

Weak reactions in few nucleon systems: pp and muon captures

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Program INT-12-3
Light nuclei from first principles
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Outline

- 1 Introduction
- 2 Theoretical ingredients
- 3 Results: muon capture
- 4 Results: pp capture
- 5 Conclusion & Outlook

Collaborators

- L.E. Marcucci - *INFN & Pisa University, Pisa (Italy)*
- L. Girlanda - *INFN & Salento University, Lecce (Italy)*
- R. Schiavilla & M. Piarulli- *Jefferson Lab. & ODU, Norfolk (VA, USA)*
- A. Kievsky - *INFN & Pisa University, Pisa (Italy)*

Reactions of interest ($A \leq 4$)

Weak captures of astrophysical interest:

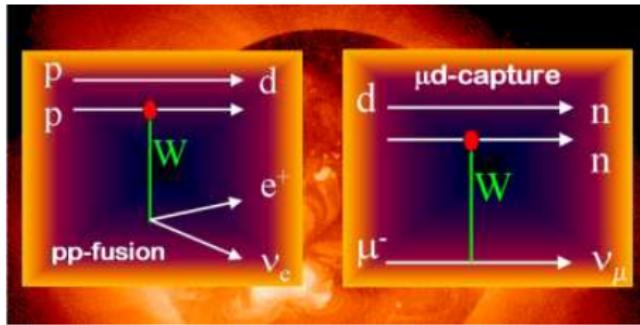
- $p + p \rightarrow {}^2\text{H} + e^+ + \nu_e$
- $p + {}^3\text{He} \rightarrow {}^4\text{He} + e^+ + \nu_e$

Muon captures:

- $\mu^- + d \rightarrow n + n + \nu_\mu$
- $\mu^- + {}^3\text{He} \rightarrow {}^3\text{H} + \nu_\mu$ (70%)
- $\mu^- + {}^3\text{He} \rightarrow n + d + \nu_\mu$ (20%)
- $\mu^- + {}^3\text{He} \rightarrow n + n + p + \nu_\mu$ (10%)

Motivations

- Calculations of reaction rates at astrophysical energies
- Muon captures: test of the theoretical models of weak current (MuSun experiment)
- Extraction of $G_{PS}(q^2) \rightarrow$ validate χPT

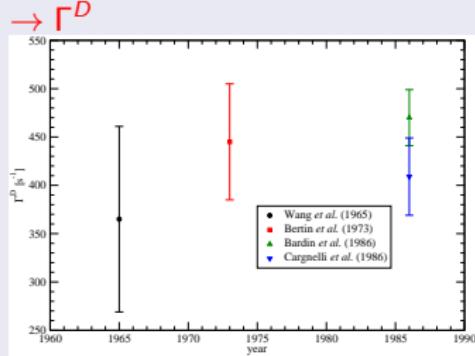


$$j^\mu = \overline{u_p} \left[F_1(q^2) \gamma^\mu + F_2(q^2) \frac{i \sigma^{\mu\nu} q_\nu}{2M_N} - G_A(q^2) \gamma^\mu \gamma^5 - G_{PS}(q^2) \frac{q^\mu \gamma^5}{2M_N} \right] u_n$$

Experimental situation



- Two hyperfine states: $f = 1/2$ and $3/2$
- Dominant capture from $f = 1/2$
 $\rightarrow \Gamma^D$



- New measurement in progress:
MuSun
- \rightarrow Kammel's talk



- Hyperfine states:
 $(f, f_z) = (1, \{\pm 1, 0\})$ and $(0, 0)$

$$\frac{d\Gamma}{d(\cos \theta)} = \frac{1}{2} \Gamma_0 [1 + A_v P_v \cos \theta + A_t P_t \left(\frac{3 \cos^2 \theta - 1}{2} \right) + A_\Delta P_\Delta]$$

$$P_v = P_{1,1} - P_{1,-1}$$

$$P_t = P_{1,1} + P_{1,-1} - 2P_{1,0}$$

$$P_\Delta = 1 - 4P_{0,0}$$

- Γ_0 = total capture rate, $A_{v,t,\Delta}$ = angular correlation parameters

- Ackerbauer et al., (1998):

$$\Gamma_0 = 1496(4) \text{ s}^{-1}$$

- Souder et al., (1998): $A_v = 0.63 \pm 0.09$ (stat.) $^{+0.11}_{-0.14}$ (syst.)

Theoretical studies: “main ingredients” (I)

NN potentials

- "Old models": Argonne V18, CD-Bonn, Nijmegen ($\chi^2 \approx 1$)
- Fit of 3N data using non-locality in P-waves (INOY [Doleschall, 2008])
- Effective field theory (EFT)
 - J-N3LO – [Epelbaum and Coll, 1998-2006]
 - N3LO – [Entem & Machleidt, 2003]

3N potentials

- "Old models": Tucson-Melbourne [Coon *et al*, 1979, Friar *et al*, 1999]; Brazil [Robilotta & Coelho, 1986]; Urbana [Pudliner *et al*, 1995]
- Effective field theory
 - at N2LO [Epelbaum *et al*, 2002], [Navratil, 2007]
 - Illinois [Pieper *et al*, 2001]
 - Under progress: N3LO, N4LO

Accurate nuclear wave functions

- Methods for $A \geq 3$: Faddeev-Yakubovsky Equations, GFMC, Variational methods (Gaussians, NCSM)
 - HH method [J. Phys. G 35, 063101 (2008)]

