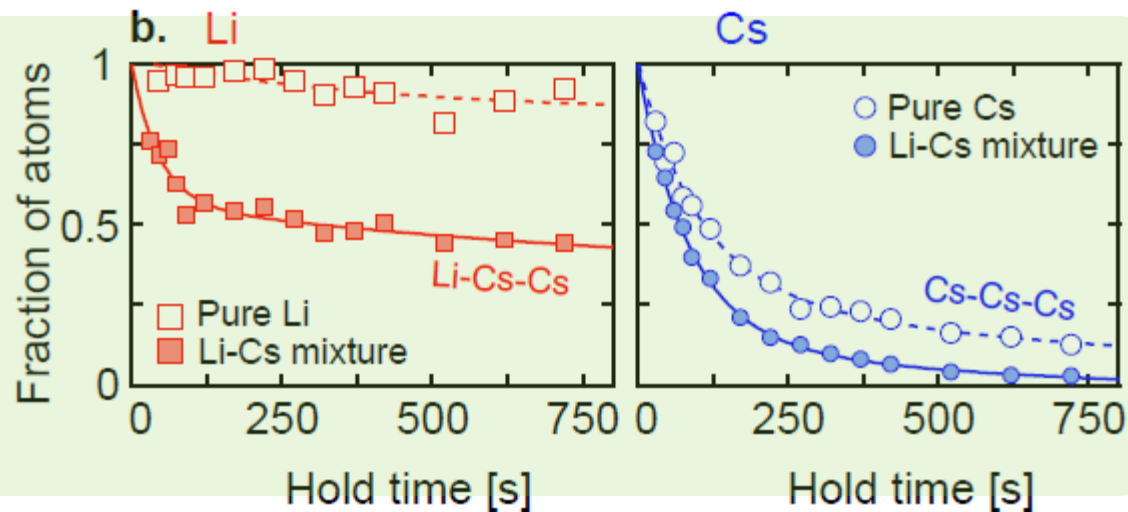
30×10^3

Trap loss measurement: Cs+Cs+Cs vs. Li+Cs+Cs

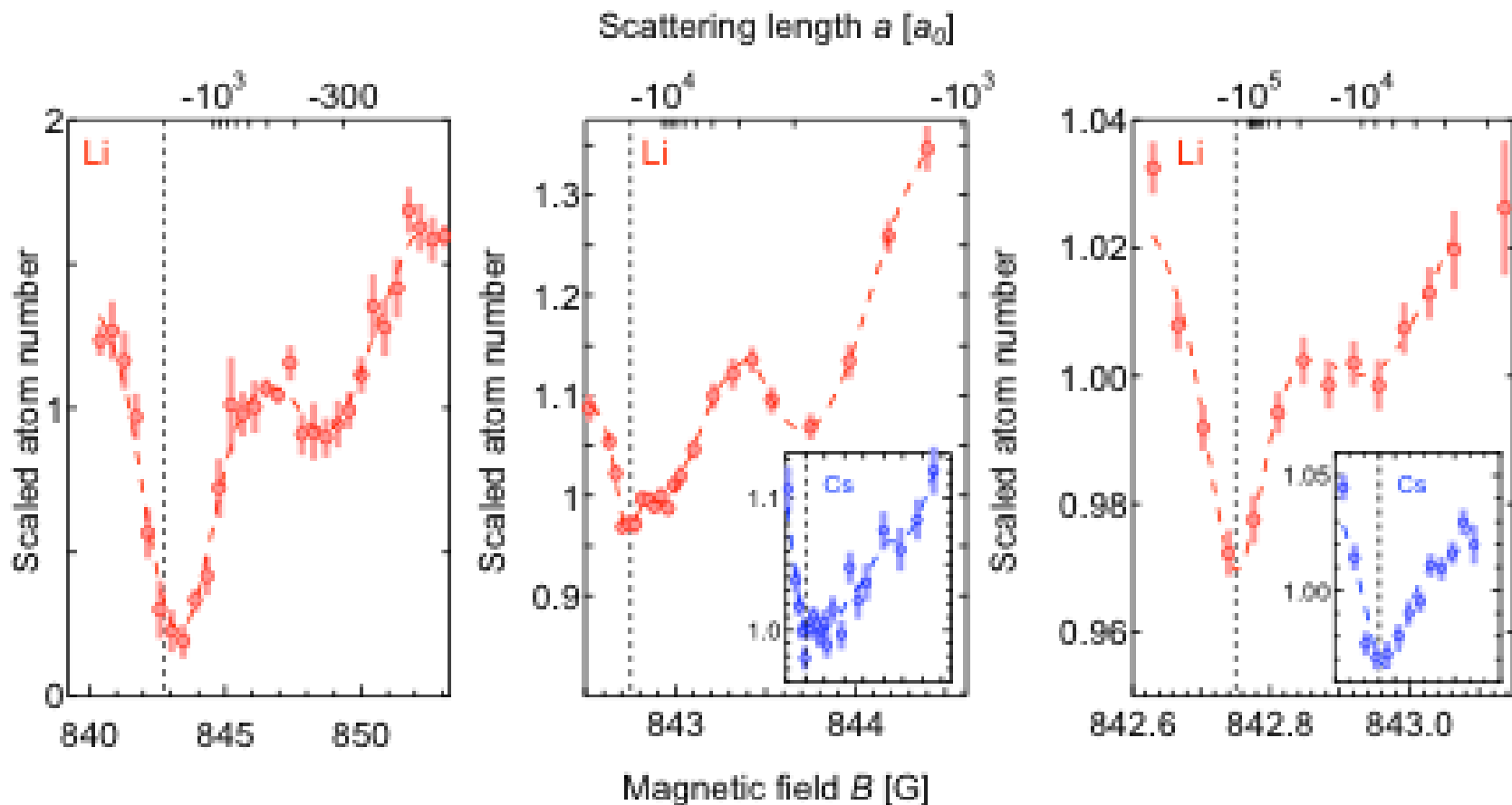
5G away
($-300 a_0$)



