

# Experimental Overview of Past and Future Studies of the *EMC Effect*

Dave Gaskell - JLab

February 14, 2013

*INT Workshop on Nuclear Structure  
and Dynamics at Short Distance*

# Outline

- Measurements of  $\sigma_A/\sigma_D$ 
  - Early measurements
  - $x$ ,  $Q^2$ , nuclear dependence, universality
- JLab results and implications
  - EMC effect and local density
  - EMC-SRC connection
  - Flavor dependence
  - Nuclear dependence of  $R=\sigma_L/\sigma_T$
- Summary

# Quarks in the Nucleus

Typical nuclear binding energies  
→ MeV while DIS scales → GeV

Naïve expectation:

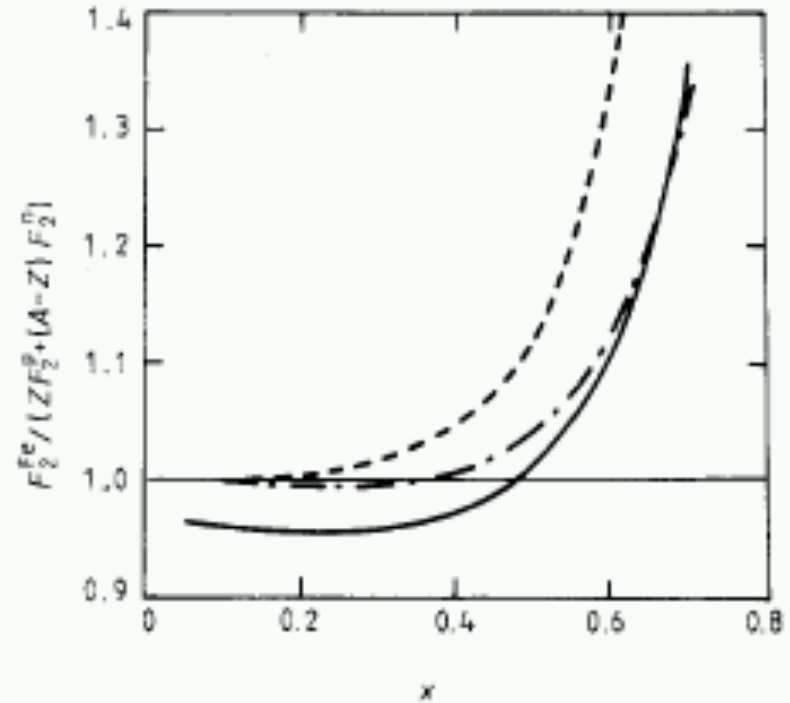
$$F_2^A(x) = ZF_2^p(x) + (A - Z)F_2^n(x)$$

More sophisticated approach  
includes effects from Fermi  
motion

$$F_2^A(x) = \sum_i \int_x^{M_A/m_N} dy f_i(y) F_2^N(x/y)$$

Quark distributions in nuclei were  
not expected to be significantly  
different (below  $x=0.6$ )

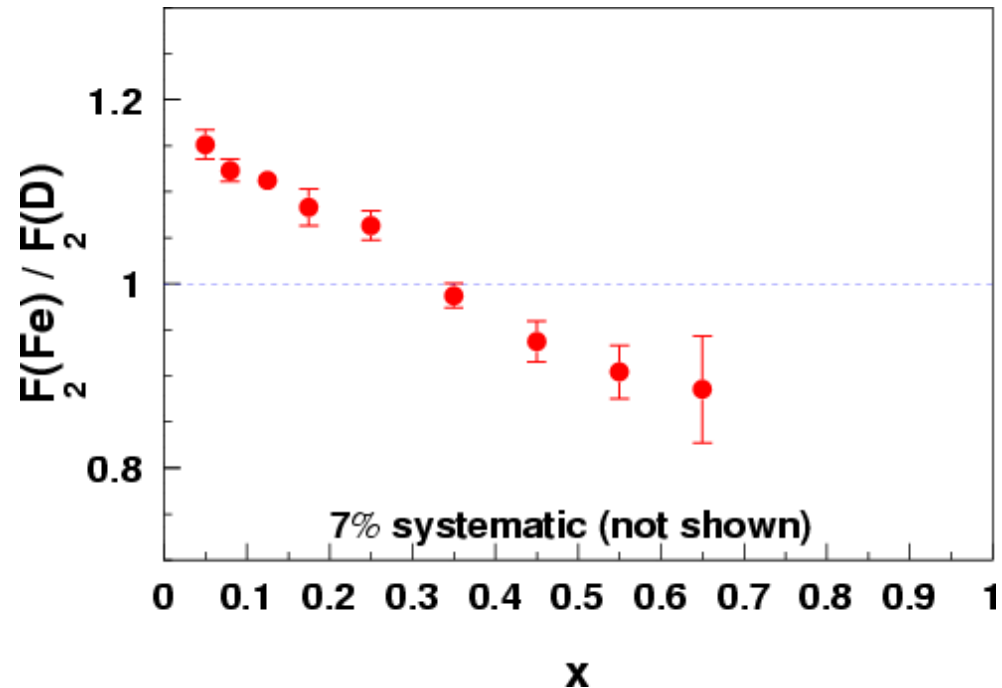
$$F_2^{Fe} / (ZF_2^p + (A - Z)F_2^n)$$



*Bodek and Ritchie*  
*PRD 23, 1070 (1981)*

# First Measurement of the EMC Effect

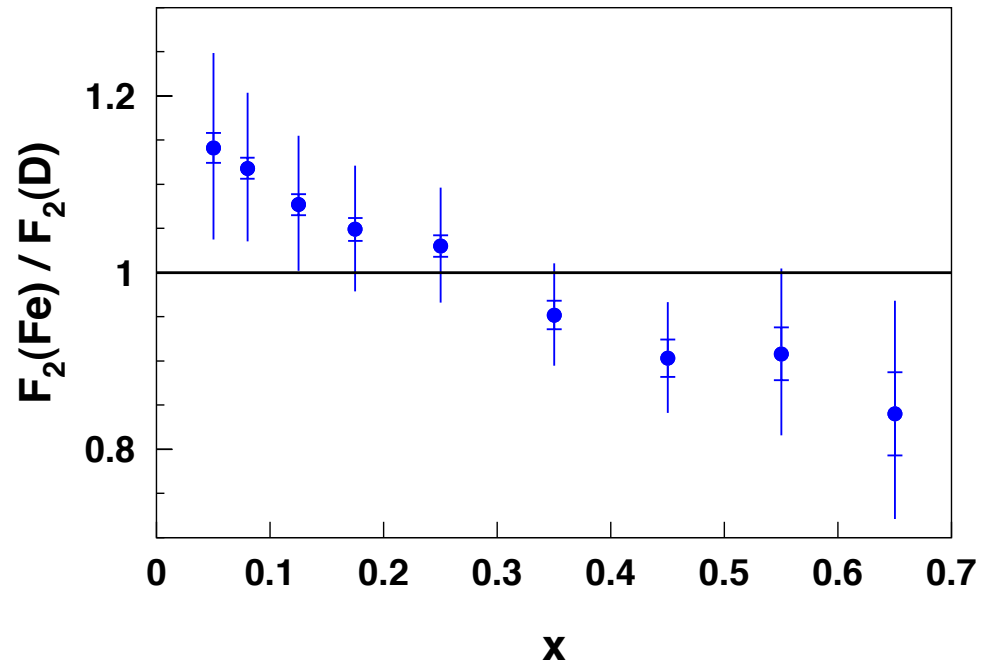
- First published measurement of nuclear dependence of  $F_2$  by the European Muon Collaboration in 1983
- Observed 2 mysterious effects
  - Significant enhancement at small  $x$  → Nuclear Pions! (see my thesis)
  - Depletion at large  $x$  → the “EMC Effect”
- Enhancement at  $x < 0.1$  later went away



*Aubert et al, Phys. Lett. B123, 275 (1983)*

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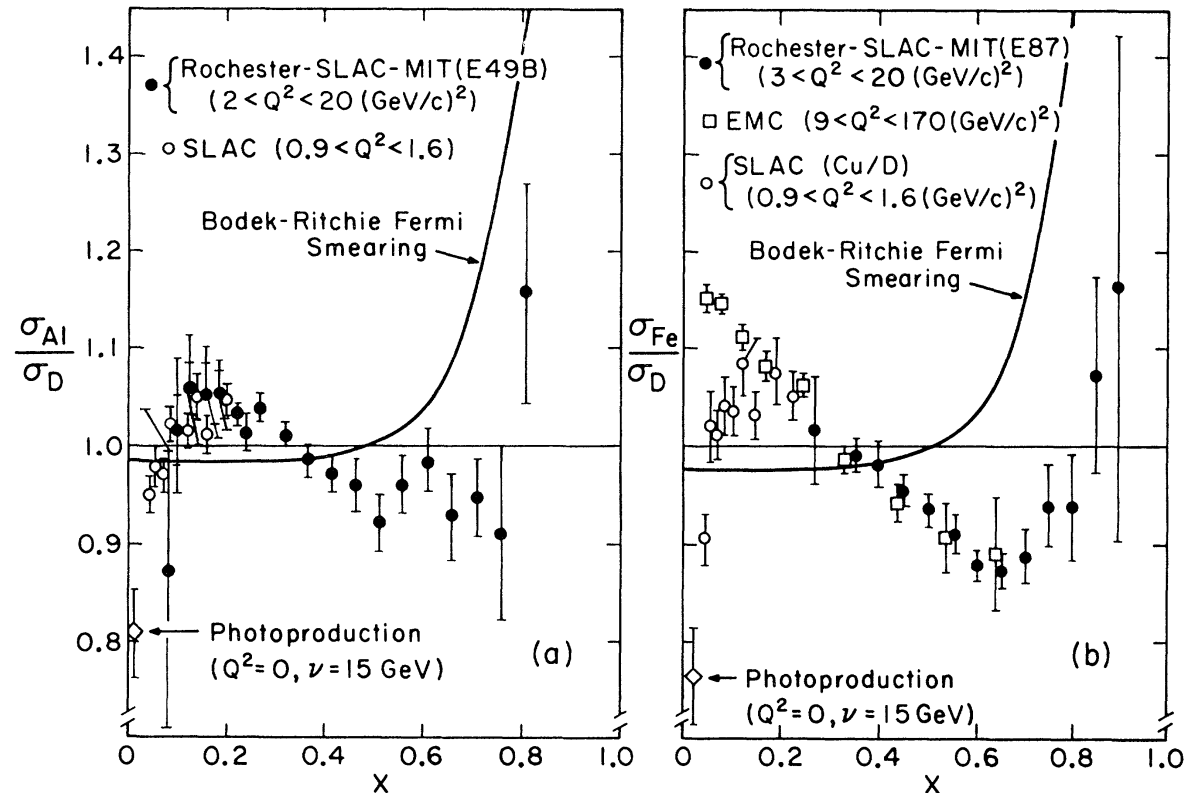


*Aubert et al, Nucl. Phys. B293, 740 (1987)*

# Confirmation of the Effect

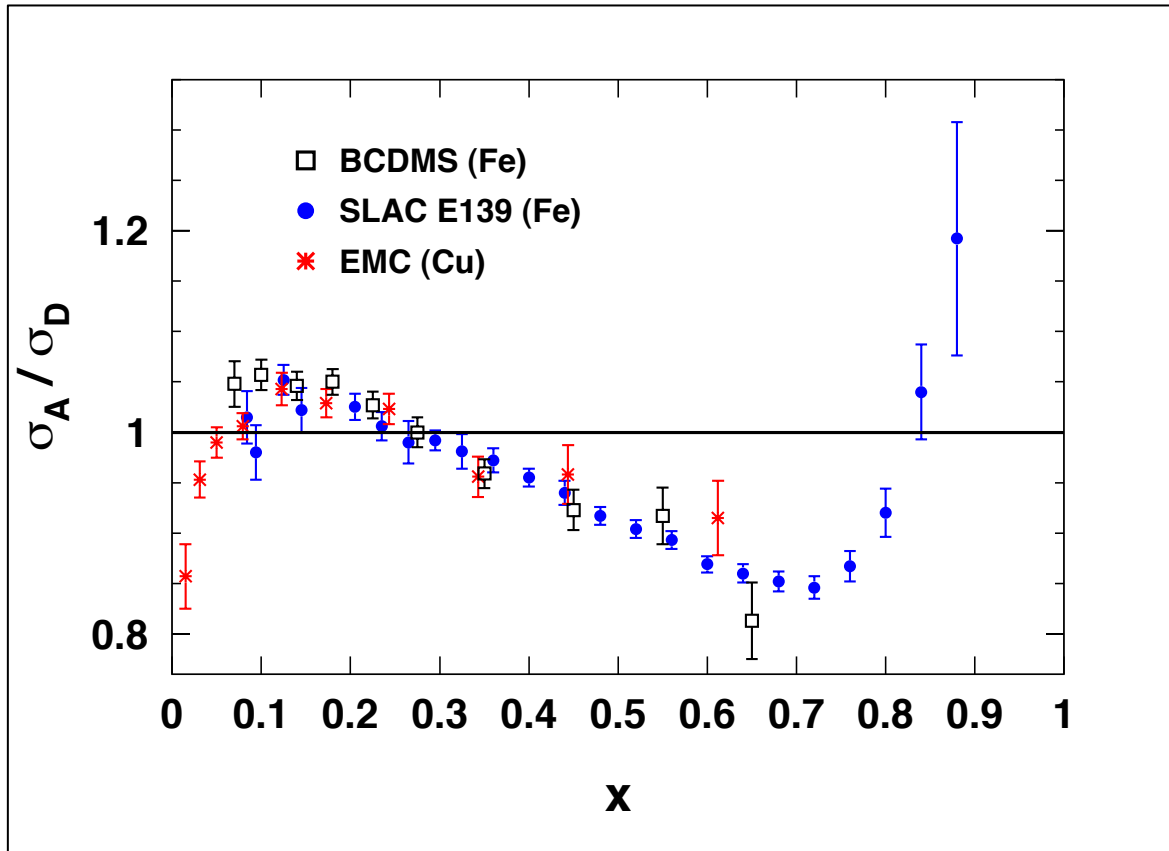
SLAC re-analysis of old solid target data used for measurements of cryotarget wall backgrounds

→ Effect for  $x > 0.3$  confirmed  
 → No large excess at very low  $x$



Bodek et al, PRL 50, 1431 (1983) and PRL 51, 534 (1983)

# Subsequent Measurements



A program of dedicated measurements quickly followed

The resulting data is remarkably consistent over a large range of beam energies and species

# EMC Effect Measurements

| Laboratory/<br>collaboration | Beam | Energy<br>(GeV) | Target  | Year          |
|------------------------------|------|-----------------|---|---------------|
| SLAC E139                    | e    | 8-24.5          | <b>D</b> , <sup>4</sup> He, Be, C, Ca, Fe, Ag, Au                         | 1994, 1984    |
| SLAC E140                    | e    | 3.75-19.5       | <b>D</b> , Fe, Au   | 1992, 1990    |
| CERN NMC                     | μ    | 90              | <sup>6</sup> <b>Li</b> , <sup>12</sup> C, <sup>40</sup> Ca                | 1992          |
|                              | μ    | 200             | <b>D</b> , <sup>4</sup> He, C, Ca   | 1991, 1995    |
|                              | μ    | 200             | Be, <b>C</b> , Al, Ca, Fe, Sn, Pb   | 1996          |
| CERN BCDMS                   | μ    | 200             | <b>D</b> , Fe   | 1987          |
|                              | μ    | 280             | <b>D</b> , N, Fe  | 1985          |
| CERN EMC                     | μ    | 100-280         | <b>D</b> , Cu   | 1993          |
|                              | μ    | 280             | <b>D</b> , C, Ca  | 1988          |
|                              | μ    | 100-280         | <b>D</b> , C, Cu, Sn  | 1988          |
|                              | μ    | 280             | H, <b>D</b> , Fe  | 1987          |
|                              | μ    | 100-280         | <b>D</b> , Fe   | 1983          |
| FNAL E665                    | μ    | 490             | <b>D</b> , Xe   | 1992          |
|                              | μ    | 490             | <b>D</b> , Xe   | 1992          |
| DESY HERMES                  | e    | 27              | <b>D</b> , <sup>3</sup> He, N, <b>Kr</b>                                  | 2000, 2003    |
| Jefferson Lab                | e    | 6               | <b>D</b> , <sup>3</sup> He, <sup>4</sup> He, Be, C, <b>Cu</b> , <b>Au</b> | 2009          |
|                              | e    | 6               | <b>D</b> , <b>C</b> , <b>Cu</b> , <b>Au</b>                               | 2004 (thesis) |

Geesaman, Saito, and Thomas, *Ann. Rev. Nucl. Sci.* 45, 337 (1995) – updated by Gaskell



# Nuclear dependence of structure functions

Experimentally, we measure cross sections (and the ratios of cross sections)

$$\frac{d\sigma}{d\Omega dE'} = \frac{4\alpha^2(E')^2}{Q^4 v} \left[ F_2(v, Q^2) \cos^2 \frac{\theta}{2} + \frac{2}{Mv} F_1(v, Q^2) \sin^2 \frac{\theta}{2} \right] \quad F_2(x) = \sum_i e_i^2 x q_i(x)$$

$$R = \frac{\sigma_L}{\sigma_T} = \frac{F_2}{2xF_1} \left( 1 + 4 \frac{M^2 x^2}{Q^2} \right) - 1 \quad \xrightarrow{\text{In the limit } R_A = R_D} \quad \boxed{\sigma_A/\sigma_D = F_2^A/F_2^D}$$

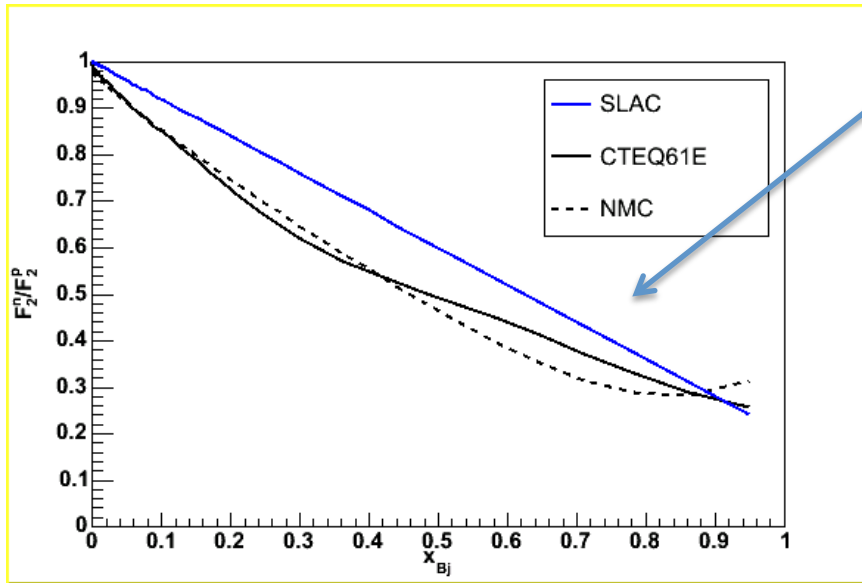
Experiments almost always display cross section ratios,  $\sigma_A/\sigma_D$

→ Often these ratios are labeled or called  $F_2^A/F_2^D$

→ Sometimes there is an additional uncertainty estimated to account for the  $\sigma \rightarrow F_2$  translation. Sometimes there is not.

# Isoscalar Corrections

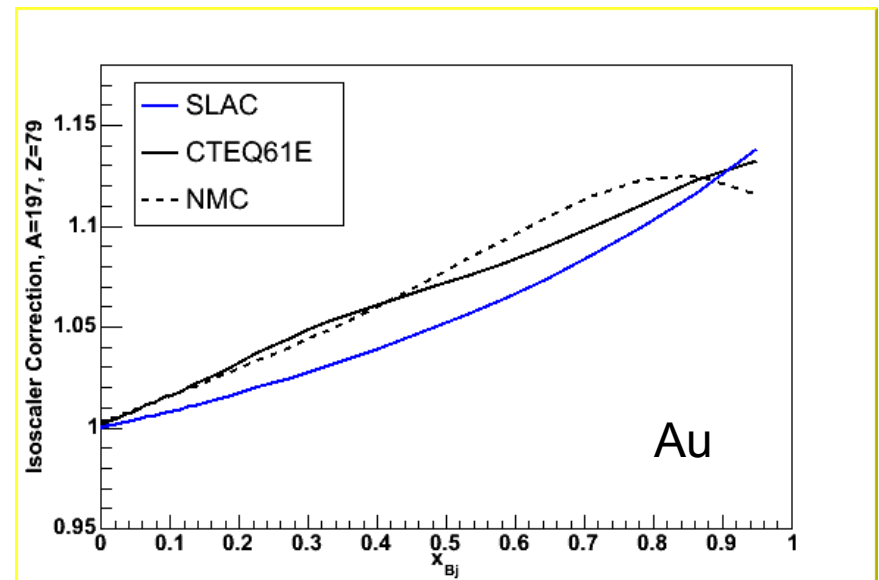
In the case of nuclei where  $N \neq Z$ , need to remove the “trivial” change in nuclear cross section due to  $\sigma_n \neq \sigma_p$   
 → Different experiments often use slightly different parameterizations/estimates for this correction