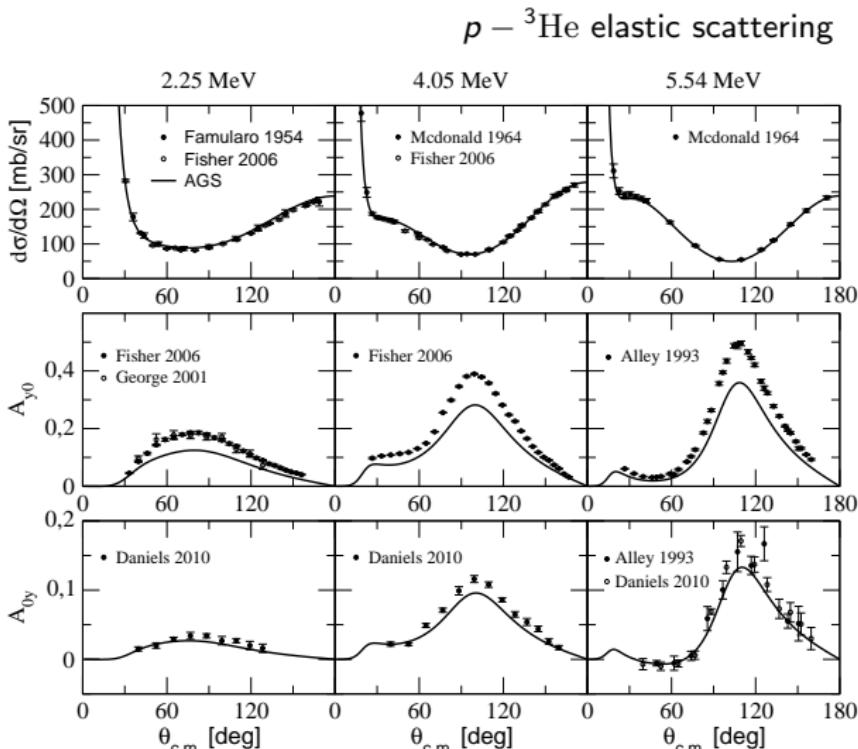
Some results for $A = 2\text{--}4$

$A = 2$	AV18	N3LO	Exp.
B_d (MeV)	2.22457	2.22456	2.224574(9)
a_{nn} (fm)	-18.487	-18.900	-18.9(4)
$^1a_{np}$ (fm)	-23.732	-23.732	-23.740(20)
$^3a_{np}$ (fm)	5.412	5.417	5.419(7)
$A = 3$	AV18/UIX	N3LO/N2LO	Exp.
$B_3\text{H}$ (MeV)	8.479	8.474	8.482
$B_3\text{He}$ (MeV)	7.750	7.733	7.718
$^2a_{nd}$ (fm)	0.590	0.675	0.645(10)
$^4a_{nd}$ (fm)	6.343	6.342	6.35(2)
$A = 4$	AV18/UIX	N3LO/N2LO	Exp.
$B_4\text{He}$ (MeV)	28.45	28.36	28.30
$^0a_{n^3\text{He}}$ (fm)	7.81	7.61	7.57(3)
$^1a_{n^3\text{He}}$ (fm)	3.39	3.37	3.36(1)

Accuracy of the calculation tested in several benchmarks

Bound states: [Kamada *et al.*, 2001] – Scattering states [MV *et al.*, 2011]

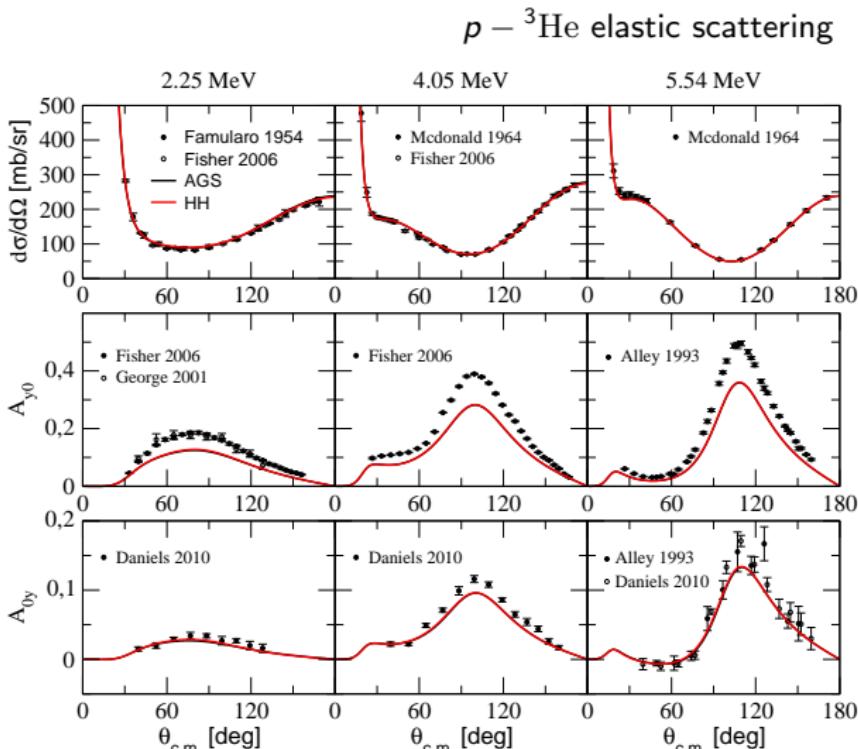
Benchmark test of 4N scattering calculations [PRC 84, 054010 (2011)]