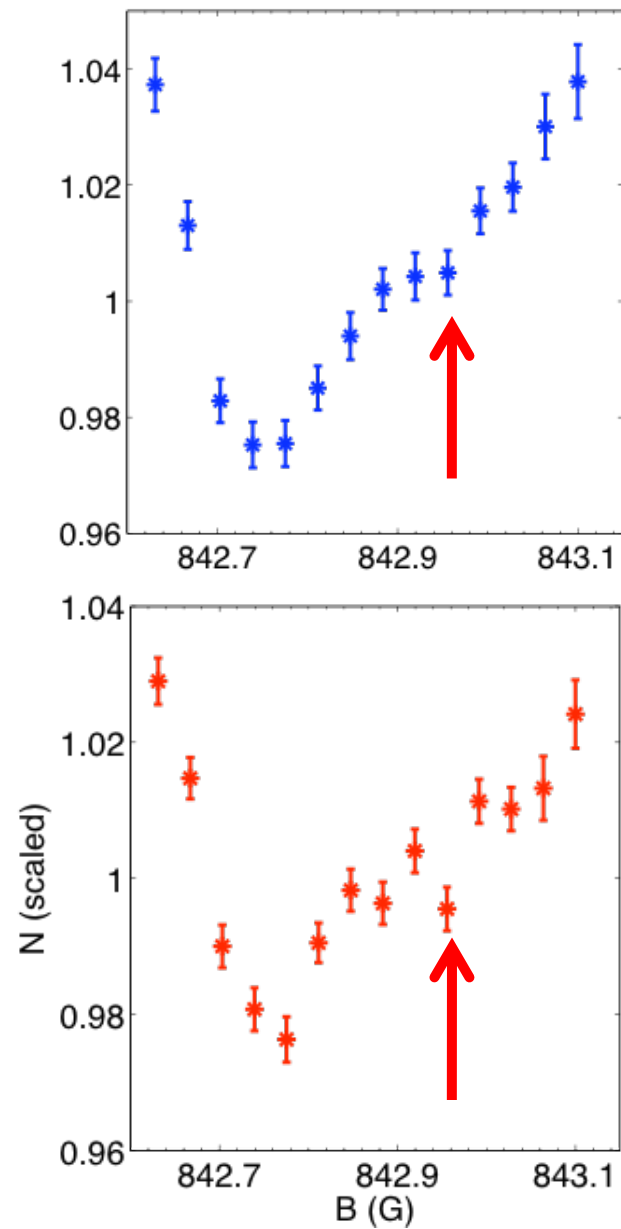
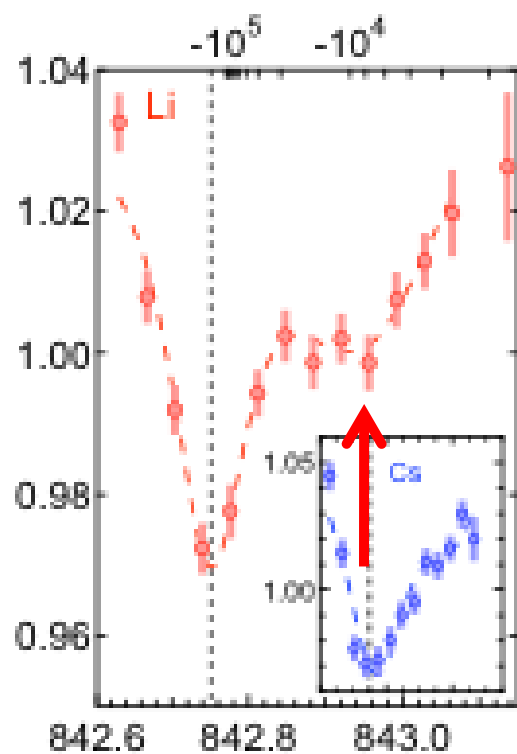
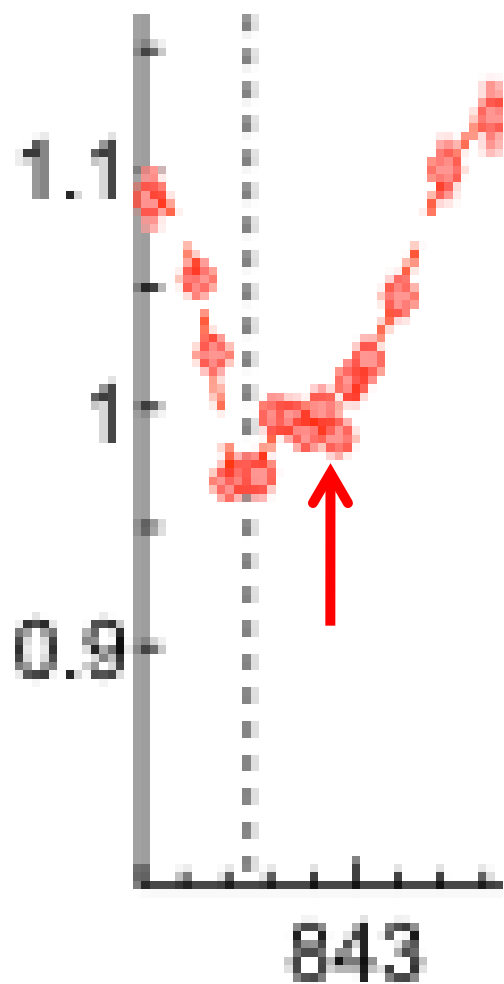
Feshbach
resonance



