

~~*From Cold Atoms to Halo Nuclei*~~

An EFT description of three-body physics

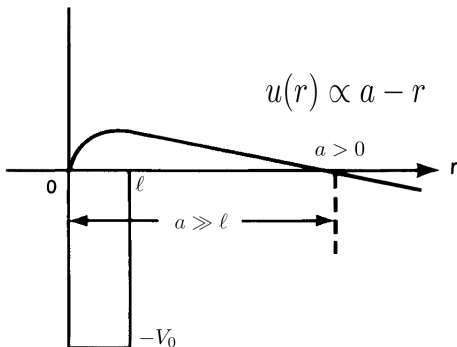
Chen Ji || TRIUMF

In collaboration with
B. Acharya, D.R. Phillips & Ch. Elster (Ohio Univ.)

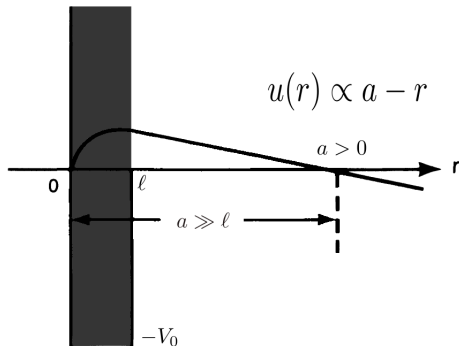
INT workshop, May 12-16, 2014



- Separation of scales:
 $a \gg \ell$
- 2-body universality:
 $B_2 = 1/ma^2$



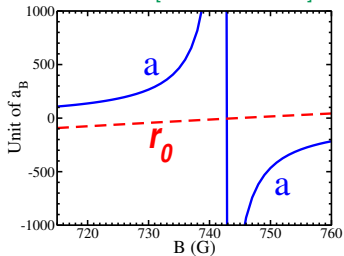
- **Separation of scales:**
 $a \gg \ell$
- **2-body universality:**
 $B_2 = 1/ma^2$
- physics at $r \sim a$ is insensitive to physics at $r \sim \ell$
- large- a physics is studied in ℓ/a expansion
- effects from SR-dynamics can be included as perturbation



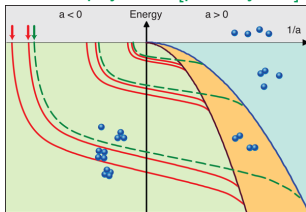
universal physics exists in systems with $\ell \ll a$

- atomic physics
 - cold atomic gases (^{133}Cs , ^7Li , ^{39}K):
 $\ell \ll a$ varies near Feshbach resonance

^7Li atoms [Gross et al. '09]



Efimov physics [pic:Esry '09]

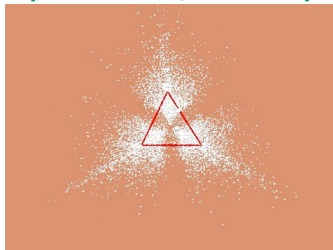


universal physics exists in systems with $l \ll a$

- atomic physics
 - cold atomic gases (^{133}Cs , ^7Li , ^{39}K):
 $l \ll a$ varies near Feshbach resonance
 - ^4He atomic clusters:
 $l_{vdw} \sim 7\text{\AA}$, $a \sim 100\text{\AA}$

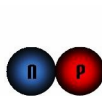
^4He trimer

[pic: Blume@physics.wsu.edu]

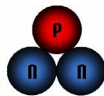


universal physics exists in systems with $l \ll a$

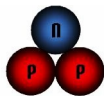
- atomic physics
 - cold atomic gases (^{133}Cs , ^7Li , ^{39}K):
 $l \ll a$ varies near Feshbach resonance
 - ^4He atomic clusters:
 $l_{vdw} \sim 7\text{\AA}$, $a \sim 100\text{\AA}$
- nuclear physics
 - few-nucleon systems (^3H , ^3He , ^4He):
 $l_{np}^t \sim 1.7\text{ fm}$, $a_{np}^t \sim 5.4\text{ fm}$



Deuterium



Tritium



Helium-3

universal physics exists in systems with $\ell \ll a$

- atomic physics

- cold atomic gases** (^{133}Cs , ^7Li , ^{39}K):
 $\ell \ll a$ varies near Feshbach resonance

- ^4He atomic clusters:**
 $\ell_{vdw} \sim 7\text{\AA}$, $a \sim 100\text{\AA}$

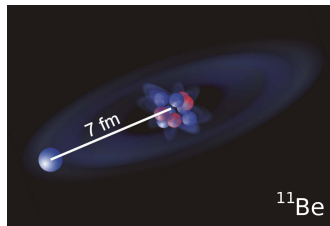
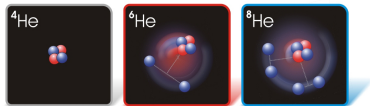
- nuclear physics

- few-nucleon systems** (^3H , ^3He , ^4He):
 $\ell_{np}^t \sim 1.7\text{ fm}$, $a_{np}^t \sim 5.4\text{ fm}$

- halo nuclei** (core + valence N)

$$^6\text{He}: S_{2n} \approx 1\text{ MeV} \quad E_c^* \approx 20\text{ MeV}$$


$$Q \sim \sqrt{m_N S_{2n}} \approx 30\text{ MeV} \quad \Lambda \sim \sqrt{m_N E_c^*} \approx 140\text{ MeV}$$



- **EFT is an approach to systems with a separation of scales**
 - atomic systems \rightarrow zero-range EFT
 - few-nucleon systems \rightarrow pionless EFT
 - halo nuclei \rightarrow halo EFT
 - **study physics in expansion of ℓ/a or Q/Λ**

- **EFT with contact interactions**


- **2-body contact (LO)**



$$= -iC_0$$

C_0 determined by a 2-body observable

- **3-body contact (LO)**



$$= -iD_0$$

D_0 determined by a 3-body observable

- EFT is an approach to systems with a separation of scales

- atomic systems \rightarrow zero-range EFT
- few-nucleon systems \rightarrow pionless EFT
- halo nuclei \rightarrow halo EFT

