

The strange way from nuclei to neutron stars: a Quantum Monte Carlo approach

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Trento, Italy



Main collaborators:

- ★ F. Pederiva (Trento, Italy)
- ☆ S. Gandolfi (LANL, US-NM)

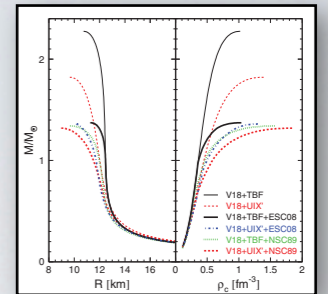
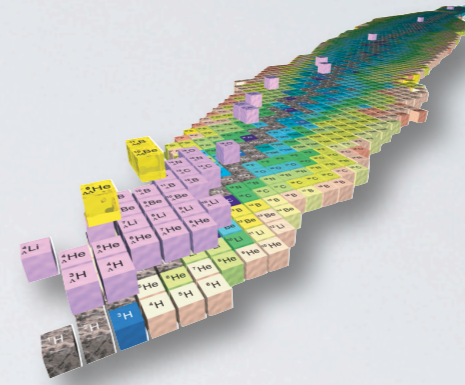


INT, Seattle - July 29, 2013

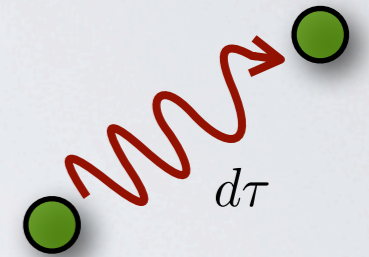
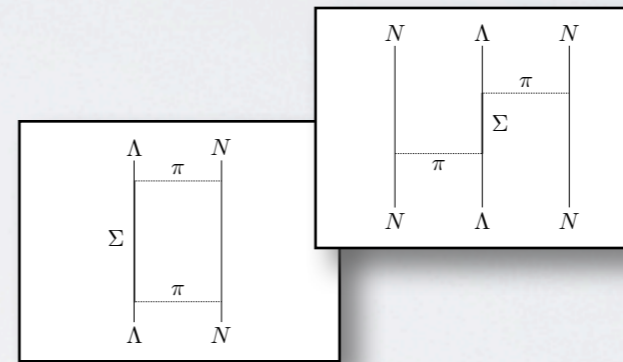


Outline

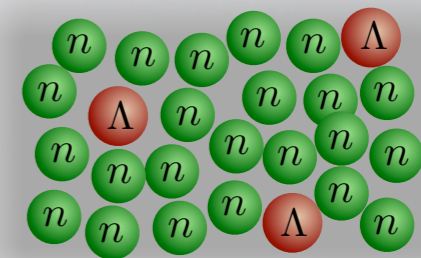
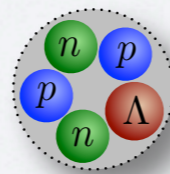
- ✓ Strangeness in nuclear physics: why?
 - terrestrial experiments: hypernuclei
 - neutron stars: theory vs. observations



- ✓ The strange AFDMC project
 - the idea
 - the hyperon-nucleon interaction
 - the AFDMC code

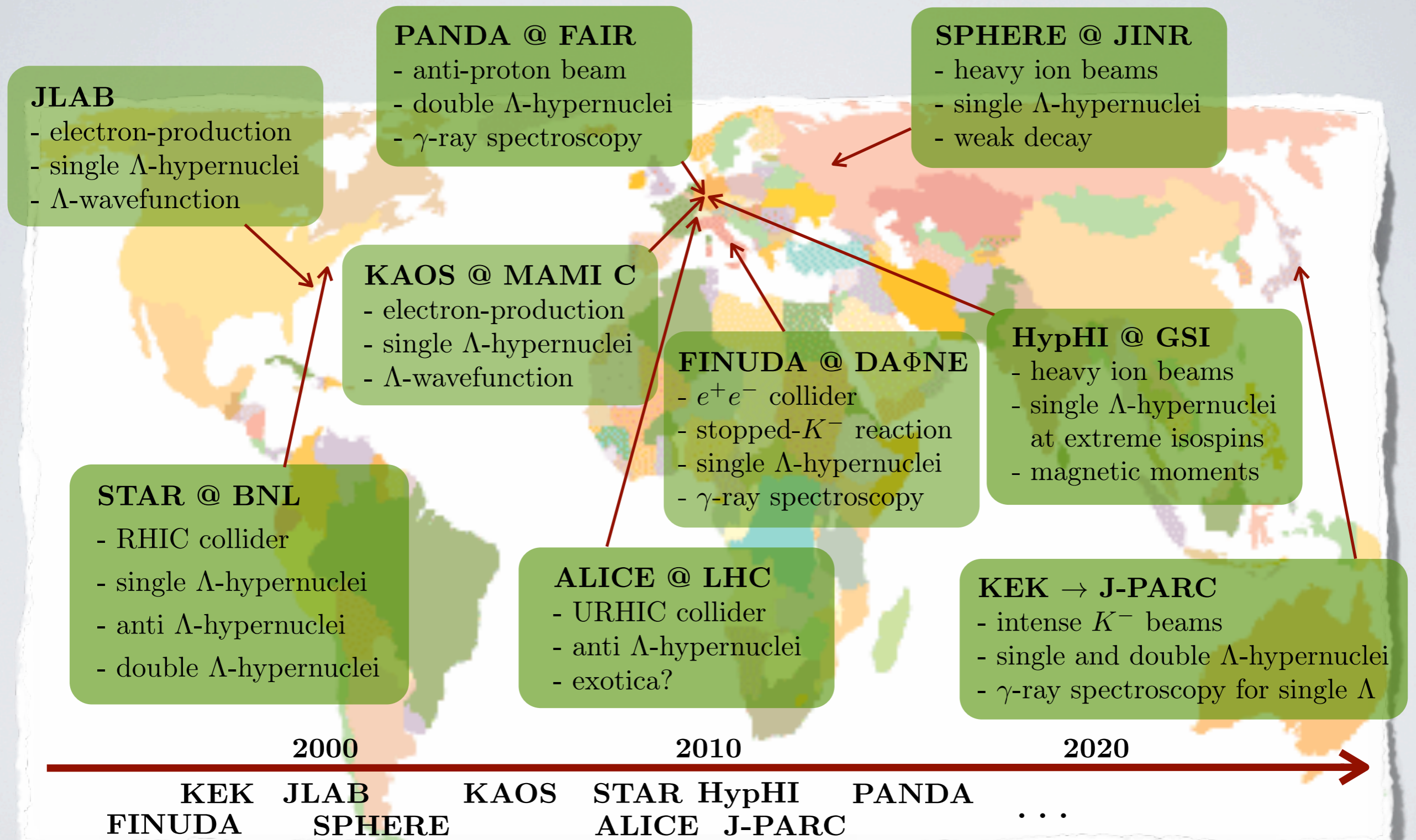


- ✓ Results
 - Λ -hypernuclei
 - Λ -neutron matter



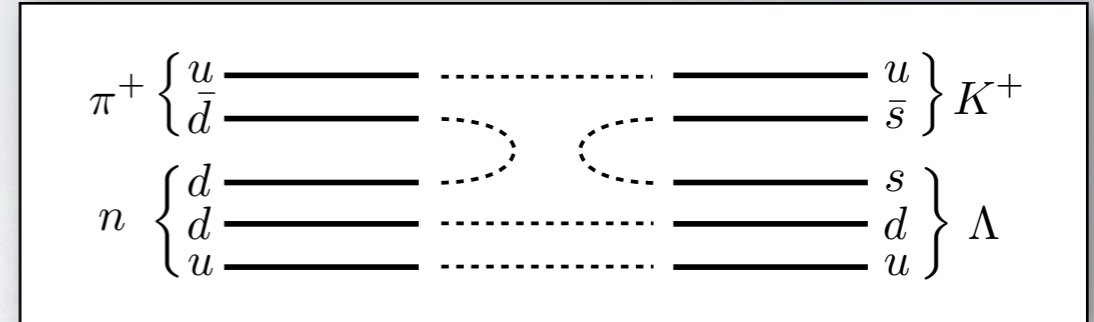
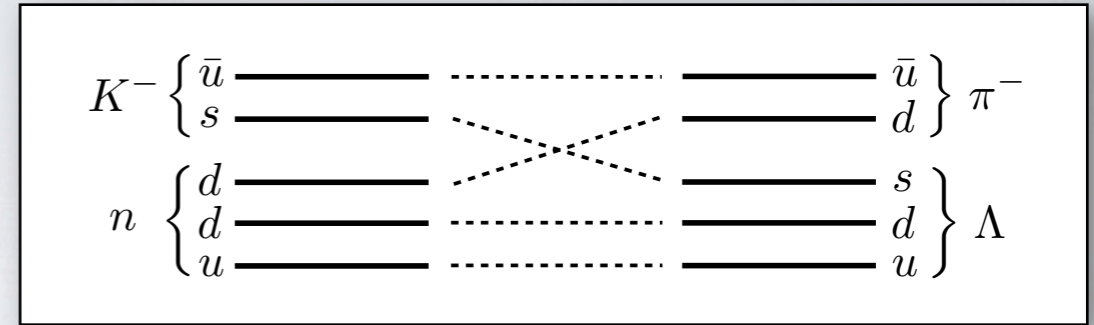
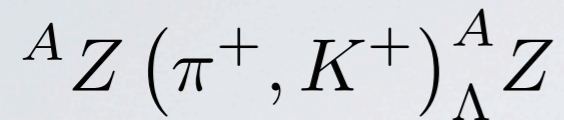
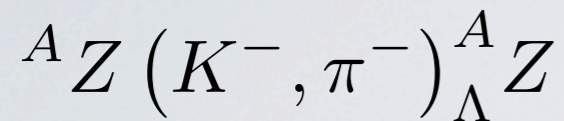
- ✓ Conclusions

Strangeness in nuclear physics: hypernuclei experiments

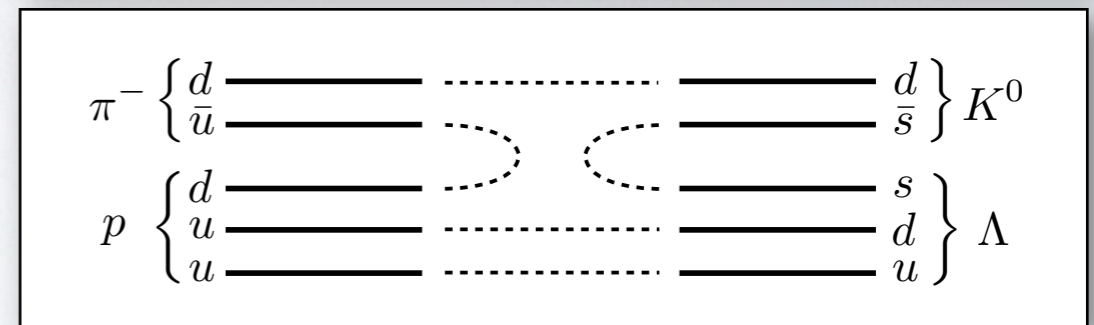
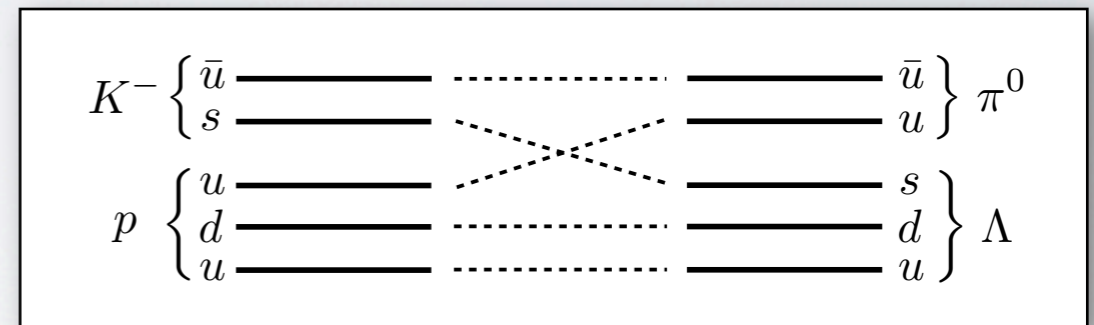
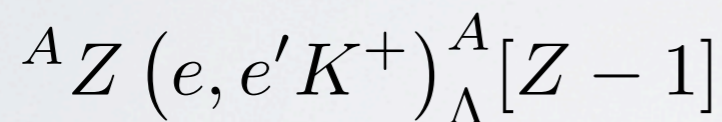
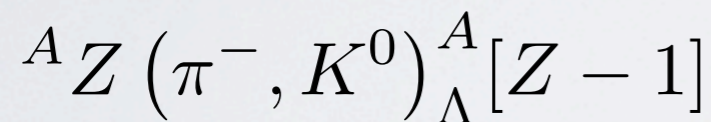
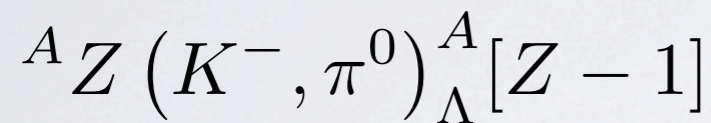


Strangeness in nuclear physics: hypernuclei experiments

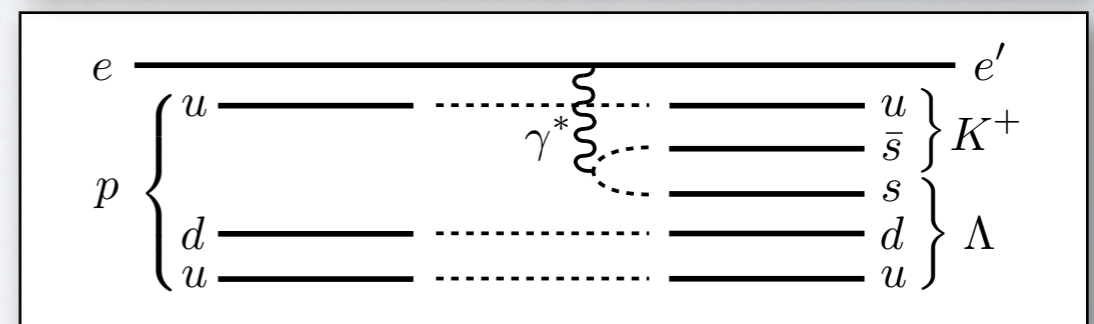
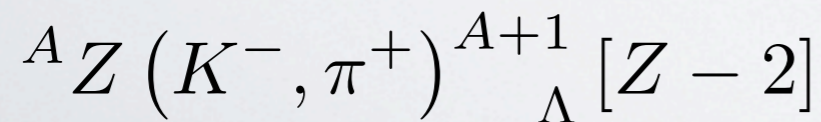
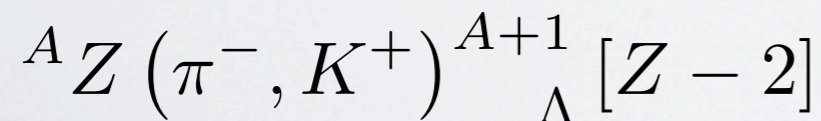
✓ Charge conserving reactions



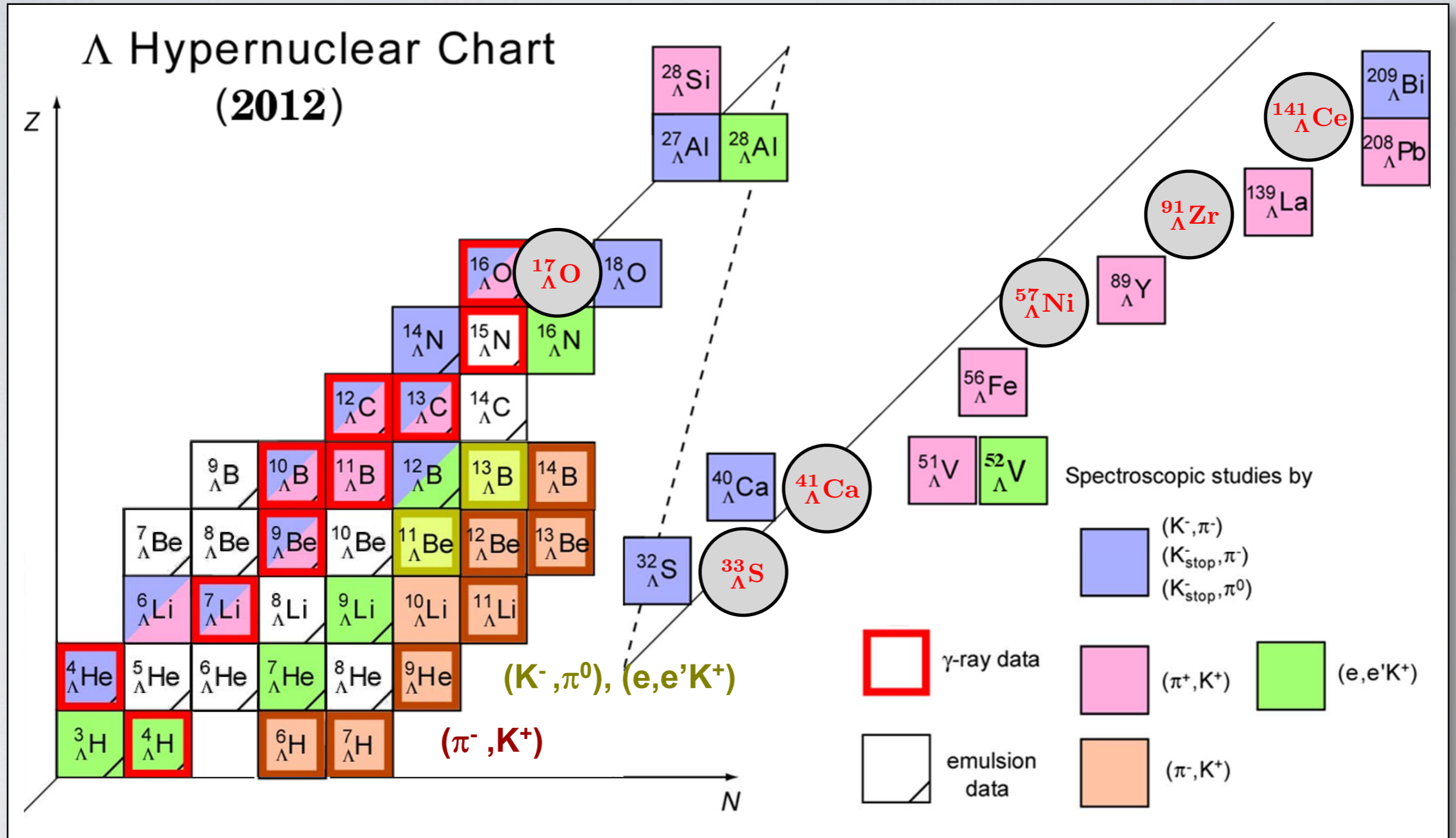
✓ Single charge exchange reactions (SCX)



✓ Double charge exchange reactions (DCX)



Strangeness in nuclear physics: hypernuclei experiments

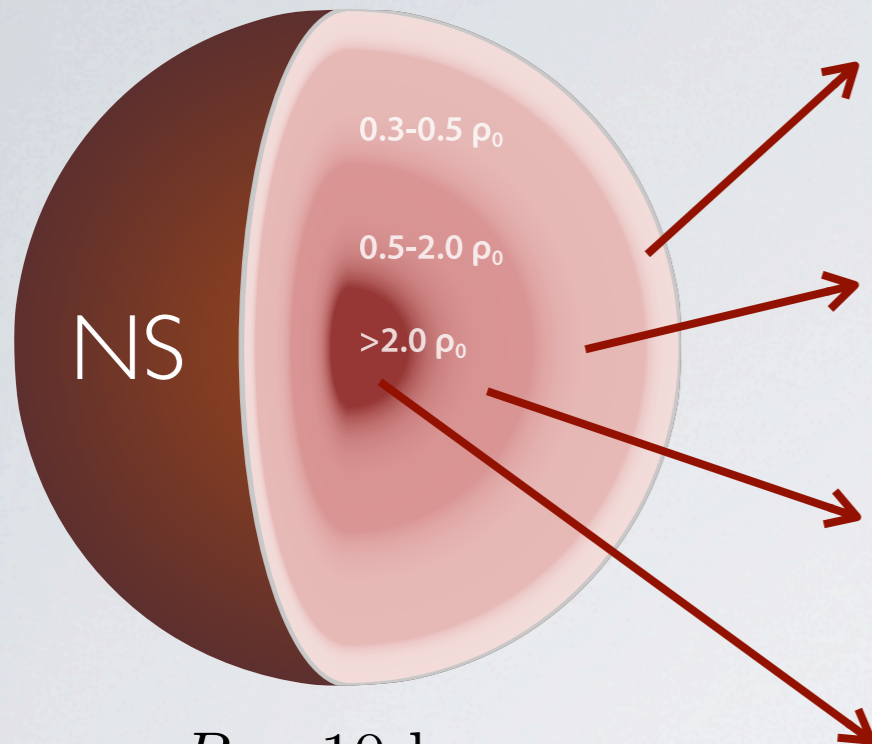


Updated from: O. Hashimoto, H. Tamura, Prog. Part. Nucl. Phys. 57, 564 (2006)

Λp scattering data

{	~ 600 low energy ($p_{\text{lab}} = 200 \div 300 \text{ Mev}/c$)
	~ 250 high energy ($p_{\text{lab}} = 300 \div 1500 \text{ Mev}/c$)

