

Neutrino Interactions with Nucleons and Nuclei

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Motivation and Contents

- Determination of neutrino oscillation parameters requires knowledge of neutrino energy
- Modern experiments use complicated nuclear targets: from Carbon to Argon
- Nuclear effects affect everything:
 - event cross section measurements
 - event identification
 - final state particles
 - neutrino energy reconstruction
 - determination of oscillation parameters

Neutrino Oscillations

- 2-Flavor Oscillation:

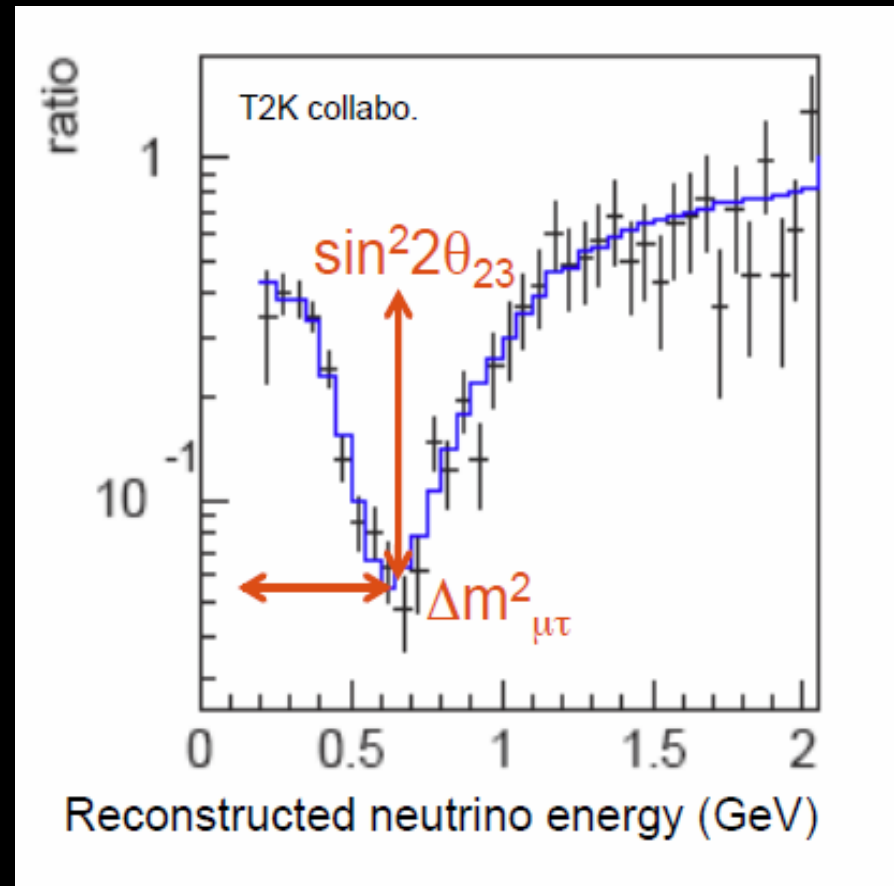
$$P(\nu_\mu \rightarrow \nu_e) = \sin^2 2\theta \sin^2 \left(\frac{\Delta m^2 L}{4E_\nu} \right)$$

Know: L , need E_ν to determine $\Delta m^2, \theta$

- 3-Flavor Oscillation: allows for CP violation

Observable Oscillation Parameters

$$P(\nu_\mu \rightarrow \nu_e) = \sin^2 2\theta \sin^2 \left(\frac{\Delta m^2 L}{4E_\nu} \right)$$



Oscillation probability

Long-Baseline Accelerator Appearance Experiments

- Oscillation probability complicated and dependent not only on θ_{13} but also:

1. CP violation parameter (δ)
2. Mass hierarchy (sign of Δm_{31}^2)
3. Size of $\sin^2\theta_{23}$

$$\begin{aligned}
 P(\nu_\mu \rightarrow \nu_e) = & 4C_{13}^2 S_{13}^2 S_{23}^2 \sin^2 \frac{\Delta m_{31}^2 L}{4E} \times \left(1 + \frac{2a}{\Delta m_{31}^2} (1 - 2S_{13}^2) \right) \\
 & + 8C_{13}^2 S_{12} S_{13} S_{23} (C_{12} C_{23} \cos \delta - S_{12} S_{13} S_{23}) \cos \frac{\Delta m_{32}^2 L}{4E} \sin \frac{\Delta m_{31}^2 L}{4E} \sin \frac{\Delta m_{21}^2 L}{4E} \\
 & - 8C_{13}^2 C_{12} C_{23} S_{12} S_{13} S_{23} \sin \delta \sin \frac{\Delta m_{32}^2 L}{4E} \sin \frac{\Delta m_{31}^2 L}{4E} \sin \frac{\Delta m_{21}^2 L}{4E} \\
 & + 4S_{12}^2 C_{13}^2 \{ C_{12}^2 C_{23}^2 + S_{12}^2 S_{23}^2 S_{13}^2 - 2C_{12} C_{23} S_{12} S_{23} S_{13} \cos \delta \} \sin^2 \frac{\Delta m_{21}^2 L}{4E} \\
 & - 8C_{13}^2 S_{13}^2 S_{23}^2 \cos \frac{\Delta m_{32}^2 L}{4E} \sin \frac{\Delta m_{31}^2 L}{4E} \frac{aL}{4E} (1 - 2S_{13}^2)
 \end{aligned}$$

⇒ These extra dependencies are both a “curse” and a “blessing”

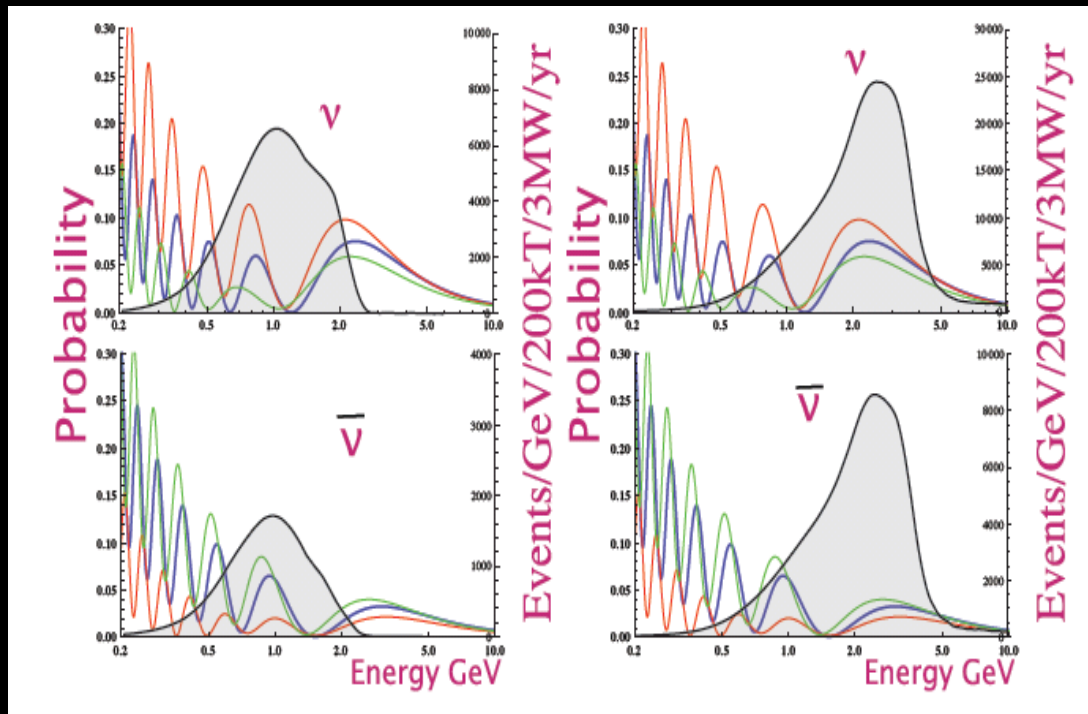
Reactor Disappearance Experiments

- Reactor disappearance measurements provide a straight forward method to measure θ_{13} with no dependence on matter effects and CP violation

$$P(\bar{\nu}_e \rightarrow \bar{\nu}_e) = 1 - \sin^2 2\theta_{13} \sin^2 \frac{\Delta m_{13}^2 L}{4E} + \textit{small terms}$$