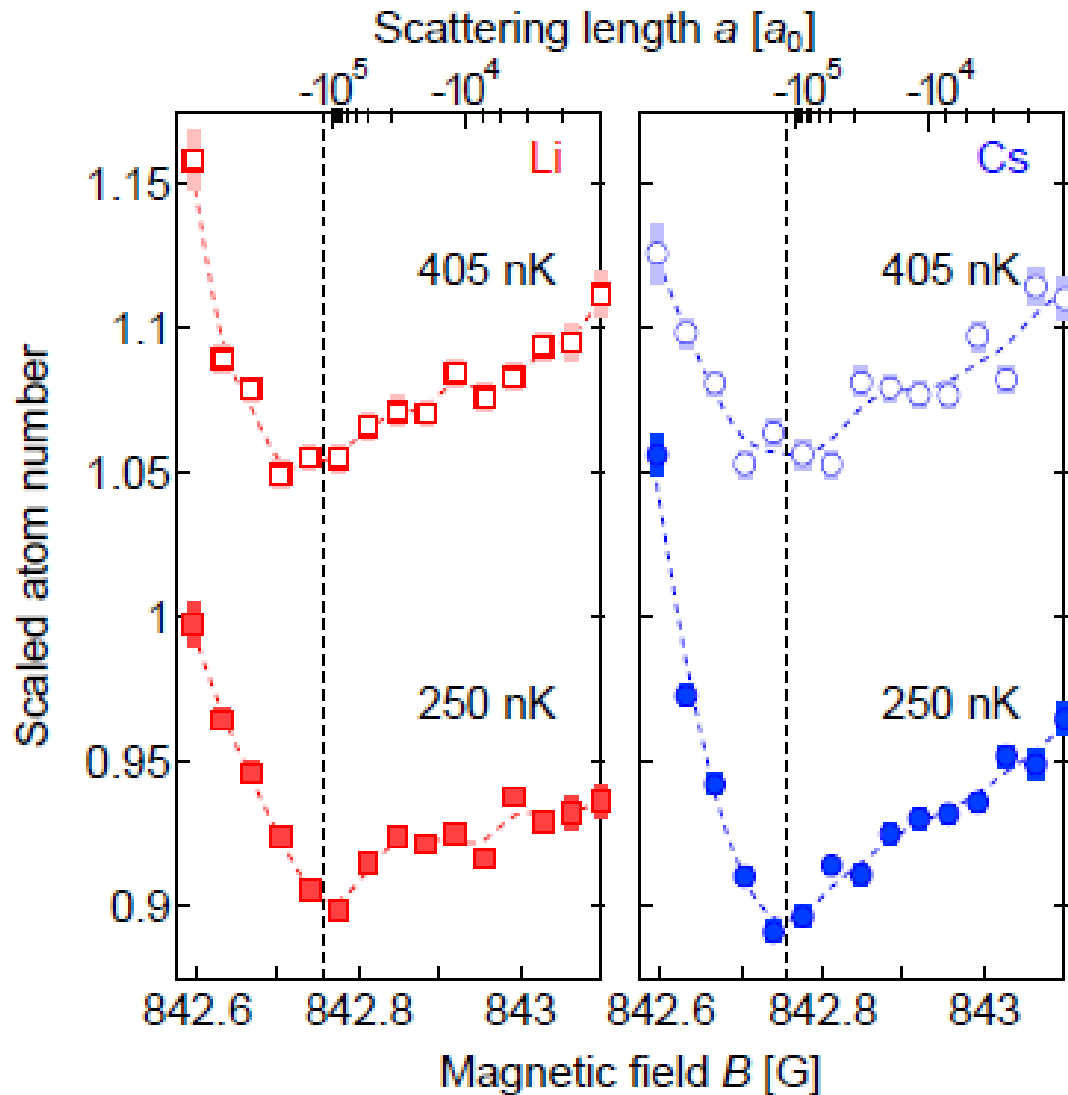
Trap loss measurement: Field scans



Finest Scans



Temperature dependence?



$\langle T \rangle = 405$ nK
209 averages

$\langle T \rangle = 250$ nK
328 averages

Result

- Three Efimov resonances:

First resonance: +5.6(2) G, $a_1 = -337(9)$ Bohr

Second resonance: +1.07(2) G, $a_2 = -1650(30)$ Bohr

Third resonance: +0.22(4) G, $a_3 = -7900(1400)$ Bohr

Feshbach: 842.75(1)G

- Scaling ratio: $a_1 : a_2 = 1 : 4.90(16)$

$$a_2 : a_3 = 1 : 4.79(87)$$

Weighted ratio: 4.85(44)

Theory: 4.88 D’Incao and B.Esry, PRA (2006)

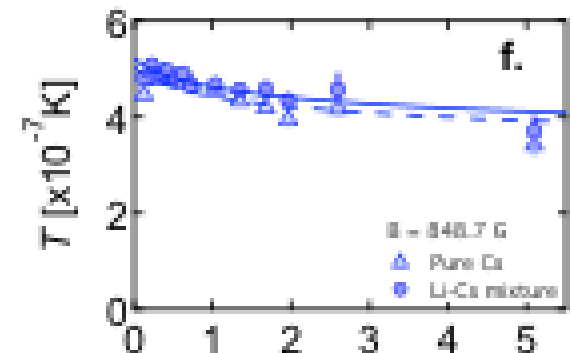
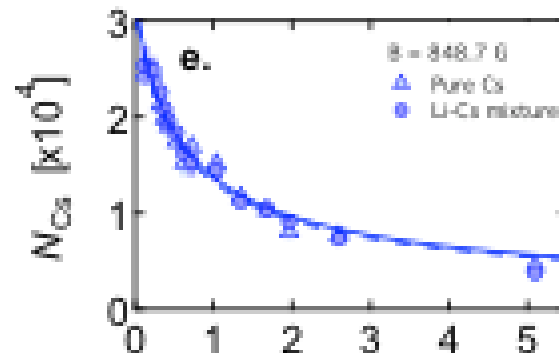
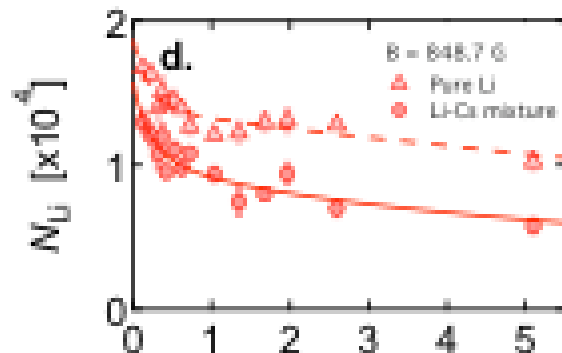
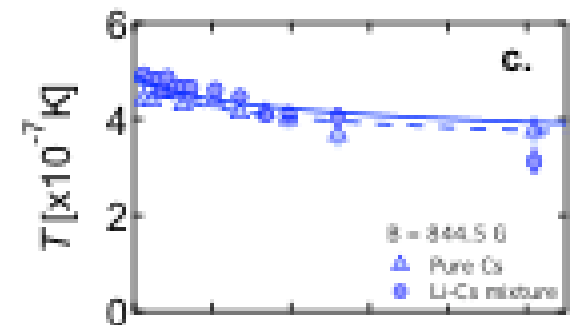
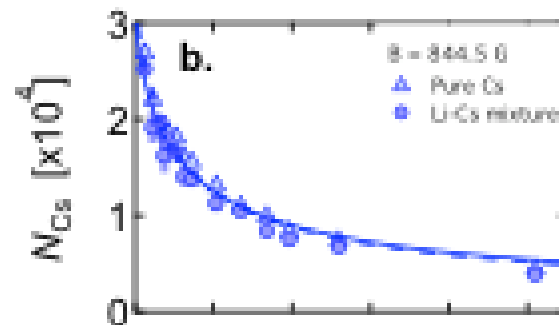
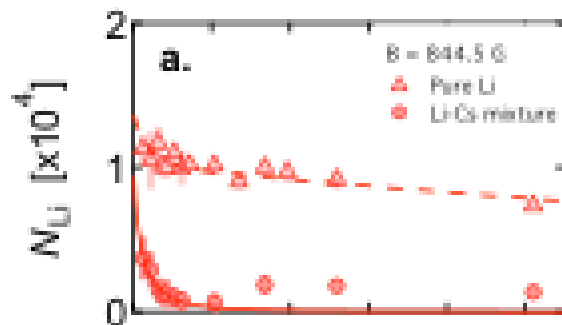
Systematics (Thanks to R. Grimm and C. Salomon)

