

Microscopic description of exotic nuclei in the Berggren basis

J. Rotureau

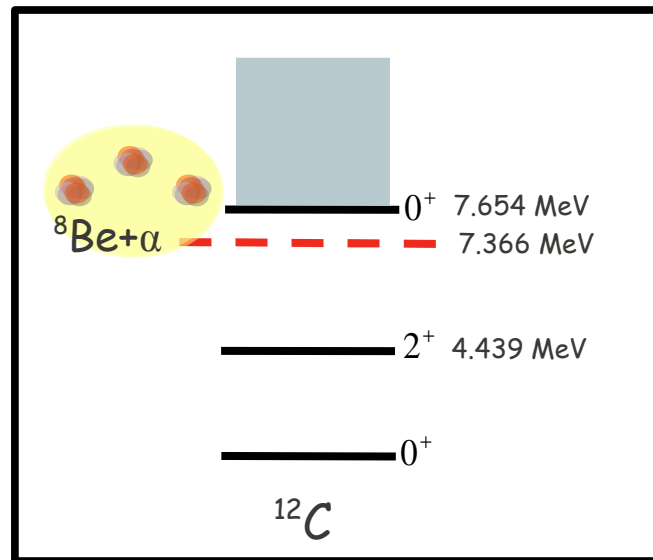
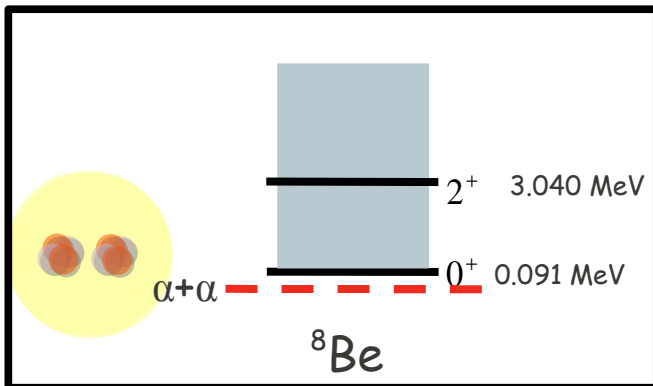
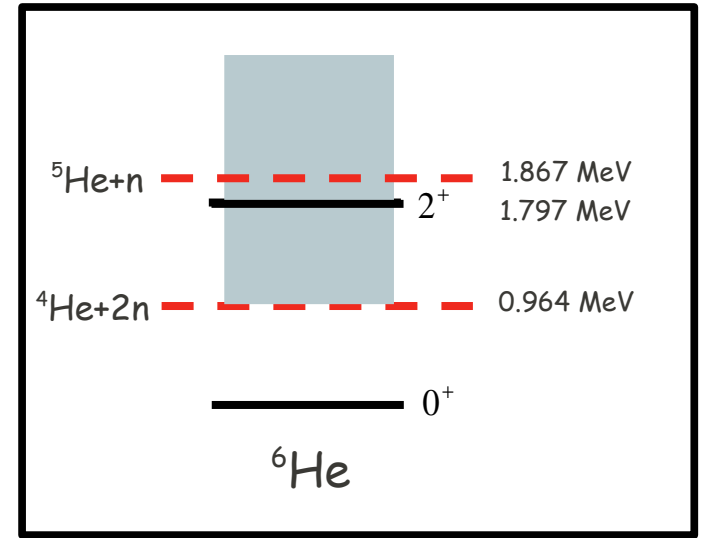
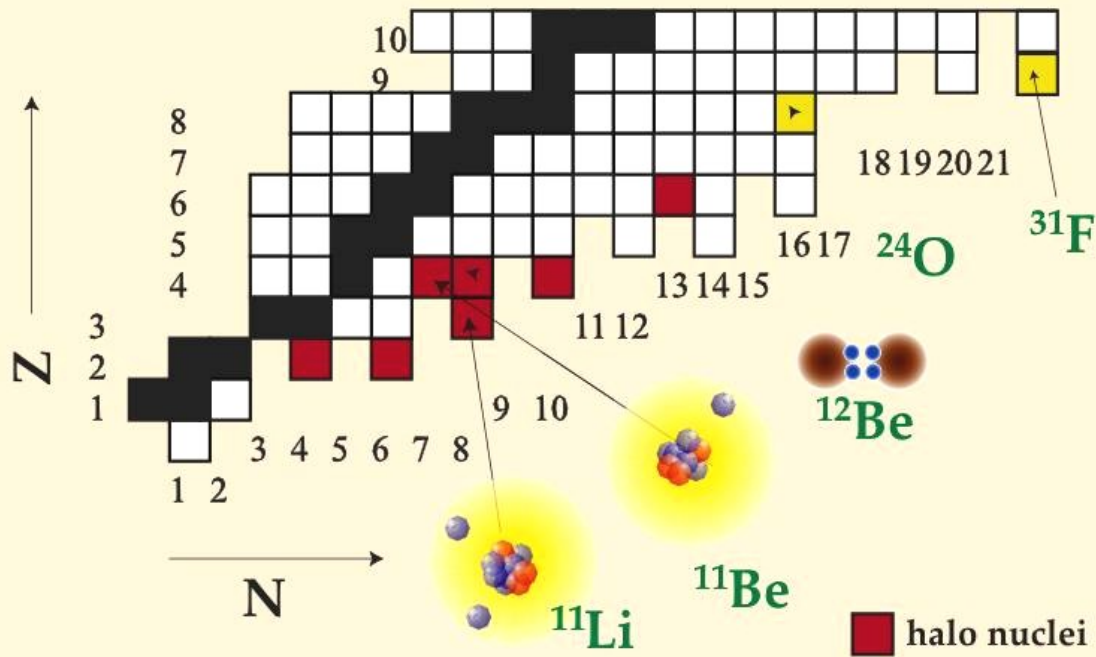


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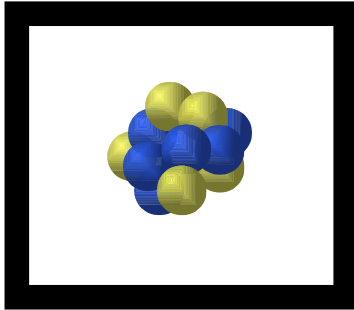


Light drip line nuclei

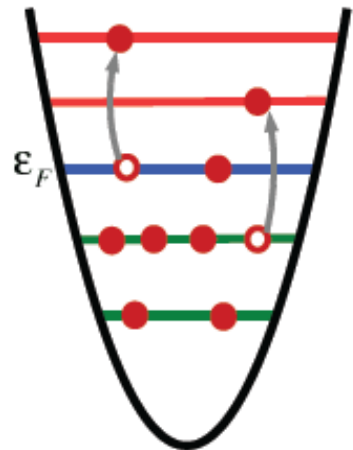


Importance of continuum in the structure of nuclei far from stability

Closed quantum systems



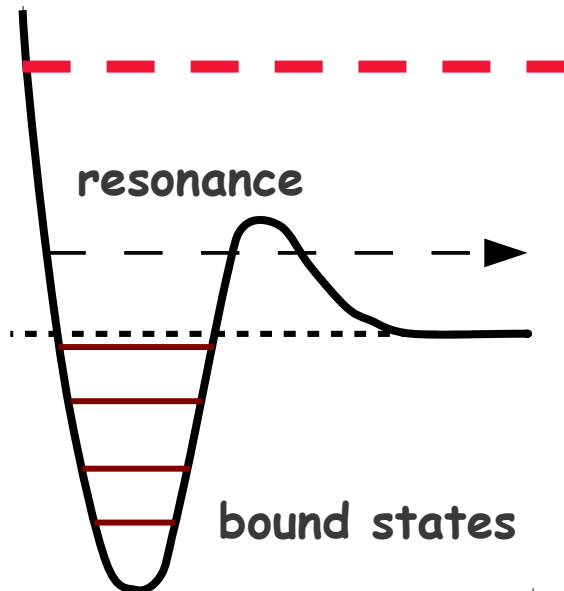
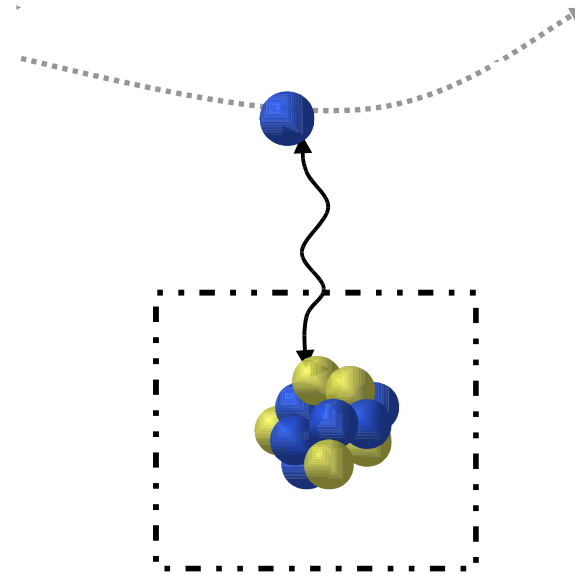
infinite well



discrete states only,
HO basis usually

exact treatment of the
c.m, analytical solution...

