

Range effects on Efimov physics in cold atoms

An Effective field theory approach

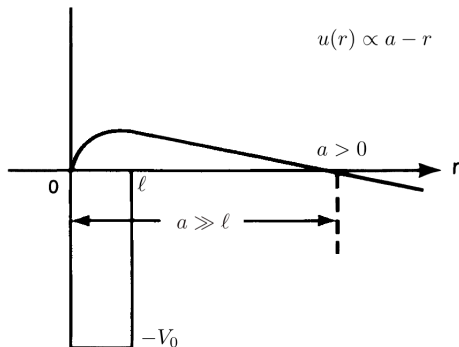
Chen Ji || TRIUMF

in collaboration with D. R. Phillips, L. Platter

INT workshop, November 7 2012



- **Separation of scales:**
 $a \gg \ell$
- **2-body S-wave universality:**
 $B_d = 1/Ma^2$



- **Separation of scales:**

$$a \gg \ell$$

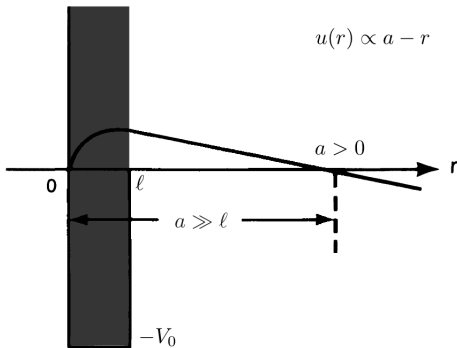
- **2-body S-wave universality:**

$$B_d = 1/Ma^2$$

- Physics at large distance is insensitive to physics at short distance

- Large-distance physics is studied in ℓ/a expansion

- Effects from SR-dynamics can be included in perturbation theory



Universal Physics exists in systems with
 $l \ll a$

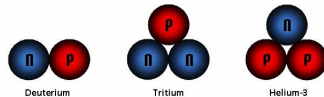
- Nuclear Physics

- **Few-nucleon systems (NN, NNN):**

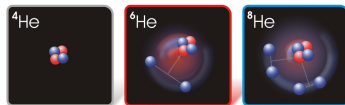
- i.e., $l_{np}^t \sim 1.7$ fm, $a_{np}^t \sim 5.4$ fm*

- **Halo nuclei (${}^6\text{He}, {}^{11}\text{Li}$)**

- $E_{\text{sep}} \ll E_{\text{core}}$



[www.theurbn.com]



[www.anl.gov]

Universal Physics exists in systems with $l \ll a$

- Nuclear Physics

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i.e., $l_{np}^t \sim 1.7 \text{ fm}$, $a_{np}^t \sim 5.4 \text{ fm}$

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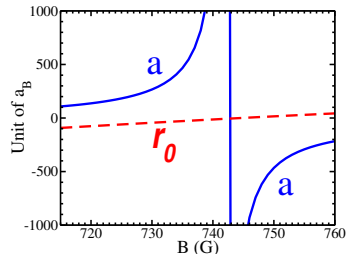
- Atomic Physics

- **Cold atomic gases (${}^{133}\text{Cs}$, ${}^7\text{Li}$, ${}^{39}\text{K}$):**

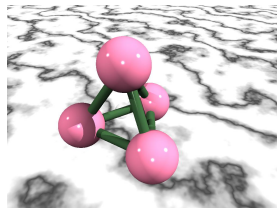
r_0 and a varies near Feshbach resonance

- **${}^4\text{He}$ atoms (dimer, trimer, tetramer):**

$l_{vdw} \sim 7 \text{ \AA}$, $a \sim 100 \text{ \AA}$



${}^7\text{Li}$ atoms [Gross et al. '09]




${}^6\text{He}$ tetramer [D. Blume]

- An approach to systems with a separation of scales
 - **Systems with $l \ll a$** \rightarrow an EFT with contact interactions
 - Few-nucleon systems \rightarrow pionless EFT
 - Halo nuclei \rightarrow halo EFT
 - Atomic systems \rightarrow short-range EFT
 - **Physical quantities are expanded in powers of r_0/a**


- **Contact interactions at LO**

- **2-body contact interaction (LO)**


$$= -iC_0$$

C_0 determined by a 2-body observable

- **3-body contact interaction (LO)**


$$= -iD_0$$

D_0 determined by a 3-body observable

- An approach to systems with a separation of scales
 - **Systems with $l \ll a$** \rightarrow an EFT with contact interactions
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 - **Physical quantities are expanded in powers of r_0/a**

