

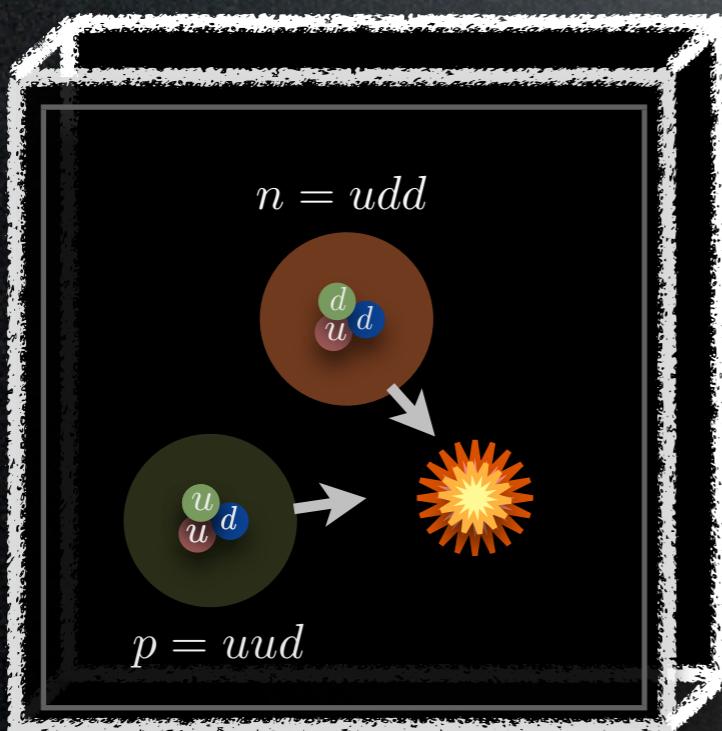
INT-13-53W
March 12, 2012

Nuclear physics in a box

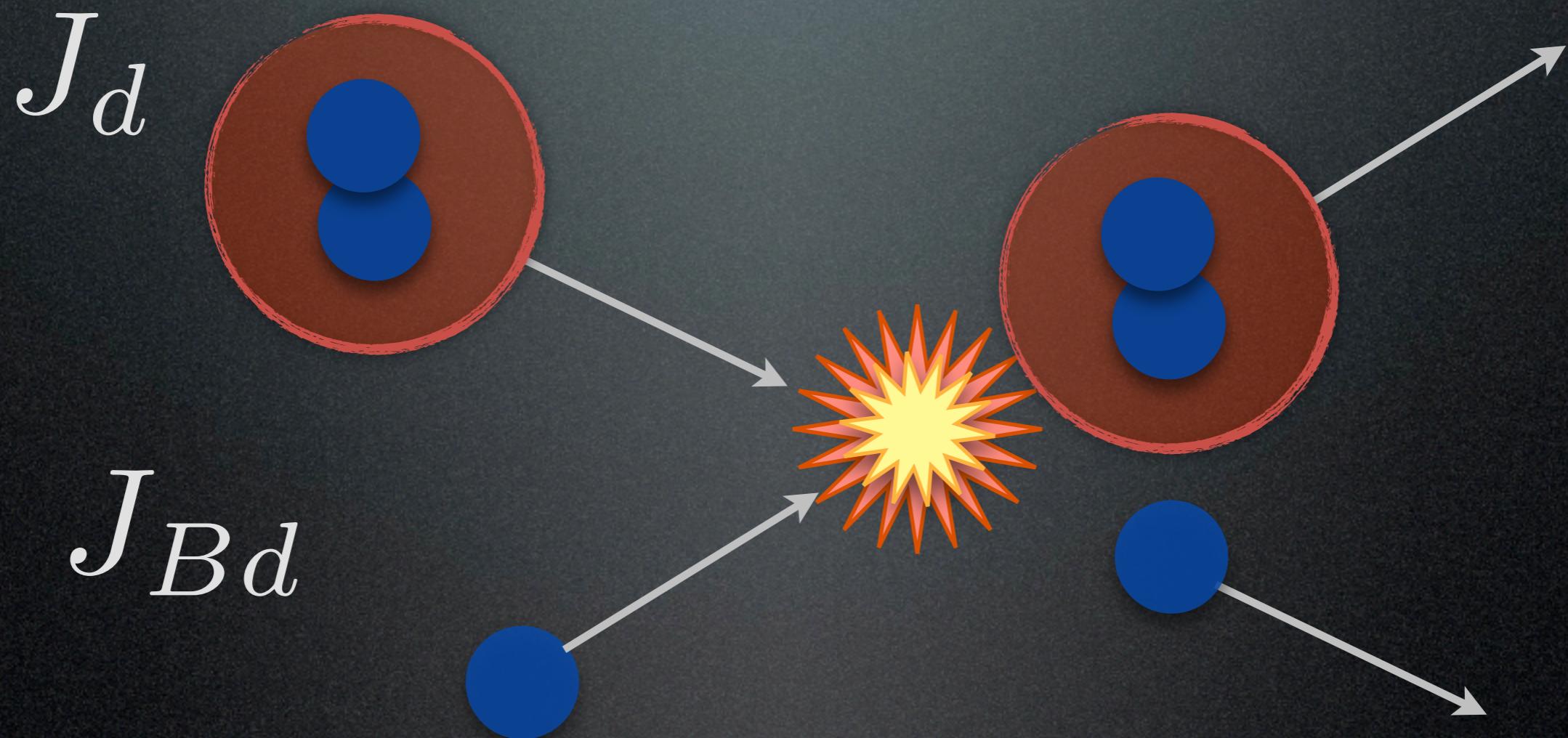
the devil is in the details

Raúl Briceño

(in collaboration with Zohreh Davoudi)

