

# **Compton Scattering at HIGS: from Giant Resonances to Spin Polarizabilities**

## **Compton@HIGS Collaboration**

- George Washington University
  - Jerry Feldman
  - Mark Sikora
- Duke University/TUNL
  - Luke Myers
  - Henry Weller
  - Mohammad Ahmed
  - Jonathan Mueller
  - Seth Henshaw
- University of Kentucky
  - Mike Kovash



# Outline

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- **What (and where) is HIGS?**

- **What have we done so far at HIGS?**

- polarized Compton scattering study of IVGQR
- elastic Compton scattering on  ${}^6\text{Li}$

- **What are we planning to do at HIGS?**

- elastic Compton scattering on deuterium
  - ✓ neutron polarizability
- polarized Compton scattering on proton
  - ✓ proton electric polarizability
- double-polarized Compton scattering on proton
  - ✓ proton spin polarizability
- double-polarized Compton scattering on  ${}^3\text{He}$ 
  - ✓ neutron spin polarizability

# **Background Information on HIGS**

# United States



# North Carolina



# Duke University



# TUNL

Triangle Universities  
Nuclear Laboratory

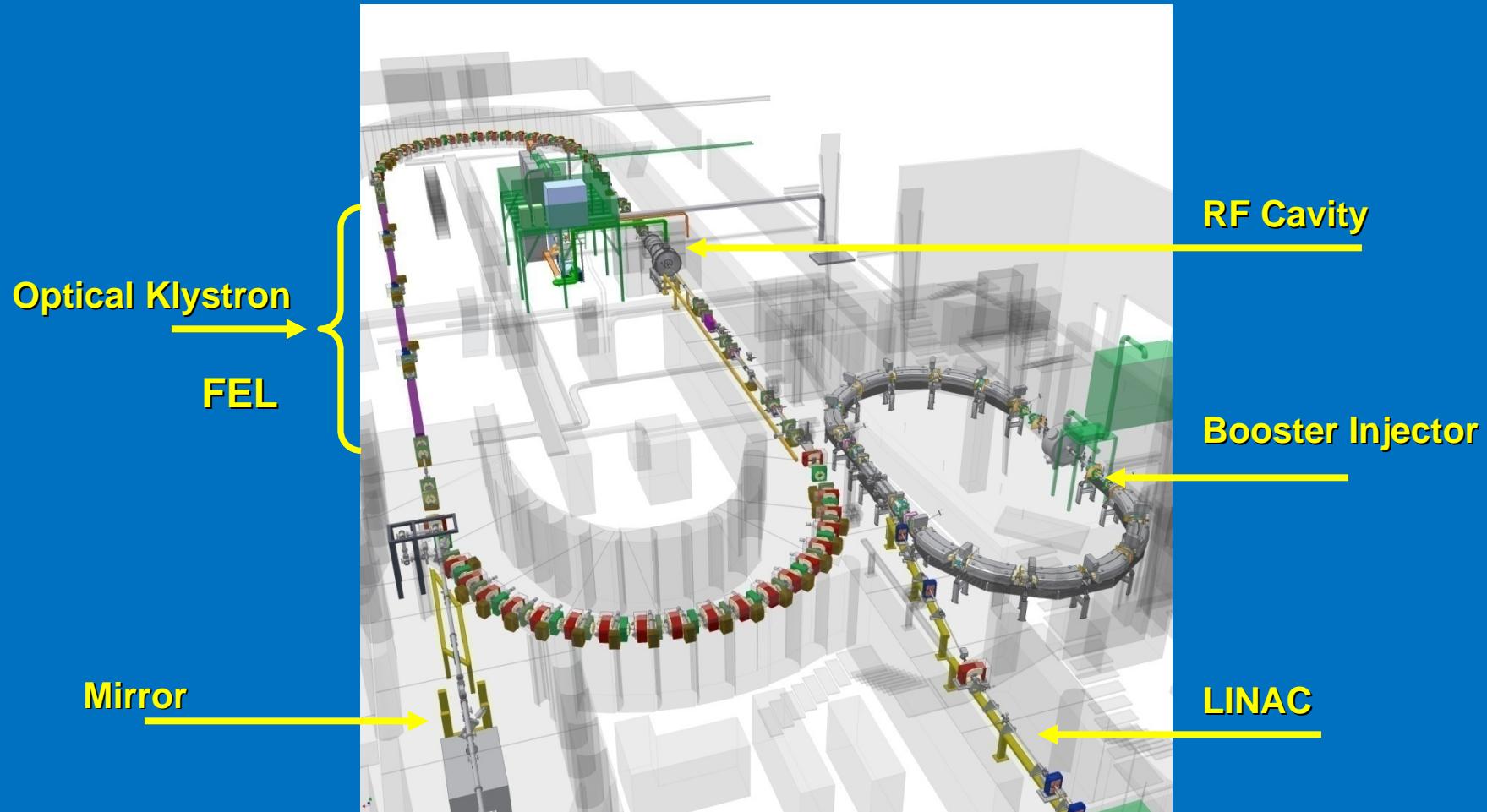


# Duke Free-Electron Laser Lab



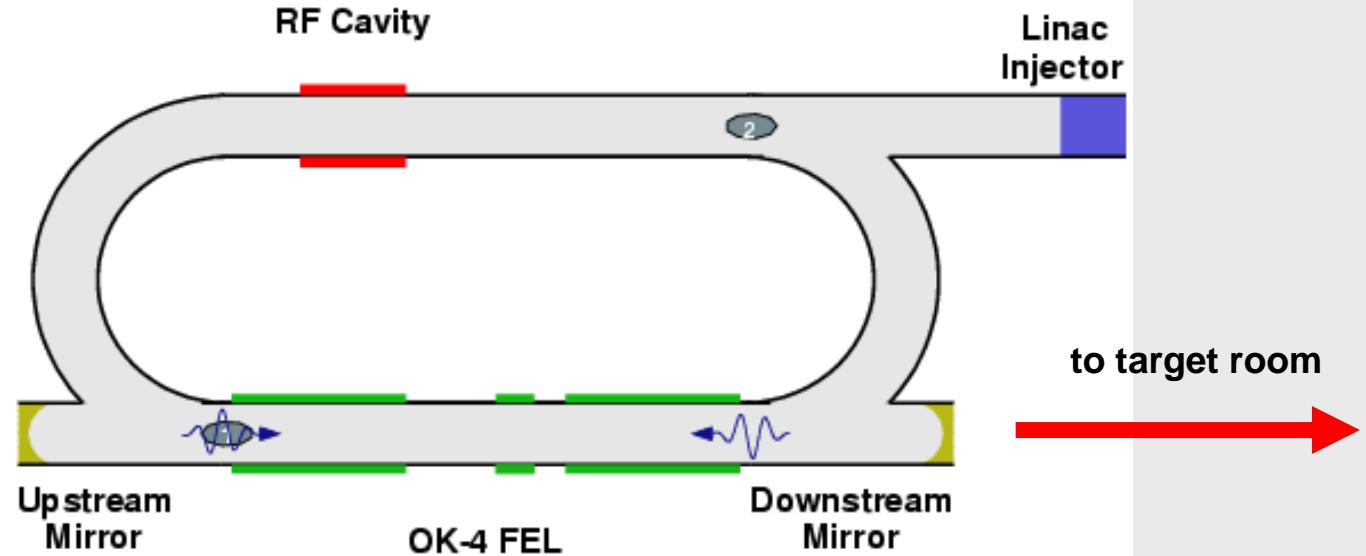
# Storage Ring and Booster

Circularly and linearly polarized  $\gamma$  rays, nearly **monoenergetic** ( $E_\gamma = 2\text{--}90 \text{ MeV}$ )  
Utilizes Compton backscattering to generate  $\gamma$  rays



# HIGS Photon Beam

## Two Bunch Mode



# HIGS Photon Beam

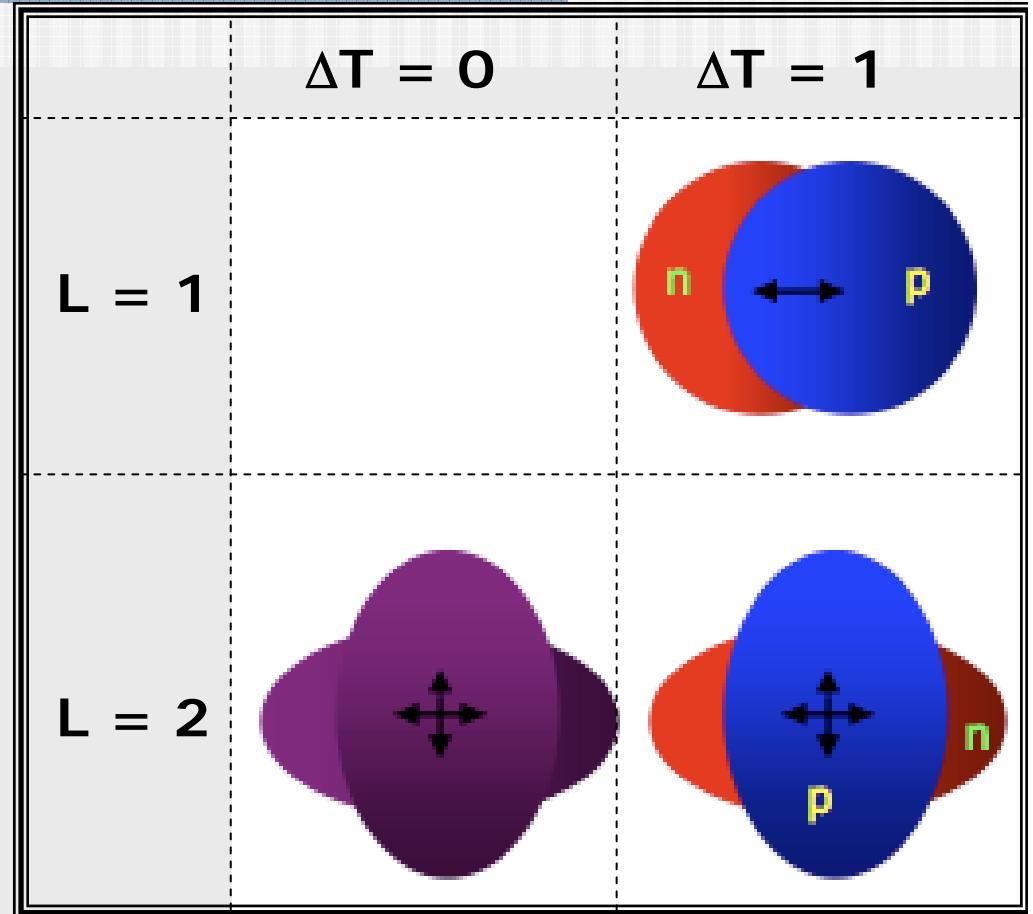
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- monoenergetic photons up to  $\sim 90$  MeV
  - energy will reach  $\sim 160$  MeV by 2015
- 100% linear or circular polarization
- high photon beam intensity
  - $\sim 10^7$  Hz at 20-60 MeV
  - $\sim 10^8$  Hz below 15 MeV
- low beam-related background
  - no bremsstrahlung typical of tagged photons

# **Polarized Compton Scattering for I VGQR Systematics**

# Giant Resonances

- collective nuclear excitations
  - GDR and ISGQR well known
  - **IVGQR poorly known**
  - photon as **isovector** probe
- 
- **use pol. photons for IVGQR**
  - map systematics vs. A
  - nuclear symmetry energy
    - ✓ neutron star eqn. of state



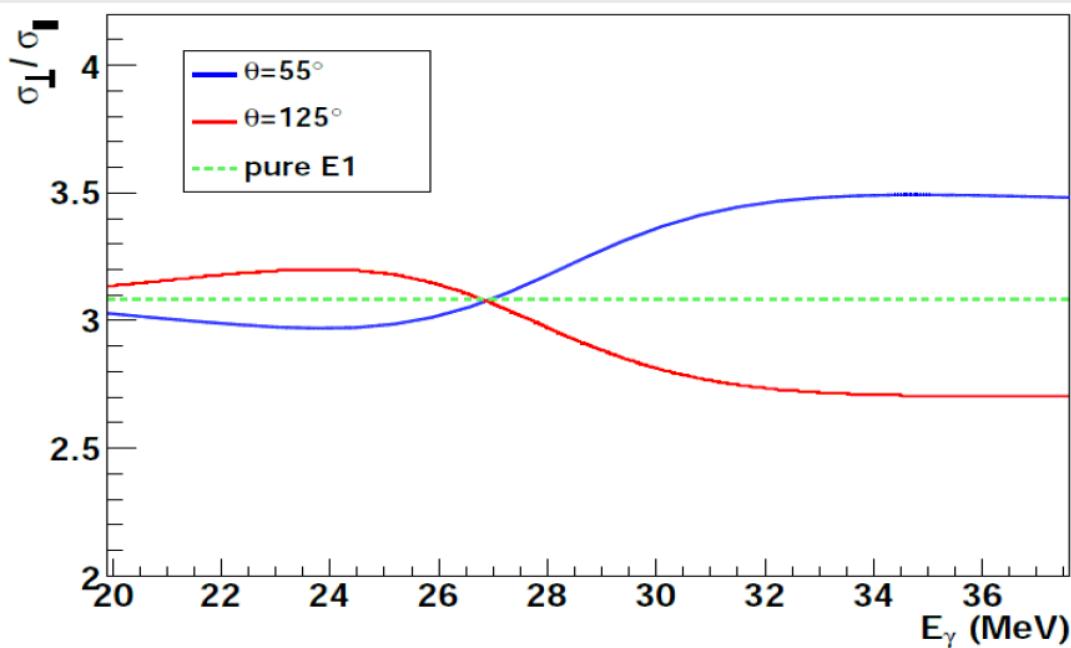
- ratio of H/V scattered photons is sensitive to E1/E2 interference
- sign difference in interference term at forward/backward angles

# Photon Asymmetry in IVGQR

$$\frac{\sigma_{\parallel}}{\sigma_{\perp}} = \cos^2 \theta + \frac{2|C^{E2}|}{|C^{E1}|} \cos(\phi_{E2} - \phi_{E1}) \left[ \cos^3 \theta - \cos \theta \right]$$

pure E1

E1/E2 interference



$$\cos(\phi_{E2} - \phi_{E1}) \begin{cases} < 0, E < E_{res} \\ > 0, E > E_{res} \end{cases}$$

$\theta$	$\cos^2 \theta$	$\cos^3 \theta - \cos \theta$
$125^\circ$	0.33	0.38
$55^\circ$	0.33	-0.38

