

# Direct association of halo dimers and trimers in ultracold atoms.

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*Electroweak properties of light nuclei, INT Seattle 7/11/2012*

# People

Bar-Ilan University, Israel



L. Kh, Eli Shwartz, Noam Gross, Zav Shotan, Olga Machtey  
Not on the picture: Aviad Schori, Cornee Ravensbergen (short term visitor)

Eindhoven University of  
Technology, The Netherlands



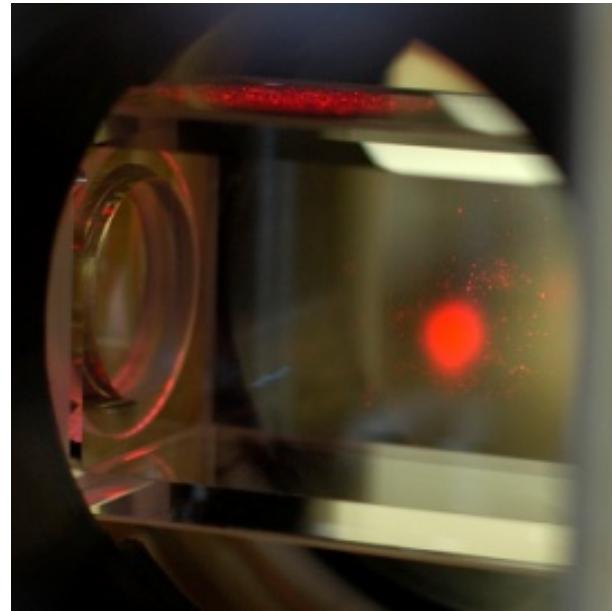
Servaas Kokkelmans

# Ultracold (Li) atoms

Magneto-optical trap

Close to the resonance (orbital electronic states) visible (laser) light – 671 nm ( $\sim 2$  eV)

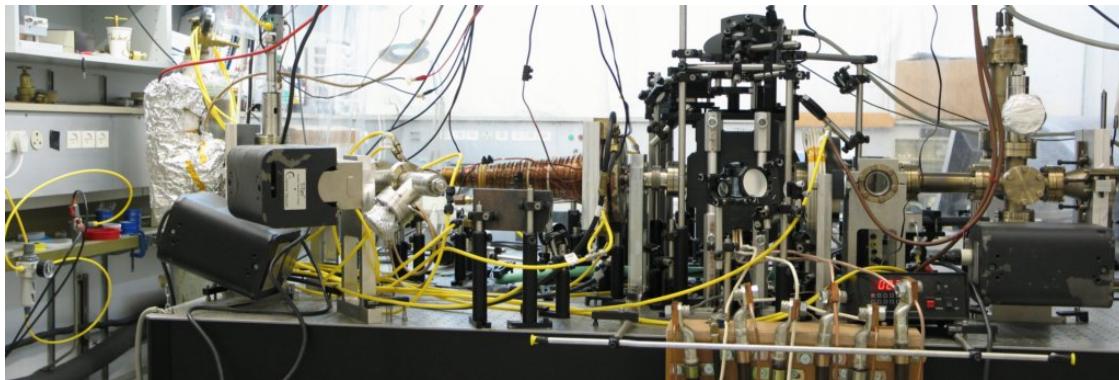
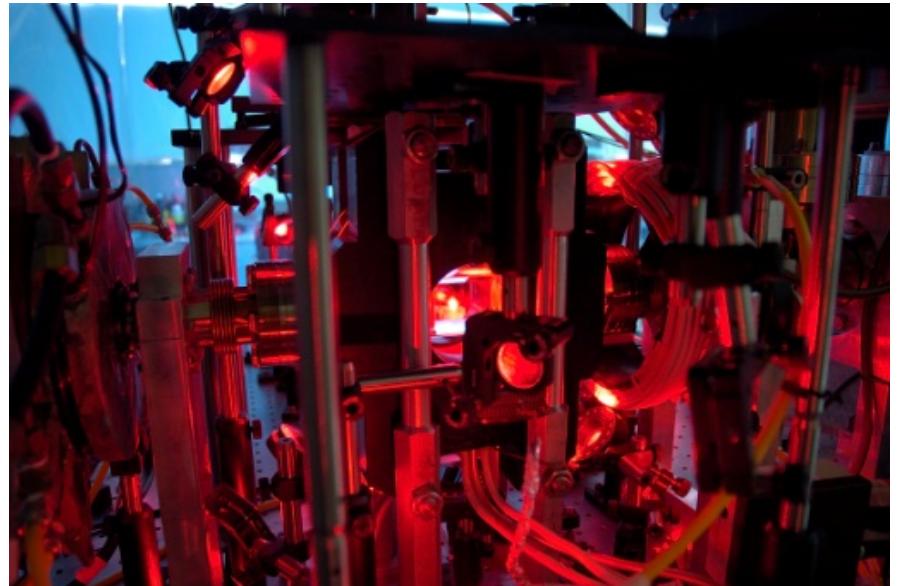
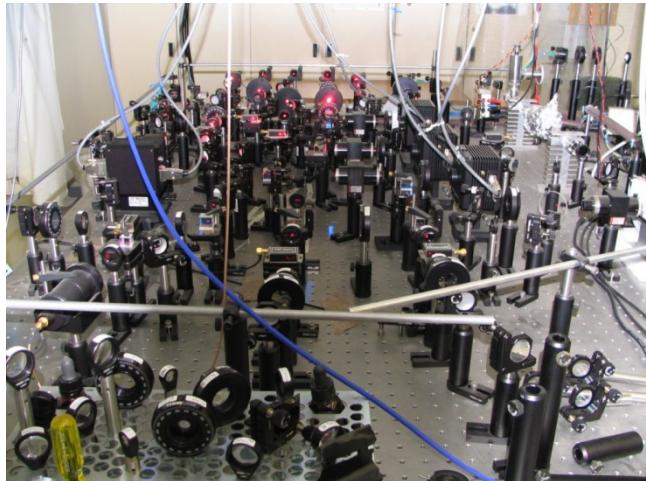
Ultrahigh vacuum  
environment



Magnetic fields

**Dissipative trap**  
 $N \sim 5 \times 10^8$  atoms  
 $n \sim 10^{10}$  atoms/cm<sup>3</sup>  
 $T \sim 300$  μK

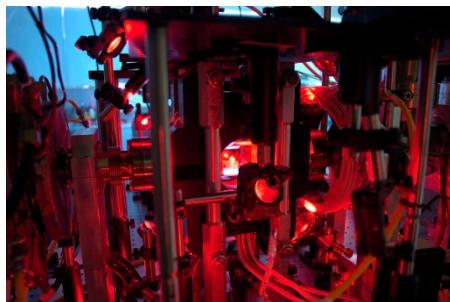
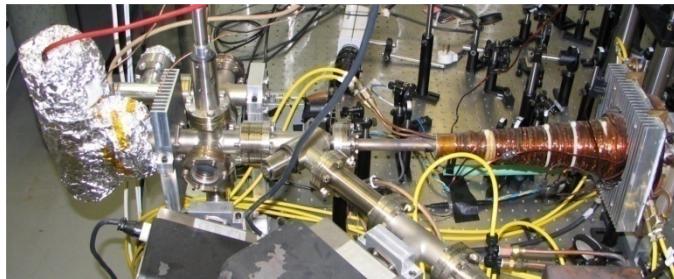
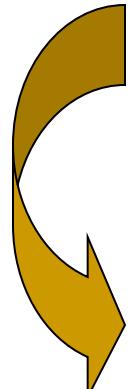
# Ultracold atoms – table-top experiment



# Ultracold (Li) atoms

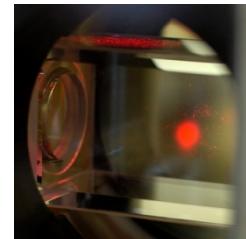
## Cooling:

Zeeman  
slower



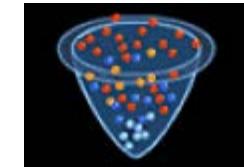
MOT  
 $\sim 10^9$  atoms

CMOT  
 $\sim 5 \times 10^8$  atoms  
( $300 \mu\text{K}$ )

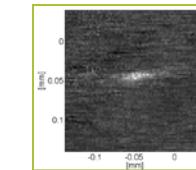


## Trapping: conservative atom trap

(our case: focus of powerful infrared laser)



## Crossed-beam optical trap



Evaporation:  
 $\sim 2 \times 10^4$  atoms  
 $\sim 1.5 \mu\text{K}$

## Typical numbers:

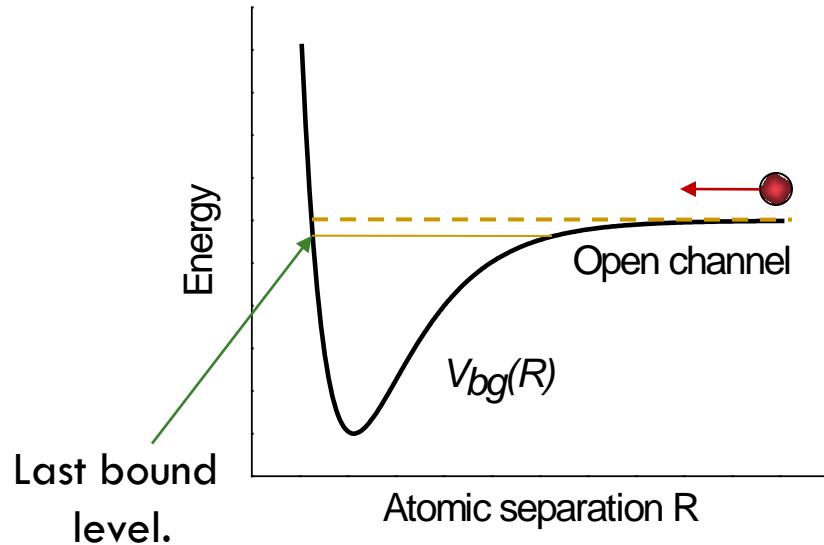
Temperature:  $\sim \mu\text{K}$

Relative velocities: few cm/sec

Collision energies: few peV

# Scattering length

At low temperatures the scattering is completely s-wave dominated.



s-wave scattering length  $a$   
is determined by the last bound state

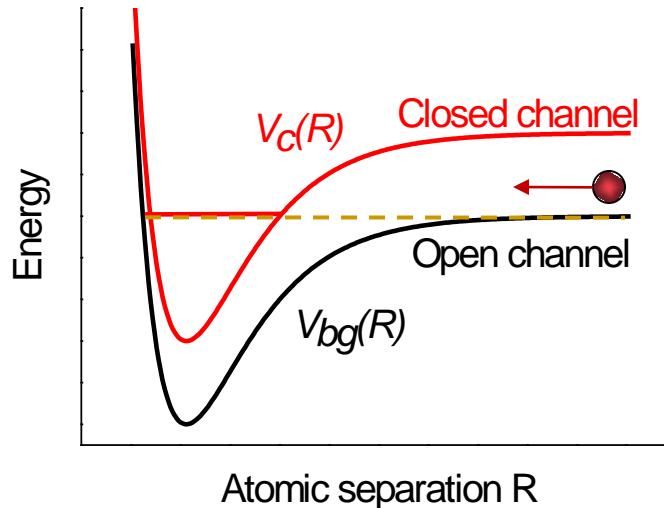
$$a = a_{bg}$$

$$\text{Li atoms: } a_{bg} = 10a_{Bohr}$$

Typical size of the interatomic potential – the van der Waals length       $l_{vdW} \sim 100a_{Bohr}$

# Feshbach resonance

Magnetic field tuning of the scattering length.