- **Contact interactions at LO**

- **2-body contact interaction (LO)**

$$\begin{array}{ccc}
 \begin{array}{c} \diagup \\ \bullet \\ \diagdown \end{array} = -iC_0 & \xrightarrow{C_0 = g^2/\Delta} & \begin{array}{c} \text{introduce a dimer field} \\ \text{---} \end{array} \begin{array}{c} \diagup \\ \diagdown \end{array} = -i\sqrt{2}g
 \end{array}$$

C_0 determined by a 2-body observable

- **3-body contact interaction (LO)**

$$\begin{array}{ccc}
 \begin{array}{c} \diagup \\ \bullet \\ \diagdown \\ \bullet \\ \diagup \\ \diagdown \end{array} = -iD_0 & \xrightarrow{D_0 = -3hg^2/\Delta^2} & \begin{array}{c} \diagup \\ \diagdown \end{array} \begin{array}{c} \diagup \\ \diagdown \end{array} \begin{array}{c} \diagup \\ \diagdown \end{array} = ih
 \end{array}$$

D_0 determined by a 3-body observable

Bedaque, Hammer, van Kolck '99

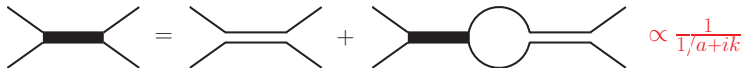
- **LO** $(r_0/a)^0$ EFT Lagrangian for 3 identical bosons

$$\mathcal{L} = \psi^\dagger \left(i\partial_0 + \frac{\nabla^2}{2m} \right) \psi - d^\dagger \left(i\partial_0 + \frac{\nabla^2}{4m} - \Delta \right) d - \frac{g}{\sqrt{2}} \left(d^\dagger \psi \psi + \text{h.c.} \right) + h d^\dagger d \psi^\dagger \psi + \dots$$

- terms with more derivatives are at higher orders $(r_0/a)^n$

- **Non-perturbative features at LO**

- atom-atom (dimer) scattering (tune g)



$$\propto \frac{1}{1/a + ik}$$

- atom-dimer scattering (tune h)



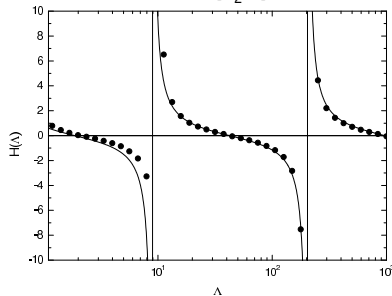
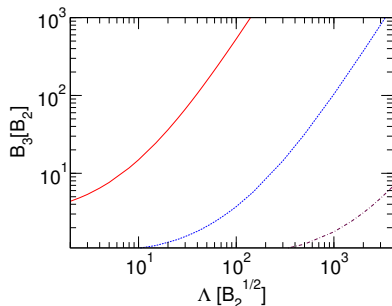
Bedaque, Hammer, van Kolck '99

- Without 3BF:

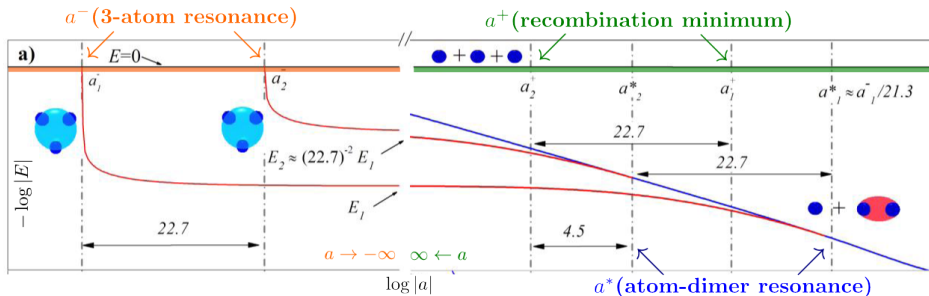
- 3-body spectrum:
 - cutoff dependent ($\Lambda \sim 1/\ell$)
 - Platter '09

- LO 3BF h :