# Phenomenological Formalism

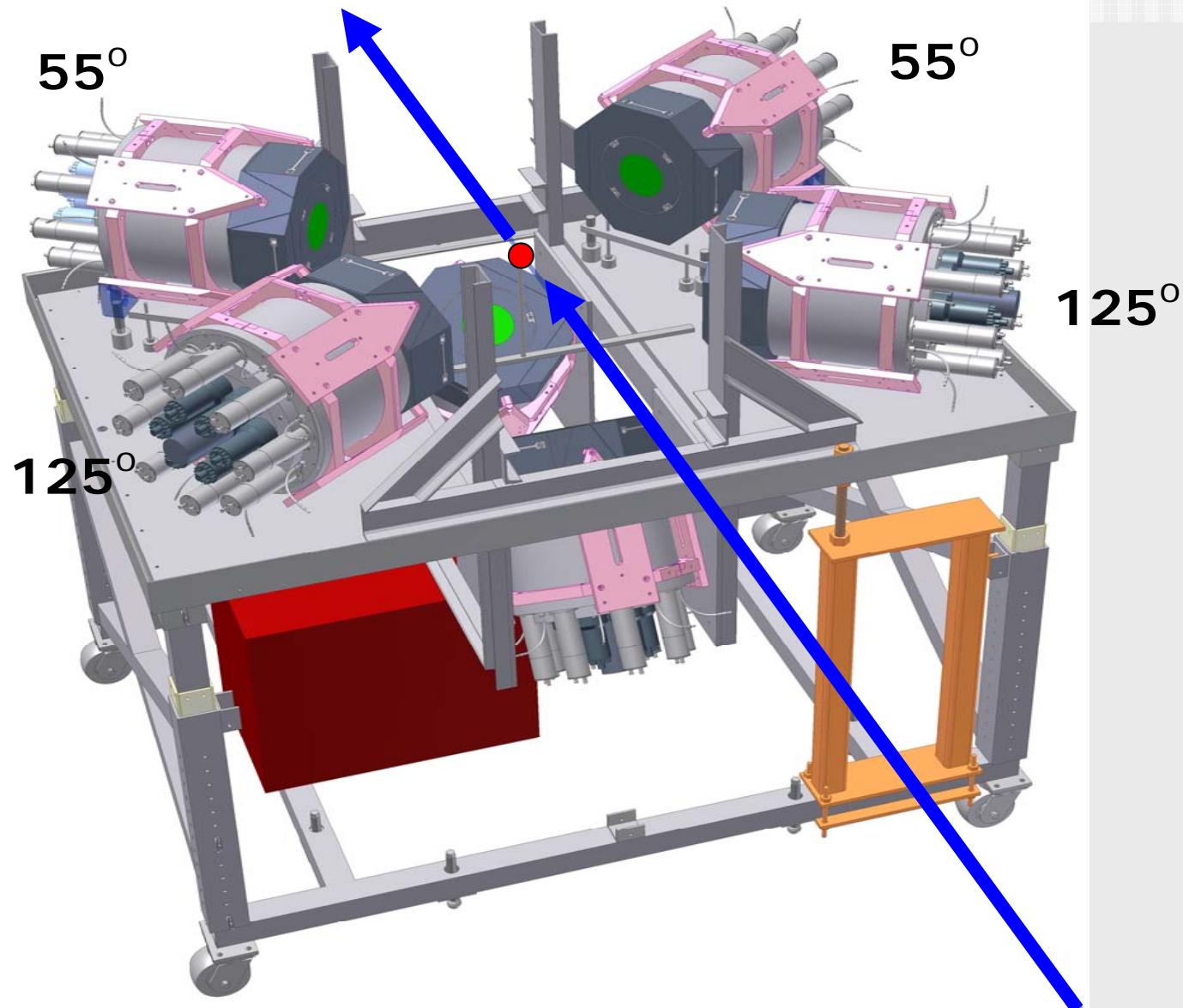
$$R(E, \theta) = R^{GR}(E, \theta) + R^{QD}(E, \theta) + R_1^{SG}(E, \theta) + R_2^{SG}(E, \theta)$$

$$\left\{ \begin{array}{l} R^{GR}(E, \theta) = f_{E1}(E)g_{E1}(\theta) + f_{E2}(E)g_{E2}(\theta) + \frac{NZ}{A}r_0[1 + \kappa_{GR}]g_{E1}(\theta) \\ \\ R^{QD}(E, \theta) = \left[ f_{QD}(E) + \frac{NZ}{A}r_0\kappa_{QD} \right] F_2(q)g_{E1}(\theta) \end{array} \right.$$

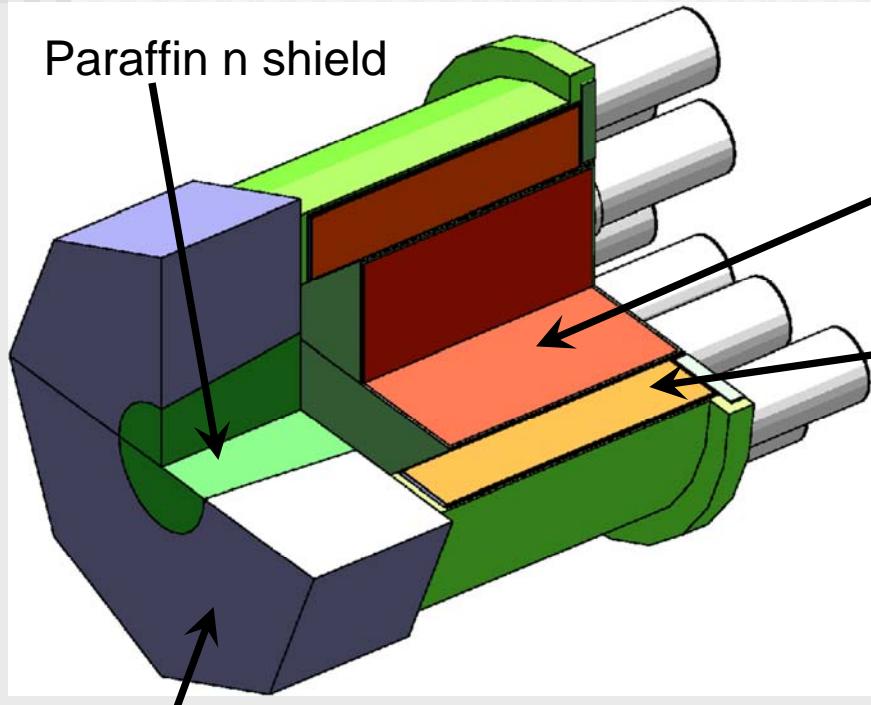
$$\left\{ \begin{array}{l} R_1^{SG}(E, \theta) = -F_1(q) \left[ Zr_0 - \left( \frac{E}{\hbar c} \right)^2 A\bar{\alpha} \right] g_{E1}(\theta) - \left[ \left( \frac{E}{\hbar c} \right)^2 A\bar{\beta} \right] g_{M1}(\theta) + O(E^4) \\ \\ R_2^{SG}(E, \theta) = -F_2(q) \frac{NZ}{A} r_0 (\kappa_{GR} + \kappa_{QD}) \end{array} \right.$$


# HINDA Array

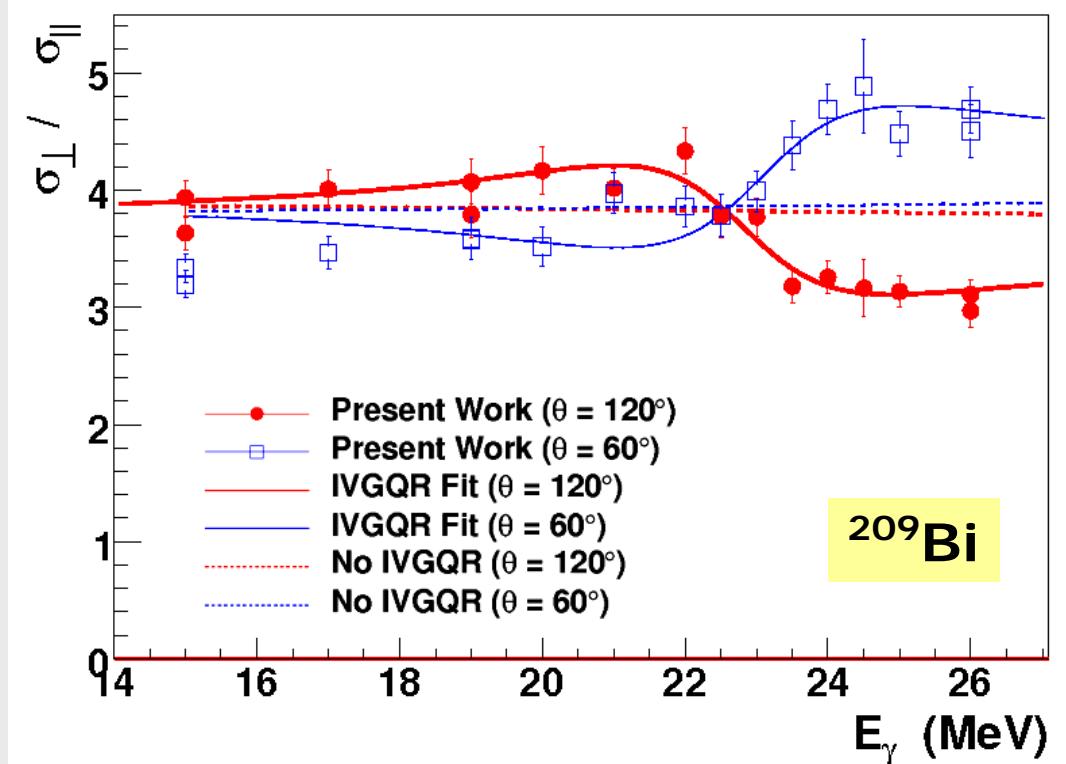
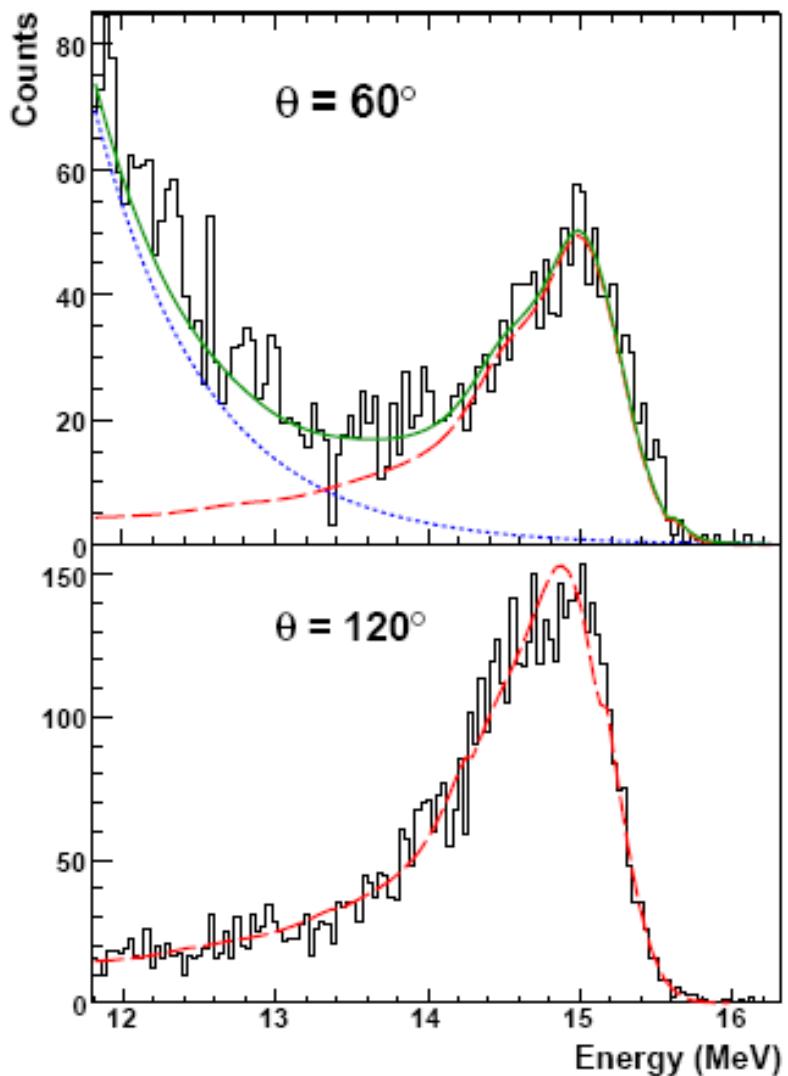
**HIGS NaI Detector Array**



# Nal Detectors



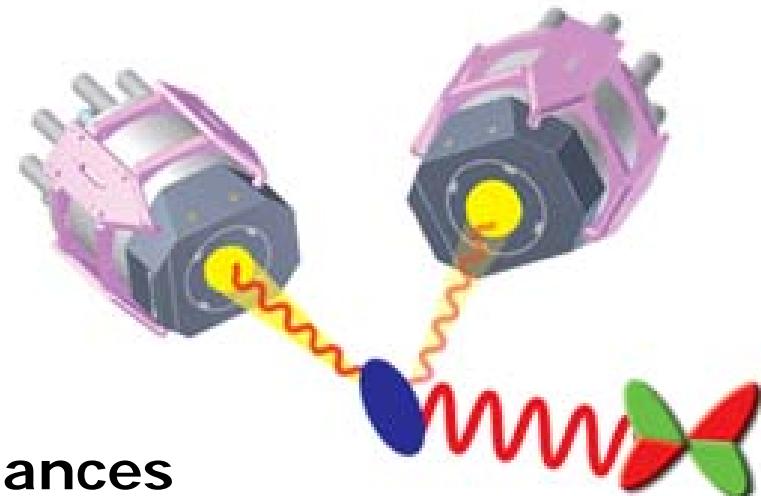
# Results for $^{209}\text{Bi}$



$E_{res}$ (MeV)	Width (MeV)	Strength (IVQEWSRs)
$23.0^{+0.13(\text{stat})}_{-0.18(\text{sys})}$	$3.9^{+0.7(\text{stat})}_{-0.6(\text{sys})}$	$0.56^{+0.04(\text{stat})}_{-0.05(\text{sys})}$

**New Method for Precise Determination of the Isovector Giant Quadrupole Resonances in Nuclei**S. S. Henshaw,<sup>1</sup> M. W. Ahmed,<sup>1,2</sup> G. Feldman,<sup>3</sup> A. M. Nathan,<sup>4</sup> and H. R. Weller<sup>1</sup><sup>1</sup>*Department of Physics and Triangle Universities Nuclear Laboratory, Duke University,  
TUNL Box 90308, Durham, North Carolina 27708-0308, USA*<sup>2</sup>*Department of Physics, North Carolina Central University, Durham, North Carolina 27707, USA*<sup>3</sup>*Department of Physics, George Washington University, Washington, D.C. 20052, USA*<sup>4</sup>*Department of Physics, University of Illinois, Urbana-Champaign, Illinois 61801, USA*

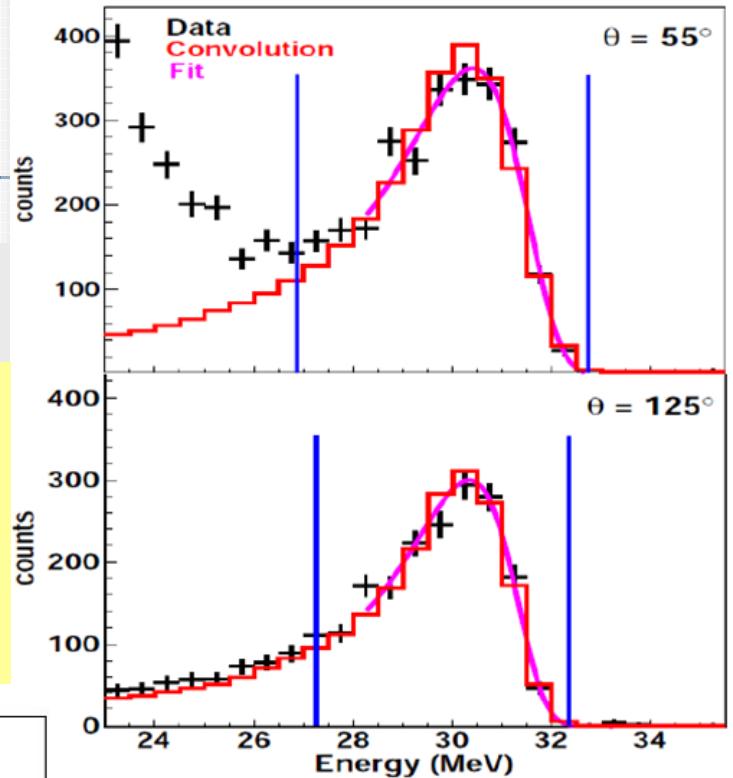
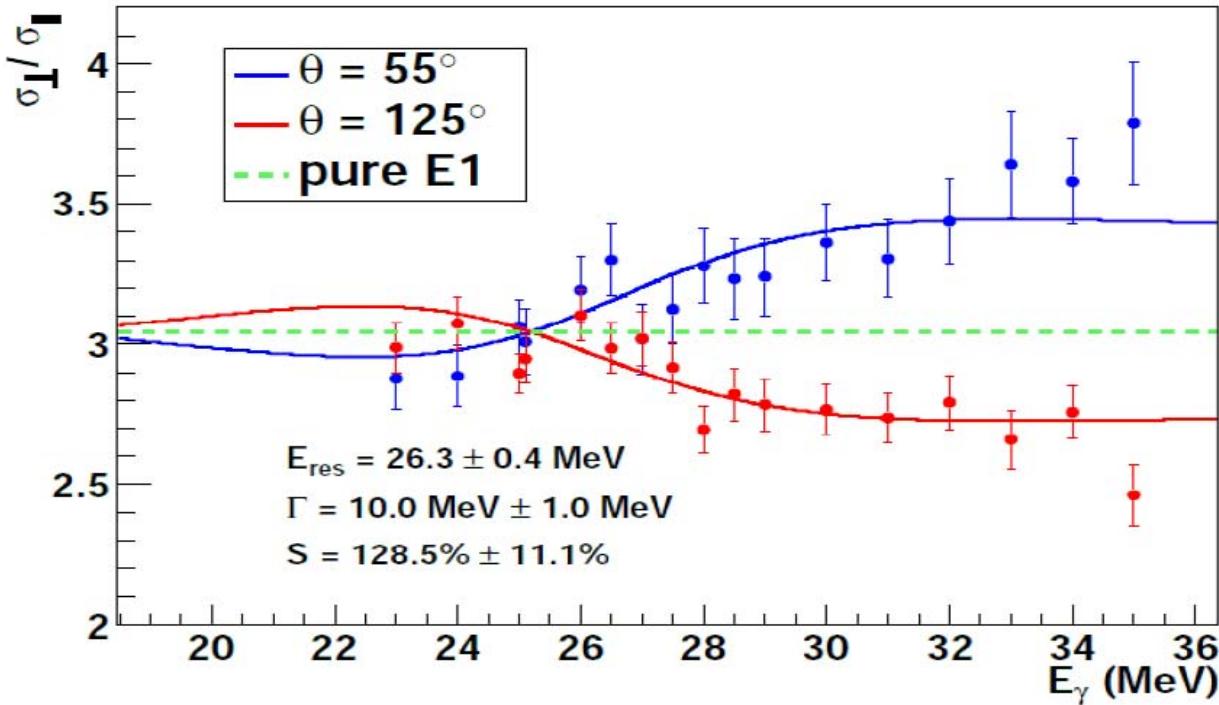
(Received 29 July 2011; published 23 November 2011)

