The No Core Shell Model for Bound, Resonant and Scattering States

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OUTLINE

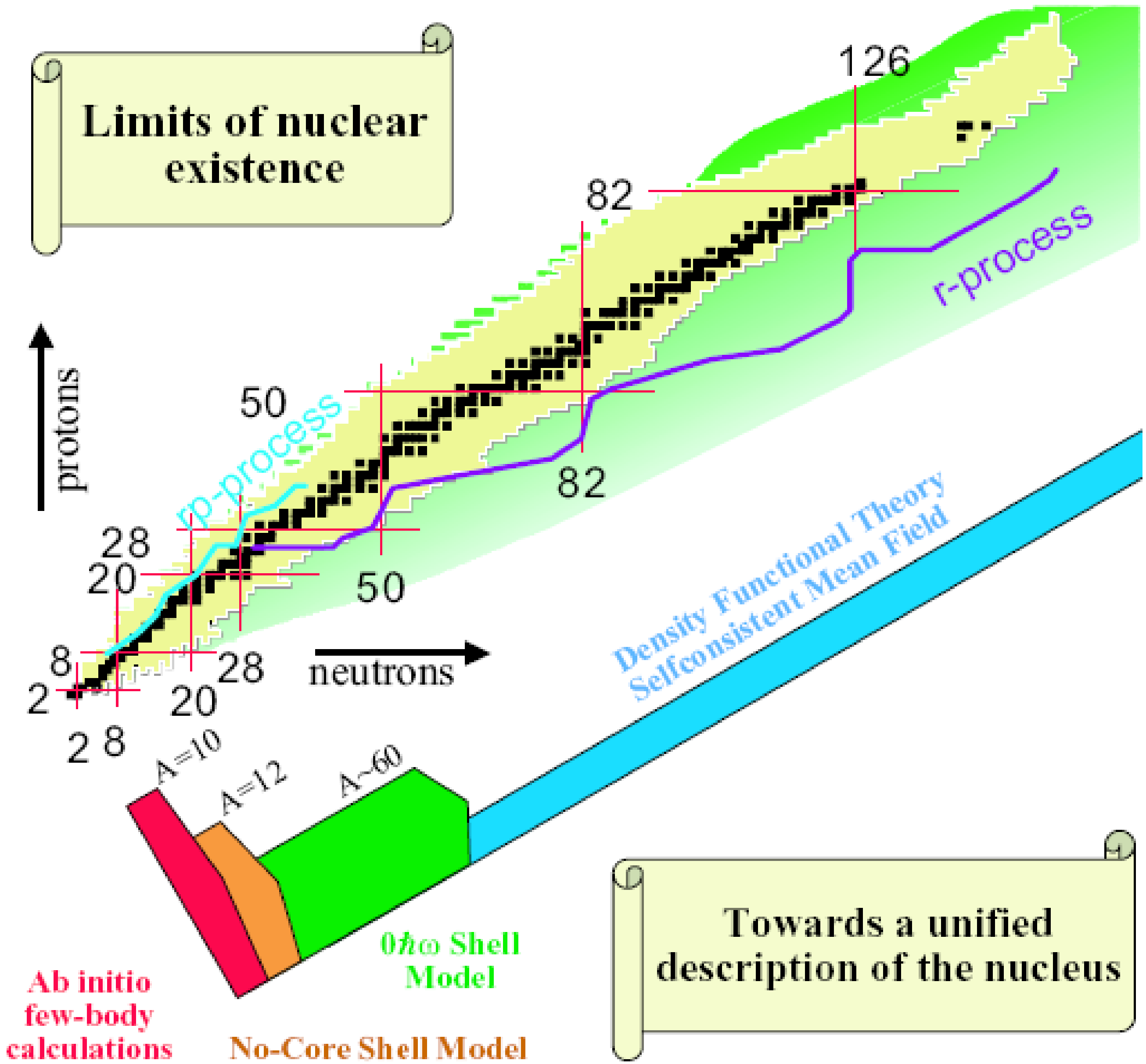
I. Introduction: No Core Gamow Shell Model (NCGSM)

