

DVCS analysis status and perspectives

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- *photon leptonproduction cross section*
- *small x_B fits (H1 and ZEUS)*
- *fixed target (HERMES, JLAB)+ small x_B fits*
- *predictions and studies for future experiments*

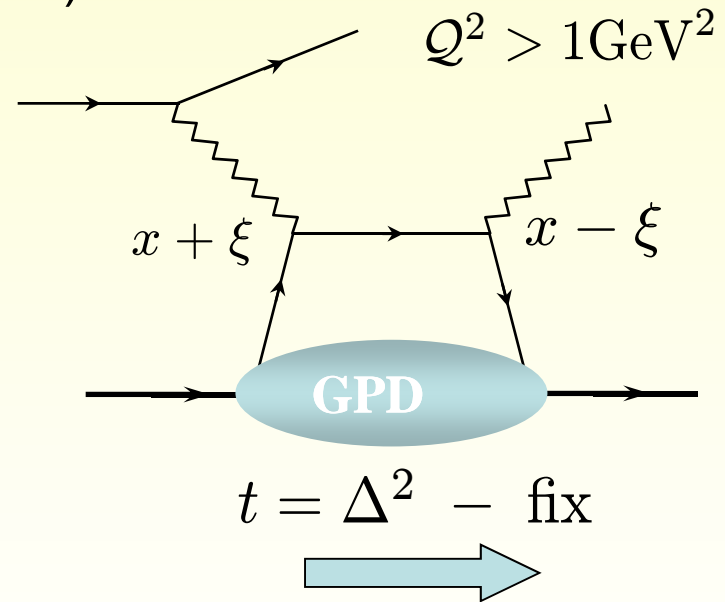
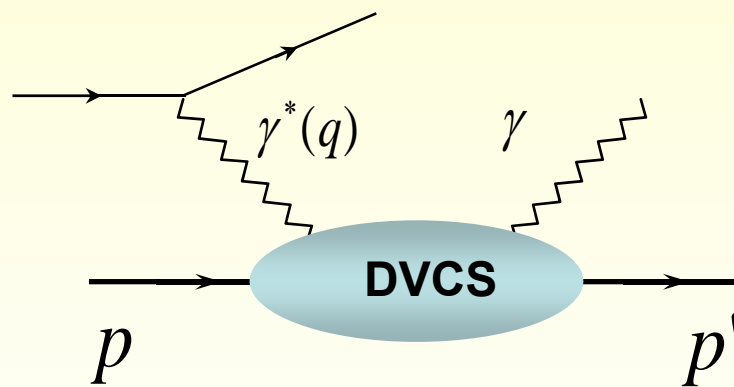
**in collaboration with K. Kumerički
E. Aschenauer and S. Fazio (EIC studies)**

GPDs embed non-perturbative physics

GPDs appear in various hard exclusive processes,

[DM et. al (90/94)
Radyushkin (96)
Ji (96)]

e.g., hard electroproduction of photons (DVCS)



$$\mathcal{F}(\xi, Q^2, t) = \int_{-1}^1 dx C(x, \xi, \alpha_s(\mu), Q/\mu) F(x, \xi, t, \mu) + O\left(\frac{1}{Q^2}\right)$$

CFF

Compton form factor

observable

hard scattering part

perturbation theory
(our conventions/microscope)

GPD

universal
(conventional)

higher twist

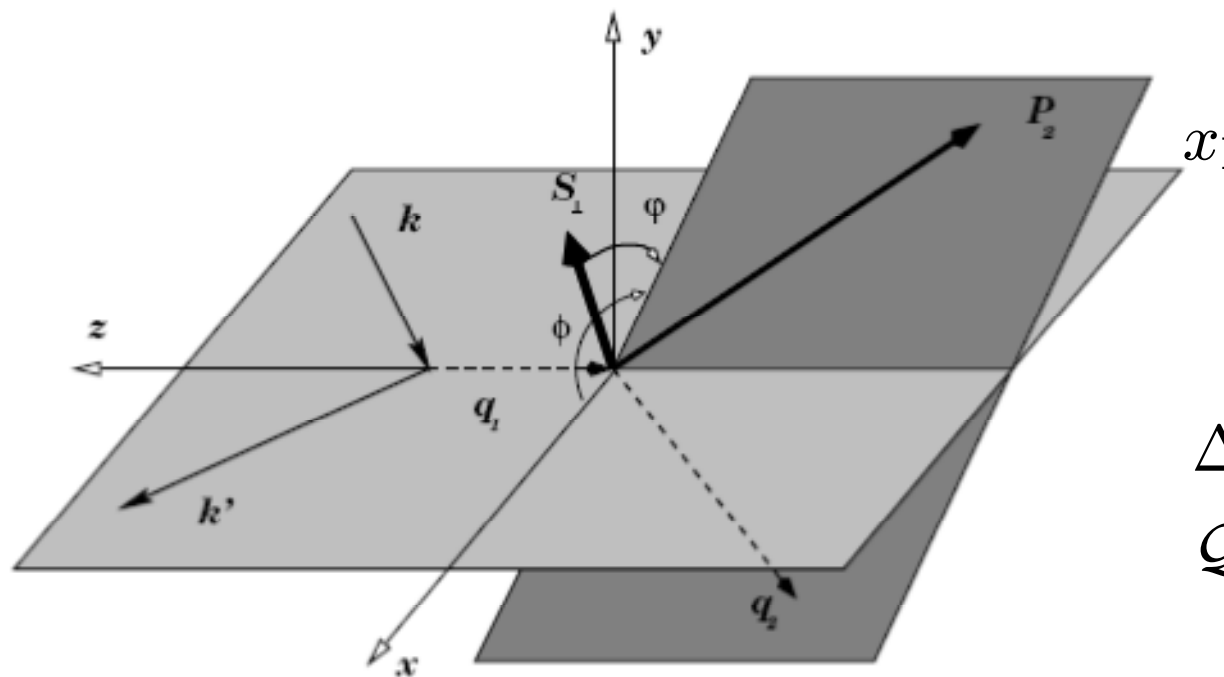
depends on
approximation

Photon leptonproduction $e^\pm N \rightarrow e^\pm N \gamma$

measured by **H1, ZEUS, HERMES, CLAS, HALL A** collaborations

planned at **COMPASS, JLAB@12GeV**, perhaps at ?? EIC,

$$\frac{d\sigma}{dx_{Bj} dy d|\Delta^2| d\phi d\varphi} = \frac{\alpha^3 x_{Bj} y}{16 \pi^2 Q^2} \left(1 + \frac{4M^2 x_{Bj}^2}{Q^2} \right)^{-1/2} \left| \frac{\mathcal{T}}{e^3} \right|^2,$$



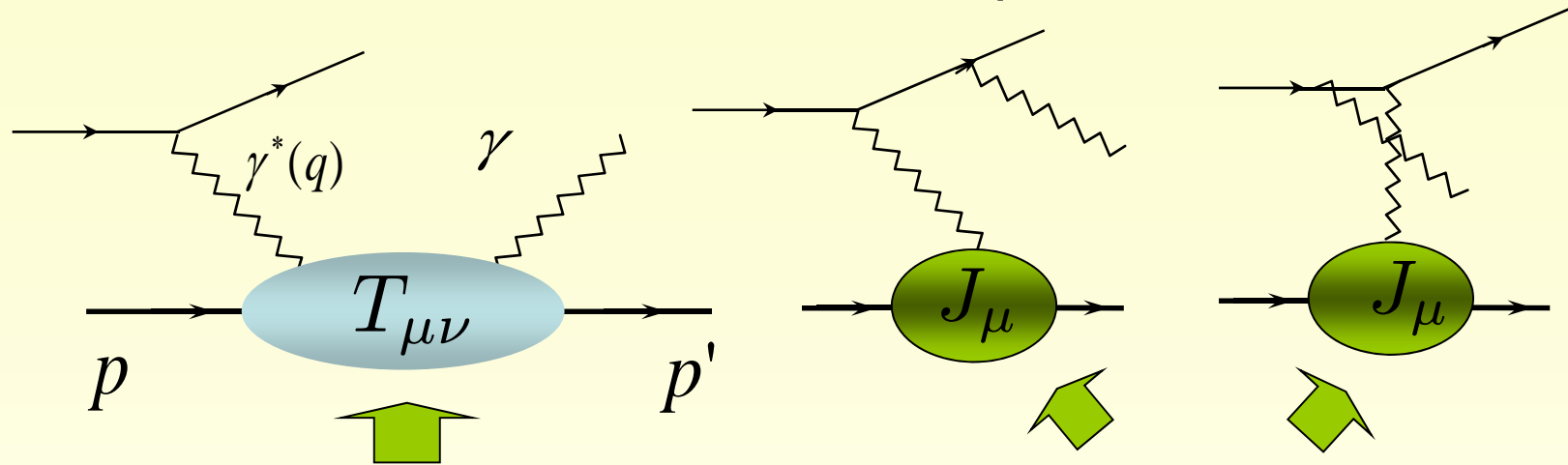
$$x_{Bj} = \frac{Q^2}{2P_1 \cdot q_1} \approx \frac{2\xi}{1 + \xi},$$

$$y = \frac{P_1 \cdot q_1}{P_1 \cdot k},$$

$$\Delta^2 = t \text{ (fixed, small),}$$

$$Q^2 = -q_1^2 (> 1\text{GeV}^2),$$

interference of *DVCS* and *Bethe-Heitler* processes



12 Compton form factors $\mathcal{H}, \mathcal{E}, \tilde{\mathcal{H}} \dots$ (helicity amplitudes) elastic form factors F_1, F_2

$$|\mathcal{T}_{\text{BH}}|^2 = \frac{e^6 (1 + \epsilon^2)^{-2}}{x_{\text{Bj}}^2 y^2 \Delta^2 \mathcal{P}_1(\phi) \mathcal{P}_2(\phi)} \left\{ c_0^{\text{BH}} + \sum_{n=1}^2 c_n^{\text{BH}} \cos(n\phi) \right\}, \quad \text{exactly known (LO, QED)}$$

$$|\mathcal{T}_{\text{DVCS}}|^2 = \frac{e^6}{y^2 Q^2} \left\{ c_0^{\text{DVCS}} + \sum_{n=1}^2 [c_n^{\text{DVCS}} \cos(n\phi) + s_n^{\text{DVCS}} \sin(n\phi)] \right\}, \quad \text{harmonics helicity ampl. } \mathbf{1:1}$$

$$\mathcal{I} = \frac{\pm e^6}{x_{\text{Bj}} y^3 \Delta^2 \mathcal{P}_1(\phi) \mathcal{P}_2(\phi)} \left\{ c_0^{\mathcal{I}} + \sum_{n=1}^3 [c_n^{\mathcal{I}} \cos(n\phi) + s_n^{\mathcal{I}} \sin(n\phi)] \right\}. \quad \text{harmonics helicity ampl.}$$

access of CFFs from measurements:

twist	sector	harmonics in \mathcal{I}				extraction of CFFs	P of Q^{-P}	Δ_{\perp}^l behavior	
	\mathcal{C} 's	unp	LP	TP_x	TP_y			unp, LP	TP
two	$\Re\mathcal{C}(\mathcal{F}), \Delta\mathcal{C}(\mathcal{F})$	c_1, c_0	c_1, c_0	c_1, c_0	$s_1, -$	over compl.	1,2	1,0	0,1
	$\Im\mathcal{C}(\mathcal{F}), \Delta\mathcal{C}(\mathcal{F})$	$s_1, -$	$s_1, -$	$s_1, -$	c_1, c_0	over compl.	1,2	1,0	0,1
three	$\Re\mathcal{C}(\mathcal{F}^{\text{eff}})$	c_2	c_2	c_2	s_2	complete	2	2	1
	$\Im\mathcal{C}(\mathcal{F}^{\text{eff}})$	s_2	s_2	s_2	c_2	complete	2	2	1
two	$\Re\mathcal{C}_T(\mathcal{F}_T)$	c_3	-	-	-	$1 \times \Re$ of 4	1	3	2
	$\Im\mathcal{C}_T(\mathcal{F}_T)$	-	s_3	s_3	c_3	$3 \times \Im$ of 4	1	3	2

three possible nucleon polarization + electron/positron beam + neglecting transversity allows to access imaginary and real part of

$$\mathcal{F} = \{\mathcal{H}, \mathcal{E}, \tilde{\mathcal{H}}, \tilde{\mathcal{E}}\}$$

$$\mathcal{F}^3 = \{\mathcal{H}^3, \mathcal{E}^3, \tilde{\mathcal{H}}^3, \tilde{\mathcal{E}}^3\}$$

twist-three offers access to quark-gluon-quark correlations

transversity arises at NLO from gluons at twist-two or at LO as a twist-four effect

$$\mathcal{F}_T = \mathcal{O}(\alpha_s, 1/Q^2)$$

relations among **harmonics** and **GPDs** are based on $1/Q$ expansion:
 (all harmonics are expressed by twist-2 and -3 GPDs)

[Diehl et. al (97)
 Belitsky, DM, Kirchner (01)]

$$\begin{cases} c_1 \\ s_1 \end{cases}^{\mathcal{I}} \propto \frac{\Delta}{Q} \text{tw-2(GPDs)} + O(1/Q^3), \quad c_0^{\mathcal{I}} \propto \frac{\Delta^2}{Q^2} \text{tw-2(GPDs)} + O(1/Q^4),$$

$$\begin{cases} c_2 \\ s_2 \end{cases}^{\mathcal{I}} \propto \frac{\Delta^2}{Q^2} \text{tw-3(GPDs)} + O(1/Q^4), \quad \begin{cases} c_3 \\ s_3 \end{cases}^{\mathcal{I}} \propto \frac{\Delta \alpha_s}{Q} (\text{tw-2})^{\mathcal{T}} + O(1/Q^3),$$

$$c_0^{\text{CS}} \propto (\text{tw-2})^2, \quad \begin{cases} c_1 \\ s_1 \end{cases}^{\text{CS}} \propto \frac{\Delta}{Q} (\text{tw-2}) (\text{tw-3}), \quad \begin{cases} c_2 \\ s_2 \end{cases}^{\text{CS}} \propto \alpha_s (\text{tw-2}) (\text{tw-2})^{\text{GT}}$$

setting up the **perturbative framework:**

[Belitsky, DM (97);

Mankiewicz et. al (97); Ji,

✓ **twist-two** coefficient functions at **next-to-leading** order [Osborne (98); Pire et. al (11)]

✓ evolution kernels at **next-to-leading** order [Belitsky, DM, Freund (01)]

✓ **next-to-next-to-leading** order in a specific conformal subtraction scheme [KMP-K & Schaefer 06]

✓ **twist-three** including quark-gluon-quark correlation at LO [Anikin, Teryaev, Pire (00); Belitsky DM (00); Kivel et. al]

✓ partially, **twist-three** sector at **next-to-leading** order [Kivel, Mankiewicz (03)]

✓ 'target mass corrections' (not well understood) [Belitsky DM (01)]

✓ **twist-four** sector [Braun, Manashov (11)]

Can one 'measure' GPDs?