**Synopsis: Ringing Nuclear Resonances**

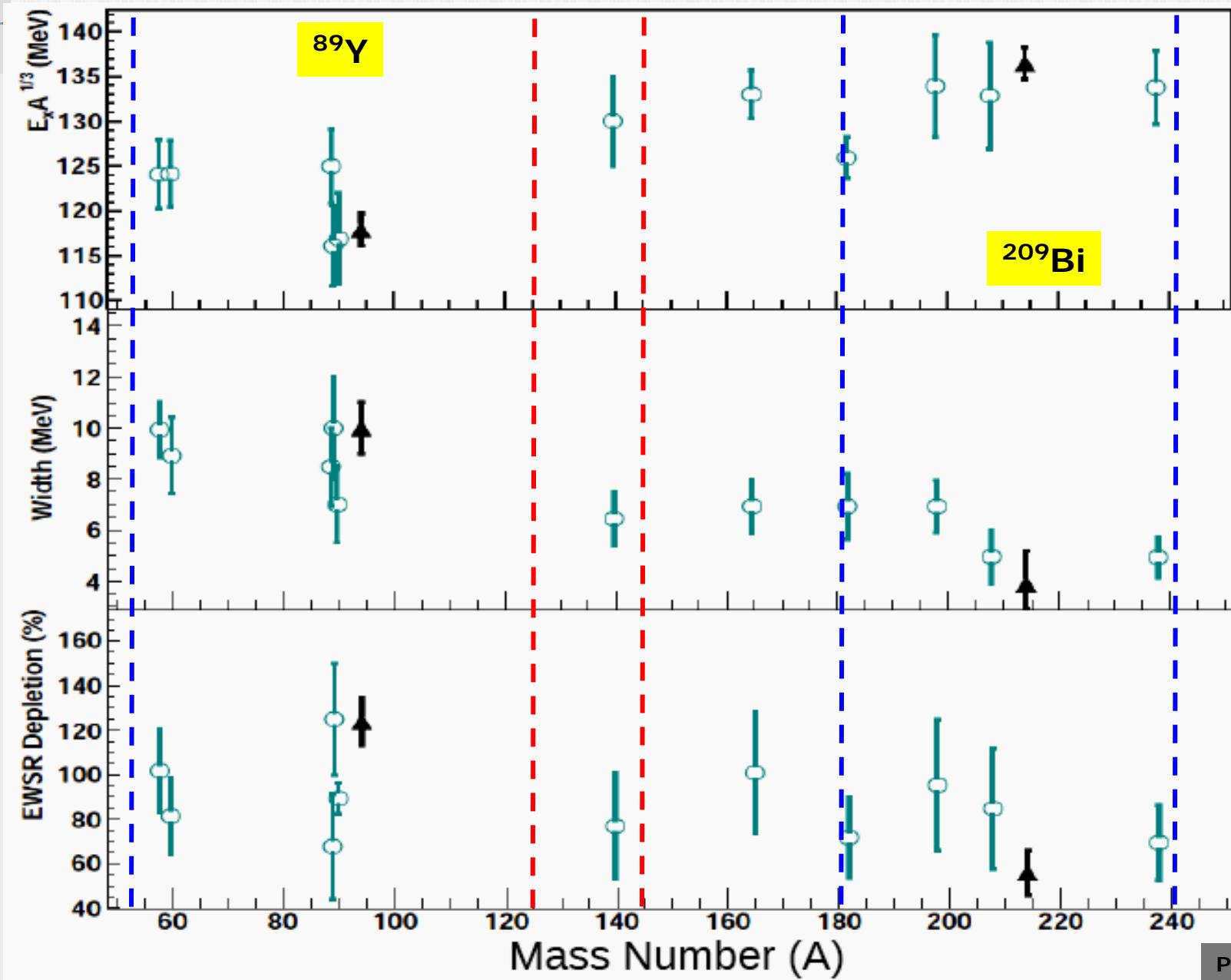
Courtesy Seth Henshaw, Duke University

# Results for $^{89}\text{Y}$

- extend measurements to  $^{89}\text{Y}$
- measure  $^{124}\text{Sn}$  early next year
- lease  $^{142}\text{Nd}$  target from ORNL for \$15k
- other tgts include  $A = 51, 181, 238$



# IVGQR Systematics



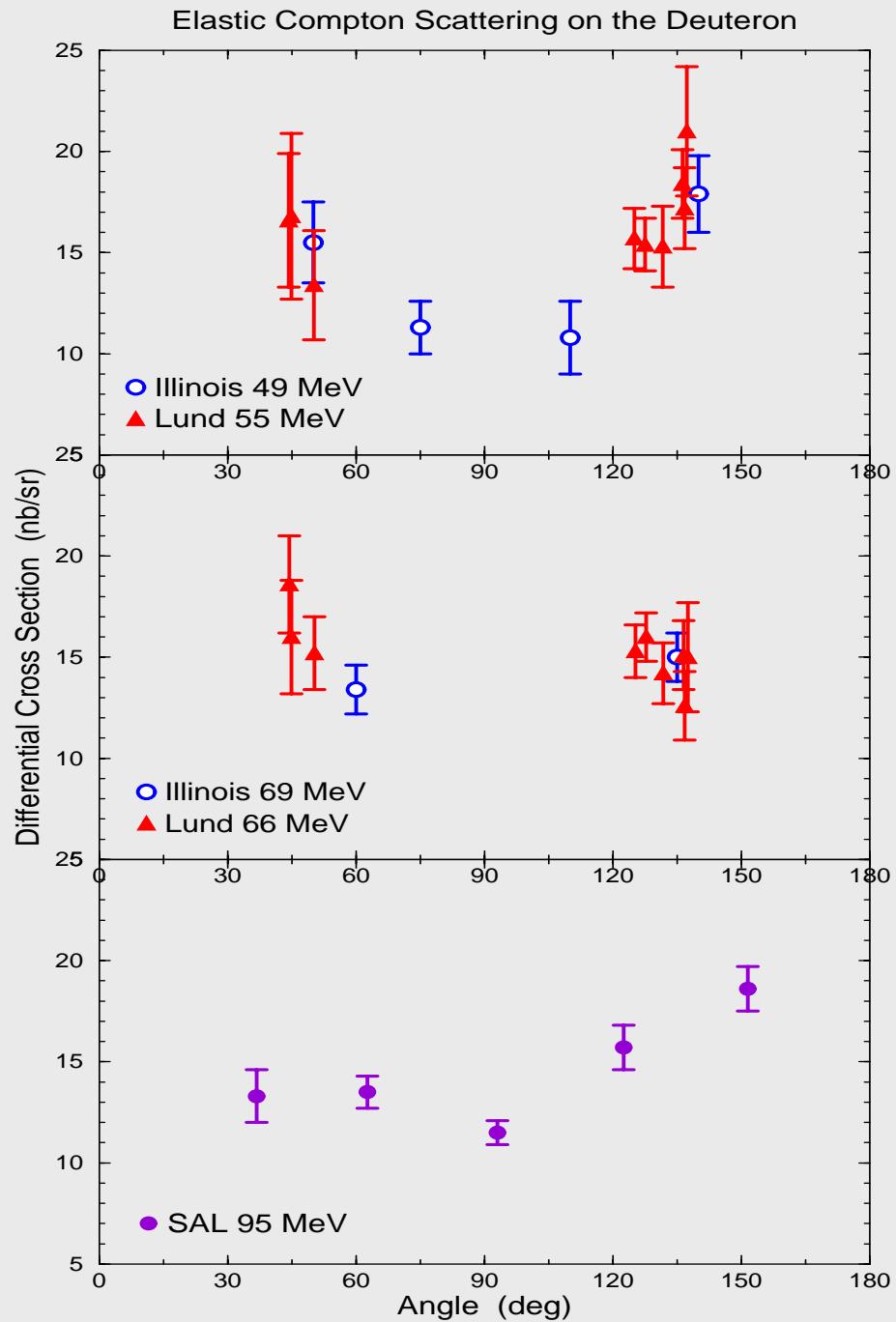
# **Compton Scattering**

## **on ${}^6\text{Li}$**

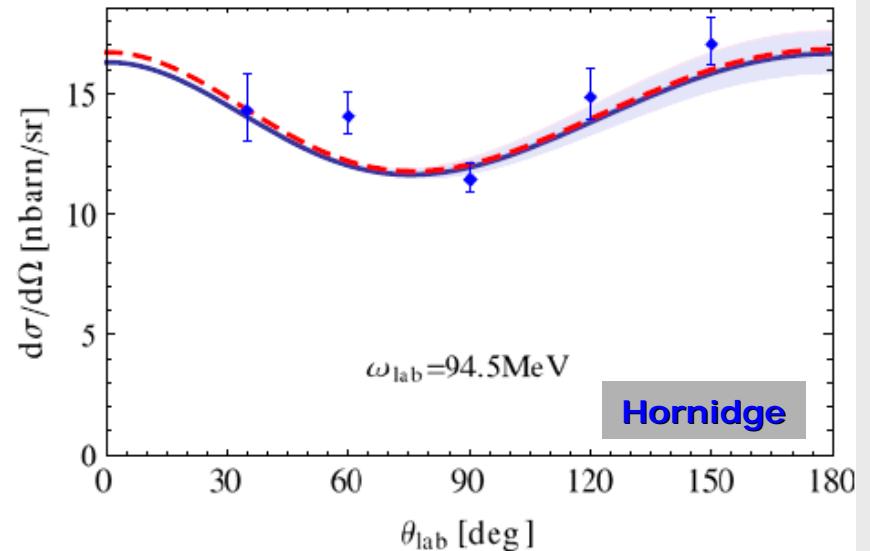
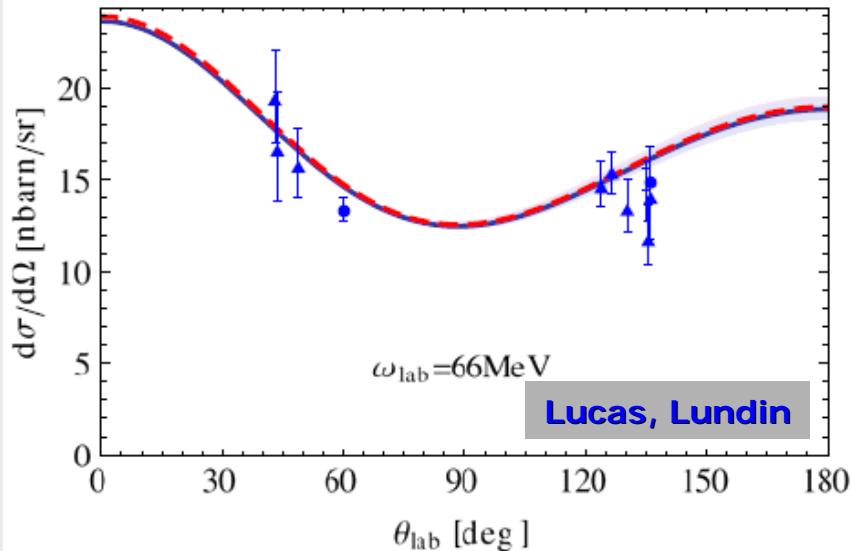
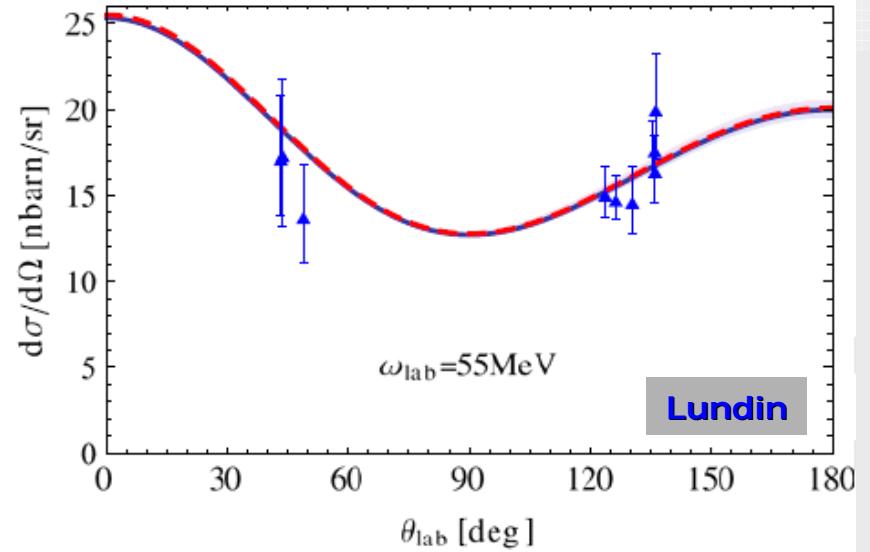
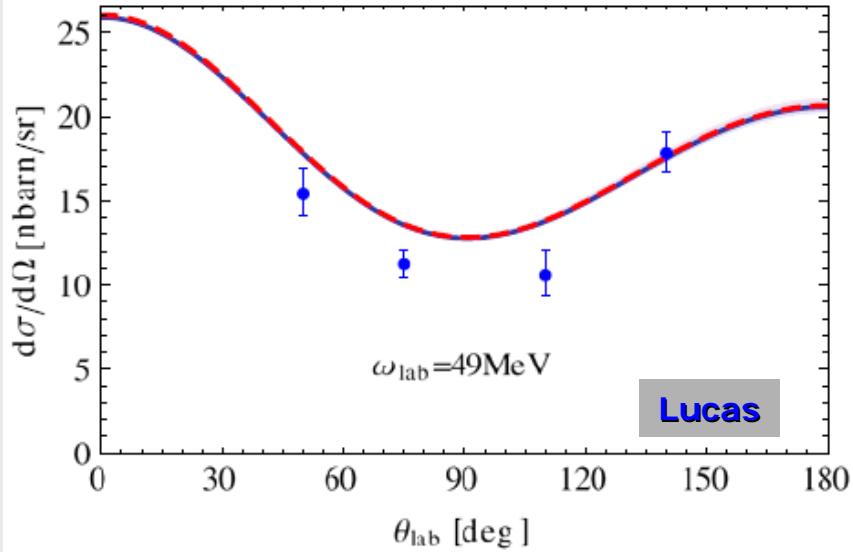
# World Data Set

$D(\gamma, \gamma)D$

- **Lucas – Illinois (1994)**  
 $E_\gamma = 49, 69$  MeV
- **Hornidge – SAL (2000)**  
 $E_\gamma = 85-105$  MeV
- **Lundin – Lund (2003)**  
 $E_\gamma = 55, 66$  MeV
- **Myers and Shonyozov  
(coming 2012)**  
**Illinois, GW, UK, Lund**  
 $E_\gamma = 58-115$  MeV



# EFT Fits to Deuteron Data



# Summary of Neutron Results

## ☐ Neutron scattering

- Schmiedmayer (91)

$$\alpha_n = 12.6 \pm 1.5(\text{stat}) \pm 2.0(\text{syst})$$

## ☐ Quasi-free Compton scattering

- Kossett (03)

$$\alpha_n = 12.5 \pm 1.8(\text{stat}) \begin{array}{l} +1.1 \\ -0.6 \end{array} (\text{syst}) \pm 1.1(\text{model})$$

$$\beta_n = 2.7 \mp 1.8(\text{stat}) \begin{array}{l} +0.6 \\ -1.1 \end{array} (\text{syst}) \mp 1.1(\text{model})$$