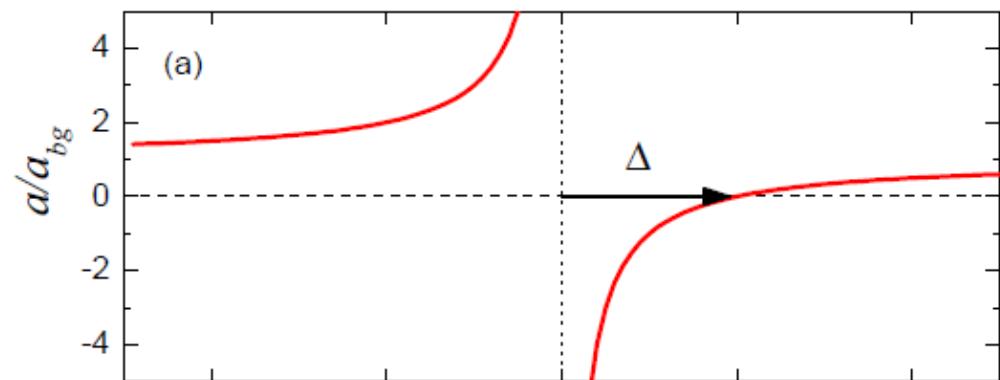


Closed channel: singlet potential  
bound state

Open channel: triplet potential  
free atoms

Different magnetic moments  
-The triplet potential depends  
on the magnetic field  
- The singlet potential does not

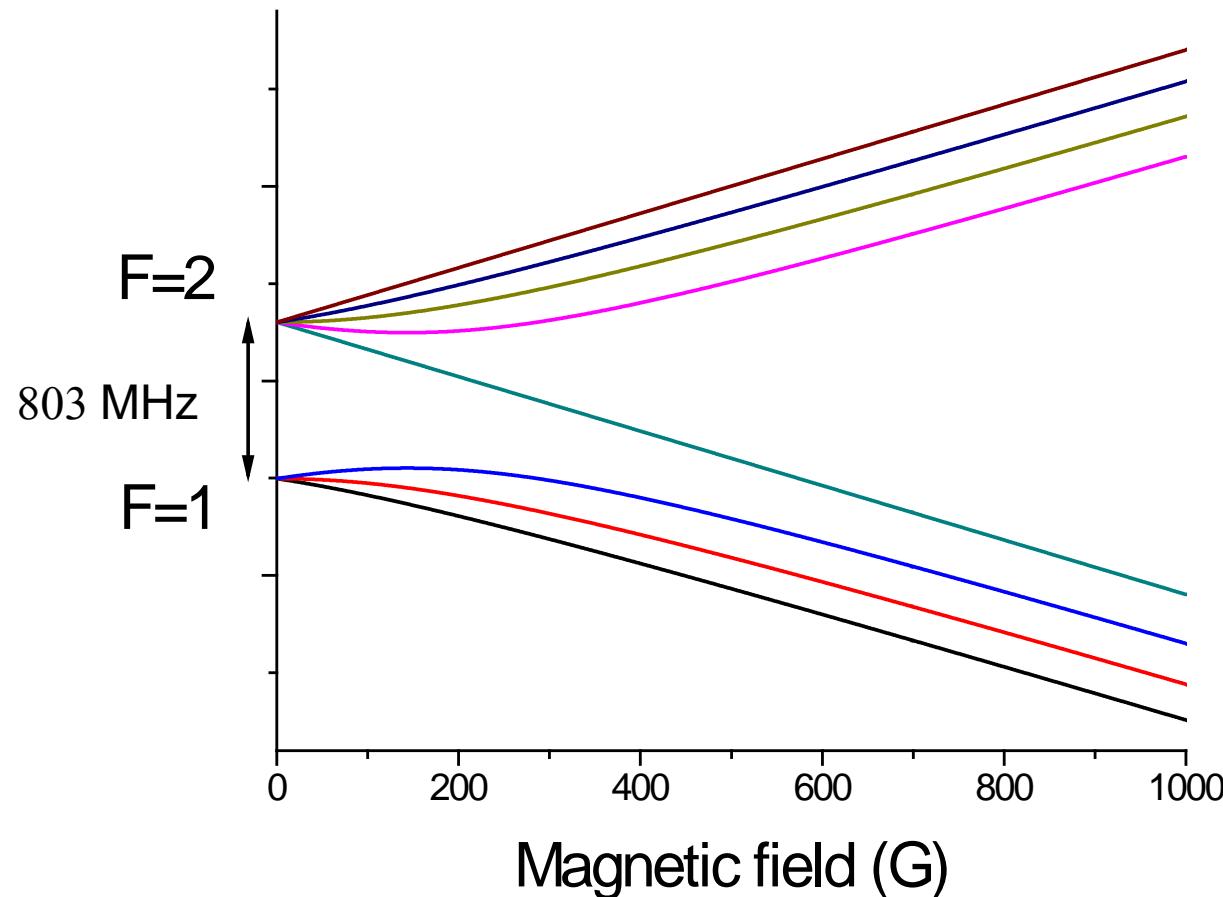
$$a = a_{bg} \left( 1 - \frac{\Delta}{B - B_0} \right)$$



# EXPERIMENTAL PLAYGROUND

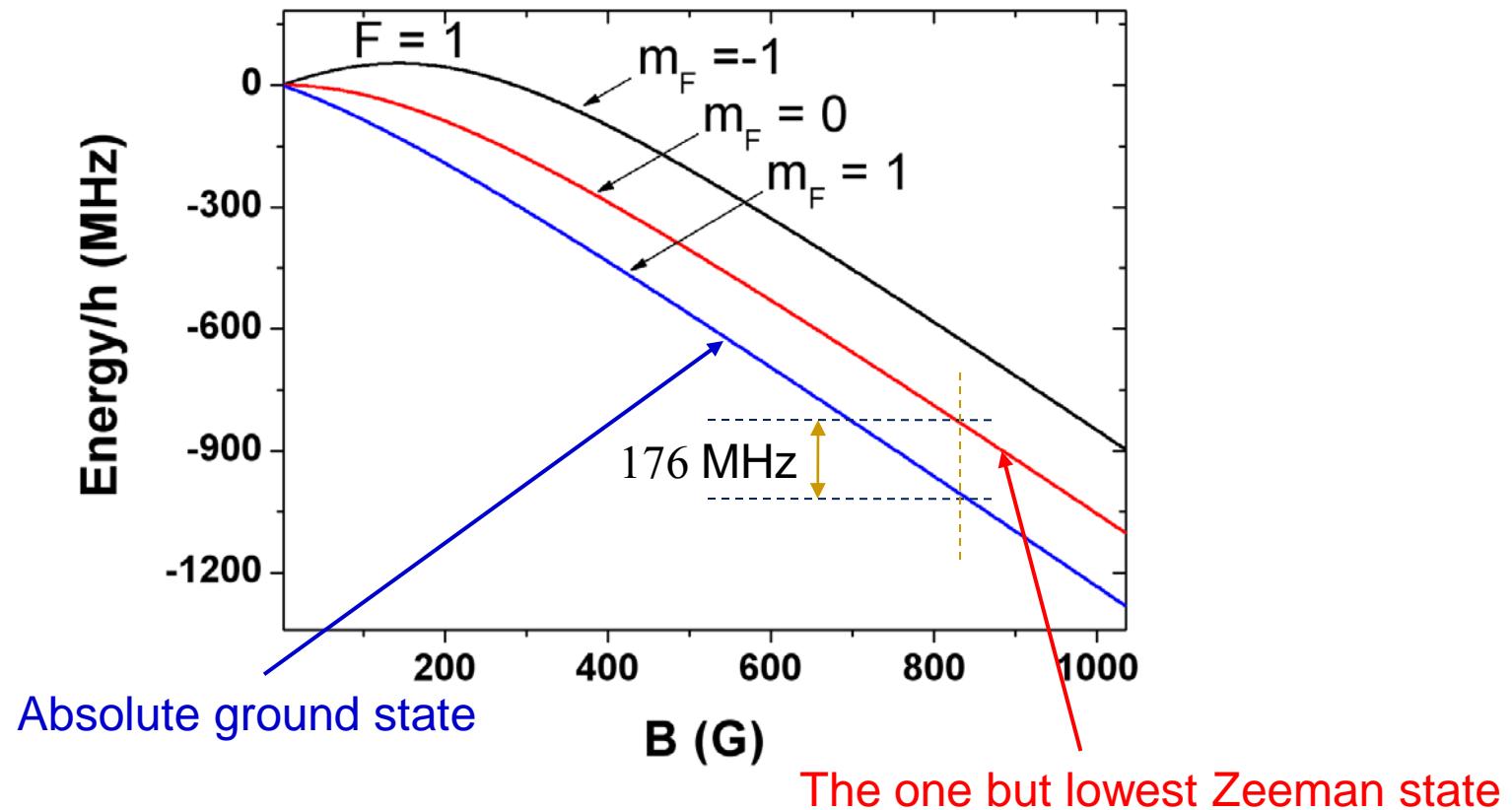
# Experimental playground - ${}^7\text{Li}$ atoms

Hyperfine structure of the ground state.



# Experimental playground - ${}^7\text{Li}$

3 identical bosons on a single nuclear-spin state.

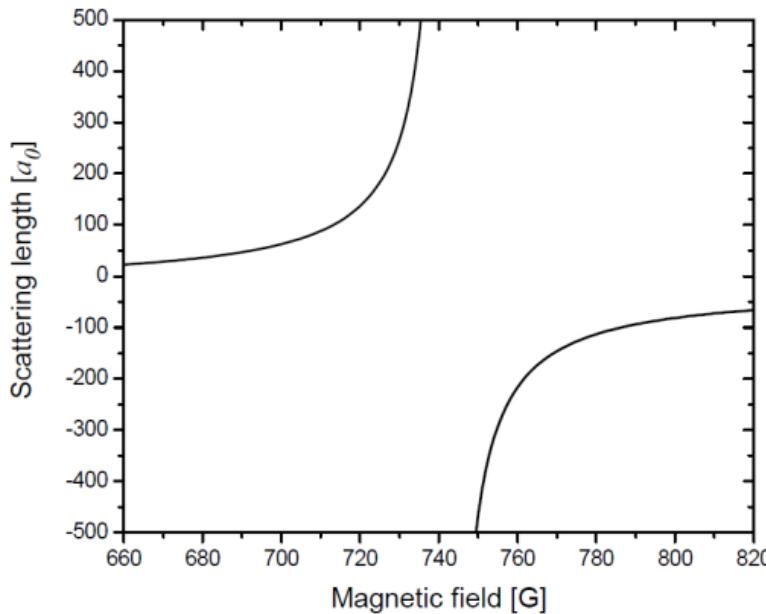


# Experimental playground - ${}^7\text{Li}$

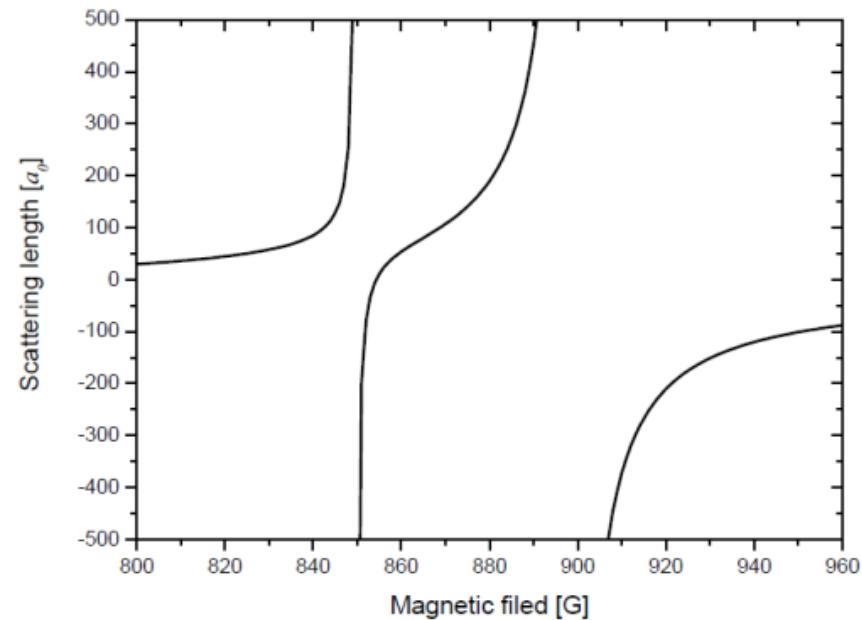
Absolute ground state

The one but lowest Zeeman state

Feshbach  
resonance

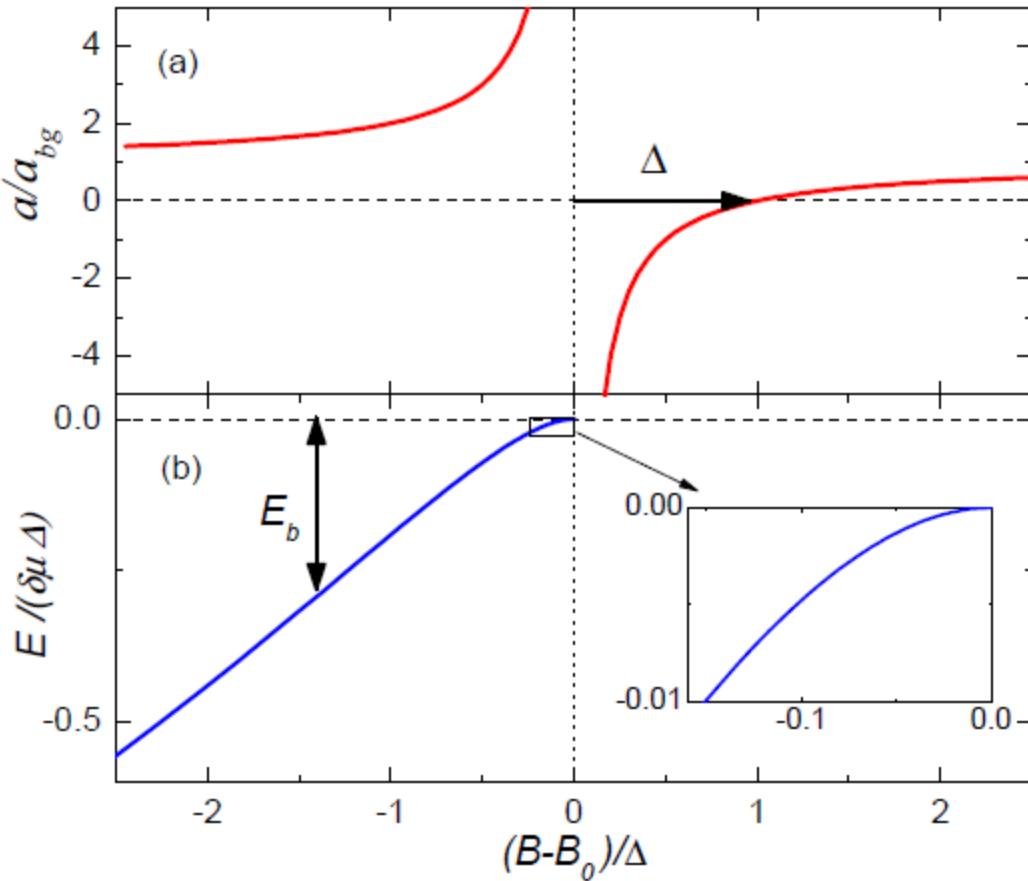


Feshbach  
resonance



# TWO-BODY UNIVERSALITY

# Feshbach molecule (universal dimer)



Feshbach molecule (universal dimer):