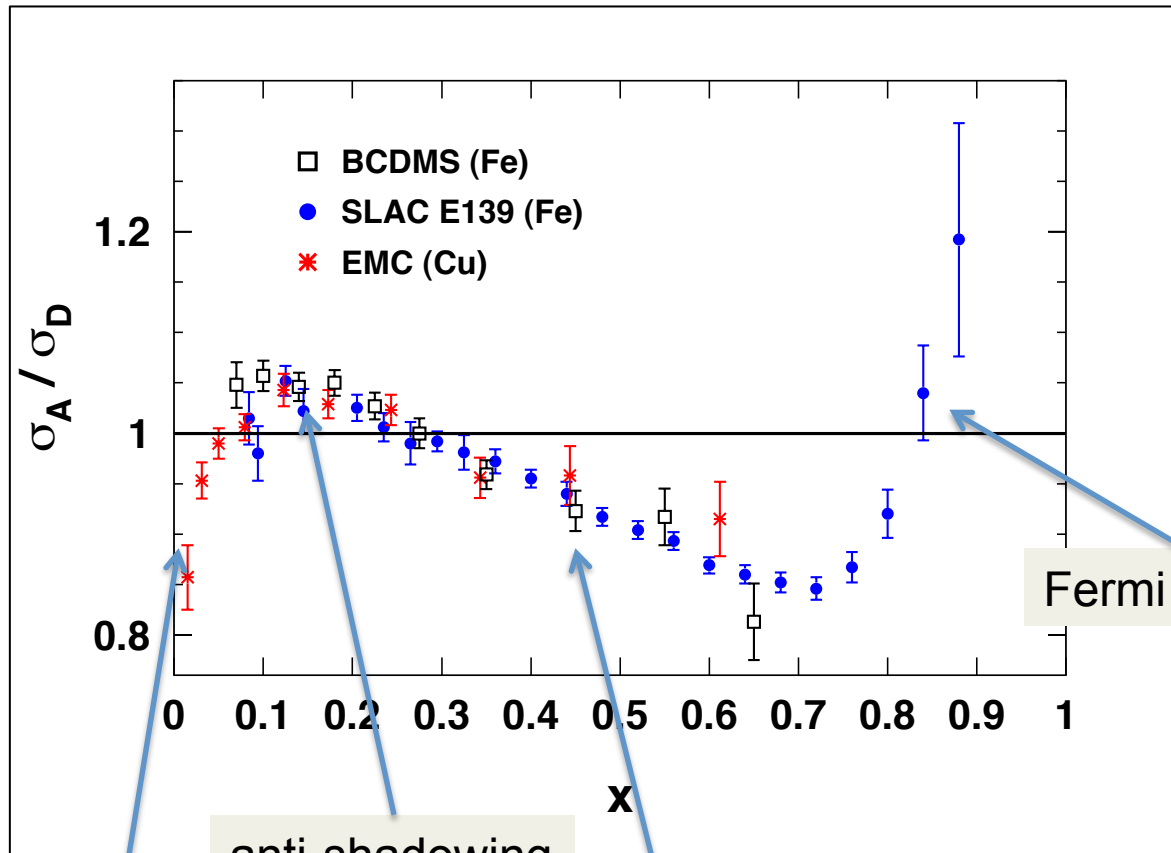
$$\frac{F_2^n}{F_2^p}$$

- SLAC param.  $(1-0.8x)$
- CTEQ
- NMC fit

Isoscalar correction applied to data



# Properties of the EMC Effect



Global properties of the EMC effect

1. Universal x-dependence

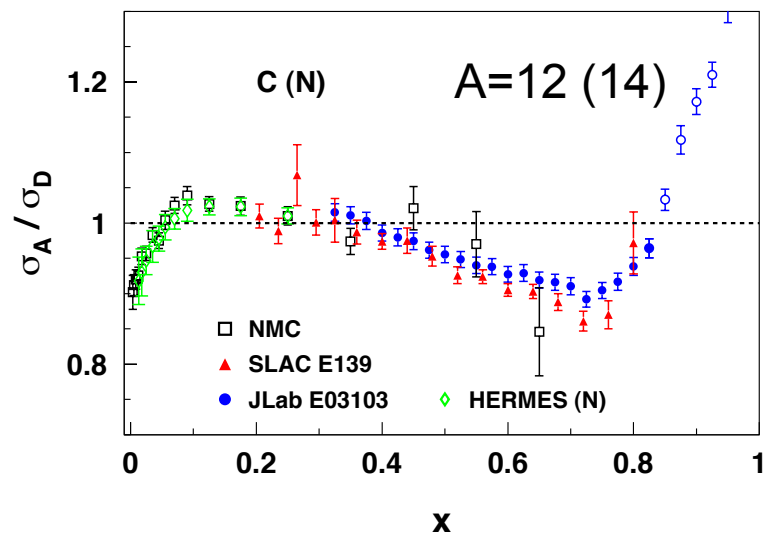
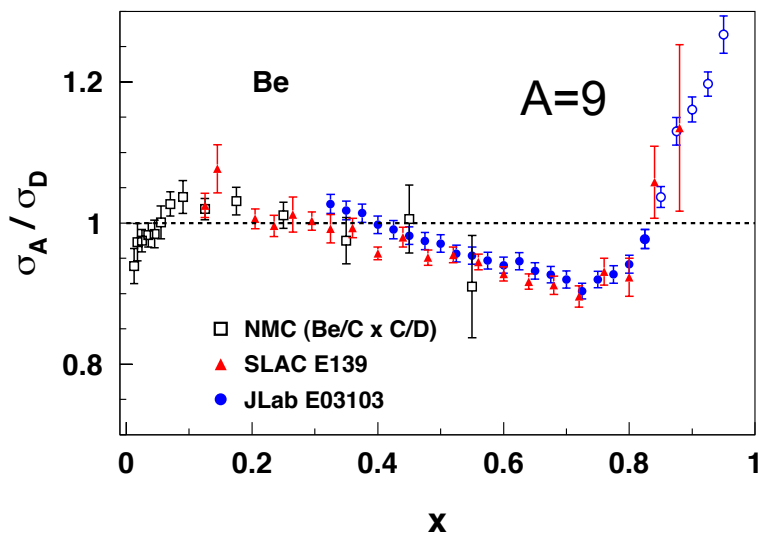
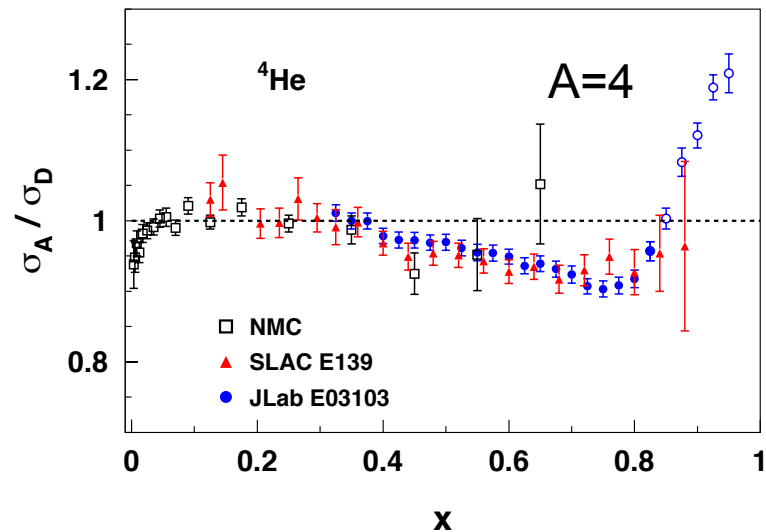
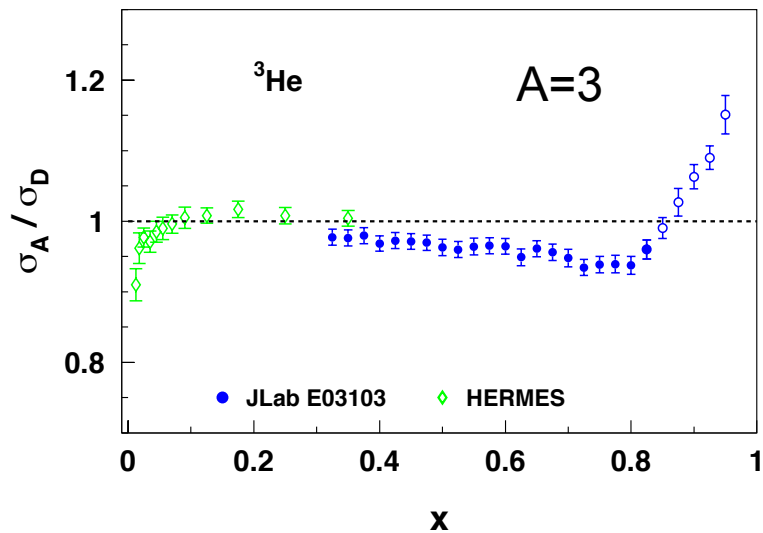
Fermi motion

anti-shadowing

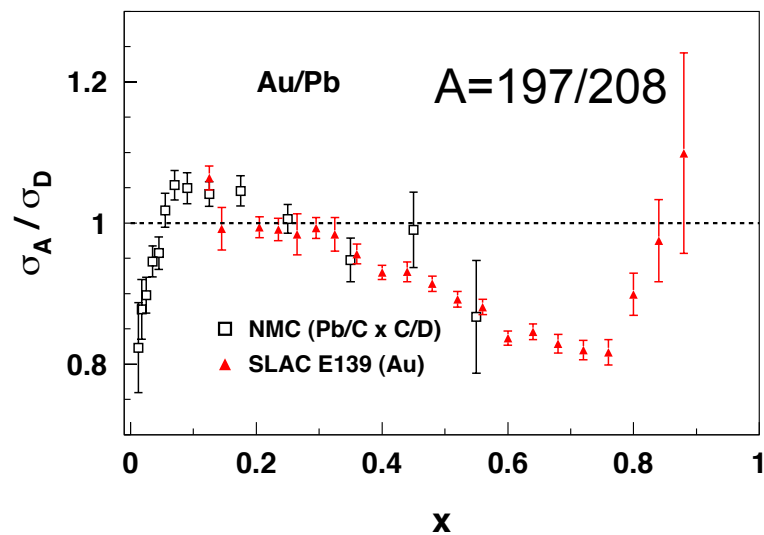
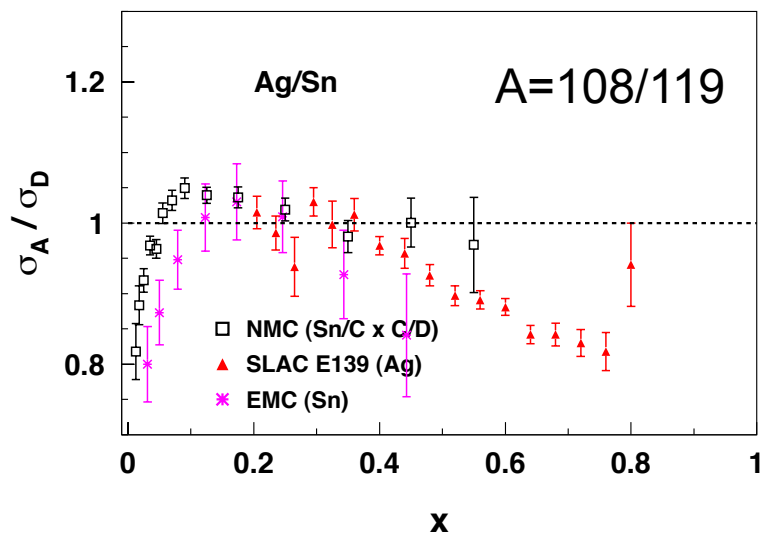
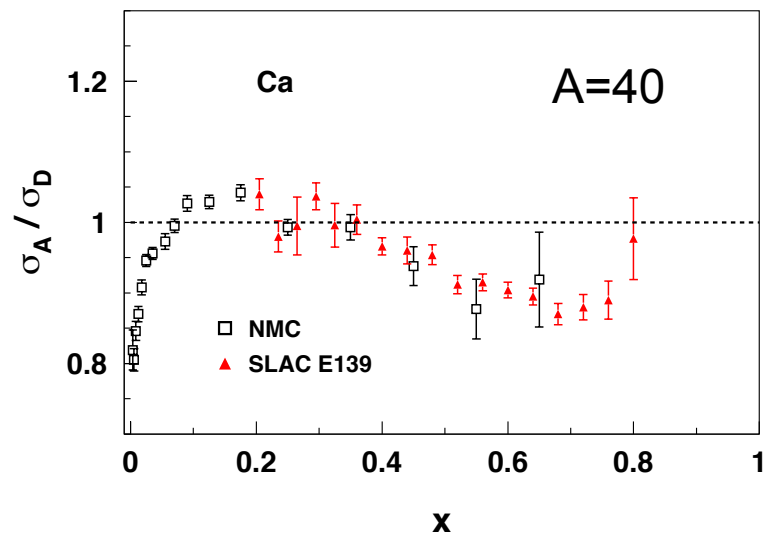
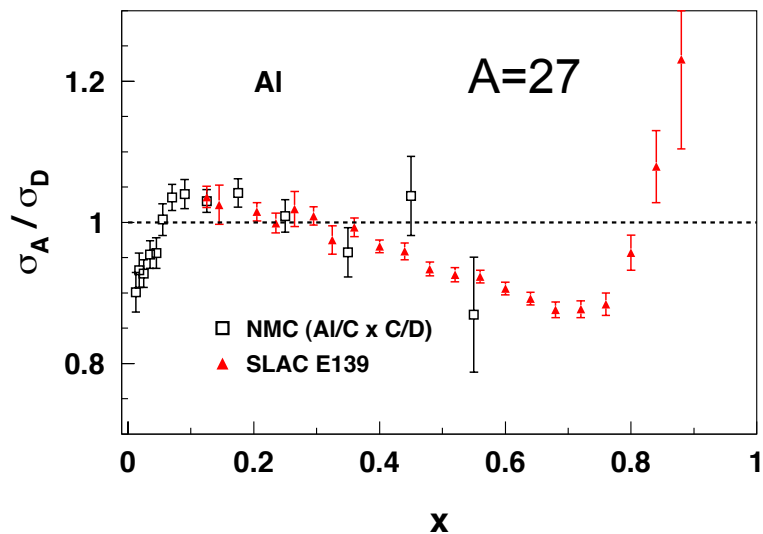
EMC-region

shadowing

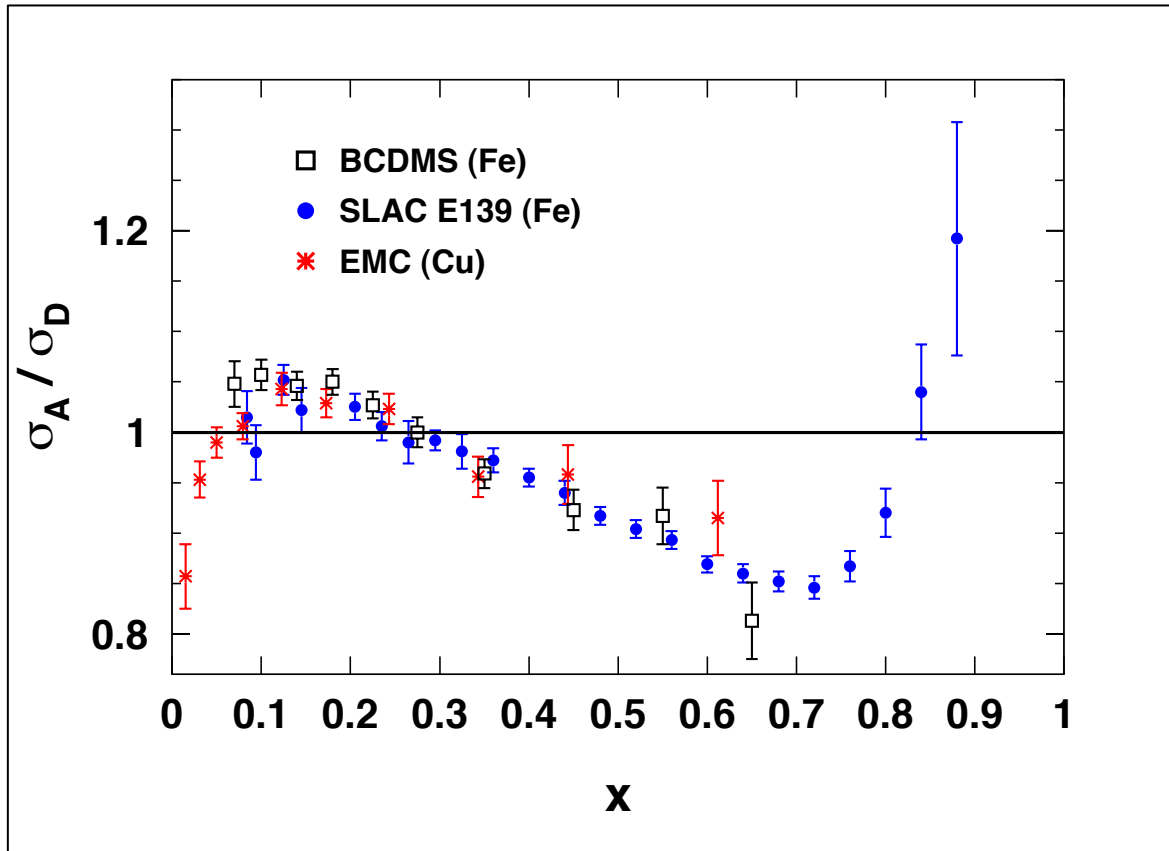
# x Dependence



# x Dependence



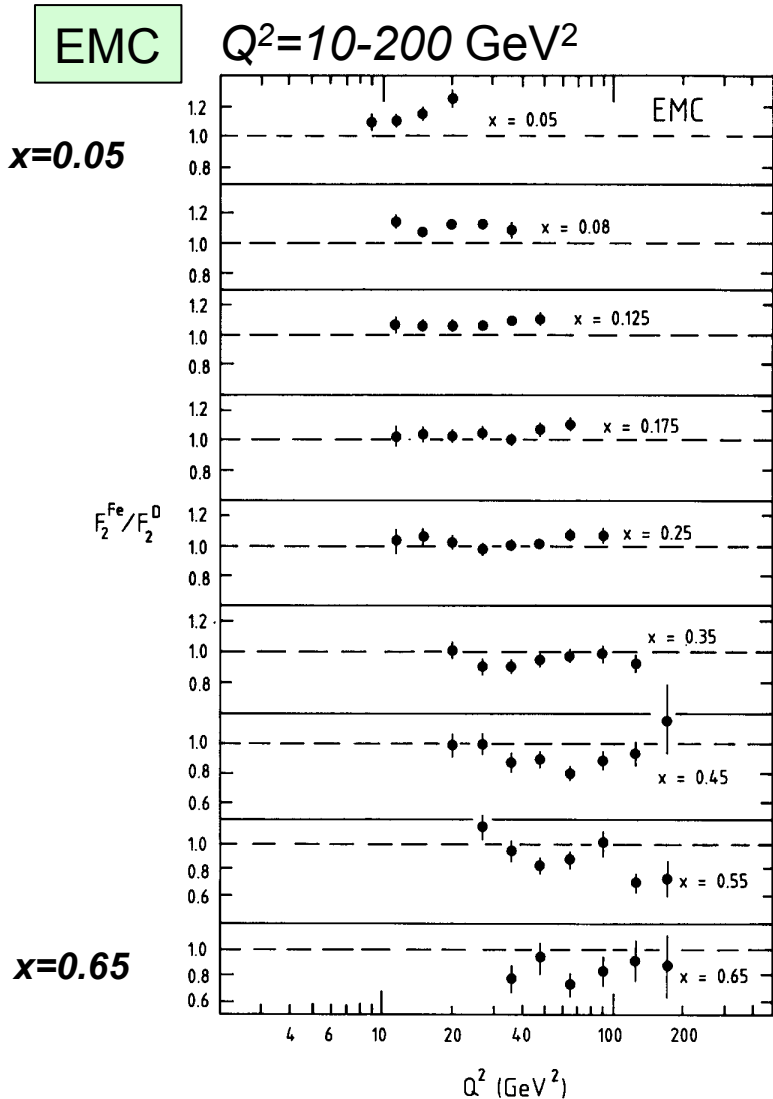
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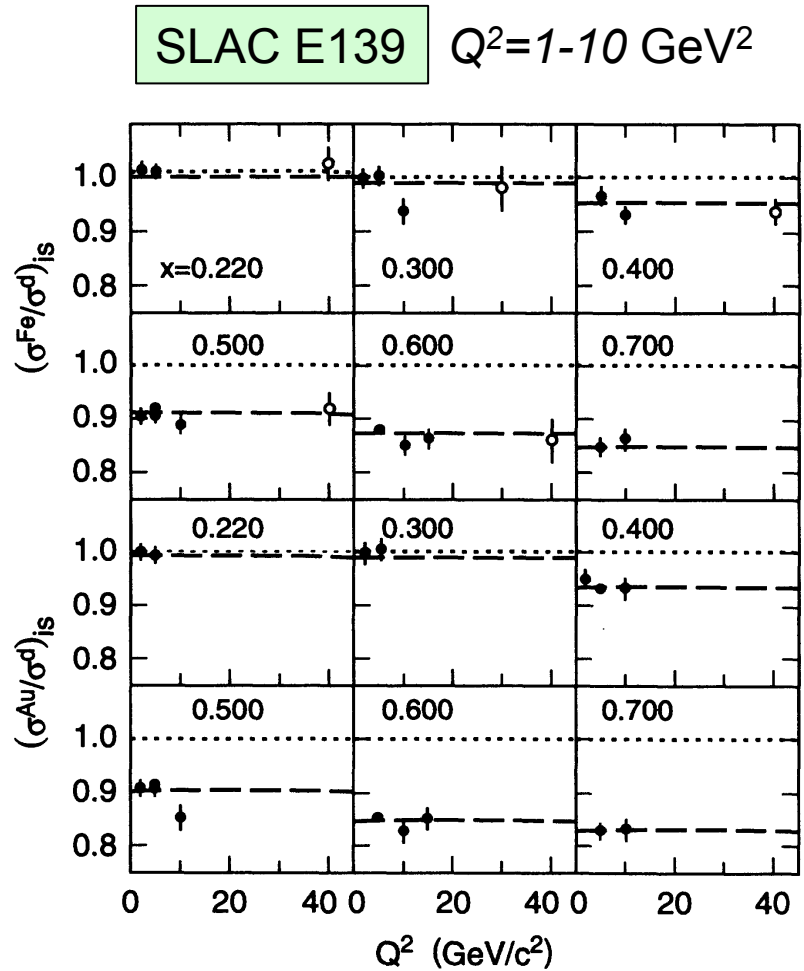
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1. Universal x-dependence
2. Little  $Q^2$  dependence\*

# Q<sup>2</sup> Dependence of the EMC Effect

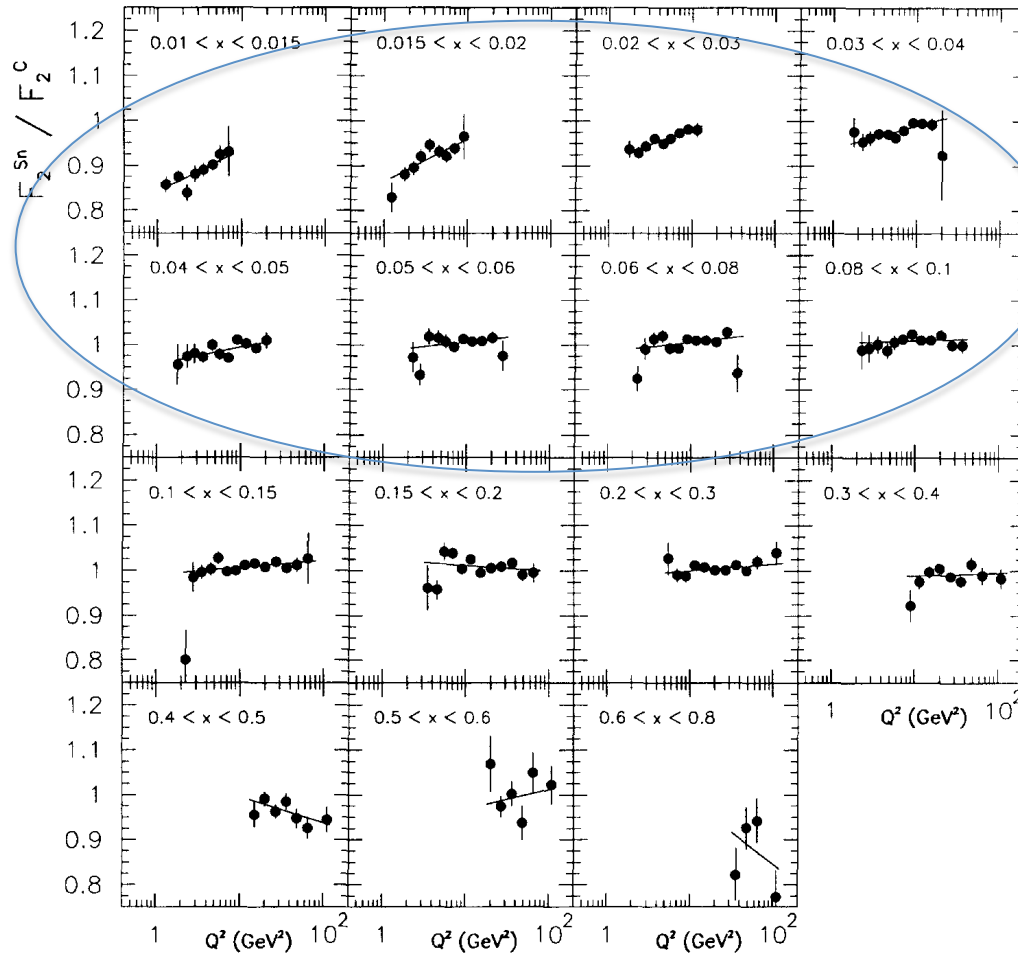


Aubert et al, Nucl. Phys. B293, 740 (1987)