- study physics in expansion of ℓ/a or Q/Λ

- EFT with contact interactions

- 2-body contact (LO) introduce a dimer field

$$\begin{array}{c} \diagup \\ \bullet \\ \diagdown \end{array} = -iC_0 \quad \xrightarrow{C_0=g^2/\Delta} \quad \begin{array}{c} \diagup \\ \text{---} \\ \diagdown \end{array} = -i\sqrt{2}g$$

C_0 determined by a 2-body observable

- 3-body contact (LO)

$$\begin{array}{c} \diagup \\ \bullet \\ \diagdown \\ \bullet \\ \diagup \\ \bullet \\ \diagdown \end{array} = -iD_0 \quad \xrightarrow{D_0=-3hg^2/\Delta^2} \quad \begin{array}{c} \diagup \\ \text{---} \\ \diagdown \end{array} = ih$$

D_0 determined by a 3-body observable

Bedaque, Hammer, van Kolck '99

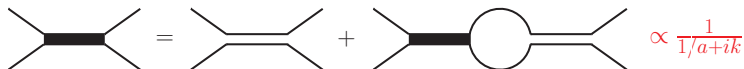
- EFT Lagrangian for 3 identical bosons at LO**

$$\mathcal{L} = \psi^\dagger \left(i\partial_0 + \frac{\nabla^2}{2m} \right) \psi - d^\dagger \left(i\partial_0 + \frac{\nabla^2}{4m} - \Delta \right) d - \frac{g}{\sqrt{2}} \left(d^\dagger \psi \psi + \text{h.c.} \right) + h d^\dagger d \psi^\dagger \psi + \dots$$

- terms with more derivatives are at higher orders

- Non-perturbative features at LO**

- particle-particle scattering (tune g)



$$\propto \frac{1}{1/a+ik}$$

- particle-dimer scattering (tune h)

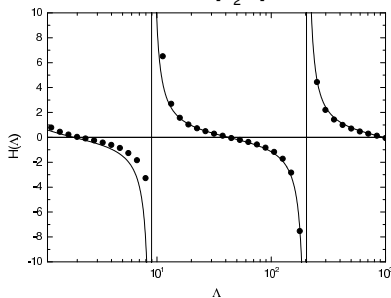
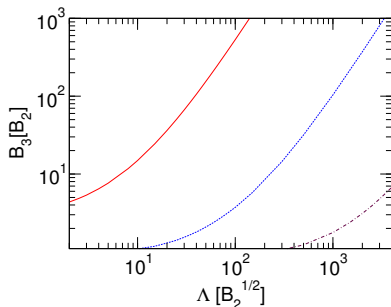


- Without 3BF:

- 3-body spectrum:
 - cutoff dependent ($\Lambda \sim 1/\ell$)
 - Platter '09

- LO 3BF h :

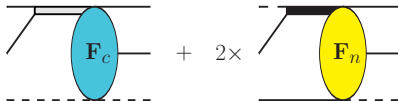
- tune $H(\Lambda) = \Lambda^2 h/2mg^2$:
 - fix one 3-body observable
- limit cycle:
 - $H(\Lambda)$ periodic for $\Lambda \rightarrow \Lambda(22.7)^n$
 - Bedaque *et al.* '00
- \rightarrow Efimov physics



- with perturbation theory:
 - study effective-range corrections (r_0/a) to 3-atom systems
 - NLO ($\sim r_0/a$) range effects in 3body Efimov spectrum
Platter, C.J., Phillips, PRA 79, 022702 (2009)
 - NLO ($\sim r_0/a$) range effects in cold-atom recombination
C.J., Platter, Phillips, EPL. 92, 13003 (2010)
C.J., Phillips, Platter, Ann. Phys. 327, 1803 (2012)
 - N²LO ($\sim r_0^2/a^2$) range effects on Helium-4 trimer
C.J., Phillips, Few Body Syst. 54, 2317 (2013)

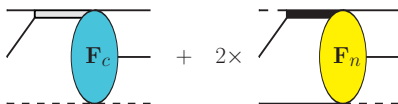
- include core + valence nucleon d.o.f. e.g., ${}^6\text{He} \rightarrow n + n + {}^4\text{He}$; ${}^{11}\text{Be} \rightarrow n + {}^{10}\text{Be}$
- include spin & isospin d.o.f.

- include core + valence nucleon d.o.f. e.g., ${}^6\text{He} \rightarrow n + n + {}^4\text{He}$; ${}^{11}\text{Be} \rightarrow n + {}^{10}\text{Be}$
- include spin & isospin d.o.f.
- 2n-halo wave functions:

$$\Psi_x(p, q) = \text{[Diagram 1]} + 2 \times \text{[Diagram 2]}$$


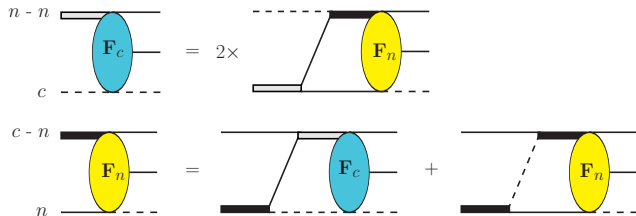
The diagram illustrates the decomposition of a 2n-halo wave function. On the left, a blue oval labeled F_c represents the core, with two lines extending upwards to represent valence nucleons. On the right, a yellow oval labeled F_n represents the core, also with two lines extending upwards. The two diagrams are separated by a plus sign and a multiplier of 2, indicating that the second diagram is multiplied by 2.