## ☐ Elastic Compton scattering

- data from Lucas (94), Hornidge (00), Lundin (03)

$$\alpha_n = 11.1 \pm 1.8 \text{ (stat)} \pm 0.4 \text{ (Baldin)} \pm 0.8 \text{ (theory)}$$

$$\beta_n = 4.1 \mp 1.8 \text{ (stat)} \pm 0.4 \text{ (Baldin)} \pm 0.8 \text{ (theory)}$$

$$\alpha_n = 11.6 \pm 1.5 \text{ (stat)} \pm 0.6 \text{ (Baldin)}$$

$$\beta_n = 3.6 \mp 1.5 \text{ (stat)} \pm 0.6 \text{ (Baldin)}$$

Hildebrandt 05

Nucleus	Energy (MeV)	Angles	Reference
D	49, 69	50°-140°	Lucas 1994
D	85-105	35°-150°	Hornidge 2000
D	55, 66	45°, 135°	Lundin 2003
D	60-115	60°, 90°, 120°, 150°	Myers, Shonyozov 2012
<sup>4</sup> He	23-70	45°, 135°	Wells 1990
<sup>4</sup> He	61	45°-150°	Proff 1999
<sup>12</sup> C	19-52	45°, 90°, 135°	Wright 1985
<sup>12</sup> C	58, 75	45°-135°	Hager 1995
<sup>12</sup> C	85-105	35°-150°	Warkentin 2001
<sup>16</sup> O	58, 75	45°-150°	Hager 1995
<sup>16</sup> O	61	50°, 135°	Proff 1999
<sup>16</sup> O	27-108	45°, 90°, 135°	Feldman 1996
<sup>40</sup> Ca	19-52	45°, 90°, 135°	Wright 1985
<sup>40</sup> Ca	58, 74	45°-150°	Proff 1999

# Experiment on ${}^6\text{Li}$ at HIGS

## □ experiment motivation

- exploit higher nuclear cross section to measure  $\alpha$  and  $\beta$ 
  - ✓ cross section scales as  $Z^2$ , so factor of 9x higher than  ${}^2\text{H}$
- solid  ${}^6\text{Li}$  target is simple
  - ✓ provided by Univ. of Saskatchewan
- no previous Compton data on  ${}^6\text{Li}$  exists      (except Pugh 1957)

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□ energies:  $E_\gamma = 60, 80 \text{ MeV}$

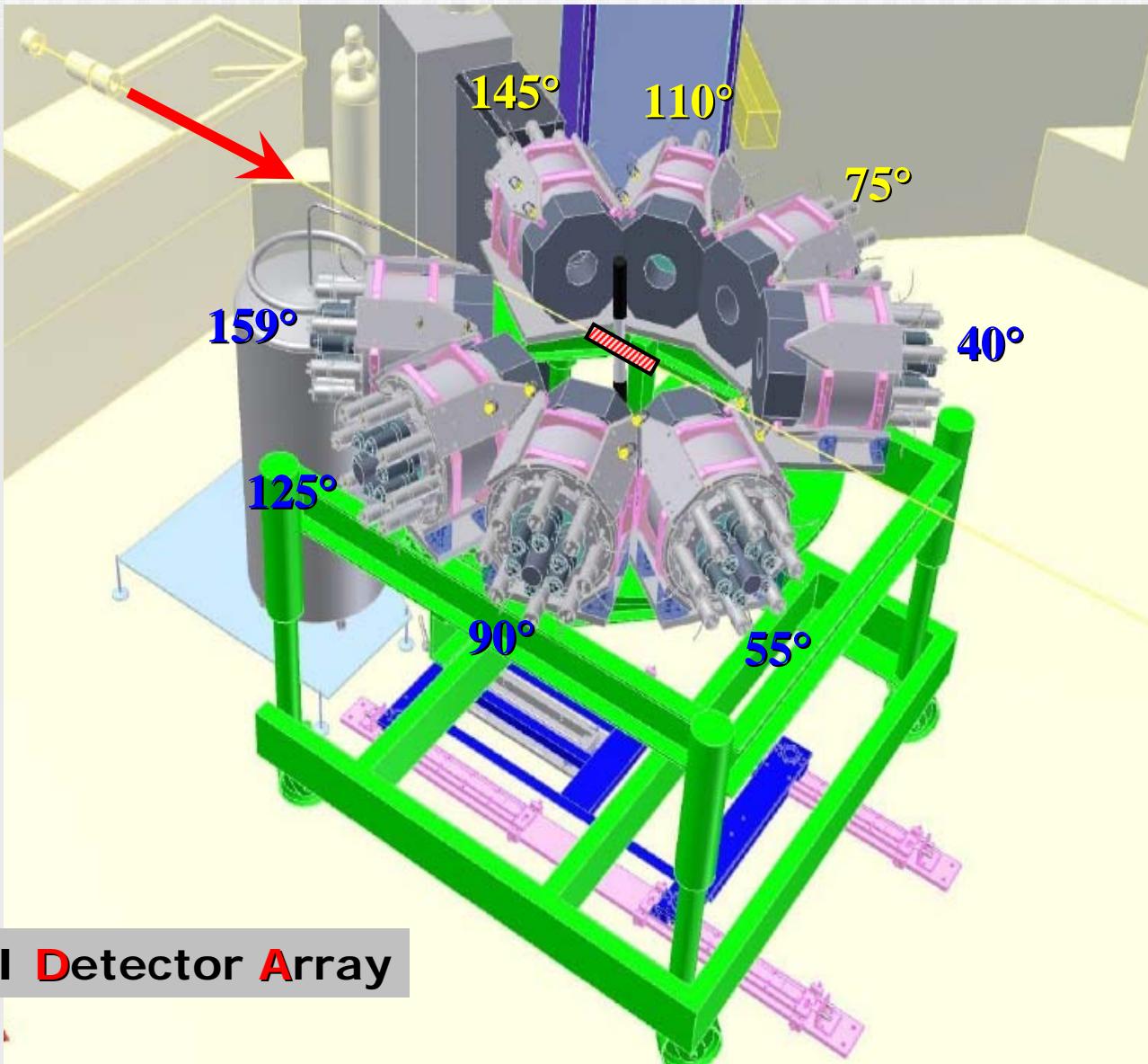
□ angles:  $\theta_\gamma = 40^\circ - 160^\circ$  ( $\Delta\theta = 17^\circ$ )

□ target: solid **12.7 cm long  ${}^6\text{Li}$**  cylinder (plus empty)

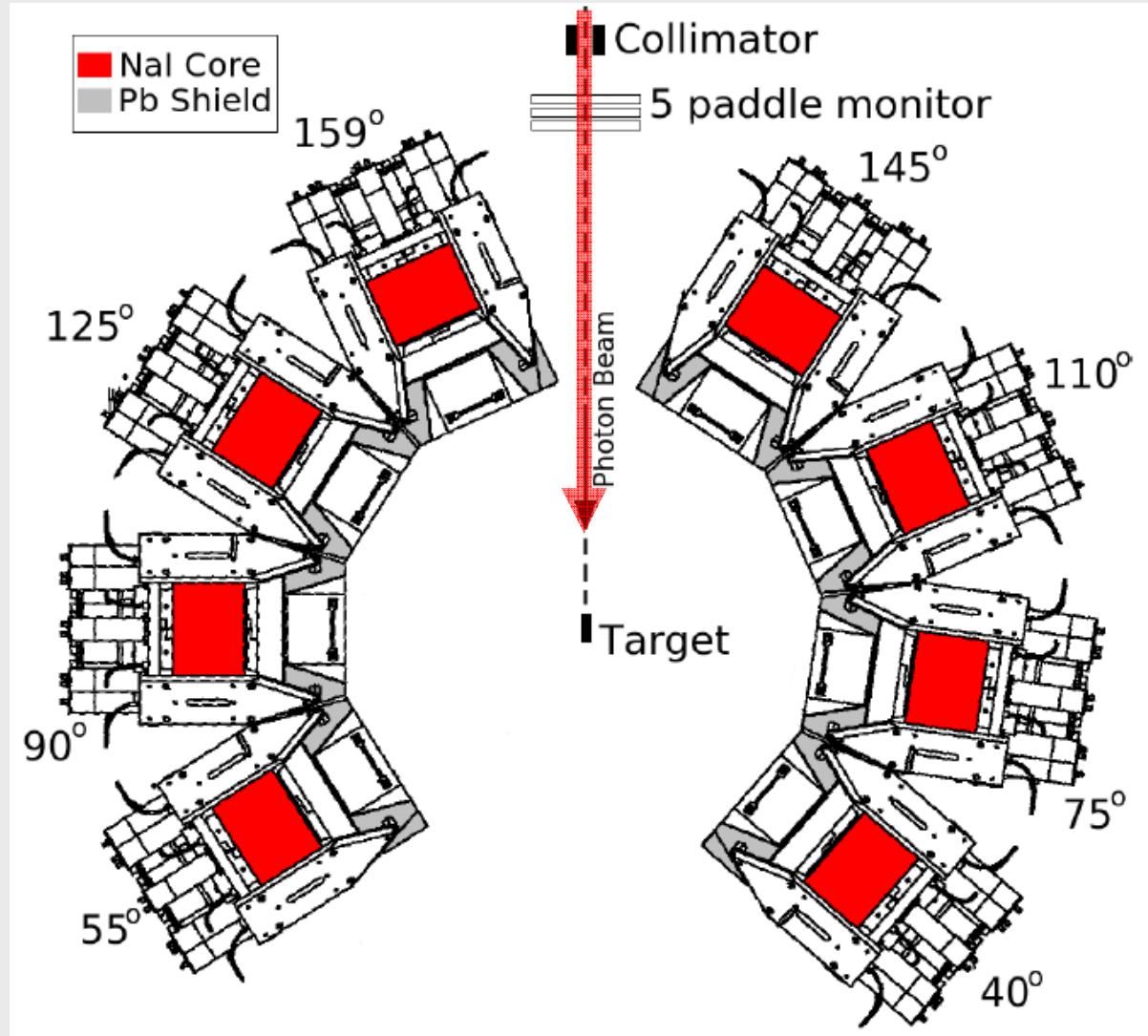
□ detectors: **eight 10"×12" NaI's** (HINDA array)

- good photon energy resolution ( $\Delta E_\gamma/E_\gamma < 5\%$ )