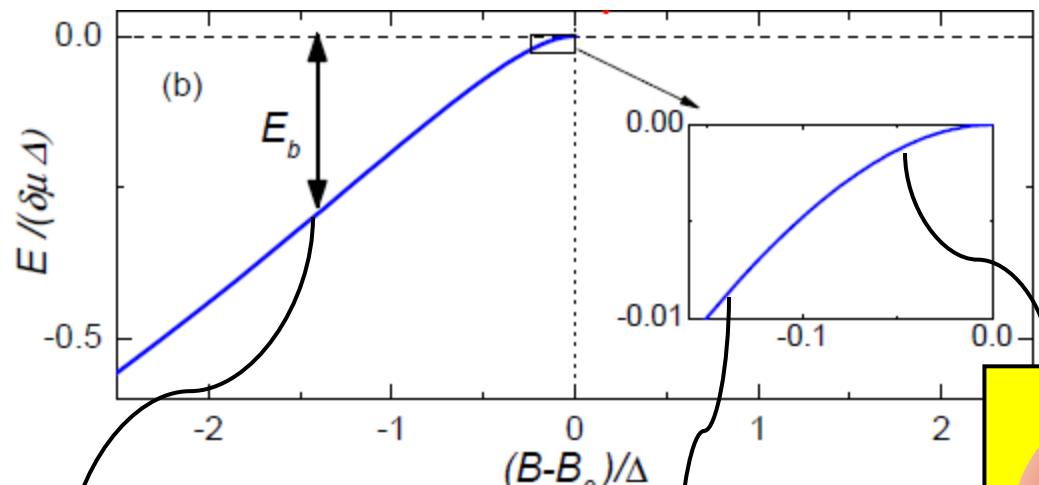
$$E_b = -\frac{\hbar^2}{ma^2}$$

Bare state (non-universal) dimer:

$$E_b = \delta\mu(B - B_0)$$

Also: deuteron, He<sub>2</sub>

# Universal dimer – quantum halo state



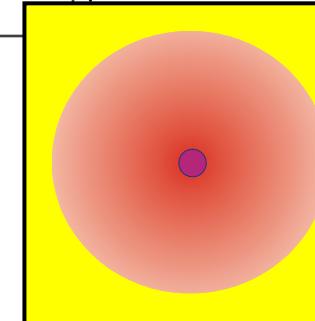
$$E_b = -\frac{\hbar^2}{ma^2}$$



The size of the bound state is that of a singlet potential:  $\sim 1.5$  nm



Progressive contamintion by the atomic continuum



A small fraction of the wave function is in the bound state. The size of the dimer is much larger than the van der Waals length.

“Quantum halo states”

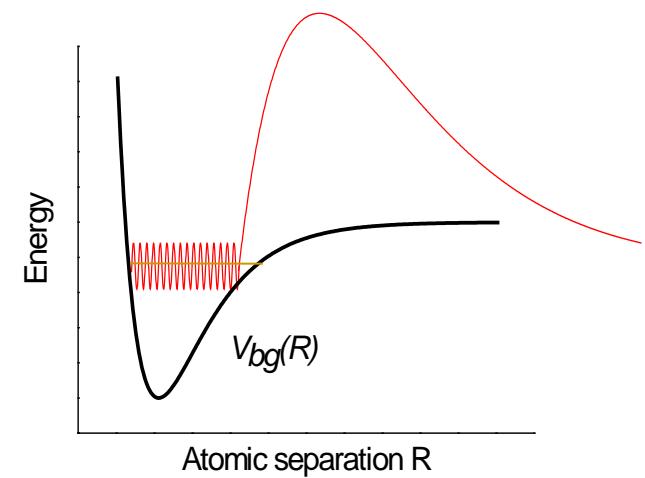
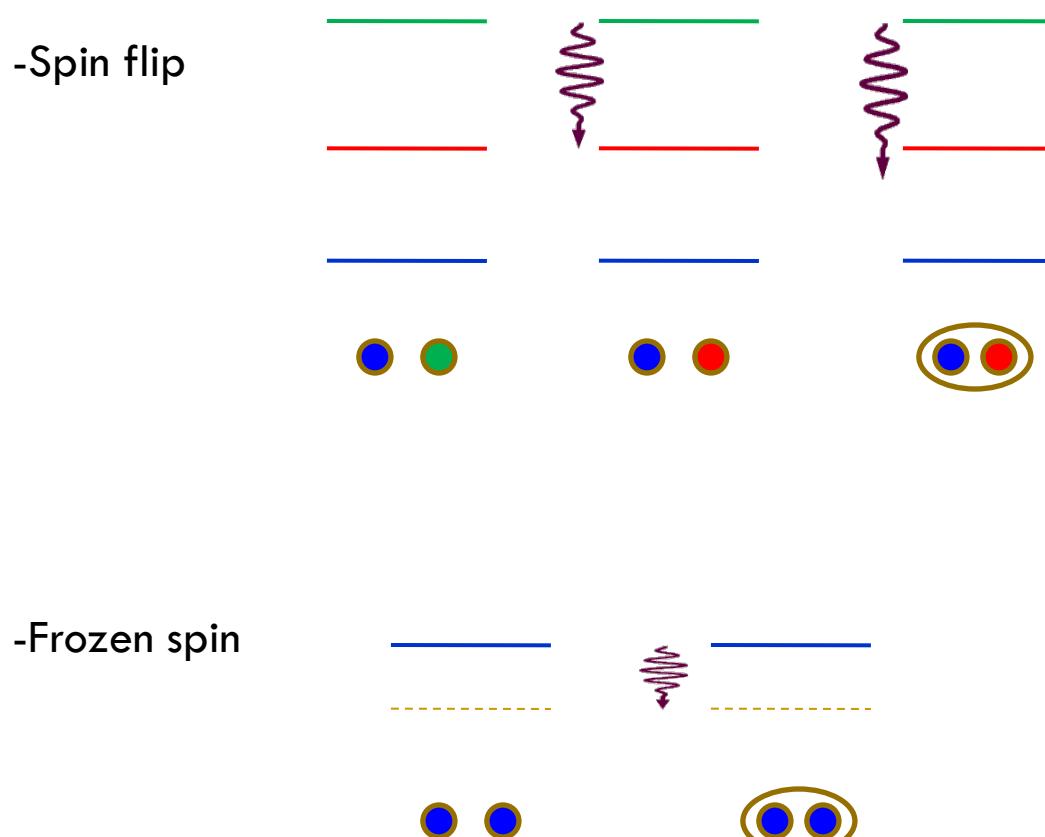
# ASSOCIATION OF HALO DIMERS

# Differenet systems

## ■ Rf association of universal dimers (partial list only) :

- 2005  $^{85}\text{Rb}$  JILA, Boulder, CO
- 2006  $^6\text{Li}$  in Innsbruck, Austria; MIT, Cambridge MA
- 2008  $^{41}\text{K}$  -  $^{87}\text{Rb}$  in Florence, Italy
- 2009  $^{40}\text{K}$  -  $^{87}\text{Rb}$  in Hannover, Germany
- 2010  $^7\text{Li}$  in BIU, Israel
- ...

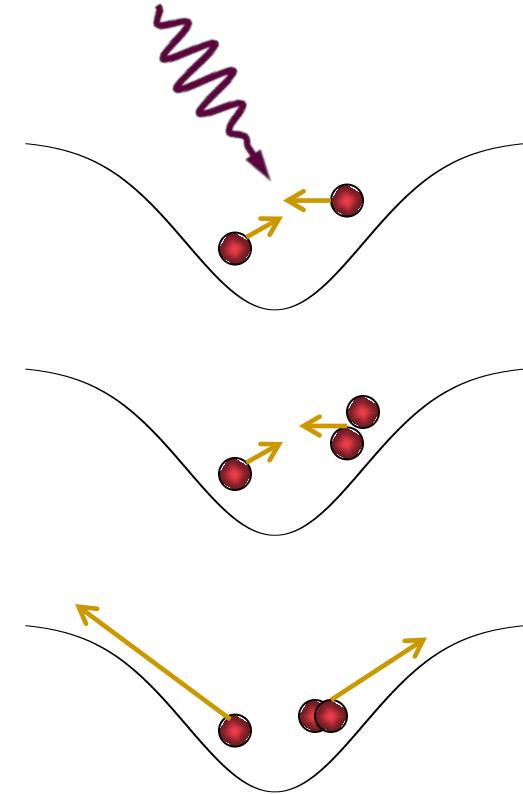
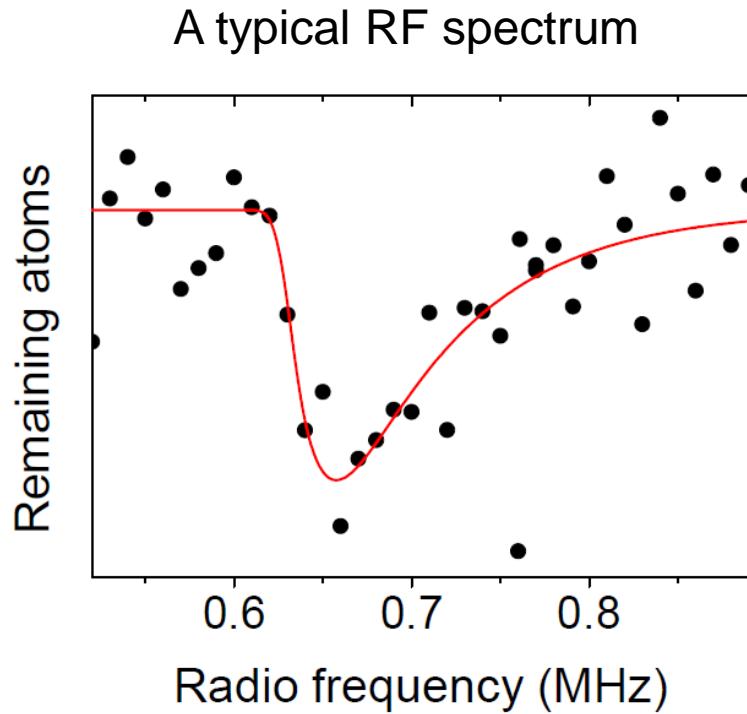
# Rf association of universal dimers