- include core + valence nucleon d.o.f. e.g., ${}^6\text{He} \rightarrow n + n + {}^4\text{He}$; ${}^{11}\text{Be} \rightarrow n + {}^{10}\text{Be}$
- include spin & isospin d.o.f.
- 2n-halo wave functions:

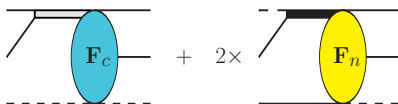
$$\Psi_x(p, q) = \text{Diagram 1} + 2 \times \text{Diagram 2}$$


- Faddeev equation

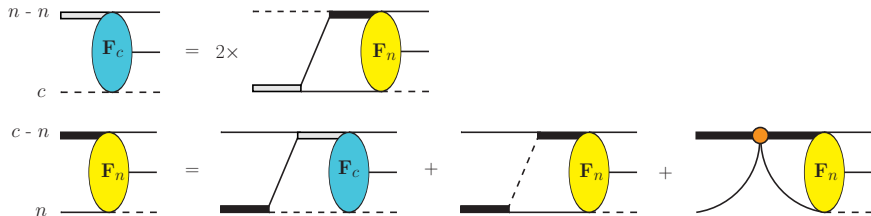
$$\begin{array}{l} n - n \\ c \end{array} \text{Diagram 1} = 2 \times \text{Diagram 2}$$

$$\begin{array}{l} c - n \\ n \end{array} \text{Diagram 3} = \text{Diagram 4} + \text{Diagram 5}$$


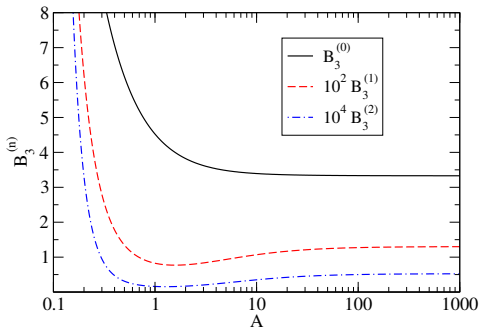
- include core + valence nucleon d.o.f. e.g., ${}^6\text{He} \rightarrow n + n + {}^4\text{He}$; ${}^{11}\text{Be} \rightarrow n + {}^{10}\text{Be}$
- include spin & isospin d.o.f.
- 2n-halo wave functions:

$$\Psi_x(p, q) = \text{Diagram 1} + 2 \times \text{Diagram 2}$$


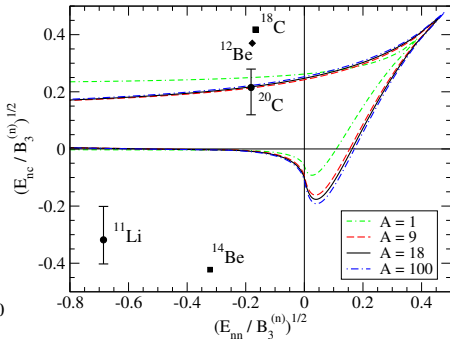
- Faddeev equation + 3-body force



- Efimov states in $a_{nn}, a_{nc} \rightarrow \infty$



- Implication of excited Efimov halo



Canham, Hammer EPJA 2008

- $1n$ halo and resonant state
 - p-wave resonance:
 - ${}^5\text{He}$ [Bertulani *et al.* '02, Bedaque *et al.* '02]

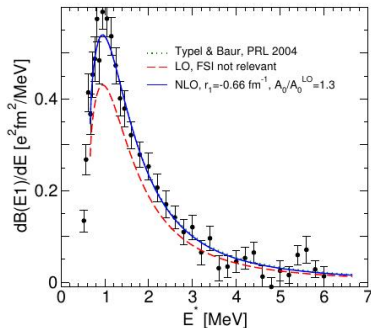
- $1n$ halo and resonant state

- p-wave resonance:

^5He [Bertulani *et al.* '02, Bedaque *et al.* '02]

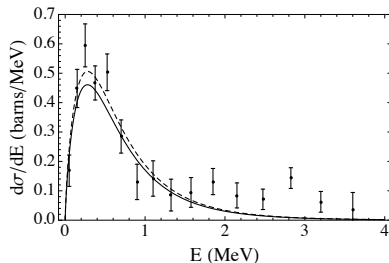
- E1 transition:

^{11}Be photo-dissociation



[Hammer, Phillips '11]

^{19}C photo-dissociation



Data: Nakamura *et al.*, RIKEN ('99,'03);
Calculation: Acharya, Phillips ('13)

- **$1p$ halo**

- ^{17}F [Ryberg, Forssén, Platter '13]

- p - α scattering [Higa, Bertulani, van Kolck '14]

- **1p halo**

^{17}F [Ryberg, Forssén, Platter '13]

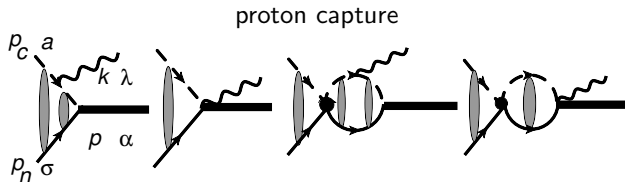
p - α scattering [Higa, Bertulani, van Kolck '14]

- **radiative nucleon captures**

$^7\text{Li} + n \rightarrow ^8\text{Li} + \gamma$ [Rupak, Higgs '11, Zhang, Nollett, Phillips '13]

$^{14}\text{C} + n \rightarrow ^{15}\text{C} + \gamma$ [Rupak, Fernando, Vaghani '12]

$^7\text{Be} + p \rightarrow ^8\text{B} + \gamma$ [Zhang, Nollett, Phillips '14]



- **2-neutron halo:**

- *n*-core in s-wave virtual/real bound state:

- ^{11}Li , ^{12}Be , ^{20}C [Canham, Hammer '08, '10]

- ^{22}C [Yamashita, Carvalho, Frederico, Tomio '11]

- ^{22}C Acharya, C.J., Phillips PLB 723 (2013)

