

Microscopic origin and universality classes of the Efimov three-body parameter



Pascal Naidon



The University
of Tokyo

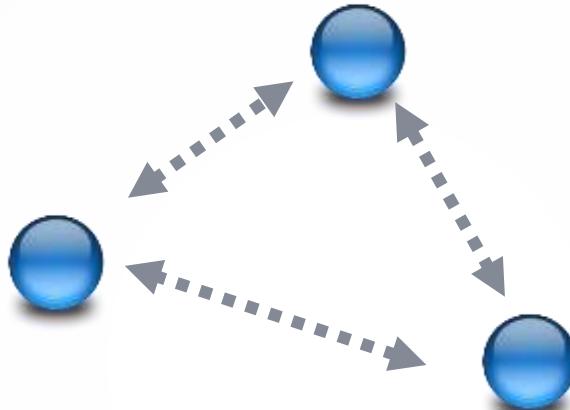


Shimpei Endo



Masahito Ueda

3 particles (bosons) with resonant two-body interactions



- Single-channel two-body interactions
- No three-body interaction

Summary

The **3-body parameter** is (mostly) determined by the **2-body correlation**.

Reason:

2-body correlation induces a **deformation** of the 3-body system.

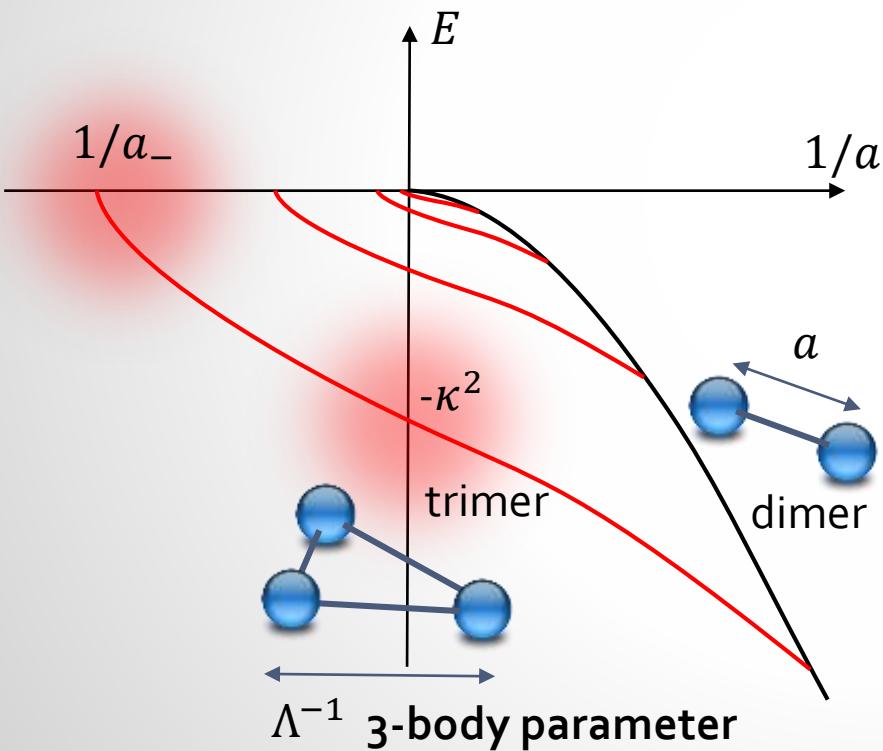
Consequences: the 3-body parameter

- is on the order of the **effective range**.
- has different universal values for **distinct classes** of interaction

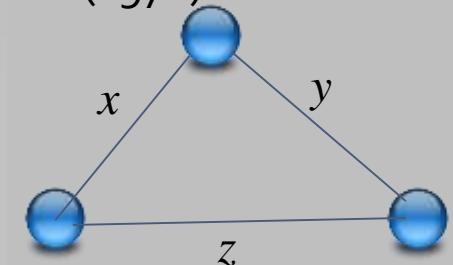
The Efimov 3-body parameter

Parameters describing particles at low energy

Scattering length a (2-body parameter)

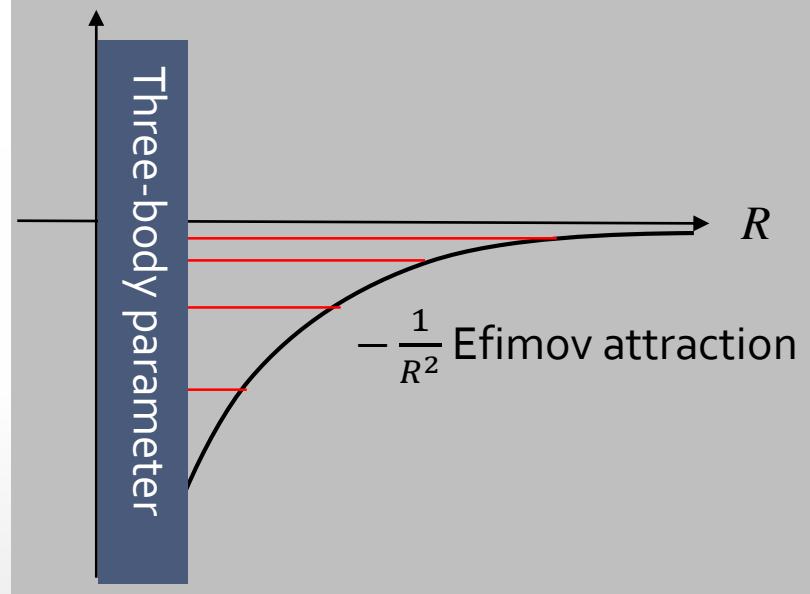


The Efimov effect (1970)



$$\text{Hyperradius } R^2 = \frac{2}{3}(x^2 + y^2 + z^2)$$

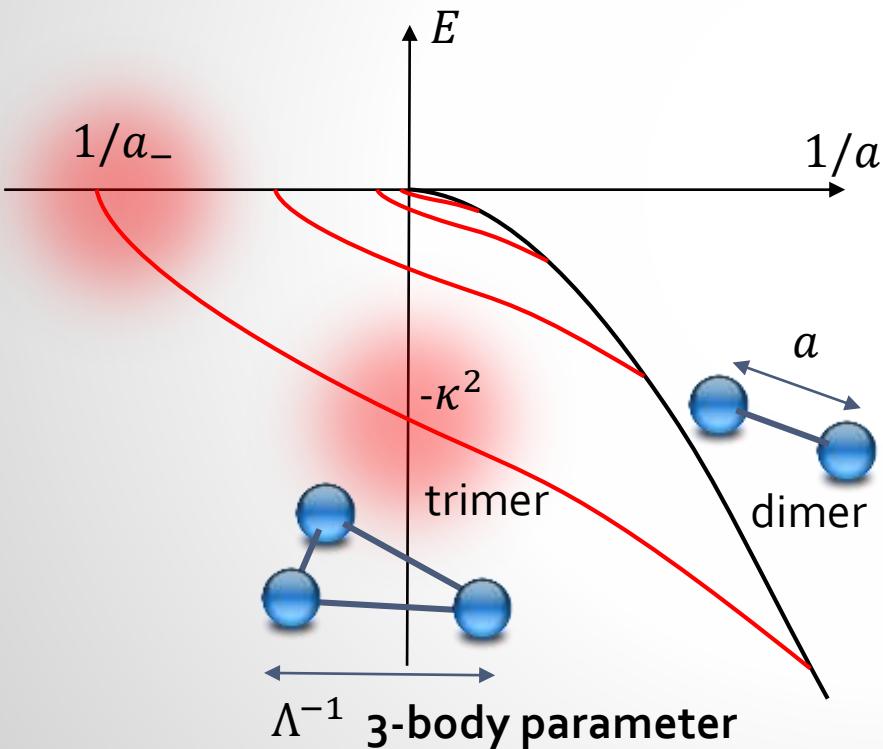
Zero-range condition with $a \rightarrow \infty$



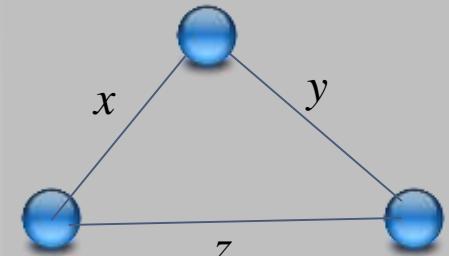
The Efimov 3-body parameter

Parameters describing particles at low energy

Scattering length a (2-body parameter)