LBNE, δ_{CP} Sensitivity

From: Bishai et al., hep-ex 12034090



From:
Bishai et al
arXiv:1203.409

$$\delta_{CP} = 0$$

$$\delta_{CP} = \pi/2$$

$$\delta_{CP} = -\pi/2$$

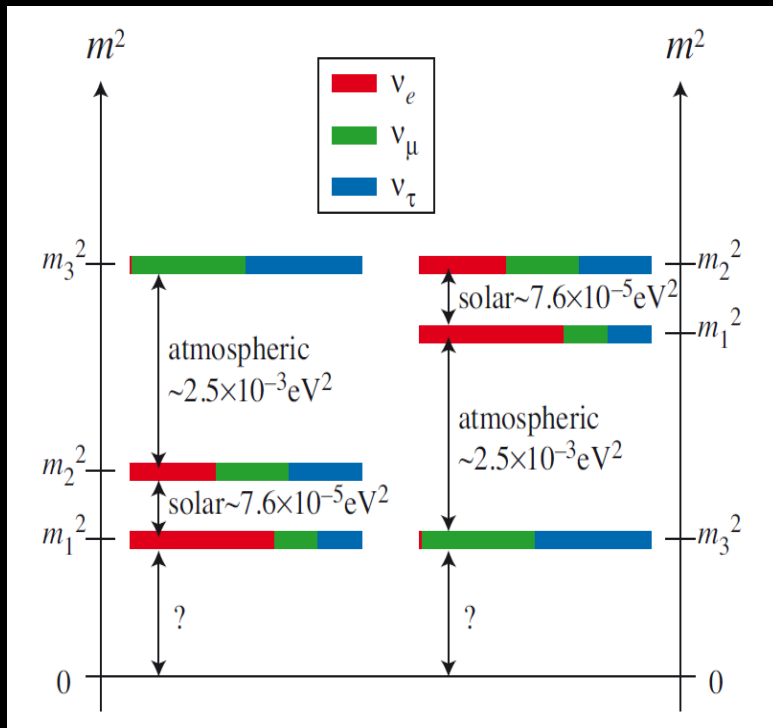
8 GeV

60 GeV

proton energy

Need energy to distinguish between different δ_{CP}

Oscillation Signal Dependence on Hierarchy and Mixing Angle



Energy has to be known better than 50 MeV

Shape sensitive to hierarchy and sign of mixing angle

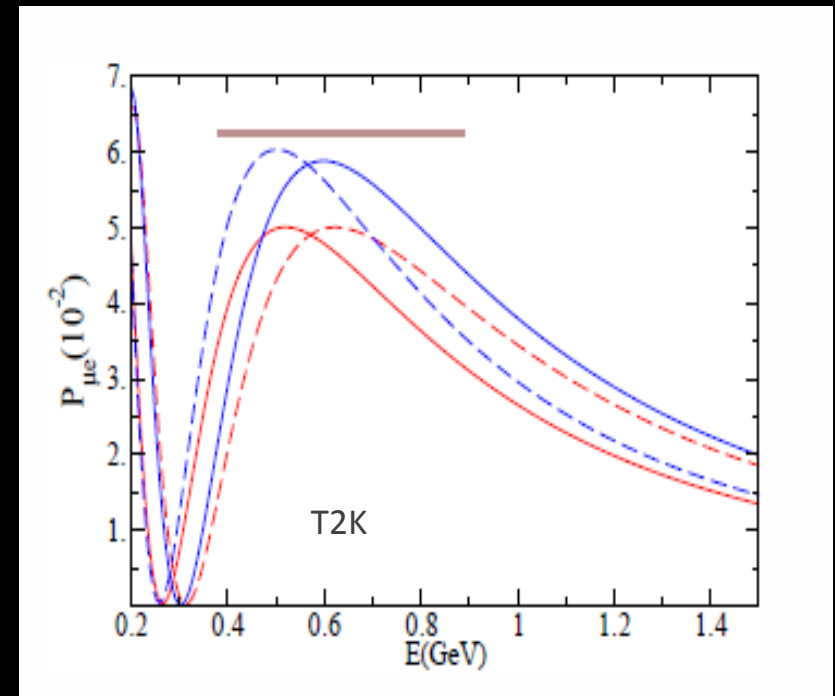


Fig. 2. $\mathcal{P}_{\mu e}$ in matter versus neutrino energy for the T2K experiment. The blue curves depict the normal hierarchy, red the inverse hierarchy. Solid curves depict positive θ_{13} , dashed curves negative θ_{13}

D.J. Ernst et al., arXiv:1303.4790 [nucl-th]

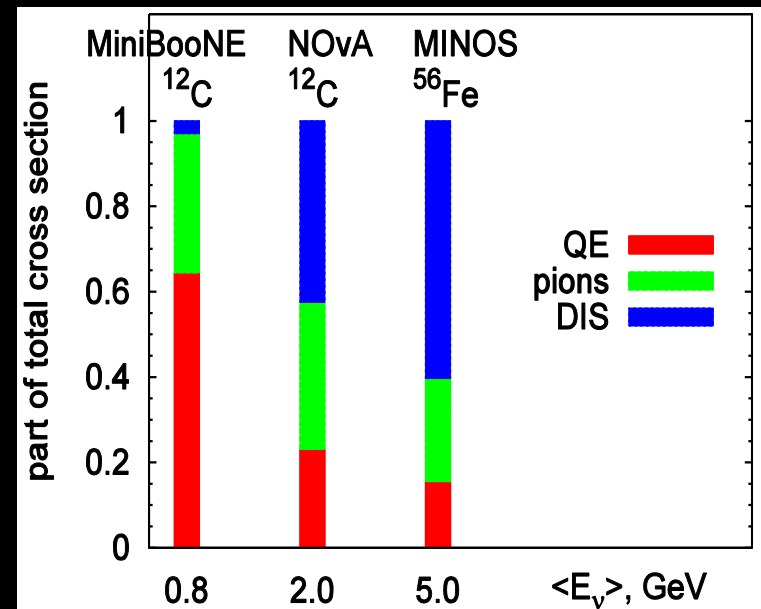
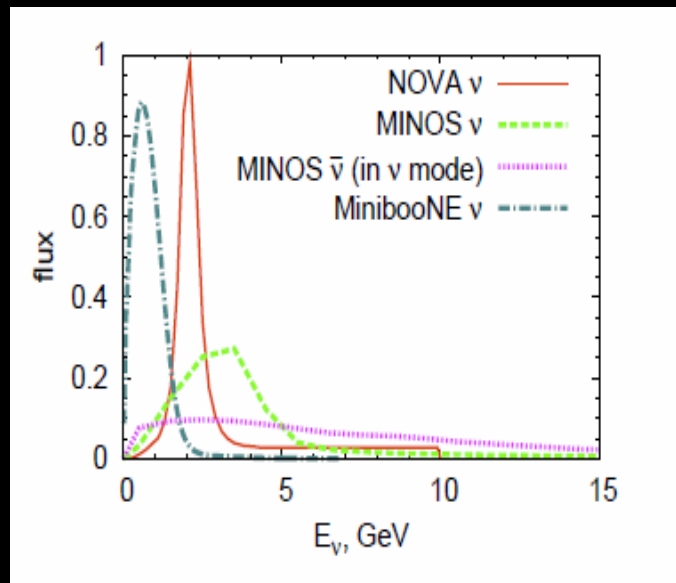
Appearance experiment

- Near detector:
 - Neutrino Flux
 - Background
 - Intrinsic ν_e
 - Neutrino energy
-
- Far detector:
 - Extrapolate Flux
 - Background
 - Neutrino energy

$$P(\nu_\mu \rightarrow \nu_e) = 1 - \sin^2 2\theta \sin^2 \left(\frac{1.27 \Delta m^2 L}{E_\nu} \right) + \text{other}$$

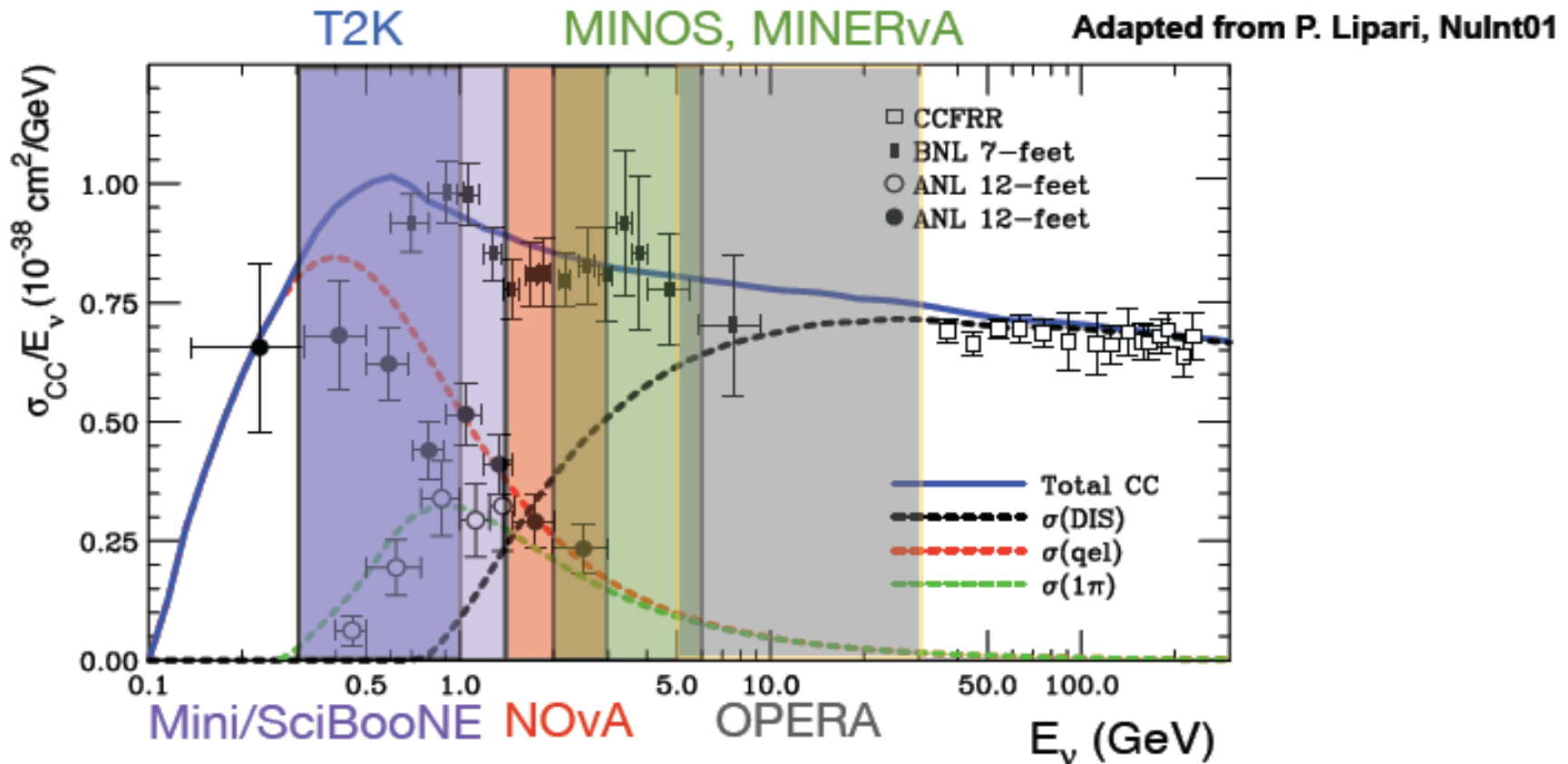
Neutrino Beams

- Neutrinos do not have fixed energy nor just one reaction mechanism



Have to reconstruct energy from final state of reaction
Different processes are entangled

Neutrino Cross-Sections



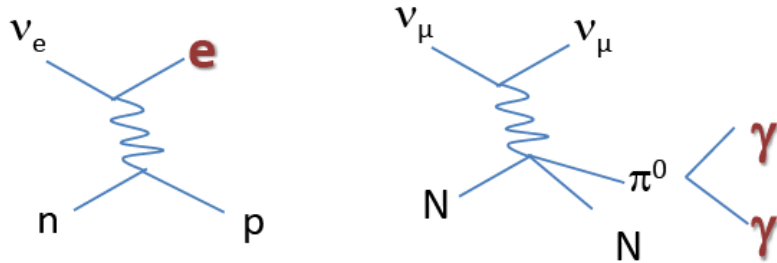
Upcoming experiments will continue to work in an “interesting” region:

- Large contributions from QE, Resonances and DIS regions
- Are these categories even sufficient ?

