

Radiative Capture Reactions in Lattice Effective Field Theory



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Nuclear Reactions from Lattice QCD, INT March 12, 2013

Outline

- Motivation
- Continuum EFT for reactions
- Lattice EFT for reactions

Reaction theory:

- Reaction theory for nuclear experiments, e.g. FRIB
- Nuclear astrophysics where data might be lacking
- Reactions are more fun than static properties

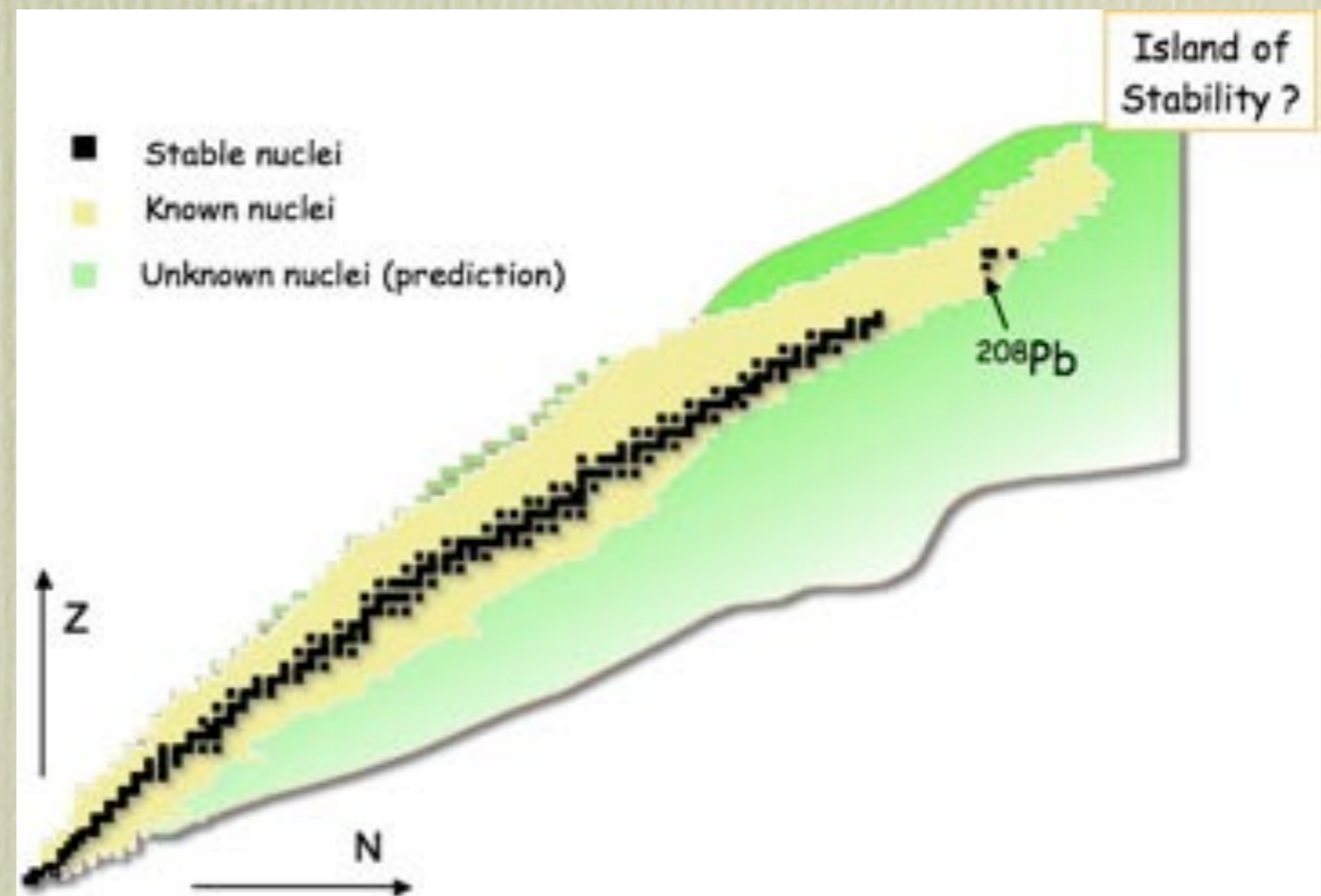
Examples: $d(p, \gamma)^3\text{He}$, $^7\text{Be}(p, \gamma)^8\text{B}$, $^{14}\text{C}(n, \gamma)^{15}\text{C}, \dots$,

the list goes on

Only halo systems for now ...

Halo systems

- Characterized small neutron/proton separation energy.
Large size.
- Interesting three-body physics: All-bound, Tango, Samba (^{12}Be), Borromean (^{11}Li)
- Exotic physics near the drip line



Some details on ${}^7\text{Li}(n, \gamma){}^8\text{Li}$

- Isospin mirror systems ${}^7\text{Li}(n, \gamma){}^8\text{Li} \leftrightarrow {}^7\text{Be}(p, \gamma){}^8\text{B}$
- Inhomogeneous BBN

Whats the theoretical error?