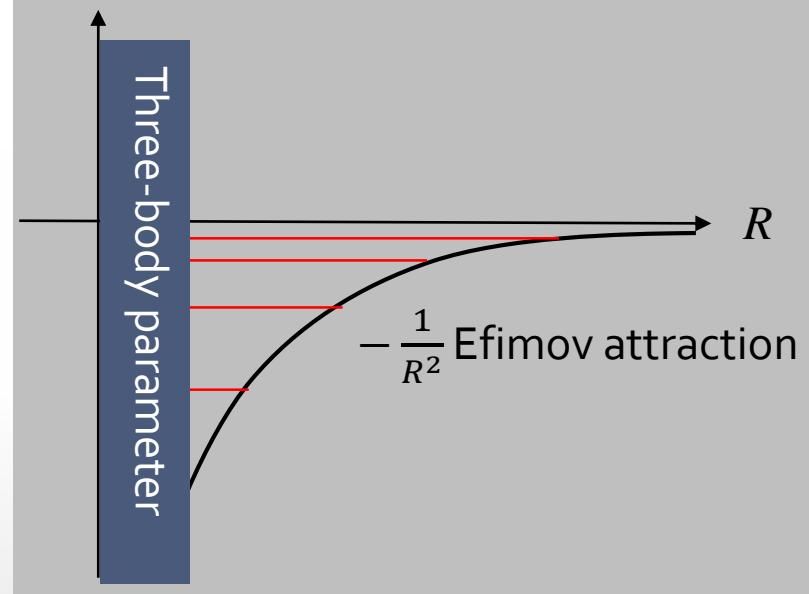


The Efimov effect



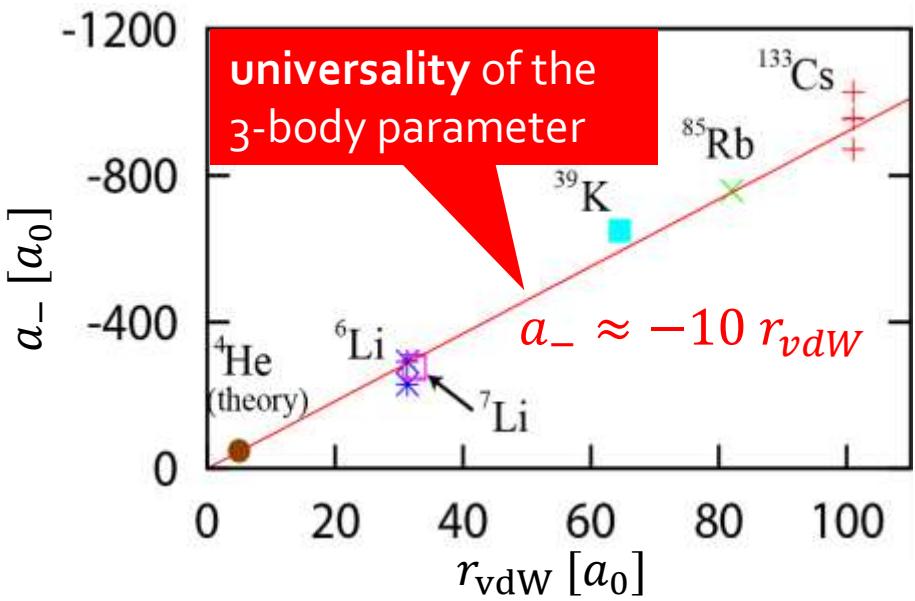
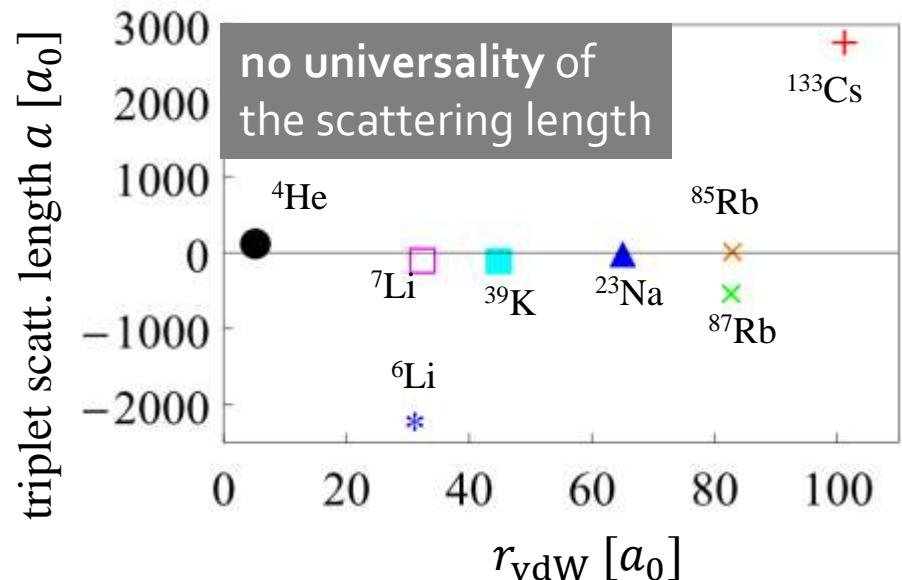
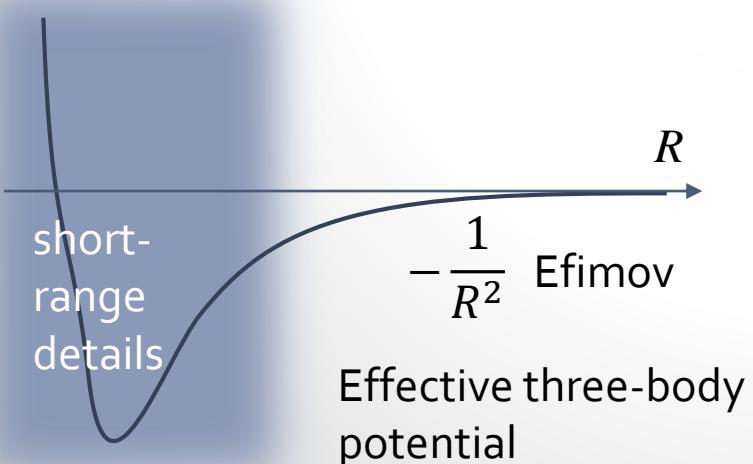
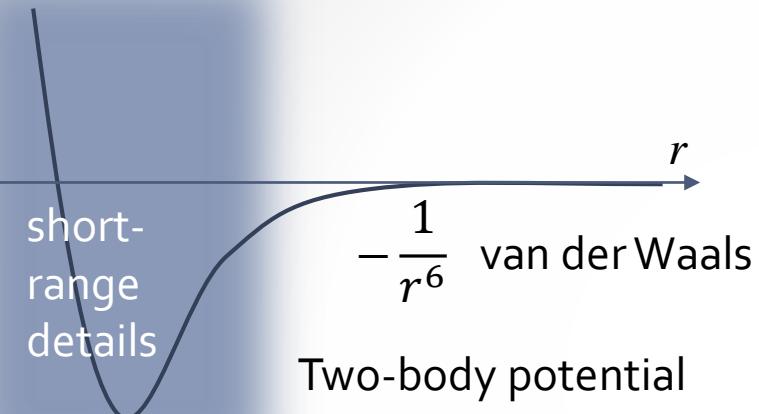
$$\text{Hyperradius } R^2 = \frac{2}{3}(x^2 + y^2 + z^2)$$