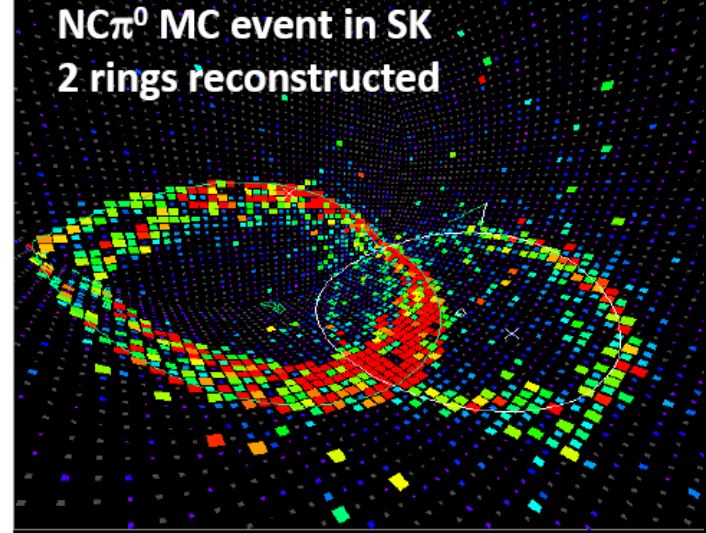
Neutrino Interactions

•for ν_e appearance

- beam ν_e
- NC π^0 events

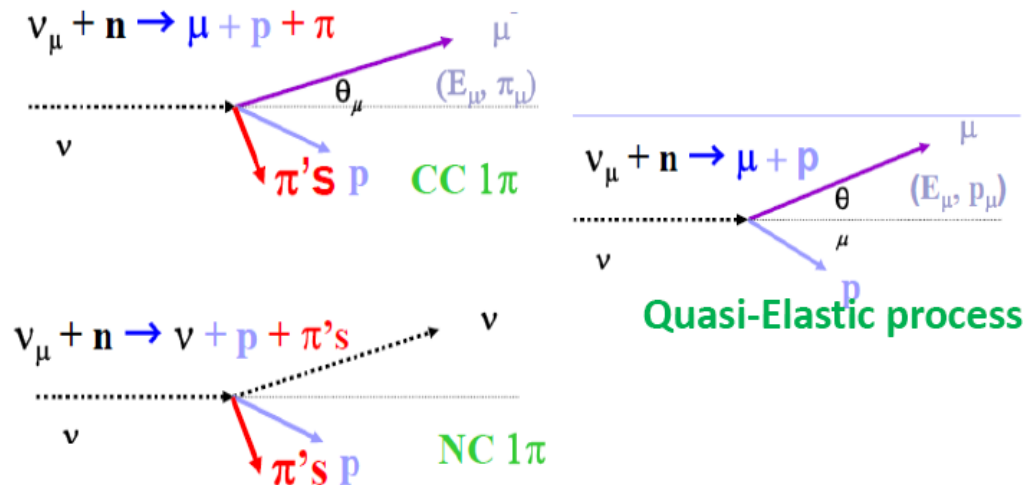


NC π^0 MC event in SK
2 rings reconstructed

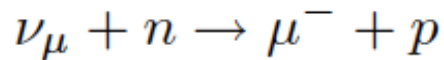


•for ν_μ disappearance (muon energy measurement)

- inelastic processes



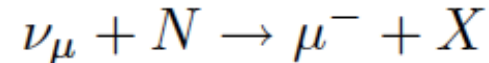
Energy reconstruction



$$E_{\nu} = E_{\nu}(E_{\mu}, \theta_{\mu})$$

Kinematic:

- Rely on underlying interaction to use relate outgoing lepton kinematics to neutrino energy
- Advantage:
 - don't need hadron reconstruction
- Disadvantages
 - energy is wrong if underlying interaction is wrong (i.e. not CCQE)
 - Nuclear effects smear resolution

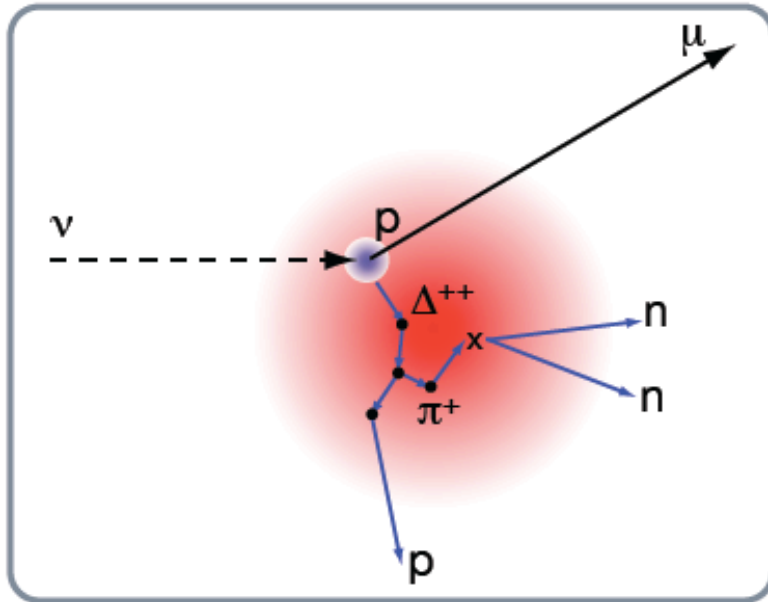


$$E_{\nu} = E_{\mu} + E_X$$

Calorimetric

- Add up the energy from the leptonic and hadronic components
- Advantages
 - No *a priori* assumption about underlying interaction
- Disadvantages
 - Relies on hadron reconstruction

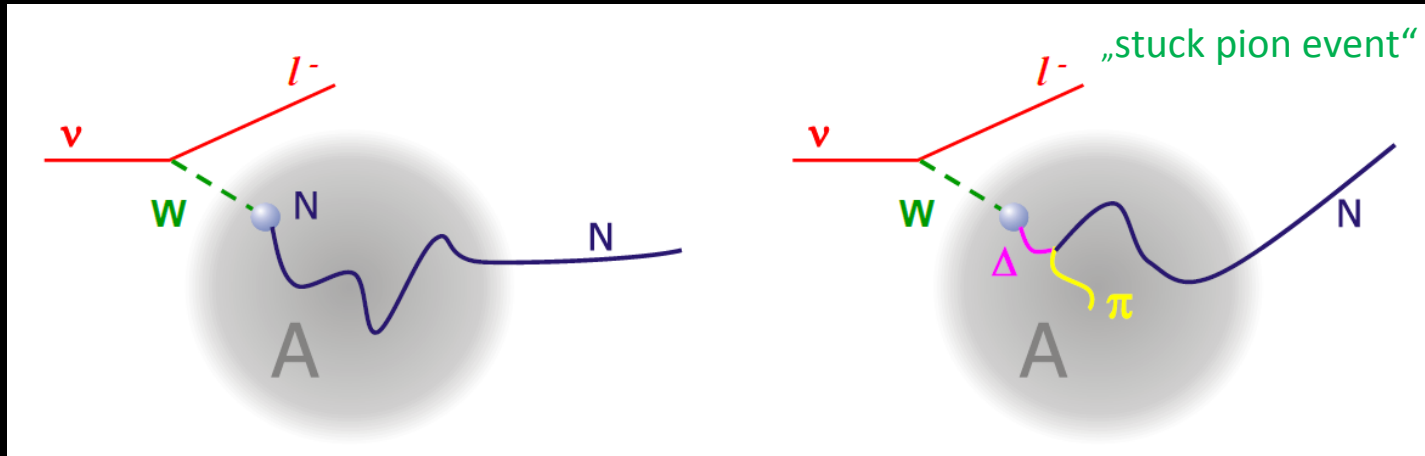
Background: Nuclear re-interactions



- Lepton kinematics shifted/smeared
- Outgoing hadronic final state (“topology”) may differ from expectation from “underlying” ν -nucleon interaction
- FSI effects may appear degenerate with hadronic interactions outside of the target nucleus.

Modeling ν interactions in nucleus

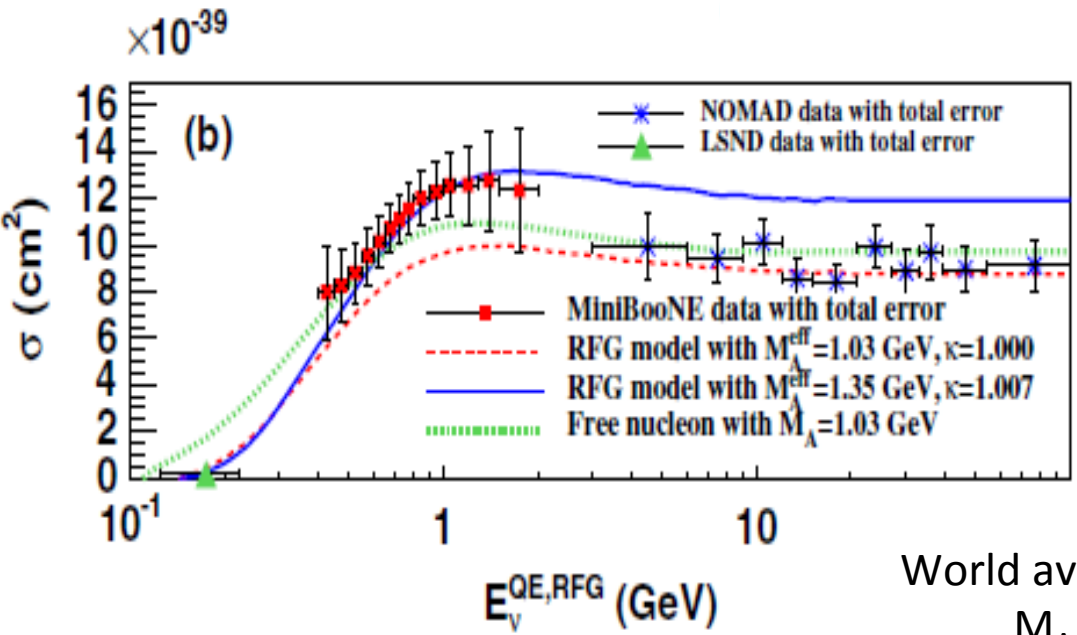
- Underlying ν -nucleon/quark interaction
 - Mode (CCQE, resonance, etc.)
 - Determine “final” state of interaction
- Initial state nucleon/quark
 - Fermi motion, binding energy
- Final state effects
 - Pauli blocking
 - Propagate hadrons within nucleus
 - Absorption, scattering, CEX, etc.