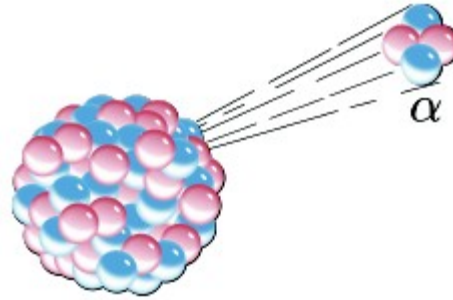
Open quantum systems (nuclei far from stability)



Gamow States

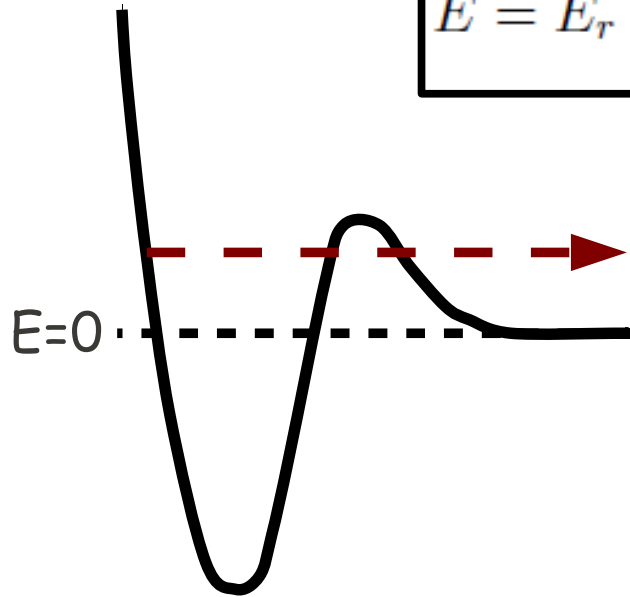
G. Gamow, Z. Phys. 51 (1928) 204

$$\tilde{E} = E_r - i\frac{\Gamma}{2}$$



$$\Psi(t, r) = e^{\frac{-i\tilde{E}t}{\hbar}} \psi(r)$$

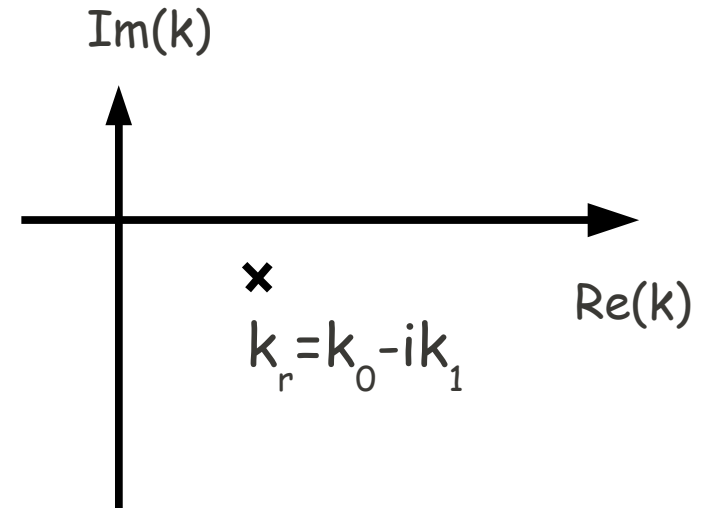
$$|\Psi(t, r)| \sim e^{-\frac{\Gamma t}{2\hbar}} e^{k_1 r}, r \rightarrow \infty$$



$$\left(-\frac{d^2}{dr^2} + \frac{2\mu}{\hbar^2} V(r) + \frac{l(l+1)}{r^2} - k^2 \right) \psi(r) = 0$$

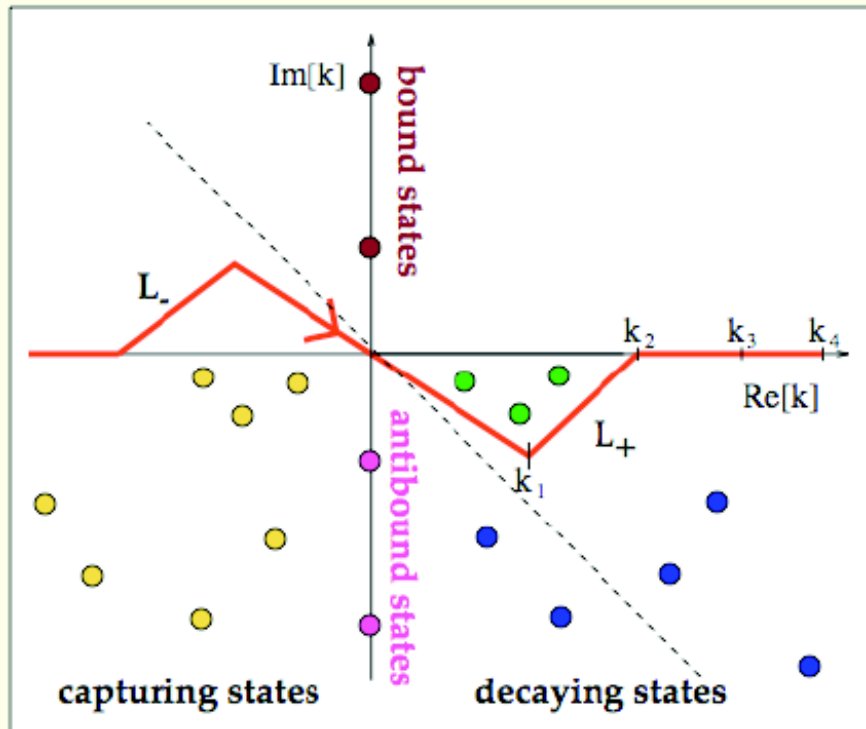
Boundary condition

$$\psi(r) \sim \mathcal{O}_l(kr) \sim e^{ikr}, r \rightarrow \infty$$



Gamow states and completeness relations

T. Berggren, Nucl. Phys. A109, 265 (1968); Nucl. Phys. A389, 261 (1982)
T. Lind, Phys. Rev. C47, 1903 (1993)



$$\sum_{n=b,r} |u_n \rangle \langle \tilde{u}_n| + \frac{1}{\pi} \int_{L_+} |u(k) \rangle \langle u(k^*)| dk = 1$$

particular case: Newton completeness relation

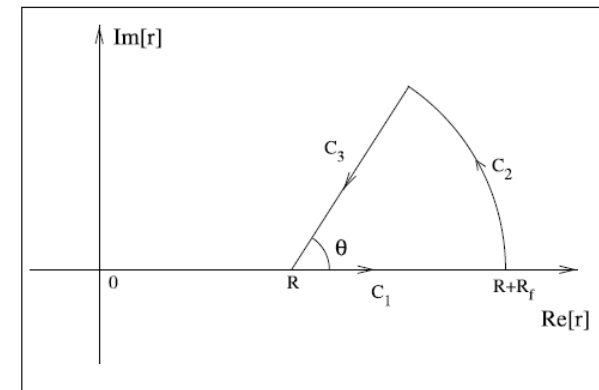
$$\sum_{n=b} |u_n \rangle \langle \tilde{u}_n| + \frac{1}{\pi} \int_R |u(k) \rangle \langle u(k^*)| dk = 1$$

Bound, resonant state

$$u(r) \rightarrow C_+ H_{l,\eta}^+(kr)$$

normalization of resonant states
with external complex scaling :

$$N_i = \sqrt{\int_0^R u_i^2(r) dr + \int_0^{+\infty} u_i^2(R+x \cdot e^{i\theta}) e^{i\theta} dx}$$

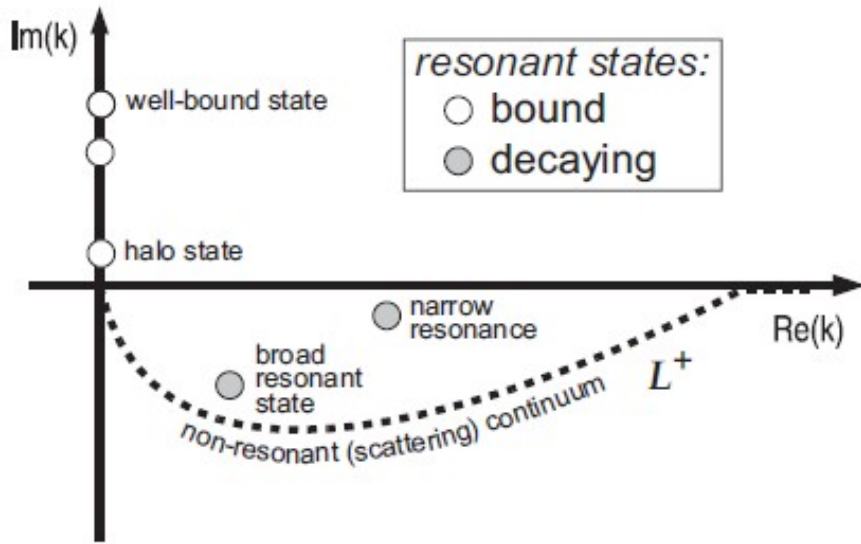


Complex scattering state

$$u(r) \rightarrow C_+ H_{l,\eta}^+(kr) + C_- H_{l,\eta}^-(kr)$$

$$C^+ C^- = \frac{1}{2\pi} \quad (\text{normalisation})$$

Gamow Shell Model



pole approximation: "0th" order approximation :

$$H^{p.a.} |\Psi^{p.a.}\rangle = E^{p.a.} |\Psi^{p.a.}\rangle$$

Many-body resonance (bound) states have the largest overlap

$$|\langle \Psi^{p.a.} | \Psi \rangle|$$

N. Michel *et al*, PRL 89 (2002) 042502; PRC67 (2003) 054311; PRC70 (2004) 064313
 G. Hagen *et al*, PRC71 (2005) 044314
 J.R *et al*, PRL 97 (2006) 110603
 N. Michel *et al*, JPG (2009) 013101
 G.Papadimitriou *et al*, PRC(R) 84 (2011) 051304

i) discretization of continuum contour

$$\sum |u_{res}\rangle \langle u_{res}| + \sum_i |u_{ki}\rangle \langle u_{ki}| \simeq 1$$

ii) construction of many-body basis

$$|SD_i\rangle = |u_{i1} \dots u_{iA}\rangle$$

iii) construction of Hamiltonian matrix

$$\langle SD_i | H | SD_j \rangle$$

(complex symmetric matrix)

iv) many-body spectrum:

bound, resonant and "spurious" continuum states

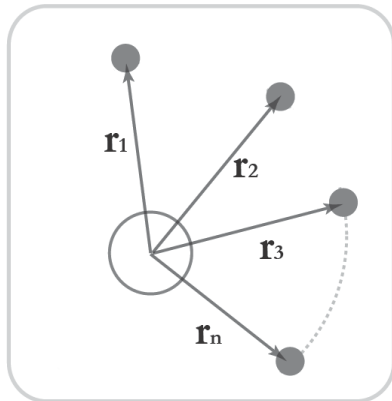
$$H = \sum_i \left[\frac{p_i^2}{2\mu} + U_i \right] + \sum_{i < j} \left[V_{ij} + \frac{1}{A_c} \vec{p}_i \vec{p}_j \right]$$

generic Hamiltonian for a GSM description of a nucleus as core+valence nucleon system

recoil term coming from the expression of H in the COSM coordinates. No spurious states