II. NCGSM Formalism

III. NCGSM: Applications to Light Nuclei

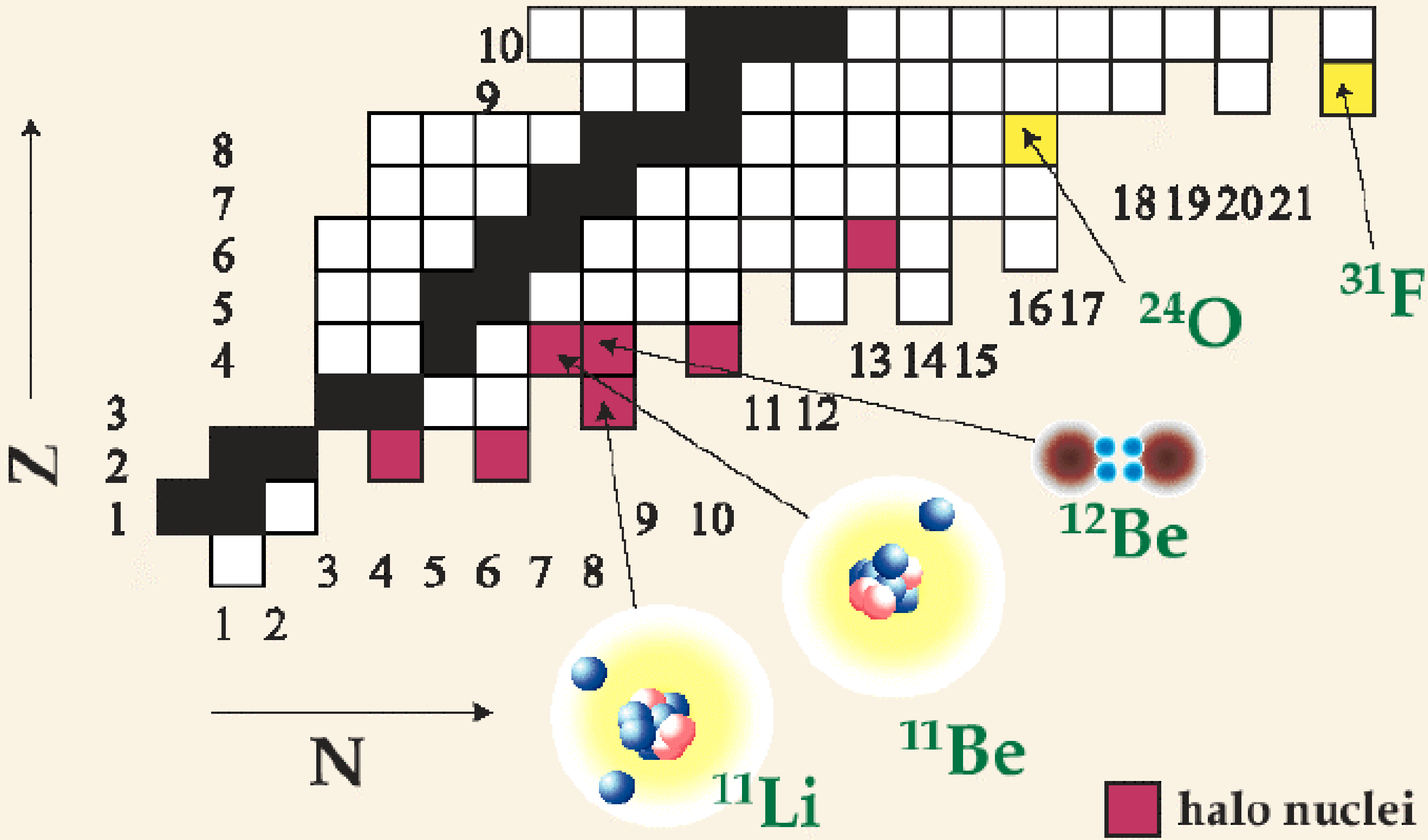
IV. Summary and Outlook

I. Introduction: No Core Gamow Shell Model

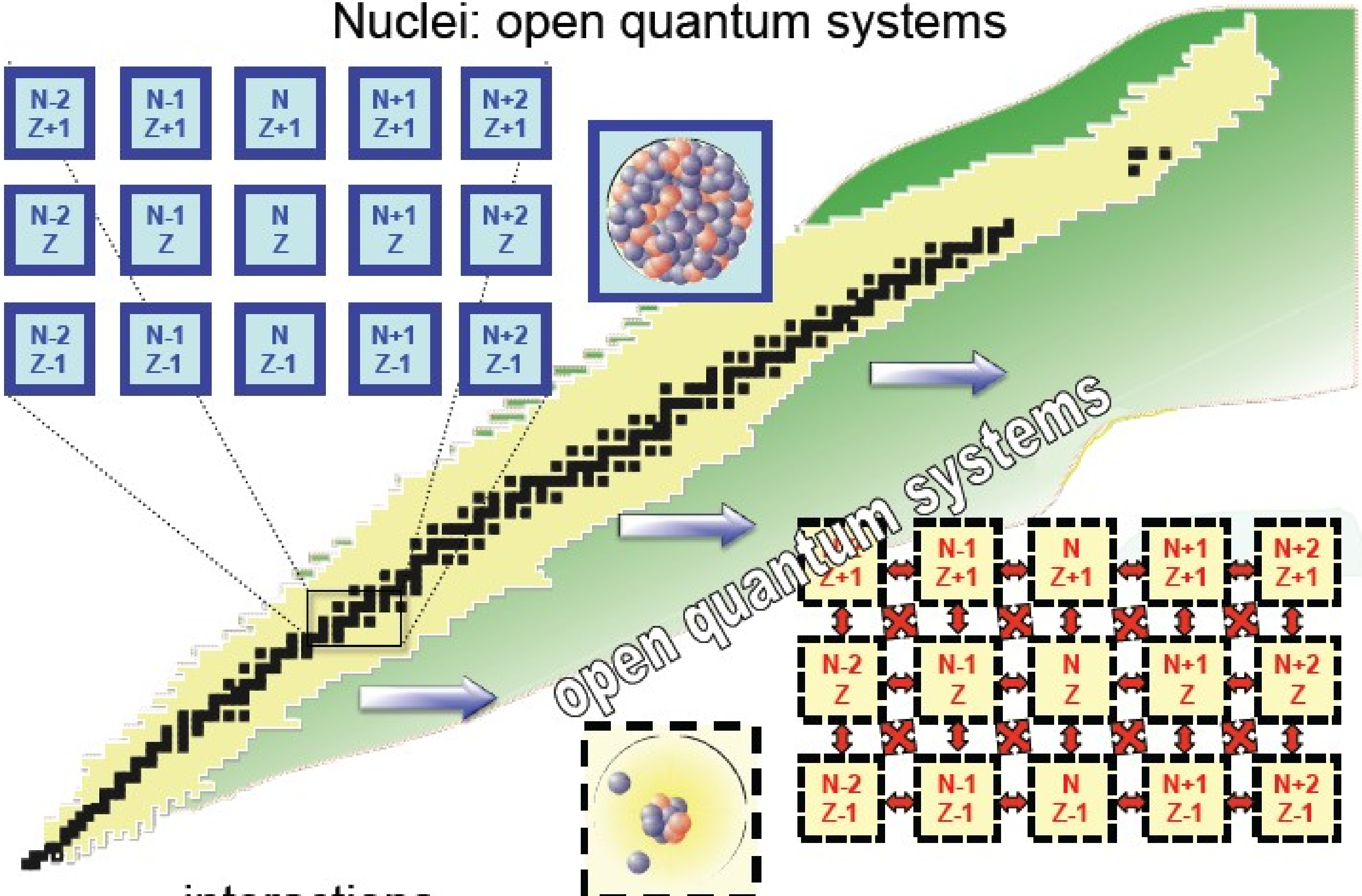


Limits of nuclear existence
protons
neutrons
rp-process
r-process
126
82
50
28
20
8
2
A=10
A=12
A~60
Density Functional Theory
Selfconsistent Mean Field
Ab initio few-body calculations
No-Core Shell Model
0ħω Shell Model
Towards a unified description of the nucleus

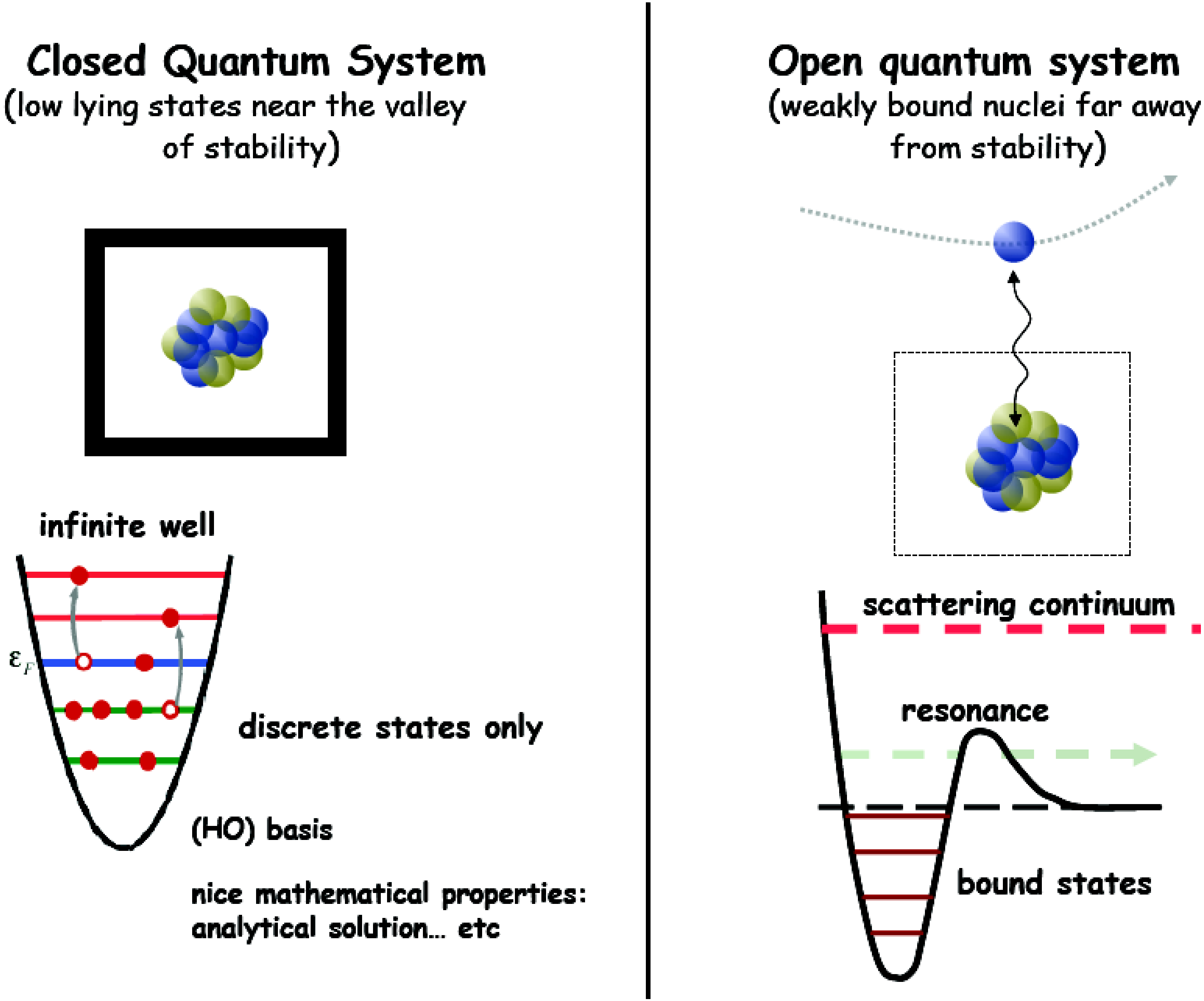
Light drip line nuclei



Nuclei: open quantum systems



interactions
correlations
many-body techniques



Closed Quantum System
(low lying states near the valley of stability)
infinite well
ε_F
discrete states only
(HO) basis
nice mathematical properties: analytical solution... etc
Open quantum system
(weakly bound nuclei far away from stability)
scattering continuum
resonance
bound states

II. NCGSM Formalism

Some selected references for the Complex Energy Gamow Shell Model

1. N. Michel, et al., Phys. Rev. C 67, 054311 (2003)

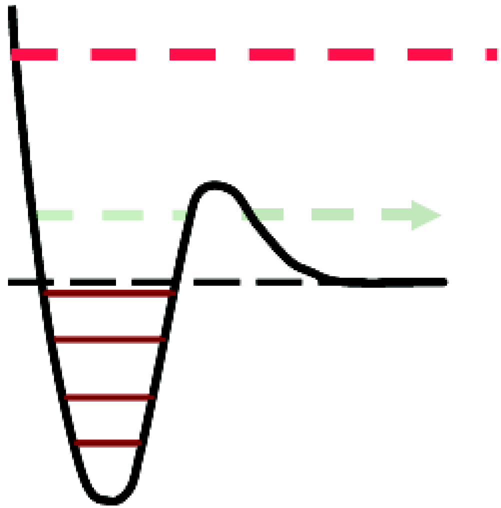
2. G. Hagen, et al., Phys. Rev. C 71, 044314 (2005)

3. J. Rotureau, et al., PRL 97, 110603 (2006)

4. M. Michel, et al., J. Phys. G: Nucl. Part. Phys. 36, 013101 (2009)

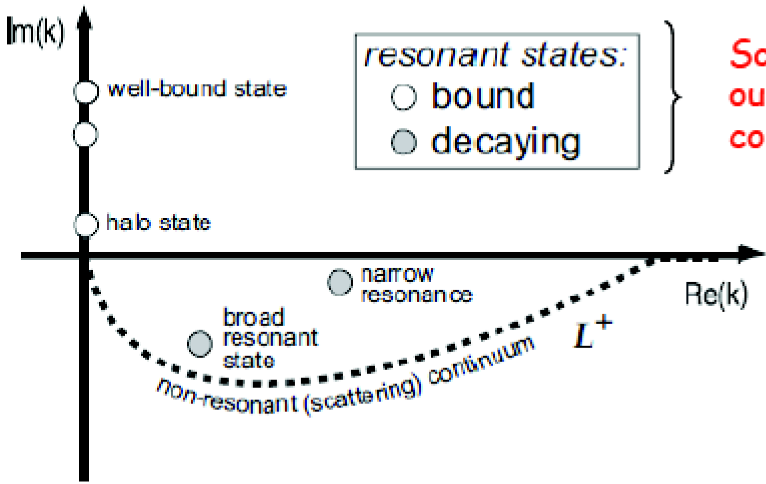
5. G. Papadimitriou, et al., PRC(R) 84, 051304 (2011)

Resonant and non-resonant states (how do they appear?)