- **2-neutron halo:**
 - n -core in s-wave virtual/real bound state:
 - ^{11}Li , ^{12}Be , ^{20}C [Canham, Hammer '08, '10]
 - ^{22}C [Yamashita, Carvalho, Frederico, Tomio '11]
 - ^{22}C Acharya, C.J., Phillips PLB 723 (2013)
 - charge radius of 2n s-wave halos [Hagen, Hammer, Platter '13]

- **2-neutron halo:**

- *n*-core in s-wave virtual/real bound state:

- ^{11}Li , ^{12}Be , ^{20}C [Canham, Hammer '08, '10]

- ^{22}C [Yamashita, Carvalho, Frederico, Tomio '11]

- ^{22}C Acharya, C.J., Phillips PLB 723 (2013)

- charge radius of 2n s-wave halos [Hagen, Hammer, Platter '13]

- heaviest 2n s-wave halo:

- ^{62}Ca [Hagen, Hagen, Hammer, Platter '13]

- **2-neutron halo:**

- n -core in s-wave virtual/real bound state:

- ^{11}Li , ^{12}Be , ^{20}C [Canham, Hammer '08, '10]

- ^{22}C [Yamashita, Carvalho, Frederico, Tomio '11]

- ^{22}C Acharya, C.J., Phillips PLB 723 (2013)

- charge radius of 2n s-wave halos [Hagen, Hammer, Platter '13]

- heaviest 2n s-wave halo:

- ^{62}Ca [Hagen, Hagen, Hammer, Platter '13]

- ^6He : n - α in p-wave resonance

- EFT + Gamow shell model [Rotureau, van Kolck '12]

- EFT + Faddeev Equations C.J., Elster, Phillips arXiv:1405.2394 (2014)

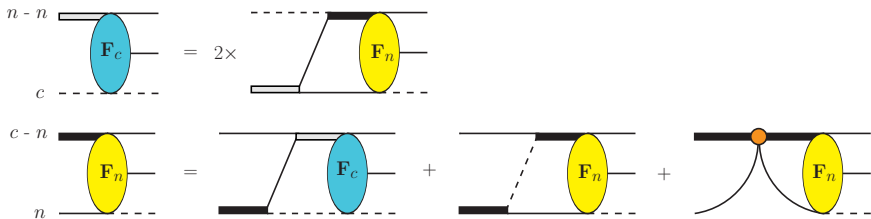
	^{20}C	^{21}C	^{22}C
bound/unbound	bound		
ground state	0^+		
binding/virtual energy	$S_{2n}=4.76$ MeV Ozawa et al. '11		
matter radius r_m	2.97(5) fm Ozawa et al. '01		

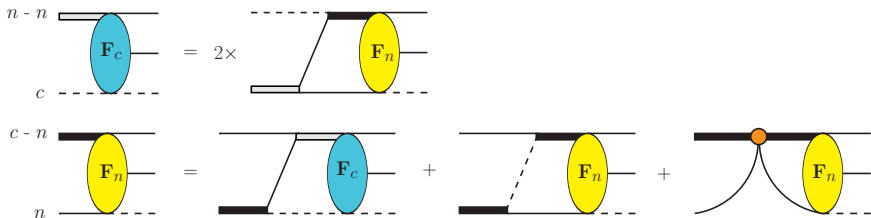
	^{20}C	^{21}C	^{22}C
bound/unbound	bound	unbound	
ground state	0^+	$S_{1/2}$	
binding/virtual energy	$S_{2n}=4.76$ MeV Ozawa et al. '11	$E_{nc} > 2.9$ MeV Mosby et al. '13 ??	
matter radius r_m	2.97(5) fm Ozawa et al. '01	—	

	^{20}C	^{21}C	^{22}C
bound/unbound	bound	unbound	bound
ground state	0^+	$S_{1/2}$	0^+
binding/virtual energy	$S_{2n}=4.76$ MeV Ozawa et al. '11	$E_{nc} > 2.9$ MeV Mosby et al. '13 ??	$S_{2n}=0.42(94)$ MeV Audi et al. '03 $S_{2n}=-0.14(46)$ MeV Gaufrey et al. '12
	matter radius r_m	2.97(5) fm Ozawa et al. '01	—

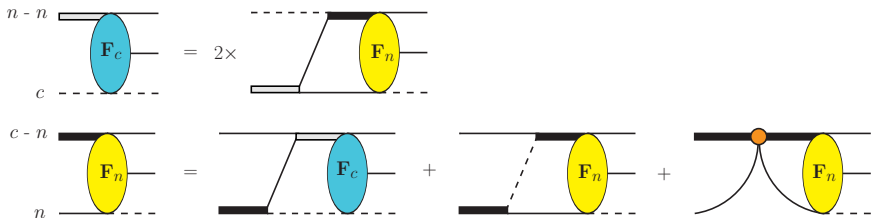
	^{20}C	^{21}C	^{22}C
bound/unbound	bound	unbound	bound
ground state	0^+	$S_{1/2}$	0^+
binding/virtual energy	$S_{2n}=4.76$ MeV	$E_{nc} > 2.9$ MeV	$S_{2n}=0.42(94)$ MeV
	Ozawa et al. '11	Mosby et al. '13 ??	Audi et al. '03 $S_{2n}=-0.14(46)$ MeV Gaufrey et al. '12
matter radius r_m	2.97(5) fm Ozawa et al. '01	—	5.4(9) fm Tanaka et al. '10

- Halo EFT: use measured ^{22}C r_m to put constraints on:
 - E_{nc} of ^{21}C ($a < 0$)
 - S_{2n} of ^{22}C





- parameterize $n - n$ and $n - {}^{20}\text{C}$ interactions:
 - 2body virtual energies: $E_{nn} = 1/(M_n a_{nn}^2)$, E_{nc}