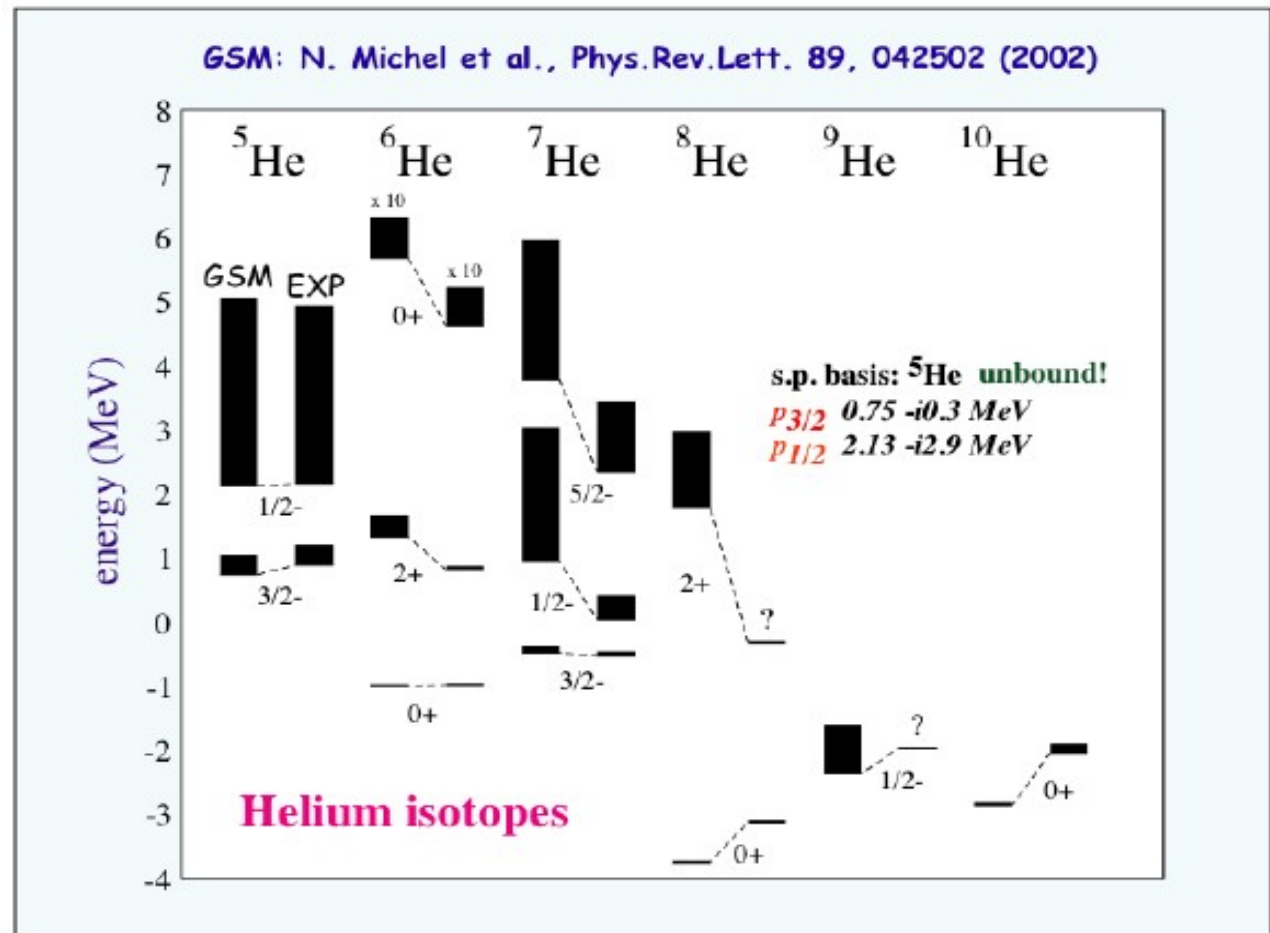
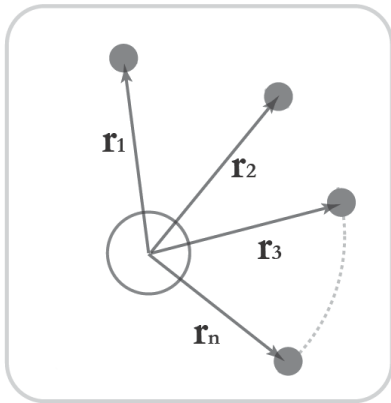
cluster Orbital Shell Model (COSM) coordinates

(Y. Suzuki et, Phys. Rev. C 38, 1 (1988))



- i) U_i core-nucleon potential
- ii) V_{ij} phenomenological, realistic NN interaction

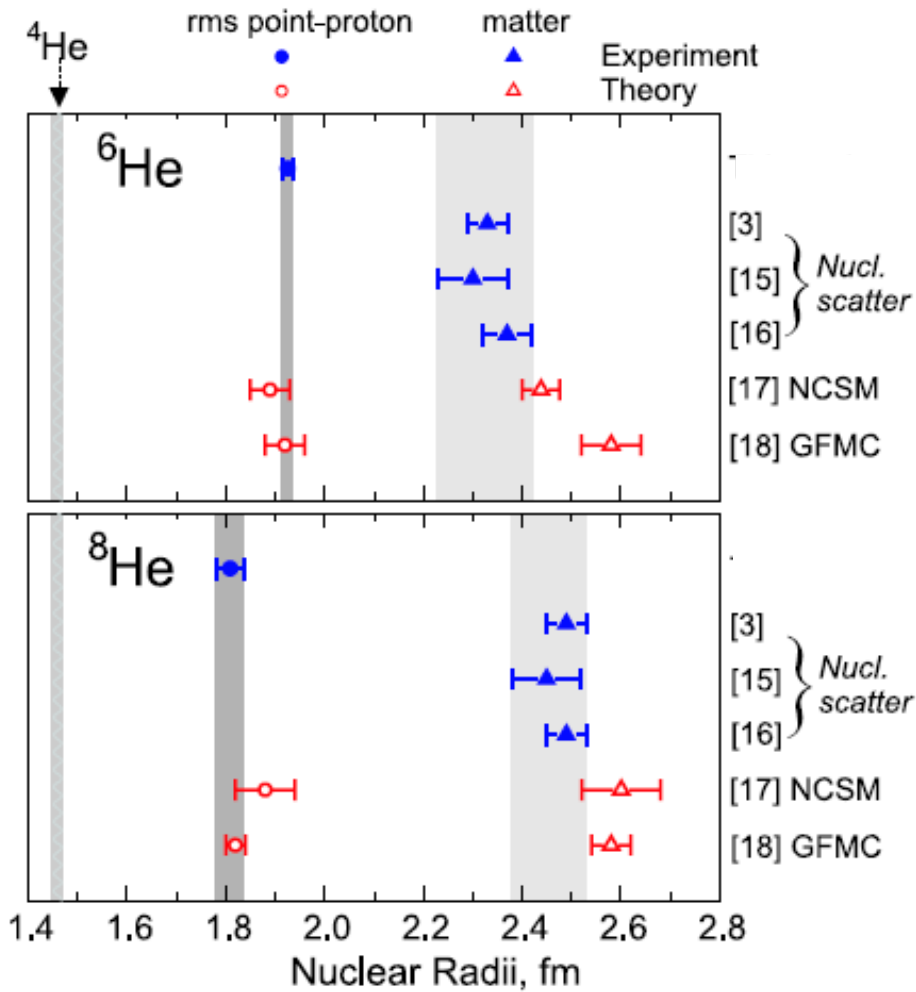
Helium chain (^4He core plus valence neutrons)



- i) Woods-Saxon potential ($^4\text{He-n}$)
- ii) two-body zero-range force (n-n)

pole approximation: $p_{3/2}$, $p_{1/2}$ resonance (^5He g.s and 1st excited state)

${}^6\text{He}, {}^8\text{He}$ charge radii



(taken from P. Muller *et al*,
PRL 99, 252501 (2007))

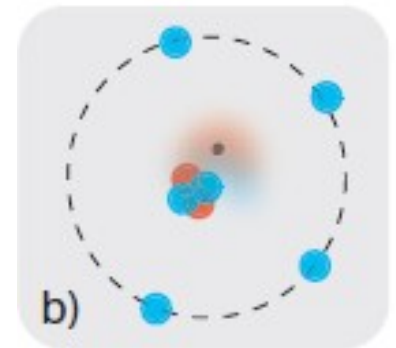
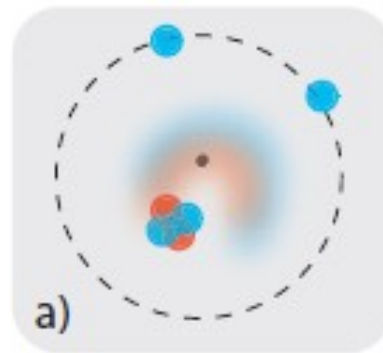
- [16] E. Caurier *et al*, PRC 73, 021302 (R) (2006)
- [17] S.C. Pieper, Riv. Nuovo Cim. 031, 709 (2008).

