

Even TMDs and Nucleon Deformation

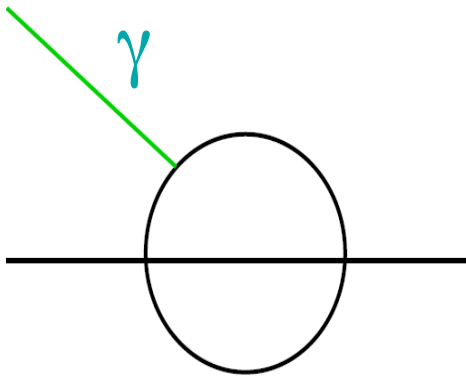
Gerald A. Miller

- **Proton form factor, model calculation-**
proton not round via **spin dependent density**
- **Model independent neutron charge density**
- **Measure shape of proton on lattice (impact parameter dependent GPD) coordinate-space probability, and in experiment (TMD): TMD is momentum-space probability**
- **GAM “Transverse Charge Densities” arXive: 1002.0355, Ann.Rev.Nucl.Part.Sci. 60 (2010) 1-25**

Ratio of Pauli to Dirac Form Factors 1995

Frank, Jennings, Miller theory, data 2000

Impulse approximation

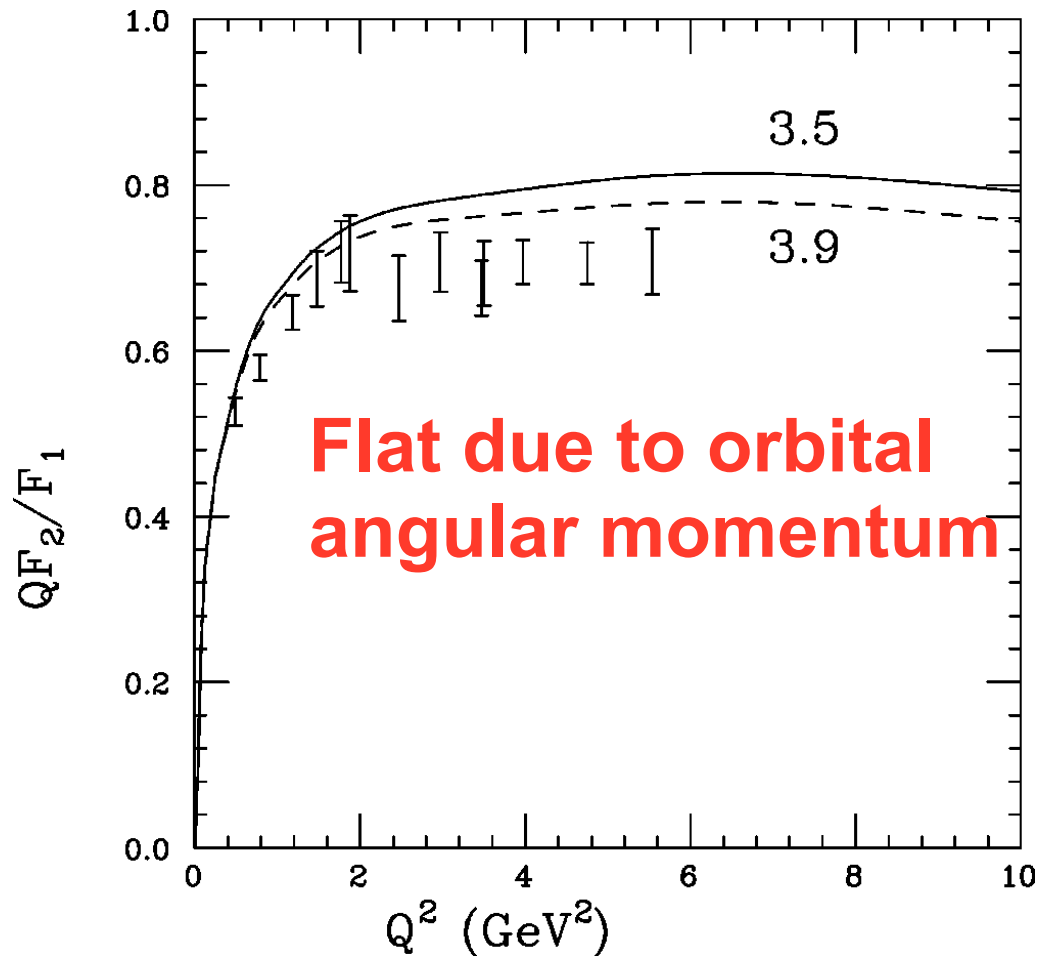


Model proton wave function $\Psi(\mathbf{k}_\perp, \mathbf{K}_\perp, \xi, \eta)$

Poincare invariant

Light front variables for boost: $\mathbf{K} \rightarrow \mathbf{K} + \eta \mathbf{q}_\perp$