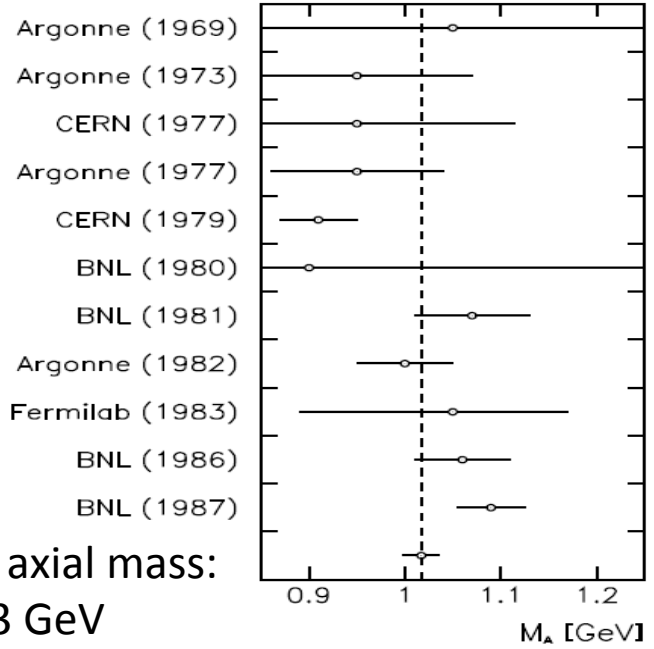
Complication to identify QE, entangled with π production

Nuclear Targets (K2K, MiniBooNE, T2K, MINOS, Minerva, ...)

MiniBooNE QE puzzle



World average axial mass:
 $M_A = 1.03 \text{ GeV}$



MiniBooNE use mineral oil (Cerenkov rings): identifies QE by muon and zero pion, corrects for 'stuck pions'

Can the nuclear effects be responsible for a higher axial mass value ?

Consider only events with no pion in final state: Cerenkov Experiments

- Experimental oscillation analyses requires QE identification (QE-like) with no pions in final state. This is why we care about QE.
- 0-pion events can involve pion production with subsequent pion absorption → ‘stuck pion events’
- Experiments remove the contribution of pionless events due to absorption according to MC models
- Definition of QE cannot distinguish between true QE (1p-1h), N^* and 2p-2h interactions

Oscillation analysis
and
Energy Reconstruction
in an ideal Long Baseline Experiment

Experimental Setup

- Ideal and perfect near detector (^{12}C or ^{16}O), 1 km, 1kton
- Far detector at 295 km, 22.5 kton
 - Oxygen
 - Carbon
- Use T2K flux, peak at 0.6 GeV, 750kW, 5 years running
- Use SK reconstruction efficiency as function of energy
- Use migration matrices produced by GiBUU(1.6) and GENIE(2.8.0)
- Muon neutrino disappearance only -> fit to atmospheric parameters

Go beyond simple case (arxiv:1311.4506)

- Use one neutrino generator (GiBUU) to simulate the nuclear effects and use another neutrino generator (GENIE) to extract the oscillation parameters
- In a real experiment the “real” effects from data will be used in the oscillation analysis together with “some” simulation of nuclear effects
- Neutrino generators are “enough” different to help understanding what will be the effect of different nuclear models on neutrino oscillation analyses

- Neglecting all FSI and multinucleon contributions, we can compute the number of events as:

$$N_i^{QE} = \sigma_{QE}(E_i)\phi(E_i)P_{\mu\mu}(E_i)$$

- However, in practice we will observe a different distribution at the detector, given by:

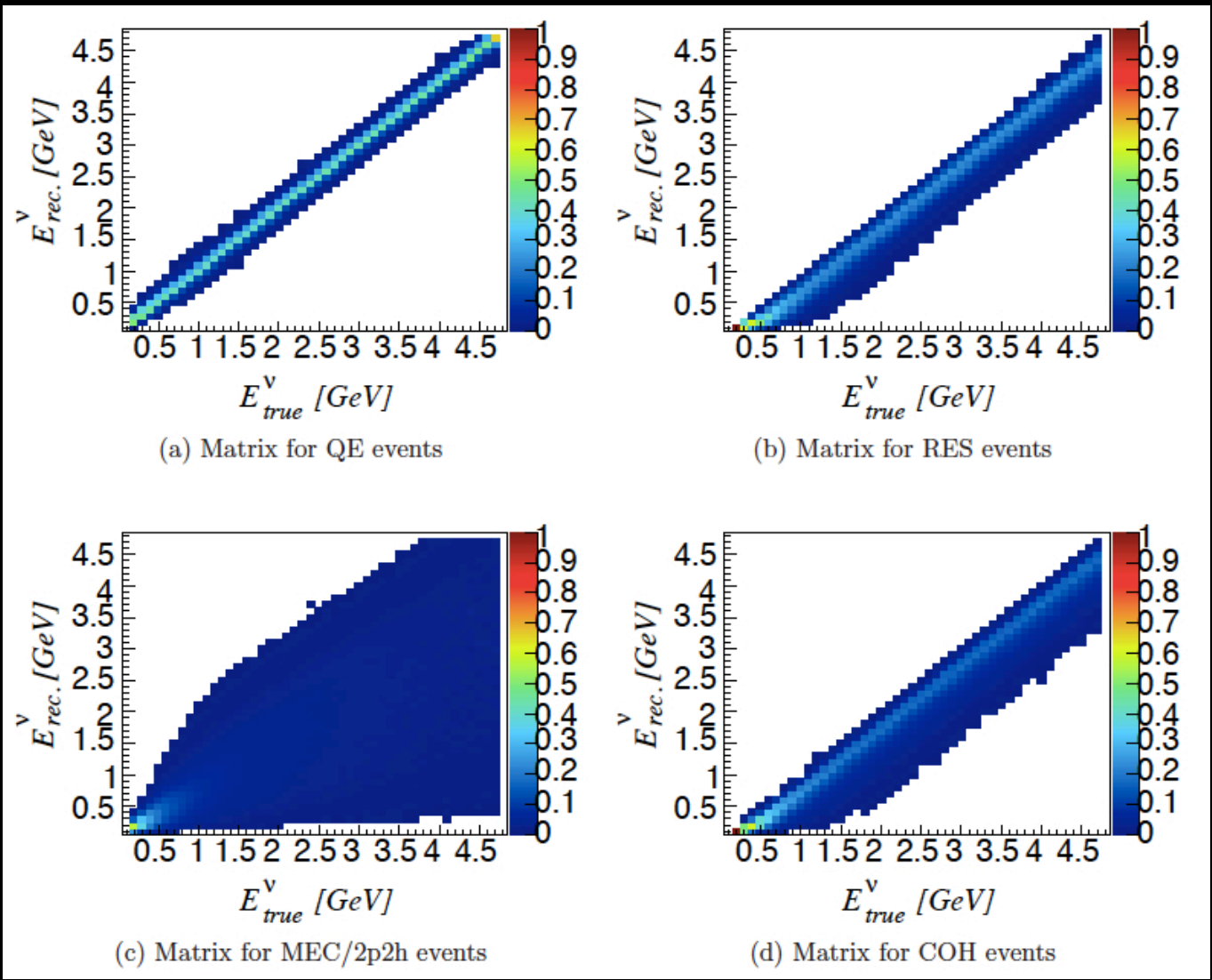
$$N_i^{QE-like} = \sum_j M_{ij}^{QE} N_j^{QE} + \sum_{non-QE} \sum_j M_{ij}^{non-QE} N_j^{non-QE}$$

- However, an intermediate situation would most likely take place:

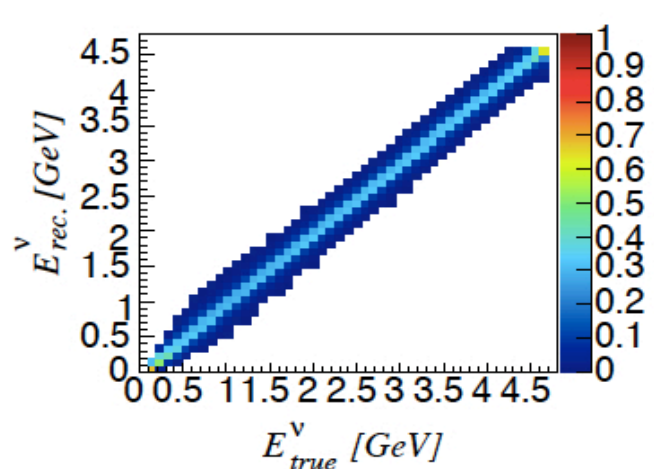
$$N_i^{test}(\alpha) = \alpha N_i^{QE} + (1 - \alpha) N_i^{QE-like}$$

Coloma and Huber, 1307.1243 [hep-ph]

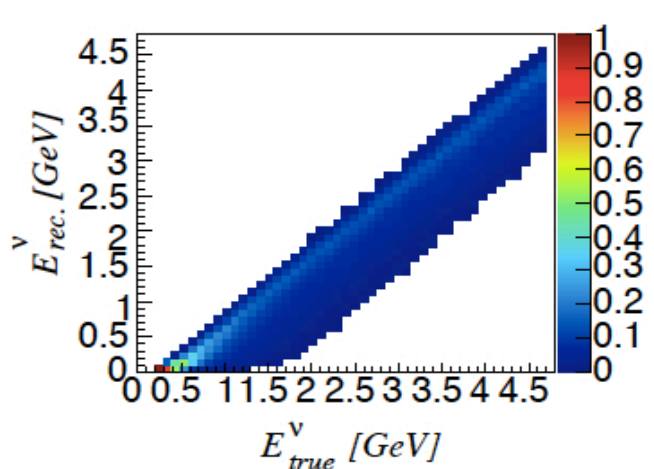
Migration matrices: GiBUU (^{16}O)



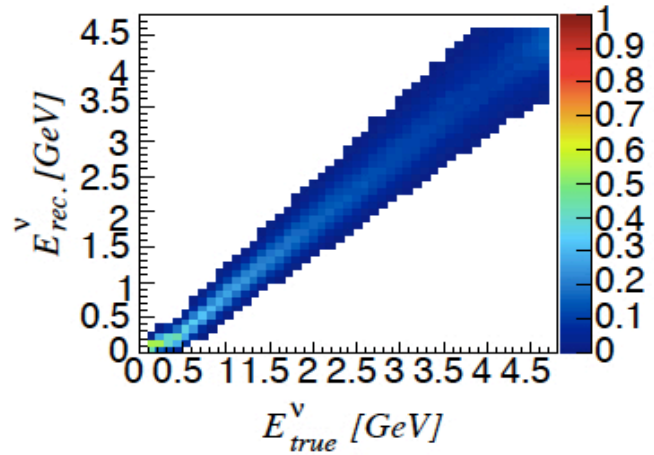
Migration matrices: GENIE (^{16}O)