- **CFF** given as **GPD convolution**:

$$\begin{aligned} \mathcal{H}(\xi, t, Q^2) &\stackrel{\text{LO}}{=} \int_{-1}^1 dx \left(\frac{1}{\xi - x - i\epsilon} - \frac{1}{\xi + x - i\epsilon} \right) H(x, \eta = \xi, t, Q^2) \\ &\stackrel{\text{LO}}{=} i\pi H^-(x = \xi, \eta = \xi, t, Q^2) + \text{PV} \int_0^1 dx \frac{2x}{\xi^2 - x^2} H^-(x, \eta = \xi, t, Q^2) \end{aligned}$$

- $H(x, x, t, Q^2)$ viewed as "**spectral function**" (s-channel cut):

$$H^-(x, x, t, Q^2) \equiv H(x, x, t, Q^2) - H(-x, x, t, Q^2) \stackrel{\text{LO}}{=} \frac{1}{\pi} \Im \mathcal{F}(\xi = x, t, Q^2)$$

- **CFFs** satisfy '**dispersion relations**'
(not the physical ones, threshold ξ_0 set to 0)

[Frankfurt et al (97)
Chen (97)
Terayev (05)
KMP-K (07)
Diehl, Ivanov (07)]

$$\Rightarrow \Re \mathcal{F}(\xi, t, Q^2) = \frac{1}{\pi} \text{PV} \int_0^1 d\xi' \left(\frac{1}{\xi - \xi'} \mp \frac{1}{\xi + \xi'} \right) \Im \mathcal{F}(\xi', t, Q^2) + \mathcal{C}(t, Q^2)$$

[Terayev (05)]

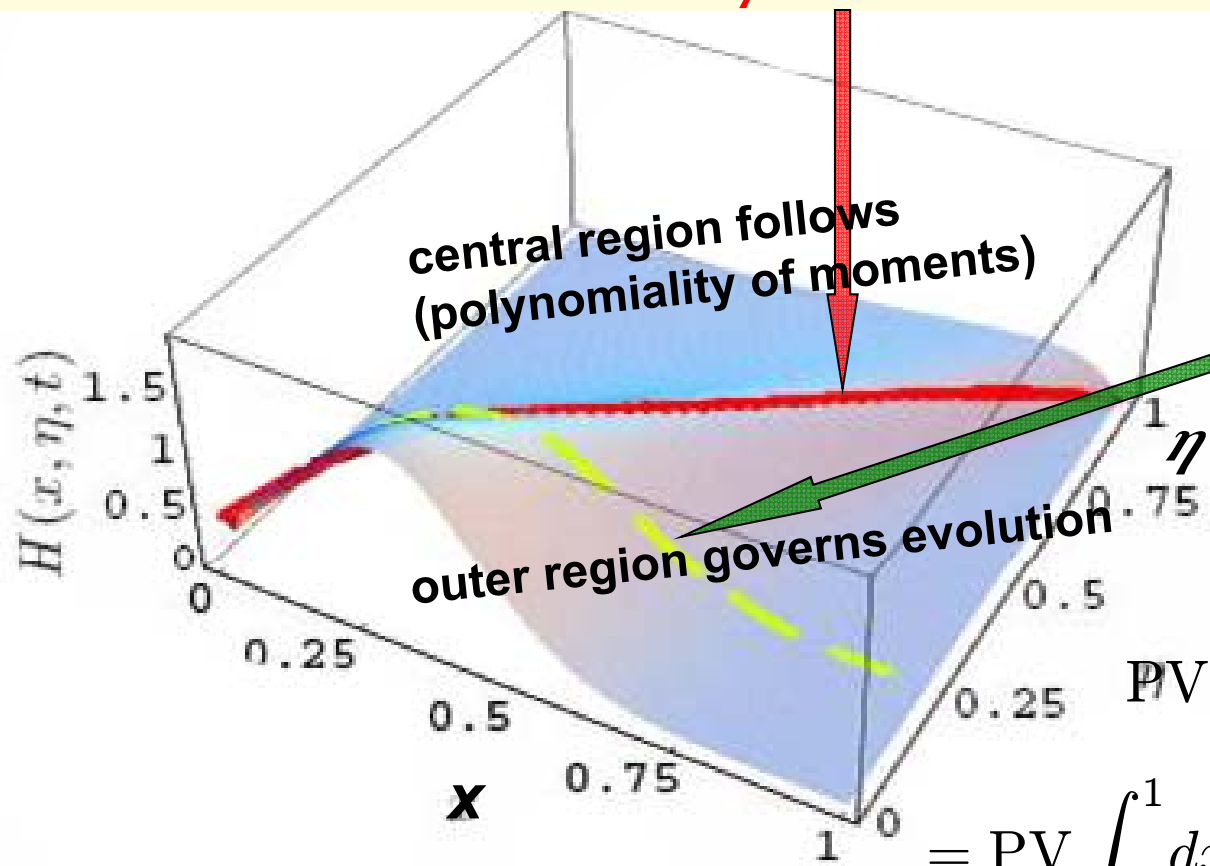
⇒ **access** to the **GPD** on the **cross-over line** $\eta = x$ (at LO)

Modeling & Evolution

outer region governs the evolution at the cross-over trajectory

$$\mu^2 \frac{d}{d\mu^2} H(x, x, t, \mu^2) = \int_x^1 \frac{dy}{x} V(1, x/y, \alpha_s(\mu)) H(y, x, \mu^2)$$

GPD at $\eta = x$ is 'measurable' (LO)



net contribution of outer + central region is governed by a sum rule:

$$\text{PV} \int_0^1 dx \frac{2x}{\eta^2 - x^2} H^-(x, \eta, t) = \text{PV} \int_0^1 dx \frac{2x}{\eta^2 - x^2} H^-(x, x, t) +_8 C(t)$$

Strategies to analyze DVCS data

(ad hoc) modeling: **VGG** code [Goeke et. al (01) based on Radyuskin's DDA]
(first decade) BKM model [Belitsky, Kirchner, DM (01) based on RDDA]
'aligned jet' model [Freund, McDermott, Strikman (02)]
Goloskokov/Kroll (05) based on RDDA (pinned down by DVMP)
'dual' model [Polyakov, Shuvaev 02; Guzey, Teckentrup 06; Polyakov 07]
" -- " [KMP-K (07) in MBs-representation]
polynomials [Belitski et al. (98), Liuti et. al (07), Moutarde (09)]

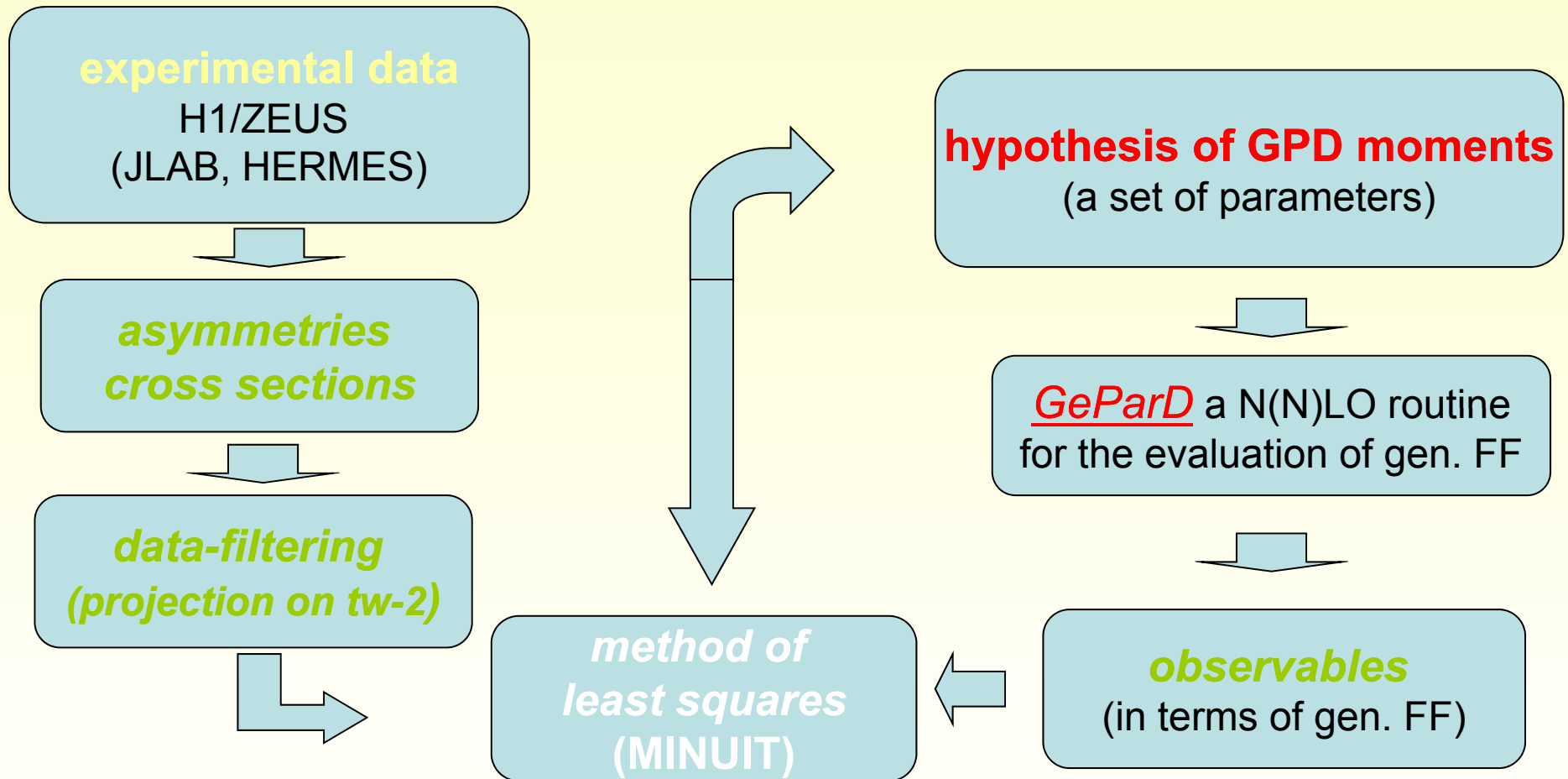
dynamical models: not applied [Radyuskin et.al (02); Tiburzi et.al (04); Hwang DM (07)]...
(respecting Lorentz symmetry)

flexible models: any representation by including *unconstrained* degrees of freedom
(for fits) **KMP-K (07/08)** for H1/ZEUS in MBs-integral-representation

Extracting CFFs from data: real and imaginary part

- i. CFF extraction with formulae [BMK (01), HALL-A (06)]
fits [Guidal, Moutarde (08...)]
neural networks [KMS (11) neural networks]
- ii. 'dispersion integral' fits [KMP-K (08), KM (08...)]
- iii. flexible GPD modeling [KM (08...)]
- vi. model comparisons **Goldstein et al. (11)** (no sea, giving up polynomiality)
Goloskokov/Kroll (07) model based on RDDA

Ready for flexible GPD model fits?



YES for small x and **we don't use it** for fixed target kinematics

- reasonable well motivated hypotheses of GPDs (moment) must be known first
- many parameters – Is a least square fit an appropriate strategy?
- some code writing is left

DVCS fits to H1 and ZEUS data

DVCS cross section measured at small $x_{Bj} \approx 2\xi = \frac{2Q^2}{2W^2 + Q^2}$

predicted by

$$\frac{d\sigma}{dt}(W, t, Q^2) \approx \frac{4\pi\alpha^2}{Q^4} \frac{W^2\xi^2}{W^2 + Q^2} \left[|\mathcal{H}|^2 - \frac{\Delta^2}{4M_p^2} |\mathcal{E}|^2 + |\tilde{\mathcal{H}}|^2 \right] (\xi, t, Q^2) \Big|_{\xi = \frac{Q^2}{2W^2 + Q^2}}$$

suppressed contributions $\ll 0.05 \gg$ relative $O(\xi)$

- LO data could not be described before **2008**
- NLO works with ad hoc GPD models [**Freund, McDermott (02)**]
results strongly depend on employed PDF parameterization

 **do a simultaneous fit to DIS and DVCS** [**KMP-K (07)**]

 **use flexible GPD models in a two-step fit** [**KM (09)**]

effective functional form at small x :

PDFs: $q^{\text{sea}}(\xi, Q) = n(Q)\xi^{-\alpha(Q)}, \quad \alpha \gtrsim 1, \quad F^{\text{sea}}(0) = 1$

GPDs: $H^{\text{sea}} = r(\eta/x = 1, Q)F^{\text{sea}}(t)\xi^{\alpha'(t, Q)}q^{\text{sea}}(\xi, Q)$

transverse distribution

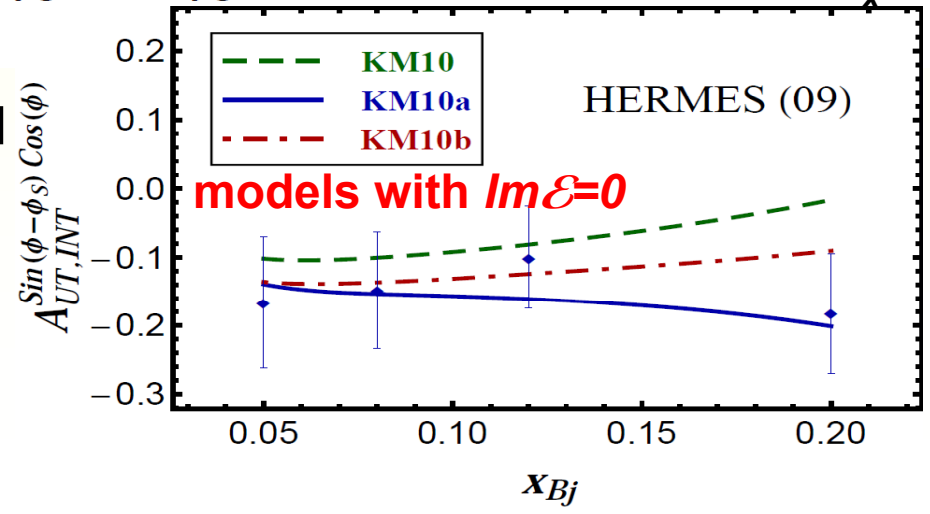
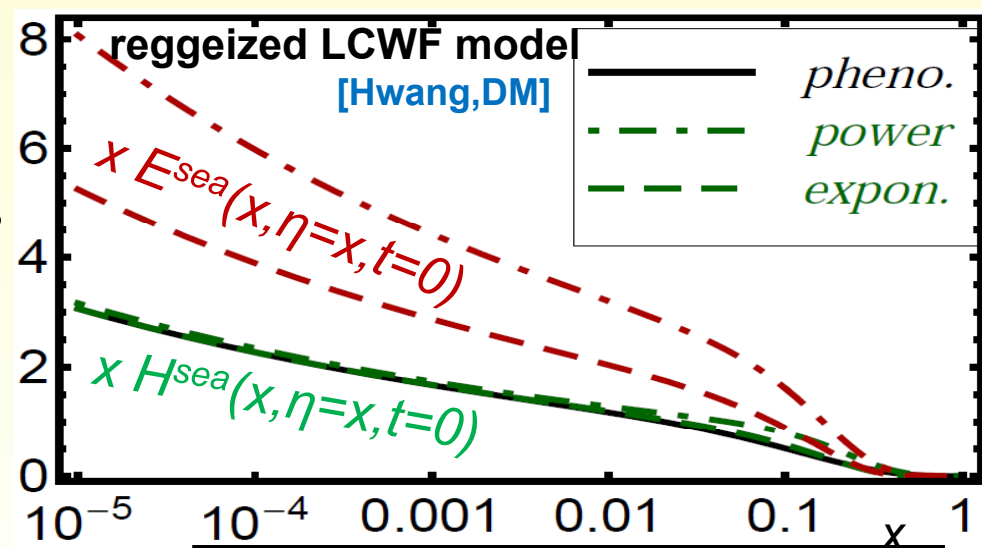
? $E(\xi, \xi, t, Q)$ *skewness*

- not seen in Regge phenomenology
- might be sizeable in instanton models
- reggeized spectator quark models
- pQCD suggests ‘pomeron’ intercept
- however, B vanishes asymptotically
- so far E is “not seen” in DVCS

