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Shell evolution and pairing in calcium isotopes with two- and three-body forces

Javier Menéndez

Institut für Kernphysik, TU Darmstadt ExtreMe Matter Institute (EMMI)

with Jason D. Holt, Achim Schwenk and Johannes Simonis

Computational and Theoretical Advances for Exotic Isotopes in the Medium Mass Region

INT Seattle, 18 April 2013

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1 [Ca ground states: pairing, shell evolution](#page-7-0)

2 [Ca excited states: shell evolution, spectra, transitions](#page-17-0)

3 [Residual 3N forces: oxygen and calcium isotopes](#page-28-0)

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The nuclear many-body approach

Big variety of nuclei in the nuclear chart, $A \sim 2...300$

Exact (*ab initio*) calculations only possible in the lightest nuclei

Poses a hard many-body problem: design approximate methods valid in different regions

In the rest of the nuclear chart, pick only the (more) relevant degrees of freedom for observable to study: idea of the Interacting Shell Model

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Medium-mass nuclei: standard view

Standard studies of medium-mass nuclei (*A* ∼ 20 − 80) are performed with theoretical approaches based on phenomenology:

- **•** Shell Model calculations use fitted interactions or modified G-matrices
- **Energy Density Functional interactions** use Skyrme, Gogny or Relativistic fitted interactions

Interactions are made to reproduce experiment for stable nuclei

- When extrapolated to exotic (neutron rich) regions results differ \Rightarrow Need to guide (ideally avoid) fits!
- Why microscopic NN interactions have to be modified? \Rightarrow Need to include 3N forces explicitly!

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Towards microscopic medium-mass nuclei calculation

Microscopic calculation of medium-mass nuclei including 3N forces

- Use Chiral Effective Field Theory (chiral EFT) interactions, includes naturally NN and 3N forces.
- Perform a renormalization group evolution to *Vlowk* interaction to enhance convergence of the MBPT calculation
- Apply Many-Body Perturbation Theory (MBPT) to 3rd order to obtain interactions to be used in Shell Model (SM) calculations
- Full diagonalizations using codes ANTOINE and NATHAN Caurier et al. RMP77 427(2005)
- \Rightarrow 3N forces are naturally included Shown necessary to reproduce light nuclei spectra
- \Rightarrow All the parameters that appear in the SM hamiltonian calculated from the input of the microscopic [in](#page-3-0)t[er](#page-5-0)[a](#page-3-0)[cti](#page-4-0)[o](#page-5-0)[n](#page-0-0) [\(](#page-0-0)[n](#page-6-0)[o](#page-7-0) [fit](#page-0-0)[s](#page-6-0)[!\)](#page-7-0)

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Chiral EFT NN+3N Forces

Systematic expansion: state-of-the-art chiral EFT forces

NN forces included up to N³LO

$3N$ forces included up to N^2LO

Weinberg, van Kolck, Kaplan, Savage, Wise, Epel[ba](#page-4-0)u[m,](#page-6-0) [K](#page-4-0)[ai](#page-5-0)[se](#page-6-0)[r,](#page-0-0) [M](#page-6-0)[ei](#page-7-0)[ßn](#page-0-0)[er](#page-6-0)[..](#page-7-0)[.](#page-0-0)

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Normal-ordered 3N Forces

Approximate treatment of 3N forces:

normal-ordered 2B: 2 valence, 1 core particle \Rightarrow (effective) Two-body Matrix Elements (TBME)

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normal-ordered 1B: 1 valence, 2 core particles \Rightarrow (effective) Single particle energies (SPE)

$$
\textcolor{red}{\text{H} \textcolor{blue}{\text{H} \times \text{H} \times \text{H}}}
$$

residual 3B:

 \Rightarrow Estimated to be suppressed by $N_{valence}/N_{core}$

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Ca isotopes: Ground-state energies

Ca calculations with respect to ⁴⁰Ca core

 $\hbar\omega = 11.48$ MeV, appropriate for Ca radii

Normal-ordered 3N force contributions to TBMEs and SPEs

3N forces provide crucial repulsion (similar to O case)

NN+3N calculations nice agreement with experimental data

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Ca isotopes: Ground-state energies

 N_{NN} Results particularly sensitive to SPEs,

NN+3N (emp) \rightarrow consider two sets:

NN+3N (MBPT) MRPT (calculated from NN+3N forces especially more neutron-rich systems, consider two sets: MBPT (calculated from NN+3N forces) Empirical (from GXPF1 interaction) to have an estimate of this uncertainty