Zero-range condition with $a \rightarrow \infty$



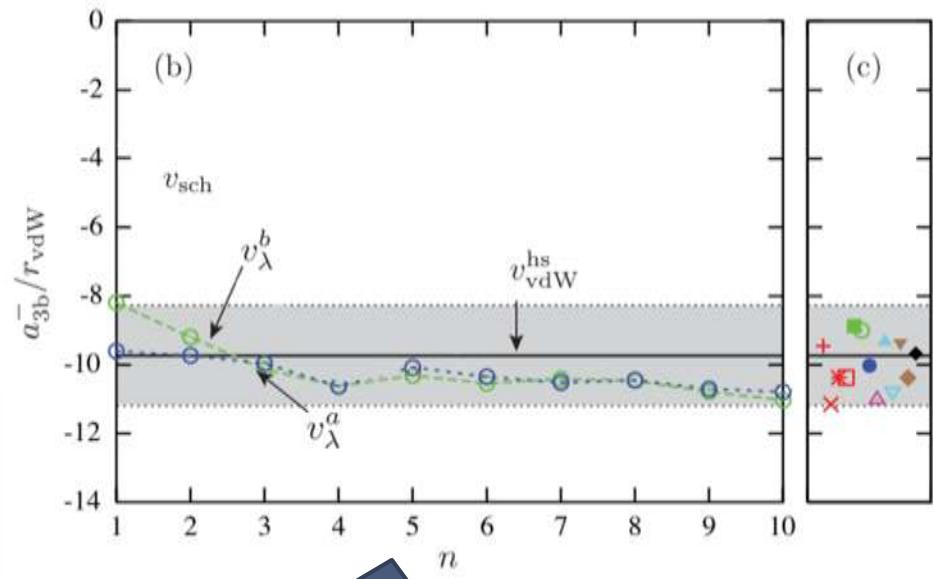
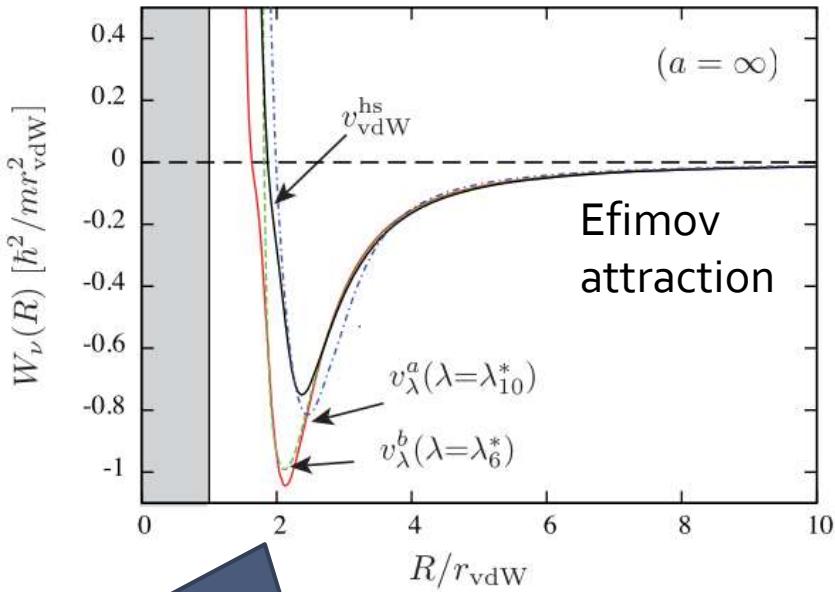
Universality for atoms

Microscopic determination?



Three-body with van der Waals interactions

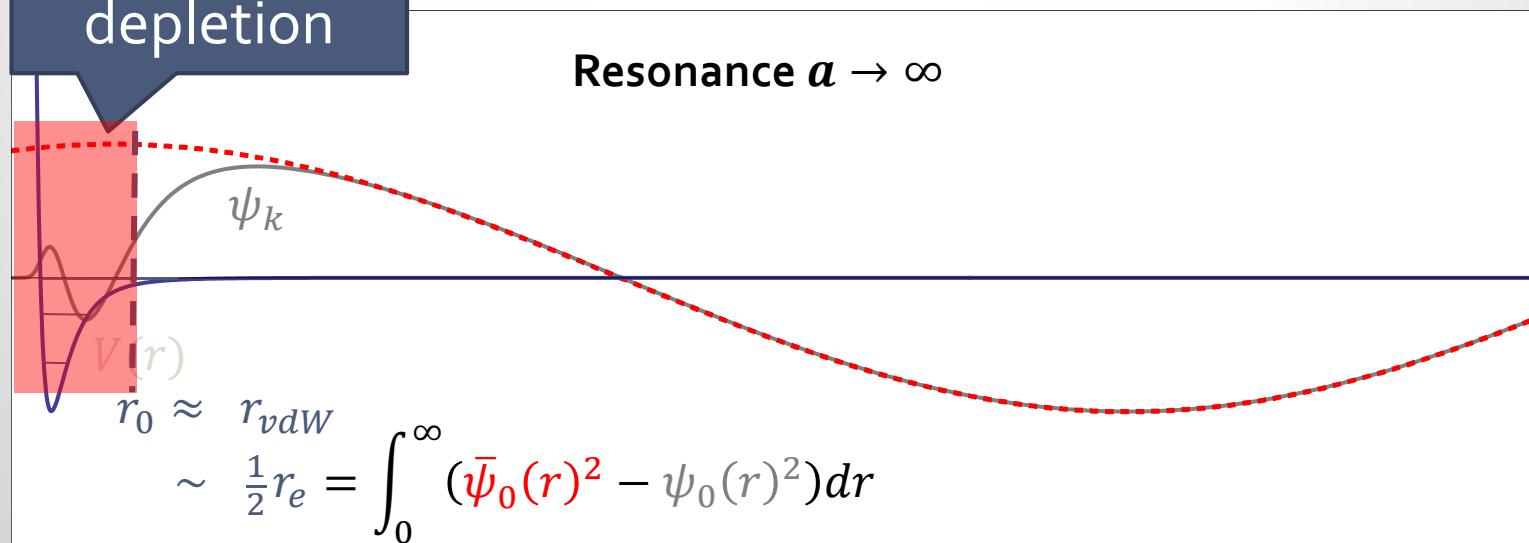
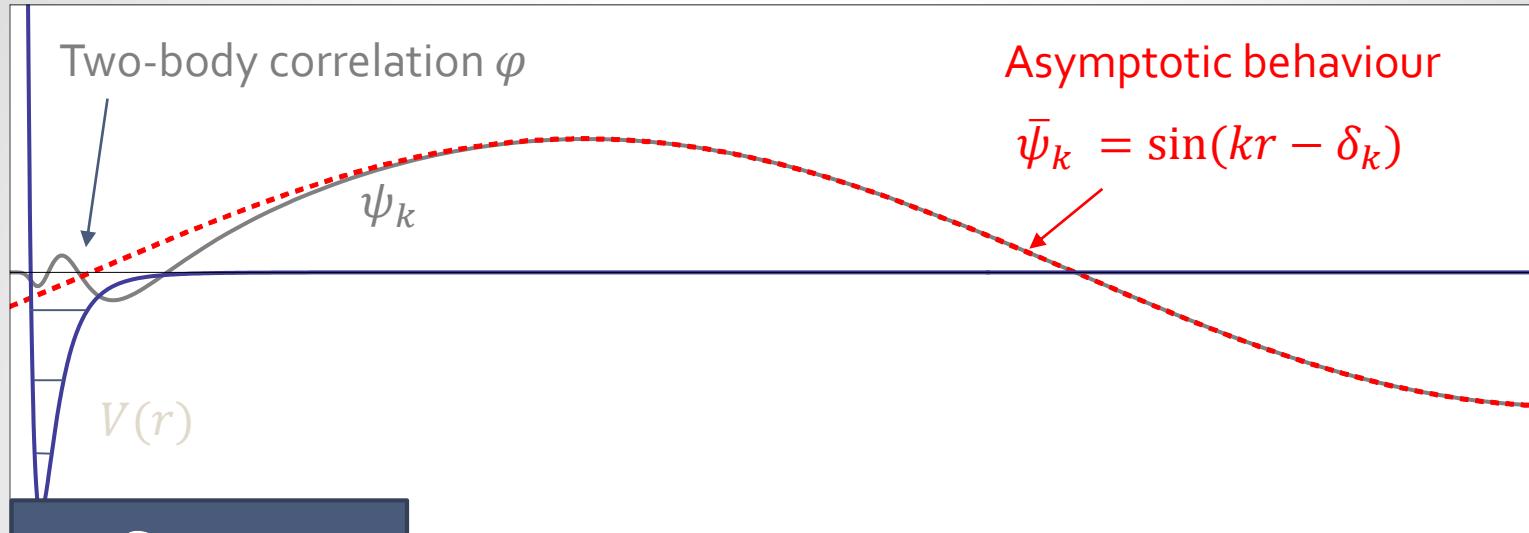
Phys. Rev. Lett. 108 263001 (2012)
J. Wang, J. D'Incao, B. Esry, C. Greene