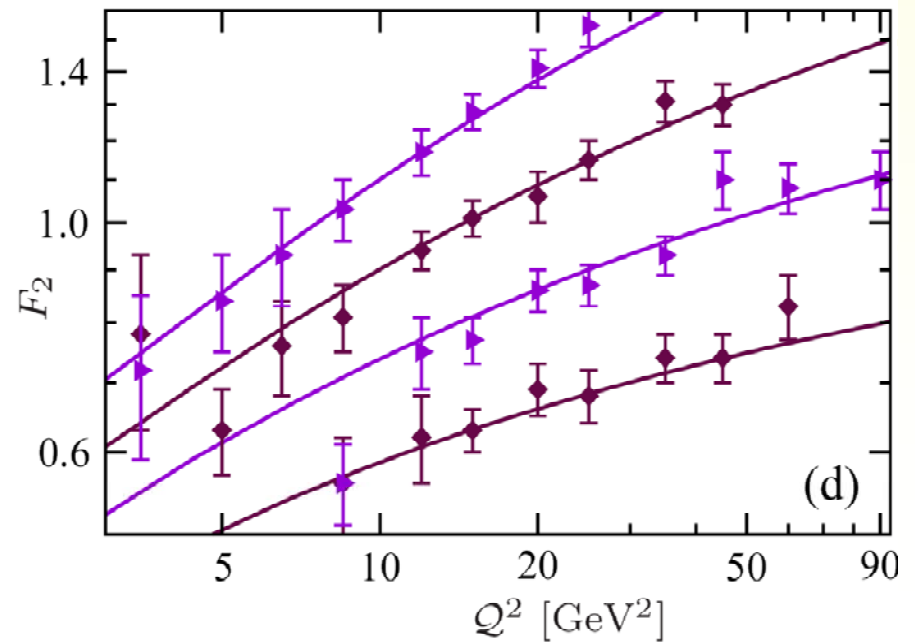
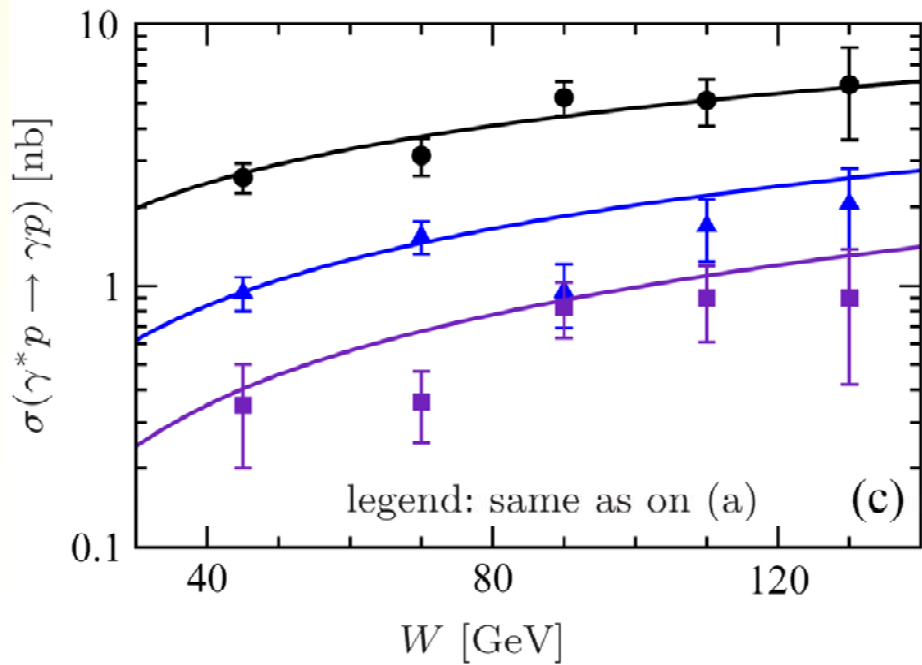
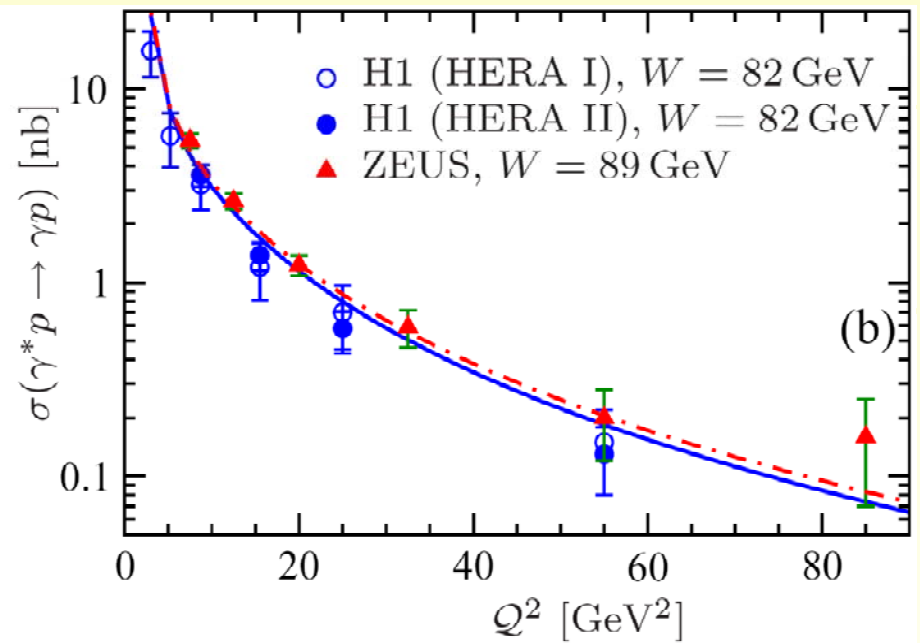
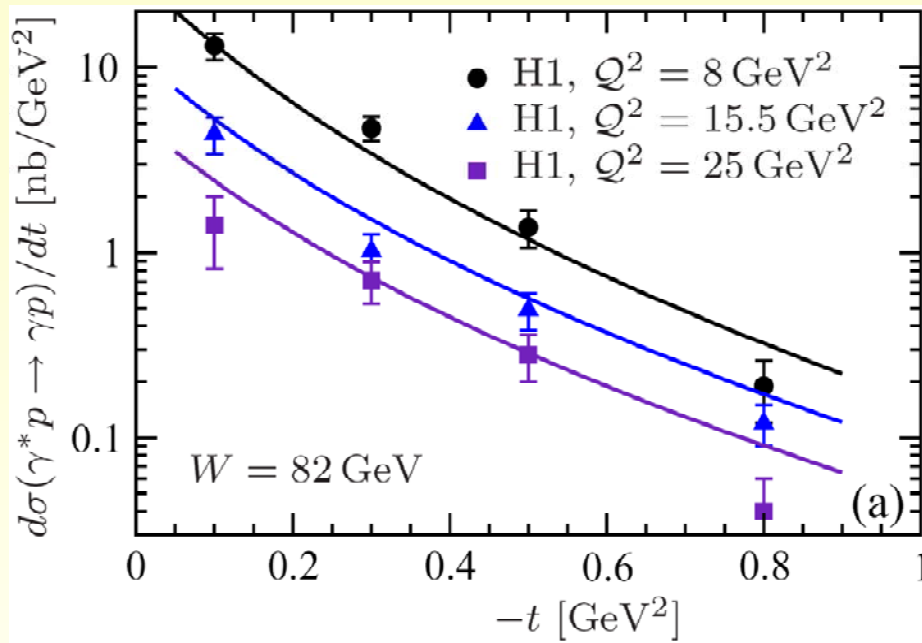
qualitative understanding of E is needed (not only for Ji’s spin sum rule)

$$B = \int_0^1 dx x E(x, \eta, t, Q)$$

transverse target spin asymmetry is **sensitive to E and sizeable at EIC**



good DVCS fits at LO, NLO, and NNLO with flexible GPD ansatz



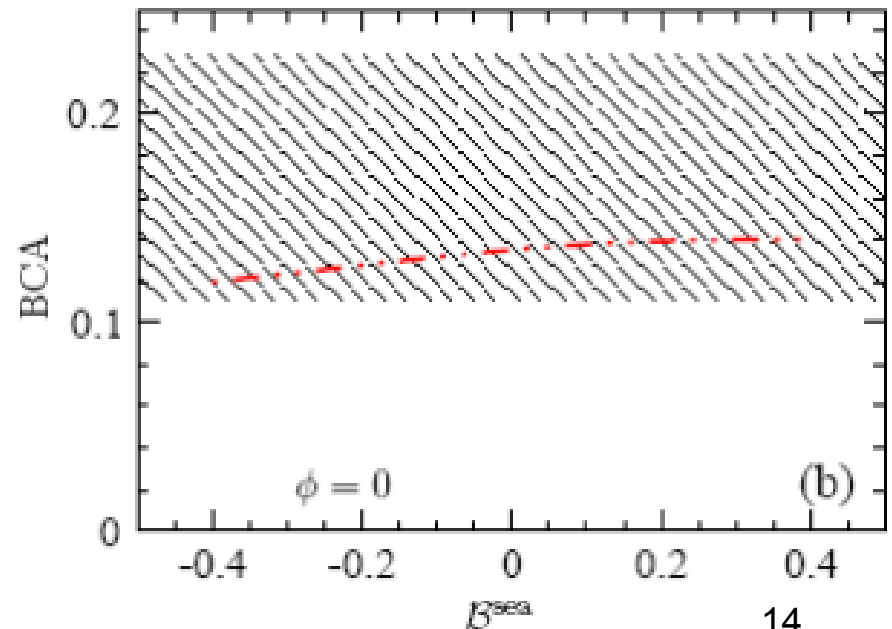
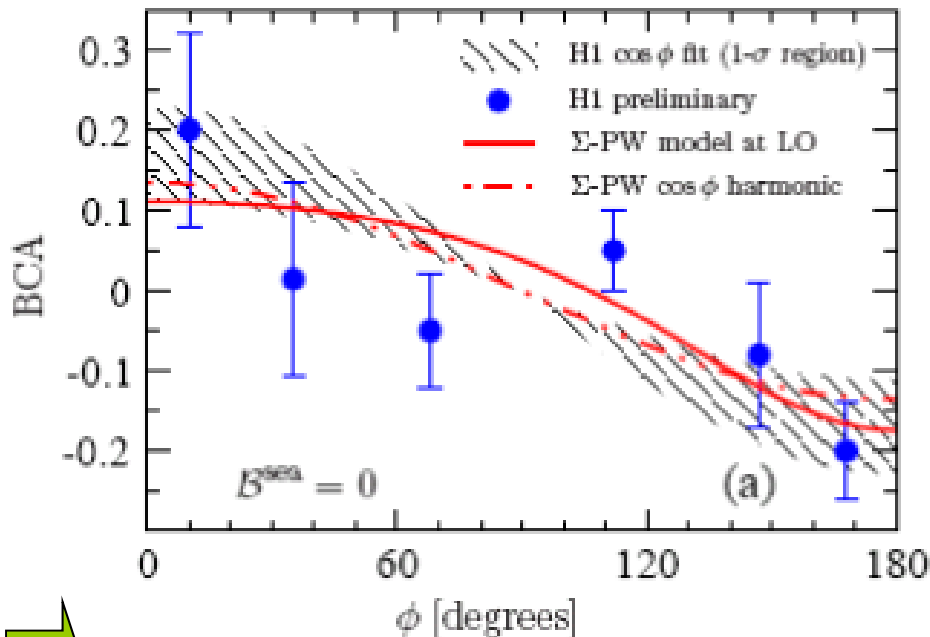
Beam charge asymmetry

$$BCA = \frac{d\sigma_{e^+} - d\sigma_{e^-}}{d\sigma_{e^+} + d\sigma_{e^-}} = \frac{\mathcal{T}_{\text{Interference}}}{|\mathcal{T}_{\text{BH}}|^2 + |\mathcal{T}_{\text{DVCS}}|^2}$$

$$\propto F_1(t) \Re \mathcal{H} + \frac{|t|}{4M^2} F_2(t) \Re \mathcal{E}$$

the unknown in Ji's nucleon spin sum rule

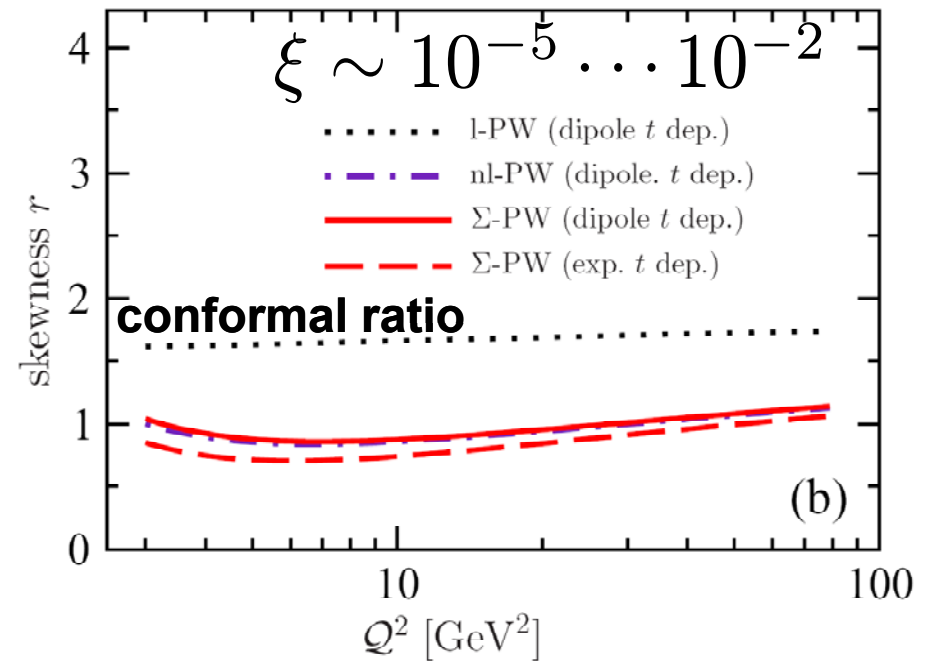
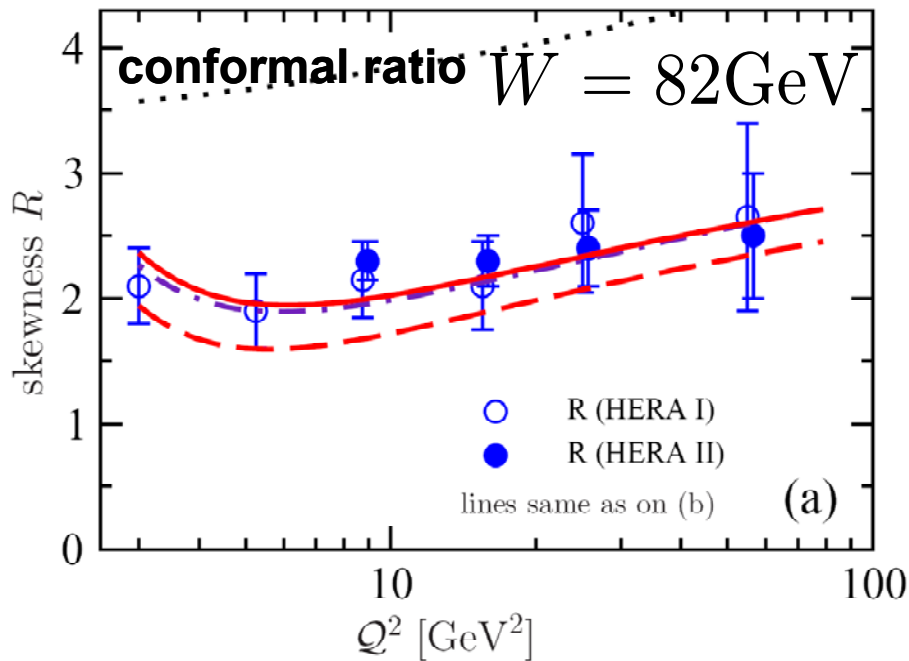
- set $E_{\text{sea}} \propto H_{\text{sea}}$ use *anomalous gravitomagnetic moment* $B_{\text{sea}} = \int_0^1 dx x E_{\text{sea}}$ as parameter



unfortunately, H1 data do not allow to access B_{sea}

quark skewness ratio from DVCS fits @ LO

$$R = \frac{\Im m A_{\text{DVCS}}}{\Im m A_{\text{DIS}}} \stackrel{\text{LO}}{=} \frac{H(\xi, \xi)}{H(2\xi, 0)} \approx 2^\alpha r \quad r = \frac{H(\xi, \xi)}{H(\xi, 0)}$$



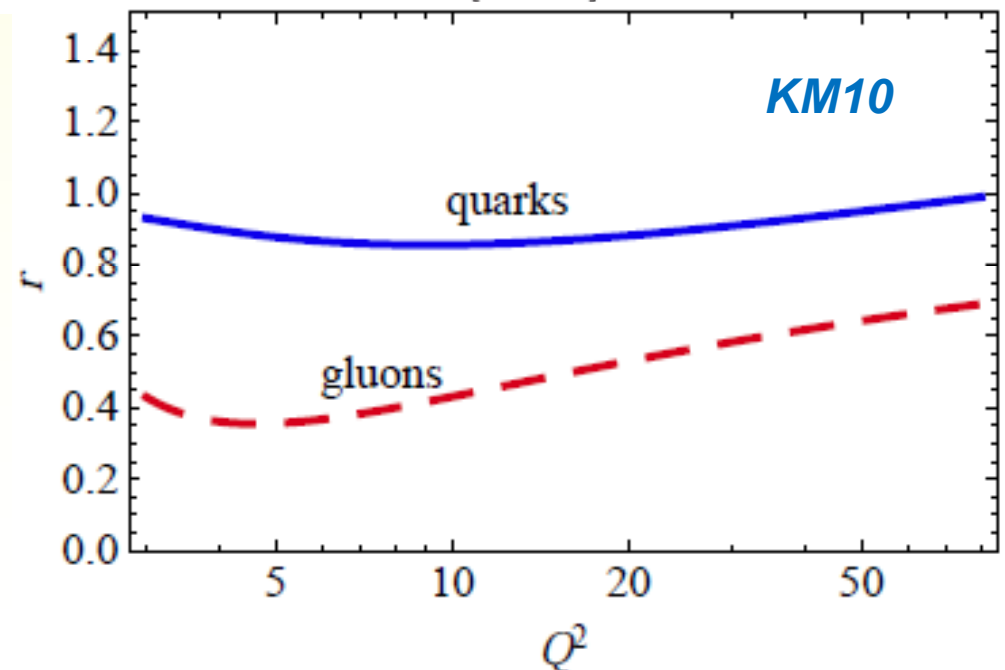
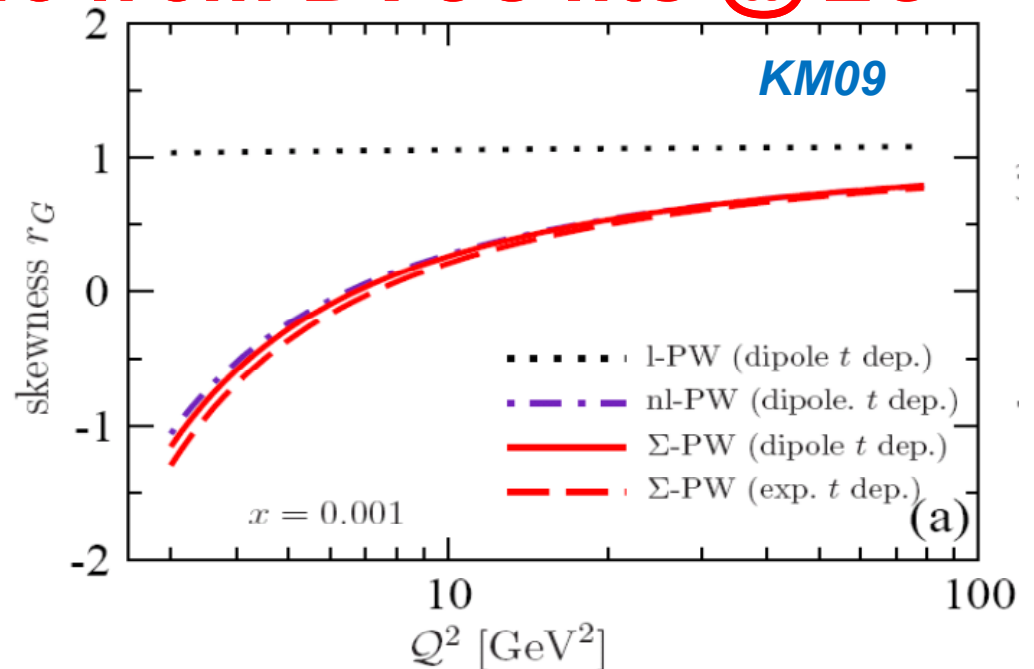
- @LO the conformal ratio $r_{\text{con}} = \frac{2^\alpha \Gamma(3/2+\alpha)}{\Gamma(3/2)\Gamma(2+\alpha)}$ is ruled out for sea quark GPD
- a generically zero-skewness effect over a large Q^2 lever arm
- scaling violation consistent with pQCD prediction
- this zero-skewness effect is non-trivial to realize in conformal space (SO(3) sibling poles are required)