$$\left(-\frac{d^2}{dr^2}+v(r)+\frac{l(l+1)}{r^2}-k^2\right)u_l(k,r)=0$$

$$k=\sqrt{\frac{2mE}{\hbar^2}}$$



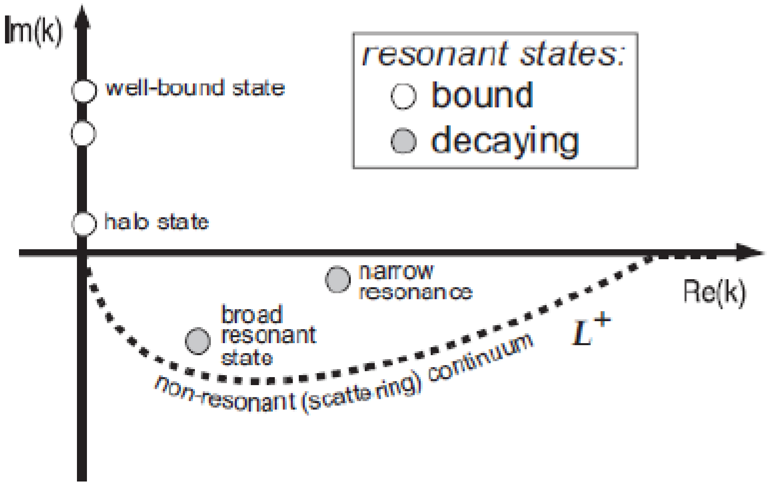
Solutions with outgoing boundary conditions

Solution of
the one-body Schrödinger equation with outgoing boundary conditions and a finite depth potential

$$u_l(k,r)\sim C_+H_l^+(k,r)\,,r\to\infty \quad bound\ \ states,\ resonances$$

$$u_l(k,r)\sim C_+H_l^+(k,r)+C\_H_l^-(k,r)\,,r\to\infty \quad scattering\ \ states$$

The Berggren basis (cont'd)



T.Berggren (1968)
NP A109, 265

The eigenstates of the 1b Shrödinger equation form a complete basis, **IF**:

we also consider the L_+ scattering states

$$\sum |u_{res}\rangle\langle u_{res}| + \int_{L^+} dk |u_k\rangle\langle u_k| = 1$$

$|u_k\rangle$ are complex continuum states along the L^+ contour (they satisfy scattering b.c)

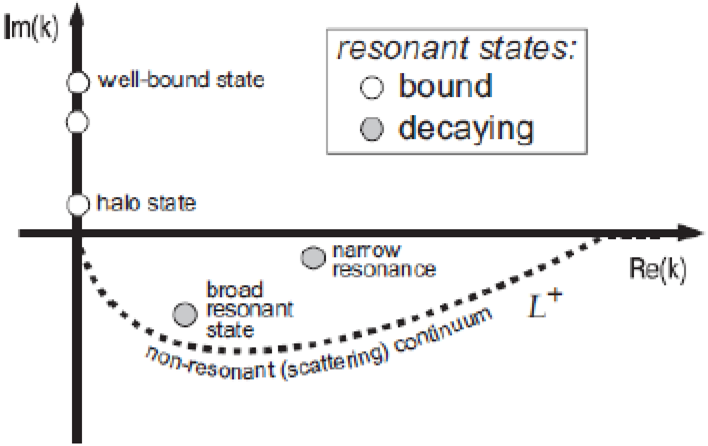
The shape of the contour is arbitrary, but it has to be below the resonance(s) position(s) (proof by T. Berggren)

In practice the continuum is discretized via a quadrature rule (e.g Gauss-Legendre):

$$\sum |u_{res}\rangle\langle u_{res}| + \sum_i |u_{ki}\rangle\langle u_{ki}| \simeq 1 \qquad \text{with} \qquad |u_k\rangle = \sqrt{\omega_i}|u_{ki}\rangle$$

Berggren's Completeness relation and Gamow Shell Model

N.Michel *et.al* 2002
PRL 89 042502



$$\sum |u_{res}\rangle\langle u_{res}| + \int_{L+} dk|u_k\rangle\langle u_k| = 1$$

resonant states (bound, resonances…)

Non-resonant Continuum along the contour

$$\sum |u_{res}\rangle\langle u_{res}| + \sum_i |u_{ki}\rangle\langle u_{ki}| \simeq 1$$

$$|SD_i\rangle = |u_{i1} \ldots\ldots u_{iA}\rangle$$

The GSM in 4 steps

Hermitian Hamiltonian

Many-body $|SD_i\rangle$ basis

Hamiltonian matrix is built (complex symmetric):

$$\langle SD|H|SD\rangle$$

Hamiltonian diagonalized

$$|\Psi\rangle = \sum_n c_n|SD_n\rangle$$

Many body correlations and coupling to continuum are taken into account simultaneously

The Density Matrix Renormalization Group (DMRG)

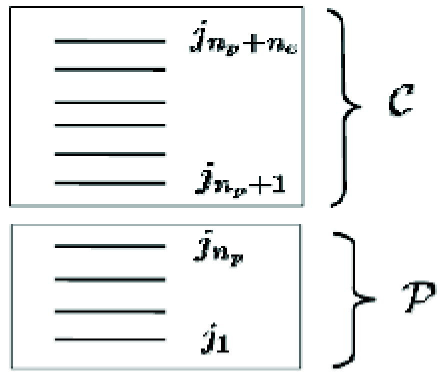
S.R White PRL 69 (1992) 2863
T.Papenbrock and D.Dean J.Phys.G 31 (2005) S1377
S.Pittel et al PRC 73 (2006) 014301
J.Rotureau et al PRC 79 (2009) 014304
J. Rotureau et al PRL 97 (2006) 110603

✓ **Truncation Method** applied to lattice models, spin chains, atomic nuclei….

GSM+DMRG

$$|\Psi\rangle^J = \sum_{p,c} \Psi_{pc} \left[|p\rangle^{J_p} |c\rangle^{J_c} \right]^J$$

DMRG → truncation in $\mathcal{C}$



P: poles
C: Continua

✓ Iterative method: In each step (N_{step}) a scattering shell is added from C.
→ Hamiltonian is diagonalized and density matrix is constructed:

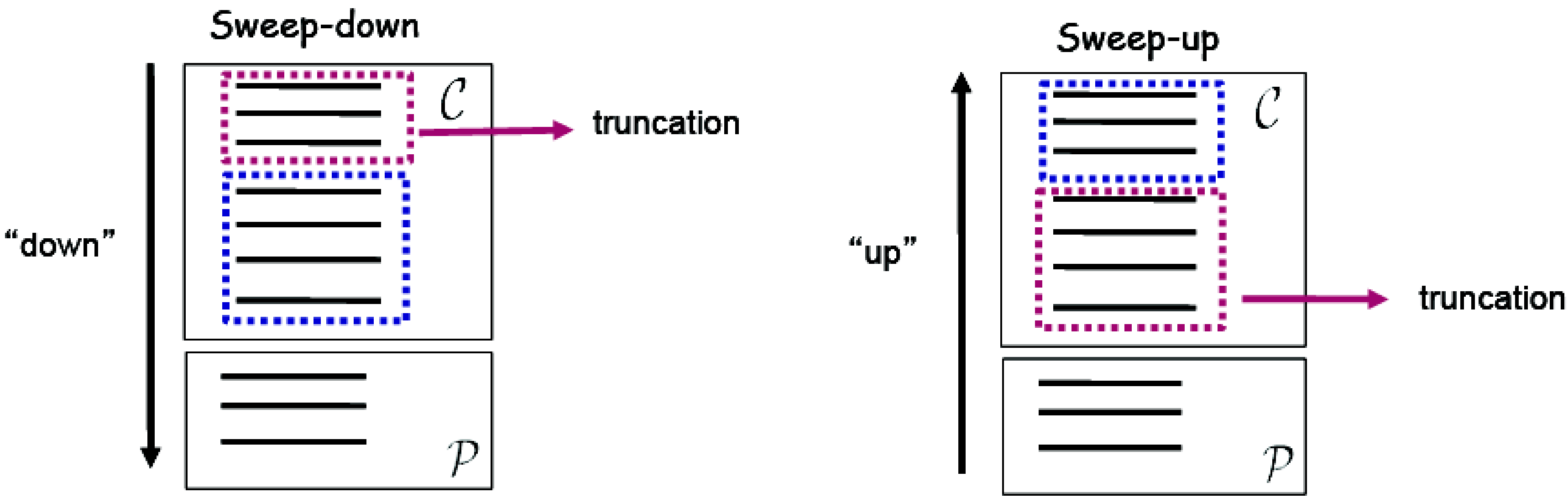
$$\rho^{J_c}_{c,c'} = \sum_p \Psi_{pc} \Psi_{pc'}$$