Strangeness in nuclear physics: neutron stars



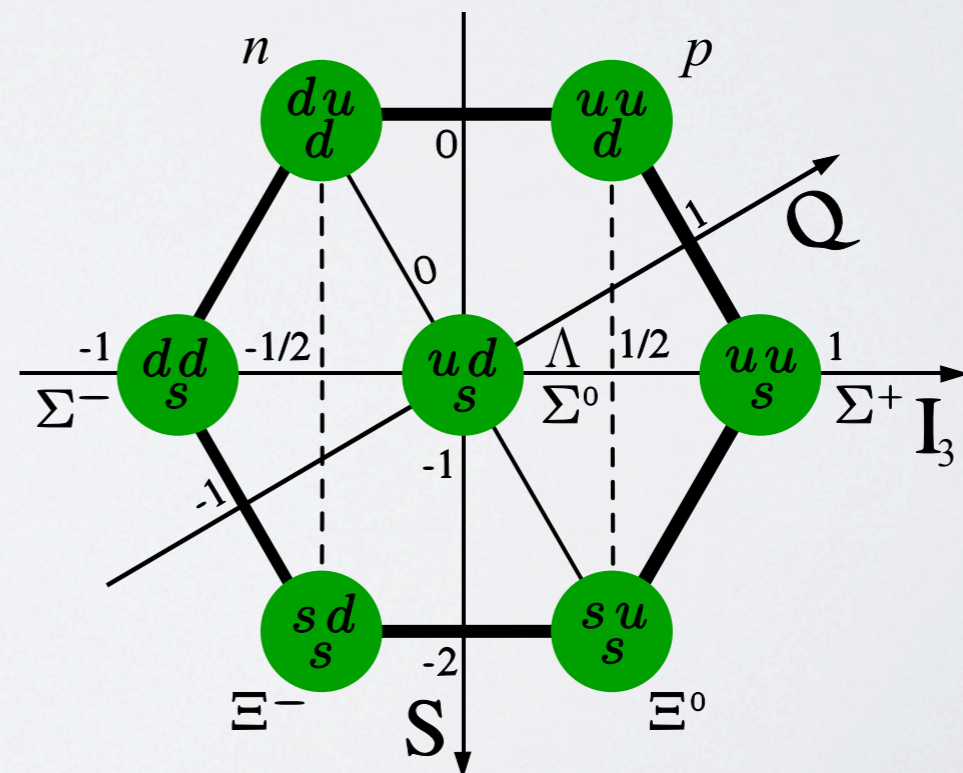
$R \sim 10 \text{ km}$
 $M \sim 1.4 M_{\odot}$

outer crust: $Z e$
 (0.3 ÷ 0.5 km)

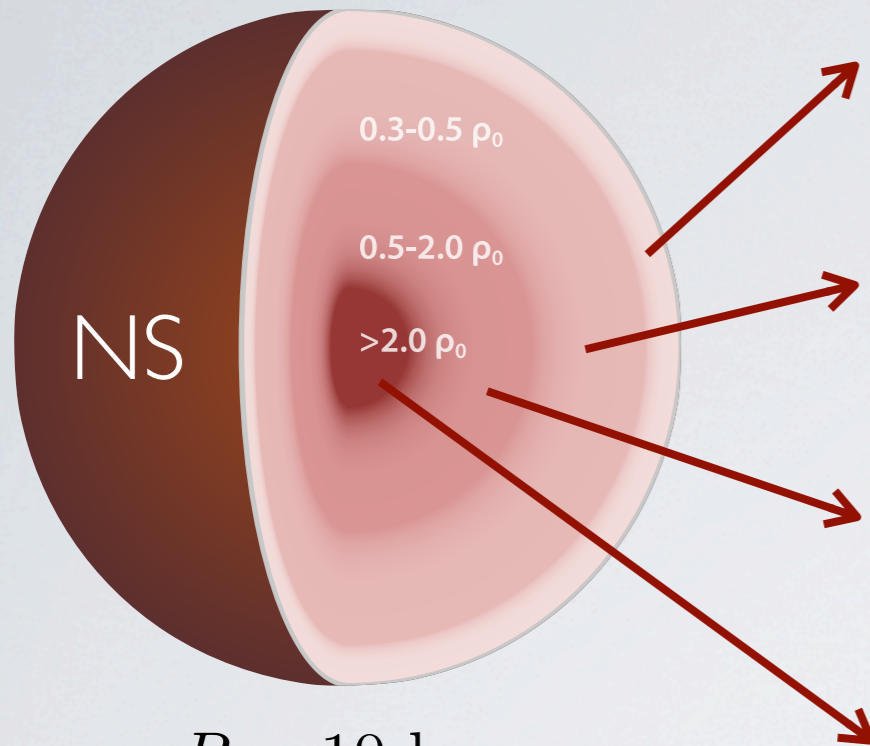
inner crust: $Z n e$
 (1 ÷ 2 km)

outer core: $n p e \mu$
 ($\sim 9 \text{ km}$)

inner core: $n p e \mu \Lambda \Sigma \Xi \pi_c K_c q_p \text{ ?}$



Strangeness in nuclear physics: neutron stars



$R \sim 10 \text{ km}$
 $M \sim 1.4 M_{\odot}$

outer crust: $Z e$
 (0.3 ÷ 0.5 km)

inner crust: $Z n e$
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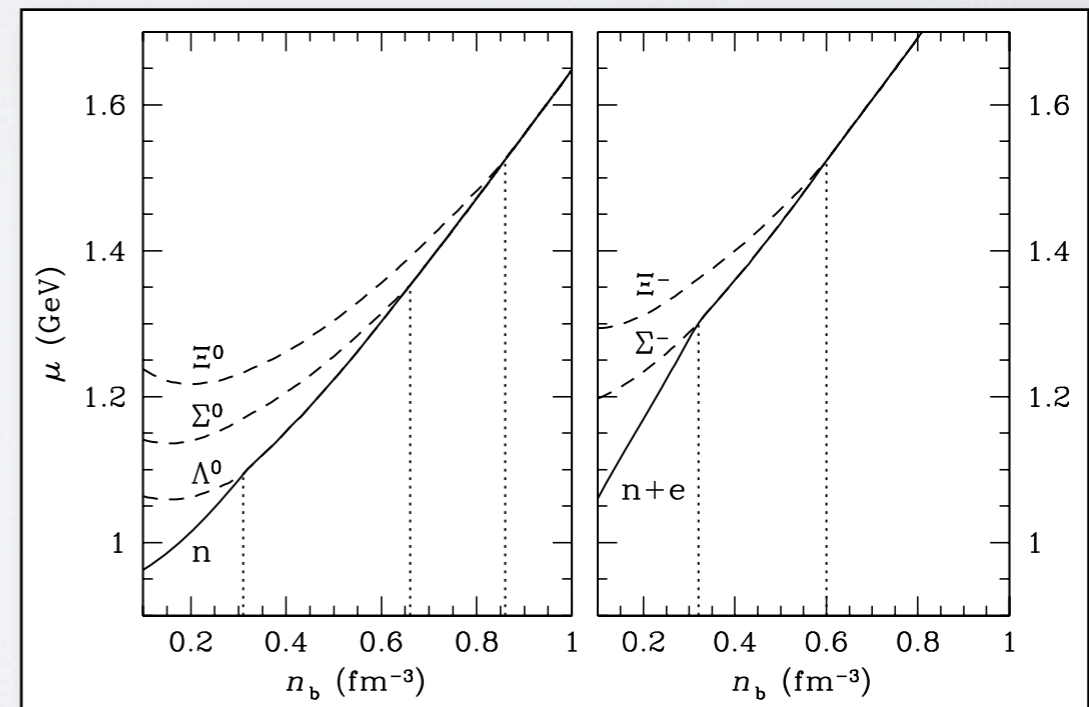
outer core: $n p e \mu$
 ($\sim 9 \text{ km}$)

inner core: $n p e \mu \Lambda \Sigma \Xi \pi_c K_c q_p ?$

$$Q = -1 : \mu_{b^-} = \mu_n + \mu_e$$

$$Q = 0 : \mu_{b^0} = \mu_n$$

$$Q = +1 : \mu_{b^+} = \mu_n - \mu_e$$

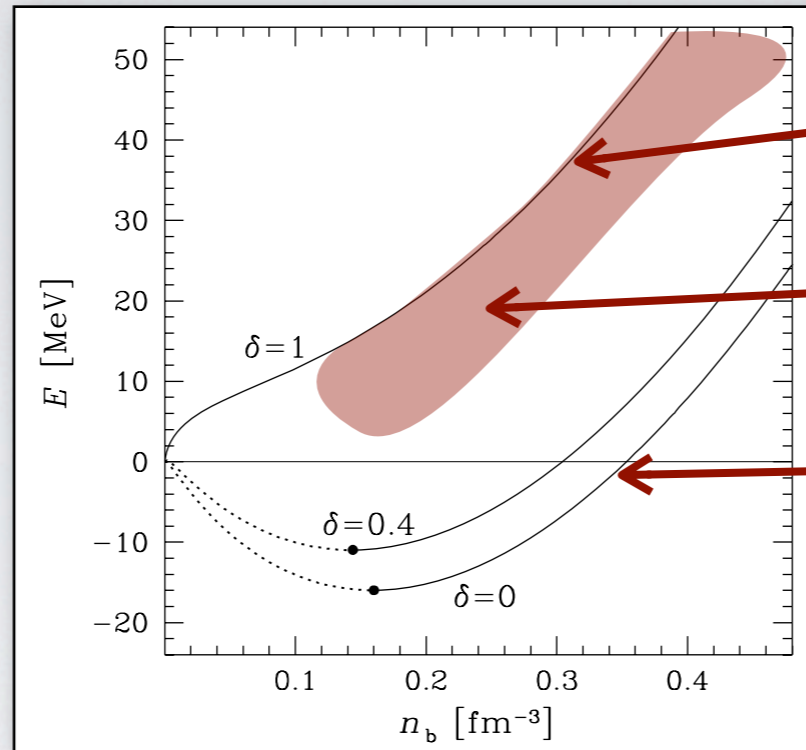


Strangeness in nuclear physics: neutron stars

$$\begin{cases} E \equiv E(n_b, \delta) \\ P = n_b^2 \frac{\partial E(n_b, \delta)}{\partial n_b} \end{cases}$$

$$n_b = n_p + n_n = A/V$$

$$\delta = \frac{n_n - n_p}{n_b}$$



pure n matter

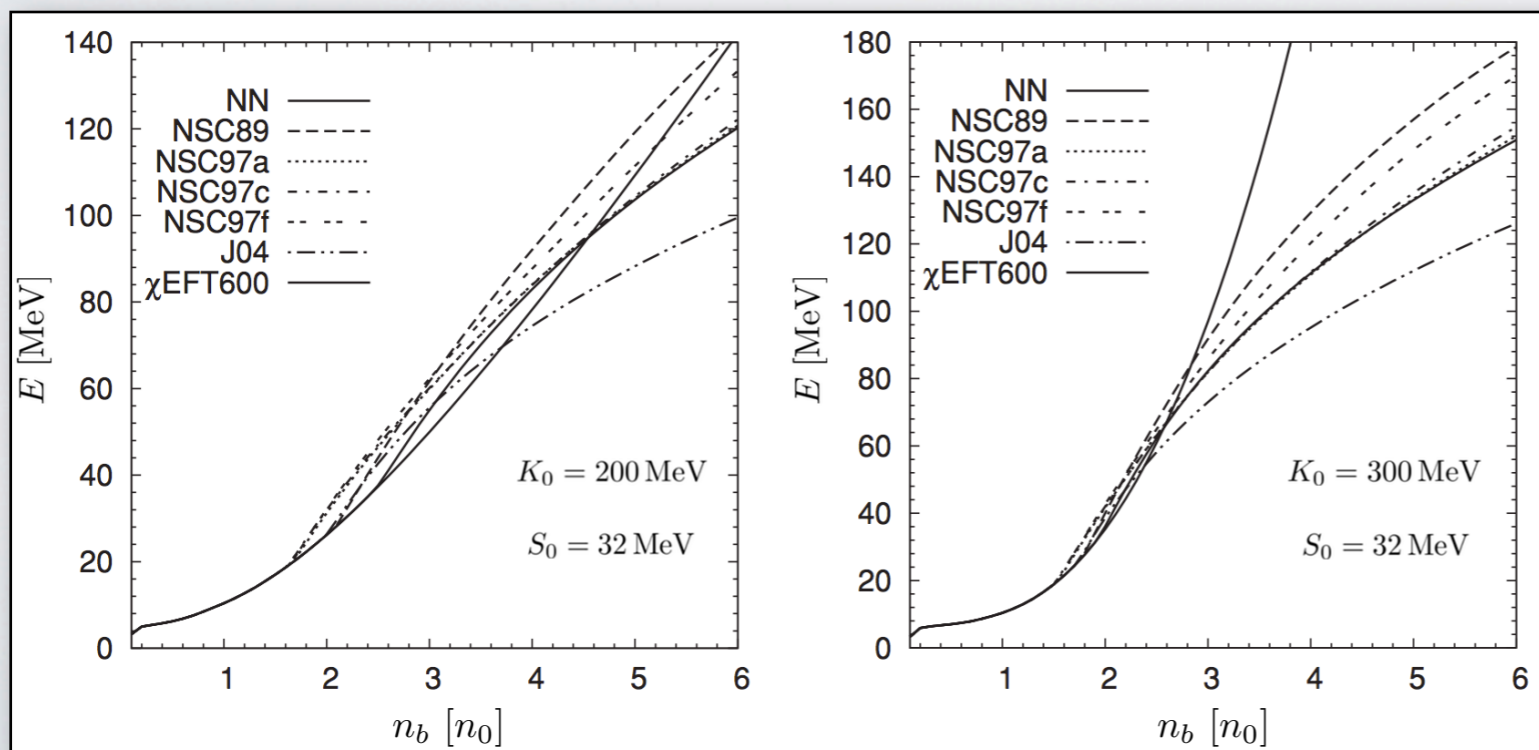
NS core

symmetric matter

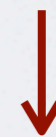
$$n_0 = 0.16 \text{ fm}^{-3}$$

$$E_0 = -16 \text{ MeV}$$

P. Haensel, A.Y. Potekhin, D.G. Yakovlev,
Neutron Stars I, Springer 2007



hyperons: softening of
the EOS

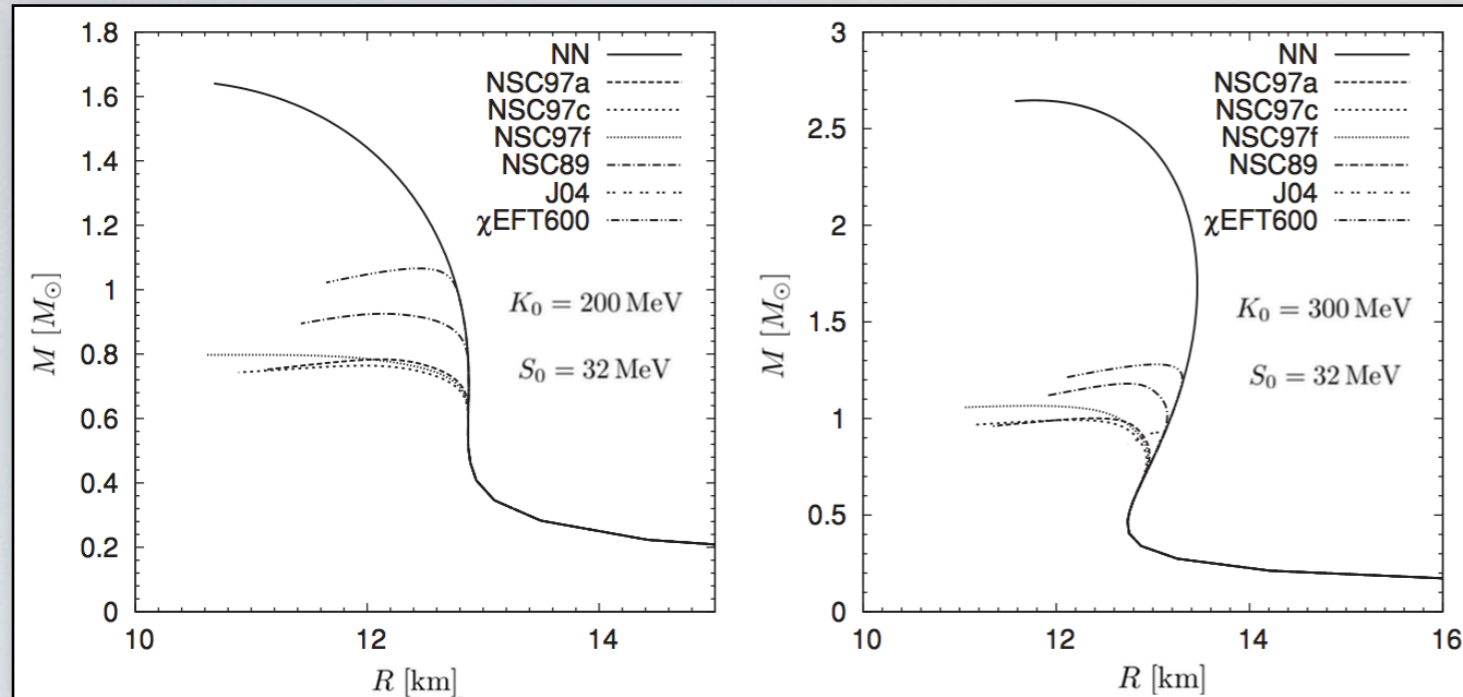


$M(R)$ & M_{max} (TOV)

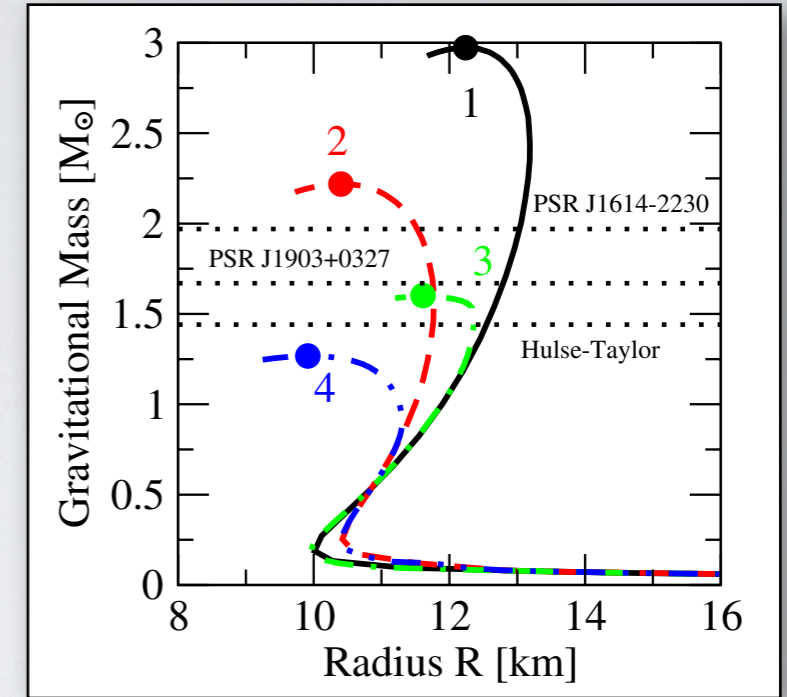
model dependent

H. Dapo, B.-J. Schaefer, and J. Wambach, Phys. Rev. C 81, 035803 (2010)

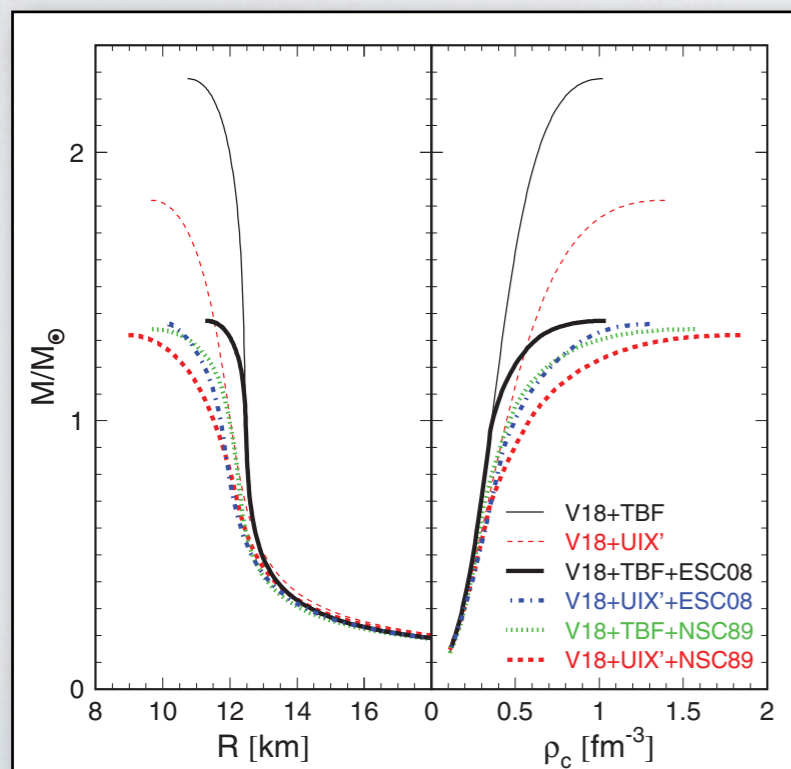
Strangeness in nuclear physics: neutron stars



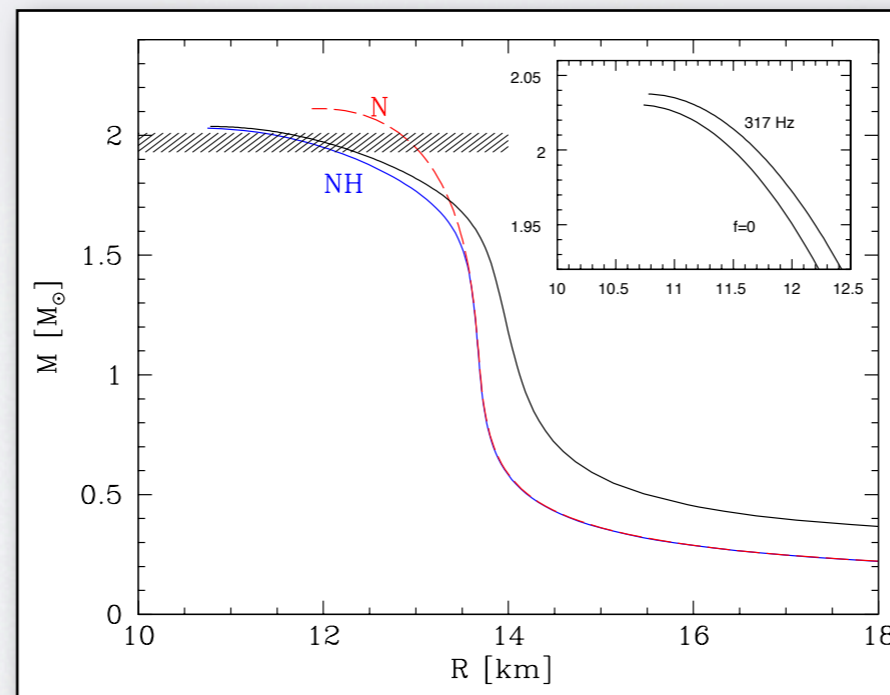
H. Ćapo, B.-J. Schaefer, J. Wambach, Phys. Rev. C 81, 035803 (2010)