Three-body repulsion at
 $R \approx 2r_{vdW}$

$$a_- \approx -11 r_{vdW}$$

Interpretation: two-body correlation

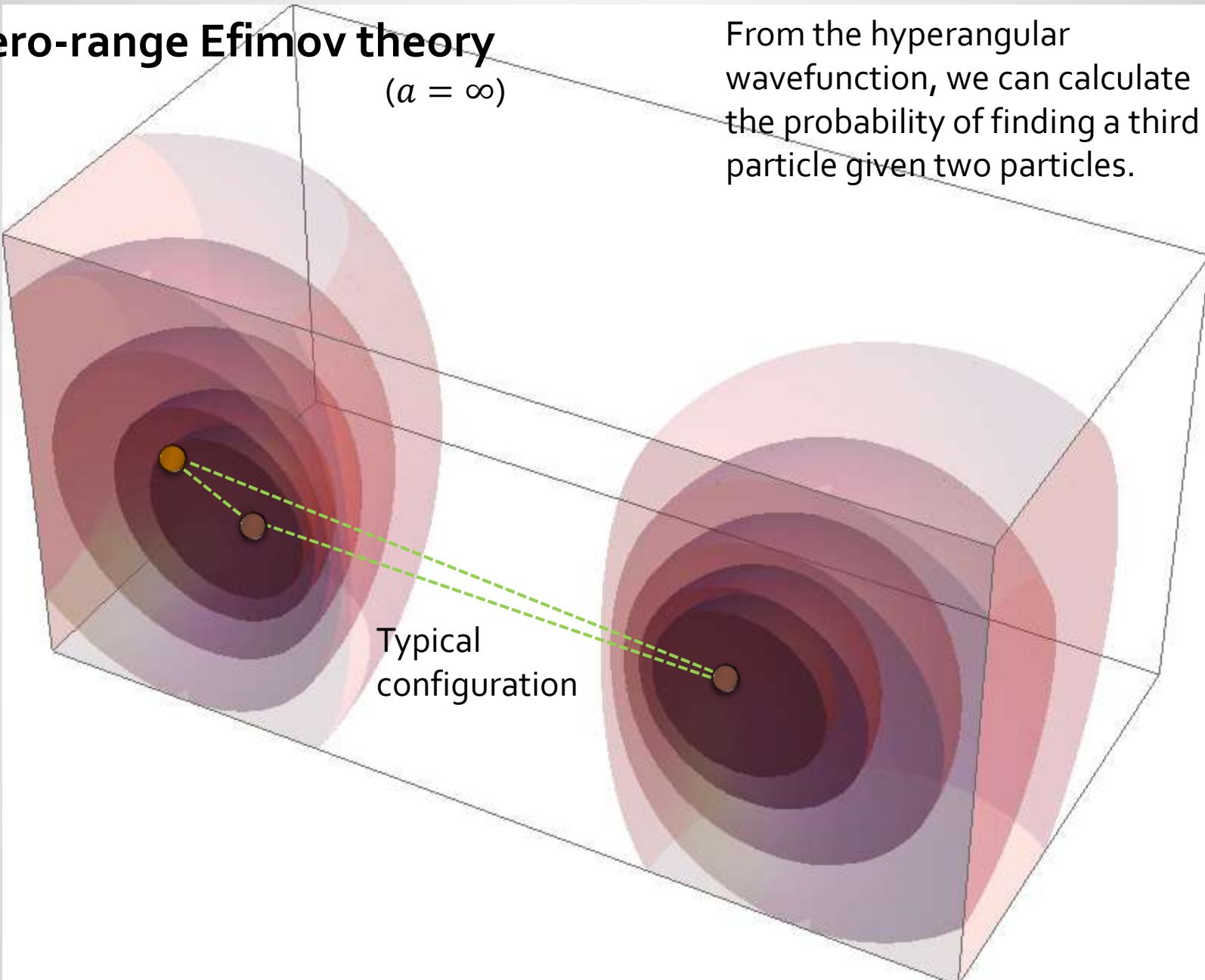


Zero-range Efimov theory

$(a = \infty)$

From the hyperangular wavefunction, we can calculate the probability of finding a third particle given two particles.