Gomez et al, Phys. Rev. D 49, 4348 (1994)

# (\*) $Q^2$ Dependence of Sn/C



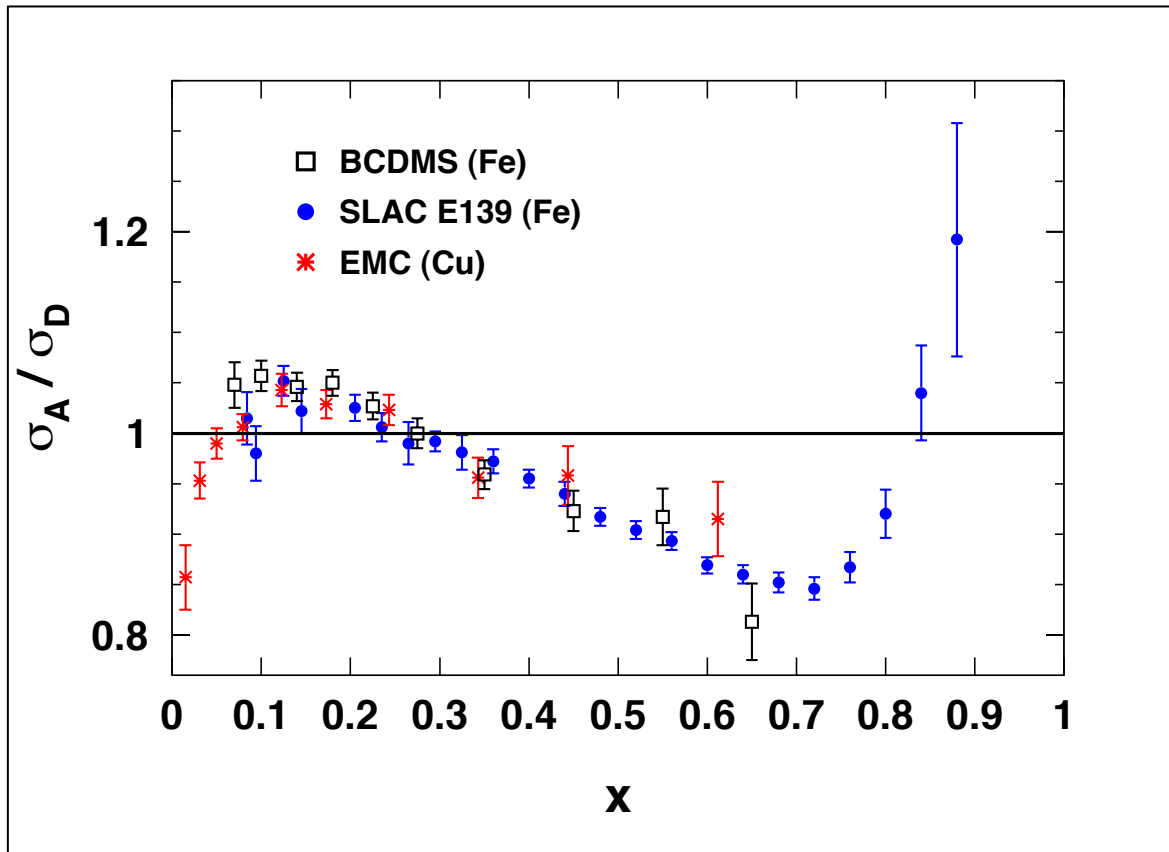
NMC measured non-zero  $Q^2$  dependence in Sn/C ratio at low small  $x$

→ This result is in some tension with other NMC C/D and HERMES Kr/D results

Arneodo et al, Nucl. Phys. B 481, 23 (1996)



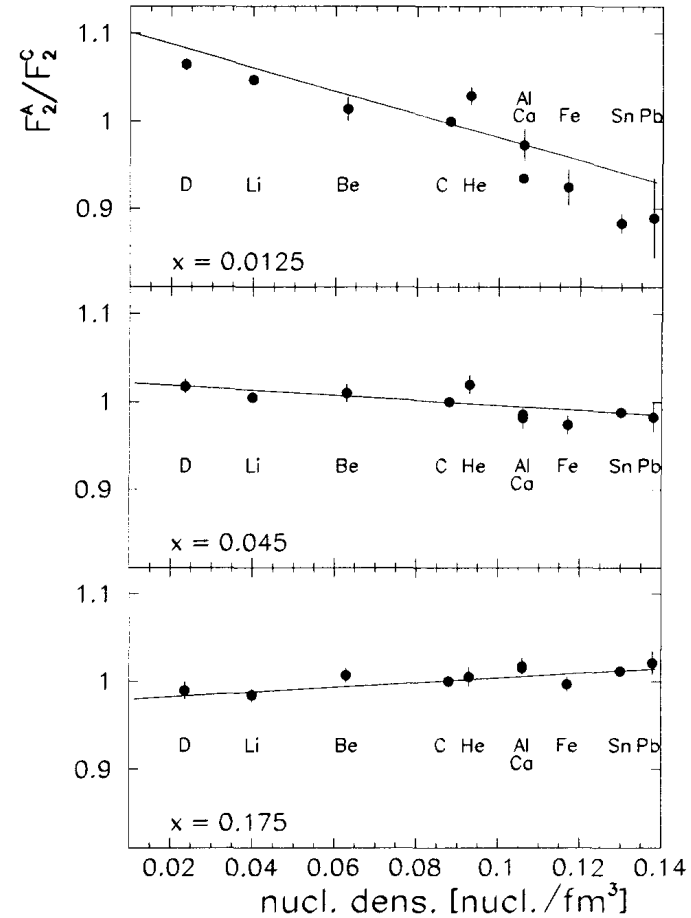
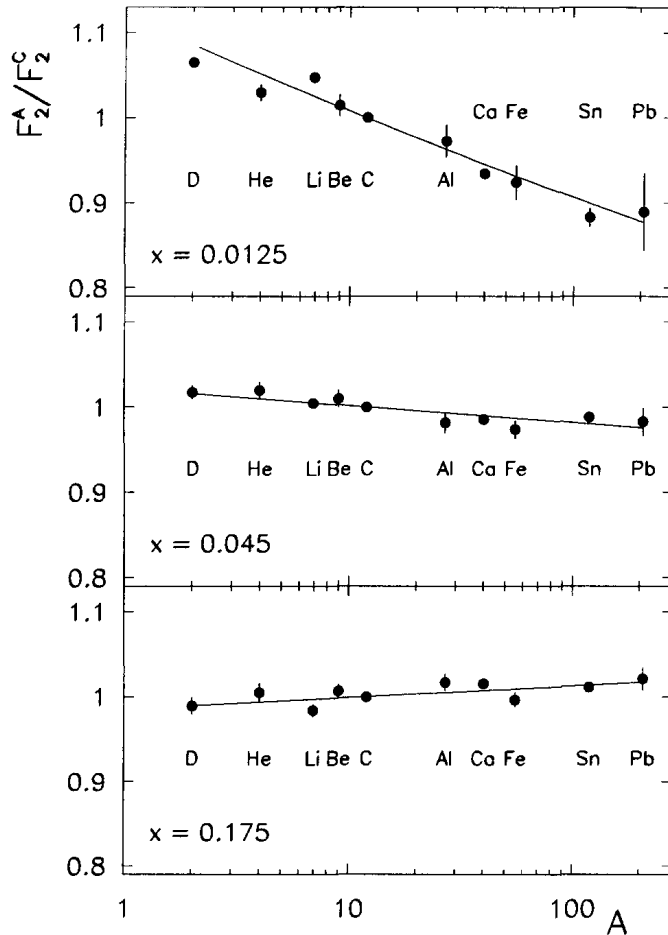
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Global properties of the EMC effect

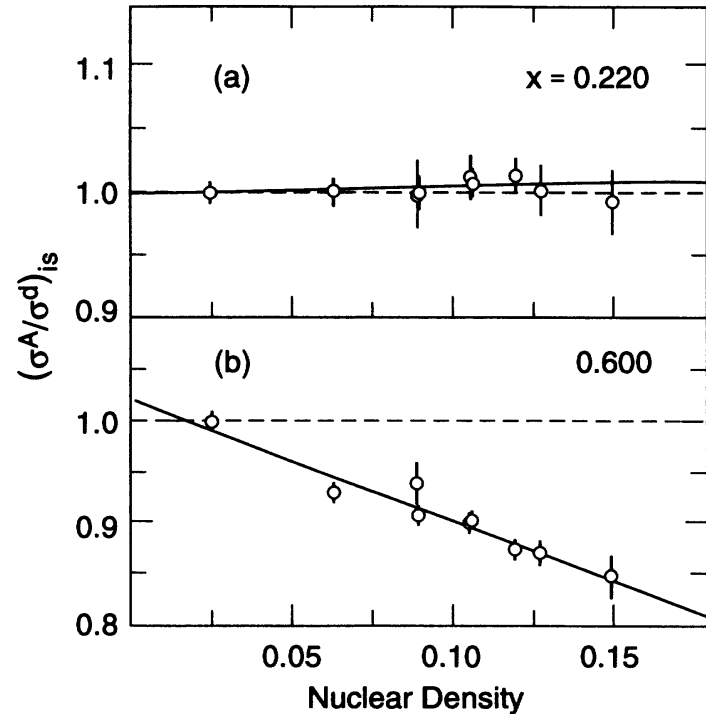
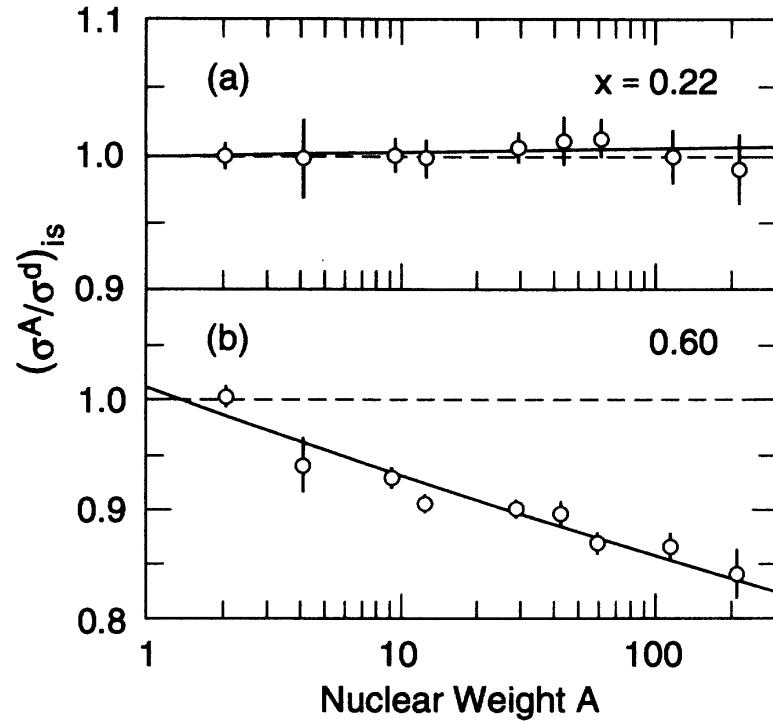
1. Universal x-dependence
  2. Little  $Q^2$  dependence
  3. EMC effect increases with  $A$
- *Anti-shadowing region shows little nuclear dependence*

# A-Dependence of EMC Effect



NMC: Arneodo et al, Nucl. Phys. B 481, 3 (1996)

# A-Dependence of EMC Effect



$$\rho = 3A/4\pi R_e^3 \quad R_e^2 = 5\langle r^2 \rangle / 3$$

$\langle r^2 \rangle$  = RMS electron scattering radius

SLAC E139: *Gomez et al, PRD 49, 4348 (1992)*

# EMC Effect Measurements at Large x

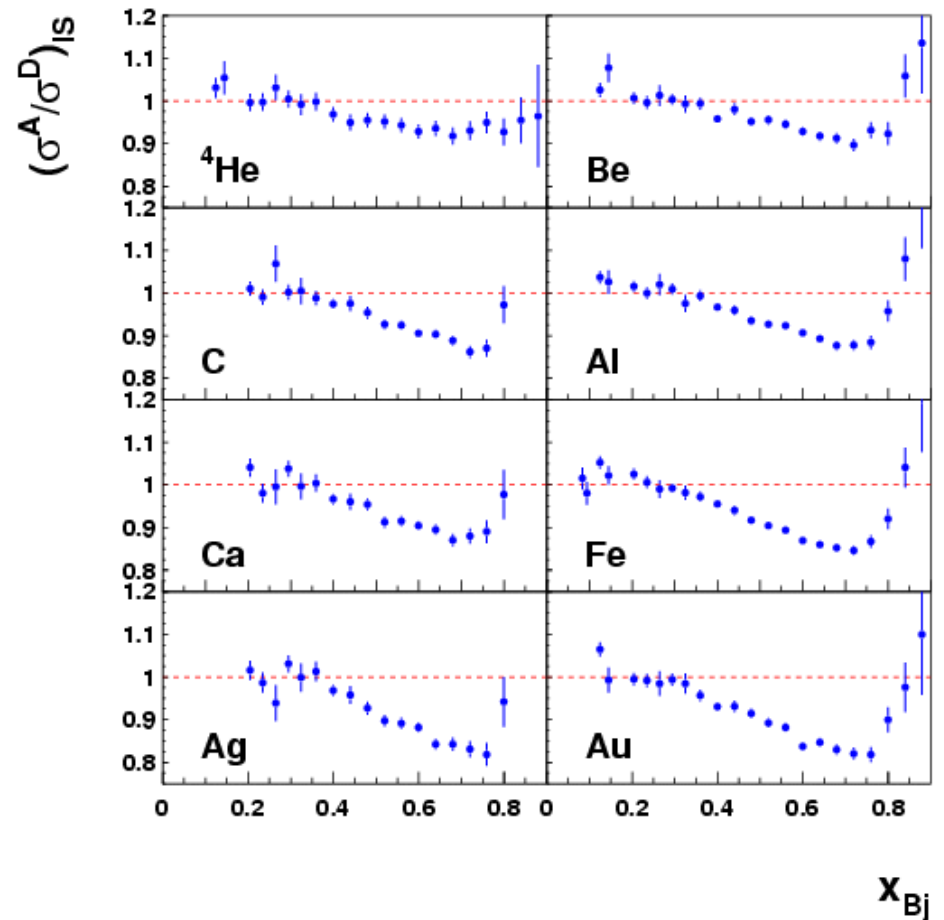
SLAC E139 provided the most extensive and precise data set for  $x > 0.2$

Measured  $\sigma_A/\sigma_D$  for  $A=4$  to 197  
→  ${}^4\text{He}$ ,  ${}^9\text{Be}$ ,  $\text{C}$ ,  ${}^{27}\text{Al}$ ,  ${}^{40}\text{Ca}$ ,  ${}^{56}\text{Fe}$ ,  
 ${}^{108}\text{Ag}$ , and  ${}^{197}\text{Au}$

→ Best determination of the  $A$  dependence  
→ Verified that the  $x$  dependence was roughly constant

Building on the SLAC data  
→ Higher precision data for  ${}^4\text{He}$   
→ Addition of  ${}^3\text{He}$   
→ Precision data at large  $x$

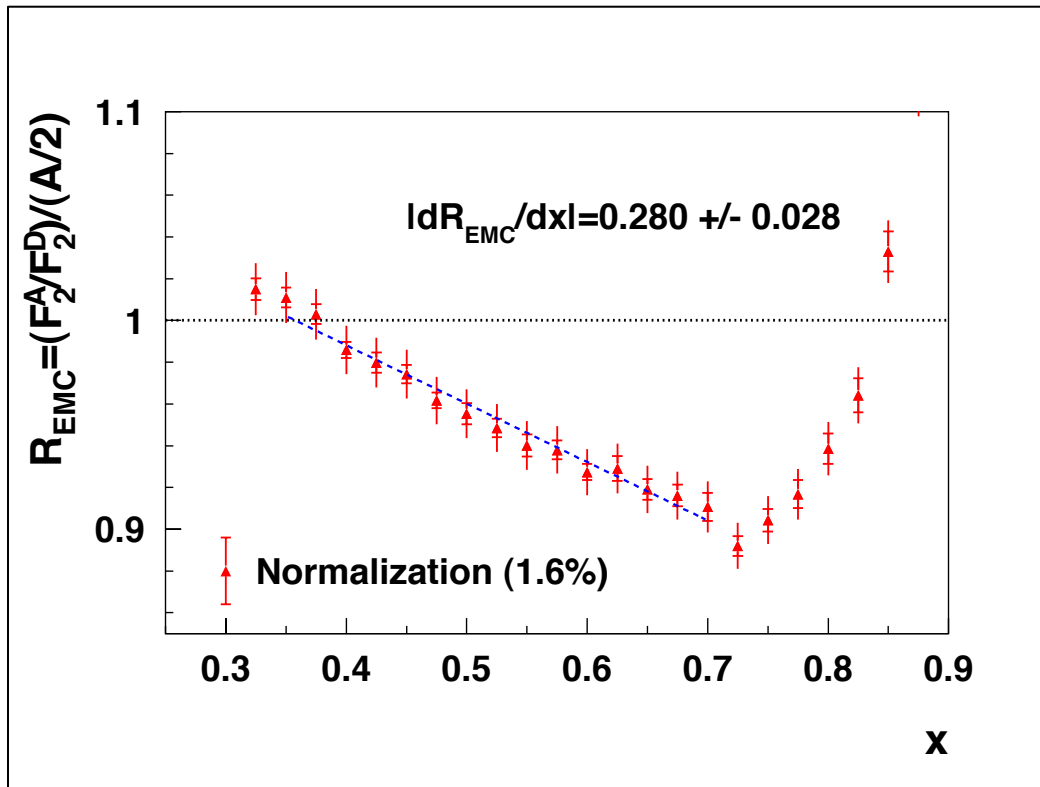
SLAC E139



# JLab E03103

E03103 in Hall C at Jefferson Lab ran Fall 2004

- Measured EMC ratios for light nuclei ( $^3\text{He}$ ,  $^4\text{He}$ , Be, and C)
- Results consistent with previous world data
- Examined nuclear dependence a la E139



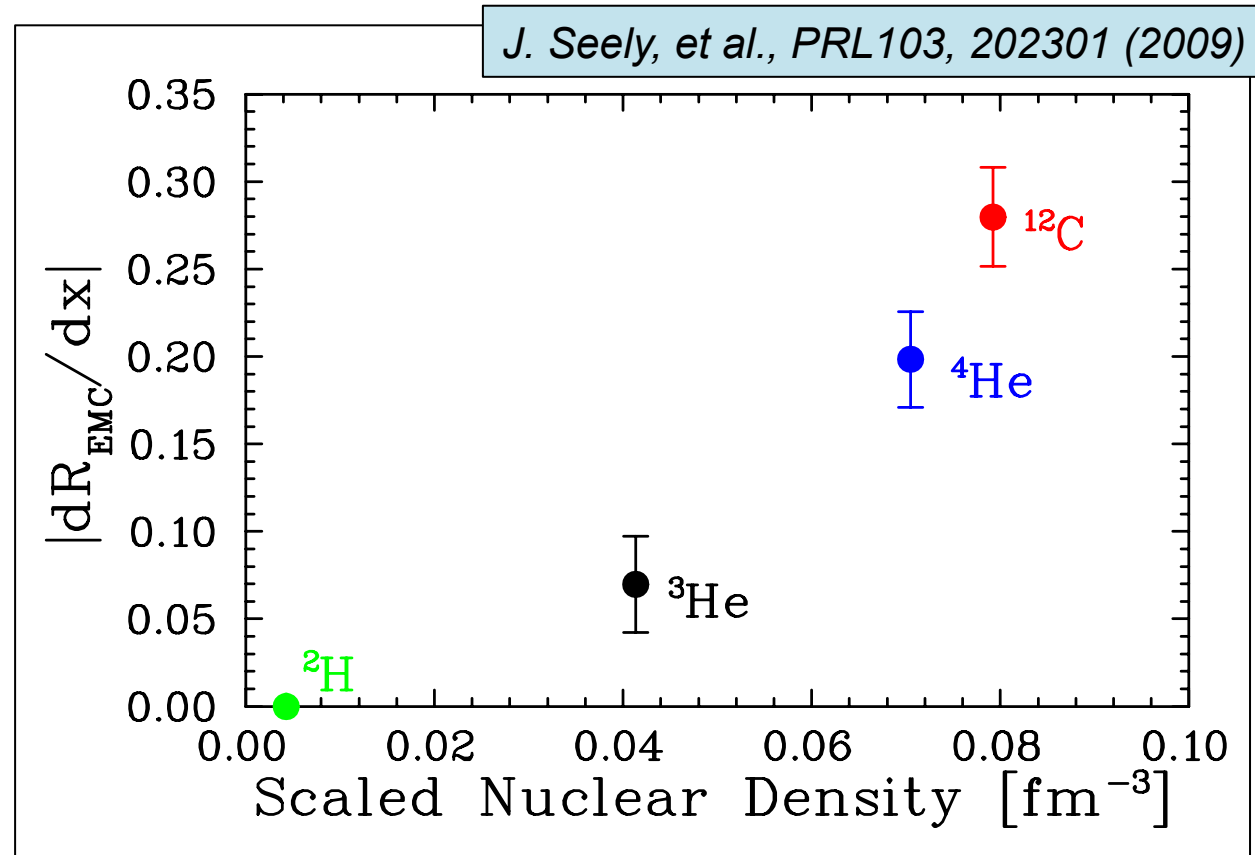
New definition of “size” of the EMC effect  
→ Slope of line fit from  $x=0.35$  to  $0.7$

Definition assumes shape of the EMC effect is universal for nuclei  
→ Data *not inconsistent* with this assumption  
→ Normalization errors mean we can only confirm this at 1-1.5% level

# JLab E03103 Results

E03103 measured  $\sigma_A/\sigma_D$   
for  $^3\text{He}$ ,  $^4\text{He}$ , Be, C