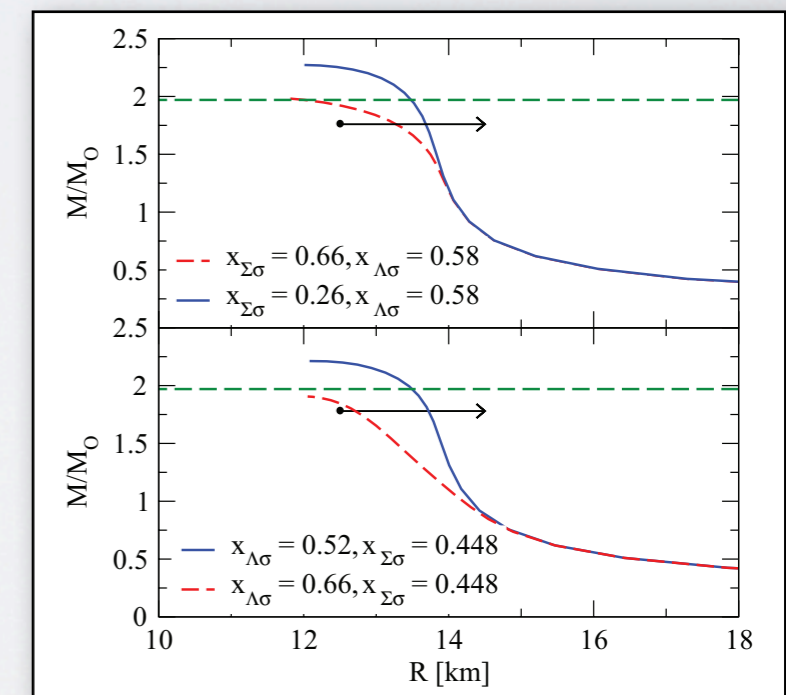
I. Vidaña, D. Logoteta, C. Providência, A. Polls, I. Bombaci, EPL 94, 11002 (2011)



H.-J. Schulze, T. Rijken, Phys. Rev. C 84, 035801 (2011)

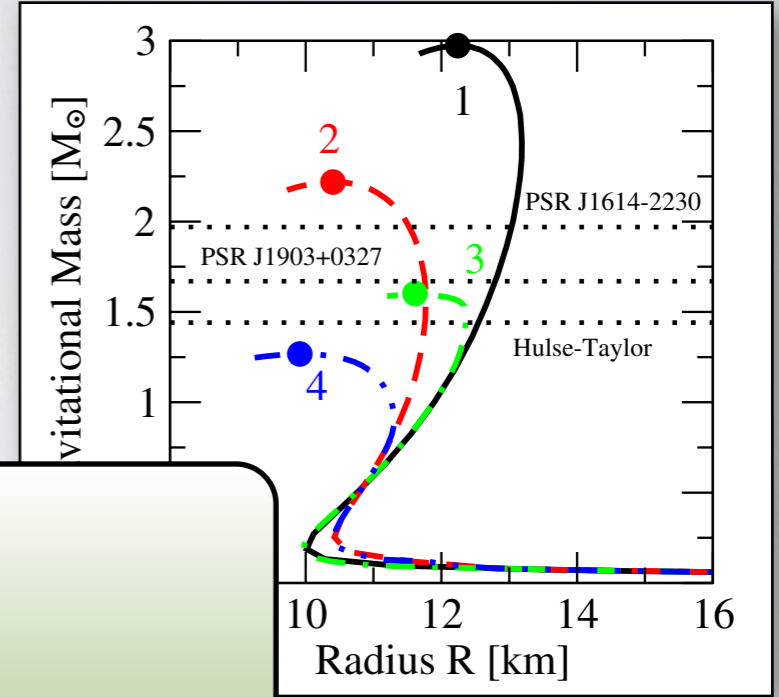
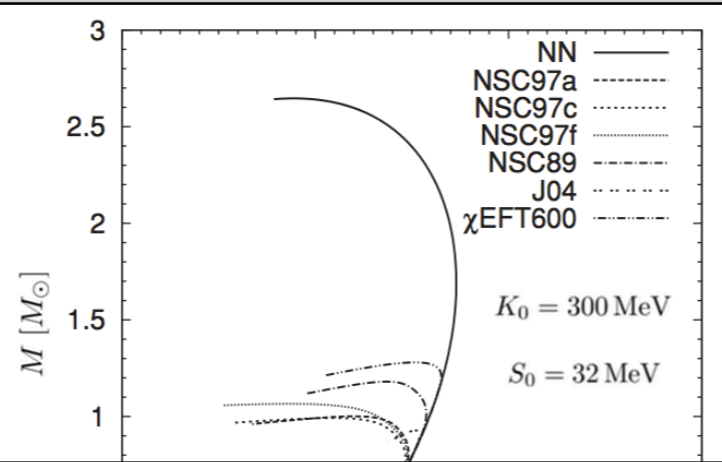
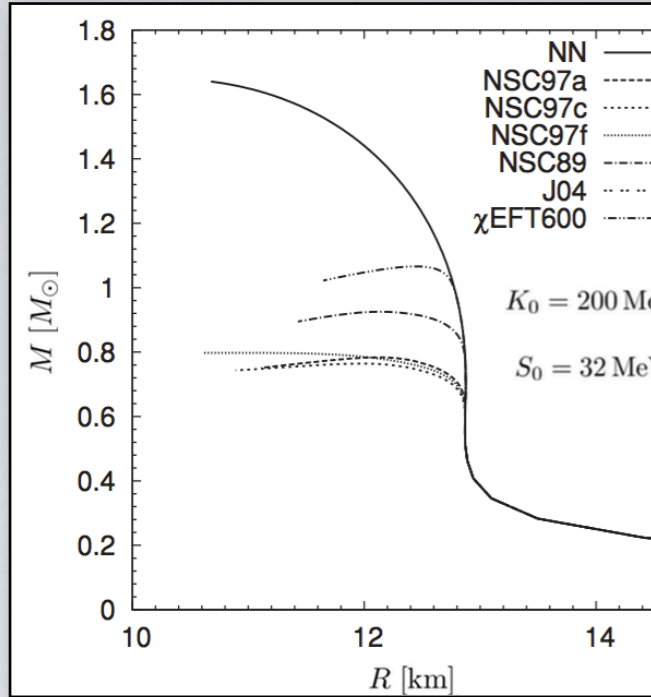


I. Bednarek, P. Haensel, J. Ł. Zdunik, M. Bejger, R. Mańka, Astron. Astrophys. 543, A157 (2012)



G. Colucci, A. Sedrakian, Phys. Rev. C 87, 055806 (2013)

Strangeness in nuclear physics: neutron stars



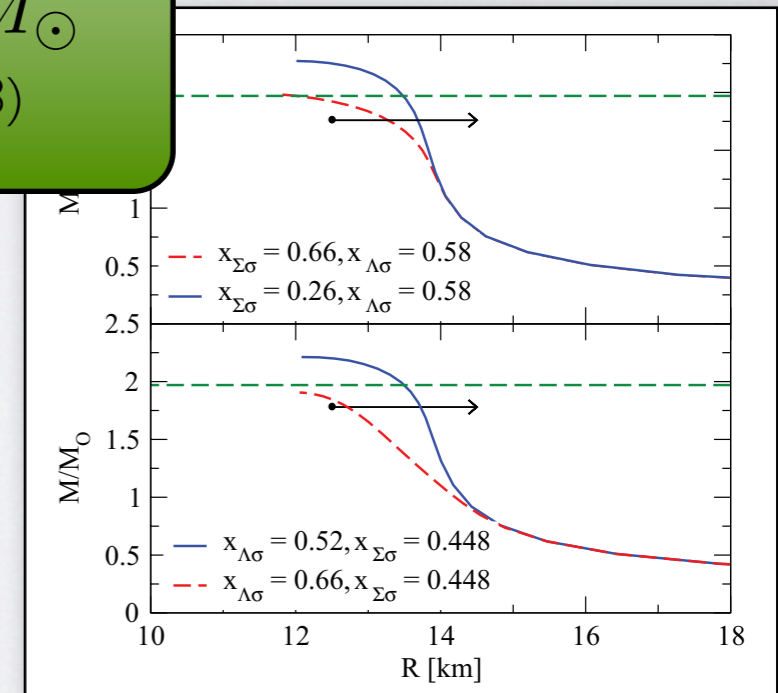
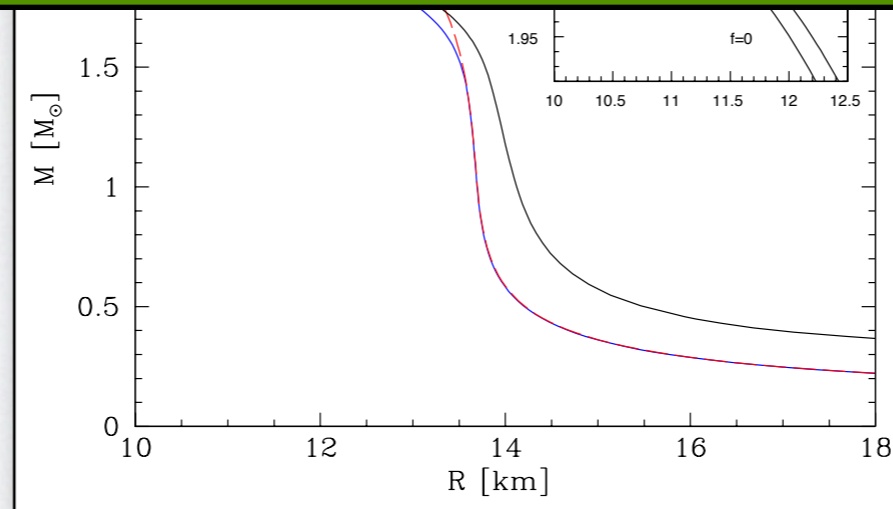
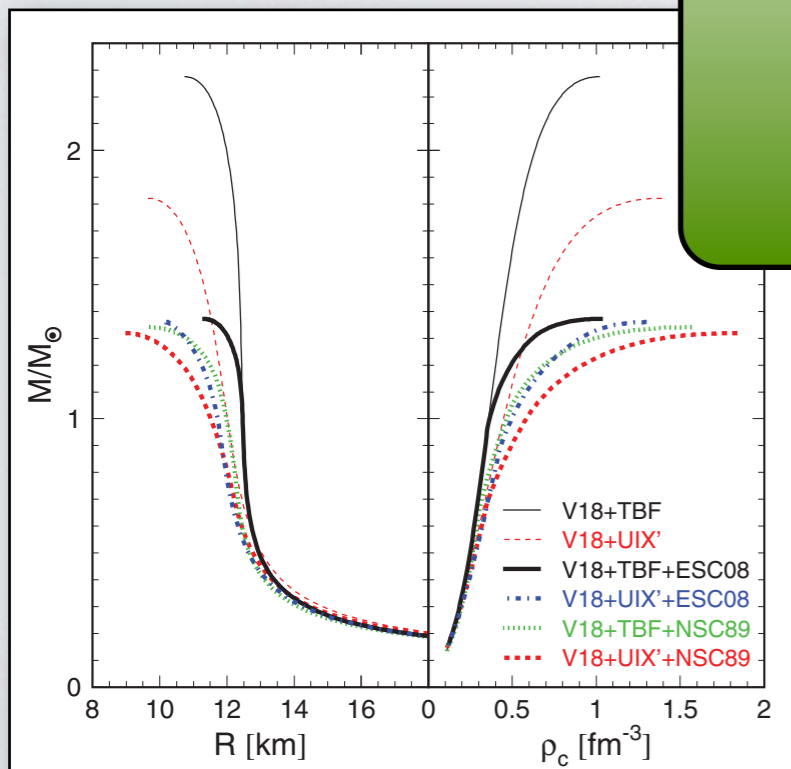
NS obs: $2 M_{\odot}$

PSR J1614-2230: $M = 1.97(4) M_{\odot}$
 P.B. Demorest et al., Nature 467, 1081 (2010)

PSR J0348+0432: $M = 2.01(4) M_{\odot}$
 J. Antoniadis et al., Science 340, 1233232 (2013)

H. Dapo, B.-J. Schaefer, J. Warr

ogoteta, C. Providência,
 paci, EPL 94, 11002 (2011)

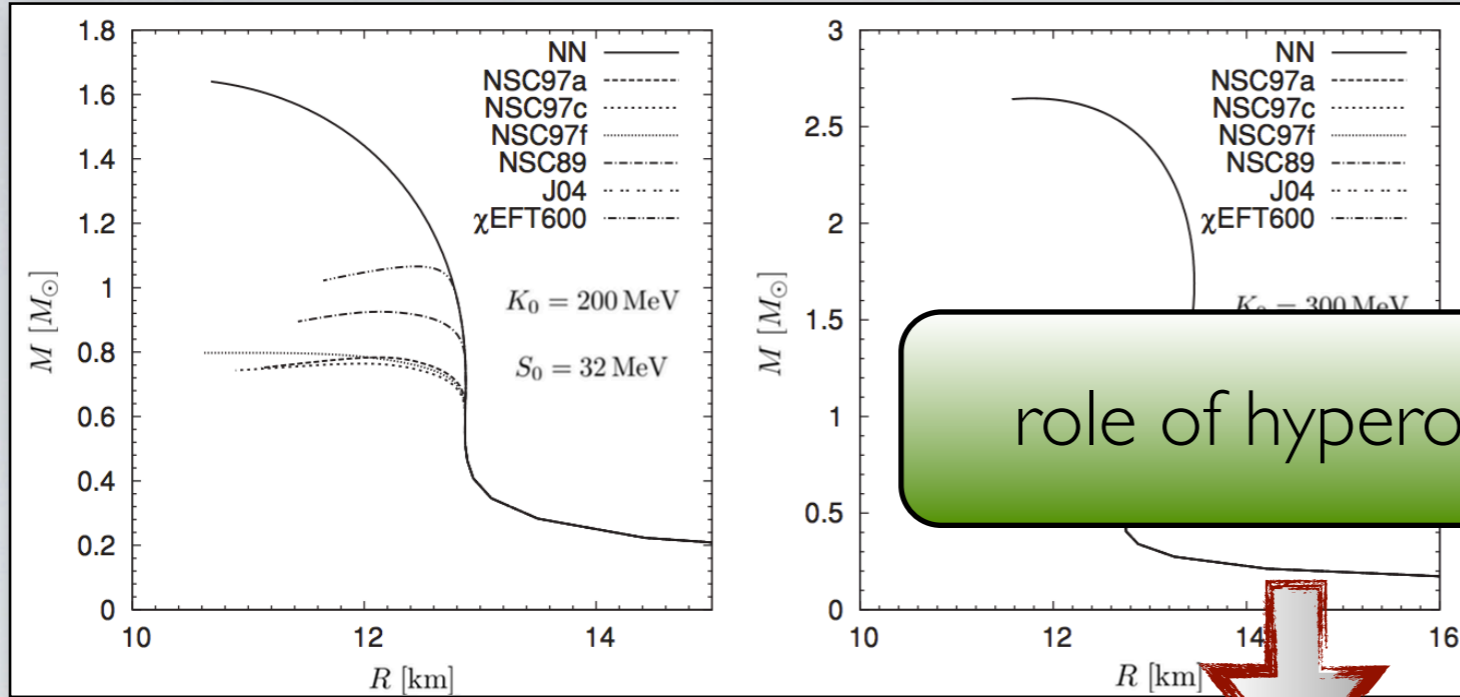


H.-J. Schulze, T. Rijken,
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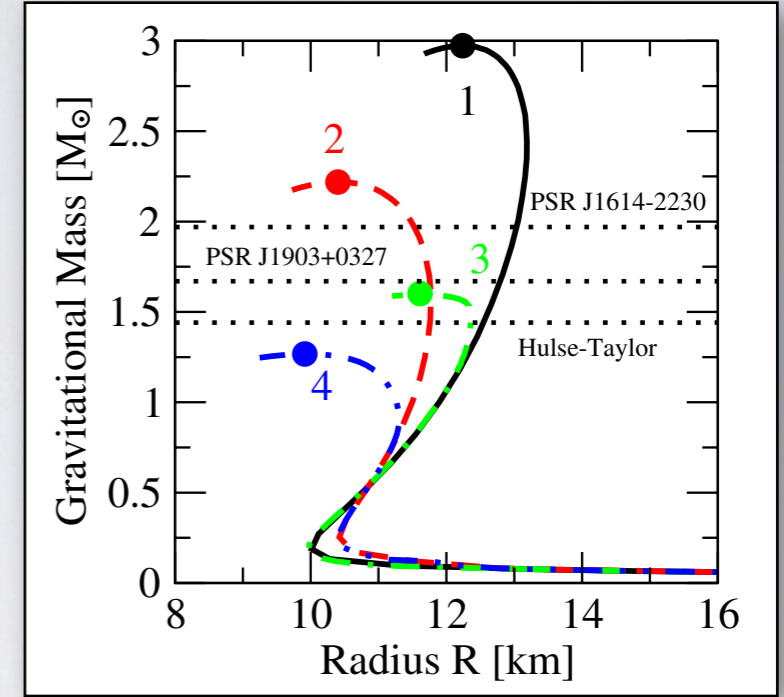
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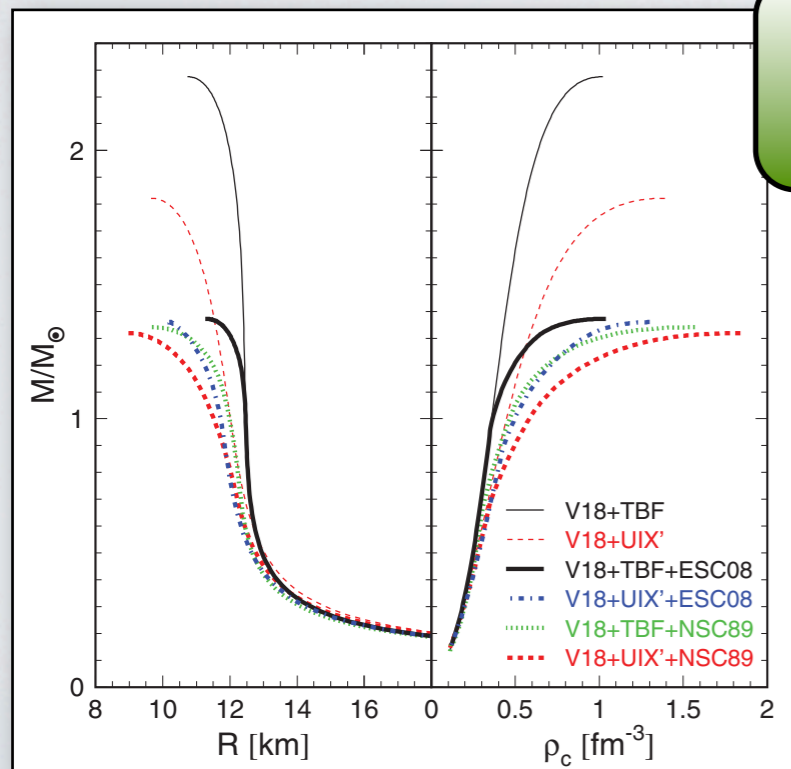
Strangeness in nuclear physics: neutron stars



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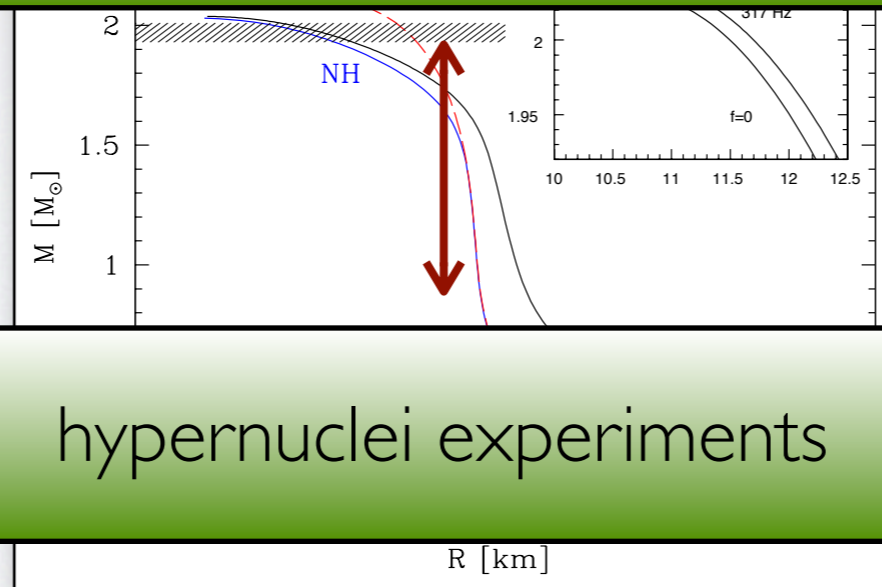


I. Vidaña, D. Logoteta, C. Providência, A. Polls, I. Bombaci, EPL 94, 11002 (2011)



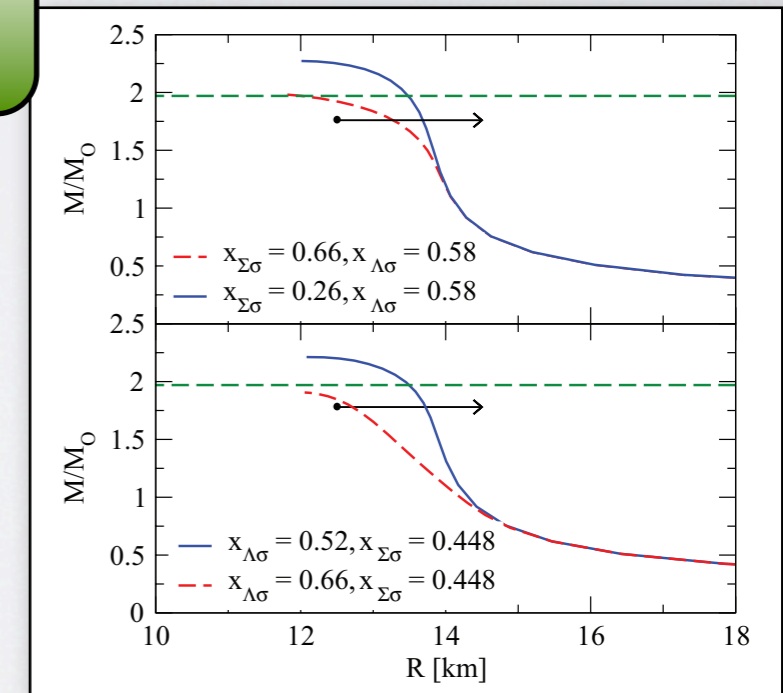
H.-J. Schulze, T. Rijken, Phys. Rev. C 84, 035801 (2011)

hyperon-nucleon interaction



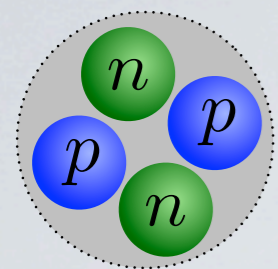
hypernuclei experiments

I. Bednarek, P. Haensel, J. L. Zdunik, M. Bejger, R. Mańka, Astron. Astrophys. 543, A157 (2012)

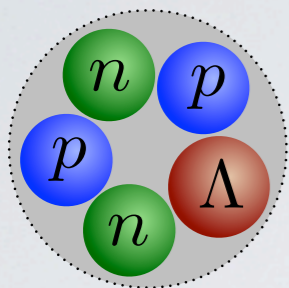


G. Colucci, A. Sedrakian, Phys. Rev. C 87, 055806 (2013)

The strange AFDMC project: the idea

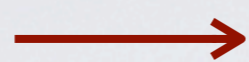


$A-1 Z$



$A_{\Lambda} Z$

nucleus



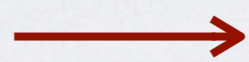
$$BE_{nuc} = \frac{\langle \psi_{nuc} | \mathcal{H}_{nuc} | \psi_{nuc} \rangle}{\langle \psi_{nuc} | \psi_{nuc} \rangle}$$