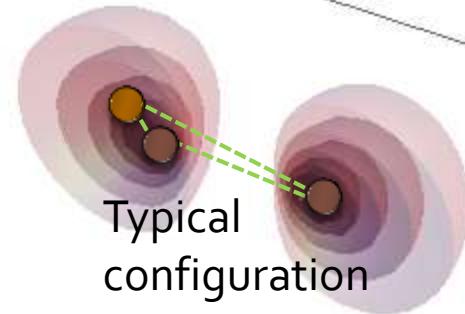
Typical configuration



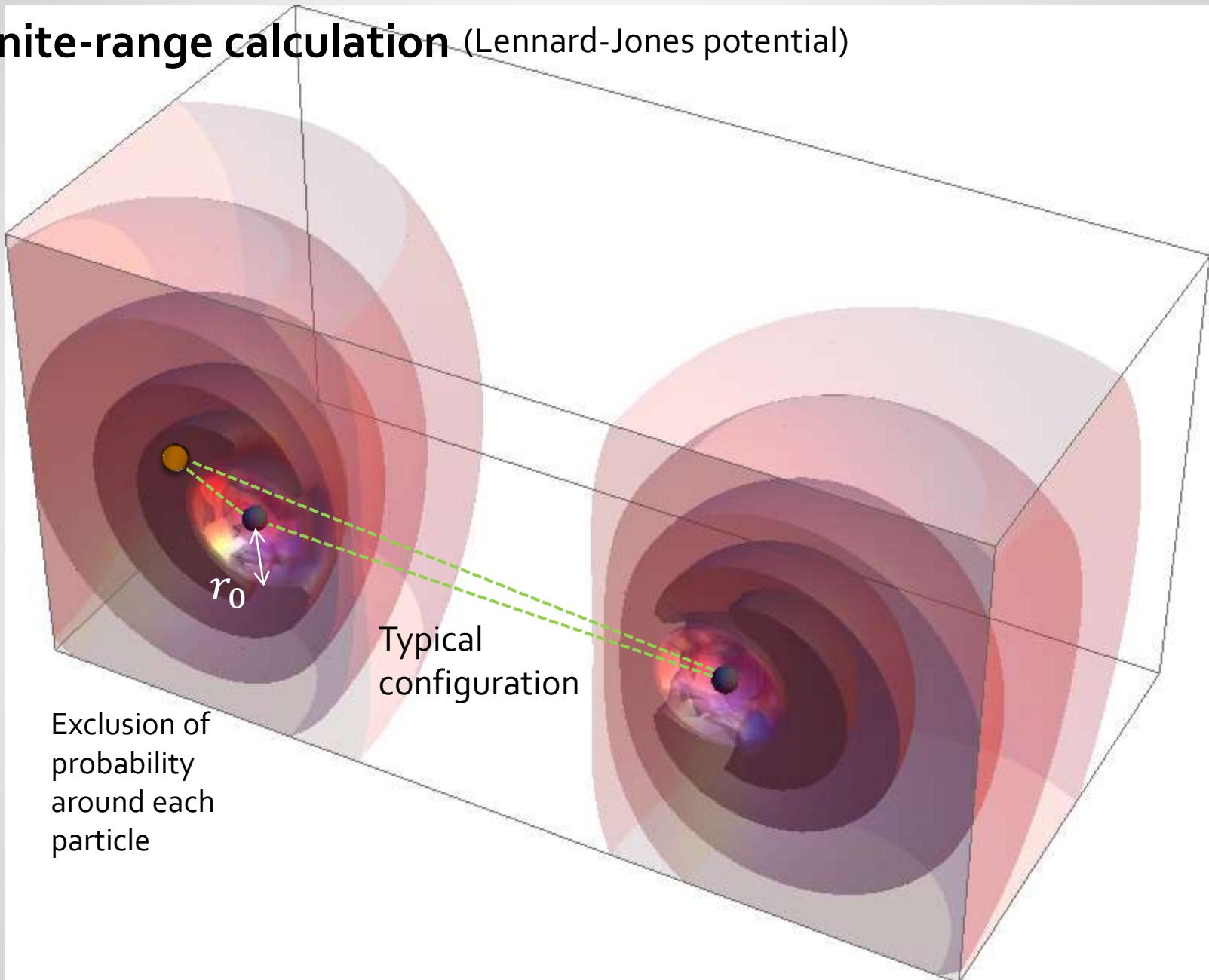
Zero-range Efimov theory

$(a = \infty)$

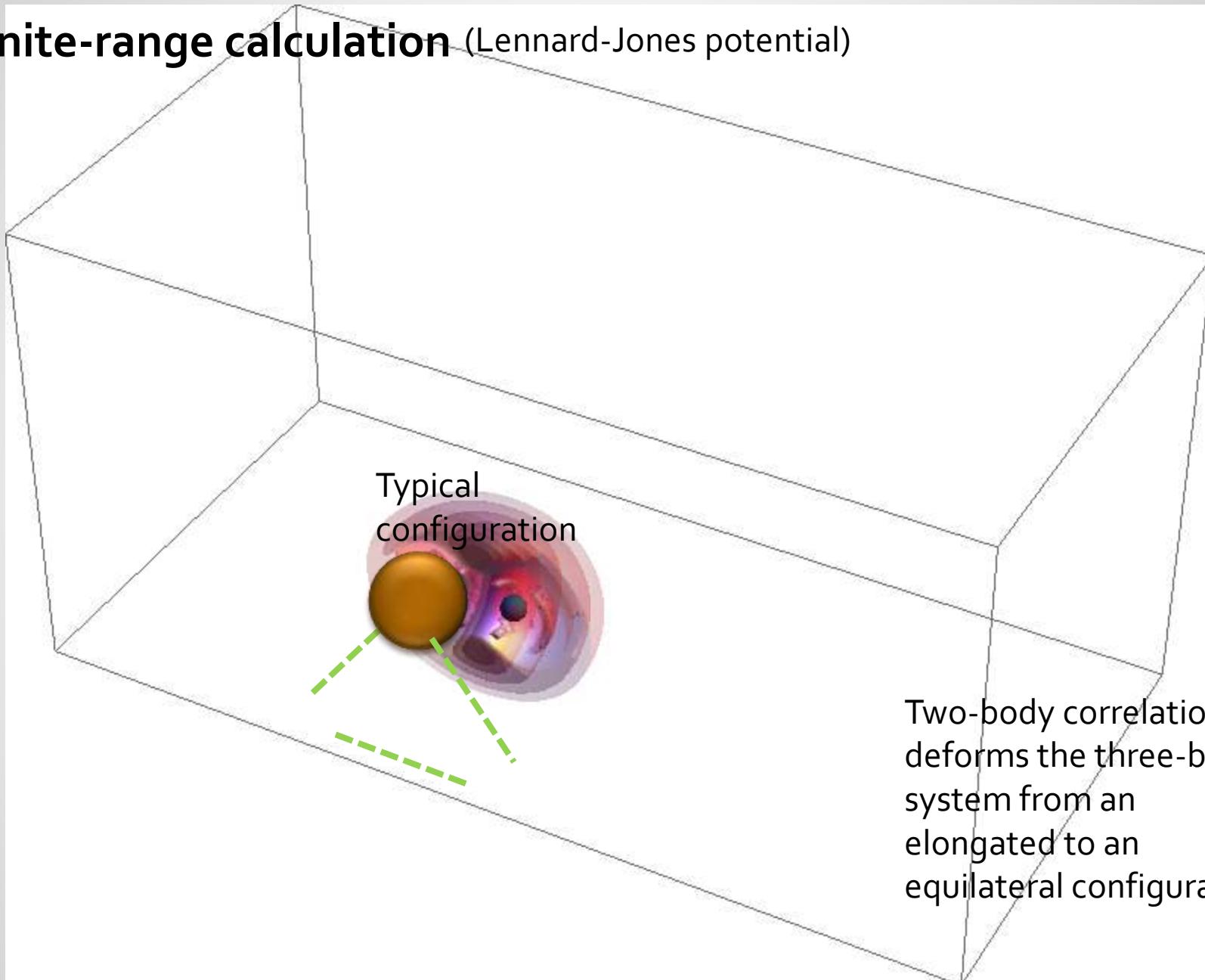
The hyperangular wavefunction is independent of the hyperradius: **the shape of the system does not change with its size.**



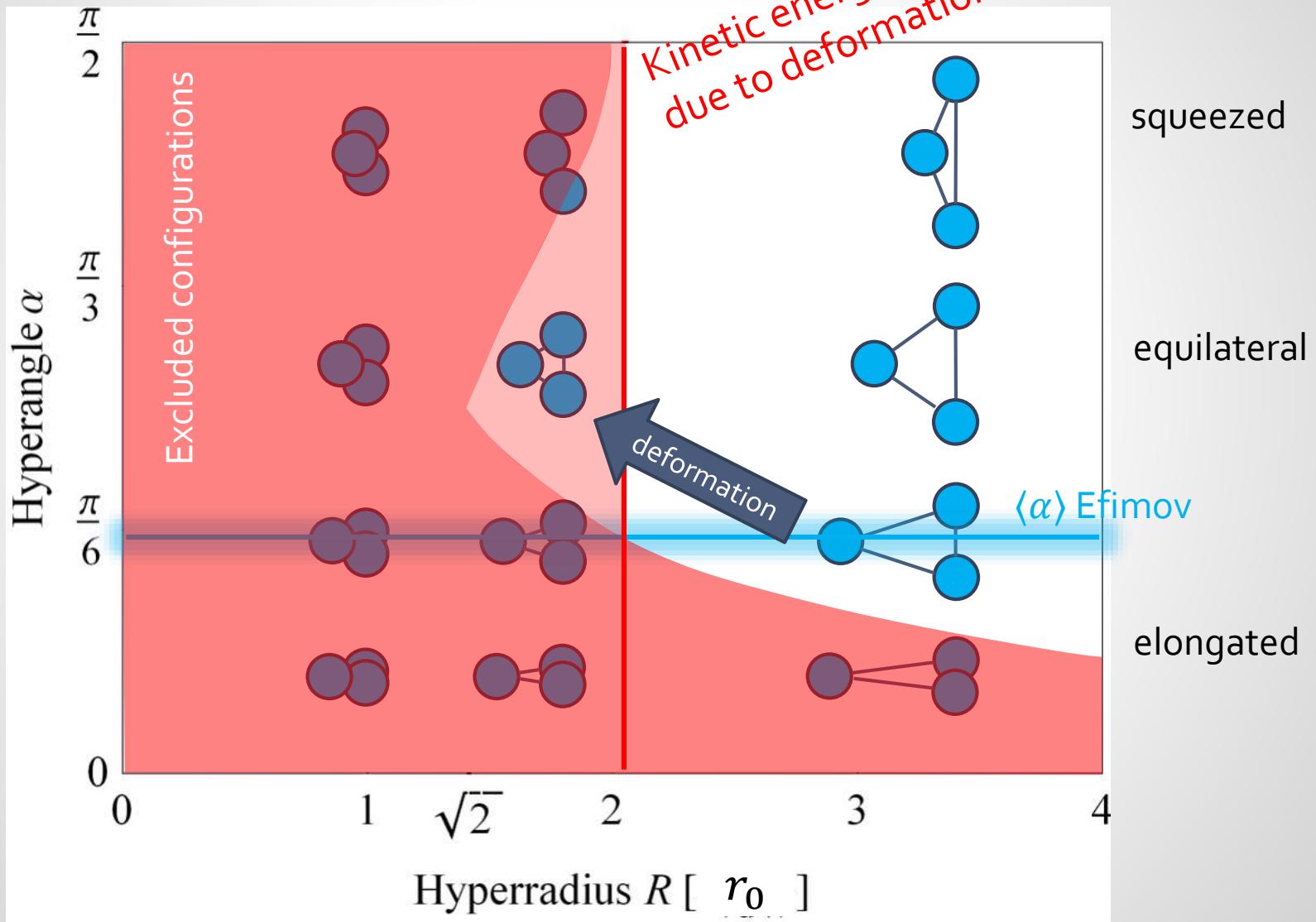
Finite-range calculation (Lennard-Jones potential)



Finite-range calculation (Lennard-Jones potential)