Gluon skewness ratio from DVCS fits @ LO

conformal ratio

$$r_{\text{con}}^G = \frac{2^{1+\alpha} \Gamma(3/2+\alpha)}{\Gamma(3/2)\Gamma(3+\alpha)} \approx 1$$

- imposed in all popular models
- accessible due to evolution
- @LO the gluonic r -ratio is smaller than the conformal one
- negative value is an artifact (still a too rigid model)
- gluonic r -ratio @LO from ρ electroproduction is 0.1...0.5
- qualitatively consistent with DVCS findings

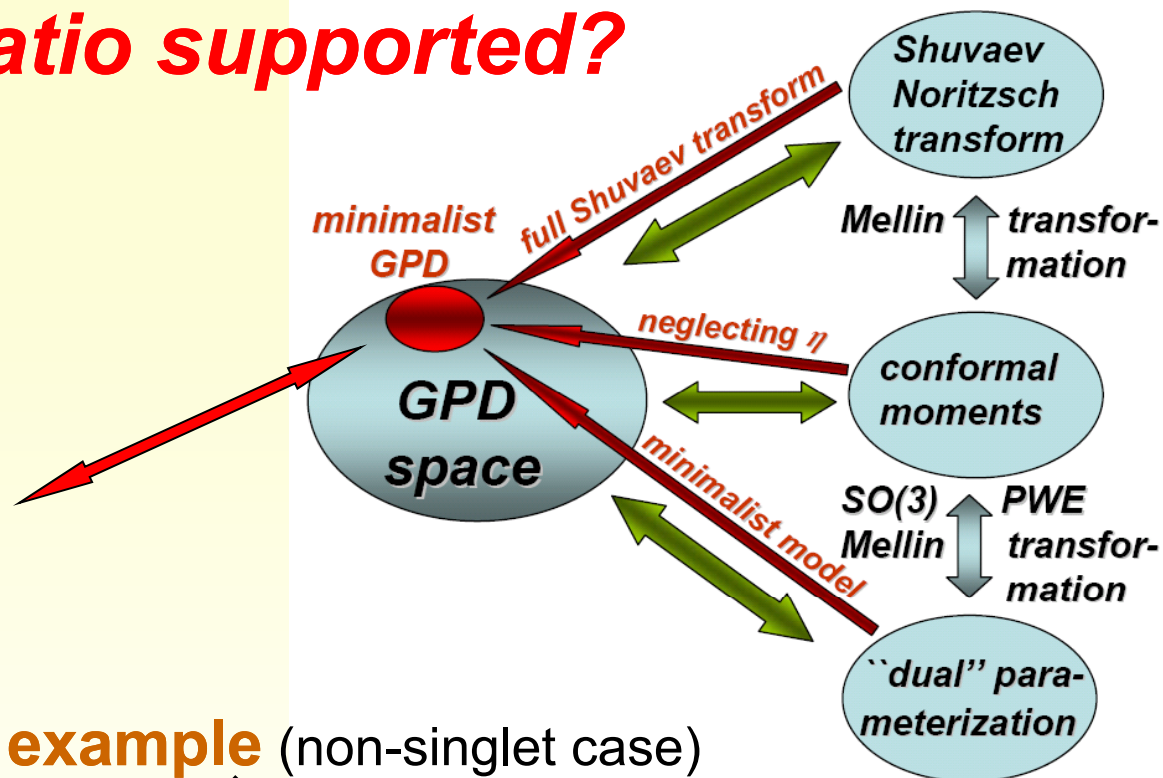


Is the conformal ratio supported?

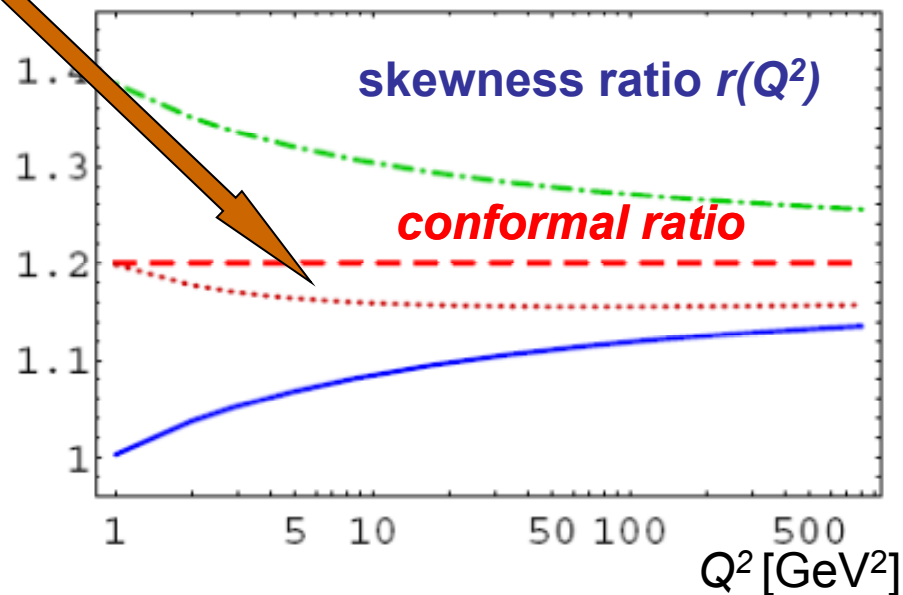
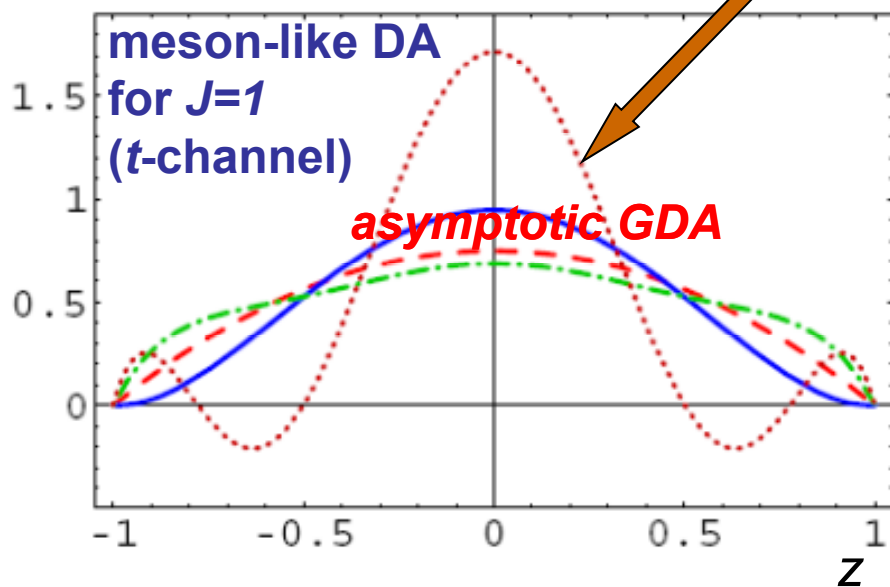
$$r = \frac{H(x, x, t=0, Q^2)}{q(x, Q^2)}$$

“erroneous small x-claim”

$$r_{\text{con}} = \frac{2^\alpha \Gamma(3/2 + \alpha)}{\Gamma(3/2) \Gamma(2 + \alpha)}$$



a counter example (non-singlet case)



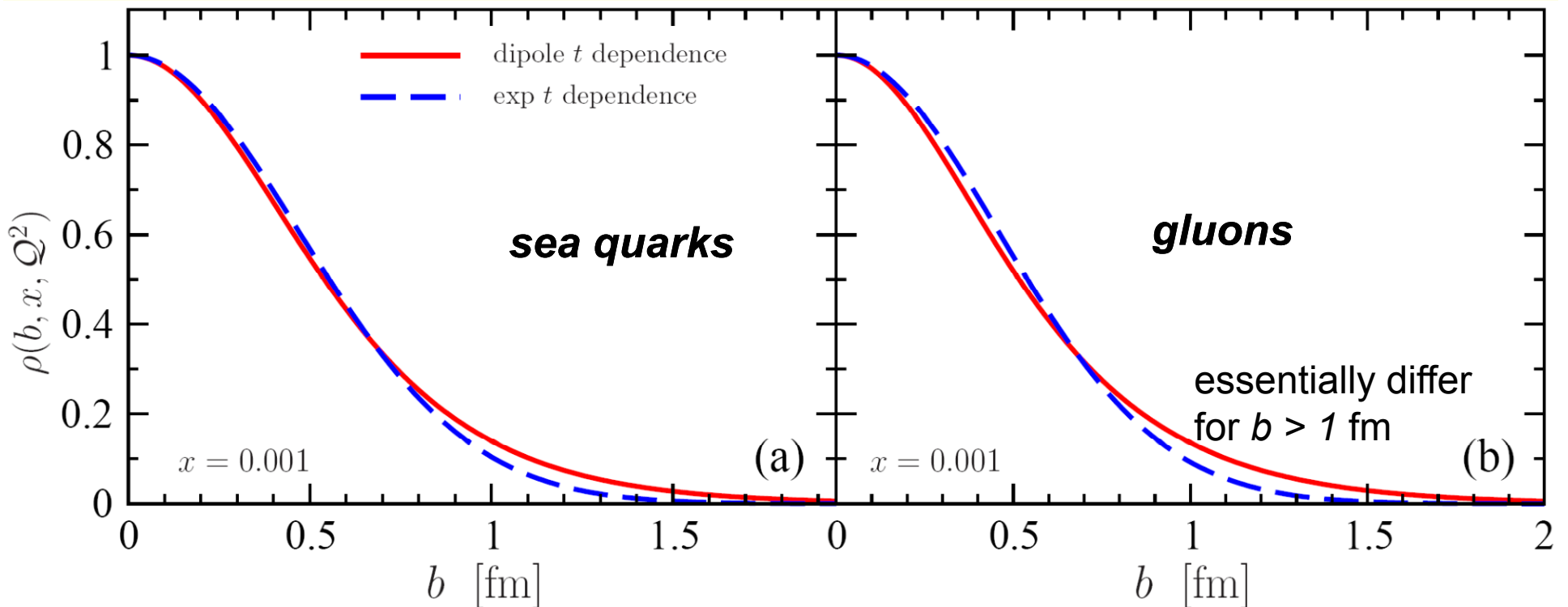
- CFF H posses ``pomeron behavior'' $\xi^{-\alpha(Q) - \alpha'(Q)t}$

- ✓ α increases with growing Q^2
- ✓ α' decreases growing Q^2

- t -dependence: exponential shrinkage is disfavored ($\alpha' \approx 0$)
- dipole shrinkage is visible ($\alpha' \approx 0.15$ at $Q^2=4 \text{ GeV}^2$)

- (normalized) profile functions

$$\rho \propto \int d^2 \vec{\Delta}_{\perp} e^{i\vec{b} \cdot \vec{\Delta}_{\perp}} H(x, 0, t = -\vec{\Delta}_{\perp}^2)$$



data set for unpolarized proton target used for KM09/10 fits

- H1/ZEUS 98 [σ , $d\sigma/dt$] +1x6 [BCA(φ)] $\langle\langle x \rangle\rangle \approx 10^{-3}$, $\langle|t|\rangle \leq 0.8 \text{ GeV}^2$
 $\langle\langle Q^2 \rangle\rangle \approx 8 \text{ GeV}^2$
- HERMES(02) 12+3 [BSA, $\sin(\varphi)$]
- HERMES(08) 12x2 [BCA, $\cos(0 \varphi)$, $\cos(\varphi)$] $0.05 \leq \langle x \rangle \leq 0.2$, $\langle|t|\rangle \leq 0.4 \text{ GeV}^2$
 12x2 [$\cos(2 \varphi)$, $\cos(3 \varphi)$] $\langle\langle Q^2 \rangle\rangle \approx 2.5 \text{ GeV}^2$
- HERMES(09,10) BSA and BCA data (included in KM10 fits)
- CLAS(07) 12x12 [BSA(φ)] $0.14 \leq \langle x \rangle \leq 0.35$, $\langle|t|\rangle \leq 0.3 \text{ GeV}^2$
 40x12 [BSA(φ)] (large $|t|$ or bad sta.) $\langle\langle Q^2 \rangle\rangle \approx 1.8 \text{ GeV}^2$
- HALL A(06) 12x24 [$\Delta\sigma(\varphi)$] $\langle x \rangle = 0.36$, $\langle|t|\rangle \leq 0.33 \text{ GeV}^2$
 3x24 [$\sigma(\varphi)$] $\langle\langle Q^2 \rangle\rangle \approx 1.8 \text{ GeV}^2$

How to analyze φ dependence?

- fit within assumed functional form **[CLAS(07)]**
- likelihood fit with respect to dominant and higher harmonics **[HERMES]**
- utilize Fourier transform (with or without additional weight) **[BMK(01)]**

 equivalent results for CLAS data with small stat. errors

Dispersion relation fits to unpolarized DVCS

- model of GPD $H(x,x,t)$ within DD motivated ansatz at $Q^2=2 \text{ GeV}^2$

fixed:
rules

PDF normalization \downarrow eff. Reage pole \downarrow large t -counting \downarrow

$$H(x, x, t) = \frac{n r 2^\alpha}{1+x} \left(\frac{2x}{1+x} \right)^{-\alpha(t)} \left(\frac{1-x}{1+x} \right)^b \frac{1}{\left(1 - \frac{1-x}{1+x} \frac{t}{M^2} \right)^p}$$