Dirac spinors carry orbital angular momentum



Flat due to orbital angular momentum

Model exists

- **lower components of Dirac spinor**
- **orbital angular momentum**
- **shape of proton?? Wigner Eckart**
no quadrupole moment
- **spin dependent densities SDD**
non-relativistic example

I: Non-Rel. $p_{1/2}$ proton outside 0^+ core

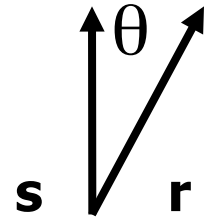
$$\langle \mathbf{r}_p | \psi_{1,1/2s} \rangle = R(r_p) \boldsymbol{\sigma} \cdot \hat{\mathbf{r}}_p |s\rangle \quad \text{Binding pot'l} \\ \text{rotationally invariant}$$

$$\rho(r) = \langle \psi_{1,1/2s} | \delta(\mathbf{r} - \mathbf{r}_p) | \psi_{1,1/2s} \rangle = R^2(r)$$

probability proton at \mathbf{r} & spin direction \mathbf{n} :

$$\rho(\mathbf{r}, \mathbf{n}) = \langle \psi_{1,1/2s} | \delta(\mathbf{r} - \mathbf{r}_p) \frac{(1 + \boldsymbol{\sigma} \cdot \mathbf{n})}{2} | \psi_{1,1/2s} \rangle$$

$$= \frac{R^2(r)}{2} \langle s | \boldsymbol{\sigma} \cdot \hat{\mathbf{r}} (1 + \boldsymbol{\sigma} \cdot \mathbf{n}) \boldsymbol{\sigma} \cdot \hat{\mathbf{r}} | s \rangle$$



$$\mathbf{n} \parallel \hat{\mathbf{s}} : \quad \rho(\mathbf{r}, \mathbf{n} = \hat{\mathbf{s}}) = R^2(r) \cos^2 \theta$$

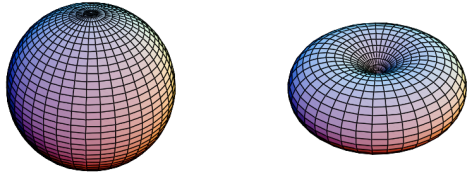
$$\mathbf{n} \parallel -\hat{\mathbf{s}} : \quad \rho(\mathbf{r}, \mathbf{n} = -\hat{\mathbf{s}}) = R^2(r) \sin^2 \theta$$

non-spherical shape depends on spin direction

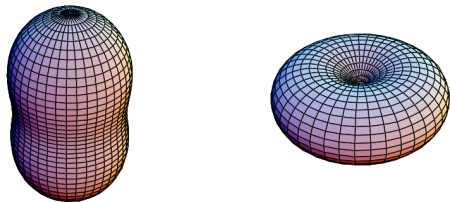
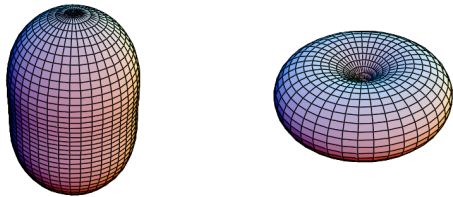
Shapes of the proton

Phys.Rev. C68 (2003) 022201

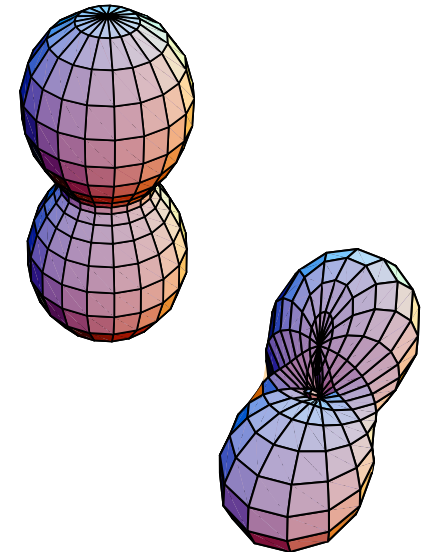
Momentum space



vectors \mathbf{n} , \mathbf{K} , \mathbf{S}



Coordinate space



Pretzelicity

How to measure? - Lattice and/or experiment

Relation between coordinate and momentum space densities? Model independent technique needed.

