INT13-3 2 Oct. 2013

Collective mass parameters in Skyrme EDF

Niigata Univ. Kenichi Yoshida Key ingredients for spontaneous fission in nuclear EDF

Collective Hamiltonian describing the fission dynamics

$$\mathcal{H} = \frac{1}{2} \sum_{kl} M_{kl}(s) \dot{s}_k \dot{s}_l + V_{\text{coll}}(s)$$

Dynamical variables $s = (s_1, s_2, \dots) = (\beta, \gamma, \beta_3, \dots)$

Collective potential $V_{coll}(s)$

Inertia functions for the collective mode $M_{kl}(s)$

Collective potential

✓ potential energy surface in multi-dimensional constrained EDF method

Gogny

relativistic





Skyrme



Lu, Zhao, Zhou, PRC85(2012)011301(R)

Staszczak, Baran, Dobaczewski, Nazarewicz, PRC85(2012)024314

✓ continuous potential energy surface toward a class-3 PES (talk by Dubray)

Microscopic collective-mass parameter

• GCM-GOA mass

exact calculation (talk by Robledo)

• ATDHFB mass

$$\mathcal{M}_{ij} = \frac{i}{2\dot{q}_i \dot{q}_j} \operatorname{Tr}(F^{i*}Z^j - F^i Z^{j*}).$$
(34)

time-odd fields neglected

Baran, Sheikh, Dobaczewski, Nazarewicz, Staszczak, PRC85(2012)024314

Cranking approximation

$$\mathcal{M}_{ij}^{C} = \frac{1}{2\dot{q}_{i}\dot{q}_{j}} \sum_{\alpha\beta} \frac{\left(F_{\alpha\beta}^{i*}F_{\alpha\beta}^{j} + F_{\alpha\beta}^{i}F_{\alpha\beta}^{j*}\right)}{E_{\alpha} + E_{\beta}}$$

Perturbative cranking

$$\mathcal{M}^{C^{p}} = \frac{1}{4} [M^{(1)}]^{-1} M^{(3)} [M^{(1)}]^{-1}$$
$$M_{ij}^{(K)} = \sum_{\alpha\beta} \frac{\langle 0|\hat{Q}_{i}|\alpha\beta\rangle\langle\alpha\beta|\hat{Q}_{j}^{\dagger}|0\rangle}{(E_{\alpha} + E_{\beta})^{K}}$$



Local QRPA method

Hinohara, Sato, Nakatsukasa, Matsuo, Matsuyanagi, PRC82(2010)064313

CHFB eq. $\delta \langle \phi(s) | \hat{H}_{CHFB} | \phi(s) \rangle = 0$ $\hat{H}_{CHFB} = \hat{H}_{HFB} - \sum_{i} \mu_{i} \hat{s}_{i}$

 $|\phi(s)\rangle$ $V_{\rm coll}(s) = \langle \phi(s) | \hat{H}_{\rm HFB} | \phi(s) \rangle$

Local harmonic approximation at each state $|\phi(s)
angle$

$$\delta\langle\phi(s)|[\hat{H}_{\rm CHFB},\hat{Q}_{\mu}(s)] - \frac{1}{i}\hat{P}_{\mu}(s)|\phi(s)\rangle = 0$$

$$\delta\langle\phi(s)|[\hat{H}_{\rm CHFB},\frac{1}{i}\hat{P}_{\mu}(s)] - C_{\mu}(s)\hat{Q}_{\mu}(s)|\phi(s)\rangle = 0$$

collective mode generated self-consistently
 time-odd effects taken into account

Collective kinetic energy

$$\mathcal{T} = \frac{1}{2} \sum_{i} (\dot{q}_{i})^{2}$$
$$= \frac{1}{2} \sum_{i} \sum_{kl} \frac{\partial q_{i}}{\partial s_{k}} \frac{\partial q_{i}}{\partial s_{l}} \dot{s}_{k} \dot{s}_{l} \equiv \frac{1}{2} \sum_{kl} M_{kl}(s) \dot{s}_{k} \dot{s}_{l}$$

Note that the infinitesimal displacement of the collective coordinates brings about a corresponding change;

$$ds_k = \sum_i \frac{\partial s_k}{\partial q_i} dq_i$$

derivative w.r.t. the collective coordinate q_i

$$\frac{\partial s_k}{\partial q_i} = \frac{\partial}{\partial q_i} \langle \phi(s) | \hat{s}_k | \phi(s) \rangle$$
$$= \langle \phi(s) | \left[\hat{s}_k, \frac{1}{i} \hat{P}_i(s) \right] | \phi(s) \rangle$$

not in need of numerical derivative

Numerical implementation with use of Skyrme EDF KY, Hinohara, PRC83(2011)061302(R)