\uparrow \uparrow \uparrow \uparrow

free:
 r -ratio at small x
sea quarks (taken from LO fits) large x -behavior p -pole mass

$$n = 0.68, \quad r = 1, \quad \alpha(t) = 1.13 + 0.15t/\text{GeV}^2, \quad m^2 = 0.5\text{GeV}^2, \quad p = 2$$

valence quarks

$$n = 1.0, \quad \alpha(t) = 0.43 + 0.85t/\text{GeV}^2, \quad p = 1$$

flexible parameterization of subtraction constant

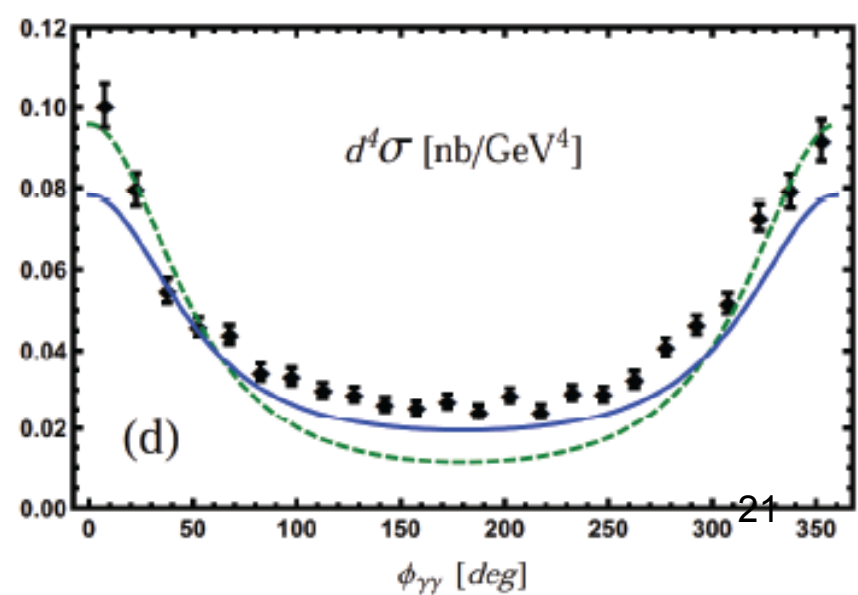
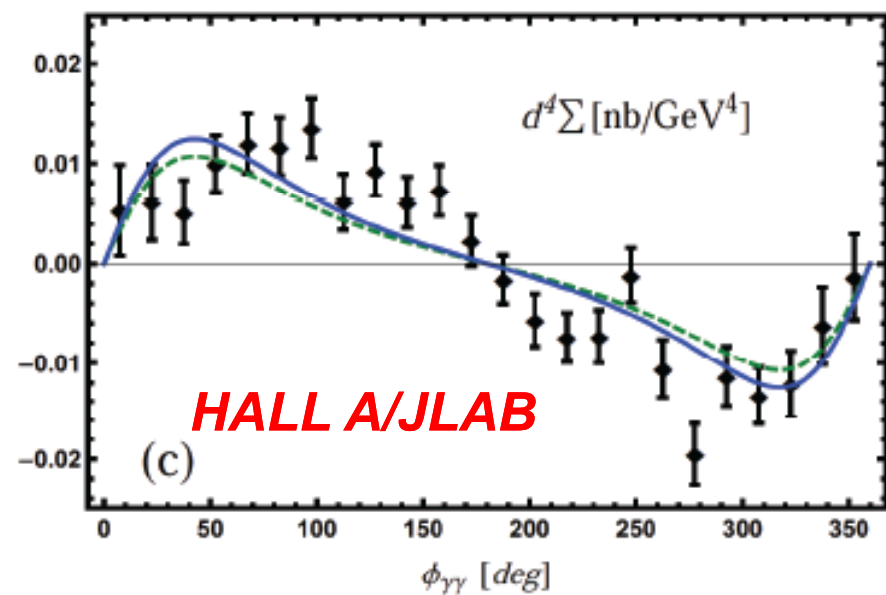
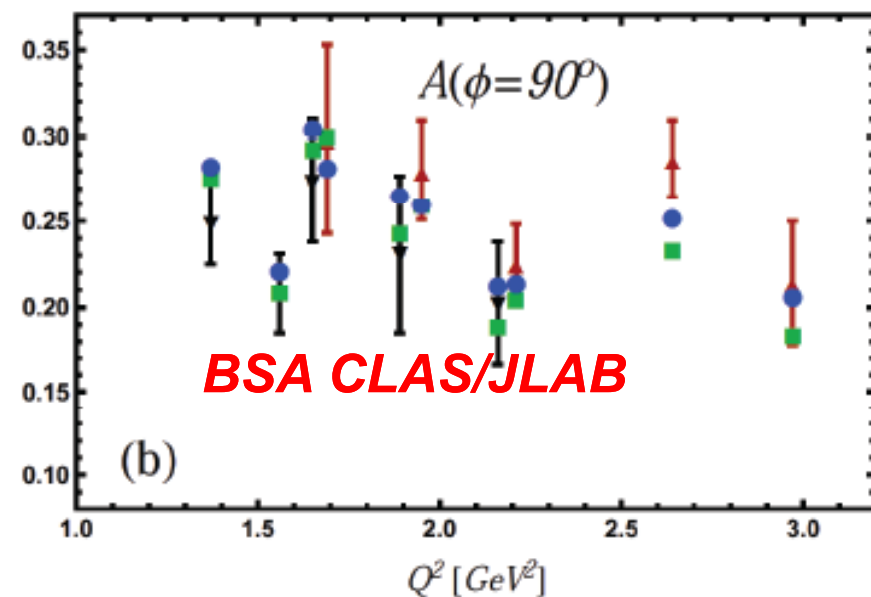
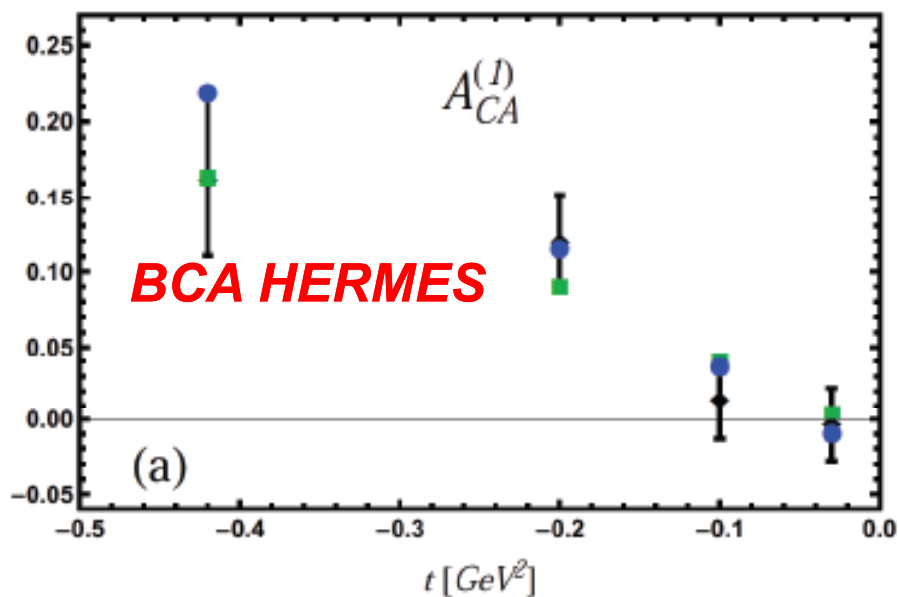
$$\mathcal{D}(t) = \frac{-C}{(1-t/M_c^2)^2}$$

+ pion-pole contribution

36 + 4 data points quality of **global fit** is good

$$\chi^2/\text{d.o.f.} \approx 1^{20}$$

Global GPD fit example: HERMES & JLAB

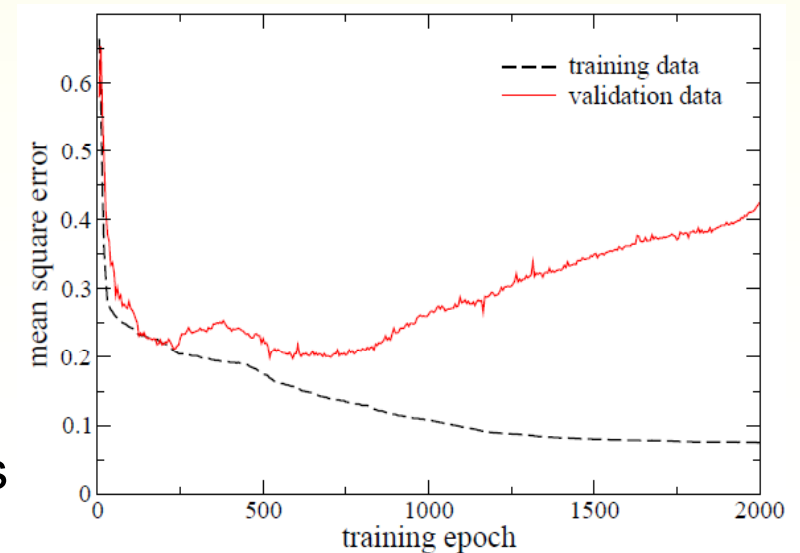
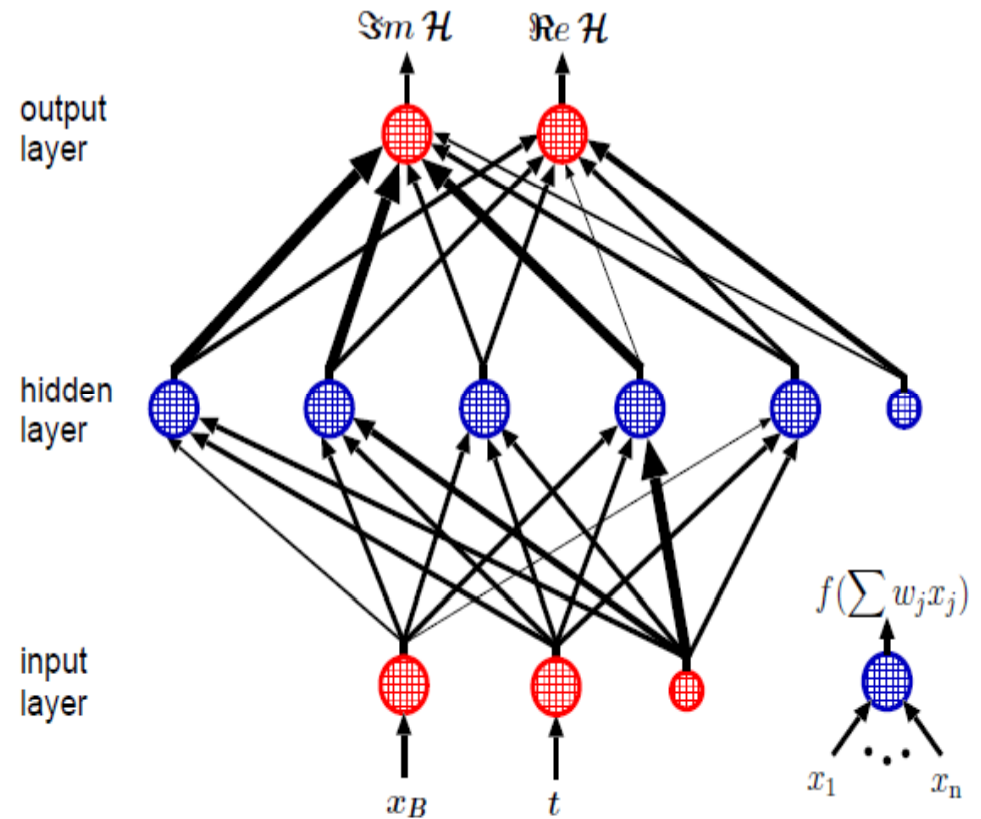


Neural Networks

- kinematical values are represented by the input layer
- propagated through the network, where weights are set randomly
- random values for $Im\mathcal{H}$ and $Re\mathcal{H}$
- calculation of χ^2
- backwards propagation
- adjusting weights so that error decreases
- repeat procedure
- taking next kinematical point

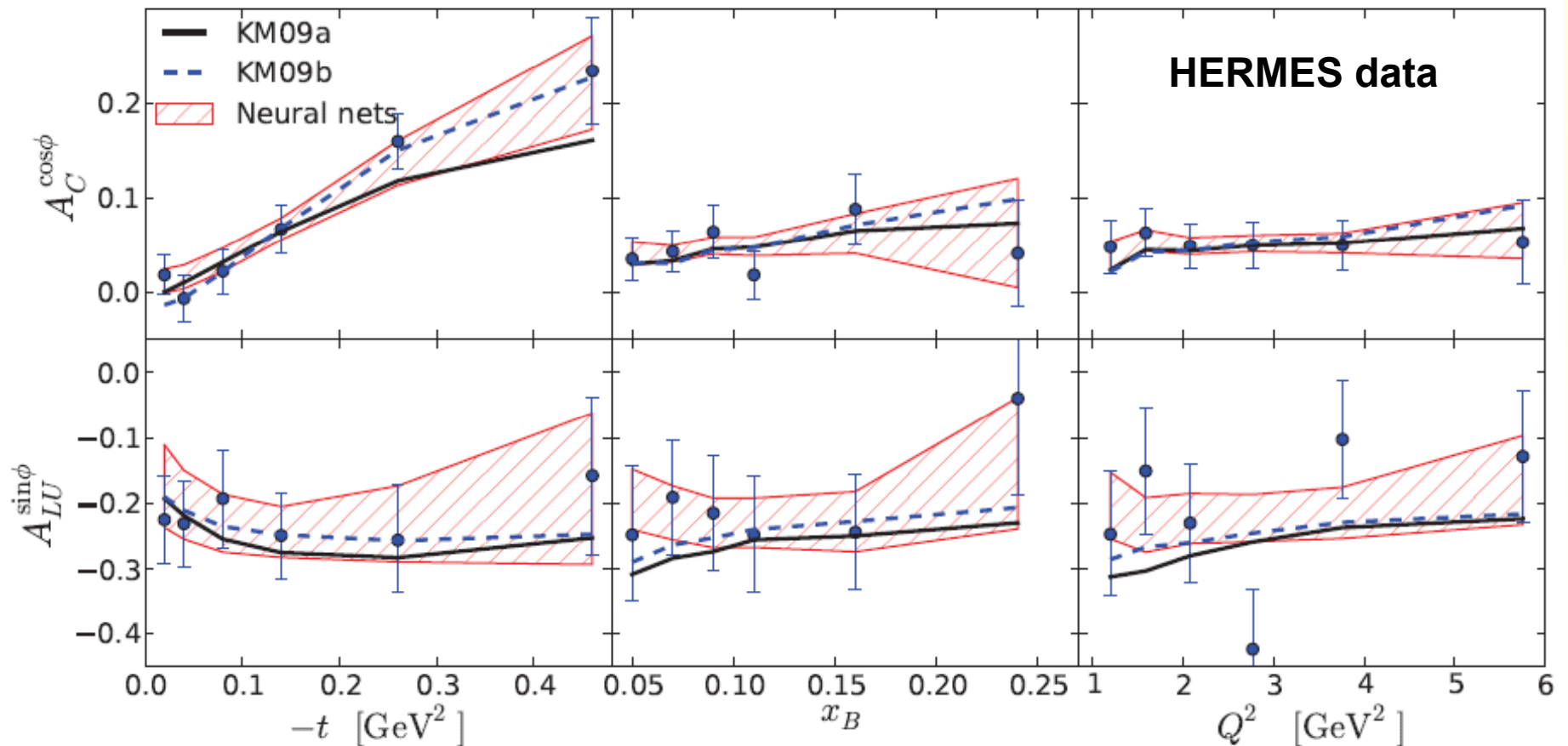
Monte Carlo procedure to propagate errors, i.e., generating a replica data set

avoiding over fitting (fitting to noise), dividing data set, taking a control example if error increases after decreasing – one stops

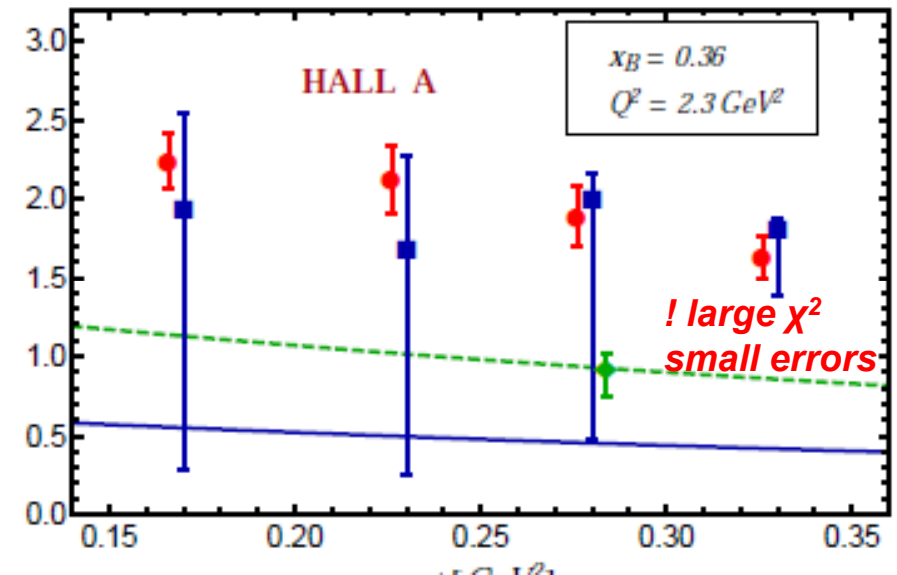
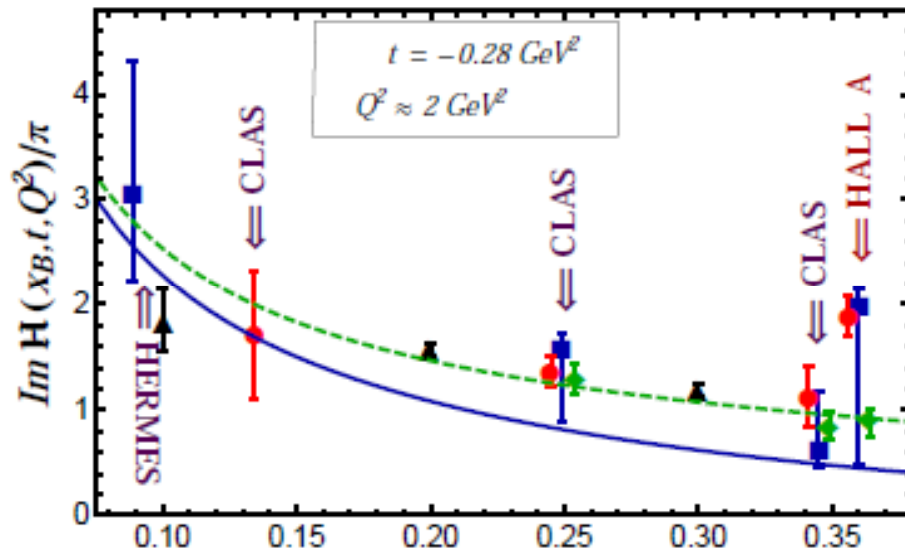


A first use of neural network fits

(ideal) tool for error propagation and quantifying model uncertainties:
so far it is used to access real and imaginary part of \mathcal{H} CFF
results are compatible to dispersion relation fits



KM09 versus CFF fits & large-x “model” fit



GUIDAL

twist-two dominance hypothesis

7 parameter fit to all harmonics of unpolarized cross section

propagated errors + “theoretical” error estimate

GUIDAL

same + longitudinal TSA

Moutarde

H dominance hypothesis within a smeared polynomial expansion

propagated errors + “theoretical” error estimate

KM

neural network within H dominance hypothesis

green (blue) curves (KM09) without (with) HALL A data (ratios)

- reasonable agreement for HERMES and CLAS kinematics
- large x-region and real part remains unsettled
- next step: NN within twist-two dominance hypothesis + ‘dispersion’ integral

How reliable is the twist-two dominance?