Model independent transverse charge density

$$J^+(x^-, \mathbf{b}) = \sum_q e_q q_+^\dagger(x^-, b) q_+(x^-, b)$$

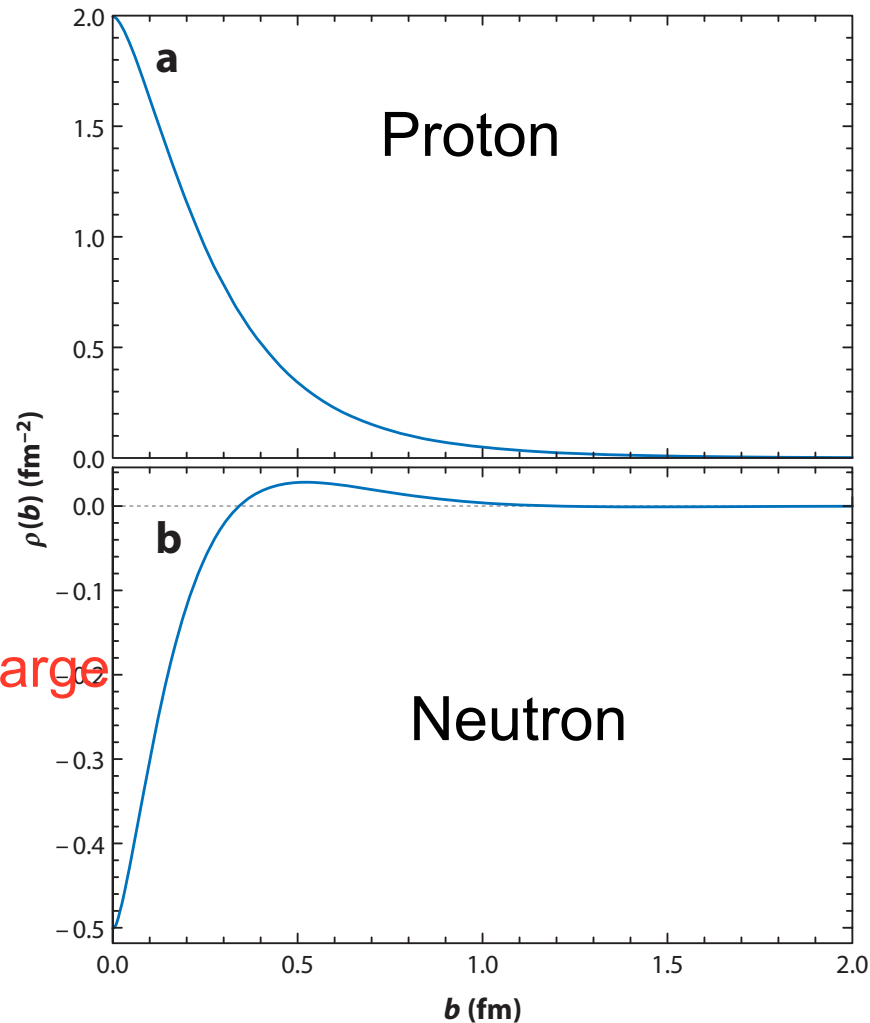
Charge Density operator IMF

$$\rho_\infty(x^-, \mathbf{b}) = \langle p^+, \mathbf{R} = \mathbf{0}, \lambda | \sum_q e_q q_+^\dagger(x^-, b) q_+(x^-, b) | p^+, \mathbf{R} = \mathbf{0}, \lambda \rangle$$

$$F_1 = \langle p^+, \mathbf{p}', \lambda | J^+(0) | p^+, \mathbf{p}, \lambda \rangle$$

$$\rho(b) \equiv \int dx^- \rho_\infty(x^-, \mathbf{b}) = \int \frac{Q dQ}{2\pi} F_1(Q^2) J_0(Qb)$$