\mathcal{H}_{NN}

ab-initio
method : AFDMC



Λ -hypernucleus



$$BE_{hyp} = \frac{\langle \psi_{hyp} | \mathcal{H}_{hyp} | \psi_{hyp} \rangle}{\langle \psi_{hyp} | \psi_{hyp} \rangle}$$

$\mathcal{H}_{NN} + \mathcal{H}_{\Lambda N(N)}$

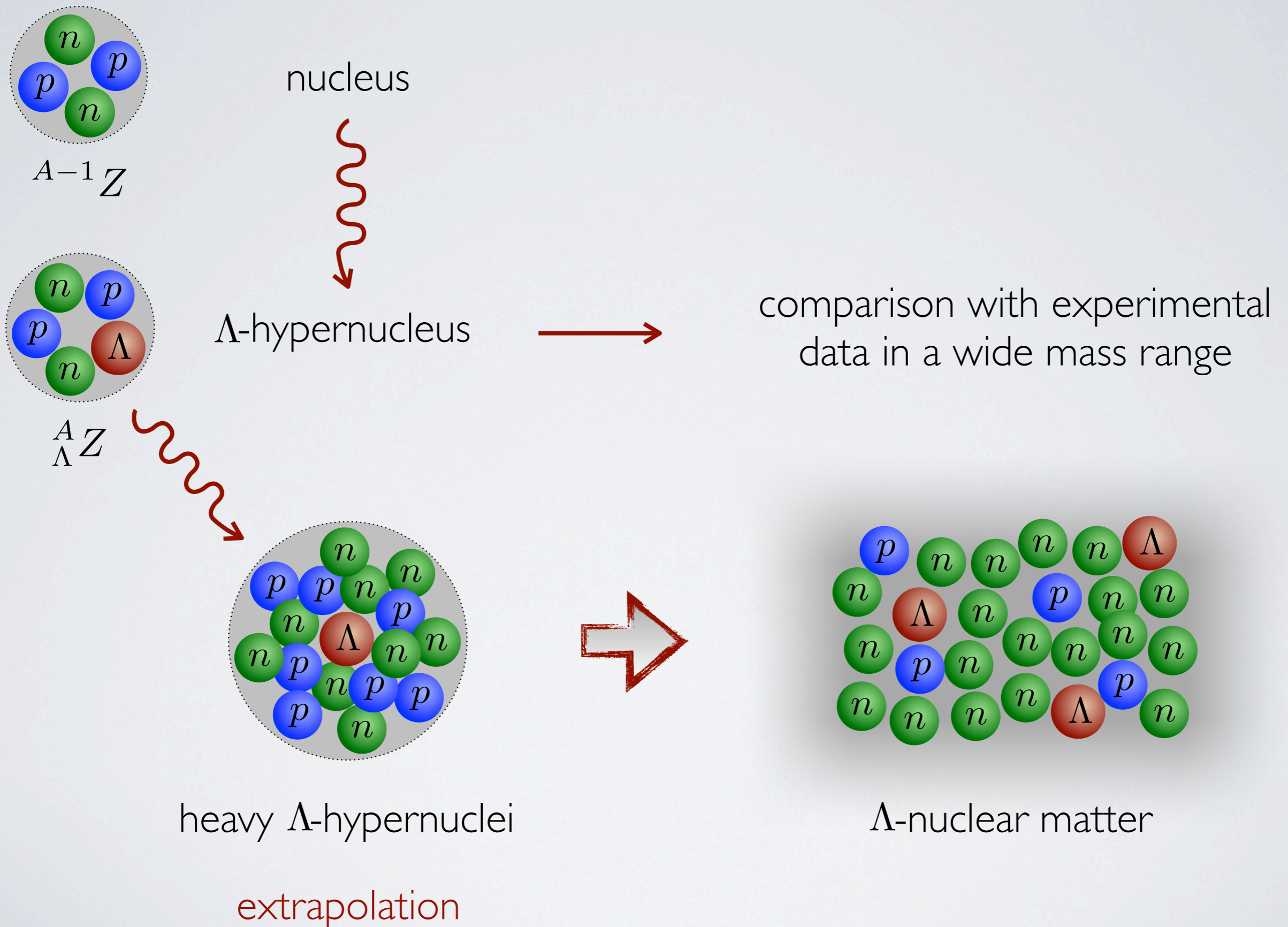
$$B_{\Lambda} = BE_{nuc} - BE_{hyp}$$

Hyp.: nuclear effects cancel at most

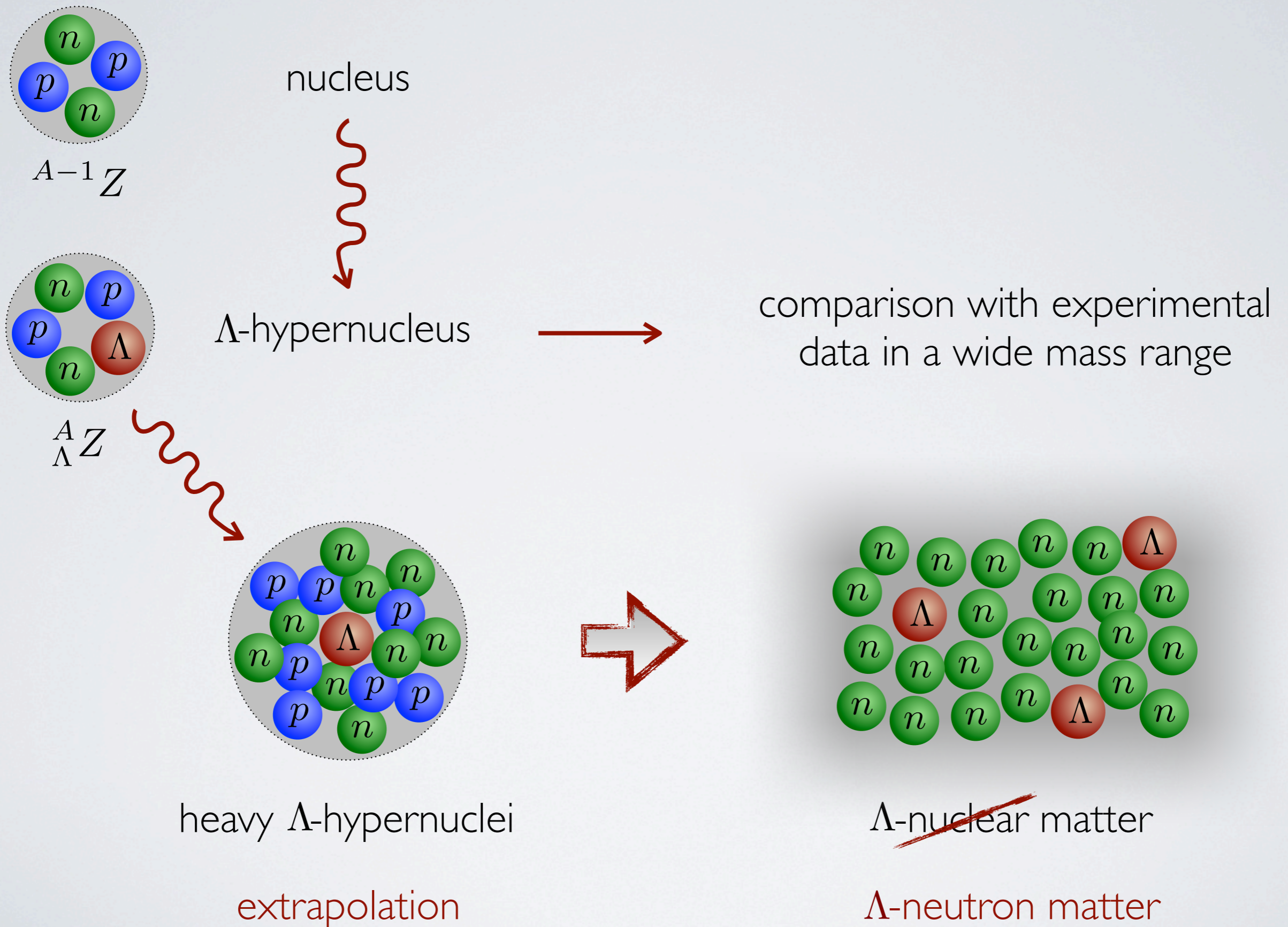


information about the hyperon-nucleon interaction

The strange AFDMC project: the idea



The strange AFDMC project: the idea



The strange AFDMC project: the hyperon-nucleon interaction

✓ Phenomenologic

- phenomenological nature
- 2-body terms
- shell & cluster model calculations

D.J. Millener, Nucl. Phys. A 1 (2013), in press

E. Hiyama, Nucl. Phys. A 1 (2013), in press

✓ Nijmegen & Jülich

- one meson exchange model
- 2-body terms, several parametrizations
- G-matrix calculations

I. Vidaña, A. Polls, A. Ramos, M. Hjorth-Jensen, Nucl. Phys. A 644, 201 (1998)

✓ χ -EFT

- derived from chiral EFT (NLO)
- 2-body terms (@ NLO)
- Faddeev-Yakubovsky calculations

A. Nogga, Nucl. Phys. A 1 (2013), in press

✓ Usmani & co.

- diagrammatic contributions, due to pion exchange
- 2-body and 3-body terms
- variational calculations

A. A. Usmani, F. C. Khanna, J. Phys. G: Nucl. Part. Phys. 35, 025105 (2008)

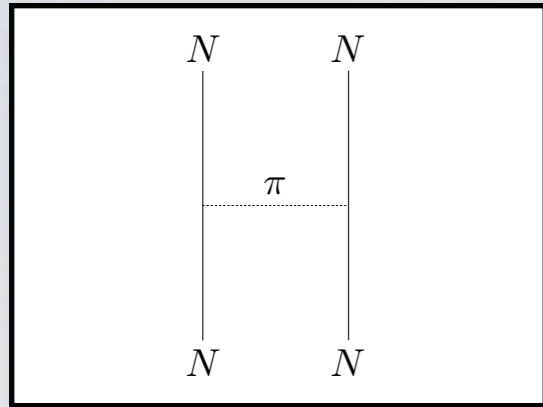
AV18+UIX like



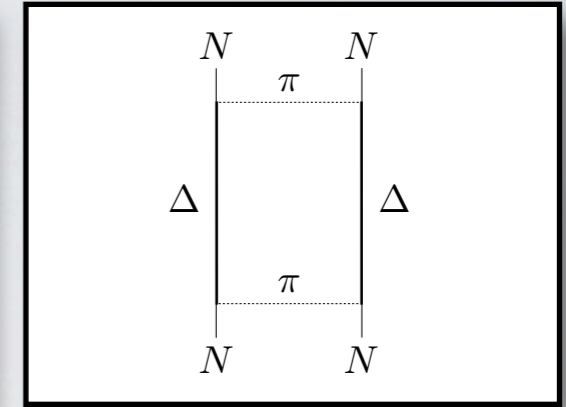
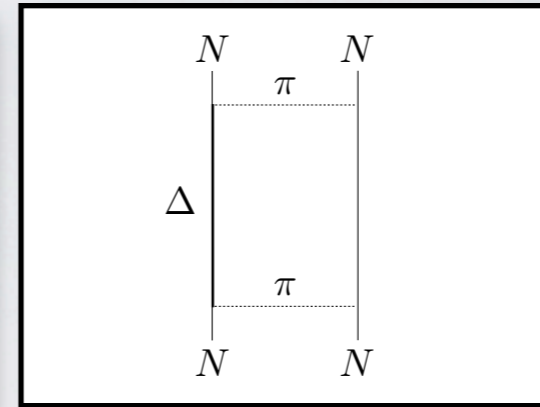
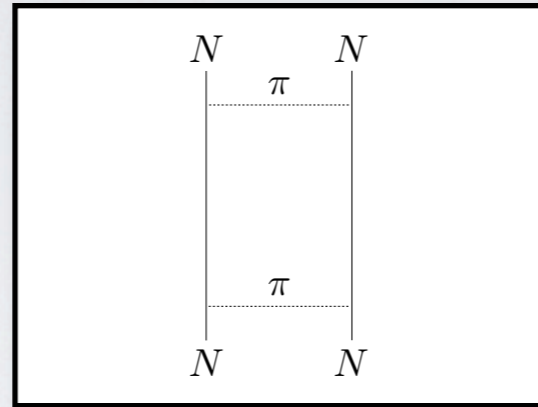
The strange AFDMC project: the hyperon-nucleon interaction

2-body

1π



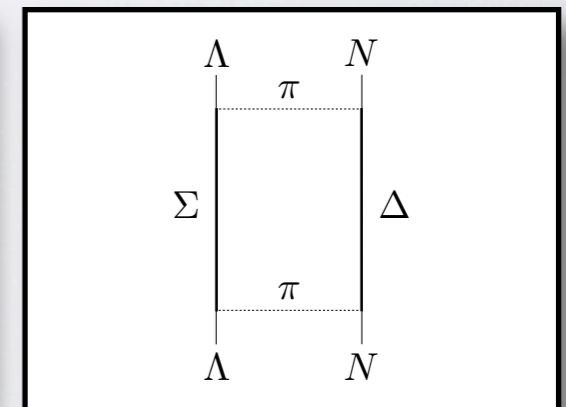
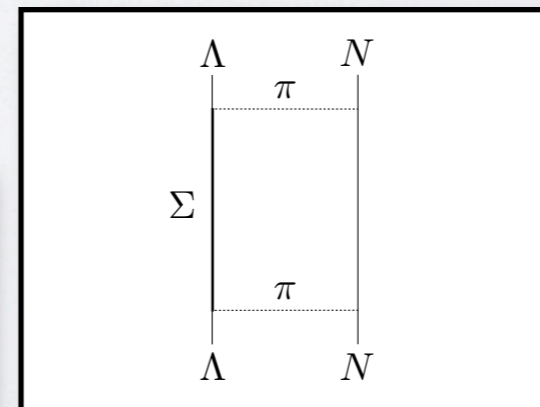
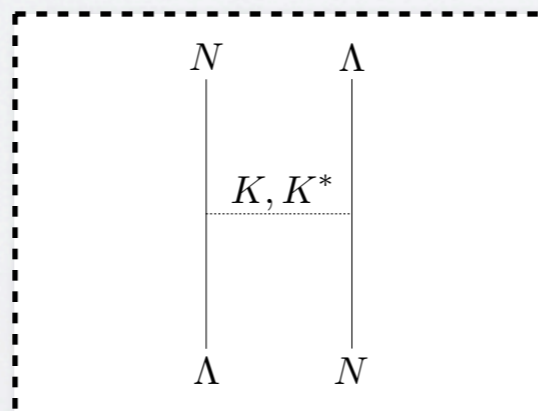
2π



$\Lambda\pi\Sigma$ vertex

forbidden

1 meson

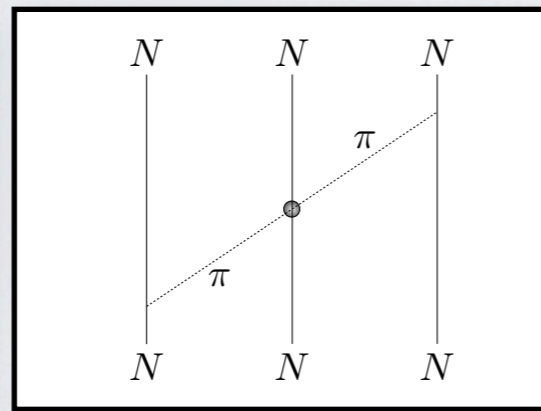
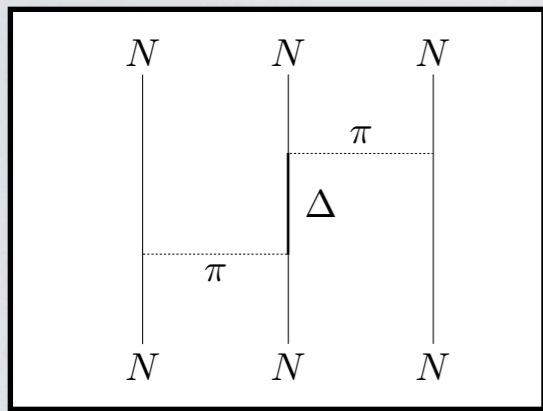


CSB ($A = 4$)

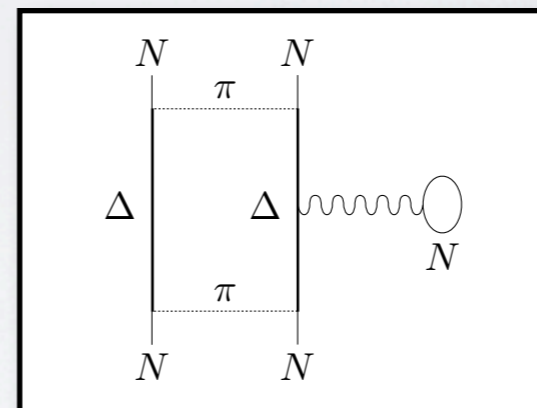
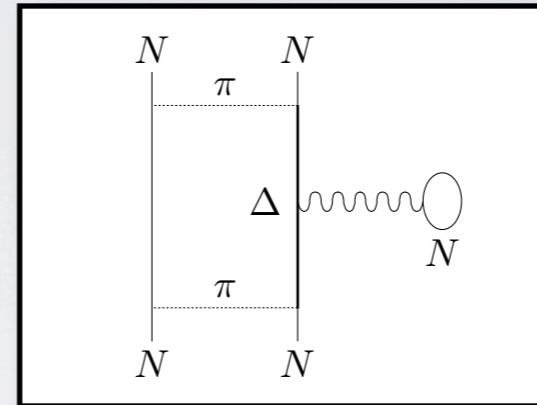
The strange AFDMC project: the hyperon-nucleon interaction

3-body

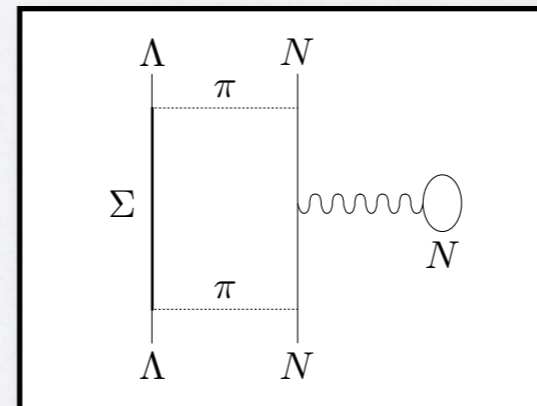
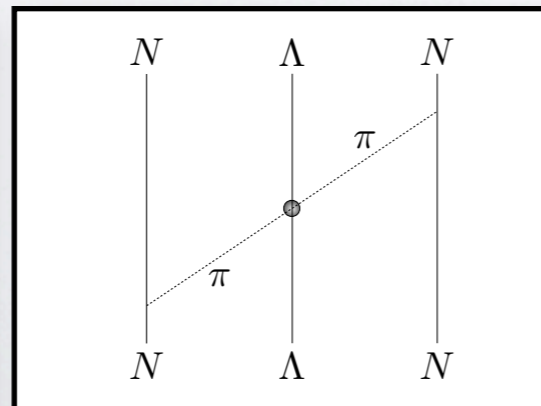
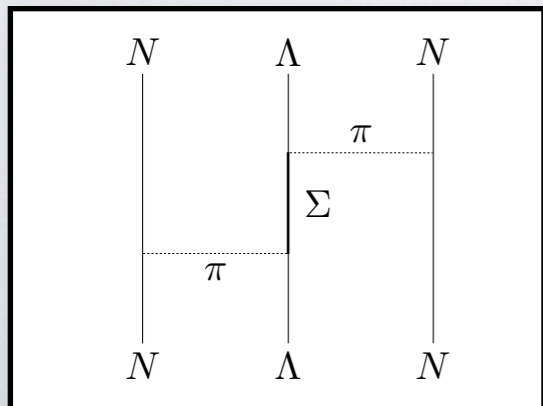
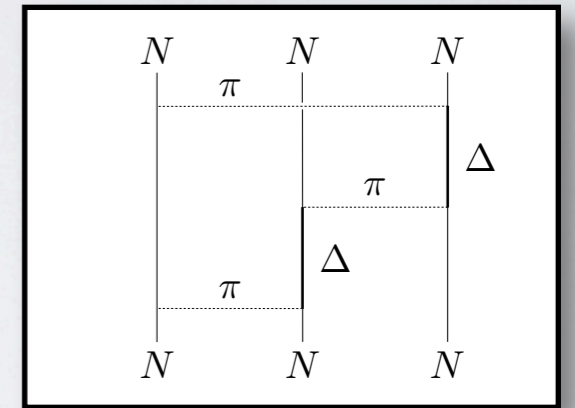
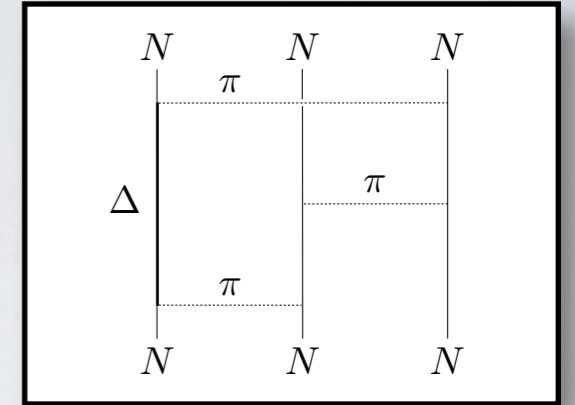
2π