Strong overlapping  
with scattering continuum

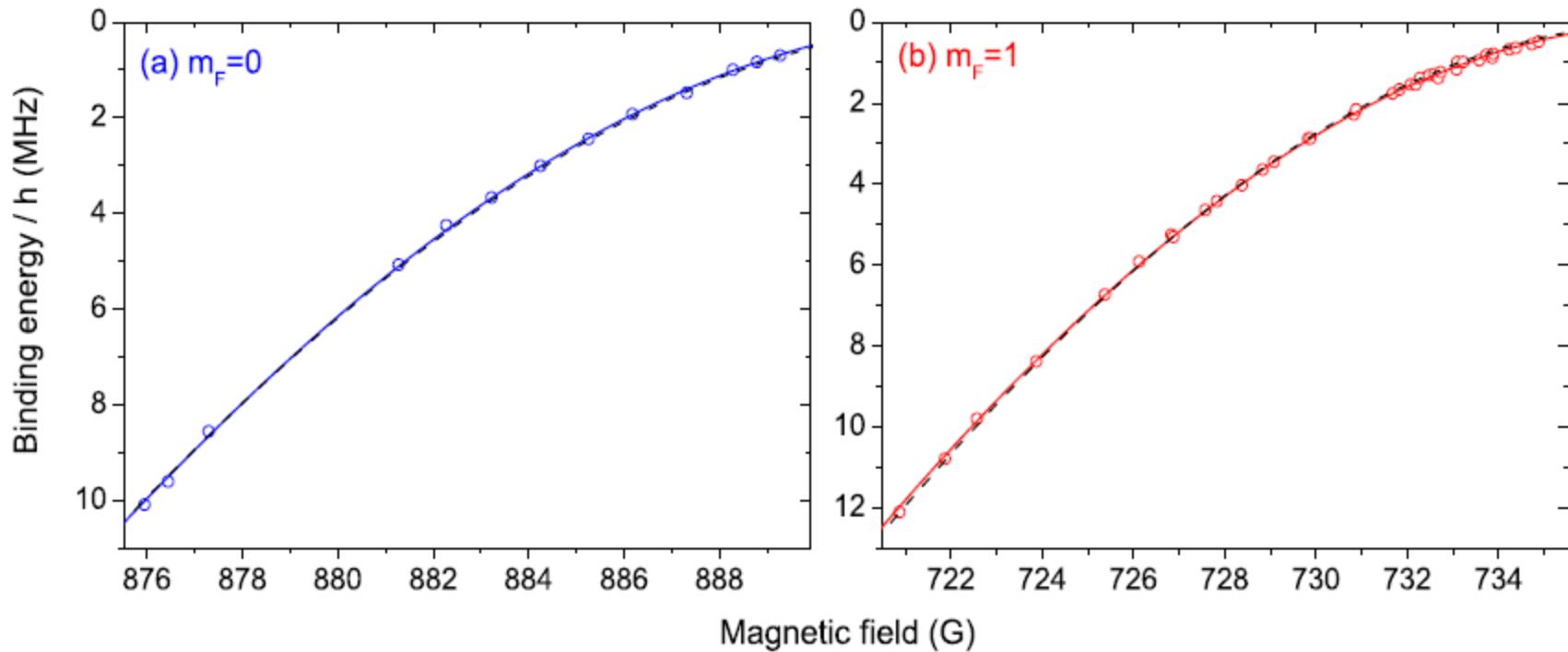
# Rf association of universal dimers

*Precise characterization of Feshbach resonances by rf-spectroscopy of universal dimers.*



# Rf association of universal dimers

*Precise characterization of Feshbach resonances by rf-spectroscopy of universal dimers.*



# Mapping between the scattering length and the applied magnetic field

*Precise characterization of Feshbach resonances by rf-spectroscopy of universal dimers.*

State	Type	$B_0$ (G)	
		Combined fit	Experimental
$ m_F = 0\rangle$	narrow	845.54	844.9(8)
$ m_F = 0\rangle$	wide	893.95(5)	893.7(4)
$ m_F = 1\rangle$	wide	737.88(2)	738.2(4)

*Improved characterization of Li inter-atomic potential.*

$$a_s = 34.33(2)a_0$$

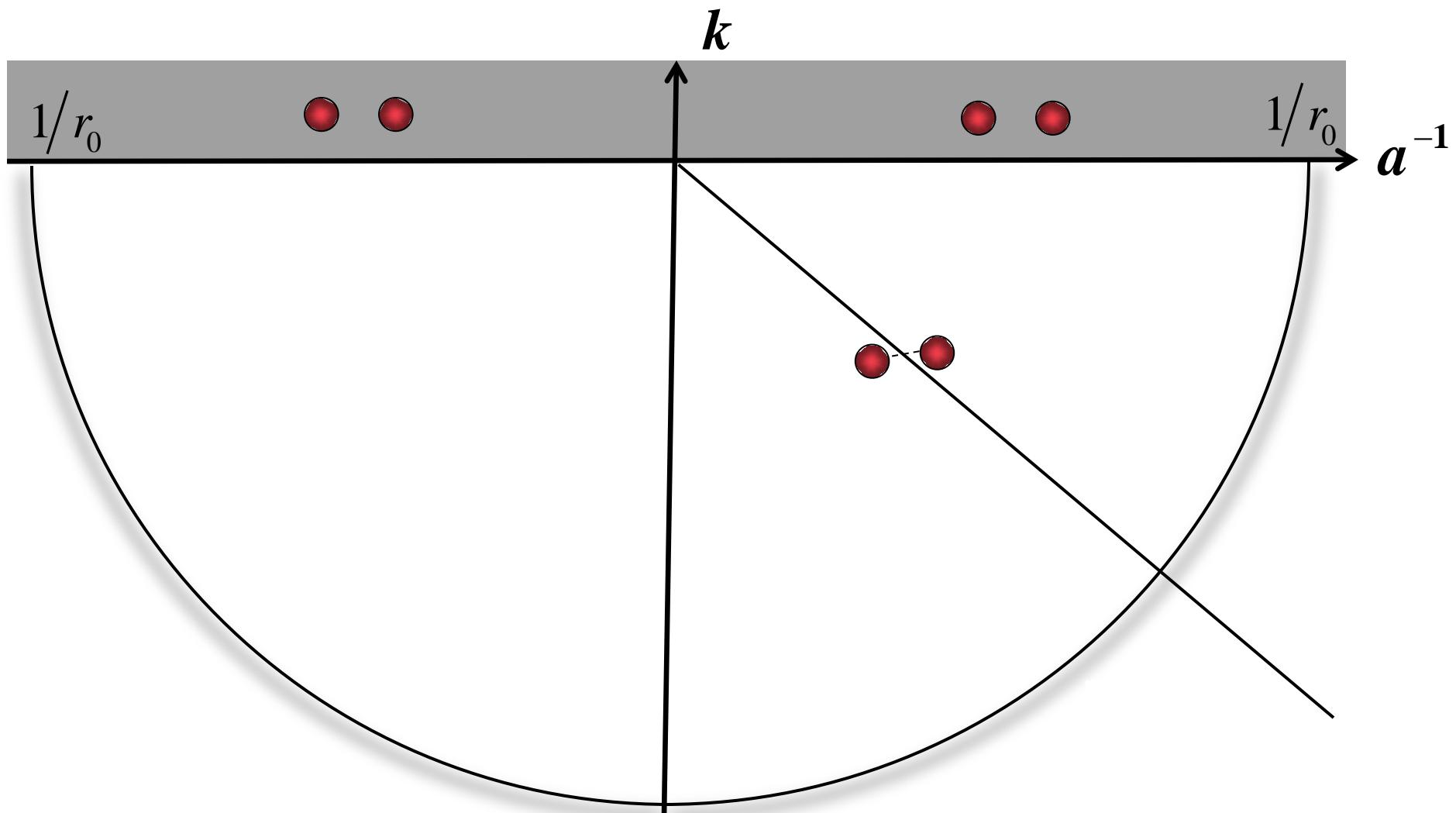
$$a_T = -26.87(8)a_0$$

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# THREE-BODY UNIVERSALITY: EFIMOV QUANTUM STATES

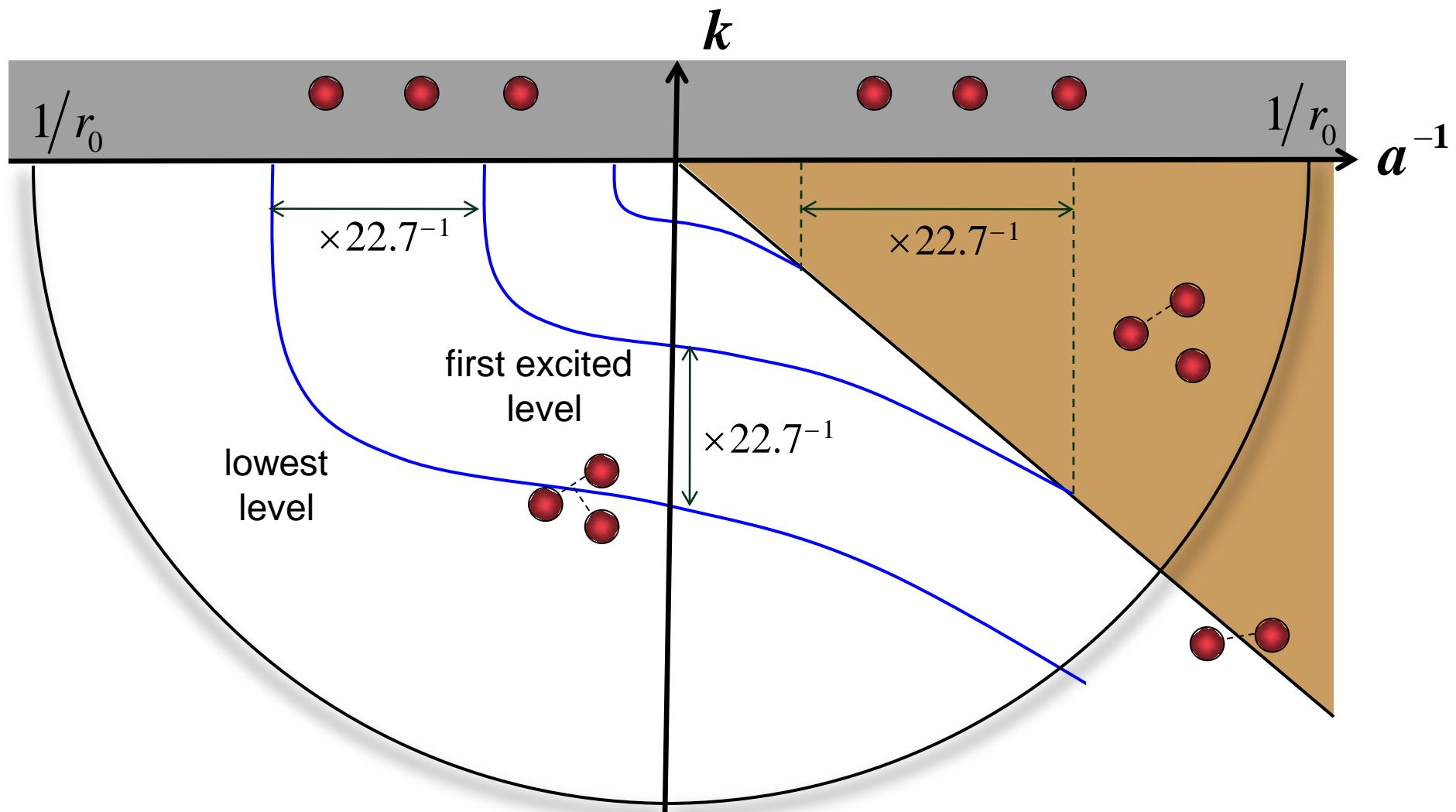
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# Quantum states near 2-body resonance