Transverse charge densities from parameterizations (Alberico)



Negative central charge density

Negative central density - GAM PRL '07

Generalized Coordinate Space Densities

$$\rho^\Gamma(\mathbf{b}) = \sum_q e_q \int dx^- q_+^\dagger(x^-, \mathbf{b}) \gamma^+ \Gamma q_+(x^-, \mathbf{b})$$

$$\Gamma = \frac{1}{2} (1 + \mathbf{n} \cdot \boldsymbol{\gamma}) \text{ gives spin - dependent density}$$

PRL 98, 222001 (2007)

PHYSICAL REVIEW LETTERS

week ending
1 JUNE 2007

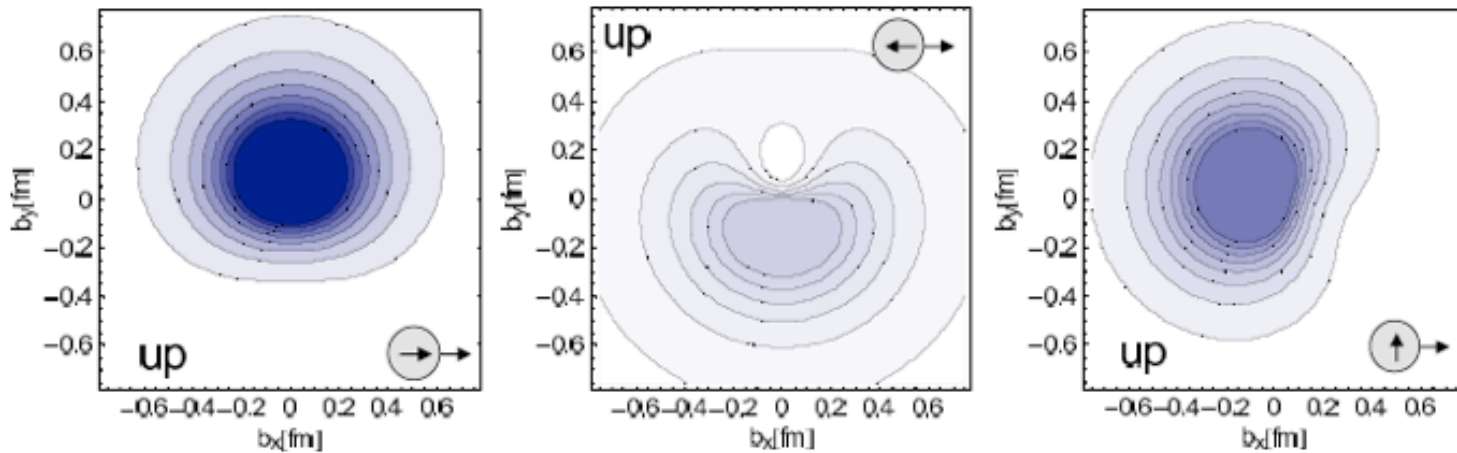
Transverse Spin Structure of the Nucleon from Lattice-QCD Simulations

M. Göckeler,¹ Ph. Hägler,^{2,*} R. Horsley,³ Y. Nakamura,⁴ D. Pleiter,⁴ P. E. L. Rakow,⁵ A. Schäfer,¹ G. Schierholz,^{6,4}
H. Stüben,⁷ and J. M. Zanotti³