# HINDA Array

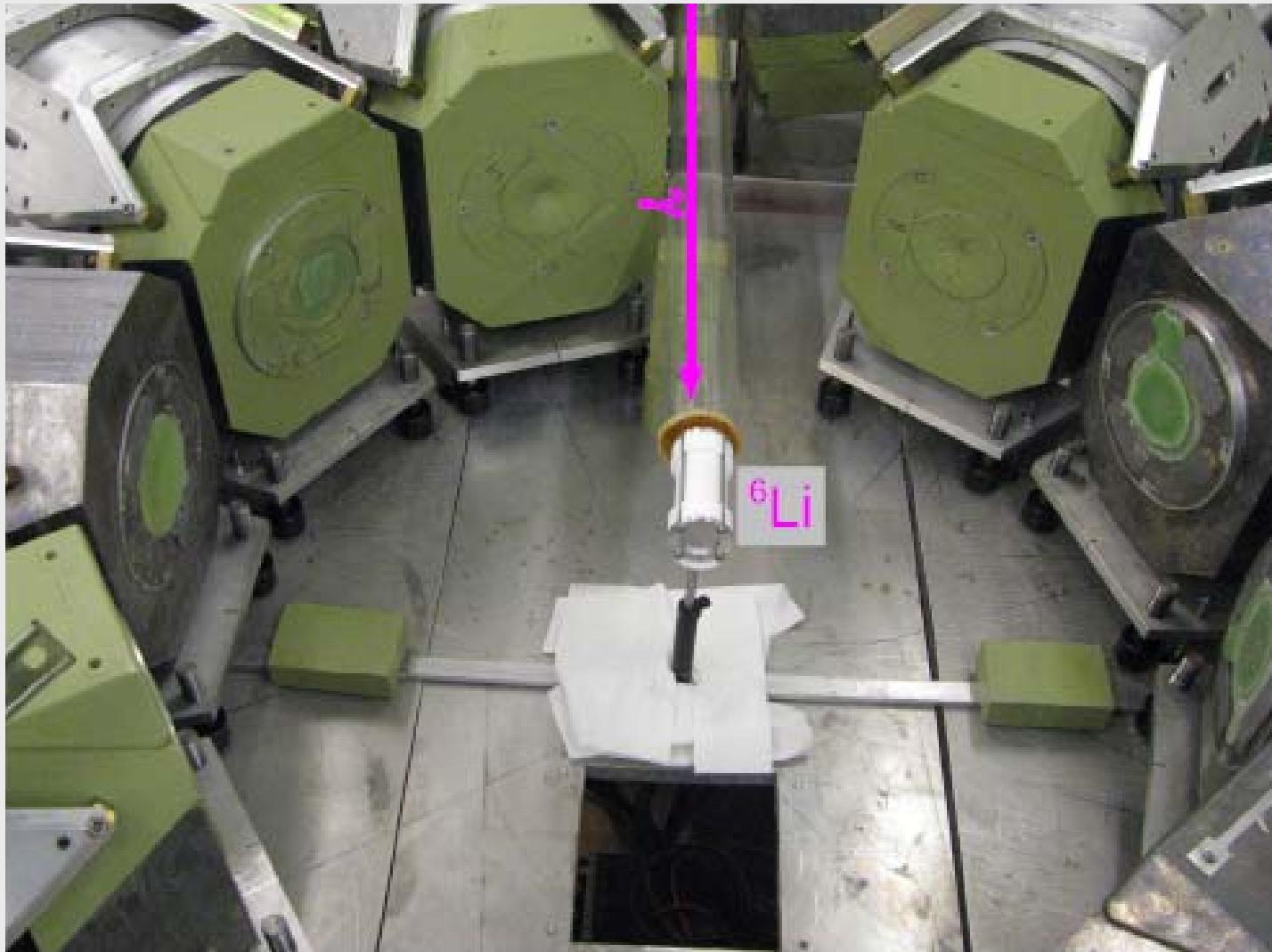


# HINDA Array



# Experimental Setup

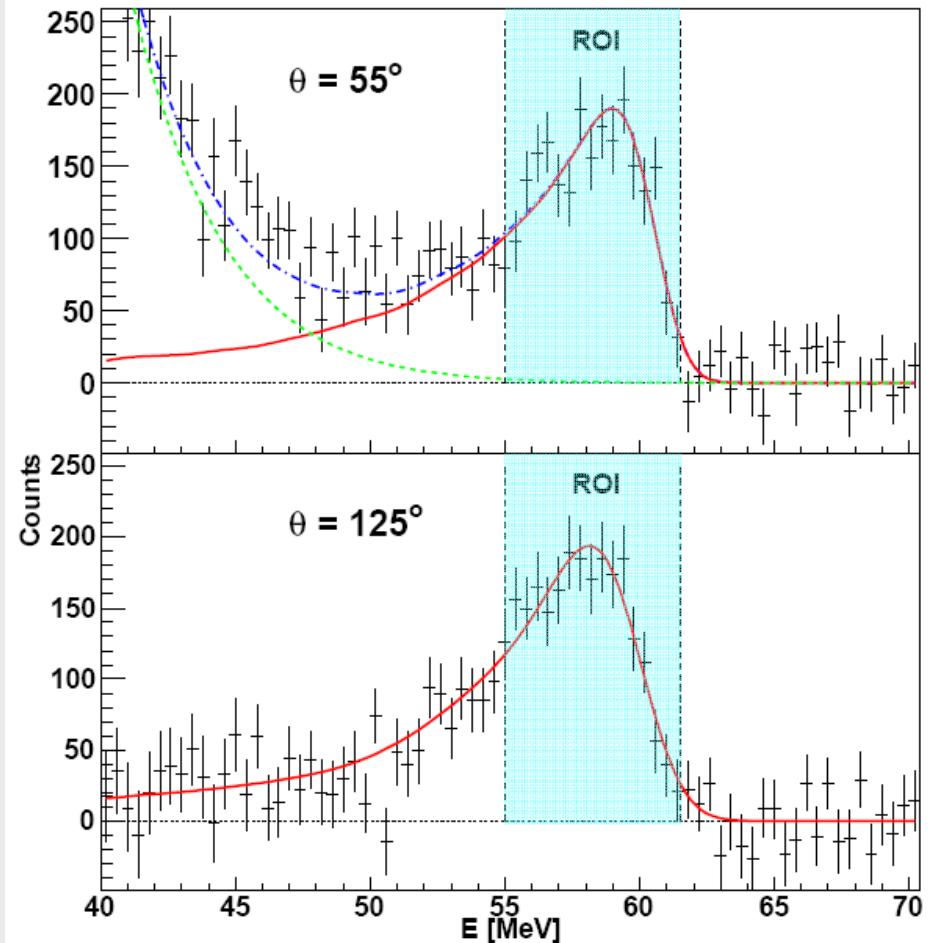
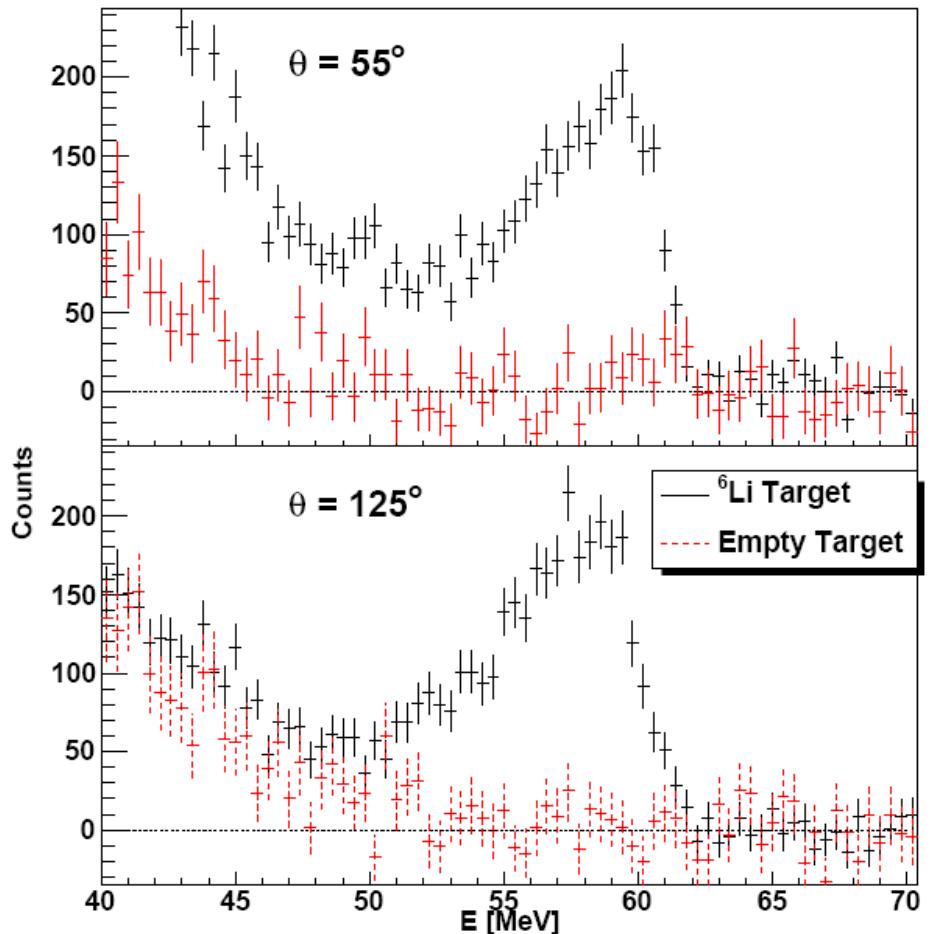
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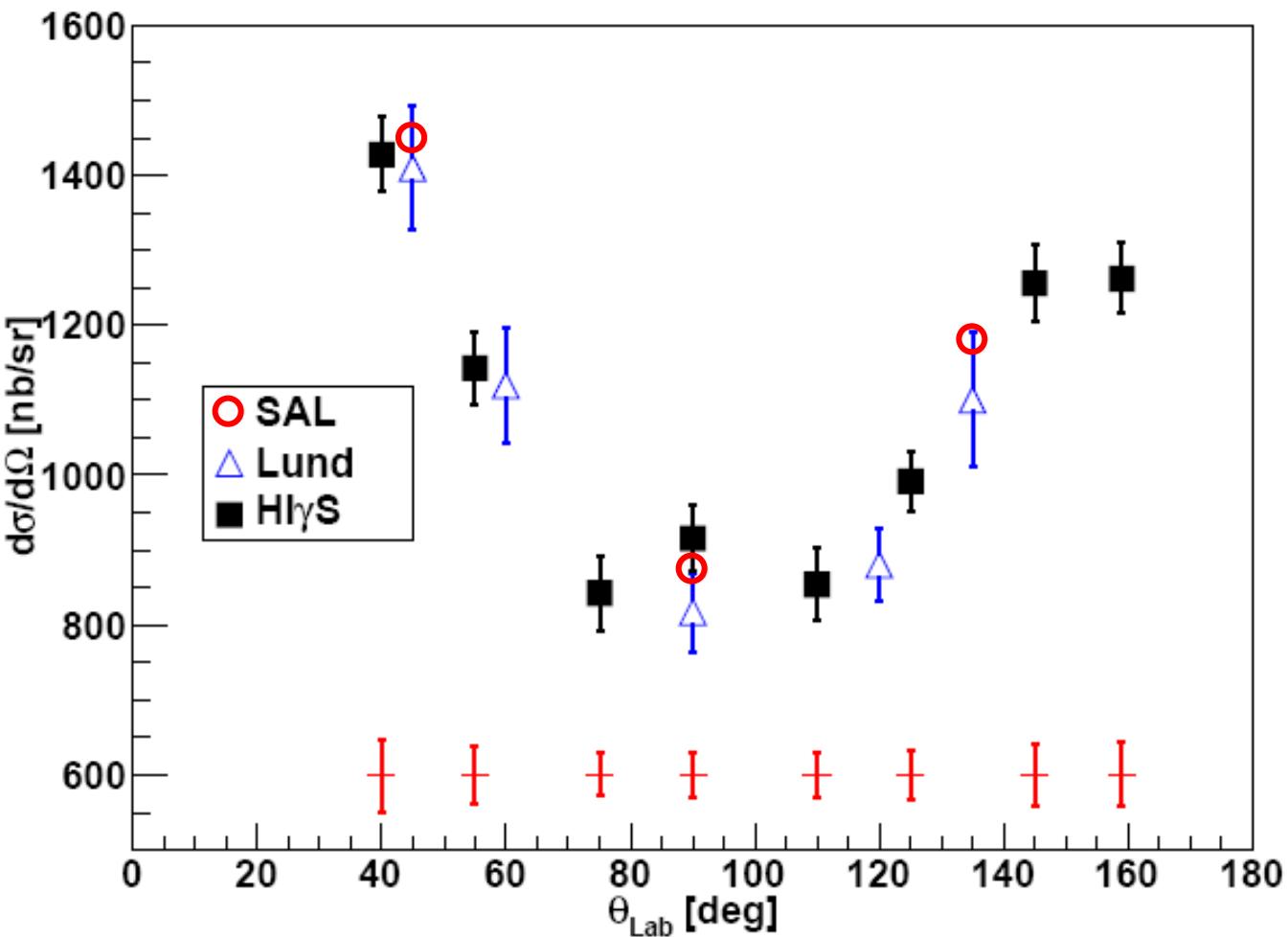
# Sample Spectra

${}^6\text{Li}(\gamma, \gamma){}^6\text{Li}$

$E_\gamma = 60 \text{ MeV}$

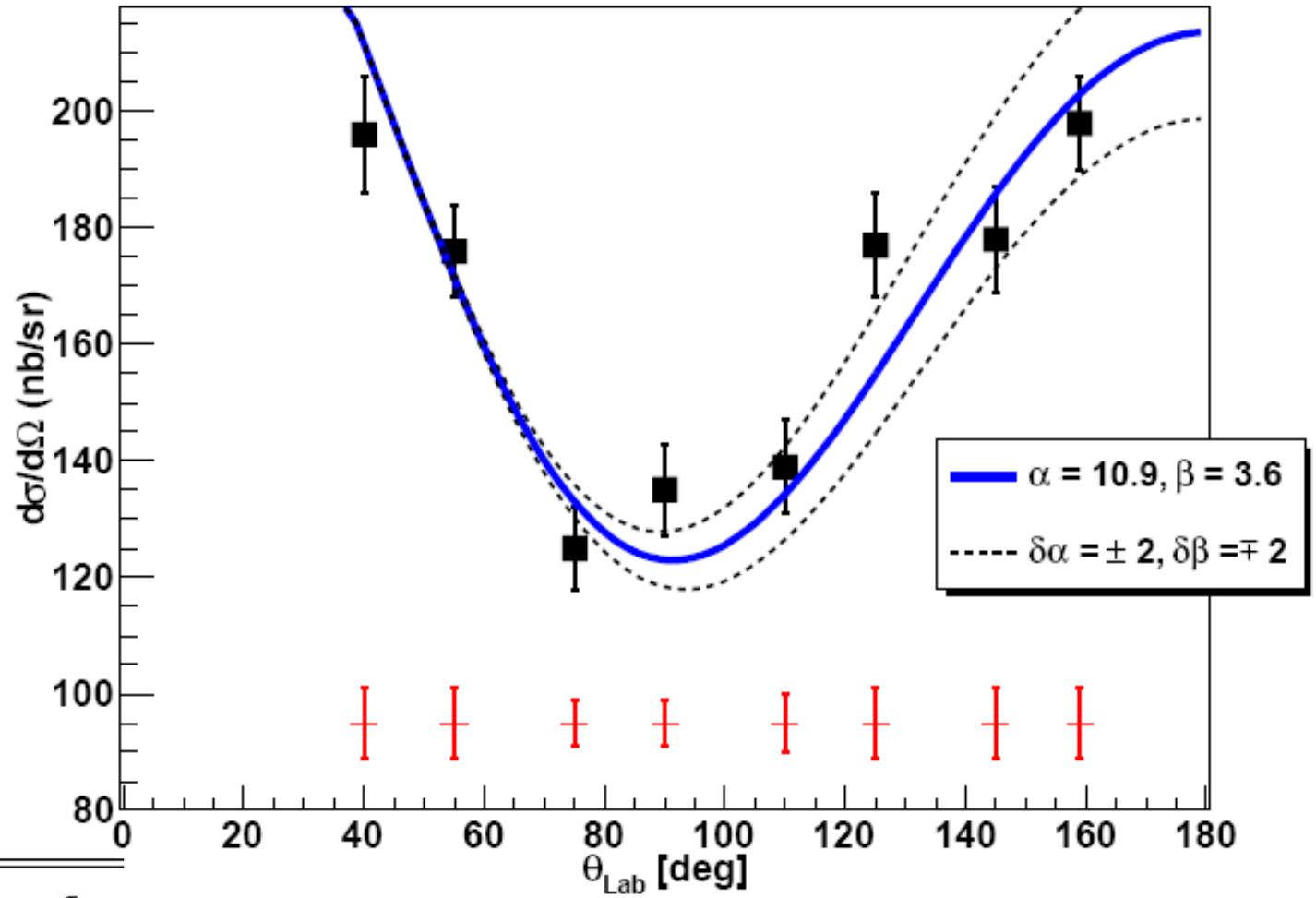


# Cross Section for $^{16}\text{O}(\gamma, \gamma)^{16}\text{O}$



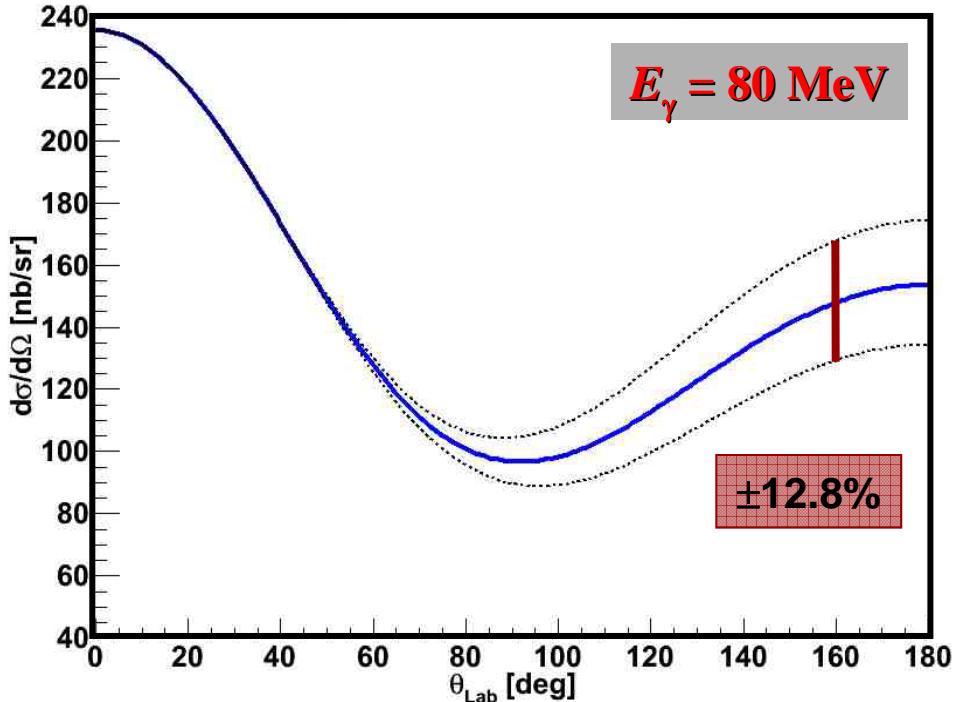
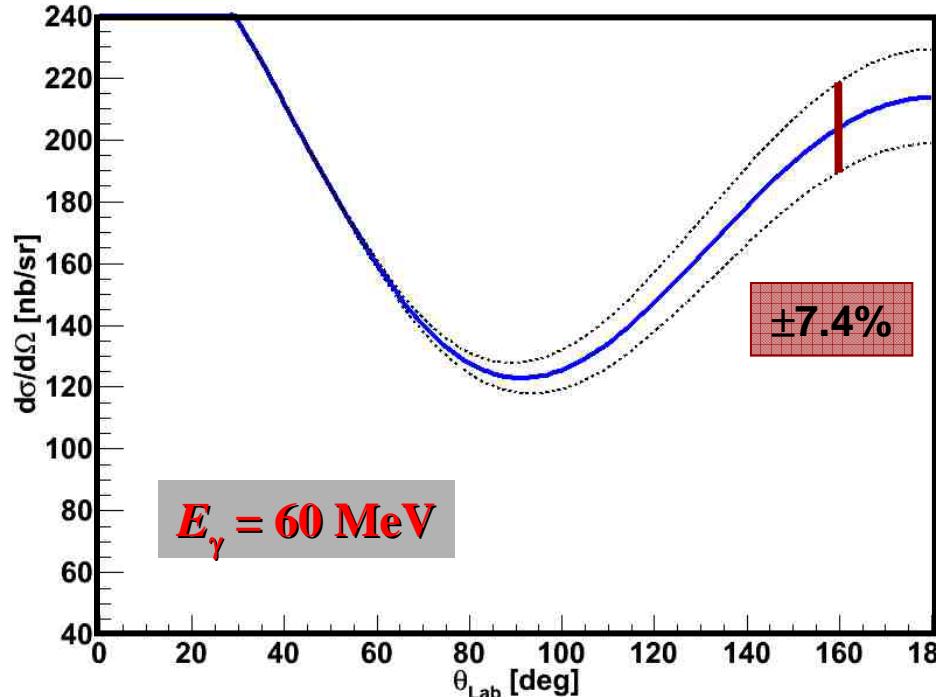
# Cross Section for ${}^6\text{Li}(\gamma, \gamma){}^6\text{Li}$

L. Myers *et al.*  
Phys. Rev. C86  
(2012)



Resonance	$E_{res}$ (MeV)	$\Gamma_{res}$ (MeV)	$\sigma_{res}$ (mb)
$E1$	25.0	6.0	7.2
$E2$	34.0	8.0	0.14
$QD$	40.0	100	1.1