Interpretation: two-body correlation



Confirmation 1: pair correlation model

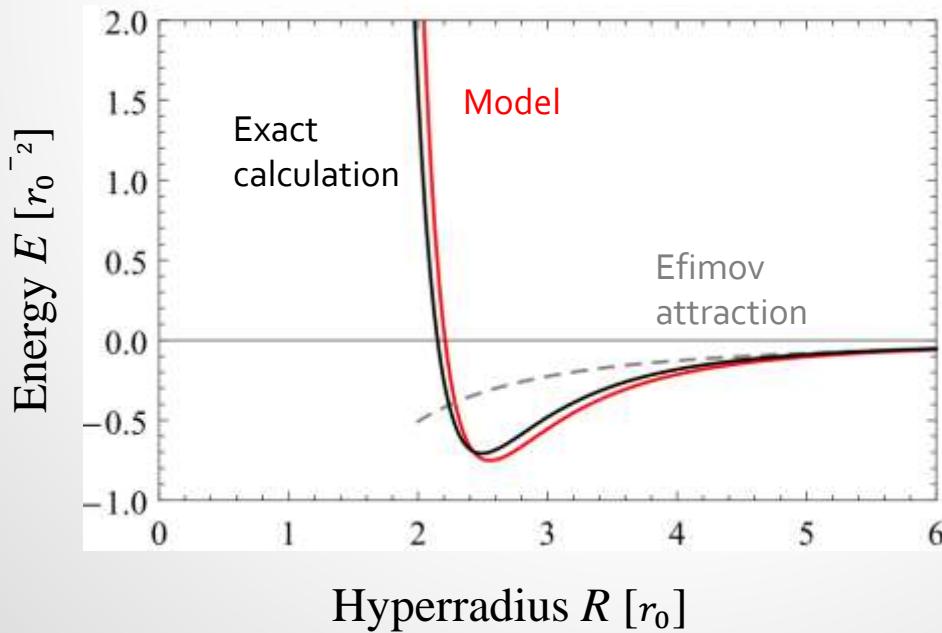
$$\Phi_{\text{Model}} = \Phi_{\text{Efimov}} \times \varphi(r_{12}) \varphi(r_{23}) \varphi(r_{31})$$

(hyperangular wave function)

(product of pair correlations)



3-body potential $U(R) = \frac{\lambda}{R^2} + \int d \cos \theta d\alpha \left| \frac{\partial \Phi}{\partial R} \right|^2$



(for a Lennard-Jones interaction with 5 bound states)

Confirmation 2: separable model

$$V = \xi |\chi\rangle\langle\chi|$$

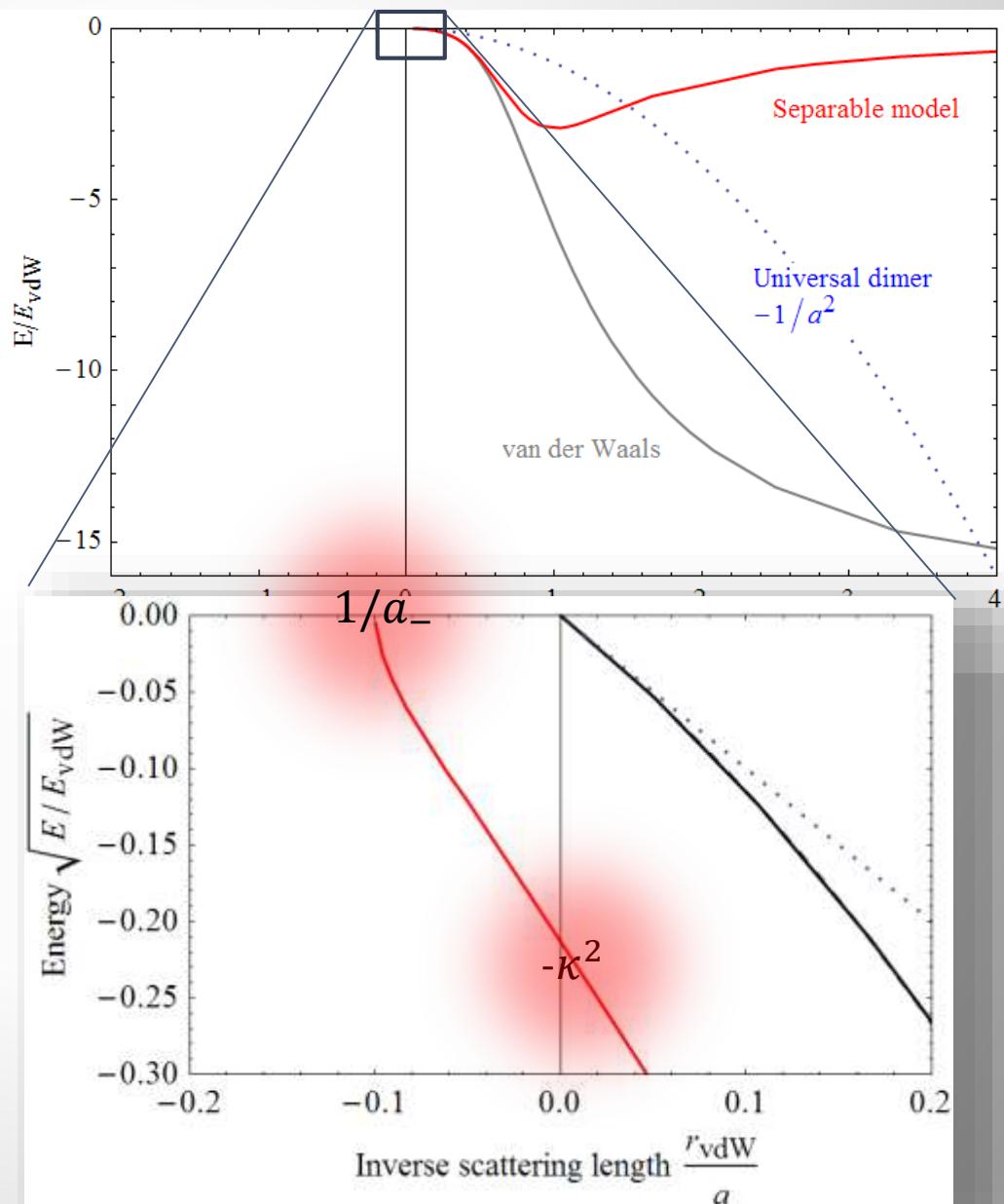
Parameterised to reproduce
exactly the two-body
correlation at zero energy.

$$\chi(q) = 1 - q \int_0^\infty (\bar{\psi}_0(r) - \psi_0(r)) \sin qr dr$$

$$\xi = 4\pi \left(\frac{1}{a} - \frac{2}{\pi} \int_0^\infty |\chi(q)|^2 dq \right)^{-1}$$

Reproduces the low-energy
2-body physics very well

- Scattering length
- Effective range
- Last bound state
-

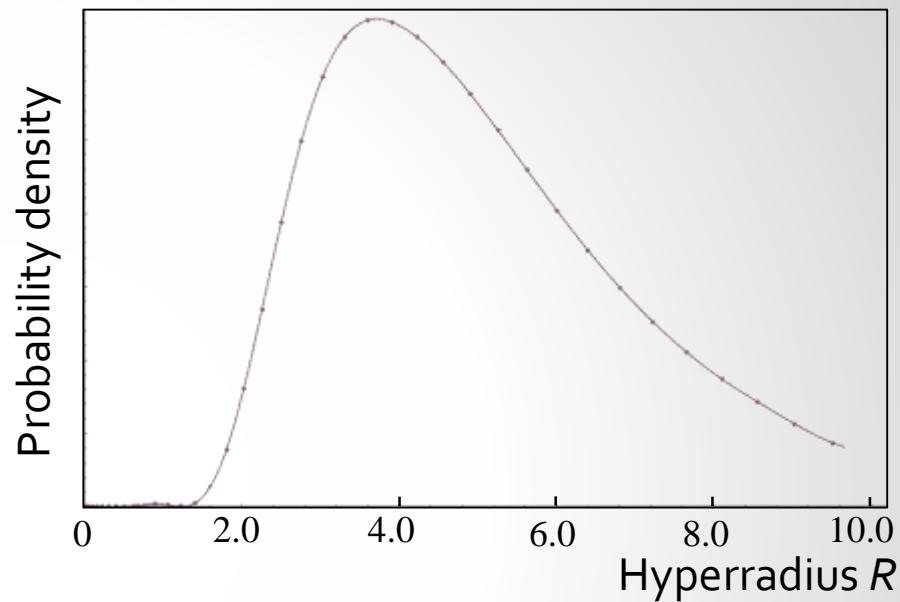
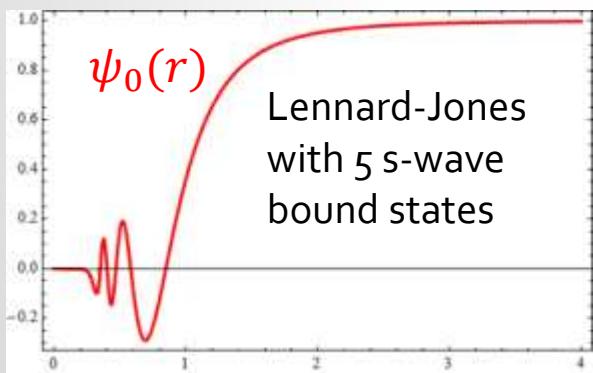