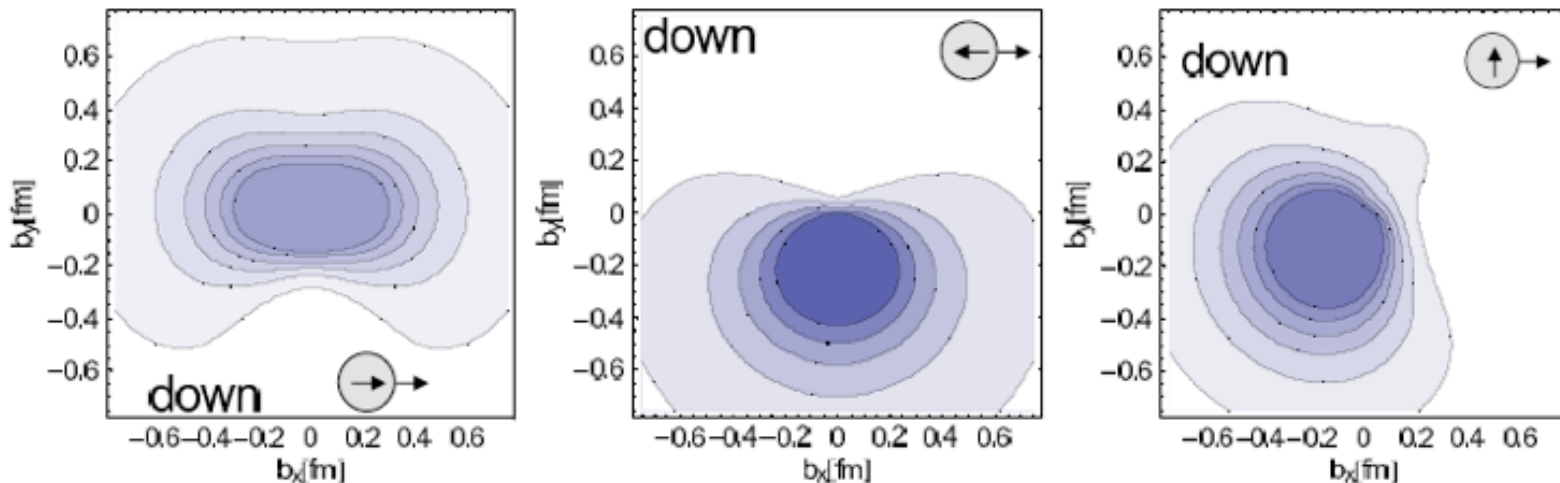
$$\begin{aligned} \rho^n &= \int_{-1}^1 dx x^{n-1} \rho(x, b_\perp, s_\perp, S_\perp) \\ &= \frac{1}{2} \left\{ A_{n0}(b_\perp^2) + s_\perp^i S_\perp^i \left(A_{Tn0}(b_\perp^2) - \frac{\Delta_{b_\perp} \tilde{A}_{Tn0}(b_\perp^2)}{4m^2} \right) \right. \\ &\quad + \frac{b_\perp^j \epsilon^{ji}}{m} [S_\perp^i B'_{n0}(b_\perp^2) + s_\perp^i \bar{B}'_{Tn0}(b_\perp^2)] \\ &\quad \left. + s_\perp^i (2b_\perp^i b_\perp^j - b_\perp^2 \delta^{ij}) S_\perp^j \frac{1}{m^2} \tilde{A}''_{Tn0}(b_\perp^2) \right\}, \end{aligned} \quad (1)$$

spin-dependent density
-depends on direction
of \mathbf{b} : proton is not round

Spin dependent densities-transverse- Lattice QCDSF, Zanotti, Schierholz...

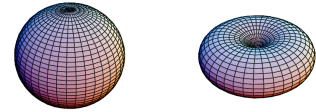


This is not zero! proton is not round



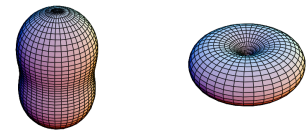
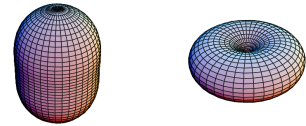
Shapes of the proton

Relate spin dependent density to experiment



Phys.Rev.C76:065209,2007

Field-theoretic spin dependent
momentum density is related to the
transverse momentum distribution h_{1T}^\perp



$$\Phi^{[\Gamma]}(x, \mathbf{K}_T) = \int \frac{d\xi^- d^2\xi_T}{2(2\pi)^3} e^{i\mathbf{K}\cdot\xi} \langle P, S | \bar{\psi}(0) \Gamma \mathcal{L}(0, \xi; n_-) \psi(\xi) | P, S \rangle \Big|_{\xi^+=0}$$

Mulders Tangerman'96

$$\Phi^{[i\sigma^{i+}\gamma^5]}(x, \mathbf{K}_T) = S_T^i h_1(x, K_T^2) + \frac{(K_T^i K_T^j - \frac{1}{2} K_T^2 \delta_{ij}) S_T^j}{M^2} h_{1T}^\perp(x, K_T^2)$$

$$\sigma^{i+}\gamma^5 \sim \gamma^0 \gamma^+ \sigma^i,$$

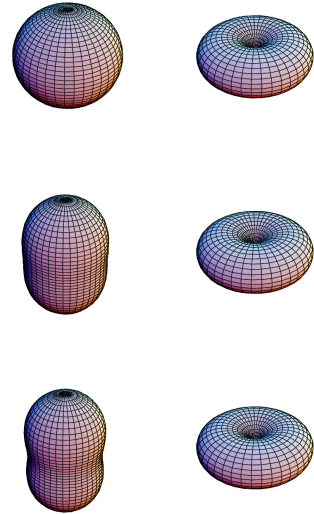
then relate equal time to $\xi^+ = 0$ by integration over x