M. Brodeur *et al.*
PRL 108, 052504 (2012)

${}^6\text{He}$ = 1.910 ± 0.011 fm

${}^8\text{He}$ = 1.835 ± 0.019 fm

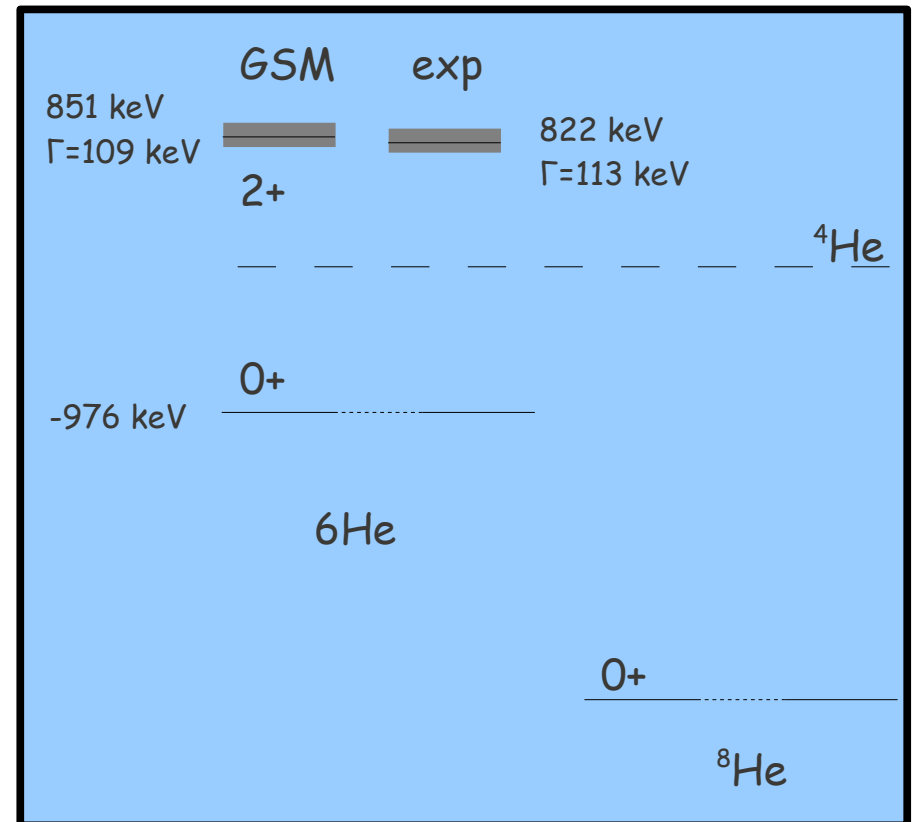
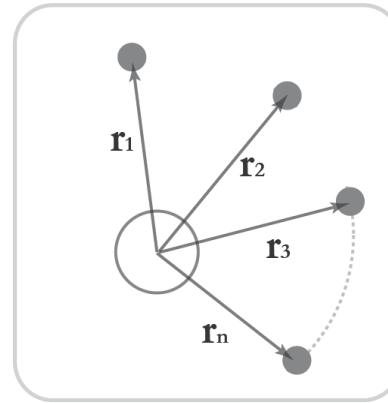
(point charge radius)



structural information on nuclear
hamiltonian and nuclear many-body
dynamics (the radial extent of the halo
nucleus is reflected in the charge radius)

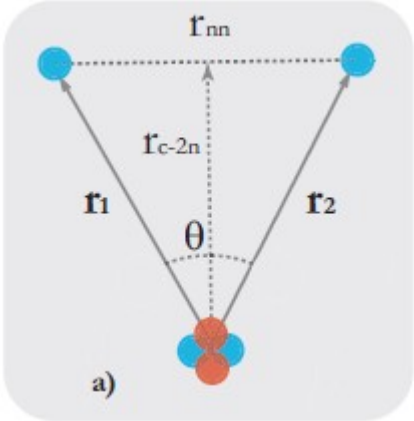
${}^6\text{He}, {}^8\text{He}$ Hamiltonian

- Woods-Saxon potential for ${}^4\text{He}$ -n (fitted to ${}^5\text{He}$ resonances)
- "Minnesota like" interaction, 2 parameters (adjusted to ${}^6\text{He}, {}^8\text{He}$ g.s.)
- $p_{3/2}$ resonance + $p_{3/2}$ complex continuum
- $p_{1/2}$ sd real continuum



${}^6\text{He}$ g.s

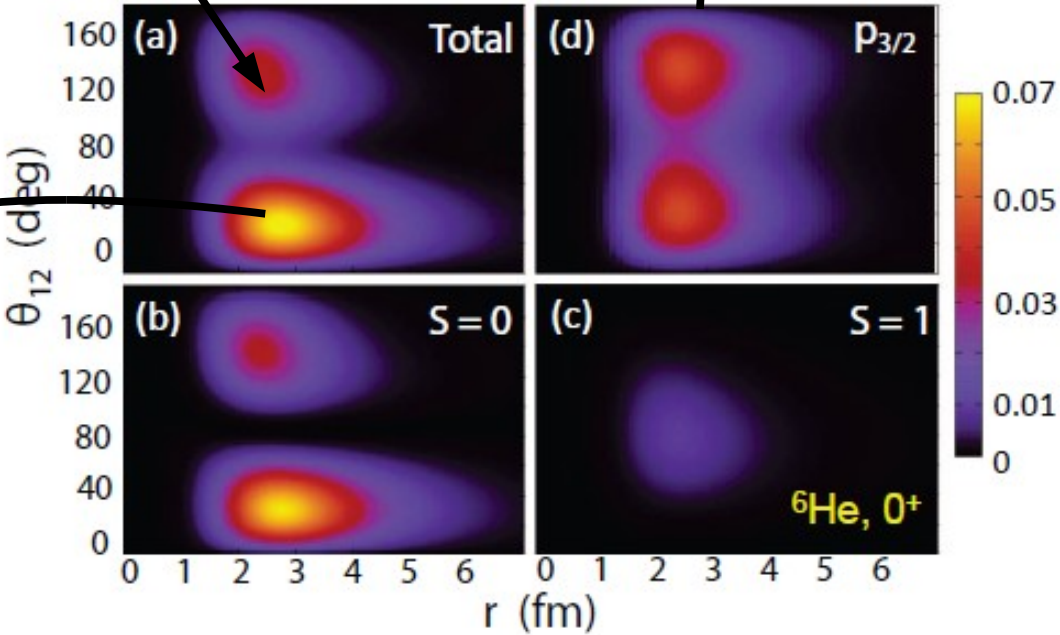
two-neutron density $\rho_{nn}(r_1=r, r_2=r, \theta)$



cigar-like configuration

$p_{3/2}$ only, no $p_{1/2}$ sd (continuum) shells

"dineutron"-like configuration



small components (~ 10 %) but dramatic effect !

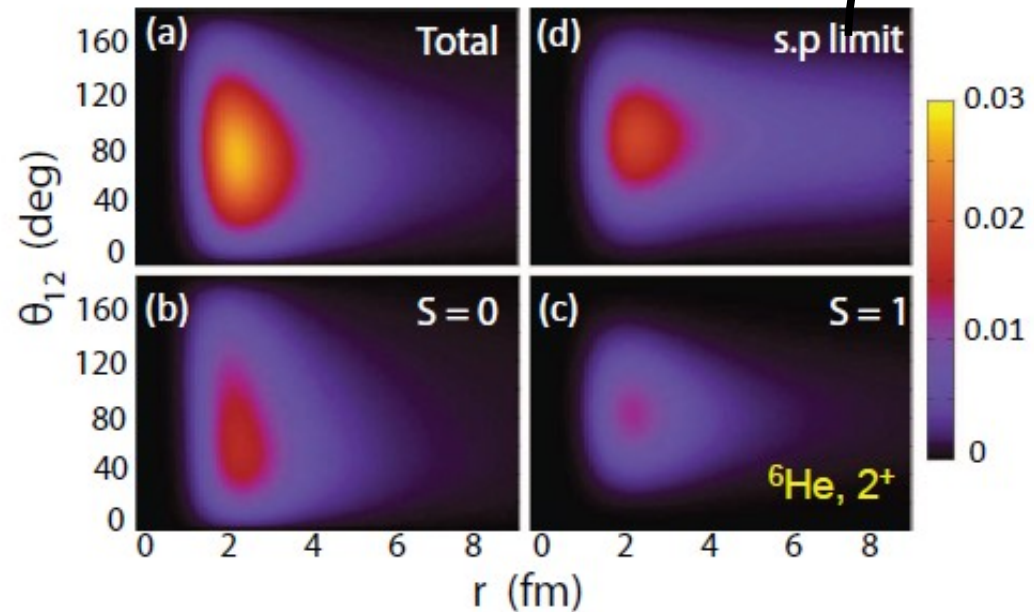
$S=0 \sim 87 \%$
 $S=1 \sim 13 \%$

${}^6\text{He}$ resonant
state $J^\pi=2^+$

two-neutron density $\rho_{nn}(r_1=r, r_2=r, \theta)$

unbound state

single particle "limit"



$S=0 \sim 66\%$
 $S=1 \sim 33\%$

G. Papadimitriou, A. T. Kruppa, N. Michel, W. Nazarewicz M. Płoszajczak
and J. R, PRC 84 (2011)

Charge radii

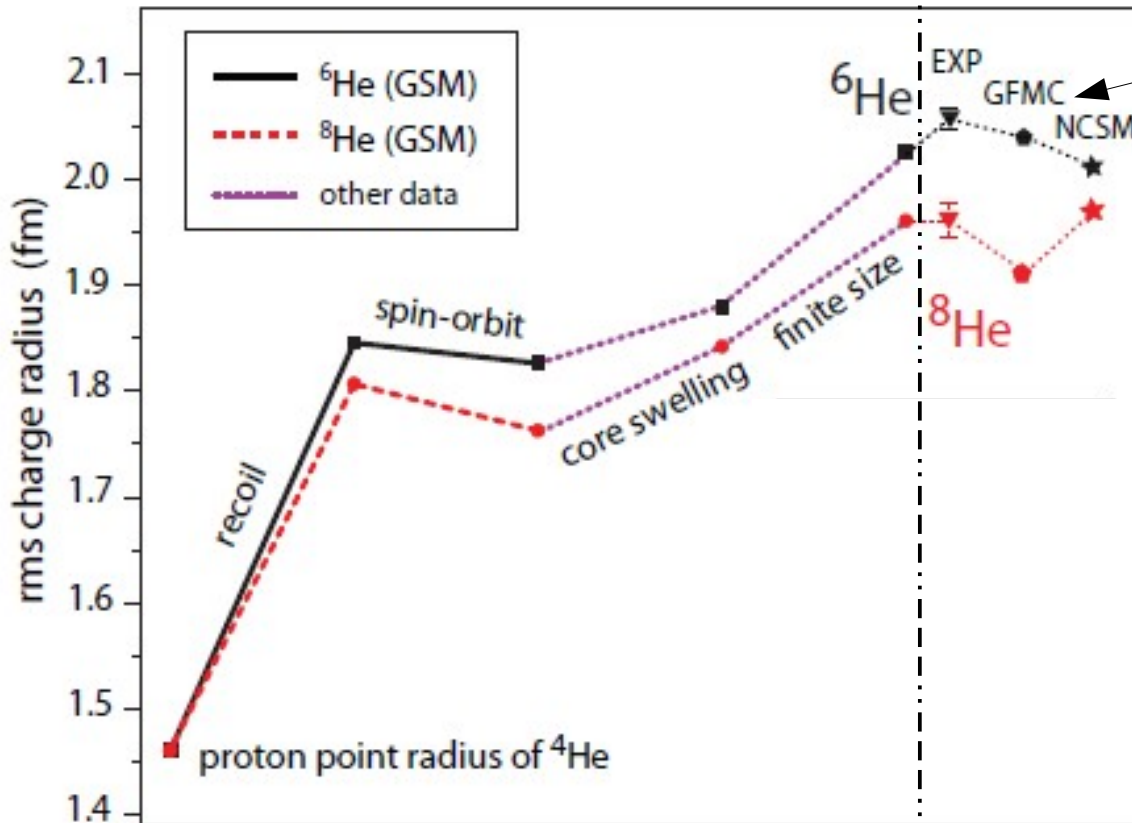
point proton radius finite size corrections

$$\langle r_{\text{ch}}^2 \rangle = \langle r_{\text{pp}}^2 \rangle + \langle R_{\text{p}}^2 \rangle + \frac{N}{Z} \langle R_{\text{n}}^2 \rangle + \frac{3}{4M_{\text{p}}^2} + \langle r^2 \rangle_{\text{so}}$$

$$\langle r_{\text{pp}}^2 \rangle = \langle r_{\text{pp}}^2(^4\text{He}) \rangle + \langle \text{recoil} \rangle$$