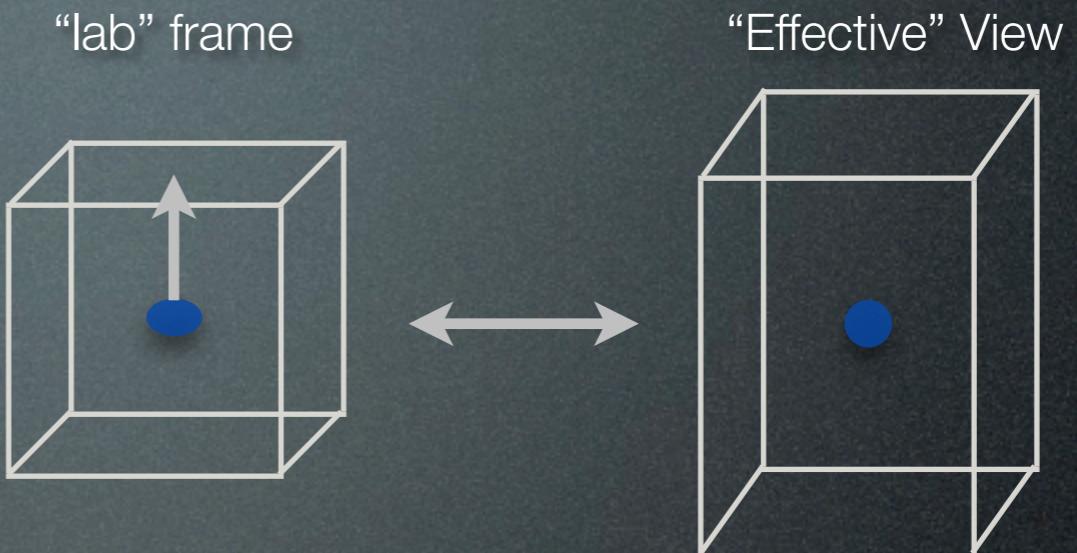


Three-Body Sector: Partial Waves & Mixing



Boosts

Symmetry is reduced:



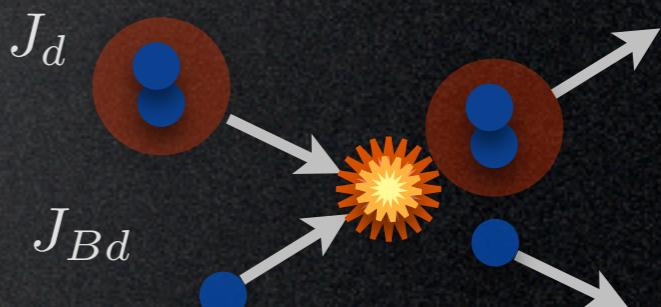
Boson-diboson CM:

- diboson is boosted: $J_d = \{0, 2, 4, \dots\}$
- Bd is unboosted: $J_{Bd} = \{0, 4, 6, \dots\}$