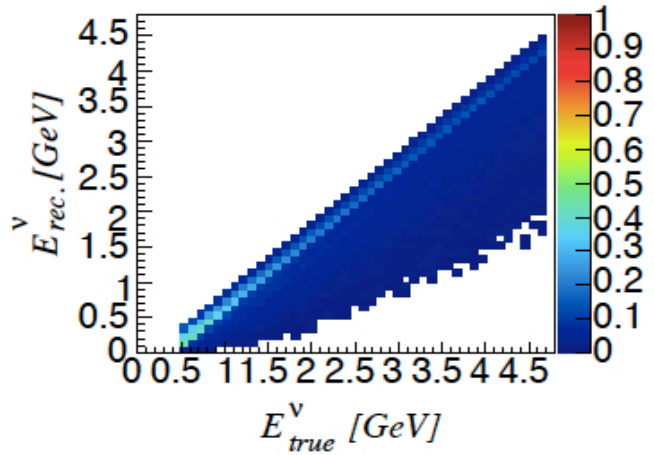
(a) Matrix for QE events



(b) Matrix for RES events

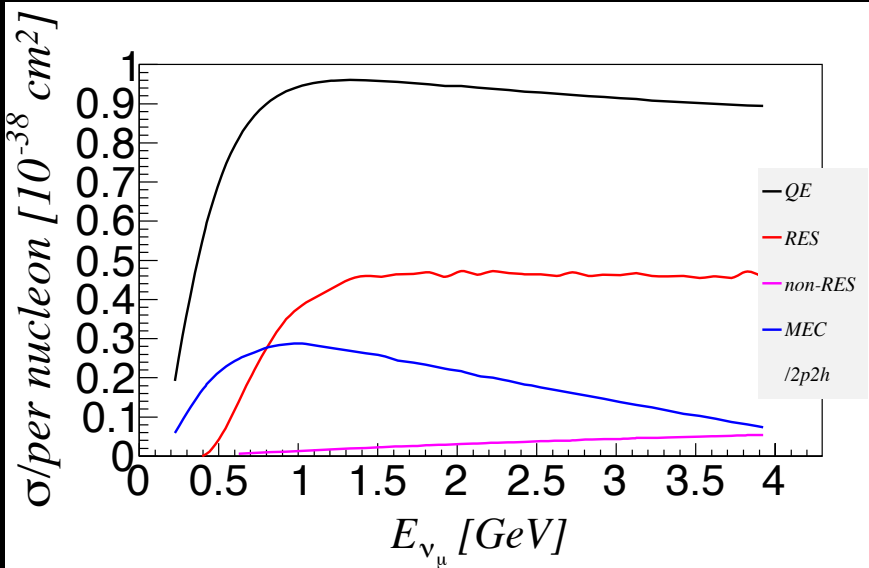


(c) Matrix for MEC/2p2h events

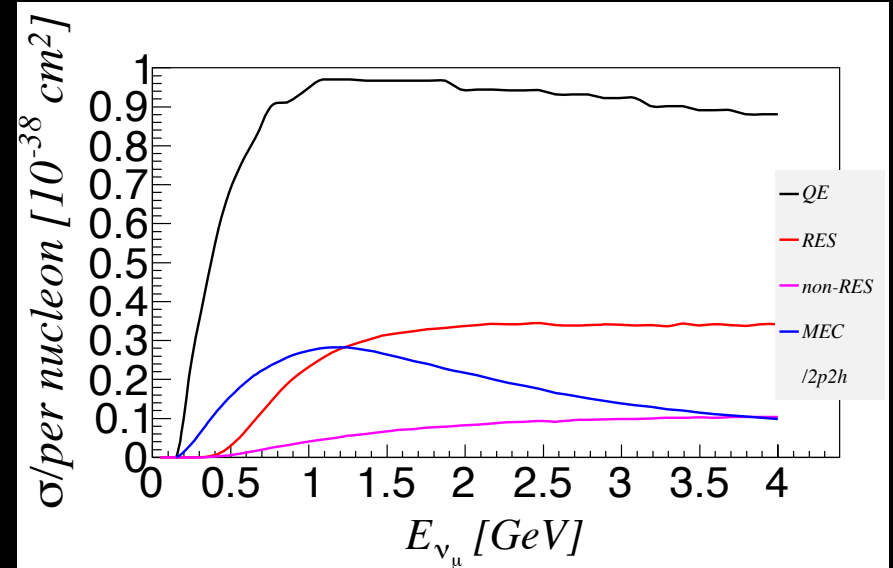


(d) Matrix for COH events

Cross-sections

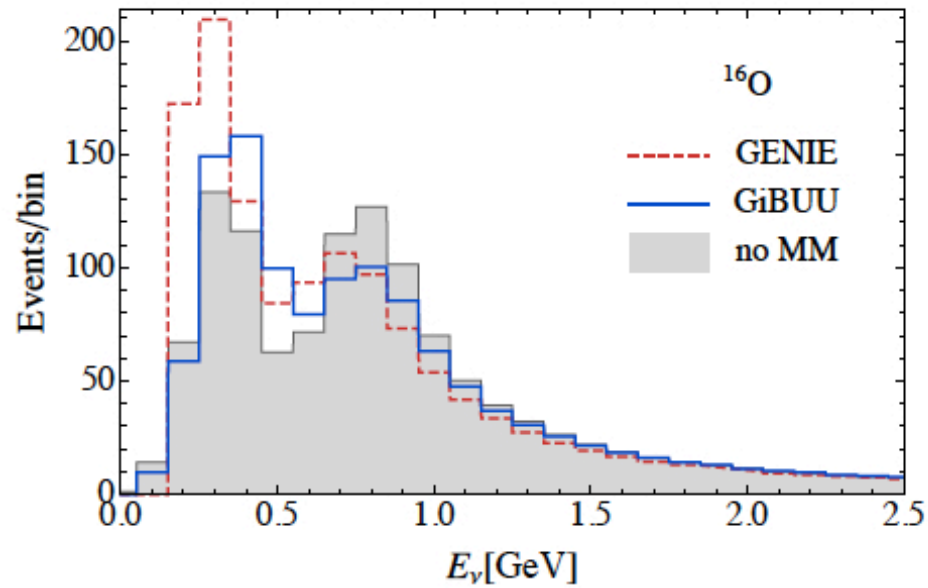


Genie 2.8.0

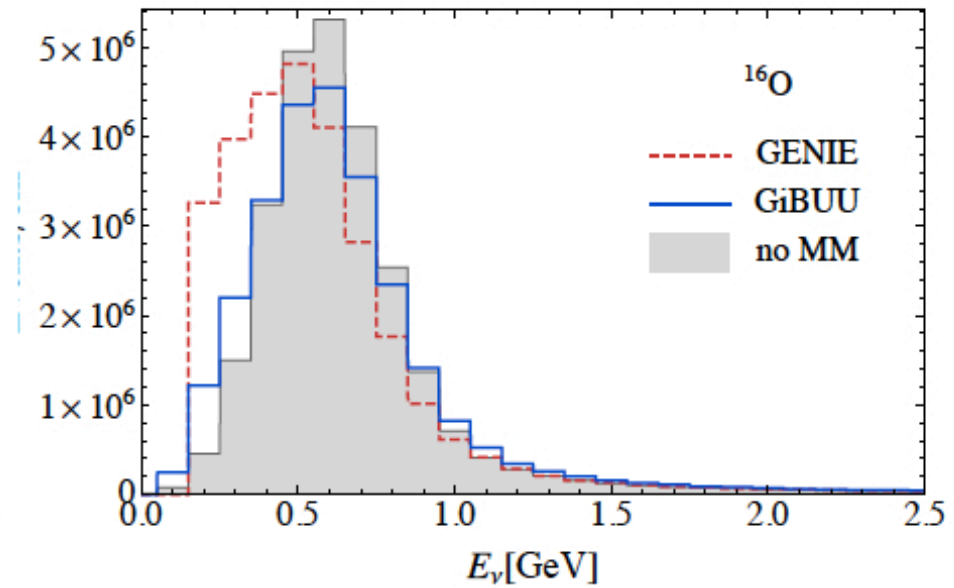


GiBUU 1.6

Event distributions

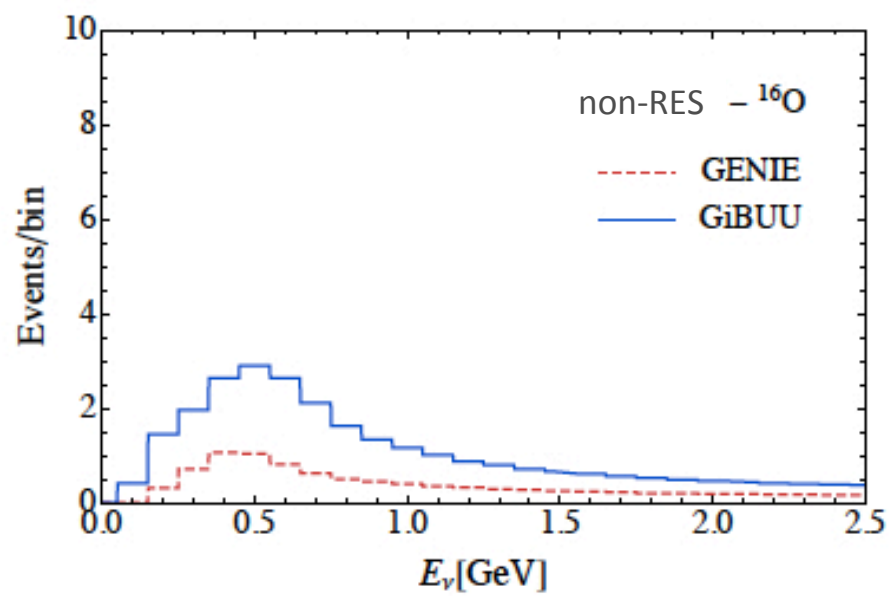
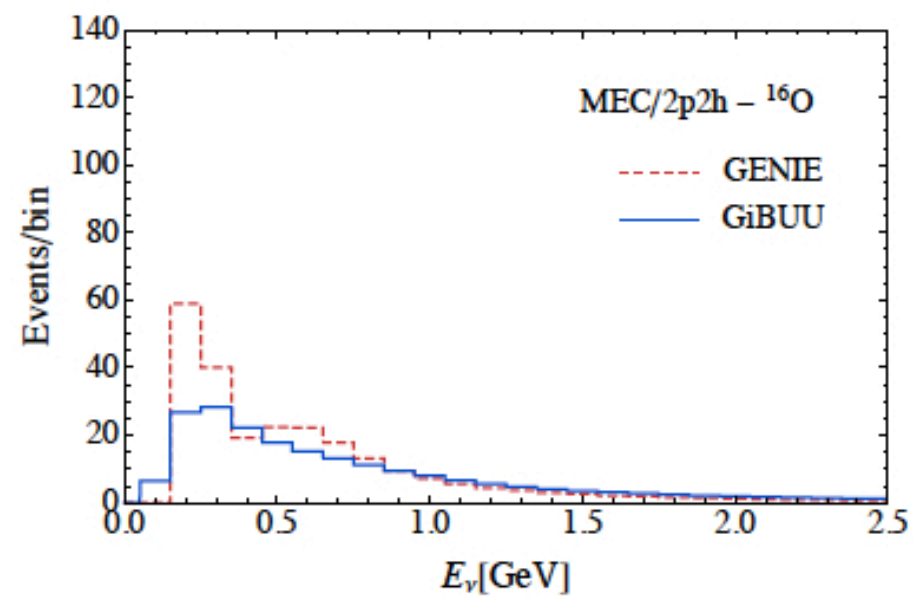
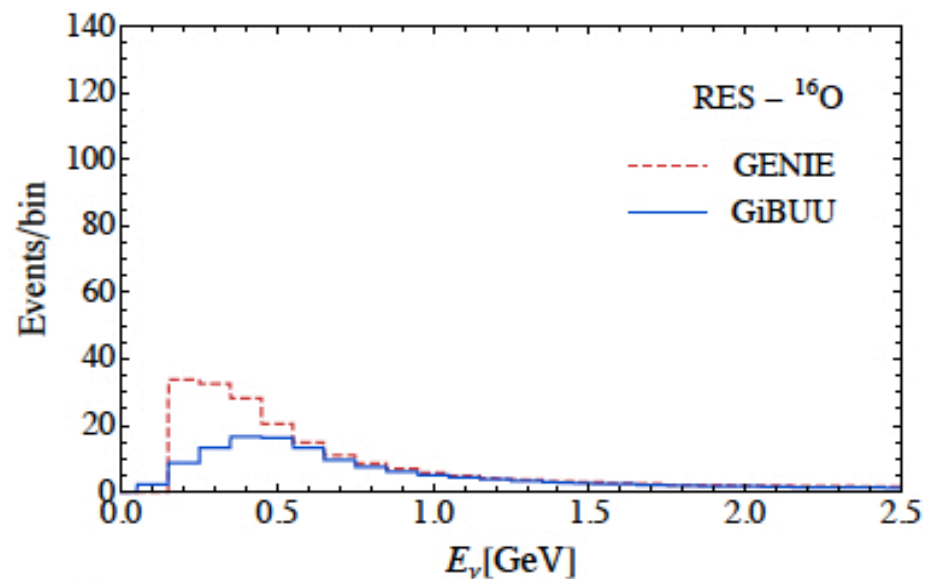
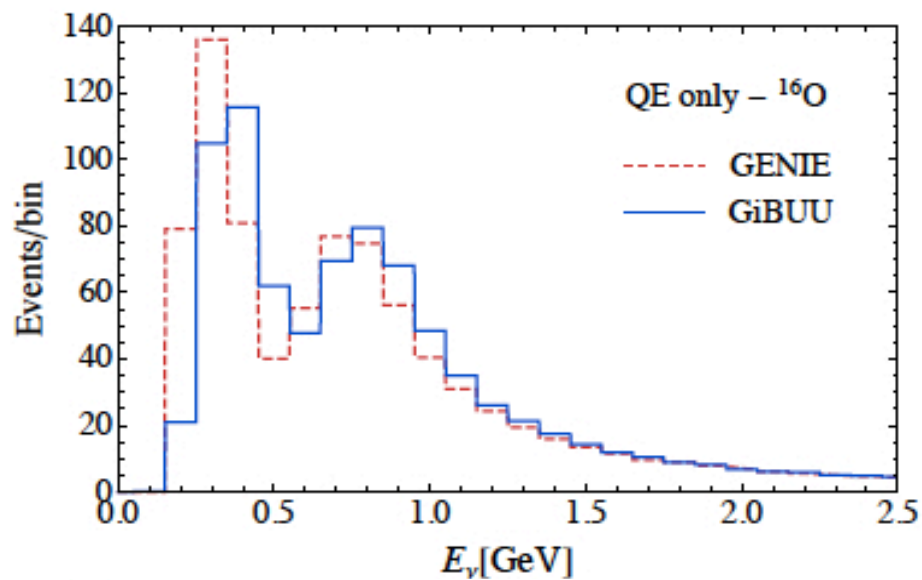