N3LO potential

AGS = Deltuva & Fonseca
FY = Lazauskas & Carbonell

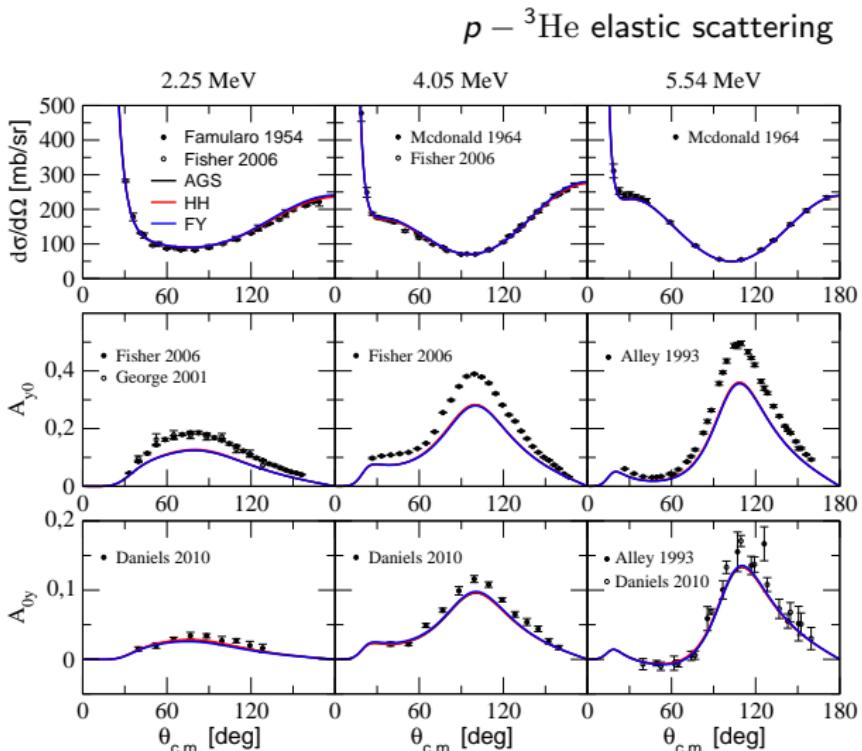
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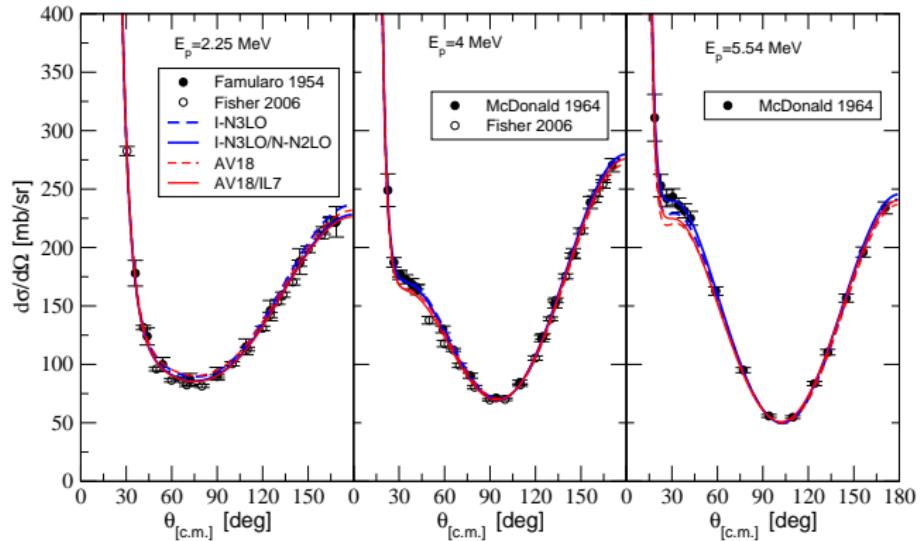
Benchmark test of 4N scattering calculations [PRC 84, 054010 (2011)]