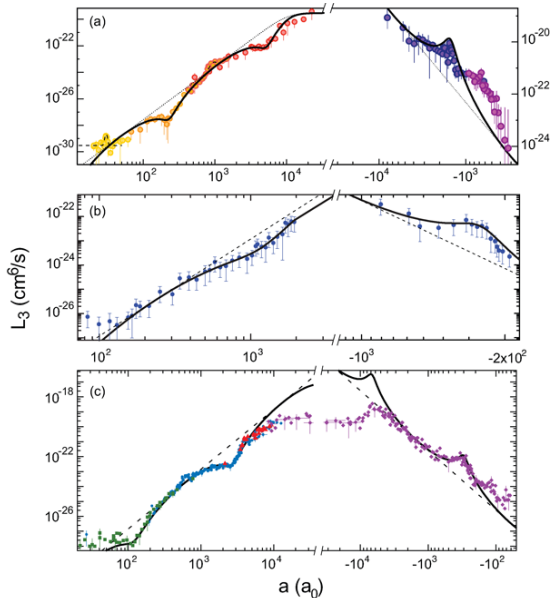
- tune $H(\Lambda) = \Lambda^2 h / 2mg^2$:
 - fix one 3-body observable
- limit cycle:
 - $H(\Lambda)$ periodic for $\Lambda \rightarrow \Lambda(22.7)^n$
 - Bedaque *et al.* '00
- scaling invariance \rightarrow Efimov physics
 - Efimov '71



- **Universal features in three-body systems (Efimov effects)**
 - 3-body spectrum: a function of scattering length a
 - **geometric spectrum**
 - $E_n = (22.7)^{-2} E_{n-1}$ in the limit $a \rightarrow \infty$
 - **universal relation of recombination features**
 - $a^* = a^+ / 4.5 = -a^- / 21.3$
 - i.e. $a_{(n)}^- = 22.7 a_{(n-1)}^-$



Zaccanti et al. '09



loss rates of free atoms in ultracold atomic gases

- ³⁹K atoms

Zaccanti *et al.* '09

- ⁷Li atoms ($|F=1, m_F=0\rangle$)

Gross *et al.* '09

- ⁷Li atoms ($|F=1, m_F=1\rangle$)

Pollack *et al.* '09

Pic. credit: Ferlaino, Grimm '10

- range effects on universal physics
 - 2-body observable: $k \cot \delta_0 = -\frac{1}{a} + \frac{r_0}{2} k^2 + \dots$
 - 3-body observables: \rightarrow in r_0/a expansion
- Previous EFT calculations of r_0/a corrections (fixed a):
 - Hammer, Mehen '01 (NLO)
 - Bedaque, Rupak, Griesshammer, Hammer '03 (N²LO, partial iteration)
 - Platter, Phillips '06 (N²LO, partial iteration)
- We perform a rigorous perturbative calculation
 - NLO (systems with variable a)
 - ultracold atomic gases (r_0/a varies near Feshbach resonance)
 - 3-body recombination
 - N²LO (systems with a fixed a)
 - ⁴He atoms ($r_0/a \sim 0.07$)
 - ⁴He trimer (bound state), atom-dimer phase shift (scattering state)

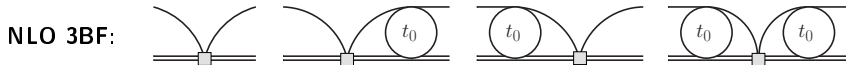
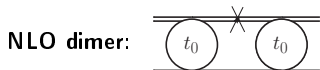
- Calculate r_0/a correction to atom-dimer amplitude
 - NLO correction to atom-atom propagator (**dimer**):

$$\text{Diagram: two parallel lines with a crossed-out double line between them} = ir_0 \frac{2\pi}{mg^2} \frac{1/a+ik}{-1/a+ik}$$

- NLO correction to atom-dimer contact term (**3BF**):

$$\text{Diagram: two parallel lines with a square contact vertex and a loop above it} = i \frac{2mg^2}{\Lambda^2} H_1(a, \Lambda)$$

- Contribution in 1st order perturbation theory:

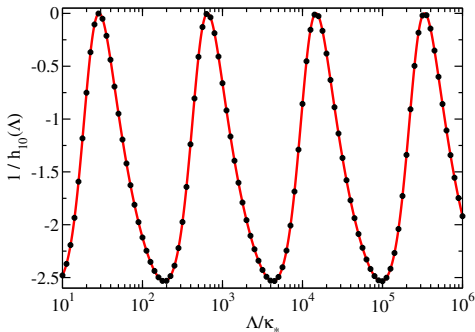


- NLO 3-body force:

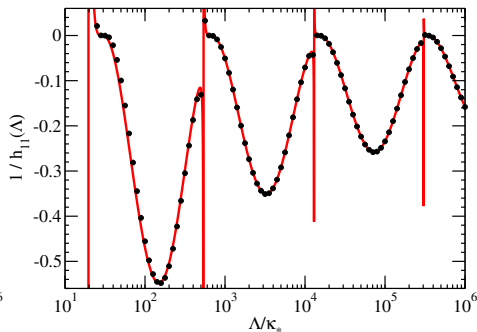
$$H_1(\Lambda) = r_0 \Lambda h_{10}(\Lambda) + r_0/a h_{11}(\Lambda)$$

- a **fixed**: h_{11} is absorbed (no additional 3-body input is needed)
- a **varies**: one additional 3-body input is needed

- $H_1(\Lambda) = r_0\Lambda h_{10}(\Lambda) + r_0/a h_{11}(\Lambda)$

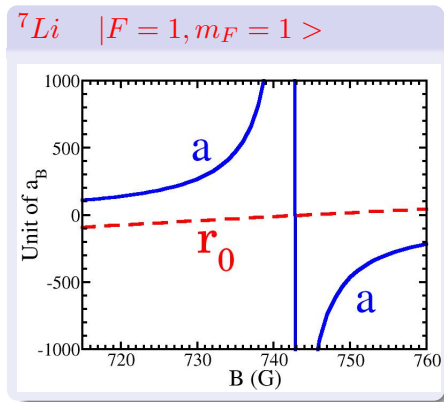
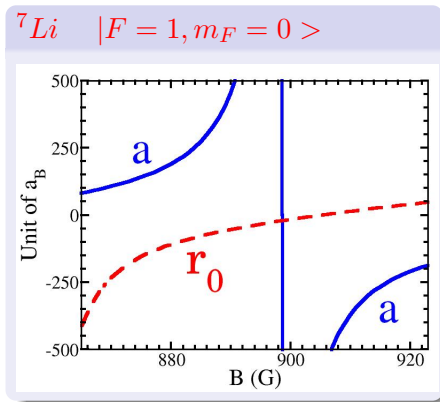


$1/h_{10}(\Lambda)$



$1/h_{11}(\Lambda)$

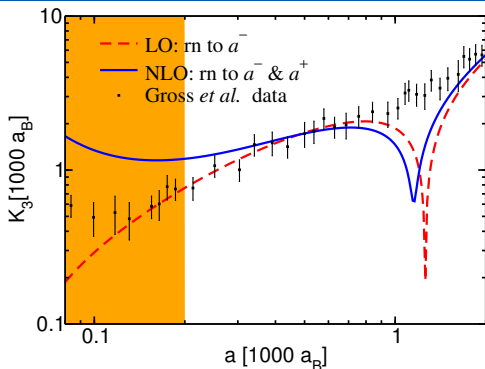
a and r_0 are functions of the magnetic field



Gross et al. '09

$$|F = 1, m_F = 0\rangle$$

- 3-body recombination rate of ${}^7\text{Li}$ atoms
 - NLO (r_0/a) range effects to recombination
 - positions are not shifted by deep-dimer at LO
- [Braaten *et al.* '08]



Experiment^{†‡}

3A res	$a_{(-)}^- [a_B]$	-264^\dagger
rec min	$a^+ [a_B]$	1160^\dagger
Ad res	$a^* [a_B]$	180^\ddagger

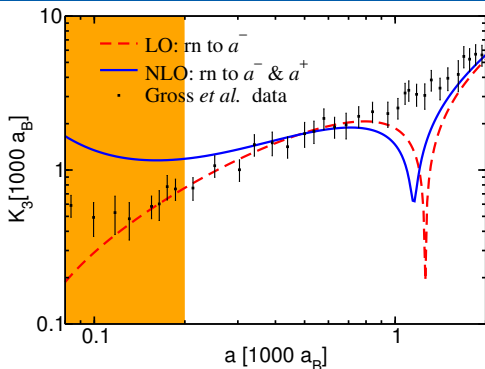
† Gross *et al.* '09

‡ Machtey *et al.* '12

★ Ji, Phillips, Platter '10

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	Experiment ^{†‡}	LO
3A res $a_{(-)}^- [a_B]$	-264^\dagger	-264
rec min $a^+ [a_B]$	1160^\dagger	1254
Ad res $a^* [a_B]$	180^\ddagger	281