dispersive



3π



forbidden

The strange AFDMC project: the hyperon-nucleon interaction

✓ 2-body ΛN interaction

- charge symmetric term: $v_{\Lambda i} = v_0(r) + v_0(r)\varepsilon(\mathcal{P}_x - 1) + \frac{1}{4}v_\sigma T_\pi^2(r)\boldsymbol{\sigma}_\Lambda \cdot \boldsymbol{\sigma}_i$

$$\left\{ \begin{array}{l} v_0(r) = v_c(r) - v_{2\pi}(r) \\ v_c(r) = W_c \left[1 + e^{\frac{r-\bar{r}}{a}} \right]^{-1} \\ v_{2\pi}(r) = \bar{v} T_\pi^2(r) \end{array} \right. \quad \left\{ \begin{array}{l} v_\sigma = v_s - v_t \\ \bar{v} = \frac{1}{4}(v_s + 3v_t) \end{array} \right.$$

parameters fitted on Λp scattering data

- charge symmetry breaking term: $v_{\Lambda i}^{\text{CSB}} = \tau_i^3 v_0^{\text{CSB}} T_\pi(r)$

parameter fitted on $A = 4$ mirror hypernuclei

The strange AFDMC project: the hyperon-nucleon interaction

✓ 3-body ΛNN interaction

$$v_{\Lambda ij} = v_{\Lambda ij}^{2\pi} + v_{\Lambda ij}^D = v_{\Lambda ij}^{PW} + v_{\Lambda ij}^{SW} + v_{\Lambda ij}^D$$

$$\left\{ \begin{array}{l} v_{\Lambda ij}^{PW} = -\frac{1}{6} C^P \{X_{i\Lambda}, X_{\Lambda j}\} \boldsymbol{\tau}_i \cdot \boldsymbol{\tau}_j \\ v_{\Lambda ij}^{SW} = C^S Z_\pi(r_{\Lambda i}) Z_\pi(r_{\Lambda j}) (\boldsymbol{\sigma}_i \cdot \hat{\mathbf{r}}_{i\Lambda} \boldsymbol{\sigma}_j \cdot \hat{\mathbf{r}}_{j\Lambda}) \boldsymbol{\tau}_i \cdot \boldsymbol{\tau}_j \\ v_{\Lambda ij}^D = W^D T_\pi^2(r_{\Lambda i}) T_\pi^2(r_{\Lambda j}) \left[1 + \frac{1}{6} \boldsymbol{\sigma}_\Lambda \cdot (\boldsymbol{\sigma}_i + \boldsymbol{\sigma}_j) \right] \end{array} \right.$$

parameters not yet fixed



fitting of the parameters to reproduce
experimental separation energies

The strange AFDMC project: the AFDMC code

✓ Diffusion Monte Carlo

$$-\frac{\partial}{\partial \tau} \psi(\mathbf{R}, \mathbf{S}, \tau) = \mathcal{H} \psi(\mathbf{R}, \mathbf{S}, \tau) \quad \rightarrow \quad \psi(\mathbf{R}, \mathbf{S}, \tau + d\tau) = e^{-\mathcal{H}d\tau} \psi(\mathbf{R}, \mathbf{S}, \tau)$$

$(\tau = it/\hbar)$ $\downarrow \tau \rightarrow \infty$

ground state

$$\psi(\mathbf{R}, \mathbf{S}, \tau + d\tau) = \int \langle \mathbf{S} \mathbf{R} | e^{-(\mathcal{H} - E_0)d\tau} | \mathbf{R}' \mathbf{S}' \rangle \langle \mathbf{S}' \mathbf{R}' | \psi(\tau) \rangle d\mathbf{R}' d\mathbf{S}'$$

$$D = \hbar^2 / 2m$$

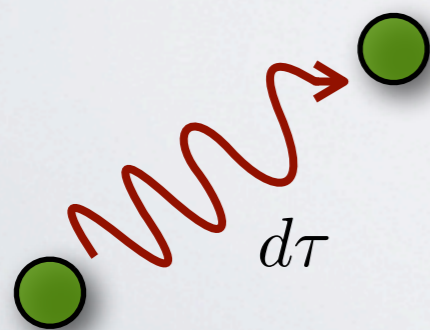
diffusion

branching

walkers

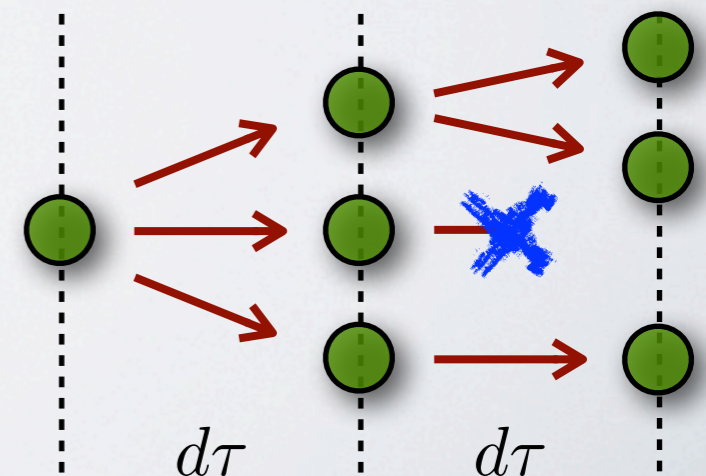
$$(4\pi D d\tau)^{\frac{3A}{2}} e^{-\frac{(\mathbf{R} - \mathbf{R}')^2}{4D d\tau}}$$

$$e^{-\left(\frac{\hat{V}(\mathbf{R}') + \hat{V}(\mathbf{R})}{2} - E_0\right) d\tau}$$



kinetic term

potential term



The strange AFDMC project: the AFDMC code

✓ Auxiliary Field

$$\mathcal{P} \sim e^{-\frac{1}{2}\lambda d\tau \mathcal{O}^2} \quad \Rightarrow \quad \psi \sim \frac{A!}{Z!(A-Z)!} 2^A \text{ terms} \quad \text{GFMC: } A \leq 12$$

Idea: Hubbard-Stratonovich transformation

$$e^{-\frac{1}{2}\lambda d\tau \mathcal{O}^2} = \frac{1}{\sqrt{2\pi}} \int dx e^{-\frac{x^2}{2} + \sqrt{-\lambda d\tau} x \mathcal{O}}$$

auxiliary field

rotation over spin-isospin configurations

potential matrices diagonalization

computational cost $\sim A^3$

The strange AFDMC project: the AFDMC code

✓ AFDMC matrices: nuclear systems

$$V_{NN}(\text{AV6}) = \sum_{i < j} \sum_{p=1}^6 v_p(r) \mathcal{O}_{ij}^p \quad \mathcal{O}_{ij}^{p=1,6} = \{1, \boldsymbol{\sigma}_i \cdot \boldsymbol{\sigma}_j, S_{ij}\} \otimes \{1, \boldsymbol{\tau}_i \cdot \boldsymbol{\tau}_j\}$$

$$V_{NN}^{sd} = \frac{1}{2} \sum_{i \neq j} \sum_{\gamma} \tau_i^{\gamma} \left(\mathcal{A}_{ij}^{[\tau]} \right) \tau_j^{\gamma} \quad A \times A$$

$$+ \frac{1}{2} \sum_{i \neq j} \sum_{\alpha\beta} \sigma_i^{\alpha} \left(\mathcal{A}_{i\alpha, j\beta}^{[\sigma]} \right) \sigma_j^{\beta} \quad 3A \times 3A$$

$$+ \frac{1}{2} \sum_{i \neq j} \sum_{\alpha\beta\gamma} \tau_i^{\gamma} \sigma_i^{\alpha} \left(\mathcal{A}_{i\alpha, j\beta}^{[\sigma\tau]} \right) \sigma_j^{\beta} \tau_j^{\gamma} \quad 3A \times 3A$$

Ex: $\sum_{j\beta} \mathcal{A}_{i\alpha, j\beta}^{[\sigma]} \psi_{n, j\beta}^{[\sigma]} = \lambda_n^{[\sigma]} \psi_{n, i\alpha}^{[\sigma]}$



$$\tilde{\mathcal{O}}_n^{[\sigma]} = \sum_{j\beta} \sigma_j^{\beta} \psi_{n, j\beta}^{[\sigma]}$$

The strange AFDMC project: the AFDMC code

- ✓ AFDMC matrices: hyper-nuclear systems

$$V_{\Lambda N} = \sum_i \left(\cancel{\mathcal{B}_{(\Lambda)i}^{[\hat{P}_x]}} \right) \hat{P}_{x\Lambda i} + \sum_i \left(\mathcal{B}_i^{[CSB]} \right) \tau_i^z$$
$$+ \sum_i \sum_{\alpha} \sigma_{\Lambda}^{\alpha} \left(\mathcal{B}_{(\Lambda)i}^{[\sigma]} \right) \sigma_i^{\alpha} \quad 1 \times (A - 1)$$

$$V_{\Lambda NN}^D = \frac{1}{2} \sum_{i \neq j} \sum_{\alpha} \sigma_{\Lambda}^{\alpha} \left(\mathcal{C}_{(\Lambda j)i}^{[\sigma]} \right) \sigma_i^{\alpha} \quad 1 \times (A - 1)$$

$$V_{\Lambda NN}^{2\pi} = \frac{1}{2} \sum_{i \neq j} \sum_{\alpha\beta\gamma} \tau_i^{\gamma} \sigma_i^{\alpha} \left(\mathcal{C}_{i\alpha,j\beta}^{[\sigma\tau]} \right) \sigma_j^{\beta} \tau_j^{\gamma} \quad 3(A - 1) \times 3(A - 1)$$

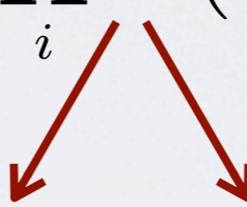


good for Hubbard-Stratonovich

The strange AFDMC project: the AFDMC code

✓ Wave functions

$$\Psi_T(R, S) = \Psi_T^N(R_N, S_N) \otimes \Psi_T^\Lambda(R_\Lambda, S_\Lambda)$$

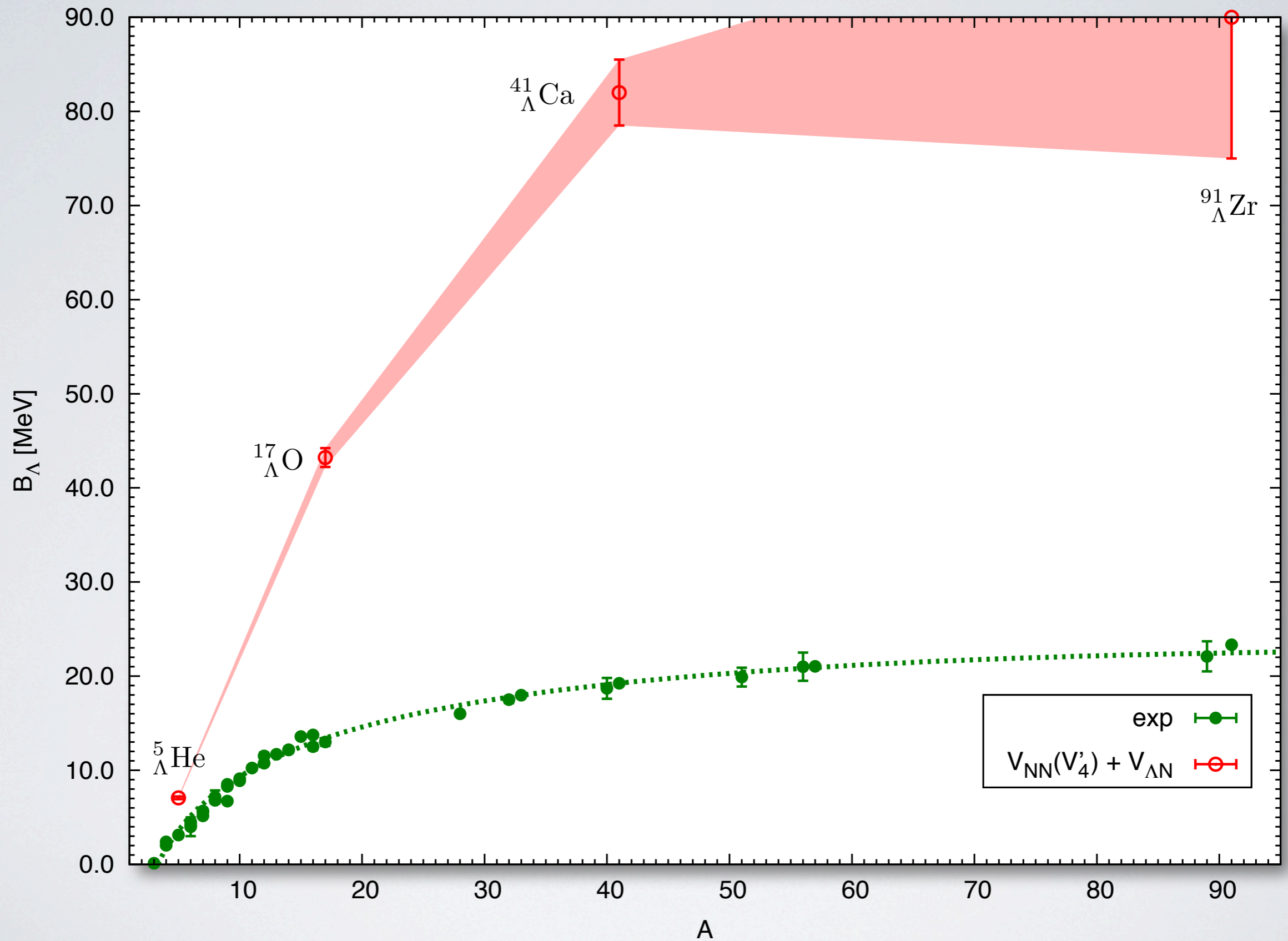
$$\Psi_T^p(R_p, S_p) = \left[\prod_{i < j} f_{ij}^c \right]_p \mathfrak{A} \left[\prod_i \phi_i(\vec{r}_i - \vec{R}_{\text{CM}}, \vec{s}_i) \right]_p$$


HF s.p. orbitals plane waves

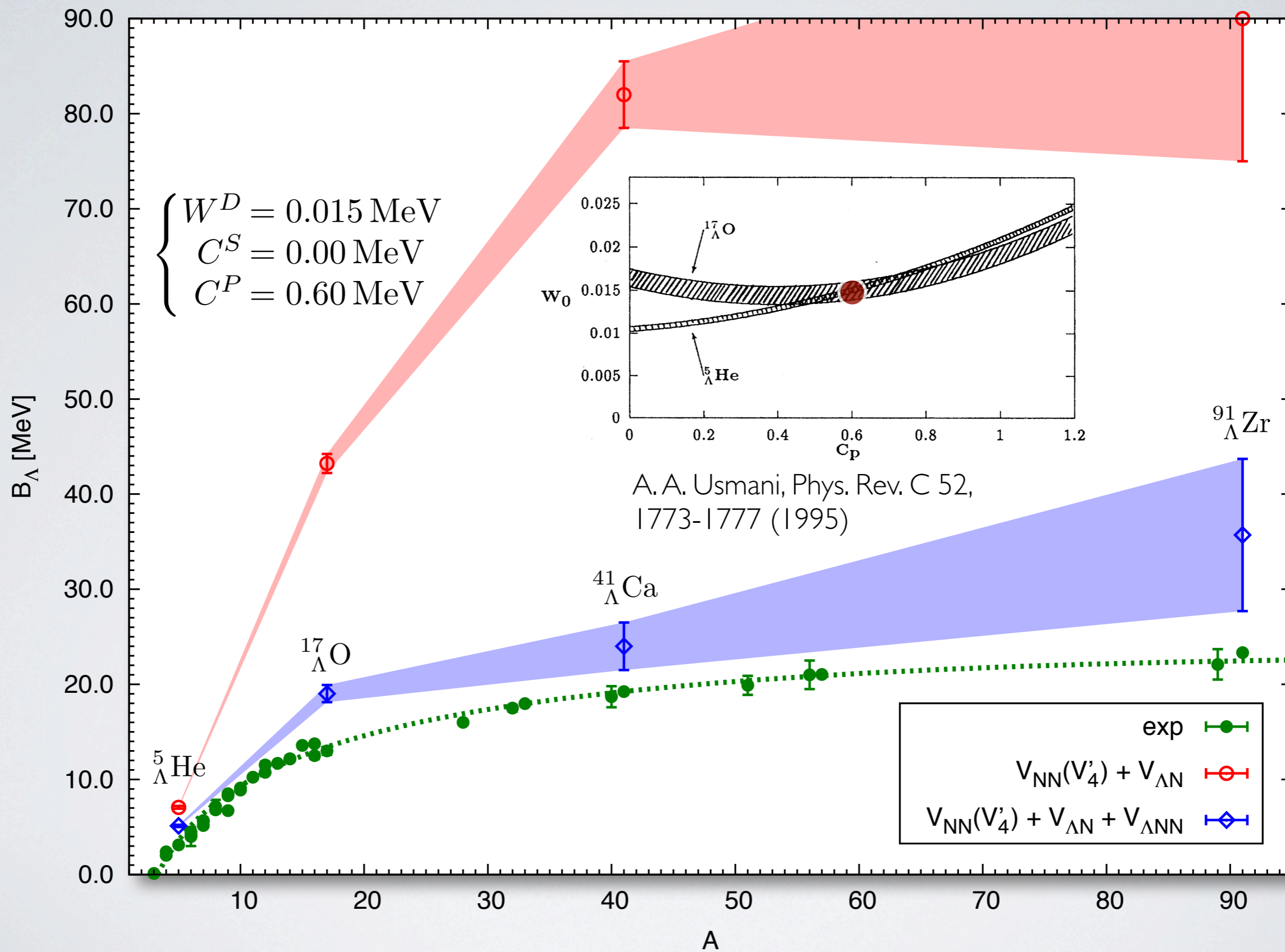
$$\vec{s}_i^N = \begin{pmatrix} a_i \\ b_i \\ c_i \\ d_i \end{pmatrix} = a_i |p \uparrow\rangle + b_i |p \downarrow\rangle + c_i |n \uparrow\rangle + d_i |n \downarrow\rangle$$

$$\vec{s}_i^\Lambda = \begin{pmatrix} u_i \\ v_i \end{pmatrix} = u_i |\uparrow\rangle + v_i |\downarrow\rangle$$

Results: Λ -hypernuclei



Results: Λ -hypernuclei



Results: Λ -hypernuclei

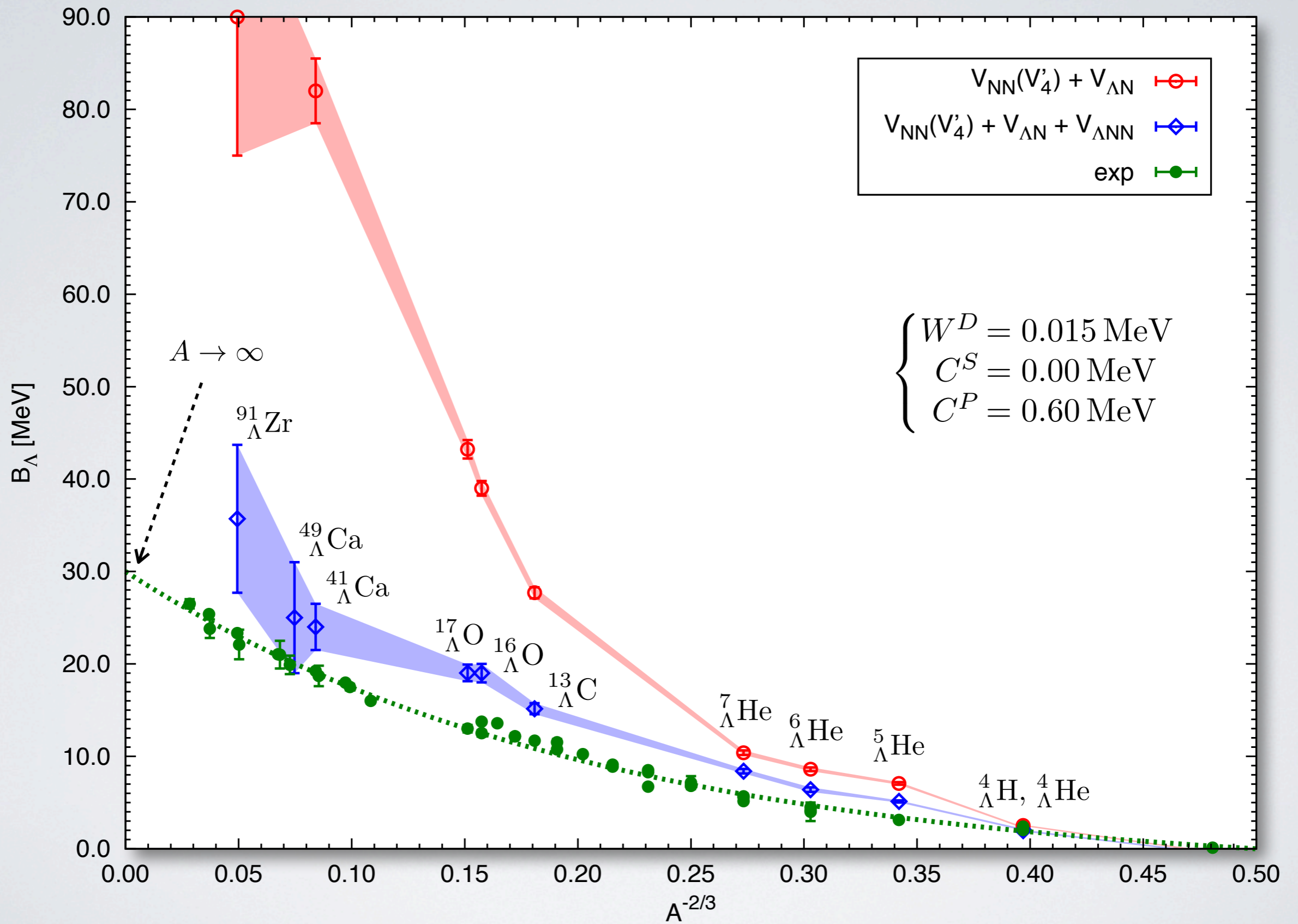
Hyp.: nuclear effects cancel at most ✓

NN potential	${}^5_{\Lambda}\text{He}$		${}^{17}_{\Lambda}\text{O}$	
	$V_{\Lambda N}$	$V_{\Lambda N} + V_{\Lambda NN}$	$V_{\Lambda N}$	$V_{\Lambda N} + V_{\Lambda NN}$
Argonne V4'	7.1(1)	5.1(1)	43(1)	19(1)
Argonne V6'	6.3(1)	5.2(1)	34(1)	21(1)
Minnesota	7.4(1)	5.2(1)	50(1)	17(2)
Expt.		3.12(2)		13.0(4)