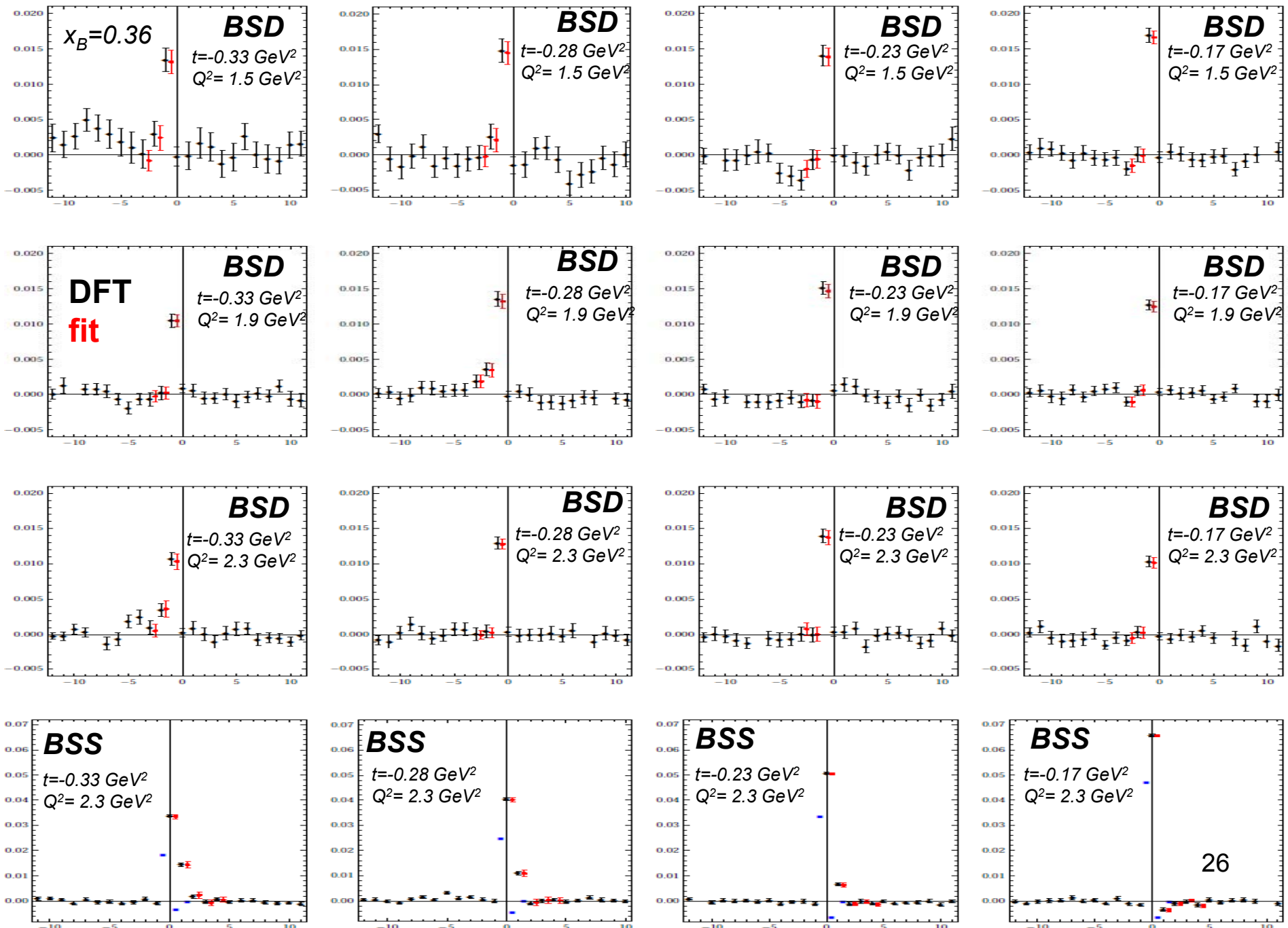
- each interference term harmonic contains twist-2, -3, and transversity (kinematical counting according to BMK + admixture effects [BM (09,10)])
- transversity (photon helicity flip by two units) might be neglected (so far not seen in data, kinematical suppressed by t'/M^2 or t'/Q^2)
- dominant $\sin/\cos(\varphi)$ harmonics of the interference term is contaminated by twist-three term from DVCS
- HALL A : twist-3 effects are hardly visible

➔ ***not conclusive***

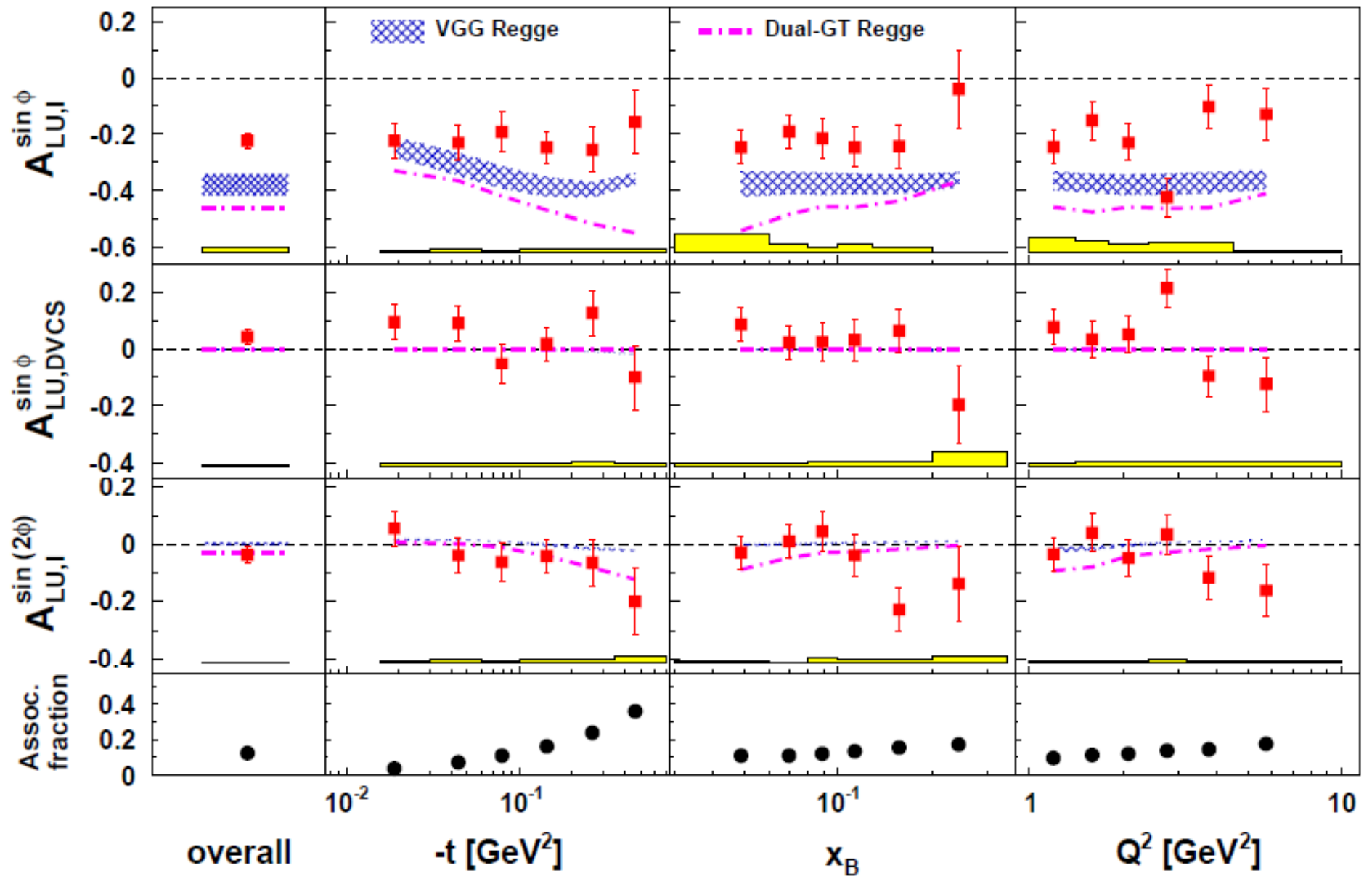
- electron/positron beam is needed for a cleaner separation
- HERMES (09) data might indicate a $\sim 4\%$ overall DVCS twist-3 BSA effect a relative 20% uncertainty for a BSA measurement within electron beam
- HERMES (10) data indicate a zero overall DVCS twist-3 BSA effect

➔ ***there might be a twist-three contamination on some 10% level***

Fourier spectra of HALL A data (12 + 4 cross section measurements versus φ)



Fourier spectra of HERMES (09) data (1996-2005)



KM10 fits to (unpolarized) DVCS

- a hybrid model: three effective SO(3) PWs for sea quarks/gluons
dispersion relations for valence
flexible pion pole contribution
still E GPD is neglected (only D-term)
- framework

leading order, including evolution for sea quarks/ gluons
twist-two dominance hypothesis within CFF convention [BM10]

- data selection (taking moments of azimuthal angle harmonics)

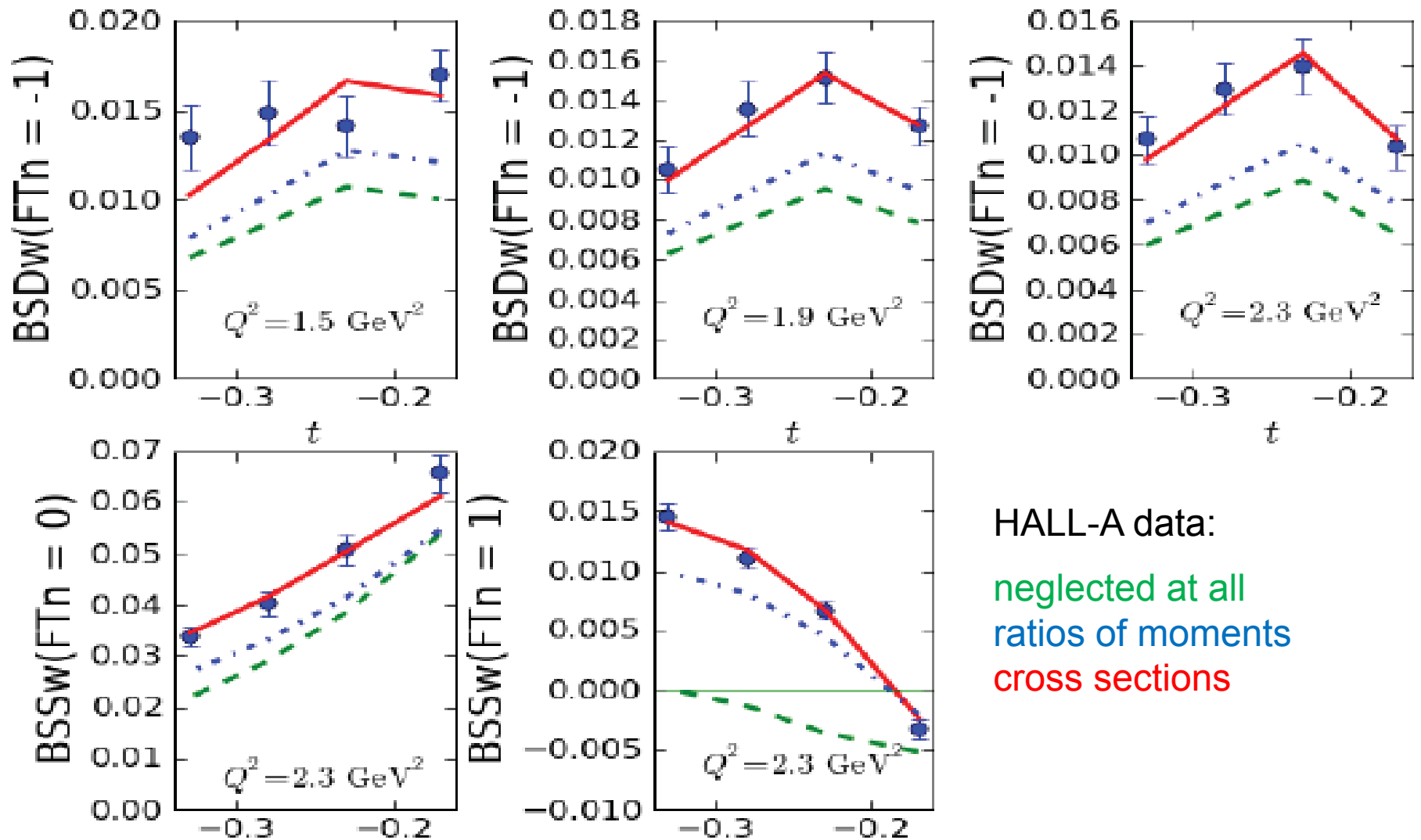
- i. neglecting,
 - ii. forming ratios of moments, or
 - iii. original HALL-A data
- neglecting large $-t$ BSA CLAS data

15 parameter fit, e.g.,
including all HALL-A data

175 data points
 $\chi^2/d.o.f. = 132/165$

```
-----  
MO2S = 0.51 +- 0.02  
SECS = 0.28 +- 0.02  
SECG = -2.79 +- 0.12  
THIS = -0.13 +- 0.01  
THIG = 0.90 +- 0.05  
  Mv = 4.00 +- 3.33 (edge)  
  rv = 0.62 +- 0.06  
  bv = 0.40 +- 0.67  
  C = 8.78 +- 0.98  
  MC = 0.97 +- 0.11  
tMv = 0.88 +- 0.24  
trv = 7.76 +- 1.39  
tbv = 2.05 +- 0.40  
rpi = 3.54 +- 1.77  
Mpi = 0.73 +- 0.37  
-----
```

- results are given as **xs.exe** on <http://calculon.phy.hr/qpd/>



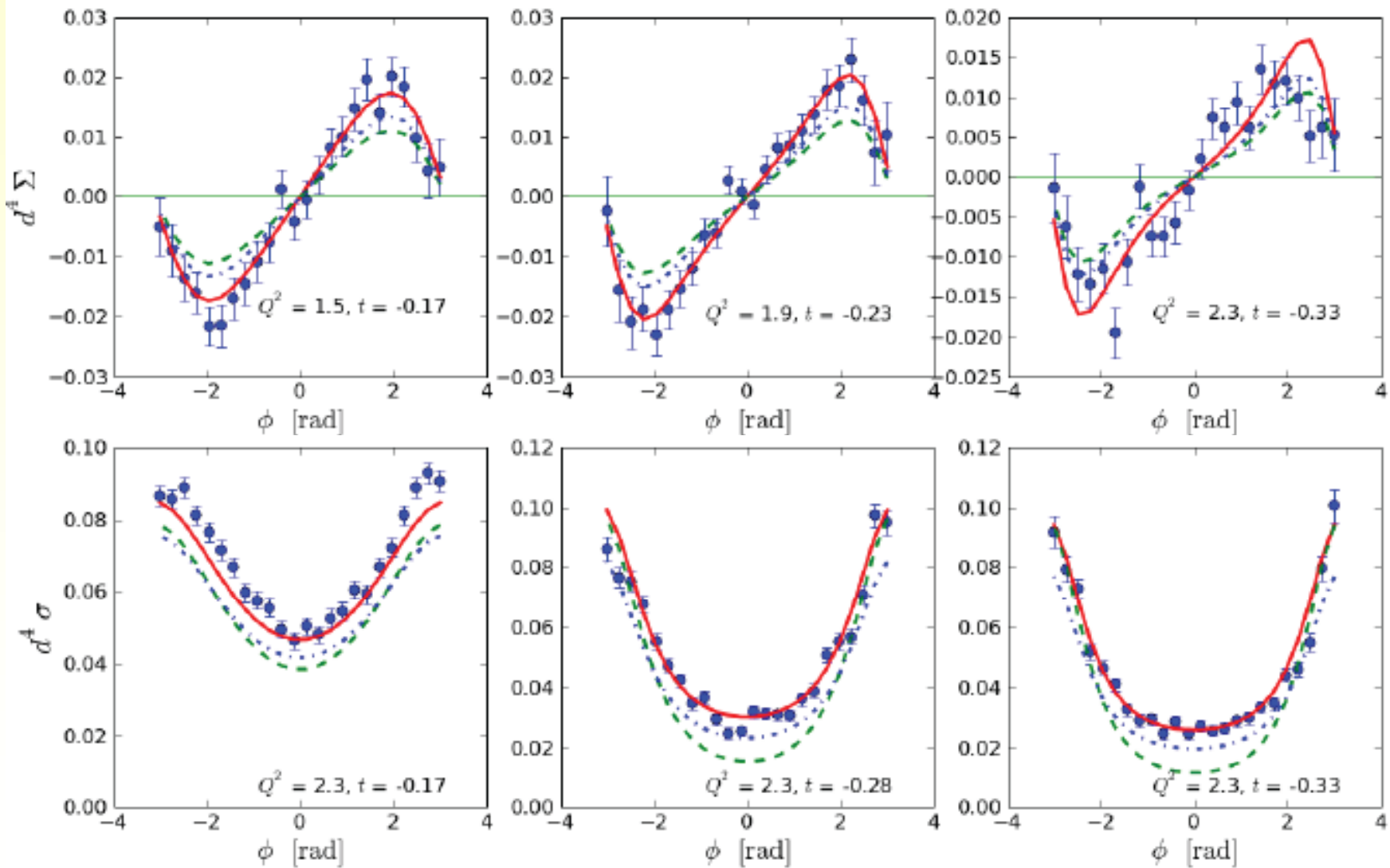
HALL-A data:

neglected at all
ratios of moments
cross sections

- fits to HALL A harmonics are fine for unexpected large \hat{H} or \check{E} contribution
- large \hat{H} KM09 scenario is excluded from longitudinal TSA (HERMES, CLAS)
- large pion pole scenario might look reasonable (cf. [\[Goloskokov and Kroll \(10\)\]](#))