Confirmation 2: separable model

$$V = \xi |\chi\rangle\langle\chi|$$

Parameterised to reproduce
exactly the two-body
correlation at zero energy.

van der Waals two-body correlation



$$a_- = -10.86(1) r_{vdW}$$

Exact: $a_- = -11(1) r_{vdW}$

Experiments: $a_- = -10(1) r_{vdW}$

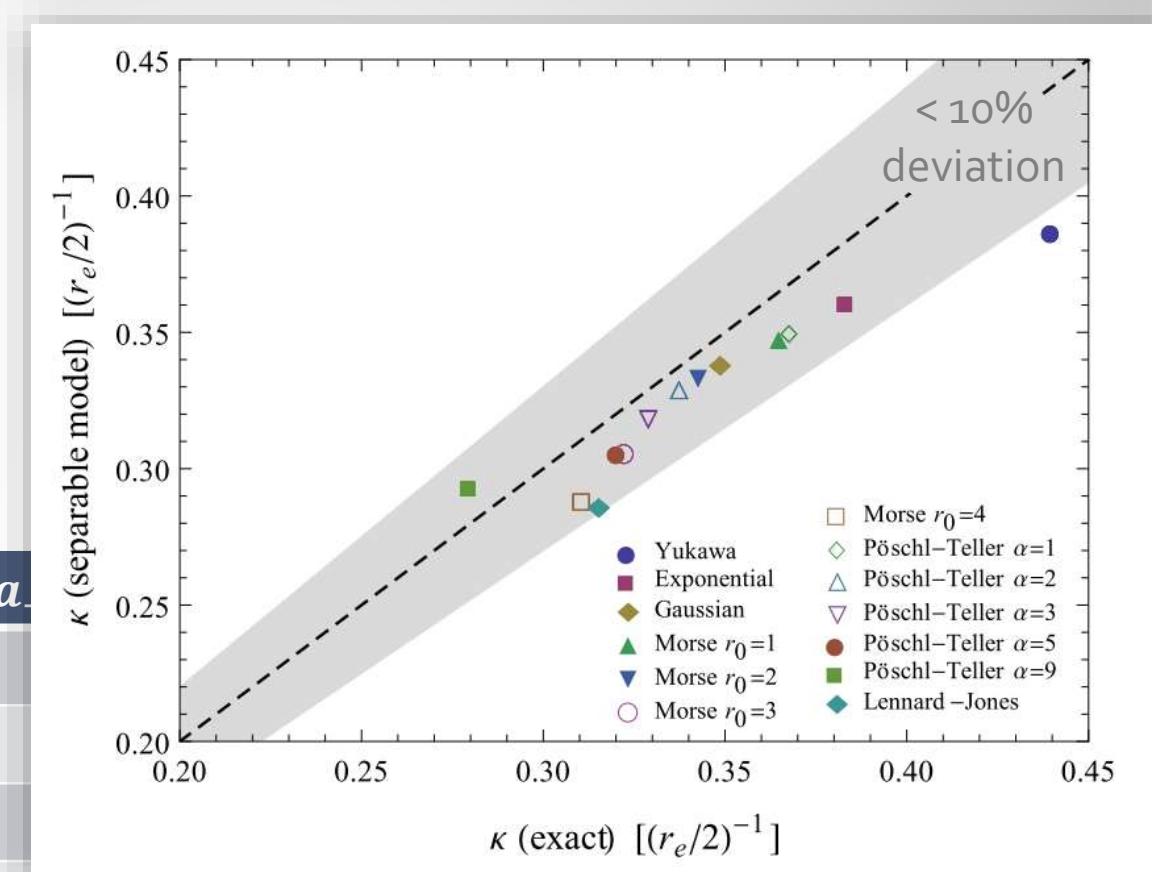
Confirmation 2: separable model

$$V = \xi |\chi\rangle\langle\chi|$$

Parameterised to reproduce
exactly the two-body
correlation at zero energy.

Other potentials

Potential	a			
Yukawa	-5.73			
Exponential	-10.7			
Gaussian	-4.27			
Morse ($r_0 = 1$)	-12.3	-12.6	0.180	0.173
Morse ($r_0 = 2$)	-16.4	-16.3	0.131	0.128
Pöschl-Teller ($\alpha = 1$)	-6.02	-6.23	0.367	0.350