→  $^3\text{He}$ ,  $^4\text{He}$ , C, EMC  
effect scales well with  
density



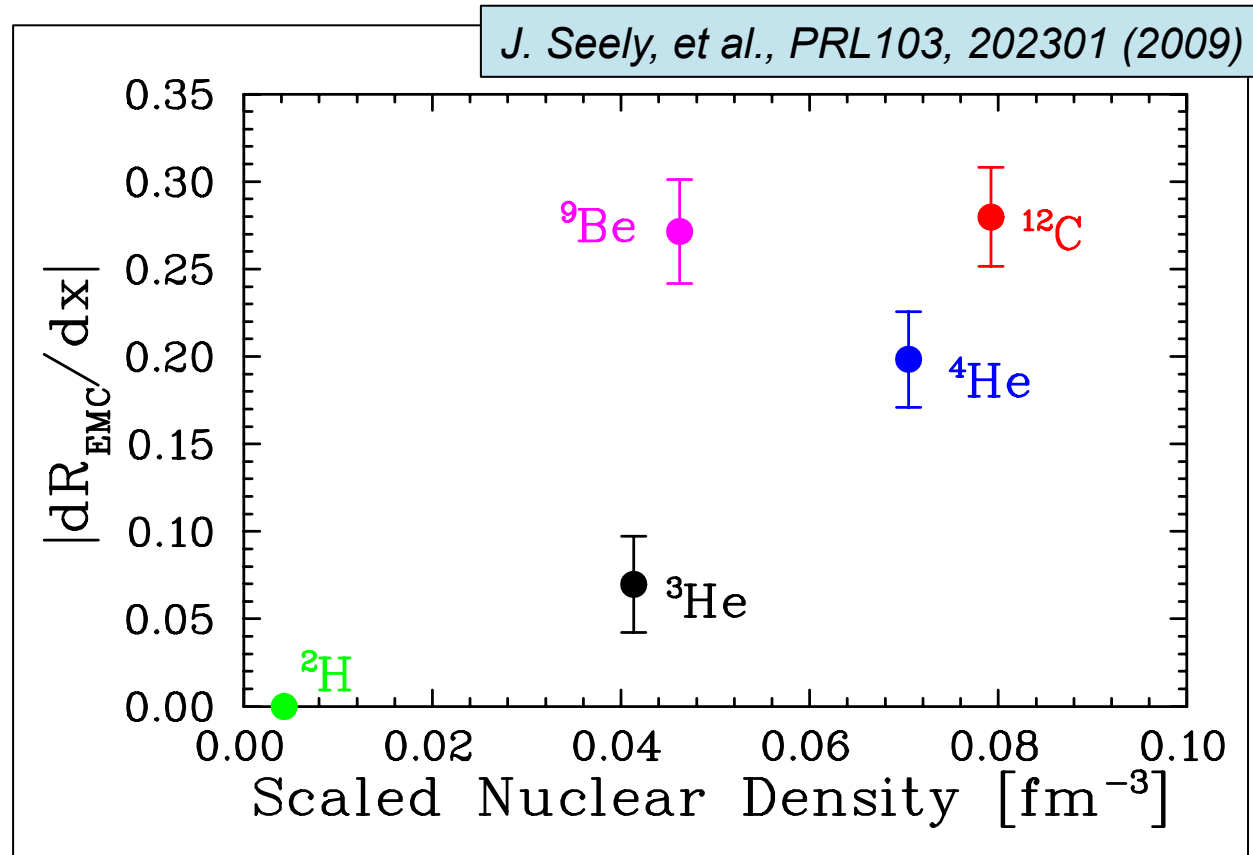
Scaled nuclear density =  $(A-1)/A \langle\rho\rangle$   
→ remove contribution from struck nucleon

$\langle\rho\rangle$  from ab initio few-body calculations  
→ [S.C. Pieper and R.B. Wiringa, *Ann. Rev. Nucl. Part. Sci* 51, 53 (2001)]

# JLab E03103 Results

E03103 measured  $\sigma_A/\sigma_D$   
for  $^3\text{He}$ ,  $^4\text{He}$ , Be, C

→  $^3\text{He}$ ,  $^4\text{He}$ , C, EMC  
effect scales well with  
density  
→ Be does not fit the  
trend



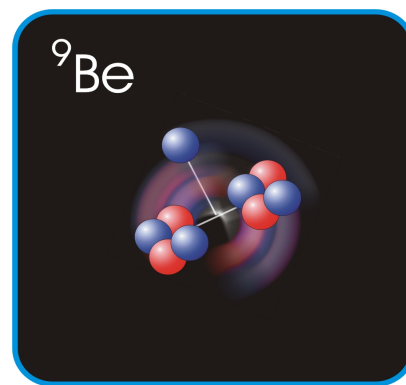
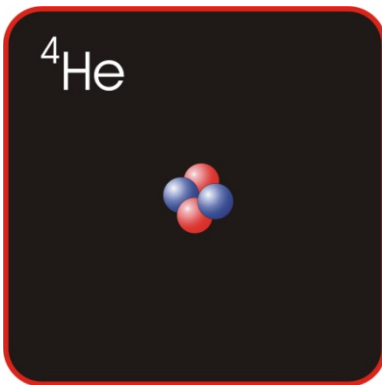
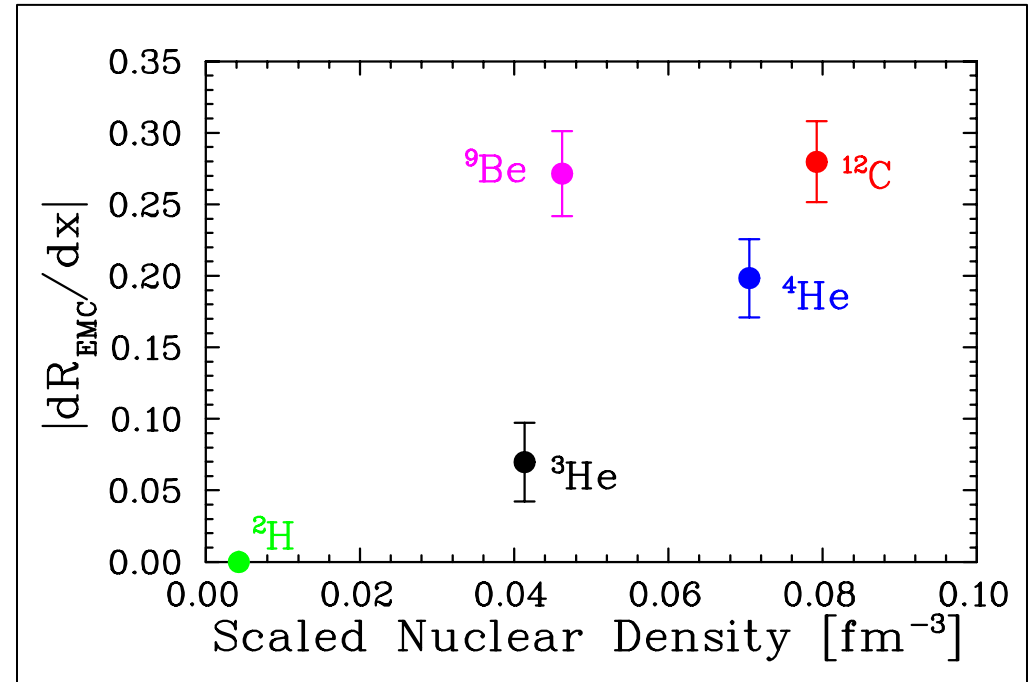
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$\langle\rho\rangle$  from ab initio few-body calculations  
→ [S.C. Pieper and R.B. Wiringa, *Ann. Rev. Nucl. Part. Sci* 51, 53 (2001)]

# EMC Effect and Local Nuclear Density

${}^9\text{Be}$  has low average density  
→ Large component of structure is  $2\alpha+n$   
→ Most nucleons in tight,  $\alpha$ -like configurations

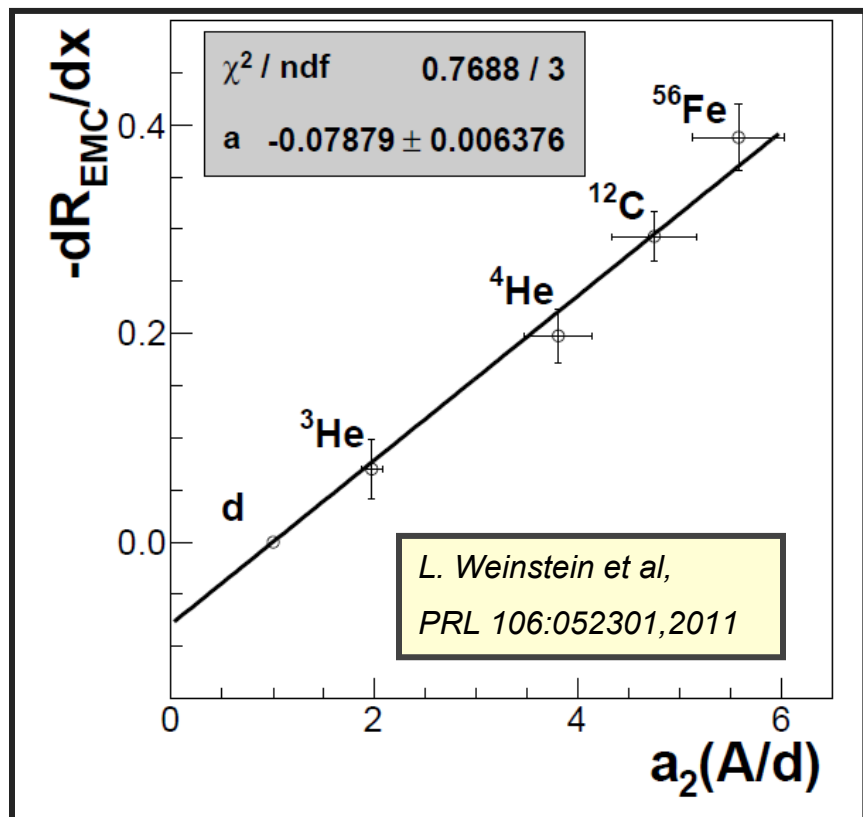
EMC effect driven by *local* rather than *average* nuclear density



“Local density” is appealing in that it makes sense intuitively – can we make this more quantitative?



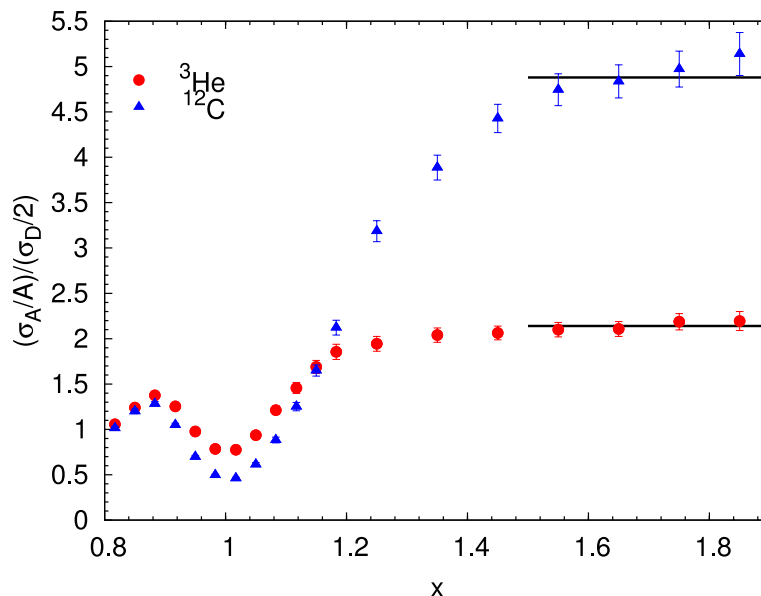
# EMC Effect and Short Range Correlations



Weinstein et al observed linear correlation between size of EMC effect and Short Range Correlation “plateau”

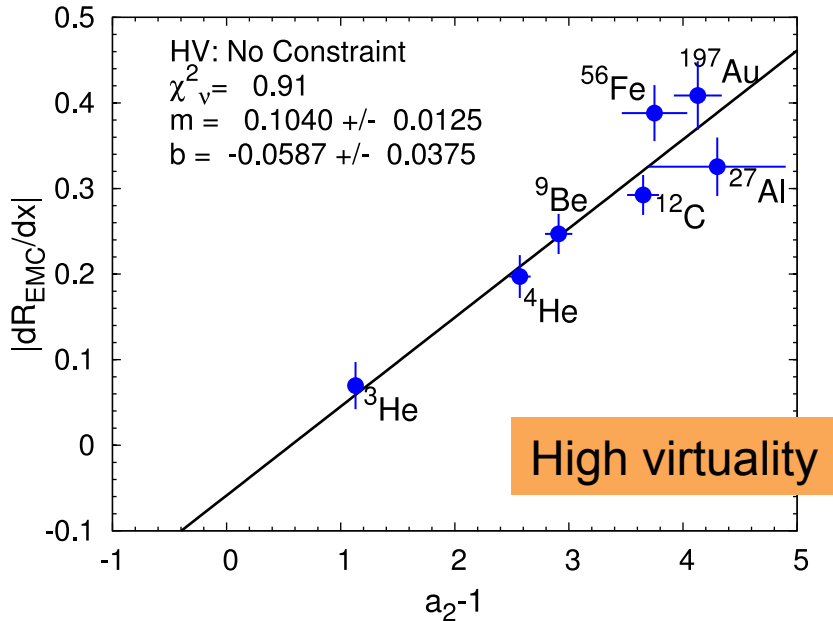
→ Observing Short Range Correlations requires measurements at  $x > 1$   
 → Reaction dynamics very different – DIS vs. QE scattering, why the same nuclear dependence?

$$\frac{2}{A} \frac{\sigma_A}{\sigma_D} = a_2(A)$$

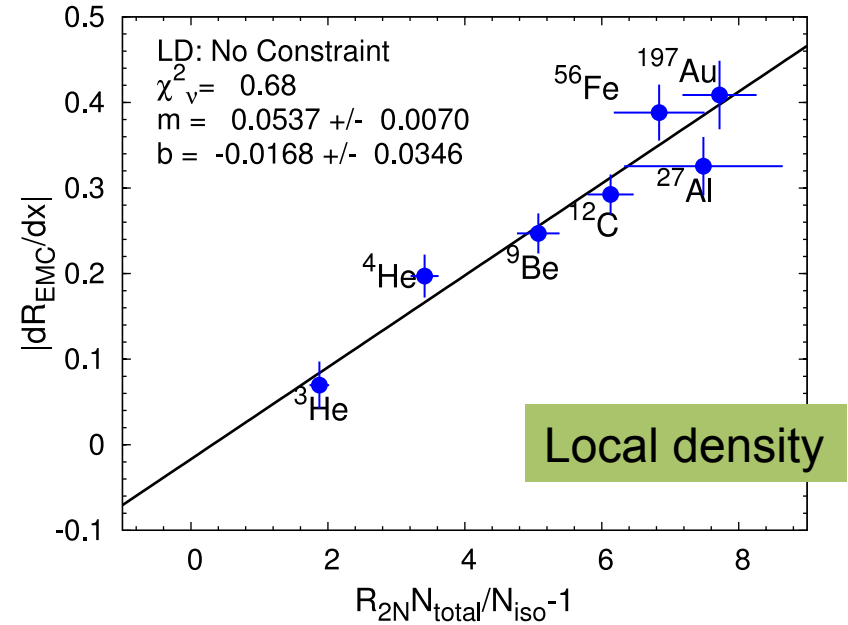


# Nuclear Dependence of EMC and SRCs

Arrington et al, PRC 86, 065204 (2012)



$a_2 \sim$  number of high momentum nucleons

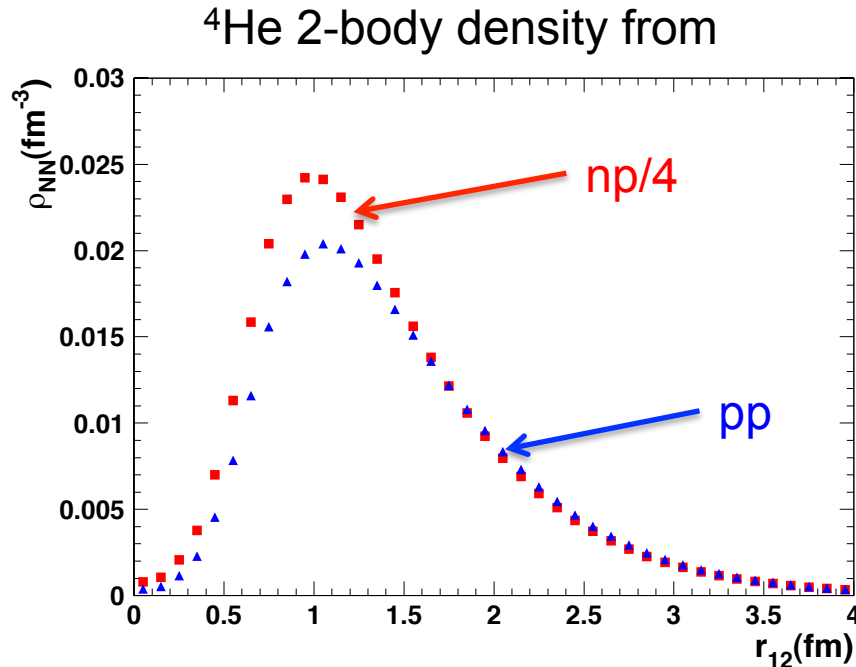


$R_{2N} \sim$  number of nucleons “close” together

Detailed study of nuclear dependence of EMC effect and SRCs (see N. Fomin’s talk from Monday) does not favor either picture

**Can we distinguish between these two pictures via some new observable?  $\rightarrow$  Flavor dependence of the EMC effect**

# Flavor dependence and SRCs



S.C. Pieper and R.B. Wiringa, *Ann. Rev. Nucl. Part. Sci.* 51, 53 (2001)

High momentum nucleons from SRCs emerge from tensor part of  $NN$  interaction –  $np$  pairs dominate

→ Probability to find 2 nucleons “close” together nearly the same for  $np$ ,  $nn$ ,  $pp$

For  $r_{12} < 1.7$  fm:

$$P_{pp} = P_{nn} \approx 0.8P_{np}$$

If EMC effect due to **high virtuality**, flavor dependence of EMC effect emerges naturally

→ If EMC effect from **local density**,  $np/pp/nn$  pairs all contribute (roughly) equally

# Flavor dependence and SRCs

High momentum nucleons in the nucleus come primarily from  $np$  pairs

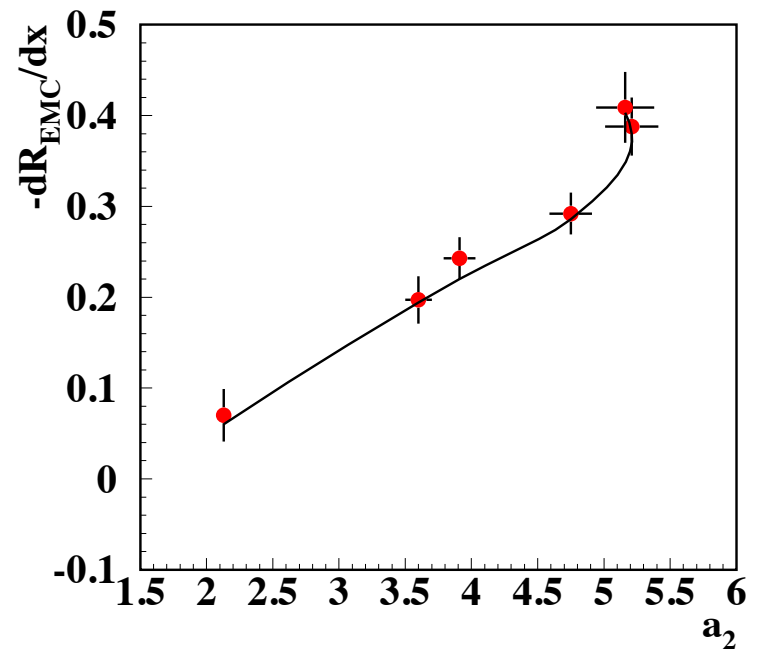
→ The relative probability to find a high momentum proton is larger than for neutron for  $N > Z$  nuclei

$$n_p^A(p) \approx \frac{1}{2x_p} a_2(A, y) n_d(p) \quad x_p = \frac{Z}{A}$$

$$n_n^A(p) \approx \frac{1}{2x_n} a_2(A, y) n_d(p) \quad x_n = \frac{A - Z}{A}$$

Probability to find SRC

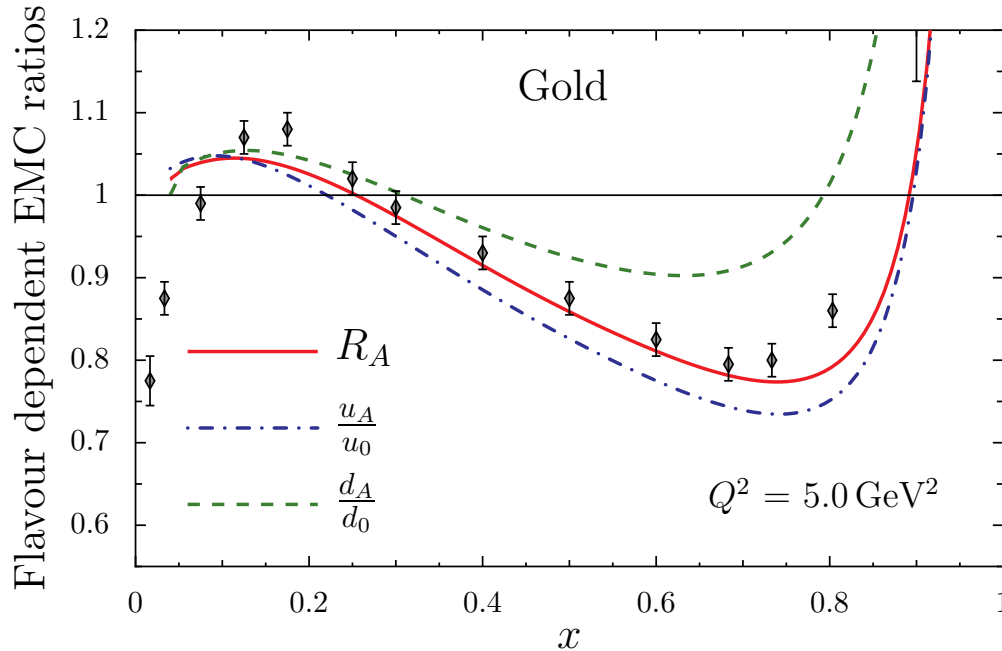
$$u_A = \frac{Z\tilde{u}_p + N\tilde{d}_p}{A} \quad d_A = \frac{Z\tilde{d}_p + N\tilde{u}_p}{A}$$



Under the assumption the EMC effect comes from “high virtuality” (high momentum nucleons), effect driven by protons (u-quark dominates) → similar flavor dependence is seen in some “mean-field” approaches

# Flavor Dependence of the EMC Effect

Mean-field calculations predict a flavor dependent EMC effect for  $N \neq Z$  nuclei



Cloët, Bentz, and Thomas, *PRL* 102, 252301 (2009)

Medium modified  
quark distributions

$$u_A = \frac{Z\tilde{u}_p + N\tilde{d}_p}{A} \quad d_A = \frac{Z\tilde{d}_p + N\tilde{u}_p}{A}$$