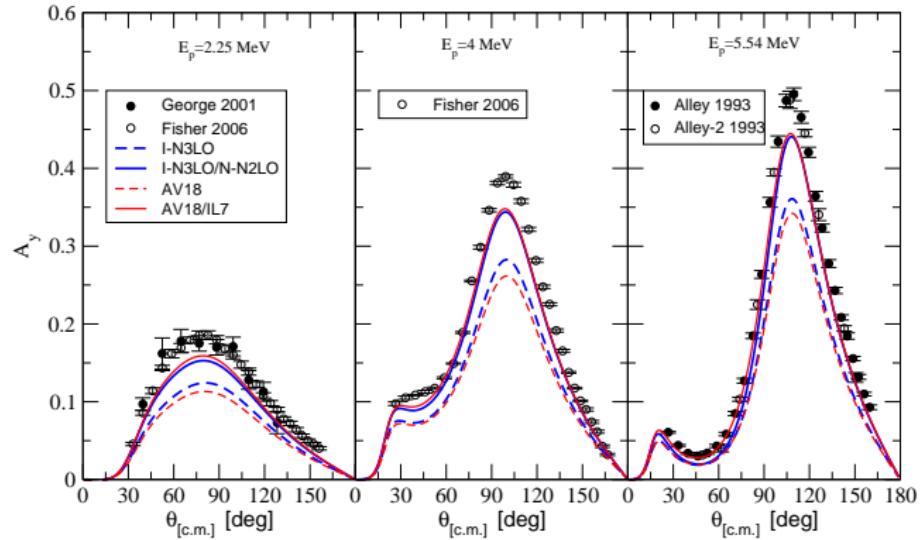
N3LO potential

AGS= Deltuva & Fonseca
FY= Lazauskas & Carbonell

$p - {}^3\text{He}$ differential cross section



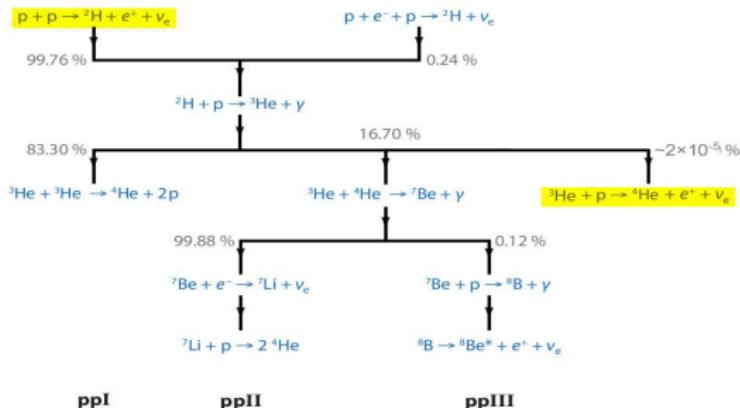
$p - {}^3\text{He}$ A_y



Theoretical studies: “main ingredients” (II)

- Nuclear weak transition operators: $[\rho^{(A,V)}, \mathbf{j}^{(A,V)}]$
 - Standard Nuclear Physics Approach - SNPA [Schiavilla *et al.*, PRC **58**, 1263 (1998); Marcucci *et al.*, PRC **63**, 015801 (2000)]
 - Chiral Effective Field Theory “hybrid” Approach - χ EFT* [Park, Min, & Rho, Phys. Rep. **233**, 341 (1993); Park *et al.*, PRC **67**, 055206 (2003)]

SNPA and χ EFT* used for $p + p \rightarrow d + e^+ + \nu_e$ and $p + {}^3\text{He} \rightarrow {}^4\text{He} + e^+ + \nu_e$



- “Less hybrid” Chiral Effective Field Theory Approach - χ EFT [Marcucci *et al.*, PRL **108**, 052502 (2012)]

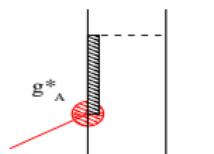
Nuclear transition operators: SNPA

- One-body operators: NRR of $j_i^\mu \rightarrow O(1/m^2)$
- Two-body operators: $\pi-$, $\rho-$, $\omega-$, ... exchanges + Δ d.o.f. excitations
 - $\rho^{(V)}$ and $\mathbf{j}^{(V)}$: CVC \rightarrow EM operators (constructed to verify current conservation)

	$\mu(^3\text{H})$	$\mu(^3\text{He})$	
1b	2.5745	-1.7634	AV18/UIX, \Rightarrow Full=1b+2b+3b
Full	2.9525	-2.1299	[Marcucci <i>et al.</i> , PRC 72, 014001 (2005)]
Exp.	2.9790	-2.1276	

- Two-body $\rho^{(A)}$: PCAC + low-energy theorem \rightarrow π -exchange and short-range terms
- Two-body $\mathbf{j}^{(A)}$: π - and ρ -exchange, $\pi\rho$ mechanism, and $\mathbf{j}^{(A)}(\Delta)$

Largest contribution to $\mathbf{j}^{(A)}(\Delta)$ from



g_A^* fit to observable: GT_{exp} of tritium β -decay

Nuclear transition operators: χ EFT at N³LO

- One-body operators \equiv SNPA
- Two-body operators: from [Park, Min, & Rho, Phys. Rep. 233, 341 (1993)], [Song *et al.*, PRC 79, 064002 (2009)]
 - Two-body $\rho^{(A)}$: soft π -exchange dominant
 - Two-body $\rho^{(V)} = 0$ at N³LO
 - Two-body $j^{(V)}$: CVC \rightarrow EM current
here $1\pi + 2\pi + \text{CT} \rightarrow$ two LECs (g_{4S} & g_{4V}) \Rightarrow from $\mu(^3\text{H} - ^3\text{He})$
 - Two-body $j^{(A)}$: $1\pi + \text{CT} \rightarrow$ one LEC (d_R) \Rightarrow from GT_{exp} of ^3H β -decay

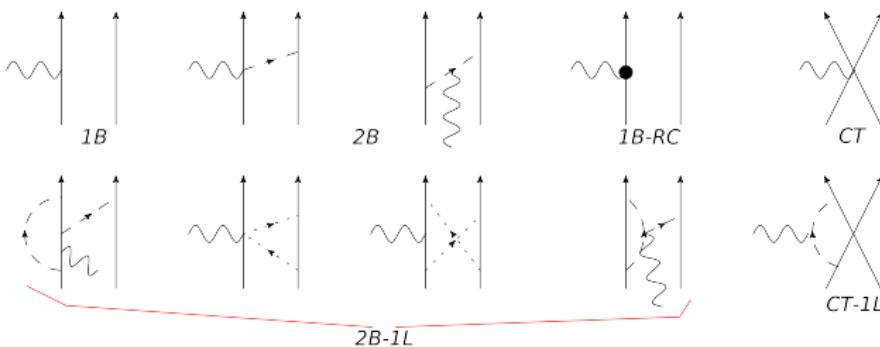
“hybrid” χ EFT* \Rightarrow AV18/UIX [N3LO/N2LO] \Rightarrow current and potentials
“uncorrelated”