Skyrme + pairing EDF: $E[arrho(m{r}), \widetilde{arrho}(m{r})]$

CHFB eq. in cylindrical coordinates assuming the axial, reflection symmetries

$$\begin{pmatrix} h(\boldsymbol{r}\sigma) - \lambda & \tilde{h}(\boldsymbol{r}\sigma) \\ \tilde{h}(\boldsymbol{r}\sigma) & -h(\boldsymbol{r}\sigma) + \lambda \end{pmatrix} \begin{pmatrix} \varphi_{1,\alpha}(\boldsymbol{r}\sigma) \\ \varphi_{2,\alpha}(\boldsymbol{r}\sigma) \end{pmatrix} = E_{\alpha} \begin{pmatrix} \varphi_{1,\alpha}(\boldsymbol{r}\sigma) \\ \varphi_{2,\alpha}(\boldsymbol{r}\sigma) \end{pmatrix}$$

Mean field (ph) and pair field (pp): β constraint

$$h = \frac{\delta E}{\delta \varrho} - \mu q_{20}, \qquad \tilde{h} = \frac{\delta E}{\delta \tilde{\varrho}}$$

LQRPA eq. in the matrix form (P-Q representation)

$$(A - B)Q_{\mu} = \frac{1}{i}P_{\mu}$$
$$(A + B)\frac{1}{i}P_{\mu} = \omega_{\mu}^{2}Q_{\mu}$$

$$(A+B)(A-B)Q_{\mu} = \omega_{\mu}^2 Q_{\mu}$$

applicable to the situation where the eigen-frequencies of the local normal modes are imaginary

 $\boldsymbol{r} = (
ho, z, \phi)$

Quadrupole collective mass $M_{\beta\beta}(\beta)$ in the LQRPA

• choose "the most collective mode" out of numerous eigenmodes

practically, the mode possessing the smallest quadrupole mass in the low-frequency region of $\omega_i^2 < 15 \text{ MeV}^2$

• calculate the quadrupole collective mass

$$M_{Q_{20}Q_{20}}(\beta) = \left(\frac{dq_i}{dQ_{20}}\right)^2 \qquad \text{Note:} \\ = \left|\langle\phi(\beta)|[\hat{Q}_{20}, \frac{1}{i}\hat{P}_i(\beta)|\phi(\beta)\rangle\right|^{-2} \qquad \eta = \sqrt{\frac{\pi}{5}}\frac{1}{\langle r^2\rangle} \\ = \left[\frac{2}{i}\sum_{\alpha\alpha'}q_{20,\alpha\alpha'}P_{i,\alpha\alpha'}(\beta)\right]^{-2} \\ = \left[\sum_{\alpha\alpha'}B_{\alpha\alpha'}(\beta)\right]^{-2}$$

 $B_{\alpha\alpha'}(\beta)$ Microscopic structure of the collective mass

CHFB calculation: MPI parallelization

use of N processors



QRPA calculation on parallel computer: MPI and BLACS

QRPA: use of N processors

matrix elements of the QRPA eq.: $A_{lphaeta\gamma\delta} \;\; B_{lphaeta\gamma\delta}$

Scalapack

2D-block cyclic distribution for load balancing

function: indxl2g for distribution subroutine: pdsyev for diagonalization

Ex. dim. : 50,000 for K=0 in ²⁴⁰Pu

w/ 512 cores

1024 cores

matrix element: 14,000 secs diagonalization: 1,400 secs inversion : 21 secs 7,000 secs 740 secs 15 secs 2048 cores

3,600 secs 680 secs 13 secs

Numerical results for ²⁴⁰Pu: Potential energy

SkM* +

pairing EDF

Yamagami, Shimizu, Nakatsukasa, PRC80(2009)064301

adjusted to pairing properties of welldeformed nuclei



Quadrupole mass parameter in ²⁴⁰Pu

strong dependence on the shell structure/ configuration

curvature of the collective potential $C_i(\beta) = \omega_i^2(\beta)$







Microscopic structure of collective mass in ²⁴⁰Pu



Enhanced collective mass



LQRPA and ATDHFB-Cranking masses in ²⁵⁶Fm

SkM* + mixed-type pairing



ATDHFB:

Baran, Sheikh, Dobaczewski, Nazarewicz, Staszczak, PRC85(2012)024314

using HFODD



Summary

Collective mass parameters for the quadrupole vib. calculated by use of the LQPRA method w/Skyrme EDF

Time-odd components are included

strongly dependent on the structure of the vacua

 $0\hbar\omega$ 2qp excitation sensitive to the shell structure more than twice at most as large as the perturbative-cranking mass

Perspective

Deeper understanding of the mass parameter microscopically in terms of the quasiparticle excitation

Benchmark is needed

among the LQRPA, ATDHFB(-Cranking), GCM-GOA masses Quadrupole mass parameters on the symmetric path way

Application to spontaneous fission dynamics LQRPA on top of the triaxial and octupole deformed states

3D-QRPA code is needed

HFODD + parallelized m-FAM may be a practical way for it(?) m-FAM: Avogadro, Nakatsukasa, PRC87(2013)014331