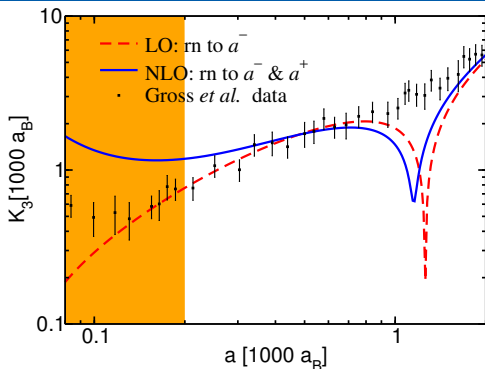
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	Experiment ^{†‡}	LO	LO
3A res $a_{(-)}^- [a_B]$	-264 [†]	-264	-244
rec min $a^+ [a_B]$	1160 [†]	1254	1160
Ad res $a^* [a_B]$	180 [‡]	281	259

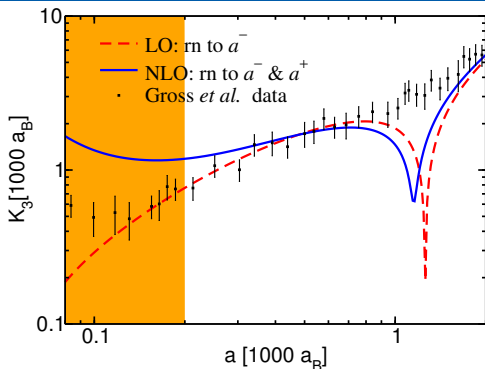
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3A res $a_{(-)}^- [a_B]$	-264 [†]	-264	-244	-264
rec min $a^+ [a_B]$	1160 [†]	1254	1160	1160
Ad res $a^* [a_B]$	180 [‡]	281	259	210(44)

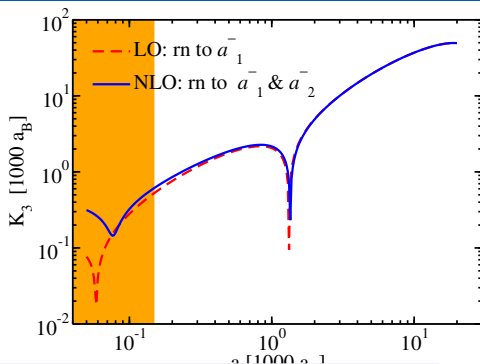
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$$|F = 1, m_F = 1\rangle$$

- ${}^7\text{Li}(m_F = 1)$ recombination [Pollack *et al.* '09]
- data disagree with universality by a factor of 2
- EFT analysis at NLO agrees with universality



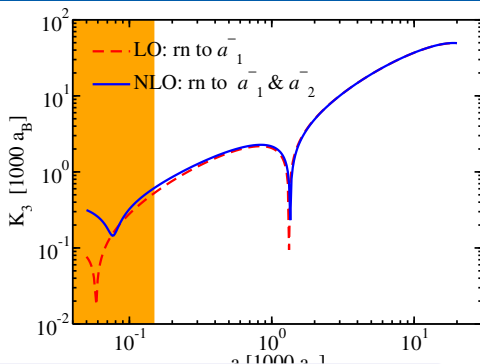
Experiment

[Pollack *et al.* '09]

3A res	$a_{(-)}^- [a_B]$	-298
3A res	$a^- [a_B]$	-6301
rec min	$a^+ [a_B]$	2676
Ad res	$a^* [a_B]$	608

$$|F = 1, m_F = 1\rangle$$

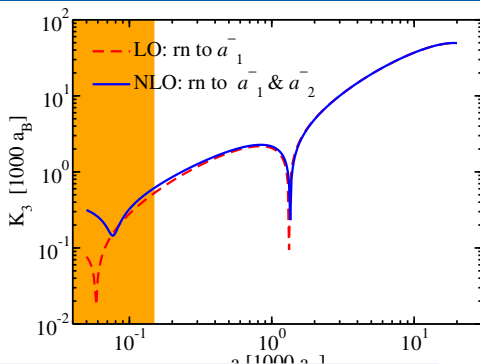
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	Experiment [Pollack <i>et al.</i> '09]	LO
3A res $a_{(-)}^- [a_B]$	-298	-298
3A res $a^- [a_B]$	-6301	-6763
rec min $a^+ [a_B]$	2676	1415
Ad res $a^* [a_B]$	608	317

$$|F = 1, m_F = 1\rangle$$

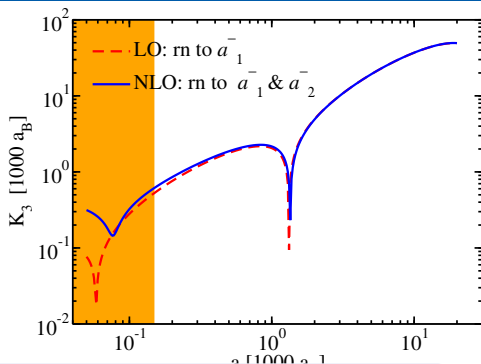
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	Experiment [Pollack <i>et al.</i> '09]	LO	LO
3A res $a_{(-)}^- [a_B]$	-298	-298	-278
3A res $a^- [a_B]$	-6301	-6763	-6301
rec min $a^+ [a_B]$	2676	1415	1319
Ad res $a^* [a_B]$	608	317	295