	$\Lambda = 500$ MeV	$\Lambda = 600$ MeV	$\Lambda = 800$ MeV
d_R	0.97(7)	1.75(8) [1.00(9)]	3.89(10)
g_{4S}	0.69(1)	0.55(1) [0.11(1)]	0.25(2)
g_{4V}	2.065(6)	0.793(6) [3.124(6)]	-1.07(1)

χ charge & current: chiral counting

Contributions from each type of current at $\mathbf{q} = \mathbf{p}_e + \mathbf{p}_\nu = 0$.

J^μ	LO	NLO	N^2LO	N^3LO	N^4LO
\mathbf{j}^A	1B	—	1B-RC	2B	1B-RC, 2B-1L and 3B
ρ^A	—	1B	2B	1B-RC	1B-RC, 2B-1L
\mathbf{j}^V	—	1B	2B	1B-RC	1B-RC, 2B-1L
ρ^V	1B	—	—	2B	1B-RC, 2B-1L and 3B



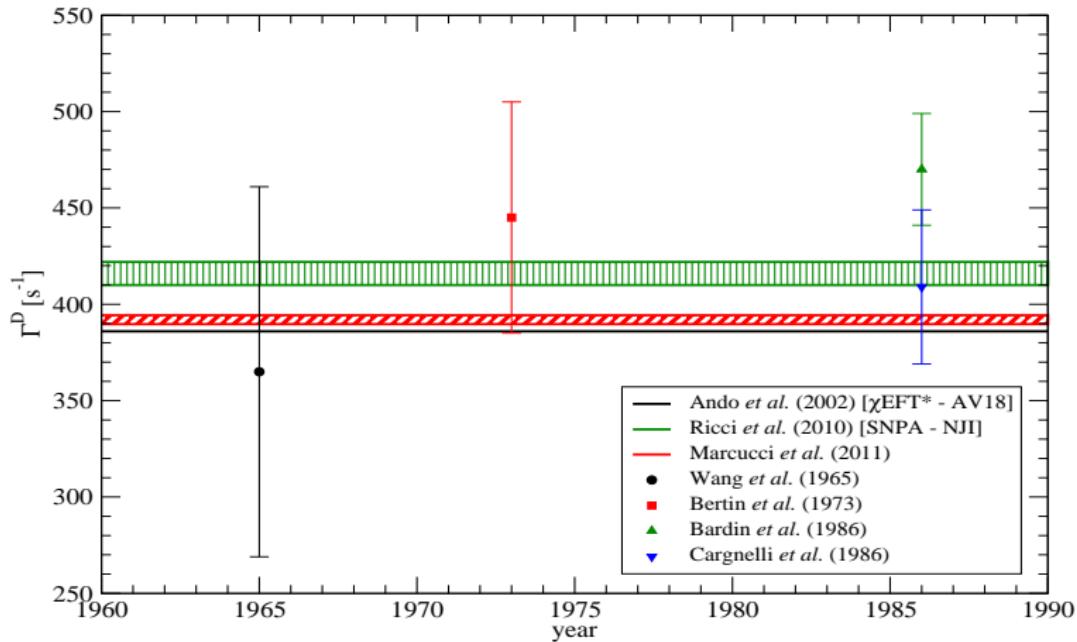
Results: $\Gamma^D(\mu^- + d)$ (SNPA and χ EFT*)

SNPA(AV18)	1S_0	3P_0	3P_1	3P_2	1D_2	3F_2	Total
$g_A=1.2654(42)$	246.6(7)	20.1	46.7	71.6	4.5	0.9	390.4(7)
$g_A=1.2695(29)$	246.8(5)	20.1	46.8	71.8	4.5	0.9	390.9(7)
χ EFT*(AV18)	1S_0	3P_0	3P_1	3P_2	1D_2	3F_2	Total
$\Lambda = 500$ MeV	250.0(8)	19.9	46.2	71.2	4.5	0.9	392.7(8)
$\Lambda = 600$ MeV	250.0(8)	19.8	46.3	71.1	4.5	0.9	392.6(8)
$\Lambda = 800$ MeV	249.7(7)	19.8	46.4	71.1	4.5	0.9	392.4(7)
χ EFT*(N3LO)	1S_0	3P_0	3P_1	3P_2	1D_2	3F_2	Total
$\Lambda = 600$ MeV	250.5(7)	19.9	46.4	71.5	4.4	0.9	393.6(7)

$$\Rightarrow \boxed{\Gamma^D = 390 \div 394 \text{ s}^{-1}}$$

Marcucci *et al.*, PRC **83**, 014002 (2011)

Comparison with data and previous calculations



Results: $\Gamma_0(\mu^- + {}^3\text{He})$ (SNPA and χ EFT*)