EFT

- Identify degrees of freedom

$$\mathcal{L} = c_0 O^{(0)} + c_1 O^{(1)} + c_2 O^{(2)} + \dots$$

Hide UV ignorance



IR explicit

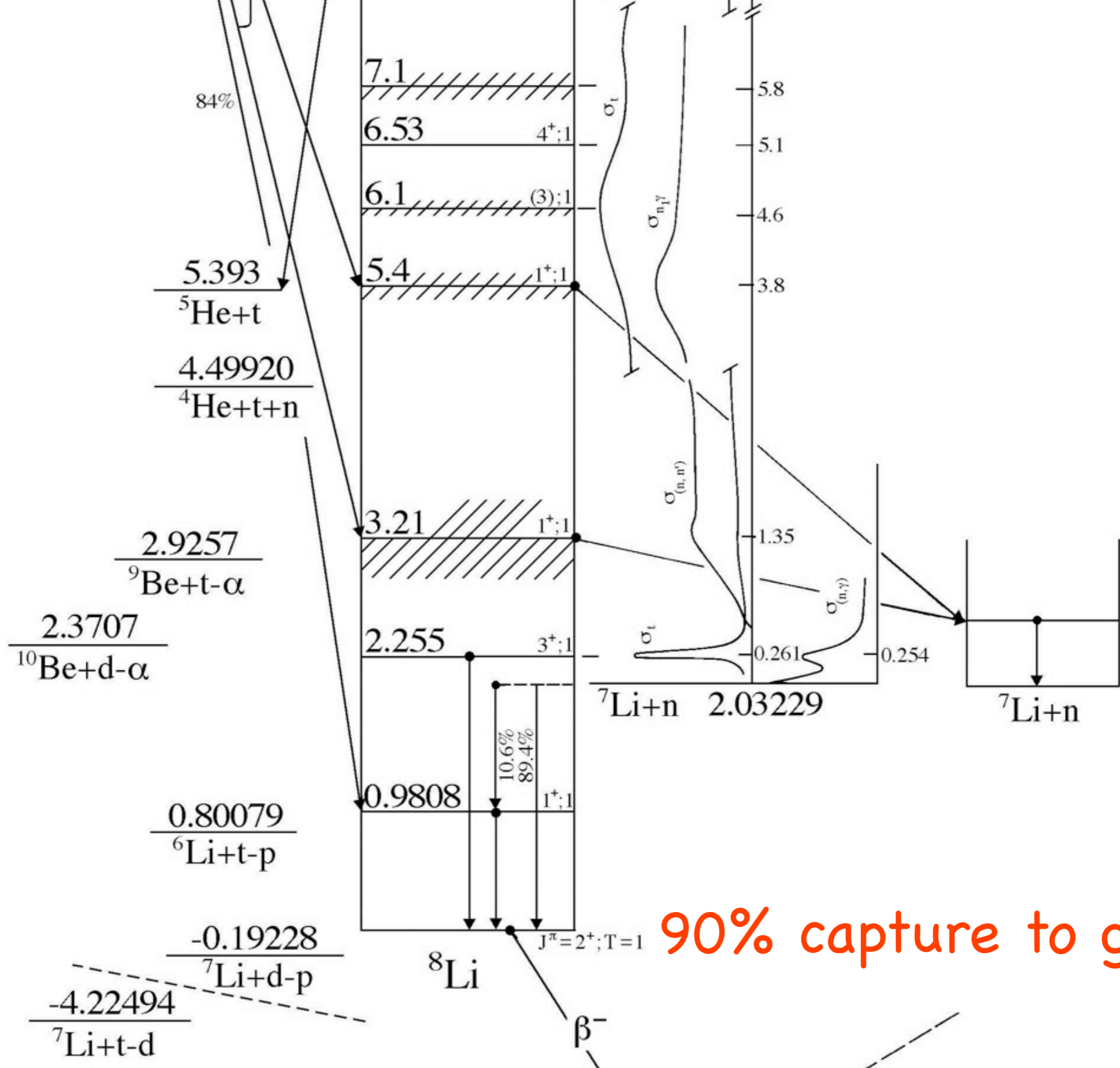


- Determine c_n from data (elastic, inelastic)

- EFT ERE + currents + relativity

Not just Ward identity





$\frac{11.747}{^5\text{He}+p+n}$

$\frac{9.9754}{^6\text{He}+p}$

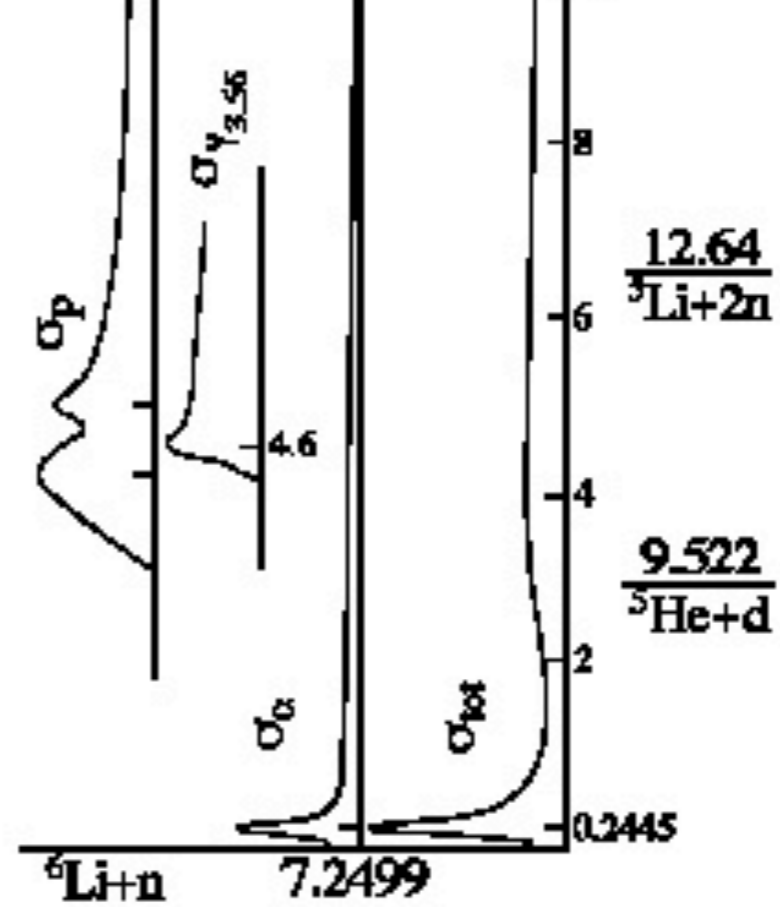
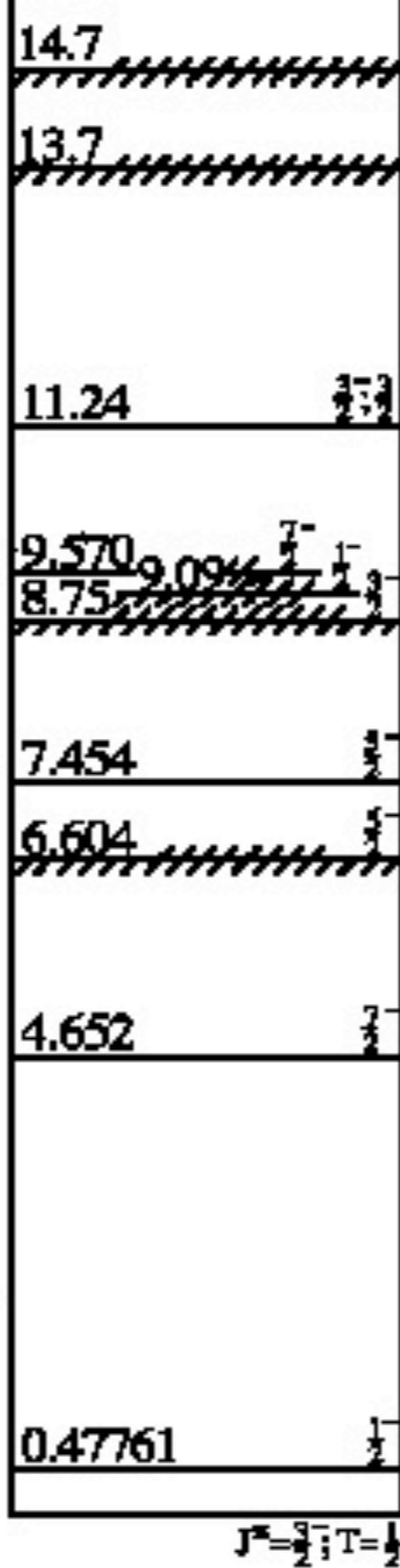
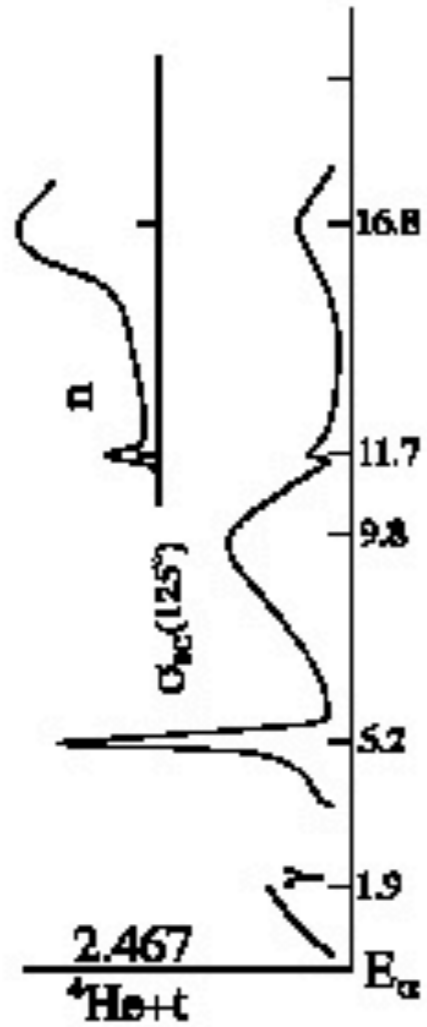
$\frac{7.1507}{^9\text{Be}+d-\alpha}$

$\frac{5.0254}{^6\text{Li}+d-p}$

$\frac{0.9927}{^6\text{Li}+t-d}$

$\frac{-17.3468}{^4\text{He}+^4\text{He}-p}$

$\frac{-133.1026}{^6\text{Li}+p-\pi^+}$



**p=84 MeV
excited state**

$\frac{2.7895}{^{10}\text{B}+n-\alpha}$

$\frac{0.8618}{^7\text{Be}}$

Cohen, Gelman, van Kolck '04

$\frac{-8.7564}{^{11}\text{B}+\alpha-^9\text{Be}}$

$\frac{-11.2023}{^9\text{Be}+p-^3\text{He}}$

^7Li

12-01

Look at E1 transition

-- Initial state: s-wave

single operator fitted to scattering length a

-- Ground state: p-wave

p-wave needs two operators (a_V, r_1)