> Calculation performed in *pf g*_{9/2} valence space

Flat behaviour towards ⁶⁰Ca does not allow clear prediction of the dripline: enlarge valence space include continuum degrees of freedom

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Nuclear Pairing Gaps

- Pairing correlations are important in nuclei
	- Odd-even mass staggering
	- Moments of Inertia
	- \bullet ...
- Theoretical pairing gaps compared to experiment via the three point mass formula:

$$
\Delta_N^3 = (-1)^N \frac{BE (N-1) + BE (N+1) - 2BE (N-1)}{2}
$$

Holt, JM, Schwenk, arXiv:1304.0434

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MBPT convergence

- **•** Succesive orders *pp*, *hh* ladder diagrams build up pairing gaps
- At third order *pp*, *hh* ladders results seem to be converged
- Third order ladders results still away from experiment Incorrect even-odd staggering (too attractive mean-field)

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Pairing, 3N forces and EDF

- At *pp*-*hh* level 3N forces reduce the pairing gaps in 200-500 keV first observed in EDF: Lesinski et al. JPhysG39 015108 (2012)
- EDF Δ_N^3 slightly closer to experiment, and more staggering No indication of shell closures

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Nuclear Pairing Gaps: NN vs NN+3N

- Full third order MBPT improves agreement with experiment
- Core-polarization effects significantly enhance pairing gaps
- Good agreement with experimental trends

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Pairing gaps and seniority

- **•** Perform seniority truncation (limited number broken $J = 0$ pairs)
- Lowest order (no broken pairs) already good approximation apart from mid-shell nuclei
- 2 broken pairs needed in mid-shell, especially when including $g_{9/2}$ orbital

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 $\Delta_n^{(3)}$ can also tell us about shell evolution

The experimental trend is very well reproduced by NN+3N forces

Theoretical results systematically 0.5 MeV higher than experiment

Shell closure at $N = 28$, sub-shell closure $N = 32$ (moderate) and no apparent $N = 34$ closure

NN forces fail to reproduce $N = 28$ or $N = 32$ predict closed $N = 34$

Phenomenological interactions also describe [dat](#page-13-0)[a v](#page-15-0)[e](#page-13-0)[ry](#page-14-0)[w](#page-6-0)[e](#page-7-0)[ll](#page-16-0) Ω

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Two-neutron separation energies

Compare S_{2n} theoretical calculations with experimental results

$$
S_{2n} = -[B(N,Z) - B(N-2,Z)]
$$

NN S_{2n} too flat: No $N = 28$ or $N = 32$ closures Points to $N = 34$ closure

When NN+3N forces are included Very good agreement with experiment

 $NN+3N$ point to $N = 28$ or $N = 32$ closures and no $N = 34$

53,54 Ca masses very recently measured at ISOLDE Very good agreement with NN+3N Nature, in pr[int](#page-15-0) ([20](#page-17-0)[1](#page-15-0)[3\)](#page-16-0) Ω

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Shell closures and 2^+_1 energies

- Correct closure at $N = 28$ when 3N forces are included
- \bullet 3N forces enhance closure at $N = 32$
- 3N forces reduce strong closure at $N = 34$ (1.7-2.2 MeV)

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Shell closures and 2^+_1 energies

Similar results from Coupled-Cluster

Hagen et al. PRL109 032502 (2012)

⁵⁴Ca sensitive to 3N interaction:

Difference 1st/3rd MBPT 3N forces

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Effective SPEs and shell evolution

Calculate effective SPEs (monopoles)

Closed-shells $N = 32$ and $N = 34$ with NN and NN+3N forces guided by ESPEs

Neither ESPEs shows closed $N = 28$

Different behaviour as phenomenological inter[act](#page-18-0)i[on](#page-20-0)[s](#page-18-0).

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Effective SPEs and shell evolution

Calculate effective SPEs (monopoles)

Closed-shells $N = 32$ and $N = 34$ with NN and NN+3N forces guided by ESPEs

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Different behaviour as phenomenological inter[act](#page-19-0)i[on](#page-21-0)[s](#page-19-0)

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Challenge: Doubly-closed nucleus ⁴⁸Ca

Spectra too compressed with NN forces only or *pf* space

2⁺ state only ∼appropriate energy in *pfg*⁹/² NN+3N calculation

 0^+_1 state too low (1st excited state) especially compared to phenomenological interactions

Importance of 3N forces

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Importance of including *g*⁹/² orbit