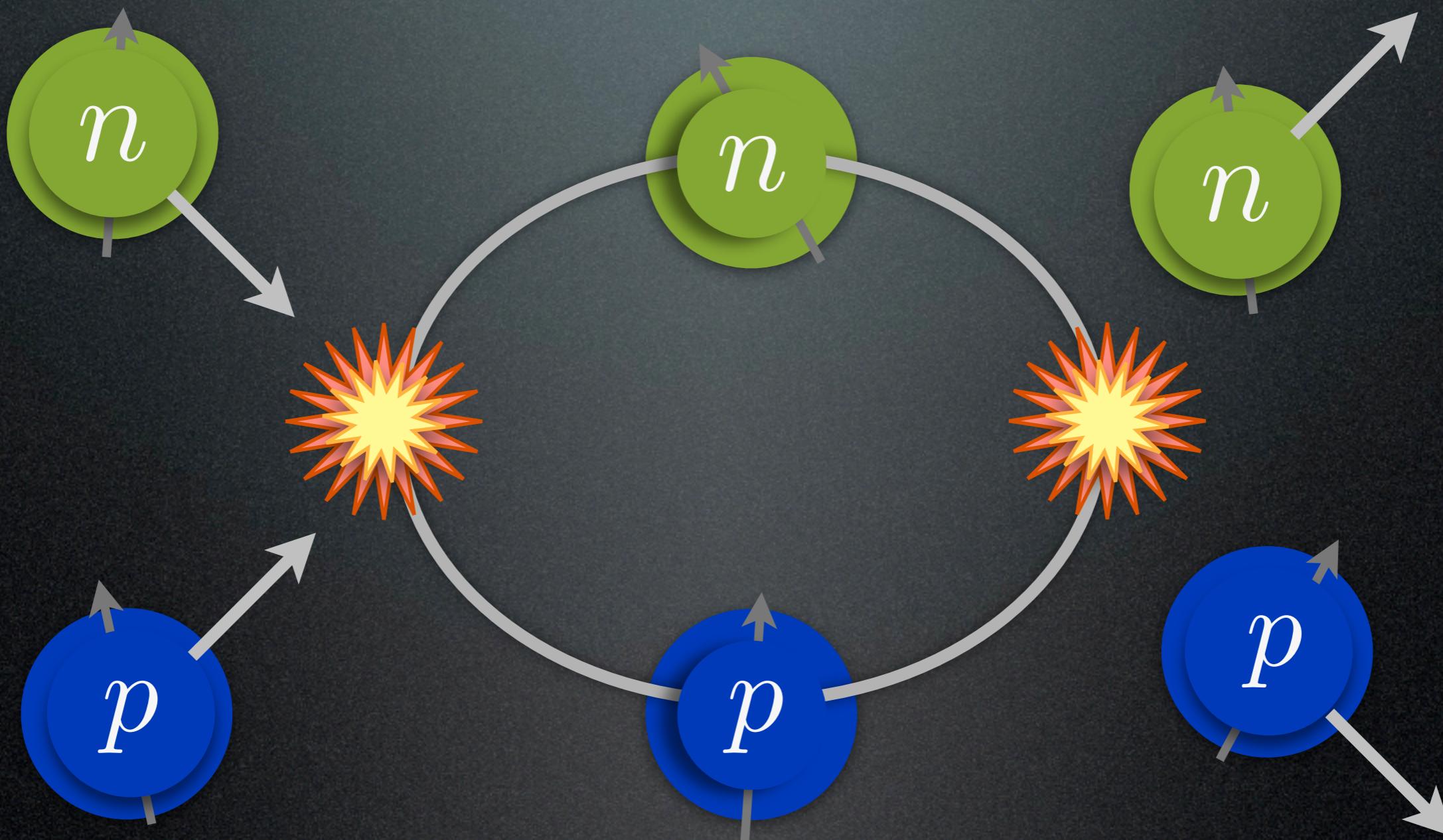
Must generalize
dimer formalism

Boson-diboson Boosted:

- $J_{Bd} = \{0, 1, 2, \dots\}$



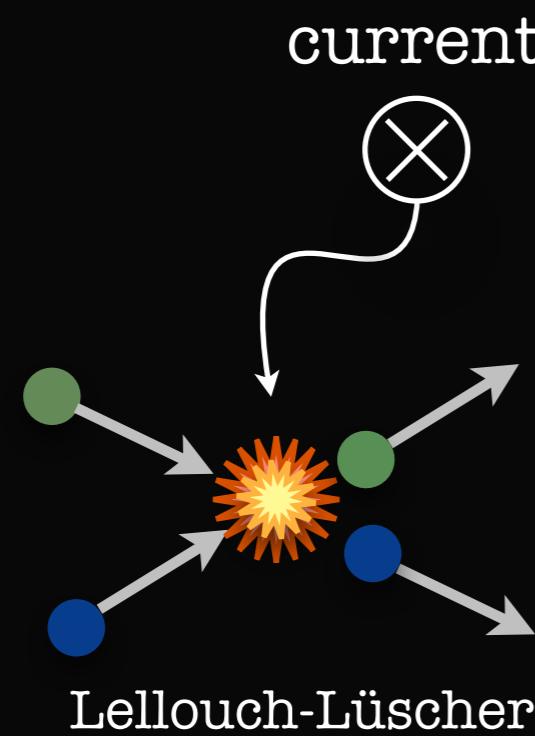
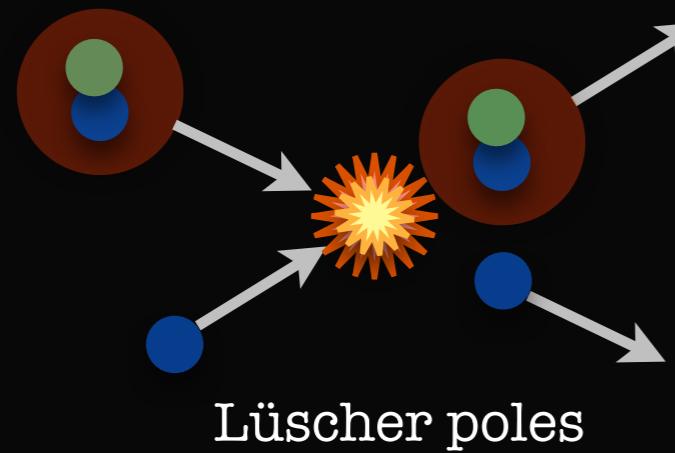
NN-Scattering