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	Experiment [Pollack <i>et al.</i> '09]	LO	LO	NLO [Ji <i>et al.</i> '10]
3A res $a_{(-)}^- [a_B]$	-298	-298	-278	-298
3A res $a^- [a_B]$	-6301	-6763	-6301	-6301
rec min $a^+ [a_B]$	2676	1415	1319	1348(151)
Ad res $a^* [a_B]$	608	317	295	356(55)

- $\frac{1}{a} = \gamma - \frac{1}{2}r_0\gamma^2$

Ji, Phillips, Platter '12

- $\frac{1}{a} = \gamma - \frac{1}{2}r_0\gamma^2$

- **NLO corrections**

$$\gamma_0 = -0.210\gamma_-^{(-)} + r_0(a_0) \cdot \mathcal{I}_0^{(-)}\gamma_-^{(-)2}$$

$$\gamma_* = -0.939\gamma_-^{(-)} + r_0(a_*) \cdot \mathcal{I}_*^{(-)}\gamma_-^{(-)2}$$

Ji, Phillips, Platter '12

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- **Linear correlation at NLO**

$$\mathcal{I}_*^{(-)} = -0.309 + 7.17 \mathcal{I}_0^{(-)}$$

Ji, Phillips, Platter '12

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- **Linear correlation at NLO**

$$\mathcal{I}_*^{(-)} = -0.309 + 7.17 \mathcal{I}_0^{(-)}$$

- **Universal relation at NLO**

$$\gamma_* = -0.939\gamma_-^{(-)} - 0.309 r_0(a_*)\gamma_-^{(-)2} + 7.17 \frac{r_0(a_*)}{r_0(a_0)} \left(\gamma_0 + 0.210\gamma_-^{(-)} \right)$$

Ji, Phillips, Platter '12

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- **NLO corrections**

$$\gamma_0 = -0.210\gamma_-^{(-)} + r_0(a_0) \cdot \mathcal{I}_0^{(-)}\gamma_-^{(-)2}$$

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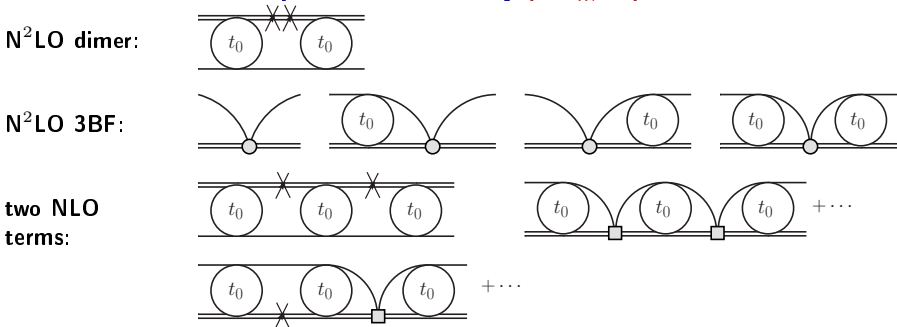
- **Universal relation at NLO**

$$\gamma_* = -0.939\gamma_-^{(-)} - 0.309 r_0(a_*)\gamma_-^{(-)2} + 7.17 \frac{r_0(a_*)}{r_0(a_0)} \left(\gamma_0 + 0.210\gamma_-^{(-)} \right)$$

$$\gamma_* = 4.47\gamma_0 - 7.02 r_0(a_*)\gamma_0^2 + 0.566 \frac{r_0(a_*)}{r_0(a_-^{(-)})} \left(\gamma_-^{(-)} + 4.76\gamma_0 \right)$$

Ji, Phillips, Platter '12

- $N^2\text{LO}$ corrections to atom-dimer scattering amplitude:
 - in 2nd order perturbation theory ($\sim r_0^2/a^2$):