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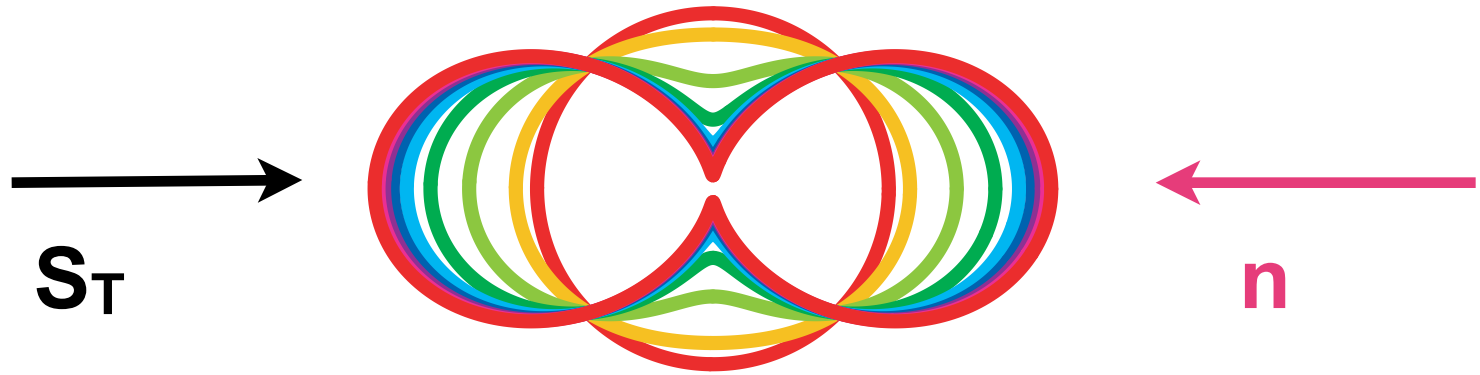
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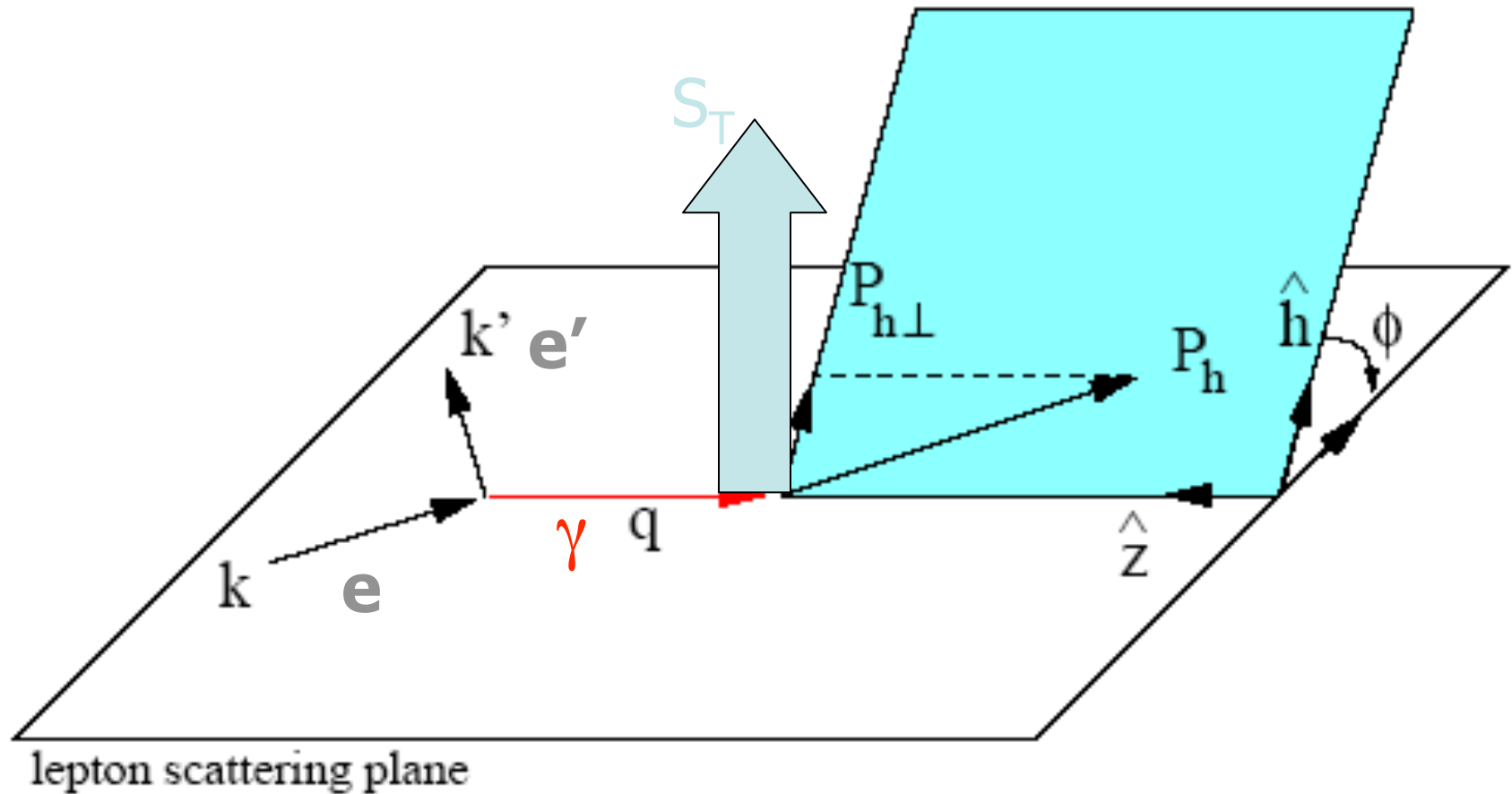
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Transverse Shapes of the Proton