We also include the excited state and the p-wave 3^+ resonance (M1 transition)

$$i\mathcal{A}(p) = \frac{2\pi}{\mu} \frac{i}{p \cot \delta_0 - ip} = \frac{2\pi}{\mu} \frac{i}{-1/a + \frac{r}{2}p^2 + \dots - ip}$$

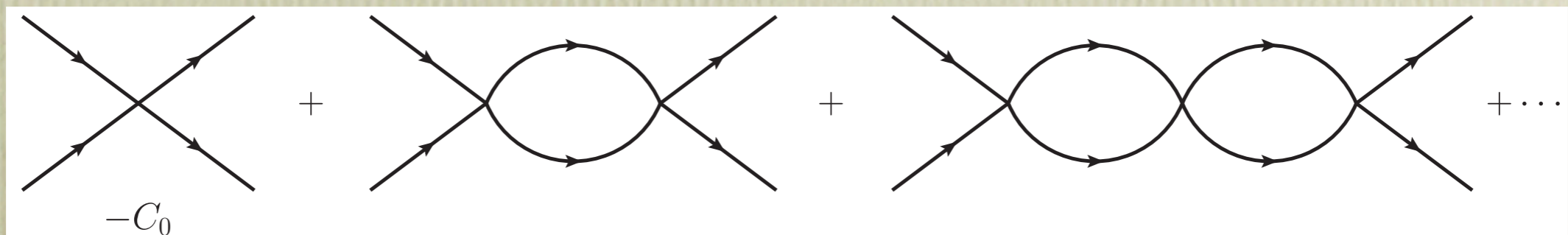
--- Natural case $a, r \sim 1/\Lambda \ll 1/p$

expand in small p , EFT perturbative

--- Large scattering length $a \gg 1/\Lambda$

$$i\mathcal{A}(p) \approx -\frac{2\pi}{\mu} \frac{i}{1/a + ip} \left[1 + \frac{1}{2} \frac{rp^2}{1/a + ip} + \dots \right]$$

EFT non-perturbative



$$i\mathcal{A}(p) = \frac{-i}{\frac{1}{C_0} + i\frac{\mu}{2\pi}p} \Rightarrow C_0 = \frac{2\pi a}{\mu}$$

Weinberg '90
 Bedaque, van Kolck '97
 Kaplan, Savage, Wise '98

p-wave $i\mathcal{A}(p) = \frac{2\pi}{\mu} \frac{ip^2}{-\frac{1}{a_V} + \frac{r_1}{2}p^2 + \dots - ip^3}$

Shallow systems

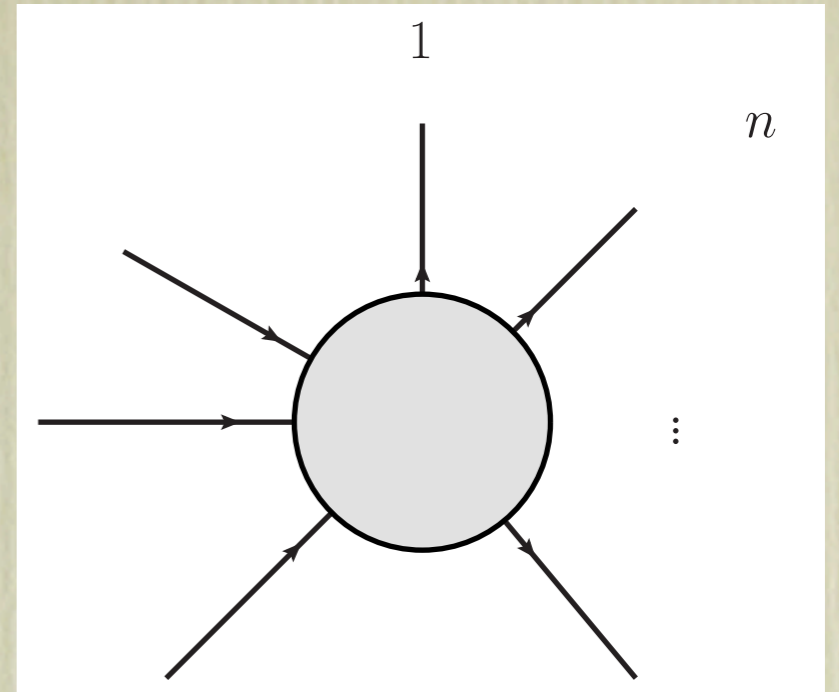
2 fine tuning Bertulani, Hammer, van Kolck '02

1 fine tuning Bedaque, Hammer, van Kolck '03

Requires two non-perturbative operators at LO

Residues and poles n-point function

LSZ reduction: $G^{(n)} \sim \prod_{i=1}^n \sqrt{Z_i}$



Consider a scalar theory

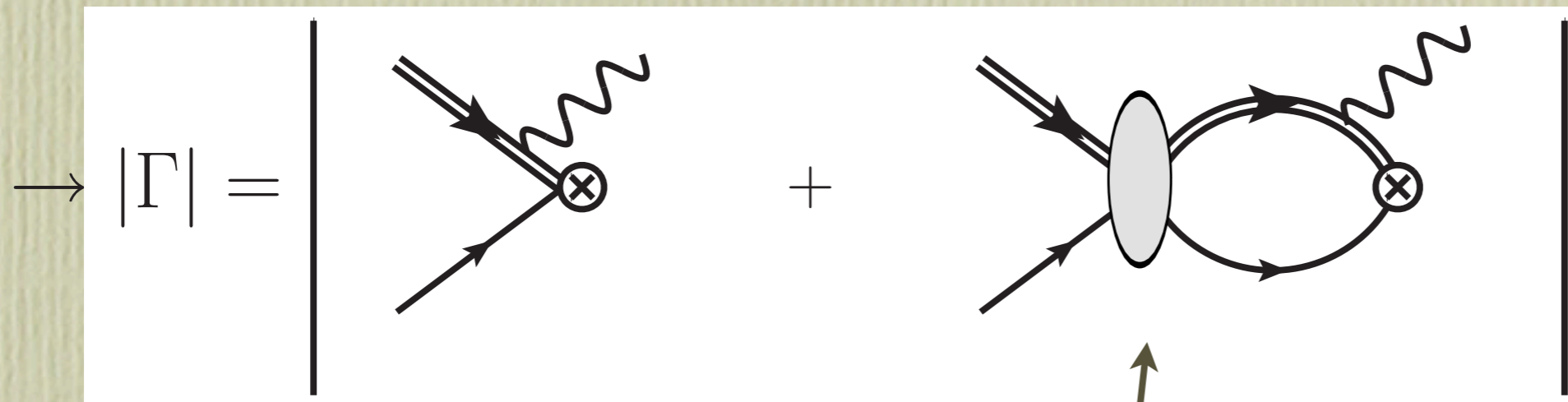
$$\begin{array}{c}
 \text{---} \rightarrow \text{---} + \text{---} \rightarrow \text{---} \text{---} \text{---} \text{---} + \text{---} \rightarrow \text{---} \text{---} \text{---} \text{---} \text{---} \text{---} + \text{---} \rightarrow \text{---} \text{---} \text{---} \text{---} \text{---} \text{---} \text{---} + \dots = \frac{Z}{p_0^2 - p^2 - m^2}
 \end{array}$$

For deuteron: $\frac{Z_d}{p_0 - \frac{p^2}{4M} + \frac{\gamma^2}{M}}, \quad Z_d = \frac{8\pi\gamma}{M^2} \frac{1}{1 - \rho\gamma}$

$$\psi_d(r) = \sqrt{\frac{\gamma}{2\pi} \frac{1}{1 - \rho\gamma}} \frac{e^{-r\gamma}}{r} = \sqrt{\frac{\mu^2}{4\pi^2} Z_d} \frac{e^{-r\gamma}}{r}$$