(a) Expected events at the far detector

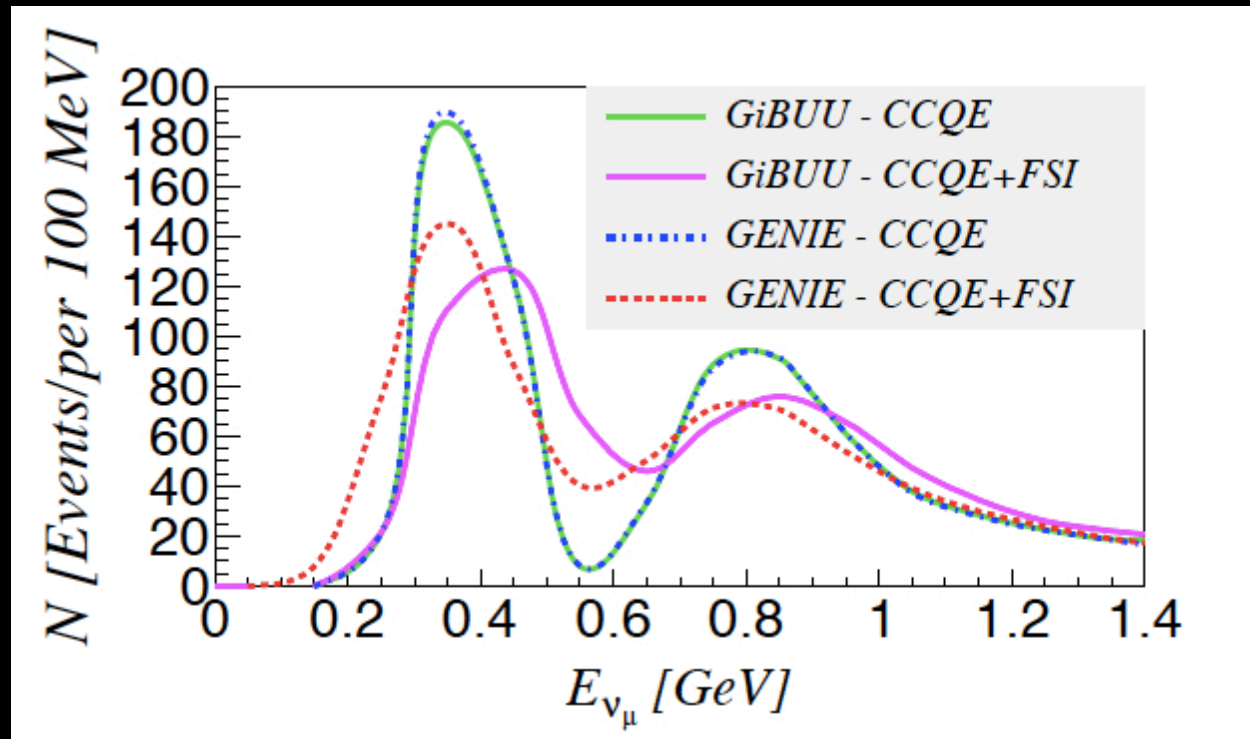


(b) Expected events at the near detector

	QE	RES	non-RES	MEC/2p2h	Total
GiBUU	870	152	32	214	1268
GENIE	877	221	11	249	1358



A surprise ...

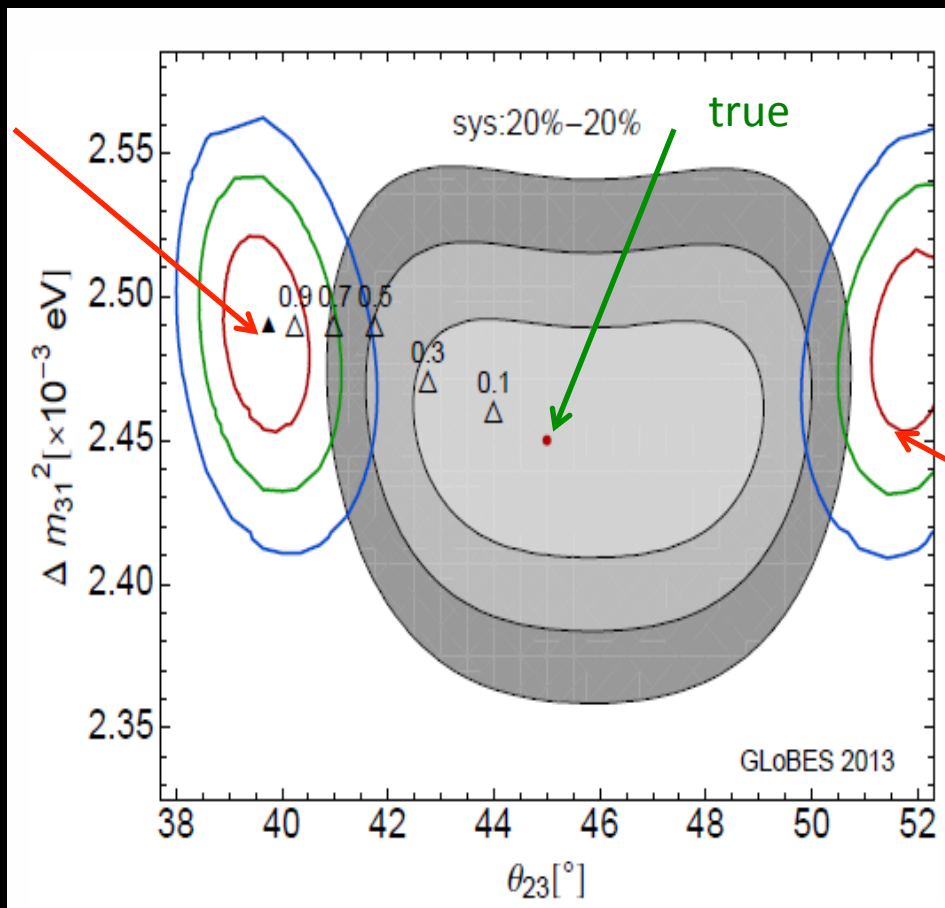


Number of events predicted as function of neutrino energy shifted by 10% for pure QE and 17% for all the QE-like events

- Due to FSI – difference is in the migration matrices
- Intrinsic model differences between GENIE and GiBUU
- Intrinsic differences in the model implementations