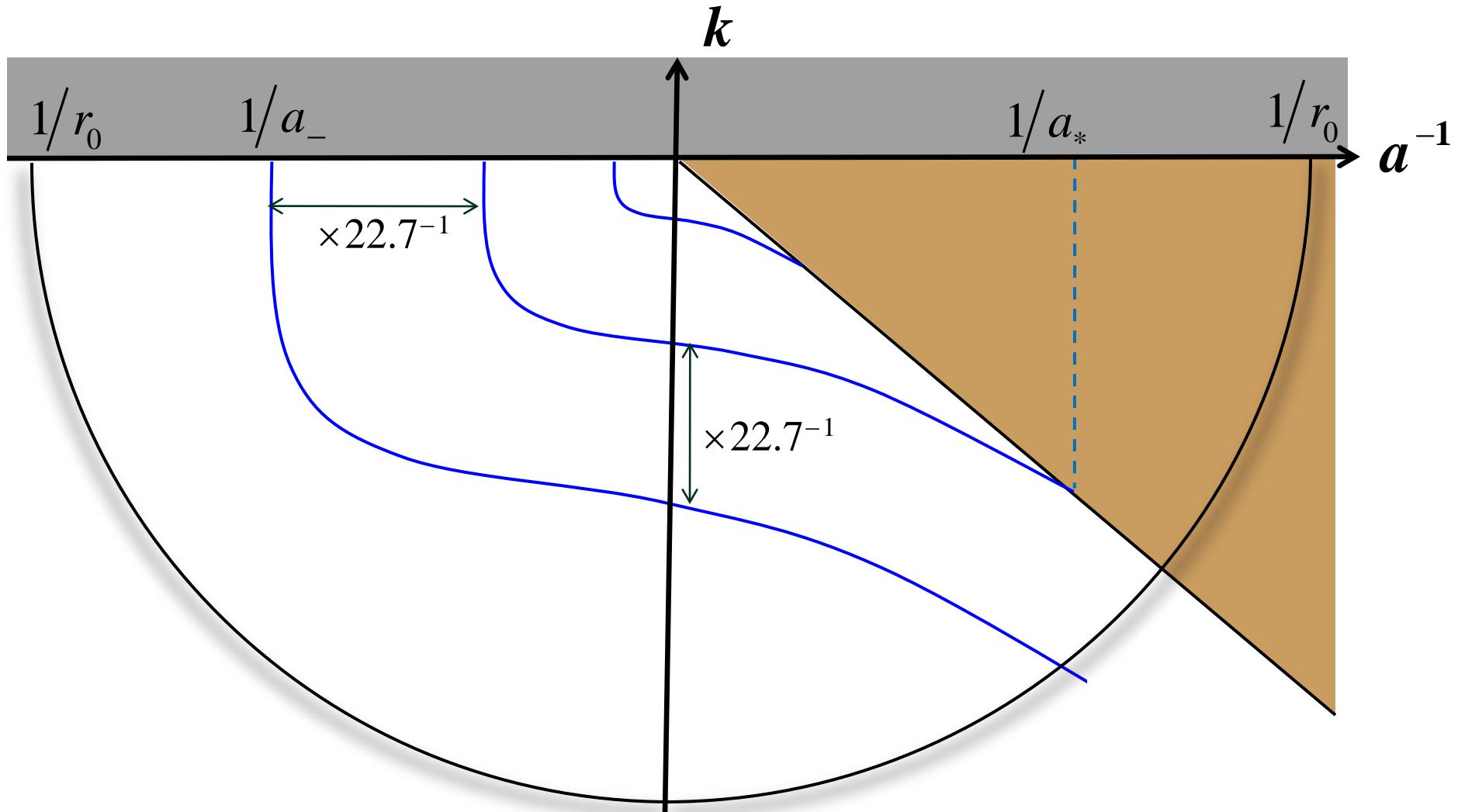
van der Waals range ( ${}^7\text{Li}$ ):  $r_0 = \left( \frac{m C_6}{16 \hbar^2} \right)^{1/4} \approx 32.5 a_0$

# Efimov scenario – universality window



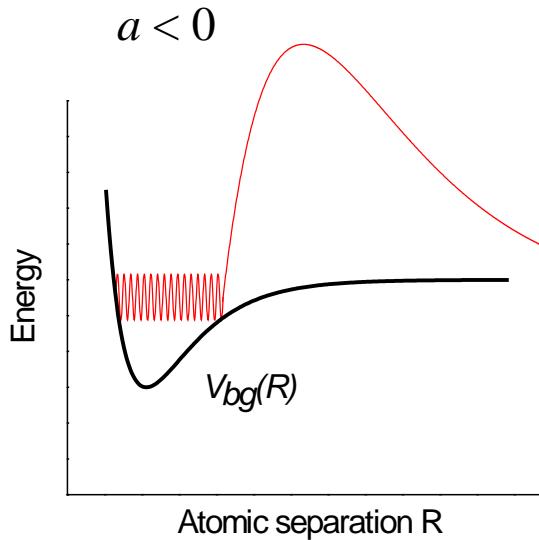
van der Waals range ( ${}^7\text{Li}$ ):  $r_0 = \left( \frac{m C_6}{16 \hbar^2} \right)^{1/4} \approx 32.5 a_0$

# Efimov scenario – 3-body parameter



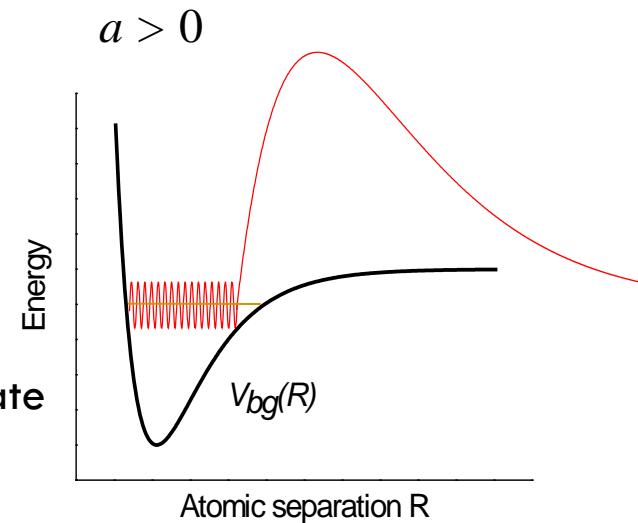
Position of a highly excited Efimov state is fixed by a 3-body parameter.

# Efimov scenario and real molecules

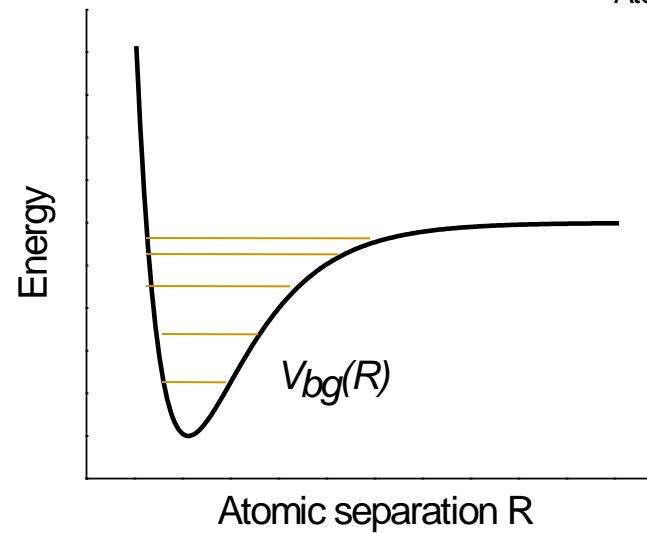


No 2-body bound states

One 2-body bound state

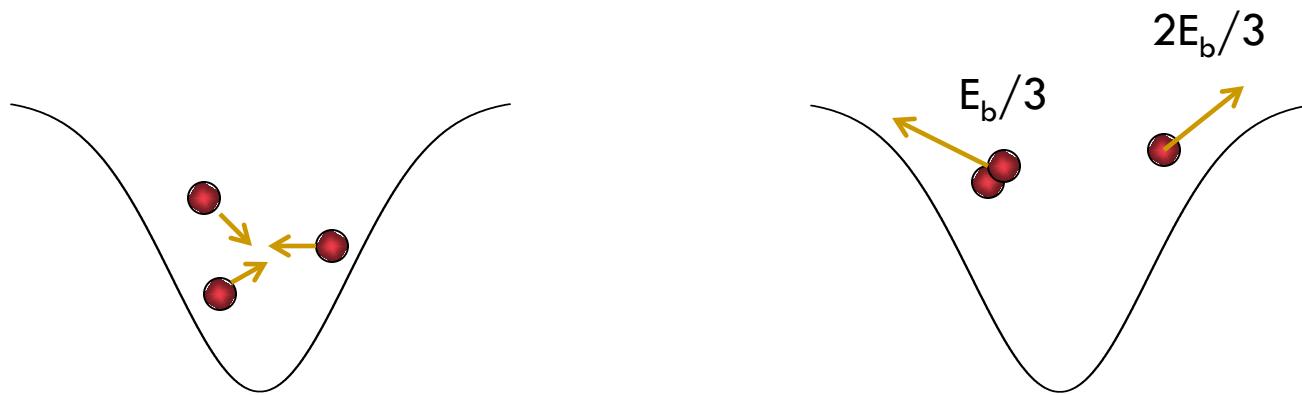


*Real molecules:*  
many deeply bound states



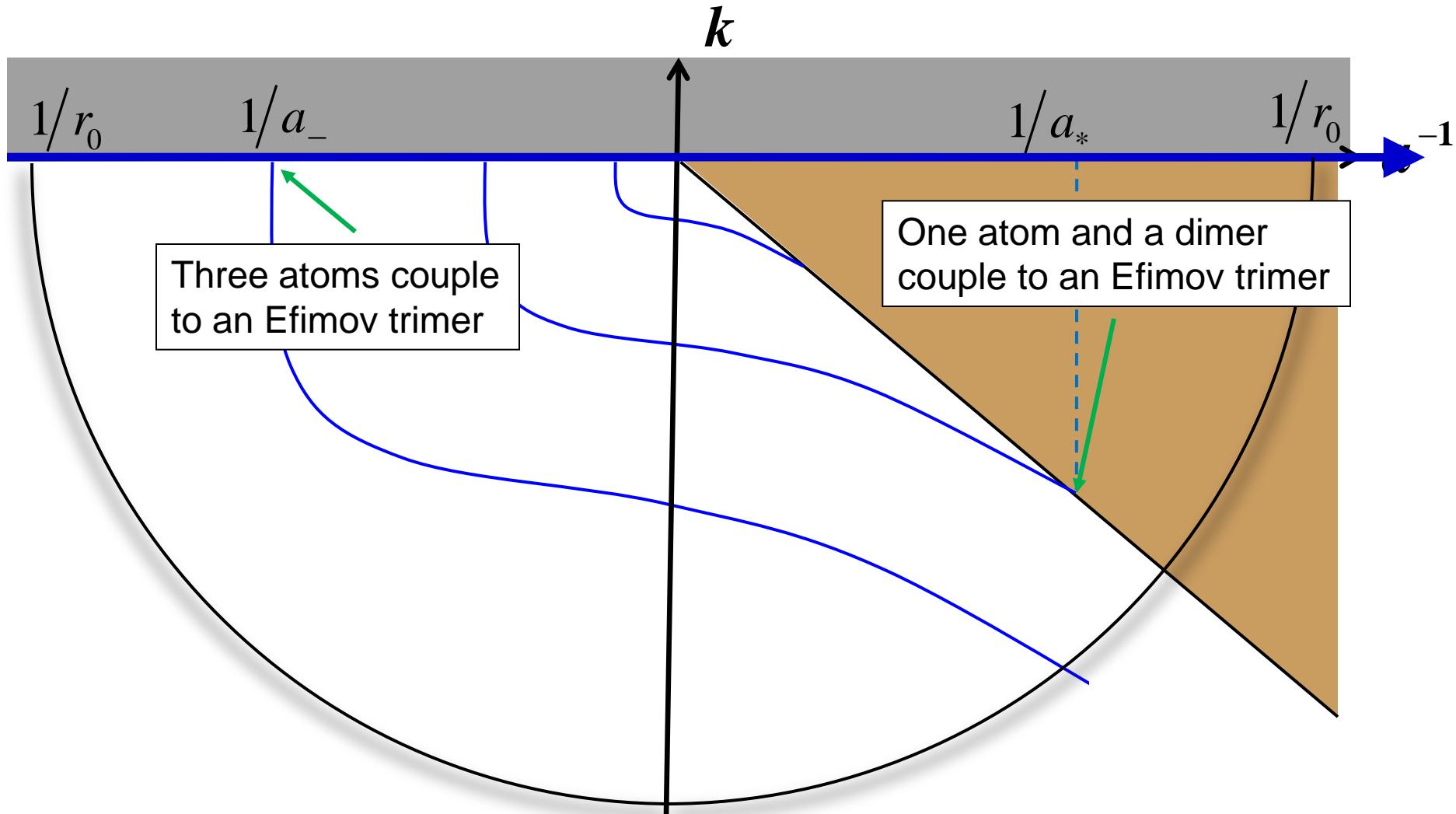
# Three-body recombination

Three body inelastic collisions result in a weakly (or deeply) bound molecule.