- parameterize $n - n$ and $n - {}^{20}\text{C}$ interactions:
 - 2body virtual energies: $E_{nn} = 1/(M_n a_{nn}^2)$, E_{nc}
- parameterize $n - {}^{21}\text{C}$ 3body contact term
 - fix $r_m[{}^{22}\text{C}]$

- two-body form factor

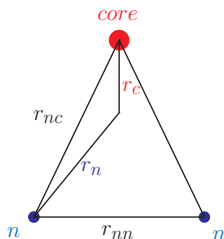
$$\mathcal{F}(k^2) = \int d\mathbf{p} \int d\mathbf{q} \Psi^\dagger(\mathbf{p}, \mathbf{q}) \Psi(\mathbf{p} - \mathbf{k}, \mathbf{q}) \approx 1 - \frac{1}{6} r_{2b}^2 k^2$$

- two-body form factor

$$\mathcal{F}(k^2) = \int d\mathbf{p} \int d\mathbf{q} \Psi^\dagger(\mathbf{p}, \mathbf{q}) \Psi(\mathbf{p} - \mathbf{k}, \mathbf{q}) \approx 1 - \frac{1}{6} r_{2b}^2 k^2$$

- matter radius (with point-like core)

$$R = \sqrt{A r_c^2 + 2 r_n^2} / (A + 2)$$



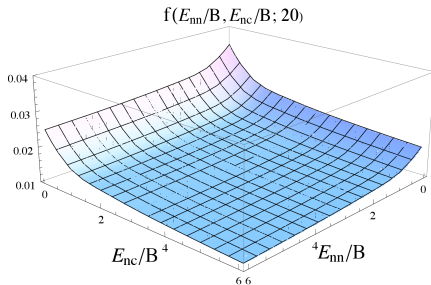
- Define a dimensionless quantity:

$$M_n B_3 \cdot R^2 \equiv f \left(\frac{E_{nn}}{B_3}, \frac{E_{nc}}{B_3}; A \right)$$

- Define a dimensionless quantity:

$$M_n B_3 \cdot R^2 \equiv f \left(\frac{E_{nn}}{B_3}, \frac{E_{nc}}{B_3}; A \right)$$

- At $A = 20$:
 - a function of E_{nn}/B_3 and E_{nc}/B_3



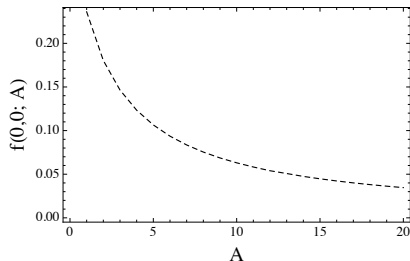
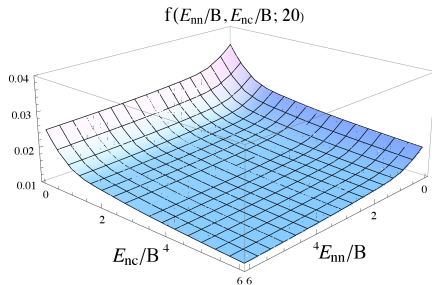
- Define a dimensionless quantity:

$$M_n B_3 \cdot R^2 \equiv f\left(\frac{E_{nn}}{B_3}, \frac{E_{nc}}{B_3}; A\right)$$

- At $A = 20$:

- a function of E_{nn}/B_3 and E_{nc}/B_3

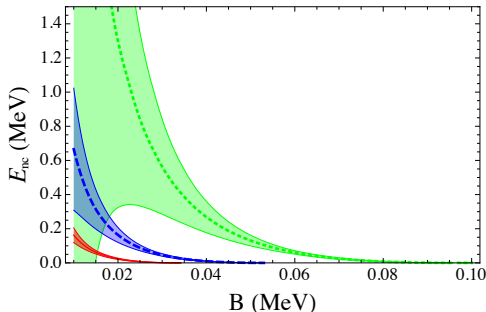
- Unitary limit $E_{nn} = E_{nc} = 0$:
 - a function of A



Acharya, C.J., Phillips PLB 2013

- Include finite size of ^{20}C

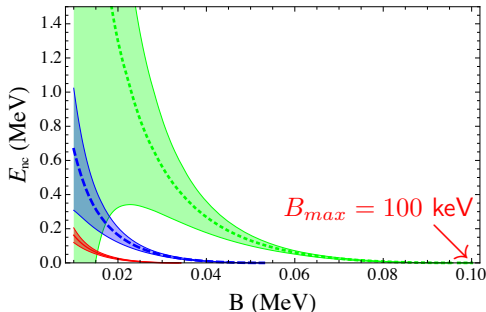
$$R^2 = r_m^2[^{22}\text{C}] - \frac{20}{22} r_m^2[^{20}\text{C}] = (5.4^{+0.9}_{-0.9})^2 - \frac{20}{22} (2.97 \pm 0.05)^2 \text{ fm}^2$$



bands: EFT power counting

- Include finite size of ^{20}C

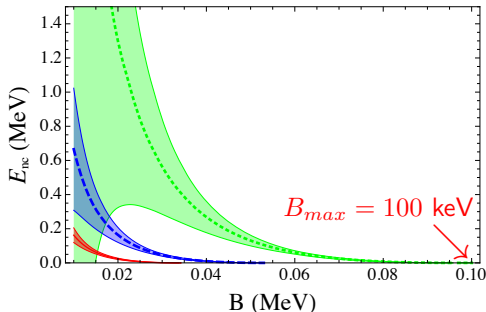
$$R^2 = r_m^2[^{22}\text{C}] - \frac{20}{22} r_m^2[^{20}\text{C}] = (5.4^{+0.9}_{-0.9})^2 - \frac{20}{22} (2.97 \pm 0.05)^2 \text{ fm}^2$$



bands: EFT power counting

- Include finite size of ^{20}C

$$R^2 = r_m^2[^{22}\text{C}] - \frac{20}{22} r_m^2[^{20}\text{C}] = (5.4_{-0.9}^{+0.9})^2 - \frac{20}{22} (2.97 \pm 0.05)^2 \text{ fm}^2$$