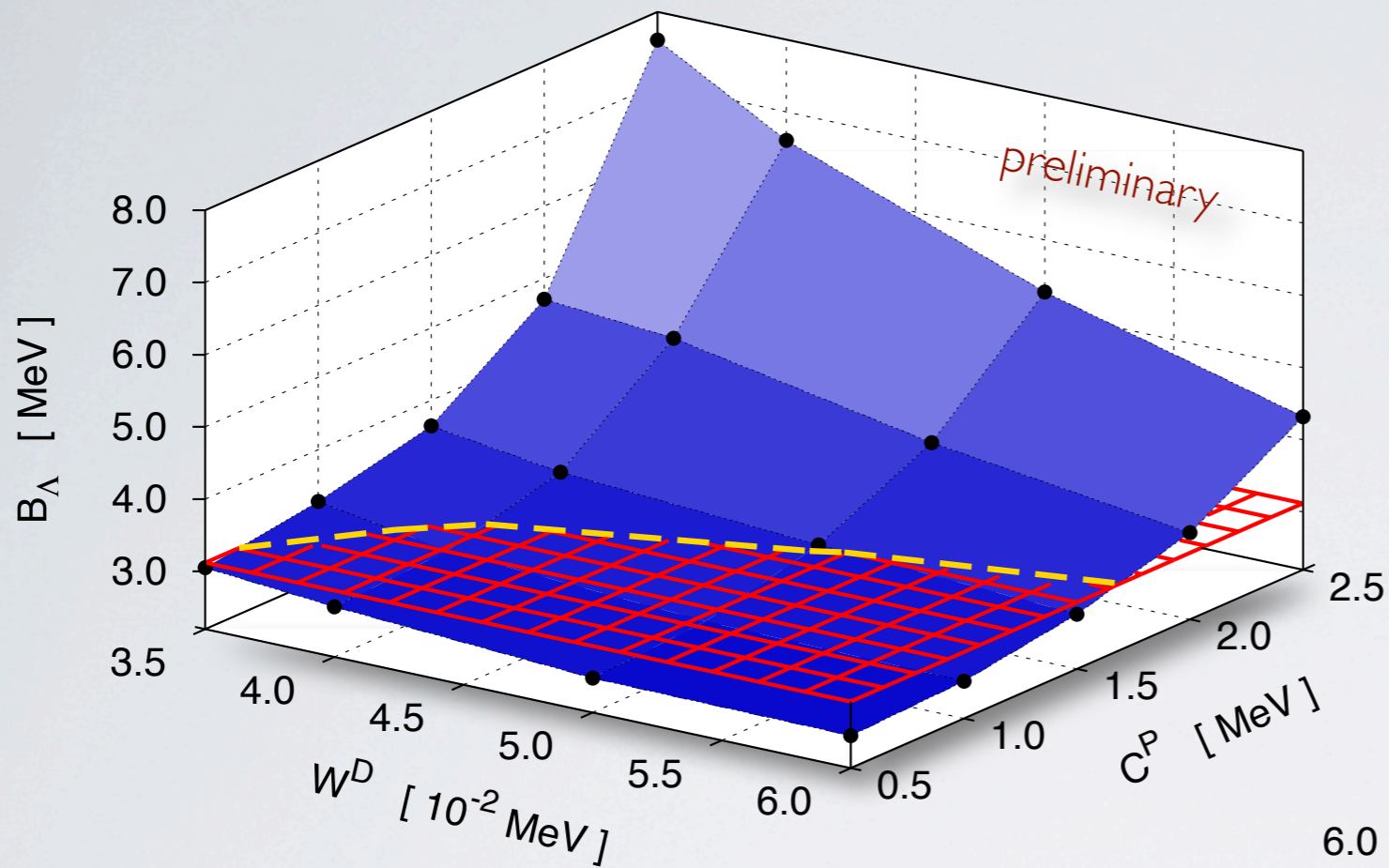
D. Lonardoni, S. Gandolfi, F. Pederiva, Phys. Rev. C 87, 041303(R) (2013)

hyperon separation energy not sensitive
to the details of nuclear interaction

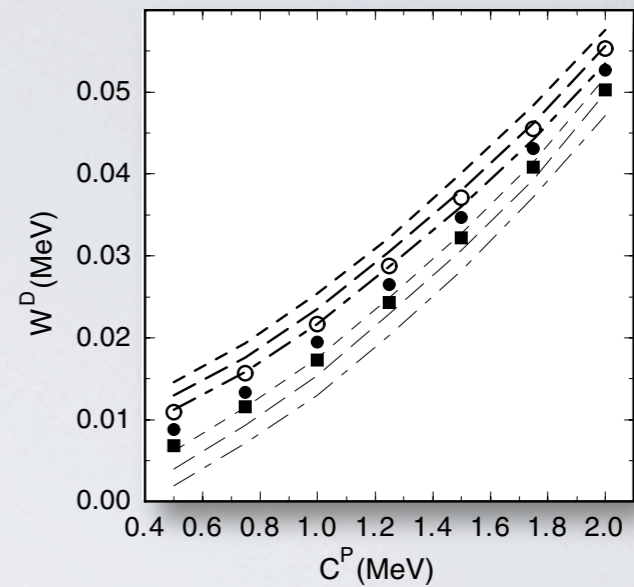
Results: Λ -hypernuclei



Results: Λ -hypernuclei



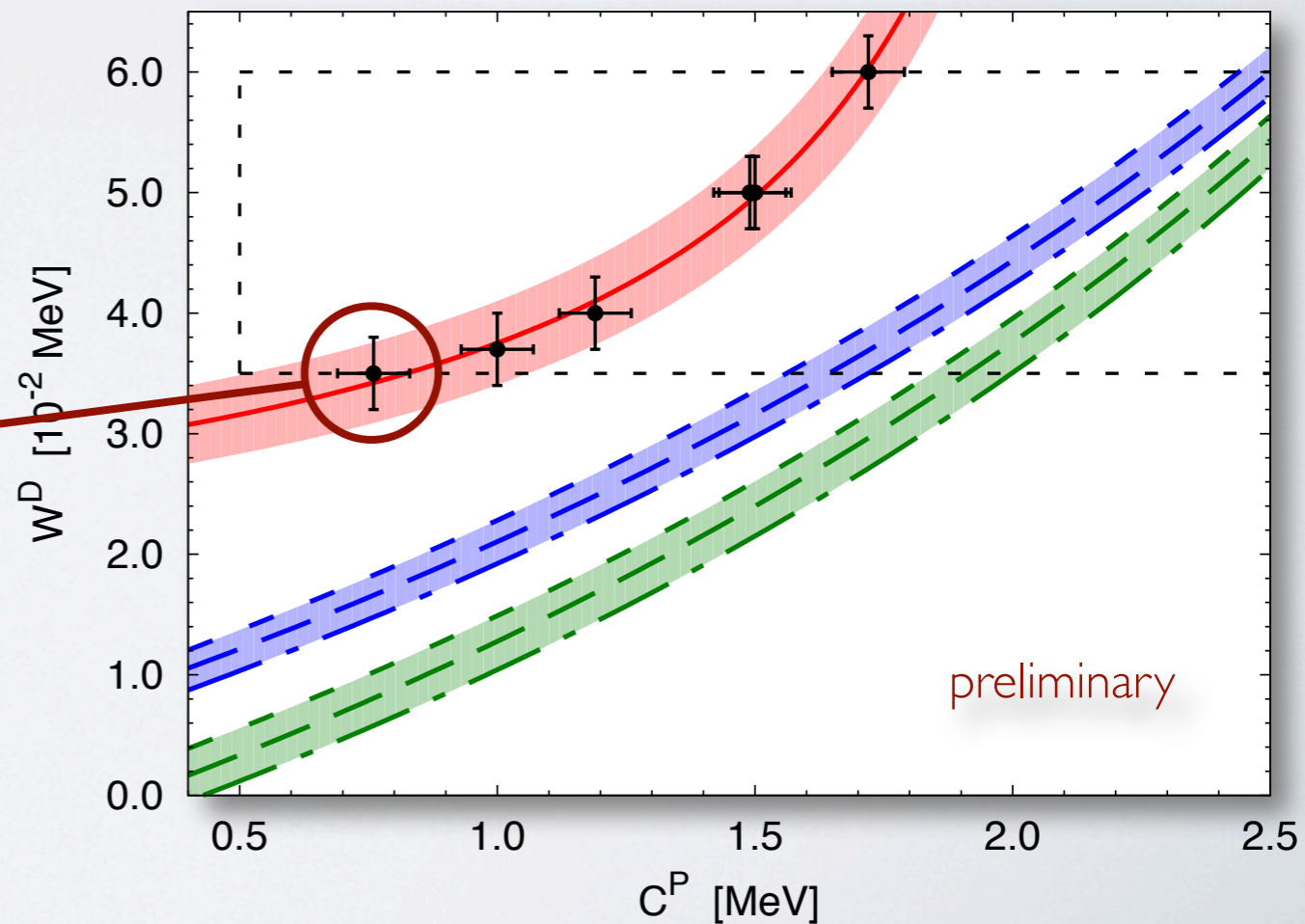
${}^5_\Lambda\text{He}$



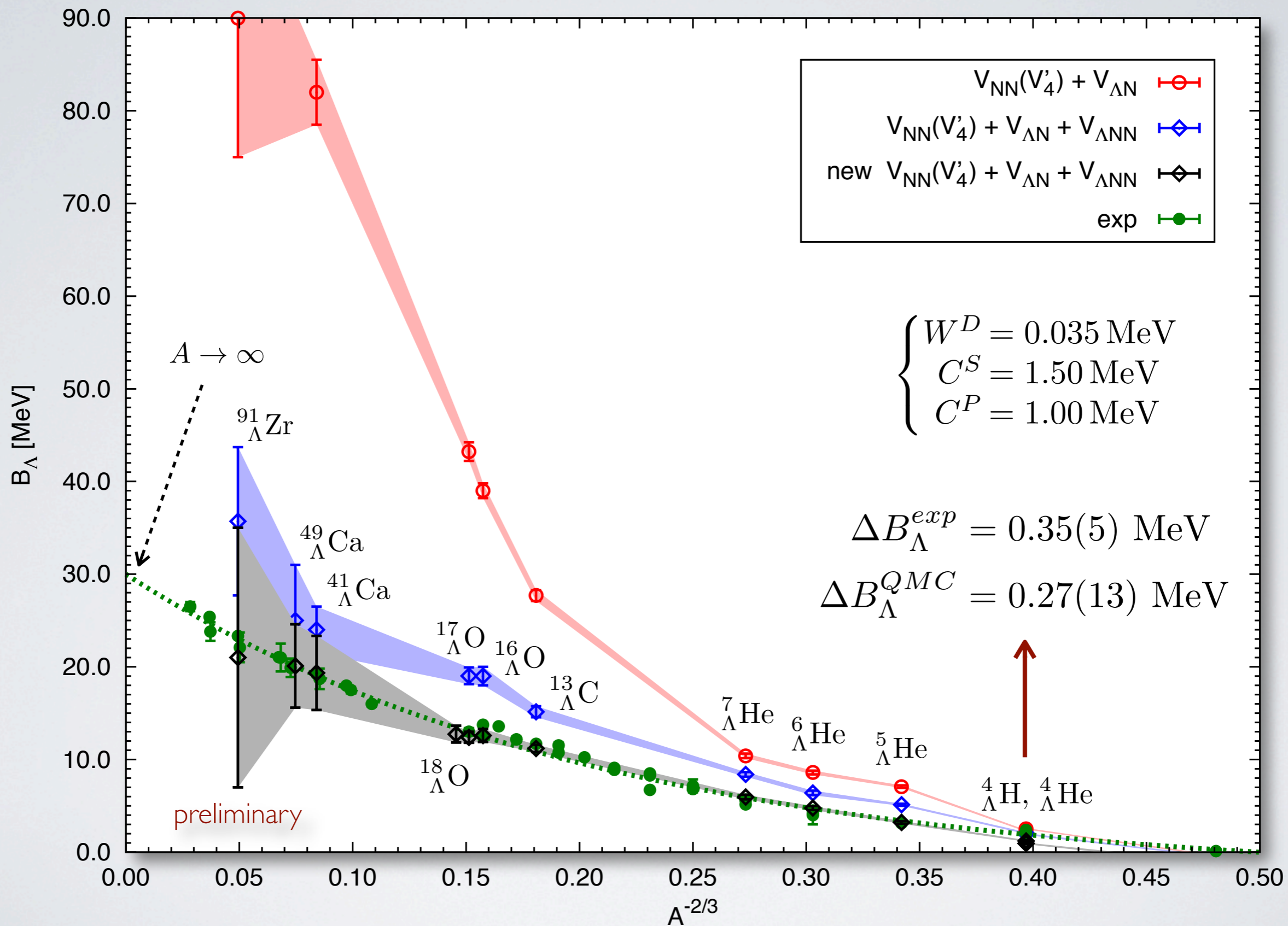
A. A. Usmani, F. C. Khanna,
J. Phys. G: Nucl. Part. Phys. 35, 025105 (2008)

$$\left\{ \begin{array}{l} W^D = 0.035 \text{ MeV} \\ C^S = 1.50 \text{ MeV} \\ C^P = 1.00 \text{ MeV} \end{array} \right.$$

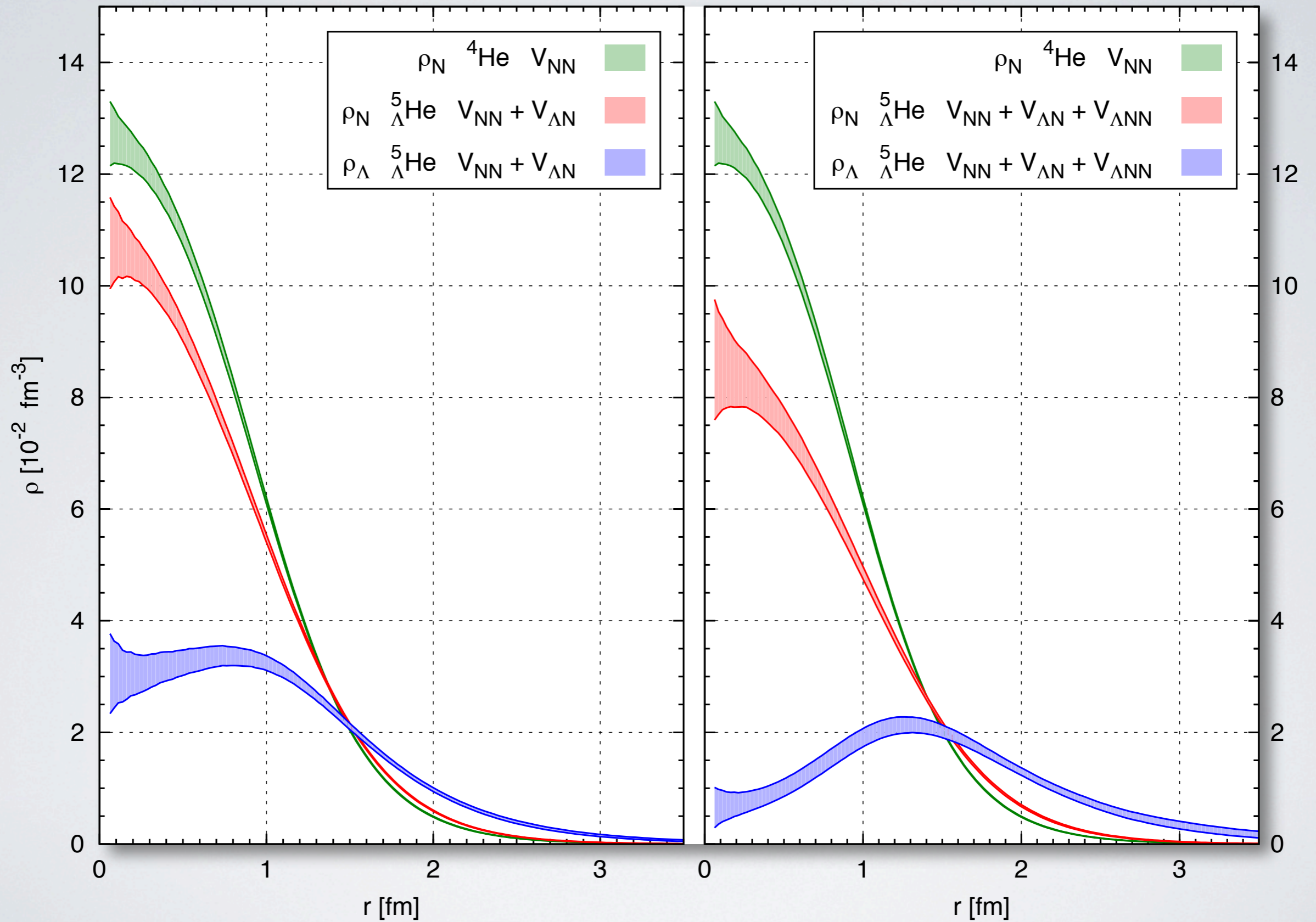
good for ${}^5_\Lambda\text{He}$ and ${}^{17}_\Lambda\text{O}$



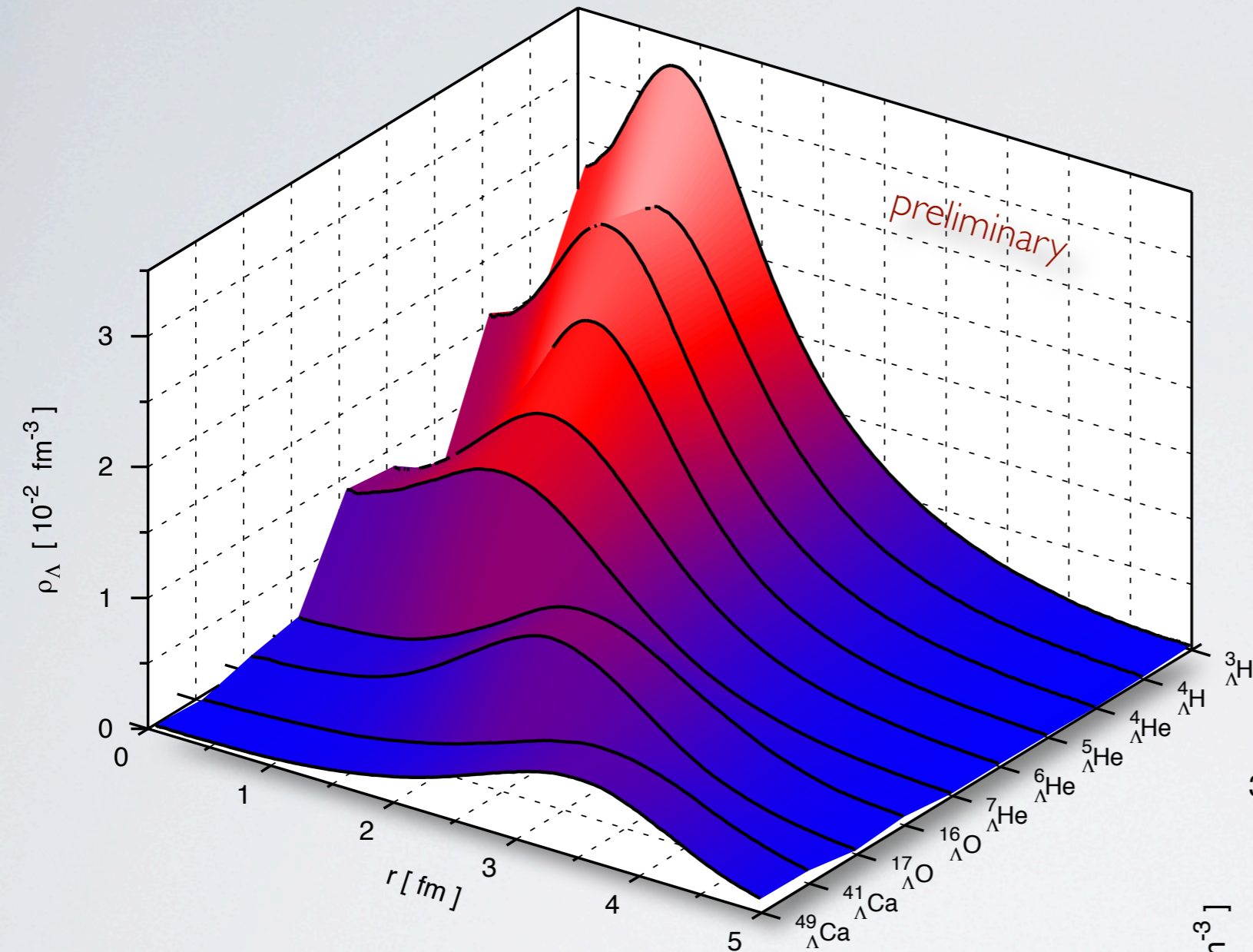
Results: Λ -hypernuclei



Results: Λ -hypernuclei

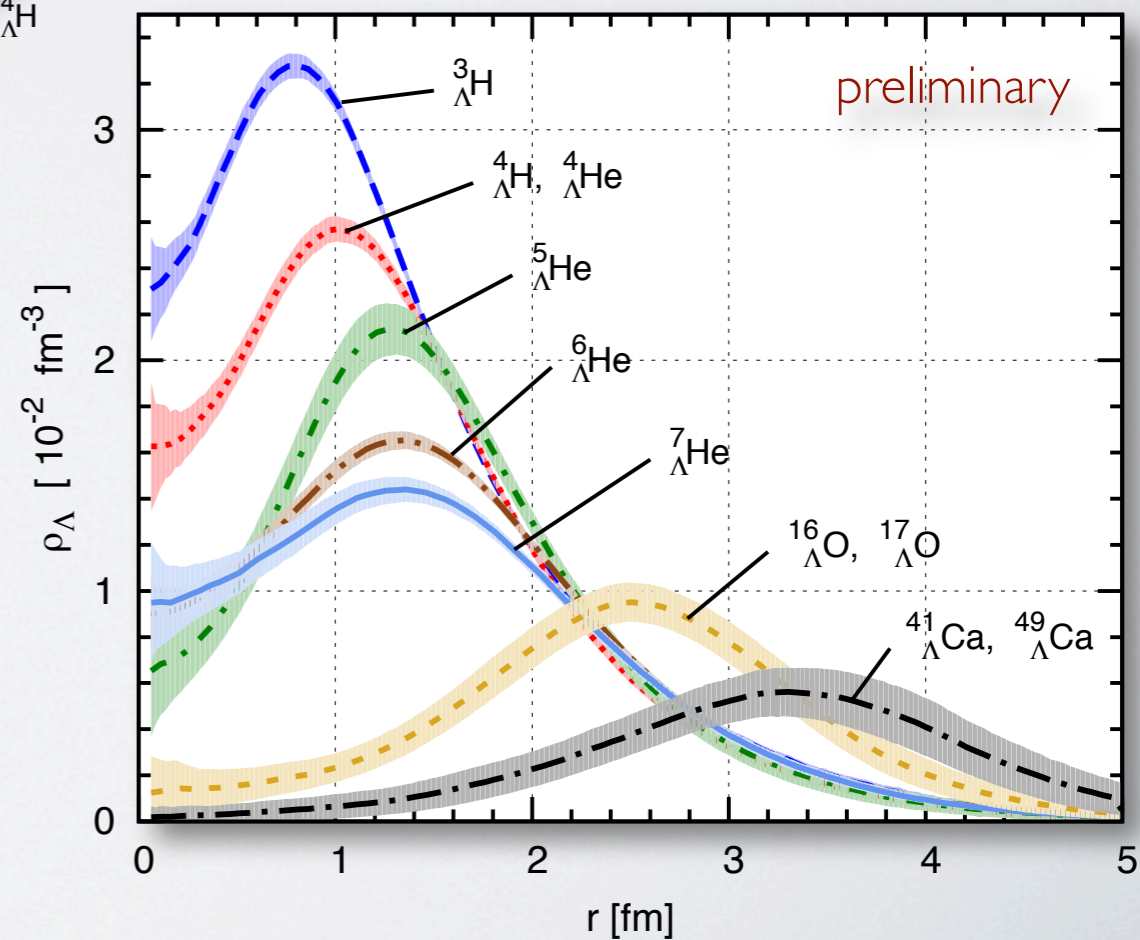


Results: Λ -hypernuclei



Λ single particle densities

significant 3-body effects



Results: Λ -neutron matter

neutron matter

$$E(n_b)$$

$\xrightarrow{\beta \text{ eq.}}$

nuclear matter

$$E(n_b, x_p)$$

Λ -neutron matter

$$E(n_b, x_\Lambda)$$

$$x_p = \frac{n_p}{n_n + n_p}$$

$$x_\Lambda = \frac{n_\Lambda}{n_n + n_\Lambda}$$

$$E(n_b, x_p) = E(n_b, \frac{1}{2}) + \Delta E$$

$$\Delta E = S(n_b)(1 - 2x_p)^2$$

?