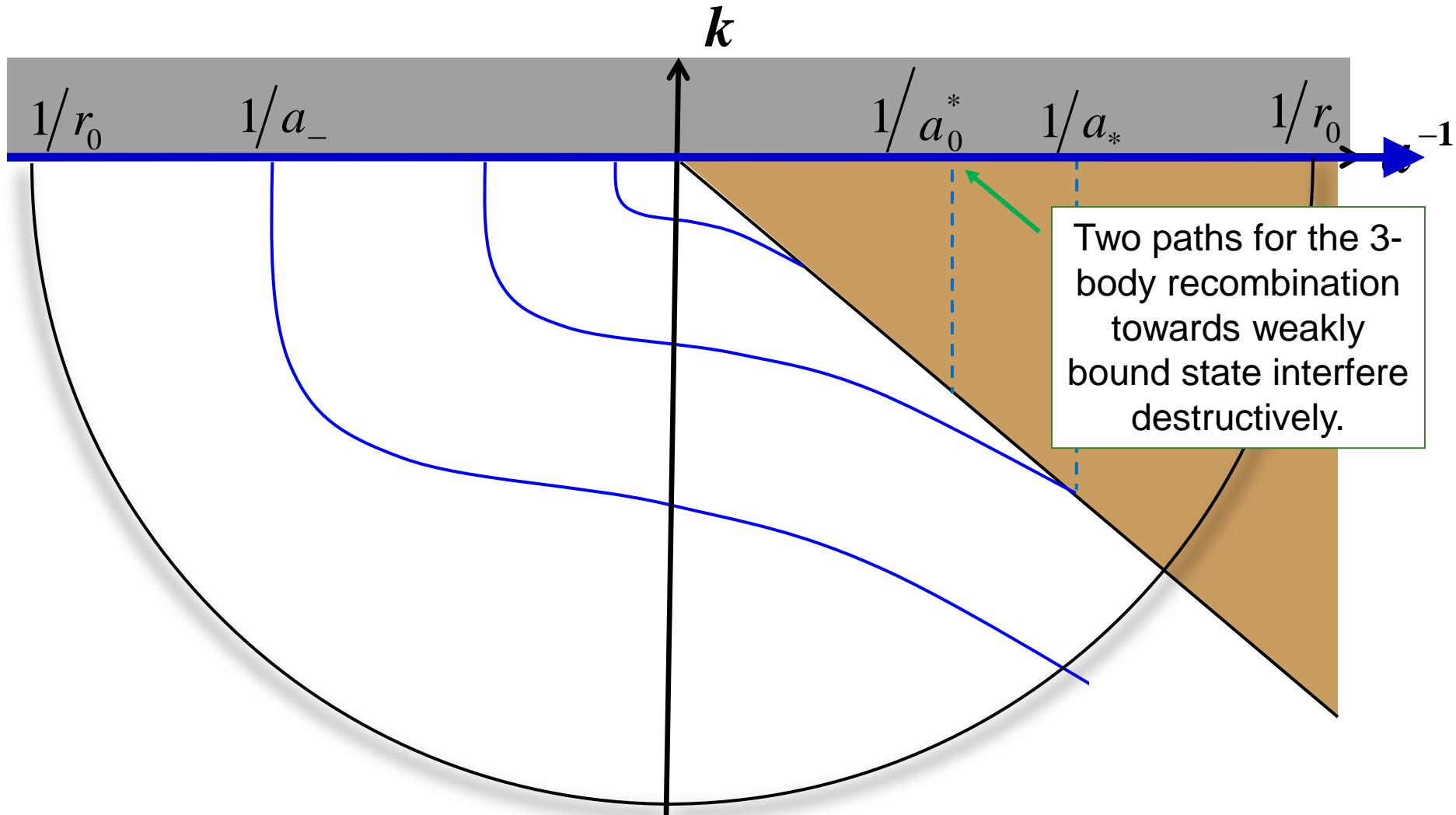
Release of binding energy causes loss which probes 3-body physics.

# Experimental observables



Experimental observable - **enhanced** three-body recombination.

# Experimental observables



Experimental observable – recombination **minimum**.

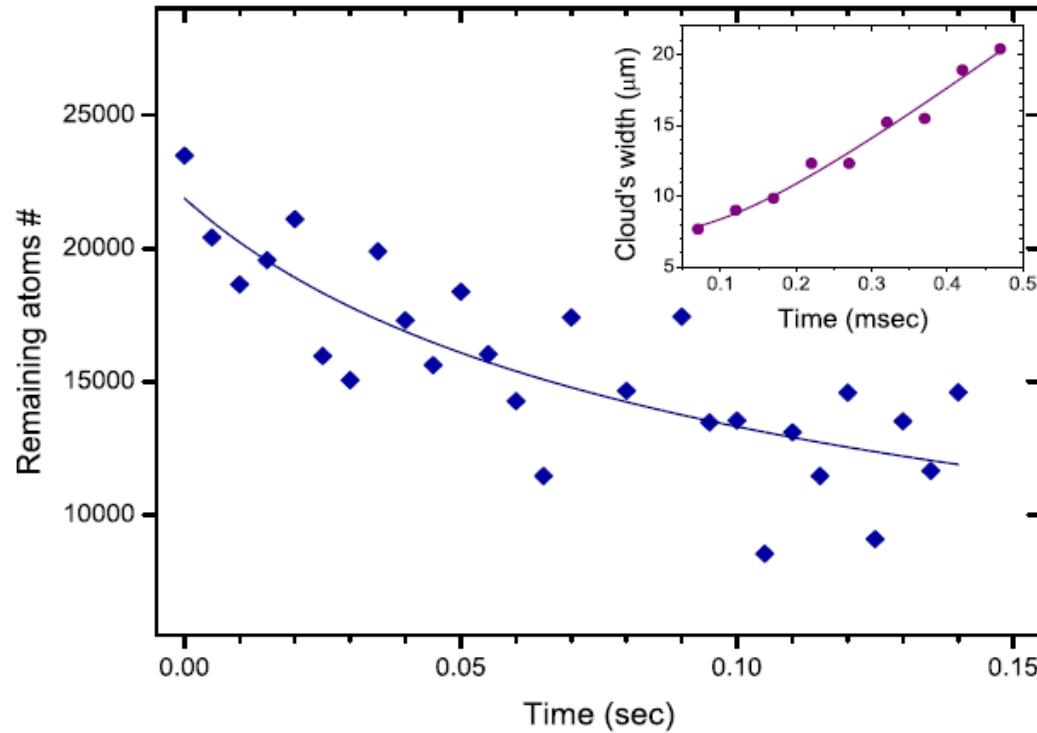
# EXPERIMENTAL RESULTS

# Different systems

- Efimov physics (and beyond) with ultracold atoms:
  - 2006 - ...  $^{133}\text{Cs}$  Innsbruck
  - 2008 – 2010  $^6\text{Li}$  3-component Fermi gas in Heidelberg, Penn State and Tokyo Univ.
  - 2009  $^{39}\text{K}$  in Florence, Italy
  - 2009  $^{41}\text{K}$  -  $^{87}\text{Rb}$  in Florence, Italy
  - 2009  $^7\text{Li}$  in Huston University, TX
  - 2009 - ...  $^7\text{Li}$  in BIU, Israel
  - 2012  $^{85}\text{Rb}$  JILA, Boulder, CO

# Experimental results

Typical set of measurements - atom number decay and temperature:



Loss rate from a trap:

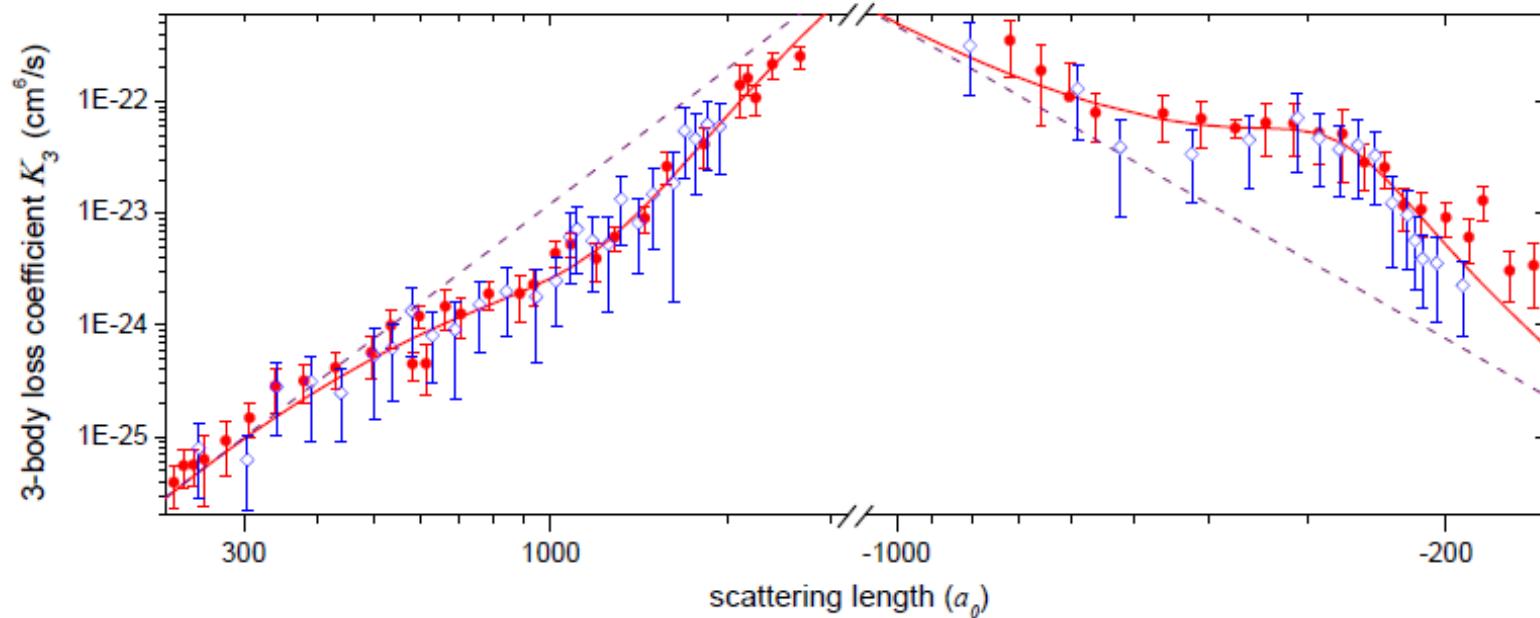
$$\dot{N} = -K_3 \langle n^2 \rangle N$$

$K_3$  – 3-body loss coefficient [cm<sup>6</sup>/sec]

# Experimental results

$a > 0$ : T = 2 – 3  $\mu\text{K}$

$a < 0$ : T = 1 – 2  $\mu\text{K}$



mf = 1; Feshbach resonance ~738G.  
mf = 0; Feshbach resonance ~894G.

# Summary of the results

## Fitting parameters to the universal theory:

State	$\eta_+$	$\eta_-$	$a_+/a_0$	$a_-/a_0$	$a_+/ a_- $
$ m_F = 0\rangle$	0.213(79)	0.180(48)	238(25)	-280(12)	0.85(11)
$ m_F = 1\rangle$	0.170(41)	0.253(62)	265(16)	-274(12)	0.97(8)

$$|a_-|/r_0 = 8.6(4) \quad \text{UT prediction:}$$

$$a_+/|a_-| = 0.96(3)$$

## Derived parameters:

### Position of recombination minima:

$$\begin{aligned} |m_F = 0\rangle & \quad a_0^* = 1134(120)a_0 \\ |m_F = 1\rangle & \quad a_0^* = 1262(76)a_0 \end{aligned}$$

### Position of atom-dimer threshold:

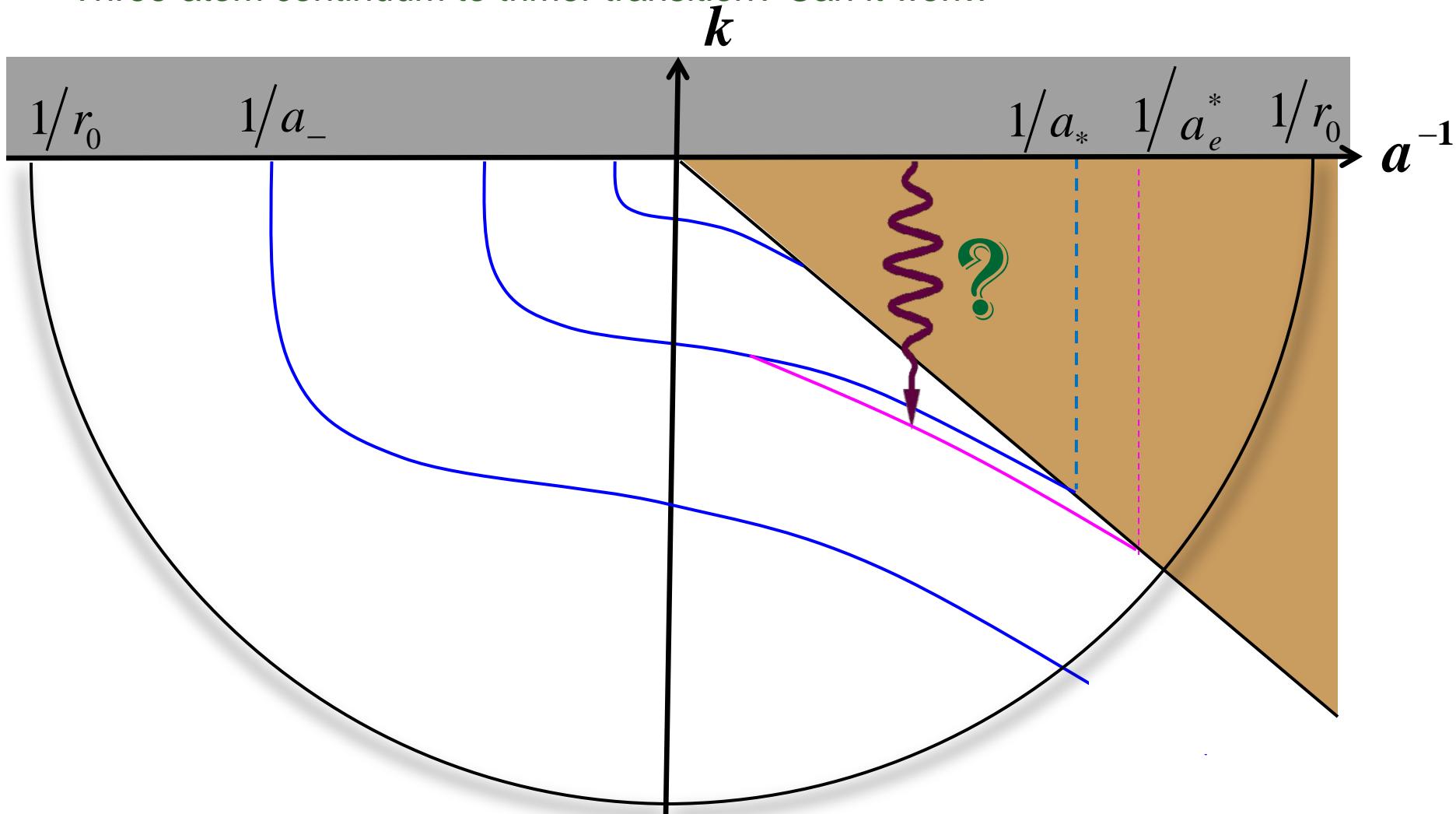
$$\begin{array}{ll} a_+ & a_- \\ a_* = 262(27)a_0 & a_* = 288(26)a_0 \\ a_* = 292(18)a_0 & a_* = 282(16)a_0 \end{array}$$