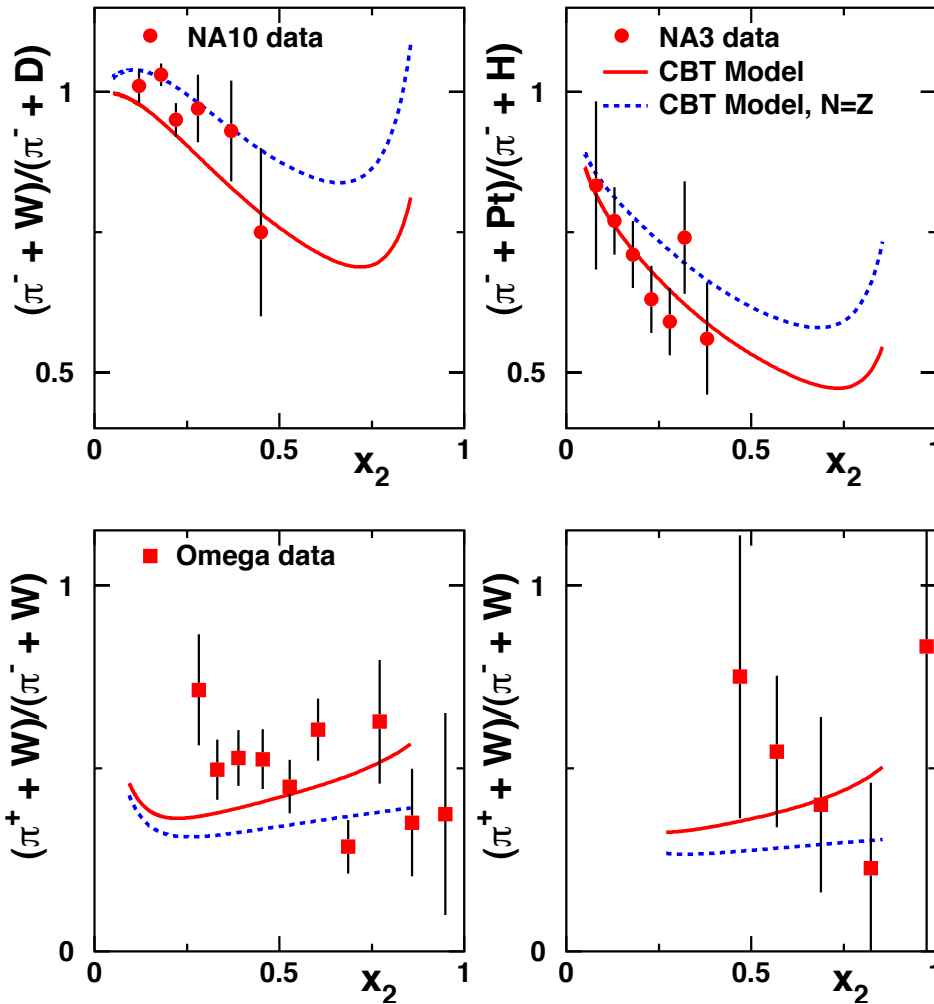
Free nucleon  
quark distributions

$$u_0 = \frac{Zu_p + Nd_p}{A} \quad d_0 = \frac{Zd_p + Nu_p}{A}$$

Isovector-vector mean field ( $\rho$ ) causes u (d) quark to feel additional vector attraction (repulsion) in  $N \neq Z$  nuclei

*Experimentally, this flavor dependence has not been observed directly*

# EMC Flavor Dependence: Pion Drell-Yan



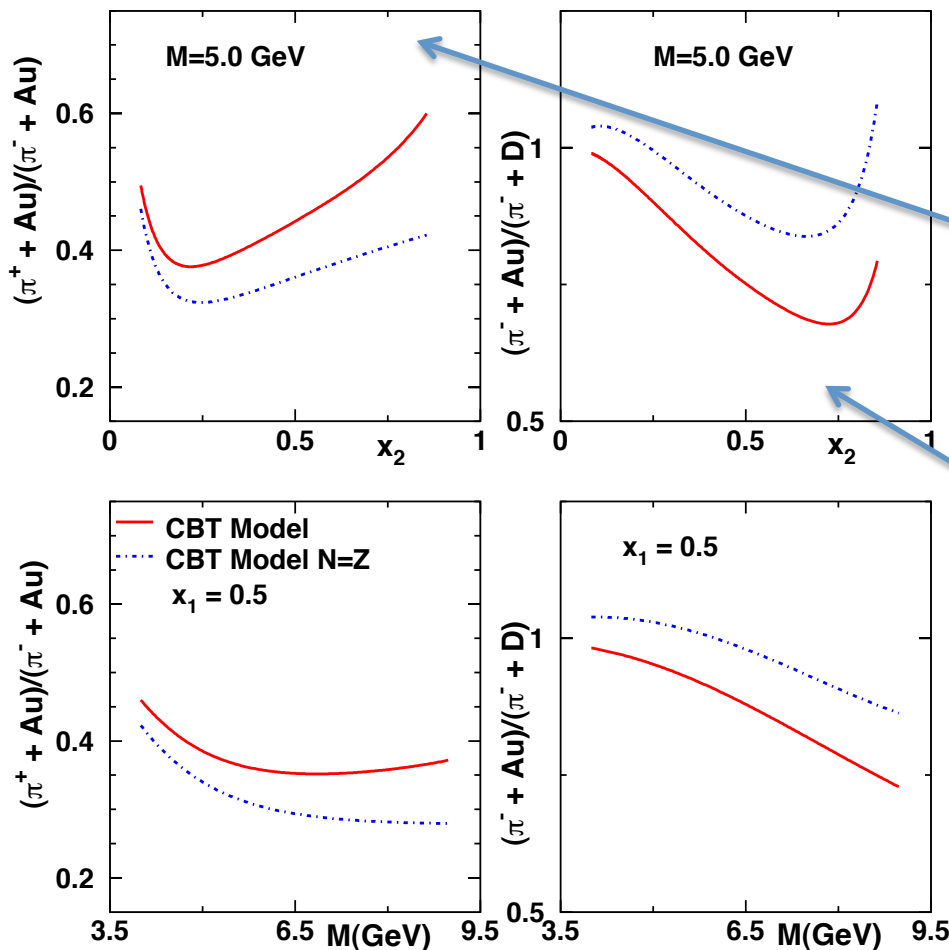
| Experiment          | Flavor Ind. | Flavor dep. |
|---------------------|-------------|-------------|
| NA3                 | 1.3         | 0.5         |
| NA10                | 0.60        | 2.5         |
| Omega (low $Q^2$ )  | 6.2         | 3.2         |
| Omega (high $Q^2$ ) | 1.4         | 0.96        |

$\chi^2/\text{DOF}$

Pion-induced Drell-Yan sensitive to potential flavor dependence, but existing data lack precision

Dutta, Peng, Cloët, DG, PRC 83, 042201 (2011)

# Pion Drell-Yan at COMPASS



160 GeV pions on gold

$$\frac{\sigma^{DY}(\pi^+ + A)}{\sigma^{DY}(\pi^- + A)} \approx \frac{d_A(x)}{4u_A(x)}$$

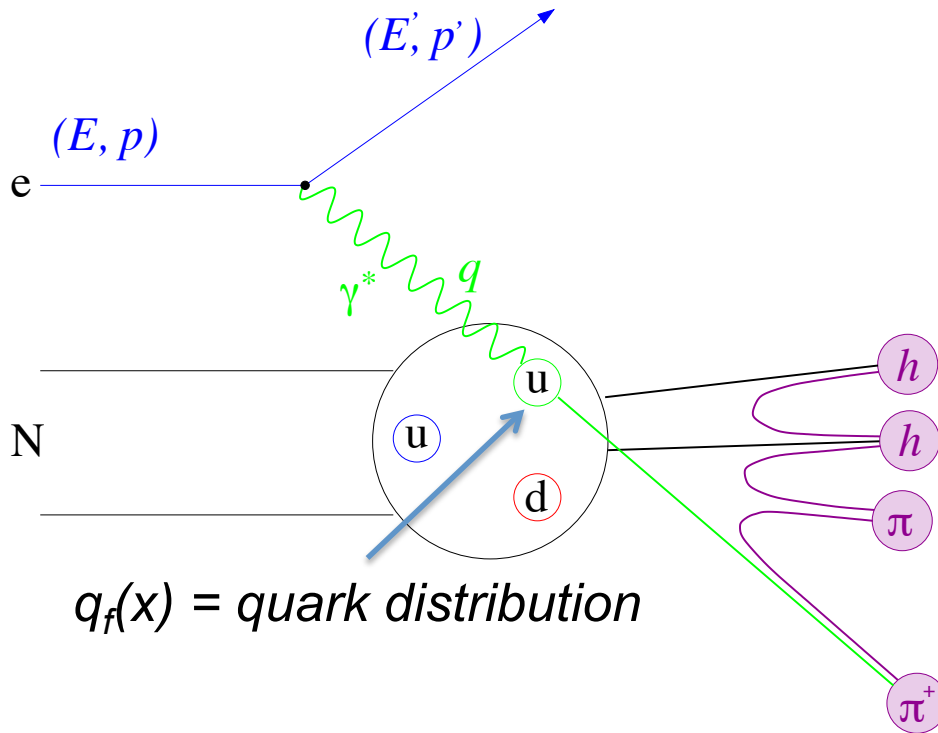
$$\frac{\sigma^{DY}(\pi^- + A)}{\sigma^{DY}(\pi^- + D)} \approx \frac{u_A(x)}{u_D(x)}$$

*Dutta et al, PRC 83, 042201 (2011)*

First measurements on NH3 (and nuclear targets) planned for 2014

$$\frac{d\sigma_{\pi^\pm A}}{dx_\pi dx_2} = \frac{4\pi\alpha^2}{9sx_\pi x_2} \sum_q e_q^2 [q_{\pi^\pm}(x_\pi)\bar{q}_A(x_2) + \bar{q}_{\pi^\pm}(x_\pi)q_A(x_2)]$$

# Semi-Inclusive DIS



Assuming factorization holds, SIDIS acts as a “flavor tag” for struck quark  
 → Similar to polarized quark distribution extractions

$D_f^h(z)$  – fragmentation function  
 quark of flavor  $f \rightarrow$  hadron  $h$

$$\frac{d\sigma}{dx dQ^2 dz} = \frac{\sum_f e_f^2 q_f(x) D_f^h(z)}{\sum_f e_f^2 q_f(x)} \left( \frac{d\sigma}{dx dQ^2} \right)$$

$x =$  fraction of proton momentum carried by quark

$z = E_{hadron}/\nu$



# Semi-Inclusive DIS

Extract flavor dependence via semi-inclusive pion yields from gold and deuterium

Super-ratio

$$\frac{Y_{Au}^{\pi^+} / Y_{Au}^{\pi^-}}{Y_D^{\pi^+} / Y_D^{\pi^-}}$$

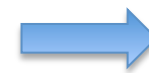
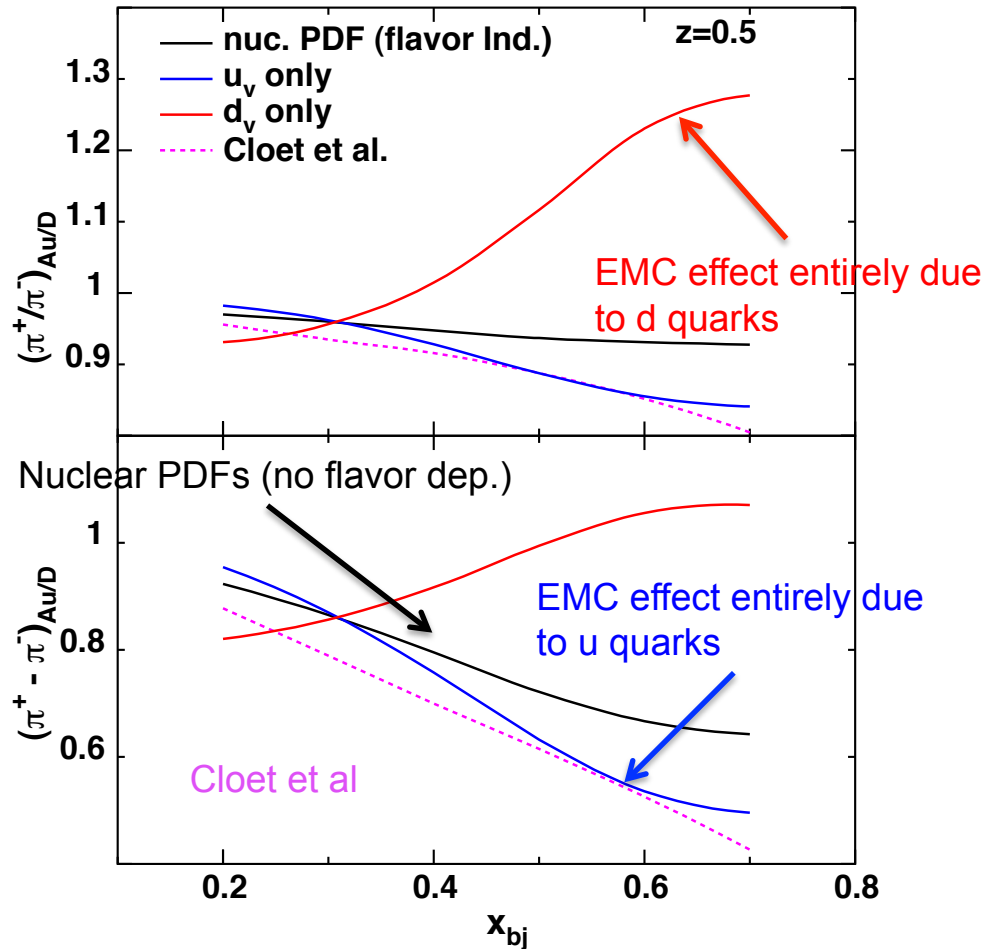
Difference ratio

$$\frac{Y_{Au}^{\pi^+} - Y_{Au}^{\pi^-}}{Y_D^{\pi^+} - Y_D^{\pi^-}}$$

Toy model:

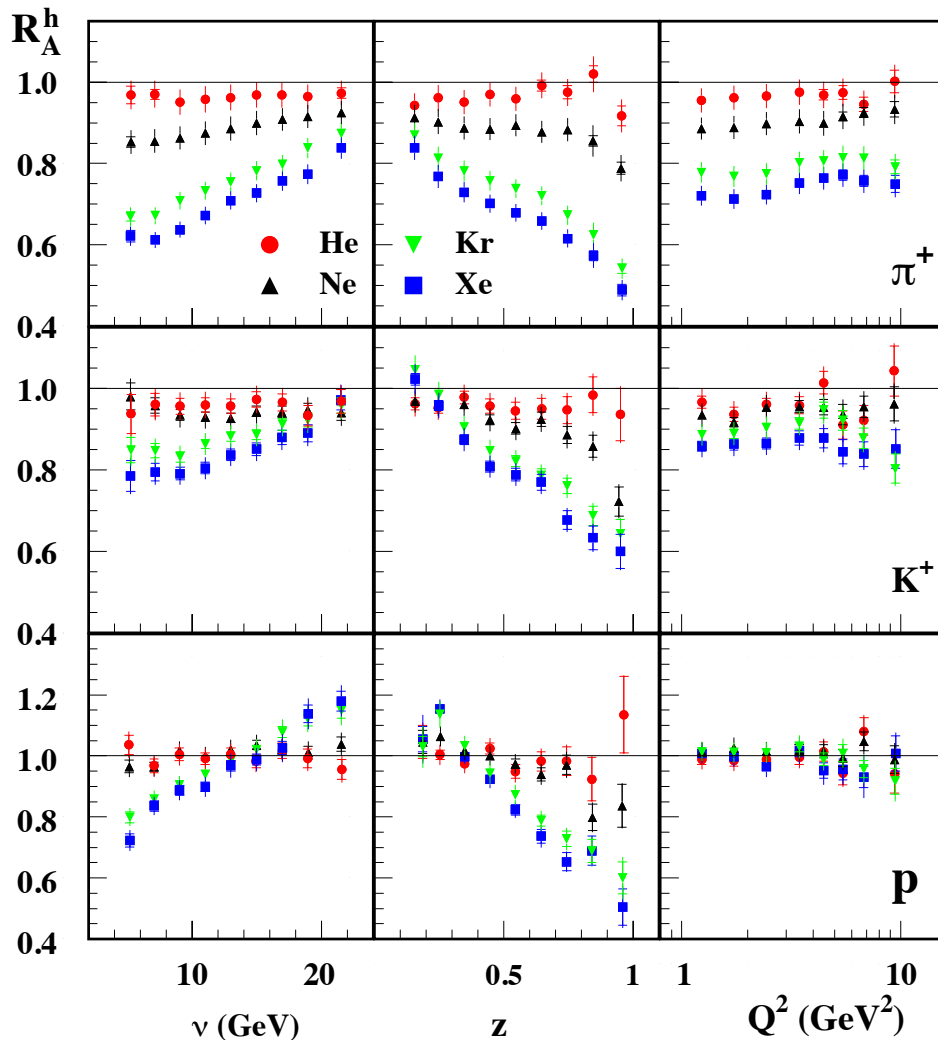
$u_V$  only: EMC effect due to modification of  $u_A$  only

$d_V$  only: EMC effect due to modification of  $d_A$  only



$F_2^A$  unchanged

# SIDIS - Interpretability



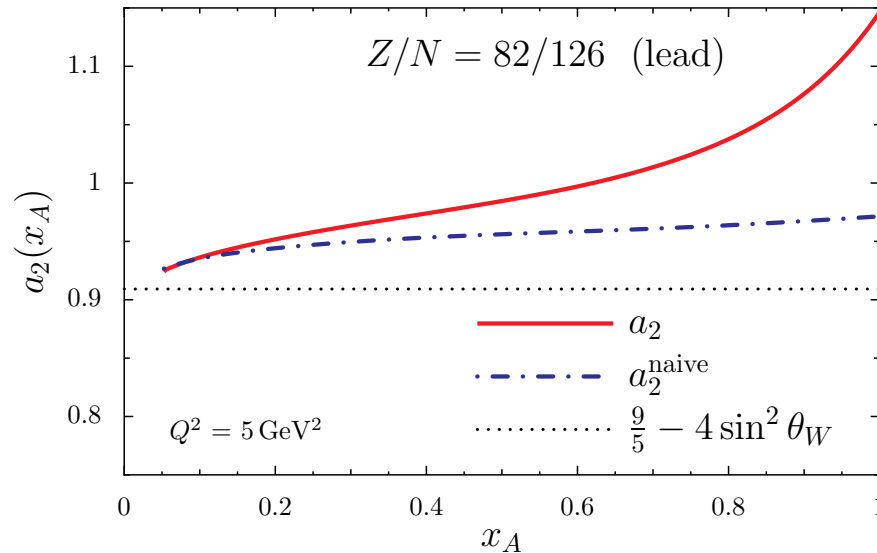
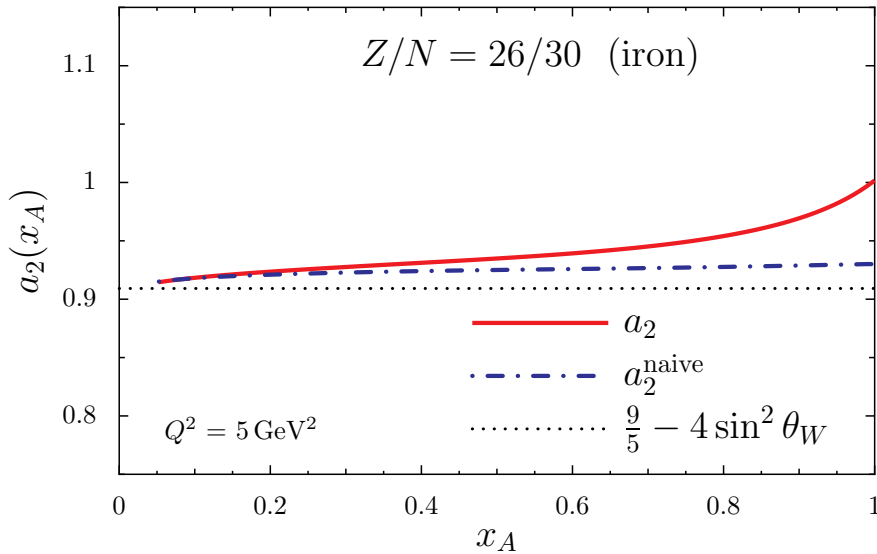
$$R_h^A(z, \nu) = \frac{\left( \frac{1}{\sigma_e} \frac{d\sigma}{dzd\nu} \right)_A}{\left( \frac{1}{\sigma_e} \frac{d\sigma}{dzd\nu} \right)_D}$$

Hadronization is modified in the nuclear medium  
 → Probability for quark  $f$  to form hadron  $h$  changes  
 → Depends on  $A$ , hadron kinematics

Complicates interpretation of SIDIS measurements of flavor dependence if effect different for  $\pi^+$  and  $\pi^-$

→ This can be checked with measurements at  $x=0.3$  (no EMC effect)

# Parity Violating DIS



Flavor dependence of EMC effect can also be explored via parity violating DIS

$$A_{PV} = \frac{G_F Q^2}{4\sqrt{2}\pi\alpha_{em}} \left[ a_2(x) + \frac{1 - (1-y)^2}{1 + (1-y)^2} a_3(x) \right]$$

suppressed

quark weak vector couplings

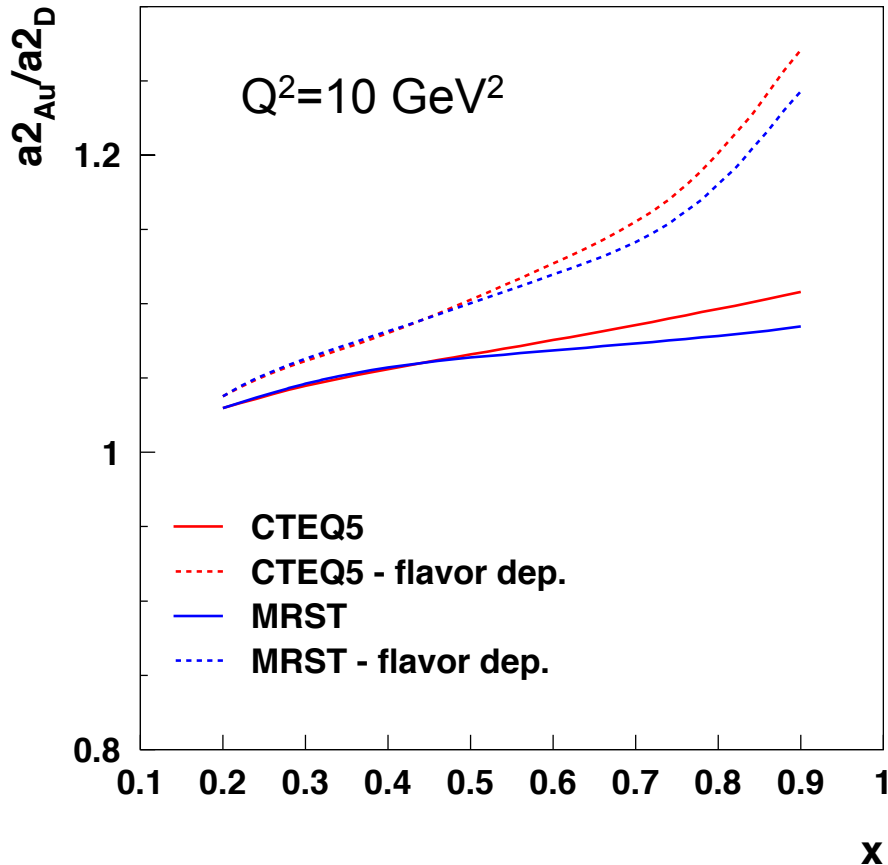
$$a_2(x) = \frac{2 \sum_q e_q g_V^q [q_A(x) + \bar{q}_A(x)]}{\sum_q e_q^2 [q_A(x) + \bar{q}_A(x)]}$$