- "swelling" of ^4He from GFMC
- recoil from GSM

nucleons charge radius +
Darwin-Foldy term + Spin-
Orbit contribution (obtained
from GSM wf)



Experimental data and other
theoretical approaches

	GSM	Exp
^6He	2.026 fm	2.059(7) fm
^8He	1.961 fm	1.959(16) fm

Density Matrix Renormalization Group (DMRG)

S. R. White, Phys. Rev. Lett. 69 (1992) 2863

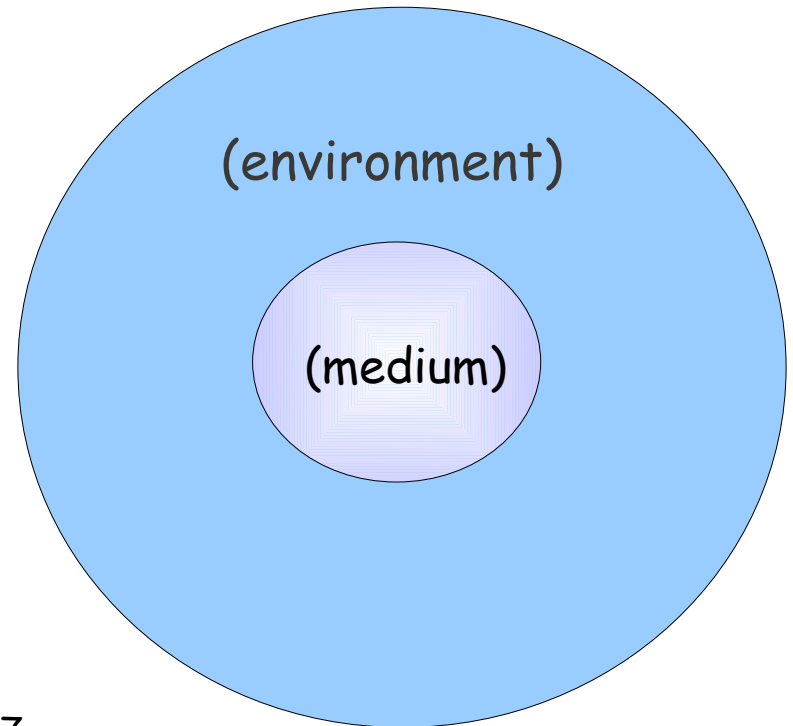
S. R. White, Phys. Rev. B 48 (1993) 10345

S.R. White et al, Phys. Rev. B 48 (1993) 3844

lattice models, spin chain, quantum dots, atomic nuclei.....

Reduction of the number of degrees of freedom + renormalization

- * Separation into a 'medium' and 'environment'
- * Truncation of degrees of freedom in the environment



Application
for nuclei

T.Papenbrock et al J.Phys.G 31 (2005) S1377

J. R et al, PRL 97 (2006) 110603

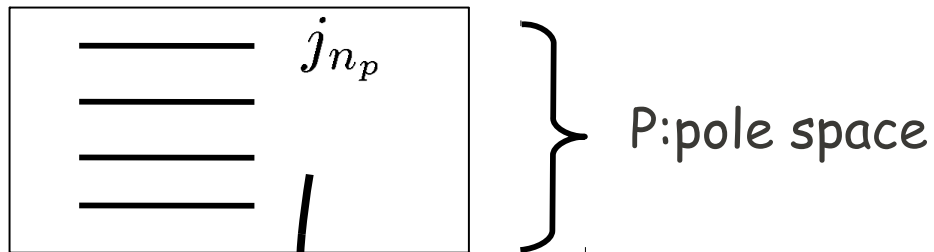
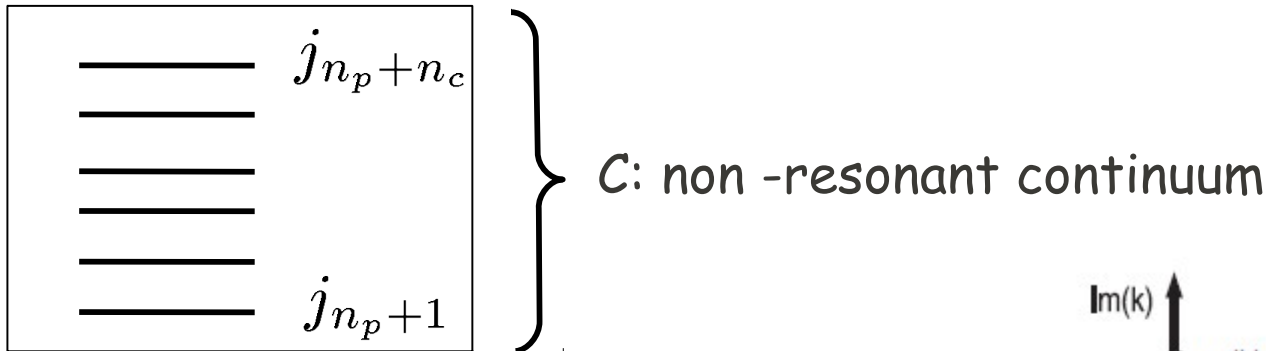
S.Pittel et al PRC 73 (2006) 014301 (R)

B. Thakur et al, Phys. Rev. C 78 (2008) 041303(R)

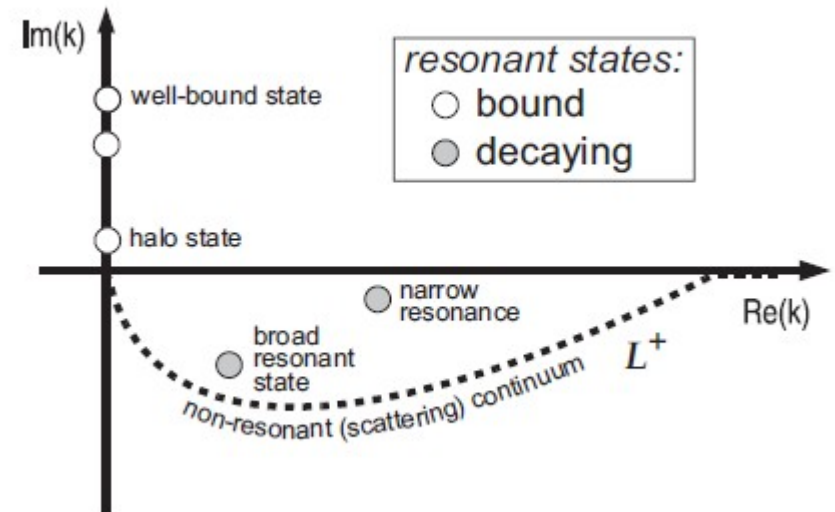
J.R et al, PRC 79 (2009) 014304

GSM+DMRG

$$|\Psi\rangle^J = \sum_{p,c} \Psi_{pc} [|p\rangle^{J_p} |c\rangle^{J_c}]^J$$



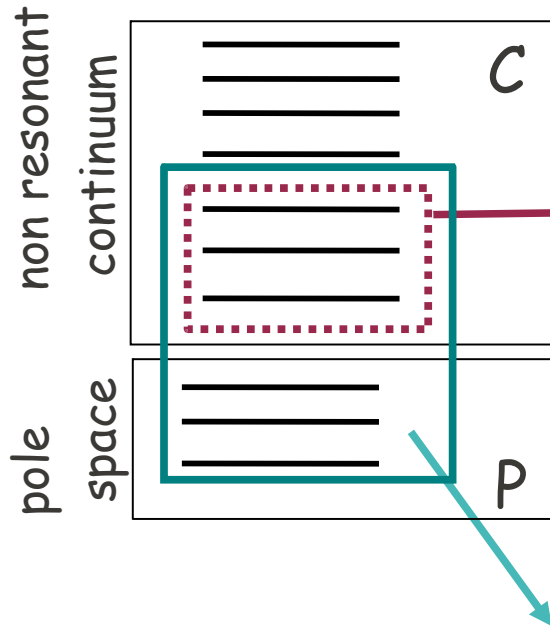
largest contribution to the GSM wave function



DMRG \longrightarrow truncation in environment C

Warm up phase

Construction of 2nd quantization operators and states in P and C



block

$|c\rangle$: states with 0,1,...n nucleons

operators : $a_i^\dagger, (a_j^\dagger a_k^\dagger)^K, [(a_i^\dagger a_j^\dagger)^{K_1} \tilde{a}_k]^{K_2} \dots$

*shells in C added one by one
one step=one shell*

* diagonalization in the superblock



$$|\Psi\rangle^J = \sum_{p,c} \Psi_{pc} (|p\rangle^{J_p} |c\rangle^{J_c})^J$$

(state with the largest overlap with the pole approx)