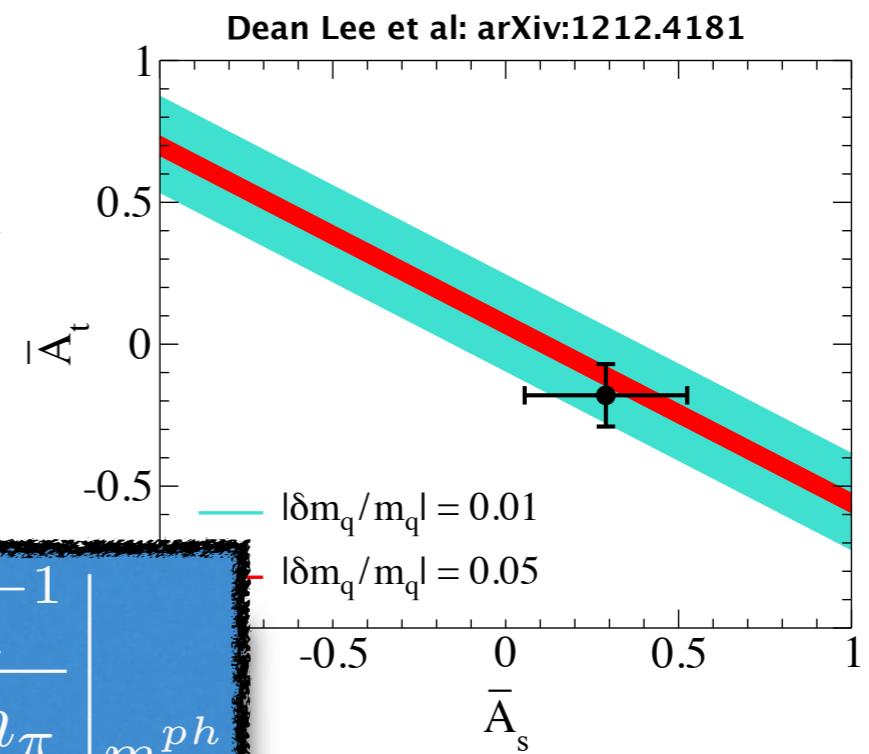
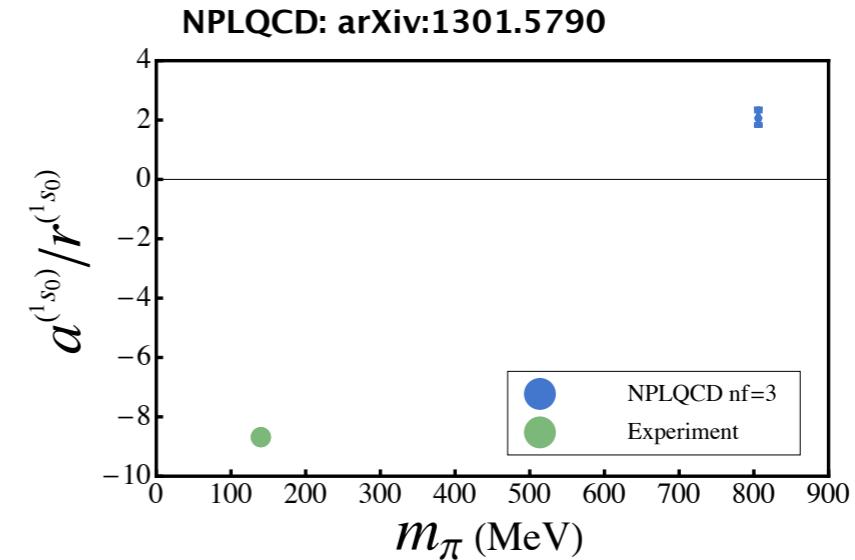
Why NN?



Pragmatic issues



Issues of Insight



$$\bar{A}_i \equiv \left. \frac{\partial a_i^{-1}}{\partial m_\pi} \right|_{m_\pi^{ph}}$$

Dimer Formalism: Scalar Sector

spherical dimer

$$(E, P) \ l_1 m_1 \quad \overline{l_2 m_2} = \frac{-i\delta_{l_1 l_2} \delta_{m_1 m_2}}{E - \frac{\mathbf{P}^2}{4m} - \Delta_{l_1} + \sum_{n=1} c_{l_1, n} \left(E - \frac{\mathbf{P}^2}{4m} \right)^{n+1}}$$

$$(E, P) \ l_1 m_1 \quad \begin{array}{c} / \\ \backslash \end{array} \quad (E - E_k, P - k) = -ig_{2,l_1} \sqrt{4\pi} Y_{l_1 m_1}^*(k^*) \ k^{*l_1} \ \delta_{1,(-1)^{l_1}}$$

$$k^* \equiv k - P/2$$

Galilean invariance

fully dressed dimer

$$l_1 m_1 \quad \overline{l_2 m_2} = \quad \begin{array}{c} / \\ \backslash \end{array} + \quad \begin{array}{c} / \\ \backslash \end{array} \quad \infty \quad = \quad \frac{i\delta_{l_1 l_2} \delta_{m_1 m_2} \ k^{*2l_1}}{k^{*2l_1+1} \cot \delta_d^{(l_1)} - ik^{*2l_1+1}}$$

$$\text{ERE: } k^{*2l+1} \cot \delta_d^{(l)} = -\frac{1}{a_l} + \frac{r_l k^{*2}}{2} + \sum_{n=2}^{\infty} \frac{\rho_{n,l}}{2n!} \ k^{*2n}$$



Finite volume result agrees with Rummukainen & Gottlieb, and Kim, Sharpe & Sachrajda

Nuclear Sector: Infinite Volume

Complicated by physical mixing

[e.g. positive parity, iso-singlet sector, $J<4$]

$$\mathcal{M}_{S=1, I=0}^{\infty, \pi=+} = \begin{pmatrix} \mathcal{M}_{J=1}^{\infty} & 0 & 0 \\ 0 & \mathcal{M}_{J=2}^{\infty} & 0 \\ 0 & 0 & \mathcal{M}_{J=3}^{\infty} \end{pmatrix}$$

~~25 x 25~~
~~18 x 18~~

$$[\mathcal{M}_{J=1}^{\infty}]_{m_j m_{j'}} = \begin{pmatrix} \mathcal{M}_{1m_j; 1m_{j'}}^S & \mathcal{M}_{1m_j; 1m_{j'}}^{SD} \\ \mathcal{M}_{1m_j; 1m_{j'}}^{DS} & \mathcal{M}_{1m_j; 1m_{j'}}^D \end{pmatrix}$$

$$[\mathcal{M}_{J=2}^{\infty}]_{m_j m_{j'}} = \mathcal{M}_{1m_j; 1m_{j'}}^D$$

$$[\mathcal{M}_{J=3}^{\infty}]_{m_j m_{j'}} = \begin{pmatrix} \mathcal{M}_{1m_j; 1m_{j'}}^D & \mathcal{M}_{1m_j; 1m_{j'}}^{DG} \\ \mathcal{M}_{1m_j; 1m_{j'}}^{DG} & \mathcal{M}_{1m_j; 1m_{j'}}^G \end{pmatrix}$$