**sum rule:**  $\alpha + \beta = 14.5$

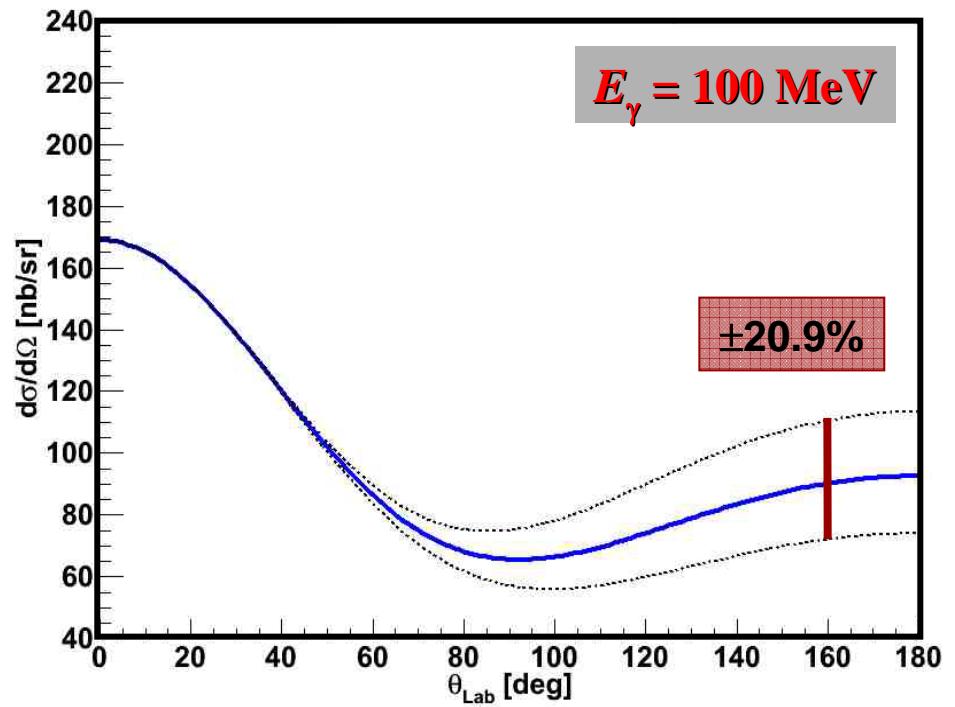


$$(\alpha, \beta) = (10.9, 3.6)$$

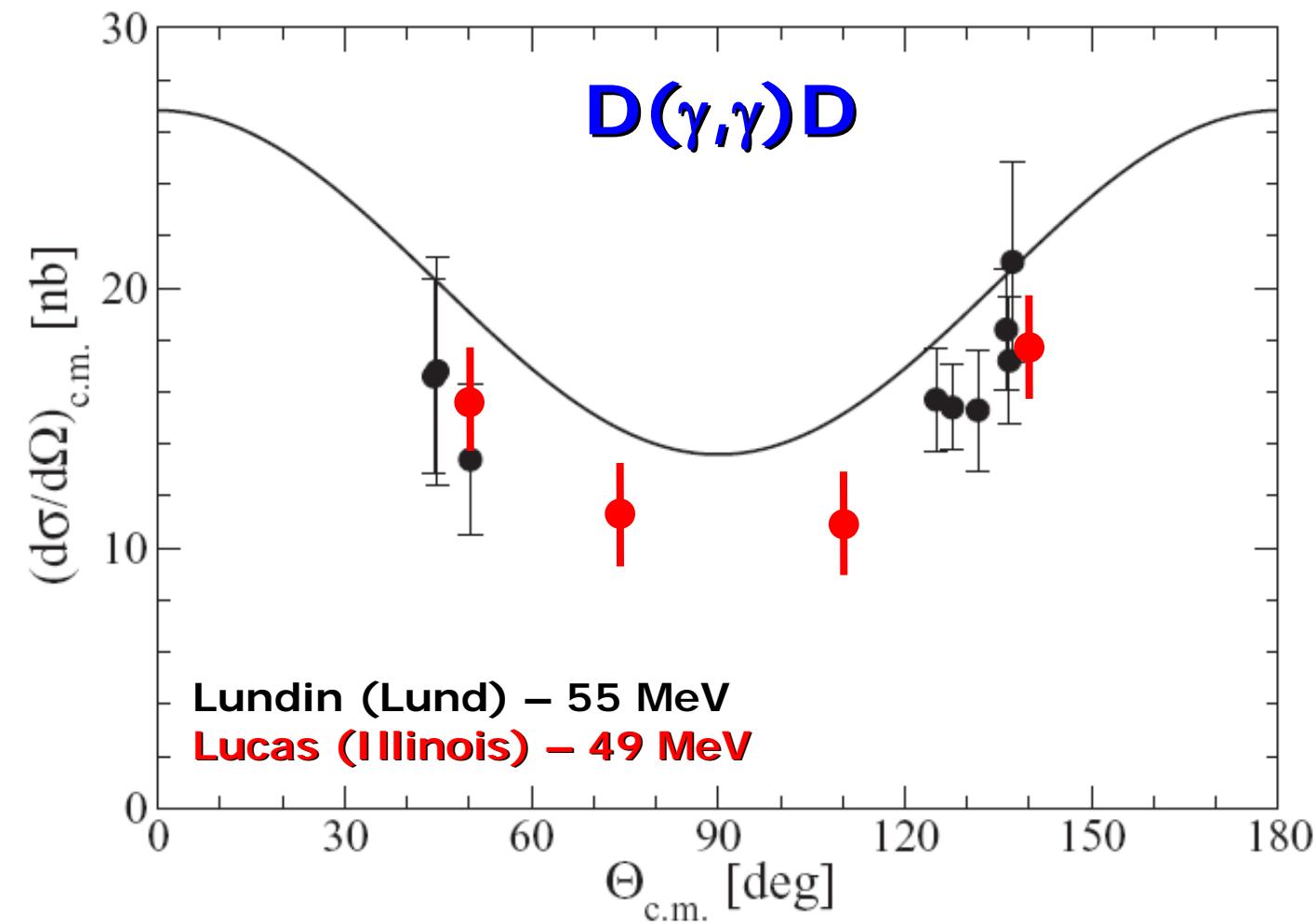
$$\Delta\alpha = \pm 2$$

$$\Delta\beta = \mp 2$$

$$\text{sum rule: } \alpha + \beta = 14.5$$



# LIT Method for Compton Scattering



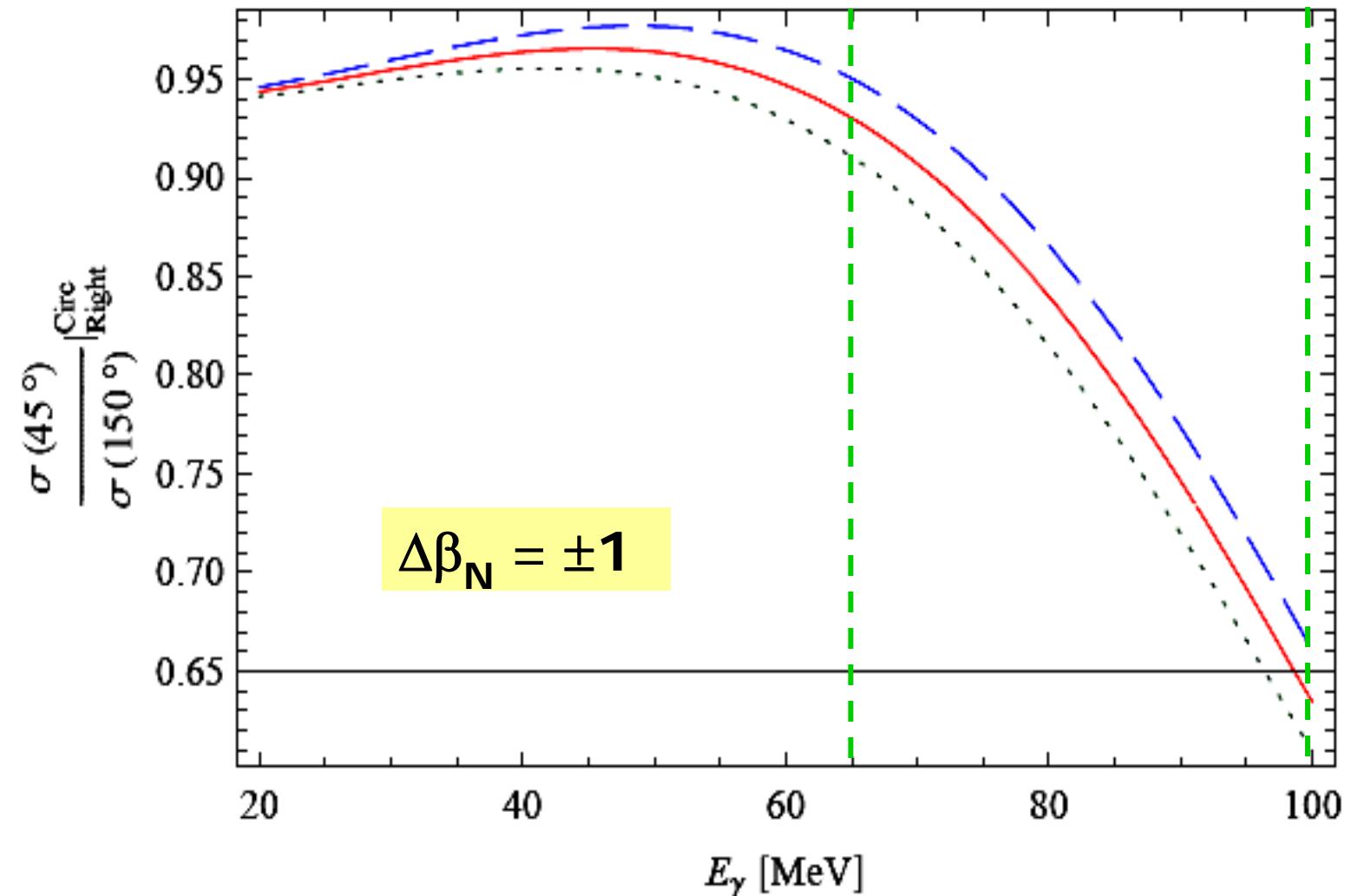
# **Compton Scattering on the Proton and Deuteron**

# Cross-Section Ratios for Deuterium

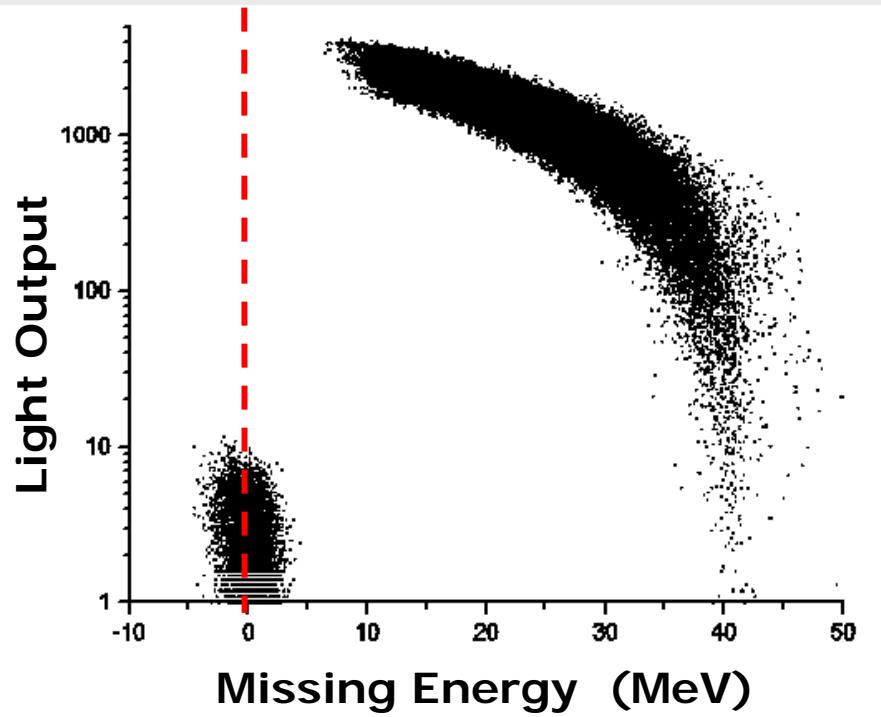
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- **proposal by Henry Weller**
- unpolarized photon beam and unpolarized deuterium tgt
  - scintillating active target (detect recoils in coincidence)
- scattering angles  $45^\circ, 80^\circ, 115^\circ, 150^\circ$  ( $E_\gamma = 65, 100$  MeV)
- requires 300 hrs (65 MeV) + 100 hrs (100 MeV)
- detectors: eight  $10'' \times 12''$  NaI's
  - arranged symmetrically left/right

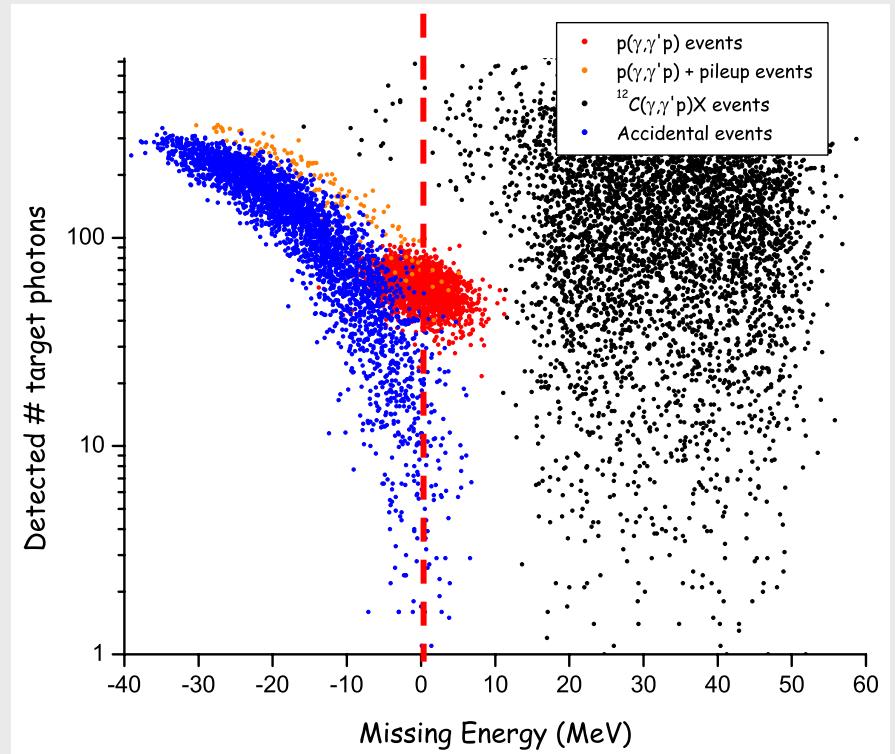
# Cross-Section Ratios for Deuterium



# Compton Scattering with scintillating target



deuteron



proton

simulations: R. Miskimen

# Sum-Rule-Independent Measurement of $\alpha_p$

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- **proposal by Mohammad Ahmed**
- linearly polarized photon beam (unpolarized target)
  - scintillating active target (detect recoils in coincidence)
- measure scattered photons at  $90^\circ$  ( $E_\gamma = 82$  MeV)
  - scattering cross section is independent of  $\beta_p$
  - extraction of  $\alpha_p$  is independent of the Baldin sum rule
  - extraction of  $\alpha_p$  is model-independent
- requires 300 hrs for 5% uncertainty in  $\alpha_p$
- detectors: four  $10'' \times 12''$  NaI's (HINDA array)
  - located left, right, up, down

# Sum-Rule-Independent Measurement of $\alpha_p$

$$\left\{ \begin{array}{l} \frac{d\sigma_{\perp}}{d\Omega} = \frac{d\sigma_{\perp}}{d\Omega}(point) - \underbrace{\left[ \frac{e^2}{Mc^2} \right] \left[ \frac{\omega'}{\omega} \right]^2 \omega \omega' (2\bar{\alpha} + 2\bar{\beta} \cos \theta)}_{d\sigma_{\perp}^{dipole}} \\ \\ \frac{d\sigma_{\parallel}}{d\Omega} = \frac{d\sigma_{\parallel}}{d\Omega}(point) - \underbrace{\left[ \frac{e^2}{Mc^2} \right] \left[ \frac{\omega'}{\omega} \right]^2 \omega \omega' (2\bar{\alpha} \cos^2 \theta + 2\bar{\beta} \cos \theta)}_{d\sigma_{\parallel}^{dipole}} \end{array} \right.$$

- The  $90^\circ$  parallel cross section is independent of  $\bar{\alpha}$  and  $\bar{\beta}$ :  $\frac{d\sigma_{\parallel}}{d\Omega}(90^\circ) = \left( \frac{d\sigma_{\parallel}}{d\Omega}(90^\circ) \right)_{point}$
- The cross section asymmetry  $A(\omega, \theta) = \frac{d\sigma_{\perp} - d\sigma_{\parallel}}{d\sigma_{\perp} + d\sigma_{\parallel}}$

$$A(\omega, 90^\circ) = \frac{C_1(\omega) - C_2(\omega)\bar{\alpha}}{C_3(\omega) + C_2(\omega)\bar{\alpha}}$$

# Nucleon Spin Polarizability

$$\begin{aligned}\mathcal{L}_{\text{pol}} = & 2\pi N^\dagger [\alpha_{E1}(\omega) \vec{E}^2 + \beta_{M1}(\omega) \vec{B}^2 + \\ & + \gamma_{E1E1}(\omega) \vec{\sigma} \cdot (\vec{E} \times \dot{\vec{E}}) + \gamma_{M1M1}(\omega) \vec{\sigma} \cdot (\vec{B} \times \dot{\vec{B}}) \\ & - 2\gamma_{M1E2}(\omega) \sigma_i B_j E_{ij} + 2\gamma_{E1M2}(\omega) \sigma_i E_j B_{ij} + \dots] N\end{aligned}$$