c.f. Yamashita et al. '11

→ $B_{max} \sim 120 \text{ keV}$

Fortune & Sherr '12

→ $B_{max} \sim 220 \text{ keV}$

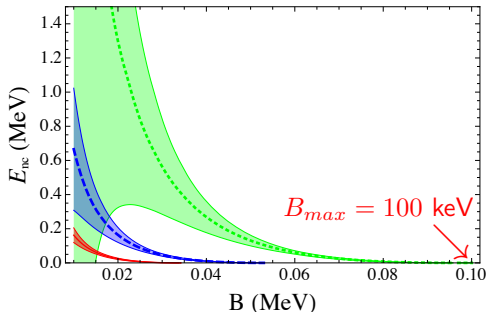
Gaufrey et al. '12

→ $B_{max} \sim 320 \text{ keV}$

bands: EFT power counting

- Include finite size of ^{20}C

$$R^2 = r_m^2[^{22}\text{C}] - \frac{20}{22} r_m^2[^{20}\text{C}] = (5.4^{+0.9}_{-0.9})^2 - \frac{20}{22} (2.97 \pm 0.05)^2 \text{ fm}^2$$



bands: EFT power counting

c.f. Yamashita et al. '11

→ $B_{max} \sim 120 \text{ keV}$

Fortune & Sherr '12

→ $B_{max} \sim 220 \text{ keV}$

Gaufrey et al. '12

→ $B_{max} \sim 320 \text{ keV}$

Mosby et al. '13

$a_{nc} < 2.8 \text{ fm} \iff E_{nc} > 2.9 \text{ MeV}$

→ $B_{max} < 20 \text{ keV}$

inconsistency in measurements?

- **experiment in ${}^6\text{He}$**

- matter radius Tanihata *et al.* '92, Alkhazov *et al.* '97, Kislev *et al.* '05
- charge radius Wang *et al.* '04, Mueller *et al.* '07
- ${}^6\text{He}$ mass Brodeur *et al.* '12

- **cluster model**

- separable potential Ghovanlou, Lehman '74
- density-dependent nn contact interaction Esbensen *et al.* '97

- ***ab initio* calculation**

- no-core shell model Navrátil *et al.* '01
- hyperspherical harmonics Bacca *et al.* '12
- Green's function Monte Carlo Pieper *et al.* '01

- **halo EFT**

- explore **universal physics** in halo nuclei
- compare **predictions** with experiments and other theories

- $n\alpha$ interaction is dominated by the ${}^2P_{3/2}$ state

$$\begin{array}{c} n \\ \alpha \end{array} \begin{array}{c} \diagup \\ \diagdown \end{array} \text{---} \text{---} \begin{array}{c} \diagdown \\ \diagup \end{array} = \frac{1}{4\pi^2\mu_{n\alpha}} \frac{\vec{p} \cdot \vec{q}}{-1/a_1 + r_1 k^2/2 - ik^3}$$

$$a_1 = -62.95 \text{ fm}^3, r_1 = -0.8819 \text{ fm}^{-1} \text{ Ardnt et al. '73}$$

- $n\alpha$ interaction is dominated by the ${}^2P_{3/2}$ state

$$\begin{array}{c} n \\ \alpha \end{array} \begin{array}{c} \diagup \\ \diagdown \end{array} \text{---} \text{---} \begin{array}{c} \diagdown \\ \diagup \end{array} = \frac{1}{4\pi^2\mu_{n\alpha}} \frac{\vec{p} \cdot \vec{q}}{-1/a_1 + r_1 k^2/2 - ik^3}$$

$$a_1 = -62.95 \text{ fm}^3, r_1 = -0.8819 \text{ fm}^{-1} \text{ Ardnt et al. '73}$$

- $r_1 \neq 0$ Nishida '12

- $n\alpha$ EFT power counting: Bedaque, Hammer, van Kolck '02

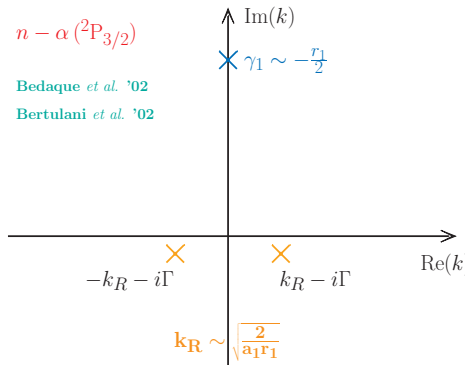
- $a_1 \sim 1/(Q^2\Lambda)$ $r_1 \sim \Lambda$
- $Q/\Lambda \sim 0.15$

- ${}^2P_{3/2}$:

shallow resonance:

$$k_R \sim Q, \quad \Gamma \sim Q^2/\Lambda$$

deep bound state: $\gamma_1 \sim \Lambda$



- $n\alpha$ EFT power counting: Bedaque, Hammer, van Kolck '02

- $a_1 \sim 1/(Q^2\Lambda)$ $r_1 \sim \Lambda$
- $Q/\Lambda \sim 0.15$

- ${}^2P_{3/2}$:

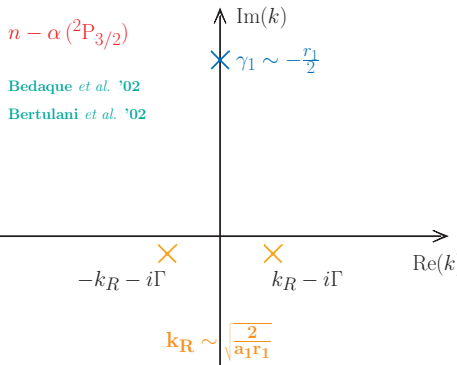
shallow resonance:

$$k_R \sim Q, \quad \Gamma \sim Q^2/\Lambda$$

deep bound state: $\gamma_1 \sim \Lambda$

- LO $n\alpha$ t-matrix ($ik^3 \rightarrow 0$)

$$\begin{aligned} t_{n\alpha} &= \frac{1}{4\pi^2\mu_{n\alpha}} \frac{\vec{p} \cdot \vec{q}}{-1/a_1 + r_1 k^2/2} \\ &= \frac{-1}{4\pi^2\mu_{n\alpha}} \frac{\vec{p} \cdot \vec{q}}{\gamma_1 (k^2 - k_R^2)} \end{aligned}$$