Rich structure MUST be reflected in the dimer formalism

Dimer Formalism: Nuclear Sector

$$d_{\nu_J \nu_I, P}^{lS} = \sum_{m_l, m_s} \langle \nu_J | \nu_l \nu_S \rangle d_{\nu_l \nu_S \nu_I, P}$$

Dimer operator with (J,I) quantum numbers

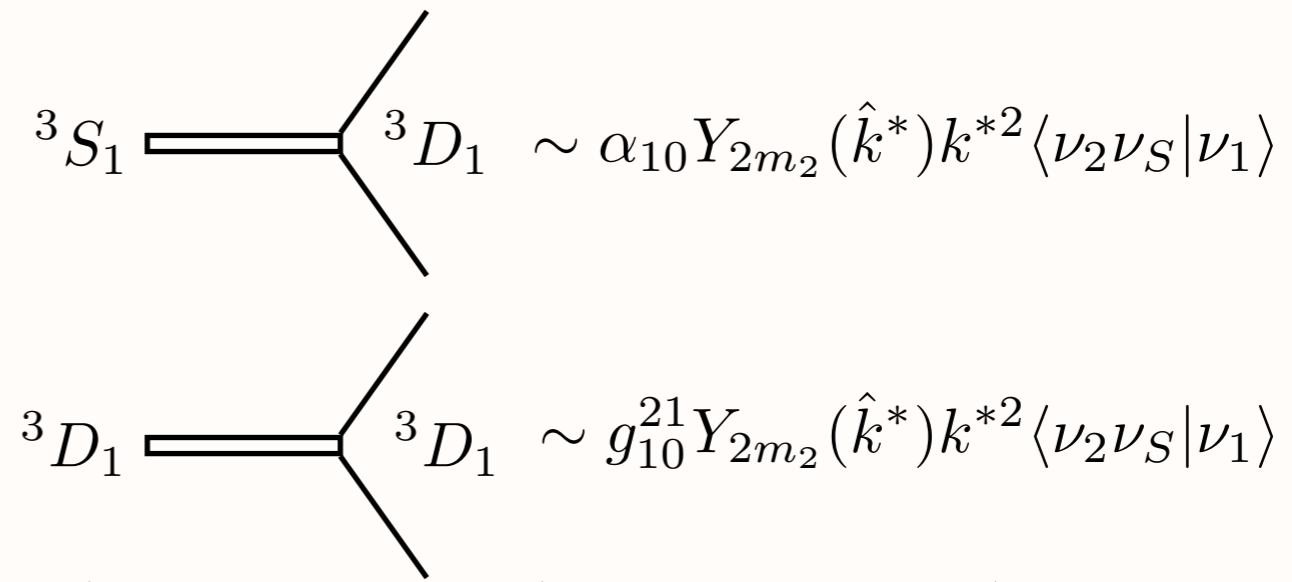
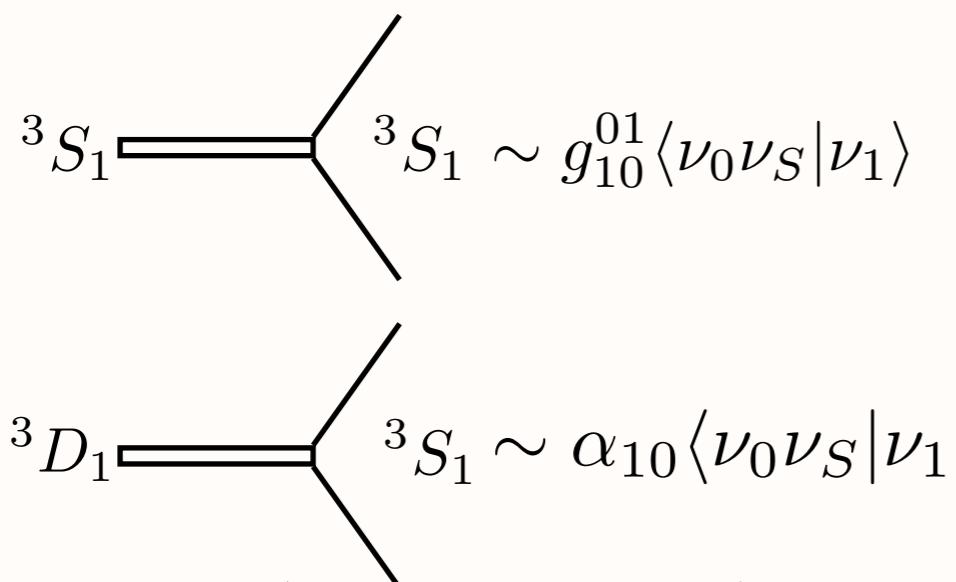
mixes different (l,S) states

convenient notation:

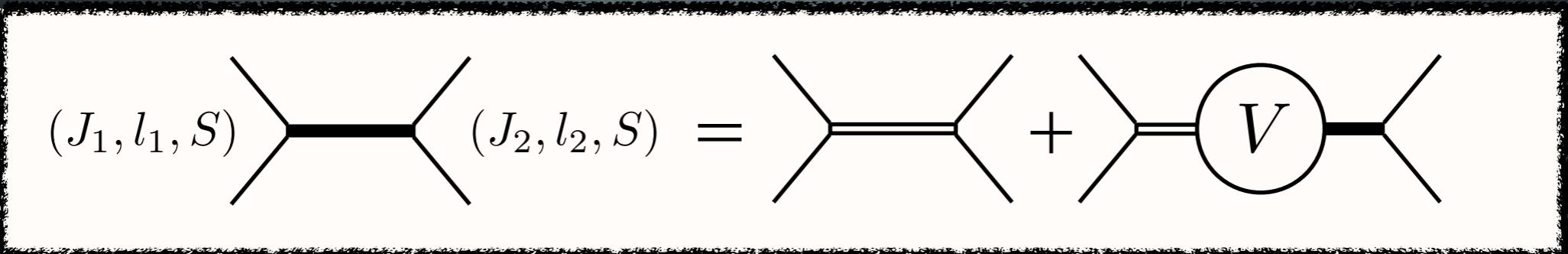
$$|\nu_X\rangle \equiv |Xm_X\rangle$$

e.g. ${}^3S_1 \longleftrightarrow {}^3D_1$

Dimer operator in the (l,S,I) - basis



Quantization Condition



independent of
dimer formalism

Holds for all partial waves,
parity & boosts

$$\det \left((\mathcal{M}^\infty)^{-1} + \delta \mathcal{G}^V \right) = 0$$

Scattering amplitude
Diagonal in J_2 -basis
mixes l states

Kinematic function of (L, E_L)
Mixes angular momentum

Quantization Condition

sanity check

$$\det \left((\mathcal{M}^\infty)^{-1} + \delta \mathcal{G}^V \right) = 0$$

$$i [\delta \mathcal{G}^V]_{\nu_J l; \nu_{J'} l'; SI} = -im \frac{P^{*l'+l}}{4\pi} \sum_{\substack{m_S, \\ m_l, m_{l'}}} \langle \nu_l \nu_S | \nu_J \rangle \langle \nu_{J'} | \nu_{l'} \nu_S \rangle \delta_{(-1)^{l+1-I-S}, 1} [F^P(P^{*2}) - iP^*]_{\nu_l, \nu_{l'}}$$

S = 0 limit

Agreement with
Gottlieb & Rummukainen
Kim, Sachrajda, Sharpe

S = 1 / 2 limit

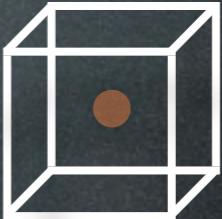
Agreement with Göckeler et al.

P=0 case by N. Ishizuka, proceeding

Boosts & Symmetry (I)

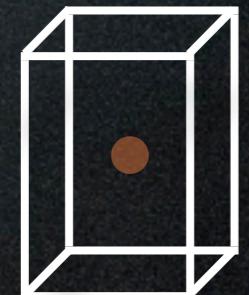
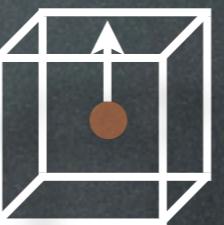
- O_h : Octahedral

$$d = (0,0,0)$$



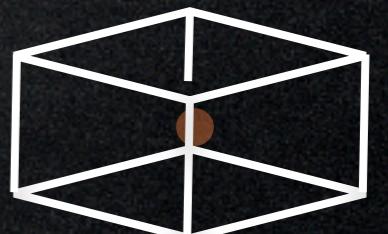
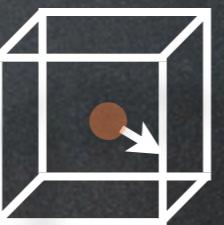
- D_{4h} : Tetragonal

$$d = (0,0,1)$$



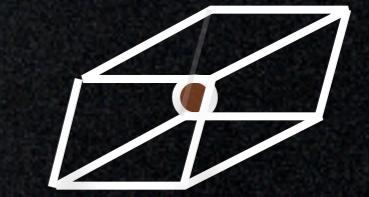
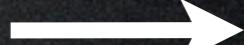
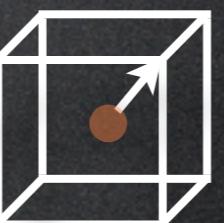
- D_{2h} : Orthorhombic

$$d = (1,1,0)$$



- D_{3h} : Trigonal

$$d = (1,1,1)$$



Boosts & Symmetry (II)

$$c_{\nu_l''}^P(x) = \left[\frac{1}{L^3} \sum_{\mathbf{k}} -\mathcal{P} \int \frac{d^3\mathbf{k}}{(2\pi)^3} \right] \frac{\sqrt{4\pi} Y_{\nu_l''}(\hat{\mathbf{k}}^*) k^{*l''}}{k^{*2} - x}$$

- Rotations: $c_{lm_l}^P = \sum_{m_{l'}=\{-l,l\}} \mathcal{D}_{m_l, m'_l}^{(l)}(R\chi) c_{lm'_l}^P$
- Parity + Isospin limit: $c_{lm_l}^P = (-1)^l c_{lm_l}^P$

- Examples [$J < 4$]:

e.g. 24 relations for O_h

- O_h : $c_{00}^P, c_{44}^P = \sqrt{\frac{5}{14}} c_{40}^P, \text{ else } c_{lm_l}^P = 0$
- D_{2h} : $c_{lm_l}^P = 0, \text{ if } l \text{ or } m_l \text{ are odd}$

$$c_{lm_l}^P = (-1)^{m_l/2} c_{l-m_l}^P$$

Lüscher (1990)