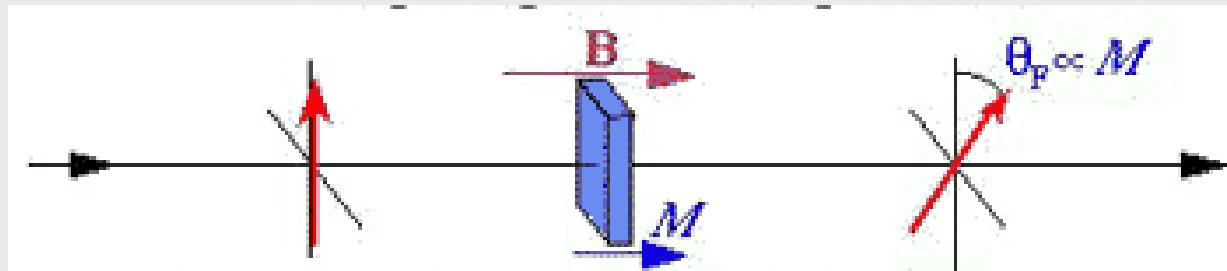
$$\gamma_0(\omega) = -\gamma_{E1E1}(\omega) - \gamma_{M1M1}(\omega) - \gamma_{E1M2}(\omega) - \gamma_{M1E2}(\omega)$$

$$\gamma_\pi(\omega) = -\gamma_{E1E1}(\omega) + \gamma_{M1M1}(\omega) - \gamma_{E1M2}(\omega) + \gamma_{M1E2}(\omega)$$

**forward and backward spin polarizabilities**

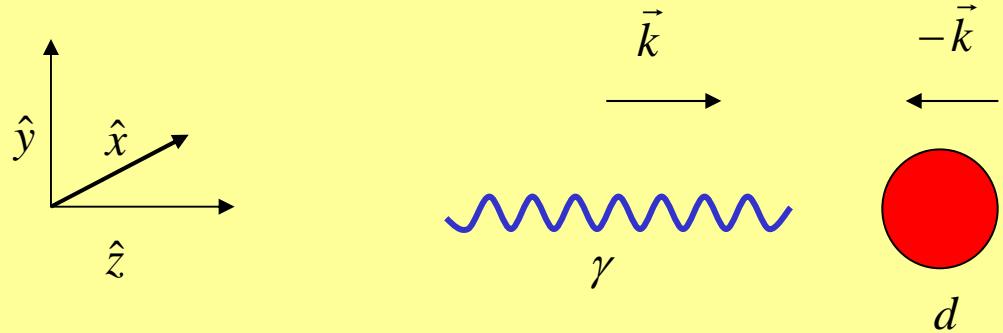
# Nucleon Spin Polarizability

- classical analogy: Faraday rotation of linearly polarized light in a spin-polarized medium



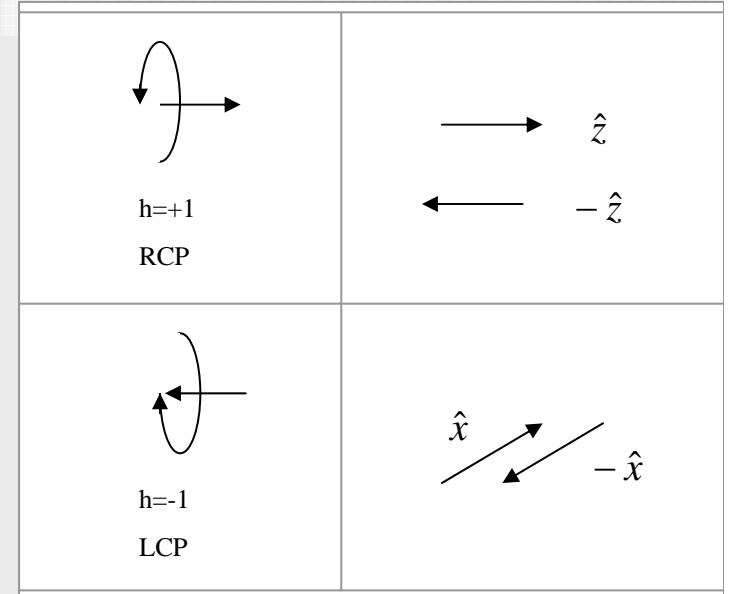
- four spin polarizabilities:  $\gamma_1, \dots, \gamma_4$ 
  - forward spin polarizability:  $\gamma_0 = \gamma_1 - \gamma_2 - 2\gamma_4$
  - backward spin polarizability:  $\gamma_\pi = \gamma_1 + \gamma_2 + 2\gamma_4$
- expt. asymmetries with circularly polarized photons
  - $\Sigma_x$ : target spin  $\perp$  photon helicity (in reaction plane)
  - $\Sigma_z$ : target spin parallel to photon helicity

# Double Polarization Observables



$$\Sigma_z = \left[ \frac{\left( \frac{d\sigma}{d\Omega} \right)_{\uparrow\uparrow} - \left( \frac{d\sigma}{d\Omega} \right)_{\uparrow\downarrow}}{Sum} \right]_{h=\pm 1}$$

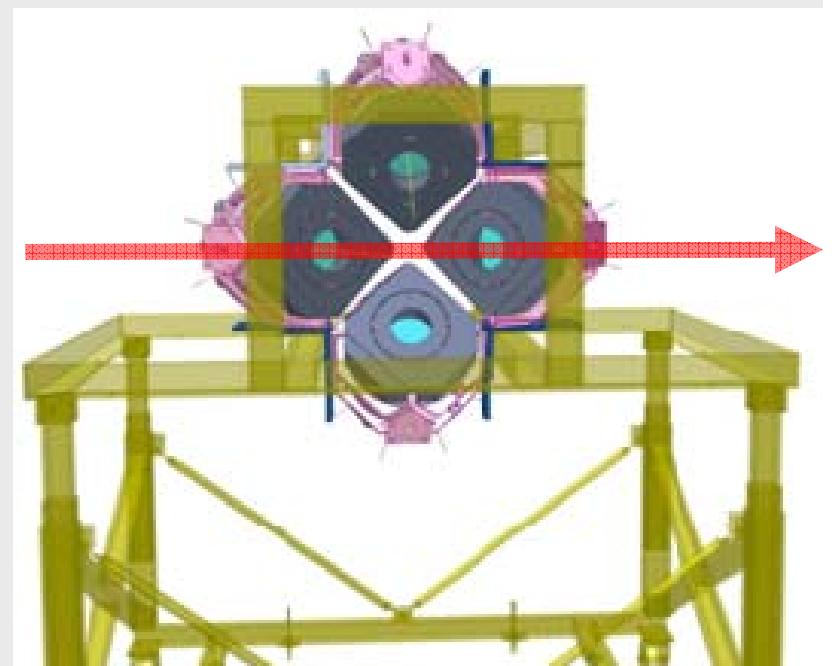
$$\Sigma_x = \left[ \frac{\left( \frac{d\sigma}{d\Omega} \right)_{\rightarrow\rightarrow} - \left( \frac{d\sigma}{d\Omega} \right)_{\rightarrow\leftarrow}}{Sum} \right]_{h=\pm 1}$$



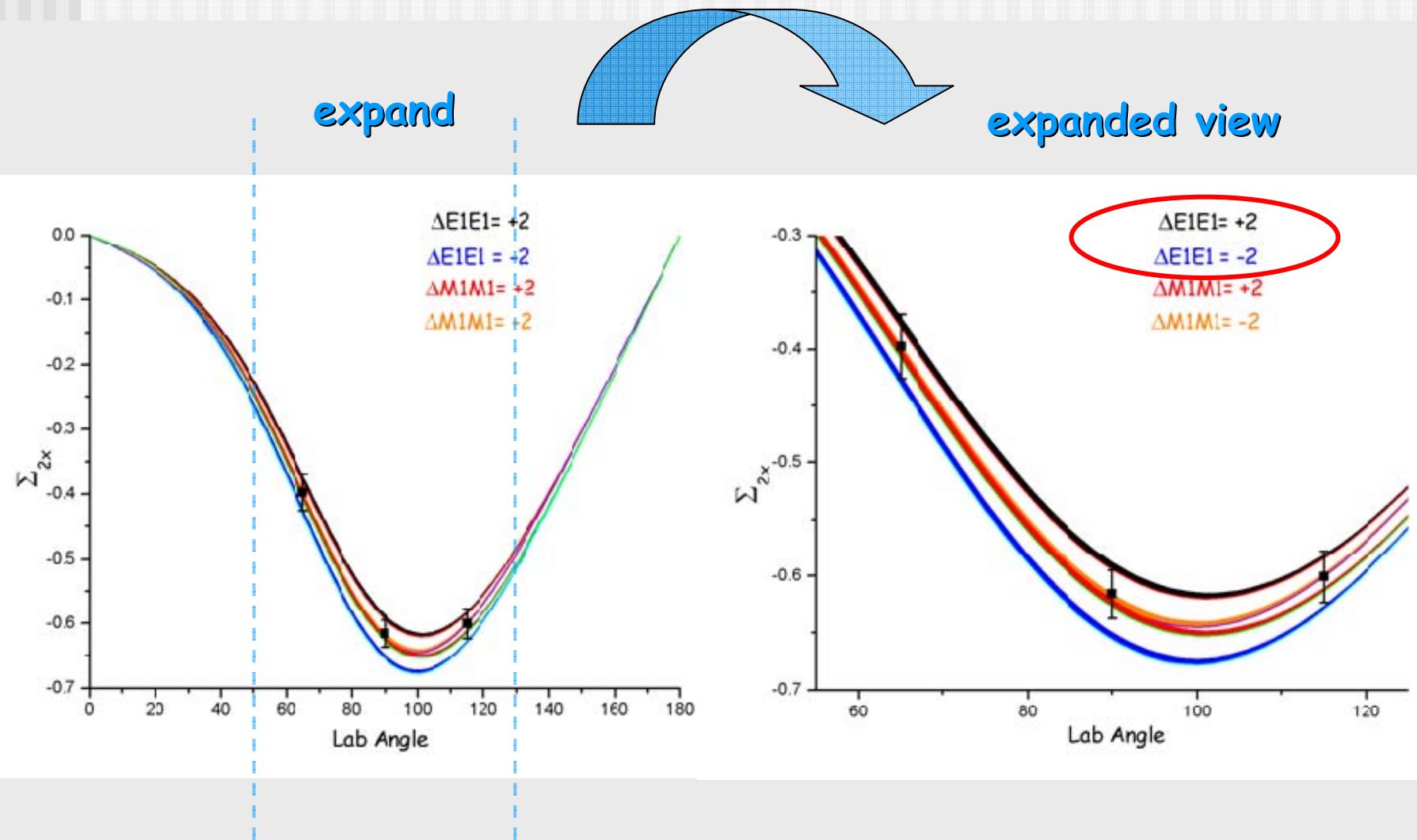
**circularly polarized beam**  
**vector polarized target**

# Spin Polarizabilities of the Proton

- ❑ proposal by Rory Miskimen
- ❑ first measurement of proton  $\gamma_{E1E1}$
- ❑ circularly polarized photon beam
  - scintillating active ***transverse polarized*** target ( $P \sim 80\%$ )
- ❑ scattering angles  $65^\circ, 90^\circ, 115^\circ$  ( $E_\gamma = 100$  MeV)
- ❑ requires 800 hrs for  $\Delta\gamma_{E1E1} = \pm 1$
- ❑ detectors: eight  $10'' \times 12''$  NaI's
  - 4 in plane, 4 out of plane