SNPA(AV18/UIX)	Γ_0
$g_A=1.2654(42)$	1486(8)
$g_A=1.2695(29)$	1486(5)
χ EFT*(AV18/UIX)	Γ_0
$\Lambda = 500 \text{ MeV}$	1487(8)
$\Lambda = 600 \text{ MeV}$	1488(9)
$\Lambda = 800 \text{ MeV}$	1488(8)
χ EFT*(N3LO/N2LO; $\Lambda=600 \text{ MeV}$)	1480(9)

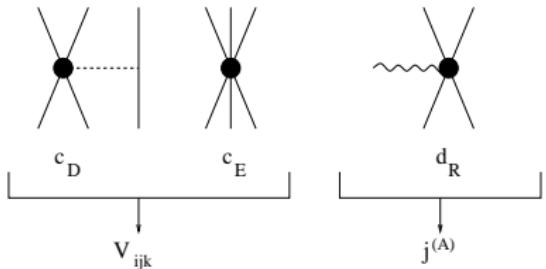
\Rightarrow

$$\Gamma_0 = 1484(13) \text{ s}^{-1}$$

To be compared with:

- $\Gamma_0(\text{exp})=1496(4) \text{ s}^{-1}$
- Marcucci *et al.*, PRC **66**, 054003 (2002) [SNPA–AV18/UIX] $\rightarrow 1484(8) \text{ s}^{-1}$
- Gazit, PLB **666**, 472 (2008) [χ EFT*–AV18/UIX] $\rightarrow 1499(16) \text{ s}^{-1}$

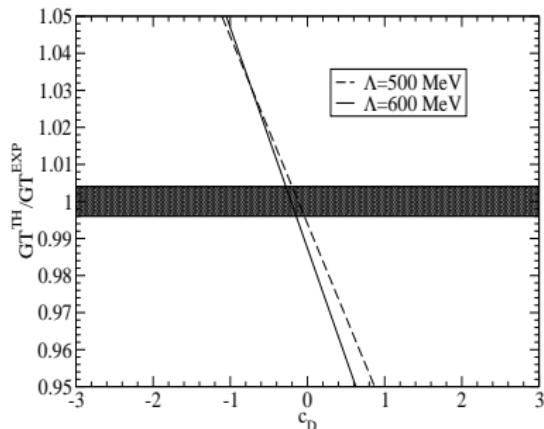
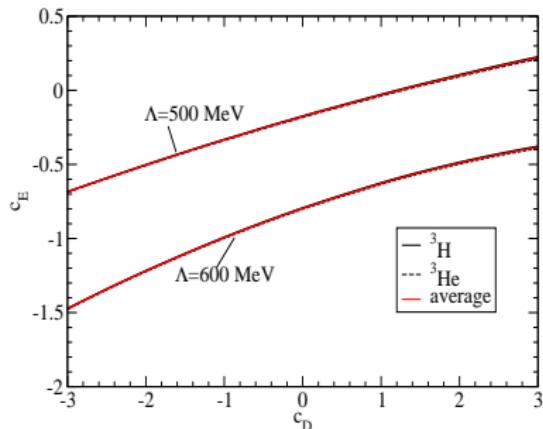
"Less hybrid" χ EFT calculation



$$d_R = \frac{M_N}{\Lambda_\chi g_A} c_D + \frac{1}{3} M_N(c_3 + 2c_4) + \frac{1}{6}$$

Gardestig and Phillips, PRL **96**, 232301 (2006)
 Gazit *et al.*, PRL **103**, 102502 (2009)

fit c_D and c_E to $B(A=3)$ and GT_{exp} with N3LO/N2LO
 $\Rightarrow \{c_D; c_E\}_{MAX}$ and $\{c_D; c_E\}_{MIN}$



Remaining LEC's: g_{4S} and g_{4V} in the vector current \Rightarrow fit to the $A = 3$ magnetic moments

	$\{c_D; c_E\}$	g_{4S}	g_{4V}
$\Lambda = 500 \text{ MeV}$	$\{-0.20; -0.208\}$	0.207 ± 0.007	0.765 ± 0.004
	$\{-0.04; -0.184\}$	0.200 ± 0.007	0.771 ± 0.004
$\Lambda = 600 \text{ MeV}$	$\{-0.32; -0.857\}$	0.146 ± 0.008	0.585 ± 0.004
	$\{-0.19; -0.833\}$	0.145 ± 0.008	0.590 ± 0.004

Radiative corrections¹ ARE included

¹ Czarnecki *et al.*, PRL **99**, 032003 (2007)

Results: $\Gamma^D(\mu^- + d)$ and $\Gamma_0(\mu^- + {}^3\text{He}) \rightarrow \chi\text{EFT}$

	1S_0	3P_0	3P_1	3P_2	Γ^D	Γ_0
IA – $\Lambda = 500$ MeV	238.8	21.1	44.0	72.4	381.7	1362
IA – $\Lambda = 600$ MeV	238.7	20.9	43.8	72.0	380.8	1360
FULL – $\Lambda = 500$ MeV	254.4(9)	20.5	46.8	72.1	399.2(9)	1488(9)
FULL – $\Lambda = 600$ MeV	255(1)	20.3	46.6	71.6	399(1)	1499(9)

$$\Gamma^D = 399(3) \text{ s}^{-1} \quad \& \quad \Gamma_0 = 1494(21) \text{ s}^{-1}$$

vs. $\Gamma^D(\text{exp}) \cdots \quad \& \quad \Gamma_0(\text{exp}) = 1496(4) \text{ s}^{-1}$