$$\begin{cases} E_n(n_b) \\ E_{n+\Lambda}(n_b) \\ x_\Lambda \ll 1 \end{cases}$$



$$E(n_b, x_\Lambda)$$



$$\begin{cases} \varepsilon(n_b, x_\Lambda) \\ \mu_n(n_b, x_\Lambda) \\ \mu_\Lambda(n_b, x_\Lambda) \end{cases}$$



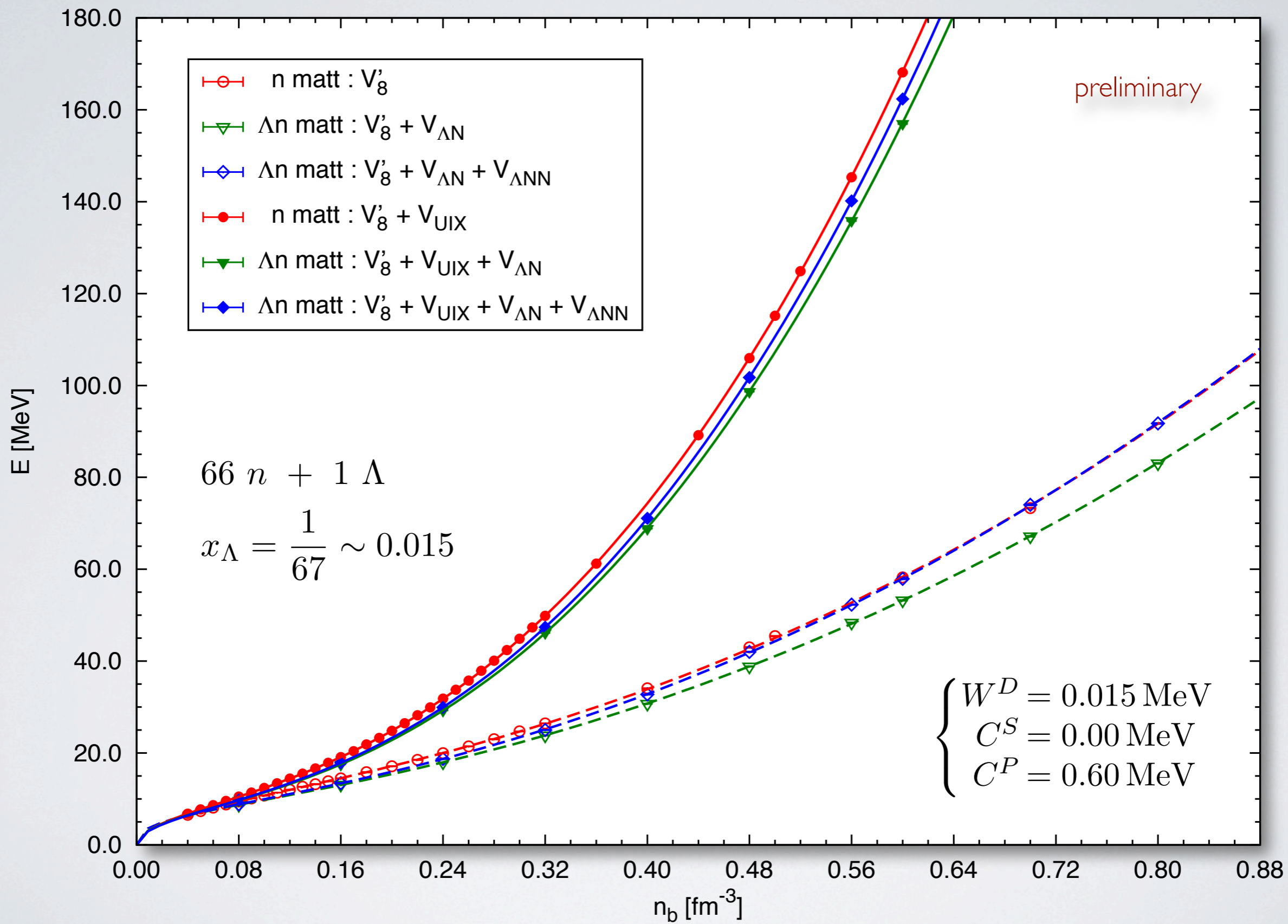
solve the TOV



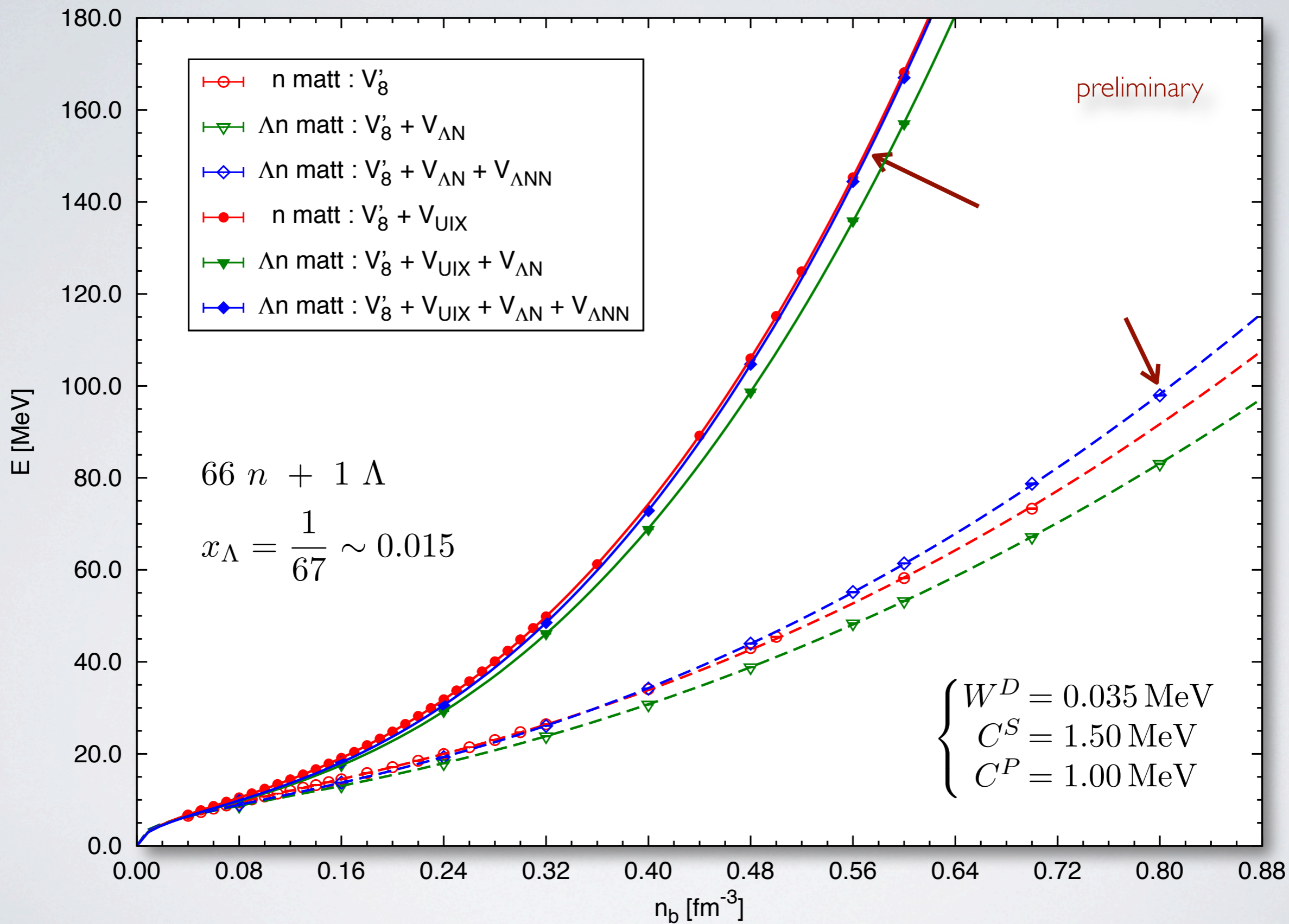
$M(R)$ & M_{\max}



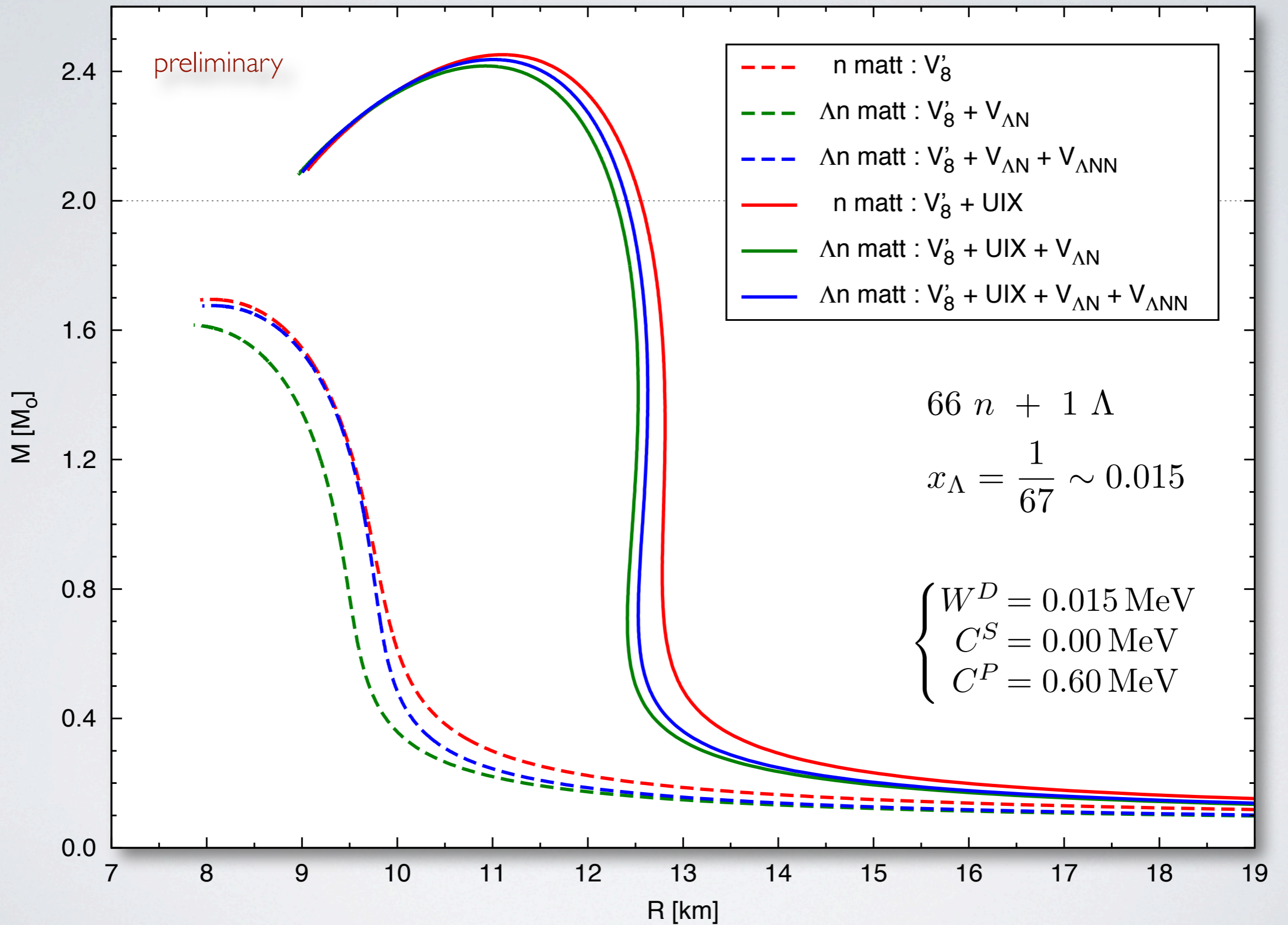
Results: Λ -neutron matter



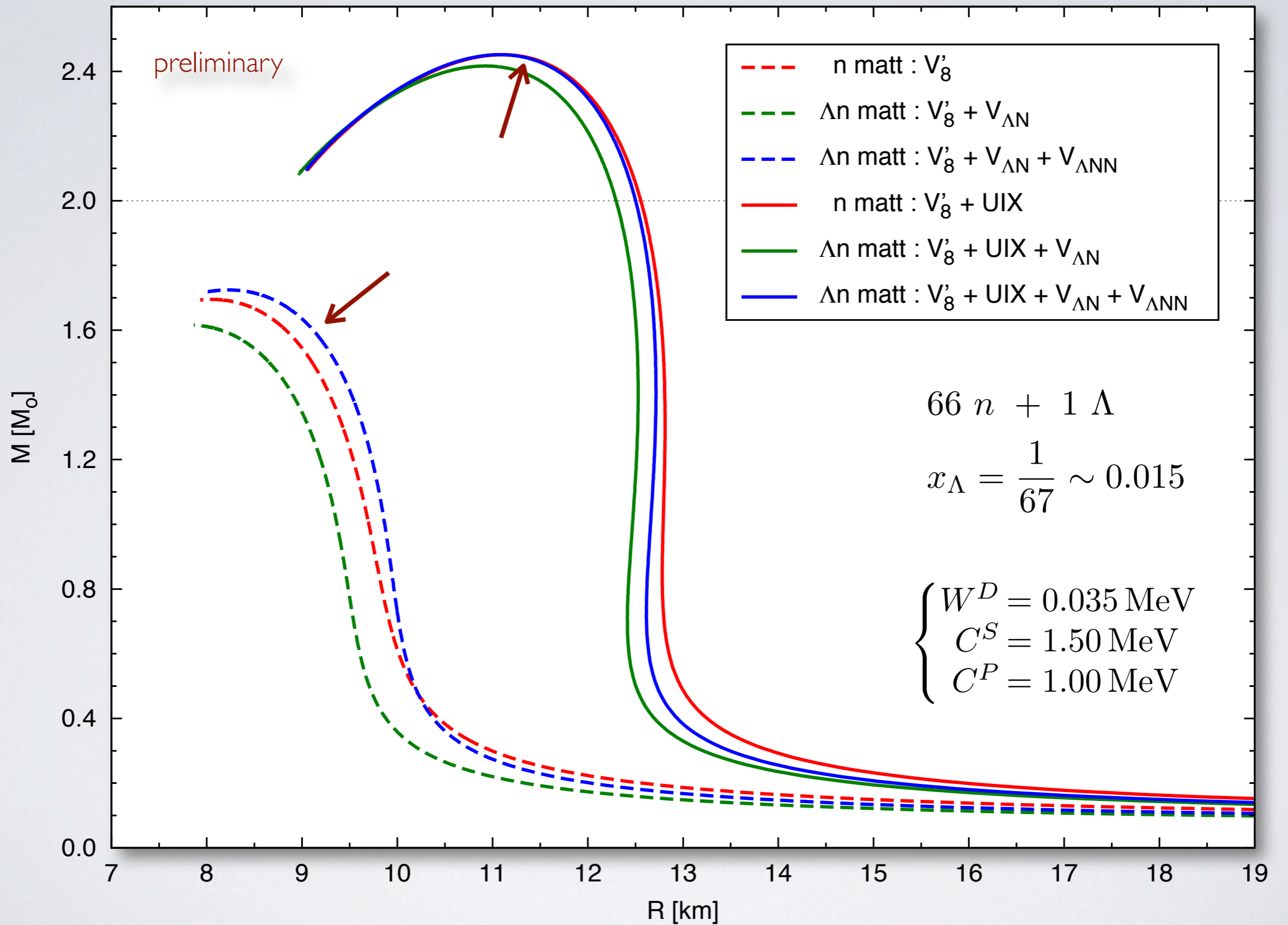
Results: Λ -neutron matter



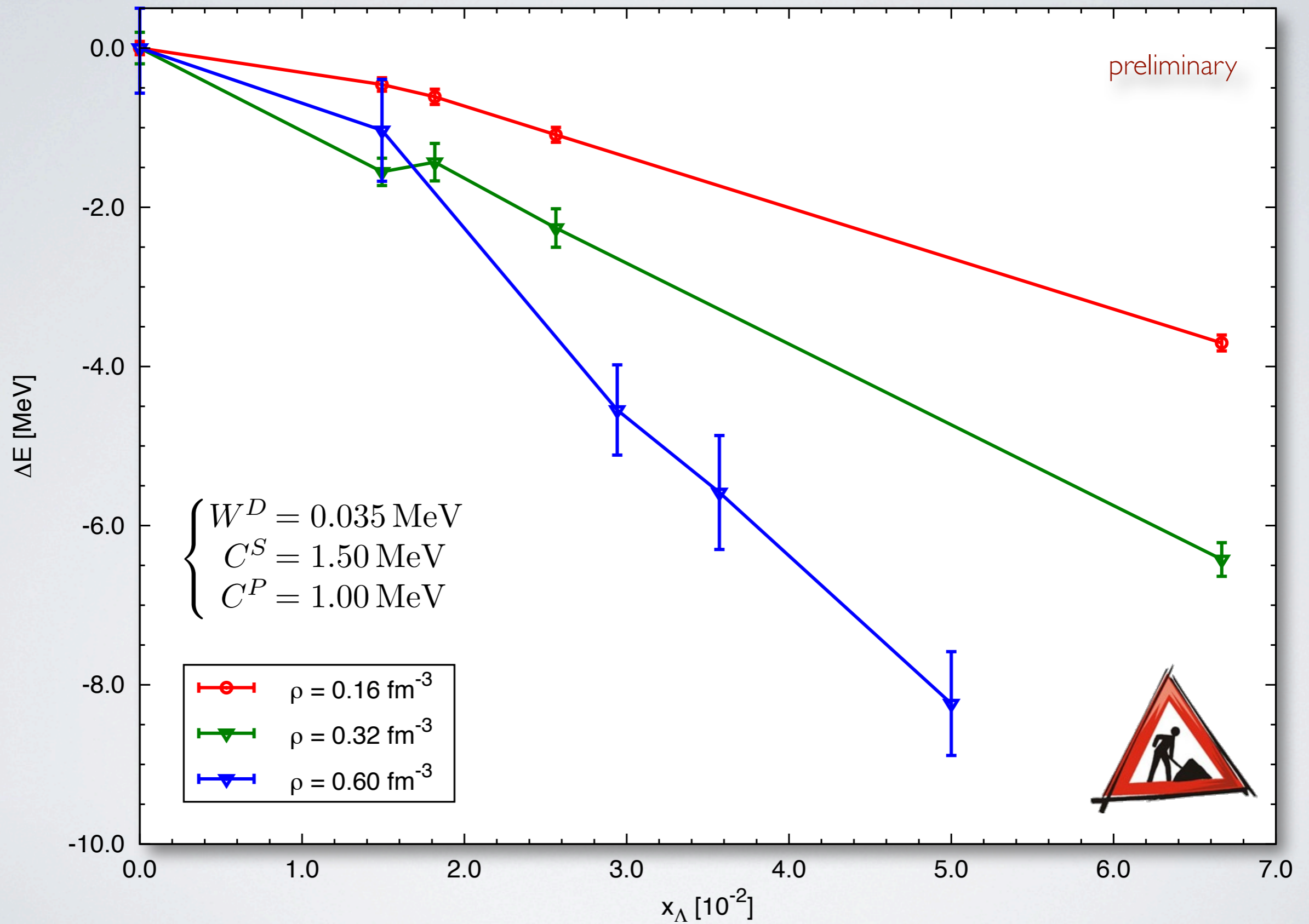
Results: Λ -neutron matter



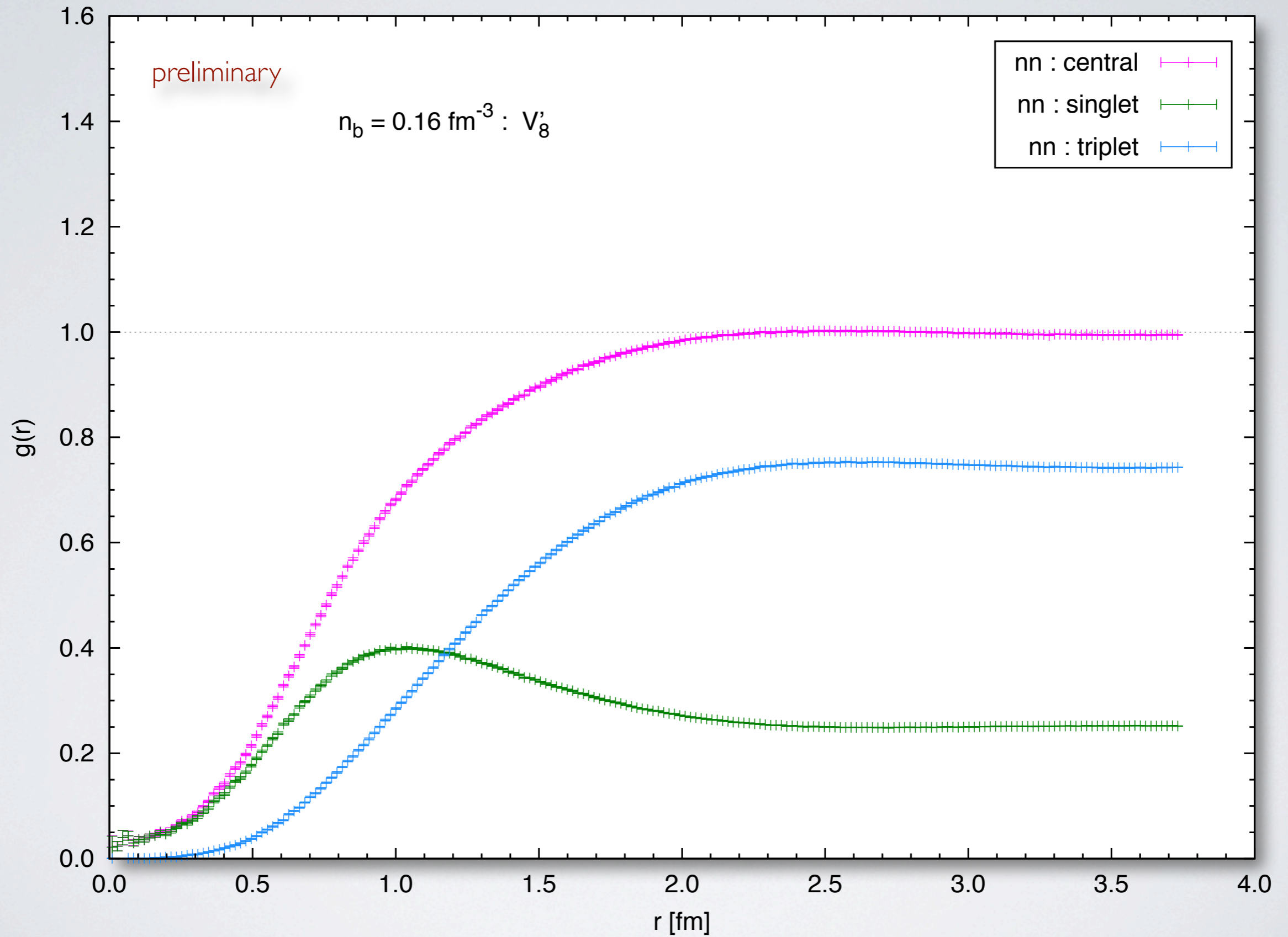
Results: Λ -neutron matter



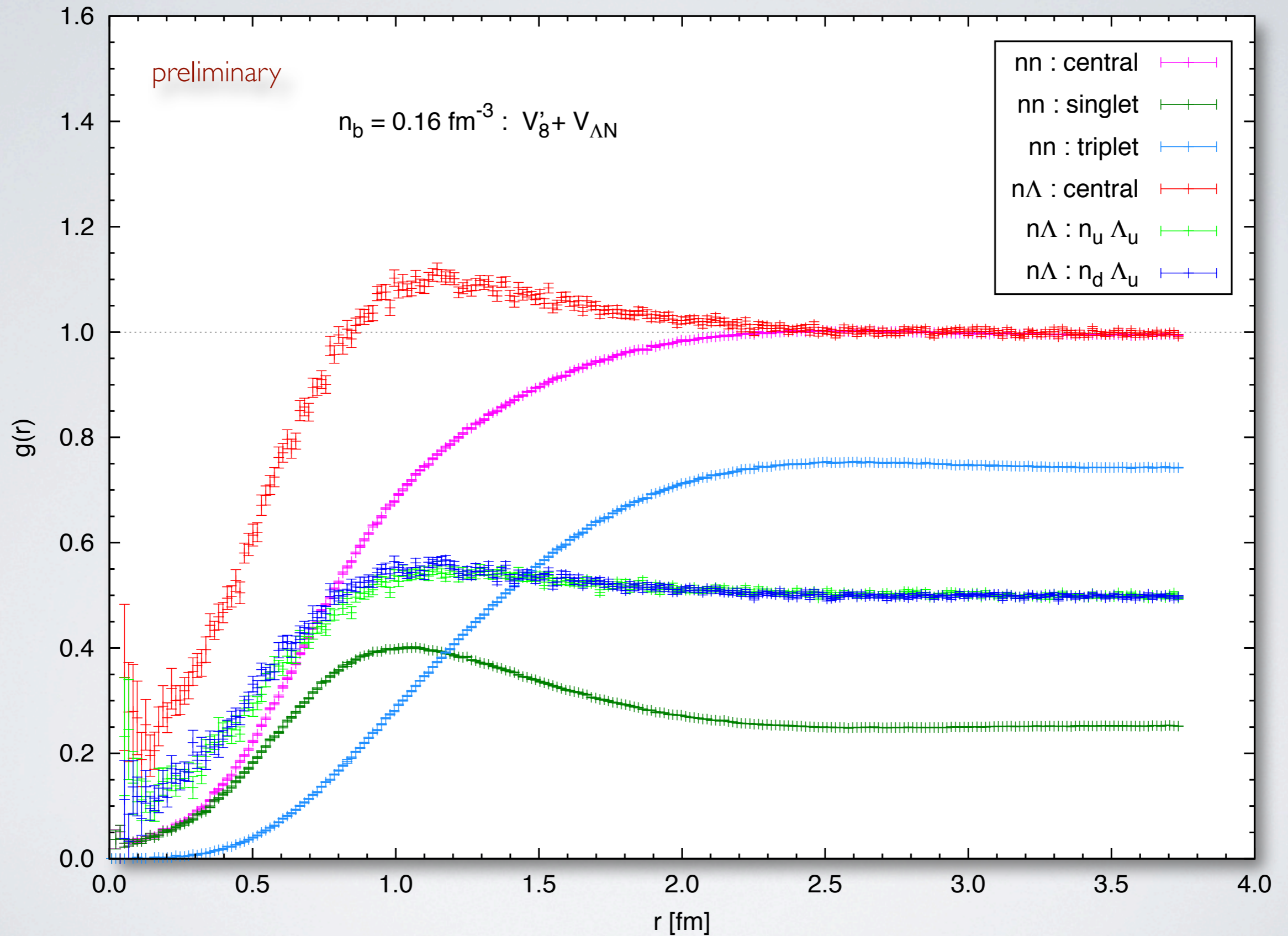
Results: Λ -neutron matter



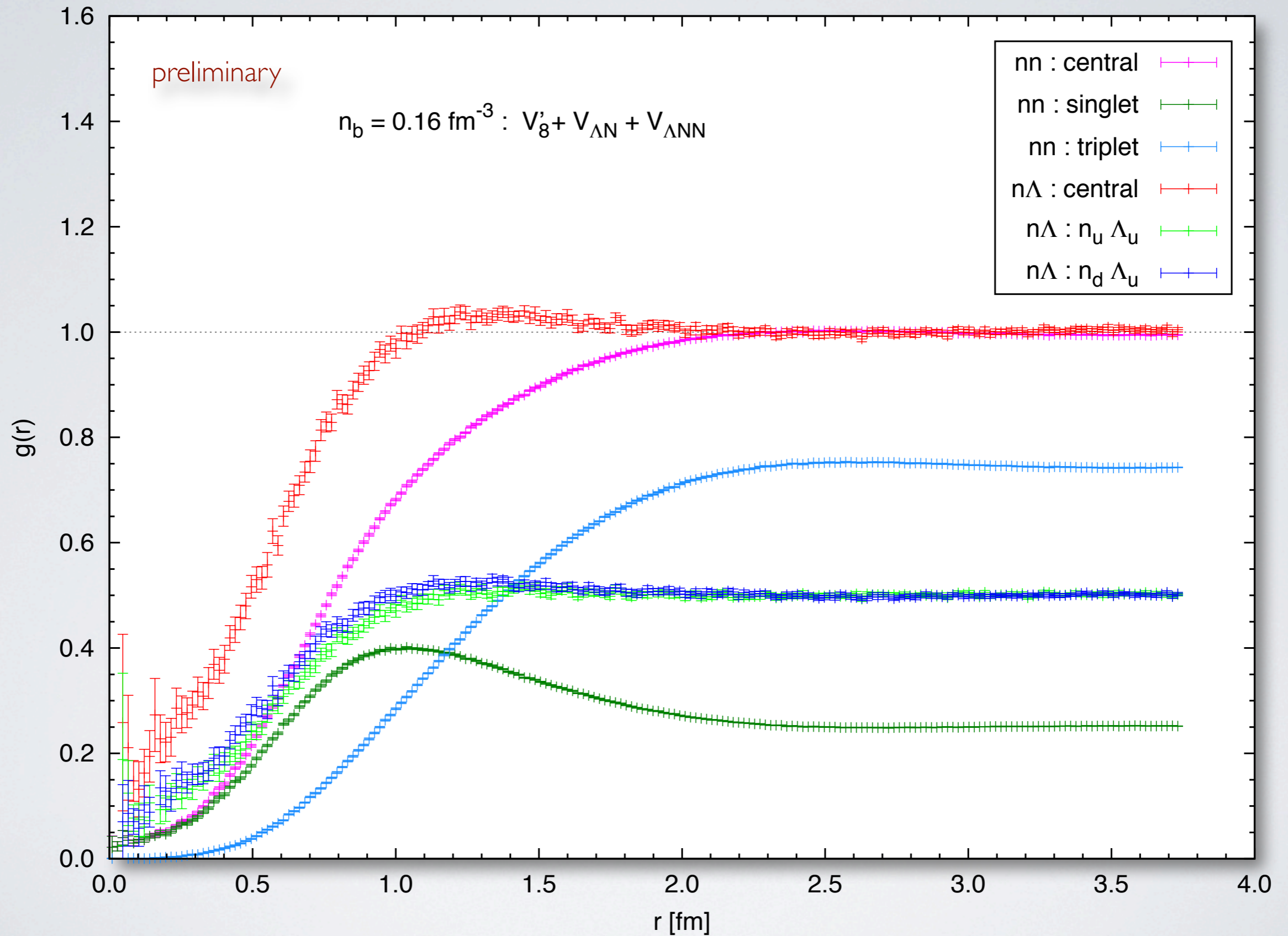
Results: Λ -neutron matter



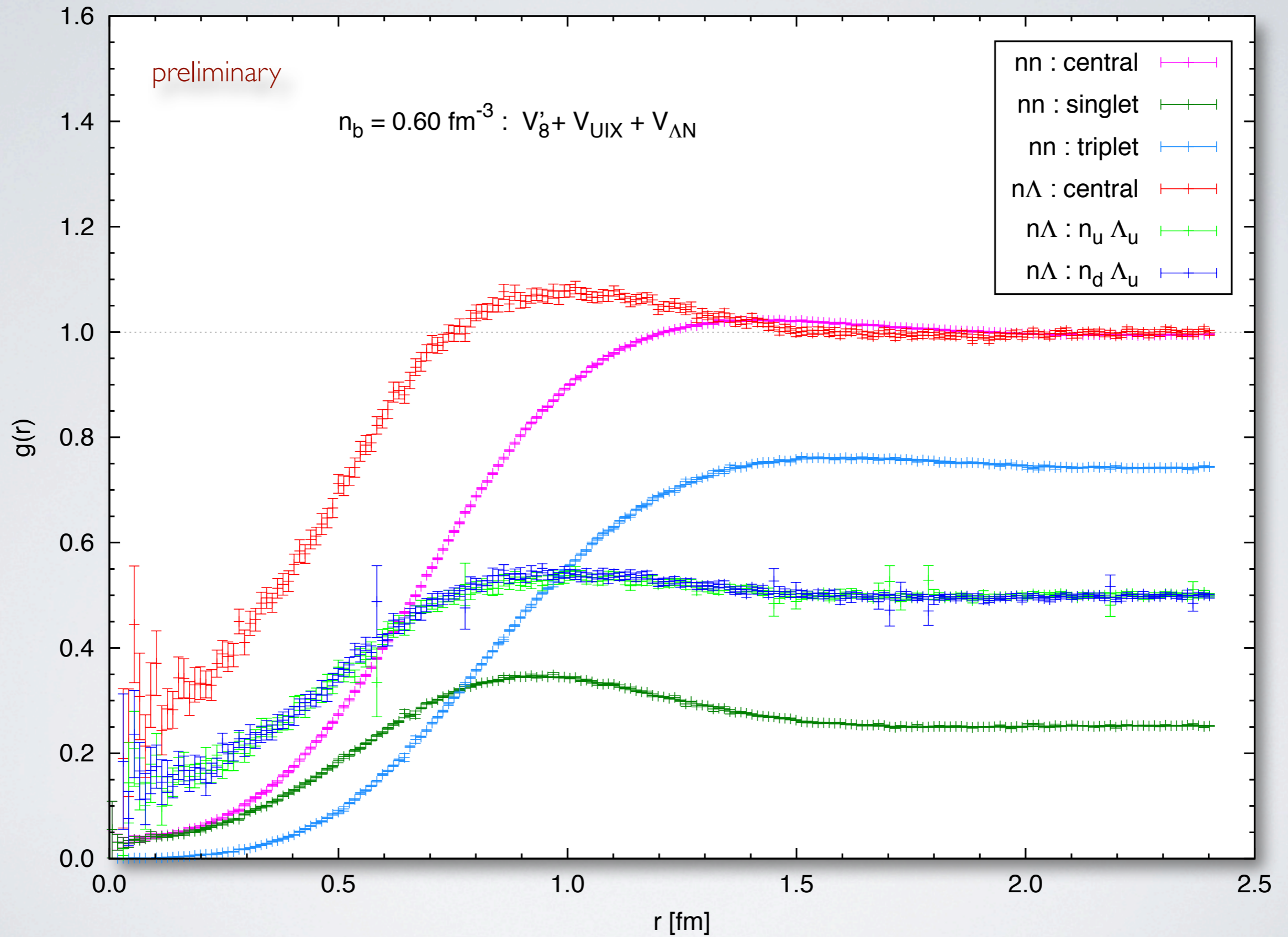
Results: Λ -neutron matter



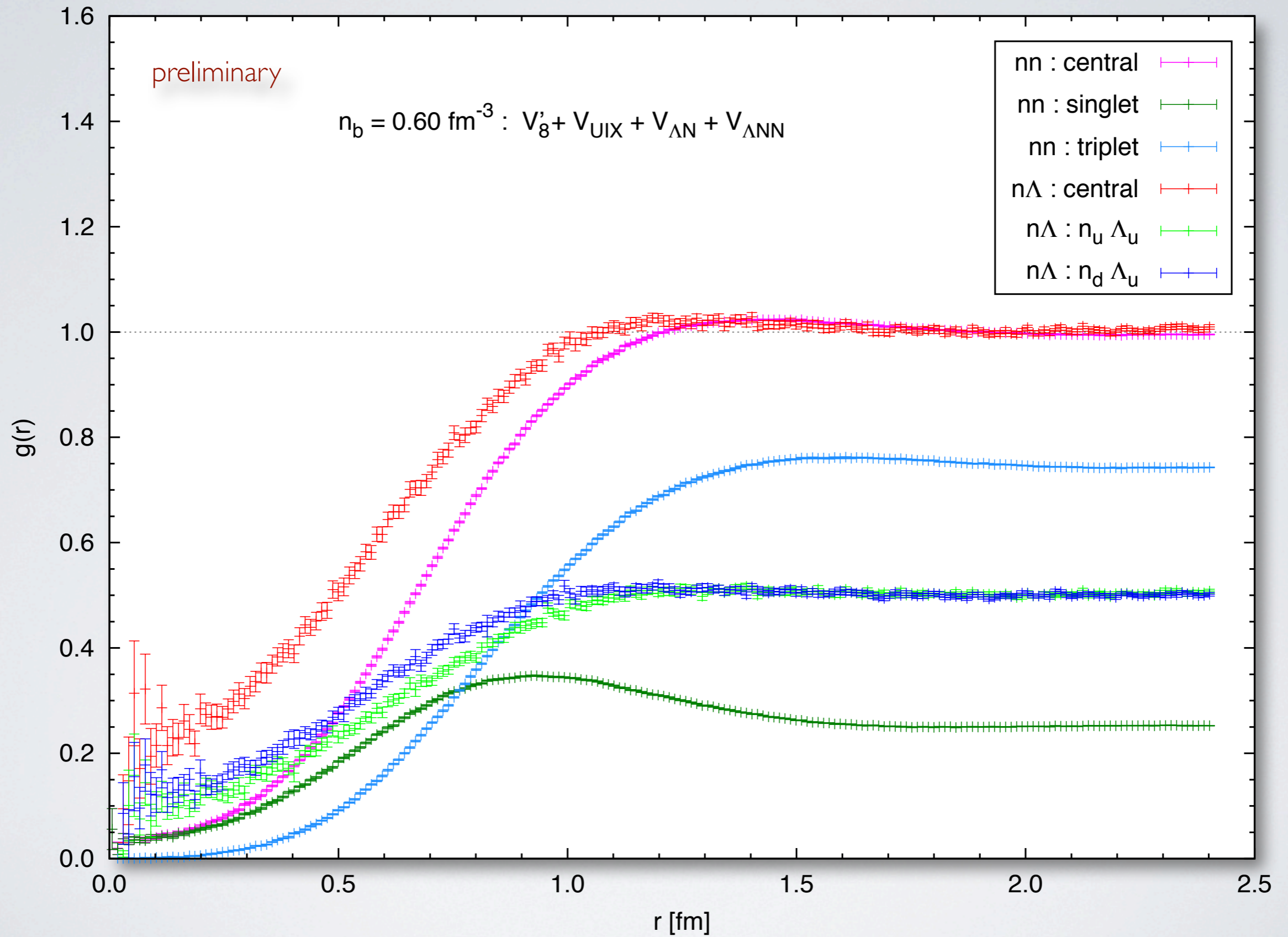
Results: Λ -neutron matter



Results: Λ -neutron matter



Results: Λ -neutron matter



Conclusions

- ✓ Extension of the AFDMC code for strange finite and infinite nuclear systems: develop of an hyperon-nucleon interaction in the Quantum Monte Carlo scheme
 - analysis of the hyperon separation energy in medium-light hypernuclei
 - analysis of the EoS of the hyperon-neutron matter at high density
- ✓ Two-body ΛN interaction not sufficient to describe the hyperon-separation energy of medium-light Λ -hypernuclei: need of a strongly repulsive three-body ΛNN interaction
- ✓ EoS for the Λ -neutron matter not too soft: chance for a NS maximum mass up to $2 M_{\odot}$ even in presence of hyperons