Measure h_{1T}^\perp : $\mathbf{e}, \uparrow \mathbf{p} \rightarrow \mathbf{e}', \pi \mathbf{X}$

H. Avakian LOI at Jlab



Cross section has term proportional to $\cos 3\phi$
 Boer Mulders '98 there are other ways to see h_{1T}^\perp

Boer Mulders, PRD57,5780

TABLE II. Leading order single spin asymmetries for the case of leptonproduction into unpolarized final states.

ABC	W	$\langle W \rangle_{ABC} \cdot [4\pi \alpha^2 s/Q^4]^{-1}$	T
OLO	$(Q_T^2/4MM_h) \sin(2\phi_h^\ell)$	$-\lambda(1-y) \sum_{a,\bar{a}} e_a^2 x_B h_{1L}^{\perp(1)a}(x_B) H_1^{\perp(1)a}(z_h)$	eo
OTO	$(Q_T/M_h) \sin(\phi_h^\ell + \phi_S^\ell)$	$ \mathbf{S}_T (1-y) \sum_{a,\bar{a}} e_a^2 x_B h_1^a(x_B) H_1^{\perp(1)a}(z_h)$	eo
OTO	$(Q_T^3/6M^2M_h) \sin(3\phi_h^\ell - \phi_S^\ell)$	$ \mathbf{S}_T (1-y) \sum_{a,\bar{a}} e_a^2 x_B h_{1T}^{\perp(2)a}(x_B) H_1^{\perp(1)a}(z_h)$	eo
OTO	$(Q_T/M) \sin(\phi_h^\ell - \phi_S^\ell)$	$ \mathbf{S}_T (1-y + \frac{1}{2}y^2) \sum_{a,\bar{a}} e_a^2 x_B f_{1T}^{\perp(1)a}(x_B) D_1^a(z_h)$	oe

There have been searches for this term

Summary

- Form factors, **GPDs**, **TMDs**, understood from unified light-front formulation, **GPD-coordinate space density**, **TMD momentum space density**
- Neutron central transverse density is negative-
- Proton is not round- lattice QCD spin-dependent-density is **not** zero
- **Experiment can whether or not proton is round by measuring**

$$h_{1T}^{\perp}$$



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The Proton