- Finite temperature shift: $< 8\%$ in $0 \sim 1\mu\text{K}$ (Y. Wang)
- Feshbach resonance position: < 10 mG (Y. Wang)
- Finite size: $E_3 = 500\text{nK}$, Cs trap freq. = 4 nK, LiCs freq. = 4.5 nK

Trap independent coefficients (a best guess)

$$\begin{aligned} \frac{dN_{\text{Li}}}{dt} &= -K_3 X(T) N_{\text{Cs}}^2 N_{\text{Li}} - A N_{\text{Li}} - B N_{\text{Li}} e^{-1/C} \\ \frac{dN_{\text{Cs}}}{dt} &= -2K'_3 X(T) N_{\text{Cs}}^2 N_{\text{Li}} - D N_{\text{Cs}}^3 T^{-3} \\ \frac{dT_{\text{Cs}}}{dt} &= -F N_{\text{Cs}}^2 T_{\text{Cs}}^{-3/2}, \end{aligned}$$

$$X(T) N_{\text{Cs}}^2 N_{\text{Li}} = \int [n_{\text{Cs}}(\mathbf{x}, T)]^2 [n_{\text{Li}}(\mathbf{x}, T)] d^3\mathbf{x},$$

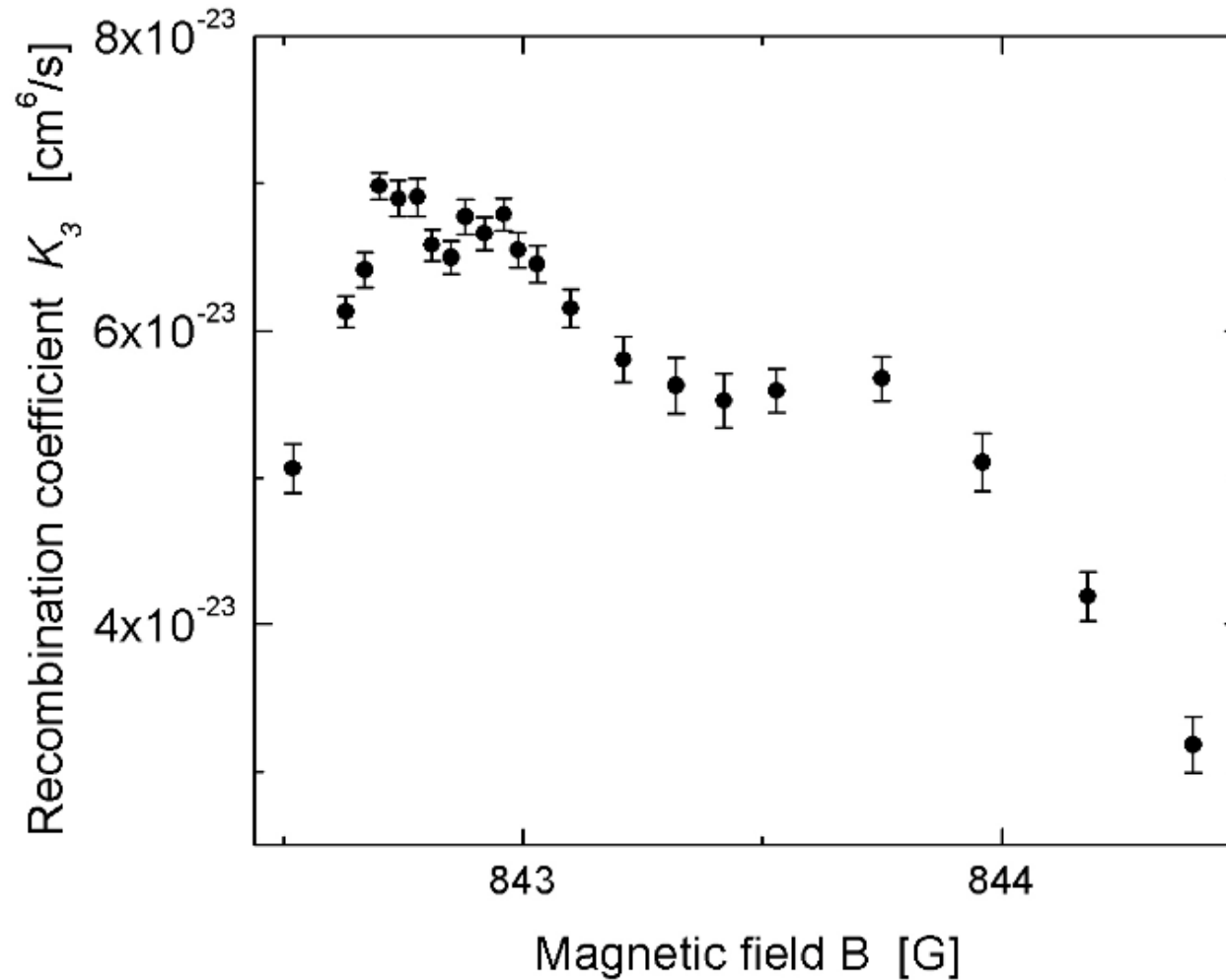


Time [s]

Time [s]