Capture Cross Section

$$\frac{d\sigma}{d\cos\theta} = \frac{1}{32\pi s} \frac{k}{p} |\Gamma|^2$$

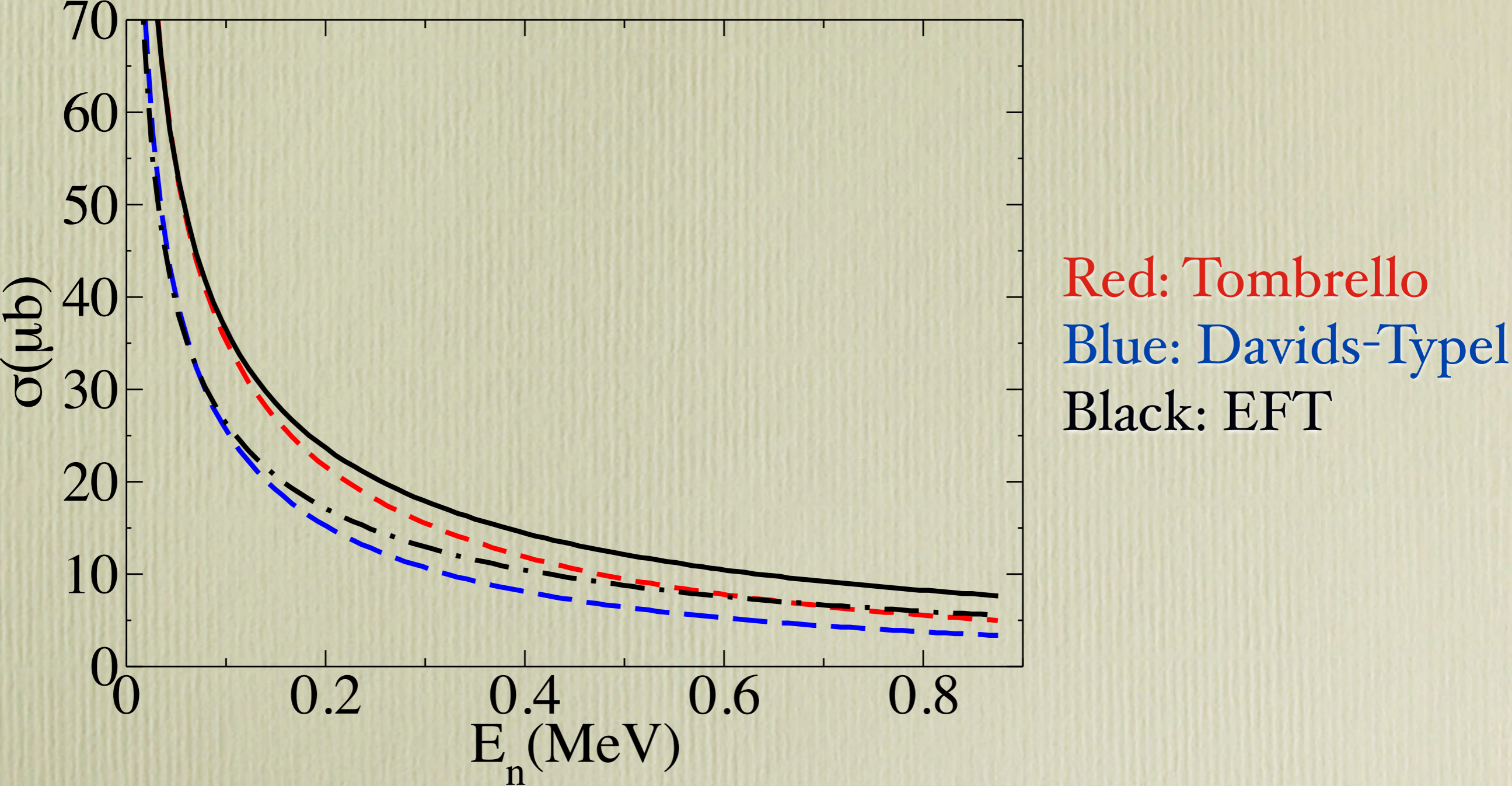


gives $1/p$ dependence

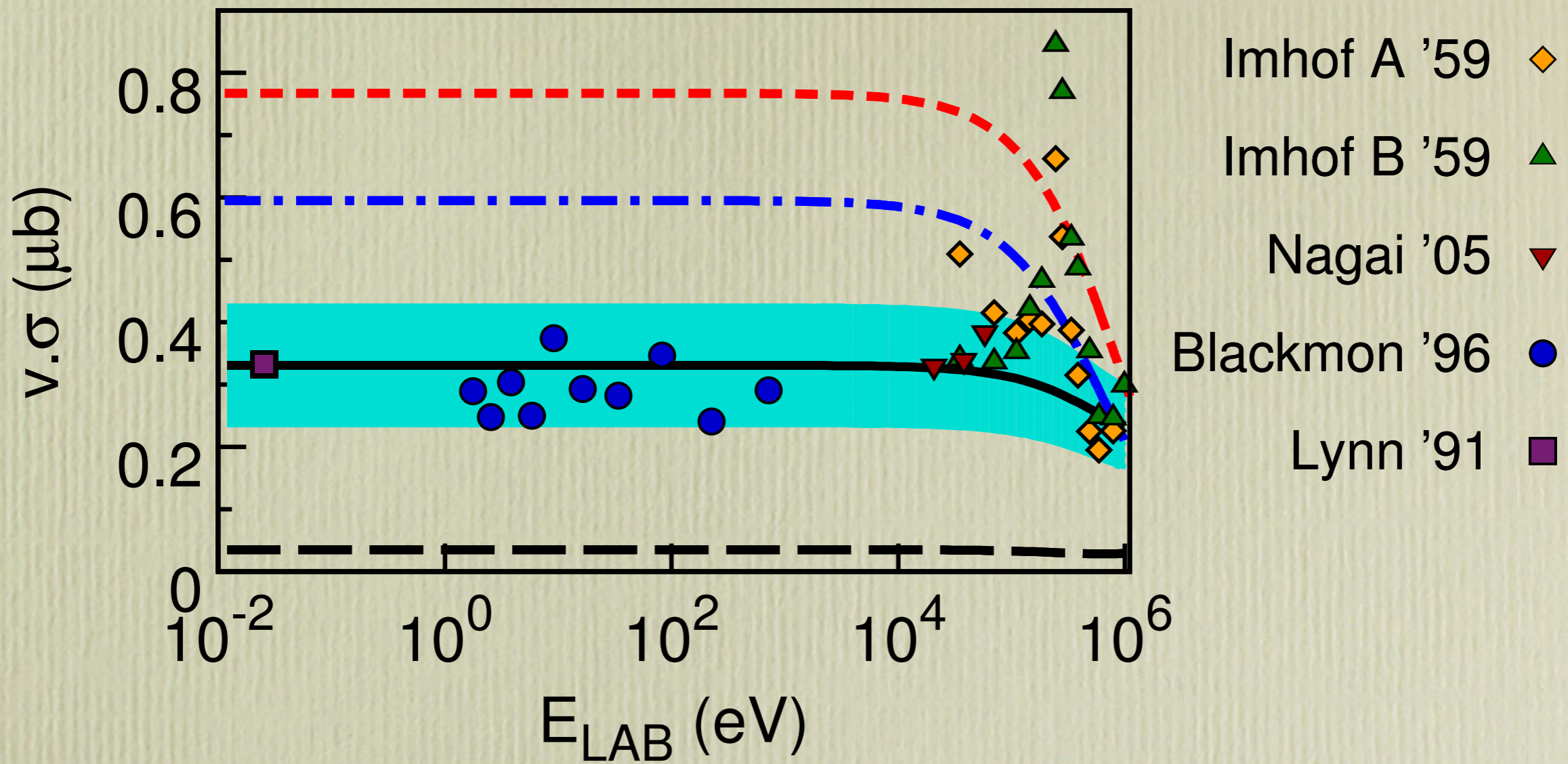
Analytic result, depends on $\sqrt{Z} = \frac{1}{\sqrt{2}} \sqrt{\frac{1}{1 + 3\gamma/r_1}}$