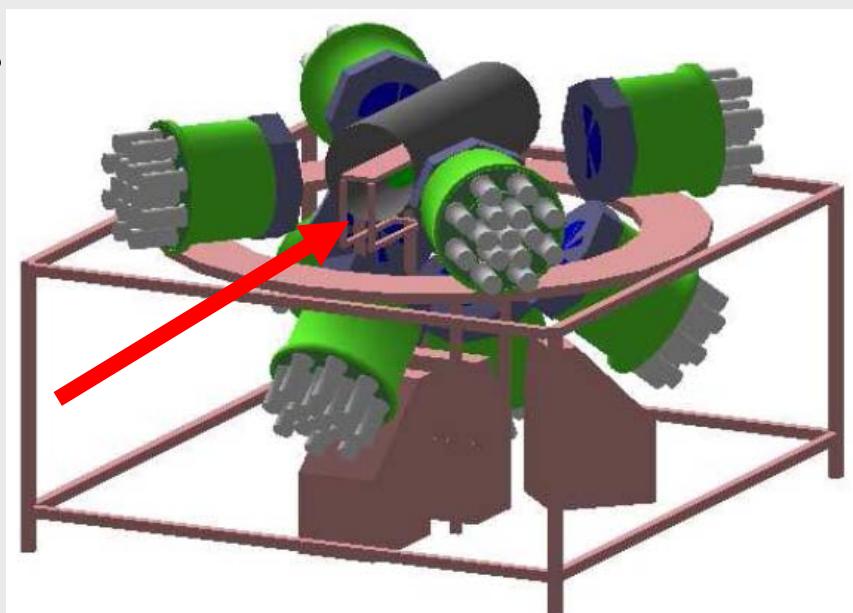


# Spin Polarizabilities of the Proton

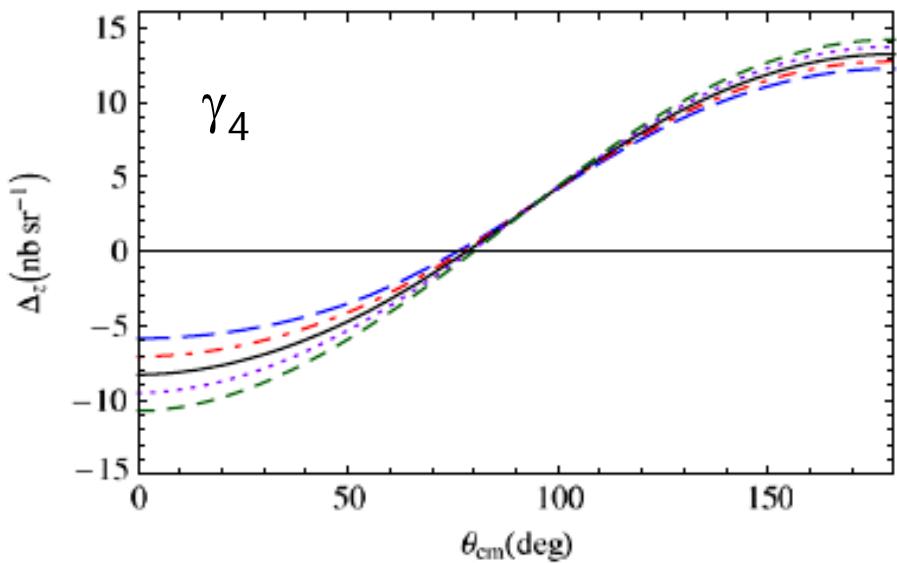
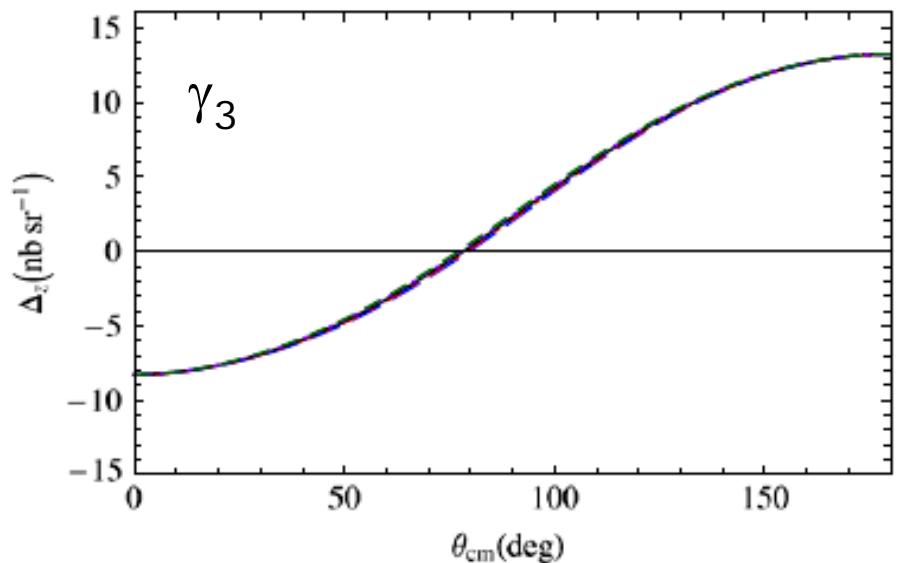
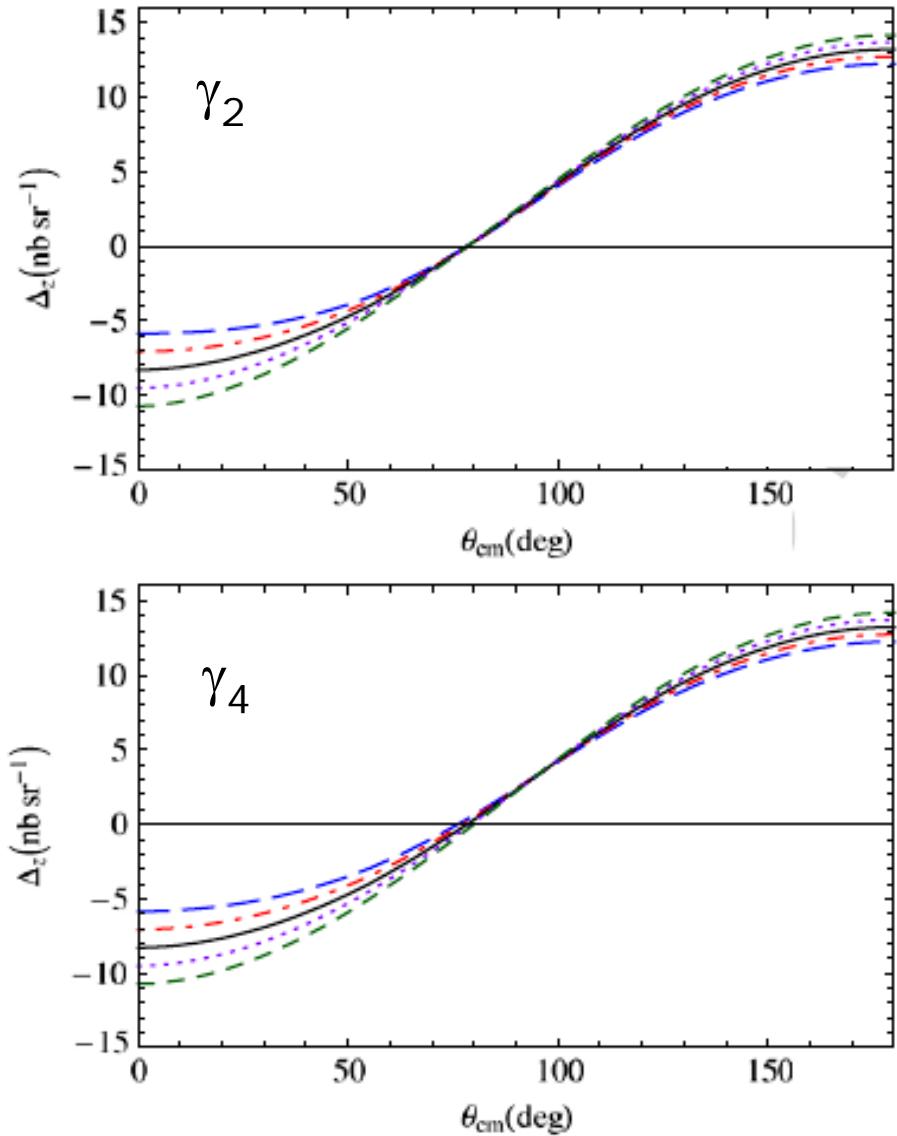
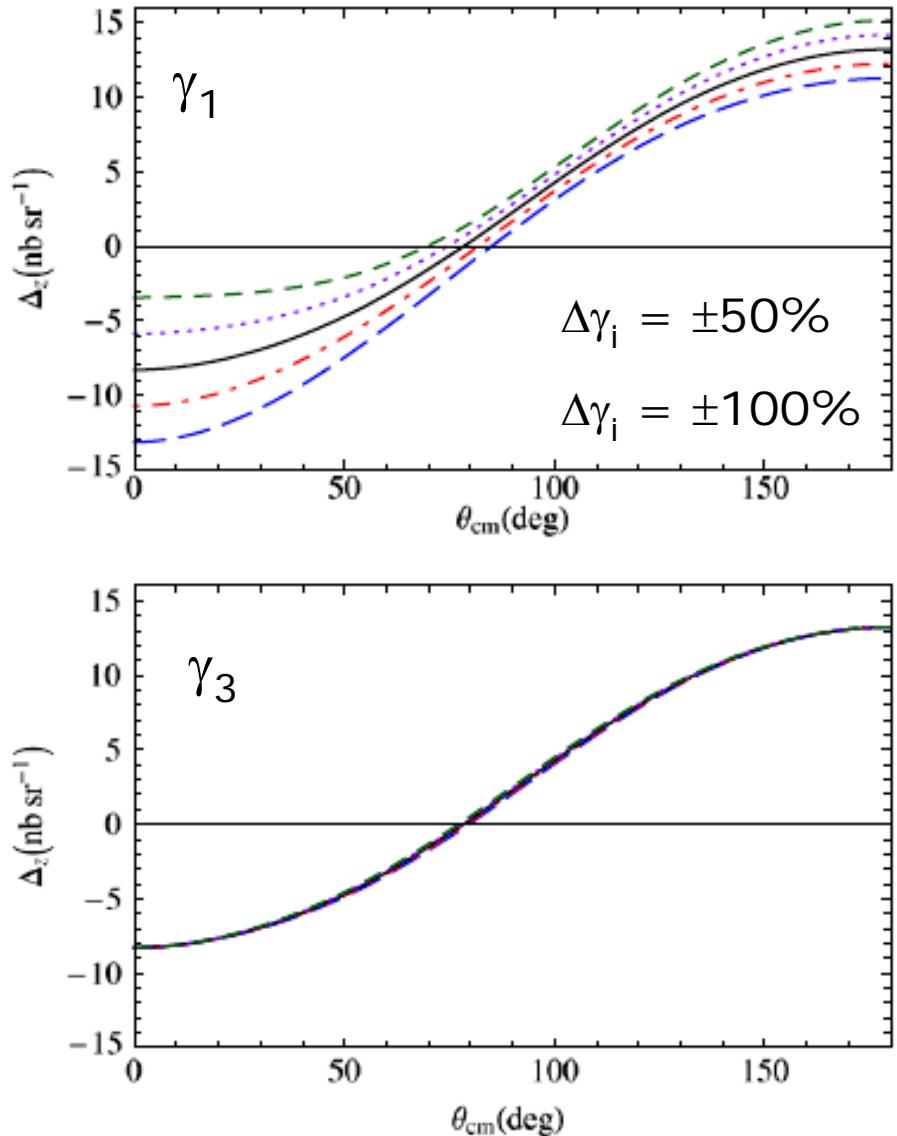


# Polarized Compton Scattering from $^3\text{He}$

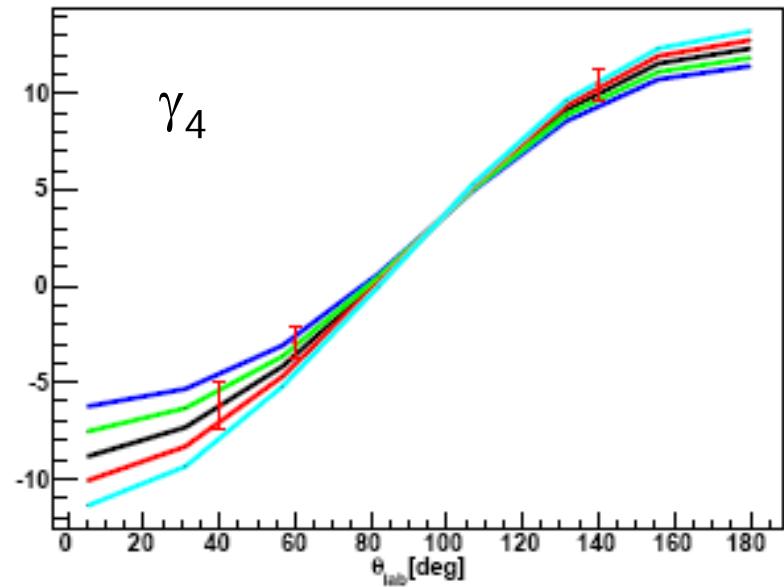
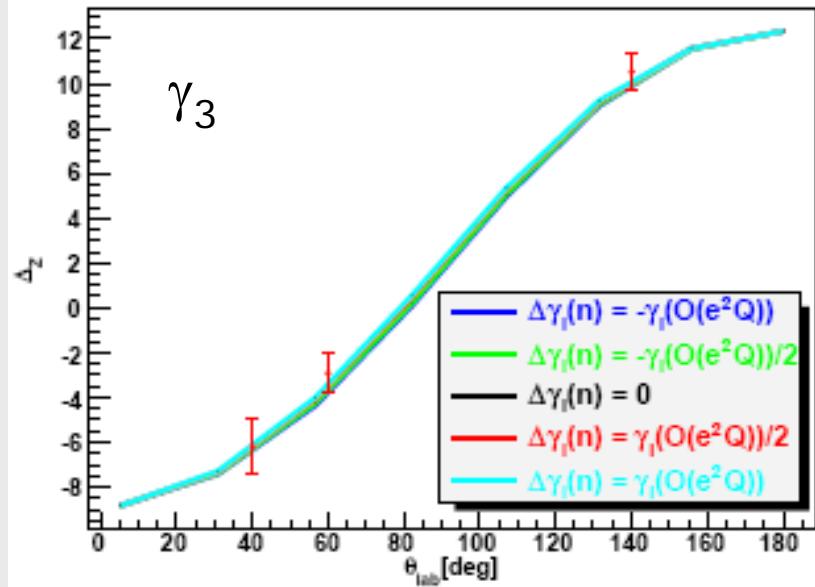
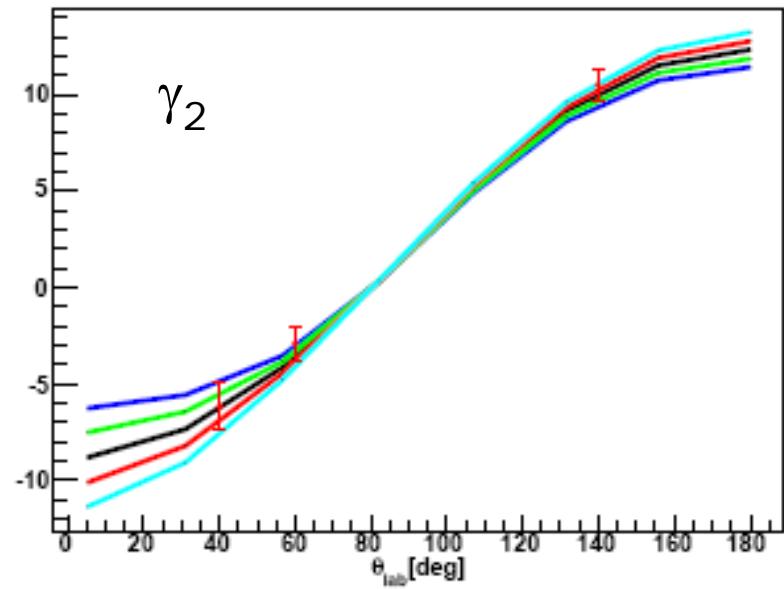
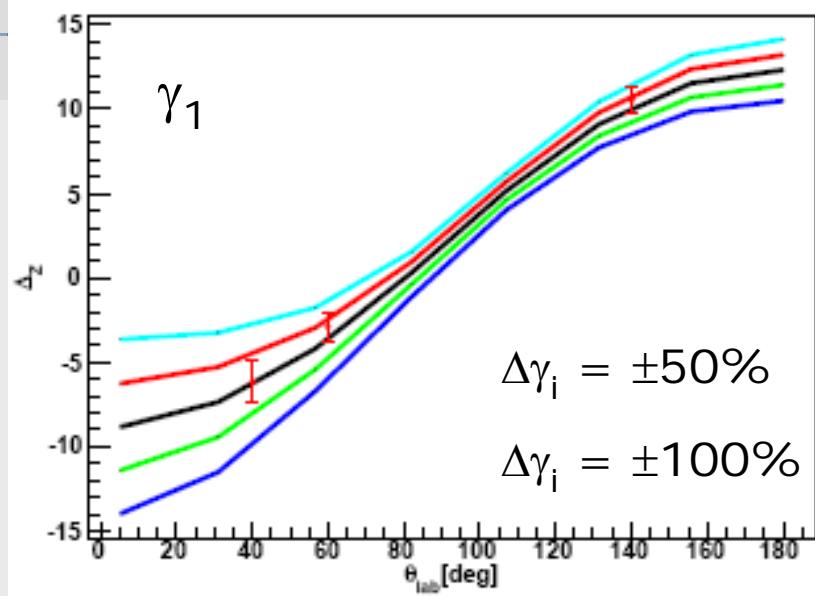
- proposal by Haiyan Gao
- first measurement of neutron spin polarizabilities
- circularly polarized photon beam
  - **longitudinally** polarized high-pressure  $^3\text{He}$  target ( $P \sim 50\%$ )
- scattering angles  $40^\circ, 60^\circ, 140^\circ$  ( $E_\gamma = 125 \text{ MeV}$ )
- requires 1200 hrs for  $\sim 20\%$  uncertainty in  $\gamma_1, \gamma_2, \gamma_4$
- detectors: eight  $10'' \times 12''$  NaI's
  - 4 in plane, 4 out of plane



# Polarized Compton Scattering from ${}^3\text{He}$



# Polarized Compton Scattering from ${}^3\text{He}$





# HIGS Capabilities for User Programs in 2012

Parameter	Value		Comments
E-beam Configuration	Symmetric two-bunch beam		High flux configuration
E-beam current [mA]	50 - 120		
Gamma-ray Energy [MeV]	1 – 100		with mirrors 1064 to 190 nm Available with existing hardware Extending wiggler current to 3.5 kA
(a) No-loss mode	Total flux [ $\gamma/s$ ]	Collimated flux ( $\Delta E/E \sim 5\%$ ) [ $\gamma/s$ ]	Both Horizontal and Circular Polarizations
1 – 3 MeV <sup>(a)</sup>	$1 \times 10^8$ – $1 \times 10^9$	$6 \times 10^6$ – $6 \times 10^7$	
3 – 5 MeV	$6 \times 10^8$ – $2 \times 10^9$	$3.6 \times 10^7$ – $1.2 \times 10^8$	
5 – 13 MeV	$4 \times 10^8$ – $4 \times 10^9$	$2.4 \times 10^7$ – $2.4 \times 10^8$	
13 – 20 MeV	$1 \times 10^9$ – $2 \times 10^9$	$6 \times 10^7$ – $1.2 \times 10^8$	
(b) Loss mode	Total flux [ $\gamma/s$ ]	Collimated flux ( $\Delta E/E \sim 5\%$ ) [ $\gamma/s$ ]	To extend mirror lifetime, circular polarization is preferred 1 <sup>st</sup> user experiment: March, 2011 190 nm, 1 <sup>st</sup> user experiment in 2013
21 – 54 MeV	$> 2 \times 10^8$ <sup>(b)</sup>	$> 1 \times 10^7$	
55 – 65 MeV	$\sim 2 \times 10^8$ <sup>(b)</sup>	$\sim 1 \times 10^7$	
<b>66 – 100 MeV</b>	<b><math>1 \times 10^8</math><sup>(b) (c)</sup></b>	<b><math>5 \times 10^6</math></b>	

<sup>(a)</sup> With present configuration of OK-5 wigglers separated by 21 m, the circular polarization is about  $\frac{1}{2}$  the values here.

<sup>(b)</sup> The flux in loss mode is mainly limited by injection rate.

<sup>(c)</sup> Thermal stability of FEL mirror may limit the maximum amount of current can be used in producing FEL lasing, thus flux.

**Highest Total Flux (2010):  $> 2 \times 10^{10} \gamma/s$  @ 9 – 11 MeV**

→ Projected energy upgrades above 100 MeV (circular pol) ←

2014: 110 – 120 MeV,  $\sim 1 \times 10^8$  g/s (total)

\* After 2015: 150+ MeV,  $\sim 1 \times 10^8$  g/s (total)

# Summary

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- **Early measurements of Compton scattering at HIGS**
  - unpolarized expt. on  ${}^6\text{Li}$
  - polarized expts. on  $A = 89\text{-}238$  for IVGQR studies
  - new proposal: low-energy (3-15 MeV) for  ${}^6\text{Li}$  nuclear pol.
- **Next generation of expts. on proton and deuteron**
  - electric and magnetic polarizability
    - ✓ unpolarized expt. on deuterium
    - ✓ **polarized** expt. on proton
  - spin polarizability
    - ✓ **double-polarized** expt. on proton
    - ✓ **double-polarized** expt. on  ${}^3\text{He}$
- **HIGS can contribute high-quality polarized data!**
  - stay tuned for further developments in the future...