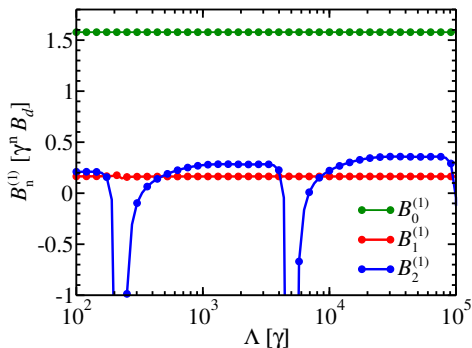
- $N^2\text{LO}$ 3-body force:

$$H_2(E, \Lambda) = r_0^2 \Lambda^2 h_{20}(\Lambda) + r_0^2 m E_3 h_{22}(\Lambda)$$

→ one additional 3-body input is needed (even when a is fixed)

c.f. Bedaque *et al.* '03 & Platter, Phillips '06

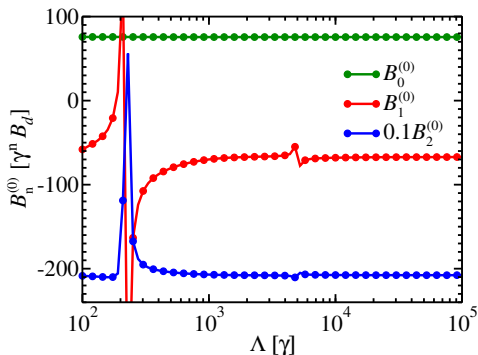
• $H_2 = r_0^2 \Lambda^2 h_{20}$



$a_{ad} \rightarrow \text{LO/NLO/N}^2\text{LO}$

$B_t^{(1)} = B_0^{(1)} + r_0 B_1^{(1)} + r_0^2 B_2^{(1)}$

• $H_2 = r_0^2 \Lambda^2 h_{20} + r_0^2 m E_t h_{22}$

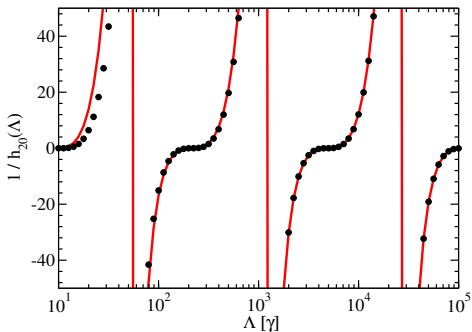


$a_{ad} \rightarrow \text{LO/NLO/N}^2\text{LO}$

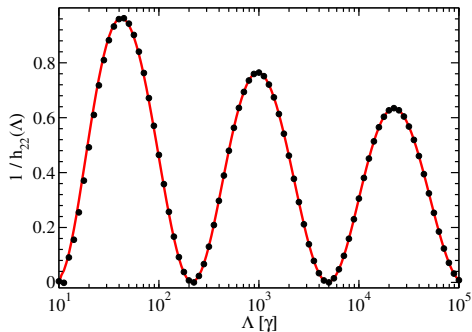
$B_t^{(1)} \rightarrow \text{N}^2\text{LO}$

$B_t^{(0)} = B_0^{(0)} + r_0 B_1^{(0)} + r_0^2 B_2^{(0)}$

- $$H_2(\Lambda) = r_0^2 \Lambda^2 h_{20}(\Lambda) + r_0^2 m E_t h_{22}(\Lambda)$$



$1/h_{20}(\Lambda)$



$1/h_{22}(\Lambda)$ [partial version]

Input	$B_t^{(1)} [B_d]$	$B_t^{(0)} [B_d]$	$a_{ad} [\gamma^{-1}]$	$r_{ad} [\gamma^{-1}]$
TTY potential	1.738	96.33	1.205	

Input		$B_t^{(1)} [B_d]$	$B_t^{(0)} [B_d]$	$a_{ad} [\gamma^{-1}]$	$r_{ad} [\gamma^{-1}]$
TTY potential		1.738	96.33	1.205	
a_{ad}	LO	1.723	97.12	1.205	0.8352
a_{ad}	NLO	1.736	89.72	1.205	0.9049
$a_{ad}, B_t^{(1)}$	N ² LO	1.738	116.9	1.205	0.9132

Input		$B_t^{(1)}$ [B_d]	$B_t^{(0)}$ [B_d]	a_{ad} [γ^{-1}]	r_{ad} [γ^{-1}]
TTY potential		1.738	96.33	1.205	
a_{ad}	LO	1.723	97.12	1.205	0.8352
a_{ad}	NLO	1.736	89.72	1.205	0.9049
$a_{ad}, B_t^{(1)}$	N ² LO	1.738	116.9	1.205	0.9132
$B_t^{(1)}$	LO	1.738	99.37	1.178	0.8752
$B_t^{(1)}$	NLO	1.738	89.77	1.201	0.9130
$B_t^{(1)}, a_{ad}$	N ² LO	1.738	115.9	1.205	0.9135