Cubic NN-Propagator

Evaluate determinant of $\left[(\mathcal{M}^\infty)^{-1} + \delta\mathcal{G}^V \right]_{\mathcal{O}_h}$

$$C_{lm} \sim Z_{lm}$$

SD				$(\mathcal{M}^\infty)^{-1}$												
b2	0	0	m1	0	0	0	0	0	0	0	0	0	0	0	0	0
0	b1	0	0	m1	0	0	0	0	0	0	0	0	0	0	0	0
0	0	b1	0	0	m1	0	0	0	0	0	0	0	0	0	0	0
m1	0	0	b2	0	0	0	0	0	0	0	0	0	0	0	0	z2
0	m1	0	0	b2	0	0	0	0	0	0	0	0	0	0	0	0
0	0	m1	0	0	b2	0	0	0	0	0	z2	0	0	0	z1	0
0	0	0	0	0	b3	0	0	0	z5	0	z6	0	0	0	z6	0
0	0	0	0	0	0	b4	0	0	0	0	0	z8	0	0	0	z9
0	0	0	0	0	0	0	b5	0	0	0	0	0	0	0	0	0
0	0	0	0	0	0	0	0	b4	0	-z9	0	0	0	-z8	0	0
0	0	0	0	0	z5	0	0	0	b3	0	-z6	0	0	0	-z6	0
0	0	0	0	z2	0	0	0	-z9	0	b6	0	0	0	z4	0	0
0	0	0	0	z6	0	0	0	-z6	0	b7	0	0	0	0	z10	0
0	0	0	z1	0	0	z8	0	0	0	0	b8	0	0	0	0	z4
0	0	0	z3	0	0	0	0	0	0	0	0	b9	0	0	0	0
0	0	0	0	z1	0	0	-z8	0	z4	0	0	0	b8	0	0	0
0	0	0	0	z6	0	0	0	-z6	0	z10	0	0	0	b7	0	0
0	0	0	z2	0	0	z9	0	0	0	0	z4	0	0	0	0	b60

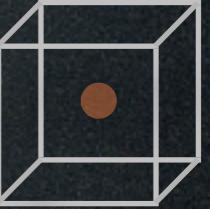
Cubic NN-Propagator

Unitary transformation: $S [\delta G^V] S^\dagger$

Diagonal blocks are diagonalized:

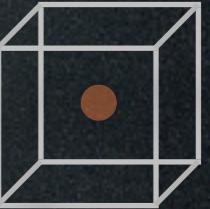
à la mode de Tom & Martin (2011)

A₂-Irrep



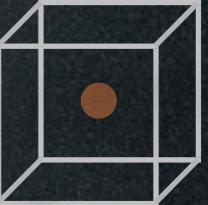
$$\frac{1}{\mathcal{M}_{J=3}^D} + \frac{imp}{4\pi} = mc_{00} - 4\sqrt{\frac{2}{35}} \frac{mc_{44}}{p^4}$$

E-Irrep

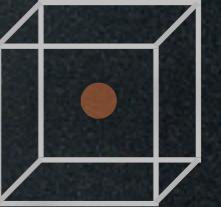


$$\left(\frac{1}{\mathcal{M}_{J=2}^D} + \frac{imp}{4\pi} - mc_{00} + 4\sqrt{\frac{2}{35}} \frac{mc_{44}}{p^4} \right)^2 = 0$$

T₂-Irrep



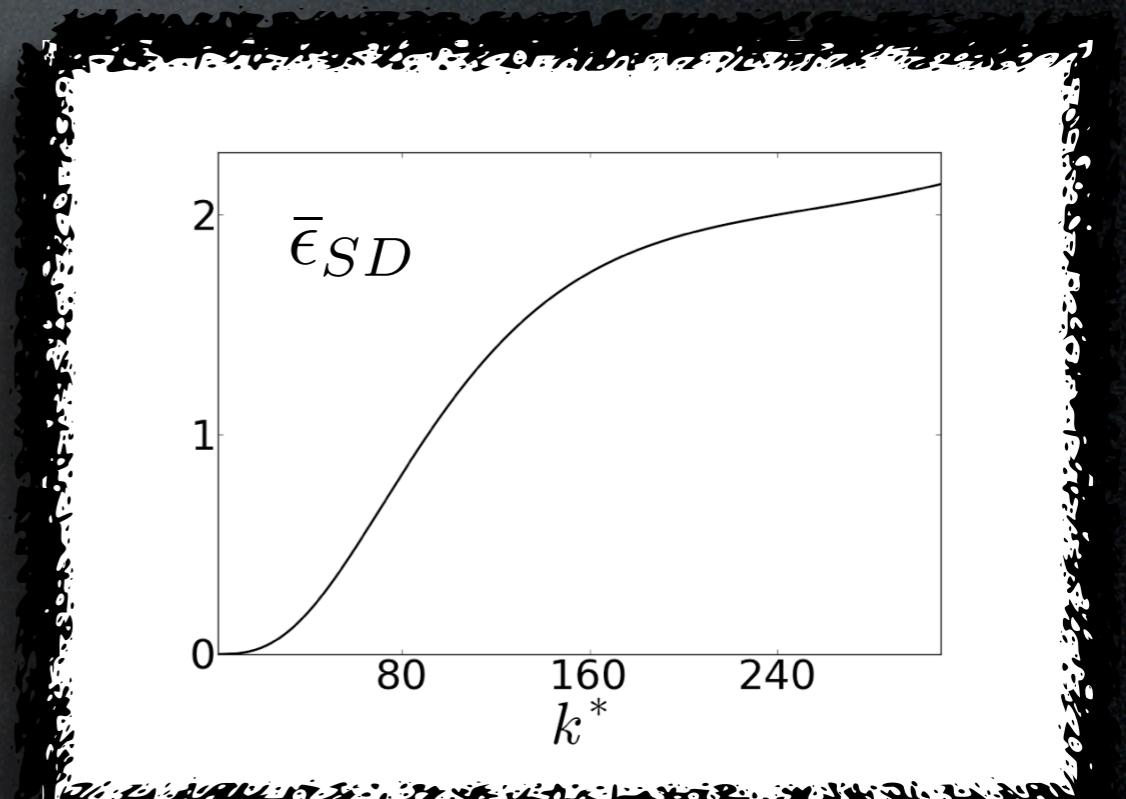
$$\det \begin{pmatrix} \frac{1}{\mathcal{M}_{J=2}^D} + \frac{imp}{4\pi} - mc_{00} - \frac{8}{3}\sqrt{\frac{2}{35}}\frac{mc_{44}}{p^4} & \frac{4}{3}\sqrt{\frac{5}{7}}\frac{mc_{44}}{p^4} \\ \frac{4}{3}\sqrt{\frac{5}{7}}\frac{mc_{44}}{p^4} & \frac{1}{\mathcal{M}_{J=3}^D} + \frac{imp}{4\pi} - mc_{00} + \frac{2}{3}\sqrt{\frac{2}{35}}\frac{mc_{44}}{p^4} \end{pmatrix} = 0$$

