Noise, sign problems and chiral symmetry breaking

Noise, sign problems, and statistics arXiv:1106.0073 [hep-lat] Michael Endres, D.K., Jong-Wan Lee, Amy Nicholson

Work in progress (!) Dorota Grabowska, D.K., Amy Nicholson

Physics motivation: can't we get beyond this cartoon??

Sign problem!



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This talk:

- From sign problem to noise in correlation functions
- The relation between noise and the pion
- Computing the distribution of noise in a model with chiral symmetry breaking
- Wild and unjustified speculations

The "sign" problem in the grand canonical approach: $Det(\not D + \mu \gamma^0)$ complex

- physics happens for µ≥ m_N/3...
- ...but sign problem starts at $\mu=m_{\pi}/2$!



P.E. Gibbs, 1986

Explanation (2-flavor QCD):

 $|\text{Det}(\not p + \mu \gamma^0)| \approx \underline{\text{isospin}}$ chemical potential

Role of phase: eliminate pion condensate for $\mu \ge m_{\pi}/2!$



Martin Savage's generalization of Lepage argument:

Consider an operator X(t) = 2 Re[baryon correlator for time t]

Then leading time behavior:

$$\langle X(t)^k \rangle \sim \begin{cases} e^{-\left(M_B + \frac{3(k-1)}{2}m_\pi\right)t} & k \text{ odd} \\ e^{-\left(\frac{3k}{2}m_\pi\right)t} & k \text{ even} \end{cases}$$

So one expects ($\kappa_n \equiv n^{\text{th}}$ cumulant):

skewness
$$= \frac{\kappa_3}{(\kappa_2)^{\frac{3}{2}}} \sim e^{-(M_B - \frac{3}{2}m_\pi)t}$$

kurtosis $= \frac{\kappa_4}{(\kappa_2)^2} \sim 1$

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QCD data (M. Savage, NPLQCD): Un-averaged A baryon correlator data

time-slice = 0



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Almost symmetric, small mean
"sign problem" big cancellations

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Challenge:

- Can we understand the correlator distribution quantitatively?
- Can we relate the distribution to the physical spectrum of the theory?
- Is it possible to use this knowledge to ameliorate the sign problem and study QCD at finite density or large baryon number?
 - either use knowledge of distribution + statistical methods to improve signal/noise, or
 - reformulate the theory in a way to improve the distribution

Some reason for optimism:

- experience simulating cold atoms at a Feshbach resonance
- analytical work (in progress) in a QFT with chiral symmetry breaking



Unitary fermions: nonrelativistic fermions with zero-range interaction, tuned to infinite scattering length (conformal system)

M. Endres, D.K., J.W. Lee, A. Nicholson, arXiv:1106.5725 [hep-lat] Lattice model:

- Short-range momentum-dependent 4-fermion interaction induced by auxiliary scalar field
- Interaction tuned to conformal fixed pt.



- Simulated up to \sim N=70 fermions on 14³ x 64 lattice
- •~1% accuracy in energies
- ~2 billion configurations for auxiliary field

Effective mass plot: 46 fermions, $V=12^3$, 40M configurations



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FOUND: Large-N NJL model in d=3

Fascinating model because it has two equivalent formulations: one with a QCD-like sign problem, one without!

Large-N NJL model in d=3

$$\mathcal{L} = N\left(\overline{\psi}_a(\partial \!\!\!/ - m)\psi_a - \frac{C}{2}\left[(\overline{\psi}_a\psi_a)^2 + (\overline{\psi}_a i\gamma_5\psi_a)^2\right]\right)$$

- a=1,...,N
- d=4 theory dimensionally reduced to d=3
- ψ = 4 component spinor (like d=4)
- γ-matrices = 4x4 (like d=4)
- Symmetry = $U(N)_V \times U(1)_A$ (approximate if m≠0)

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Two equivalent formulations with auxiliary fields:

$$\mathcal{L} = N\left(\frac{1}{2C}\left(\sigma^{2} + \pi^{2}\right) + \overline{\psi}_{a}\left[\partial - m + \sigma + i\pi\gamma_{5}\right]\psi_{a}\right) \qquad \mathbf{O}_{\mathbf{h}}$$
$$\mathcal{L} = N\left(\frac{1}{C}\operatorname{Tr}\left(V_{\mu}V_{\mu} + A_{\mu}A_{\mu}\right) + \overline{\psi}_{a}\left[\partial + i\psi + A\gamma_{5}\right]_{ab}\psi_{b}\right) \qquad \mathbf{A}_{\mathbf{h}}$$
$$(obtained by Fierz rearrangement of 4-fermion interaction)$$

For even N, the σ/π formulation has no sign problem at $\mu \neq 0$: $(\partial \!\!\!/ + \sigma + i\pi\gamma_5 + \mu\gamma_1)^* = C(\partial \!\!\!/ + \sigma + i\pi\gamma_5 + \mu\gamma_1)C$ $\det(\partial \!\!\!/ + \sigma + i\pi\gamma_5 + \mu\gamma_1)^N = \text{real, positive}$

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...but the equivalent A/V formulation has a QCD-like sign problem: $\det \left[\partial \!\!\!/ + i V \!\!\!/ + A \!\!\!/ \gamma_5 + \mu \gamma_1 \right]^N = \text{complex!}$

(Can give a Splittorff-Verbaarschot argument for why it is complex, relating phase fluctuations of determinant to pion condensation in 2N flavor theory, just as in QCD.)