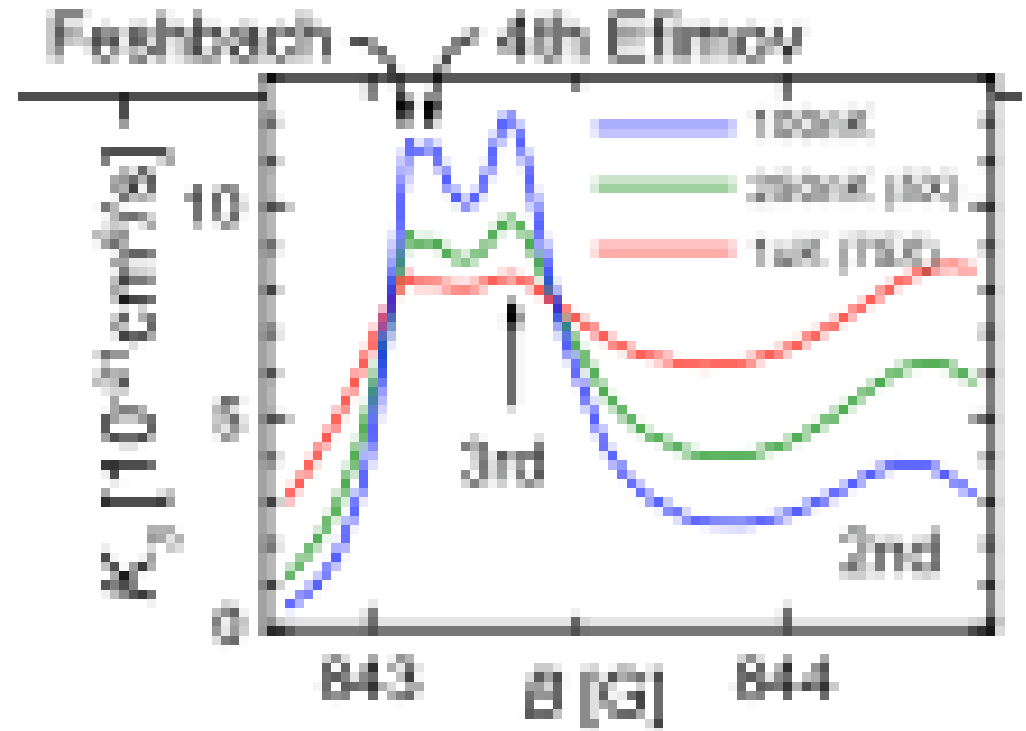
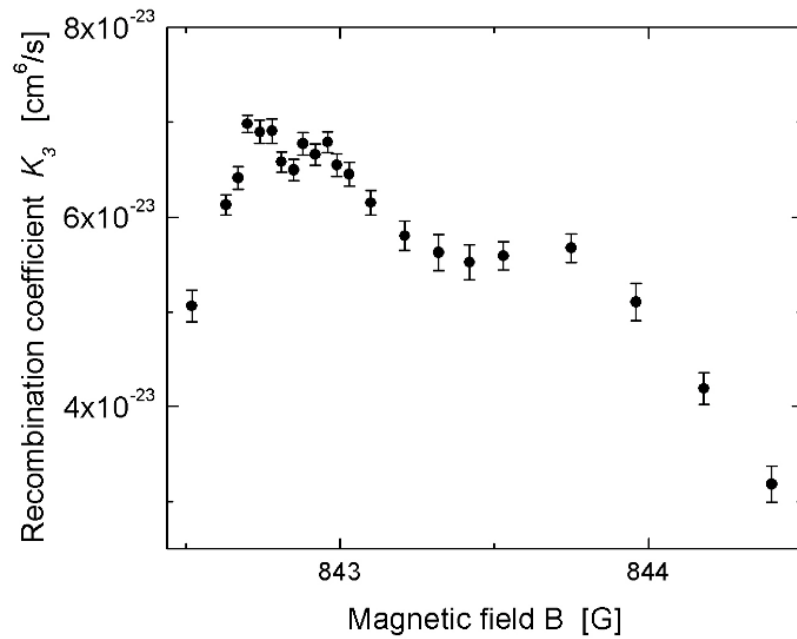
Time [s]

Trap independent coefficients (a best guess)

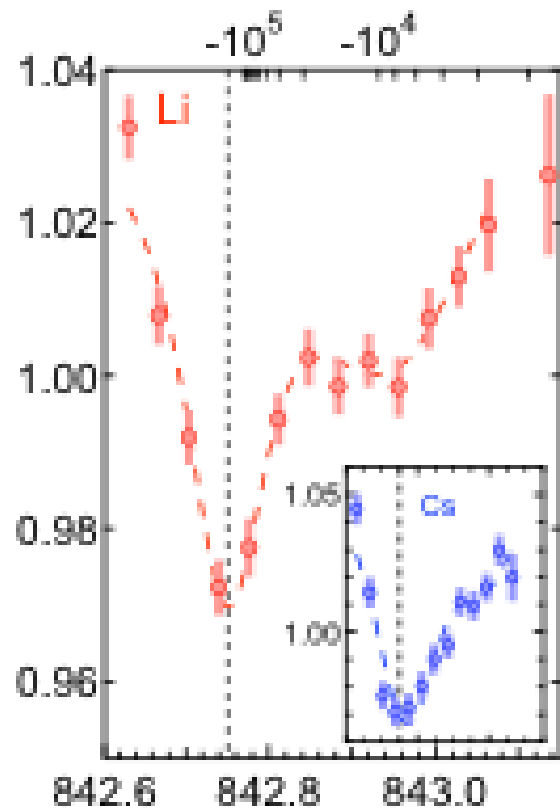


Theory:

(Yujun Wang, NIST and Kansas State)



Scattering length determination



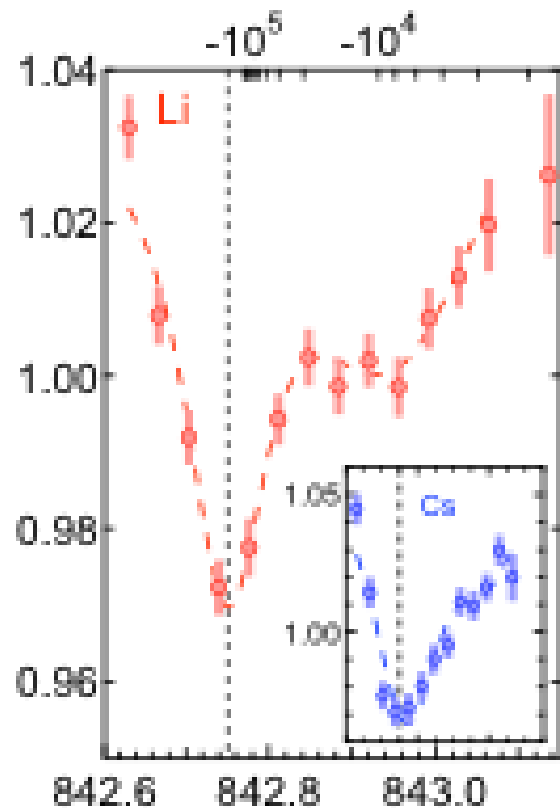
FB: 842.75(1) G
E3: +0.22(4) G
E2: +1.07(2) G
E1: +5.60(20) G