Avoids complications due to hadronization issues

**CBT model predicts 5% effect at  $x=0.6$**

# Measuring Flavor Dependence with PVDIS

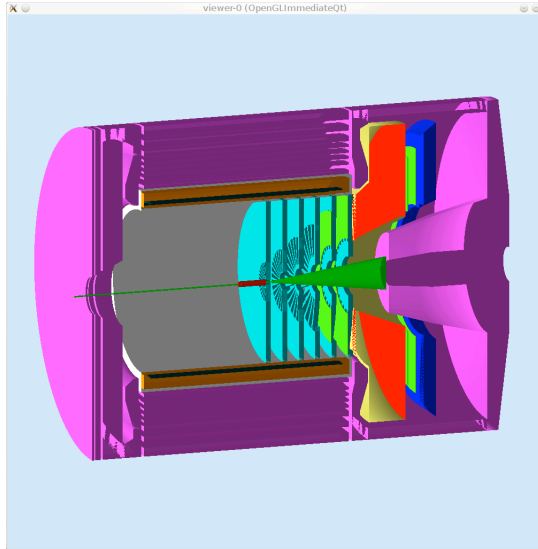
Au/<sup>2</sup>H



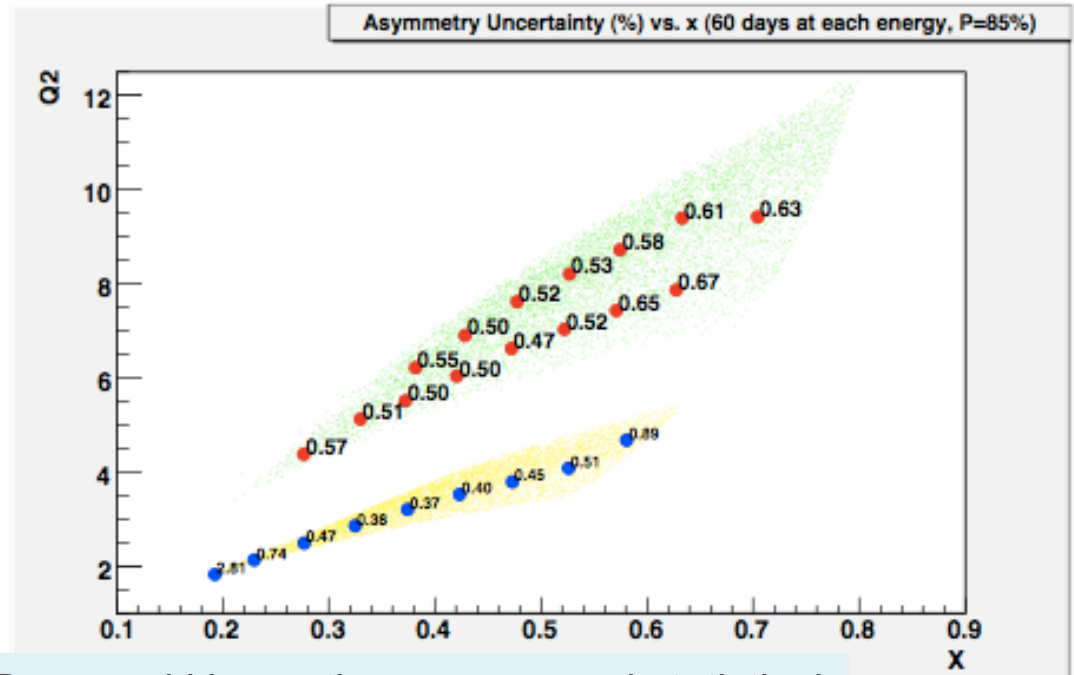
Experimentally – simpler to measure super-ratio  
→ Certain systematics are reduced (beam polarization)  
→ Less sensitivity to absolute value of weak vector couplings

Note that even the “no flavor dependence” calculation not identically 1.0  
→ Must compare experimental result to the “naïve” estimate  
→ Naïve estimate has some dependence on nucleon PDFs  
→ May be non-negligible contribution to uncertainty

# PVDIS at JLab



New solenoidal spectrometer



Proposed kinematic coverage and statistical precisions

SOLID experiment at JLab (P. Souder, spokesperson) – use PVDIS to look for physics beyond Standard Model,  $d/u$  at large  $x$   
→ awarded 169 days for H and D running  
→ no time for solid target running (flavor dependent EMC) requested yet

# Flavor Dependence with inclusive DIS

Several alternatives for accessing flavor dependence of EMC effect

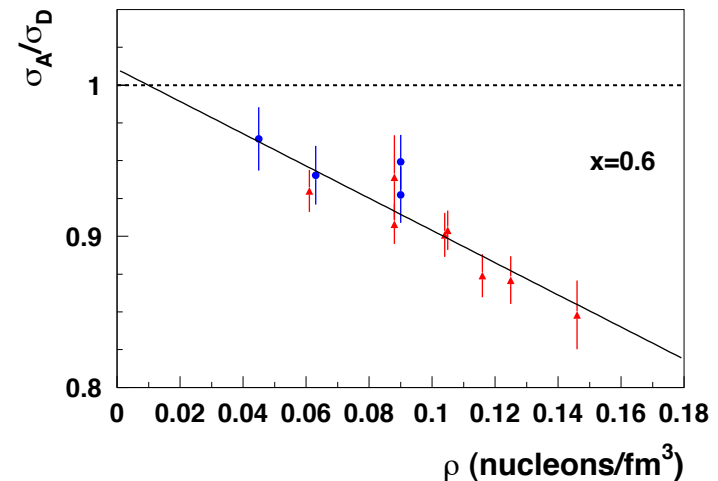
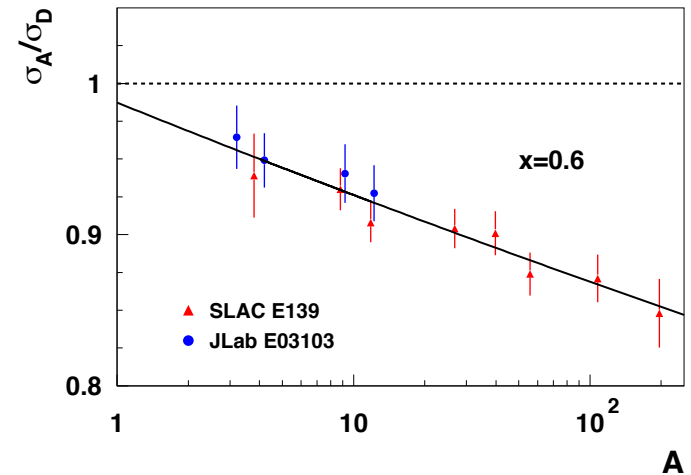
→ Pion DY @ COMPASS: sufficient statistical precision at large  $x$ ?

→ SIDIS @ JLab: hadron attenuation and factorization concerns

→ PVDIS @ JLab: SOLID experiment requires significant \$, long time scale

*Would like something “easy” that can be done on a short time scale*

Inclusive DIS on nucleus with same  $A$  and  $\rho$  but different ratio  $N/Z$



# Flavor dependence from $^{40}\text{Ca}$ and $^{48}\text{Ca}$

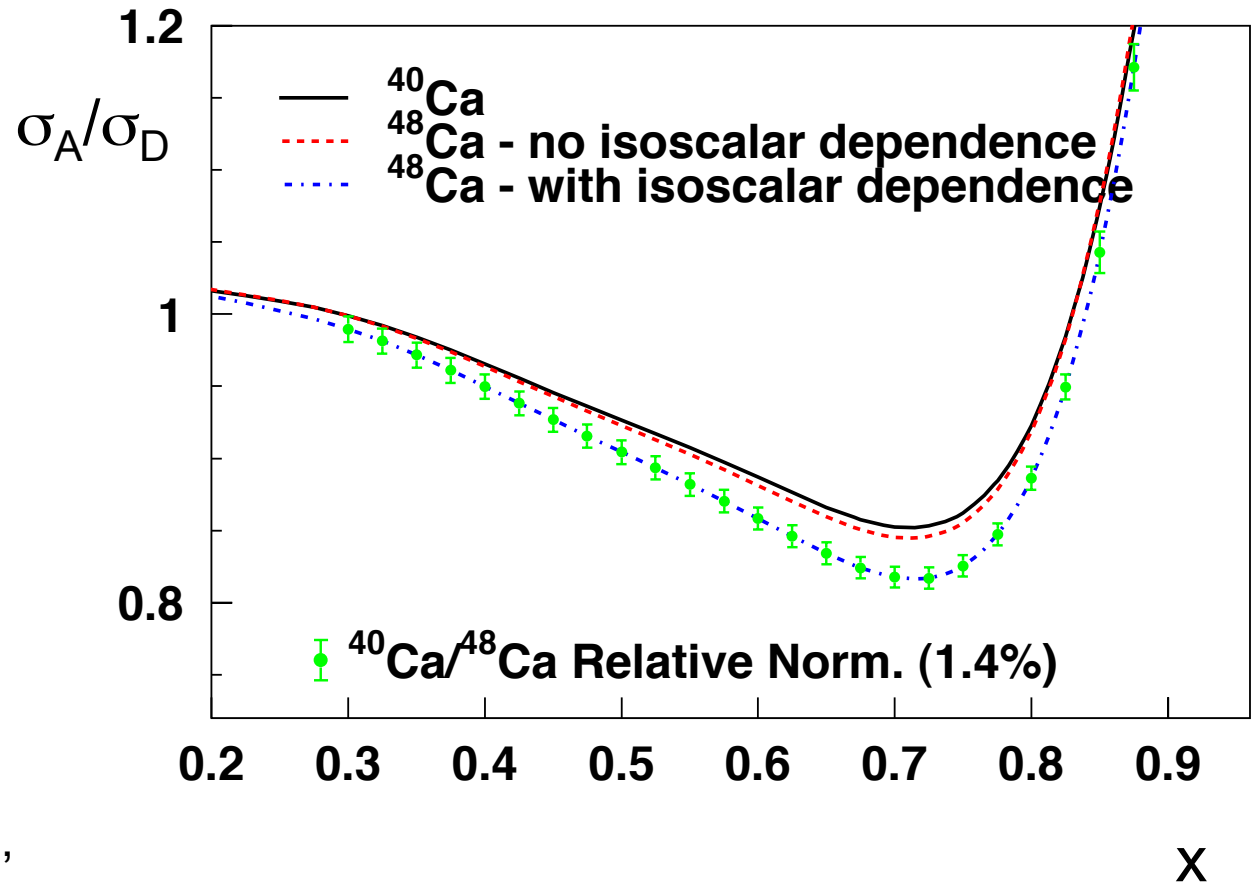
CBT model predicts a  
~3% effect for  $^{48}\text{Ca}$  at  
 $x=0.6$   
 $\rightarrow N/Z = 1.4$

Assuming no flavor  
dependence, difference  
between  $^{40}\text{Ca}$  and  $^{48}\text{Ca}$   
should be less than < 1%

Will be measured at JLab  
@ 12 GeV

E12-10-008

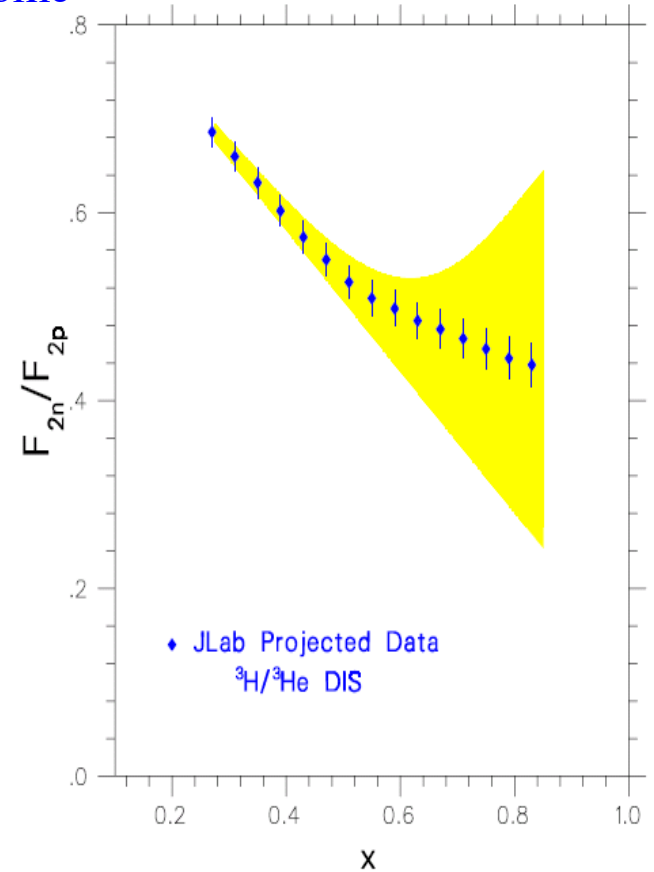
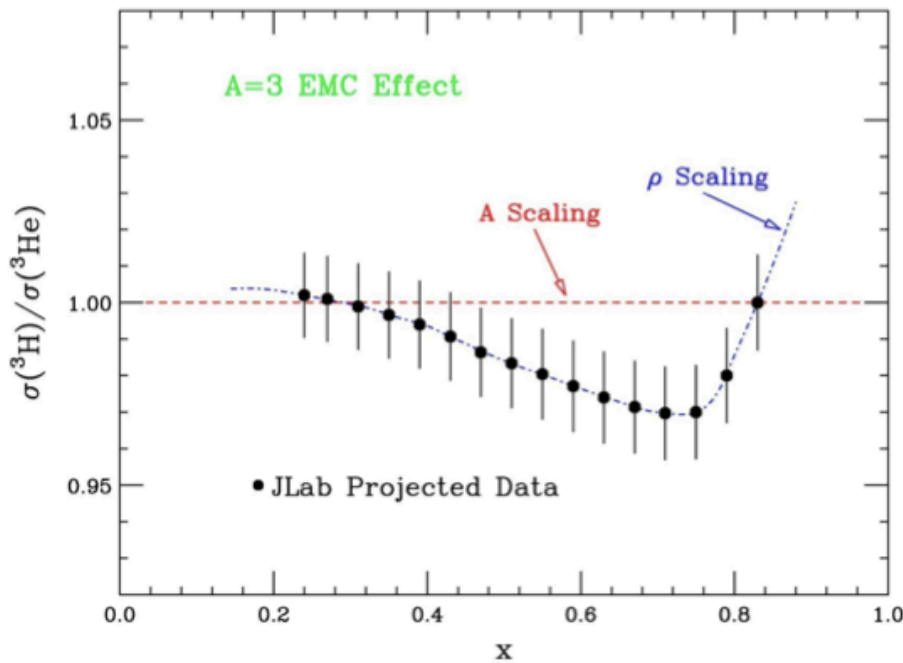
Spokespersons: Arrington,  
Gaskell, Daniel



# E12-06-118: The MARATHON experiment

Spokespeople: G. Petratos, J. Gomez, R. Holt, R. Ransome

EMC effect with A=3 mirror nuclei



→ “Free” n/p (d/u) ratios extracted using “known” corrections to difference in EMC effect in  $^3\text{He}/^3\text{H}$ ; additional flavor dependence could impact extraction



# E12-10-008 and E12-06-105

Hall C experiments will provide more inclusive data

→ E12-06-105  $x > 1$

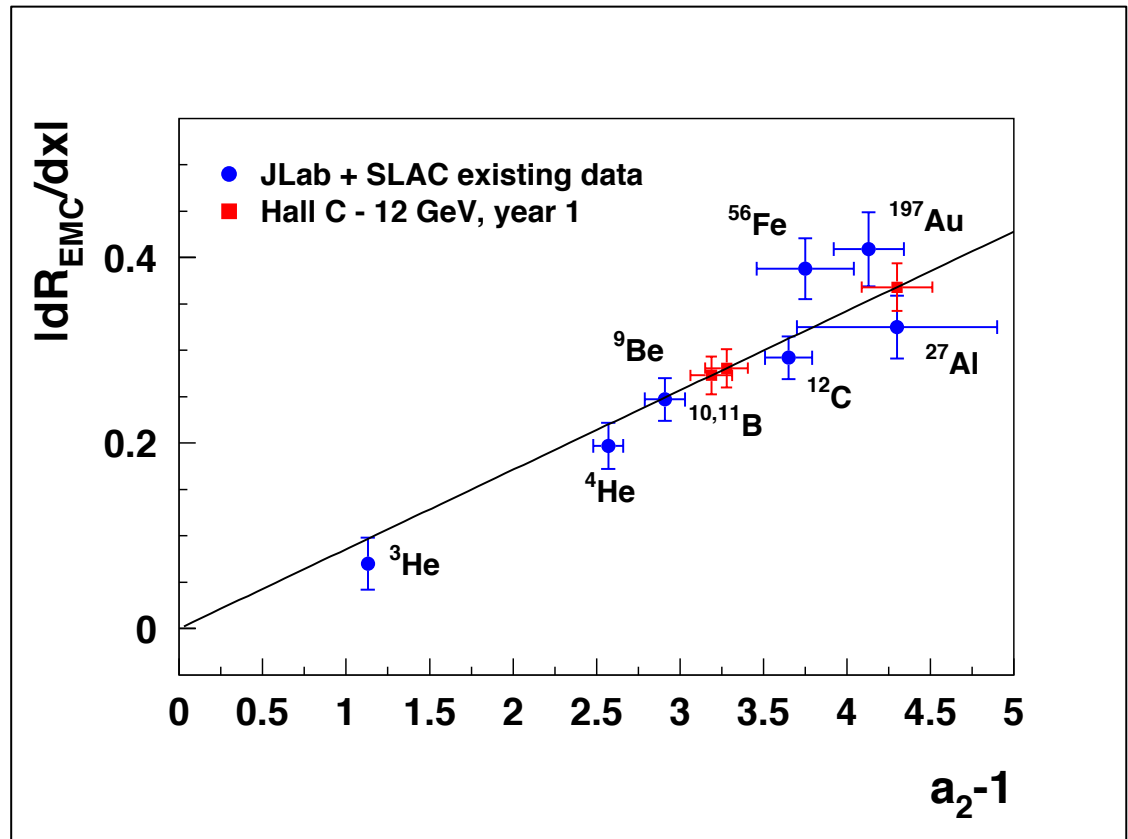
→ E12-10-008 *EMC Effect*

Will provide additional data on light and medium-heavy targets

→  $^2\text{H}$ ,  $^3\text{He}$ ,  $^4\text{He}$

→  $^6\text{Li}$ ,  $^7\text{Li}$ , Be,  $^{10}\text{B}$ ,  $^{11}\text{B}$ , C

→ Al,  $^{40}\text{Ca}$ ,  $^{48}\text{Ca}$ , Cu



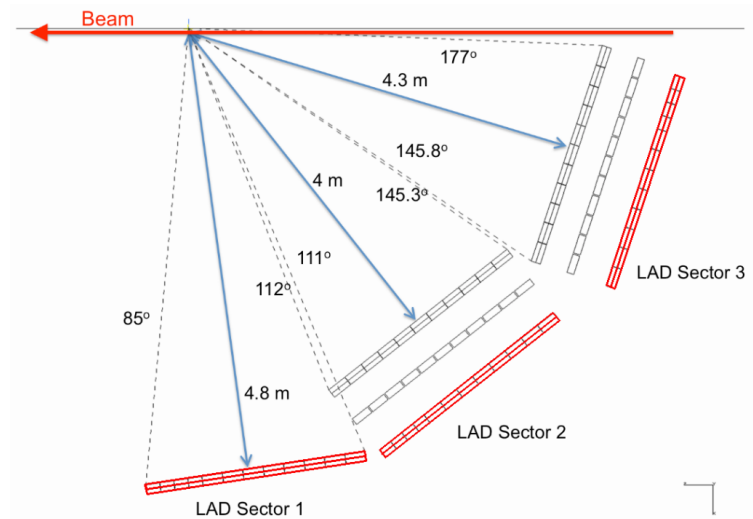
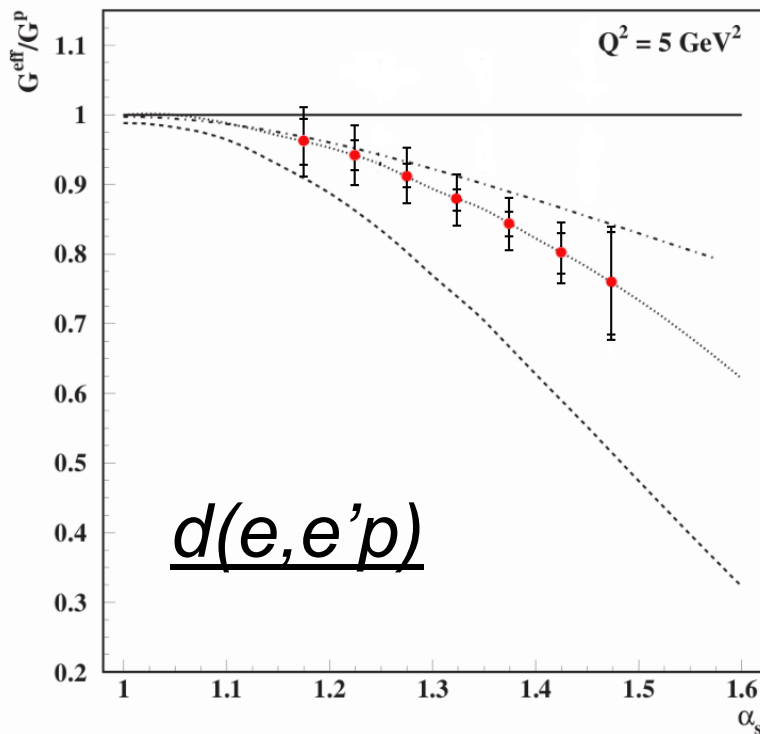
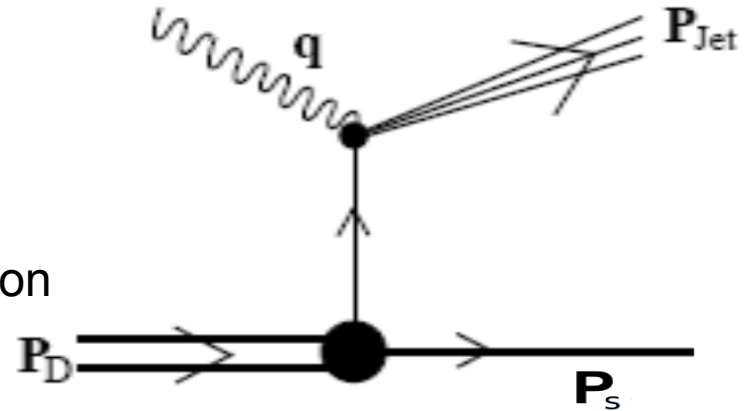
First running in Hall C after completion of 12 GeV Upgrade will include a few days for EMC/ $x > 1$  measurements on  $^{10}\text{B}$ ,  $^{11}\text{B}$ , and Al (parasitic)

# E12-11-107: In-Medium Structure Functions

Measure structure function of high momentum nucleon in deuterium by tagging the spectator

→ Final state interactions cancelled by taking double ratios

→ Requires new, large acceptance proton/neutron detector at back angles



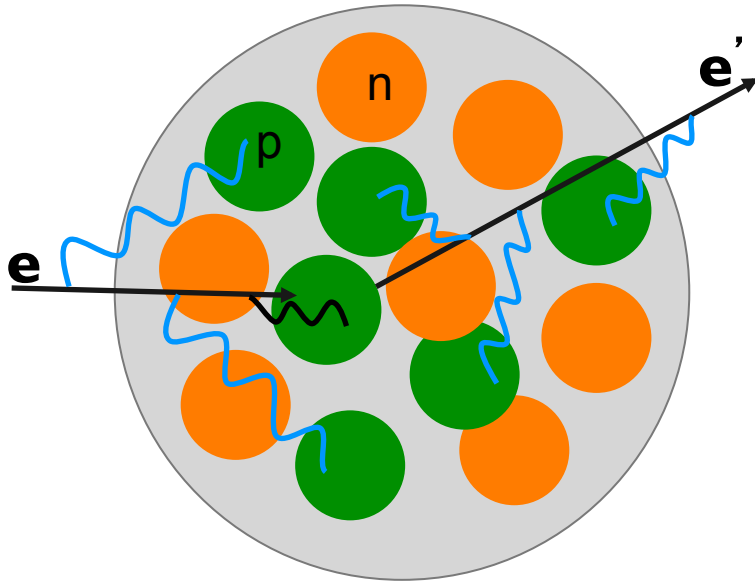
Spokespersons: O. Hen, L. Weinstein, S. Gilad, S. Wood