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B(M1) Transition in ⁴⁸Ca

B(M1) strength in ⁴⁸Ca too fragmented in *pf* space Phenomenological calculations reproduce experimental concentration

In the extended $pfg_{9/2}$ space NN forces also fragmented strength

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NN+3N calculation in *pfg*⁹/² very good agreement with experiment

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49Ca, ⁵¹Ca neutron rich spectra

Spectrum compressed unless *pfg*_{9/2} NN+3N

Correct $(1/2)^-$ energy, (but too low $(5/2)^-$), possibility to assign experimental spins

Similar quality comparable to phenomenological interactions

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Spectroscopic factors

Have a look at the spectroscopic factors (Shell Model calculation)

Occupancy of the $f_{7/2}$ orbital lower than phenomenological models (and naively expected)

Extended $g_{9/2}$ orbital takes [miss](#page-23-0)[in](#page-25-0)[g](#page-23-0) [oc](#page-24-0)[c](#page-25-0)[u](#page-16-0)[p](#page-17-0)[a](#page-27-0)[n](#page-28-0)[c](#page-16-0)[y](#page-17-0) $(1 + 1)$

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45Ca, ⁴⁷Ca light isotopes spectra

Test our interaction with light Ca isotopes

Results in *pf*–*pfg*⁹/² spaces and based on NN–NN+3N interactions compared to standard phenomenological interactions

Agreement with experiment similar to phenomenological interactions

s[d](#page-28-0) degrees of freedom might still play a role fo[r s](#page-24-0)[om](#page-26-0)[e](#page-27-0) [e](#page-25-0)[x](#page-26-0)[ci](#page-16-0)[t](#page-17-0)ed [s](#page-16-0)t[a](#page-27-0)[te](#page-28-0)[s](#page-0-0) 290

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B(E2) Transition Strengths

B(E2)s in reasonable agreement with experiment (order of magnitude)

Similar quality as phenomenological interactions (very close to KB3G)

⁴⁶Ca: *sd* degrees of freedom?

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53Ca, ⁵⁴Ca neutron rich spectra:

Phenomenological interactions different results in neutron rich nuclei

MBPT: prediction, more controlled

Explore sensitivity to theoretical uncertainties (difference with respect to 3N 1st order results)

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Residual 3N Forces

In the most neutron-rich isotopes, 3N forces between 3 valence neutrons (suppressed by *Nvalence*/*Ncore*) can give a relevant contribution

Exact treatment of residual 3N forces would demand *V* ³*^N* diagonalization

Expected small correction compared to other 3N contributions Evaluated 1st order perturbation theory:

 $\langle\Psi|V^{3N}|\Psi\rangle$

Residual 3N forces expected more important in O isotopes

Neutron Number (*N*)

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Residual 3N forces in O isotopes

Residual 3N forces small and repulsive correction

Increase with the number of neutrons

Negligible up to ²⁴O, up to ∼1 MeV in ²⁸O

Important in the dripline region: flat energy be[hav](#page-28-0)[io](#page-30-0)[ur](#page-28-0)

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Residual 3N Forces around O dripline

Caesar, Simonis et al, arXiv:1209.0156

Compare to experimental energies of unbound ²⁵O and ²⁶O (relative to 24 O)

Without residual 3N forces O dripline predicted at ²⁶O by ∼ 300 keV (3N forces to 3rd order in MBPT)

Residual 3N forces drive dripline to ²⁴O in agreement with experiment

Unbound ²⁵O and ²⁶O also in very good agreement with MSU and GSI measurements

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Residual 3N forces in Ca isotopes

Smaller that in the O case: 20 vs 8 core-neutrons

Negligible for the discussion of S_{2n} , Δ_n or spectra

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Summary and Outlook

Microscopic calculation based on chiral EFT (NN+3N forces) and MBPT gives good agreement with experimental two-neutron separation energies, pairing gaps and excitation spectra for calcium isotopes:

- Experimental trends in S_{2n} 's and $\Delta_n^{(3)}$'s nicely reproduced
- S_{2n} 's and $\Delta_n^{(3)}$ together with 2⁺ energies establish shell closures: $N = 28$ appears, $N = 32/34$ enhanced/reduced by 3N forces
- Predicted spectra for Ca neutron rich isotopes
- Residual 3N forces included and found relevant in O isotopes

Outlook:

Explore uncertainties in the theoretical calculation

L[e](#page-32-0)e-Suzuki transformation from *pf g*_{9/2} to *pf* s[pac](#page-31-0)e

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