- 3-body parameter is the same across the region of  $|a| \rightarrow \pm\infty$
- 3-body parameter is the same for both nuclear-spin sublevels.

**ASSOCIATION OF HALO  
TRIMERS.**

# Rf association of Efimov trimers

Three-atom continuum to trimer transition? Can it work?

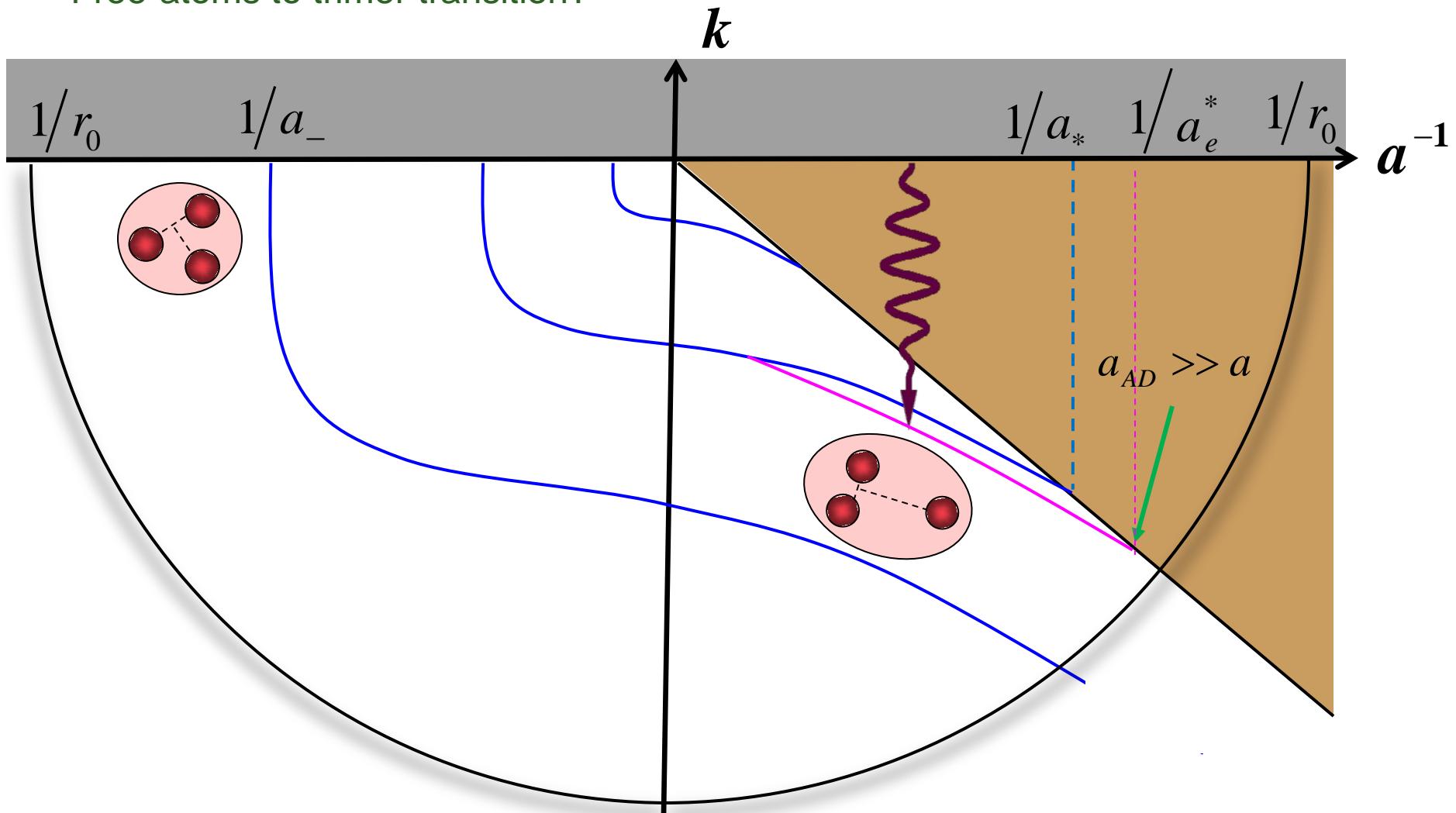


See also rf association of Efimov trimers in three-component Fermi gas:

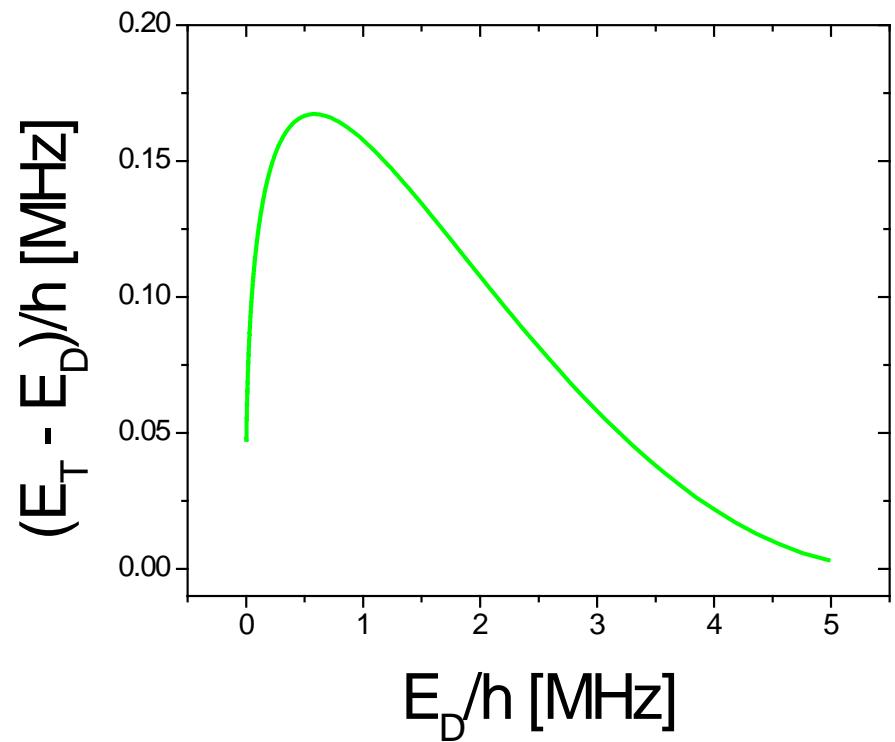
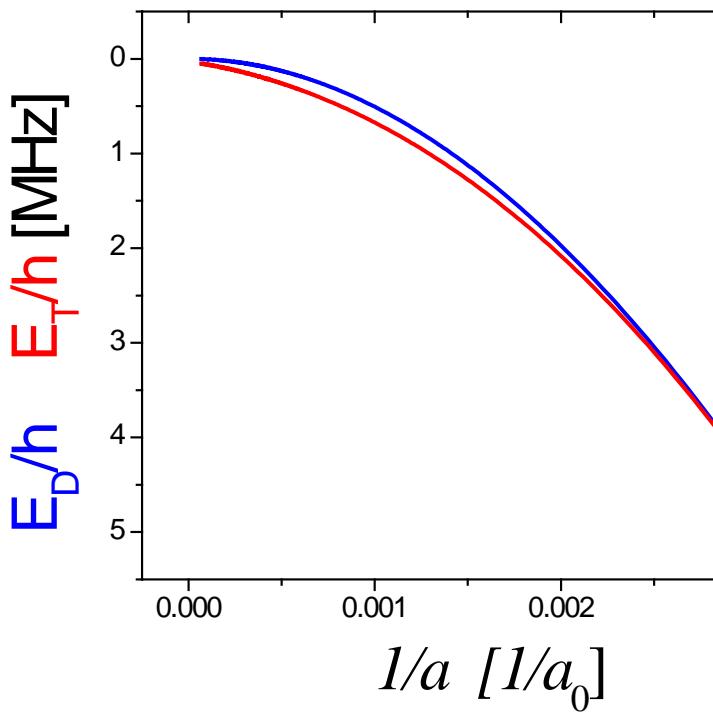
T. Lompe, T.B. Ottenstein, F. Serwane, A.N. Wenz, G. Zurn, S. Jochim , Science **330**, 940 (2010).

# Rf association of Efimov trimers

Free-atoms to trimer transition?

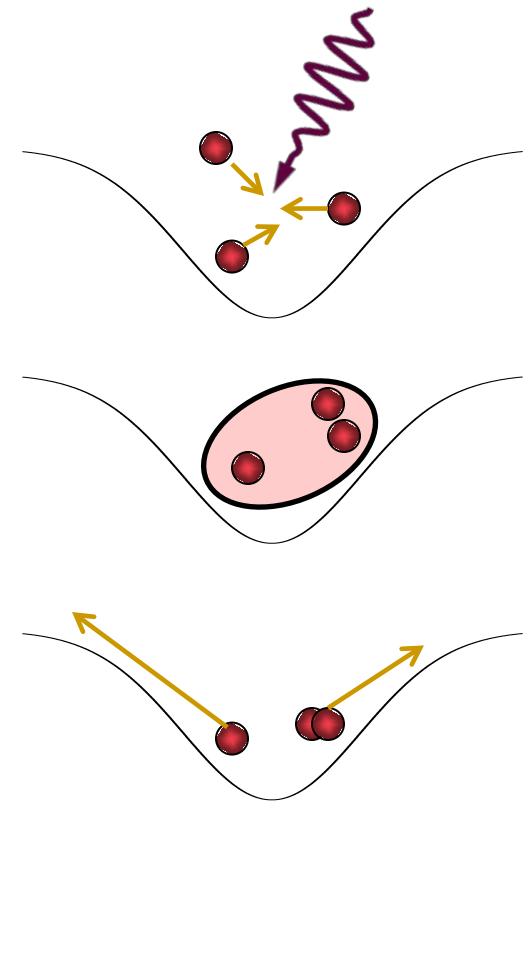
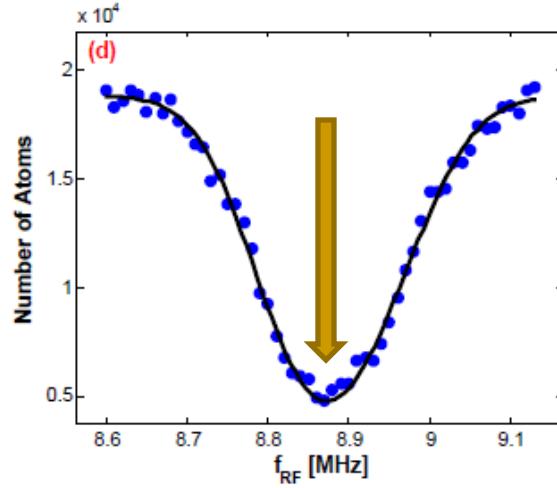
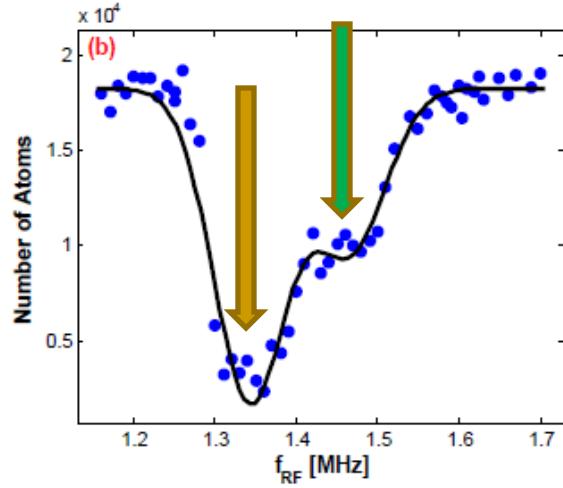
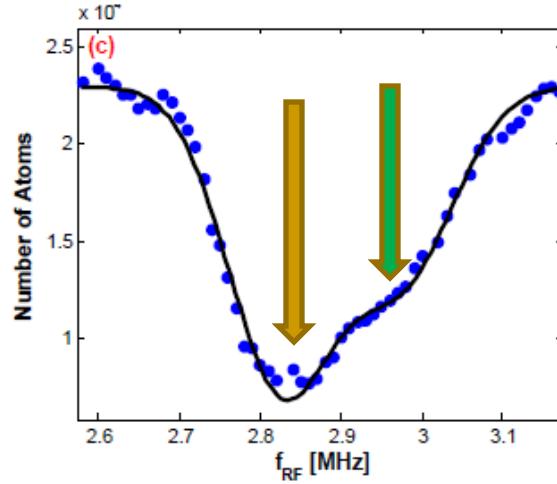
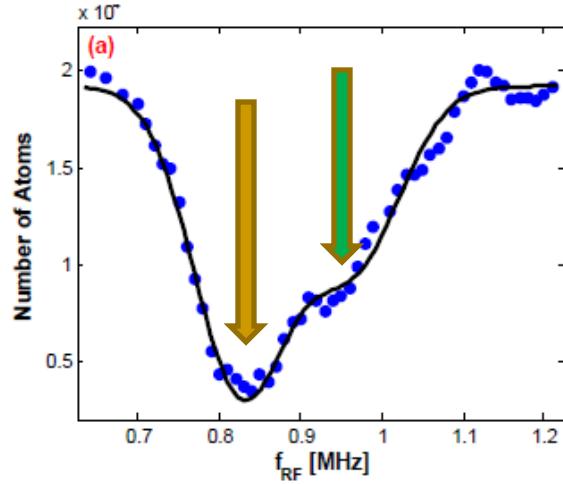