# Light to Heavy Nuclei

- New JLab data, new method of characterizing “size” of EMC effect gave insight into nuclear dependence of EMC effect.
  - Same dependence observed for  $A/D$  ratios at  $x > 1$
  - Correlation between EMC effect and SRCs
  - Local density vs. high virtuality  $\rightarrow$  flavor dependence?
- Some interesting effects have also been observed for heavy targets

# JLab E03103 – Heavy Targets

E03-103 also measured EMC ratios for Cu and Au – analysis at the relatively low 6 GeV beam energy complicated by **Coulomb Corrections**



Electrons scattering from nuclei can be accelerated/decelerated in the Coulomb field of the nucleus

→ This effect is NOT part of the hadronic structure of the nucleus we wish to study

→ Important to remove/correct for apparent changes in the cross section due to Coulomb effects

In a very simple picture – Coulomb field induces a change in kinematics in the reaction

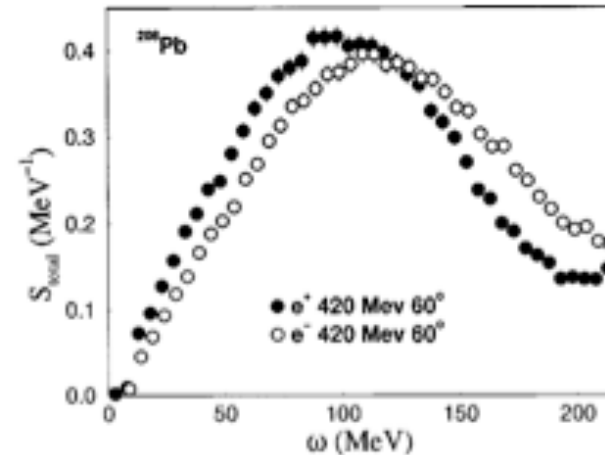
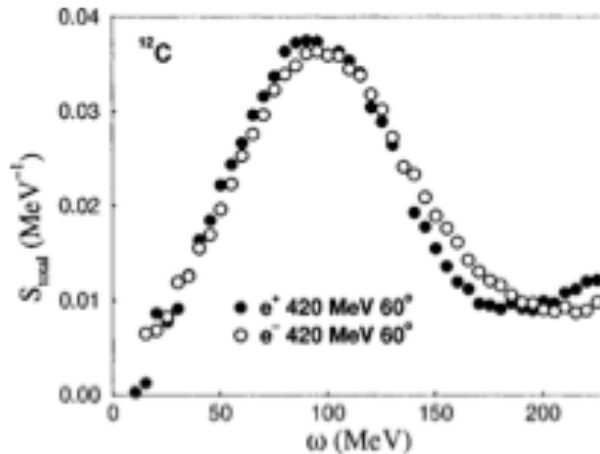
$$E_e \rightarrow E_e + V_0$$

$$E_{e'} \rightarrow E_{e'} - V_0$$

$$V_0 = 3\alpha(Z-1)/2R \quad \leftarrow \text{Electrostatic potential energy at center of nucleus}$$

# Coulomb Corrections in QE Processes

Importance of Coulomb Corrections in quasi-elastic processes well known



Gueye et al., PRC60, 044308 (1999)

Distorted Wave Born Approximation calculations are possible – but difficult to apply to experimental cross sections

→ Instead use **E**ffective **M**omentum **A**pproximation (**EMA**) tuned to agree with DWBA calculations

EMA:  $E_e \rightarrow E_e + V_0$      $E_e' \rightarrow E_e' - V_0$     with “focusing factor”  $F^2 = (1 - V_0/E)$   
 $V_0 \rightarrow (4/5)V_0$ ,  $V_0 = 3\alpha(Z-1)/2R$      $V_0 = 10$  MeV for Cu, 20 MeV for Au

[Aste et al, *Eur.Phys.J.A26:167-178,2005*, *Europhys.Lett.67:753-759,2004*]

# E03103: EMC Effect in Gold

$\sigma_A/\sigma_D$  for Gold  
A=197 Z=79

SLAC E-139

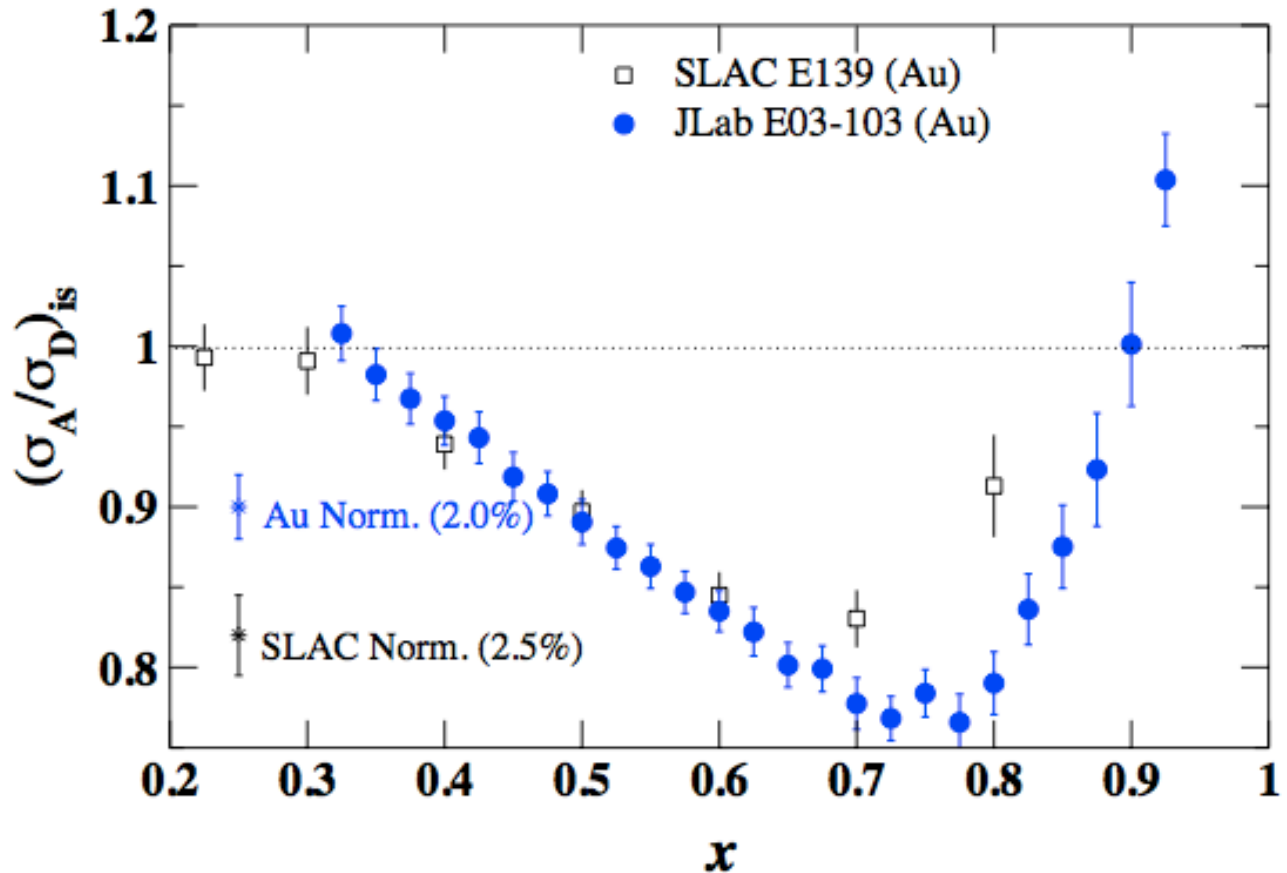
$E_e \sim 8-25$  GeV

$E_e' \sim 4-8$  GeV

JLab E03-103

$E_e \sim 6$  GeV

$E_e' \sim 1-2$  GeV



No Coulomb Corrections applied

# E03103: EMC Effect in Gold

$\sigma_A/\sigma_D$  for Gold  
A=197 Z=79

SLAC E-139

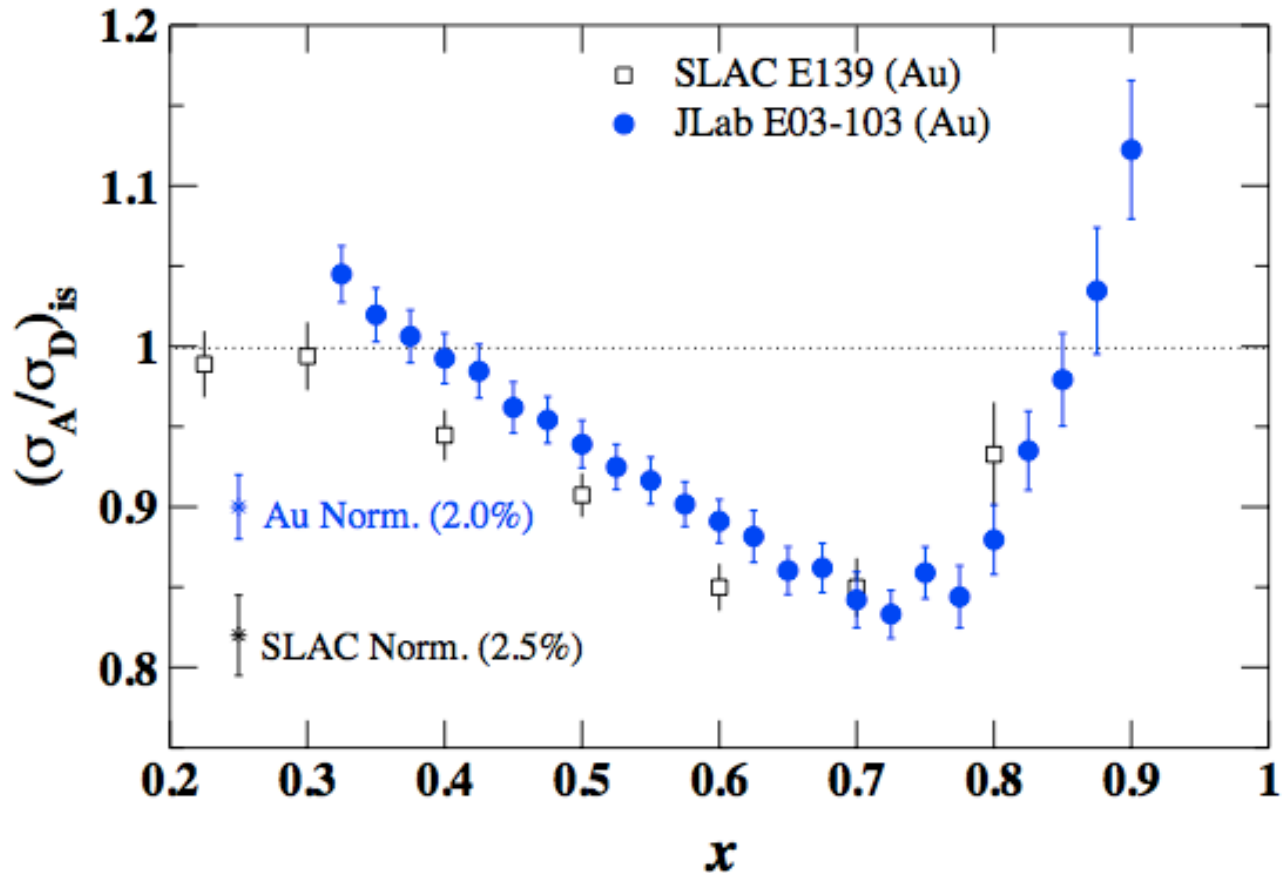
$E_e \sim 8-25$  GeV

$E_e' \sim 4-8$  GeV

JLab E03-103

$E_e \sim 6$  GeV

$E_e' \sim 1-2$  GeV



with Coulomb Corrections (both data sets)

# $R_A - R_D$

E03103 shows good agreement with E139 data for smaller  $A$   
→ agreement not as good for heavier targets. Why?

$$\frac{d\sigma}{d\Omega dE'} = \frac{4\alpha^2(E')^2}{Q^4 v} \left[ F_2(v, Q^2) \cos^2 \frac{\theta}{2} + \frac{2}{Mv} F_1(v, Q^2) \sin^2 \frac{\theta}{2} \right]$$

$$F_2(x) = \sum_i e_i^2 x q_i(x) \quad \leftarrow \text{Quark distribution functions}$$

$$\frac{d\sigma}{d\Omega dE'} = \Gamma \left[ \sigma_T(v, Q^2) + \varepsilon \sigma_L(v, Q^2) \right] \quad F_1 \propto \sigma_T \quad F_2 \text{ linear combination of } \sigma_T \text{ and } \sigma_L$$

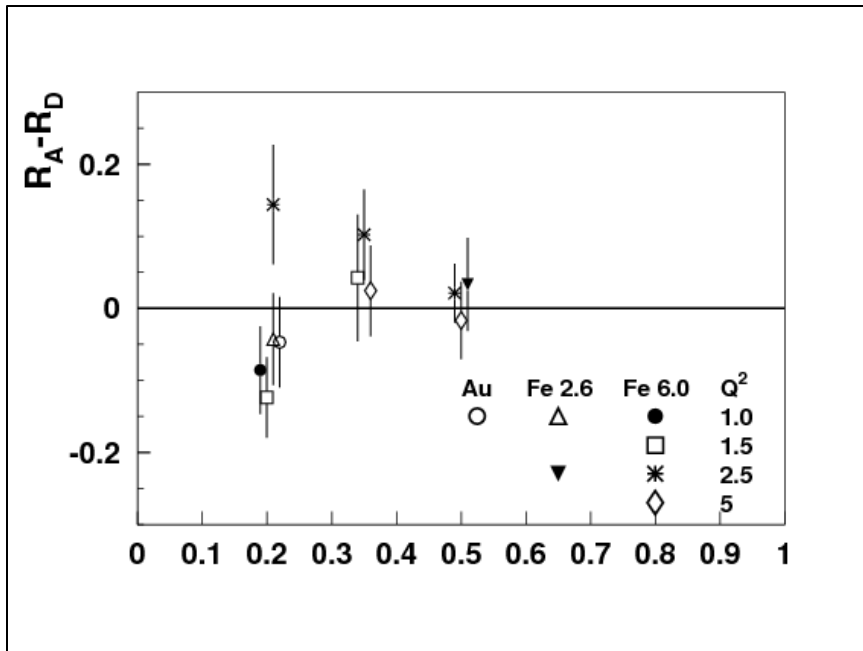
Measurements of EMC effect often assume  $\sigma_A/\sigma_D = F_2^A/F_2^D$   
→ this is true if  $R = \sigma_L/\sigma_T$  is the same for A and D

E139 data mostly at large  $\varepsilon$  – JLab data at small  $\varepsilon$  → if  $R_A \neq R_D$ , this might explain the difference

→ Motivated us to re-examine earlier experiments that measured nuclear dependence of  $R$



# SLAC E140: $R_A - R_D$



E140 measured  $\epsilon$  dependence of cross section ratios  $\sigma_A/\sigma_D$  for

$x=0.2, 0.35, 0.5$

$Q^2 = 1.0, 1.5, 2.5, 5.0 \text{ GeV}^2$

Iron and Gold targets

$R_A - R_D$  consistent with zero within errors

[E140 Phys. Rev. D 49 5641 (1993)]

**No Coulomb corrections were applied**

Large  $\epsilon$  data:  $E_e \sim 6-15 \text{ GeV}$   $E_e' \sim 3.6-8 \text{ GeV}$

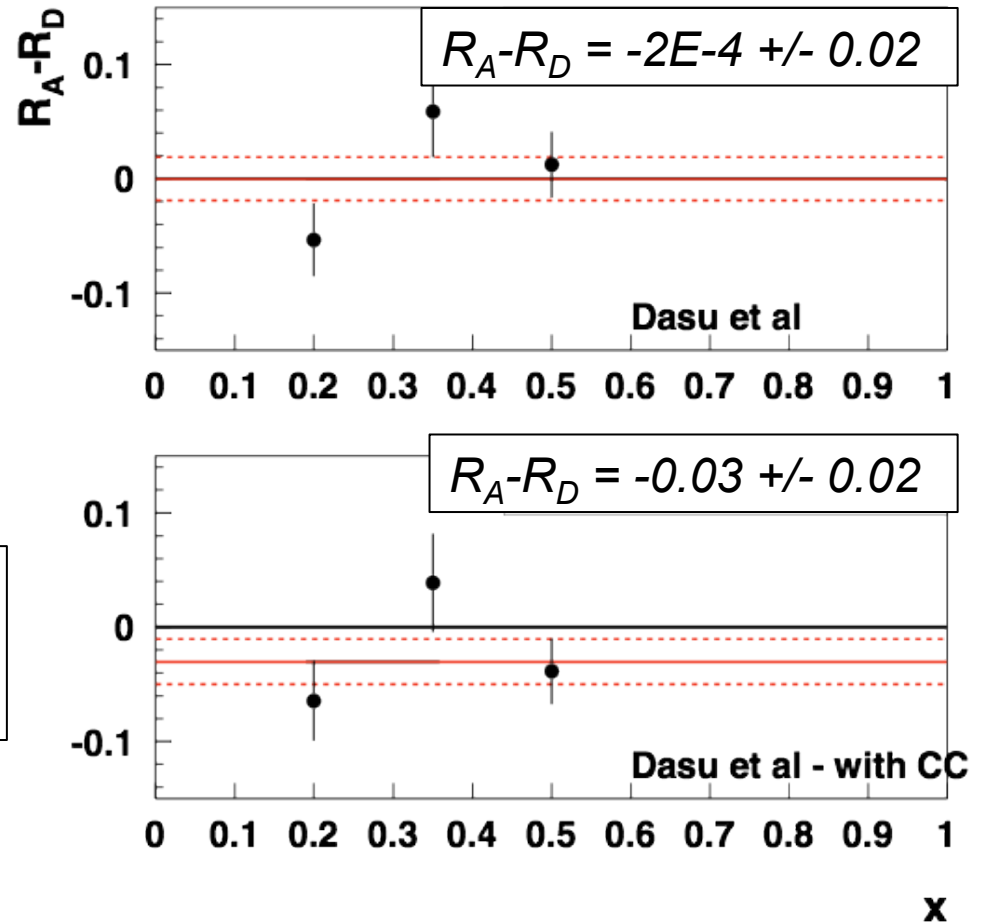
Low  $\epsilon$  data:  $E_e \sim 3.7-10 \text{ GeV}$   $E_e' \sim 1-2.6 \text{ GeV}$

# $R_A - R_D$ : E140 Re-analysis

Re-analyzed E140 data using Effective Momentum Approximation for published “Born”-level cross sections

→ Total consistency requires application to radiative corrections model as well

Including Coulomb Corrections yields result  $1.5 \sigma$  from zero when averaged over  $x$



# $R_A - R_D$ at $x=0.5$

Interesting result from E140 re-analysis motivated more detailed study

→  $x=0.5$ ,  $Q^2=5 \text{ GeV}^2$

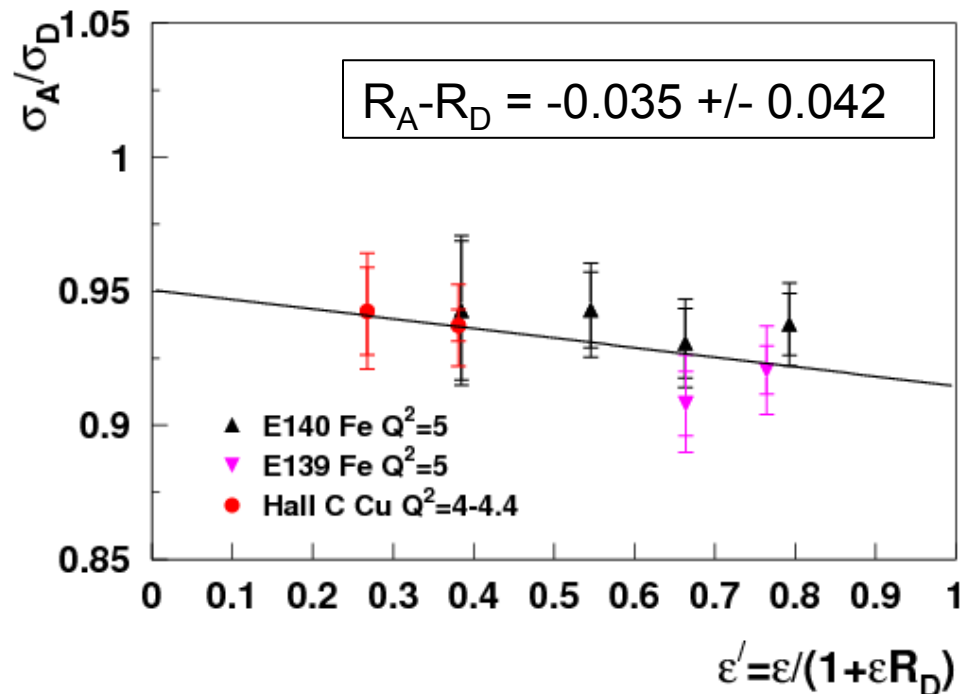
→ Include E139 Fe data

→ Include JLab data

Cu,  $Q^2=4-4.4 \text{ GeV}^2$

Normalization uncertainties between experiments treated as extra point-to-point errors