How to read the plots

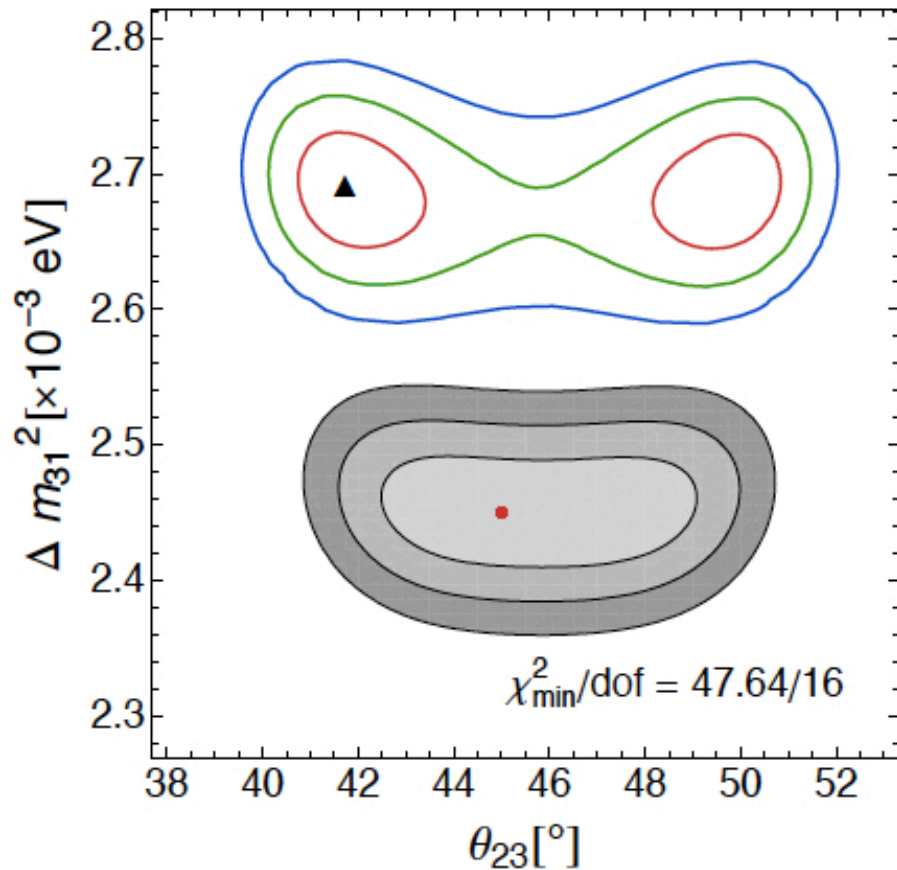
reconstructed
from naive
QE dynamics



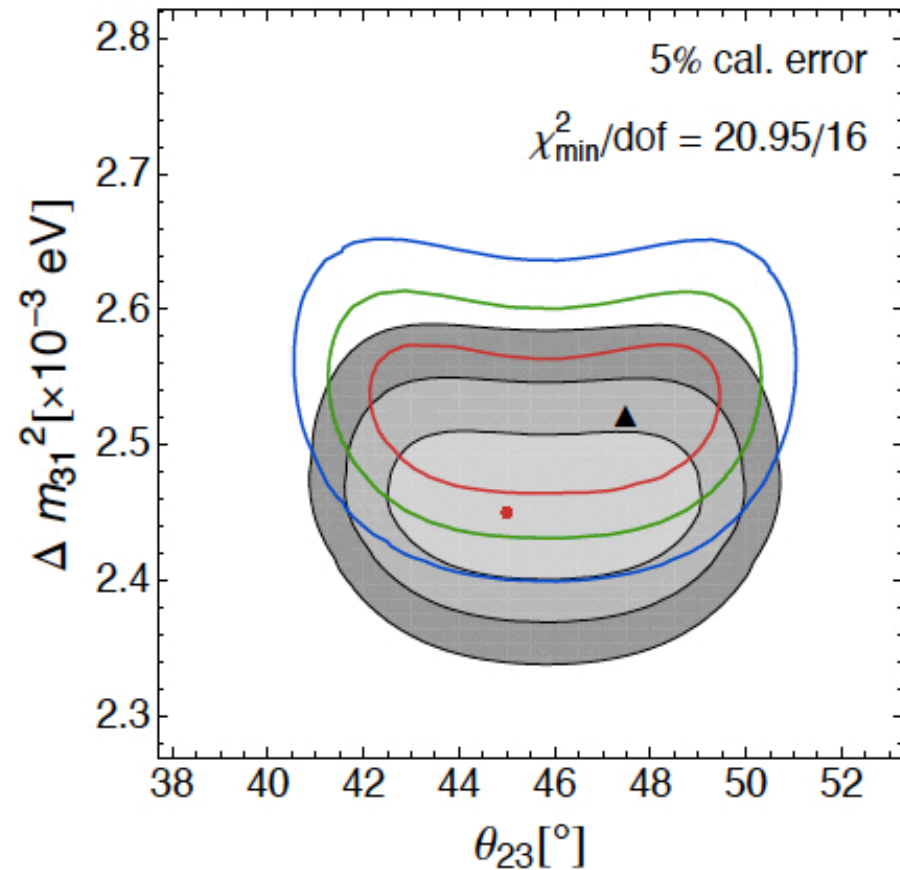
P. Coloma, P. Huber,
arXiv:1307.1243, July 2013
Analysis based on GiBUU

1, 2 and 3 σ allowed regions

Simulating with GiBUU and extracting oscillations with GENIE: with and without calibration error