Need r_1 at leading order

“Effective range” contribution



Rupak, Higa; PRL 106, 222501 (2011)



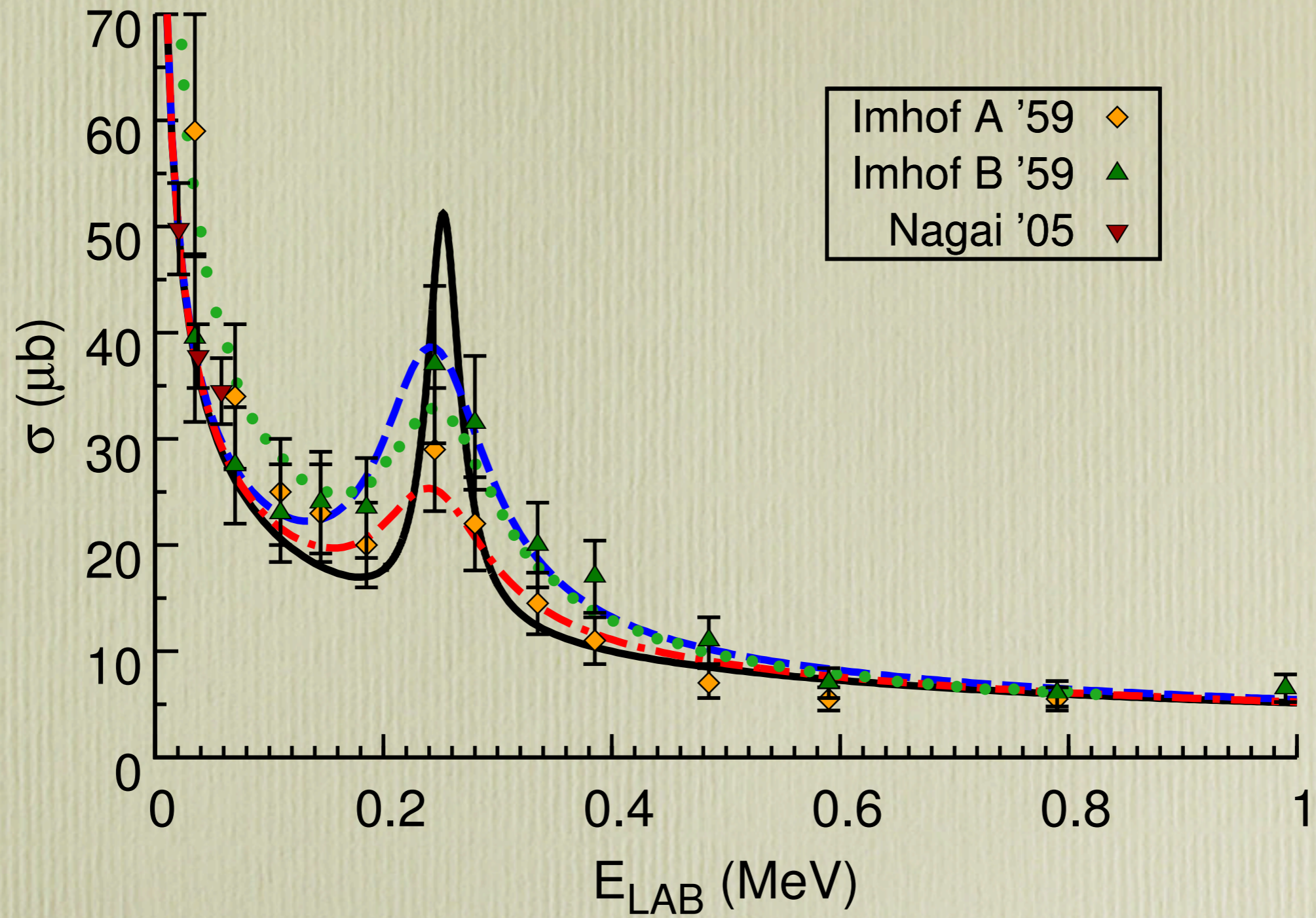
Red: Tombrello

Blue: Davids-Typel

Black: EFT

Rupak, Higa; PRL 106, 222501 (2011)

Fernando, Higa, Rupak; EPJA 48, 24 (2012)



Fernando, Higa, Rupak; EPJA 48, 24 (2012)

Lessons learned

--- Tuning potential to reproduce bound state energy is not sufficient to get the wave function renormalization constant.

--- In the strong sector directly applies to ${}^7\text{Be}(p, \gamma){}^8\text{B}$

General problem : How to constrain low-energy nuclear theory?

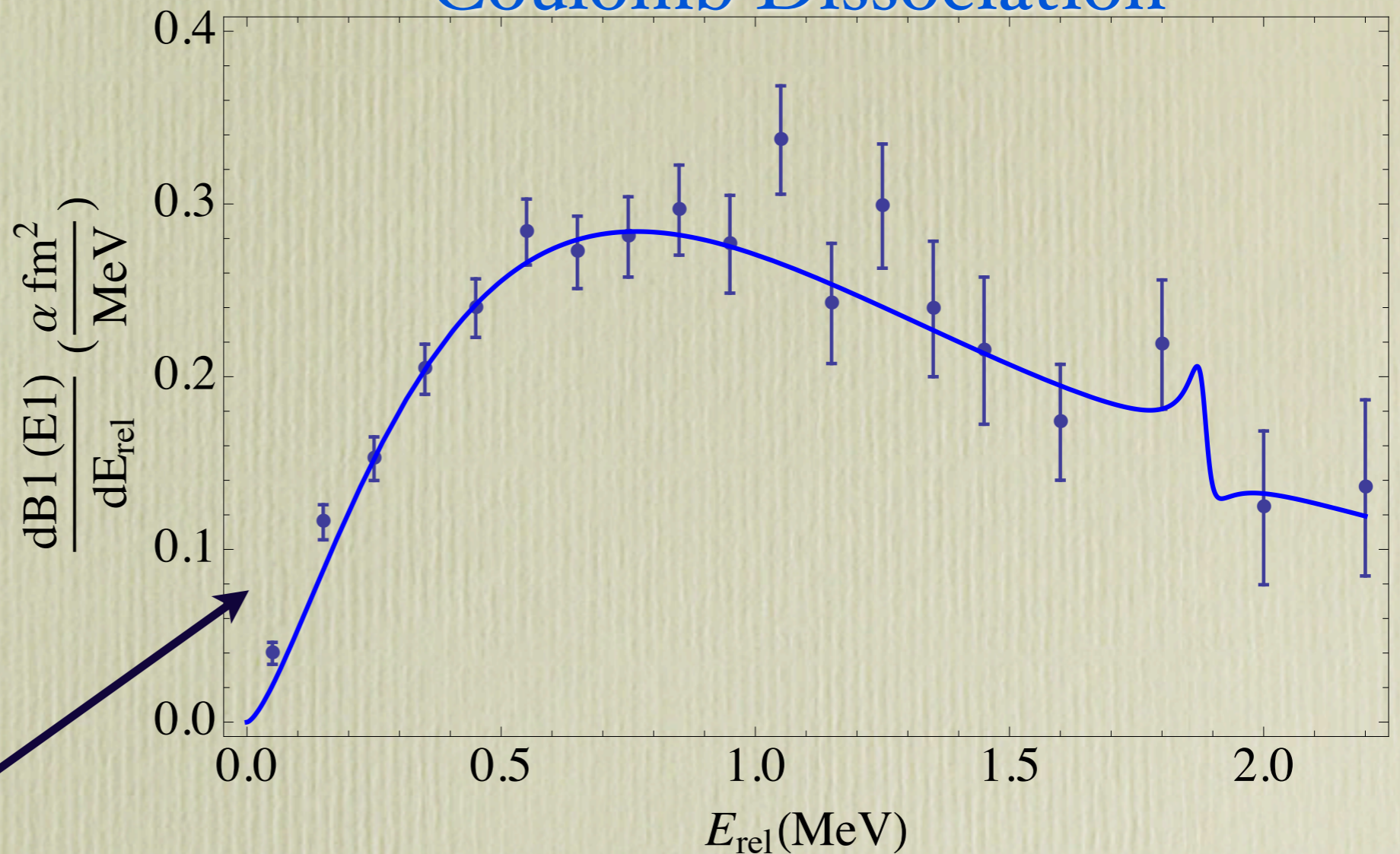
Neutron Capture/Coulomb Dissociation on Carbon-14