- $n\alpha$ EFT power counting: Bedaque, Hammer, van Kolck '02

- $a_1 \sim 1/(Q^2\Lambda)$ $r_1 \sim \Lambda$
- $Q/\Lambda \sim 0.15$

- ${}^2P_{3/2}$:

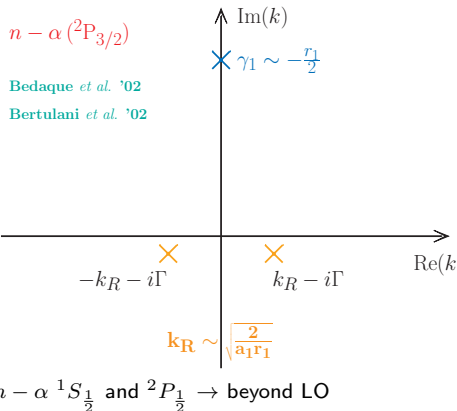
shallow resonance:

$$k_R \sim Q, \quad \Gamma \sim Q^2/\Lambda$$

deep bound state: $\gamma_1 \sim \Lambda$

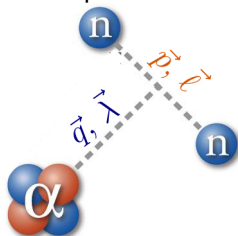
- LO $n\alpha$ t-matrix ($ik^3 \rightarrow 0$)

$$\begin{aligned} t_{n\alpha} &= \frac{1}{4\pi^2\mu_{n\alpha}} \frac{\vec{p} \cdot \vec{q}}{-1/a_1 + r_1 k^2/2} \\ &= \frac{-1}{4\pi^2\mu_{n\alpha}} \frac{\vec{p} \cdot \vec{q}}{\gamma_1 (k^2 - k_R^2)} \end{aligned}$$

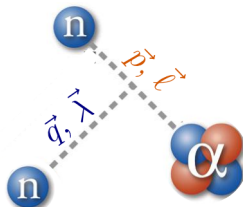


- Jacobi-momentum

α spectator



n spectator

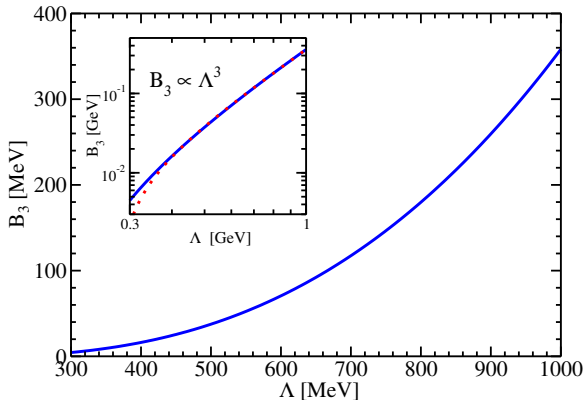


spin-orbit coupling for ${}^6\text{He}$ ($J = 0^+$)

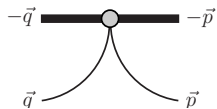
pair, spec	pair	spectator	total L, S	total J
nn, α	$\ell = 0, s_1 = 0$	$\lambda = 0, s_2 = 0$	$L = 0, S = 0$	$J = 0^+$
$n\alpha, n$	$\ell = 1, s_1 = \frac{1}{2}$	$\lambda = 1, s_2 = \frac{1}{2}$	$L = 0, S = 0$	
			$L = 1, S = 1$	

- without $nn\alpha$ 3-body force:

- S_{2n} is strongly cutoff dependent: $S_{2n} \sim \Lambda^3$ ← need 3body force!

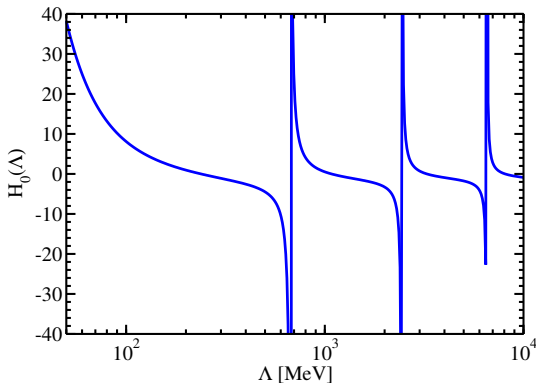


- p-wave 3BF:

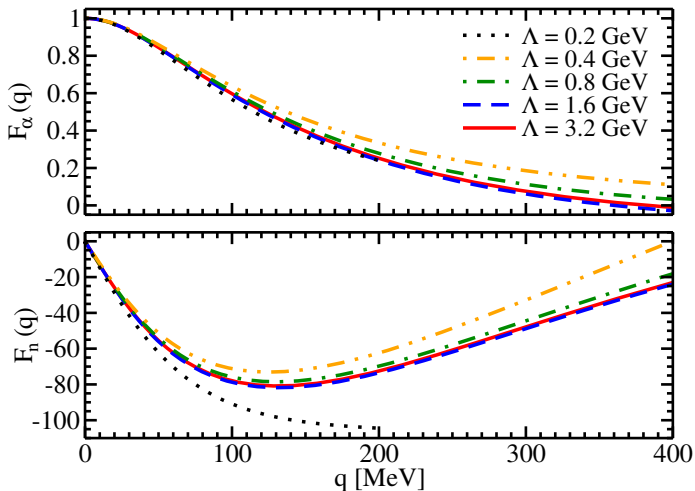


$$= M_n qp \frac{H(\Lambda)}{\Lambda^2}$$

- reproduce $S_{2n} = 0.973\text{MeV}$
- log oscillation
- No limit cycle (c.f. 3-body in S-wave)