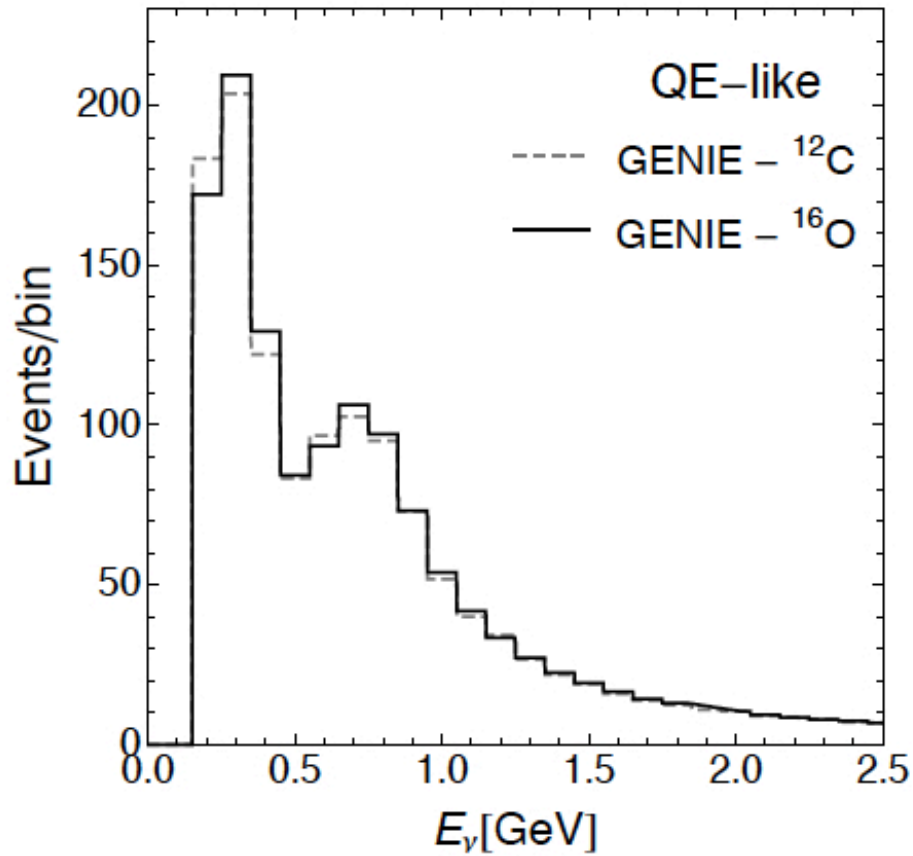


(a) No calibration error

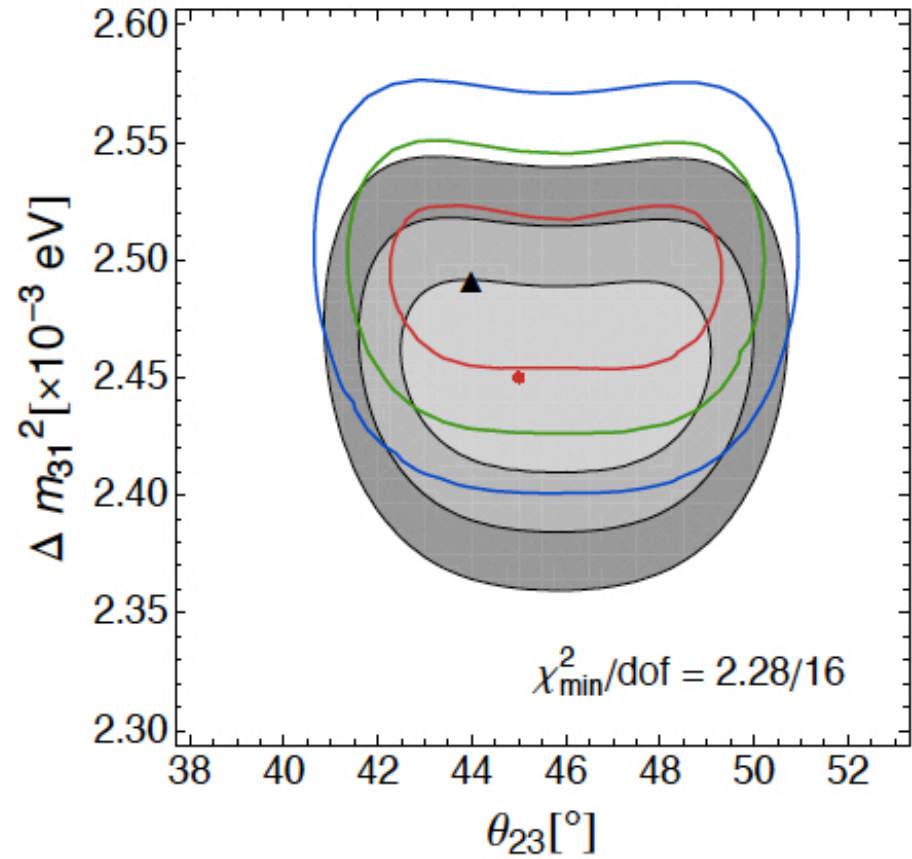


(b) 5% calibration error

Carbon vs Oxygen

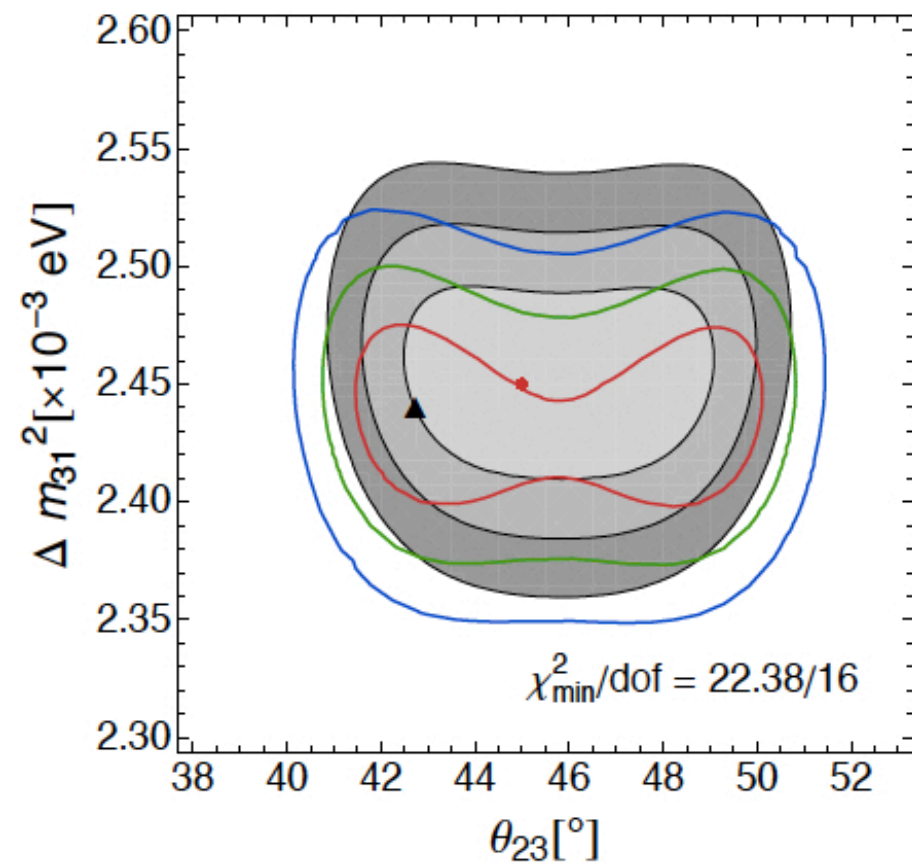


(a) QE-like event distributions

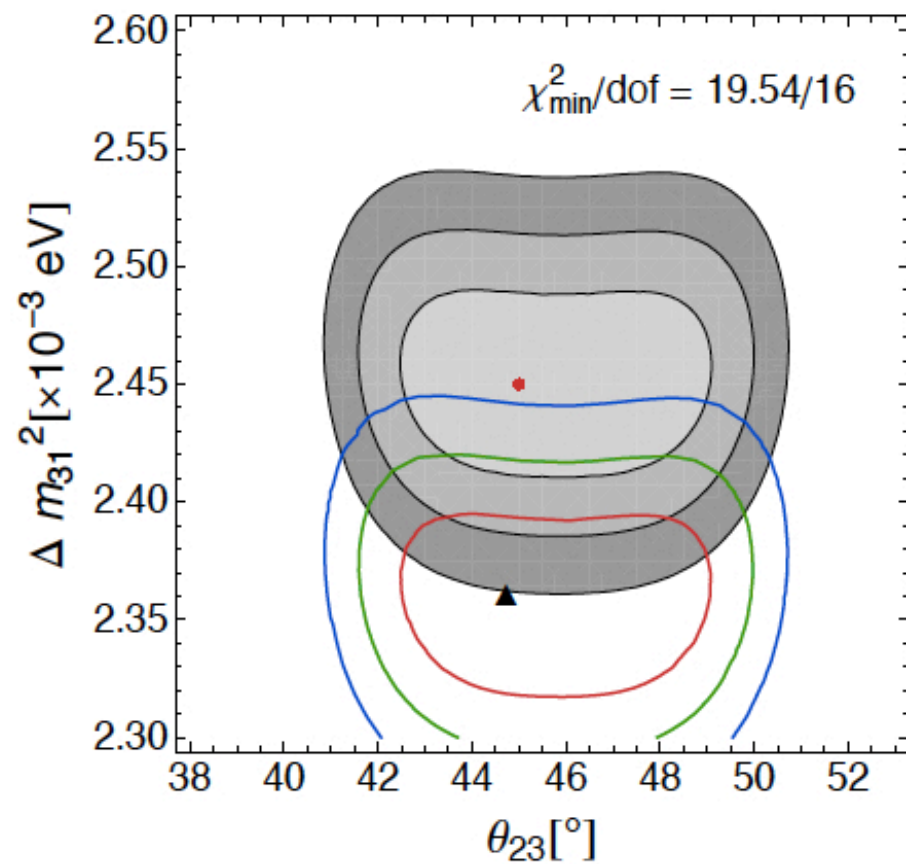


(b) Confidence regions

With and without MEC/2p2h



(a) Results using GiBUU matrices



(b) Results using GENIE matrices

Summary of results

Input “true” Values

$$\begin{aligned}\theta_{12} &= 33.2^\circ & \Delta m_{21}^2 &= 7.64 \times 10^{-5} \text{ eV}^2 \\ \theta_{13} &= 9^\circ & \Delta m_{31}^2 &= 2.45 \times 10^{-3} \text{ eV}^2 \\ \theta_{23} &= 45^\circ & \delta &= 0^\circ\end{aligned}$$

Fitted Values

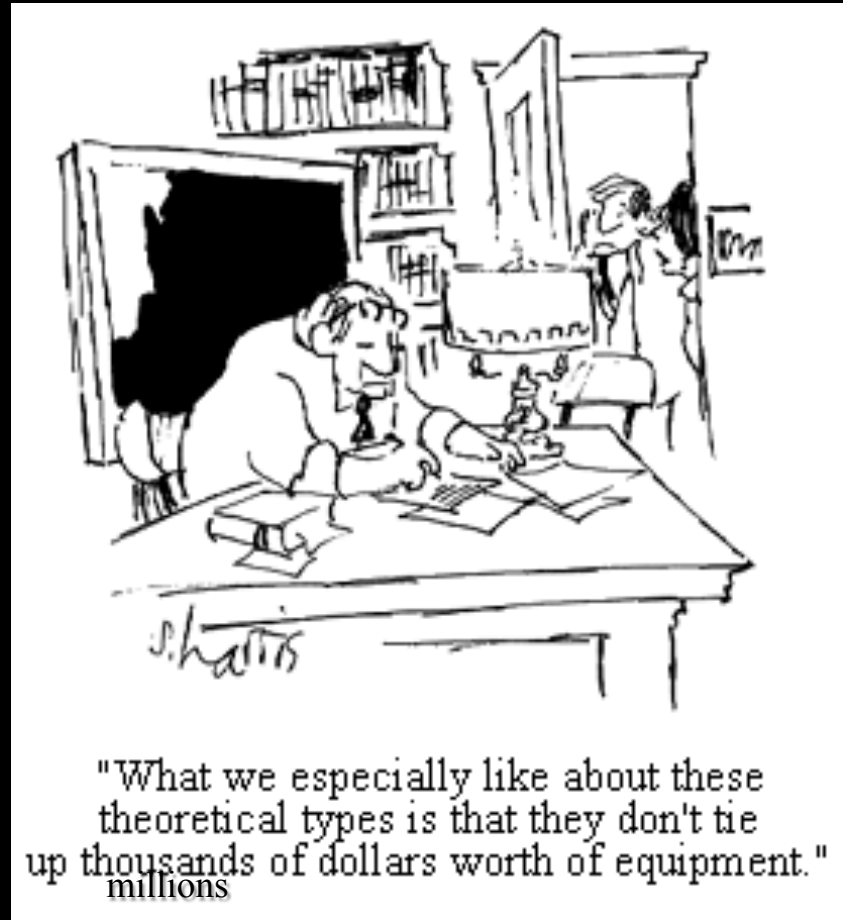
True	Fitted	$\theta_{23,min}$	$\Delta m_{31,min}^2 [\text{eV}^2]$	χ_{min}^2	σ_a
GENIE (^{16}O)	GENIE (^{12}C)	44°	2.49×10^{-3}	2.28	–
GiBUU (^{16}O)	GENIE (^{16}O)	41.75°	2.69×10^{-3}	47.64	–
		47°	2.55×10^{-3}	20.95	5%
GiBUU (^{16}O)	GiBUU (^{16}O) w/o MEC	42.5°	2.44×10^{-3}	22.38	–
GENIE (^{16}O)	GENIE (^{16}O) w/o MEC	44.5°	2.36×10^{-3}	19.54	–

Conclusions

- Energy reconstruction essential for precision determination of neutrino oscillation parameters and neutrino-hadron cross sections
- Impact on neutrino oscillation experiments due to nuclear models, what they are and how they are implemented is not negligible (order 10%)
 - comparing systematically generators is important
 - neutrino event generators use almost same data set so there are correlations that are non-negligible
 - using wrong models affect neutrino oscillation parameters determination

- In future extend the case to CP violation:
 - neutrino vs anti-neutrino cross-section, : do we have reliable event generators for anti-neutrino ?
- MEC/2p2h: heavily tuned to MiniBooNE data in both GiBUU and GENIE.
 - the contribution of MEC/2p2h is the same between Carbon and Oxygen, should it be ?
- Energy reconstruction requires reliable event generators, of same quality as experimental equipment
- Precision era of neutrino physics requires much more sophisticated generators and a dedicated effort in theory
- Theorists-phenomenologists and experimentalists need to work together: NuSTEC

Generators are a crucial part
of any experiment!
Must be of same quality as the
experimental equipment itself!
Needed resources are relatively
small, but still not available



Thanks

- Omar: for everything (where should I start ?)
- Ulrich and Olga:
 - very interesting discussion
 - GiBUU support
- Patrick and Pilar:
 - GIOBES
 - oscillation analysis
 - great discussions

Neutrino generators and oscillation analyses: discussion and questions

- Systematic errors
 - Uncertainties in input cross sections
 - Mis-identification of reaction mechanisms
 - Generator-specific numerical implementation
- What to do next
- Systematical comparison of generators
- Could and should we extend these studies and the nuclear theories to heavier targets (Ar) ?

Backup

Energy shift

$$N[E] \rightarrow N[(1 + a) E]$$

Modify the number of events as function of energy introducing a calibration error “a” and additional pull term is added to the χ^2 of the fit.