HALL A ϕ -dependence

- ϕ -dependence is described (if we fit to it)



DVCS perspectives

existing data

including longitudinal
and transverse
polarized proton data

new data

HERMES
(recoil detector data)

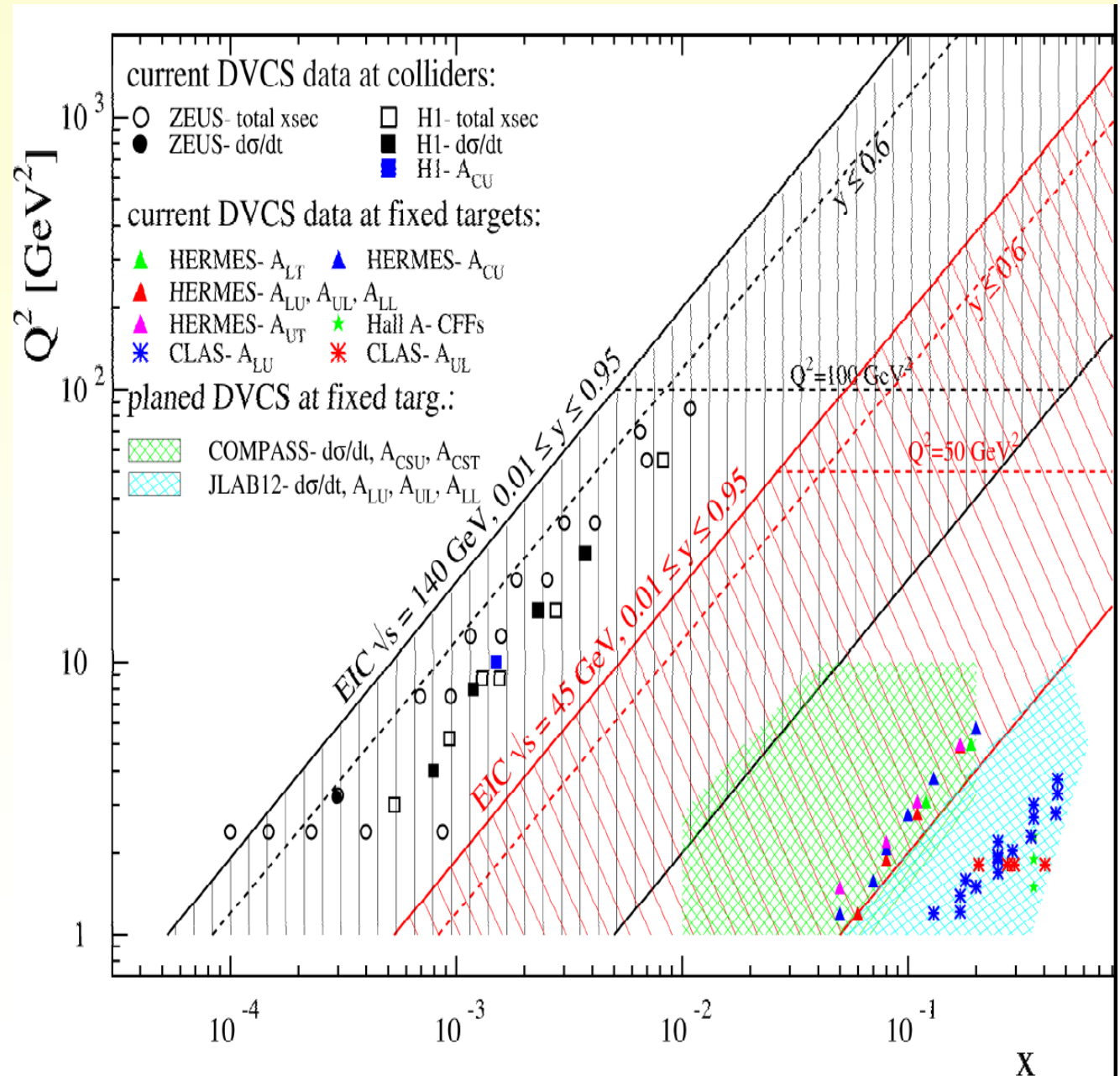
JLAB
(longitudinal TSA,
cross sections)

planned

COMPASS II, JLAB 12

proposed

EIC



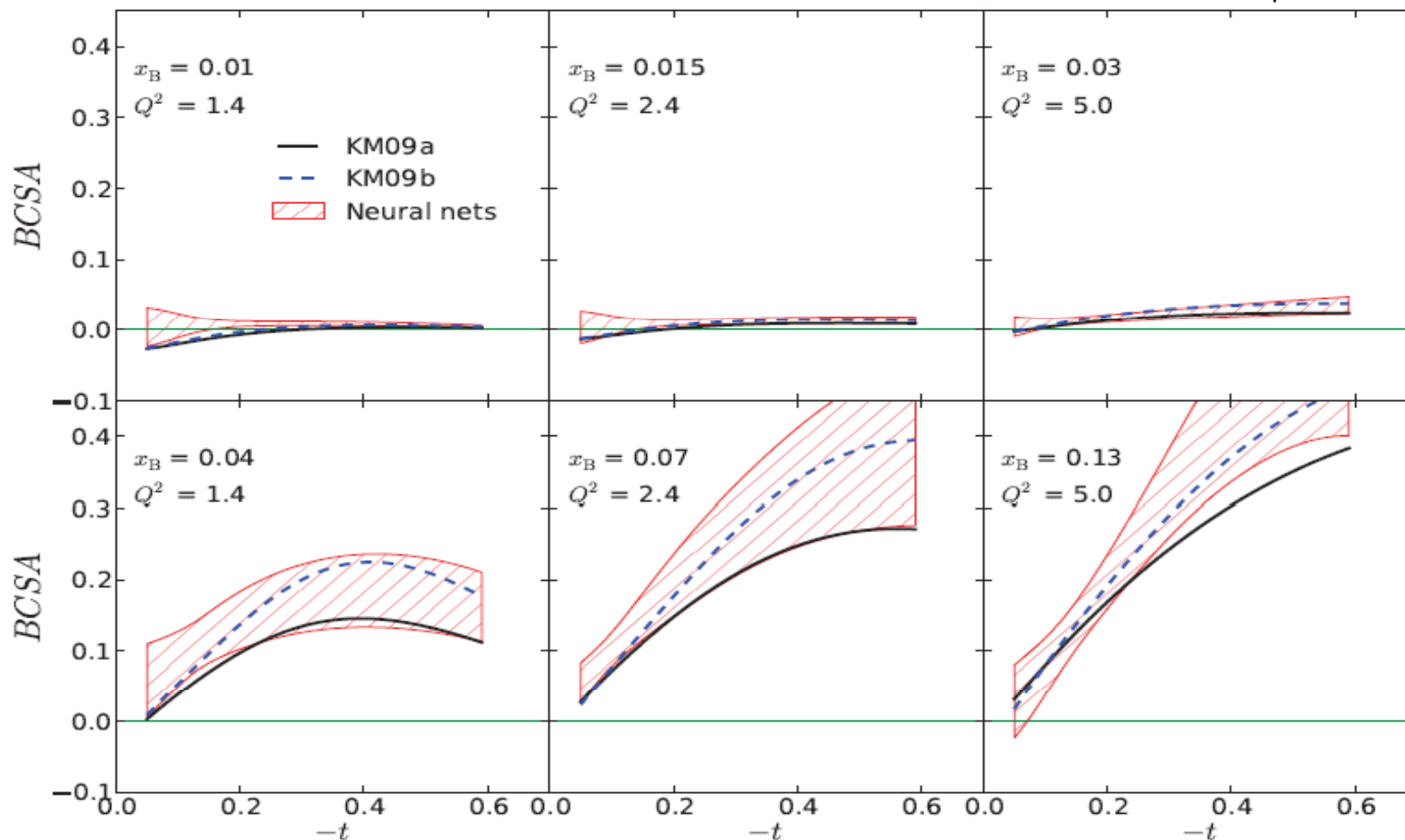
Predictions for Compass II

fixed target, polarized muon beam (~ 200 GeV)

cross sections (t -dependence), transverse polarized target (access to E GPD)

beam charge-spin asymmetry
(dominated by real part)

$$A_{\text{BCSA}} = \frac{d\sigma^{\uparrow\uparrow} - d\sigma^{\downarrow-}}{d\sigma^{\uparrow\uparrow} + d\sigma^{\downarrow-}}$$

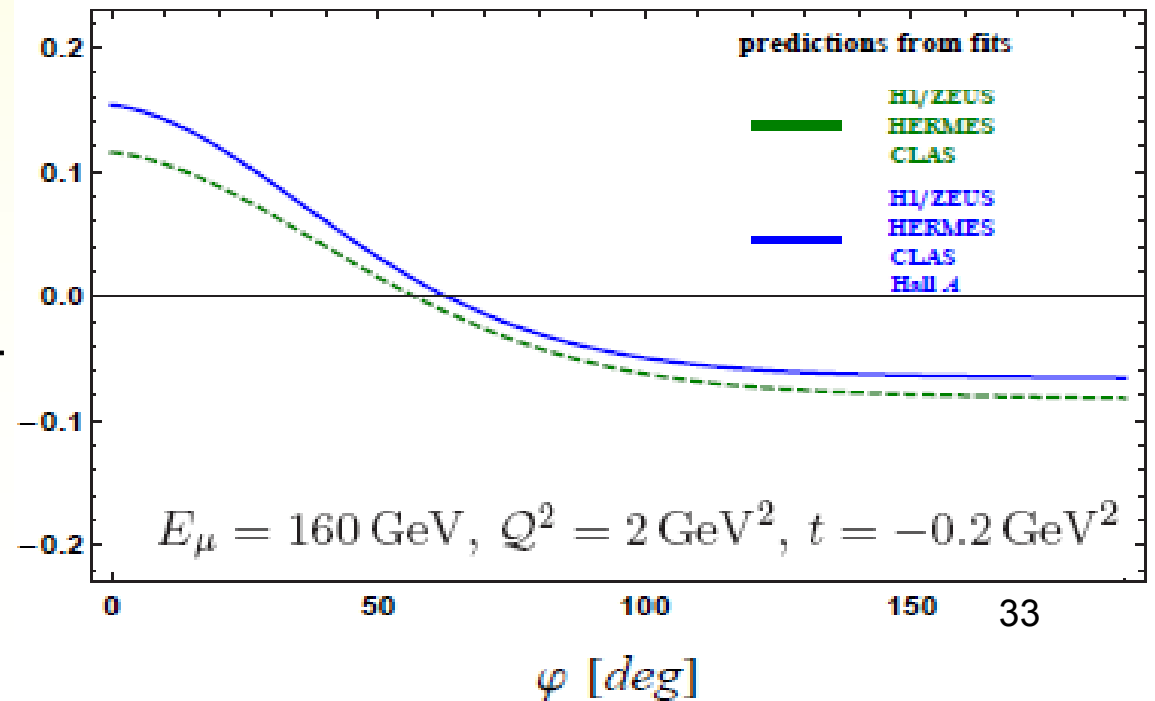
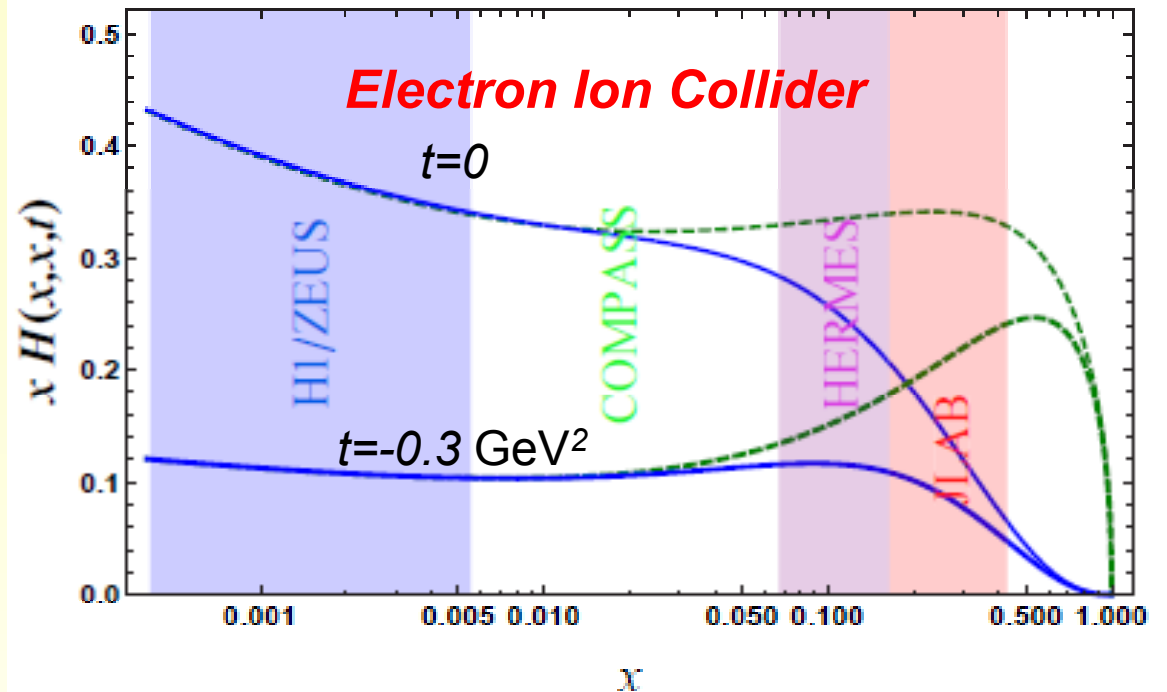


- extracting GPD from present collider and fixed target DVCS data

$$H(x,x,t, Q^2=2 \text{ GeV}^2)$$

- prediction for COMPASS

$$A_{\text{BCSA}} = \frac{d\sigma^{\uparrow\uparrow} - d\sigma^{\downarrow\downarrow}}{d\sigma^{\uparrow\uparrow} + d\sigma^{\downarrow\downarrow}}$$



EIC potential for DVCS

to address angular momentum, 3D picture, (effective) nucleon wave function within the GPD framework new DVCS experiments with

large kinematical coverage, high luminosity, and dedicated detectors are needed to quantify CFFs and GPDs on the cross-over line (and outer region,

- disentangling CFFs at small(er) x
cross sections

beam spin, target spin, and double spin flip experiments

$$BSA \propto y \left\{ F_1(t) H(\xi, \xi, t, Q^2) - \frac{t}{4M^2} F_2 E(\xi, \xi, t, Q^2) \right\}$$

$$TSA_T \propto \frac{\sqrt{-t}}{4M^2} \left\{ F_1(t) E(\xi, \xi, t, Q^2) - F_2(t) H(\xi, \xi, t, Q^2) \right\}$$

$$TSA_L \propto \left\{ F_1(t) \tilde{H}(\xi, \xi, t, Q^2) + \xi (F_1 + F_2)(t) H(\xi, \xi, t, Q^2) \right\}$$

- off neutron another possibility to access GPD E
- separation of twist-2 and twist-3 induced harmonics requires positron beam
- time-like region (a new field to study)
- off nuclei (has its own interest)