- truncation with the density matrix :

$$\rho^{J_c}_{c,c'} = \sum_p \Psi_{pc} \Psi_{pc'}$$

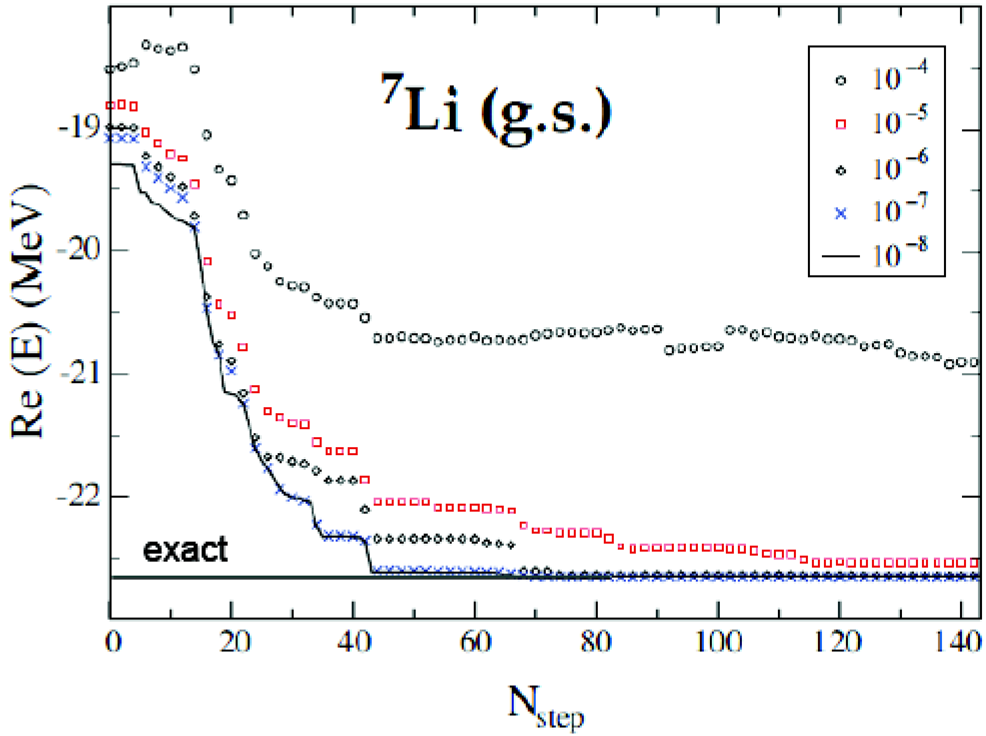
→ N_{opt} states that correspond to the largest eigenvalues of the density matrix are kept

- The process is reversed…
- In each step (shell added) the Hamiltonian is diagonalized and N_{opt} states are kept.
- Iterative method to take into account all the degrees of freedom in an effective manner.
- In the end of the process the result is the same with the one obtained by "brute" force diagonalization of H.



Density Matrix Renormalization Group - Examples - (GSM with a ^{4}He core)

J.Rotureau et al PRC 79 (2009) 014304



^{7}Li: 3 nucleons outside ^{4}He.
Max dim in DMRG: ~1400
19% of the full space space

$$\left| 1 - \mathcal{R}e \left(\sum_{i=1}^{N_\rho} w_i \right) \right| < e$$

Small ε → more states of ρ are kept in each step

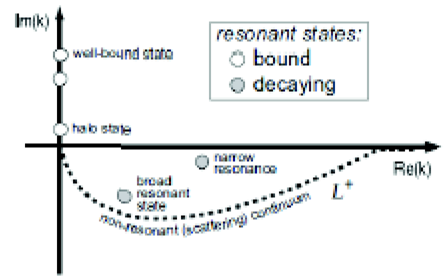
$$\sum_{\alpha} w_{\alpha} = 1$$

Gamow Shell Model in an ab-initio framework

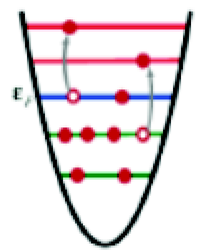
$$H = \frac{1}{A}\sum_{i<j}^{A} \frac{(\vec{p}_i - \vec{p}_j)^2}{2m} + V_{NN,ij} + \ldots \quad (1)$$

- Only NN forces at present
  → Argonne V18, (Wiringa, Stoks, Schiavilla PRC 51, 38, 1995)
  → N^3LO (D.R.Entem and R. Machleidt PRC(R) 68, 041001, 2003)
  → V_{lowk} technique used to decouple high/low momentum nodes. Λ_{Vlowk} = 1.9 fm^{-1}
  (S. Bogner et al, Phys. Rep. 386, 1, 2003)

- Basis states
  → s- and p- states generated by the HF potential



  → l > 1 H.O states

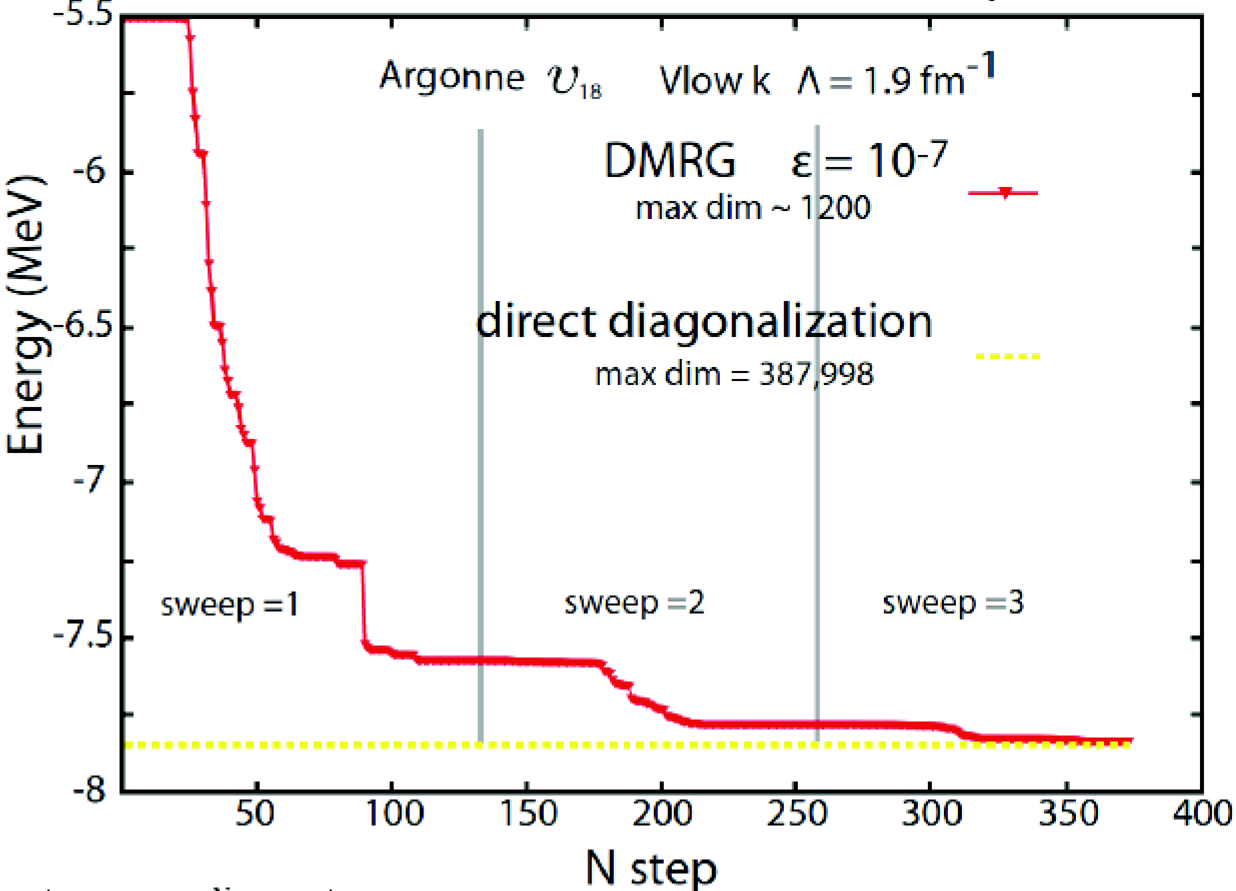


- Diagonalization of (1) → Applications to ^3H, ^4He, ^5He

III. NCGSM: Applications to Light Nuclei

Results

G.Papadimitriou, J.Rotureau, B. Barrett, N. Michel, M.Ploszajczak arXiv:1301.7140
Submitted at Phys.Rev.C



- 2 neutrons
- 1 proton
- Pole space A:0s1/2 (p/n)
- Continuum space B:
  - p3/2,p1/2,s1/2 real energy continua
  - d5/2-d3/2 H.O states
- 130 s.p. states total