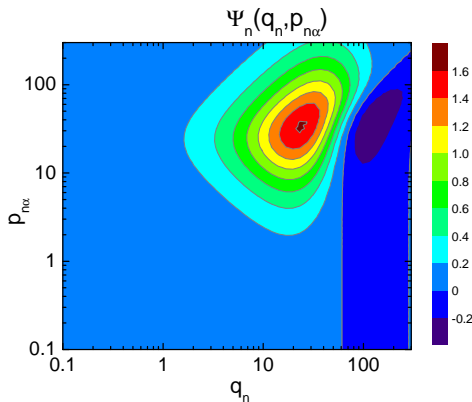
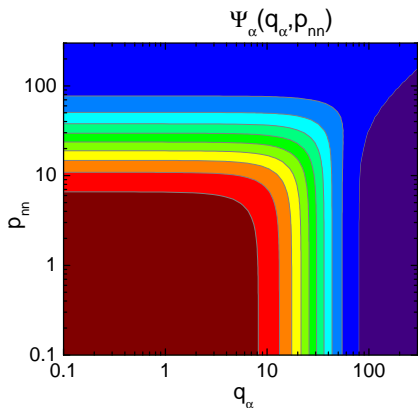
$F_\alpha(\alpha, nn)$ and $F_n(n, \alpha n)$: Λ independent



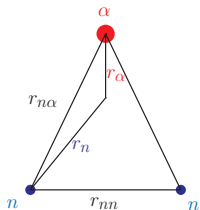
C.J., Elster, Phillips arXiv:1405.2394 (2014)

● $|\Psi\rangle$ in $\alpha - nn$ basis

● $|\Psi\rangle$ in $n - \alpha n$ basis



momenta in units of MeV



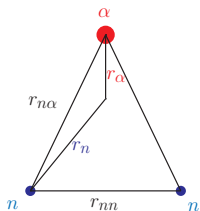
[Preliminary]

$$r_{nn} = 3.77 \text{ fm} \pm 20\%$$

$$r_{n\alpha} = 4.33 \text{ fm} \pm 20\%$$

- $\pm 20\%$ uncertainty is from higher-order EFT

		EFT [fm]	Exp [fm]
He-4	$r_{pp}[\alpha]$	—	1.455(1)
	$r_m[\alpha]$	—	1.455(1)
He-6			



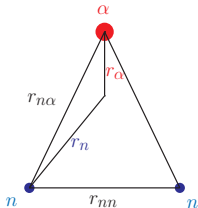
[Preliminary]

$$r_{nn} = 3.77 \text{ fm} \pm 20\%$$

$$r_{n\alpha} = 4.33 \text{ fm} \pm 20\%$$

- $\pm 20\%$ uncertainty is from higher-order EFT

		EFT [fm]	Exp [fm]
He-4	$r_{pp}[\alpha]$	—	1.455(1)
	$r_m[\alpha]$	—	1.455(1)
He-6	r_α	1.30(26)	—
	r_n	3.21(64)	—



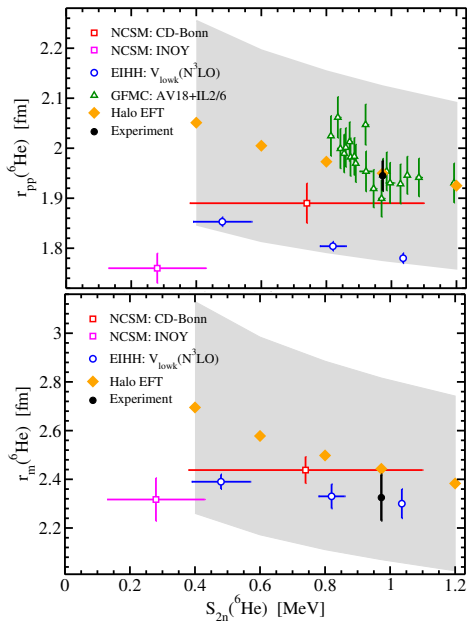
[Preliminary]

$$r_{nn} = 3.77 \text{ fm} \pm 20\%$$

$$r_{n\alpha} = 4.33 \text{ fm} \pm 20\%$$

- $\pm 20\%$ uncertainty is from higher-order EFT

		EFT [fm]	Exp [fm]
He-4	$r_{pp}[\alpha]$	—	1.455(1)
	$r_m[\alpha]$	—	1.455(1)
He-6	r_α	1.30(26)	—
	r_n	3.21(64)	—
	$r_{pp}[^6\text{He}] = \sqrt{r_\alpha^2 + r_{pp}^2[\alpha]}$	1.95(17)	1.938(23), 1.953(22)
	$r_m[^6\text{He}] = \sqrt{\frac{1}{3}(2r_\alpha^2 + r_n^2 + 2r_m^2[\alpha])}$	2.44(37)	2.33(4), 2.30(7), 2.37(5)



[Preliminary]

● He-6 point-proton radius

● He-6 matter radius

c.f. Bacca, Barnea & Schwenk, '12

- Halo nuclei are ideal places to study universal physics
 - ^{22}C : n -core in s -wave resonance
 - ^6He : n -core in p -wave resonance
- with halo EFT analysis:
 - s - / p -wave 3body force is needed for proper renormalization
 - predict probability density distribution, matter radii, point-proton radii
 - use universality to constrain experiments & other theories
- further work
 - study $^6\text{He } 2^+$ resonance state
 - extension to ^{11}Li
 - higher-order effects: range corrections, more partial waves, ...