Impact of EIC data to extract GPD H

two simulations from S. Fazio for DVCS cross section ~ 650 data points

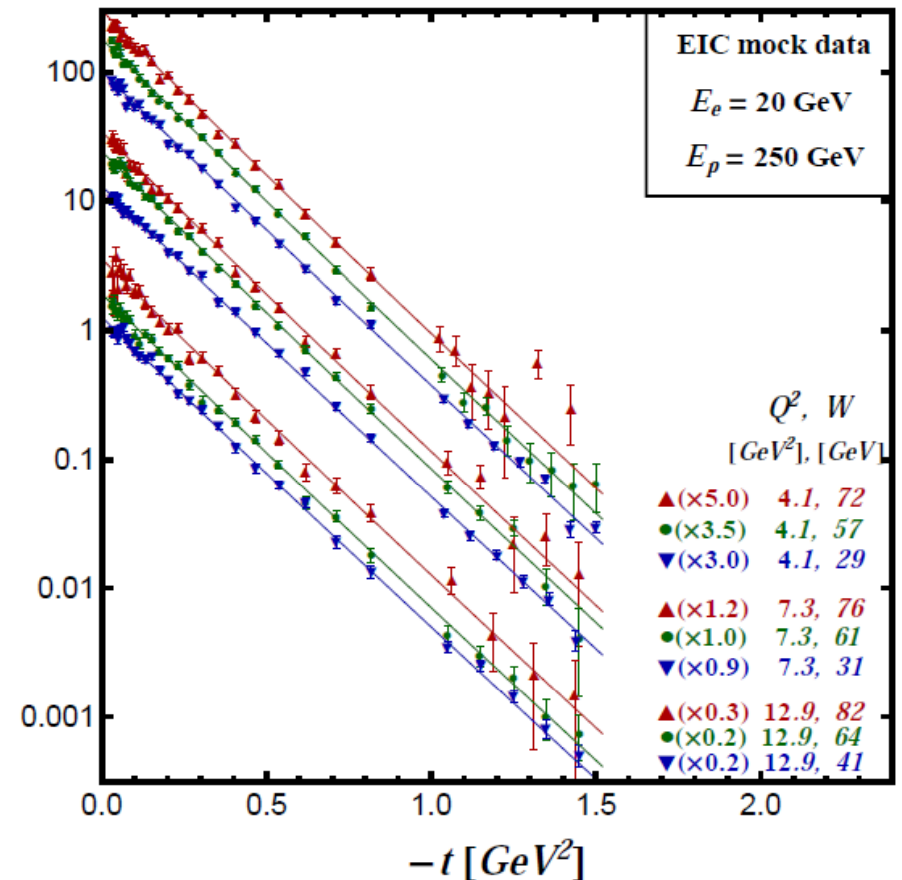
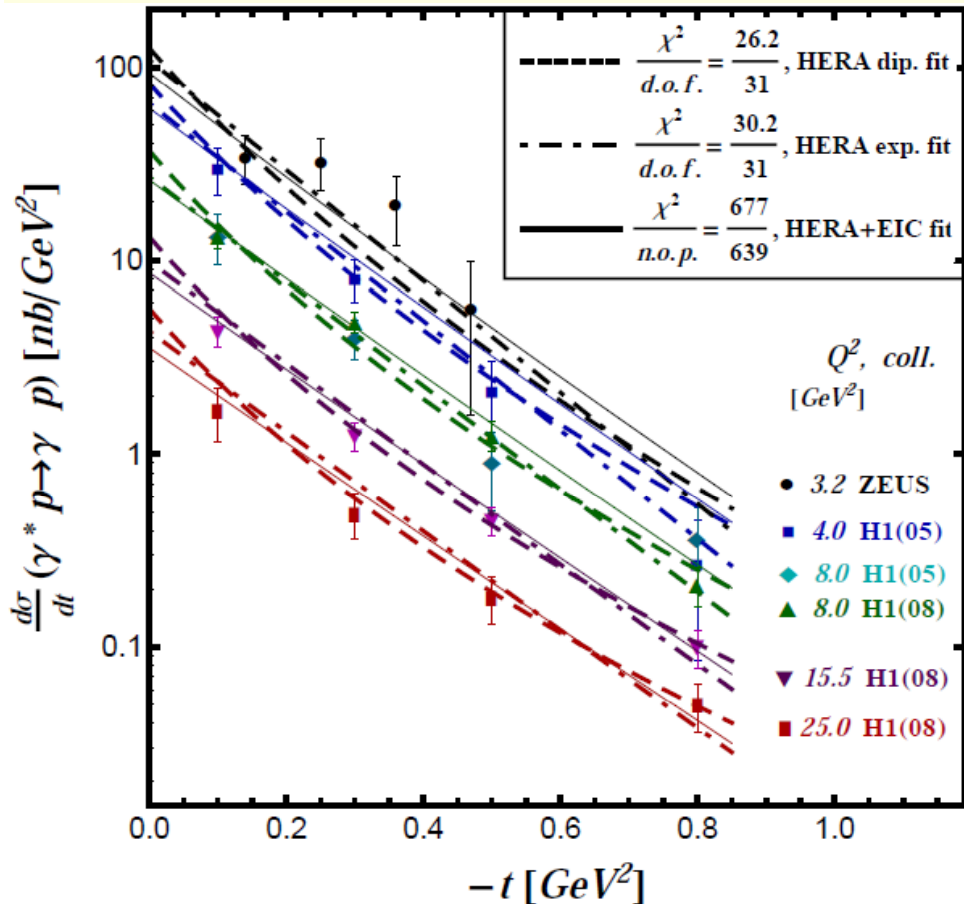
$-t < \sim 0.8 \text{ GeV}^2$ for $\sim 10/\text{fb}$

$1 \text{ GeV}^2 < -t < 2 \text{ GeV}^2$ for $\sim 100/\text{fb}$ (cut: $-t < 1.5 \text{ GeV}^2$, $4 \text{ GeV}^2 < Q^2$ to ensure $-t < Q^2$)

mock data are re-generated with GeParD

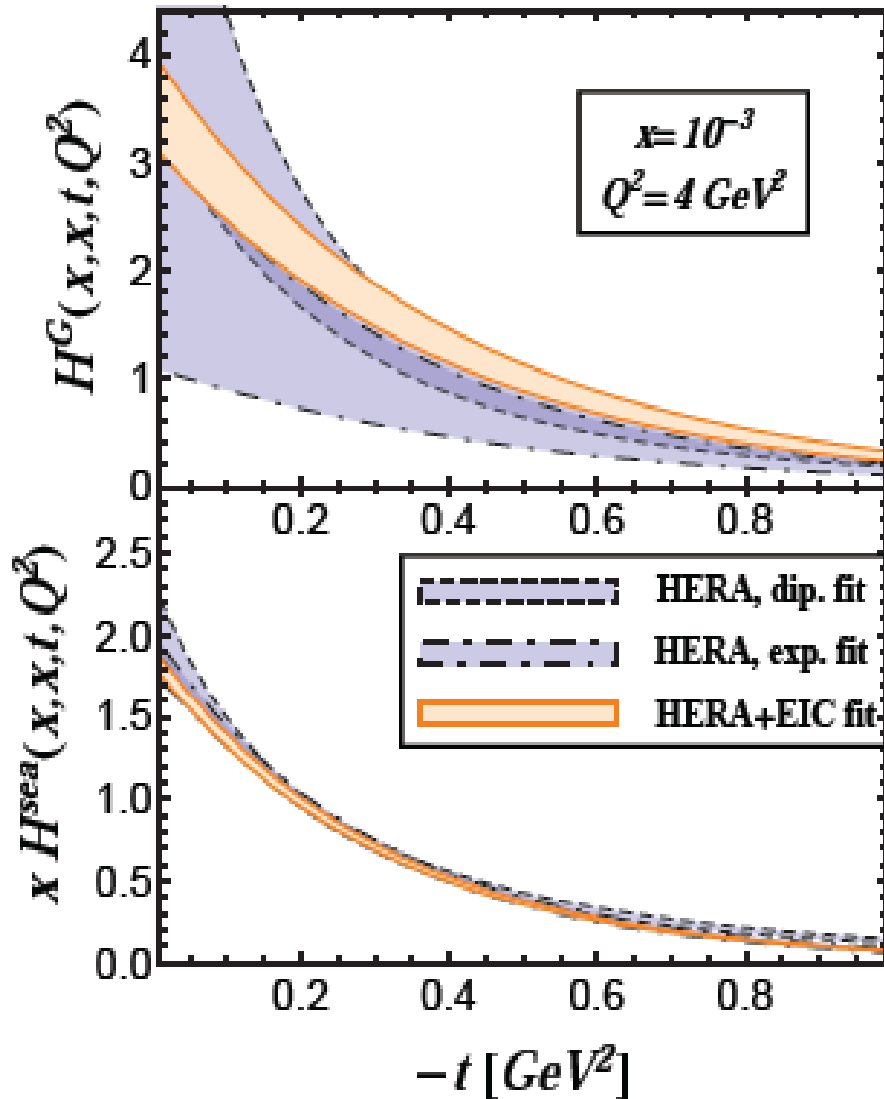
statistical errors rescaled

5% systematical errors added in quadrature, 3% Bethe-Heitler uncertainty



Imaging (probabilistic interpretation)

$$q(x, \vec{b}, \mu^2) = \frac{1}{4\pi} \int_0^\infty d|t| J_0(|\vec{b}| \sqrt{|t|}) H(x, \eta = 0, t, \mu^2)$$

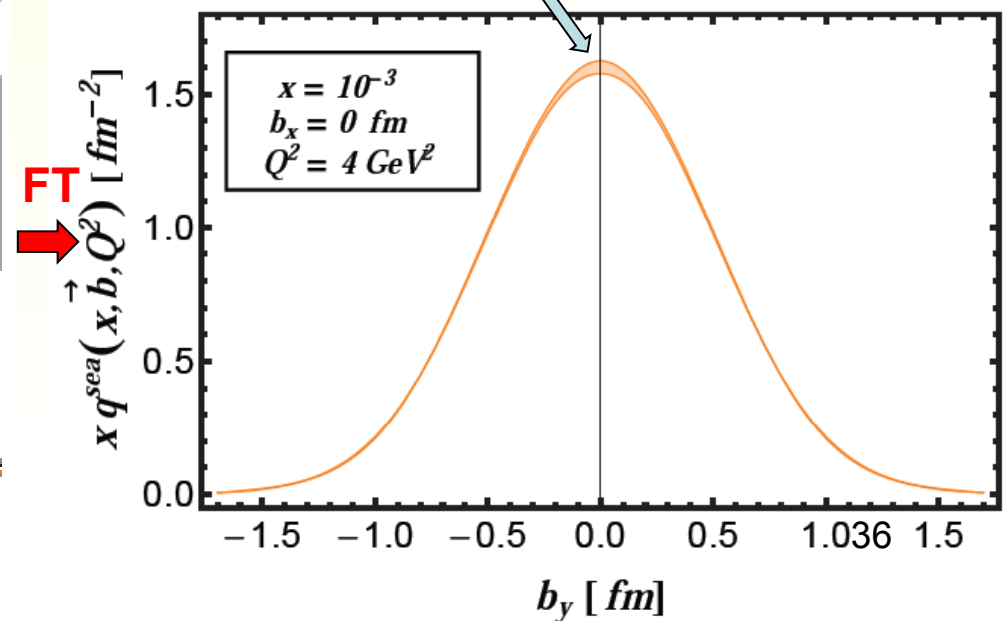


skewness effect vanishes ($s_2, s_4 \rightarrow 0$)

- reduce fit uncertainties
- increase model uncertainties

extrapolation errors for $-t \rightarrow 0$
(large b uncertainties – small effect)

extrapolation errors into $-t > 1.5 \text{ GeV}^2$
(small b uncertainties)



Single transverse target spin asymmetry

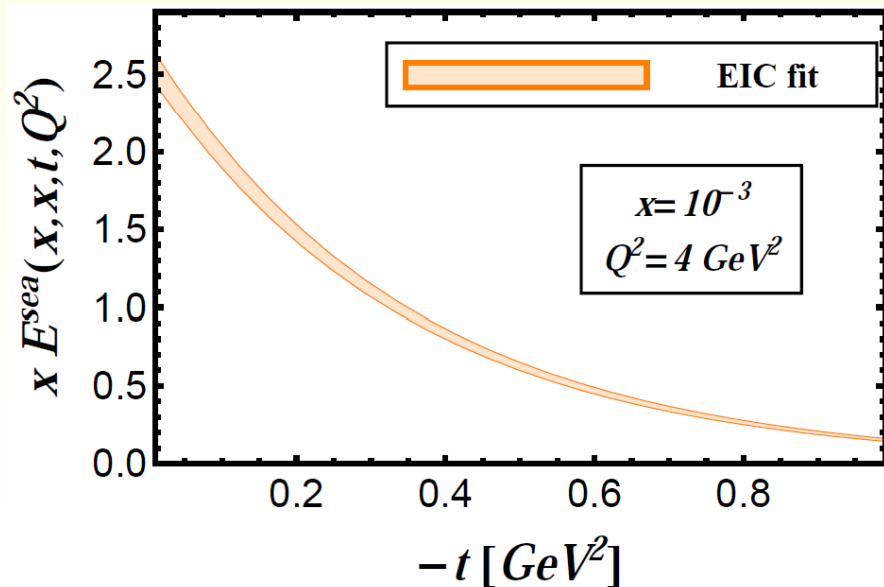
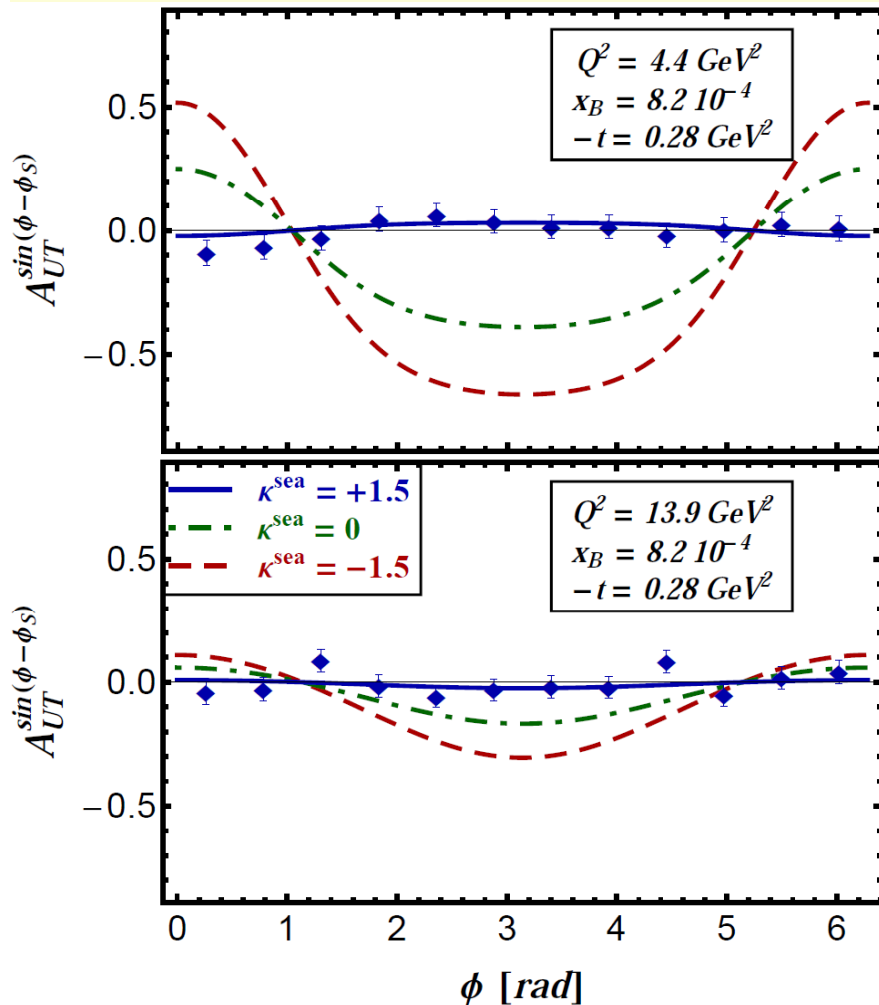
20x250 2x5/fb mock data

(~1200 data points with statistical errors + 5% systematics at cross section level)

flexible GPD model for E^{sea} and E^G

normalization (and t -dependency) of E^{sea} is reasonable constraint

E^G is essentially unconstrained



also imaging of $q^{\uparrow sea}$ is possible [see Franck's talk]

Summary

GPDs are intricate and (thus) a promising tool

- to reveal the transverse distribution of partons
- to address the spin content of the nucleon
- providing a bridge to LCWFs modeling & non-perturbative methods (lattice)

hard exclusive leptonproduction

- DVCS is widely considered as a theoretical clean process
- it is elaborated in NLO and offers a new insight in QCD
- possesses a rich structure, allowing to access various CFFs/GPDs
- new experiments (high luminosity machines and dedicated detectors) are desired to quantify exclusive (and inclusive) QCD phenomena

technology

software tools for global GPD fits have been developed for demonstration

? global QCD fits (inclusive + exclusive)