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Same theory, different formulation! As we will see: having an explicit pion field cures the sign problem

Look at these two formulations and compute the probability distribution for measurements of a single fermion correlator.

Probability distribution for an observable $X(\phi)$ -- e.g. a correlation function -- where ϕ is a stochastic field:

$$\mathcal{P}(x) = \mathcal{N} \int [\mathrm{d}\phi] \, e^{-S[\phi]} \, \delta(X[\phi] - x)$$

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Better: look at the characteristic function $\Phi_X(s)$:

$$\Phi_X(s) = e^{-W(s)} = \mathcal{N} \int [\mathrm{d}\phi] \, e^{-S[\phi] + isX[\phi]}$$

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$$W(s) = \text{ sum of connected diagrams} = -\sum_n \frac{(is)^n}{n!} \kappa_n$$

$$\mathbf{n}^{\text{th}} \text{ cumulant of } \mathbf{X}$$

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I. The σ/π theory (chiral limit)

- Integrate out fermions
- Minimize effective action for mesons at leading order in $1\!/\!N$ (dim reg)
- find chiral symmetry breaking $<\sigma>=f$ at coupling $C=-\pi/f$
- Compute π,σ propagators

$$G_{\sigma}(k) = -\left[\frac{-\frac{f}{\pi} + \frac{(4f^2 + k^2)\cot^{-1}(2f/k) + 2fk}{2\pi k}}{2\pi k}\right]^{-1},$$

$$G_{\pi}(k) = -\left[\frac{k\cot^{-1}(2f/k)}{2\pi}\right]^{-1}.$$

• Compute cumulants for the single fermion correlator:

$$X_{\Gamma} = \langle \mathbf{p} = 0, t = T | \operatorname{Tr} \Gamma \left(\partial \!\!\!/ + f + \sigma + i\pi\gamma_5 \right)^{-1} | \mathbf{p} = 0, t = 0 \rangle$$

E.g. $\Gamma = (1 + \gamma_1)/2$



Can show:

- $K_n = sum of graphs with n white vertices$
- these graphs proportional to $N^{-(n-1+L)}$, L= # loops

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Leading in 1/N:



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<u>Preliminary</u> results for σ/π formulation

- K_n fall off as $e^{-nfT}/N^{(n-1)}$, f= fermion mass... pion mass doesn't appear
- No Lepage problem because no pion propagating in the "s-channel"
- 1/N corrections go as T/N: 1/N expansion breaks down at large T
- Cumulants κ'_n for log[X] fall off as N⁻⁽ⁿ⁻¹⁾
- K'_2 grows proportional to T
- We think corrections to $\kappa^{\prime}{}_{n}$ go as 1/N, not T/N

Would imply that the fermion correlator X has a long-tailed distribution, close to log-normal, similar to what was seen for unitary fermions



II. The *A*/*V* theory (chiral limit)

No mean field for chiral symmetry breaking: solve Schwinger-Dyson eq. instead (exact in large N). Power counting looks just like large-N QCD, but no nonlinear meson couplings before fermions integrated out.

To see π,σ mesons, must solve 4-pt function







$\dots = G_{\sigma}(k), G_{\pi}(k)$ found in σ/π formulation

Can now compute cumulants for the complex fermion correlator $X_{\Gamma}(T)$ in the A/V theory:

$$\Phi(s) = e^{-W(s)} = \mathcal{N} \int [\mathrm{d}\phi] \, e^{-S[\phi] + i(sX[\phi] + \bar{s}\bar{X}[\phi])}$$
$$W(s) = -\sum_{m,n=1}^{\infty} \frac{(is)^m (i\bar{s})^n}{m! \, n!} \, \kappa_{m,n}$$

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Distribution of fermion correlator at late time T

 σ/π formulation, no sign problem



A/V formulation, QCD-like sign problem



might be tractable with cumulant expansion

very hard

Same theory; σ/π formulation has explicit pion which does not behave as a fermion bound state

Speculations about QCD?

- Sign problem is not a "fermion sign problem" it is directly attributable to chiral symmetry breaking and the existence of a light pion which is a q q* bound state
- If there is a way to include a mean field for the pion in QCD without changing the theory*, sign problem might be ameliorated or solved.

* Adding fundamental pion must also involve adding meson in t-channel to make gluon attraction between qq* weaker, so pion does not appear as bound state. Looks like quark model. Can this be done exactly without causing more problems??

Main conclusion to take away:

Since there exists a theory with a QCD-like sign problem at finite density which has another formulation without a sign problem

...and we understand why...

... maybe there is hope for QCD

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