work in progress



*Thank you for
your attention !!*

Backup: the hyperon-nucleon interaction

$$T_\pi(x) = \left(1 + \frac{3}{m_\pi x} + \frac{3}{(m_\pi x)^2} \right) Y_\pi(x) \xi(x)$$

$$Y_\pi(x) = \frac{e^{-m_\pi x}}{m_\pi x} \xi(x) \quad \xi(x) = 1 - e^{-cx^2}$$

$$Z_\pi(x) = \frac{x}{3} [Y_\pi(x) - T_\pi(x)]$$

$$X_{\Lambda i} = (\boldsymbol{\sigma}_\Lambda \cdot \boldsymbol{\sigma}_i) Y_\pi(r_{\Lambda i}) + S_{\Lambda i} Y_\pi(r_{\Lambda i})$$

$$S_{\Lambda i} = 3 (\boldsymbol{\sigma}_\Lambda \cdot \hat{\boldsymbol{r}}_{\Lambda i}) (\boldsymbol{\sigma}_i \cdot \hat{\boldsymbol{r}}_{\Lambda i}) - \boldsymbol{\sigma}_\Lambda \cdot \boldsymbol{\sigma}_i$$

constant	value	unit
m_π	138.03899	MeV
W_c	2137	MeV
\bar{r}	0.5	fm
a	0.2	fm
v_s	6.33, 6.28, 6.23	MeV
v_t	6.09, 6.04, 5.99	MeV
\bar{v}	6.15(5)	MeV
v_σ	0.24	MeV
c	2.0	fm ⁻²
ε	0.1 ÷ 0.38	–
W^D	0.002 ÷ 0.058	MeV
C^P	0.5 ÷ 2.5	MeV
C^S	~ 1.5	MeV
C^{CSB}	-0.050(5)	MeV

Backup: the hyperon-hyperon interaction

$$v_{\lambda\mu} = \sum_{k=1}^3 \left(v_0^{(k)} + v_{\sigma}^{(k)} \boldsymbol{\sigma}_{\lambda} \cdot \boldsymbol{\sigma}_{\mu} \right) e^{-\mu^{(k)} r_{\lambda\mu}^2}$$

$\mu^{(k)}$	0.555	1.656	8.163
$v_0^{(k)}$	-10.67	-93.51	4884
$v_{\sigma}^{(k)}$	0.0966	16.08	915.8

E. Hiyama, M. Kamimura, T. Motoba, T. Yamada, Y. Yamamoto,
Phys. Rev. C 66, 024007 (2002)

$$V_{\Lambda\Lambda} = \frac{1}{2} \sum_{\lambda \neq \mu} \sum_{\alpha} \sigma_{\lambda}^{\alpha} \left(\mathcal{D}_{\lambda\mu}^{[\sigma]} \right) \sigma_{\mu}^{\alpha}$$