Within the resonance width,
scattering length $\sim 1/(B-B_0)$

$$(E1 - FB) / (E2 - FB) : 5.2(2)$$

$$(E2 - FB) / (E3 - FB) : 4.9(9)$$

Scattering length determination



FB : E4?	842.75(1)	G
E3:	+0.22(4)	G
E2:	+1.07(2)	G
E1:	+5.60(20)	G

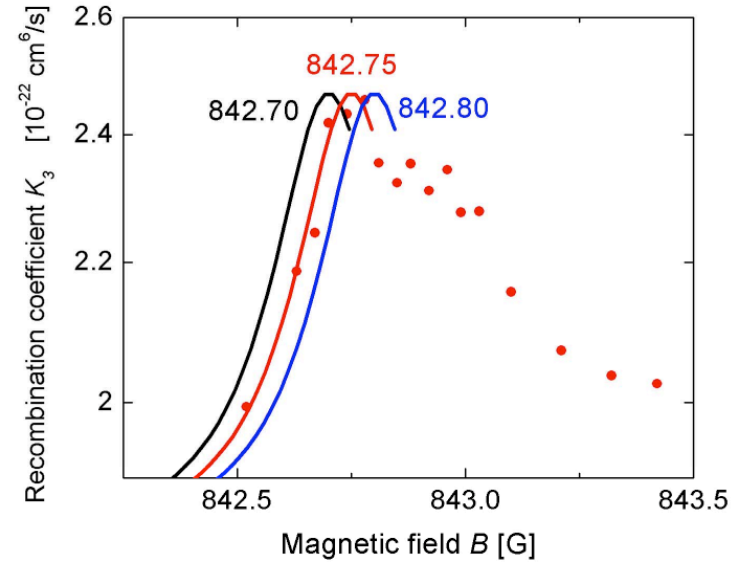
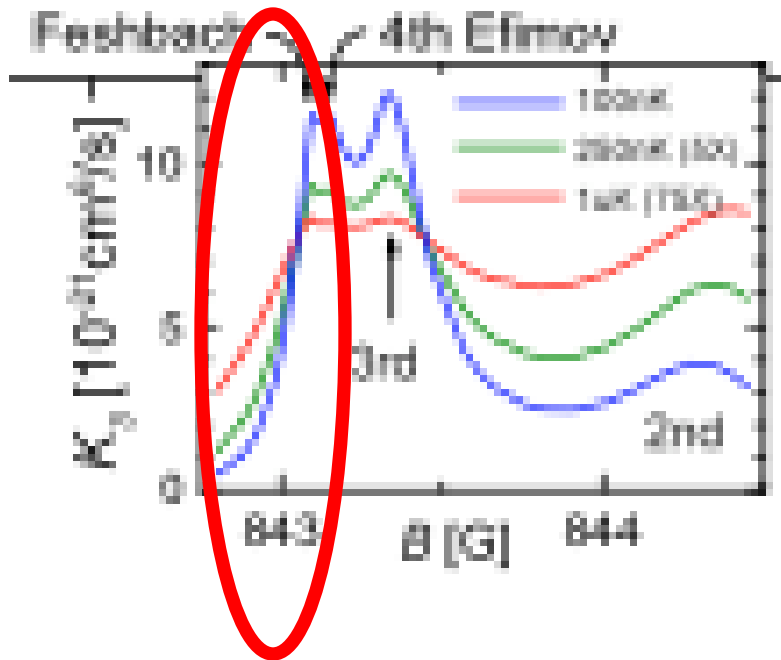
Within the resonance width,
scattering length $\sim 1/(B-B_0)$

$$(E1 - E2) / (E2 - E3) : 5.3(4)$$

$$(E2 - E3) / (E3 - E4) : 3.9(7)$$

Back to Theory:

(Yujun Wang, NIST and Kansas State)



$$K_3 \propto \text{const} + \frac{1}{(B - B_0)^2 + \Delta B^2}$$

The Crew

PI



Cheng Chin



Eric Hazlett

Postdocs



Shih-Kuang
Tung

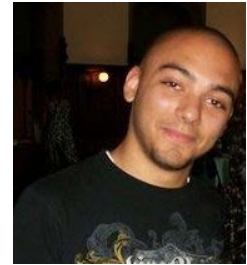


Karina
Jiménez-García

Undergrads



Nicholas
Kowalski



Dylan
O.A.B.
Sabulsky

Grads



Logan
Clark



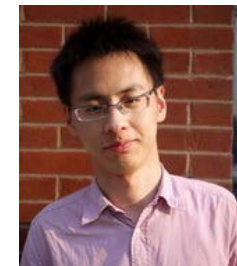
Gustaf
Downs



Lei Fang

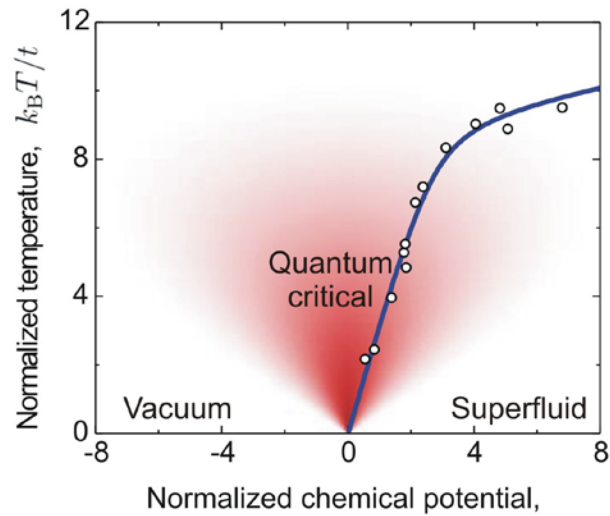


Jacob
Johansen



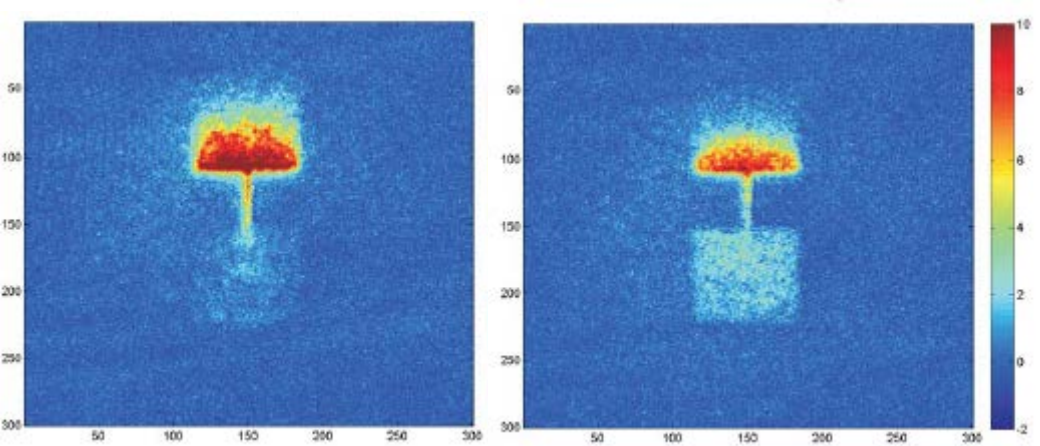
Li-Chung
Ha

Quantum criticality



Science 335, 1070 (2012)