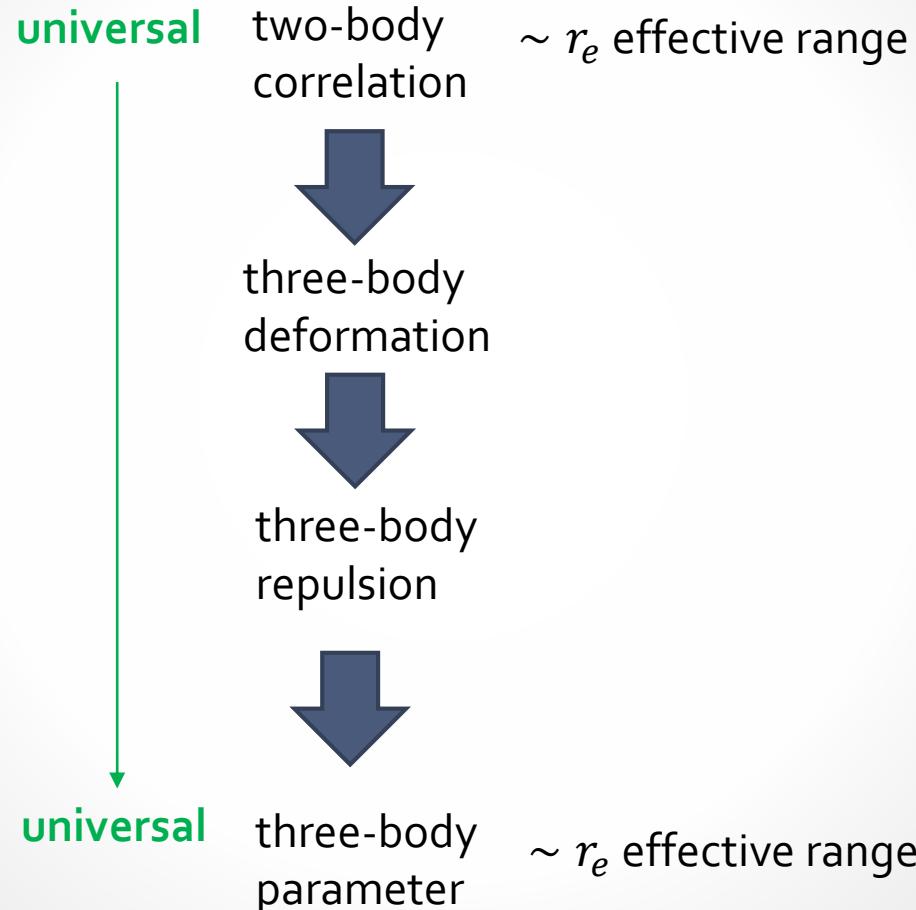


Separable
model

S. Moszkowski, S. Fleck, A. Kriek, L. Theuyl, J.-M. Richard,
and K. Varga, Phys. Rev. A 62, 032504 (2000).

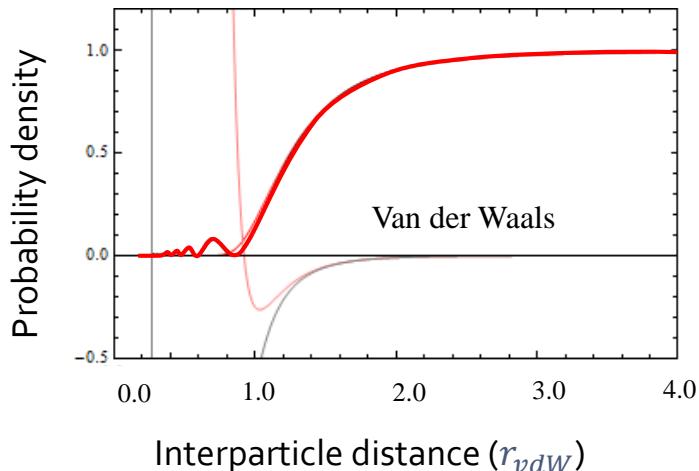
Exact
calculations

Summary

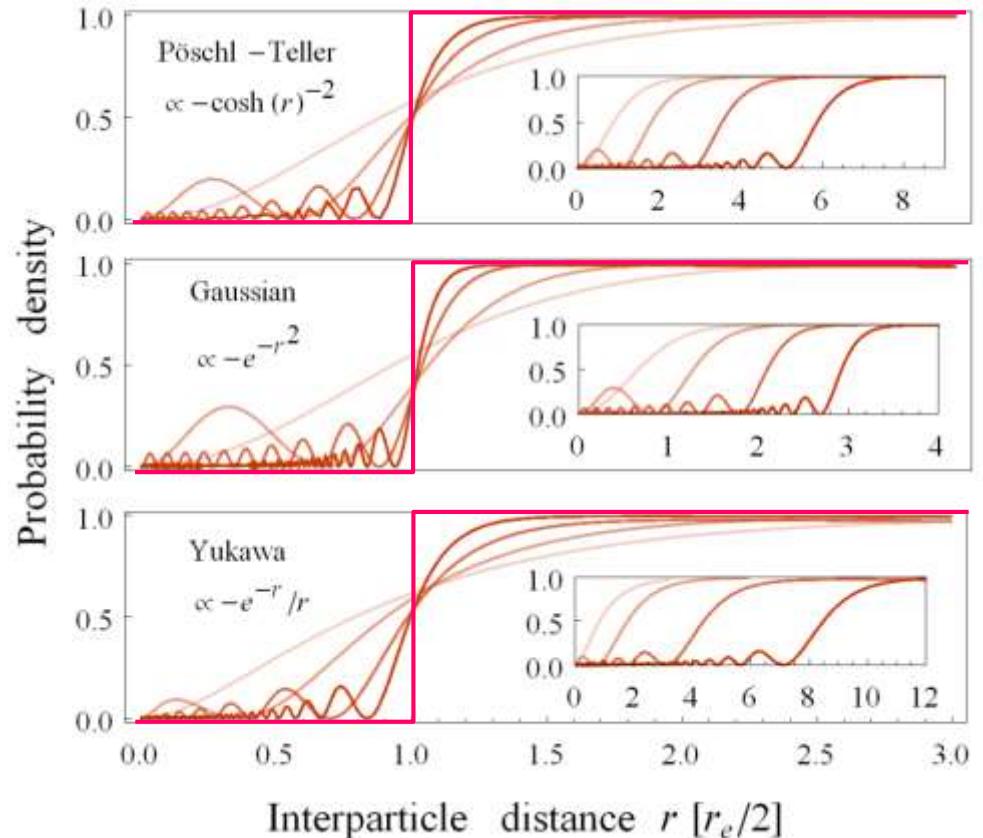


Two-body correlation universality classes

Power-law tails $\propto -\frac{1}{r^n}$ ($n > 3$)



Faster than Power-law tails



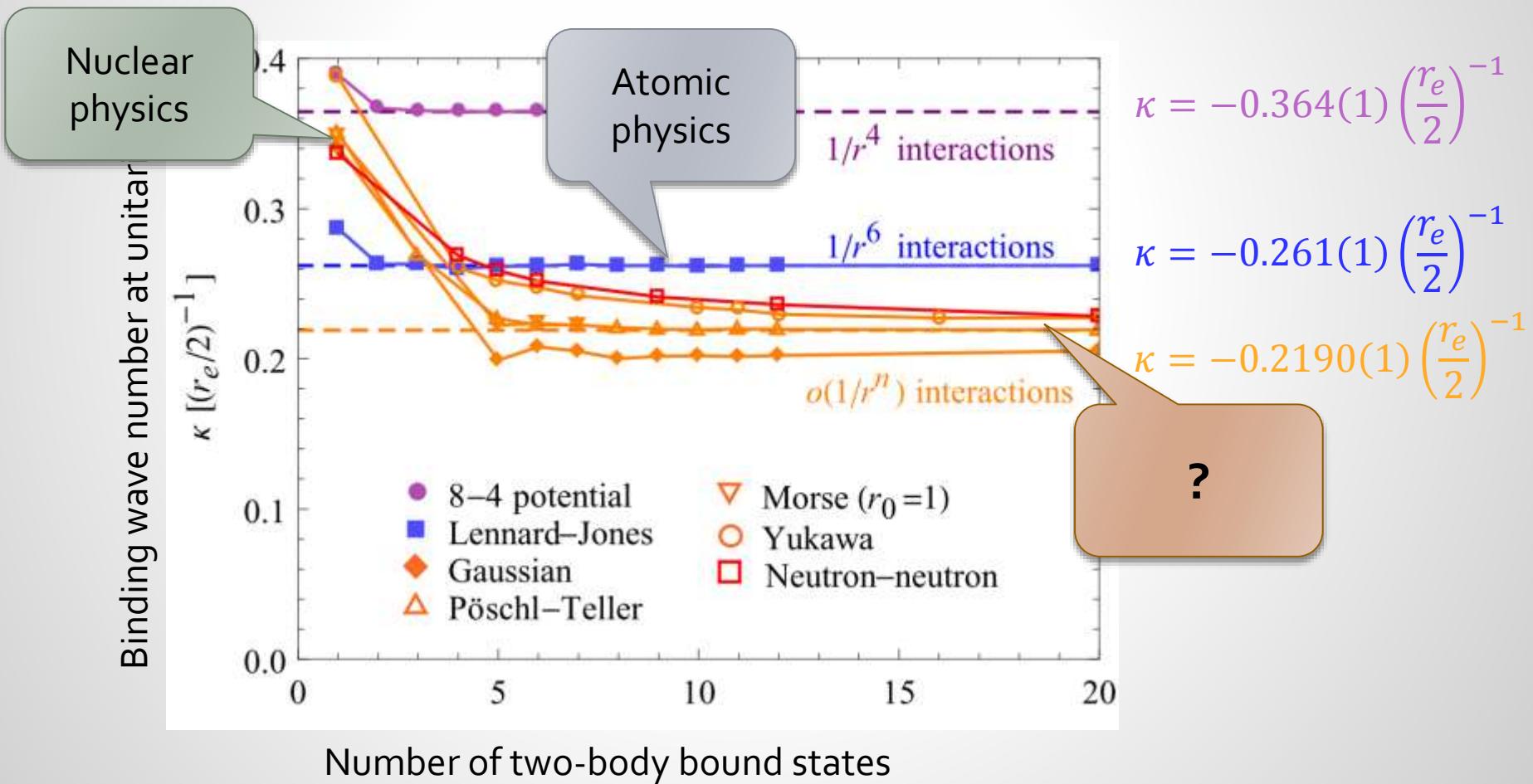
Universal correlation

$$\psi_0(r) = \Gamma\left(\frac{n-1}{n-2}\right)(r/r_n)^{1/2} J_{1/(n-2)}(2(r/r_n)^{-(n-2)/2})$$

Step function
correlation limit

Separable model

3-body parameter in units of the *two-body effective range*
(= size of two-body correlation)



Summary

The **3-body parameter** is (mostly) determined by the **2-body correlation**.

Reason:

2-body correlation induces a **deformation** of the 3-body system.

Consequences: the 3-body parameter

- is on the order of the **effective range**.
- has different universal values for **distinct classes** of interaction

P. Naidon, S. Endo, M. Ueda, arXiv:1208.3912

P. Naidon, S. Endo, M. Ueda, PRL 112, 105301 (2014)

A photograph of a road lined with cherry blossom trees in full bloom. The trees have dense clusters of white flowers. The road is paved and leads towards a bridge in the distance. The sky is clear and blue. A small signpost is visible on the left side of the road.

Thank you