Coulomb Dissociation

Power counting

$$a_1 = -n_1/Q^3$$

$$r_1 = 2n_2Q$$

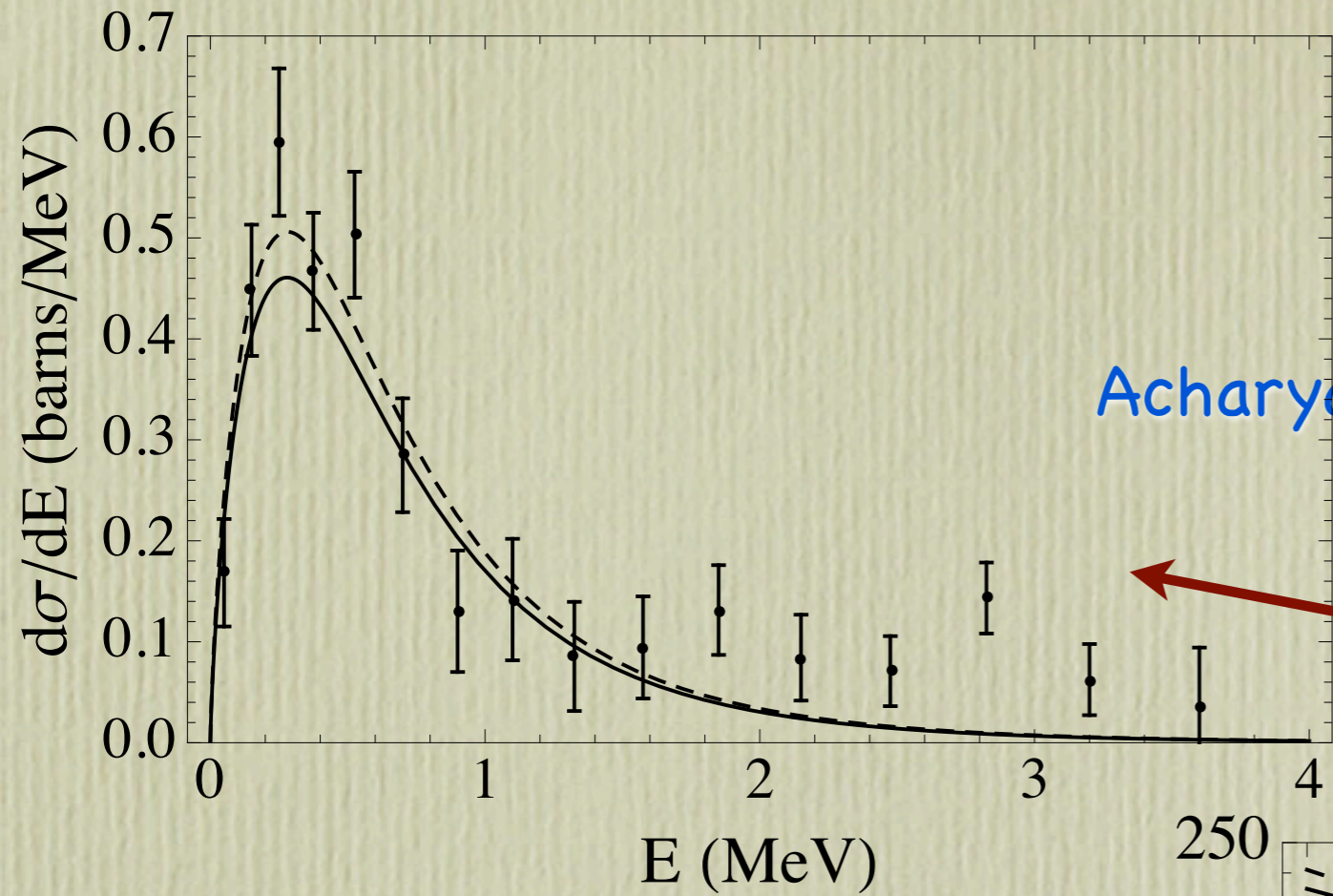


$n_1=0.7$, $n_2=1$,
 $Q=40 \text{ MeV}$

Rupak, Fernando, Vaghani, PRC 86, 044608 (2012)

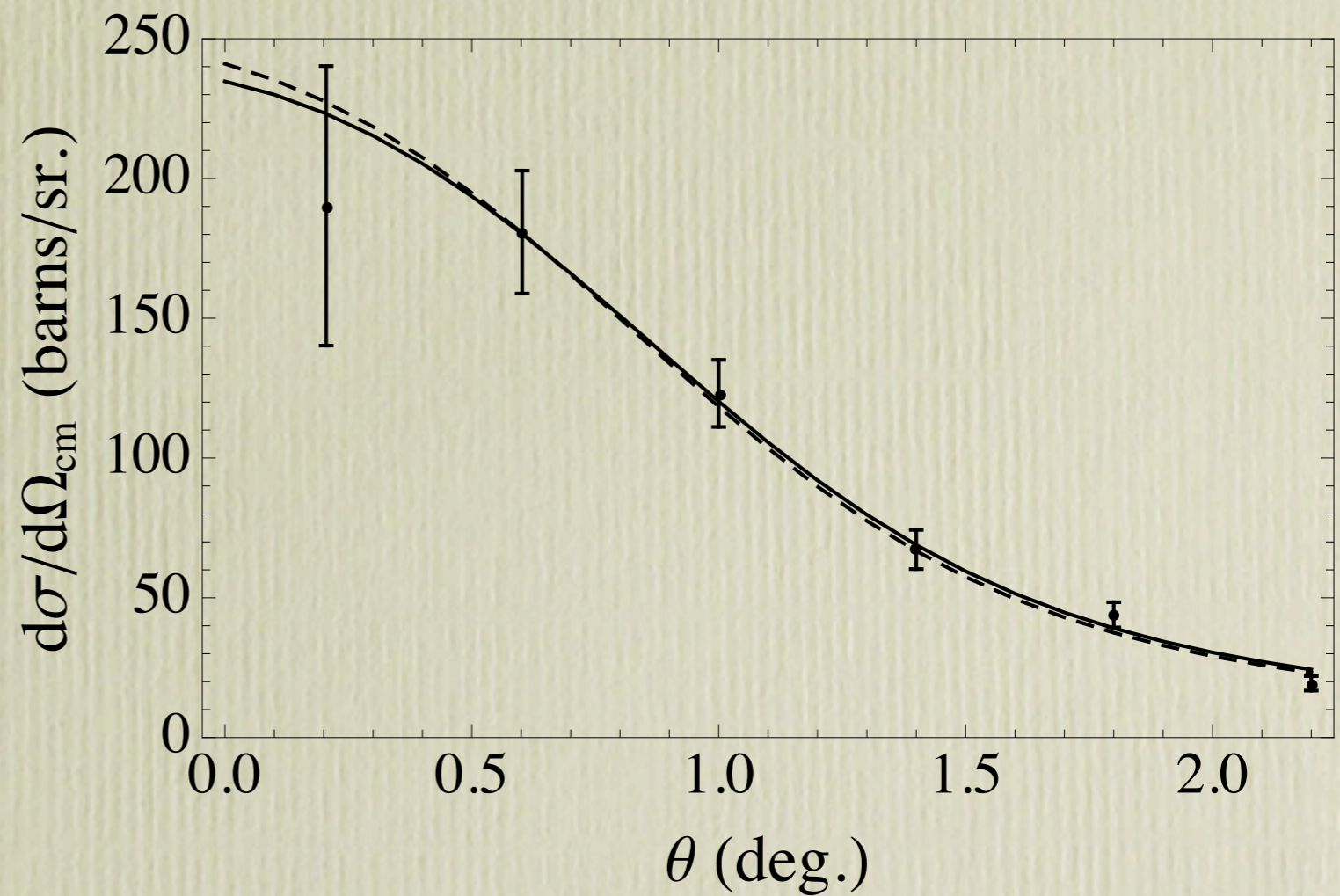
Data: T. Nakamura et al., PRC, 79, 035805 (2009)