Input		$B_t^{(1)}$ [B_d]	$B_t^{(0)}$ [B_d]	a_{ad} [γ^{-1}]	r_{ad} [γ^{-1}]
TTY potential		1.738	96.33	1.205	
a_{ad}	LO	1.723	97.12	1.205	0.8352
a_{ad}	NLO	1.736	89.72	1.205	0.9049
$a_{ad}, B_t^{(1)}$	N ² LO	1.738	116.9	1.205	0.9132
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- **Difference in 2 renormalization schemes (LO→NLO→N²LO):**
 - atom-dimer effective range r_{ad} : 5% → 0.9% → 0.02%
 - ground-state trimer $B_t^{(0)}$: 2% → 0.07% → 0.9%

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- **Difference in 2 renormalization schemes (LO→NLO→N²LO):**

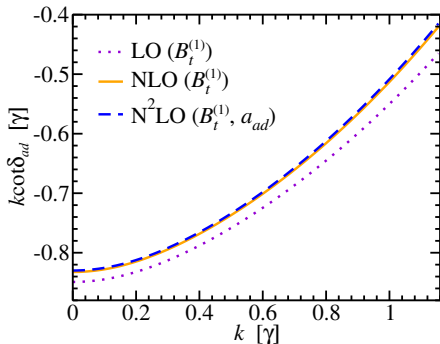
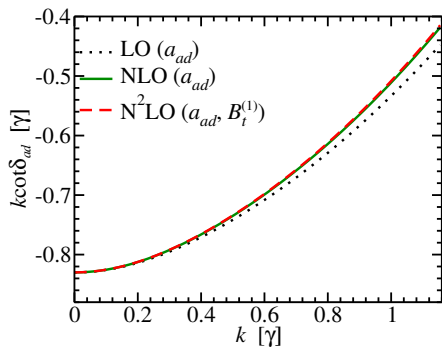
- atom-dimer effective range r_{ad} : 5% → 0.9% → 0.02%
- ground-state trimer $B_t^{(0)}$: 2% → 0.07% → 0.9%

- **Compare with TTY:**

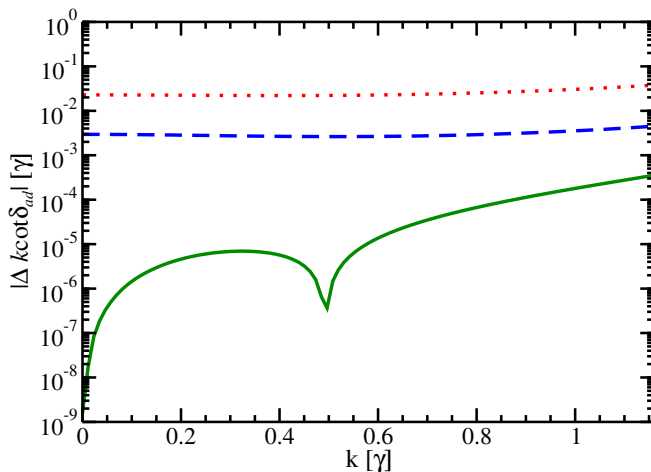
- $\Delta B_t^{(0)} \sim 20\%$
 - $r_0^2 B_t^{(0)} \sim 0.5$
 - EFT expansion needs to be corrected for deep trimer with U_{vdw} [Gao '98]

- a_{ad} → LO/NLO/N²LO
- $B_t^{(1)}$ → N²LO

- $B_t^{(1)}$ → LO/NLO/N²LO
- a_{ad} → N²LO



- LO → NLO → N²LO



- We studied effective-range corrections on Efimov physics in perturbation theory
- NLO for varying a :
 - recombination in cold atoms
 - $H_1(\Lambda) = r_0\Lambda h_{10}(\Lambda) + r_0/a h_{11}(\Lambda)$
 - one additional 3-body input is needed for NLO renormalization
- N²LO for fixed a :
 - He-4 trimer
 - $H_2(E, \Lambda) = r_0^2\Lambda^2 h_{20}(\Lambda) + r_0^2mE_3 h_{22}(\Lambda)$
 - one additional 3-body data is needed for N²LO renormalization
- Range corrections are also important in nuclear physics