Comparison between Γ_0 and $\Gamma_0(\text{exp}) \rightarrow$

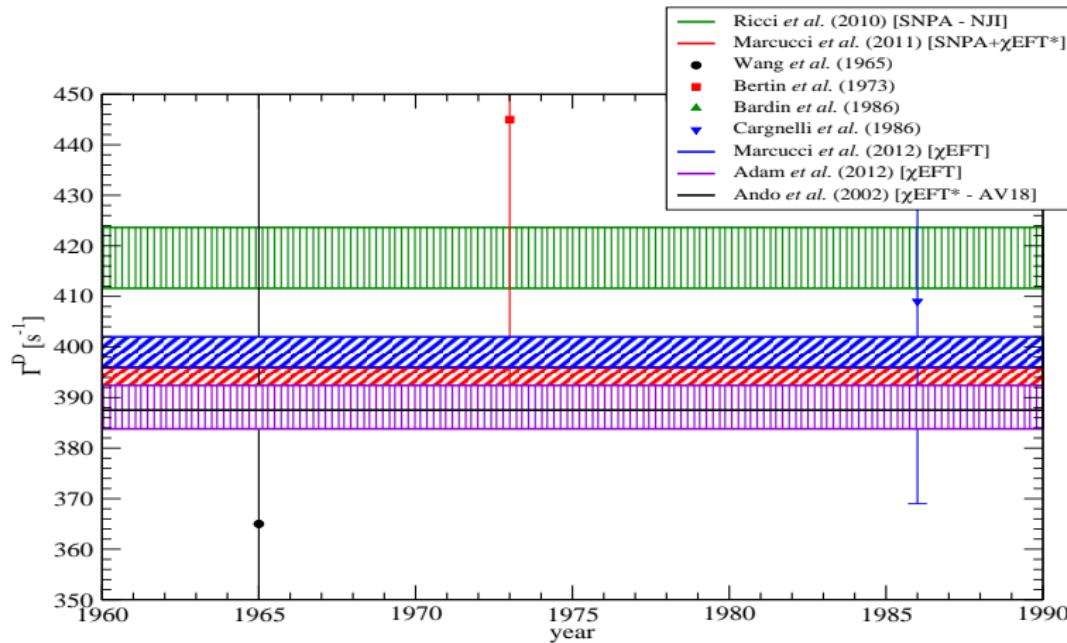
$$G_{PS} = 8.2 \pm 0.7$$

vs.

$$G_{PS}^{\chi\text{PT}} = 7.99 \pm 0.20$$

Marcucci *et al.*, PRL 108, 052502 (2012)

Comparison with data and previous calculations



$p + p \rightarrow d + e^+ + \nu_e$ astrophysical factor

$$\sigma(E) = \frac{1}{(2\pi)^3} \frac{G_V^2}{v} m_e^5 f(E) \sum_M |\langle d, M | \mathbf{A}_- | pp \rangle|^2$$

$$S_{11}(E) = S_{11}(0) + S'_{11}(0)E + \frac{1}{2} S''_{11}(0)E^2 + \dots$$

Goal: < 1% accuracy

- Dominant contribution from the 1S_0 wave
- P -wave contribution: $\sim 1\%$
- Two-body contribution: $\sim 1\%$

pp wave function

- EM interaction:
 $V_{C1} + V_{C2} + V_{DF} + V_{VP} + \dots$
- $V_{VP} \sim \exp(-2m_e r)$: sizeable effect at low energies
- Necessity to solve the Schroedinger equation up to 1,000 fm
- 1% effect

SNPA: [Schiavilla *et al.*, 1998], χ EFT*: [Park *et al.*, 2003]

$$S_{11}(0) = 3.94(1 \pm 0.0015 \pm 0.0010 \pm \epsilon) \text{ (only the } {}^1S_0 \text{ wave)}$$

errors from uncertainties in g_A , fit of the tritium β -decay, etc; ϵ “systematic error” (?)

See also the review paper: [E. G. Adelberger *et al.*, Rev. Mod. Phys. 83, 195 (2011)
[arXiv:1004.2318]]

Results with the ‘Less hybrid’ χ EFT

N3LO potential – PRELIMINARY

d_R , g_{4S} , g_{4V} fixed using the N3LO/N2LO wave functions - only V_{C1} EM int.
calculation performed between $0 < E < 10$ keV

	$S_{11}(0) [\times 10^{-25} \text{ MeV b}]$	$S'_{11}(0)/S_{11}(0) [\text{MeV}^{-1}]$	
	1S_0	$S + P$	1S_0
IA(500)	3.96	3.98	11.16
IA(600)	3.94	3.96	11.17
FULL(500)	4.025(5)	4.052(5)	11.17
FULL(600)	4.007(5)	4.033(5)	11.17

Summary:

1S_0

$$S_{11}(0) = 4.00 \div 4.03 \times 10^{-25} \text{ MeV b}$$

$$\frac{S'_{11}(0)}{S_{11}(0)} = 11.17 \text{ MeV}^{-1}$$

All waves

$$S_{11}(0) = 4.03 \div 4.06 \times 10^{-25} \text{ MeV b (1\%)}$$

$$\frac{S'_{11}(0)}{S_{11}(0)} = 11.68 \text{ MeV}^{-1} (4\%)$$

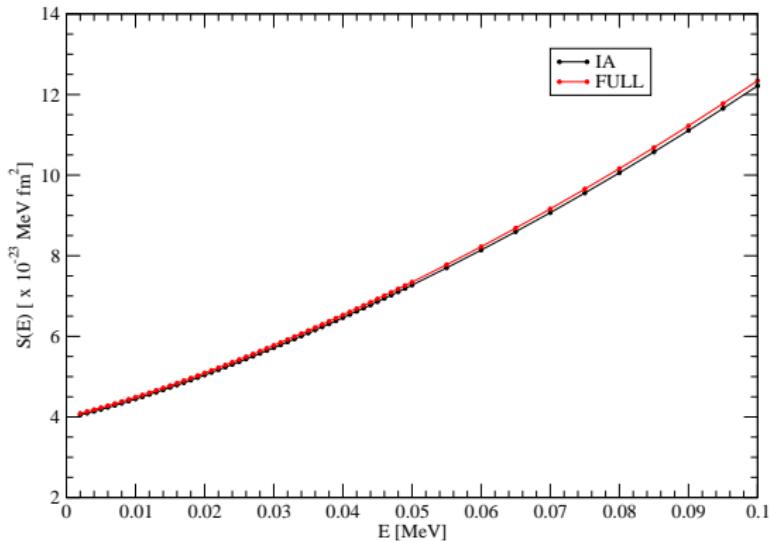
Pionless EFT at N²LO [Wei *et al.*, 2012]

$$S_{11}(0) = (3.99 \pm 0.14)10^{-25} \text{ MeV b}, S'_{11}(0)/S_{11}(0) = (11.3 \pm 0.1) \text{ MeV}^{-1}$$

New interest: $S''_{11}(0)$

$S(E)$ calculated in the range 0 – 100 keV

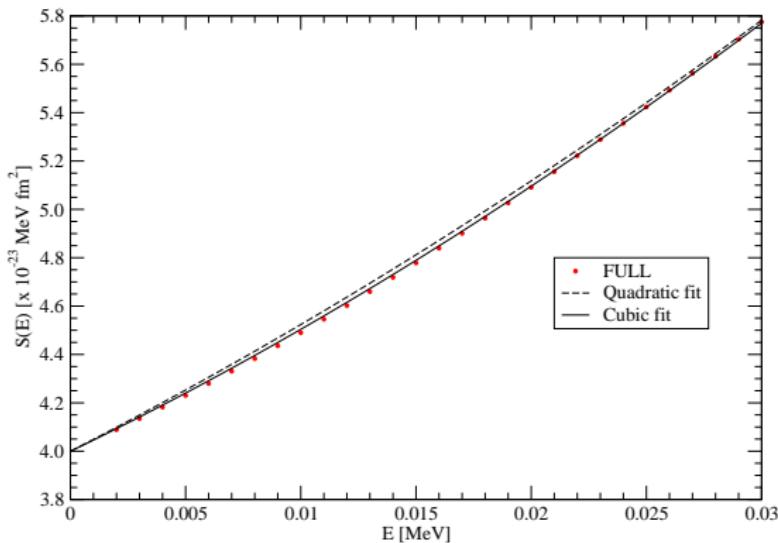
Effect of two-body currents



Calculation with AV18 - only the 1S_0 wave

$S(E)$ calculated in the range 0 – 100 keV

Test of the quadratic approximation



Calculation with AV18 - only the 1S_0 wave

Small effects

Quadratic vs. cubic fit

$$S(E) = S(0) + a_0 E + a_1 E^2$$

$$S(E) = S(0) + a_0 E + a_1 E^2 + a_2 E^3$$

	$S(0) \times 10^{23}$ [MeV fm ²]	$S'(0)/S(0)$ [MeV ⁻¹]	$S''(0)/S(0)$ [MeV ⁻²]
quadratic fit	4.00	12.23	175.0
cubic fit	4.00	11.47	233.3
Wei <i>et al.</i>	3.99 ± 0.14	11.3 ± 0.1	170 ± 2

Effect of the EM interactions $V_{C2} + V_{DF} + V_{VP} + \dots$

Calculation performed so far only for the AV18 interaction

	$S(0) \times 10^{23}$ [MeV fm ²]	$S'(0)/S(0)$ [MeV ⁻¹]	$S''(0)/S(0)$ [MeV ⁻²]
AV18+ V_{C1}	4.03	11.59	226.5
AV18+ V_{EM}	4.00	11.47	233.3

Conclusions and outlook

- Extensive theoretical work on muon capture in SNPA, χ EFT* and χ EFT
- $\Gamma_0(\mu^- + {}^3\text{He})$: nice agreement theory vs. experiment
- $\Gamma^D(\mu^- + d)$:
 - some **discrepancies** among different theoretical works
 - more accurate experimental results** → MuSun
- New refined calculation of the $p p$ fusion up to 100 keV (in progress)
- In the future:
 - χ EFT → $\mu^- + {}^3\text{He} \rightarrow n + d + \nu_\mu$
 $\mu^- + {}^3\text{He} \rightarrow n + n + p + \nu_\mu$
 - χ EFT → reactions of astrophysical interest
 - $p + {}^3\text{He} \rightarrow {}^4\text{He} + e^+ + \nu_e$
 - $p + d \rightarrow {}^3\text{He} + \gamma$
 - $d + d \rightarrow {}^4\text{He} + \gamma$
 -
- New model for \mathbf{j}^A, ρ^A consistent with the EM charge/current derived by **Pastore et al.** (→ Pastore's talk)