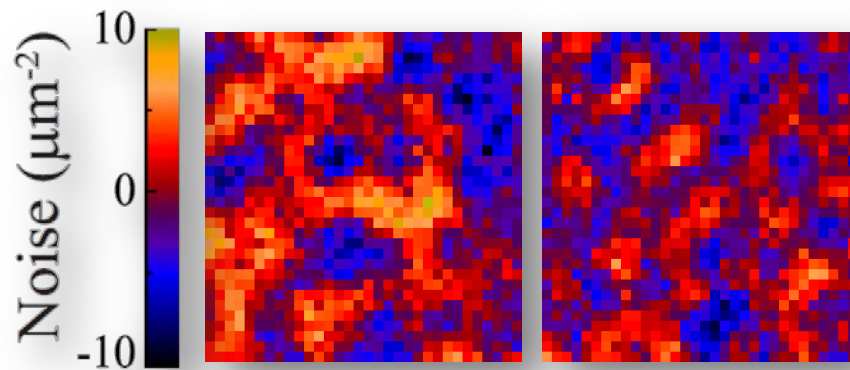
Transport



Postdoc positions available

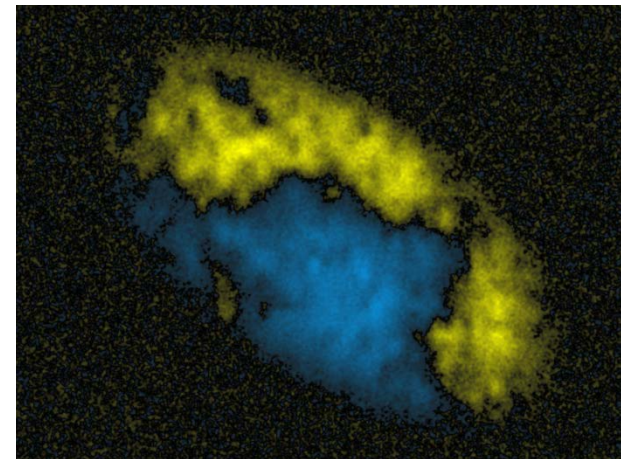
Arbitrary potential

Sakharov oscillations



Science 341, 1213 (2013)

Postdoc positions available



Nature Physics 9, 769 (2013)

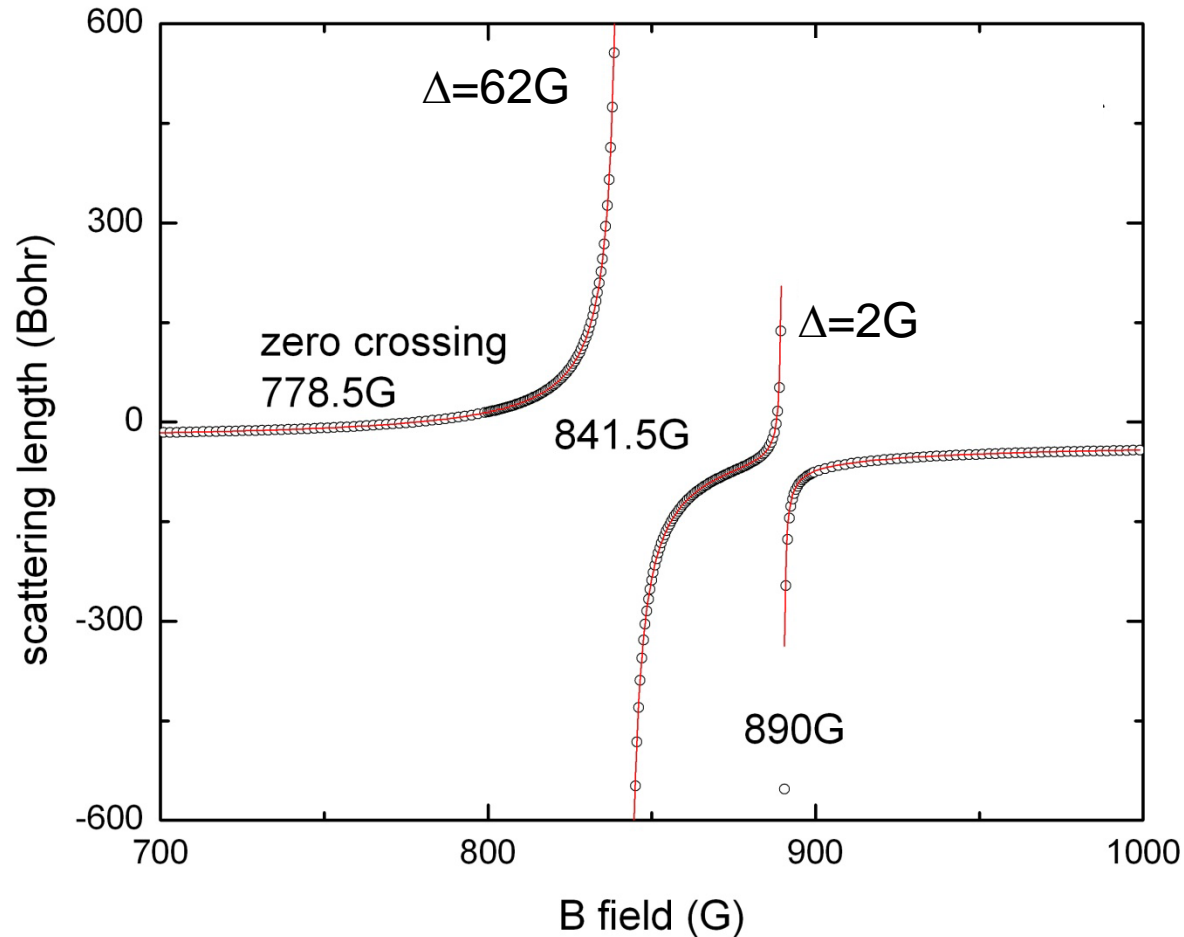
Which resonance should we choose?

TABLE I. Scattering length dependence for three-body collision rates in BBX and XYZ systems. For BBX systems both a_{BX} and a_{BB} scattering lengths are resonant, while for XYZ systems only a_{XY} and a_{XZ} are resonant. The notation $|a|$ indicates $a < 0$, and no entry indicates that the associated process is not possible. Expressions for M and P are given in Eqs. (2)–(5).