* singular value decomposition

* diagonalization of the density matrix :

$$\rho_{c,c'}^{J_c} = \sum_p \Psi_{pc} \Psi_{pc'}$$

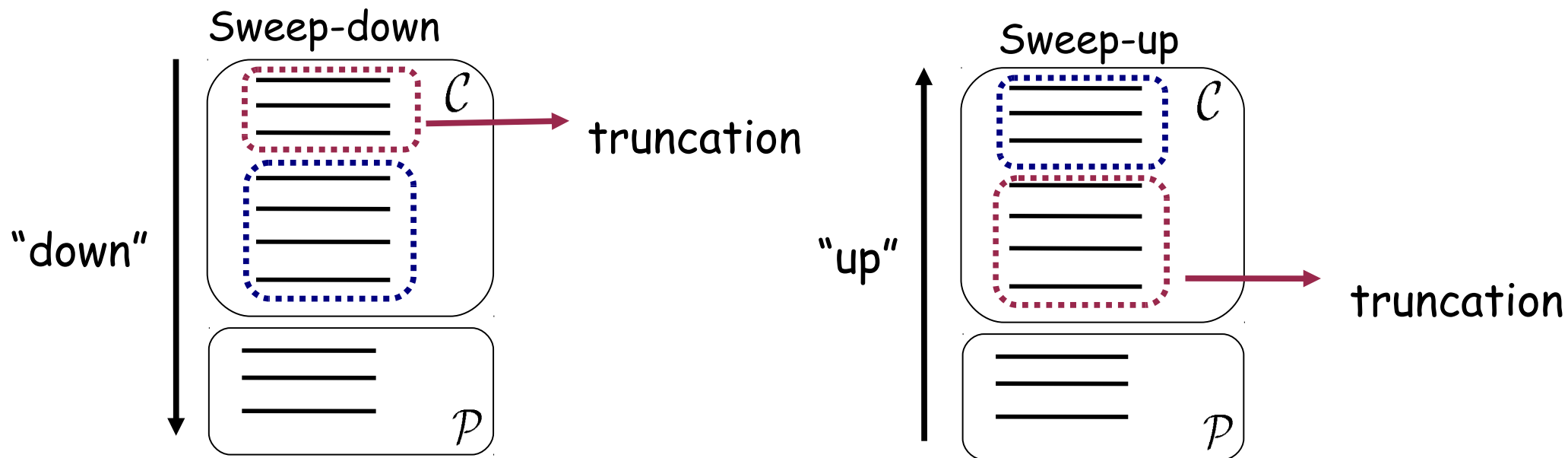


eigenstates with "largest" eigenvalues are kept.

Eigenvalues of the density are probabilities :

$$\sum_{\alpha} w_{\alpha} = 1$$

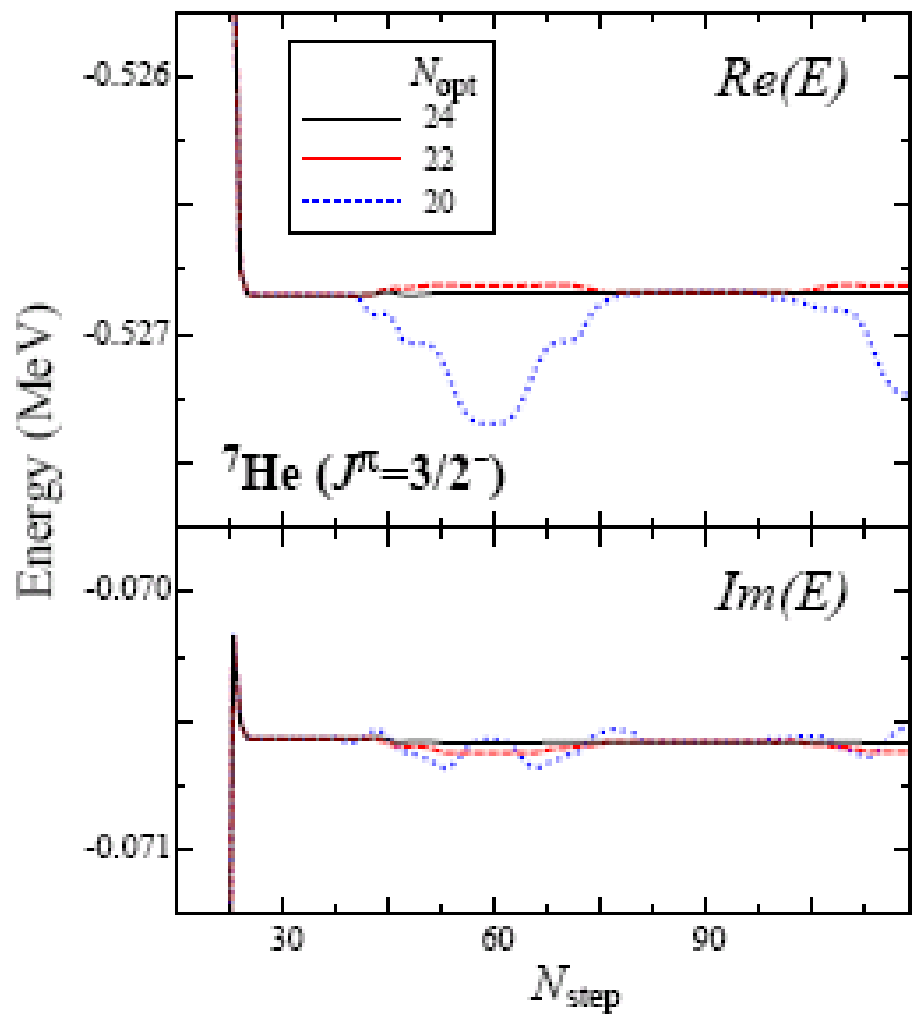
Sweeping phase



${}^7\text{He}$ g.s.

${}^4\text{He}$ core + 3 neutrons

Convergence of the energy as a function of DMRG iteration



- × pole space : $0p_{3/2}, 0p_{1/2}$
- × continuum space : $p_{3/2}, p_{1/2}$
(30 shells each)

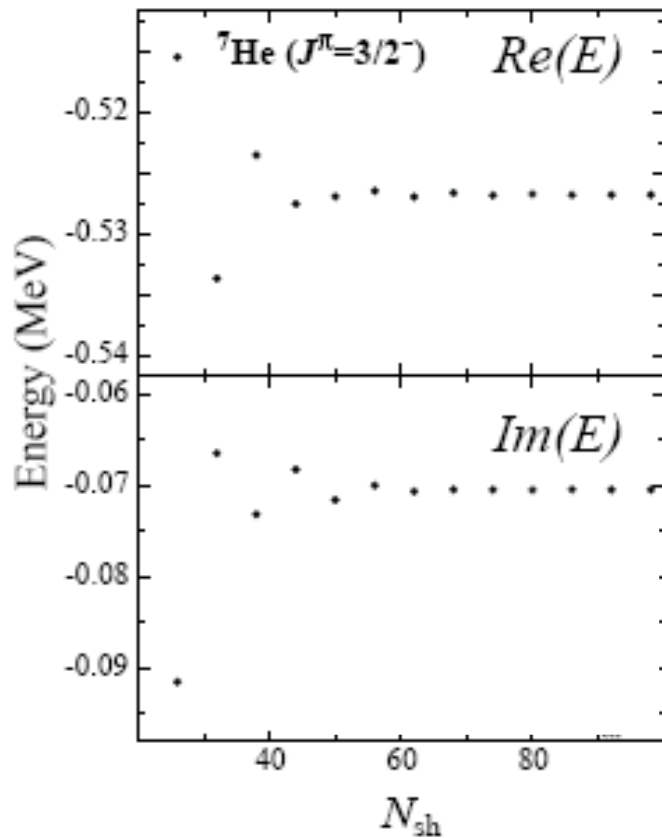
Woods-Saxon + Surface Gaussian two-body interaction :

$$V_{ij}^{J,T} = V_0(J, T) \exp \left[- \left(\frac{r_1 - r_2}{\mu} \right)^2 \right] \delta(|r_1| + |r_2| - 2R_0)$$

Shell Model dimension=83948
largest matrix in DMRG=1143

${}^7\text{He}$ g.s.

Convergence of the real (top) and imaginary part (bottom) of the g.s. energy as a function of the total number of shells



Shell Model dimension
6149 \longleftrightarrow 332171

DMRG dimension:
941 \longleftrightarrow 1001

Very good scaling
with number of shells !

DMRG truncation at $N_{\text{opt}}=22$

Ab-Initio calculations in the Berggren basis

$$H = \frac{1}{A} \sum_{i < j}^A \frac{(\vec{p}_i - \vec{p}_j)^2}{2m} + V_{NN,ij}$$

i) NN potential:

* AV18 (R.B. Wiringa et al PRC 51 (1995) 38)

* N³LO (D.R. Entem et al PRC(R) 68 (2003) 041001)

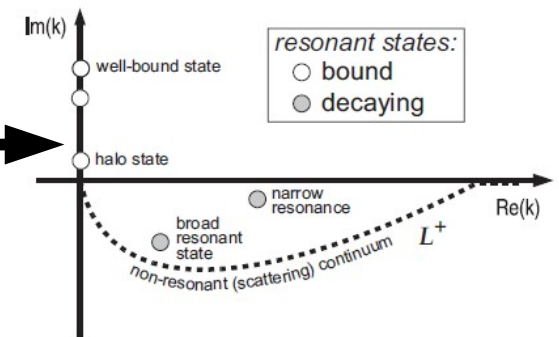
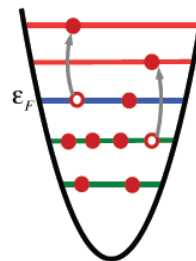
(For comparison with Faddeev, Faddeev-Yakubovsky and Coupled Cluster)

} softened by $v_{\text{low-k}}$ with $\Lambda = 1.9 \text{ fm}^{-1}$
(S. Bogner et al, Phys. Rep. 386 (2003) 1)

ii) single particle states:

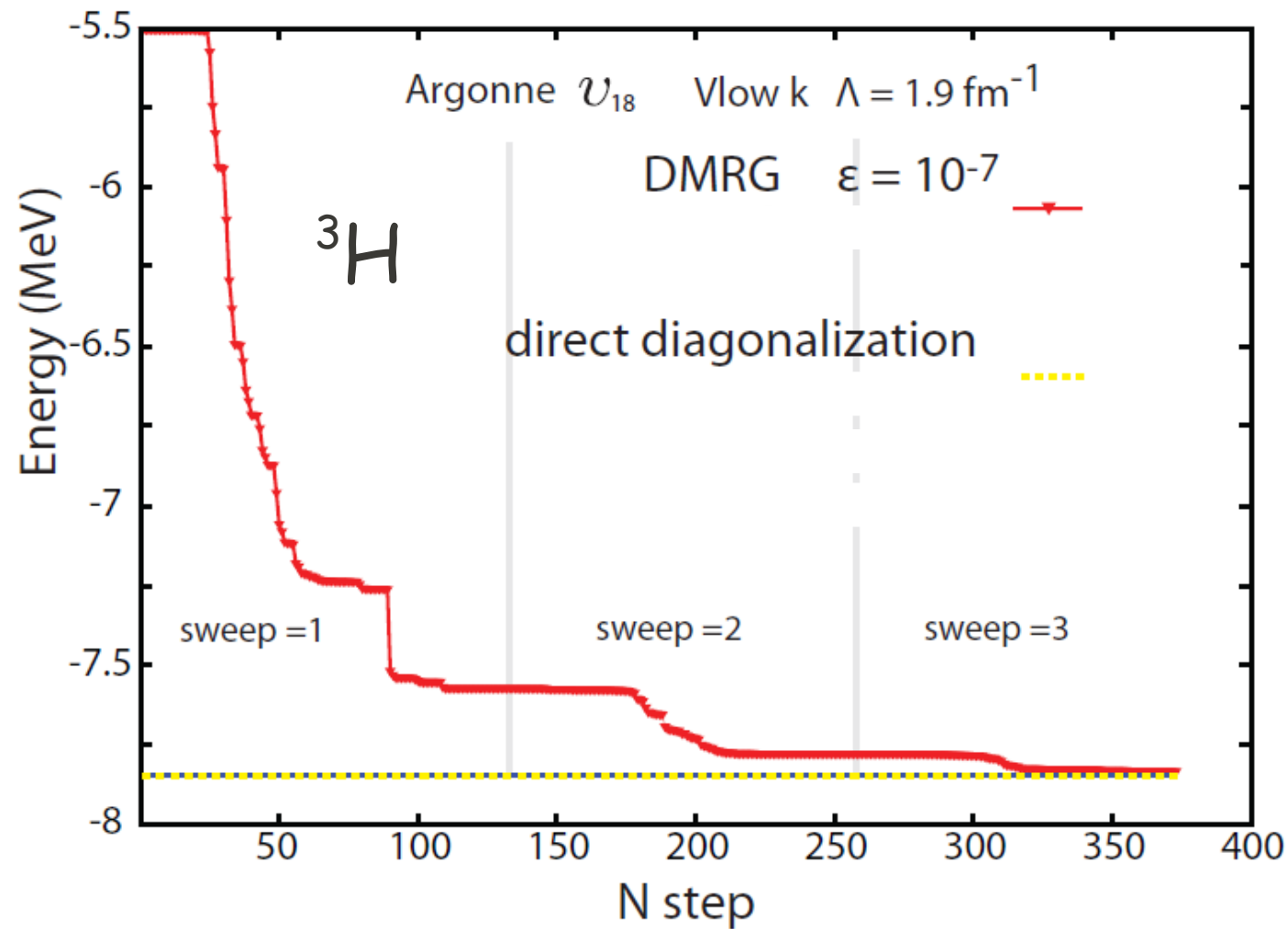
a) s- and p-shells from Hartree-Fock potential