Coulomb Dissociation of Carbon-19



Acharya & Phillips, arxiv:1302.4762

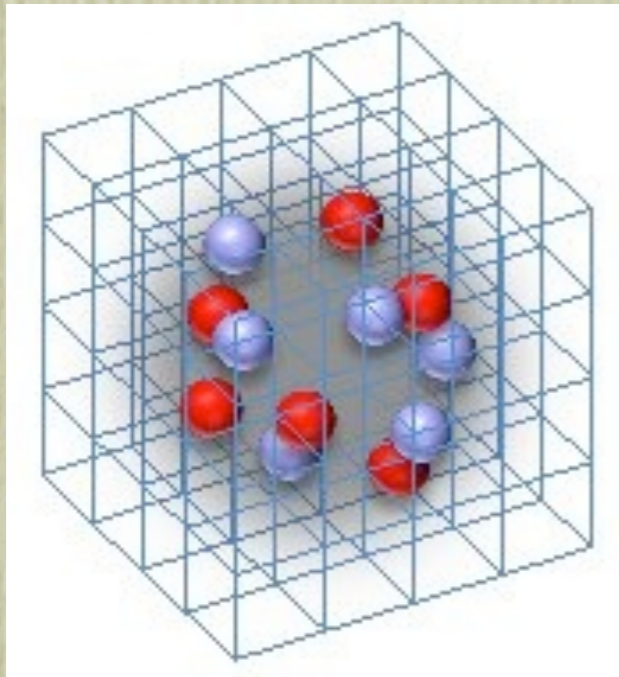
(a, r) fitted



Lattice EFT for Halo Nuclei

- Interested in $a(b, \gamma)c$
- Need interaction between clusters
- Calculate capture with cluster interaction.
Many possibilities --- traditional methods,
continuum EFT, lattice method

Nuclear Lattice Effective Field Theory collaboration



Evgeny Epelbaum,
Hermann Krebs,
Timo Lahde
Dean Lee
Ulf-G. Meissner

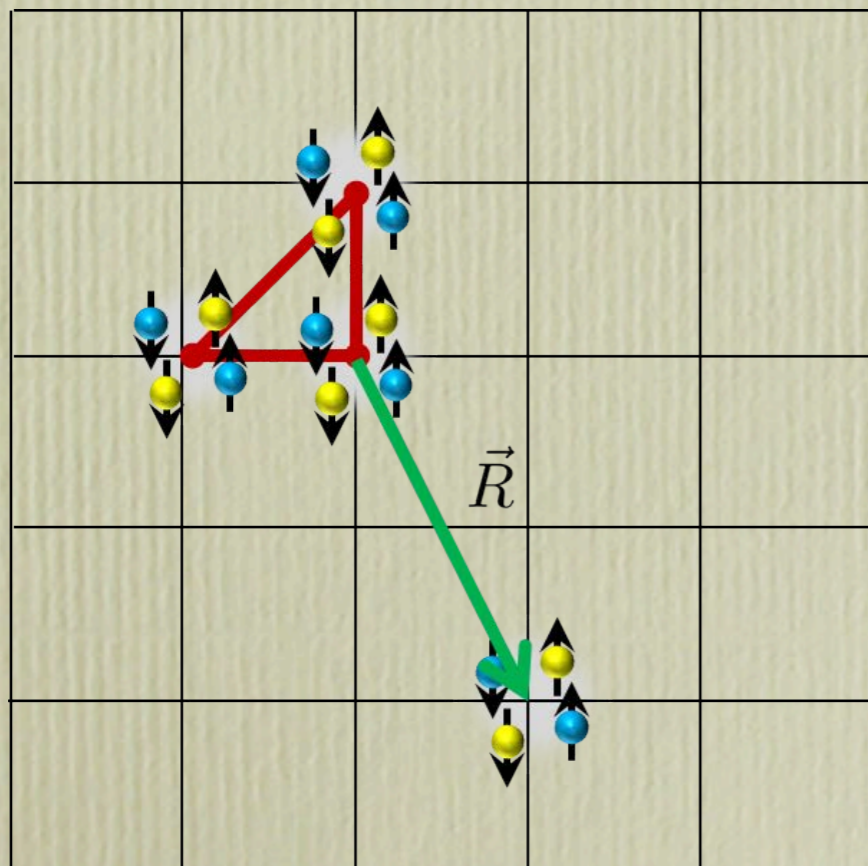
Adiabatic Hamiltonian

Microscopic Hamiltonian $L^{3(A-1)}$

Adiabatic Hamiltonian for the clusters L^3

-- acts on the cluster c.m. and spins

Blume, Greene 2000



Lee, Pine, Rupak

1D toy atom-dimer problem

Microscopic Hamiltonian: -2.130490, -2.130490, 0.1189620,
0.1189620, ...

Adiabatic Hamiltonian: -2.130505, -2.130493, 0.1189604,
0.1189781, ...

That was in 2012, Can do better now

n-d scattering in quartet channel

Microscopic Hamiltonian: 7.152, 23.37, 23.37, 23.37,
29.61, 29.61, 40.34, ...

Adiabatic Hamiltonian: 7.166, 23.42, 23.42, 23.42,
29.74, 29.74, 40.49, ...

Warm up $p(n, \gamma)d$

Write $\langle \psi_B | O_{EM} | \psi_i \rangle$ using retarded Green's function

$$\mathcal{M}(\epsilon) = \left(\frac{p^2}{M} - E - i\epsilon \right) \sum_{\mathbf{x}, \mathbf{y}} \psi_B^*(\mathbf{y}) \langle \mathbf{y} | \frac{1}{E - \hat{H}_s + i\epsilon} | \mathbf{x} \rangle e^{i\mathbf{p} \cdot \mathbf{x}}$$

Exact analytic continuum result

$$\mathcal{M}_C(\epsilon) = \frac{1}{p^2 + \gamma^2} - \frac{1}{(1/a + ip_\epsilon)(\gamma - ip_\epsilon)}, \quad p_\epsilon = \sqrt{p^2 + iM\epsilon}$$