	$a_{BX} \ll a_{BB}$	$a_{BX} \gg a_{BB}$	$ a_{BX} \gg a_{BB}$	$ a_{BX} \ll a_{BB}$
$BB^* + X \rightarrow BB^* + X$...	$P_{s_0}^*(\frac{a_{BX}}{a_{BB}})M_{s_0}(\frac{a_{BB}}{r_0})a_{BX}$
$\rightarrow BB + X, BX + B$	$P_{s_0}(\frac{a_{BX}}{r_0})a_{BX}$	$P_{s_0}^*(\frac{a_{BX}}{a_{BB}})a_{BX}$
$BB^* + X \rightarrow BX^* + B$	$M_{s_0}(\frac{a_{BB}}{r_0})a_{BX}^2/a_{BB}$
$\rightarrow BB + X, BX + B$	a_{BX}^2/a_{BB}	$P_{s_0}^*(\frac{a_{BB}}{r_0})a_{BB}$	$P_{s_0}(\frac{a_{BB}}{r_0})a_{BB}$	$P_{s_0}(\frac{a_{BB}}{r_0})a_{BX}^2/a_{BB}$
$B + B + X \rightarrow BX^* + B$	$M_{s_0}(\frac{a_{BX}}{r_0})a_{BX}^2a_{BB}^2$	$M_{s_0}^*(\frac{a_{BX}}{a_{BB}})a_{BX}^4$
$\rightarrow BB^* + X$	a_{BB}^4	$M_{s_0}(\frac{a_{BB}}{r_0})a_{BX}^4$	$P_{s_0}^*(\frac{a_{BX}}{a_{BB}})M_{s_0}(\frac{a_{BB}}{r_0})a_{BX}^4$	a_{BB}^4
$\rightarrow BB + X, BX + B$	$a_{BX}^2a_{BB}^2$	a_{BX}^4	$P_{s_0}(\frac{a_{BB}}{a_{BB}})a_{BX}^4$	$a_{BX}^2a_{BB}^2$
	$a_{BX} \ll a_{BB} $	$a_{BX} \gg a_{BB} $	$ a_{BX} \gg a_{BB} $	$ a_{BX} \ll a_{BB} $
$BB^* + X \rightarrow BB + X, BX + B$	$P_{s_0}(\frac{a_{BX}}{r_0})a_{BX}$	$P_{s_0}^*(\frac{a_{BB}}{r_0}, \frac{a_{BX}}{a_{BB}})a_{BX}$
$B + B + X \rightarrow BX^* + B$	$M_{s_0}(\frac{a_{BX}}{r_0})a_{BX}^2a_{BB}^2$	$M_{s_0}^*(\frac{a_{BB}}{r_0}, \frac{a_{BX}}{a_{BB}})a_{BX}^4$
$\rightarrow BB + X, BX + B$	$a_{BX}^2a_{BB}^2$	a_{BX}^4	$P_{s_0}^*(\frac{a_{BB}}{r_0}, \frac{a_{BX}}{a_{BB}})a_{BX}^4$	$P_{s_0}(\frac{a_{BB}}{r_0})a_{BX}^2a_{BB}^2$
	$a_{XY} \gg a_{XZ}$	$ a_{XY} \gg a_{XZ}$	$a_{XY} \gg a_{XZ} $	$ a_{XY} \gg a_{XZ} $
$XY^* + Z \rightarrow XZ^* + Y$	$M_{s_0}(\frac{a_{XZ}}{r_0})a_{XZ}^2/a_{XY}$
$\rightarrow XY + Z, XZ + Y, \dots$	a_{XZ}^2/a_{XY}	...	$P_{s_0}(\frac{a_{XZ}}{r_0})a_{XZ}^2/a_{XY}$...
$XZ^* + Y \rightarrow XY + Z, XZ + Y, \dots$	$P_{s_0}(\frac{a_{XZ}}{r_0})a_{XZ}$	$P_{s_0}(\frac{a_{XZ}}{r_0})a_{XZ}$
$X + Y + Z \rightarrow XY^* + Z$	a_{XY}^4	...	a_{XY}^4	...
$\rightarrow XZ^* + Y$	$M_{s_0}(\frac{a_{XZ}}{r_0})a_{XZ}^2a_{XY}^2$	$M_{s_0}(\frac{a_{XZ}}{r_0})a_{XZ}^2a_{XY}^2$
$\rightarrow XY + Z, XZ + Y, \dots$	$a_{XZ}^2a_{XY}^2$	$a_{XZ}^2a_{XY}^2$	$a_{XZ}^2a_{XY}^2$	$P_{s_0}(\frac{a_{XZ}}{r_0})a_{XZ}^2a_{XY}^2$

Choose the one with $a_{BB} < 0$. J. P. D’Incao and B. D. Esry, PRL 103 083202 (2009)

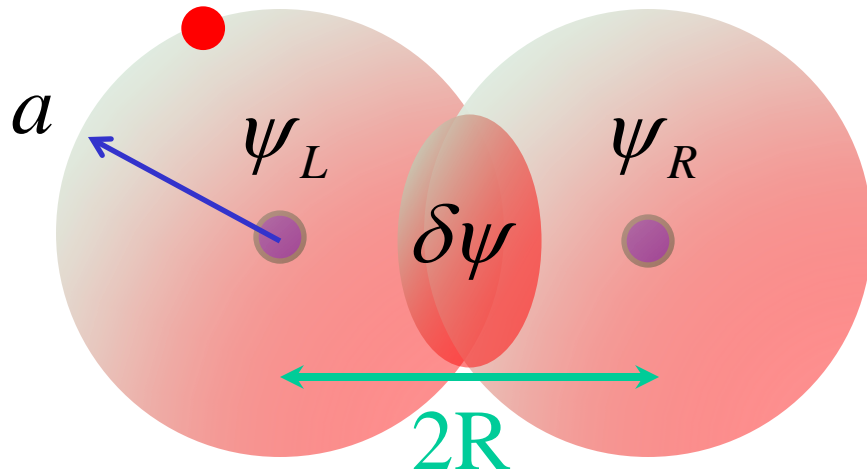
Broad (62G) ${}^6\text{Li}$ - ${}^{133}\text{Cs}$ Feshbach resonance



Theory: Y. Wang and P. Julienne (NIST)

S.K. Tung et al., PRA 87, 010702 (2013)

Picture of Efimov potential



$$\psi_{\pm}(r) = \psi_L(r) \pm \psi_R(r)$$

$$-\frac{\hbar^2}{2m} \dot{E}''(r) + V \dot{E}(r) = E \dot{E}(r) \Rightarrow \frac{\hbar^2}{2m} \frac{\delta \dot{E}(r)}{R^2} + V_{\text{eff}} \delta \dot{E}(r) = 0$$

$$\Rightarrow V_{\text{eff}} \propto -\frac{\hbar^2}{2mR^2} \quad \text{when } R < |a|$$

$$\approx 0 \quad \text{when } R > |a|$$

Geometric scaling of Efimov states