b) for $l > 1$, shells of the Harmonic Oscillator



iii) Resolution with DMRG

Calculations of ³H, ⁴He and ⁵He



Shells

$0s_{1/2}(p) : E = -9.231 \text{ MeV}$
 $0s_{1/2}(n) : E = -11.765 \text{ MeV}$

$s_{1/2}, p_{3/2}, p_{1/2}, s_{1/2}$ real energy continua

$d_{5/2}, d_{3/2}$ H.O states

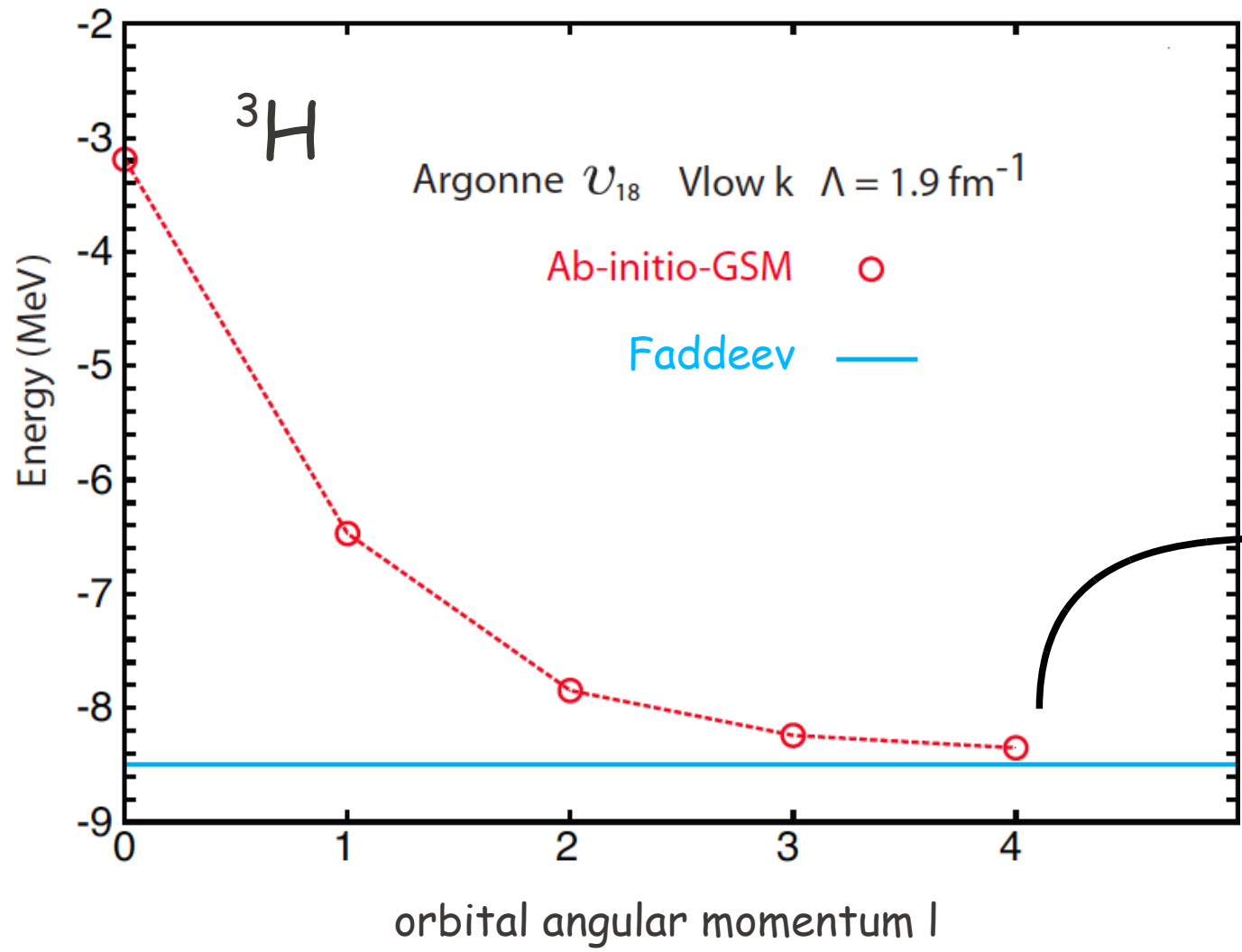
130 s.p. in total

GSM full diagonalisation: dim= 123,835
 DMRG : dim~ 1200

$$E_{\text{exact}} = -7.840 \text{ MeV}$$

$$E_{\text{DMRG}} = -7.832 \text{ MeV}$$

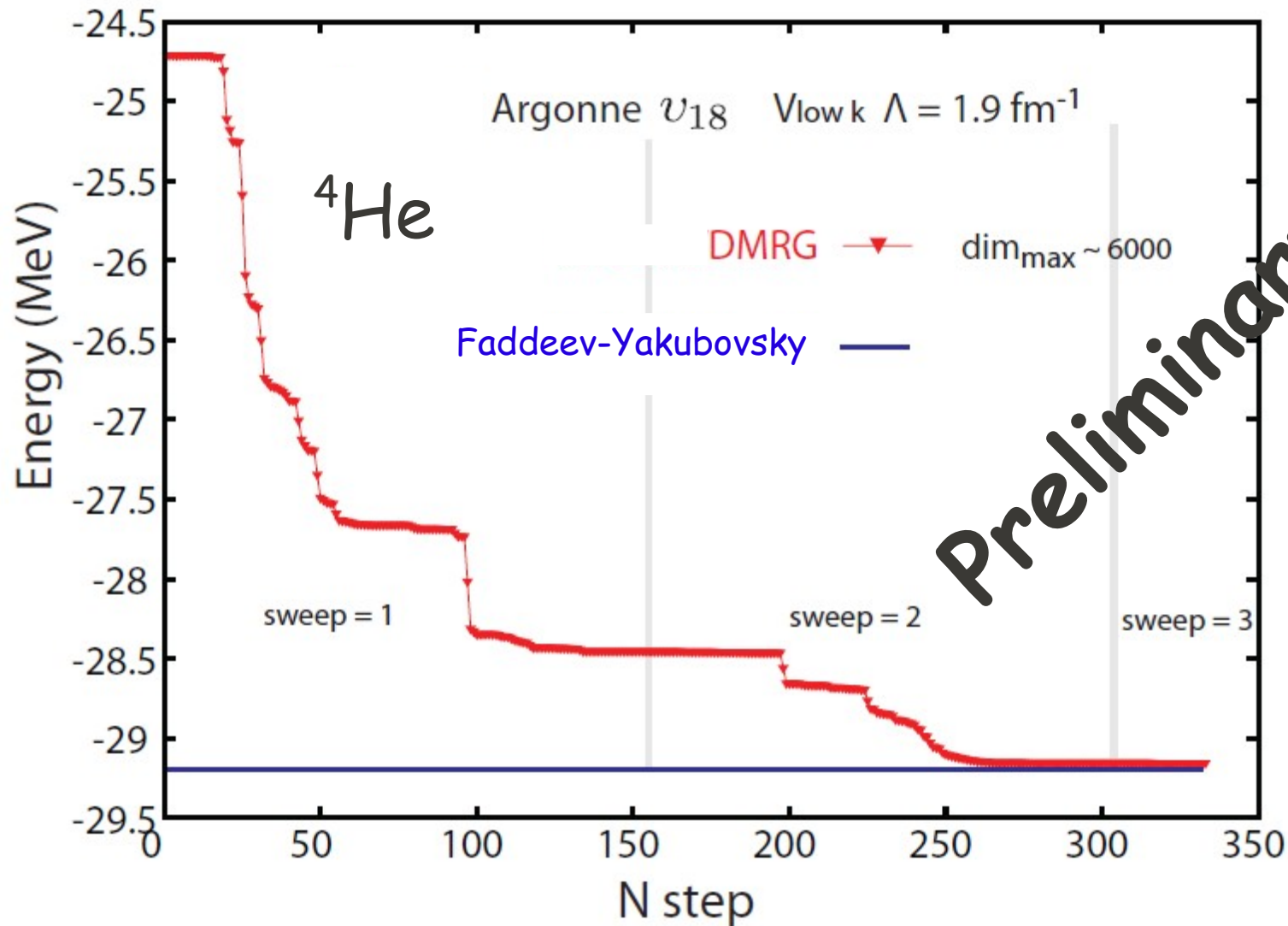
Preliminary



$E_{\text{Ab-Initio GSM}} = -8.390 \text{ MeV}$

$E_{\text{Faddeev}} = -8.470 (2) \text{ MeV}$

Faddeev result from Nogga et al, PRC 70 (2004) 061002, 2004



Preliminary

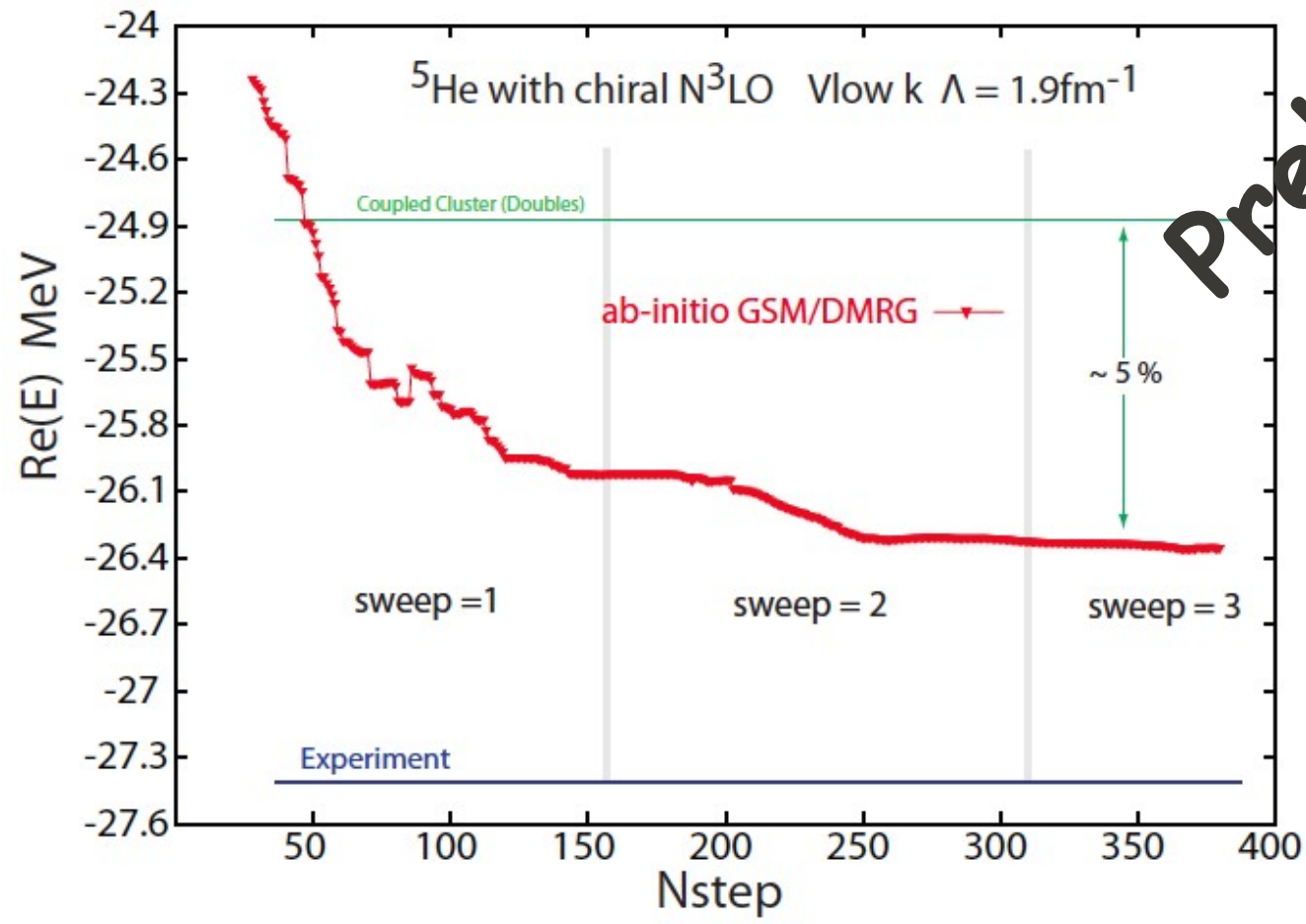
- Shells
- $0s_{1/2}(\text{p}) : E = -23.304 \text{ MeV}$
 - $0s_{1/2}(\text{n}) : E = -24.334 \text{ MeV}$
 - $s_{1/2}, p_{3/2}, p_{1/2}, s_{1/2}$ real energy continua
 - d, f, g H.O states

156 s.p. in total

GSM full dim = 6,230,512
 DMRG : dim ~ 6000

$E_{\text{ab-initio GSM}}$	$= -29.15 \text{ MeV}$
$E_{\text{Faddeev-Yakubovsky}}$	$= -29.19 (5) \text{ MeV}$

Preliminary



Shells

$0s_{1/2}(p) : E = -23.291 \text{ MeV}$

$0s_{1/2}(n) : E = -23.999 \text{ MeV}$

$0p_{3/2}(n) : E (1.194, -0.633)$

$p_{3/2}(n)$ complex contour
(discretized)

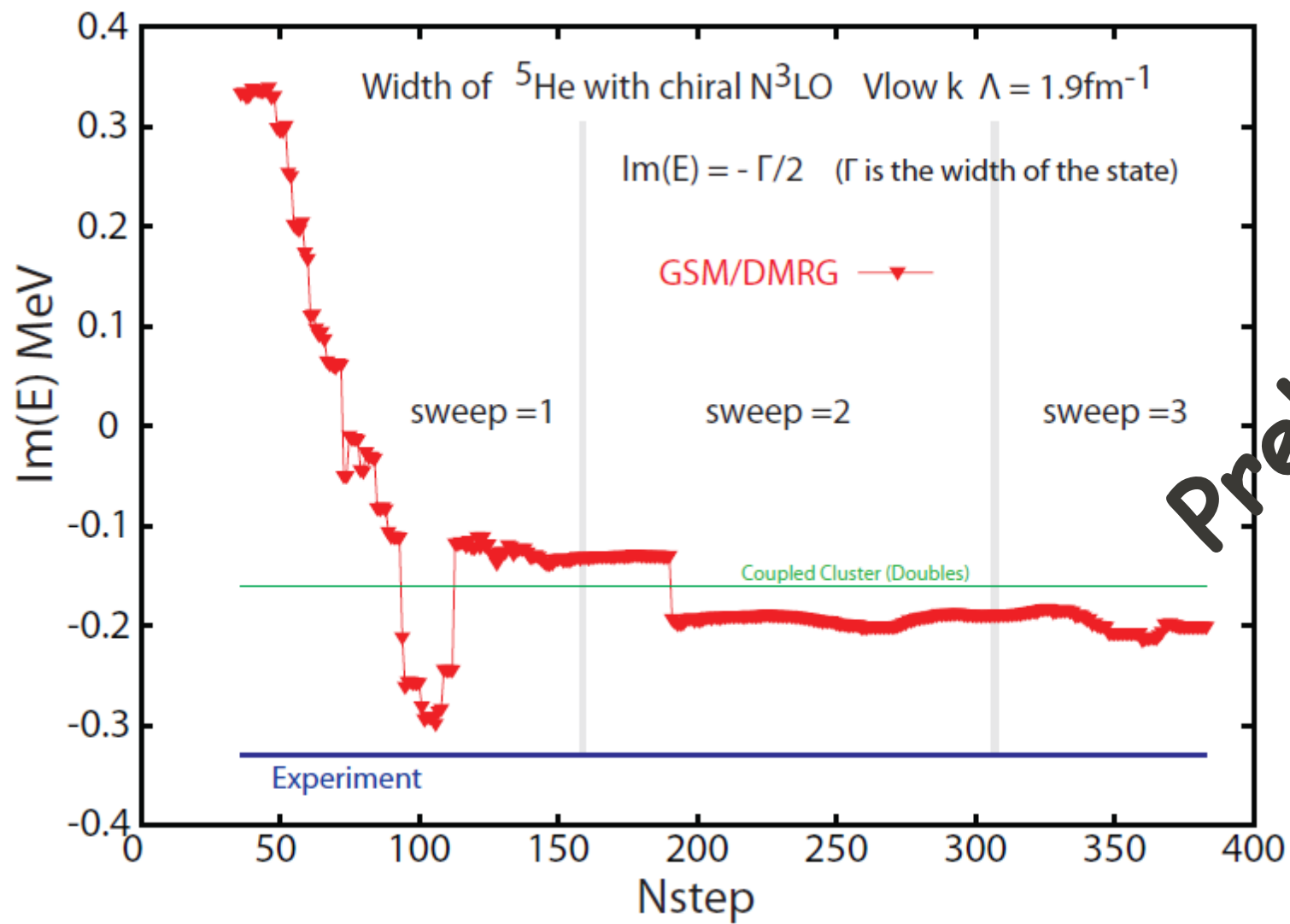
$s_{1/2}, p_{3/2}, p_{1/2}, s_{1/2}$ real energy
continua

d, f, g H.O states

157 s.p. in total

GSM full dim = 1,379,196,439

DMRG : dim $\sim 1.10^5$



Preliminary

Microscopic description of exotic nuclei in the Berggren basis

(Shell Model approach with coupling to the continuum)

i) Gamow Shell Model for helium isotopes, charge radius

ii) Ab-Initio approaches for (exotic) light nuclei with DMRG

Perspectives:

^{11}Li description as 7 nucleons above ^4He core, Oxygen isotopes with ^{22}O as a core, Ab-Initio description of Hydrogen chain.....



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W. Nazarewicz, University of Tennessee/ORNL
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Grand Accélérateur National d'Ions Lourds

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