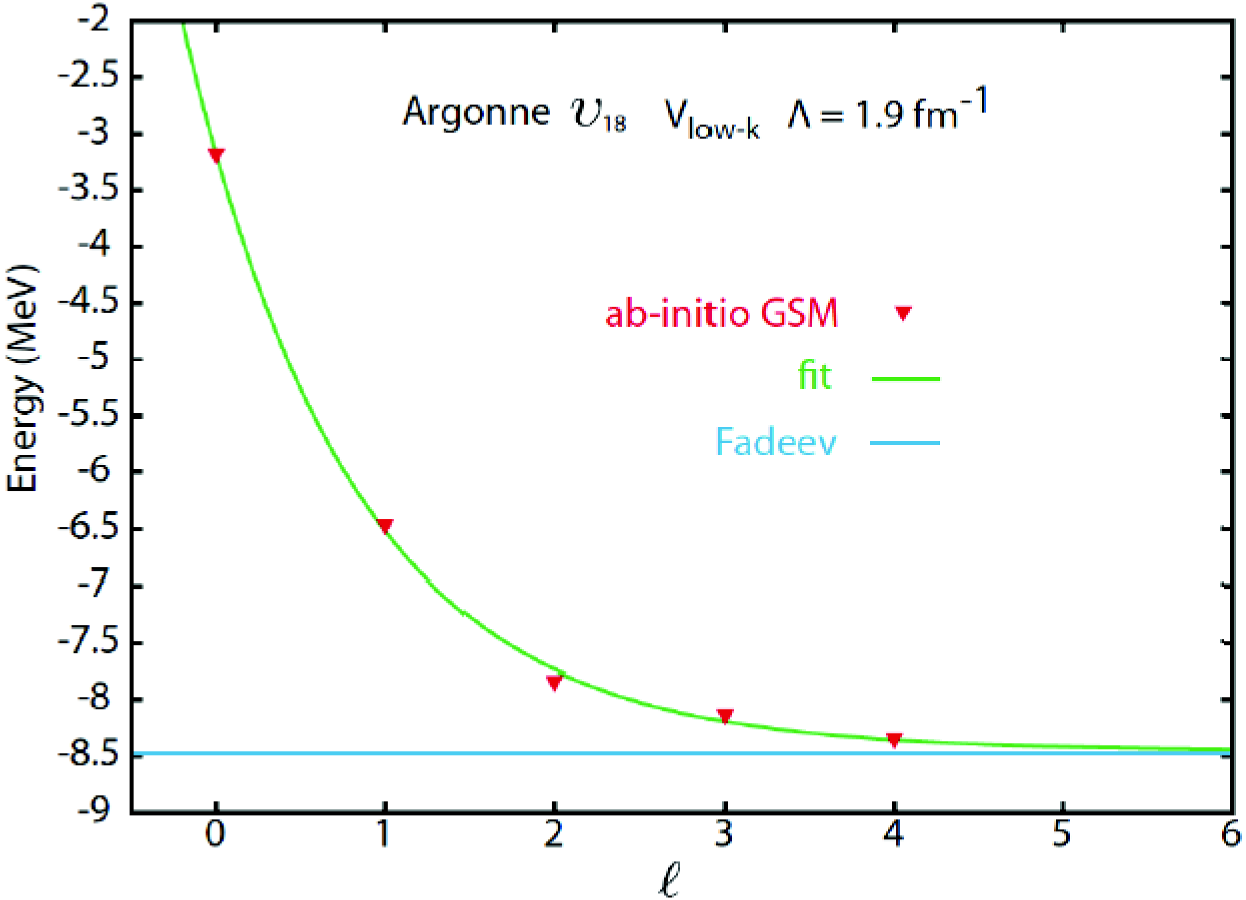
$$\left|1 - \mathcal{R}e\left(\sum_{i=1}^{N_\rho} w_i\right)\right| < e$$

$$\sum_\alpha w_\alpha = 1$$

E_{exact} = -7,840 MeV
$E_{DMRG\ (\varepsilon = 10^{-7})}$ = -7,832 MeV
$E_{DMRG\ (\varepsilon = 10^{-6})}$ = -7,820 MeV

Results: Triton

G.Papadimitriou, J.Rotureau, B. Barrett, N. Michel, M.Ploszajczak arXiv:1301.7140
Submitted at Phys.Rev.C



$E_{ab-initio}$ = -8.39 MeV

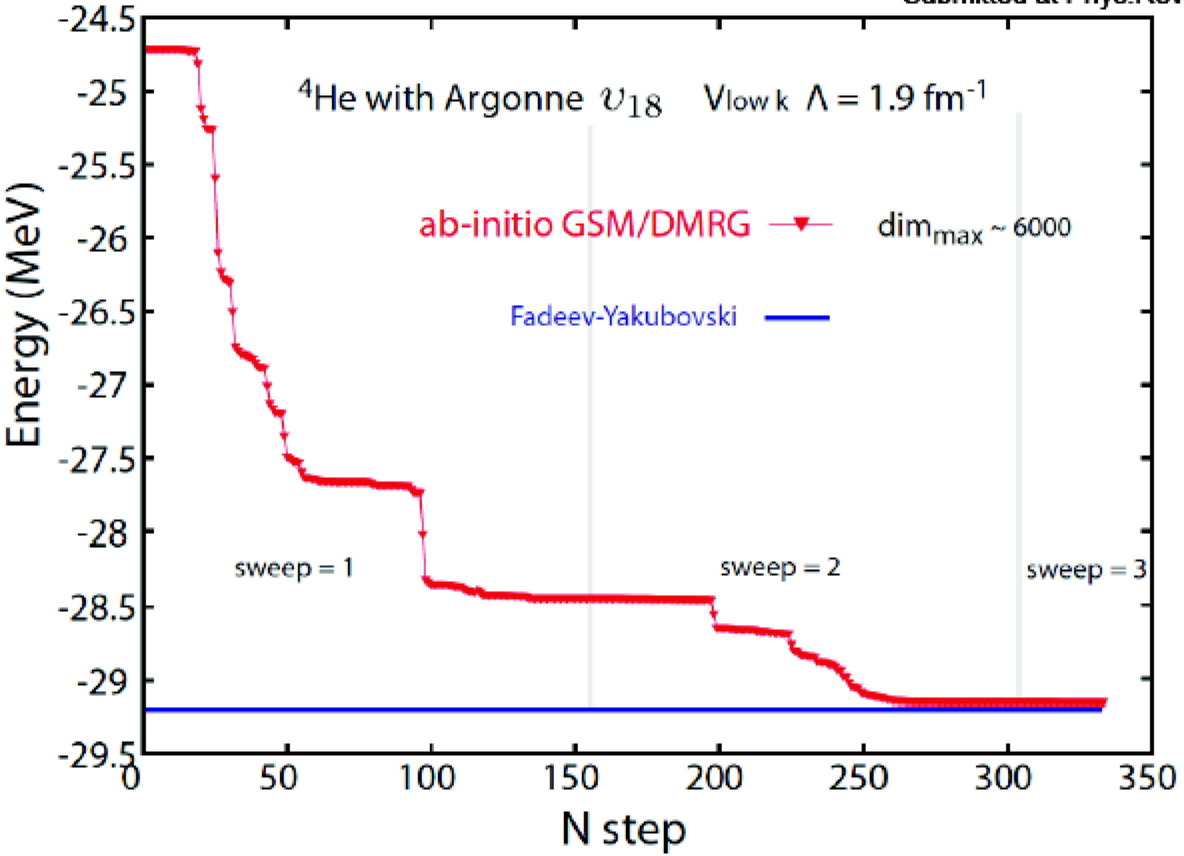
E_{extrap} = -8.44 +- 0.08 MeV

E_{Fadeev} = -8.47 MeV

Fadeev result from (Nogga, Bogner, Schwenk, PRC 70,061002, 2004)

Results: ^{4}He against Fadeev-Yakubovsky

G.Papadimitriou, J.Rotureau, B. Barrett, N. Michel, M.Ploszajczak arXiv:1301.7140
Submitted at Phys.Rev.C



- 2 neutrons
- 2 protons
- Pole space A:0s1/2 (p/n)
- Continuum space B:
  - p3/2,p1/2,s1/2 real energy continua
  - d5/2-d3/2, f5/2-f7/2, g7/2-g9/2 — H.O states
- 156 s.p. states total

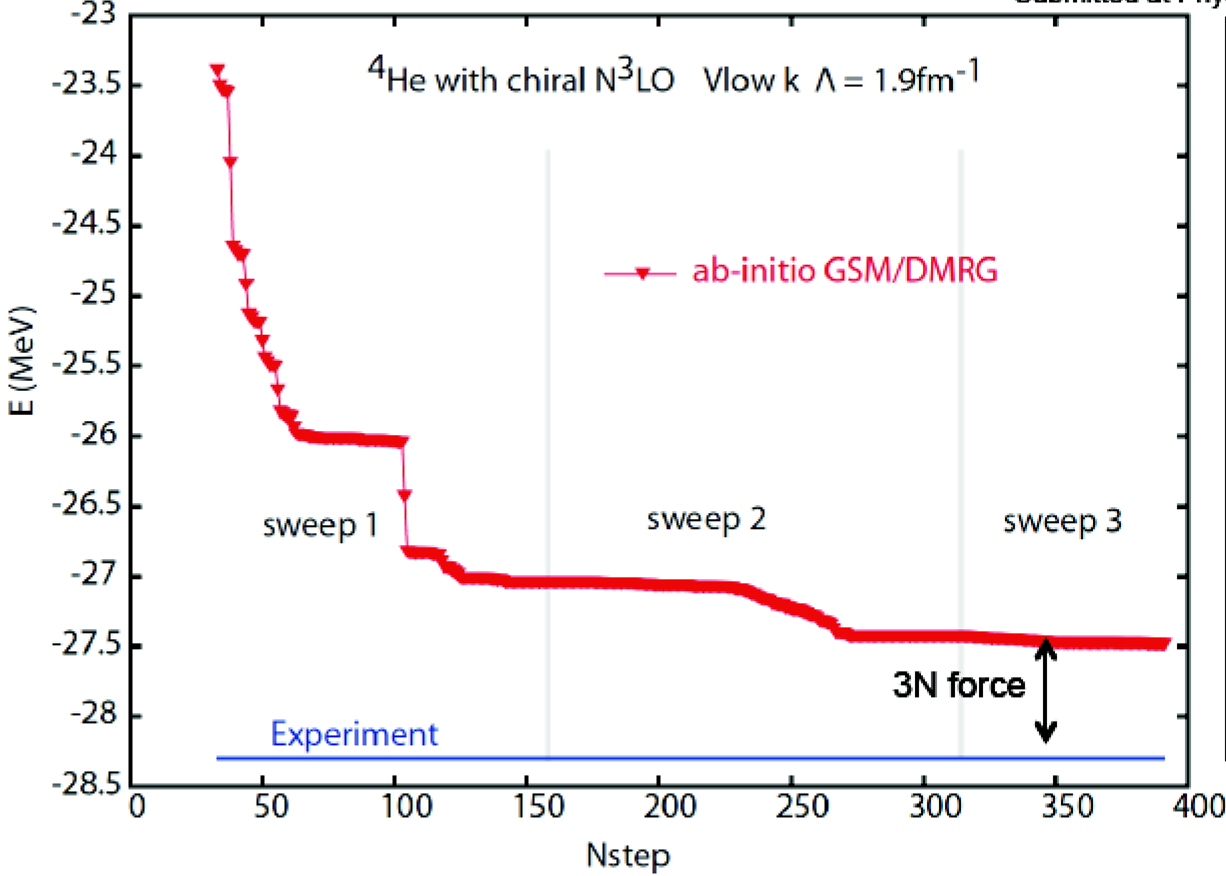
Dim for direct diagon: 119,864,088

$E_{ab\text{-}initio}$ = -29.15 MeV

E_{FY} = -29.19 MeV

Results: ^{4}He with chiral N^3LO

G.Papadimitriou, J.Rotureau, B. Barrett, N. Michel, M.Ploszajczak arXiv:1301.7140
Submitted at Phys.Rev.C