No Coulomb Corrections → combined analysis still yields  $R_A - R_D \sim 0$



*No Coulomb Corrections*

# $R_A - R_D$ at $x=0.5$

Interesting result from E140 re-analysis motivated more detailed study

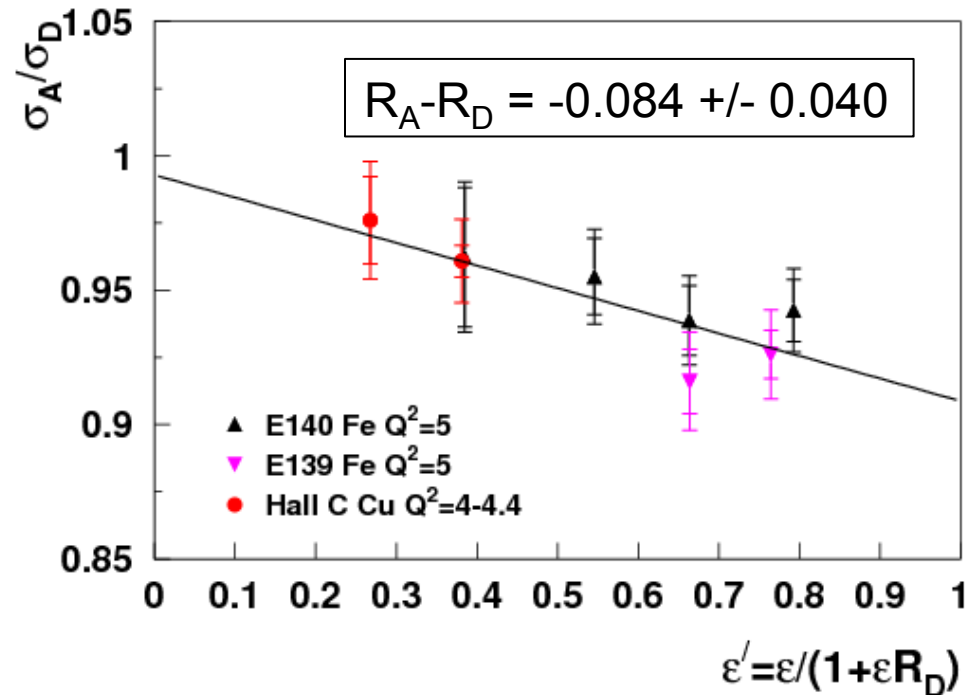
→  $x=0.5$ ,  $Q^2=5 \text{ GeV}^2$

→ Include E139 Fe data

→ Include JLab data

Cu,  $Q^2=4-4.4 \text{ GeV}^2$

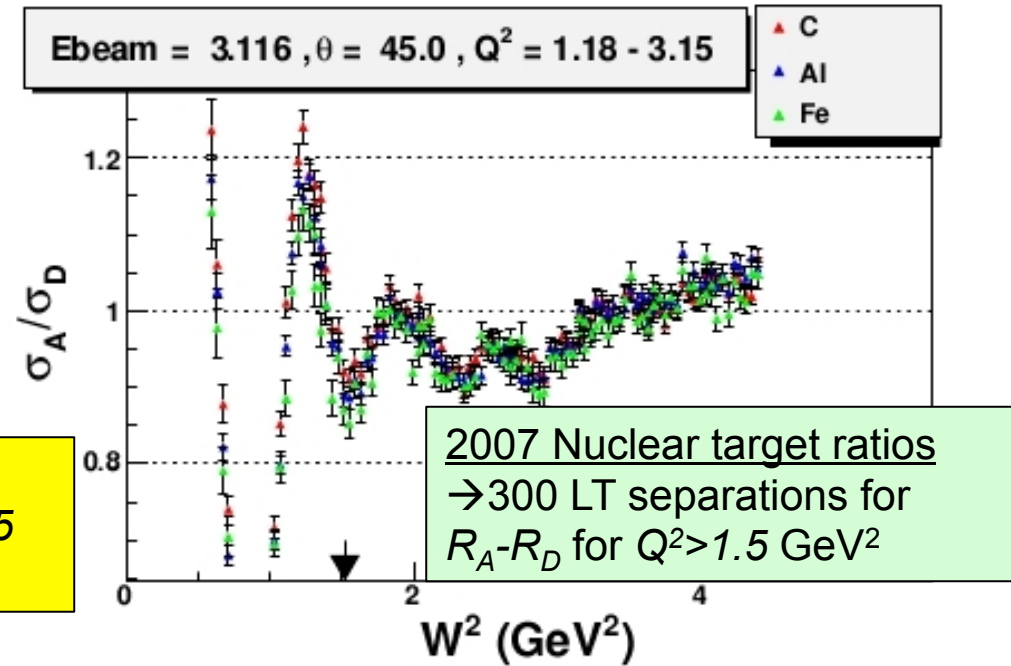
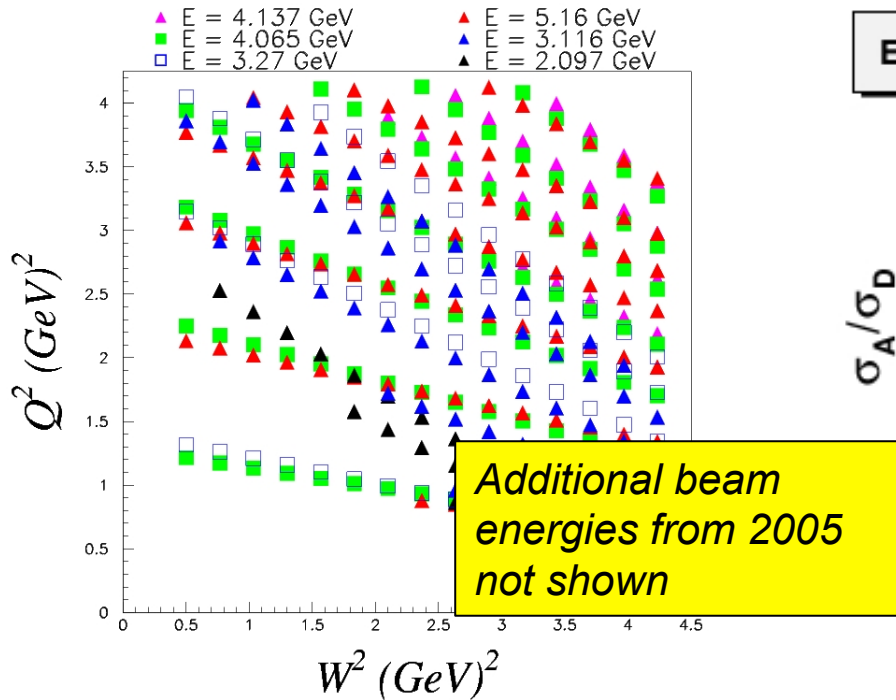
Normalization uncertainties between experiments treated as extra point-to-point errors



*with Coulomb Corrections*

Application of Coulomb Corrections →  $R_A - R_D$   $2 \sigma$  from zero

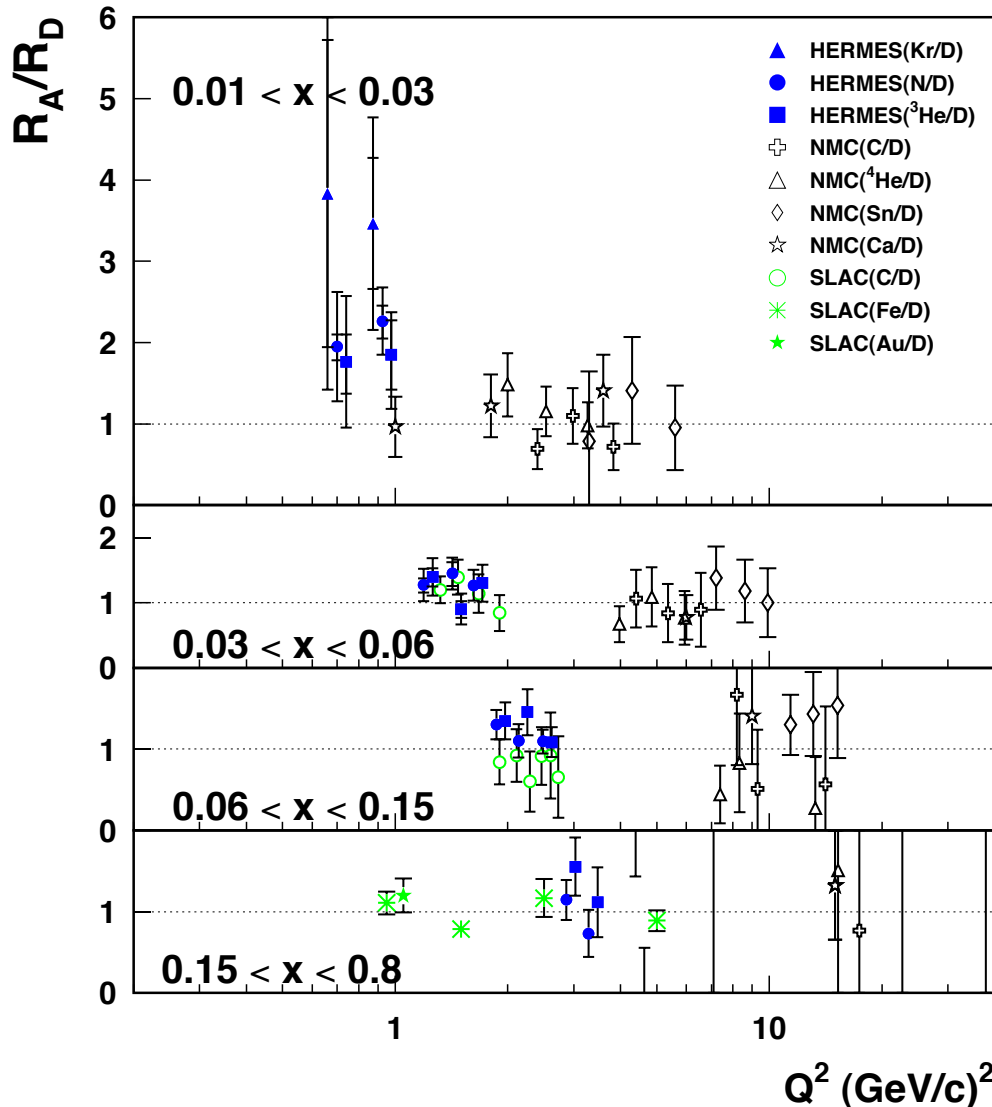
# JLab Hall C E02-109/E04-001/E06-009



- Precision extraction of separated structure functions on D, Al, C, Fe/Cu
- Search for nuclear effects in  $F_L$ , R
- Neutron and p-n moment extractions (compare to lattice calculations)
- Allow study of quark-hadron duality for neutron, nuclei separated structure functions

**$F_2$ ,  $F_L$ , R on Deuterium and heavier targets**

# World Data on $R_A/R_D$



SLAC E140: *PRD* 49, 5641 (1994)

$R_A-R_D$  for Fe, Au

Only true Rosenbluth separated data

NMC:

*Phys. Lett. B* 294, 120 (1992)

$R_{Ca}-R_C$

*Nucl. Phys. B* 481, 23 (1996)

$R_{Sn}-R_C$

Multiple beam energies,  $R_A-R_C$   
extracted using  $Q^2$  dep. fit at fixed  $x$

HERMES:

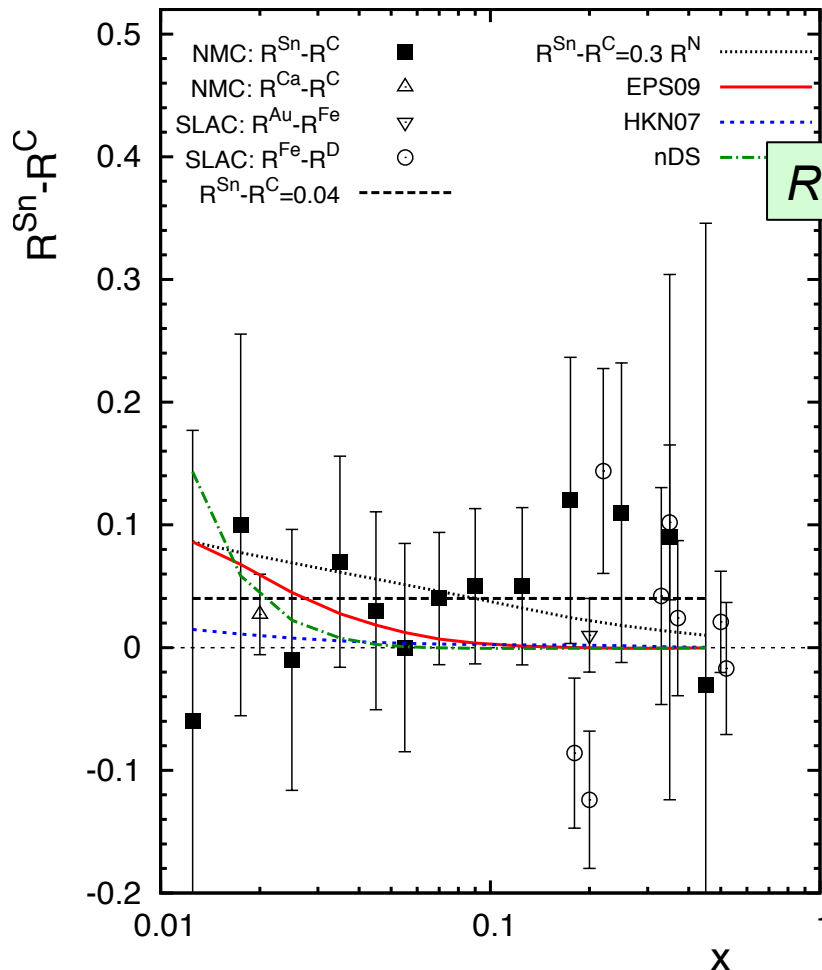
*Phys Lett. B* 567, 339 (2003)

$R_A/R_D$  for Kr, N, <sup>3</sup>He

Fit  $\epsilon$  dependence at fixed  $x$  for  
single beam energy (changing  $Q^2$ )

# Other Hints of non-zero $R_A - R_D$

NMC results for  $R_{Sn} - R_C$  systematically larger than zero



$$R_{Sn} - R_C = 0.040 \pm 0.026 \text{ (stat)} \pm 0.020 \text{ (sys)}$$

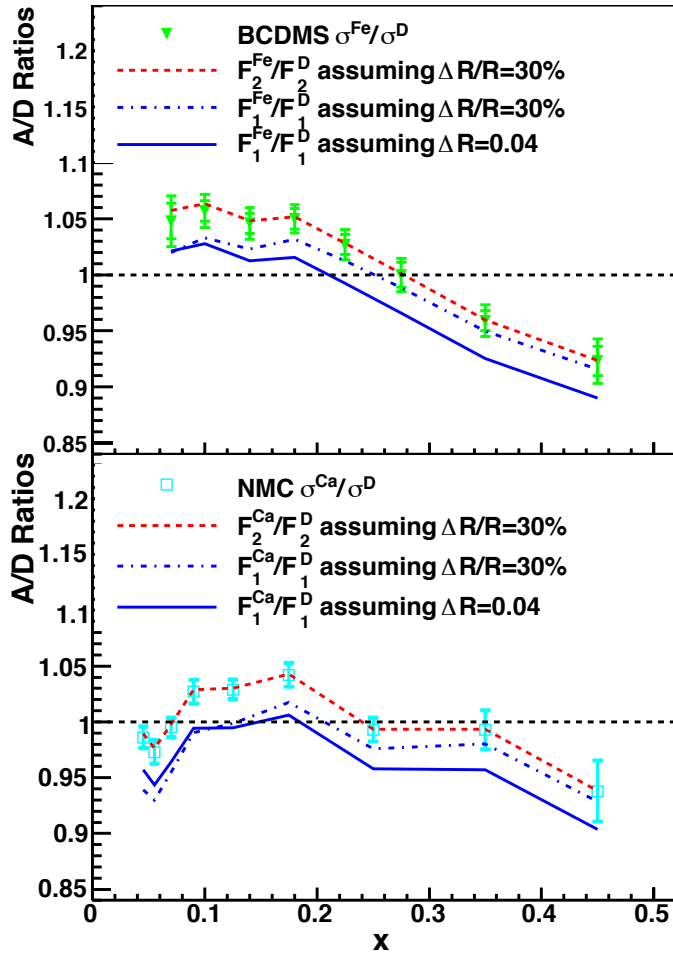
→ Averaged over  $x=0.0125 - 0.45$

→  $\langle Q^2 \rangle = 10 \text{ GeV}^2$

What are the consequences for A/D ratios for  $F_1$  and  $F_2$  if this is true?

*V. Guzey et al, PRC 86 045201 (2012)*

# Consequences of $R_A - R_D > 0$



$$\frac{\sigma_A}{\sigma_D} = \frac{F_1^A(x)}{F_1^D(x)} \left[ 1 + \frac{\epsilon(R_A - R_D)}{1 + \epsilon R_D} \right]$$

$F_1$  ratio purely transverse

Anti-shadowing disappears for  $F_1$  ratio,  
remains for  $F_2$

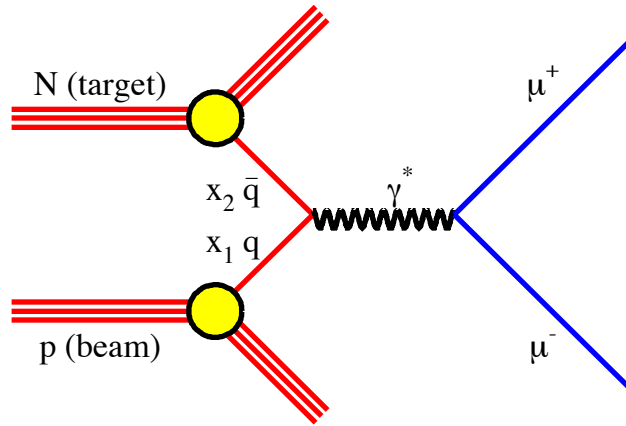
Anti-shadowing from longitudinal photons?

*More discussion in Thia Keppel's talk next week*

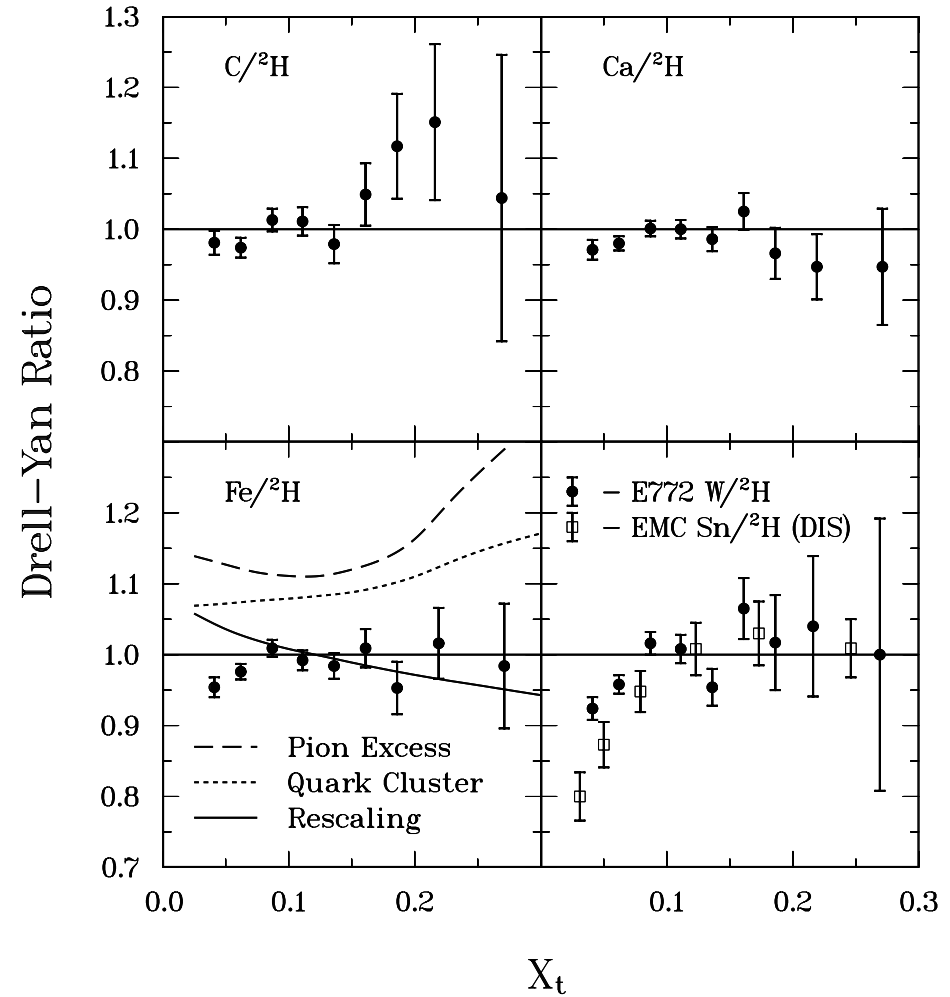
V. Guzey et al, PRC 86 045201 (2012)



# A Dependence of Anti-quark Distributions

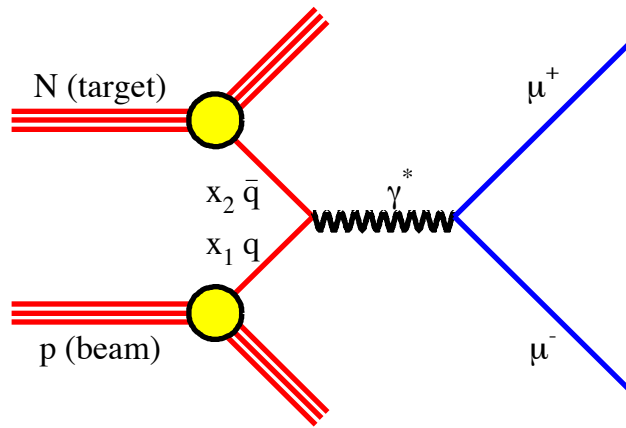


- Drell-Yan process sensitive to anti-quark distributions in the target
- E772 measured no  $A$  dependence over limited  $x$  range, with limited precision
- E906 will measure up to  $x=0.4$

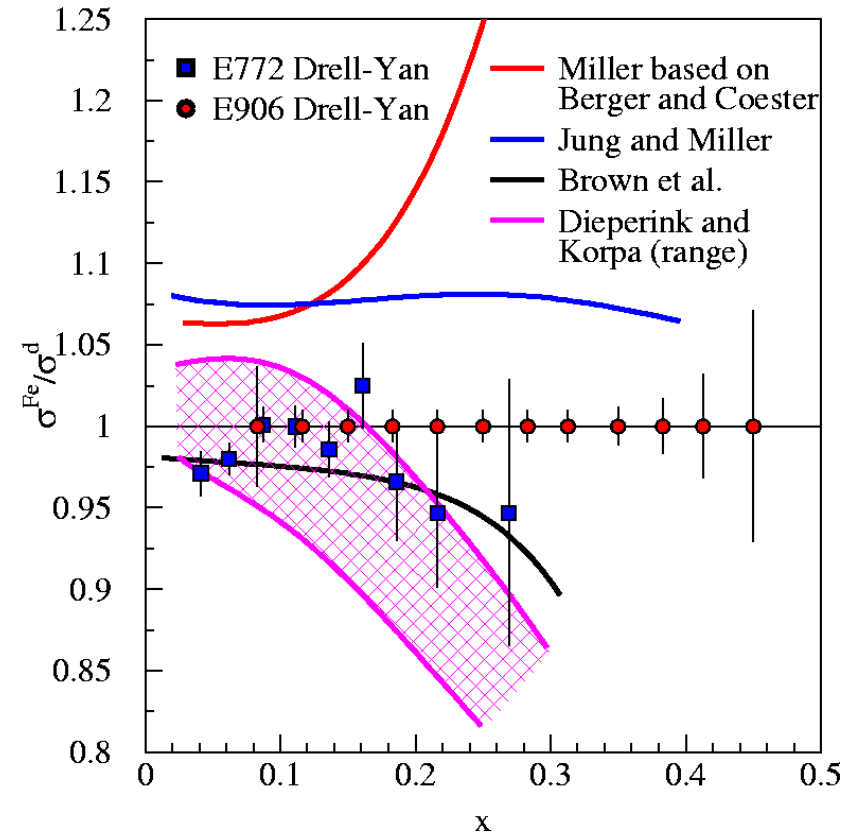


*D.M. Alde et al., PRL64: 2479 (1990)*

# A Dependence of Anti-quark Distributions



- Drell-Yan process sensitive to anti-quark distributions in the target
- E772 measured no  $A$  dependence over limited  $x$  range, with limited precision
- E906 will measure up to  $x=0.4$



E906 underway ...

# Nuclear Dependence of $R$

- Conventional wisdom was that there was little or no difference between  $R$  in heavy nuclei and free nucleon
- Recent JLab data suggests  $R_A - R_D < 0$  at large  $x$ 
  - Alternatively, Coulomb Corrections are not under control
  - Better calculations and/or experimental tests needed
- Re-examination of high energy NMC data suggests  $R_A - R_D > 0$ 
  - How can this be consistent with JLab + SLAC data?
  - $Q^2$  dependent? Problems with either data set?
- More data is needed – a systematic study over large range of  $Q^2$  and  $x$

# Summary

- The EMC effect has been with us for 30 years and motivated intense experimental (and theoretical) study
- Amazingly, it seems there is still much to learn
  - What is the link between SRCs and the EMC effect?
  - Does the EMC effect depend on quark flavor?
  - Does  $\sigma_A/\sigma_D = F_2^A/F_2^D$  for all  $x$  and  $Q^2$
- Many of these questions will be addressed at JLab after the 12 GeV upgrade
- Issues I did not discuss
  - Polarized EMC effect
  - Low  $x$  measurements  $\rightarrow$  EIC
  - Several other processes that aim to quantify the modification of nucleons in the nucleus