# Energy levels

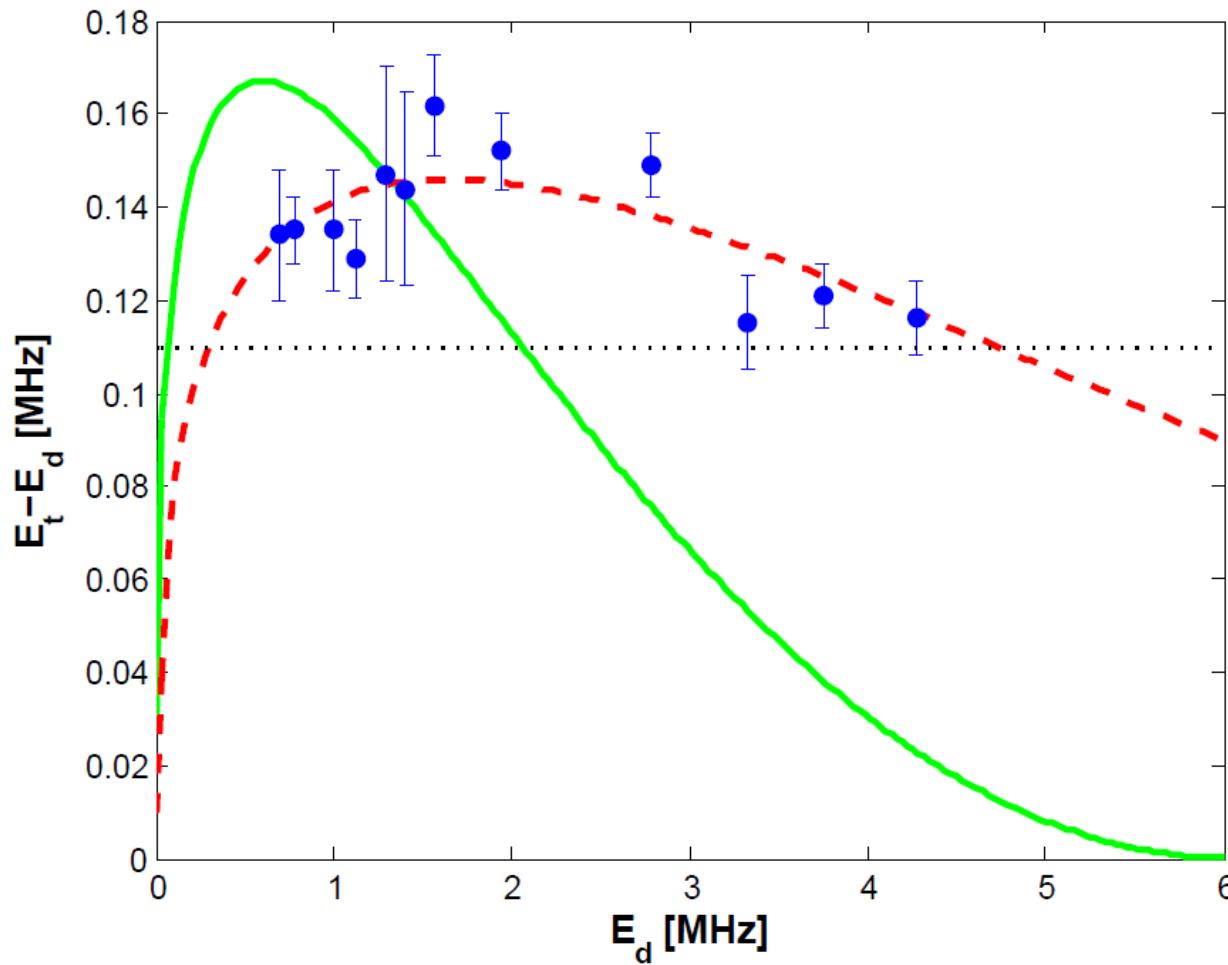


# Rf scans

Remaining atoms after rf-pulse at different magnetic fields.



# Trimer-dimer energy difference



Estimation:  
 $a_* \sim 180a_0$

# Beyond universality theory

## Beyond universality in three-body recombination: An effective field theory treatment

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<sup>2</sup> Helmholtz-Institut für Strahlen- und Kernphysik, Universität Bonn - D-53115, Bonn, Germany, EU

<sup>3</sup> Department of Physics, Ohio State University - Columbus, OH 43120, USA

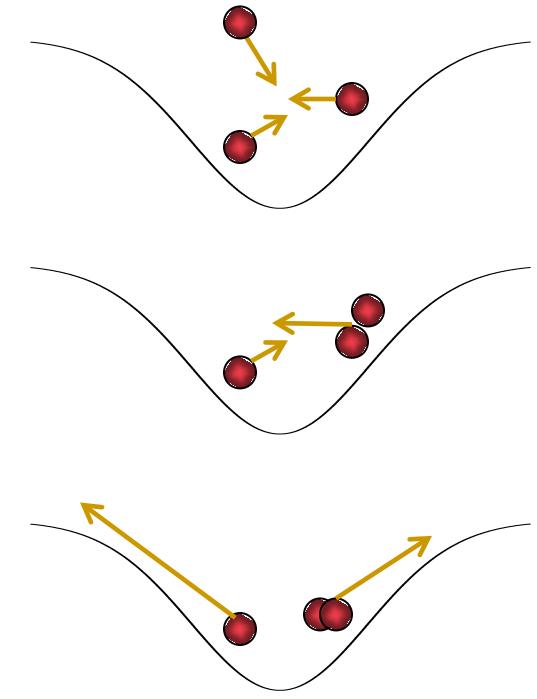
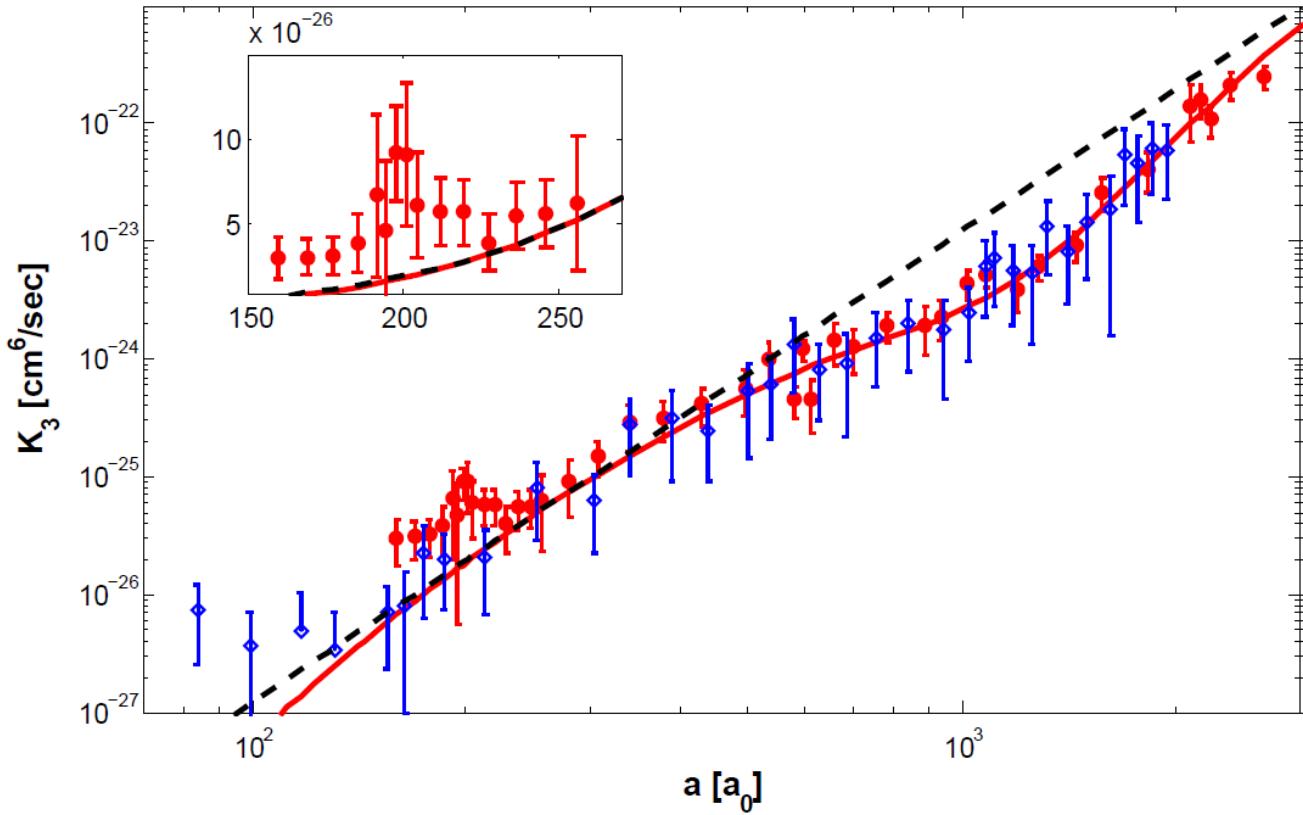
<sup>4</sup> Institute for Nuclear Theory, University of Washington - Seattle, WA 98102, USA

**Abstract** – We discuss the impact of a finite effective range on three-body systems interacting through a large two-body scattering length. By employing a perturbative analysis in an effective field theory well suited to this scale hierarchy we find that an additional three-body parameter is required for consistent renormalization once range corrections are considered. This allows us to extend previously discussed universal relations between different observables in the recombination of cold atoms to account for the presence of a finite effective range. We show that such range corrections allow us to simultaneously describe the positive and negative scattering-length loss features observed in recombination with  ${}^7\text{Li}$  atoms by the Bar-Ilan group. They do not, however, significantly reduce the disagreement between the universal relations and the data of the Rice group on  ${}^7\text{Li}$  recombination at positive and negative scattering lengths.

Prediction including finite effective range:  $a_* = (210 \pm 44)a_B$

# 3-body recombination data

Avalanche resonance.



# Conclusions and outlook

- Ultracold atoms are extremely suitable to study few body universal physics and halo quantum states.
- Remarkable progress (experimental and theoretical) has been achieved in recent years.
- There are still many open questions to be studied.