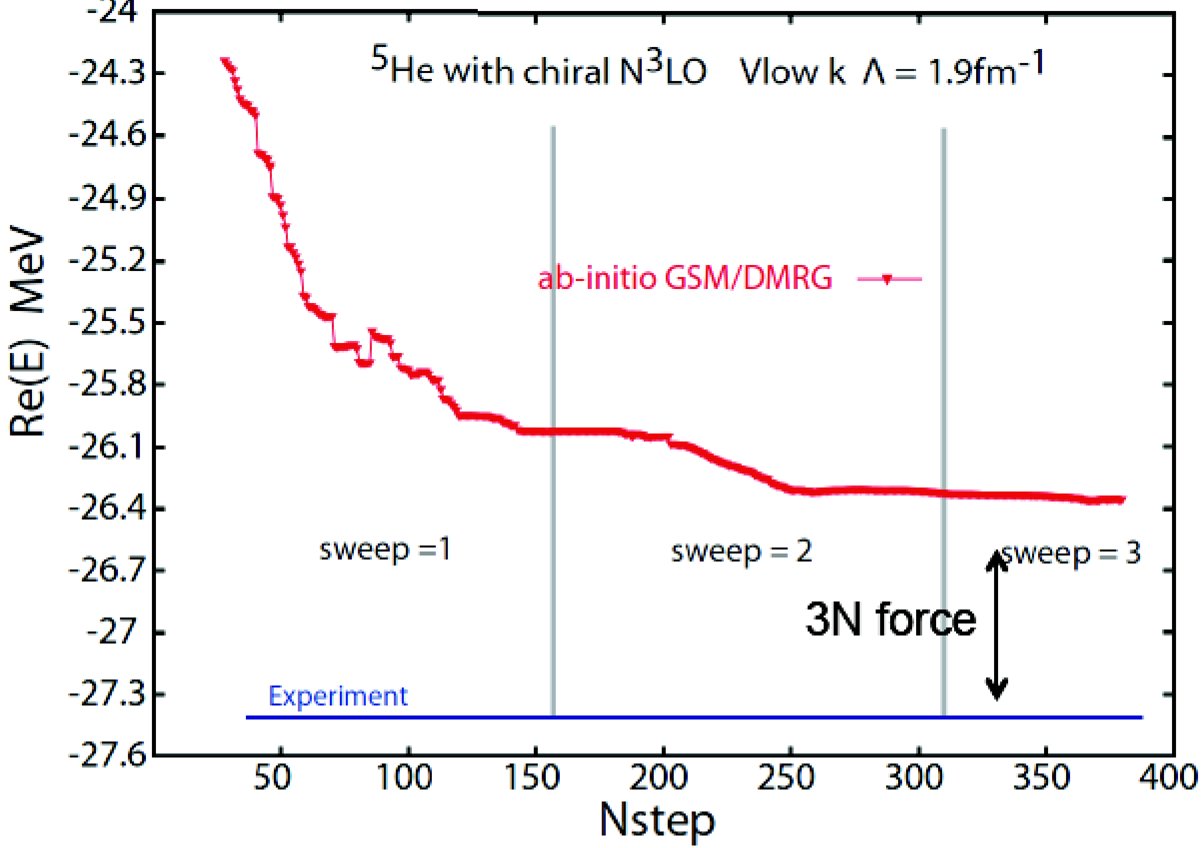


- 2 neutrons
- 2 protons
- Pole space A:0s1/2 (p/n)
- Continuum space B:
  - p3/2,p1/2,s1/2 real energy continua
  - d5/2-d3/2, f5/2-f7/2, g7/2-g9/2 — H.O states
- 156 s.p. states total

3N force arises from the renormalization of the NN interaction.

Results: ^{5}He with chiral N^{3}LO (real part)

G.Papadimitriou, J.Rotureau, B. Barrett, N. Michel, M.Ploszajczak arXiv:1301.7140
Submitted at Phys.Rev.C



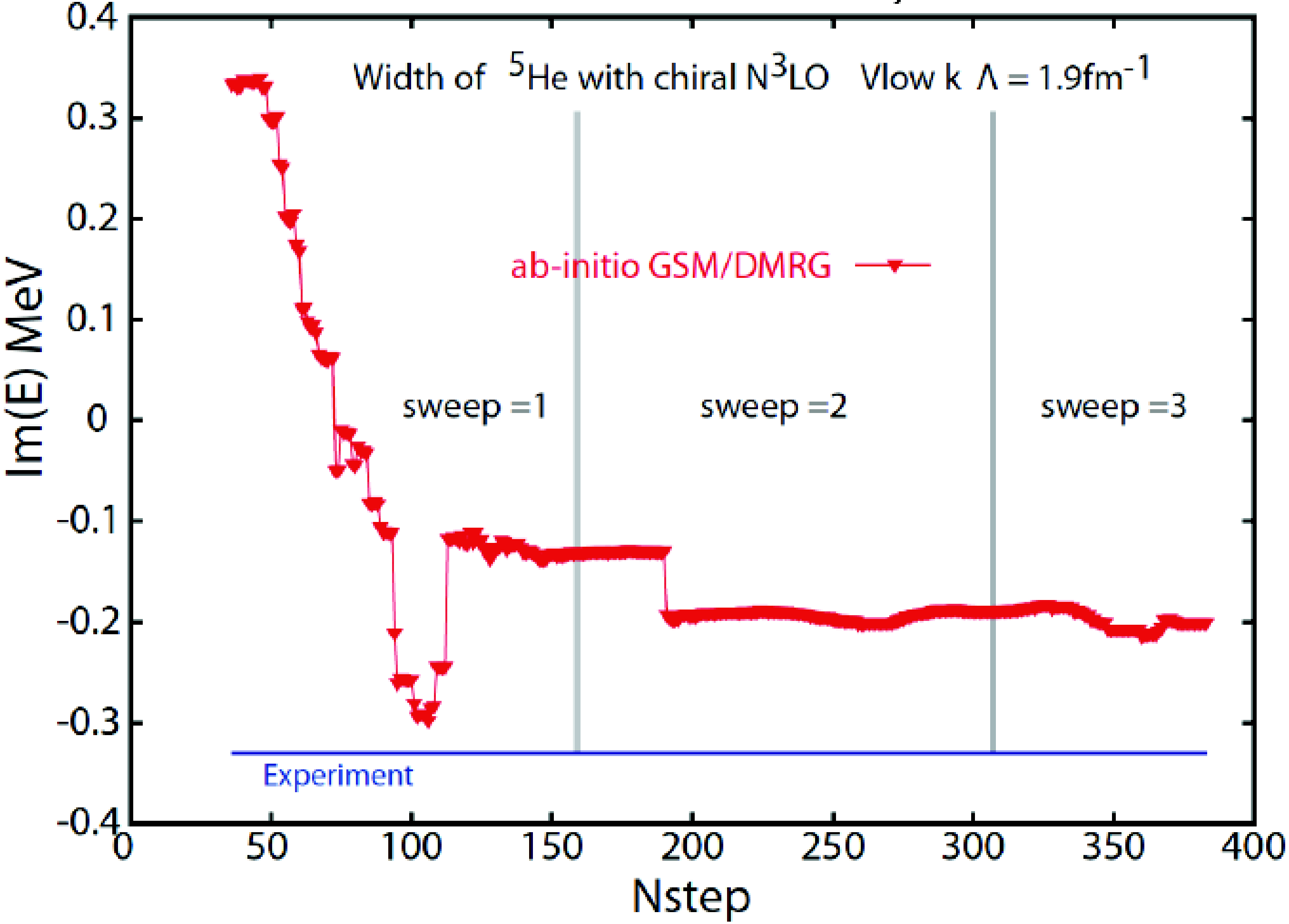
- 3 neutrons
- 2 protons
- Pole space A:0s1/2 (p/n) + 0p3/2 n resonant state
- Continuum space B:
  - p3/2 complex continuum
  - p1/2-s1/2 real continua
  - d5/2-d3/2, f5/2-f7/2, g7/2-g9/2 } H.O states
- 157 s.p. states total

Dim for direct diagon: $3*10^9$

DMRG dim ~ 10^5

Results: ^{5}He imaginary part (width) with chiral N^3LO

G.Papadimitriou, J.Rotureau, B. Barrett, N. Michel, M.Ploszajczak arXiv:1301.7140
Submitted at Phys.Rev.C