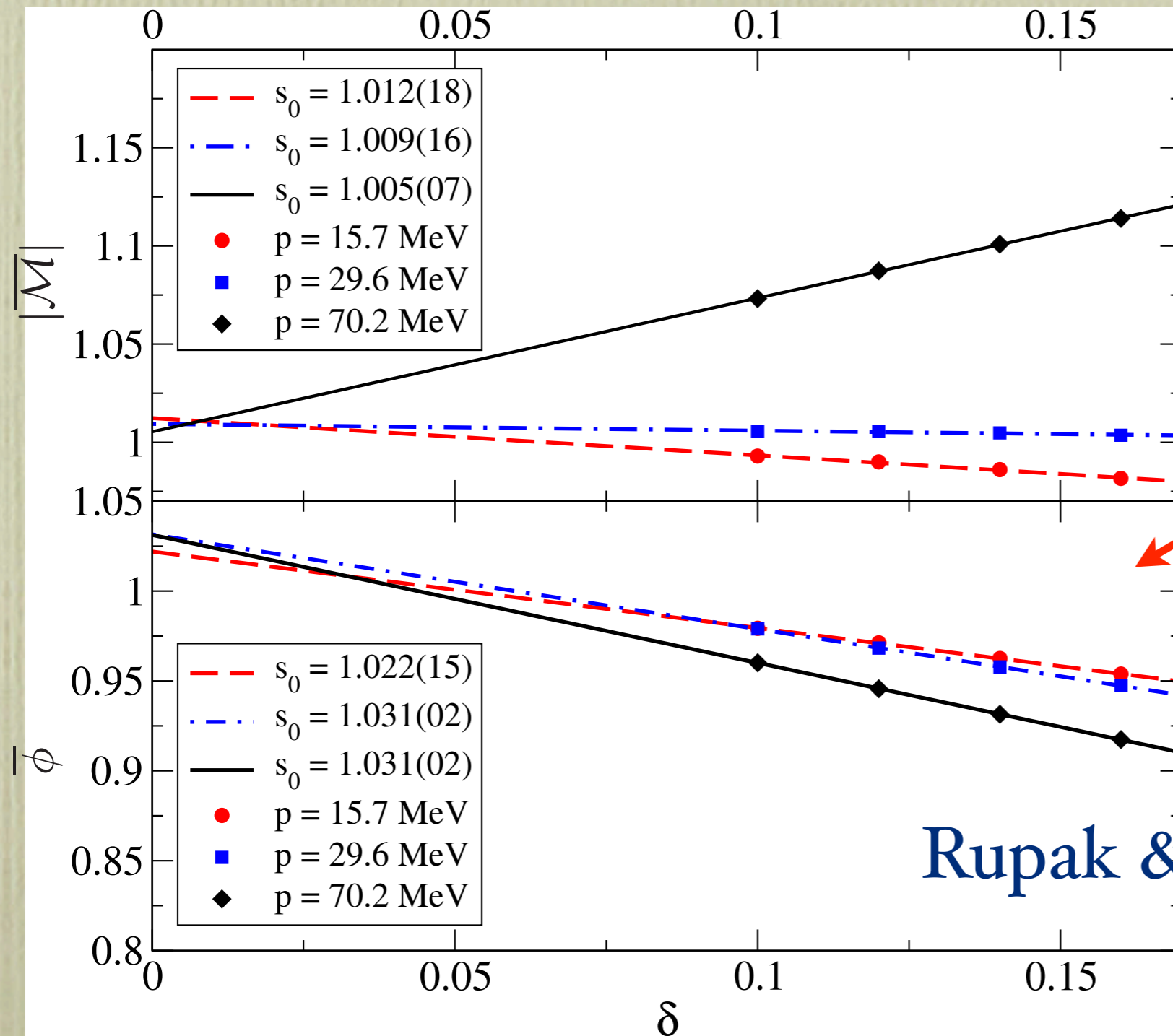
When $\epsilon \rightarrow 0^+$, \mathcal{M}_C reduces to known M1 result

Rupak, 2000

Lee & Rupak

Lattice EFT results

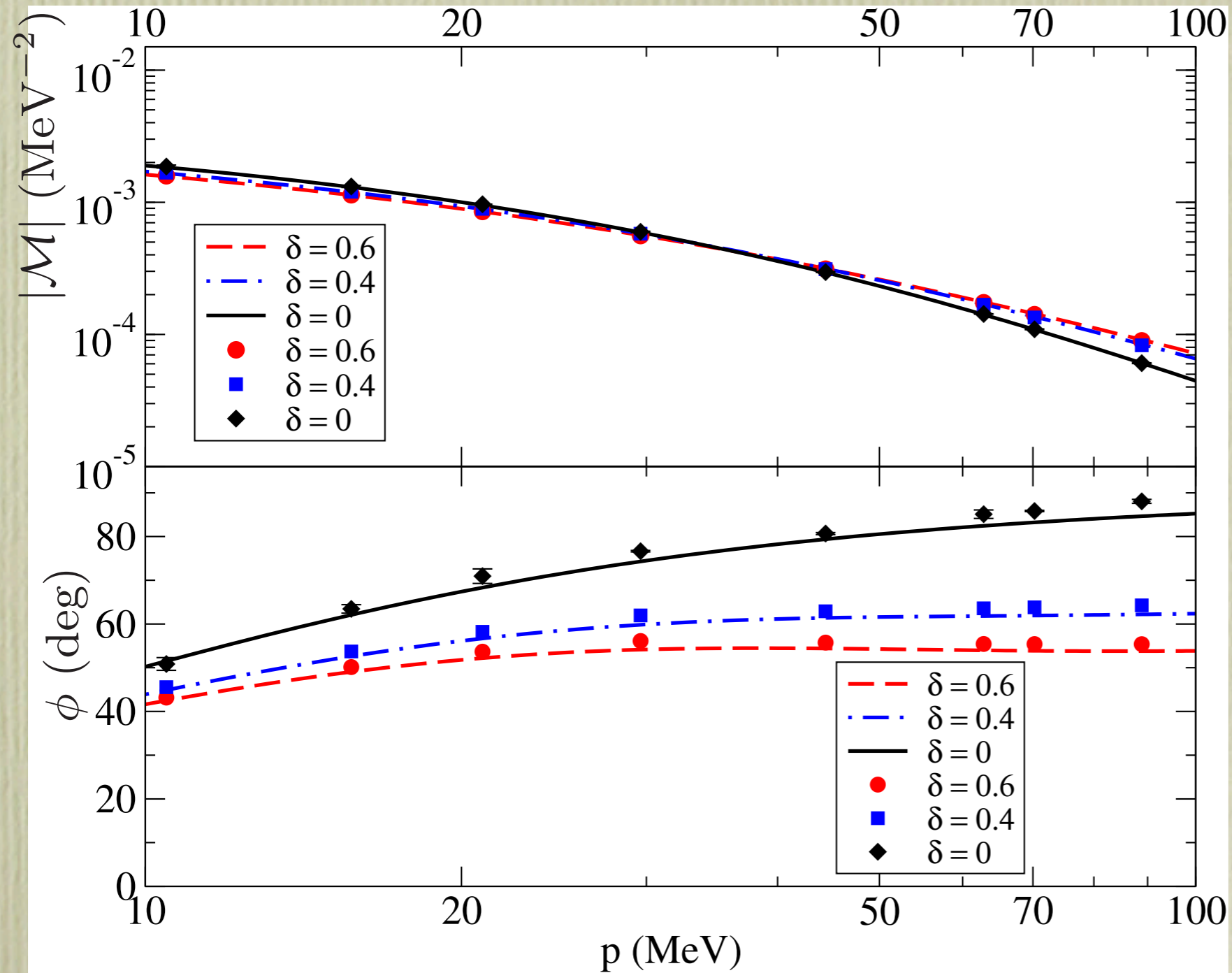
Magnitude, argument normalized to continuum



$$\delta = \epsilon M / p^2$$

Rupak & Lee, arXiv:1302.4158

Continuum extrapolation



Lattice QCD to lattice EFT

- Constrain Hamiltonian in elastic channels
- Electroweak currents --- [Detmold, Savage 2004](#)
 1. Fit pionless EFT at unphysical pion mass.
 2. Match observables in pionless and chiral EFT.
 3. Extrapolate to physical pion mass.
- Coulomb effect (in discussion)

Conclusions

- Capture reactions ${}^7\text{Li}(n, \gamma){}^8\text{Li}$, ${}^{14}\text{C}(n, \gamma){}^{15}\text{C}$, etc.
- In progress ${}^7\text{Be}(p, \gamma){}^8\text{B}$
- Capture reactions in lattice EFT

Thank you