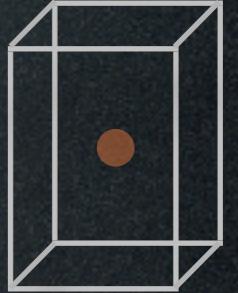


T₁-Irrep

$$\det \begin{pmatrix} \frac{\mathcal{M}_{J=1}^D}{\det \mathcal{M}_{J=1}} + \frac{imp}{4\pi} - mc_{00} & \frac{\mathcal{M}_{J=1}^{SD}}{-\det \mathcal{M}_{J=1}} & 0 \\ \frac{\mathcal{M}_{J=1}^{SD}}{-\det \mathcal{M}_{J=1}} & \frac{\mathcal{M}_{J=1}^S}{\det \mathcal{M}_{J=1}} + \frac{imp}{4\pi} - mc_{00} & 4\sqrt{\frac{3}{35}} \frac{mc_{44}}{p^4} \\ 0 & 4\sqrt{\frac{3}{35}} \frac{mc_{44}}{p^4} & \frac{1}{\mathcal{M}_{J=3}^D} + \frac{imp}{4\pi} - mc_{00} - \frac{2}{3}\sqrt{\frac{2}{35}} \frac{mc_{44}}{p^4} \end{pmatrix} = 0$$

- Free states split
- In CM the SD-mixing $\sim \sin(2\bar{\epsilon}_{SD})$
- Nearly uncoupled at physical point
- Unphysical m_π ?





Tetragonal

- Five irreps
- Two irreps with S-wave: A_2, E
- Quantization condition for A_2 -irrep:

$$\det \begin{pmatrix} -\frac{2c_{20}m}{\sqrt{5}p^2} + \frac{ipm}{4\pi} - mc_{00} + \frac{2\mathcal{M}_{J=1}^S + 2\sqrt{2}\mathcal{M}_{J=1}^{SD} + \mathcal{M}_{J=1}^D}{3 \det \mathcal{M}_{J=1}} & -\frac{-\sqrt{2}\mathcal{M}_{J=1}^S + \mathcal{M}_{J=1}^{SD} + \sqrt{2}\mathcal{M}_{J=1}^D}{3 \det \mathcal{M}_{J=1}} & \frac{9c_{20}m}{7\sqrt{5}p^2} + \frac{4c_{40}m}{7p^4} \\ -\frac{-\sqrt{2}\mathcal{M}_{J=1}^S + \mathcal{M}_{J=1}^{SD} + \sqrt{2}\mathcal{M}_{J=1}^D}{3 \det \mathcal{M}_{J=1}} & \frac{2c_{20}m}{\sqrt{5}p^2} + \frac{ipm}{4\pi} - mc_{00} + \frac{\mathcal{M}_{J=1}^S - 2\sqrt{2}\mathcal{M}_{J=1}^{SD} + 2\mathcal{M}_{J=1}^D}{3 \det \mathcal{M}_{J=1}} & \frac{2\sqrt{2}c_{40}m}{7p^4} - \frac{6\sqrt{\frac{2}{5}}c_{20}m}{7p^2} \\ \frac{9c_{20}m}{7\sqrt{5}p^2} + \frac{4c_{40}m}{7p^4} & \frac{2\sqrt{2}c_{40}m}{7p^4} - \frac{6\sqrt{\frac{2}{5}}c_{20}m}{7p^2} & -\frac{8c_{20}m}{7\sqrt{5}p^2} - \frac{2c_{40}m}{7p^4} + \frac{ipm}{4\pi} - mc_{00} + \frac{1}{\mathcal{M}_{J=3}^D} \end{pmatrix} = 0$$

Physical mixing

vs.

Symmetry reduction

- E-irrep is 5D: mixes $J=1, 2, \& 3$

Conclusion

- Generalization of dimer formalism
 - Scalar/Nuclear sectors
- Quantization condition for NN in a box
 - Arbitrary partial wave, parity, & boost
- Positive parity irreps for $J<4$
 - Boosts: $(0,0,0)$, $(0,0,1)$, $(1,1,0)$, $(1,1,1)$
- Negative parity sector under way
- Good interpolating operators needed!

Thanks!

- Martin, David, INT
- Speakers
- Participants

under
construction

please come back later...