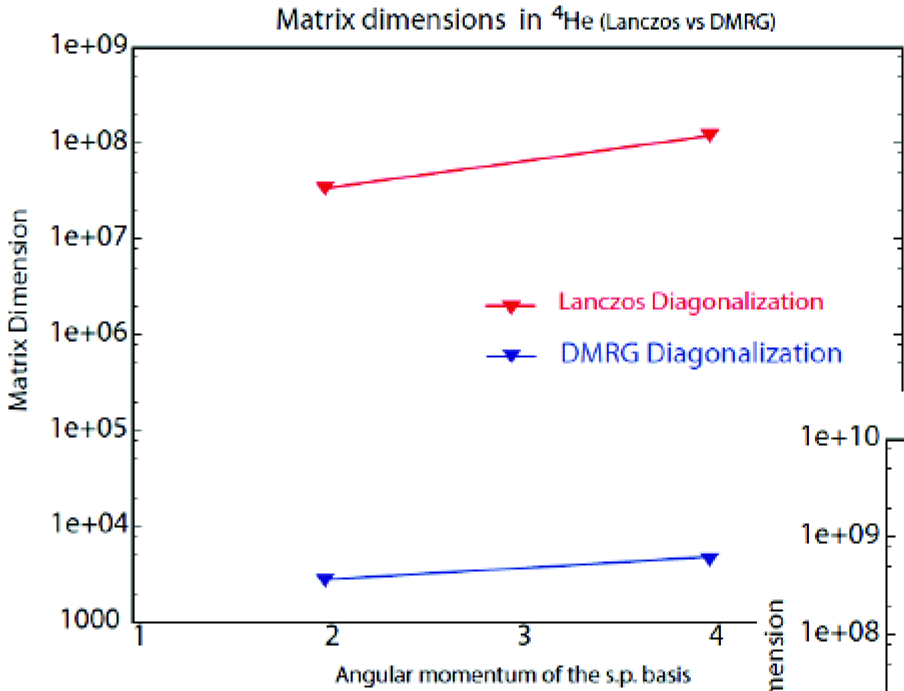
S_{1n} = -1.20 MeV

S_{1n} (exp) = -0.89 MeV

Unbound character of ^{5}He reproduced within an ab-initio framework

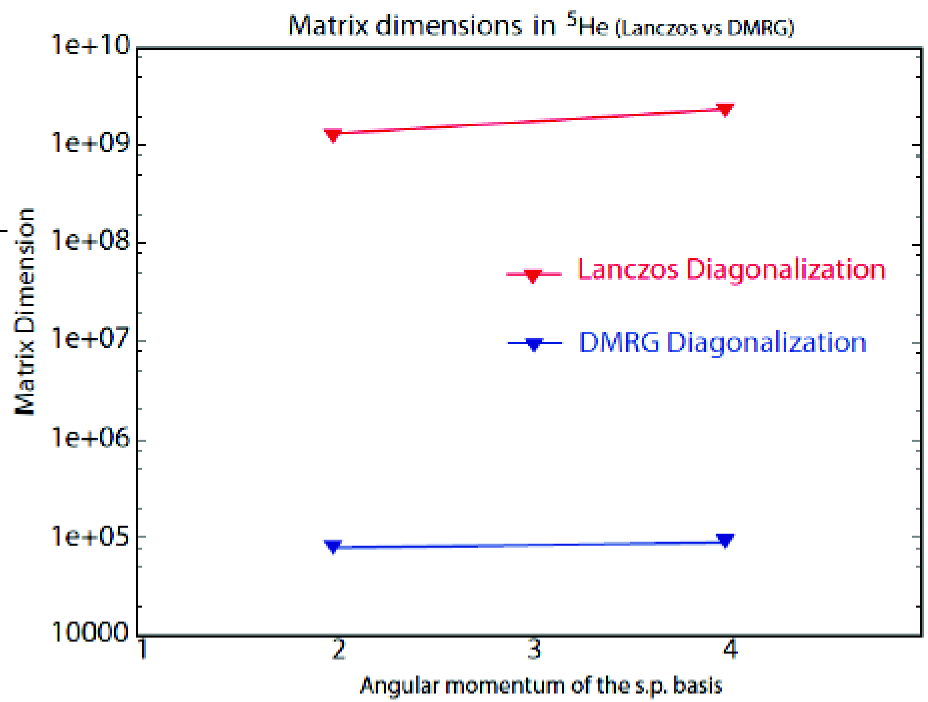
Satisfactory agreement of the **width** with experiment

Dimension comparison



→ Lanczos: "brute" force diagon of H.

→ DMRG: Diagon of H in the space where only the most important degrees of "freedom" are considered



IV. Summary and Outlook

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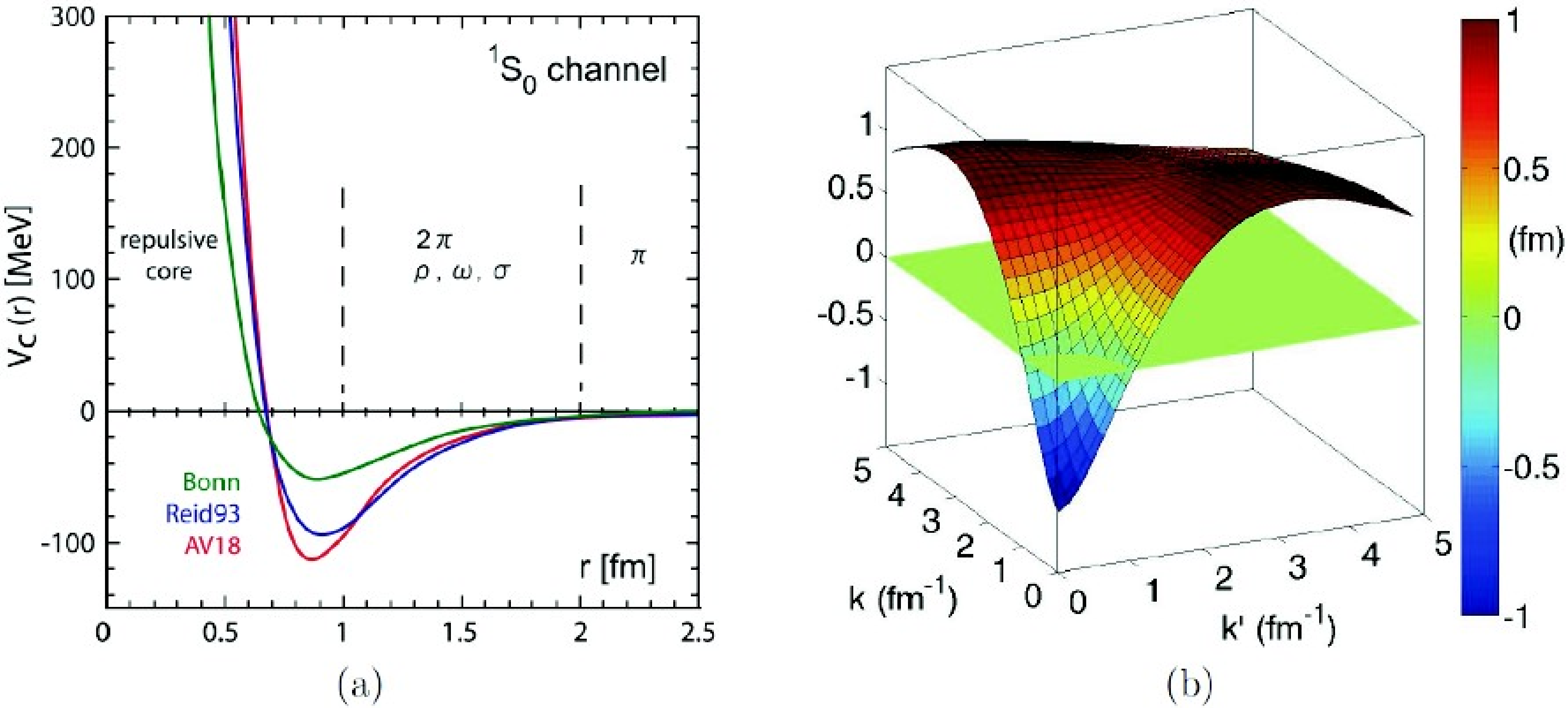
1. The Berggren basis is appropriate for calculations of weakly bound/unbound nuclei.

2. Berggren basis has been applied successfully in an ab-initio GSM framework --> No Core Gamow Shell Model for weakly bound/unbound nuclei.

3. Diagonalization with DMRG makes calculations feasible for heavier nuclei using Gamow states.

4. Future applications to heavier nuclei and to nuclei near the driplines.

Realistic two-body potentials in coordinate and momentum space



(a) (b)

Repulsive core makes calculations difficult

→ Need to decouple high/low momentum modes

✓ Achieved by $V_{low\text{-}k}$ or Similarity RG approaches (e.g. SRG)

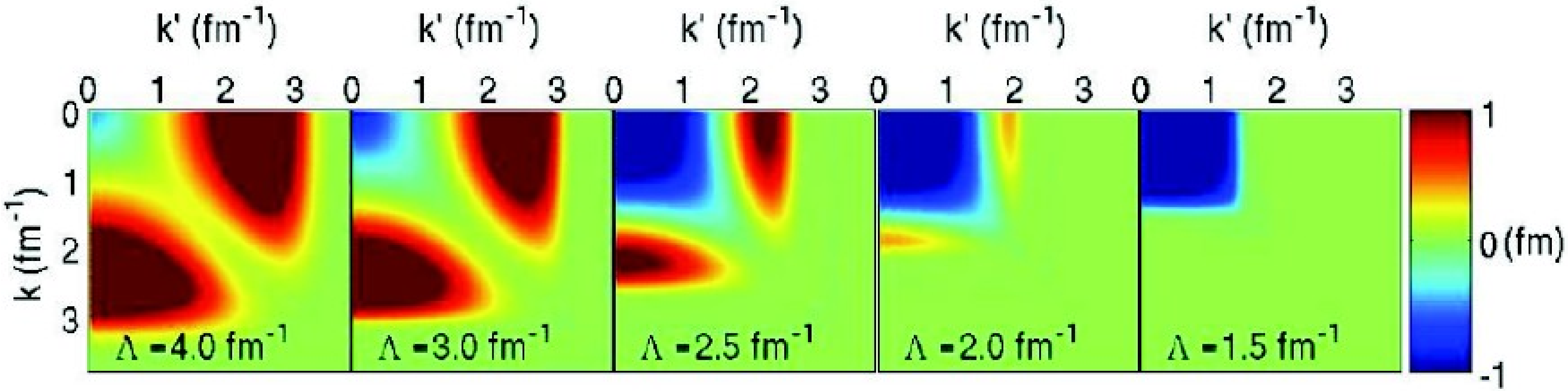
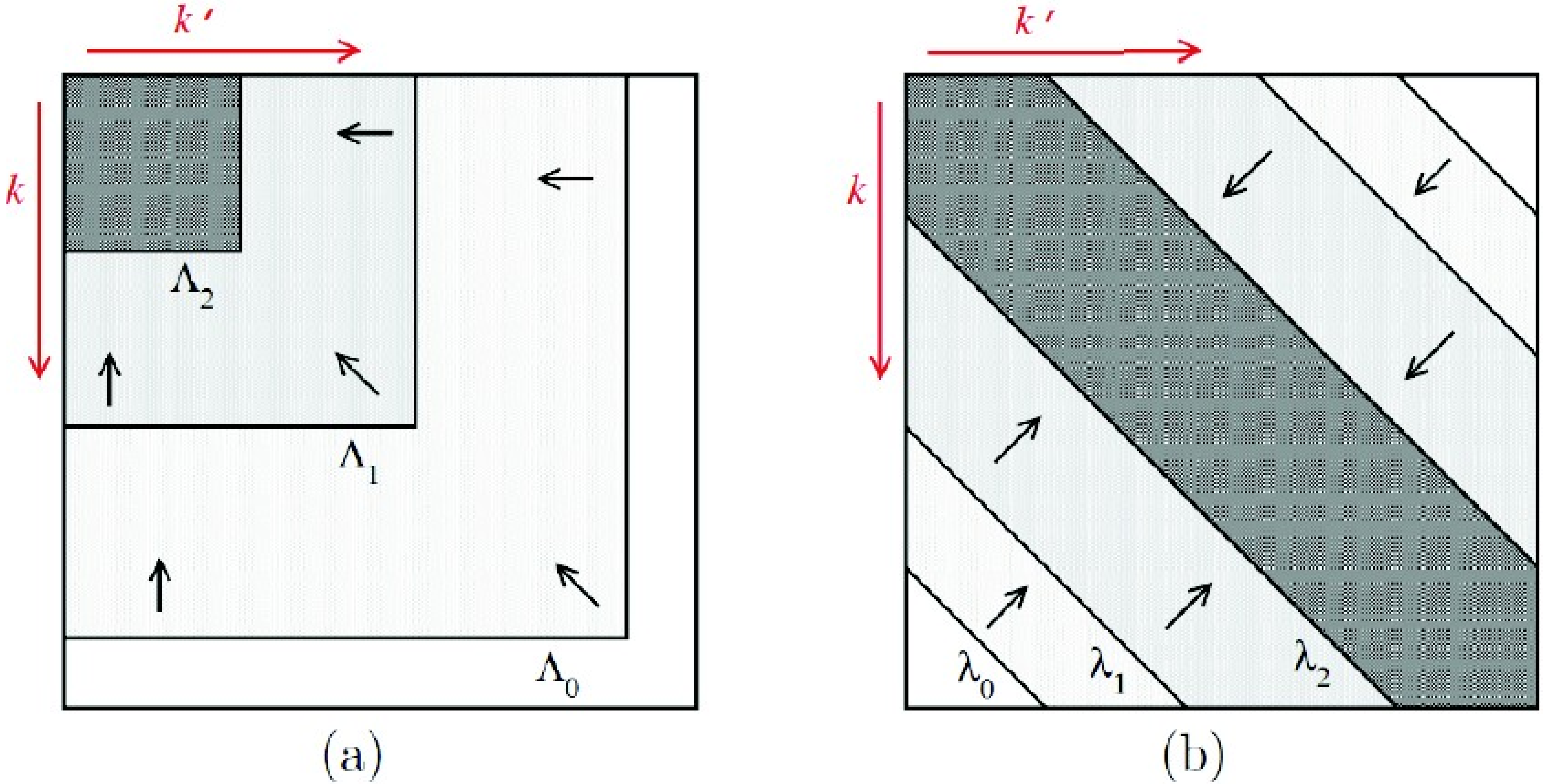


Fig. from S. Bogner et al Prog.Part.Nucl.Phys.65:94-147,2010

→ Observable physics is preserved (e.g. NN phase shifts) AND calculations become easier (work with the relevant degrees of freedom)

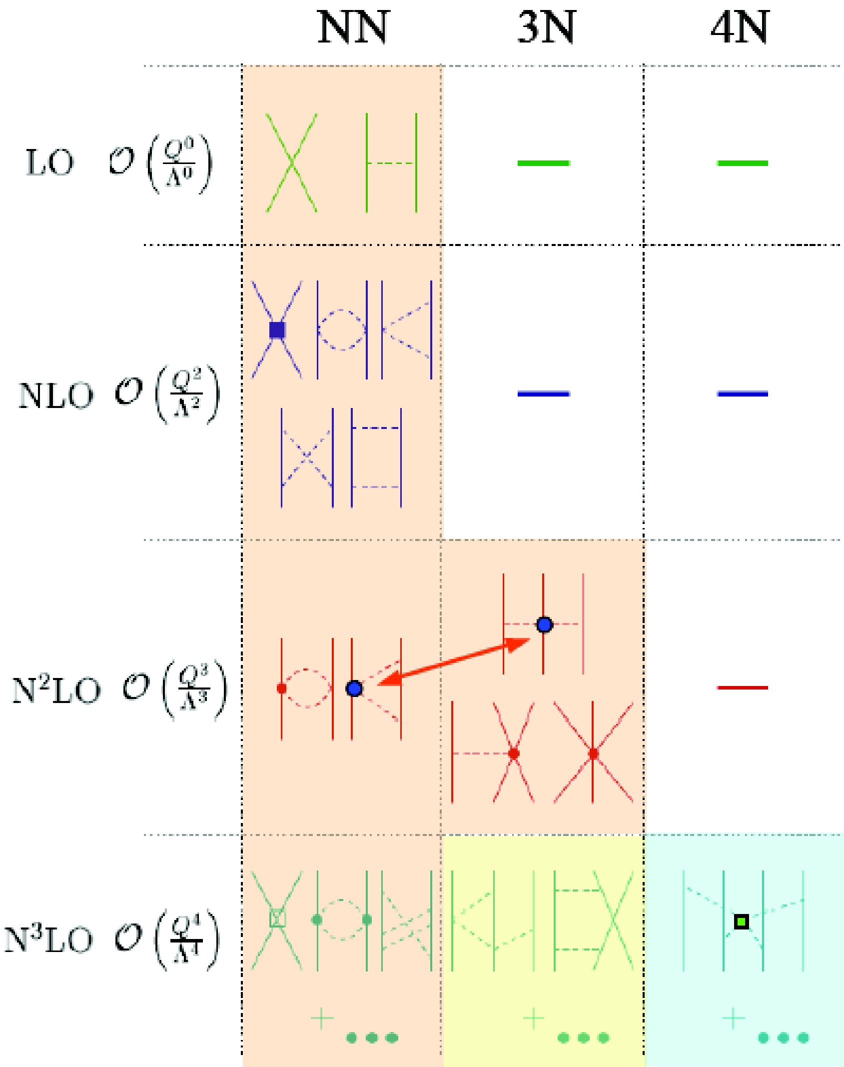
→ One has to deal with “induced” many-body forces…

Illustration on how the high momentum nodes are integrated out in the Vlowk (a) and in the SRG (b) RG methods



Chiral effective field theory (EFT) for nuclear forces

Separation of scales: low momenta $\frac{1}{\lambda} = Q \ll \Lambda_b$ breakdown scale Λ_b

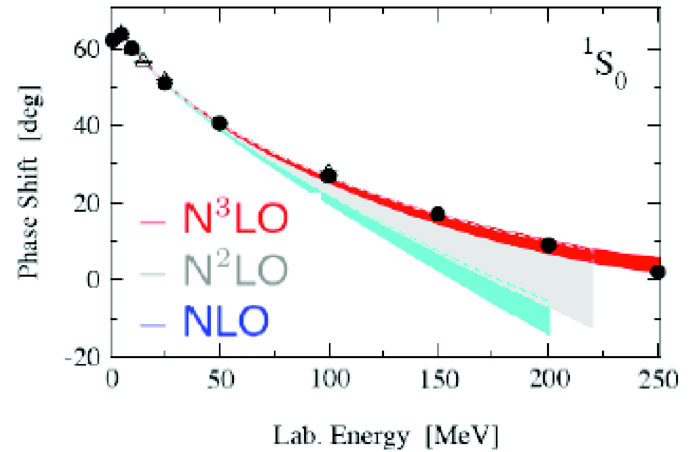


explains pheno hierarchy:
NN > 3N > 4N > …

NN-3N, πN, ππ, electro-weak,…
consistency

3N,4N: 2 new couplings to N^{3}LO!

theoretical error estimates



Weinberg, van Kolck, Kaplan, Savage, Wise, Epelbaum, Meissner, Nogga, Machleidt,…

A. Schwenk